



Interpreting the IceCube events by decaying dark matter hints and constraints

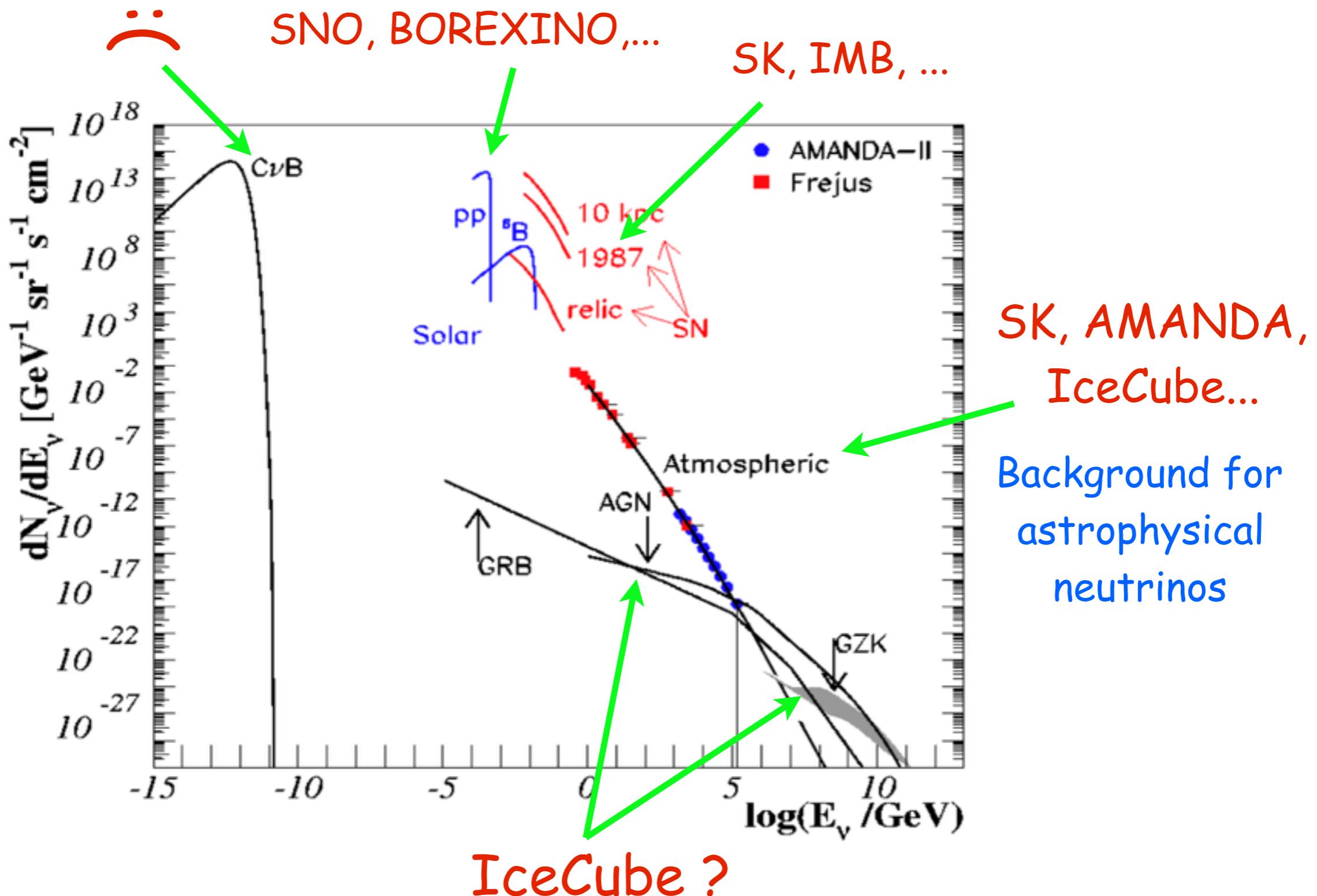
Arman Esmaili

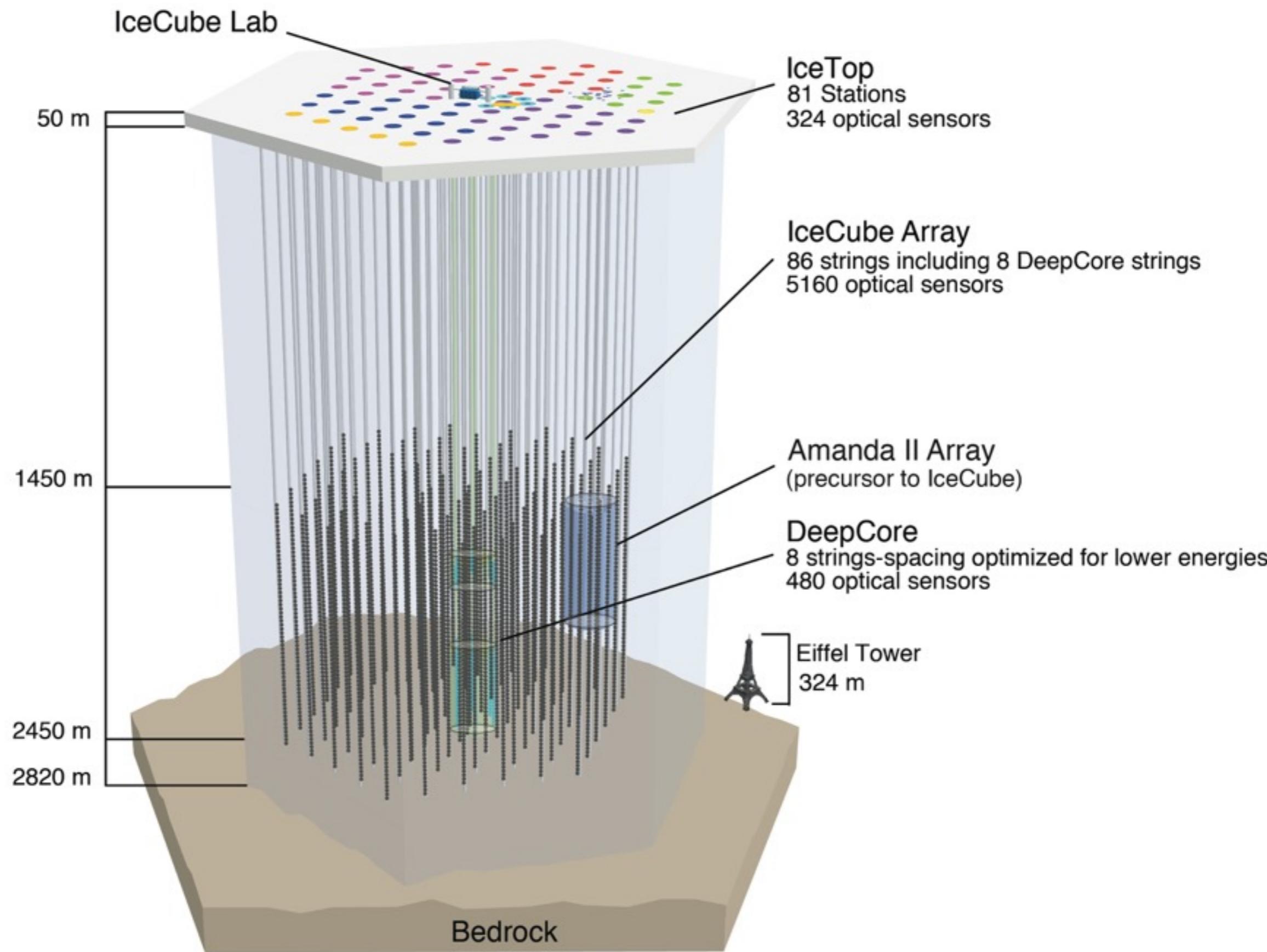
Laboratori Nazionali del Gran Sasso (LNGS)
Theory Group

18/Dec/2015

DSU @ YITP

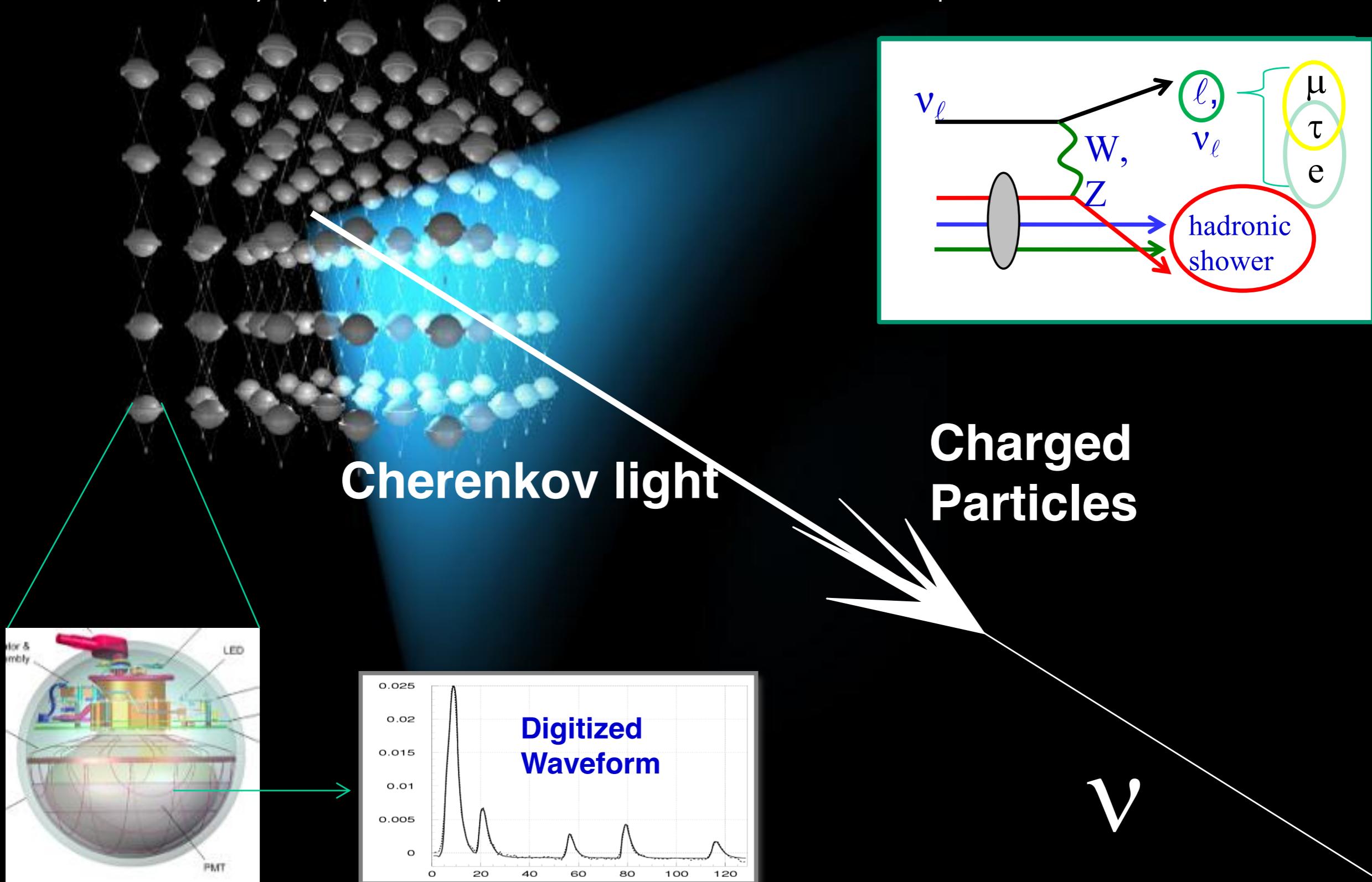
Neutrino Sky





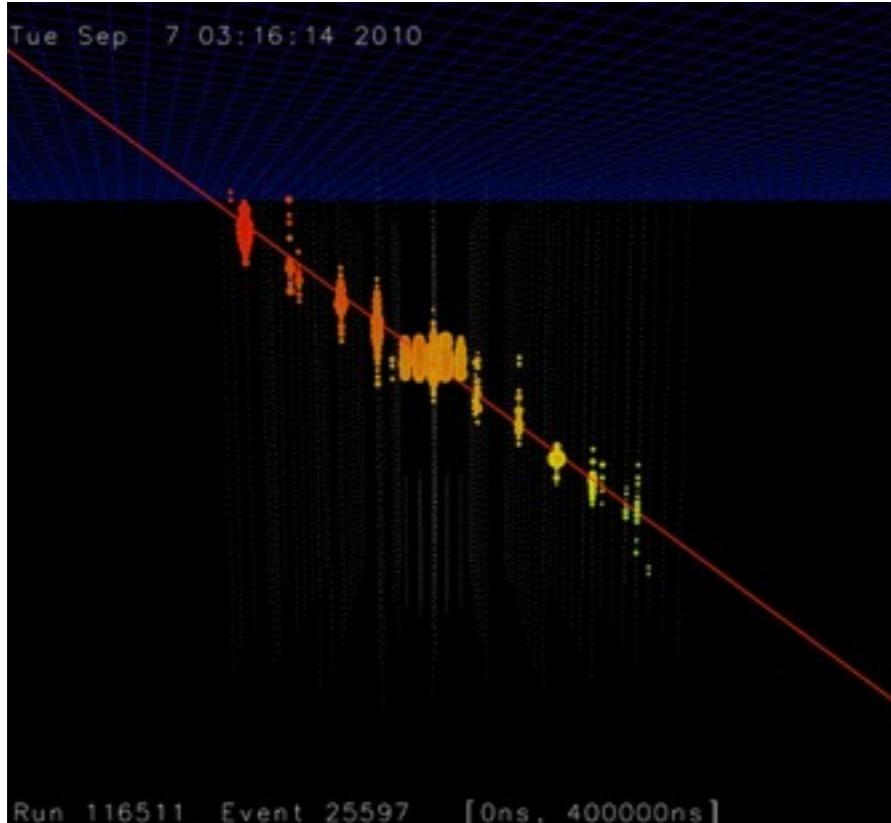
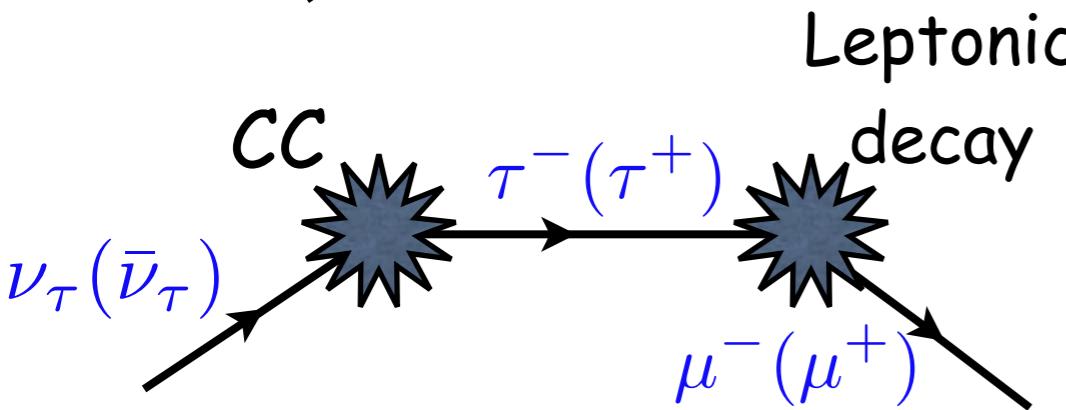
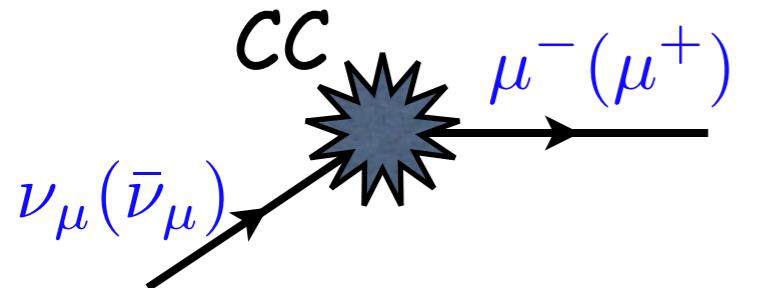
Detection Principle

An array of photomultiplier tubes + Dark and transparent material

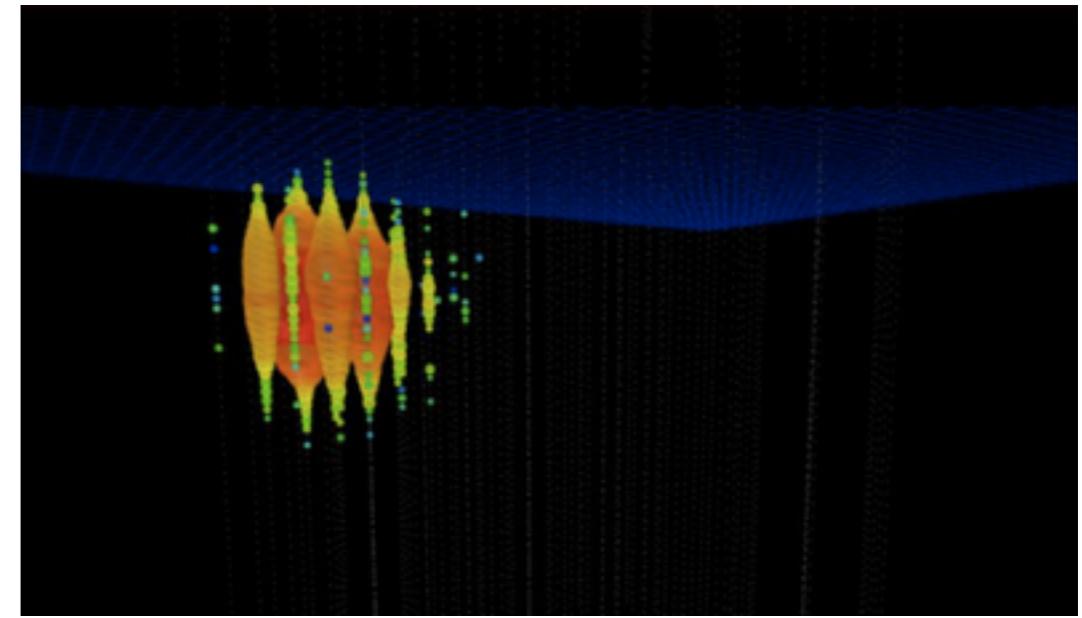
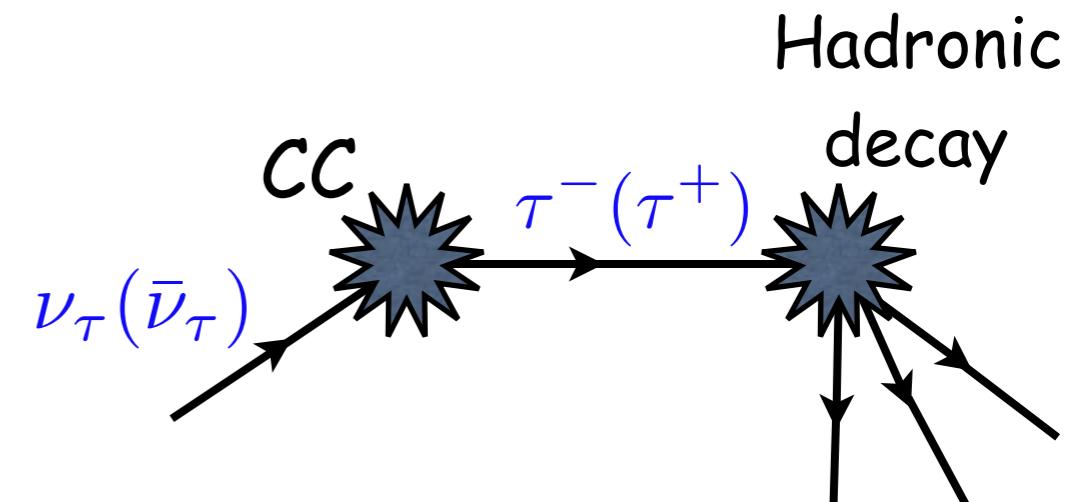
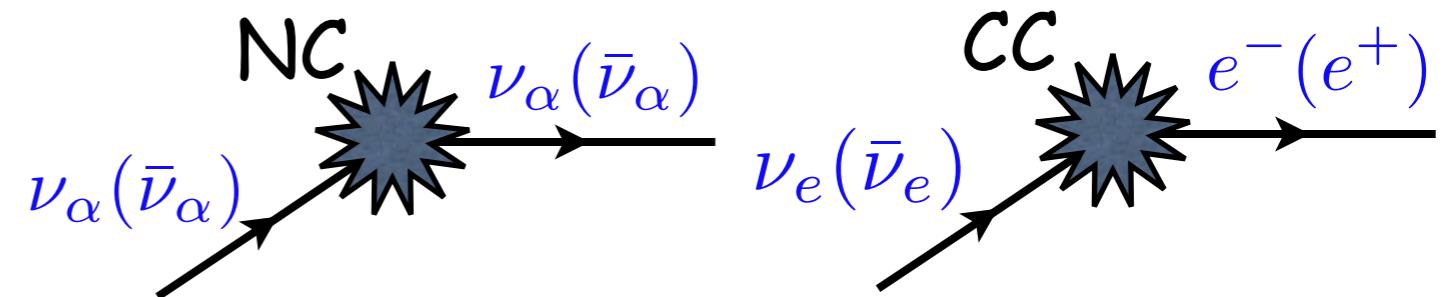


Flavoring at IceCube

muon-track events



cascade events



figures from
IceCube
website

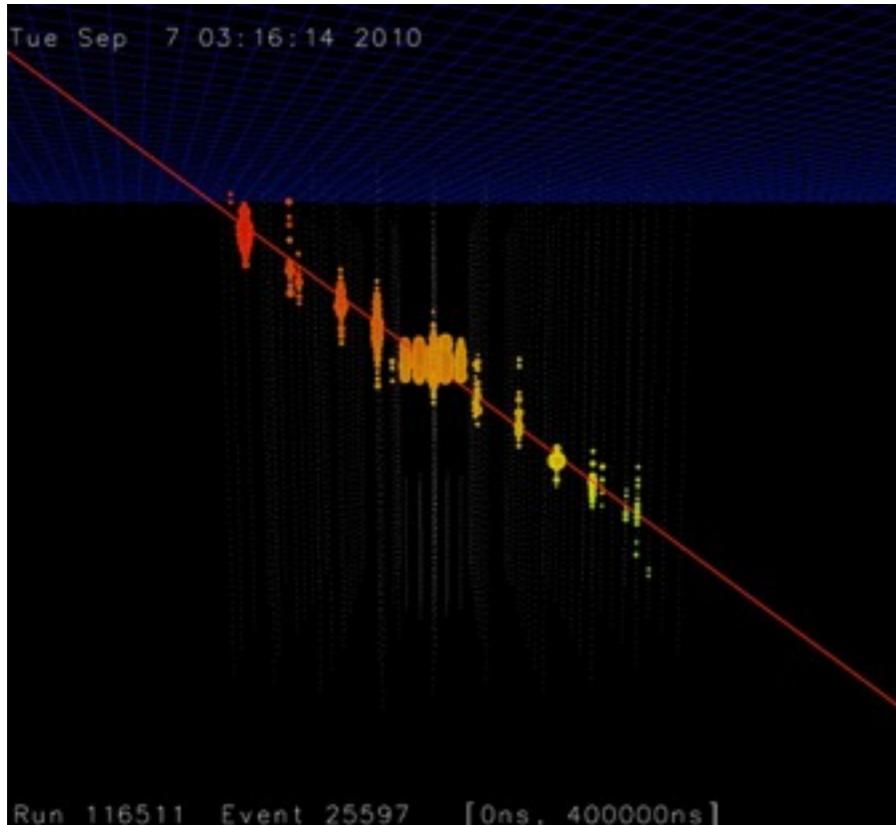
Flavoring at IceCube

muon-track events

great angular
resolution ($< 1^\circ$)

ν_τ

moderate energy
resolution ($\sigma_E \sim E$)



cascade events

$\nu_\alpha(\bar{\nu}_\alpha)$

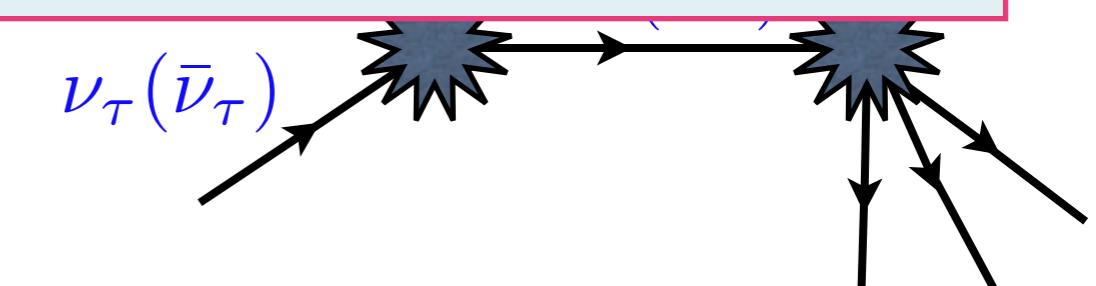
poor angular resolution
($< 10^\circ - 20^\circ$)

great energy resolution
($\sigma_E \sim 0.15 \times E$)

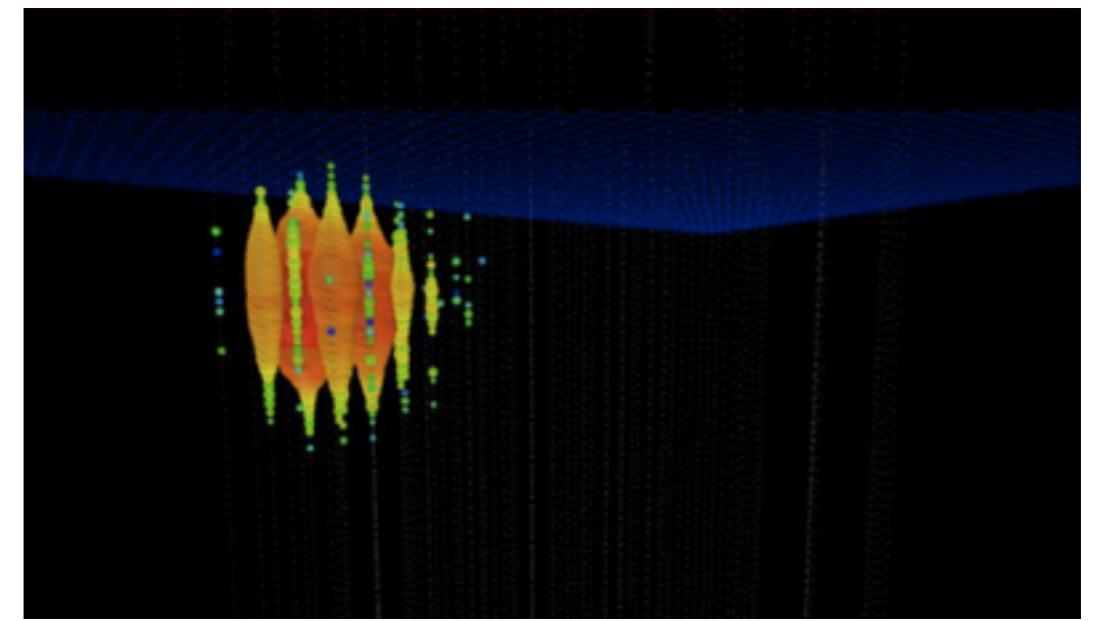
$\nu_\tau(\bar{\nu}_\tau)$

(e^+)

nic
y

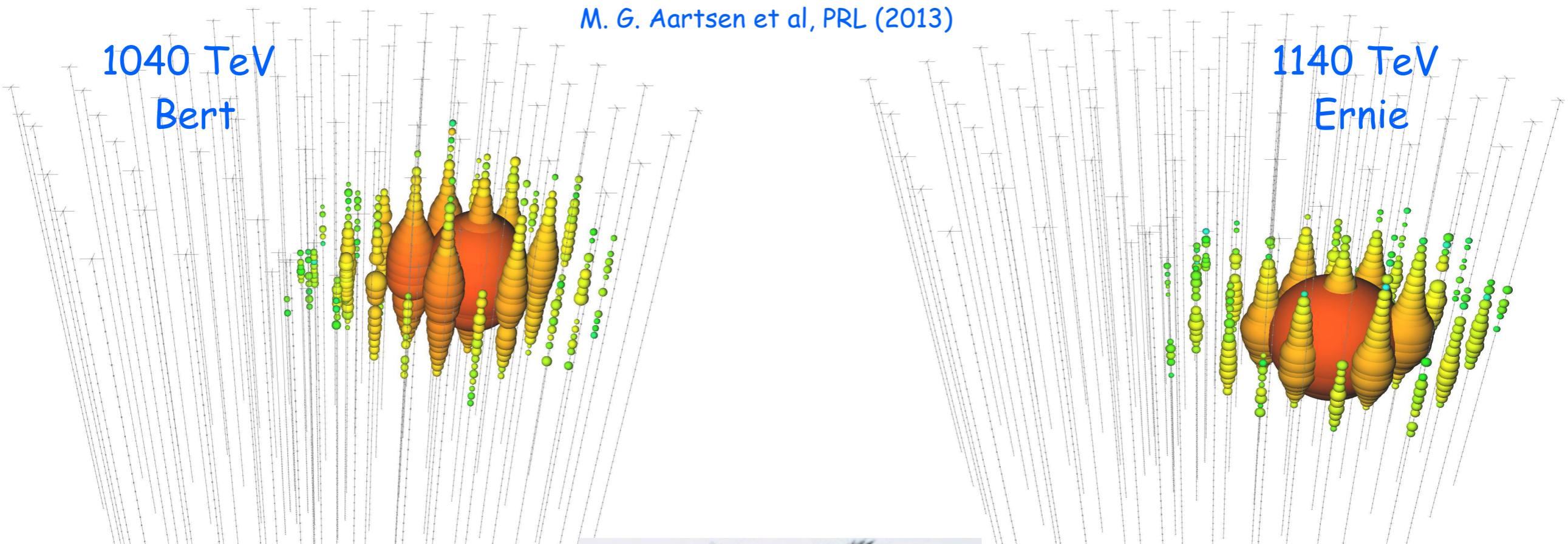


figures from
IceCube
website



Mission for IceCube began !

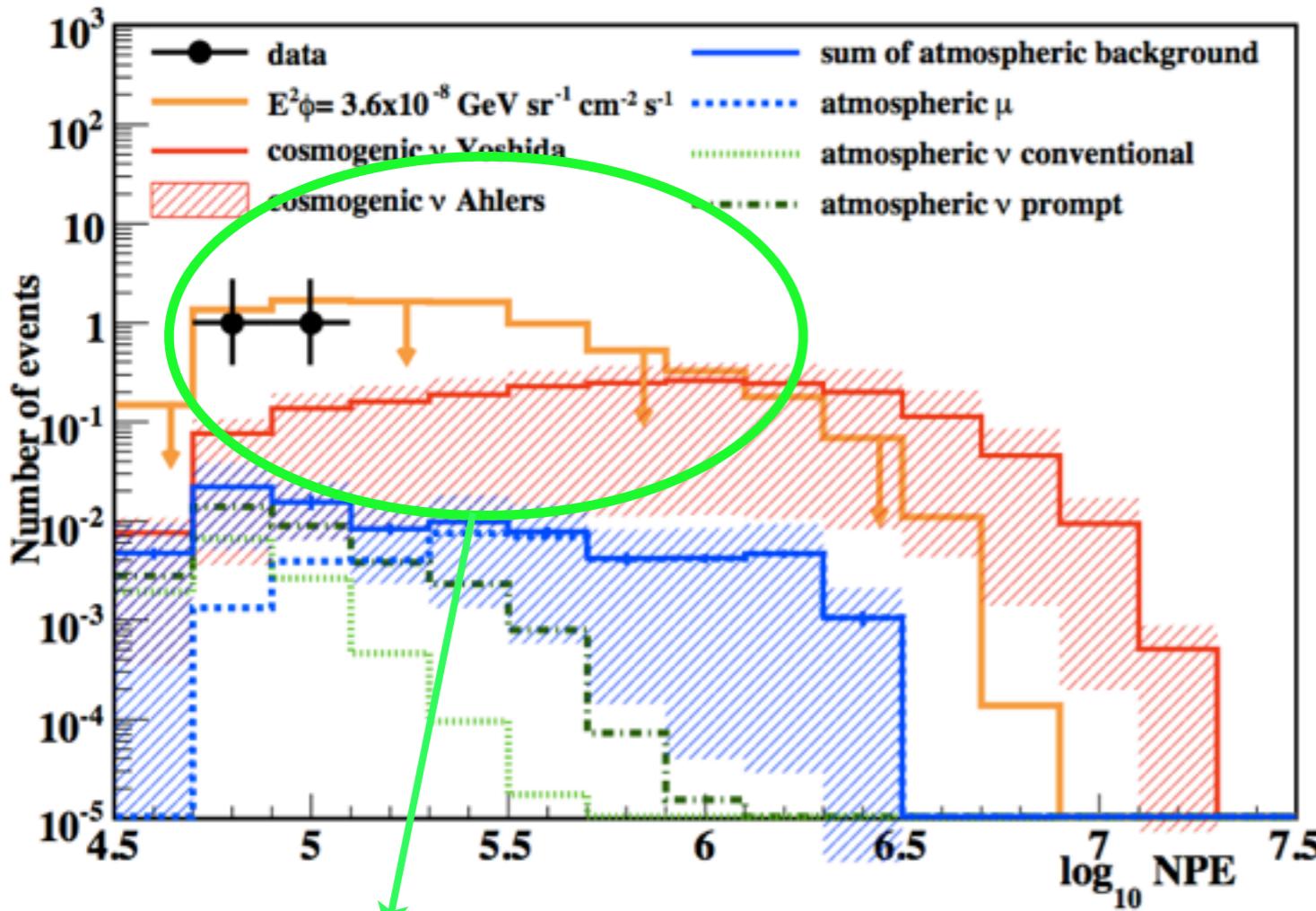
- ✓ The two PeV cascade events, 616 days livetime



Mission for IceCube began !

