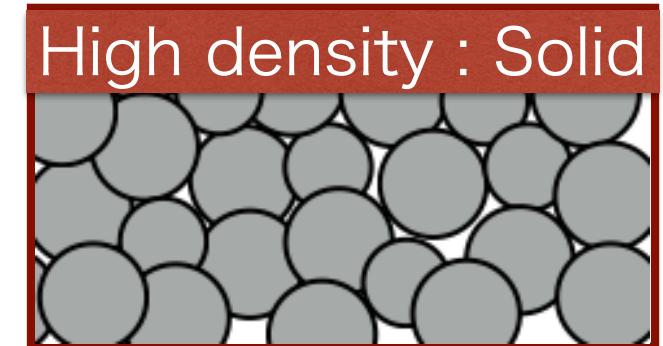
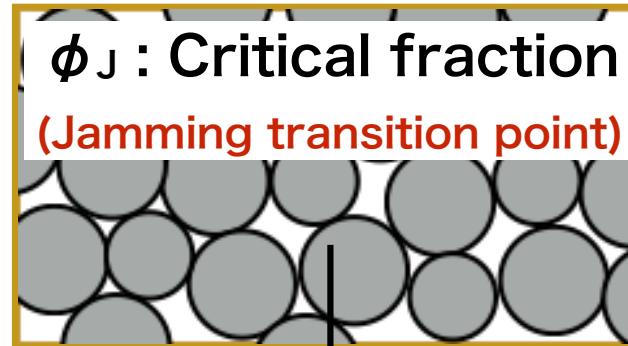
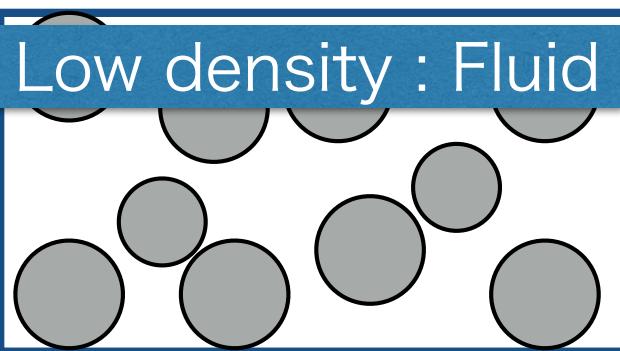


# Avalanche contribution to shear modulus of granular materials under oscillatory shear

**Effect of friction between particles**

Michio Otsuki (Shimane Univ.),  
Hisao Hayakawa (Kyoto Univ.)

# Jamming transition



$\phi$  : packing fraction

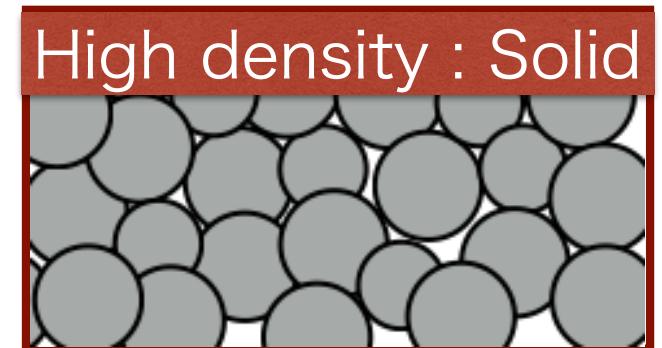
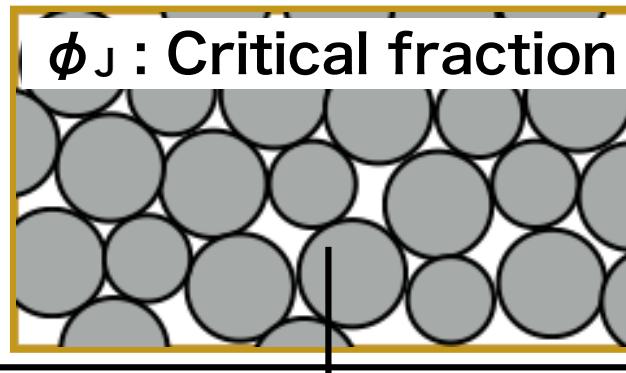
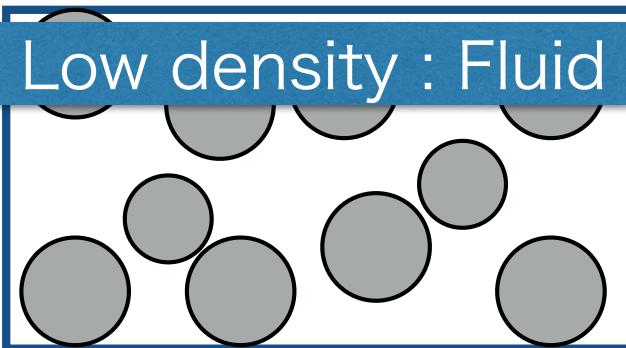
$\phi < \phi_J$  Granular materials  
flow like fluids.



$\phi > \phi_J$  Granular materials  
have rigidity like  
solids.

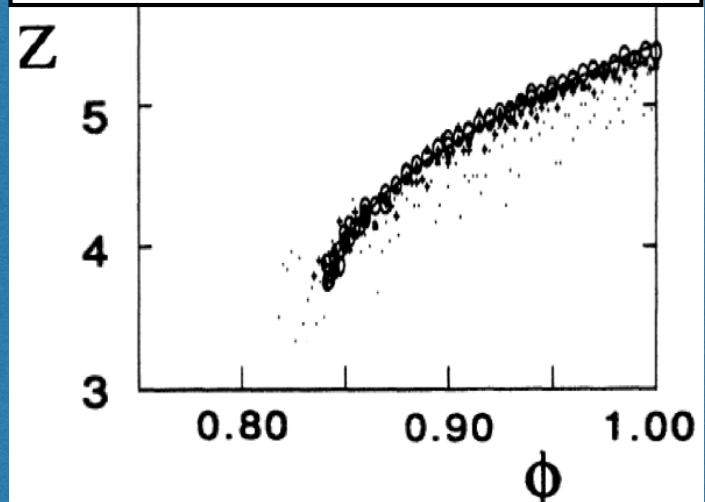


# Critical behaviors : frictionless grains

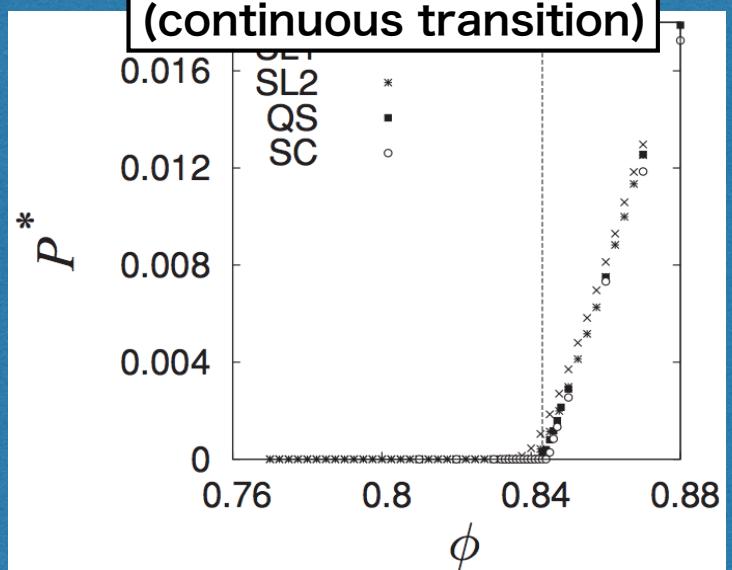


$\phi$  : packing fraction

Coordination number :  $Z$   
(discontinuous transition)



Pressure :  $P$   
(continuous transition)



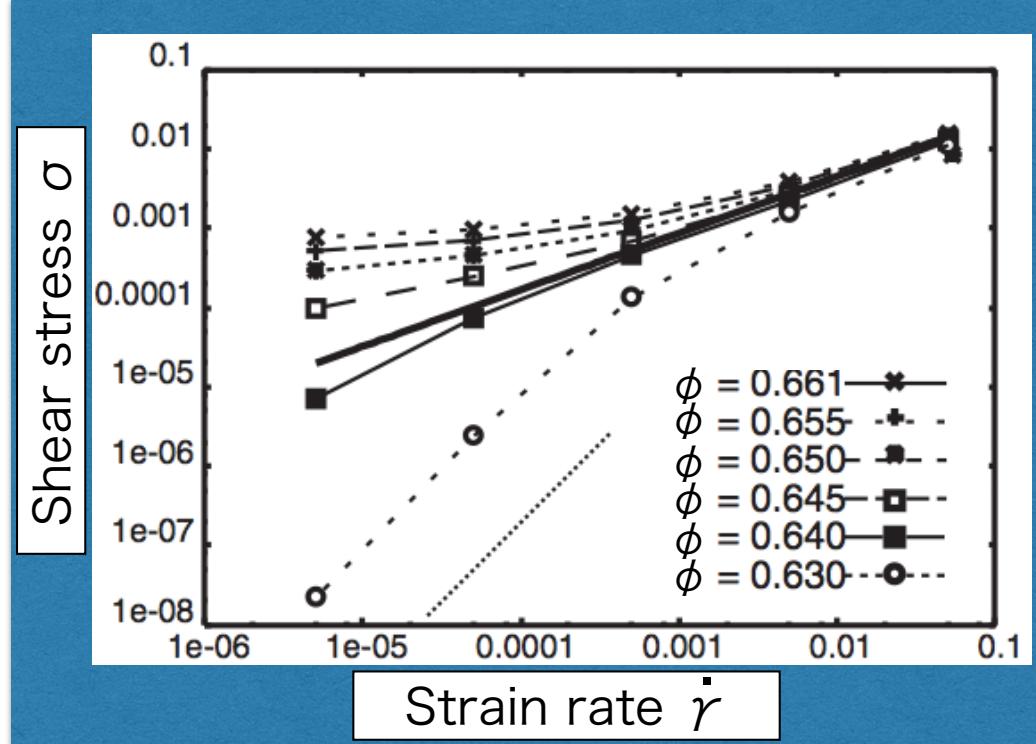
$$Z = \begin{cases} 0, & \phi < \phi_J \\ Z_c + a(\phi - \phi_J)^{1/2}, & \phi > \phi_J \end{cases}$$

$$P = \begin{cases} 0, & \phi < \phi_J \\ b(\phi - \phi_J), & \phi > \phi_J \end{cases}$$

# Jamming under steady shear

## Frictionless grains

Hatano, Otsuki, Sasa, JPSJ (2007)



$$\phi < \phi_J$$

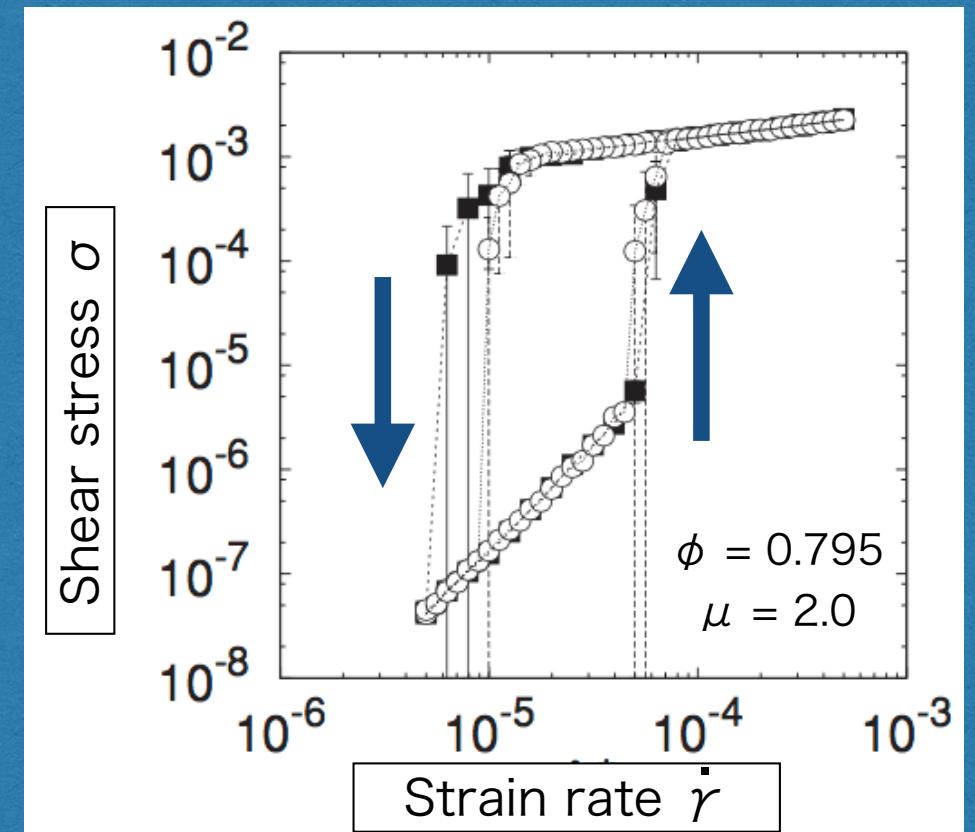
$$\sigma \propto \dot{\gamma}^2$$

$$\phi > \phi_J$$

$$\sigma = \text{const.}$$

## Frictional grains

Otsuki, Hayakawa, PRE (2011)

