

Benchmark models for Dark Matter searches at the LHC

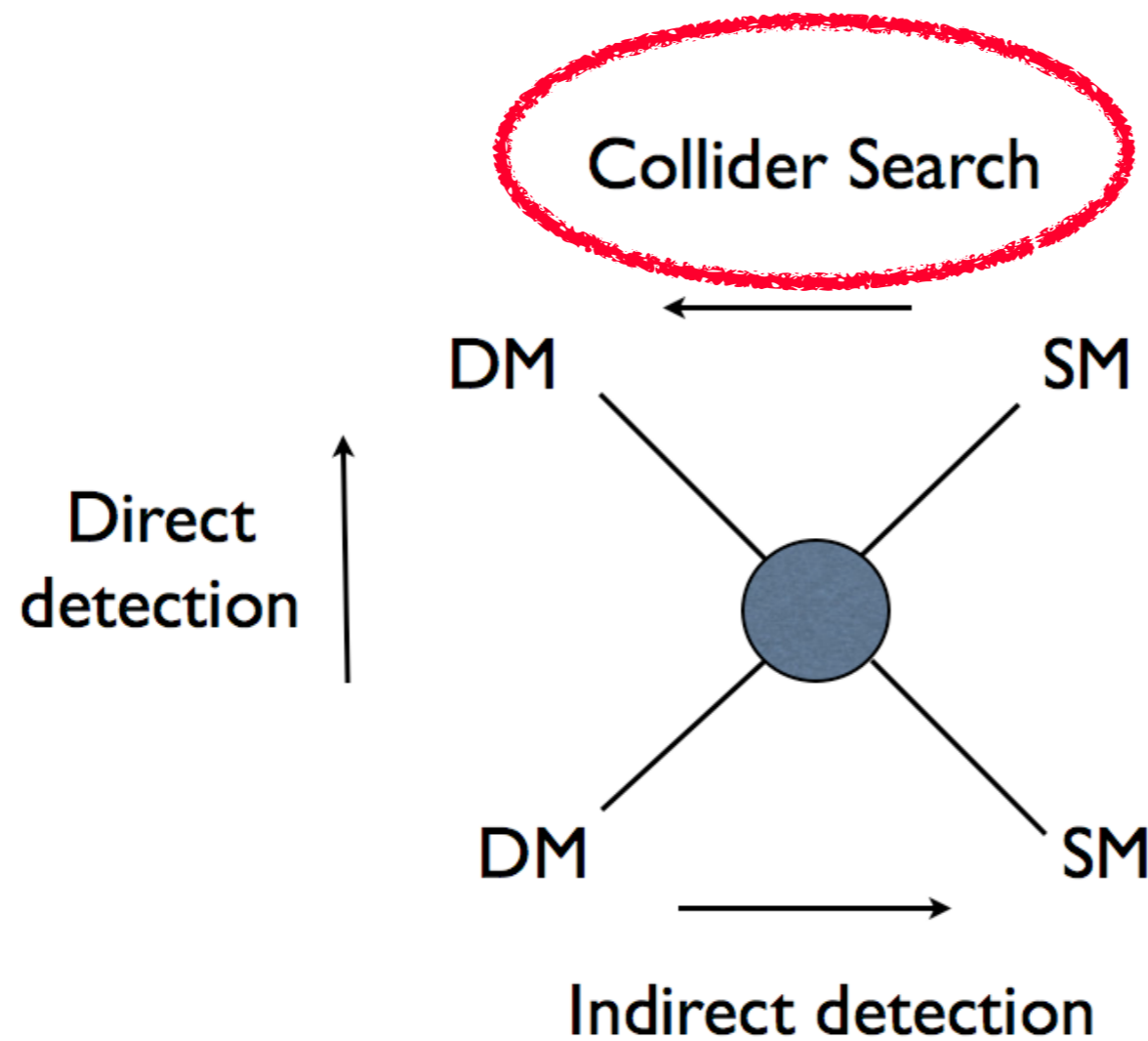
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in preparation

with M. Papucci, K. Zurek (LBL)
and A. Vichi (CERN)

DSU 2015, Kyoto

Detecting Dark Matter Signals



Two approaches

Top-down

bottom-up

Supersymmetry

Effective Field Theory

Extra dimensions

Simplified models

etc.

relatively well known

relatively new

Bottom-up approach: Effective Field Theory

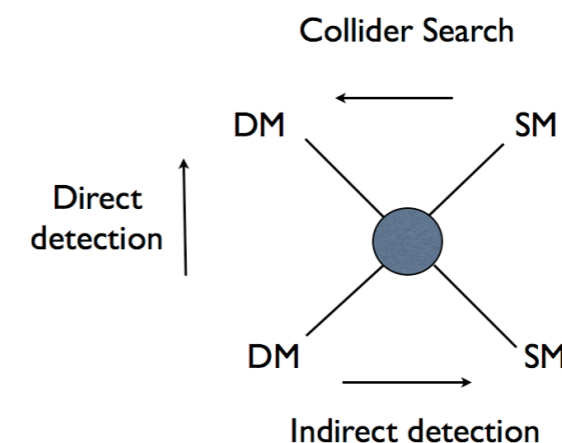
Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

Name	Operator	Coefficient
M1	$\bar{\chi}\chi\bar{q}q$	$m_q/2M_*^3$
M2	$\bar{\chi}\gamma^5\chi\bar{q}q$	$im_q/2M_*^3$
M3	$\bar{\chi}\chi\bar{q}\gamma^5q$	$im_q/2M_*^3$
M4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	$m_q/2M_*^3$
M5	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/2M_*^2$
M6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/2M_*^2$
M7	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
M8	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/8M_*^3$
M9	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^3$
M10	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/8M_*^3$

Name	Operator	Coefficient
C1	$\chi^\dagger\chi\bar{q}q$	m_q/M_*^2
C2	$\chi^\dagger\chi\bar{q}\gamma^5q$	im_q/M_*^2
C3	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu q$	$1/M_*^2$
C4	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu\gamma^5q$	$1/M_*^2$
C5	$\chi^\dagger\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^2$
C6	$\chi^\dagger\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^2$
R1	$\chi^2\bar{q}q$	$m_q/2M_*^2$
R2	$\chi^2\bar{q}\gamma^5q$	$im_q/2M_*^2$
R3	$\chi^2 G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/8M_*^2$
R4	$\chi^2 G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^2$

[Goodman et al. '10]

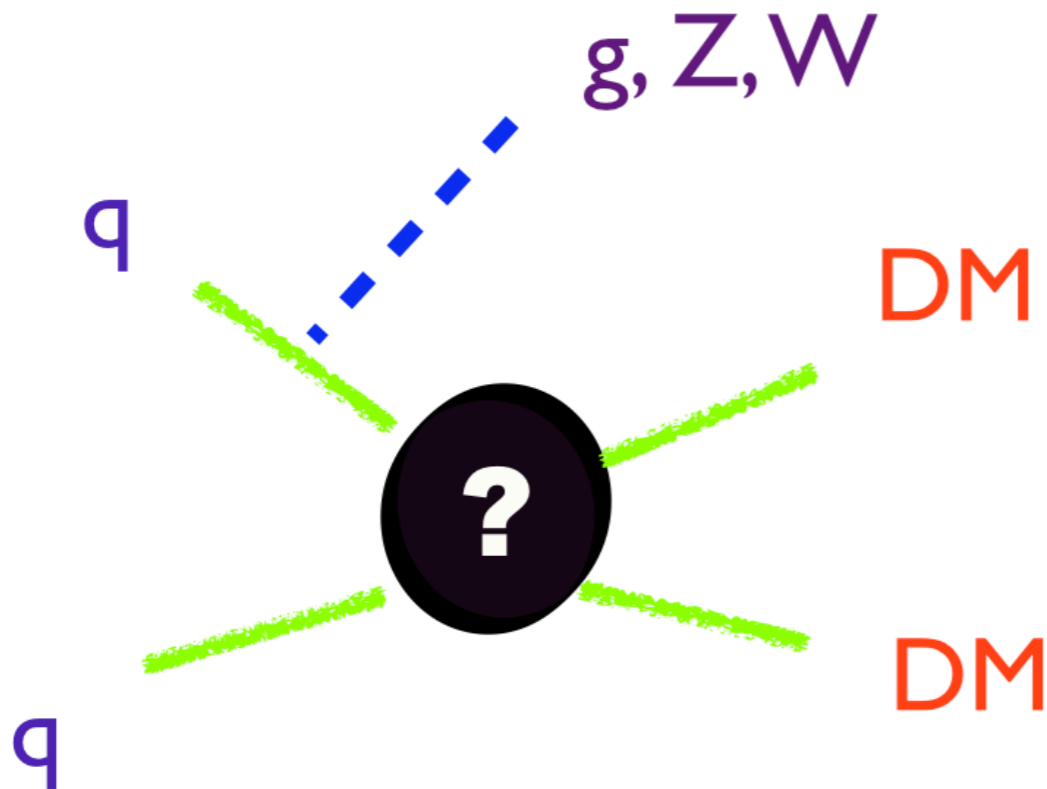
- DM is an SM singlet
- DM is odd under a Z2 symmetry
- DM is either scalar or fermion