✓ The two PeV cascade events, 616 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Phys. Rev. Lett. 111 (2013), [arXiv:1304.5356]



it is like a cut-off at \sim PeV

demands more statistics

expected bkg. (conventional+prompt)
 $\sim 0.08(-0.057)(+0.041)$ sys.

excess of events $\sim 2.8\sigma$

GZK ? too low energy, more events
should be seen in higher energies

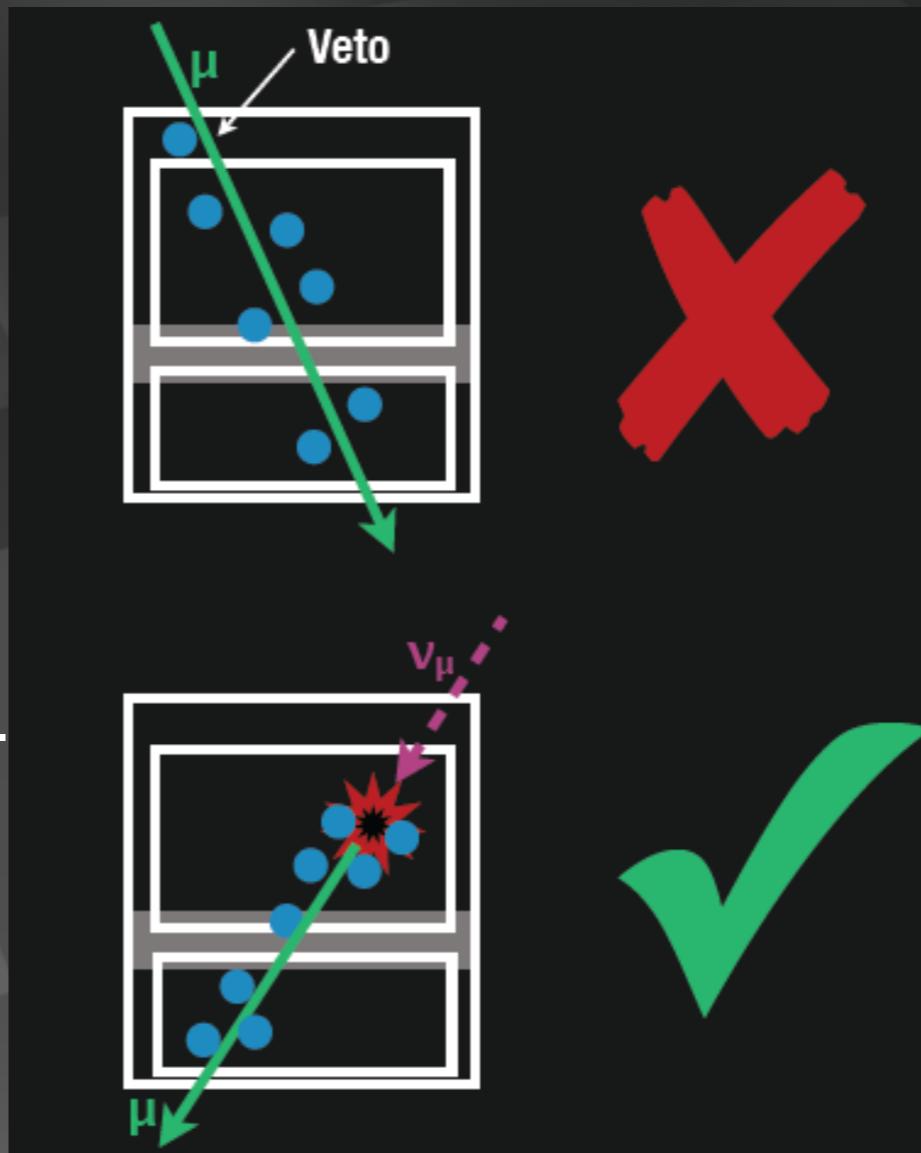
astrophysical ? an E^{-2} spectrum
would give ~ 9 more events in
higher energies

flavor composition ?
NC of ν_a or CC of ν_e

isotropy ?

HESE analysis

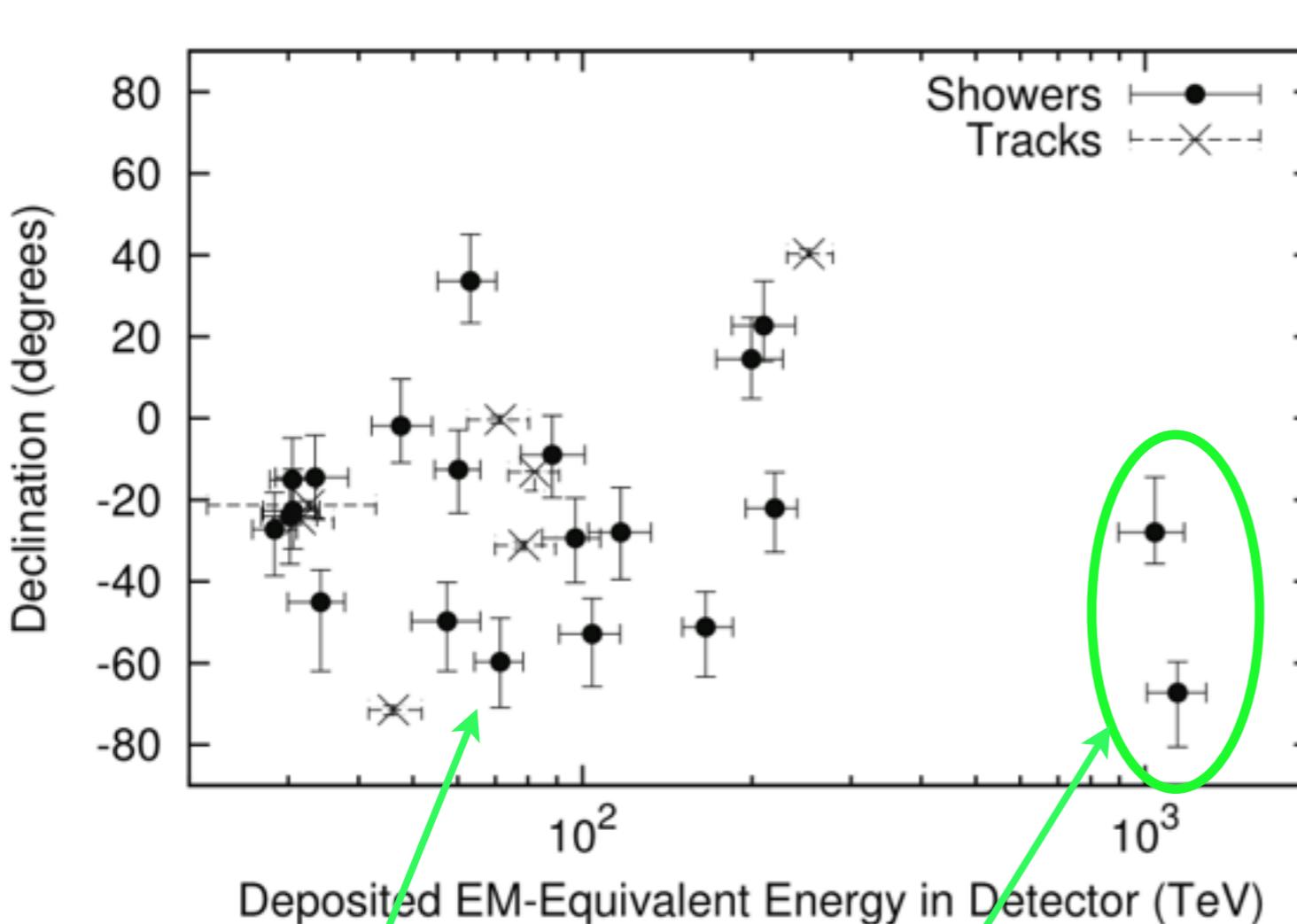
- ✓ select events interacting inside the detector only
- ✓ no light in the veto region
- ✓ veto for atmospheric muons and neutrinos (which are typically accompanied by muons)
- ✓ energy measurement: total absorption calorimetry



Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



The whole family!

previous PeV
cascade events
(Bert and Ernie)

✓ 26 more events

✓ all the new events are lower in energy

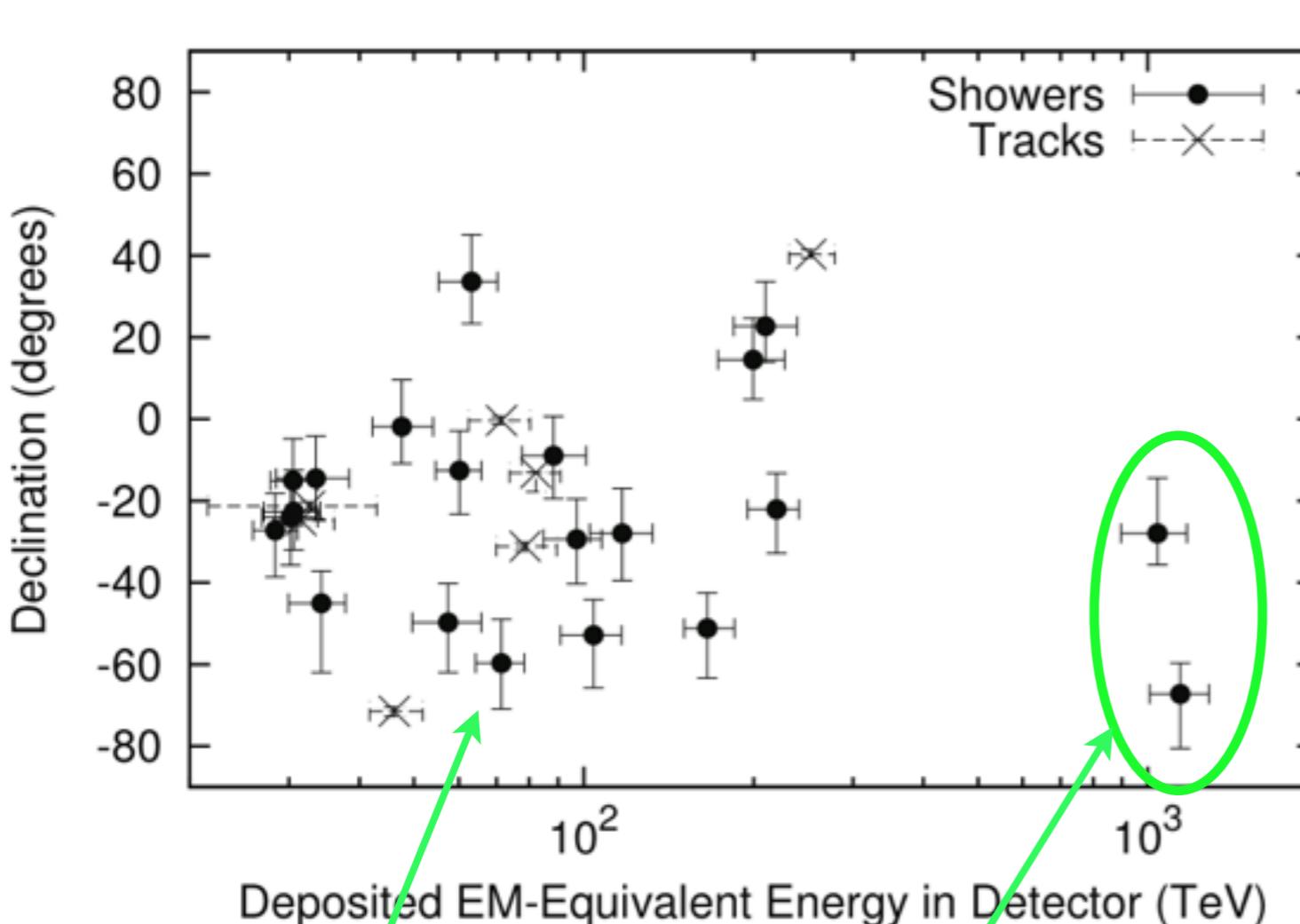
✓ expected bkg. (conventional+prompt)
 $\sim 10.6(+4.5)(-3.5)$ sys.

excess of events $\sim 4.3\sigma$

Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



The whole family!

previous PeV
cascade events
(Bert and Ernie)

which one?

atmospheric ?

astrophysical ?

or something else ?

✓ 26 more events

✓ all the new events are lower in energy

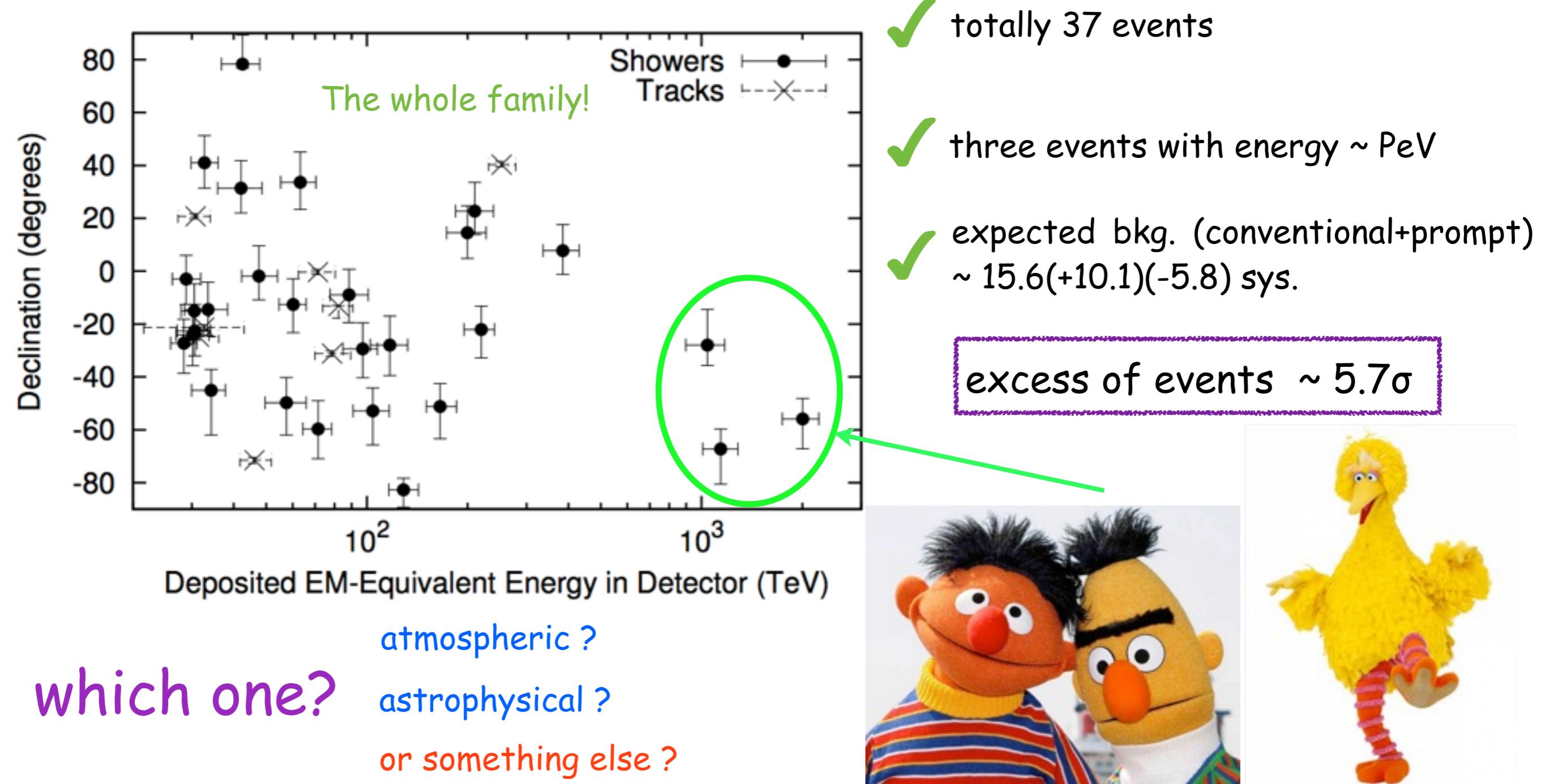
✓ expected bkg. (conventional+prompt)
~ 10.6(+4.5)(-3.5) sys.

excess of events ~ 4.3 σ

Mission for IceCube began !

✓ Looking for lower energy contained events, 988 days livetime

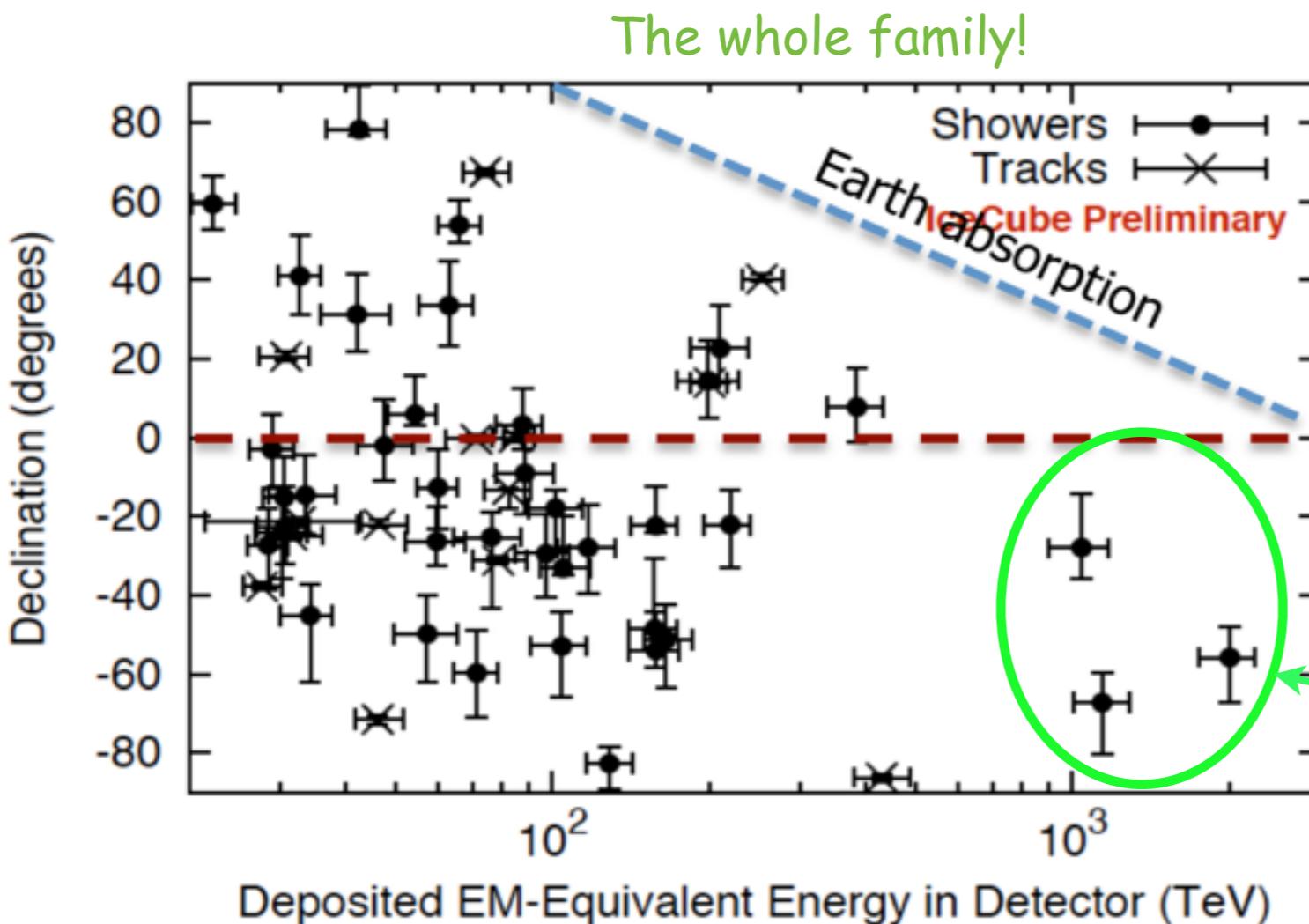
M. G. Aartsen et al. [IceCube Collaboration],
PRL 113 (2014), [arXiv:1405.5303]



Mission for IceCube began !

✓ Looking for lower energy contained events, 1347 days livetime

IPA 2015



which one?

atmospheric ?

astrophysical ?

or something else ?

✓ totally 54 events

✓ still three events with energy ~ PeV

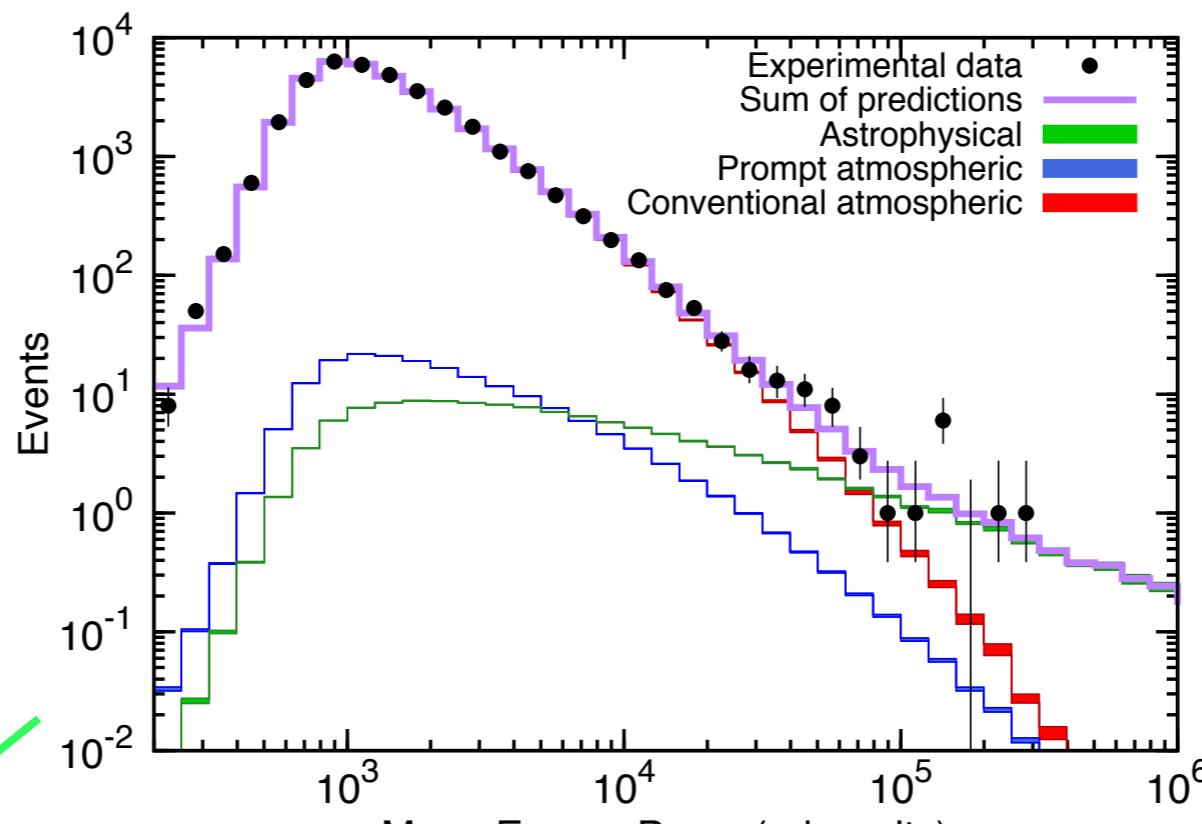
4 years of data

excess of events ~ 7 σ



Mission for IceCube began !

✓ Looking for muon-track events, 660 days livetime



$$\Phi(E_\nu) = 9.9_{-3.4}^{+3.9} \times 10^{-19} \left(\frac{E_\nu}{100 \text{ TeV}} \right)^{-2} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}, \text{ fitting the index: } 2.2 \pm 0.2$$

lowering the energy
threshold of HESE
analysis to 1 TeV

$$\Phi(E_\nu) = 2.06_{-0.3}^{+0.4} \times 10^{-18} \left(\frac{E_\nu}{100 \text{ TeV}} \right)^{-2.46 \pm 0.12} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}$$

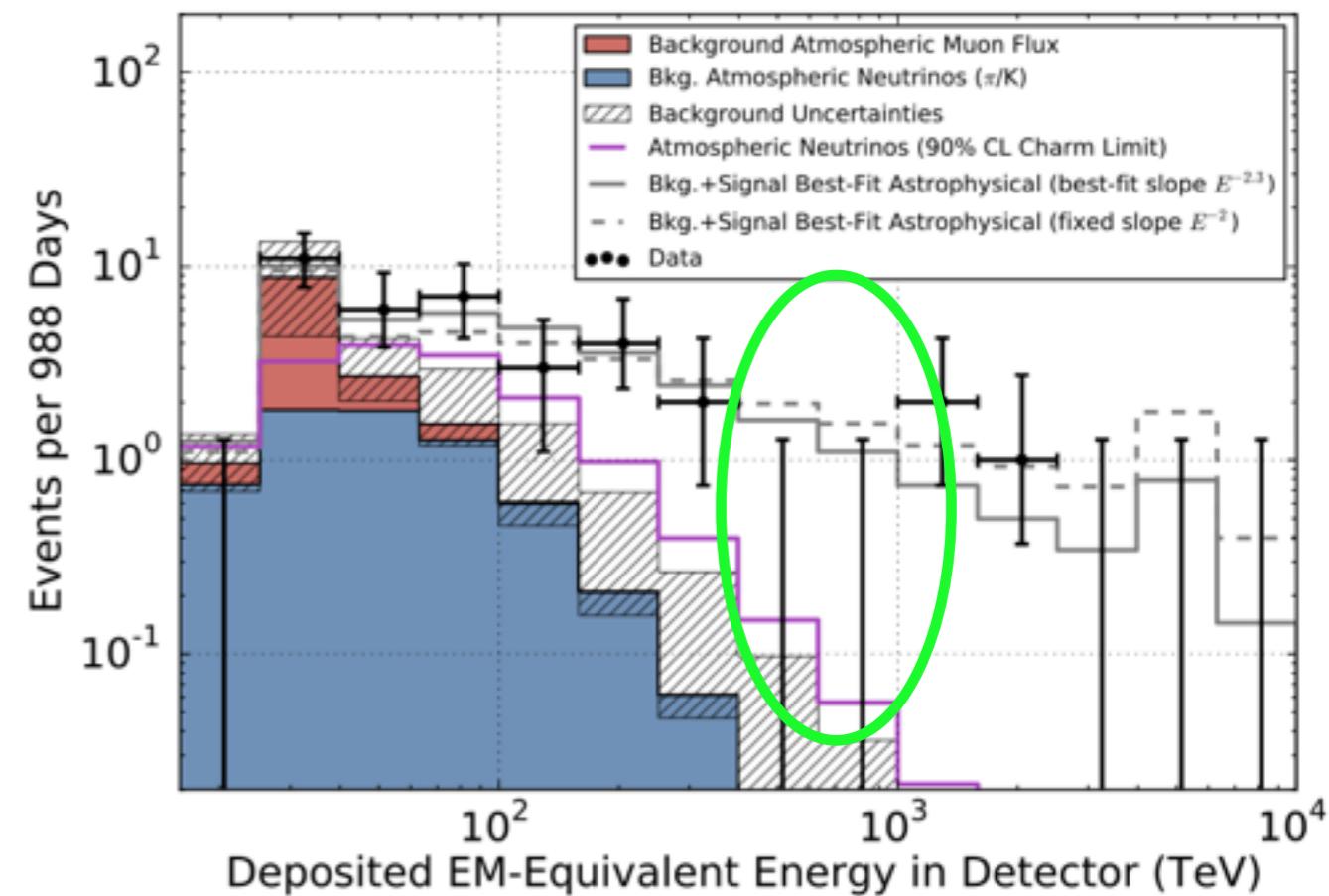
IceCube data

Some features of the observed spectrum

IceCube data

Some features of the observed spectrum

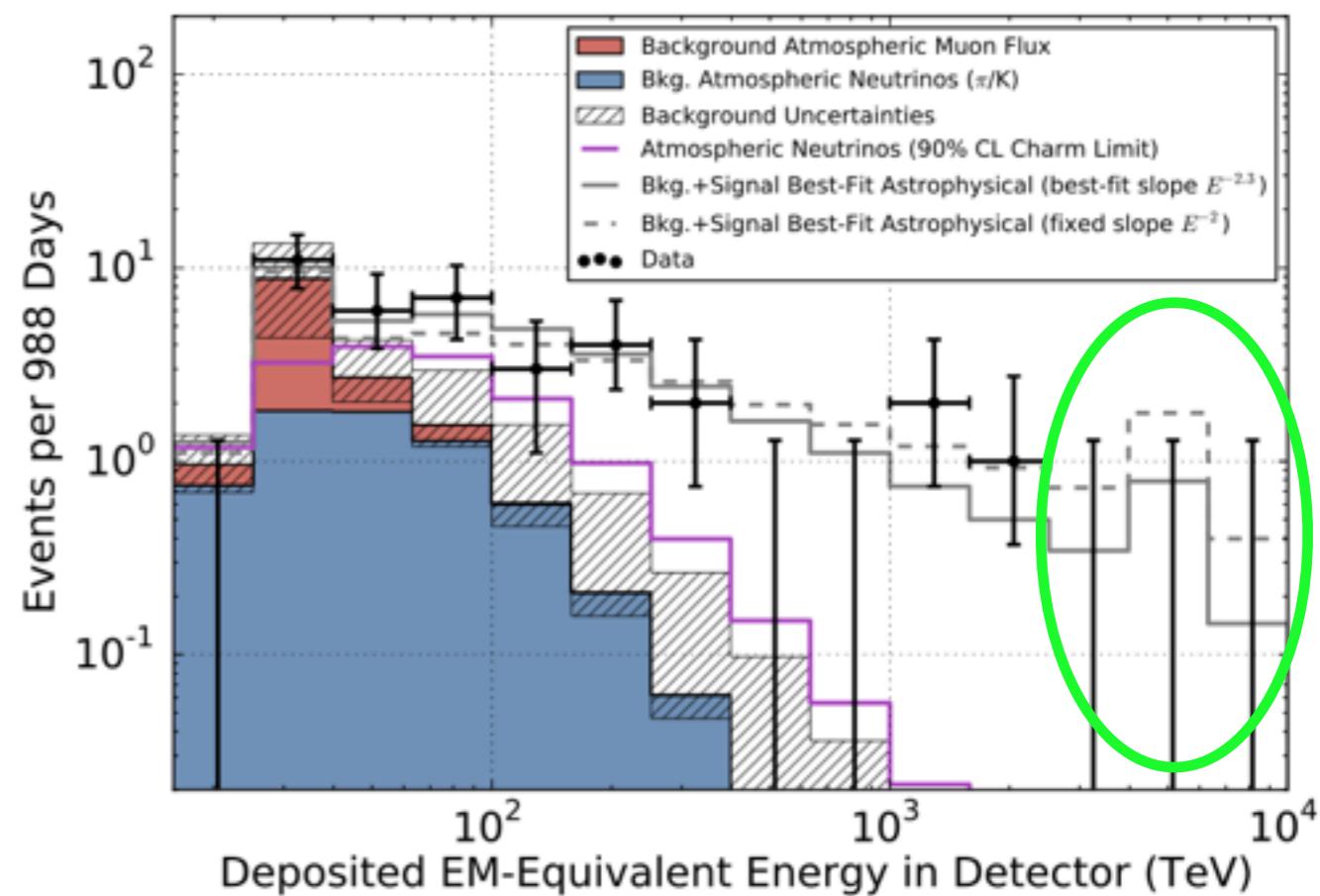
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV



IceCube data

Some features of the observed spectrum

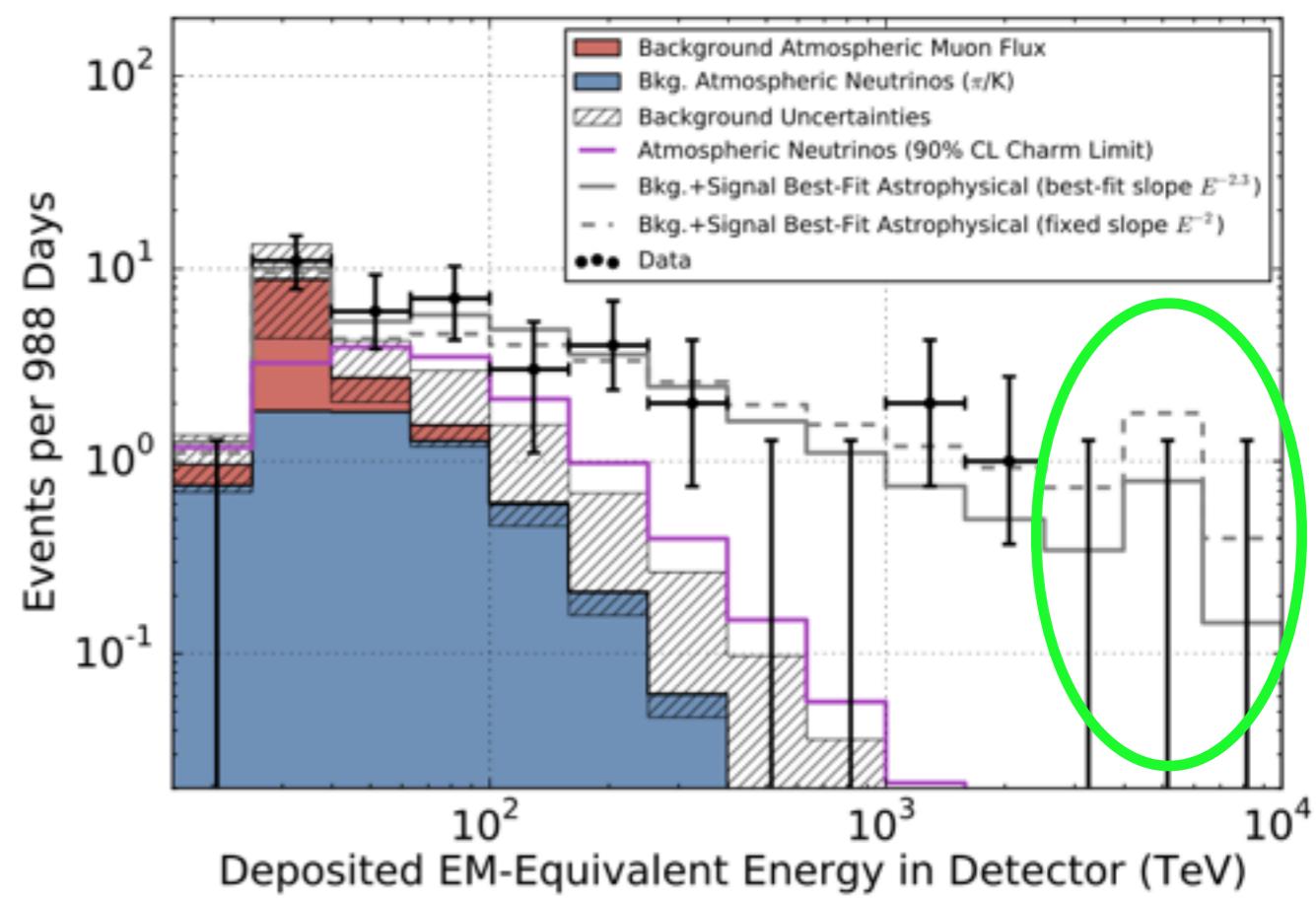
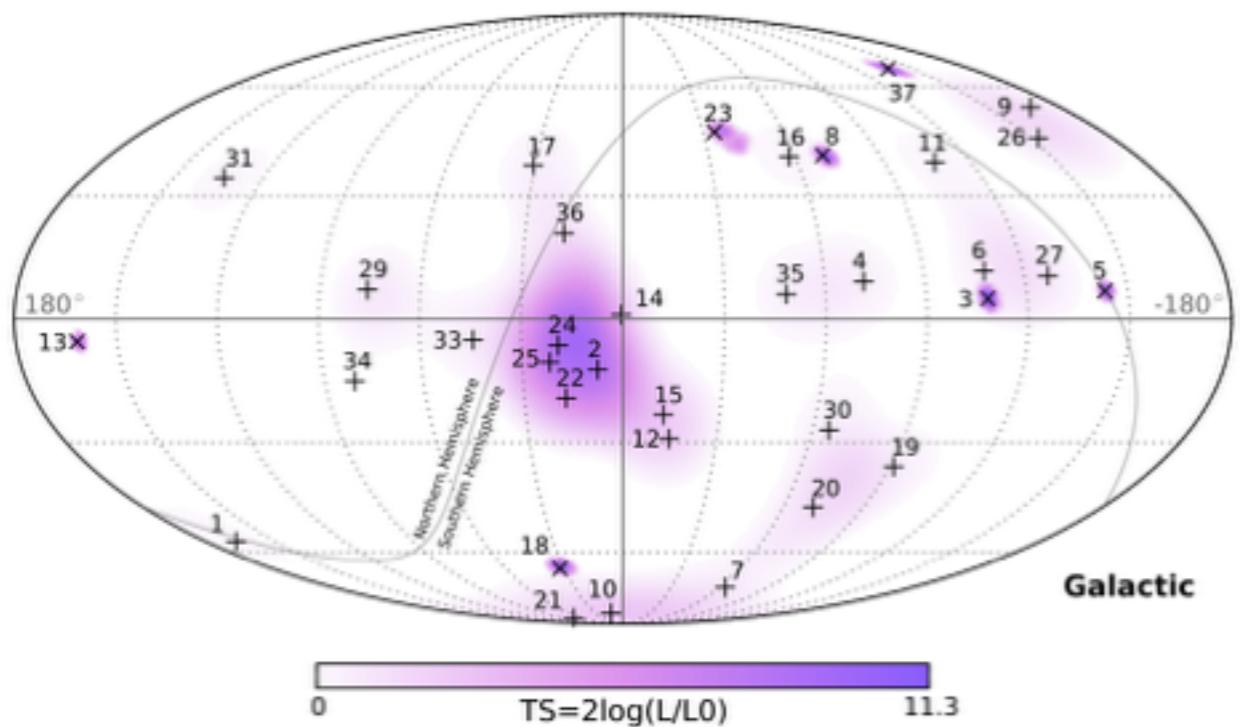
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV