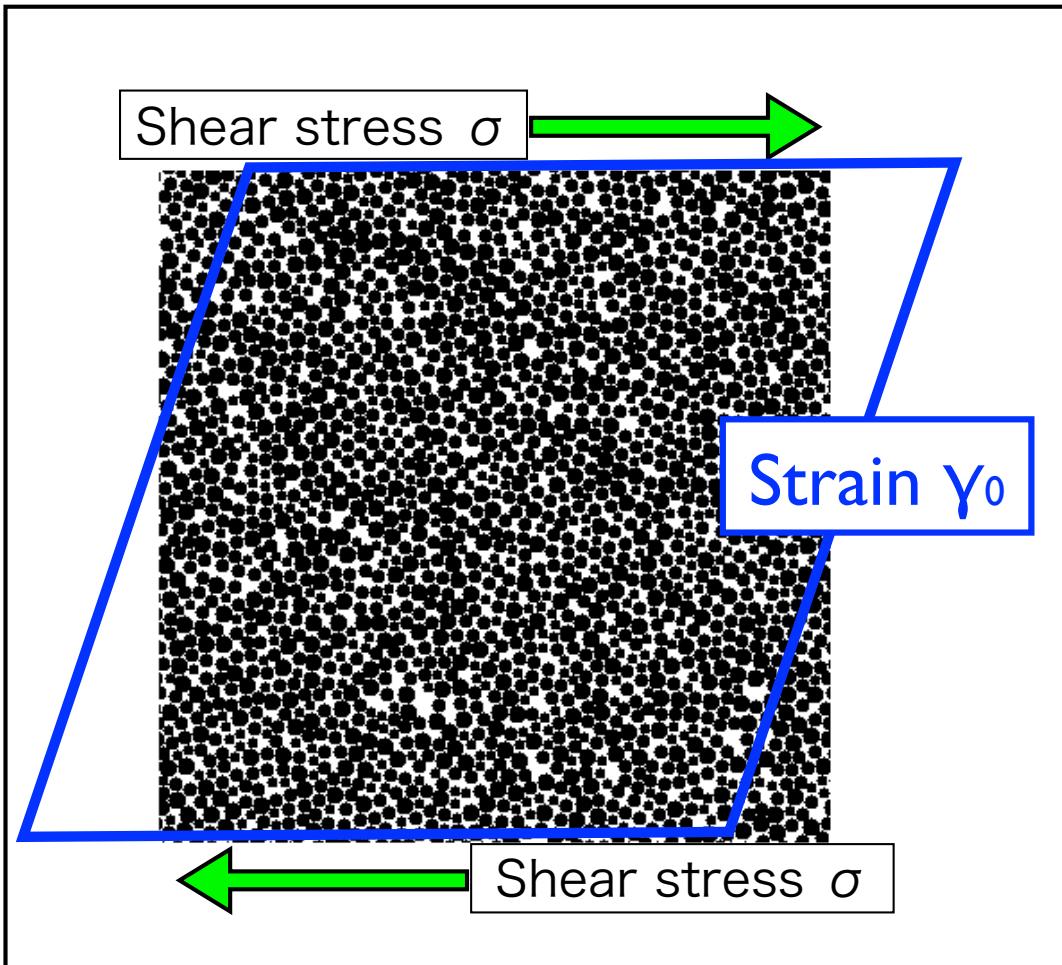
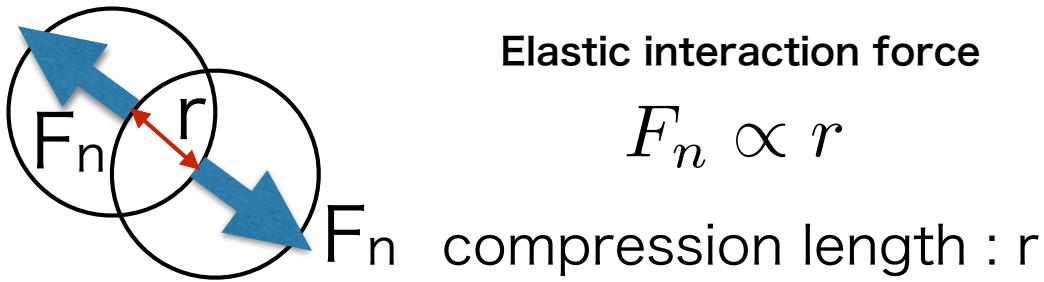


Strong shear thickening and hysteresis  
(They disappear in frictionless limit.)

Continuous transition

Discontinuous transition

# Shear modulus : frictionless grains

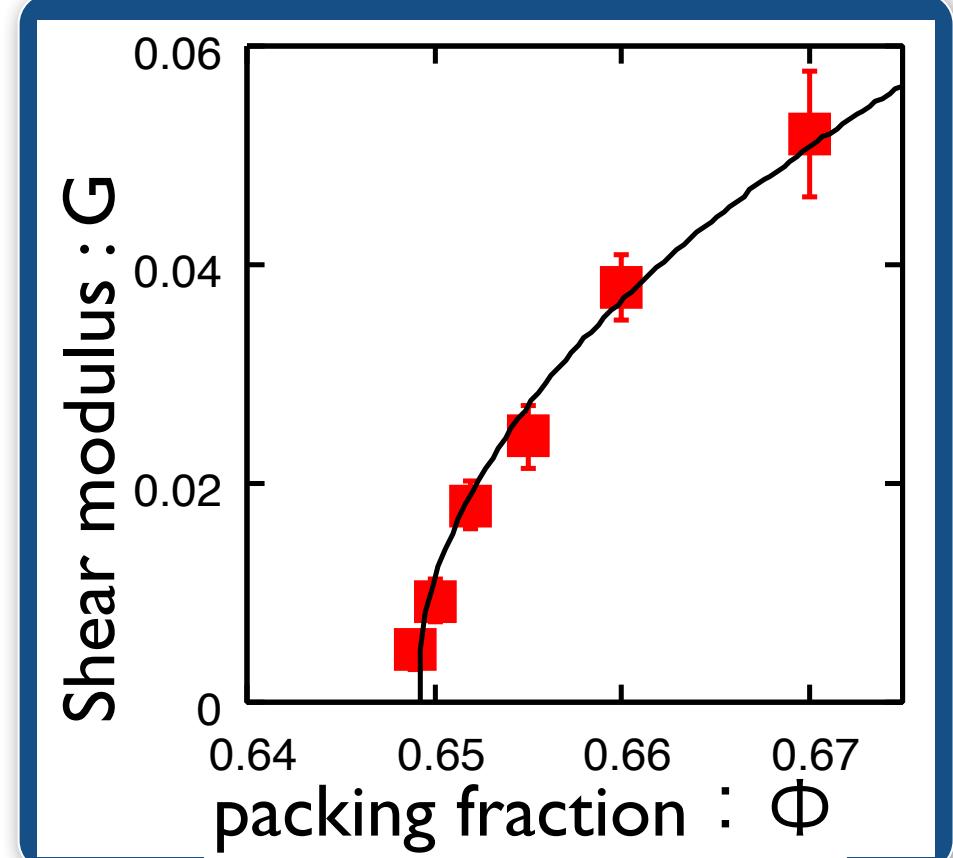


Shear modulus :  $G = \sigma / \gamma_0$

$G \propto (\phi - \phi_J)^{1/2}$

Linear response

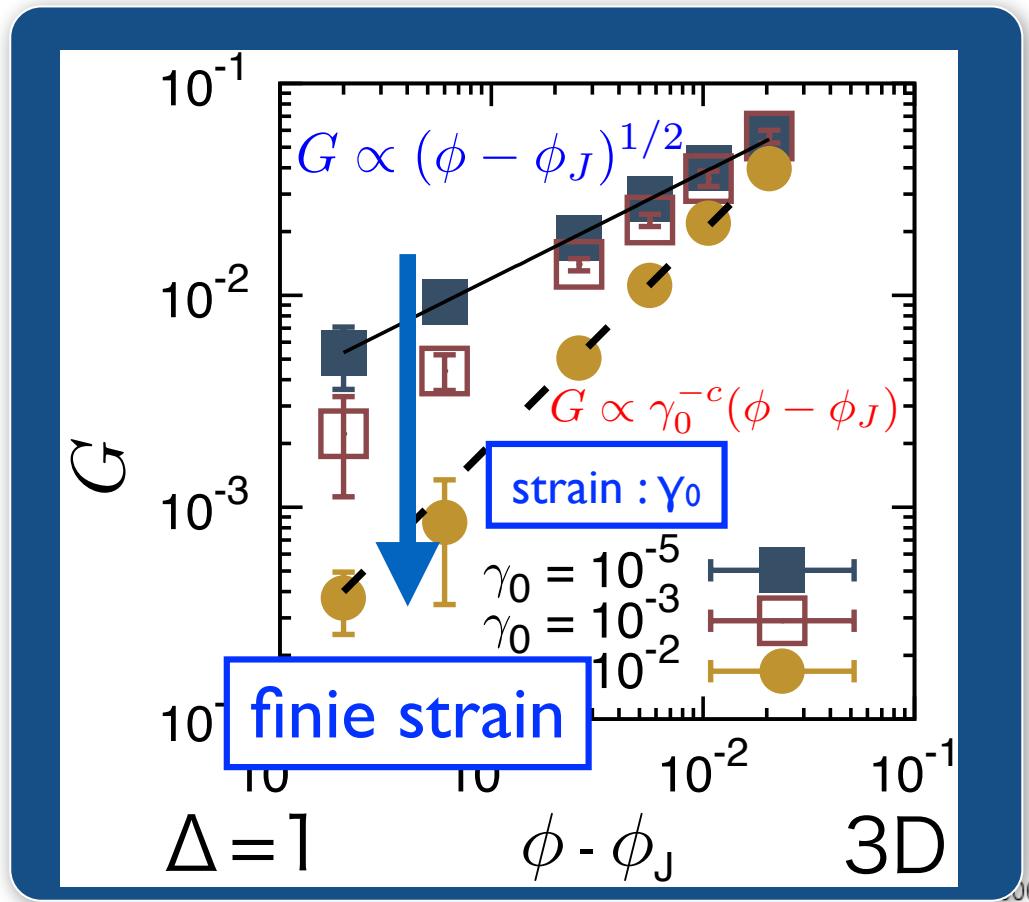
C. S. O'Hern et al., PRE 68, 011306 (2003)



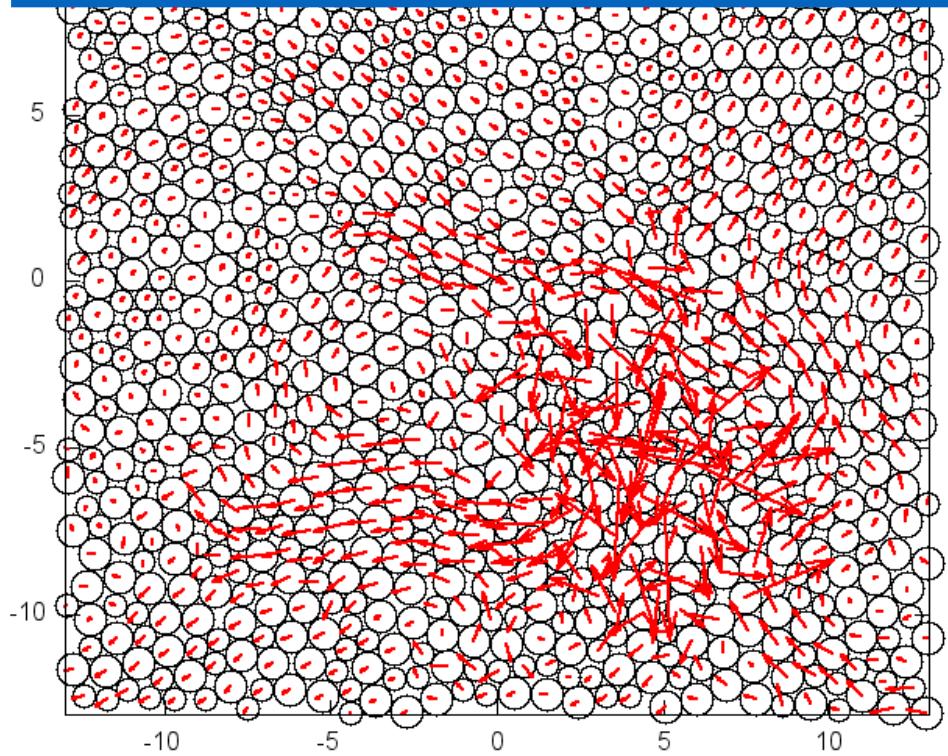
[extensively small strain  $\gamma_0 = 10^{-5}$ ]