Bottom-up approach: Effective Field Theory

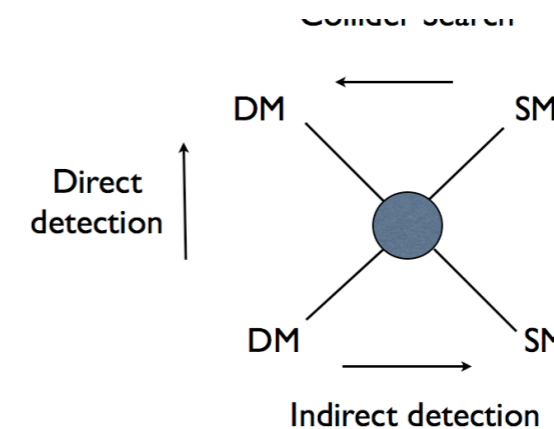
Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\chi\bar{q}q$	$i/m_q/M_*^3$

Name	Operator	Coefficient
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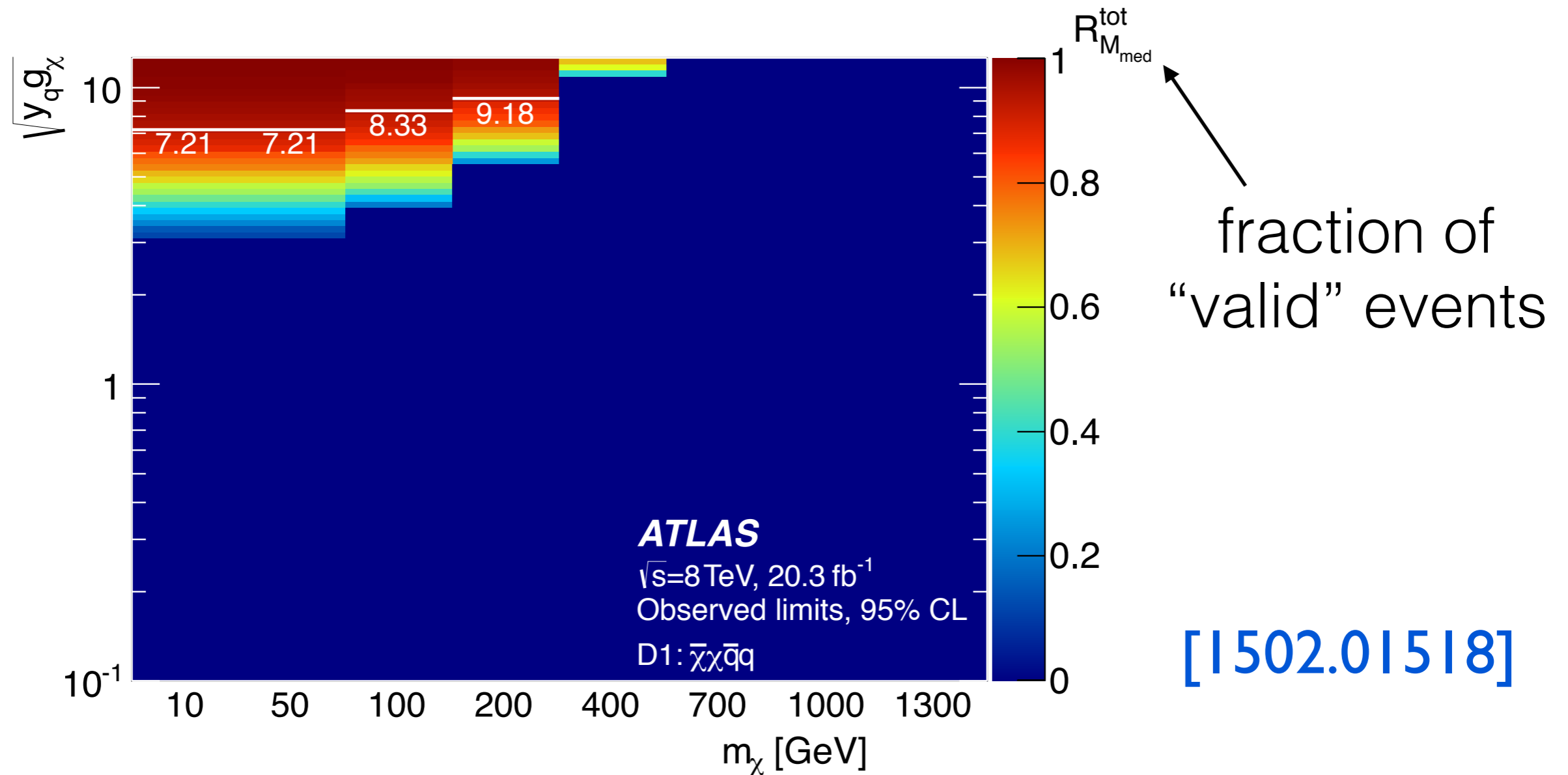
Signature:
Mono-events
with missing
energy
at the LHC

- DM is an SM singlet
- DM is odd under a Z_2 symmetry
- DM is either scalar or fermion



However, in many cases, Effective Field Theory
breaks down at LHC energy scales

Q (momentum transfer) $>$ M (mediator mass)

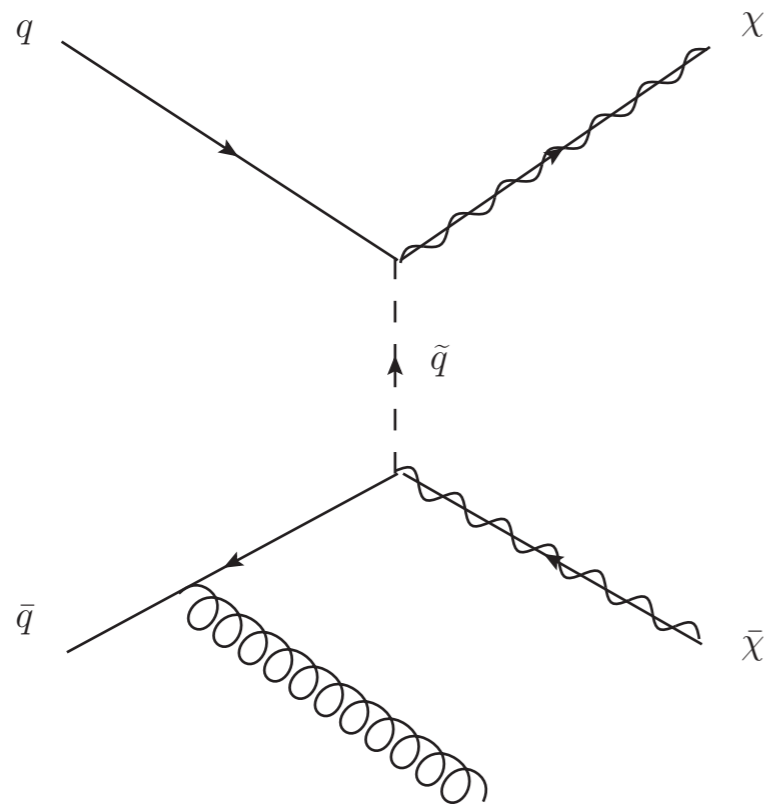


(s-channel mediator)

UV-complete models, i.e. simplified DM models are more appropriate.