IceCube data

Some features of the observed spectrum

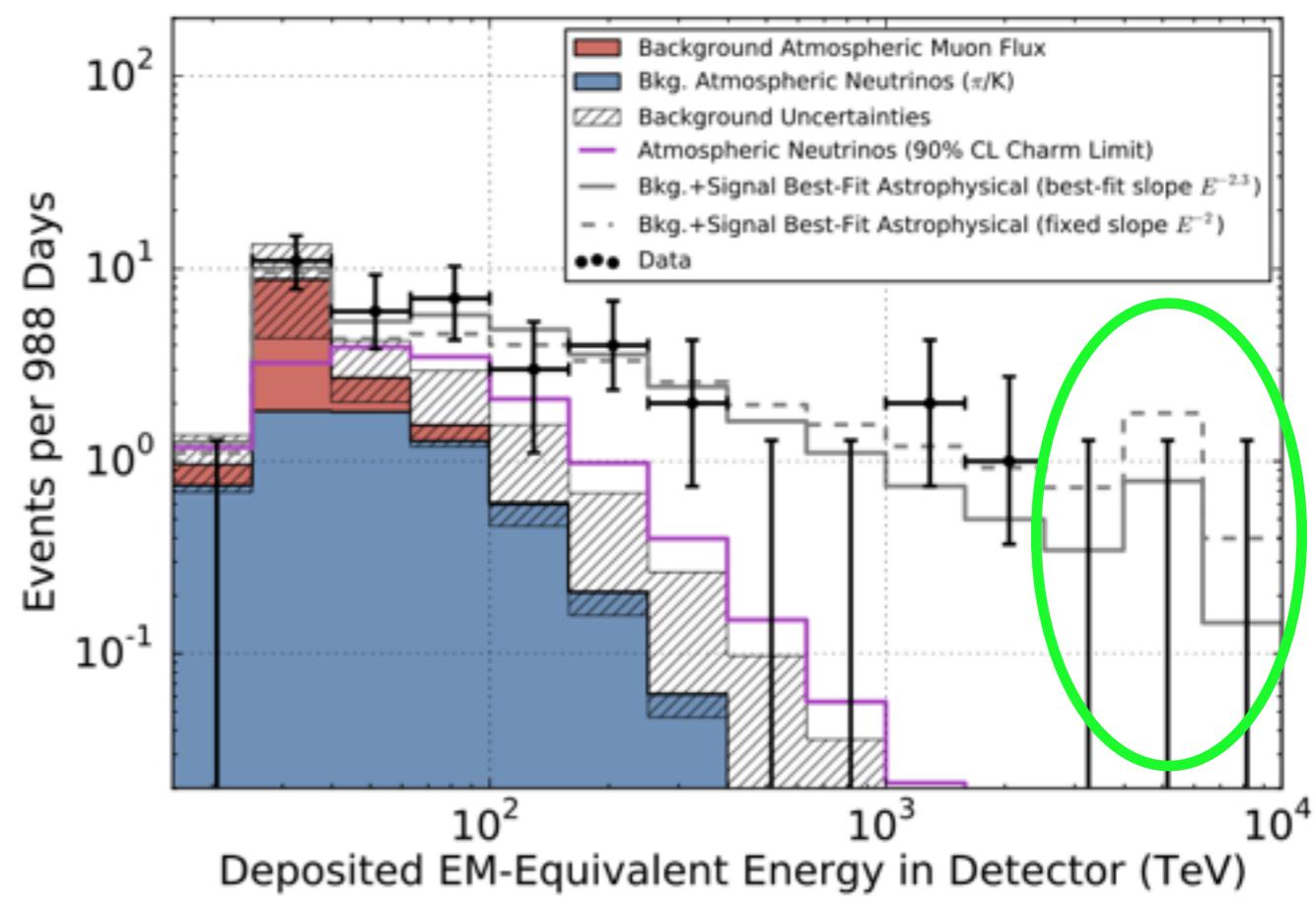
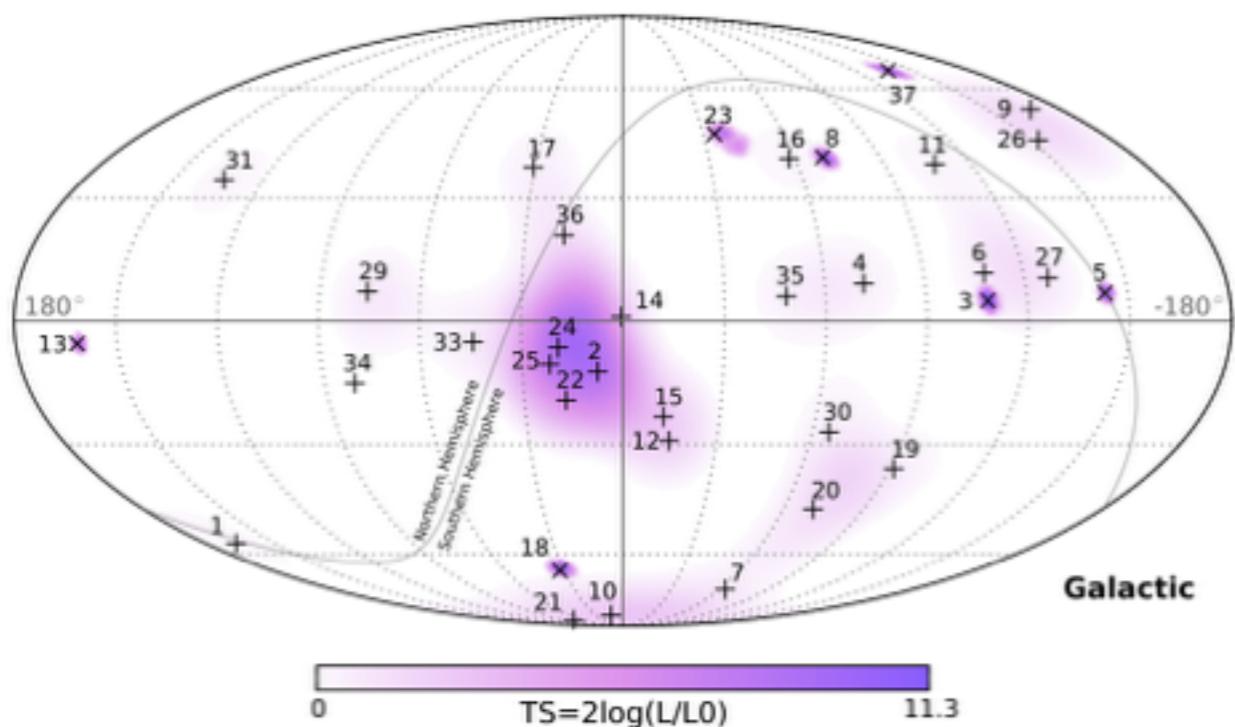
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)



IceCube data

Some features of the observed spectrum

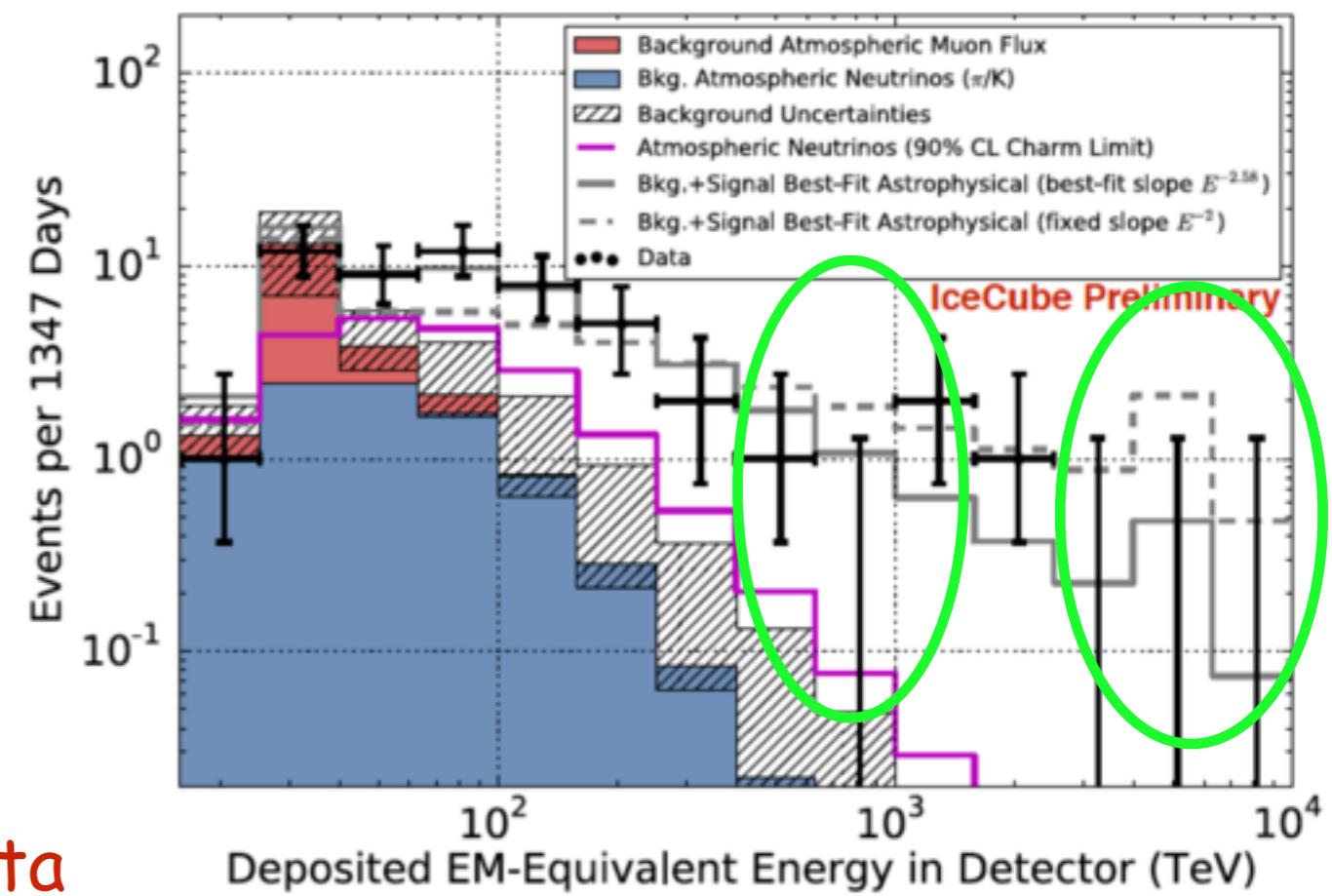
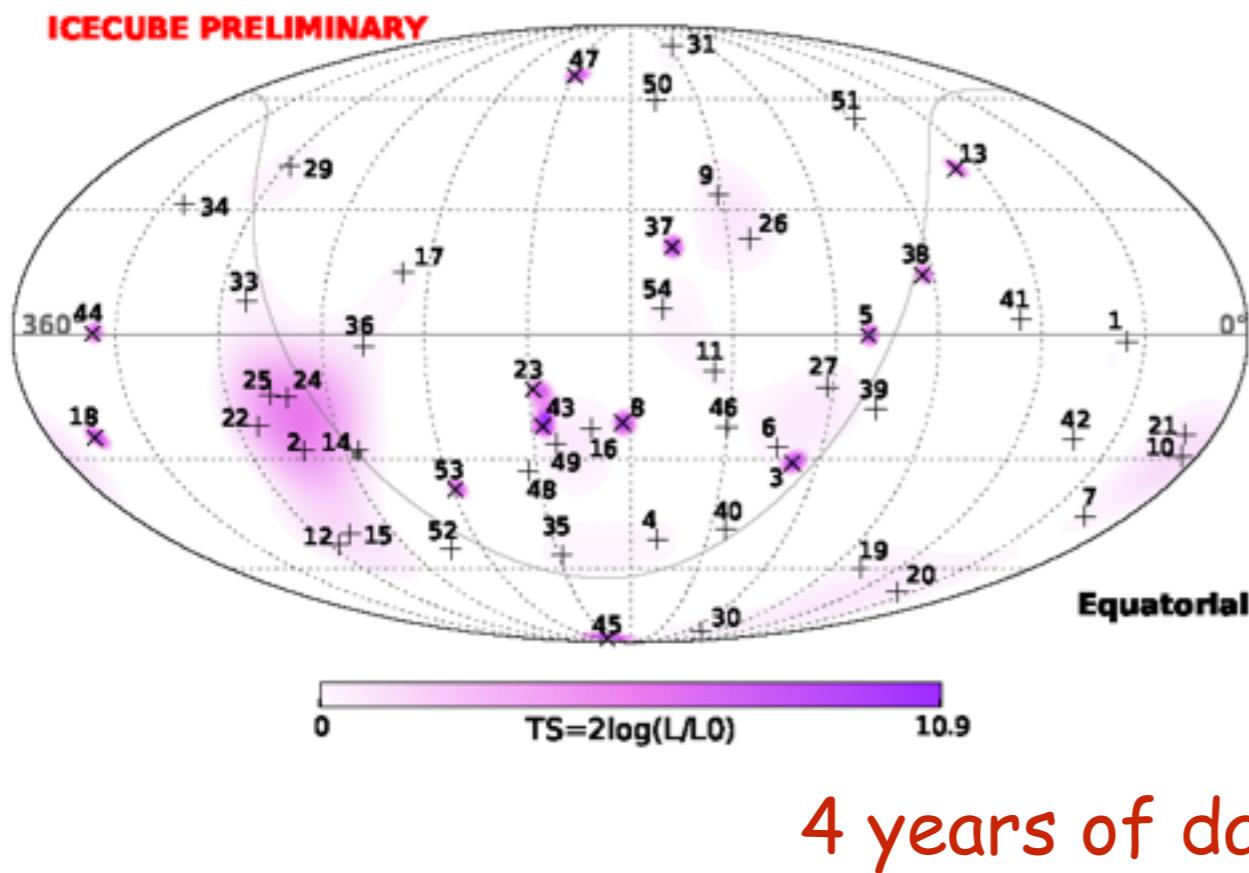
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
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- ⚠ none of the above-mentioned issues are significant



IceCube data

Some features of the observed spectrum

- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
 - ✓ cut-off in events: no events observed with energy > 2 PeV
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Interpretations of IceCube data

✓ "Conventional" interpretations of IceCube data

Cosmic ray sources

GRBs

Galaxy clusters

Star-forming galaxies

AGNs

Fermi bubbles

Galactic Center activities

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For a review

L. A. Anchordoqui, V. Barger, I. Cholis, H. Goldberg, D. Hooper, A. Kusenko, J. G. Learned and D. Marfatia et al., Journal of High Energy Astrophysics 1-2, 1 (2014) [arXiv:1312.6587 [astro-ph.HE]].

- M. D. Kistler, T. Stanev and H. Yuksel, arXiv:1301.1703 [astro-ph.HE]
K. Murase and K. Ioka, Phys. Rev. Lett. 111, no. 12, 121102 (2013) [arXiv:1306.2274 [astro-ph.HE]].
K. Murase, M. Ahlers and B. C. Lacki, Phys. Rev. D 88, no. 12, 121301 (2013) [arXiv:1306.3417 [astro-ph.HE]].
L. A. Anchordoqui, H. Goldberg, M. H. Lynch, A. V. Olinto, T. C. Paul and T. J. Weiler, arXiv:1306.5021 [astro-ph.HE].
R. Laha, J. F. Beacom, B. Dasgupta, S. Horiuchi and K. Murase, Phys. Rev. D 88, 043009 (2013) [arXiv:1306.2309 [astro-ph.HE]].
S. Razzaque, Phys. Rev. D 88, 081302 (2013) [arXiv:1309.2756 [astro-ph.HE]].
C. Y. Chen, P. S. Bhupal Dev and A. Soni, Phys. Rev. D 89, no. 3, 033012 (2014) [arXiv:1309.1764 [hep-ph]].
M. Ahlers and K. Murase, Phys. Rev. D 90, 023010 (2014) [arXiv:1309.4077 [astro-ph.HE]].
I. Tamborra, S. Ando and K. Murase, JCAP 1409, no. 09, 043 (2014) [arXiv:1404.1189 [astro-ph.HE]].
M. Kachelriess and S. Ostapchenko, Phys. Rev. D 90, 083002 (2014) [arXiv:1405.3797 [astro-ph.HE]].
M. Ahlers and F. Halzen, arXiv:1406.2160 [astro-ph.HE].
Y. Bai, A. J. Barger, V. Barger, R. Lu, A. D. Peterson and J. Salvado, Phys. Rev. D 90, 063012 (2014) [arXiv:1407.2243 [astro-ph.HE]].
A. Bhattacharya, R. Enberg, M. H. Reno and I. Sarcevic, arXiv:1407.2985 [astro-ph.HE].
C. Lunardini, S. Razzaque, K. T. Theodoseau and L. Yang, Phys. Rev. D 90, 023016 (2014) [arXiv:1311.7188 [astro-ph.HE]].

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Interpretations of IceCube data

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Cosmic ray sources

GRBs

Galaxy clusters

Star-forming galaxies

AGNs

Fermi bubbles

Galactic Center

- M. D. Kistler, T. Stanev and H. Yuksel, arXiv:1301.1703 [astro-ph.HE]
- K. Murase and K. Ioka, Phys. Rev. Lett. 111, no. 12, 121102 (2013) [arXiv:1306.2274 [astro-ph.HE]].
- K. Murase, M. Ahlers and B. C. Lacki, Phys. Rev. D 88, no. 12, 121301 (2013) [arXiv: 1306.3417 [astro-ph.HE]].
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- S. Razzaque, Phys. Rev. D 88, 081302 (2013) [arXiv:1309.2756 [astro-ph.HE]].
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- M. Ahlers and F. Halzen, arXiv:1406.2160 [astro-ph.HE].
- Y. Bai, A. J. Barger, V. Barger, R. Lu, A. D. Peterson and J. Salvado, Phys. Rev. D 90, 063012 (2014) [arXiv:1407.2243 [astro-ph.HE]].
- A. Bhattacharya, R. Enberg, M. H. Reno and I. Sarcevic, arXiv:1407.2985 [astro-ph.HE].
- C. Lunardini, S. Razzaque, K. T. Theodoreau and L. Yang, Phys. Rev. D 90, 023016 (2014) [arXiv:1311.7188 [astro-ph.HE]].

Some class of sources are starting to be excluded

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Interpretations of IceCube data

✓ "New Physics" interpretations of IceCube data

Lepto-quarks

Y. Ema, R. Jinno and T. Moroi, Phys. Lett. B 733, 120 (2014) [arXiv:1312.3501 [hep-ph]].

K. Ioka and K. Murase, PTEP 2014, (2014) [arXiv:1404.2279 [astro-ph.HE]].

K. C. Y. Ng and J. F. Beacom, Phys. Rev. D 90, 065035 (2014) [arXiv:1404.2288 [astro-ph.HE]].

M. Ibe and K. Kaneta, Phys. Rev. D 90, 053011 (2014) [arXiv:1407.2848 [hep-ph]].

V. Barger and W. Y. Keung, Phys. Lett. B 727, 190 (2013) [arXiv:1305.6907 [hep-ph]].

B. Feldstein, A. Kusenko, S. Matsumoto and T. T. Yanagida, Phys. Rev. D 88, no. 1, 015004 (2013) [arXiv:1303.7320 [hep-ph]].

Y. Bai, R. Lu and J. Salvado, arXiv:1311.5864 [hep-ph].

A. Bhattacharya, M. H. Reno and I. Sarcevic, JHEP 1406, 110 (2014) [arXiv:1403.1862 [hep-ph]].

J. Zavala, Phys. Rev. D 89, 123516 (2014) [arXiv:1404.2932 [astro-ph.HE]].

A. Bhattacharya, R. Gandhi and A. Gupta, arXiv:1407.3280 [hep-ph].

C. Rott, K. Kohri and S. C. Park, arXiv:1408.4575 [hep-ph].

T. Higaki, R. Kitano and R. Sato, JHEP 1407, 044 (2014) [arXiv:1405.0013 [hep-ph]].

A. Esmaili and P. D. Serpico, JCAP 1311, 054 (2013), arXiv:1308.1105

A. Esmaili, S. K. Kang and P. D. Serpico, JCAP 1412, 054 (2014), arXiv:1410.5979

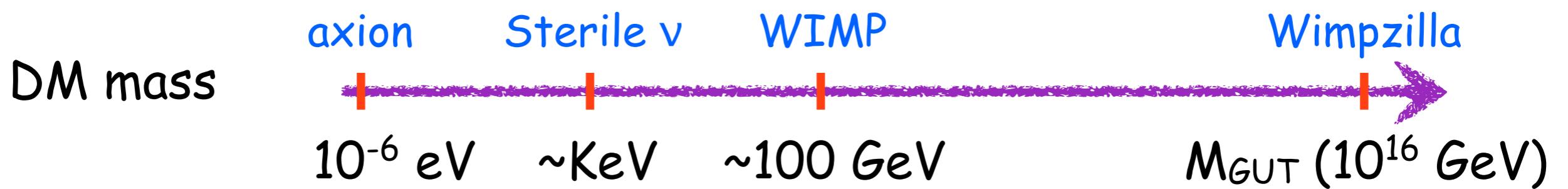
A. Esmaili and P. D. Serpico, JCAP 1510 (2015), arXiv:1505.06486



The idea

decaying
dark matter

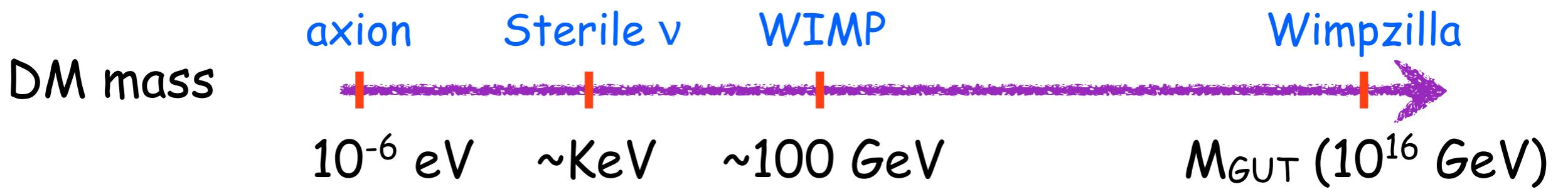
$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$



The idea

decaying
dark matter

$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$



caution: streetlight effect

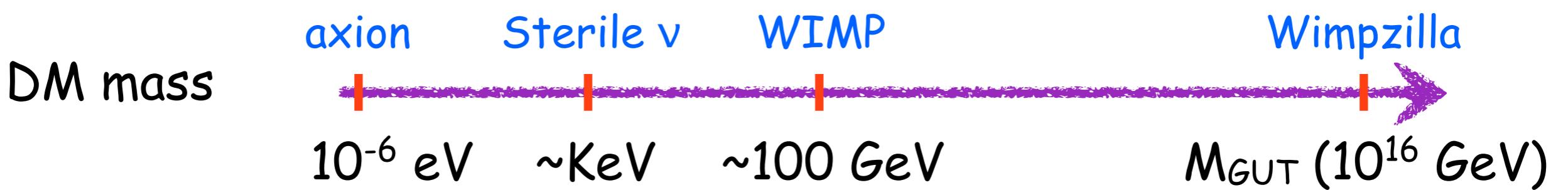


Mulla
Nasreddin



The idea

decaying
dark matter



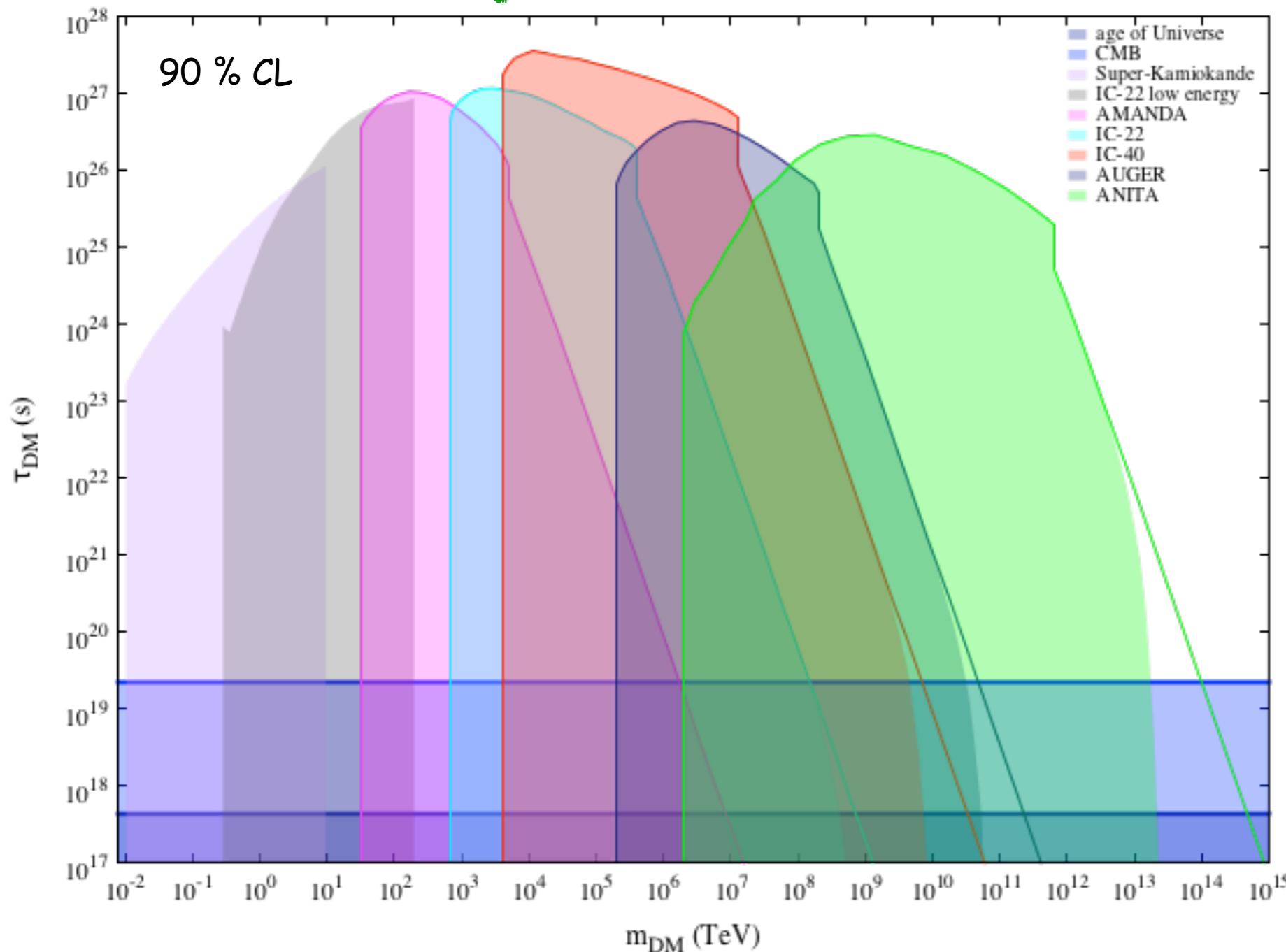
The idea

decaying
dark matter

$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$

$$\tau_{\text{DM}} \sim \mathcal{O}(10^{28} \text{ s})$$

A.E., A. Ibarra and O. L. G. Peres
JCAP (2012) [arXiv: 1205.5281]

