

Dark Matter Searches at the LHC

Masahiro Morii
Harvard University

Dark Side of the Universe 2015
Kyoto, 14–18 Dec 2015

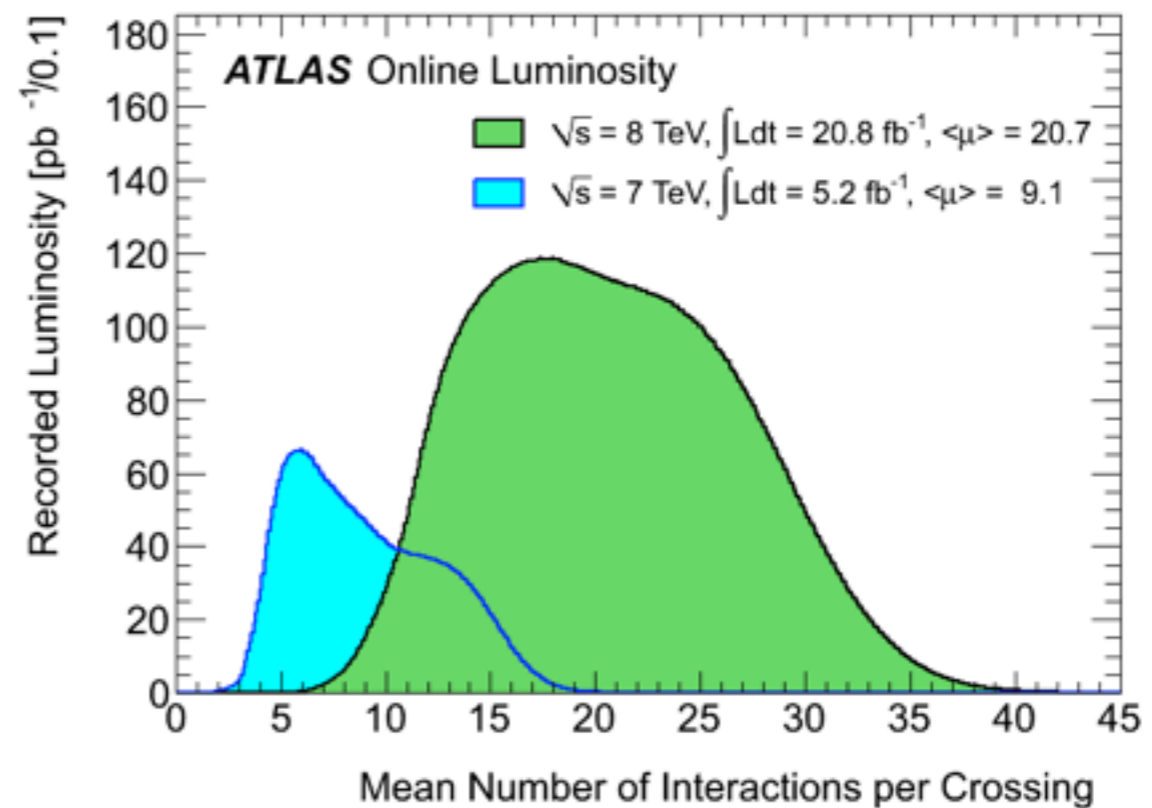
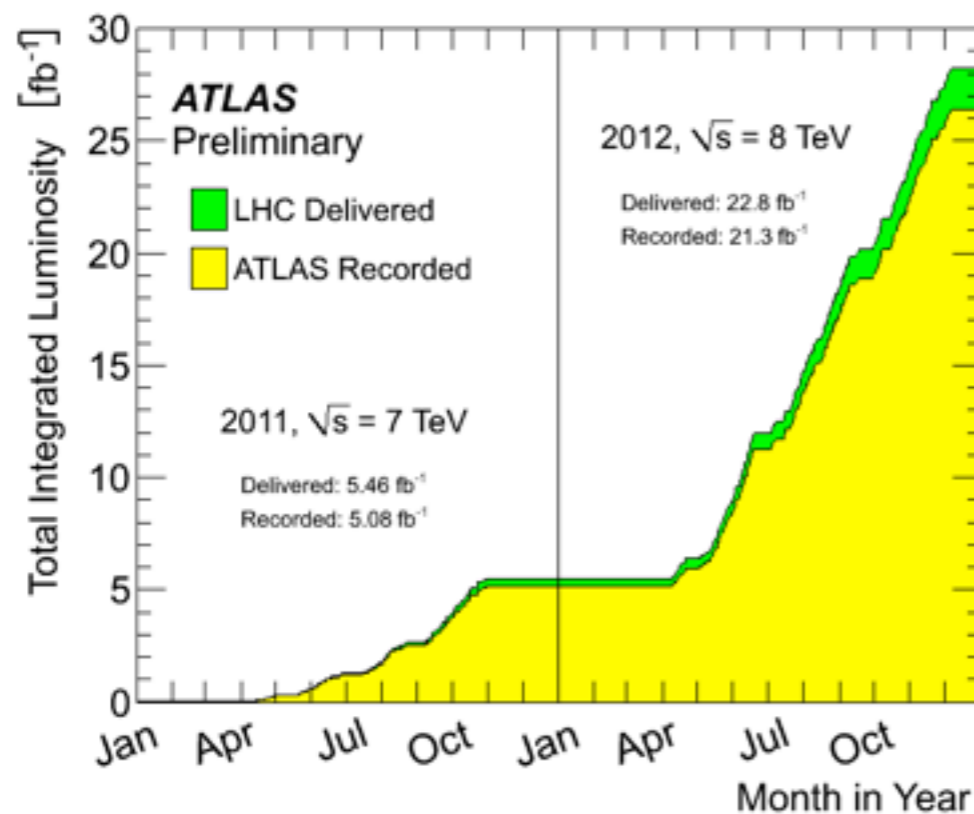
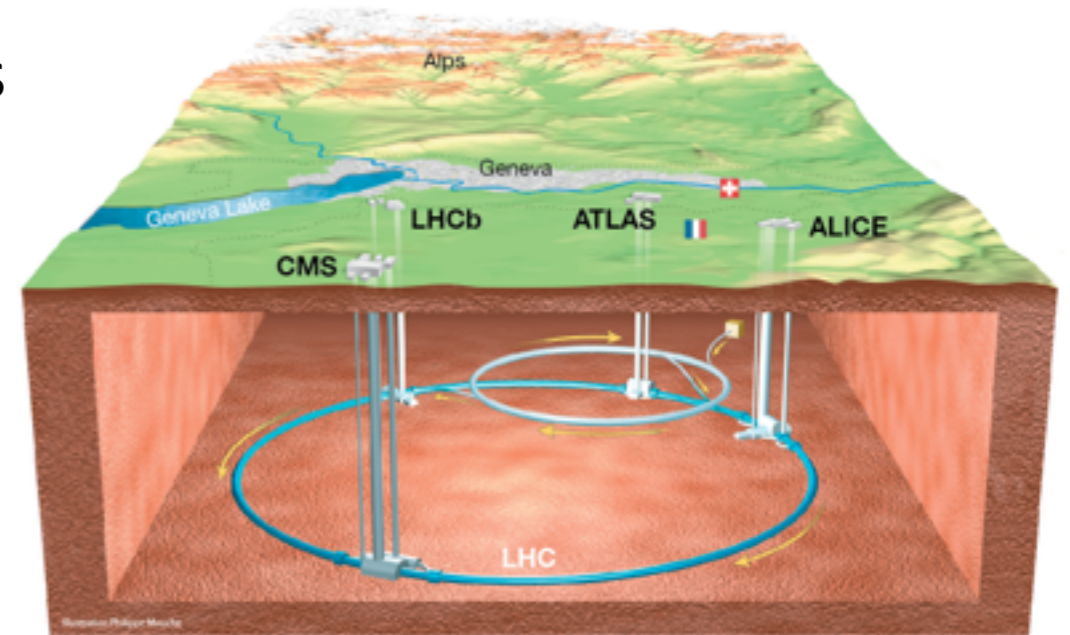
Large Hadron Collider