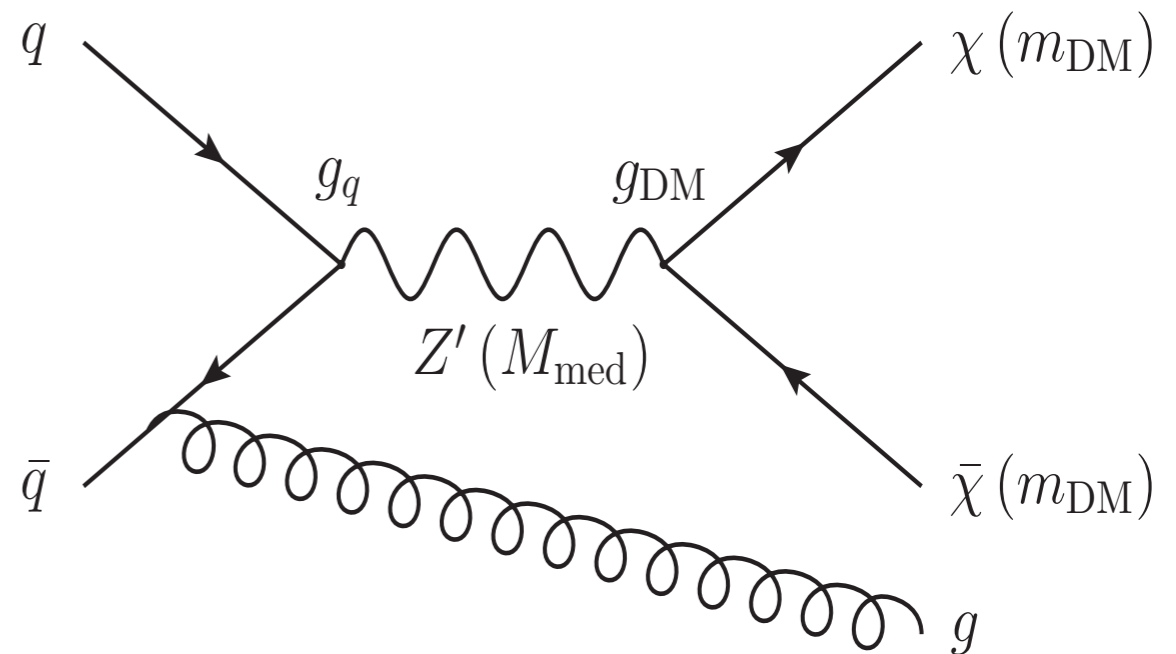
e.g.

squark-neutralino model



[Papucci et al. '14]

Z' model



[Buchmueller et al. '14]

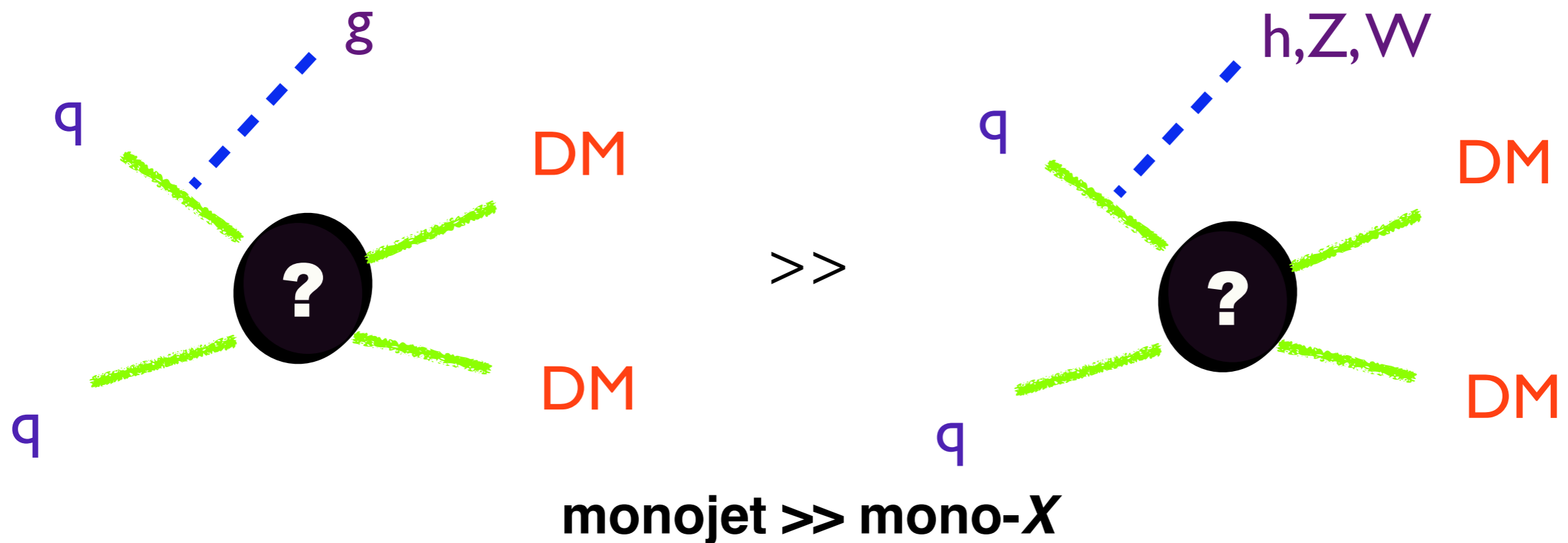
are the benchmark models to study monojet signatures

Simplified models related to monojet are well known.

How about simplified models for other mono-particle, i.e. mono- X , which is also actively looked for at the LHC?

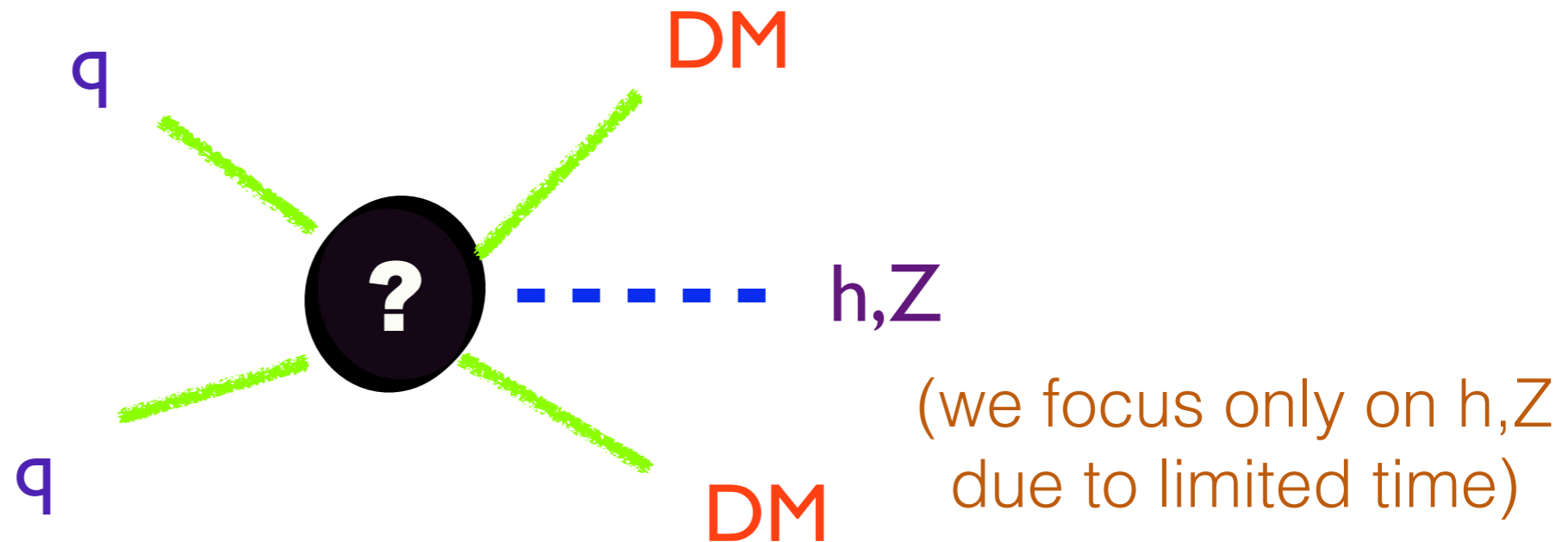
(W, Z boson, photon, Higgs, top, bottom)

firstly, colorless X from initial state cannot be dominant over jet from initial state at the **LHC**



simply because QCD processes are more important at the LHC

Simplified models of mono- X for colorless X
must have X radiated from the *final state*