# Shear modulus under finite strain

MO and H. Hayakawa, PRE (2014)



Displacement of grains under shear  
(slip avalanches)



Infinitesimal strain :  $G \propto (\phi - \phi_J)^{1/2}$

Linear response

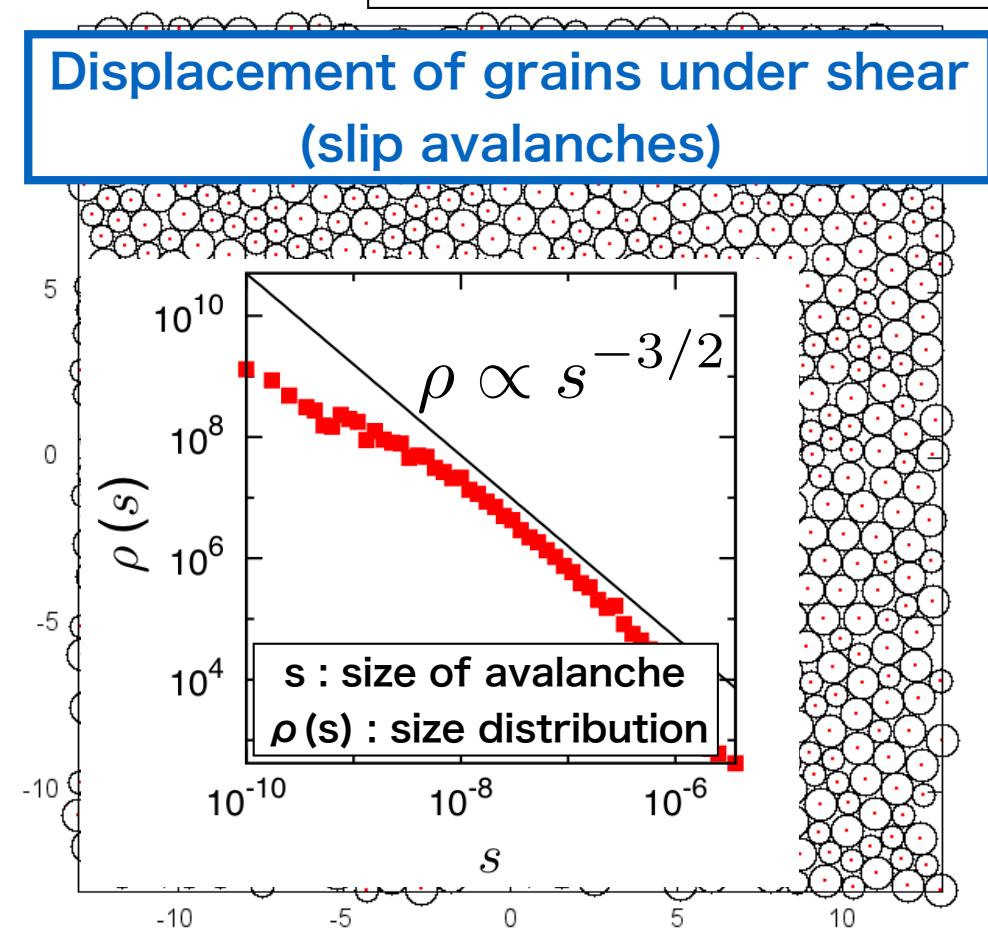
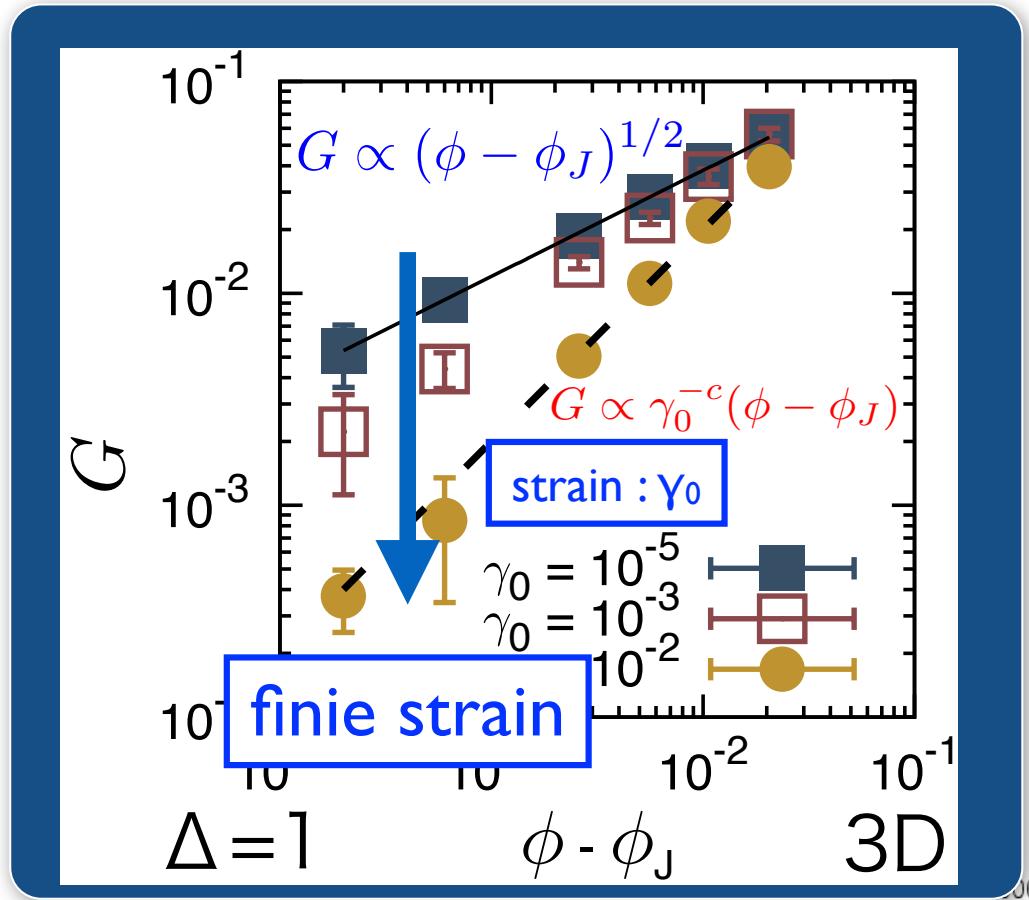
Finite strain :  $G \propto \gamma_0^{-c}(\phi - \phi_J)$

Nonlinear response

Origin of different scaling laws : slip avalanches

# Shear modulus under finite strain

MO and H. Hayakawa, PRE (2014)



Infinitesimal strain :  $G \propto (\phi - \phi_J)^{1/2}$

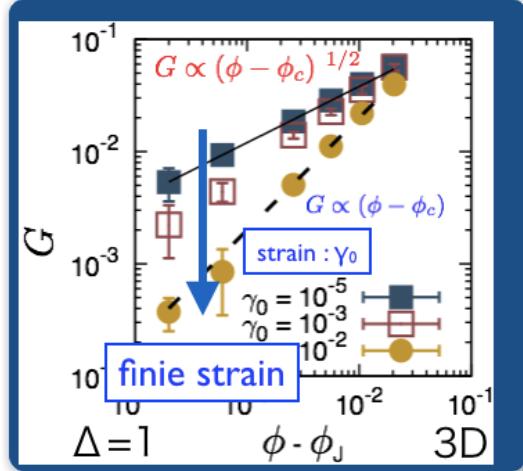
Linear response

Finite strain :  $G \propto \gamma_0^{-c}(\phi - \phi_J)$

Nonlinear response

Origin of different scaling laws : slip avalanches

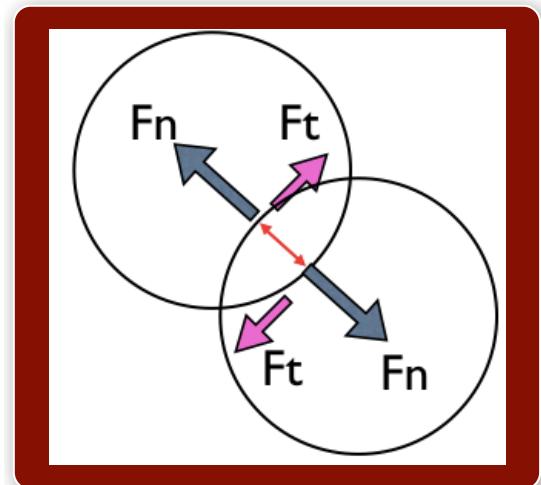
# Purpose : shear modulus of frictional grains



## Shear modulus of frictionless grains

Recent studies : including nonlinear elasticity

Otsuki and Hayakawa, PRE (2014)  
Coulais, Seguin, and Dauchot, PRL (2014),  
Goodrich, Liu, Sethna, arXiv : 1510.03469  
Nakayama, Yoshino, Zhamponi, arXiv:1512.06544  
Boschan, Vagberg, Somfai, Tighe, arXiv : 1601.00068



## Shear modulus of frictional grains

Linear elasticity :

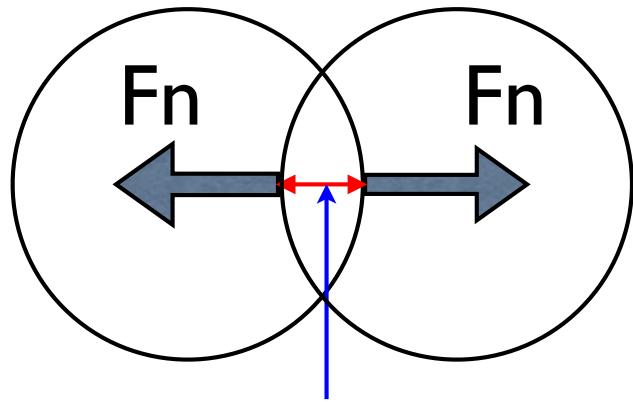
Somfai, van Hecke, Ellenbroek, Shundyak, van Saarloos, PRE (2007)  
Magnanimo, La Ragione, Jenkins, Wang, Makse, EPL 81, 34006 (2000)

There is no studies on non-linear elasticity.

We numerically study the effect of friction on G.