LHC collides protons with center-of-mass energies up to 14 TeV

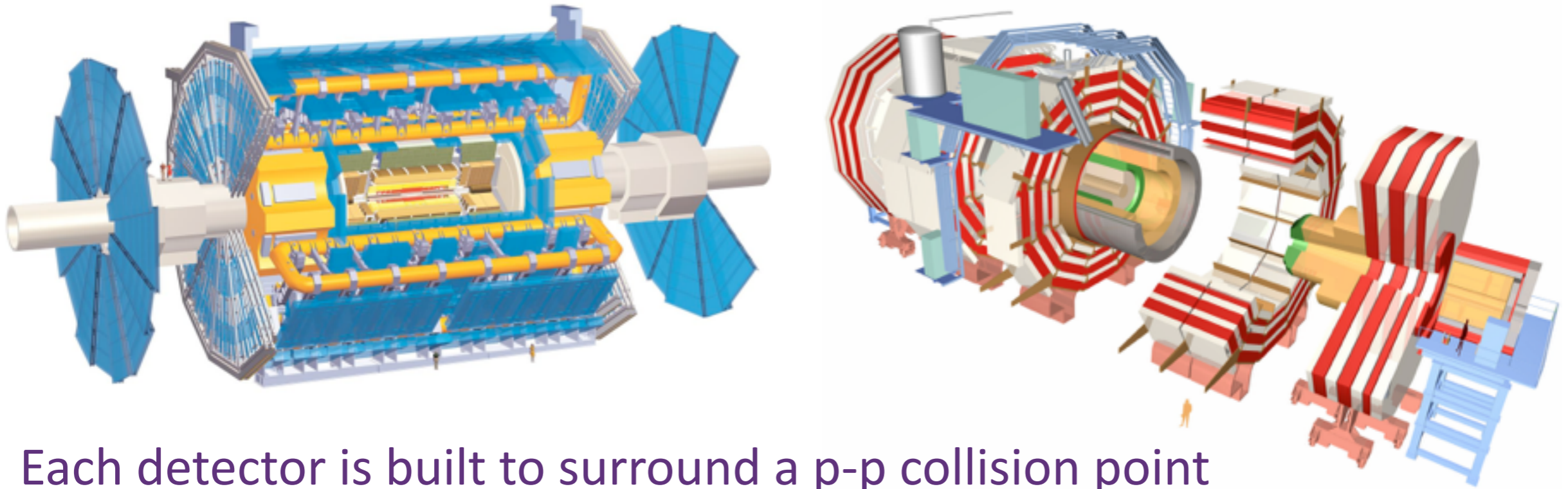
- Counter-circulating beams of >1000 bunches of $\sim 10^{11}$ protons
- Bunches cross every 25 ns (or 50 ns) producing up to 40 p-p collisions

Run 1: 7 TeV (2010–11) and 8 TeV (2012)

- All results presented today use 8 TeV data



ATLAS and CMS detectors

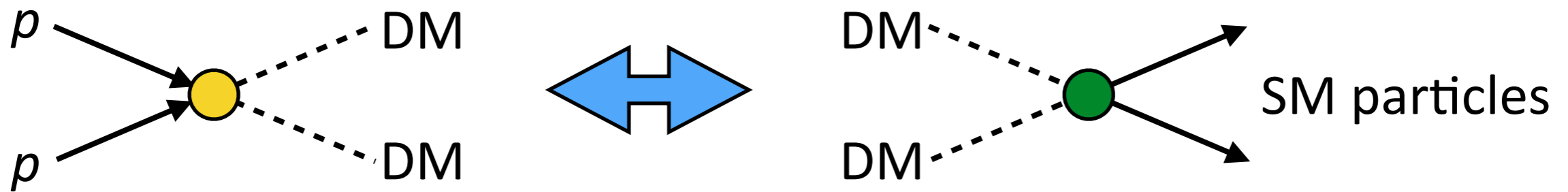


Each detector is built to surround a p-p collision point

- Near- 4π coverage to detect all collision products (photons, leptons, hadrons) **except for neutrinos and DM**
- Energy/momentum measurement of all particles \rightarrow Inference of undetected particle(s) from momentum conservation in the transverse plane
 - ▶ p_T^X = transverse momentum of X (particle or group of particles)
 - ▶ **MET** = negative sum of p_T of all detected particles in an event
- All DM searches look for events with large values of MET