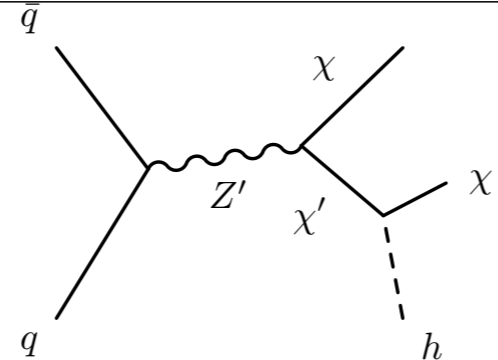
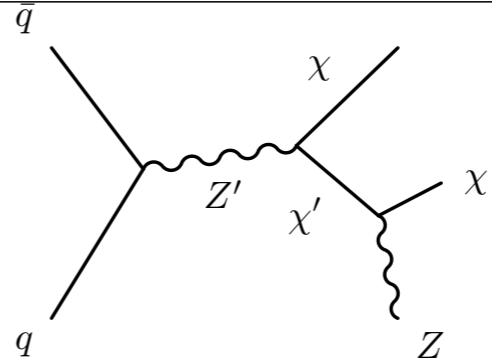
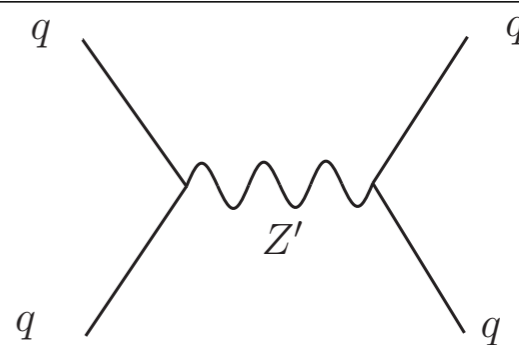
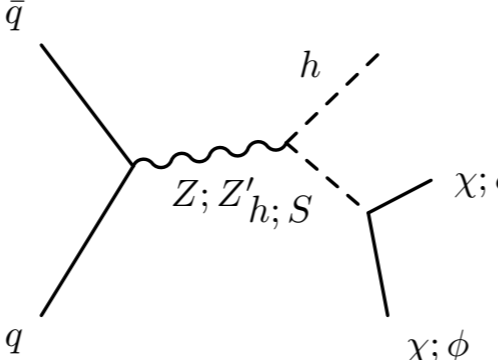
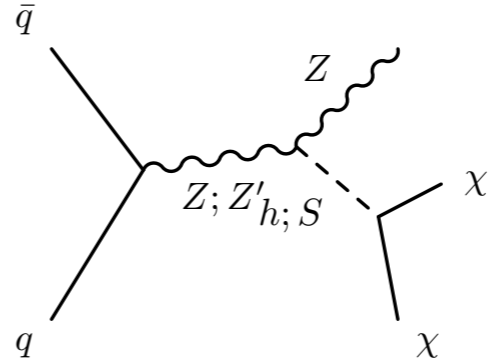
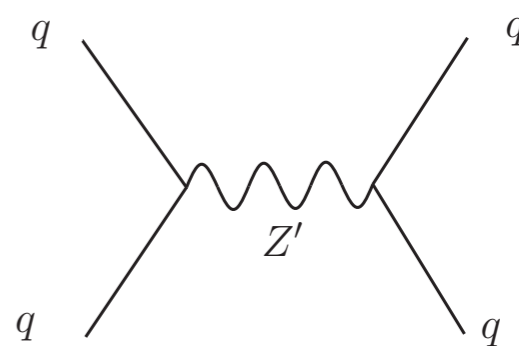
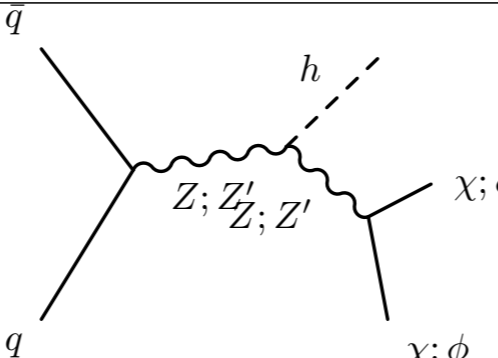

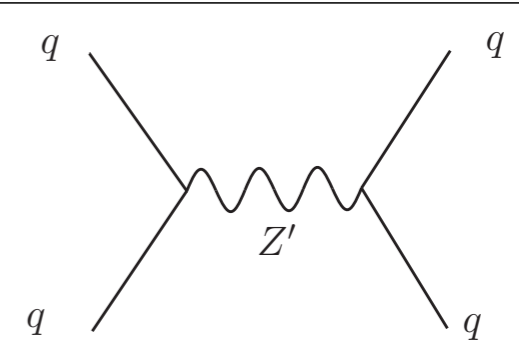
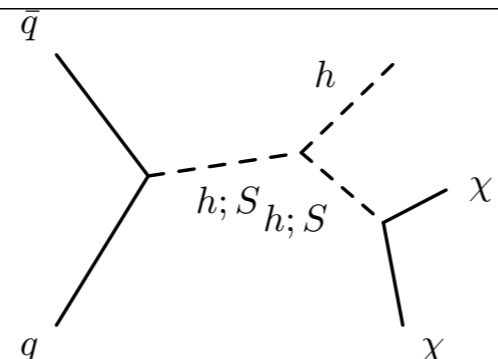
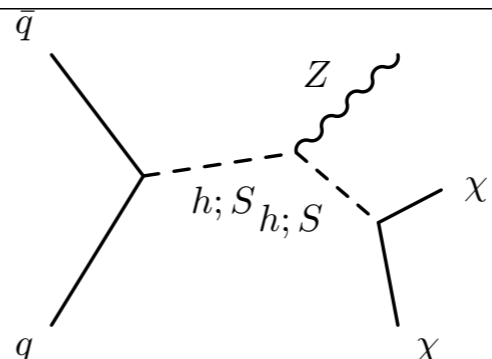
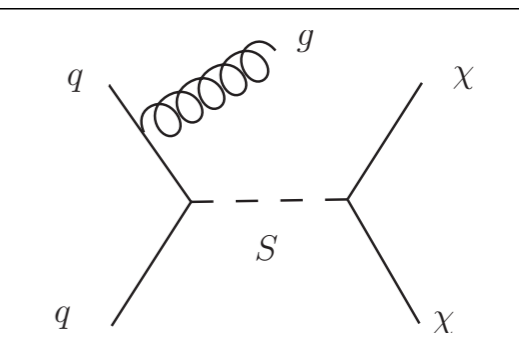


We come up with a minimal set of simplified models by imposing the following requirements:

- DM is an SM singlet and is produced in pair
- tree-level topology
- minimal number of mediators

mono- X simplified models I


 s-channel

Model	mono- h	mono- Z	direct constraints
Inelastic DM			
2HDM			
s-channel vector			
s-channel scalar			

constraints from direct searches

Model	mono- h	mono- Z	direct constraints
Inelastic DM			
2HDM			
s-channel vector			
s-channel scalar			