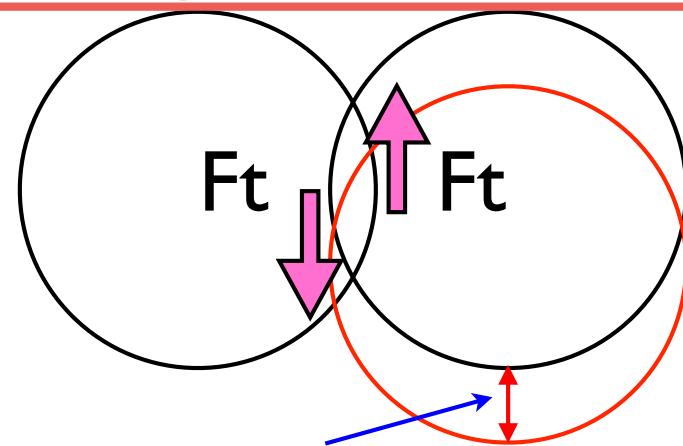
# 2D model of frictional grains

## Normal force



$r$  : compression length

## Tangential force



$\delta$  : tangential displacement

### Normal force

$$F_n = kr - \eta \dot{r}$$

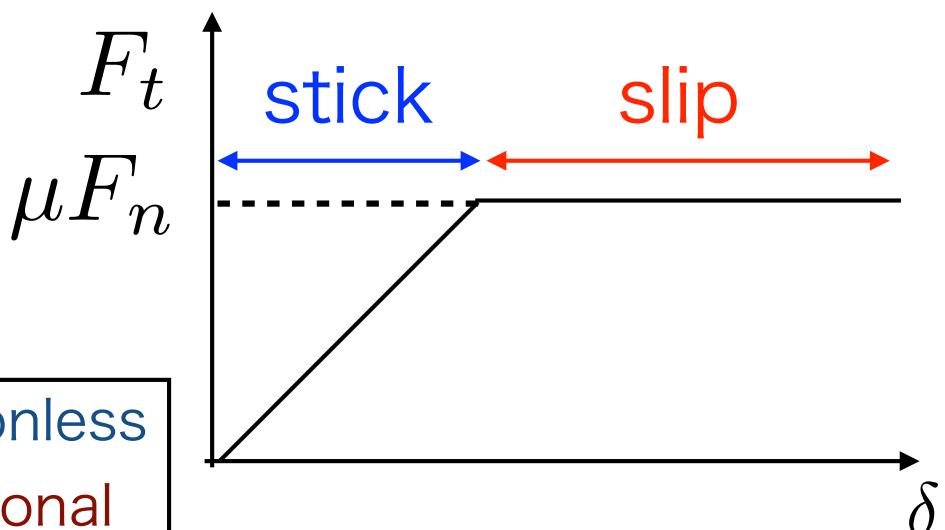
### Tangential force

$$|F_t| = k_t \delta, \quad F_t < \mu F_n$$

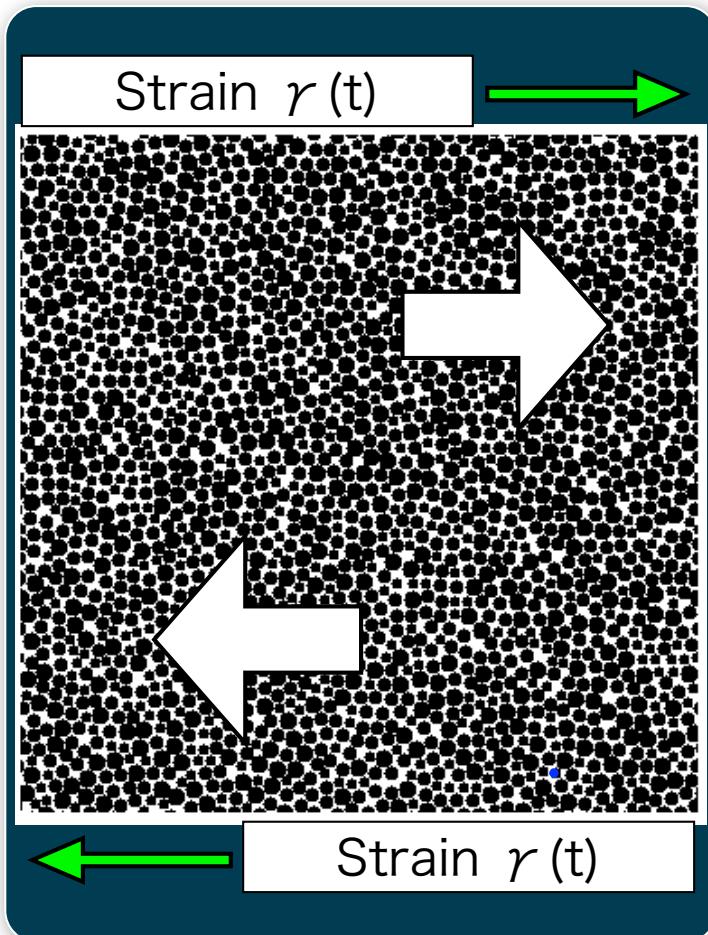
$$|F_t| = \mu F_n, \quad \text{otherwise}$$

$\mu$  : friction coefficient

$\mu = 0$  : frictionless  
 $\mu > 0$  : frictional



# Oscillatory shear



- Oscillatory shear strain :

$$\gamma(t) = \gamma_0(1 - \cos \omega t)$$

- Frequency :  $\omega$

Quasi-static limit :  $\omega \rightarrow 0$

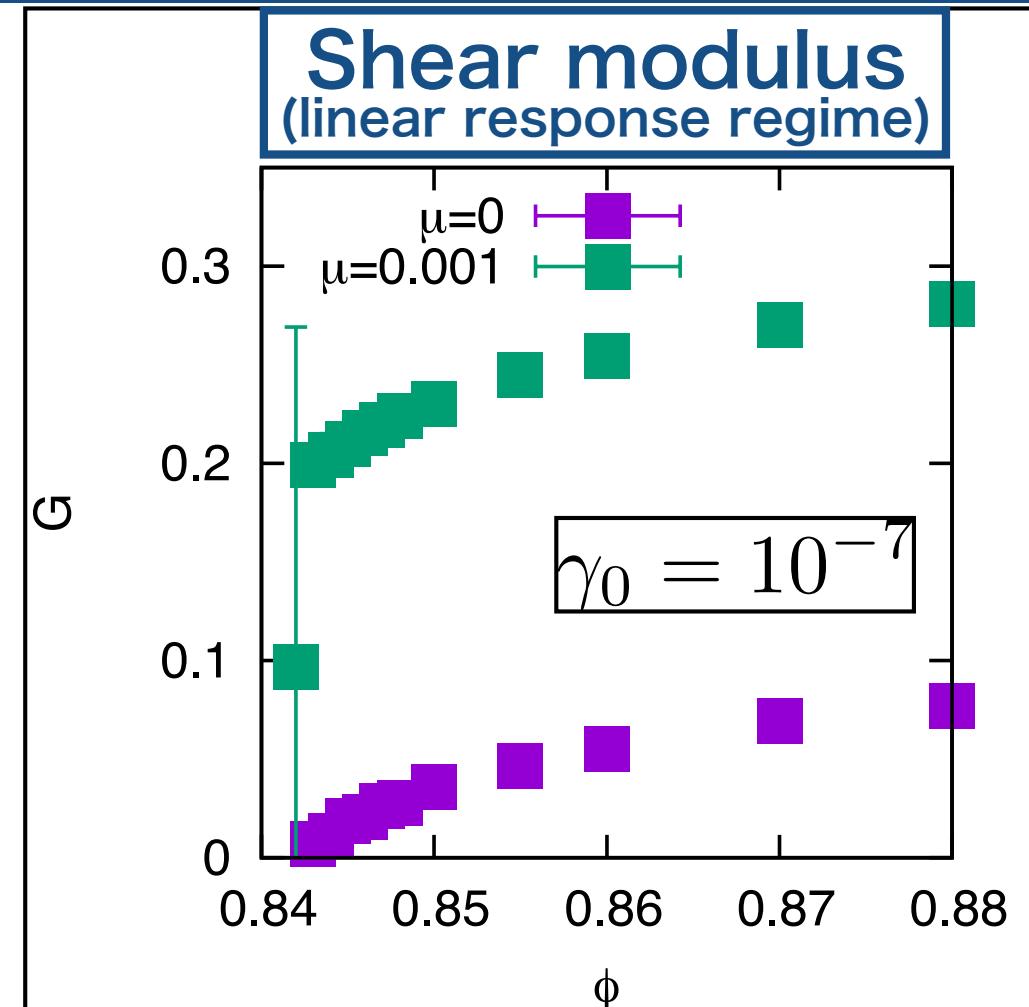
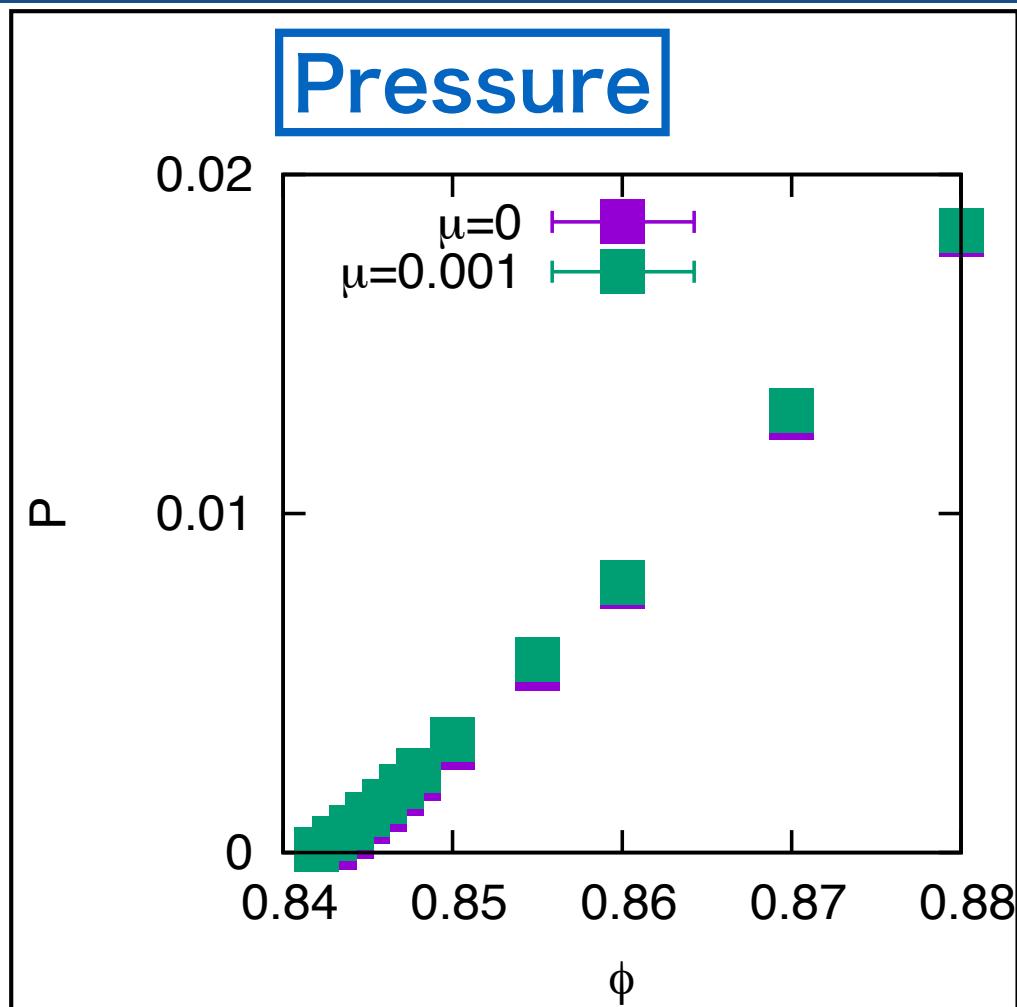
- **Strain amplitude :  $\gamma_0$**

- Shear stress :  $\sigma(t)$

## Shear modulus (storage modulus)

$$G(\gamma_0, \phi) = \frac{\omega}{\pi} \int_0^{2\pi/\omega} dt \frac{\sigma(t) \cos(\omega t)}{\gamma_0}$$