DM production at the LHC

LHC (hopefully) pair-produce DM particles in p-p collisions

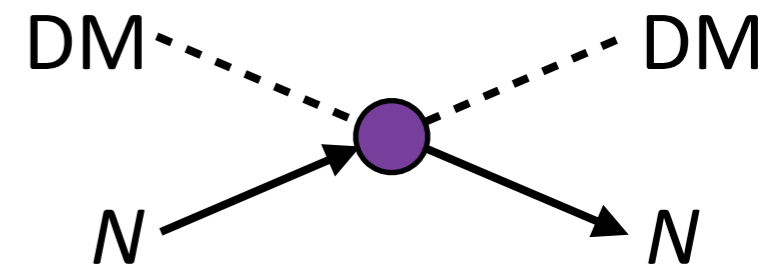


Production cross-section $\sigma(pp \rightarrow \chi\bar{\chi})$ is **related** to the annihilation cross-section $\sigma(\chi\bar{\chi} \rightarrow \text{anything})$, but it's **model-dependent**

- Protons are made (mostly) of light quarks
- It also differs from the nucleon-DM elastic scattering cross-section measured by direct detection experiments

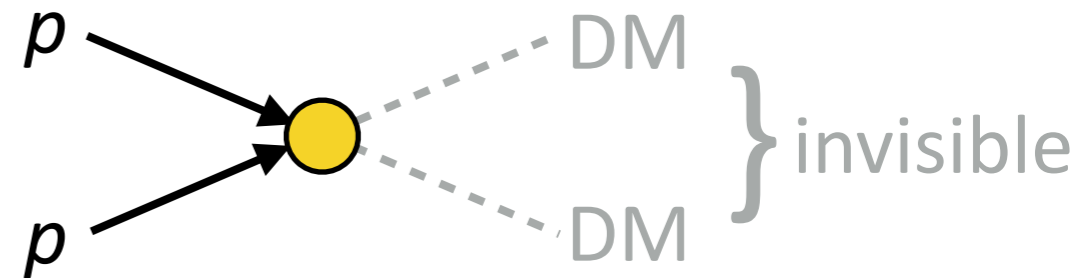
Interpretation gets complicated

- Pay attention to theory talks after the coffee break...



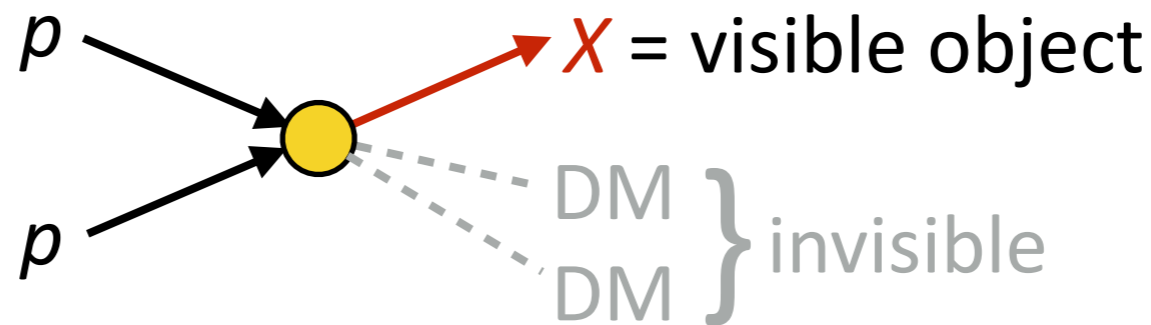
DM detection at the LHC

Experiments cannot see the final state with only DM particles



- How do we record and count such events?