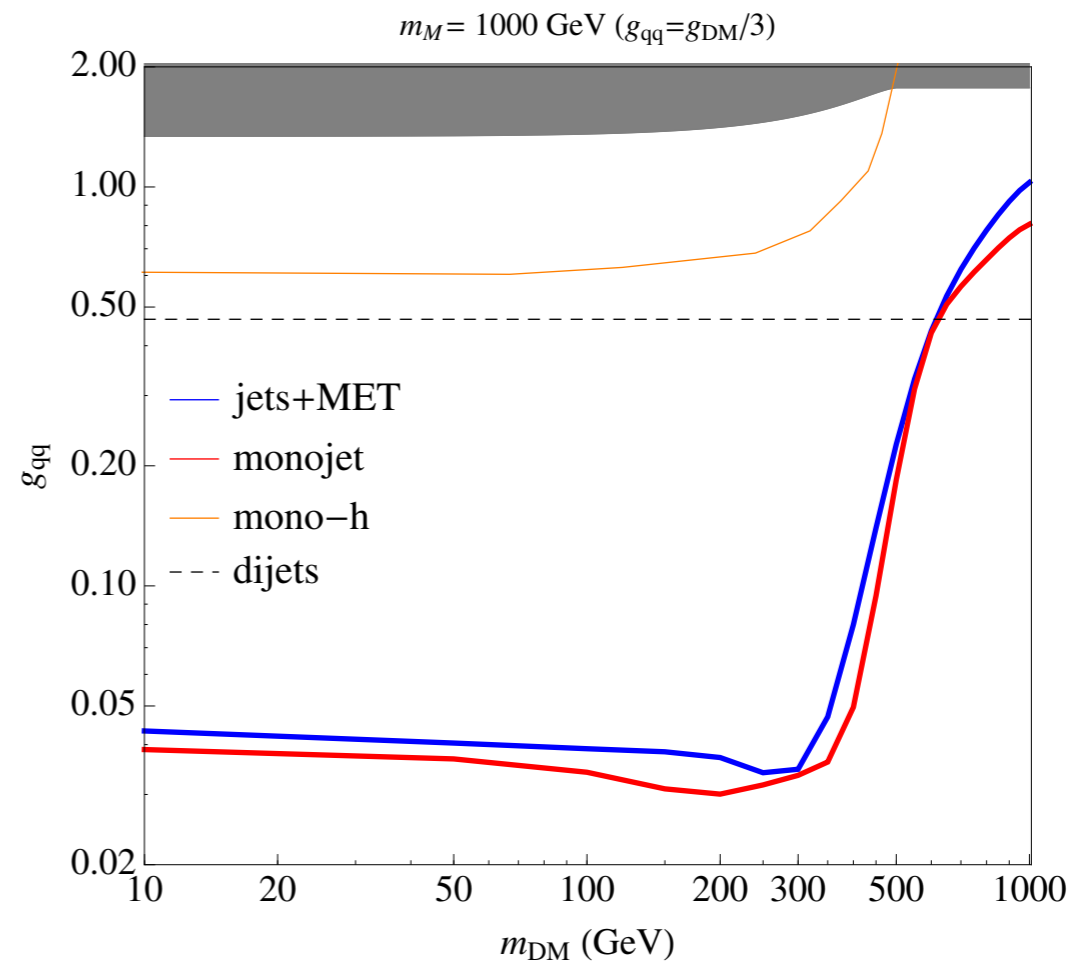
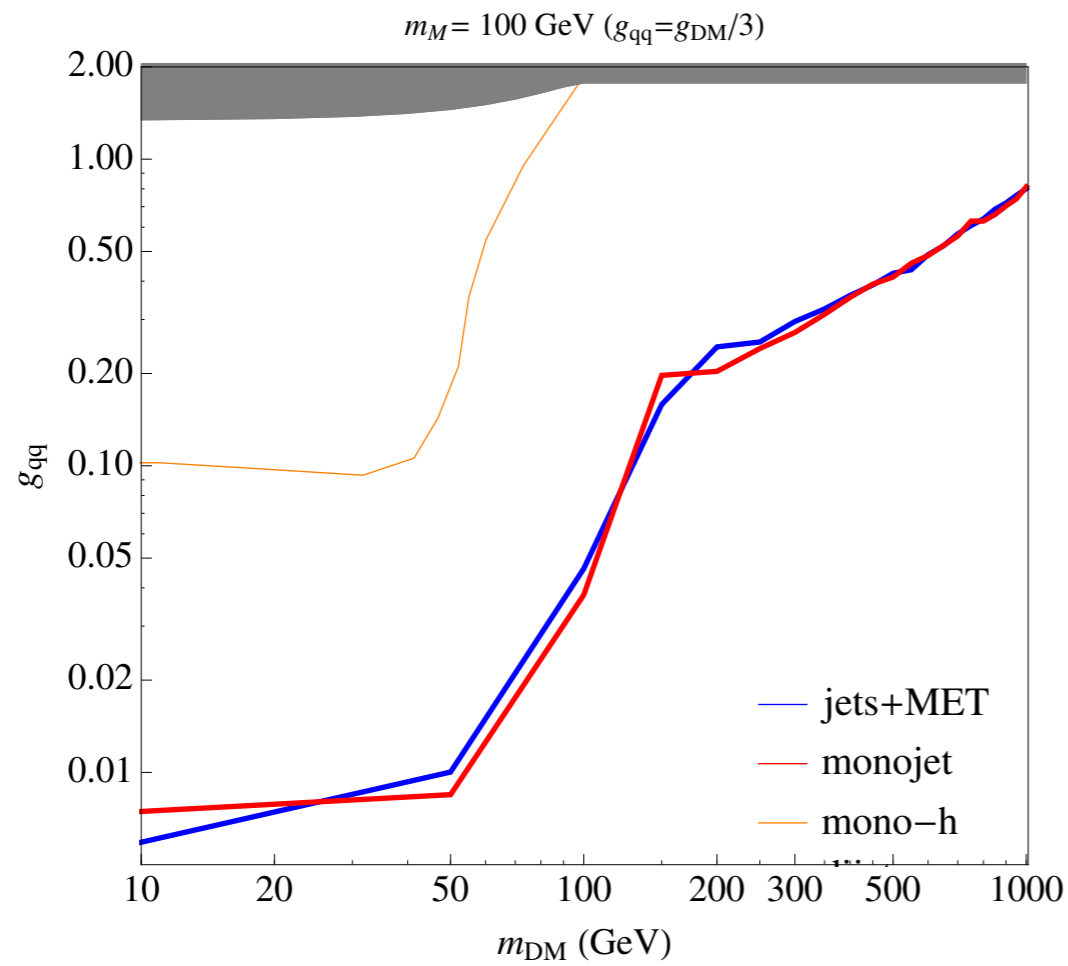
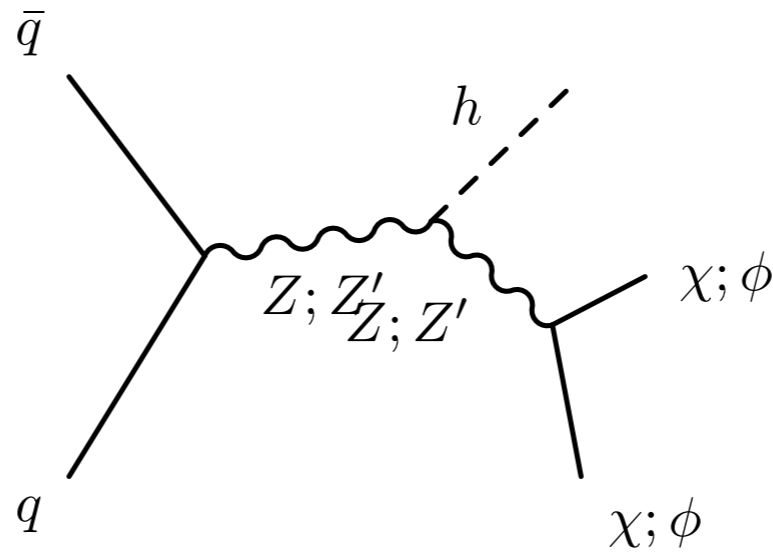
dijet

dijet

dijet

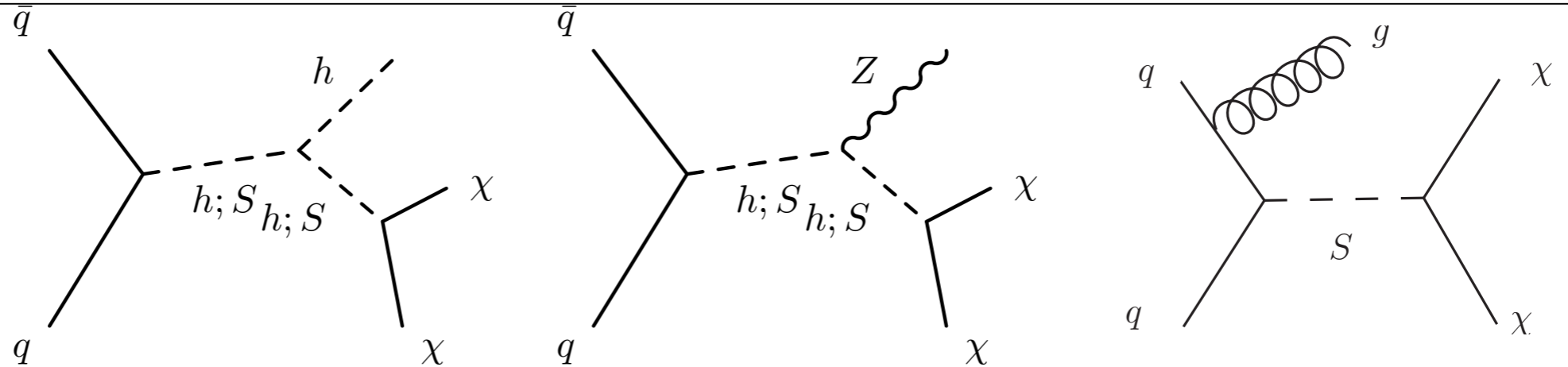
monojet

s-channel vector



Monojet, dijet, jet + MET searches are more powerful than mono-Higgs

s-channel scalar



no searches can constrain the couplings
up to perturbative level

see also [Carpenter et al.'13]

Model	mono- h	mono- Z	direct constraints
Inelastic DM			
2HDM			

mono- h and mono- Z can overcome dijet constraints -> good benchmark models for mono- h and mono- Z

(will discuss more later if time allows)

mono- X simplified models II

t -channel



Model	mono- h	mono- Z	direct constraints
Squarks/sbottoms			
Inelastic squark			

mono- h and mono- Z are weak compared to squark searches (jets + MET etc.)