# Effect of friction

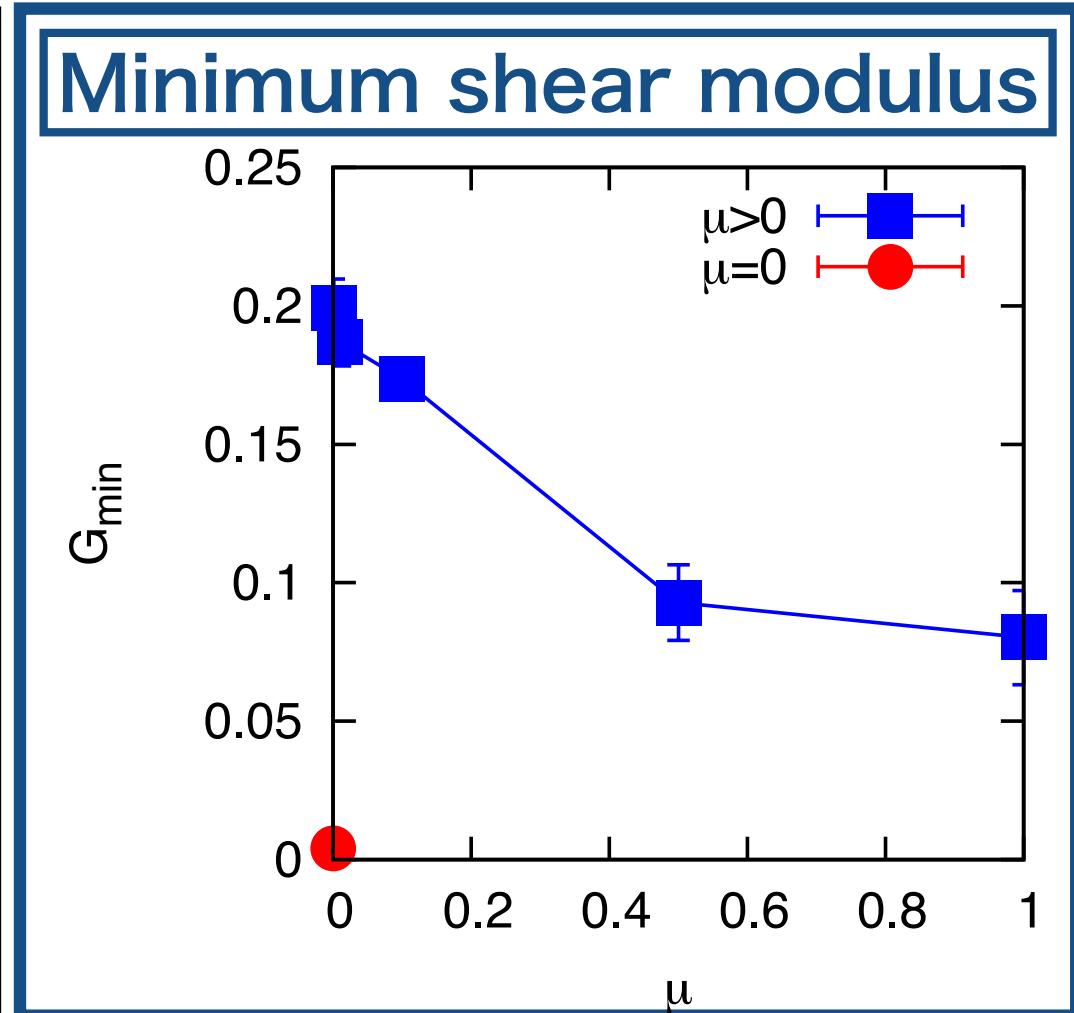
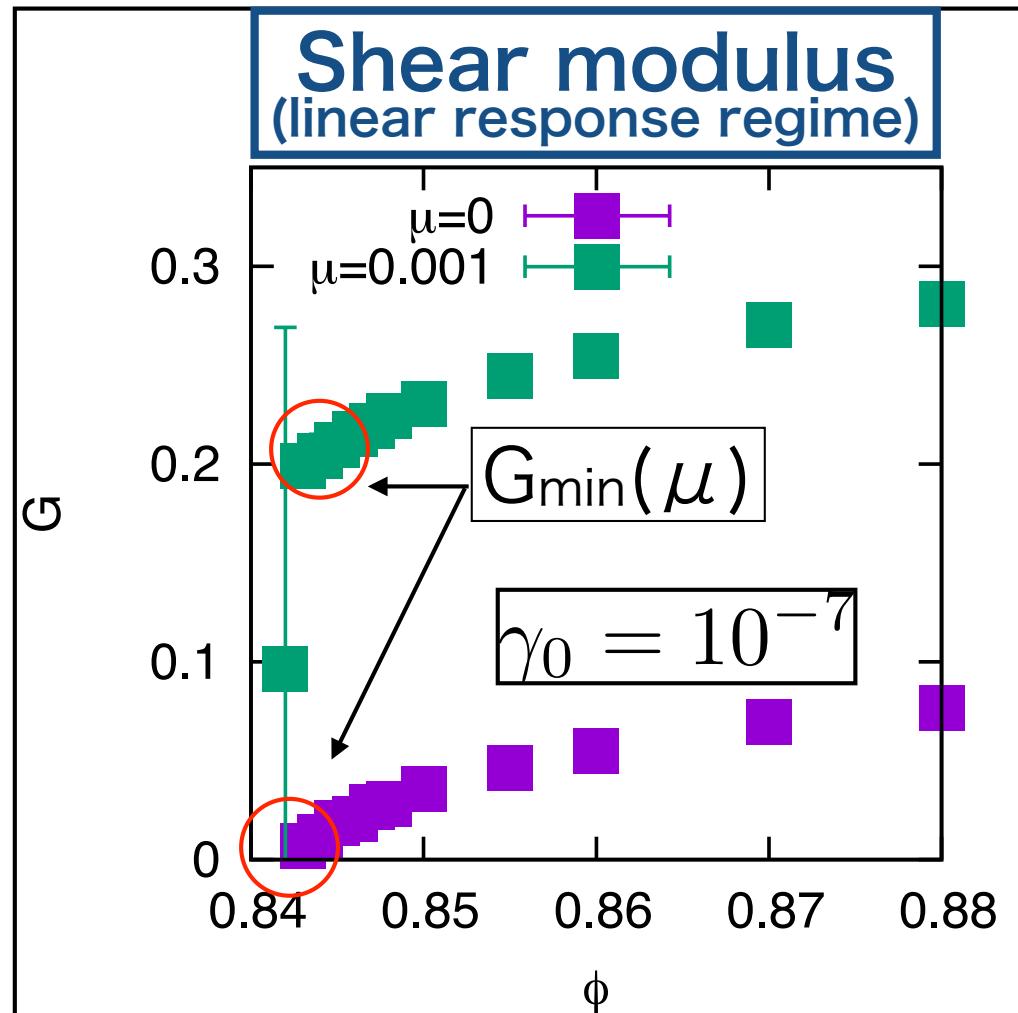


**Pressure :** Weak dependence on  $\mu$

c.f. Somfai et al., PRE (2007)

**Shear modulus :** Strong dependence on  $\mu$   
Discontinuous change at  $\phi_J$

# $\mu$ -dependence of minimum shear modulus



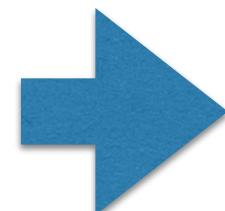
$\mu = 0$

$G_{\min} = 0$

$\mu > 0$

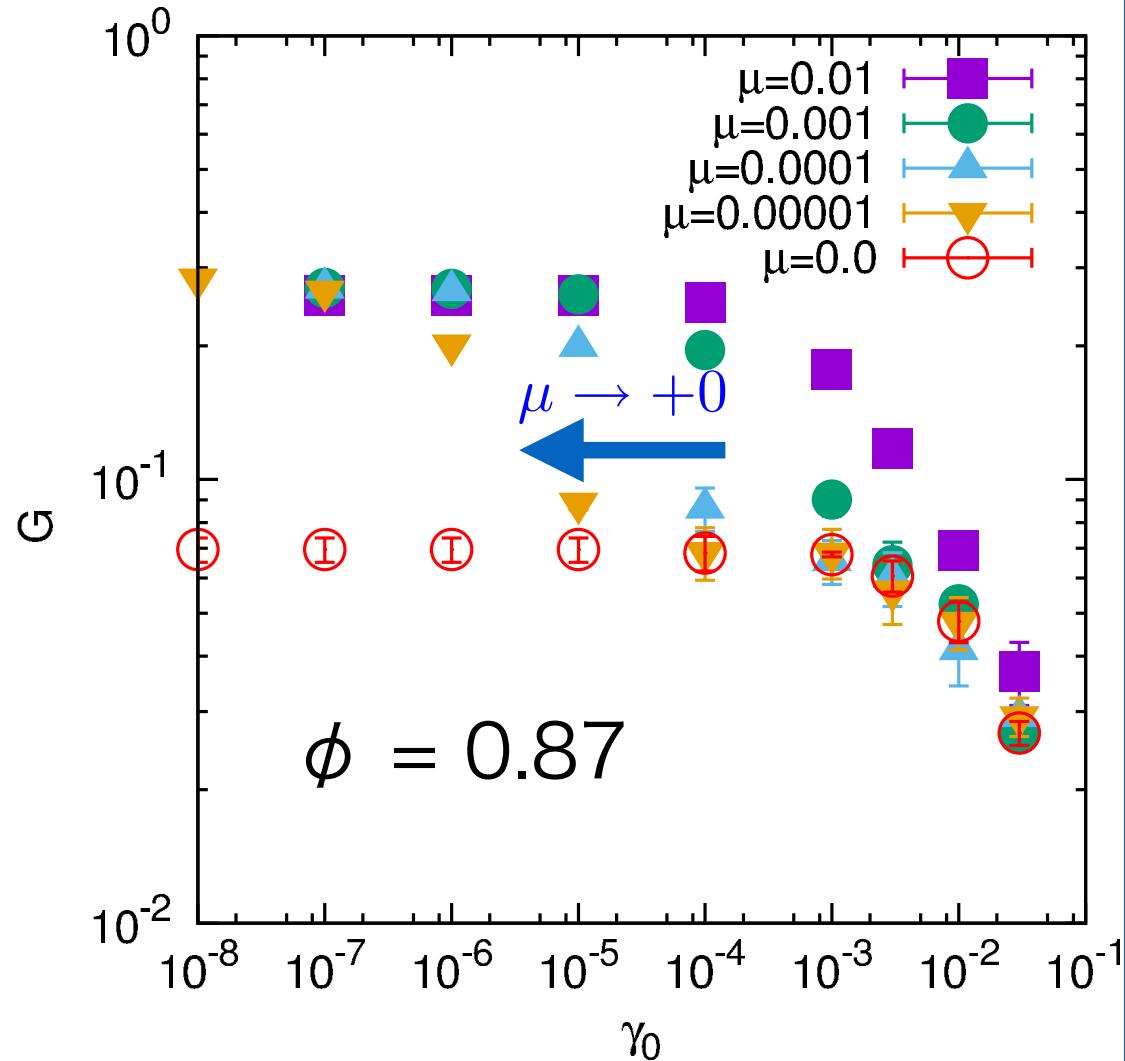
$G_{\min} > 0$  (even in the limit  $\mu \rightarrow +0$ )

G for  $\mu > 0$  is completely different from G for  $\mu = 0$  ?



Origin ?

# $\gamma_0$ -dependence of shear modulus



## $\gamma_0$ -dependence :

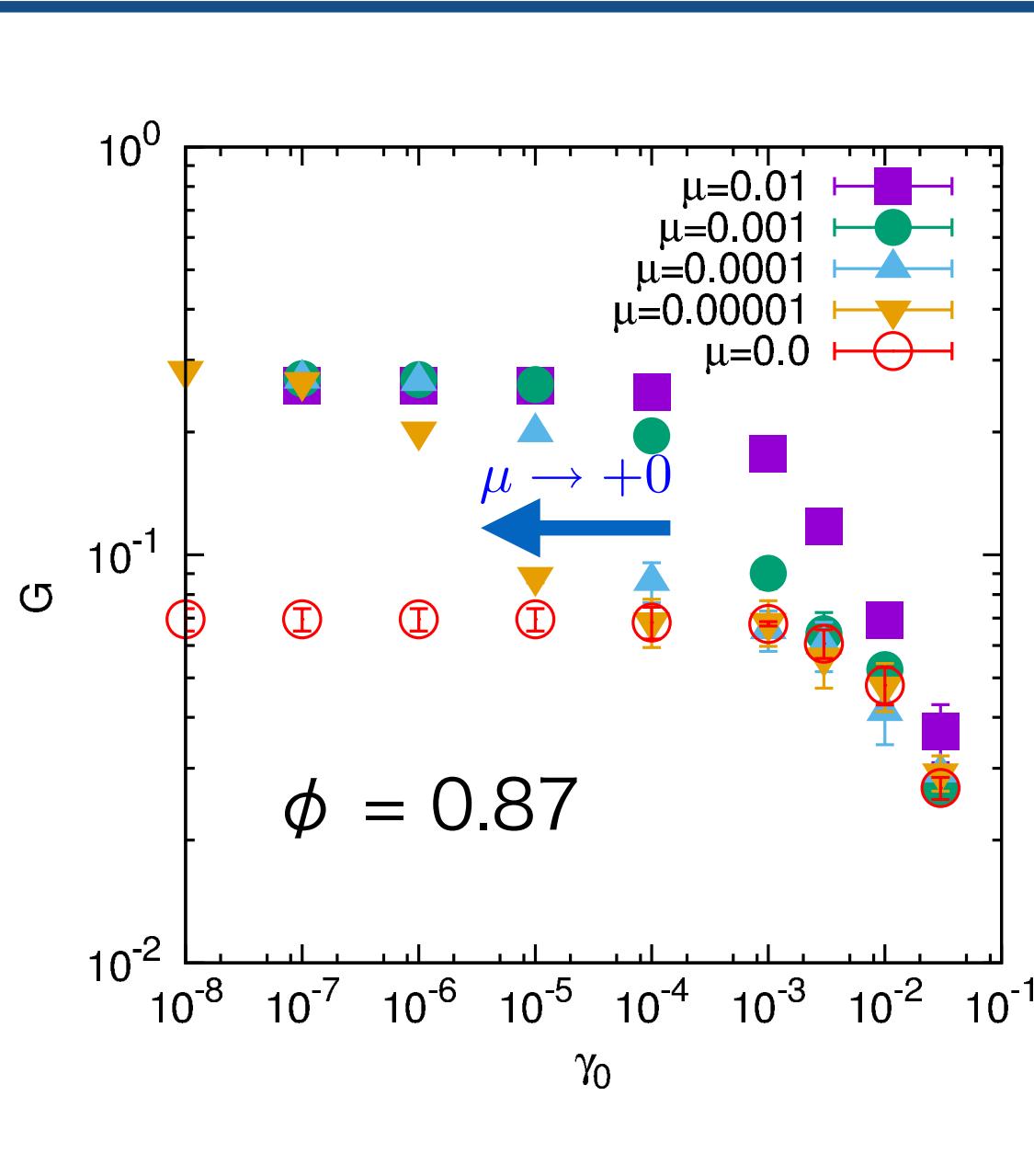
- For small  $\gamma_0$ ,  $G$  is constant.  
**(Linear response regime)**
- As  $\gamma_0$  increases,  $G$  decreases.  
**(Nonlinear response regime)**