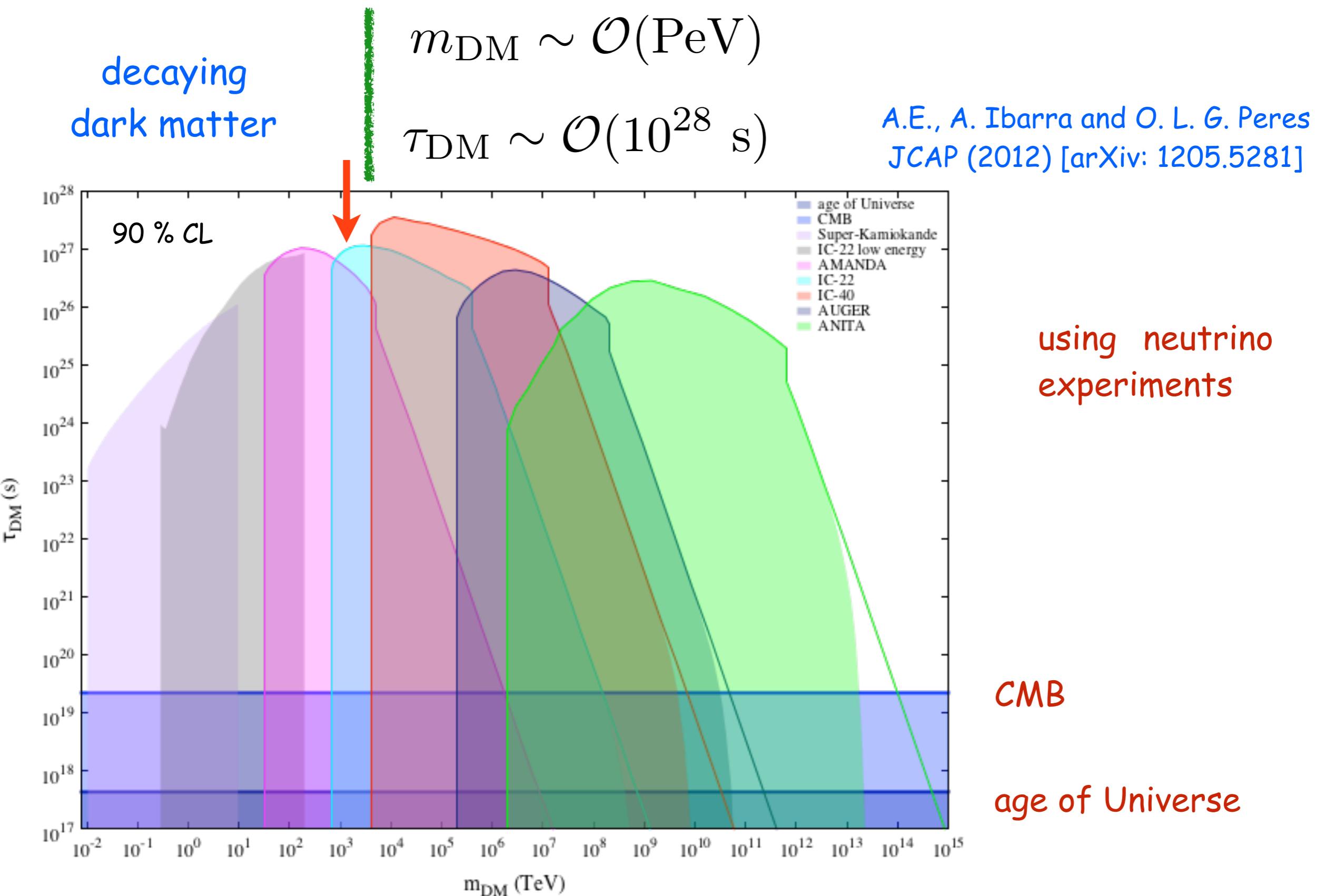


using neutrino
experiments

CMB

age of Universe

The idea



The idea

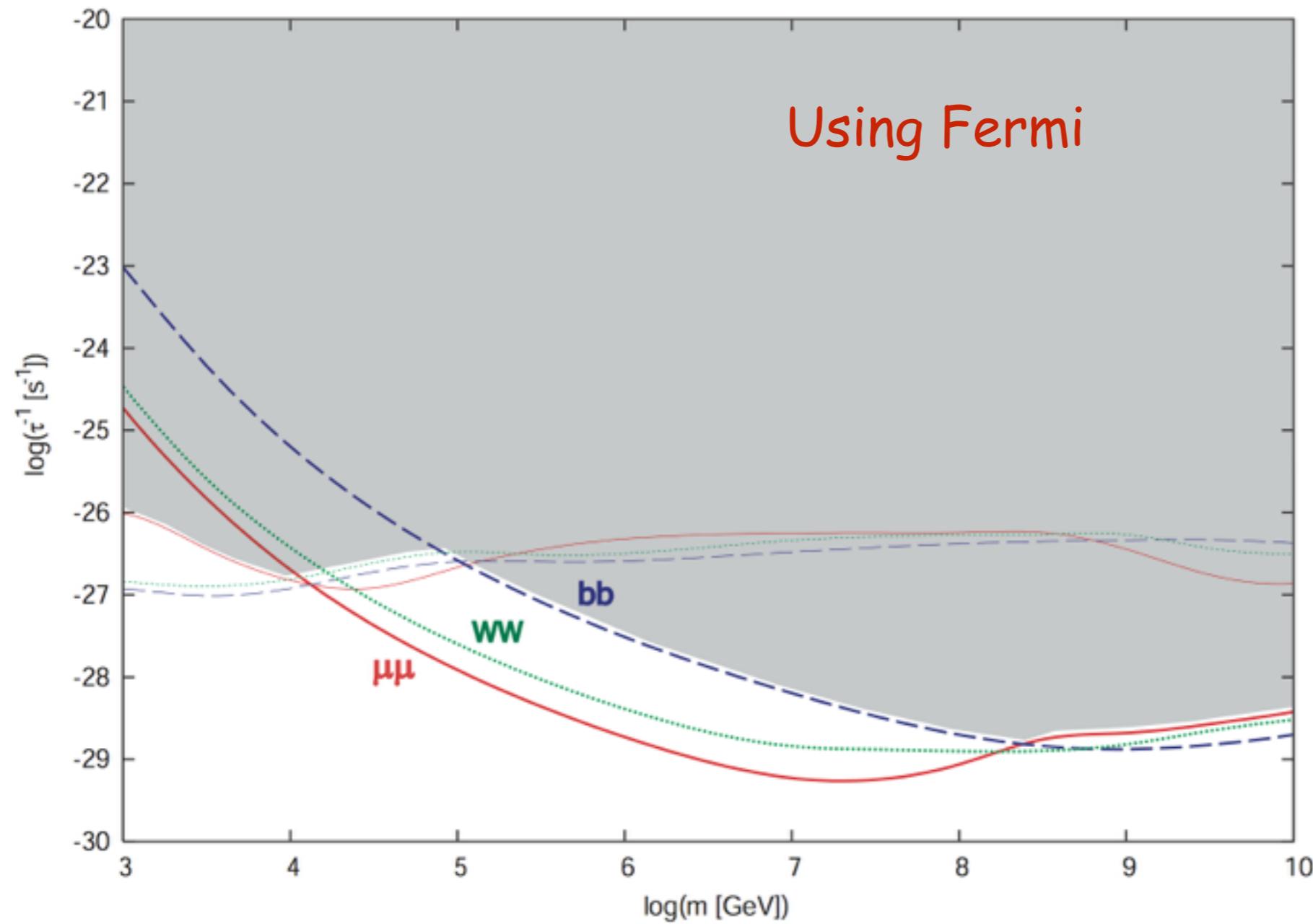
decaying
dark matter

$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$

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Using Fermi

K. Murase and J. Beacom
JCAP (2012) [arXiv: 1206.2595]



The idea

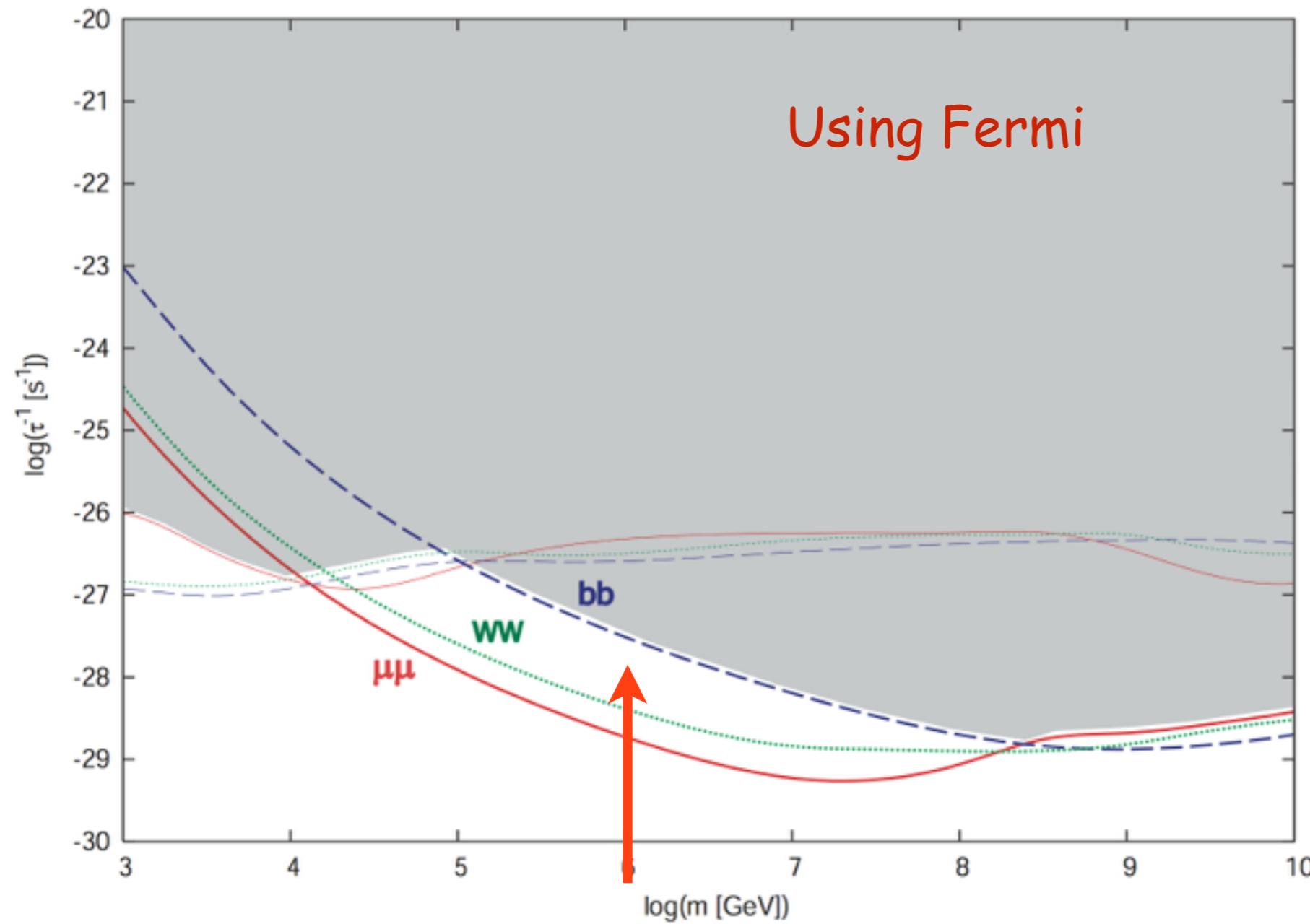
decaying
dark matter

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Using Fermi

K. Murase and J. Beacom
JCAP (2012) [arXiv: 1206.2595]



The idea

decaying
dark matter

From model-building
point of view

$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$

$$\tau_{\text{DM}} \sim \mathcal{O}(10^{28} \text{ s})$$

R-parity violating gravitino
hidden sector gauge boson
singlet fermion in extra-dim

Sterile neutrino

Sterile neutrino
+ U(1) gauge

mirror dark matter

EFT leptophilic DM

scalar + sterile neutrino

Feldstein, Kusenko, Matsumoto, Yanagida
PRD (2013) [arXiv: 1303.7320]

Higaki, Kitano and Sato
JHEP (2014) [arXiv: 1405.0013]

Ko and Tang
PLB (2015) [arXiv: 1508.02500]

Berezhiani [arXiv: 1506.09040]

Boucenna et. al.,
[arXiv: 1507.01000]

Dudas, Mambrini, Olive,
PRD [arXiv: 1412.3459]

The 2nd approach of
Mambrini's talk

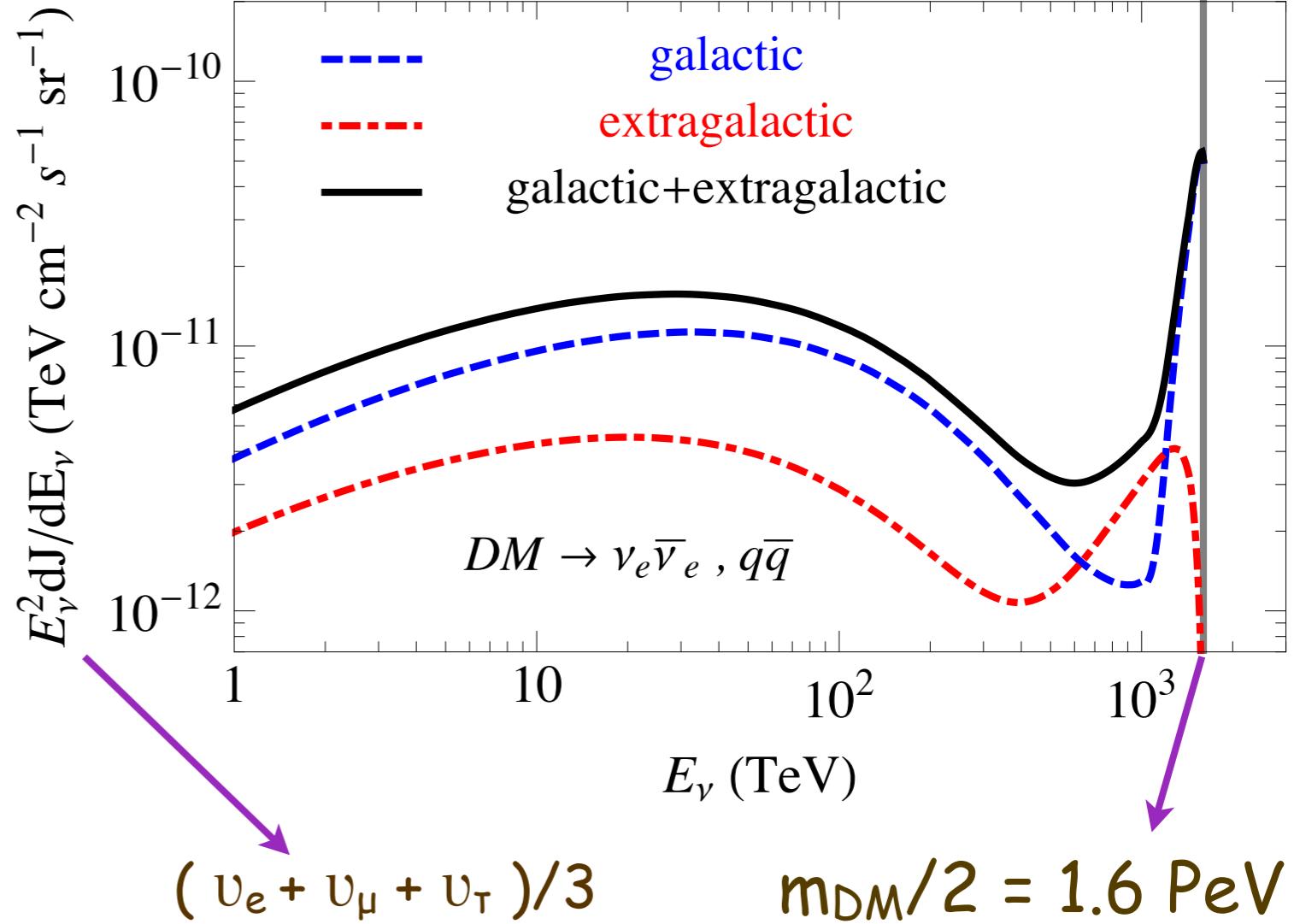
The idea

A. E., Pasquale D. Serpico, JCAP (2013) [arXiv:1308.1105]

decaying
dark matter

$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$

$$\tau_{\text{DM}} \sim \mathcal{O}(10^{28} \text{ s})$$



Galactic

$$\frac{d\Phi_{\text{gal}}}{dE_\nu} = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_{\text{DM}}[r(s, l, b)]$$

extra-Galactic

$$\frac{d\Phi_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [E_\nu(1+z)]$$

The idea

A. E., Pasquale D. Serpico, JCAP (2013) [arXiv:1308.1105]

decaying dark matter

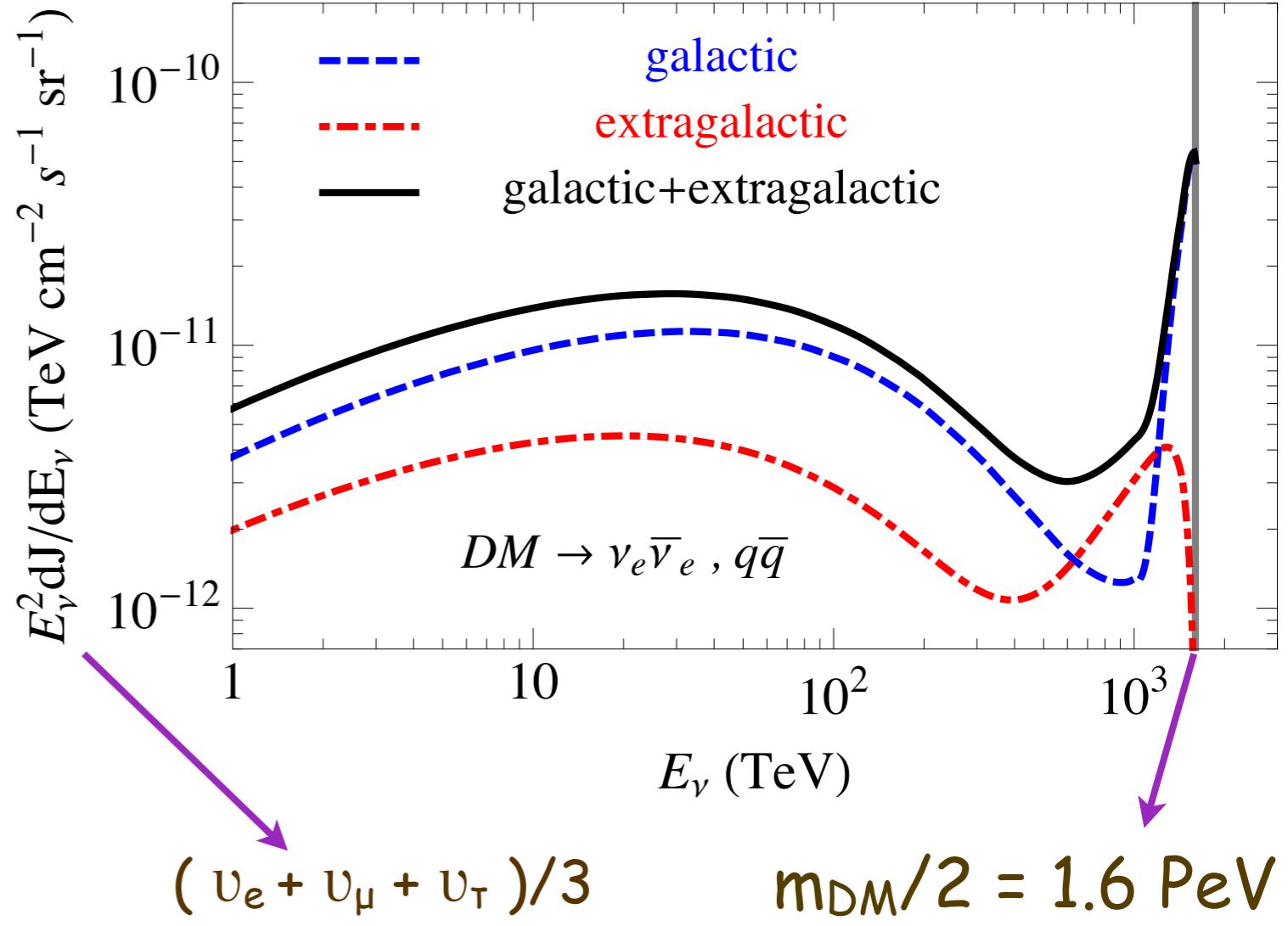
$$m_{\text{DM}} \sim \mathcal{O}(\text{PeV})$$

$$\tau_{\text{DM}} \sim \mathcal{O}(10^{28} \text{ s})$$

quarks

$$\frac{dN_\nu}{dE_\nu} = (1 - b_H) \left. \frac{dN_\nu}{dE_\nu} \right|_S + b_H \left. \frac{dN_\nu}{dE_\nu} \right|_H$$

neutrinos, charged leptons



Galactic

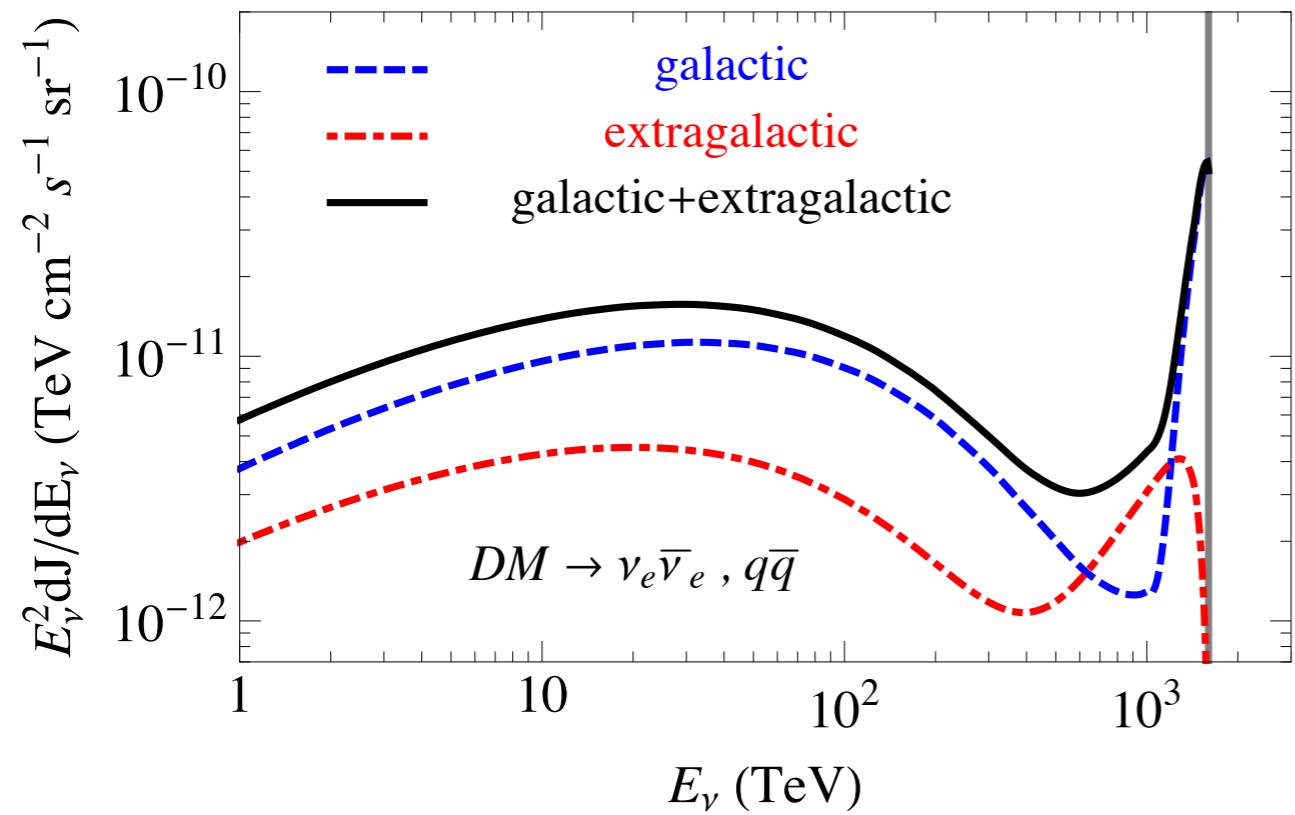
$$\frac{d\Phi_{\text{gal}}}{dE_\nu} = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_{\text{DM}}[r(s, l, b)]$$

extra-Galactic

$$\frac{d\Phi_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [E_\nu(1+z)]$$



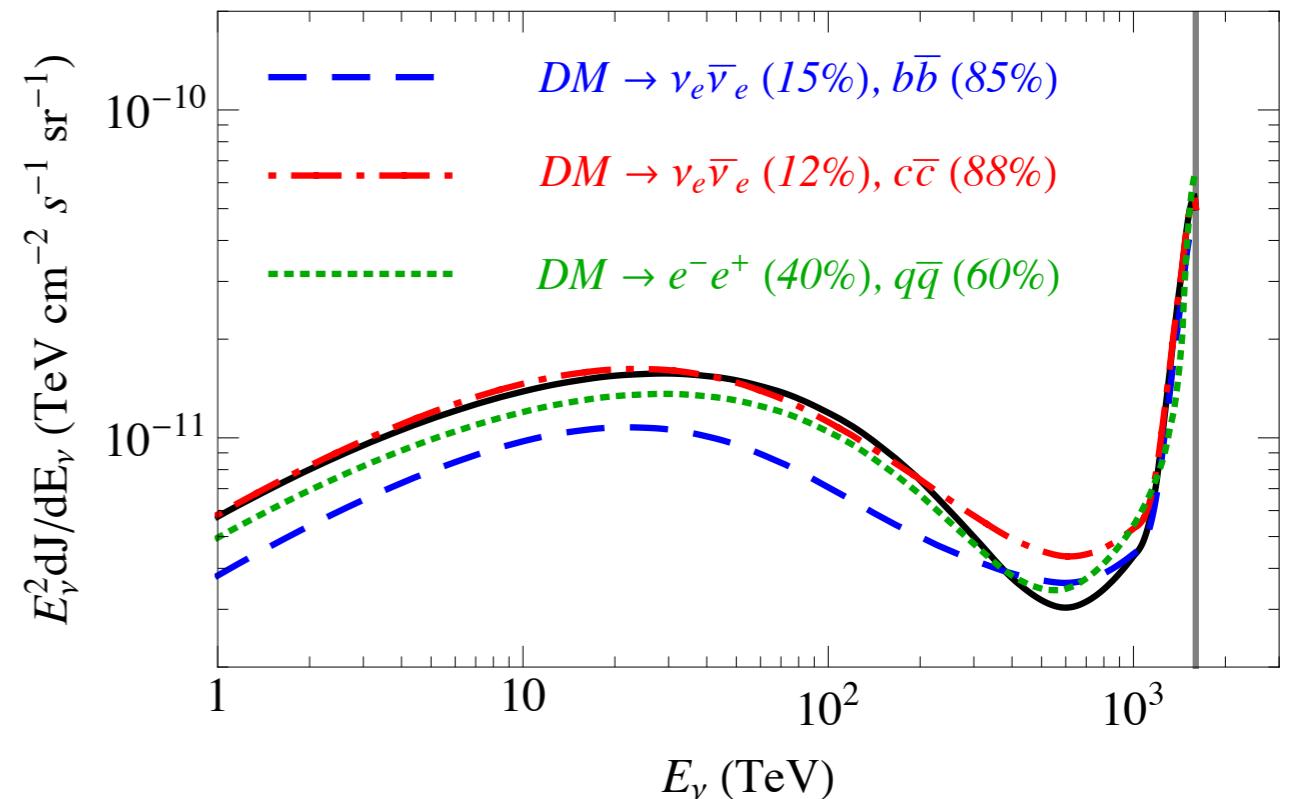
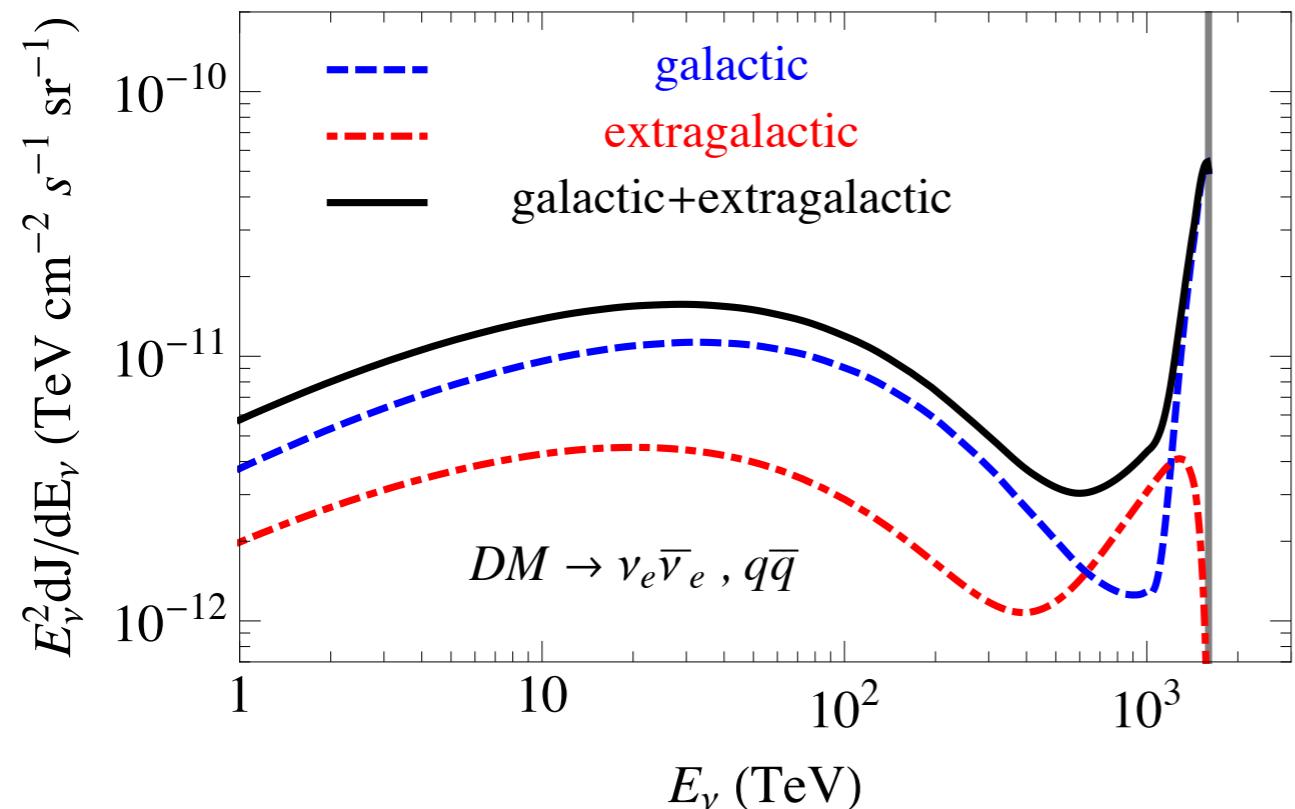
fine-tuned decay channels ?





fine-tuned decay channels ?

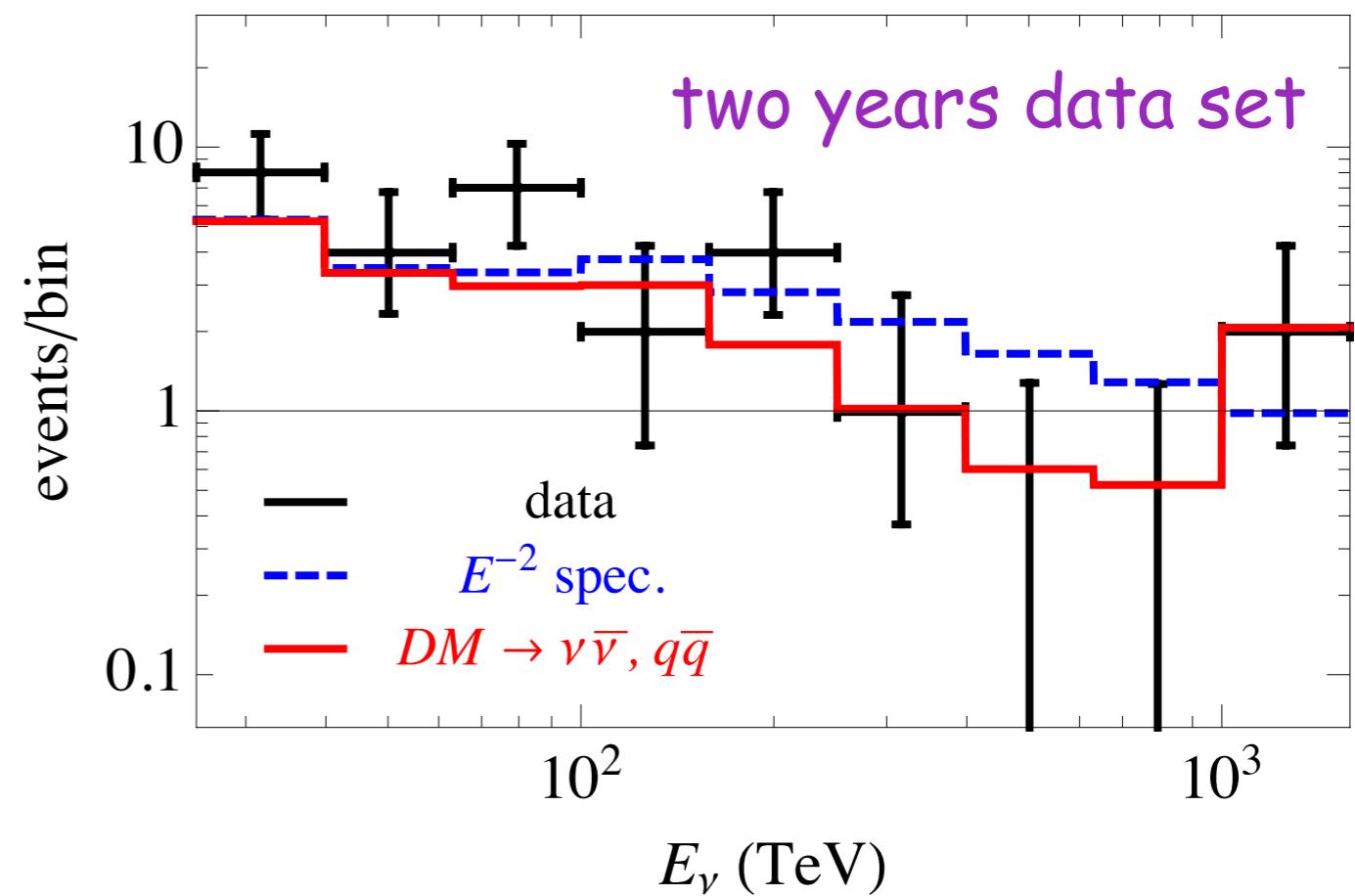
$$\tau_{DM} = (1-3) \times 10^{27} \text{ s}$$



the intriguing features are generic

Confronting with energy distribution of IceCube data

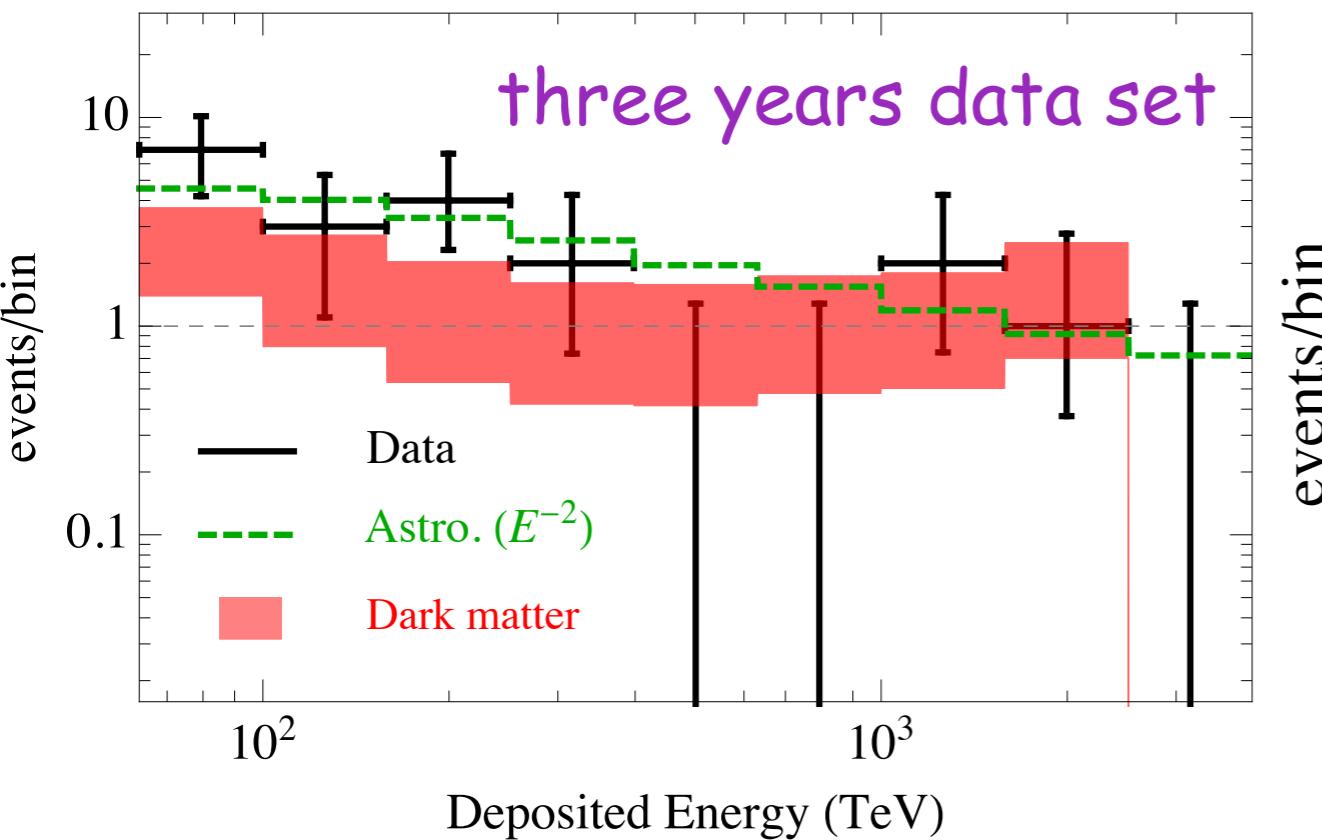
$$b_H = 0.12 \text{ and } \tau_{DM} = 2 \times 10^{27} \text{ s}$$



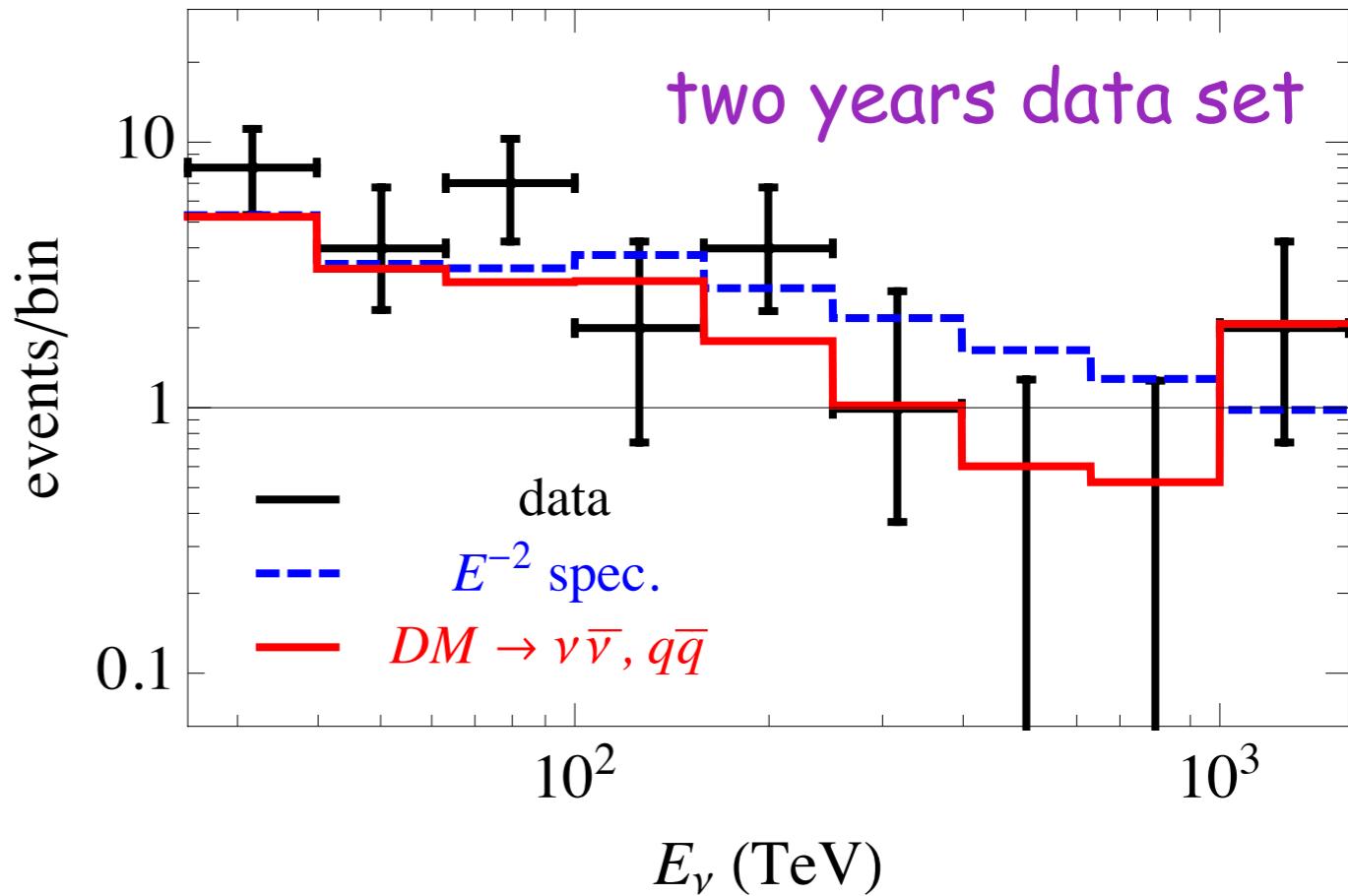
A. E. and P. D. Serpico,
JCAP (2013) [1308.1105]