We have studied a comprehensive list
of mono- X simplified models.
The following are the benchmark models
giving dominant mono- X signals at the LHC

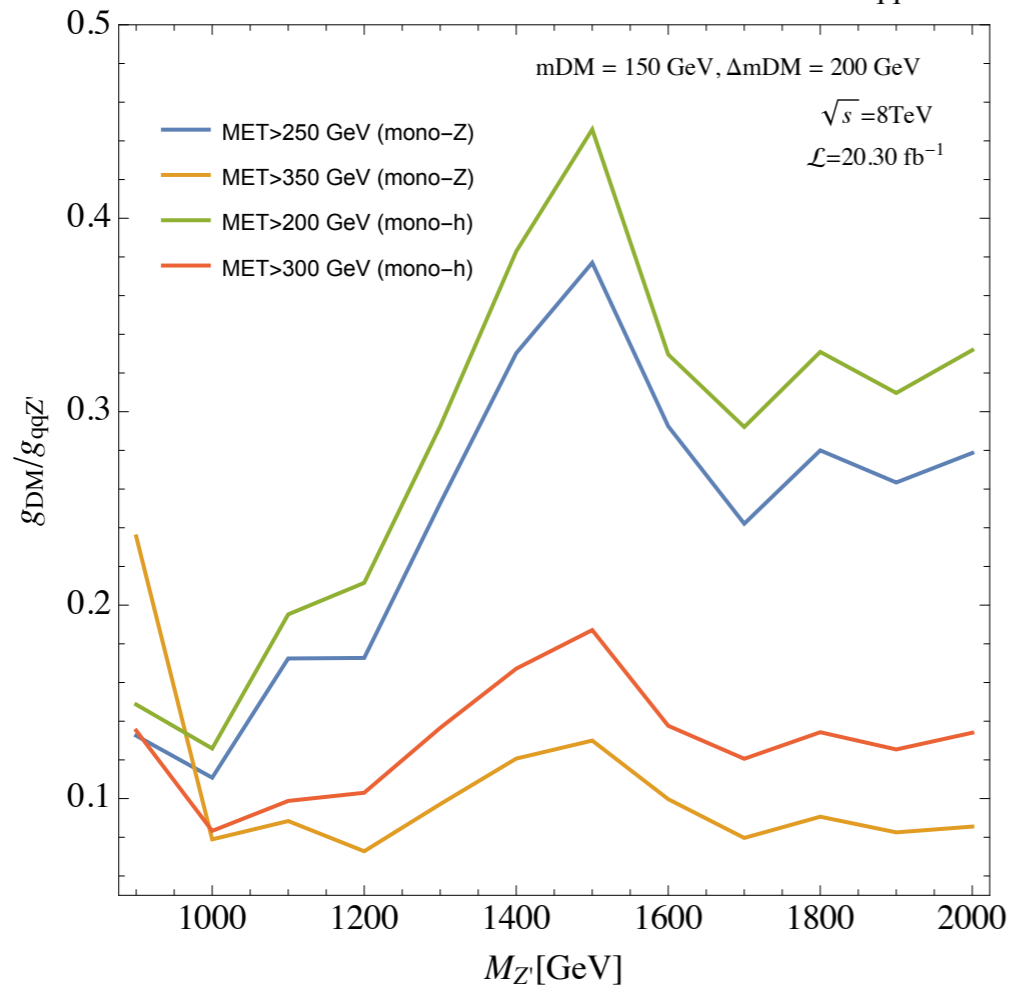
Search	Model where it matters
mono- h	Inelastic DM; 2HDM
mono- z	Inelastic DM; 2HDM
mono-jet	squark mediated production, compressed spectrum; Z'
mono- b	sbottom mediated production, compressed spectrum
mono- t	RPV stops; non-resonant mono- t

TABLE I: Summary of results: for each mono- X search we list the models where the analysis can exclude part of the parameter space not already ruled out by some other search.

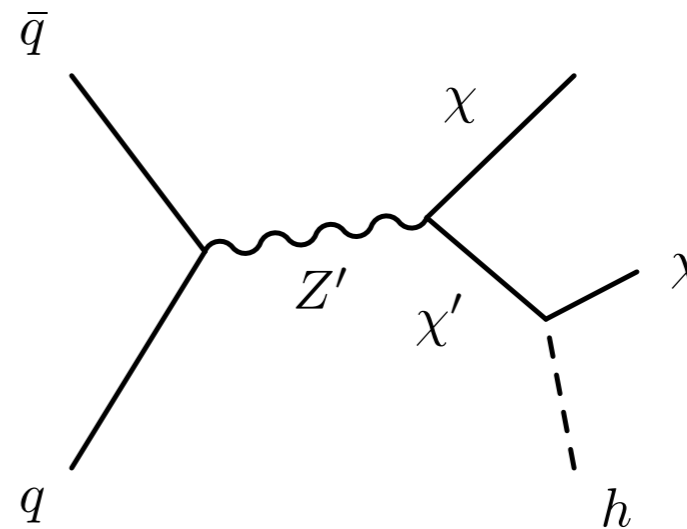
Thank you!

Backups

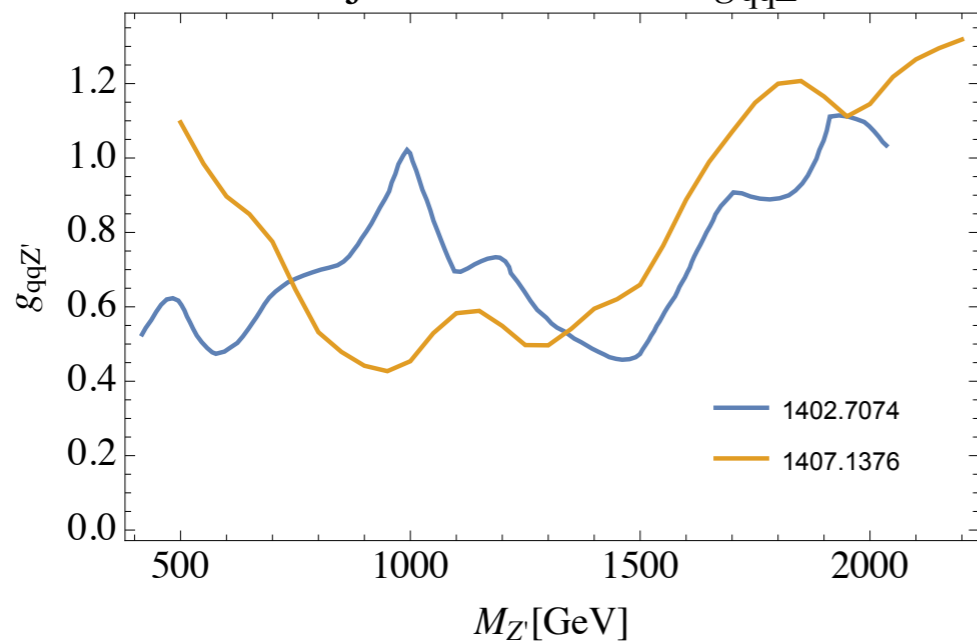
Mono-Z/h exclusion limits VS $g_{\text{DM}}/g_{\text{qq}Z'}$