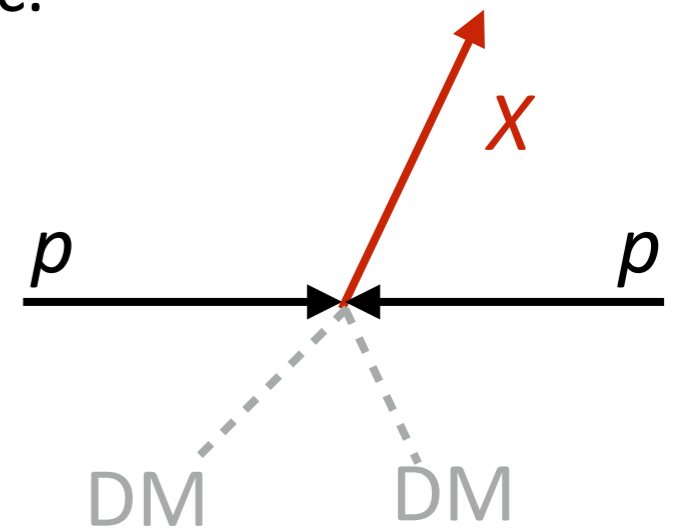
We need ≥ 1 extra particles produced with the DM particles



- We call it “mono- X ” with $X = \text{jet, photon, Higgs, } W, Z, \text{ etc.}$

Net momentum of the outgoing particles seems to violate momentum conservation \rightarrow “MET”

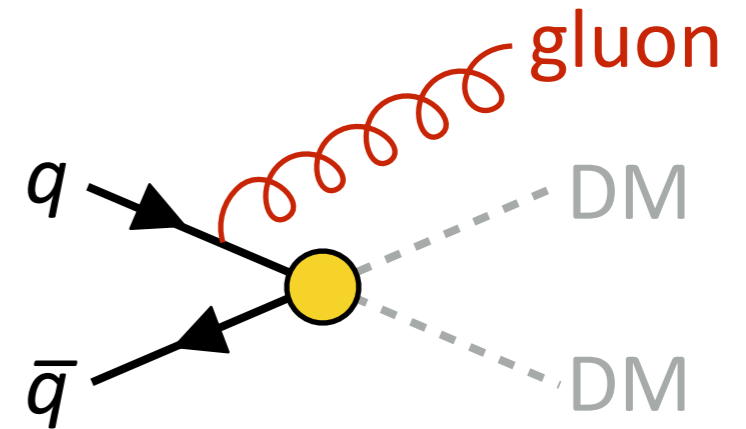
- Candidate events contain $X + \text{MET}$ and little else



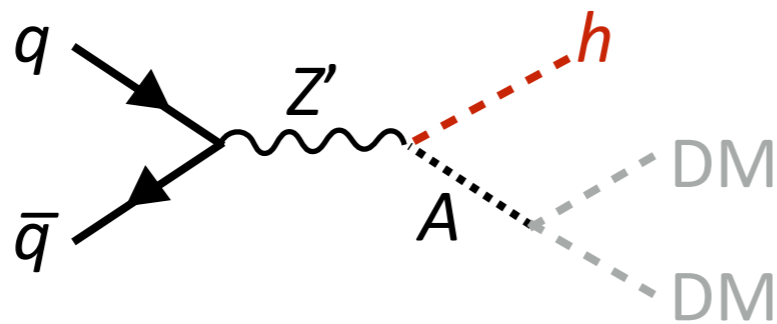
Where does “X” come from?