## Linear response regime :

- $G$  for  $\mu > 0$  does not depend on  $\mu$ .
- $G$  for  $\mu > 0$  differs from  $G$  for  $\mu = 0$ .
- As  $\mu \rightarrow +0$ , the linear response regime shrinks.

$$\lim_{\mu \rightarrow +0} \lim_{\gamma_0 \rightarrow 0} G(\mu, \gamma_0, \phi) \neq G(\mu = 0, 0, \phi)$$

# $\gamma_0$ -dependence of shear modulus



## $\gamma_0$ -dependence :

- For small  $\gamma_0$ ,  $G$  is constant.  
**(Linear response regime)**
- As  $\gamma_0$  increases,  $G$  decreases.  
**(Nonlinear response regime)**

## Linear response regime :

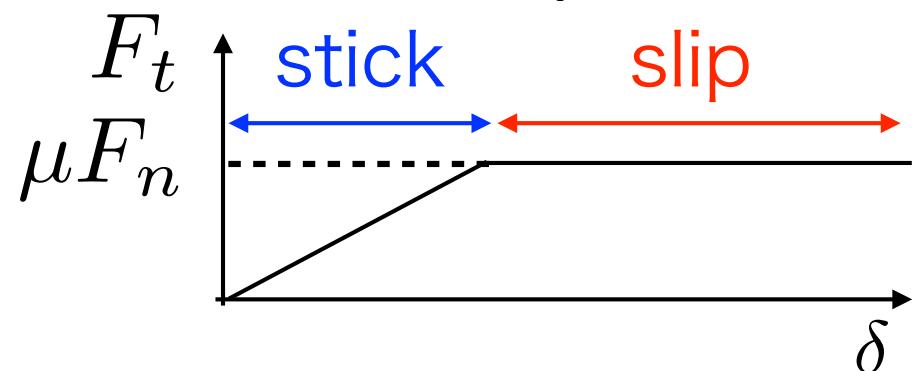
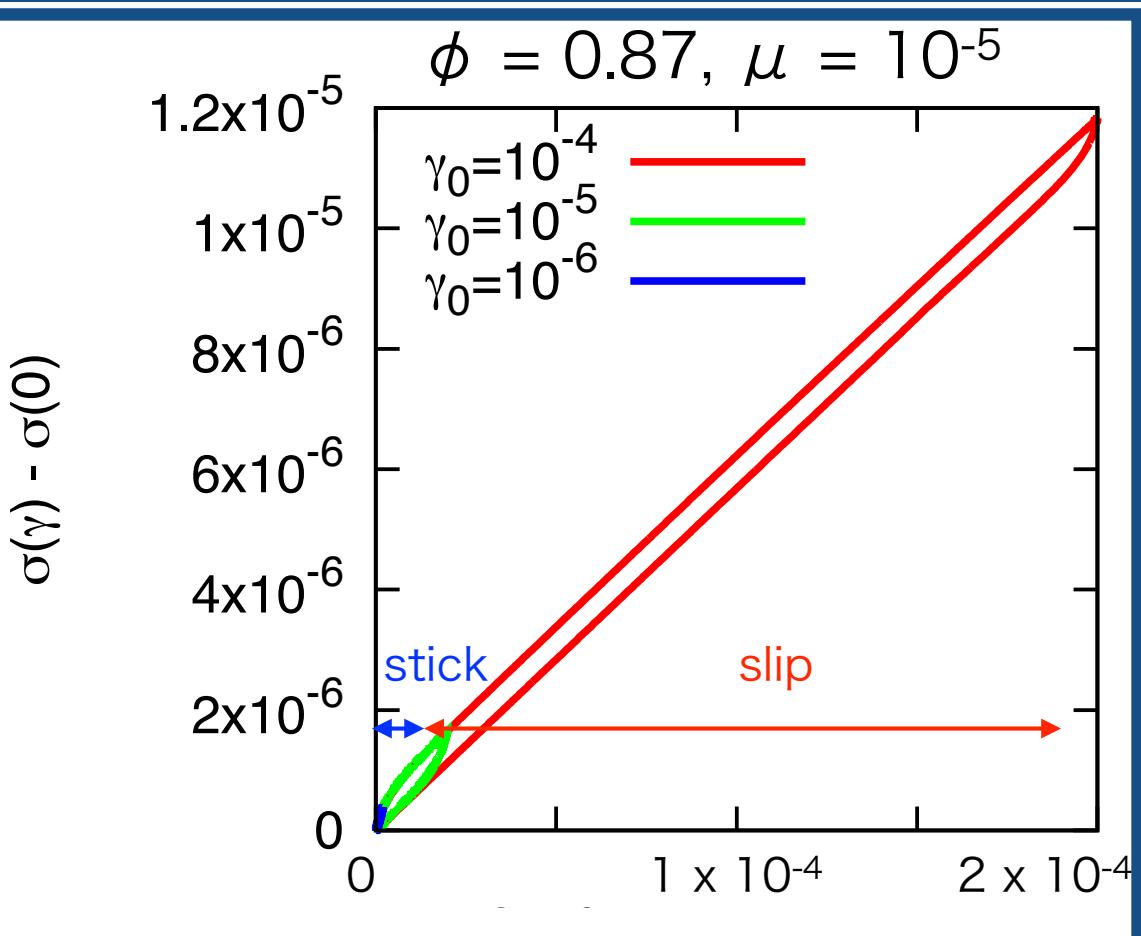
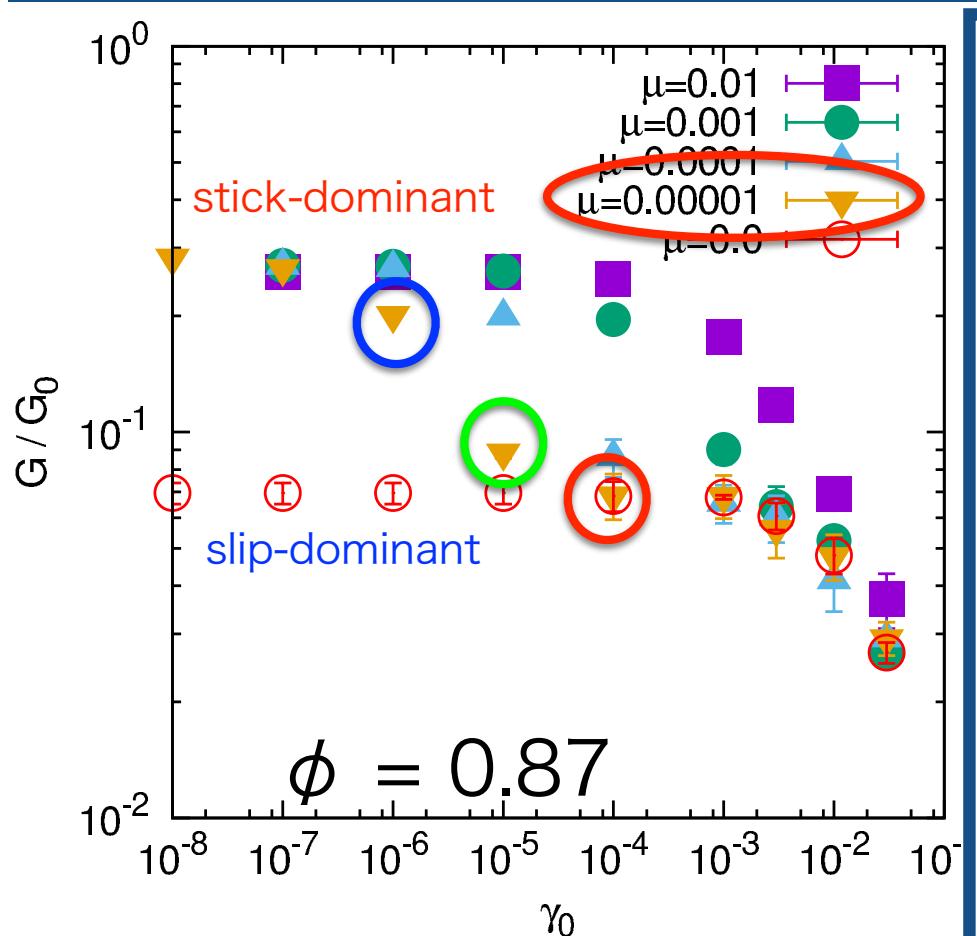
$$\lim_{\mu \rightarrow +0} \lim_{\gamma_0 \rightarrow 0} G(\mu, \gamma_0, \phi) \neq G(\mu = 0, 0, \phi)$$

## Nonlinear response regime :

- $G$  converges to  $G$  for  $\mu=0$  as  $\mu \rightarrow 0$ .
- $G$  has a plateau as  $\mu \rightarrow +0$ .