Inelastic DM

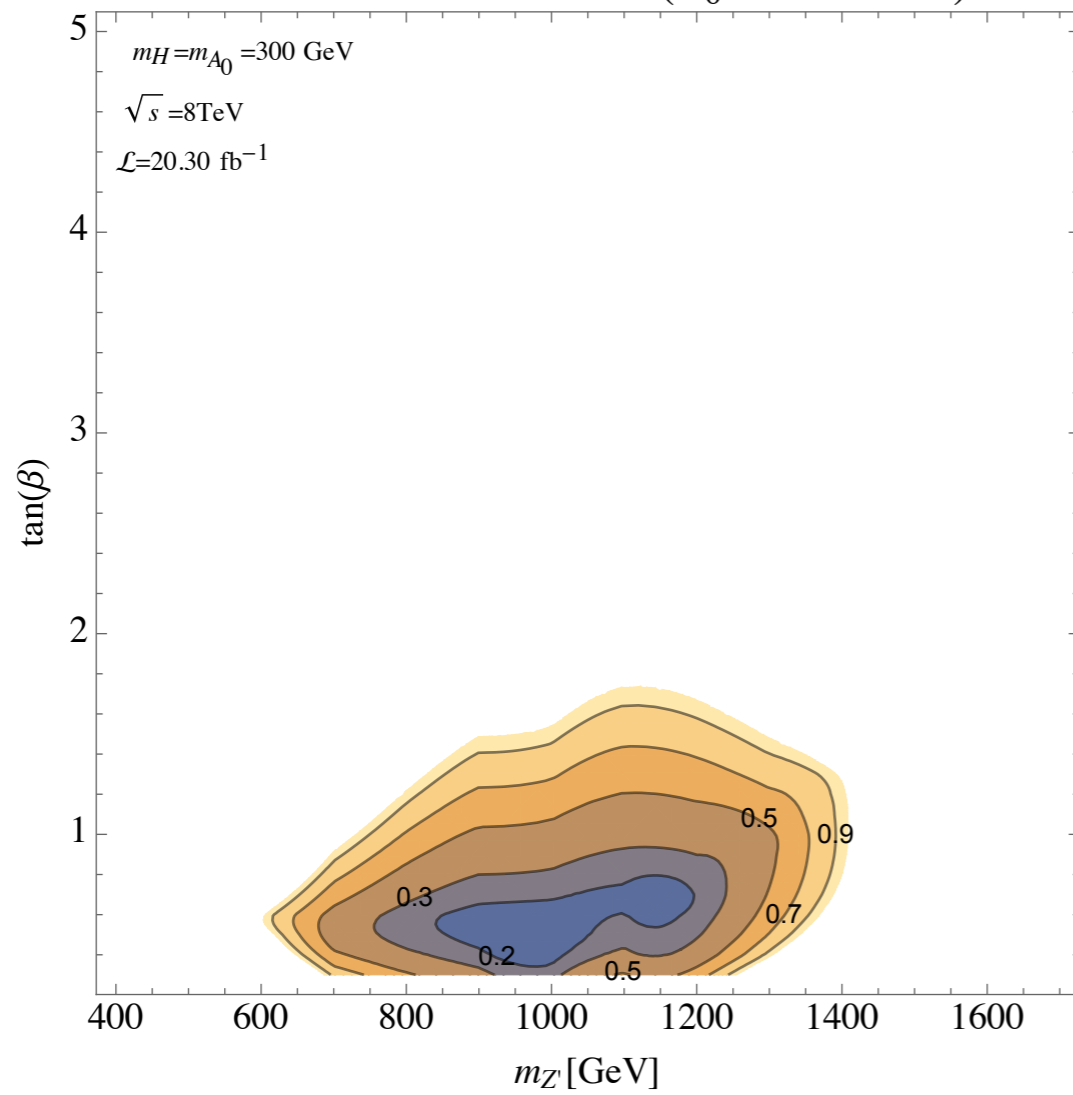


Dijet Constraints on $g_{\text{qq}Z'}$

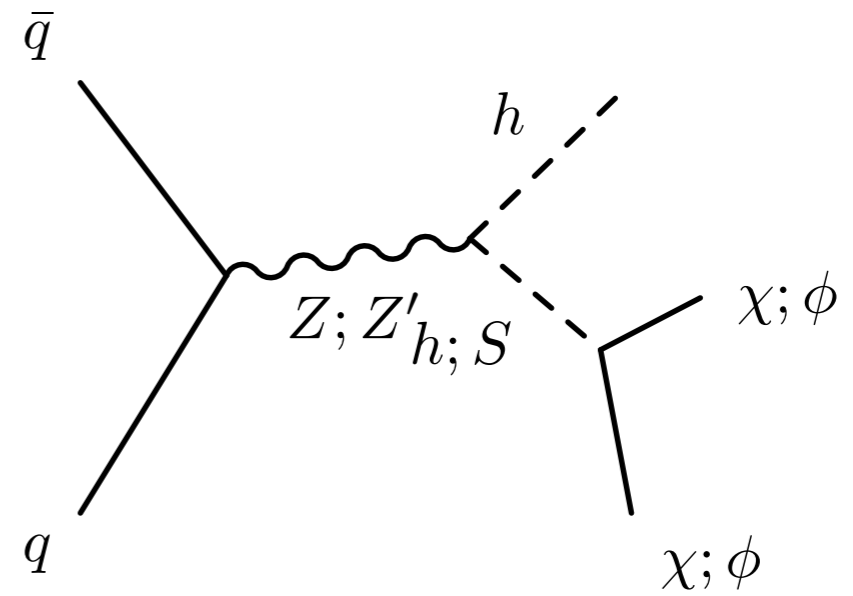


Dijet constraint

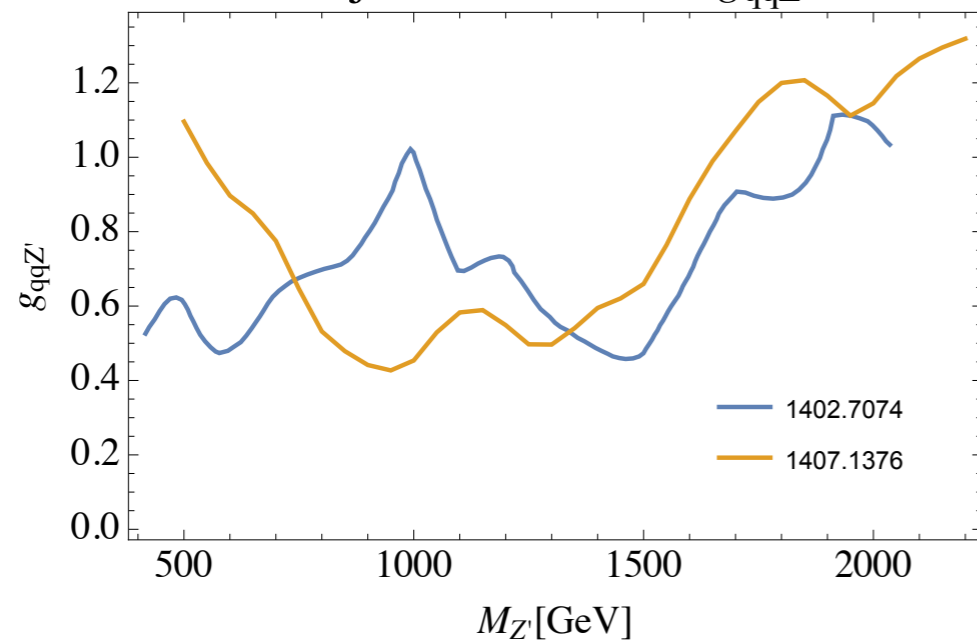
Exclusion limits on BR($A_0 \rightarrow$ Invisible)



2HDM

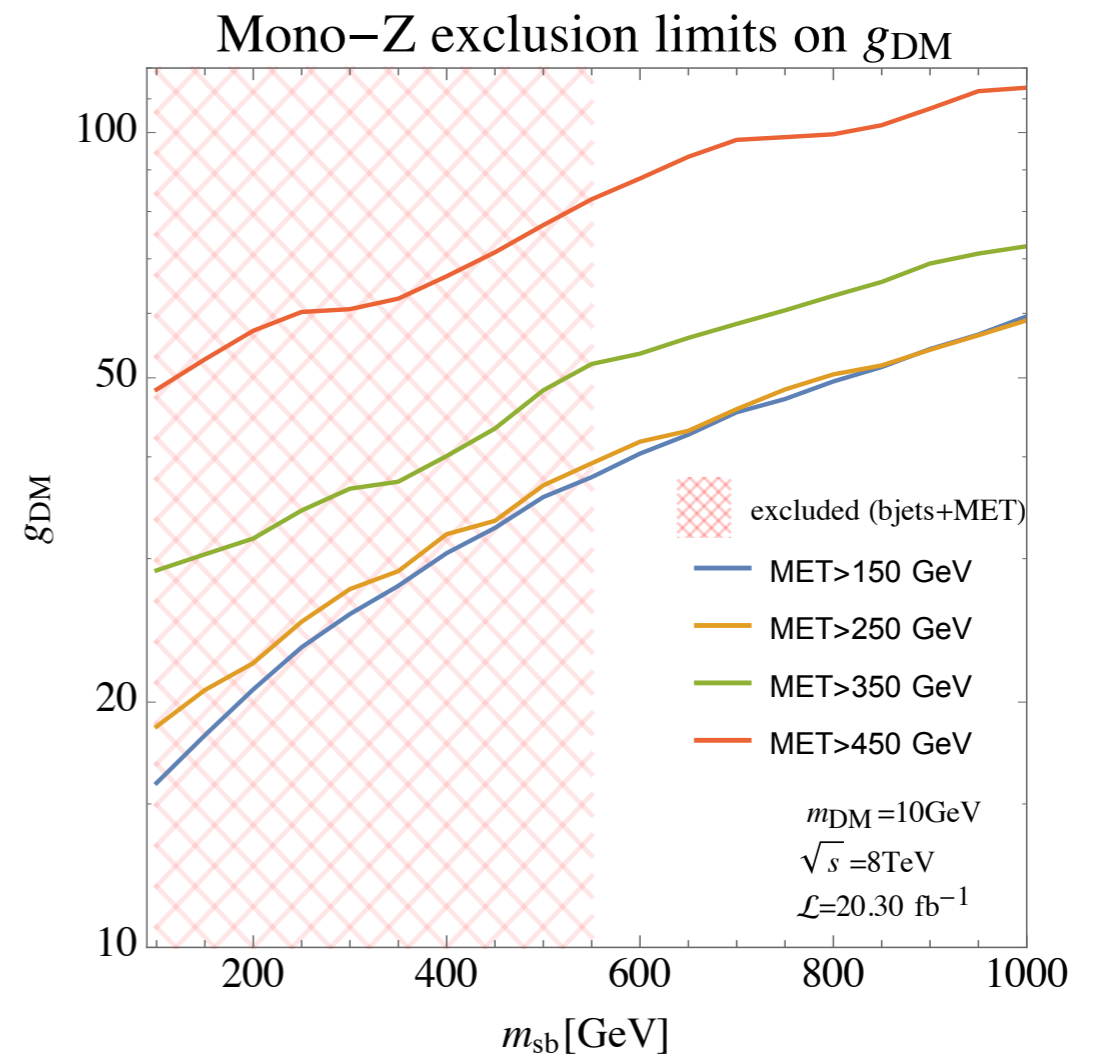
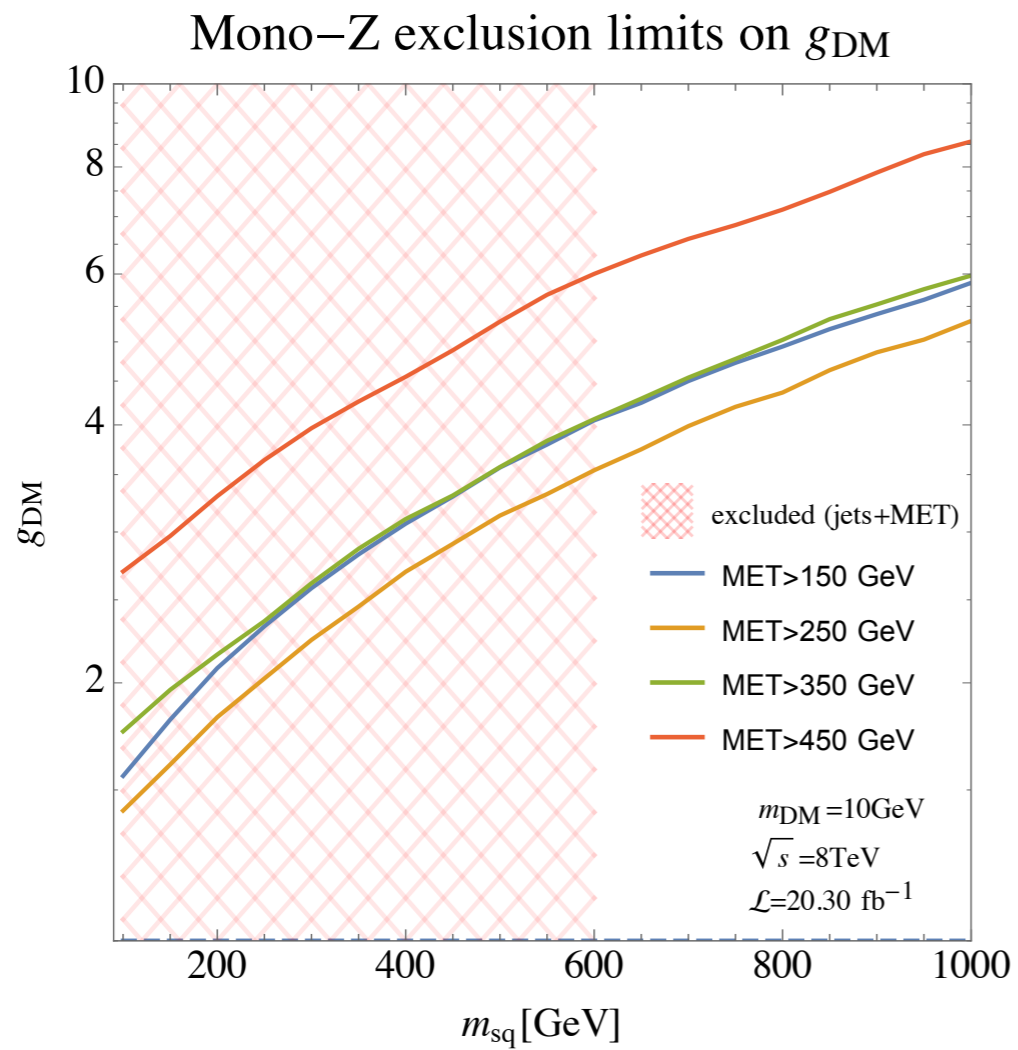