Confronting with energy distribution of IceCube data

sterile neutrino model



$b_H = 0.12$ and $\tau_{DM} = 2 \times 10^{27}$ s



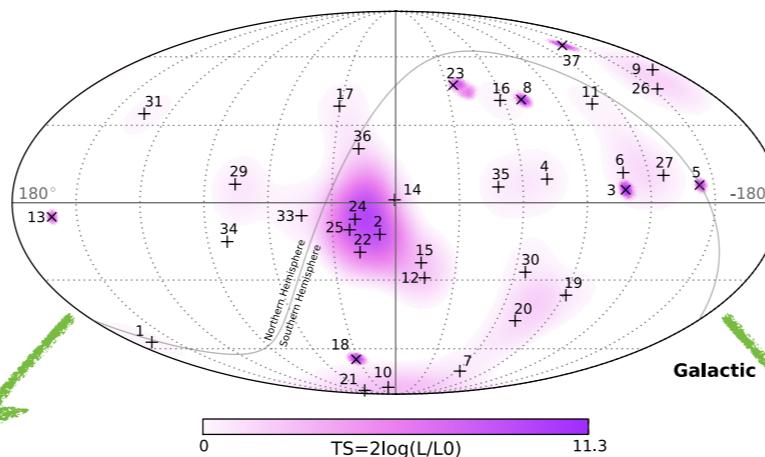
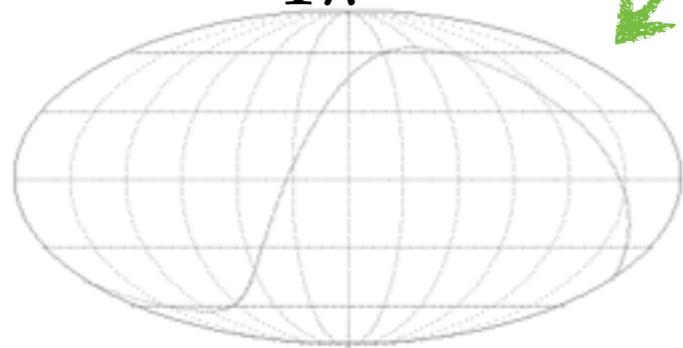
A. E., S. K. Kang and P. D. Serpico,
JCAP (2014) [1410.5979]

A. E. and P. D. Serpico,
JCAP (2013) [1308.1105]

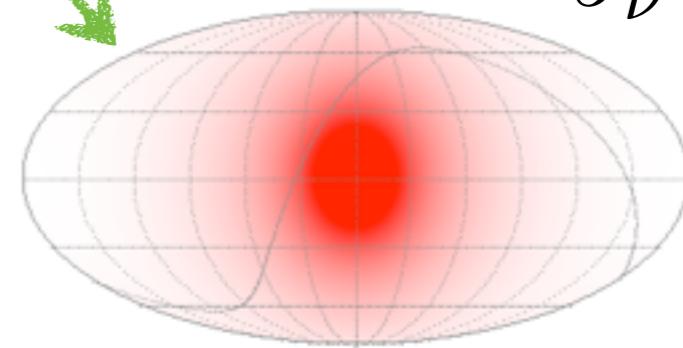
Angular distribution of neutrinos from decaying DM

✓ We would compare

$$p^{\text{iso}} = \frac{1}{4\pi}$$



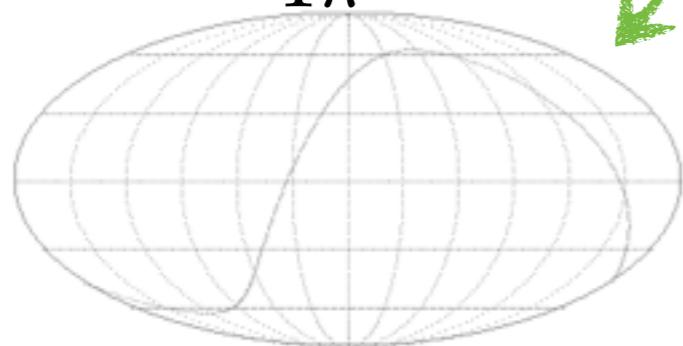
$$p^{\text{DM}} = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl}$$



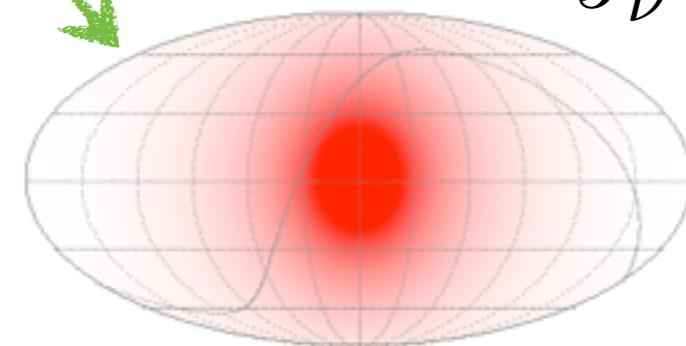
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$$p^{\text{DM}} = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl}$$



PDF of data

$$p_i(b, l) = \frac{1}{2\pi\sigma_i^2} \exp\left[-\frac{|\vec{x} - \vec{x}_i|^2}{2\sigma_i^2}\right]$$

"flat sky" approximation

PDF of isotropic dis.

$$p^{\text{iso}} = \frac{1}{4\pi}$$

$$p^{\text{DM}}(b, l) = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl} = \frac{\int_0^\infty \rho[r(s, b, l)] ds + \Omega_{\text{DM}} \rho_c \beta}{4\pi(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$

Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Test
Statistics

Number of signal events

$$TS_{\text{like}} = 2 \sum_{i=1}^N (\ln f_i - \ln p_i^{\text{iso}}) = 2 \ln \left(\prod_{i=1}^N f_i \right) - 2N \ln \left(\frac{1}{4\pi} \right)$$

$$f_i = \int p_i(b, l) p^{\text{DM}}(b, l) \cos(b) db dl = \frac{1}{2\pi\sigma_i^2} \int e^{-\frac{|\vec{x}_i - \vec{x}|^2}{2\sigma_i^2}} p^{\text{DM}}(b, l) \cos(b) db dl$$

N = 35 ? too optimistic!

Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

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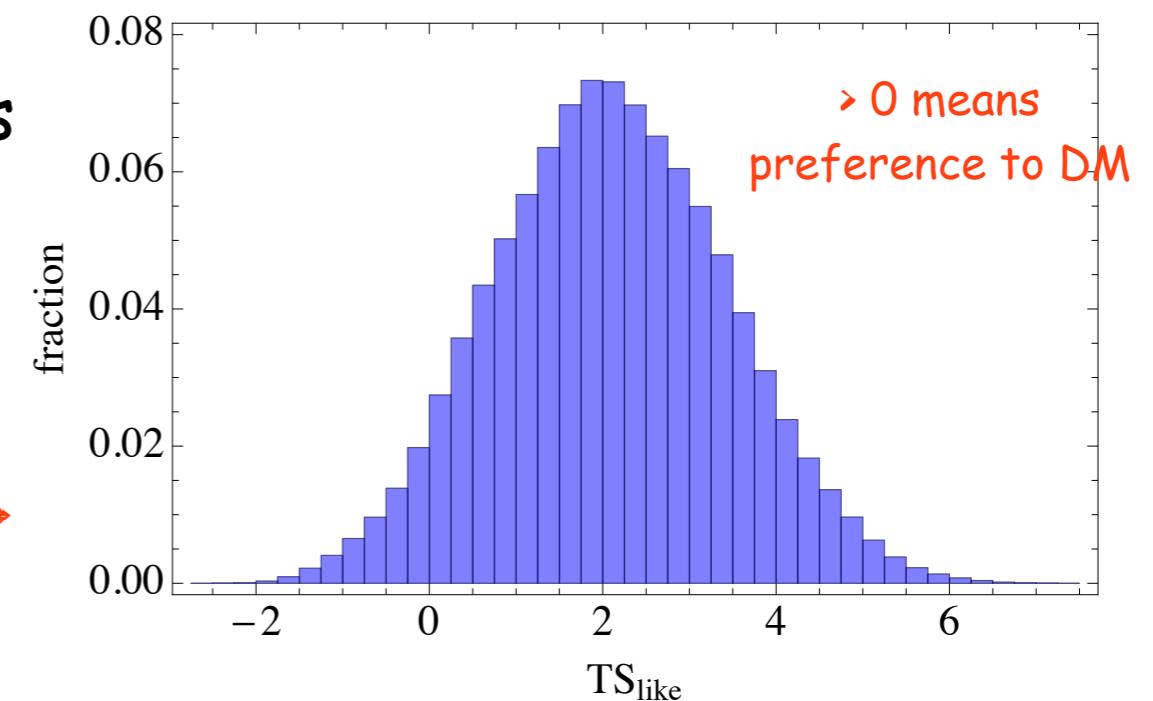
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$N = 35$? too optimistic!

✓ let's assume $N_b = 15$ and all the events with $E > 150$ TeV are signal events

→ $\binom{26}{15}$ ways of selecting the bkg events among the low energy events

Distribution of TS_{like} for all these realizations
(mean value = 2.1)



Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Test Statistics

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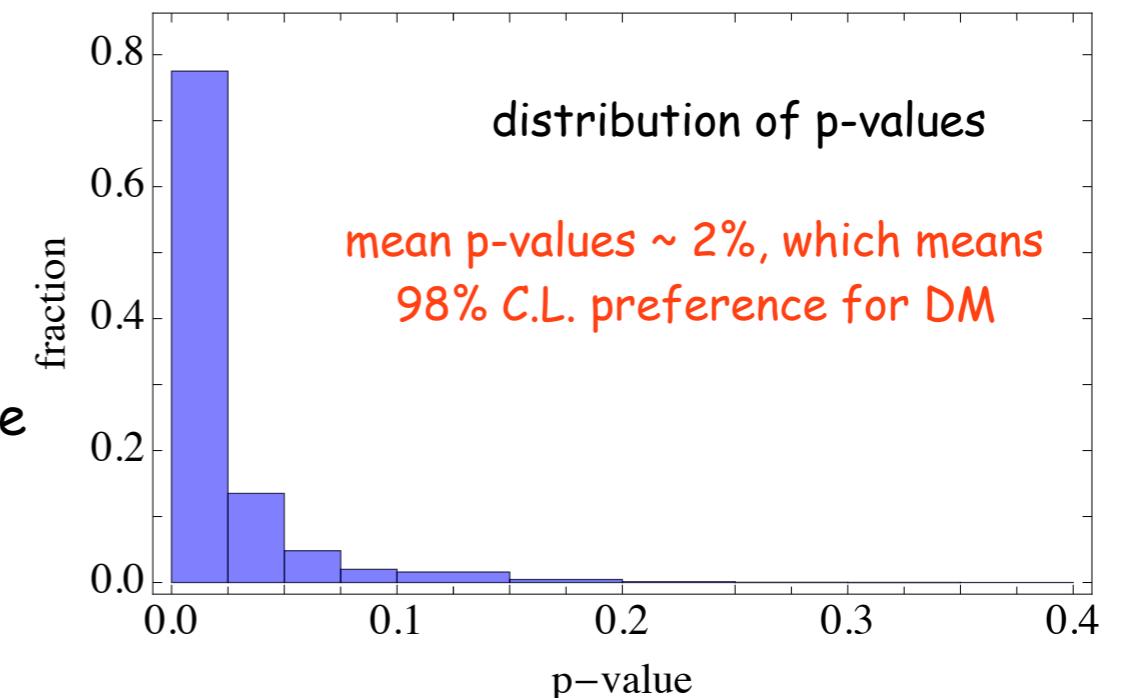
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Quantifying the preference

generating a sample (10^5) of isotropically distributed set of 20 events

p-value

→ for each realization of bkg choosing, p-value is the fraction of generated events which have smaller TS_{like} than the one computed by observed data



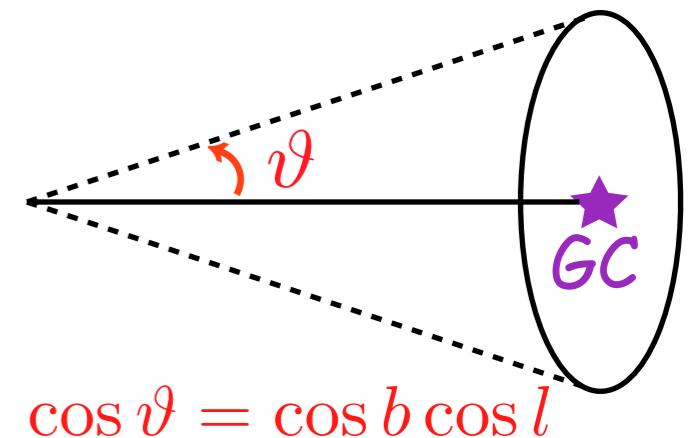
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$

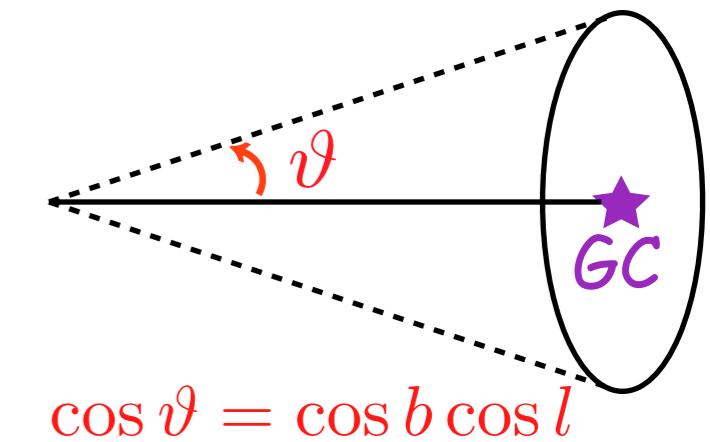


Angular distribution of neutrinos from decaying DM

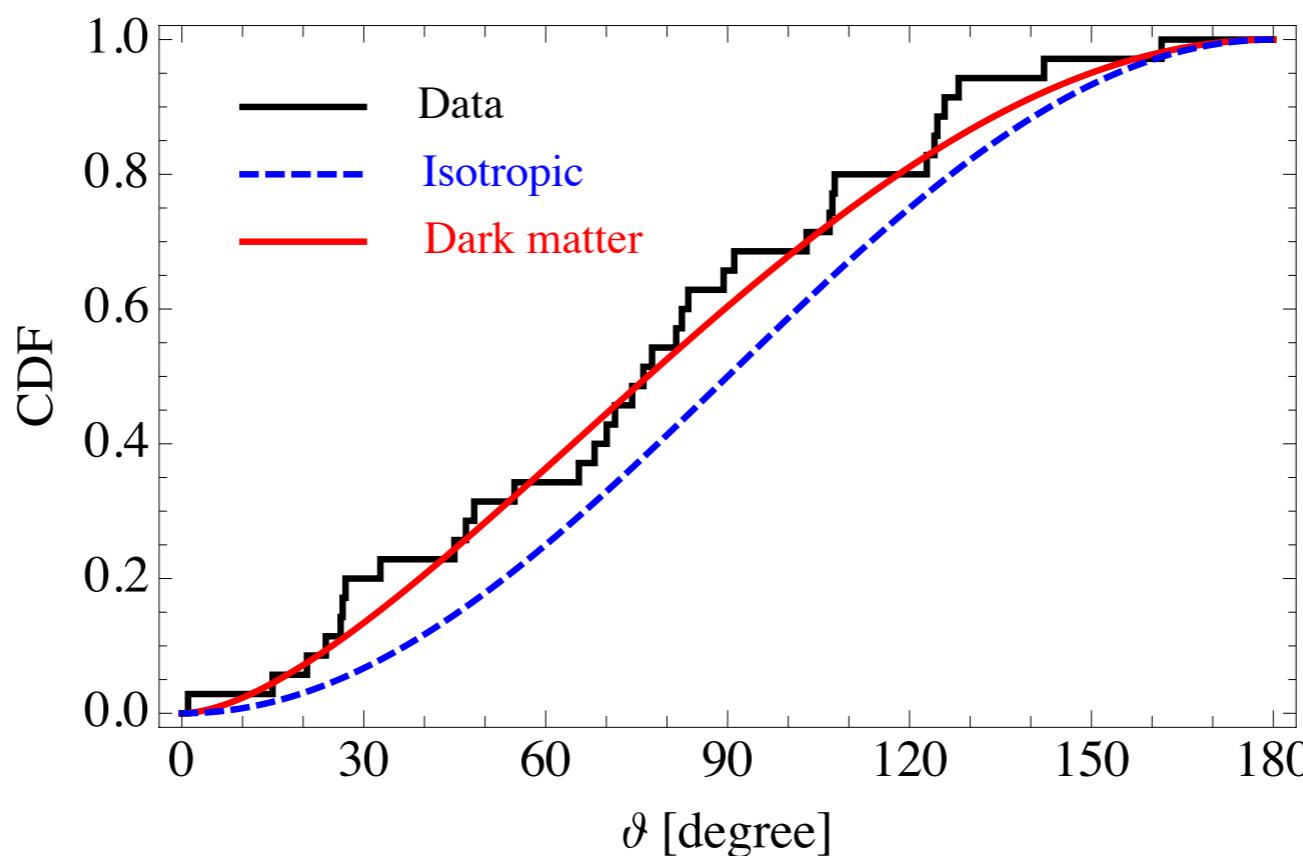
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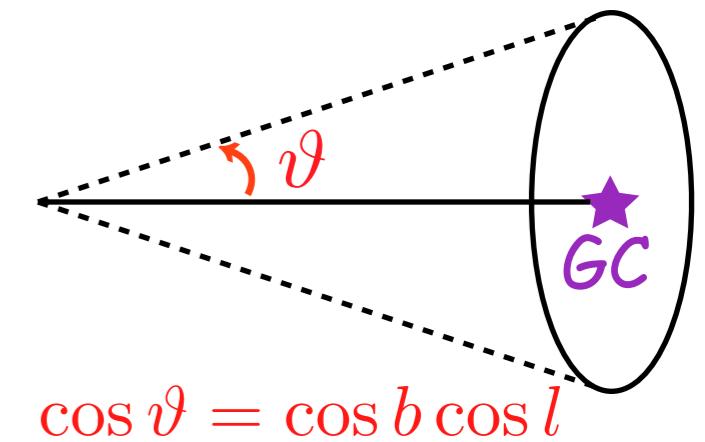


Angular distribution of neutrinos from decaying DM

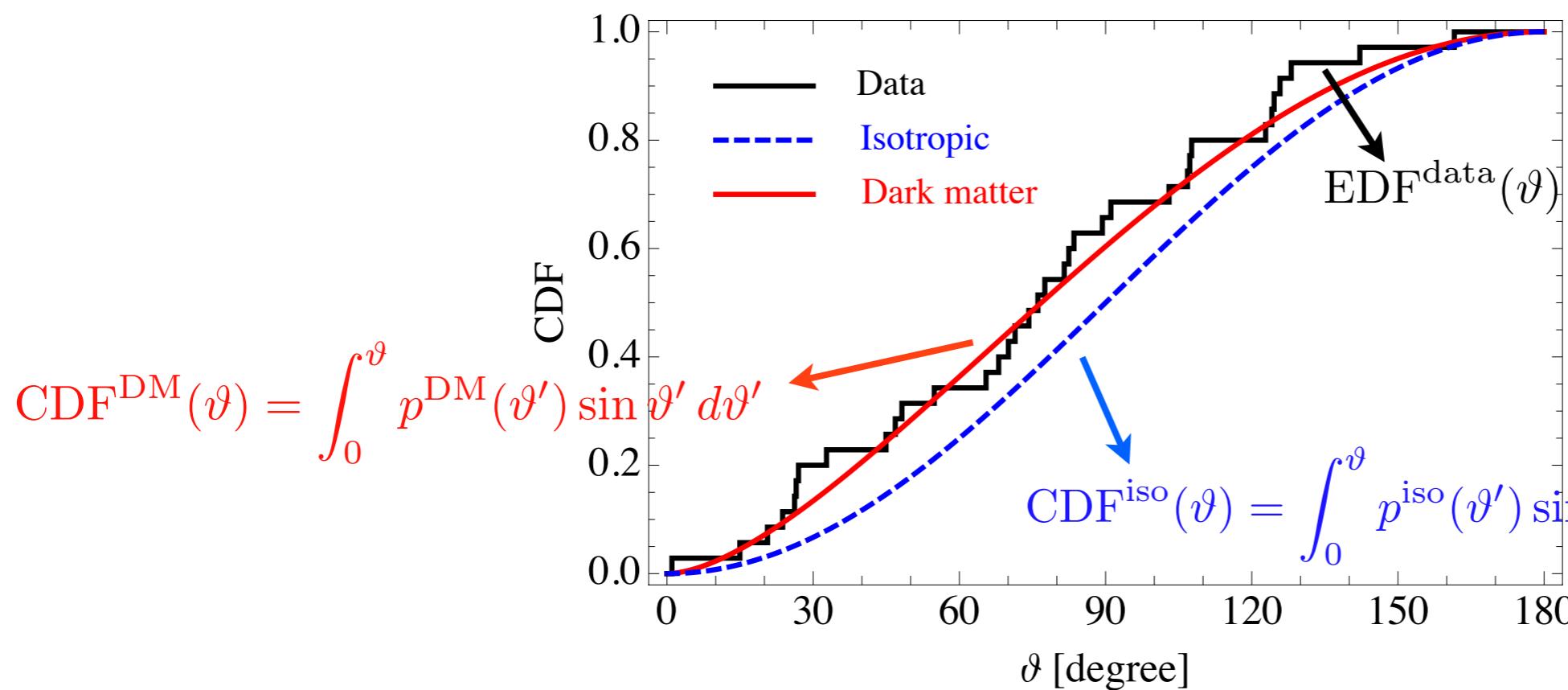
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$$\text{EDF}^{\text{data}}(\vartheta) = \frac{1}{N} \sum_{i=1}^N \Theta(\vartheta - \vartheta_i)$$

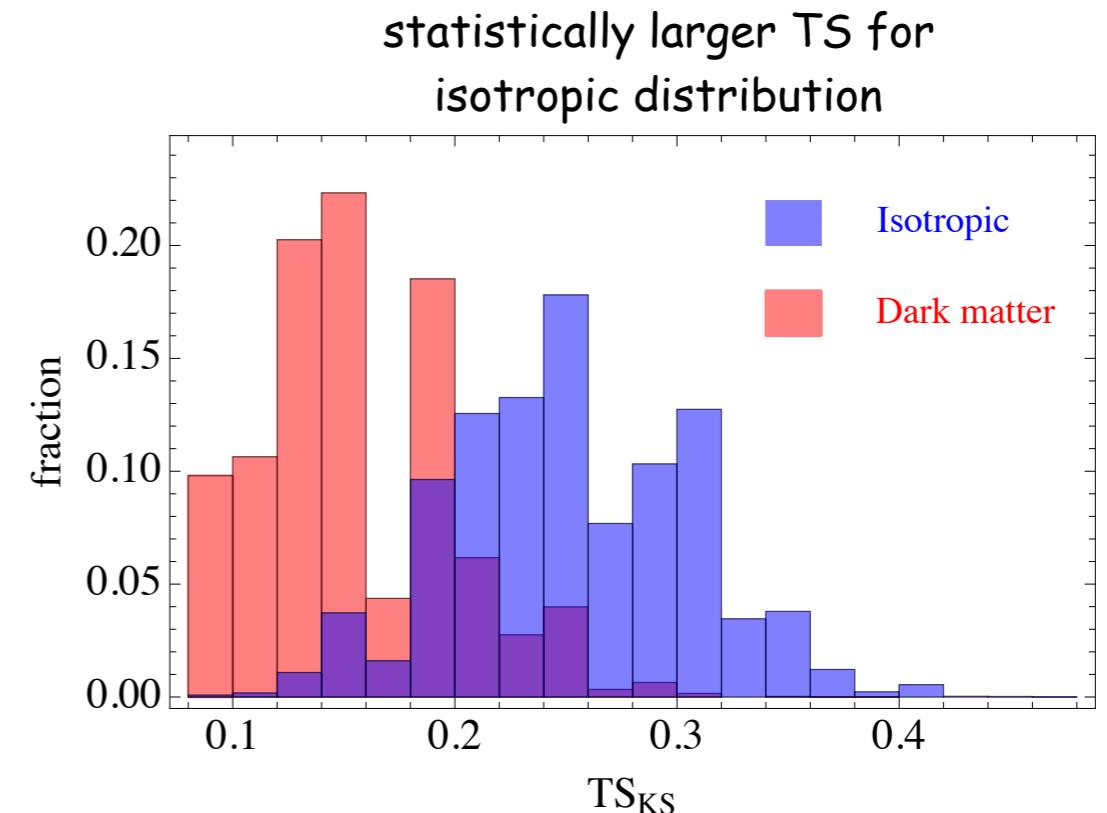
$$\text{CDF}^{\text{iso}}(\vartheta) = \int_0^\vartheta p^{\text{iso}}(\vartheta') \sin \vartheta' d\vartheta' = \frac{1 - \cos \vartheta}{2}$$

Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test:

Test Statistics

$$TS_{KS} = \max_{1 \leq i \leq N} \left\{ CDF^{DM}(\vartheta_i) - \frac{i-1}{N}, \frac{i}{N} - CDF^{DM}(\vartheta_i) \right\}$$

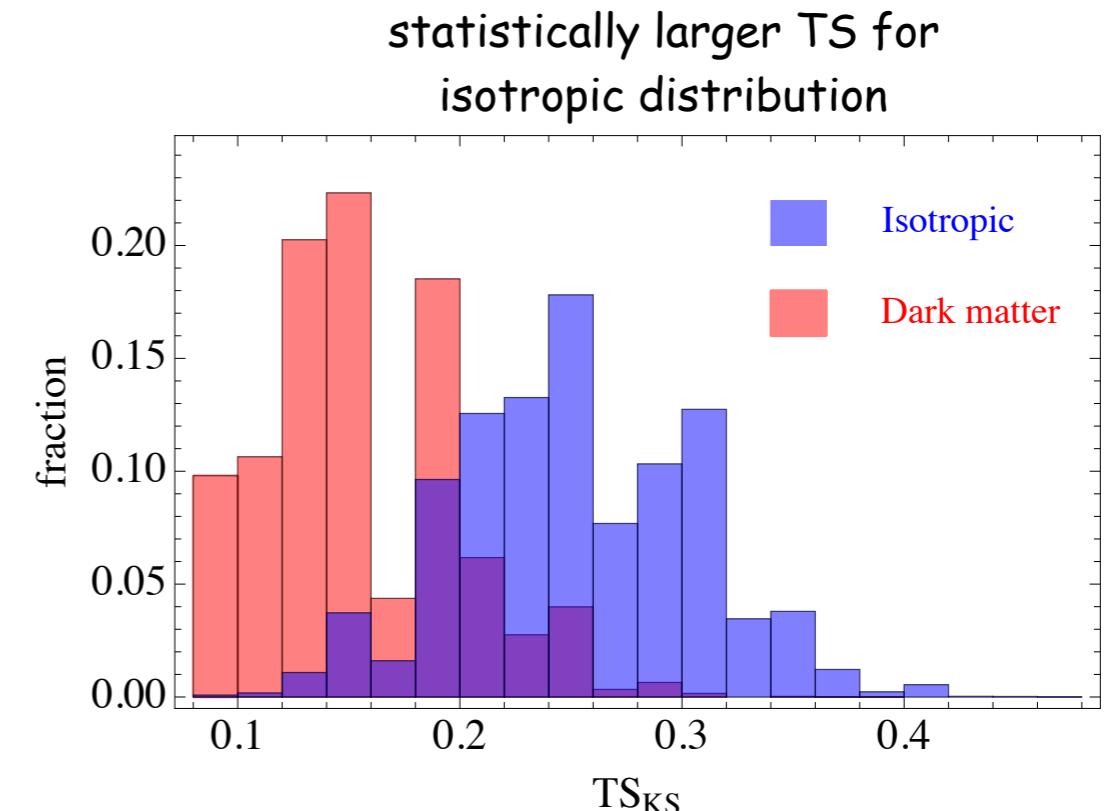


Angular distribution of neutrinos from decaying DM

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again, generating a sample (10^5) of isotropically distributed set of 20 events

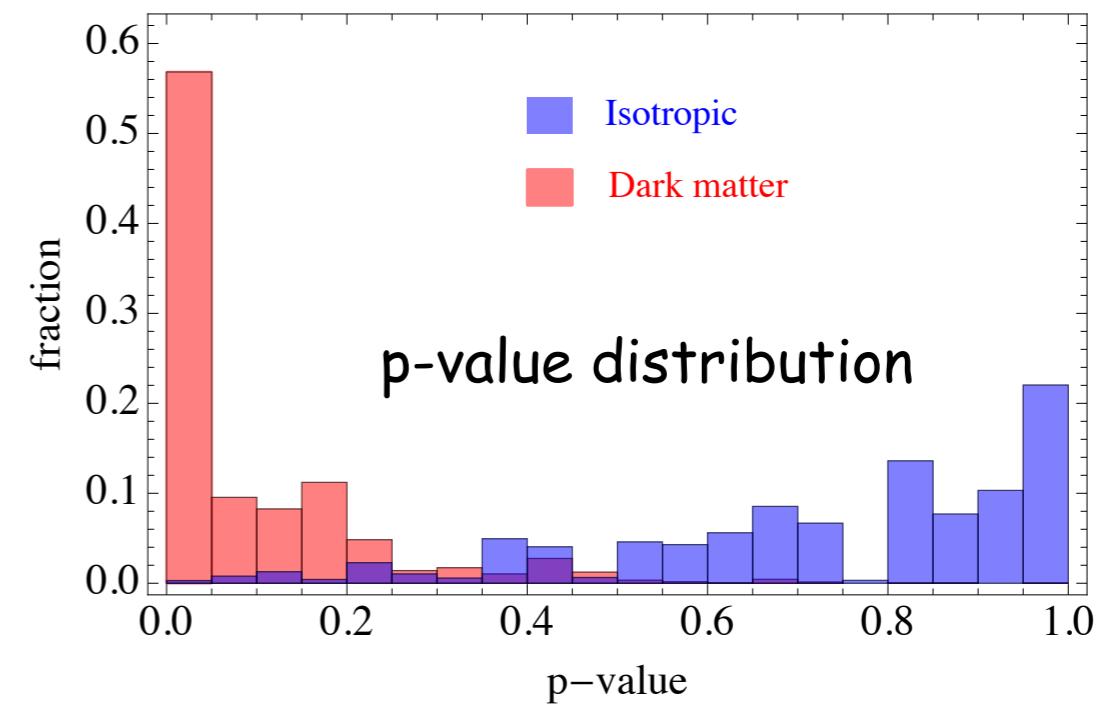


on the average, 10% of generated isotropic sample have smaller TS_{KS} than the values obtained for data vs DM dis.

for data vs isotropic dis. it is 73%



less than 2σ preference for DM dis.