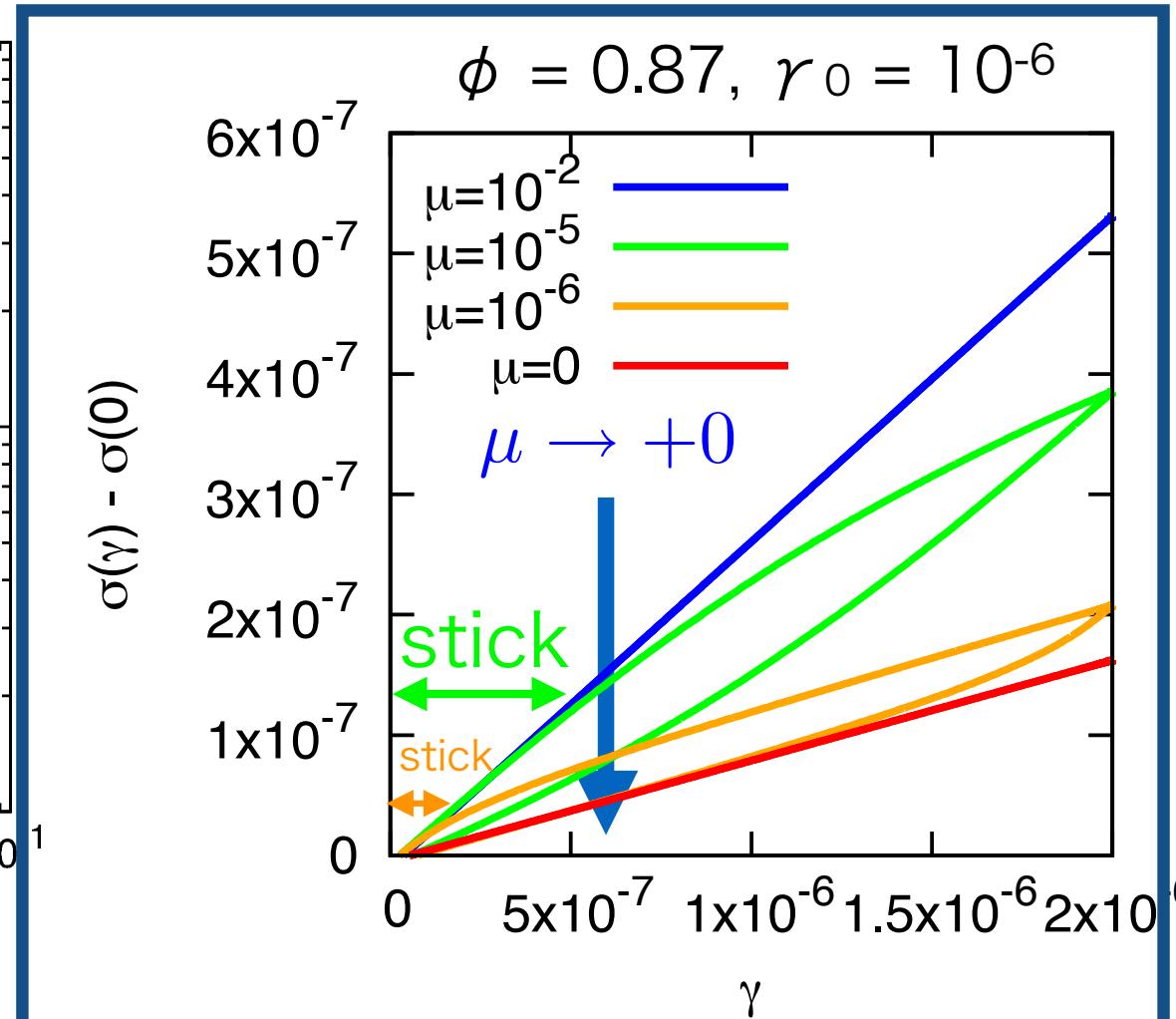
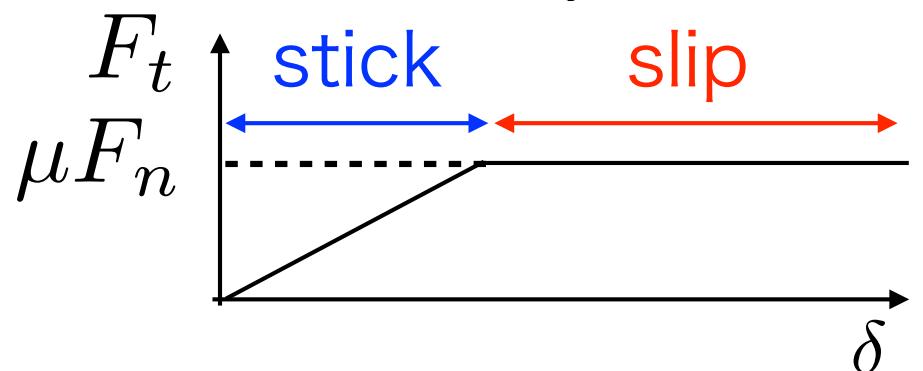
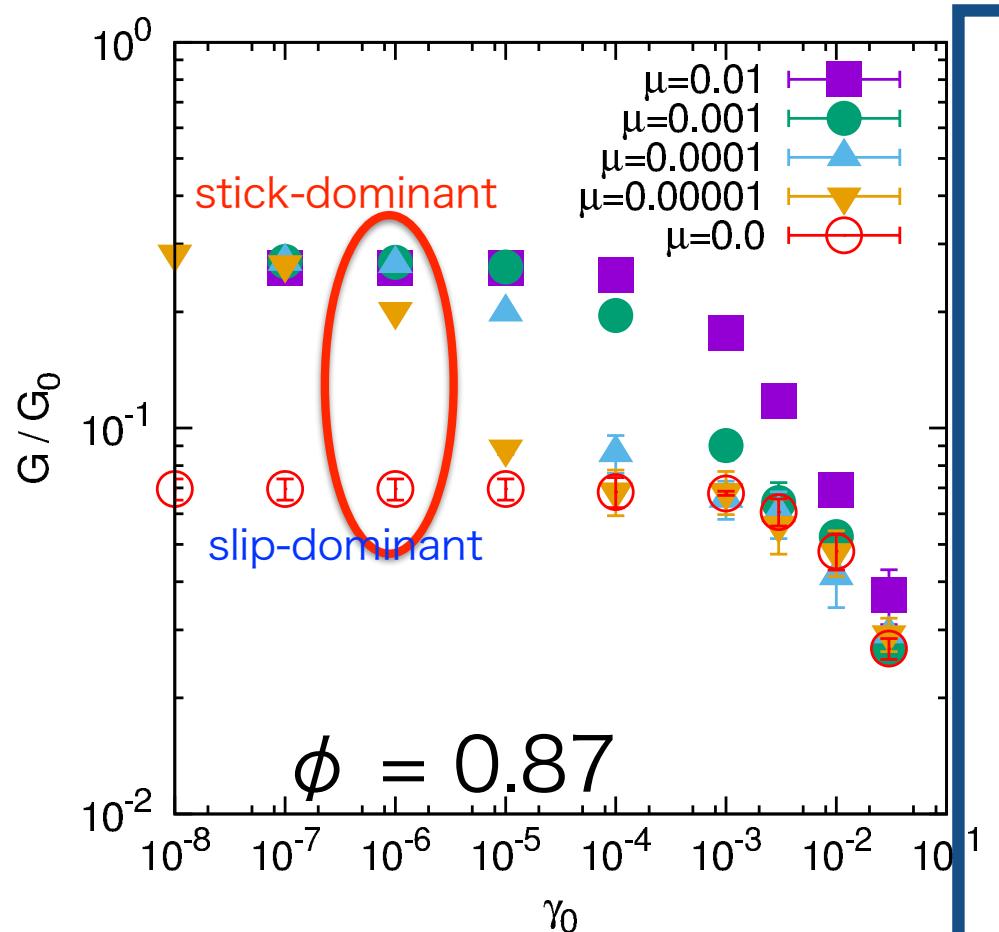
$$\lim_{\gamma_0 \rightarrow 0} \lim_{\mu \rightarrow +0} G(\mu, \gamma_0, \phi) = G(\mu = 0, 0, \phi)$$

# $\gamma_0$ -dependence of $\sigma$ - $\gamma$ relation



$\gamma_0 = 10^{-6}$  : Linear (stick region)  
 $\gamma_0 = 10^{-5}$  : Nonlinear with loop (slip region)  
 $\gamma_0 = 10^{-4}$  : Nonlinear, another linear region  
 (origin of a plateau region in  $G$  -  $\gamma_0$  plot)

# $\mu$ -dependence of $\sigma$ - $\gamma$ relation

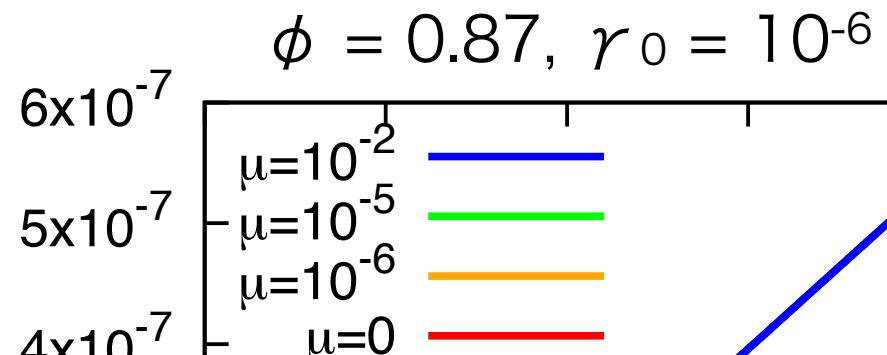
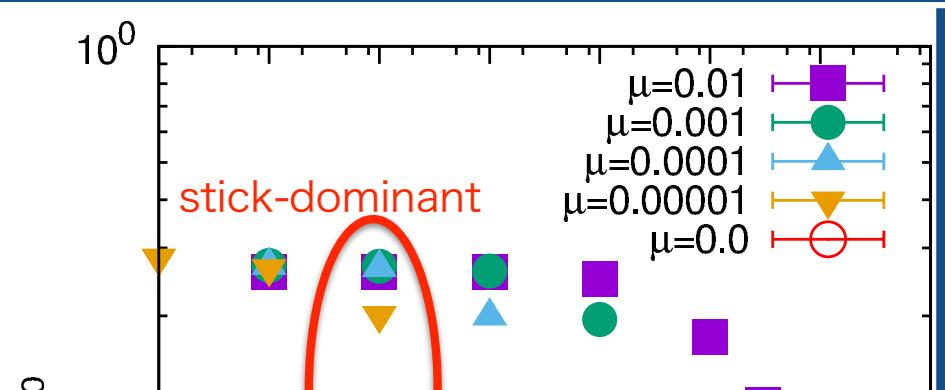


$\mu = 10^{-2}$  : Linear response

$\mu = 10^{-5}$  : Nonlinear response with loop

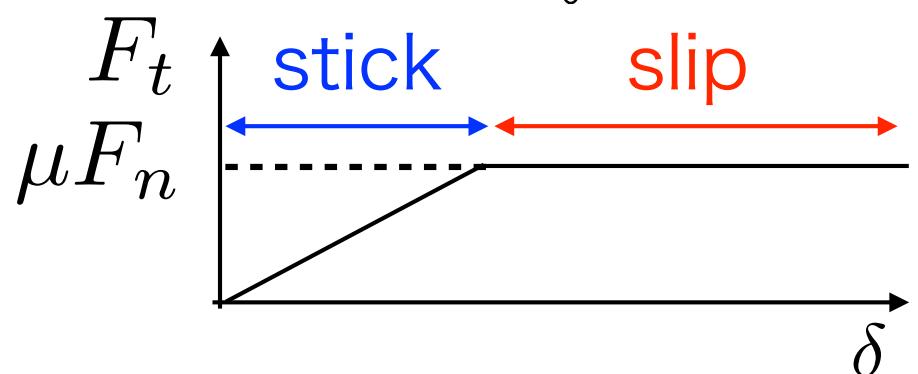
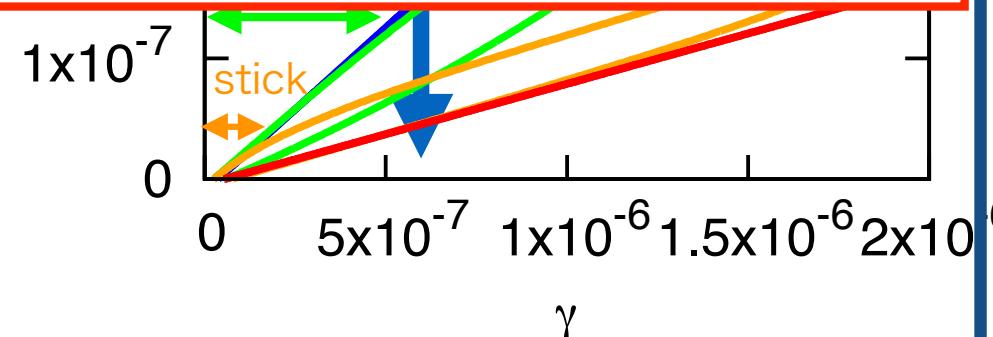
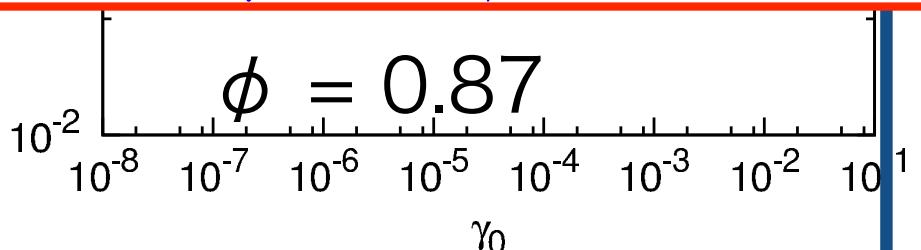
$\mu = 10^{-6}$  : Shrink of linear response regime, another linear region ([the angle] = [G for  $\mu=0$ ])

# $\mu$ -dependence of $\sigma$ - $\gamma$ relation



## Origin of singular behavior

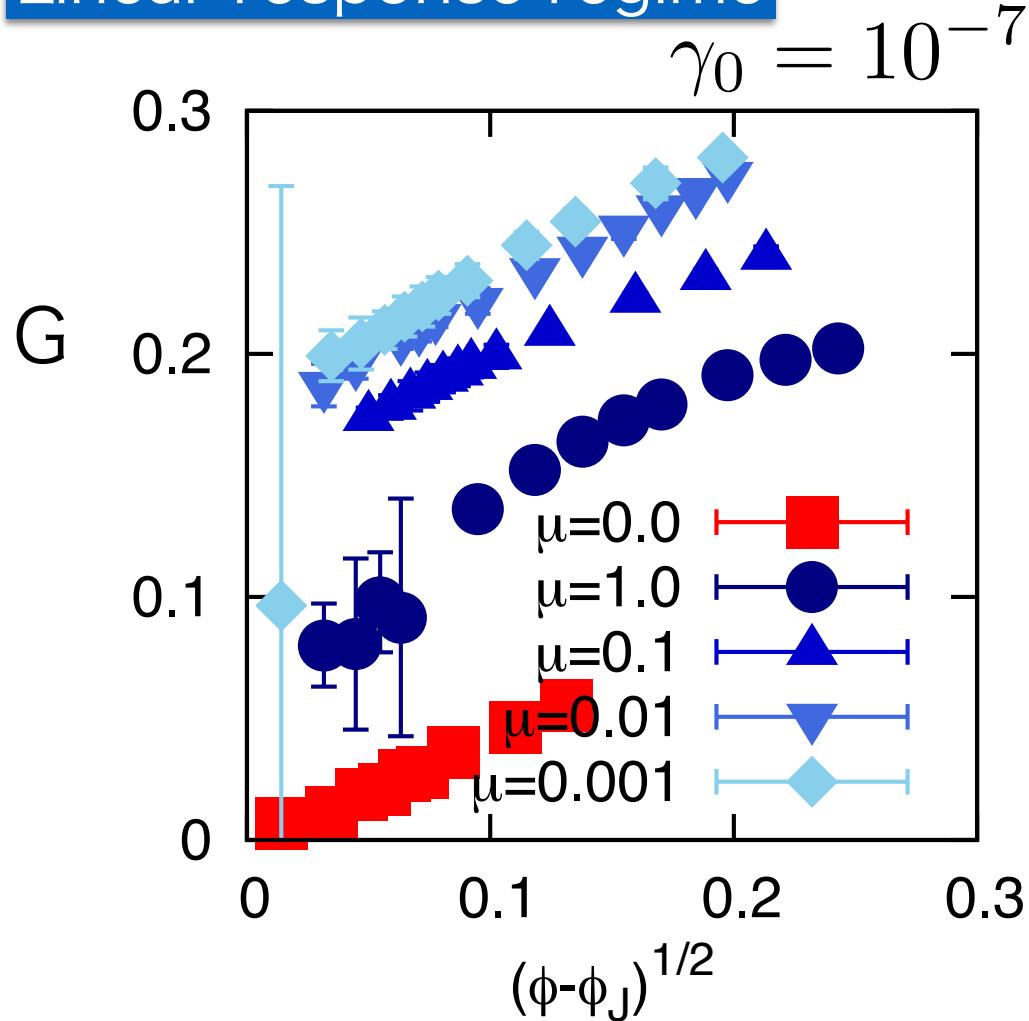
$$\lim_{\mu \rightarrow +0} \lim_{\gamma_0 \rightarrow 0} G(\mu, \gamma_0, \phi) \neq G(\mu = 0, 0, \phi)$$