Dijet Constraints on $g_{qqZ'}$

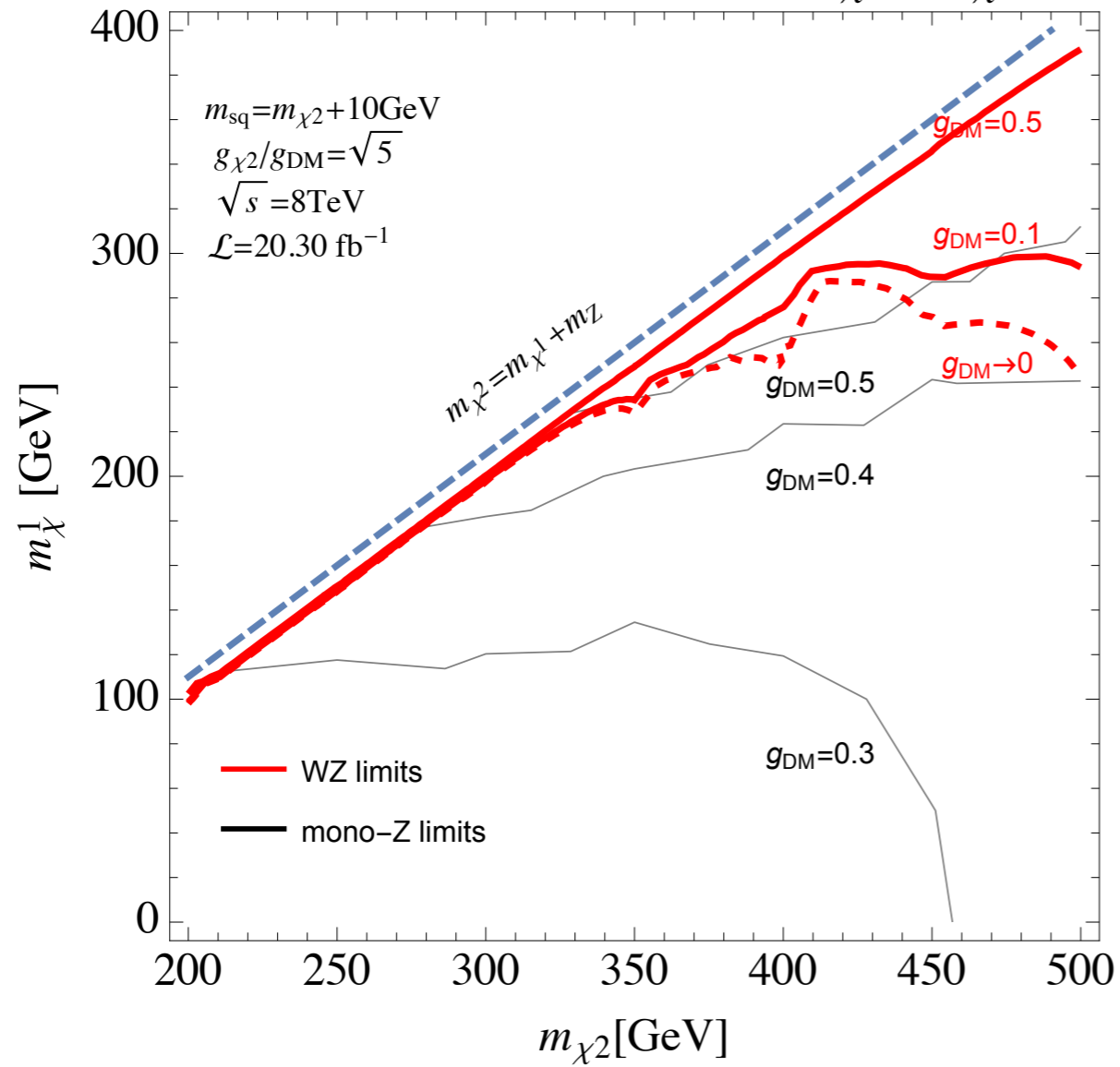


Dijet constraint

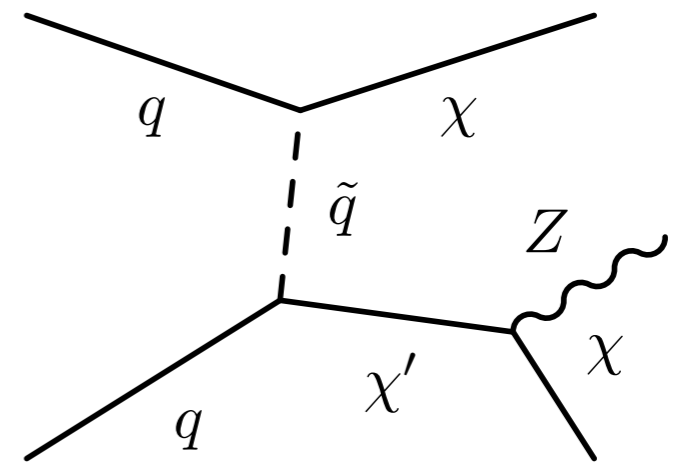
mono-Z for squark/sbottom model



Mono-Z exclusion limits on $m_{\chi_2} - m_{\chi_1}$ plane



mono-Z
(inelastic squark)



sbottom model (mono-b & b-jets + MET)