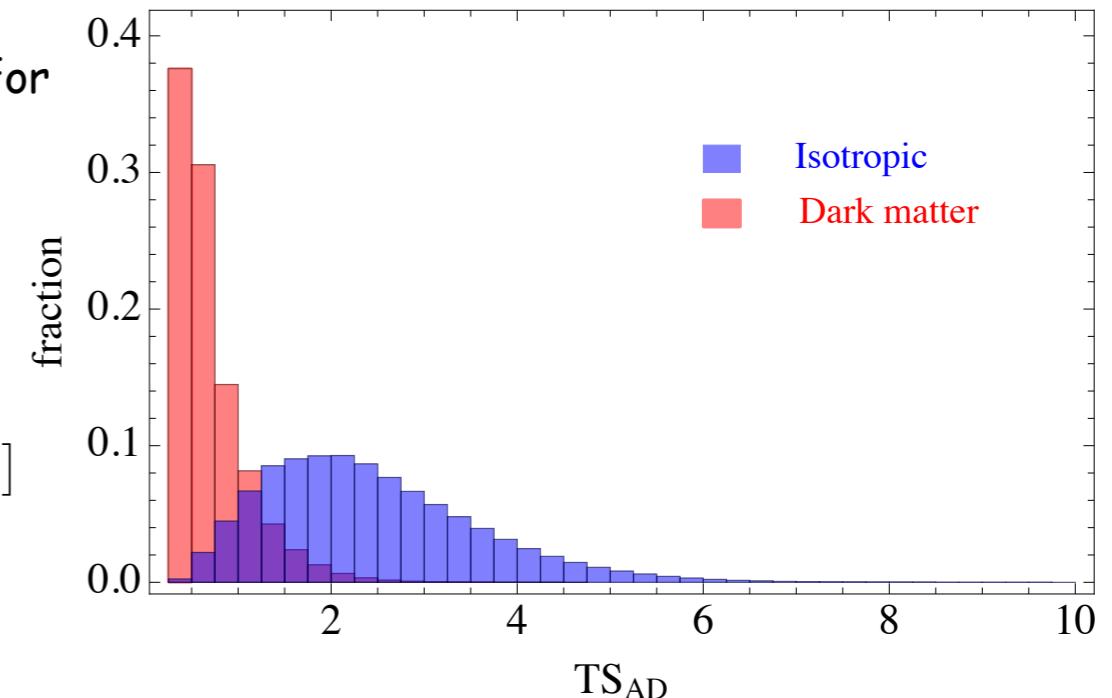
Angular distribution of neutrinos from decaying DM

✓ Anderson-Darling test: a powerful non-parametric test, especially sensitive to the end points

Test Statistics

$$TS_{AD} = -N - \frac{1}{N} \sum_{i=1}^N (2i - 1) [\ln(CDF^{DM}(\vartheta_i)) + \ln(1 - CDF^{DM}(\vartheta_{N+1-i}))]$$

statistically larger TS for isotropic distribution



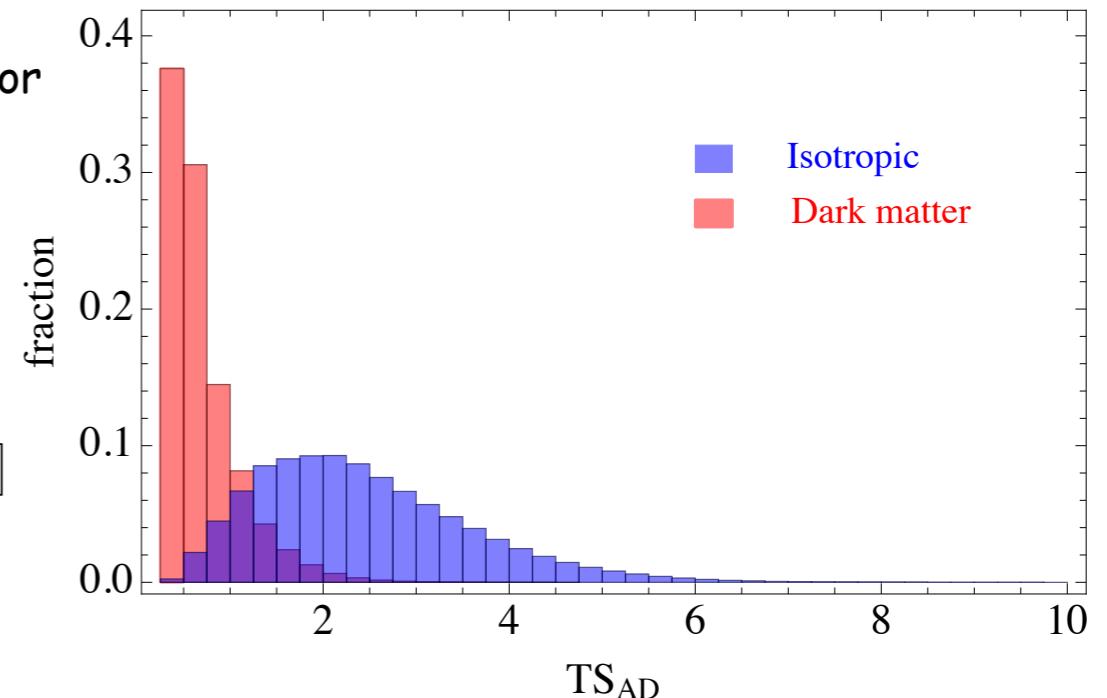
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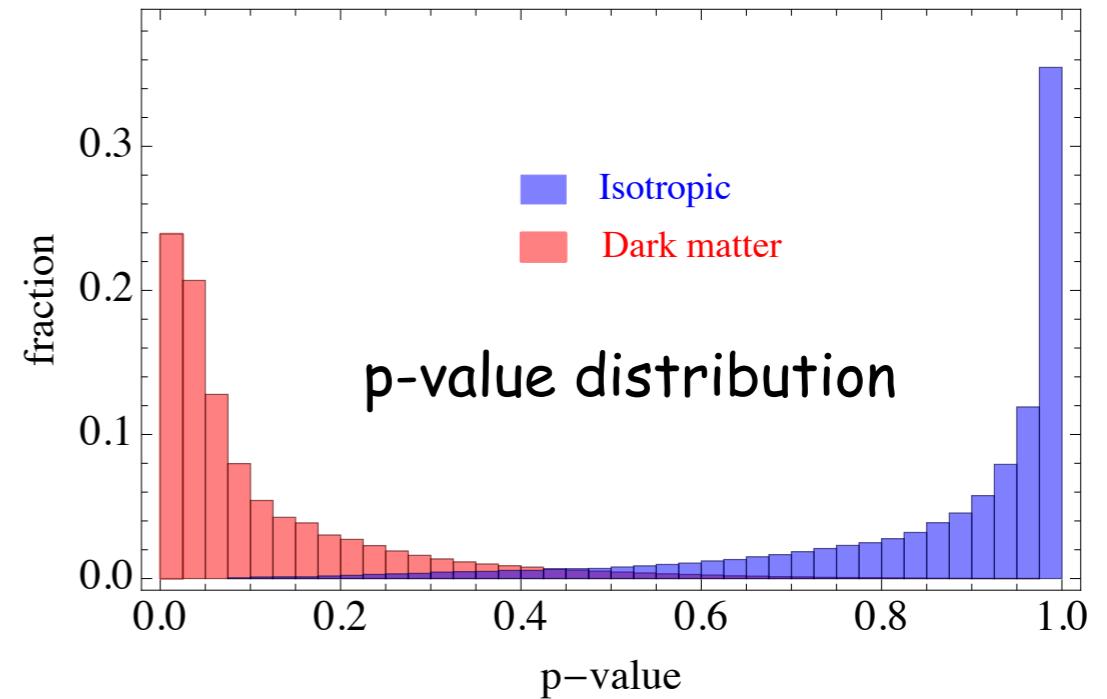
again, generating a sample (10^5) of isotropically distributed set of 20 events



on the average, 11% of generated isotropic sample have smaller TS_{KS} than the values obtained for data vs DM dis.
for data vs isotropic dis. it is 86%



less than 2σ preference for DM dis.



Gamma ray bounds

Universe is opaque for
gamma-rays with $E > 1 \text{ TeV}$



cascades develop: gamma-ray
interaction with interstellar
radiation field and CMB



gamma-rays populate at
lower energies $< 10^{(2-3)} \text{ GeV}$

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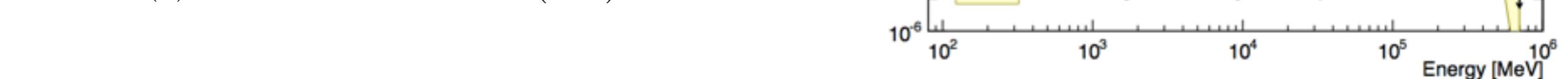
✓ Isotropic diffuse gamma-ray background by Fermi-LAT

integrated energy density

$$\omega_\gamma = \frac{4\pi}{c} \int_{E_1}^{E_2} E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} dE_\gamma \lesssim 4.4 \times 10^{-7} \text{ eV/cm}^3$$

$$E_1 \sim \mathcal{O}(1) \text{ GeV}$$

$$E_2 \sim \mathcal{O}(100) \text{ GeV}$$



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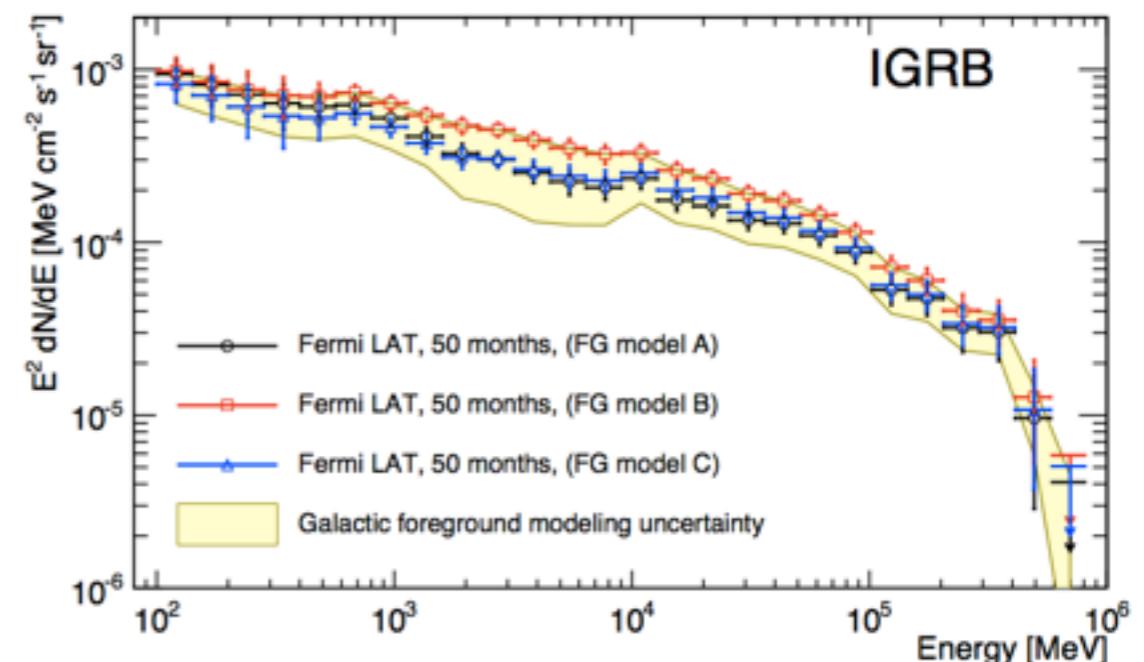
M. Ackermann et al. [The Fermi LAT Collaboration], arXiv:1410.3696 [astro-ph.HE].

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total electromagnetic energy budget
(NH case)

$$\frac{4\pi}{c} \int \sum_{i=\text{gal, extragal}} \left[E_\gamma \left(\frac{d\varphi_\gamma}{dE_\gamma} \right)^i + E_e \left(\frac{d\varphi_{e^\pm}}{dE_e} \right)^i \right] dE \simeq 5.2 \times 10^{-8} \text{ eV/cm}^3$$



Gamma ray bounds

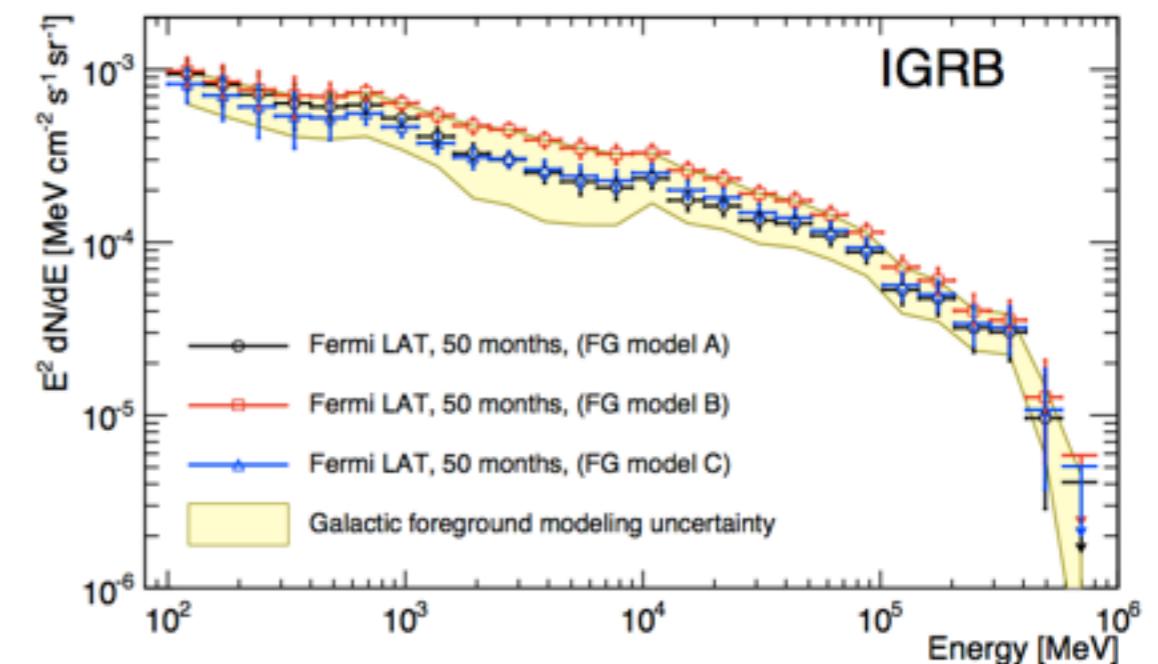
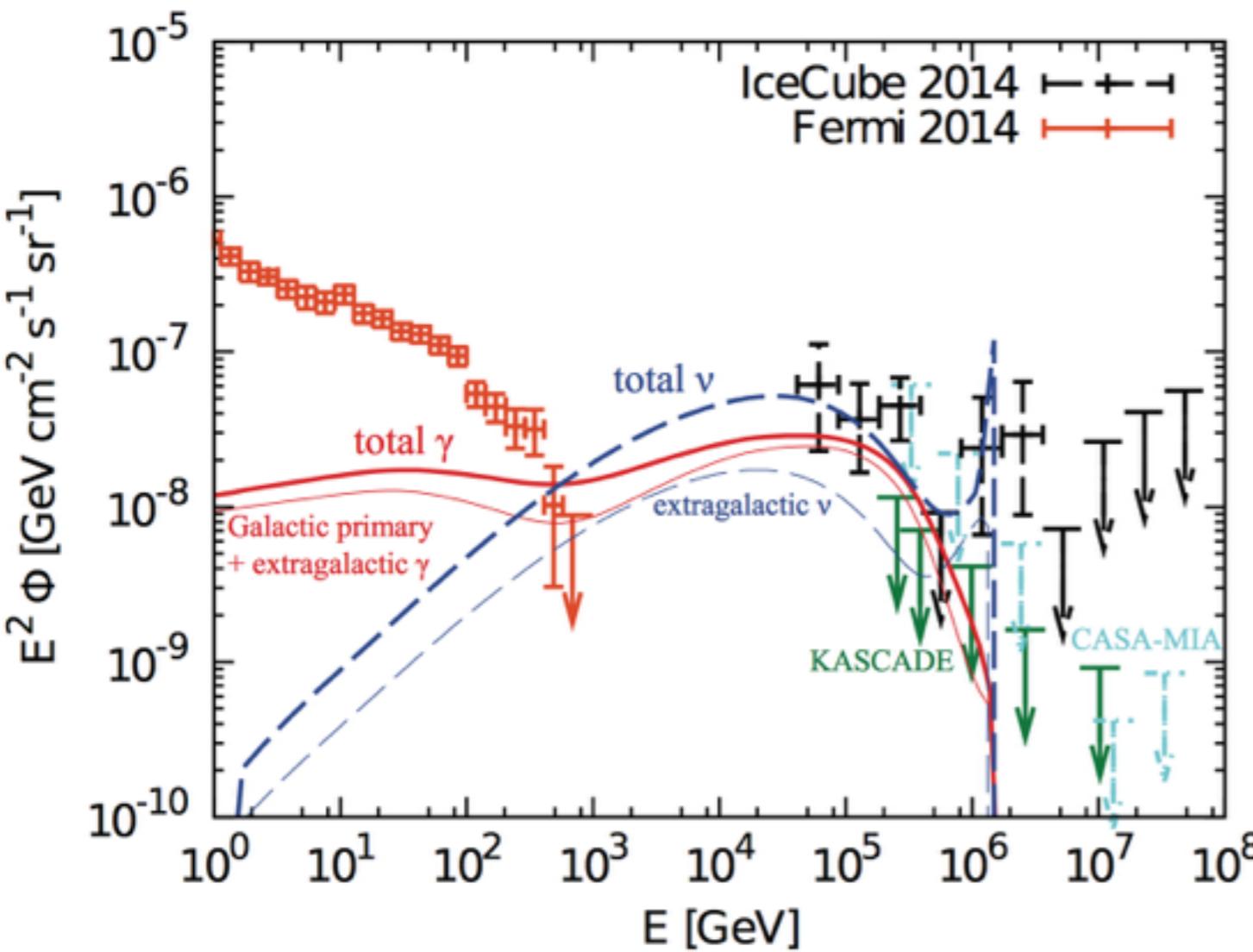
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Murase, Laha, Ando, Ahlers,
arXiv:1503.04663

Gamma ray bounds

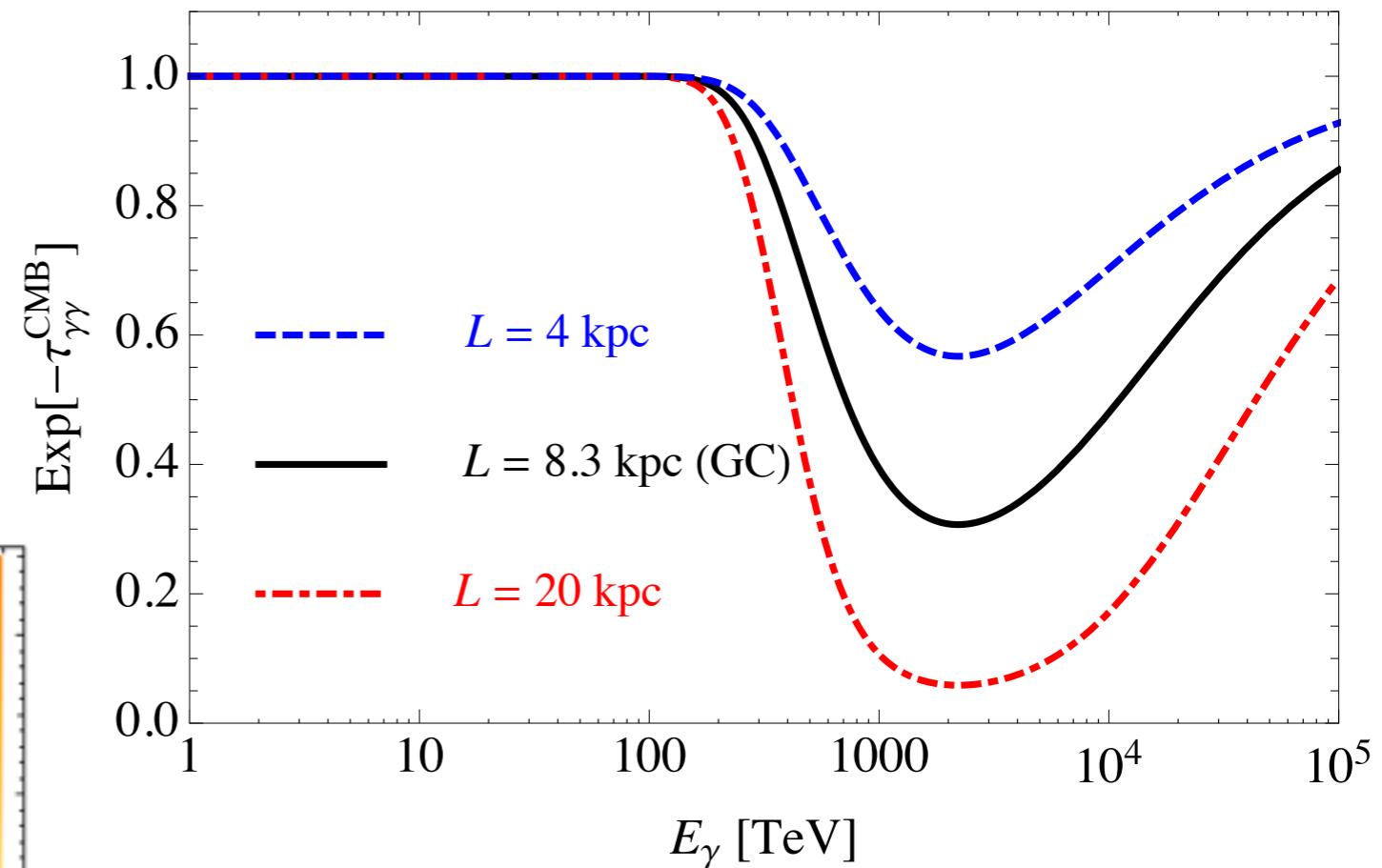
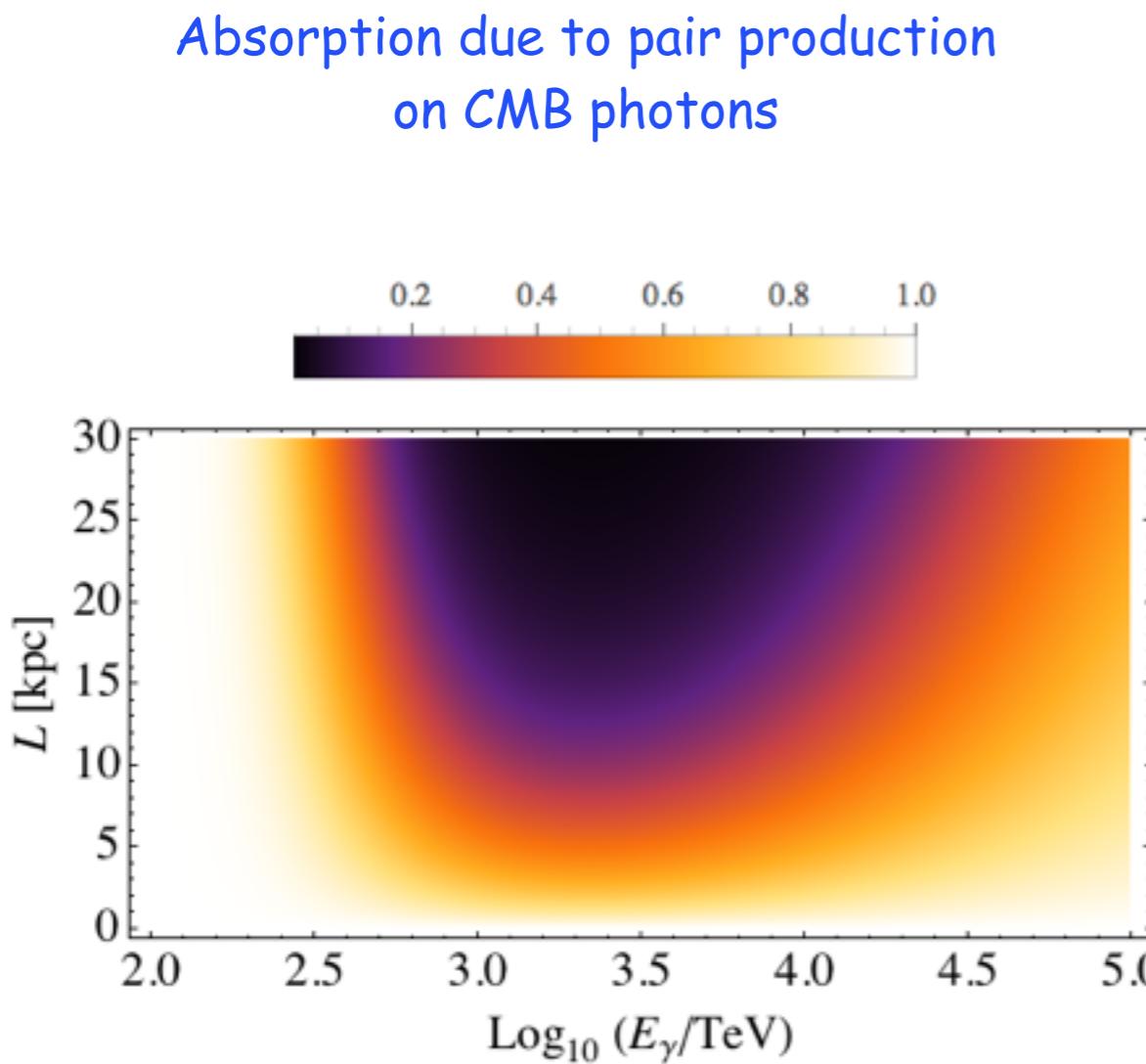
✓ Galactic component

at \sim PeV, the absorption length of gamma-rays
are comparable to Galactic distances



neither full absorption or cascade
development, nor full transparency

A. E. and P. Serpico, arXiv:1505.06486



Absorption at \sim PeV

Gamma ray bounds

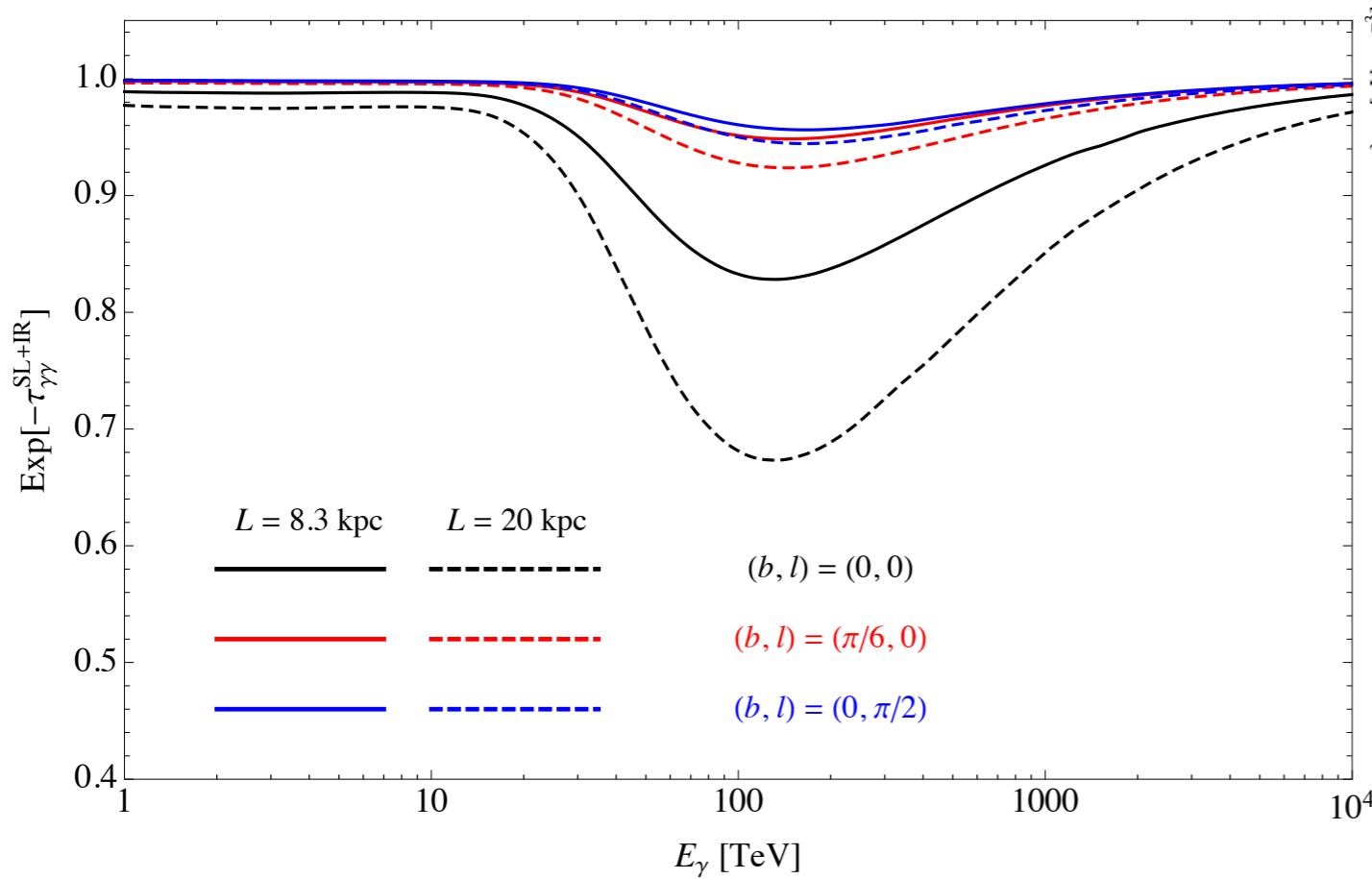
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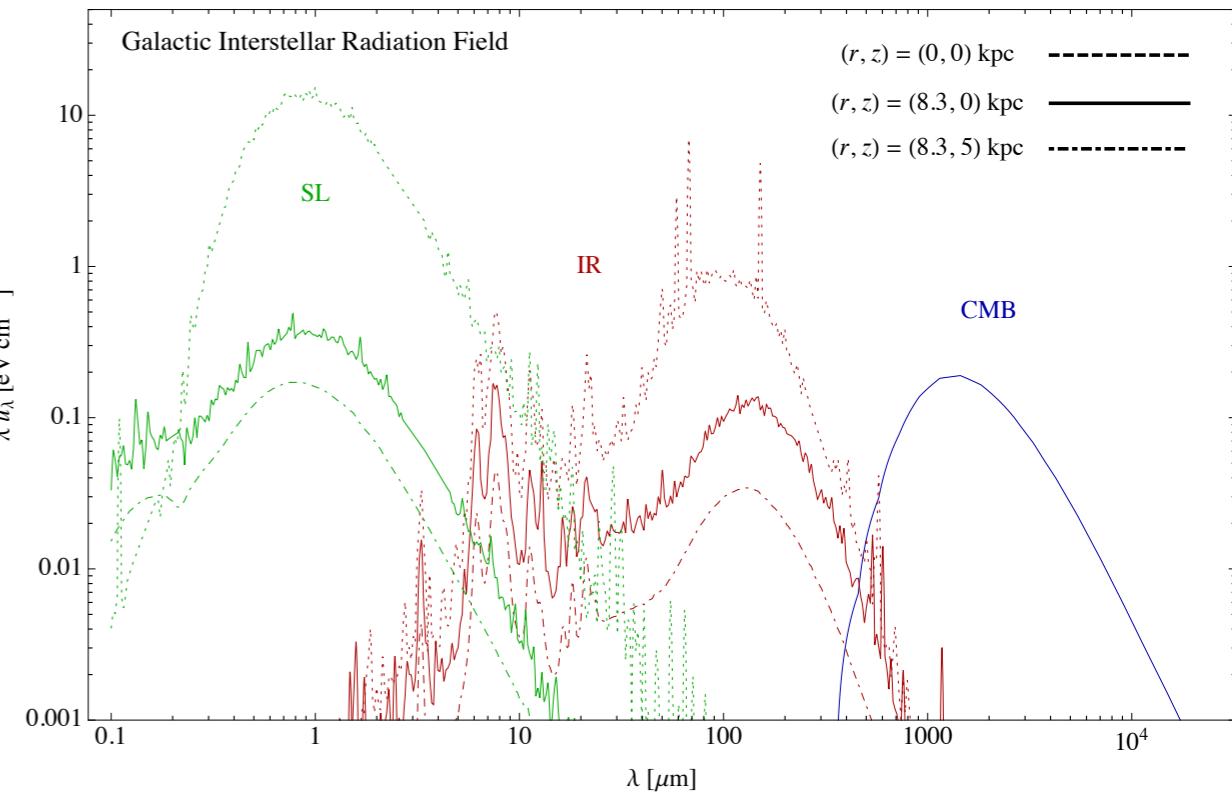


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Absorption due to pair production
on SL+IR photons



A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486



Absorption at $\sim 100 \text{ TeV}$

Gamma ray bounds

✓ Galactic component

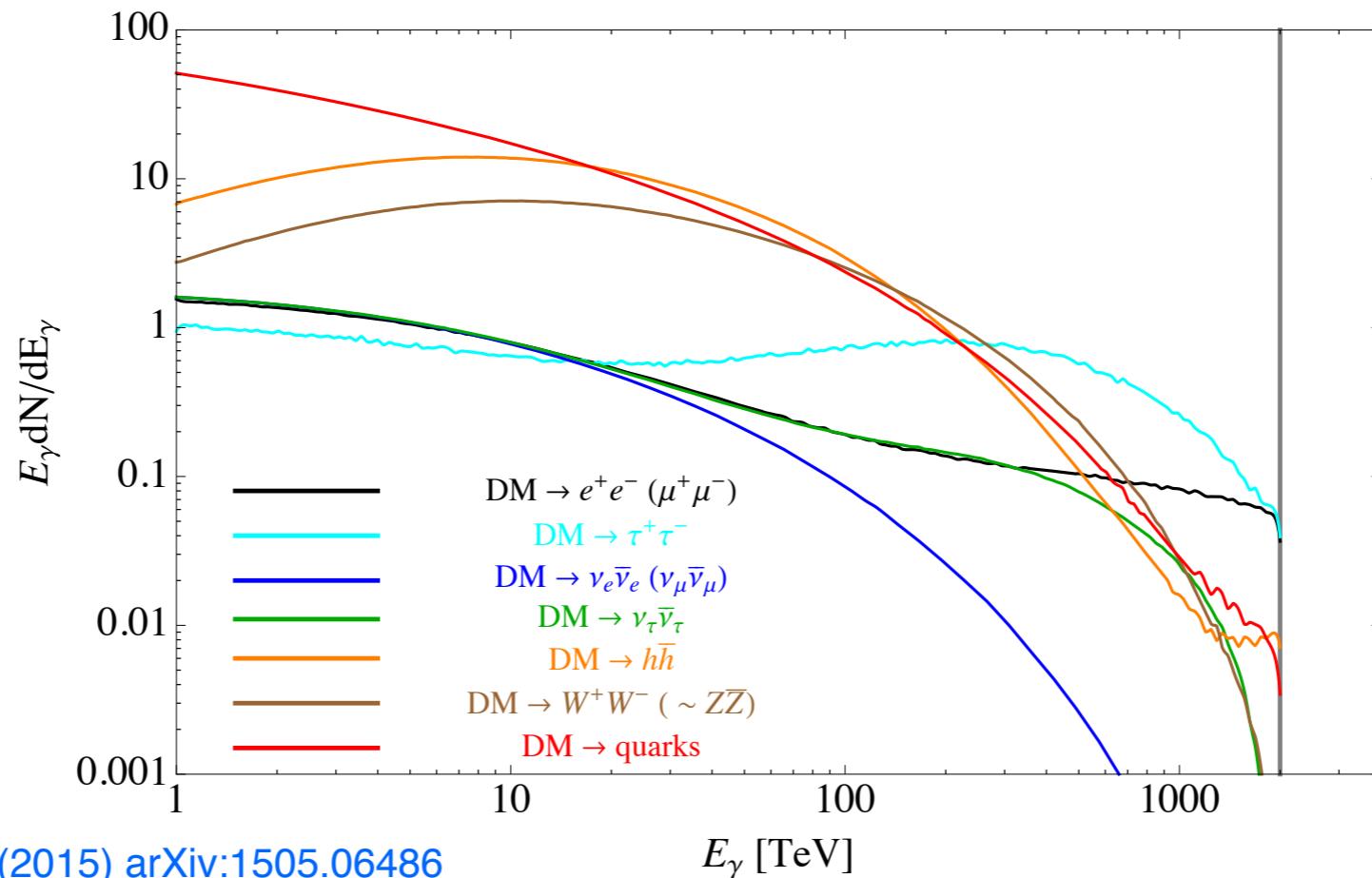
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Prompt component

$$\frac{d\Phi_\gamma}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\gamma}{dE_\gamma}(E_\gamma) \int_0^\infty \rho_h[\varrho(s, b, l)] e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} ds$$