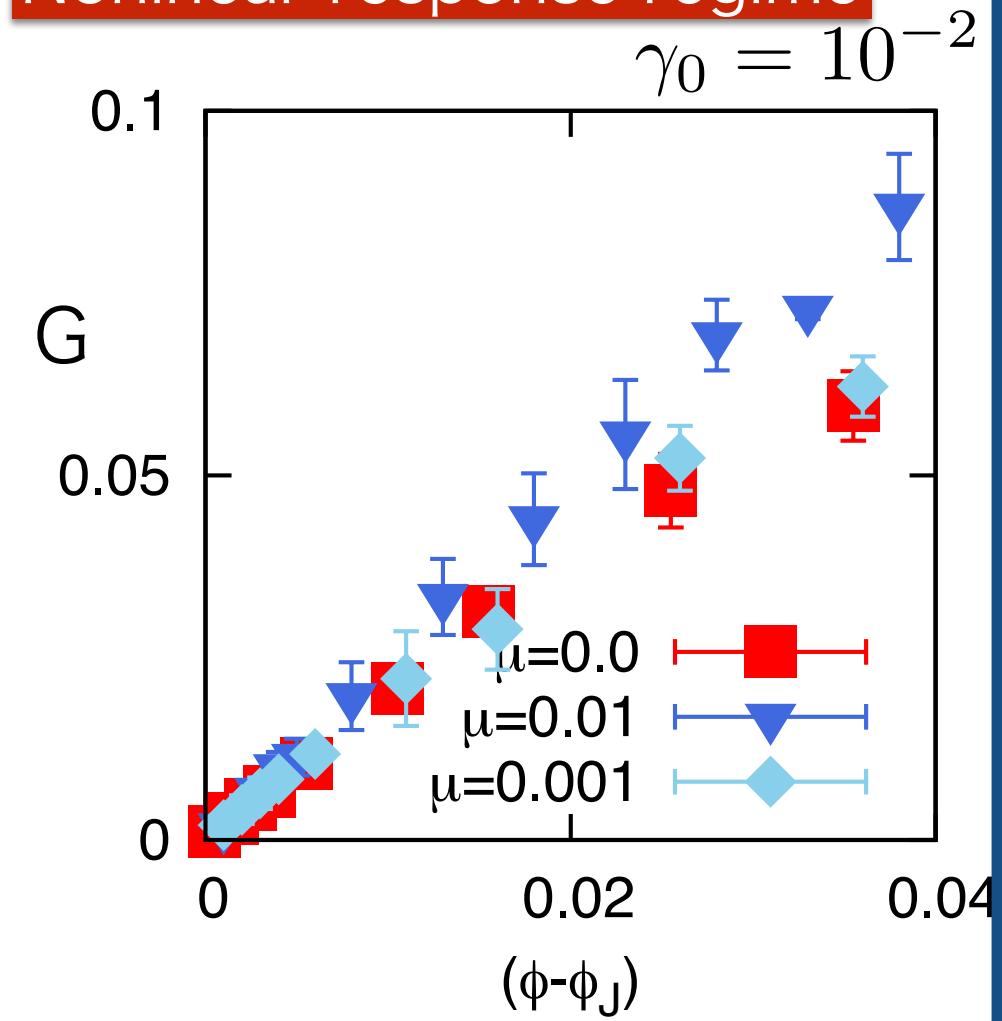
$\mu = 10^{-2}$  : Linear response  
 $\mu = 10^{-2}$  : Nonlinear response with loop  
 $\mu = 10^{-2}$  : Shrink of linear response regime,  
 another linear region ([the angle] = G for  $\mu=0$ )

# Scaling of G

Linear response regime



Nonlinear response regime



$$G = G_0(\mu) + A(\phi - \phi_J)^{1/2}$$

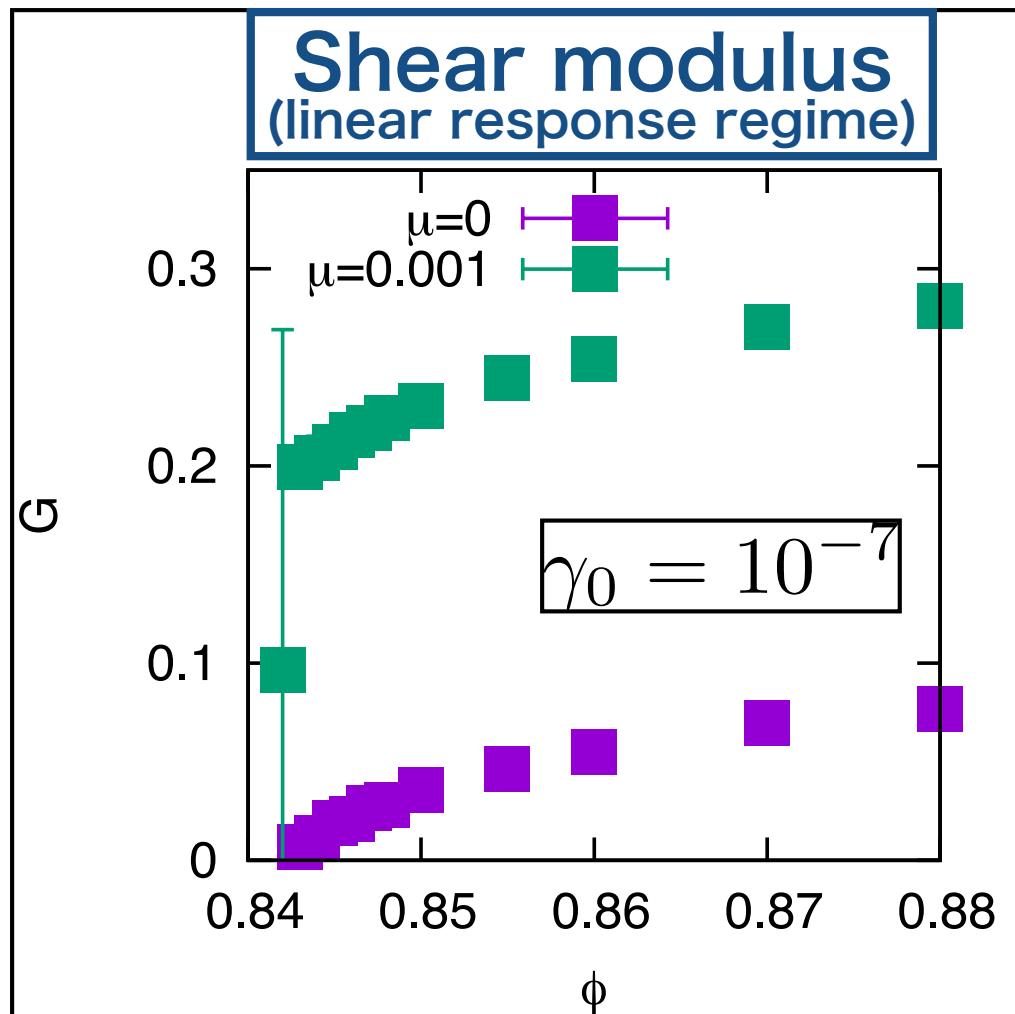
Somfai, et al., PRE (2007)

$$G \propto \gamma_0^{-c} (\phi - \phi_J)$$

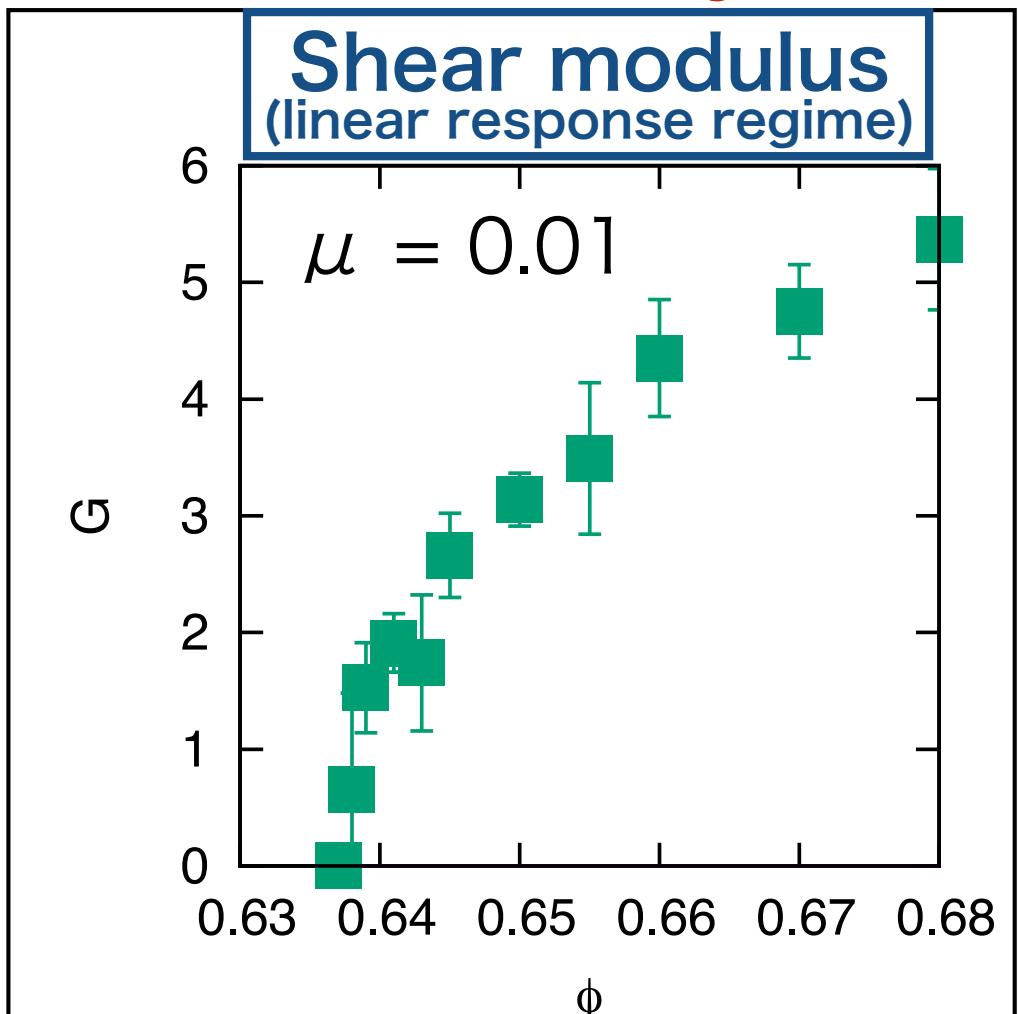
consistent with the result of  $\mu=0$ .

# Discussion : Dependence on dimension

2D



3D **(Preliminary result)**



Continuous transition in 3D ?

# Summary

- Purpose : shear modulus of frictional granular materials.
- Infinitesimal friction changes the shear modulus in linear response regime.
- The shear modulus for  $\mu \rightarrow +0$  in nonlinear response regime is consistent with that of frictionless grains.