Gamma ray bounds

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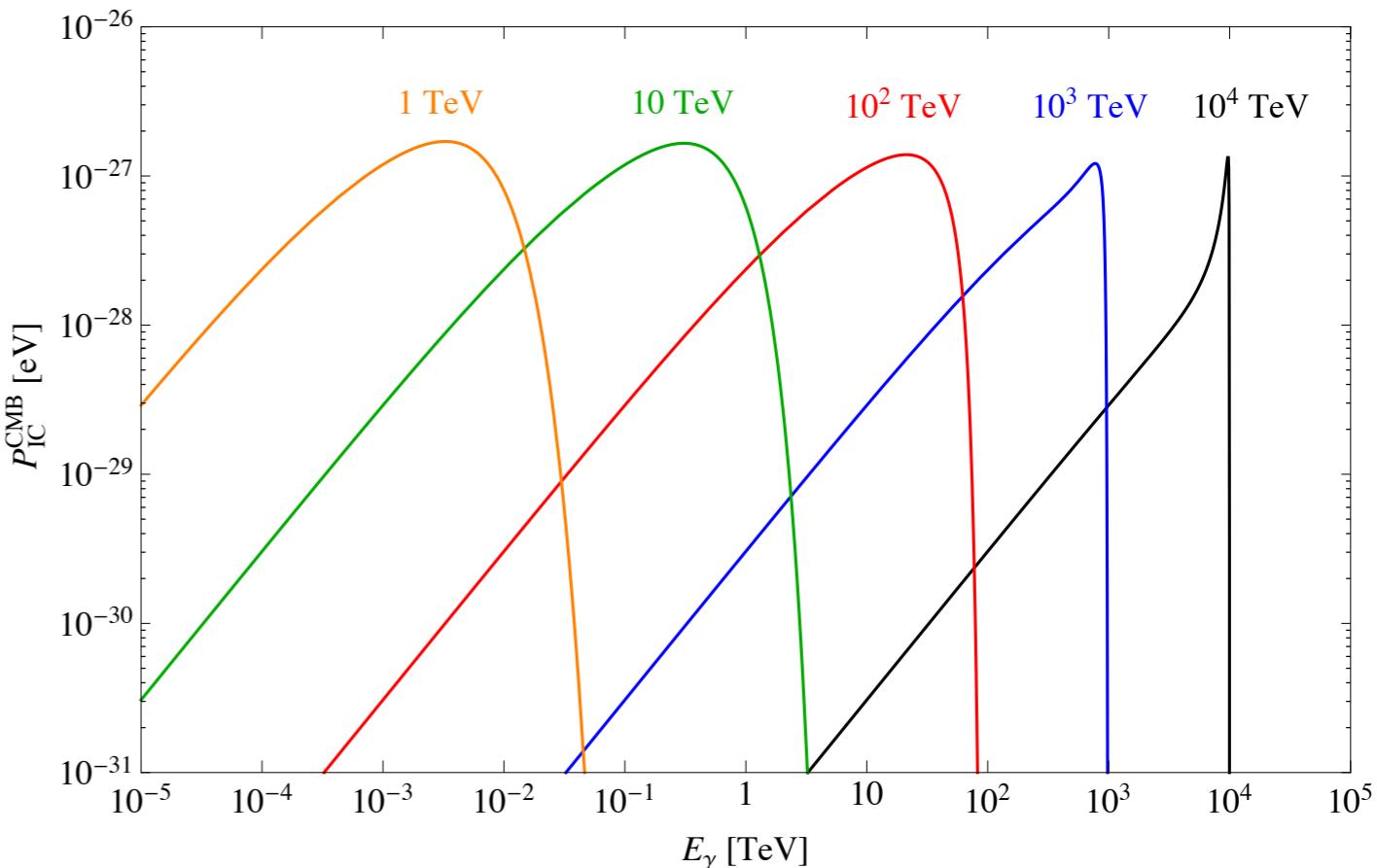


neither full absorption or cascade
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inverse-Compton
component

$$\frac{d\Phi_{\text{IC}}}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi E_\gamma} \int_0^\infty ds e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} \int_{m_e}^{m_{\text{DM}}/2} dE_e \frac{dn_e}{dE_e}(E_e, \varrho) P_{\text{IC}}(E_e, E_\gamma, \varrho)$$

$$P_{\text{IC}} = P_{\text{IC}}^{\text{CMB}} + P_{\text{IC}}^{\text{SL+IR}}$$



A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486

Gamma ray bounds

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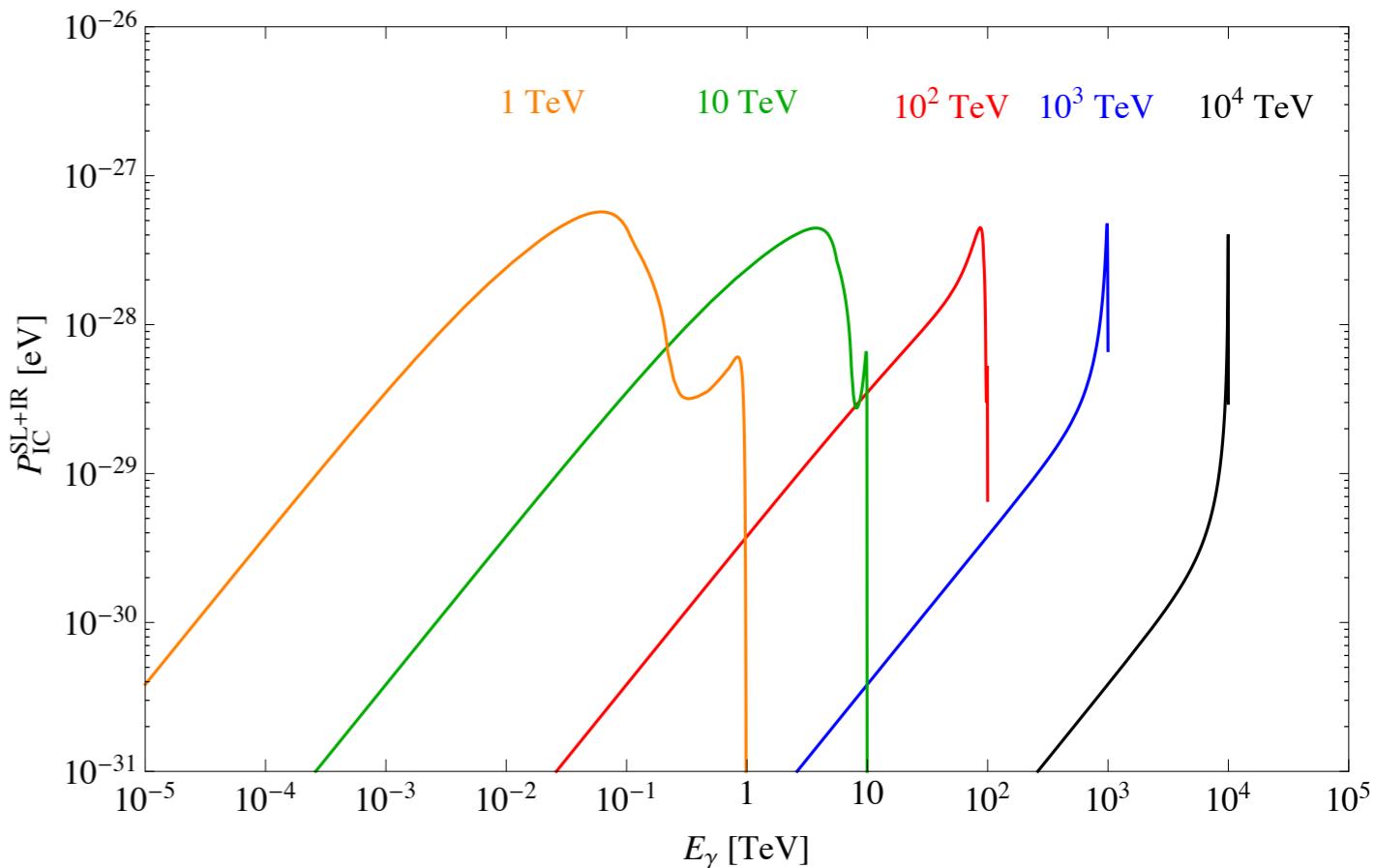


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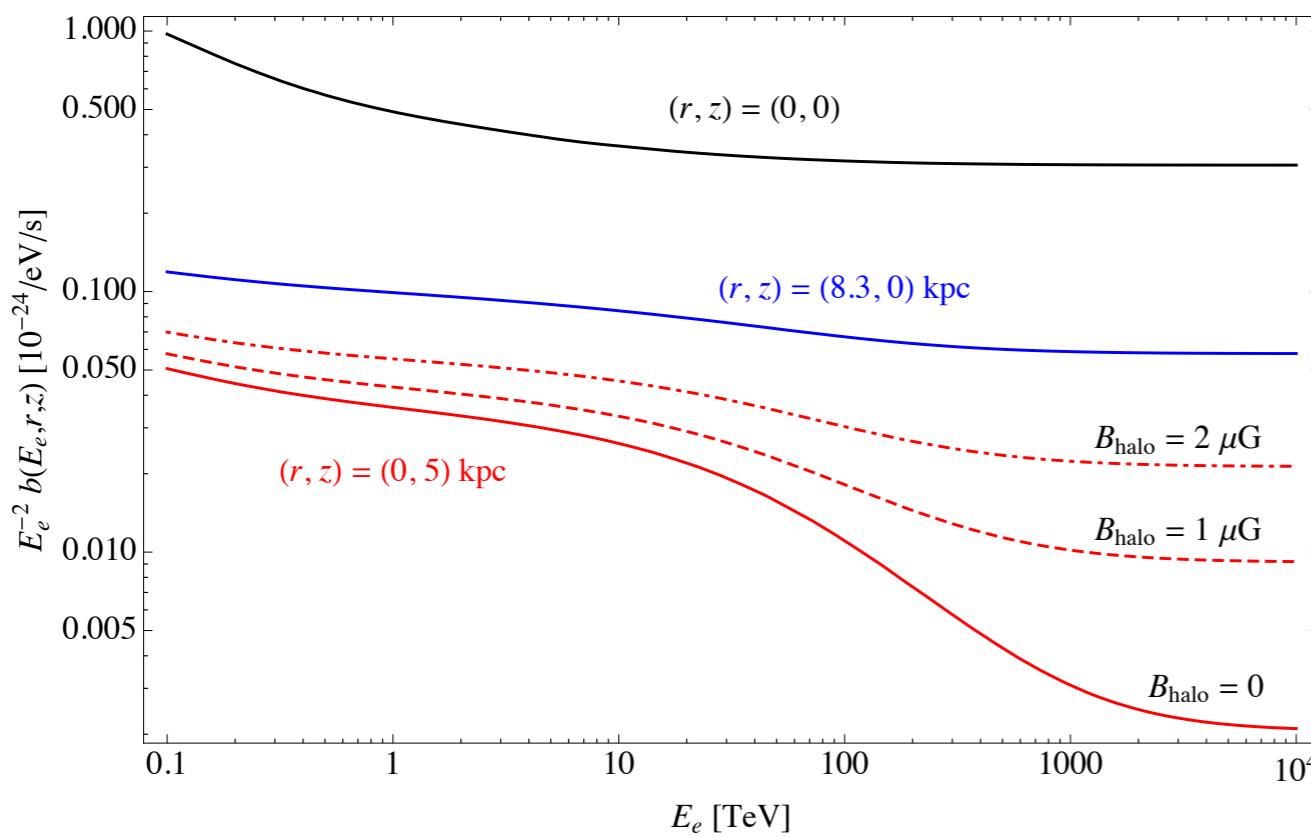


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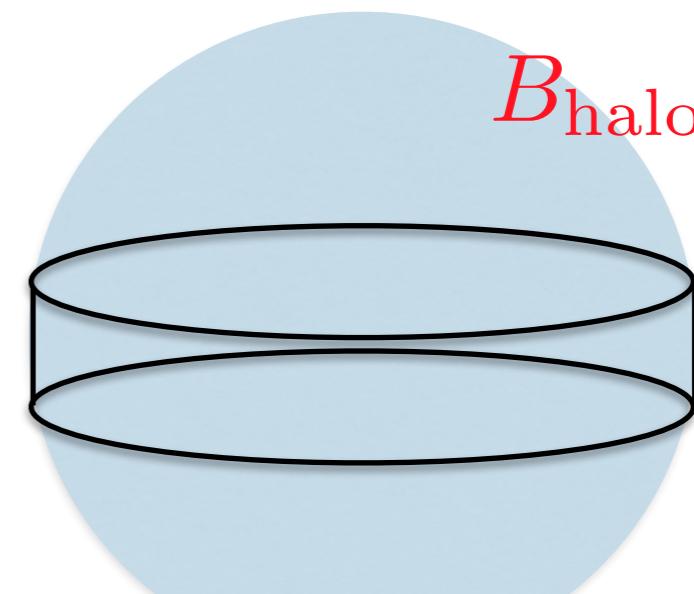
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$$\frac{dn_e}{dE_e}(E_e, \vec{x}) = \frac{1}{m_{\text{DM}} \tau_{\text{DM}}} \frac{\rho_h(\vec{x})}{b(E_e, \vec{x})} \int_{E_e}^{m_{\text{DM}}/2} \frac{dN_e}{dE'_e}(E'_e) I_{\text{diff}}(E_e, E'_e, \vec{x}) dE'_e$$



$$b(E_e, \vec{x}) \equiv -\frac{dE_e}{dt} = b_{\text{IC}}(E_e, \vec{x}) + b_{\text{syn}}(E_e, \vec{x})$$

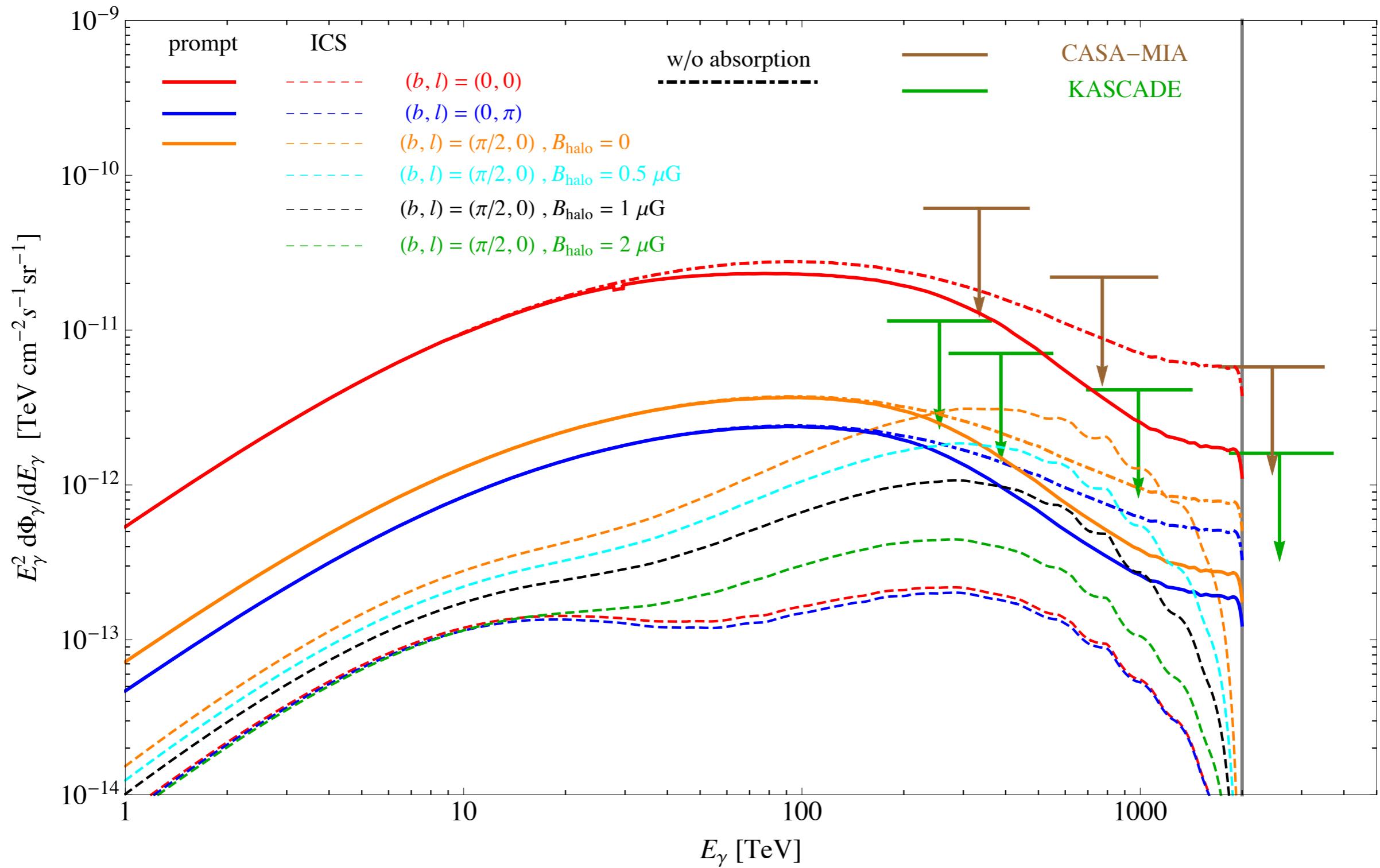


$$B_{\text{reg}}(\vec{x}) = B_0 \exp \left[-\frac{|r - R_\odot|}{r_B} - \frac{|z|}{z_B} \right]$$

Gamma ray bounds

✓ Galactic component

$$\tau_{\text{DM}} = 10^{28} \text{ s} \quad \text{and} \quad m_{\text{DM}} = 4 \text{ PeV}$$

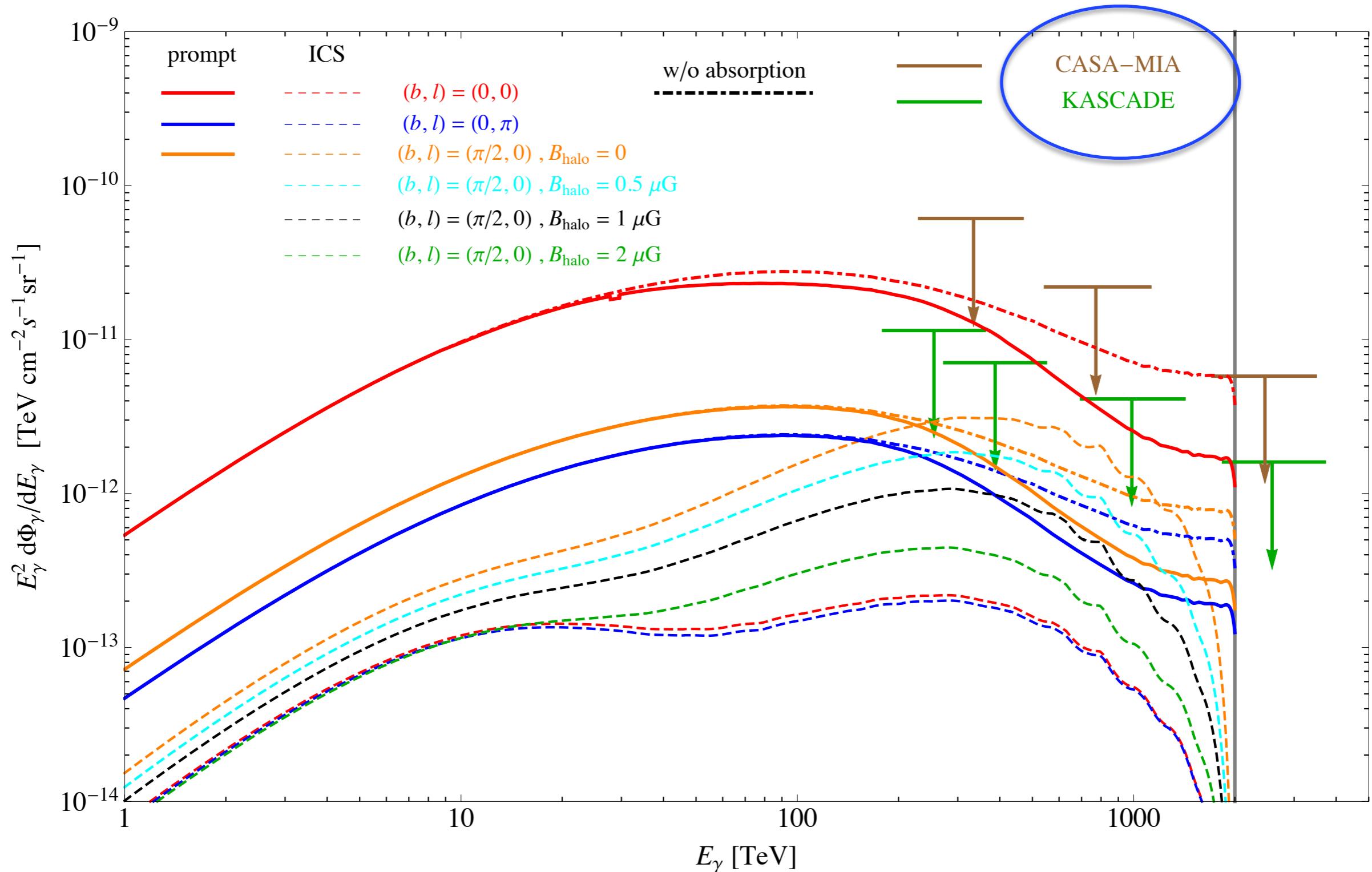


A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486

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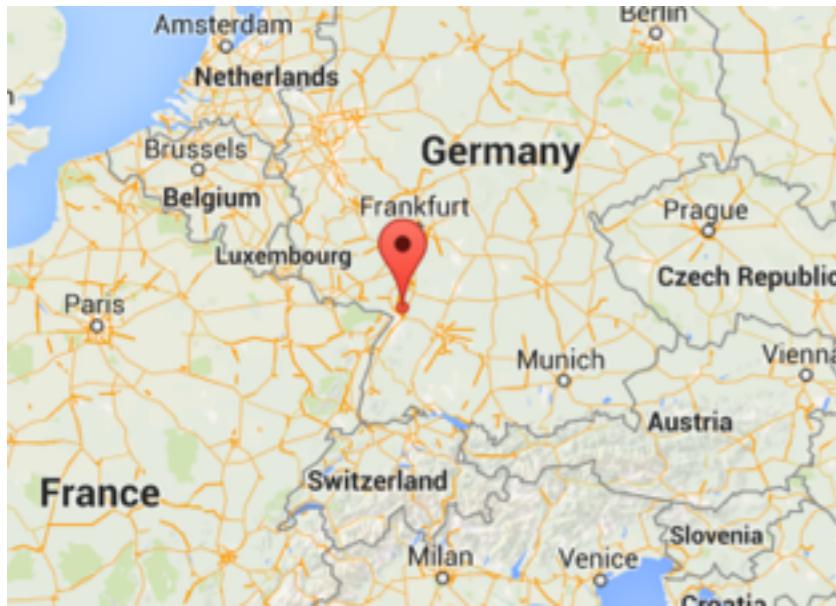
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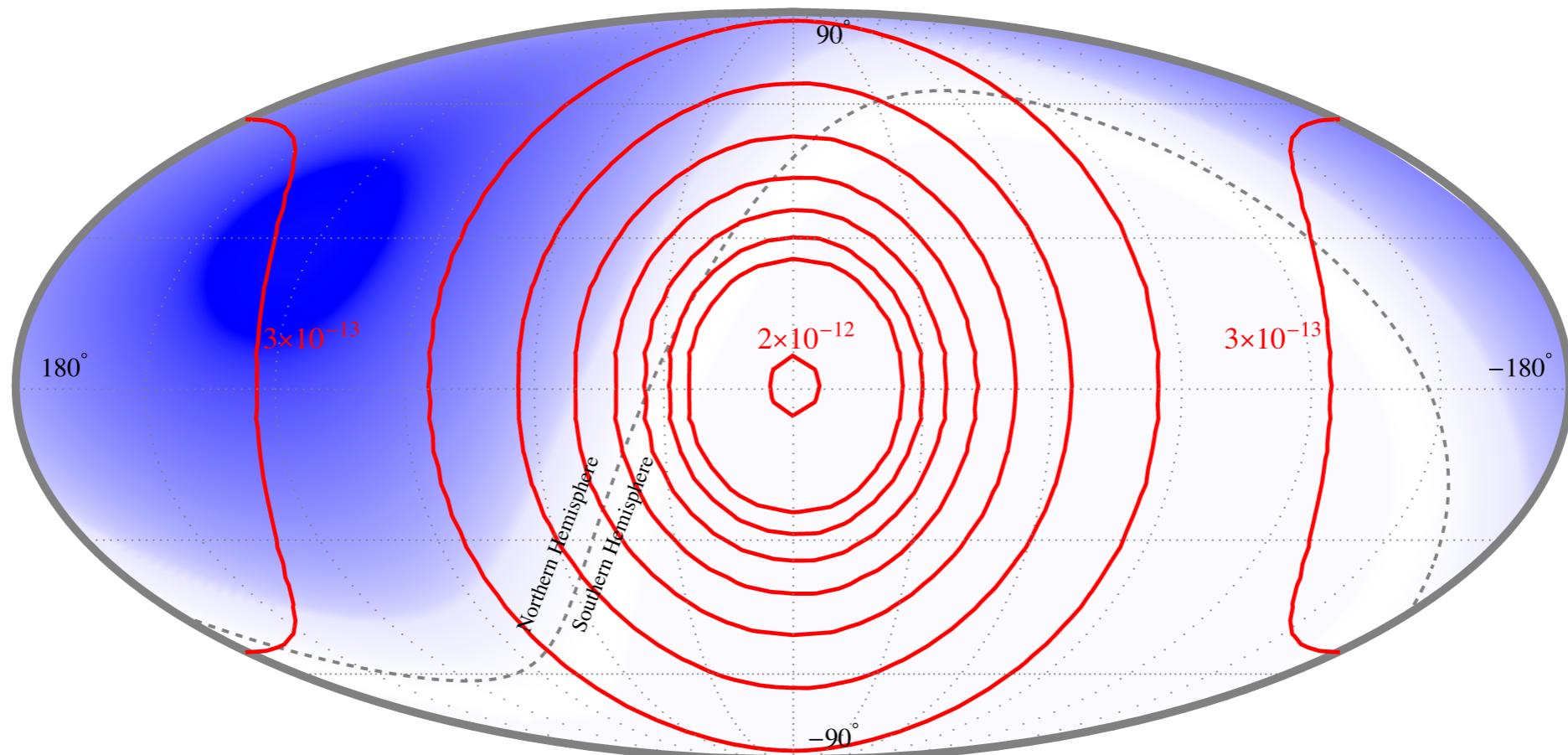
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Gamma ray bounds

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KASCADE

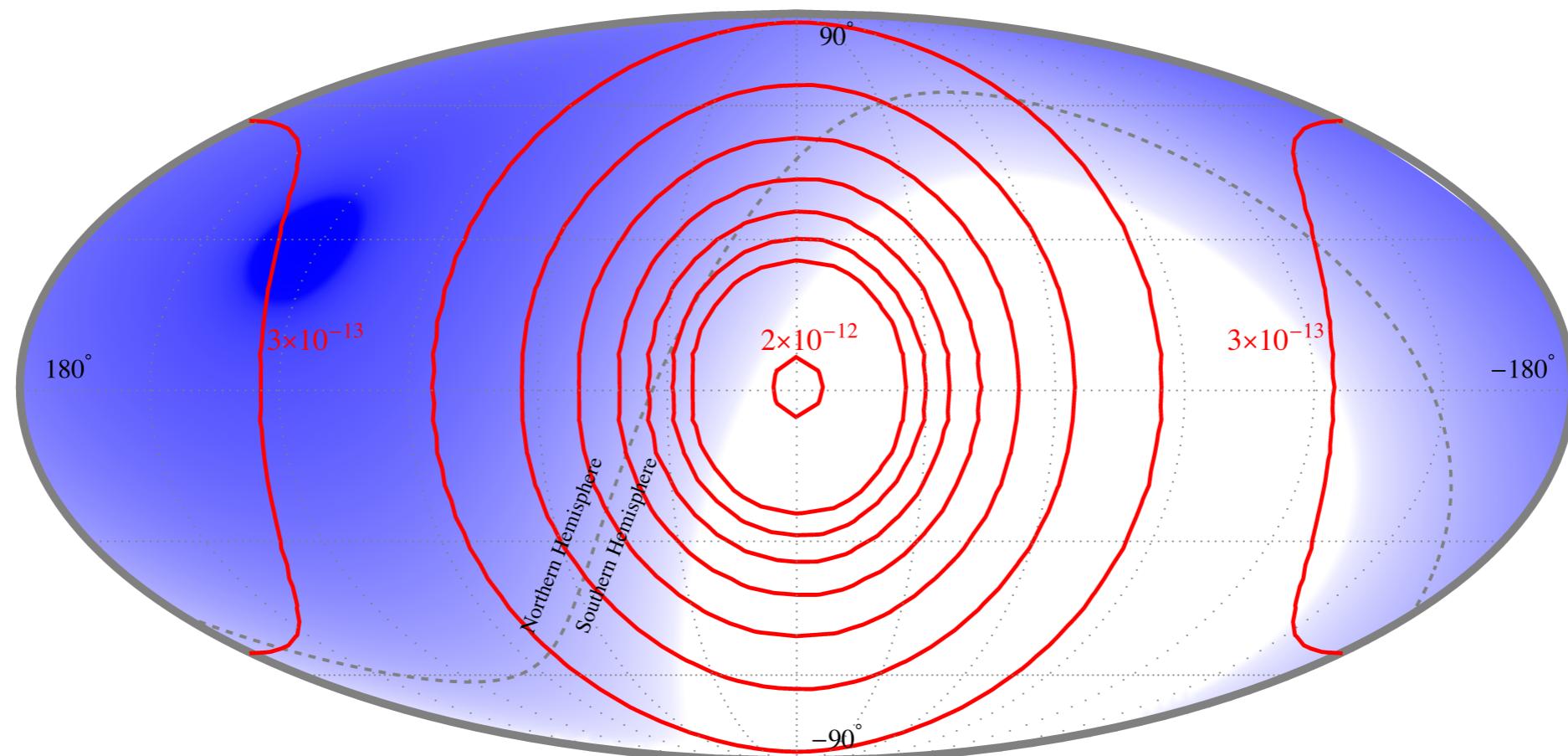


A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486

Gamma ray bounds

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CASA-MIA



A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486

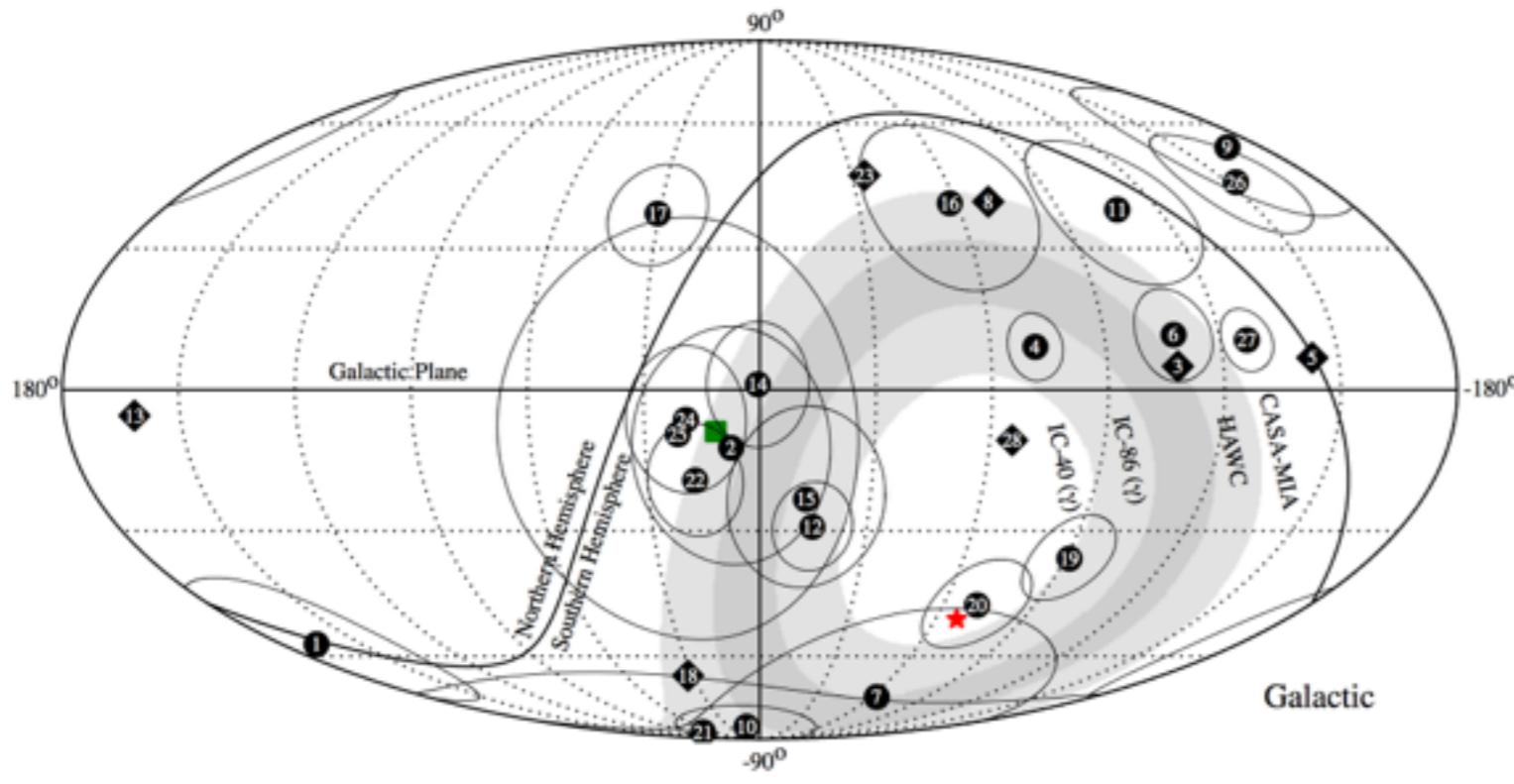
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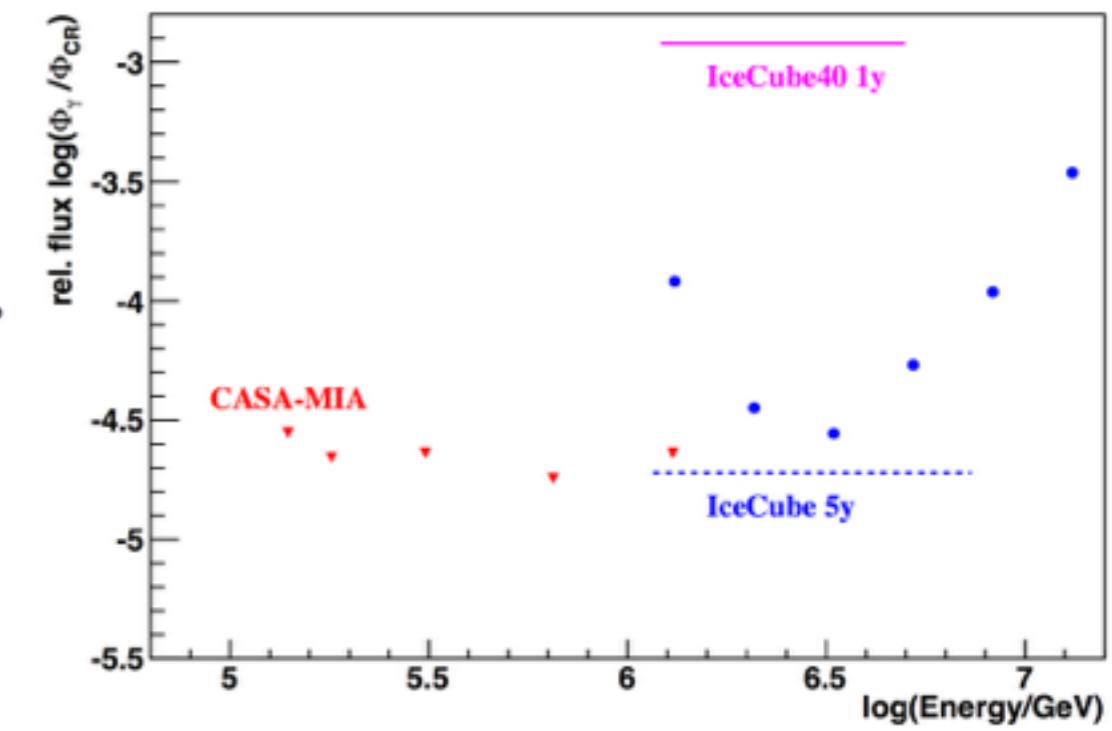
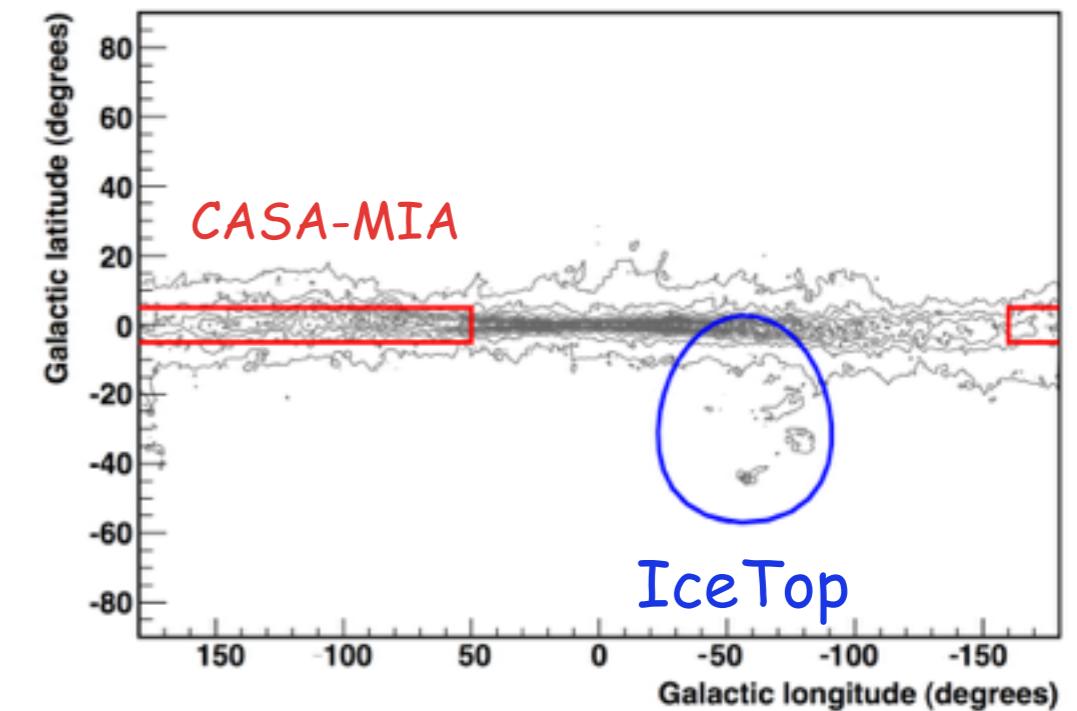
Future experiments



M. Ahlers and K. Murase,
PRD (2014), arXiv:1309.4077



IceTop



Gamma ray bounds

✓ Galactic component

Anisotropy

$$a_\gamma = \frac{\left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{GC}} - \left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{anti-GC}}}{\frac{d\Phi_{\text{CR}}}{dE}}$$

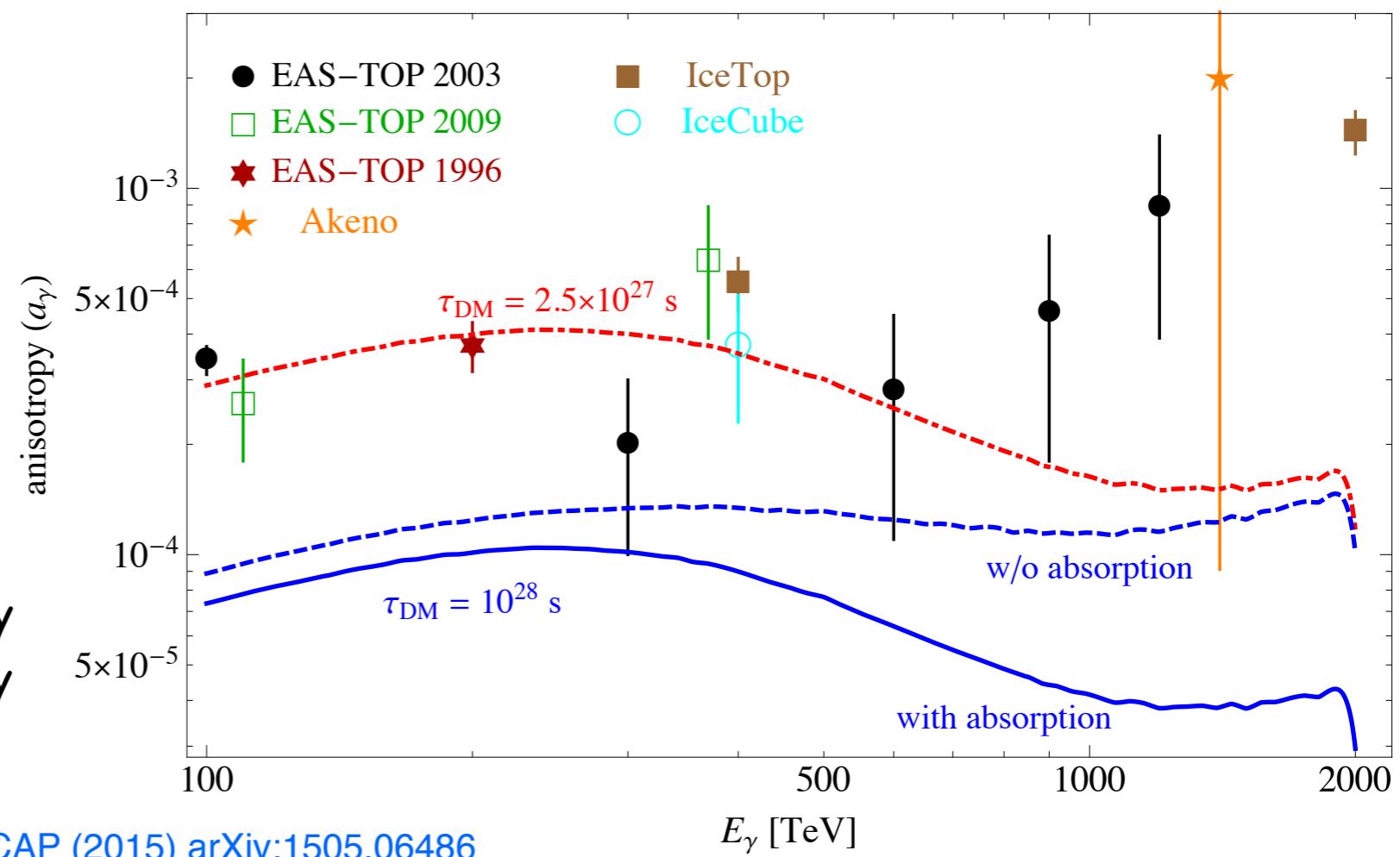
Total CR flux

✓ No need to γ /hadron discrimination

✓ Absorption suppress the anisotropy

✓ The bound 2.5×10^{27} s can be set

✓ Adding the phase info of anisotropy would improve the limits significantly



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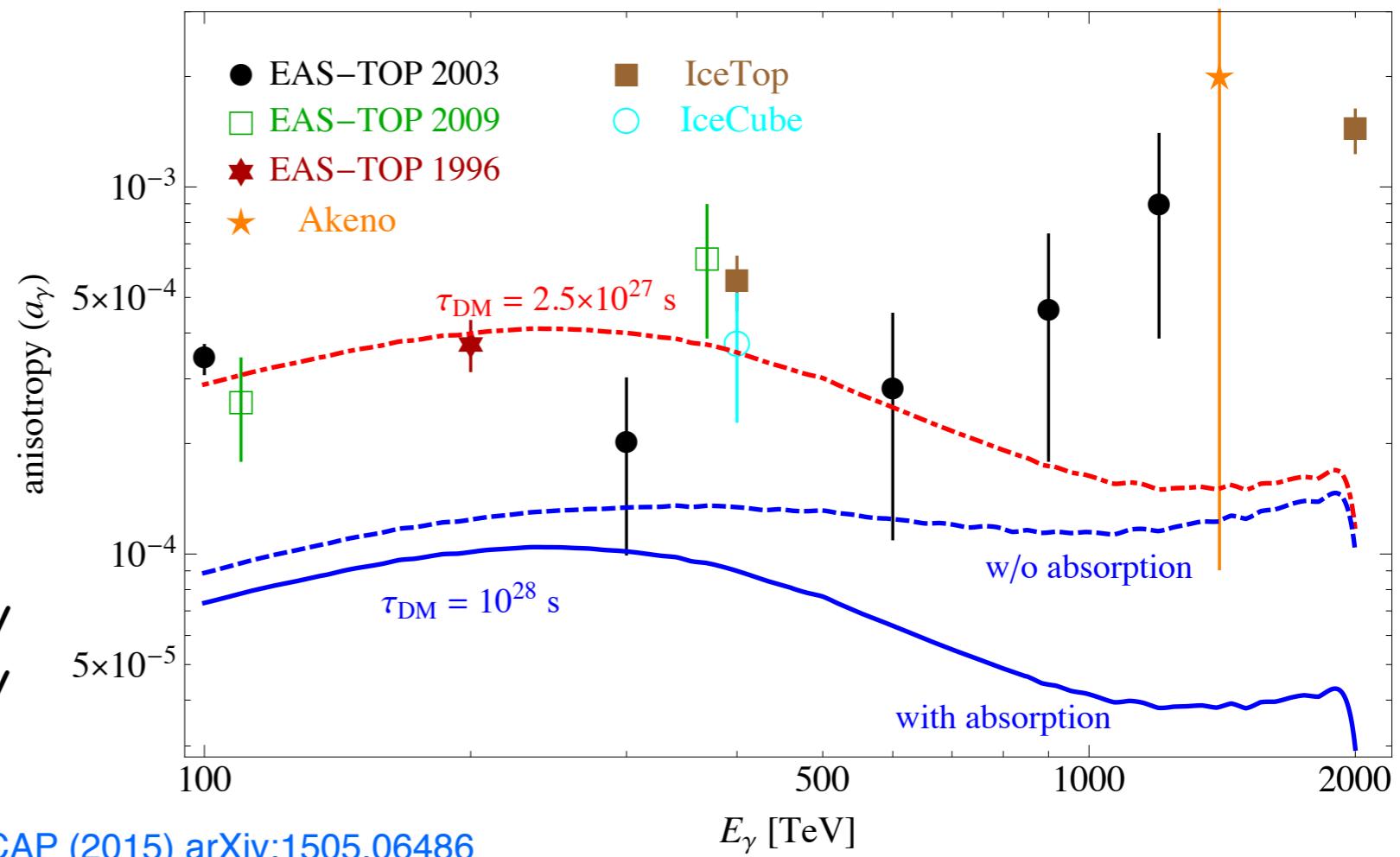
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A. E. and P. Serpico, JCAP (2015) arXiv:1505.06486

conclusions

- ✓ The excess of events observed by IceCube in the energy range $\sim 30 \text{ TeV} - 2 \text{ PeV}$ is an evidence for astrophysical flux or other "New Physics" induced fluxes

- ✓ Several features of the observed events motivate us for a DM interpretation: cut-off at $\sim 2 \text{ PeV}$, a mild dip in the $(400 - 1000) \text{ TeV}$ and anisotropy.

- ✓ We argued that a PeV-scale decaying DM, with generic decay channels, can naturally explain these features. The required lifetime is allowed by the current limits. Both the energy and angular distributions mildly prefer DM interpretation.

- ✓ With more statistics in the next few years, the DM interpretation of IceCube events can be tested. The gamma-ray flux expected in this scenario can be detected by the next generation of EAS detectors. Also, anisotropy measurements in the CR flux would be constraining.

conclusions



Thank you !

Confronting with energy distribution of IceCube data

three years data set

SM sector



Dark sector

portal type:

$$\mathcal{L}_{\text{protoal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

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A. Falkowski, J. Juknevich and J. Shelton
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heavy sterile neutrino, DM candidate

T. Higaki, R. Kitano and R. Sato, JHEP (2014)
arXiv:1405.0013

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UV completion:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_{B-L}$$

$$m_\phi \sim 10^{13} \text{ GeV}$$

"Higgs" field ϕ_{B-L} plays the role of inflaton

$$T_R \sim 10^7 \text{ GeV}$$

Confronting with energy distribution of IceCube data

three years data set

Leptogenesis: $\phi \rightarrow N_2 N_2$ $M_2 \sim 10^{12}$ GeV $\xrightarrow{\text{green arrow}} \frac{n_B}{s} \sim 10^{-10}$

DM abundance: $\Omega_{N_1} \simeq 0.2 \left(\frac{M_1}{4 \text{ PeV}} \right)^3 \left(\frac{T_R}{3 \times 10^7 \text{ GeV}} \right)^{-1}$

DM lifetime: $\tau_{N_1} \simeq 8 \times 10^{28} \text{ s} \left(\frac{M_1}{1 \text{ PeV}} \right)^{-1} \left(\frac{10^{-29}}{|y_N|^2} \right)$

DM decay channels: $\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 1}|^2$ NH

$\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 3}|^2$ IH

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three years data set

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arXiv:0908.1790 [hep-ph].

✓ d=4: $\mathcal{O}_{\text{DM}} \rightarrow N$

production mechanism:

$$m_\phi \gg m_N \quad \text{inflaton decay}$$

$$m_\phi \ll m_N \quad \text{freeze-in}$$

$$g\phi NN, \ g \simeq 10^{-6}$$

Confronting with energy distribution of IceCube data

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arXiv:0908.1790 [hep-ph].



$d = 5$:

$$\mathcal{O}_{\text{DM}} \rightarrow \chi \phi$$

singlet fermion and scalar
(Asymmetric DM)



$d = 6$:

other portals



For $d > 4$ there are more freedom in branching ratios. We have shown that for the most constrained model ($d=4$) a good fit to the data can be obtained. Obviously better fits can be achieved for $d > 4$.

Constraining DM properties

✓ DM lifetime

contribution of DM to the events in each bin should be smaller than N_{limit}

bin #	$\log_{10}(E_\nu/\text{TeV})$	$N_{\text{astro}}(E_\nu^{-2} \div E_\nu^{-2.3})$	N_{data}	$N_{\text{limit}} (E_\nu^{-2} \div E_\nu^{-2.3})$	N_{limit}
#1	1.4 – 1.6	9.46 ÷ 10	11	7.8 ÷ 7.46	16.6
#2	1.6 – 1.8	4.31 ÷ 5.3	6	6.53 ÷ 5.87	10.5
#3	1.8 – 2.0	4.55 ÷ 5.68	7	7.41 ÷ 6.58	11.8
#4	2.0 – 2.2	3.97 ÷ 4.82	3	3.98 ÷ 3.73	6.68
#5	2.2 – 2.4	3.32 ÷ 3.56	4	5.15 ÷ 5.01	8.00
#6	2.4 – 2.6	2.59 ÷ 2.42	2	3.65 ÷ 3.71	5.32
#7	2.6 – 2.8	1.96 ÷ 1.62	0	2.3 ÷ 2.3	2.3
#8	2.8 – 3.0	1.55 ÷ 1.1	0	2.3 ÷ 2.3	2.3
#9	3.0 – 3.2	1.2 ÷ 0.74	2	4.31 ÷ 4.64	5.32
#10	3.2 – 3.4	0.92 ÷ 0.5	1	3.3 ÷ 3.51	3.89
#11	3.4 – 3.6	0.73 ÷ 0.35	0	2.3 ÷ 2.3	2.3
#12	3.6 – 3.8	1.72 ÷ 0.76	0	2.3 ÷ 2.3	2.3

Poisson statistics:

at $q\%$ C.L.

$$\frac{q}{100} = \frac{\int_0^{N_{\text{limit}}^i} L(N_{\text{data}}^i, N) dN}{\int_0^{\infty} L(N_{\text{data}}^i, N) dN}$$

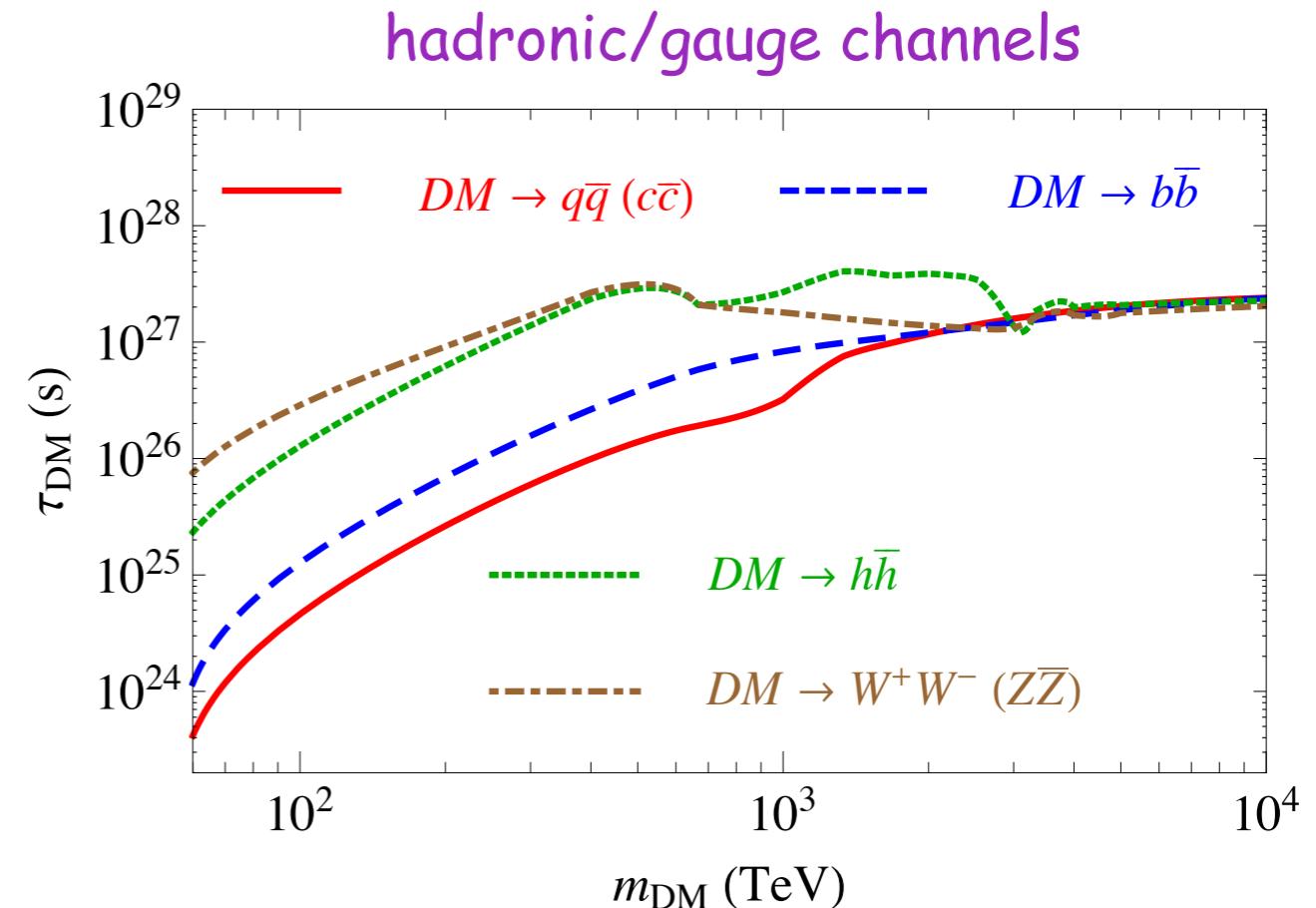
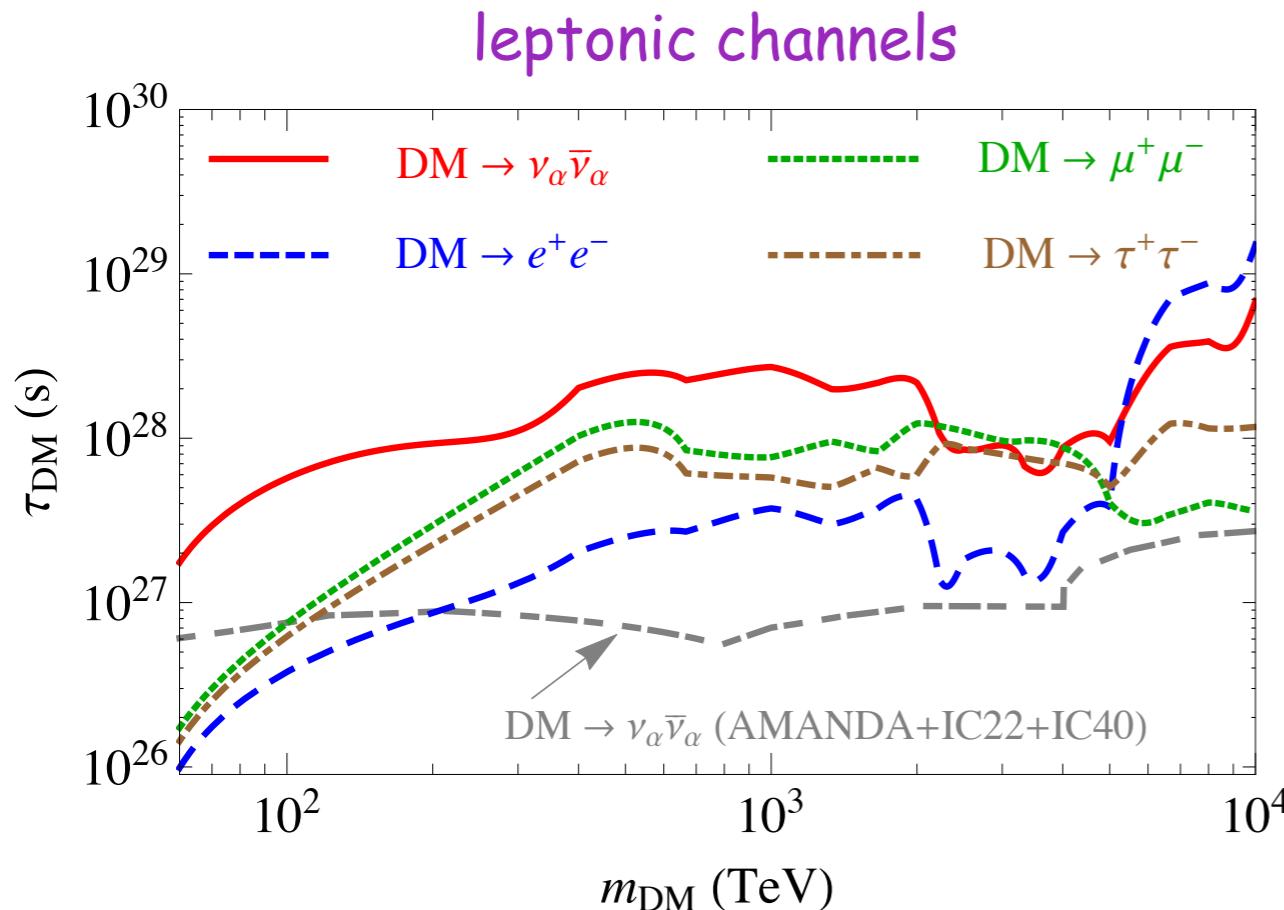
$$L(N_{\text{data}}^i, N) = \frac{(N + N_{\text{astro}}^i)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-(N + N_{\text{astro}}^i)}$$

or

$$L(N_{\text{data}}^i, N) = \frac{(N)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-N}$$

Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

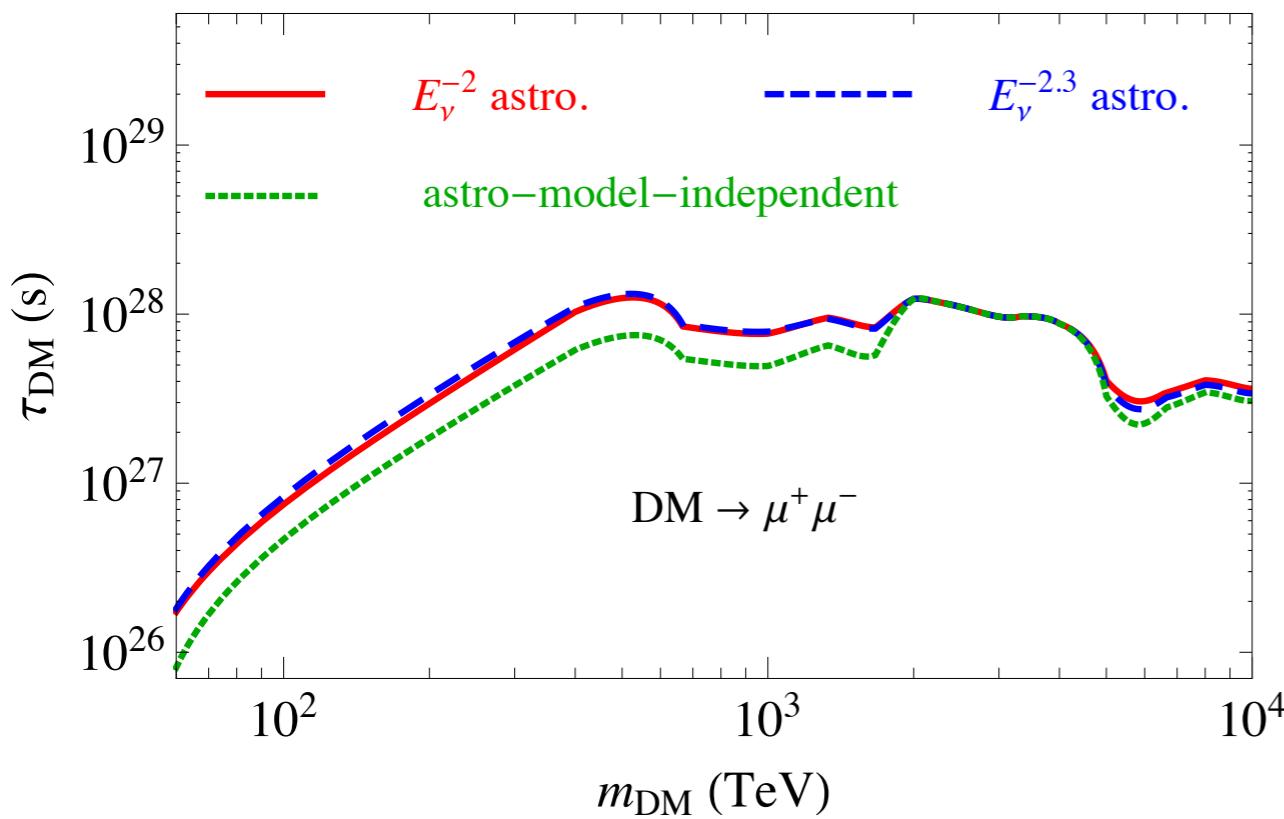
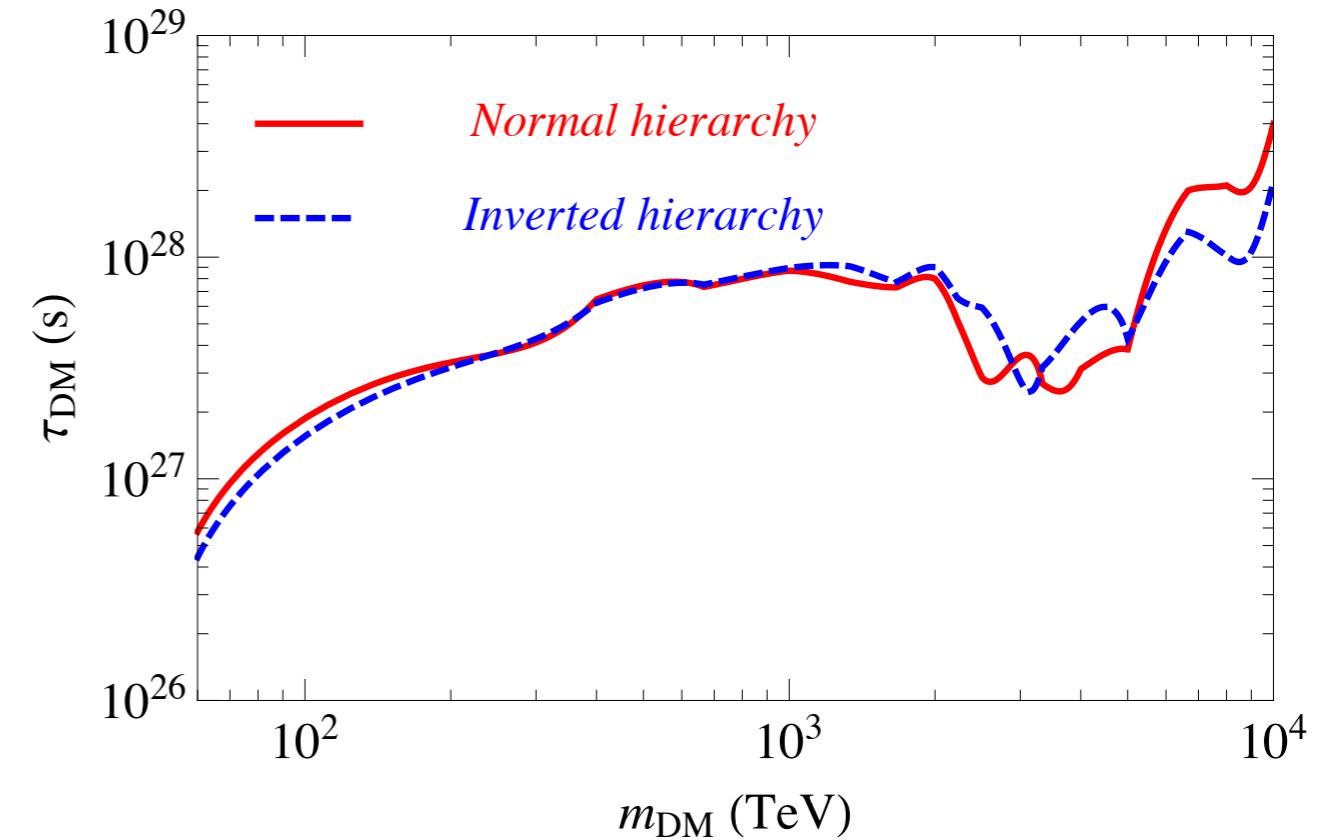


- ✓ at least one order of magnitude stronger lower limit on the DM lifetime, in the relevant DM mass range
- ✓ for a specific model, different channels should be scaled according to the corresponding branching ratios

Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

NH and IH cases



dependence on the astro.
model?

Constraining DM properties

✓ Annihilation cross section

The lower part (< 100 TeV) of the observed spectrum can be used to probe $\langle\sigma v\rangle$

Constraining DM properties

✓ Annihilation cross section

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The isotropic components of neutrino flux from DM annihilation:

The residual isotropic flux from the Galactic halo (anti-GC direction)

$$\frac{dJ_{\text{iso}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{1}{4\pi m_{\text{DM}}^2} \frac{dN}{dE_\nu} (\text{l.o.s.})_{\text{anti-GC}} \quad \text{where } (\text{l.o.s.})_{\text{anti-GC}} = \int_0^\infty \rho^2 [r(s, b=0, l=\pi)] ds$$

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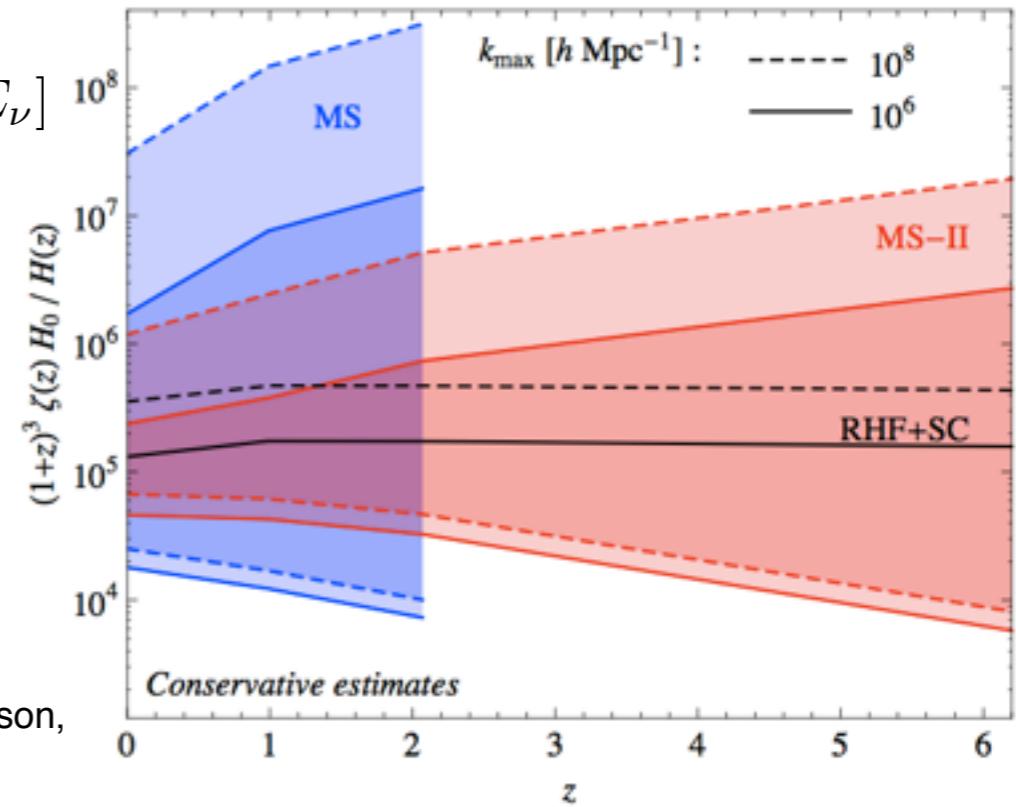
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The cosmic flux from all redshift

$$\frac{dJ_{\text{cos}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{\Omega_{\text{DM}}^2 \rho_c^2}{4\pi m_{\text{DM}}^2} \frac{c}{H_0} \int_0^\infty \frac{(1+z)^3 \zeta(z) dz}{\sqrt{\Omega_m(1+z)^3 + \Omega_\Lambda}} \frac{dN}{dE_\nu} [(1+z) E_\nu]$$

$\zeta(z)$ flux multiplier (DM clustering)



E. Sefusatti, G. Zaharijas, P. D. Serpico, D. Theurel and M. Gustafsson,
Mon. Not. Roy. Astron. Soc. (2014) [arXiv:1401.2117].

Constraining DM properties

✓ upper limits on annihilation cross section $\langle\sigma v\rangle$ (90% C.L.)

minimum ÷ maximum value used for $\zeta(z)$ unit of $\langle\sigma v\rangle$ is $10^{-22} \text{ cm}^3 \text{s}^{-1}$

m_{DM} $\text{DM} + \text{DM} \rightarrow$	100 TeV	50 TeV	30 TeV
$\nu_\alpha \bar{\nu}_\alpha$	1.39 ÷ 0.22	1.21 ÷ 0.36	2.44 ÷ 0.88
$q\bar{q}$	489 ÷ 84.5	1427 ÷ 299	9934 ÷ 4603
$b\bar{b}$	185 ÷ 30.4	517 ÷ 106	3514 ÷ 1621
$c\bar{c}$	592 ÷ 100	1708 ÷ 348	11218 ÷ 5215
e^+e^-	14.7 ÷ 2.38	17.8 ÷ 5.06	41.3 ÷ 14.2
$\mu^+\mu^-$	4.47 ÷ 0.65	9.06 ÷ 1.6	23.7 ÷ 9.23
$\tau^+\tau^-$	5.84 ÷ 0.93	10.9 ÷ 2.3	28.5 ÷ 10.8
$h\bar{h}$	21.2 ÷ 3.36	53.4 ÷ 9.49	177 ÷ 76.5
$Z\bar{Z}$	11.9 ÷ 2.05	18.1 ÷ 4.09	40.7 ÷ 16.3
W^+W^-	14.4 ÷ 2.4	23.7 ÷ 4.96	54.5 ÷ 22.3

✓ for some final states (neutrinos, charged leptons) the limit is a bit stronger than the unitary bound