X may be radiated by the incoming quarks

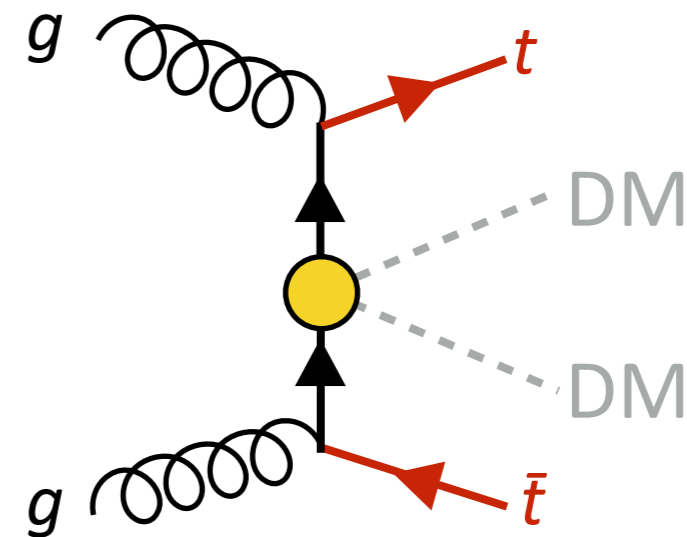
- X can be a gluon, photon, W or Z boson
- If $q\bar{q} \rightarrow \chi\bar{\chi}$ happens, $q\bar{q} \rightarrow \chi\bar{\chi} + (g/\gamma/W/Z)$ is guaranteed to happen
- If $X = \text{gluon}$, it turns into a “jet” and creates a “mono-jet” signature



X may be more closely connected to DM production, e.g.



- Cross-section depends on the exact model and the couplings
- May be large even if $q\bar{q} \rightarrow \chi\bar{\chi}$ is small



Mono- X results from Run 1

- All results use full statistics ($\sim 20 \text{ fb}^{-1}$) of 8 TeV collision data

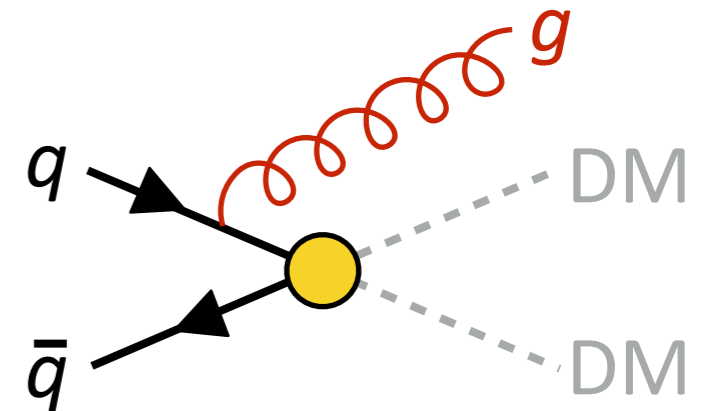
"X" in mono-X	Expt.	Reference	arXiv
jet	ATLAS	Eur. Phys. J. C 75 (2015) 299	1502.01518
	CMS	Eur. Phys. J. C 75 (2015) 235	1408.3583
≥ 2 jets	CMS	CMS-PAS-EXO-14-004	
photon	ATLAS	Phys. Rev. D 91 (2015) 012008	1411.1559
	CMS	submitted to Phys. Lett. B	1410.8812
$W (\rightarrow \ell\nu)$	ATLAS	JHEP 09 (2014) 037	1407.7494
	CMS	Phys. Rev. D 91 (2015) 092005	1408.2745
$Z (\rightarrow \ell\ell)$	ATLAS	Phys. Rev. D 90 (2014) 012004	1404.0051
	CMS	submitted to Phys. Rev. D	1511.09375
W or $Z (\rightarrow q\bar{q})$	ATLAS	Phys. Rev. Lett. 112 (2014) 041802	1309.4017
	CMS	CMS-PAS-EXO-12-055	
$t\bar{t}$ (1 lepton)	CMS	JHEP 06 (2015) 121	1504.03198
$t\bar{t}$ (2 leptons)	CMS	CMS-PAS-B2G-13-004	
t or b -quark jet	ATLAS	Eur. Phys. J C75 (2015) 92	1410.4031
$H (\rightarrow \gamma\gamma)$	ATLAS	Phys. Rev. Lett. 115 (2015) 131801	1506.01081
$H (\rightarrow b\bar{b})$	ATLAS	submitted to Phys. Rev. D	1510.06218

Mono-jet search

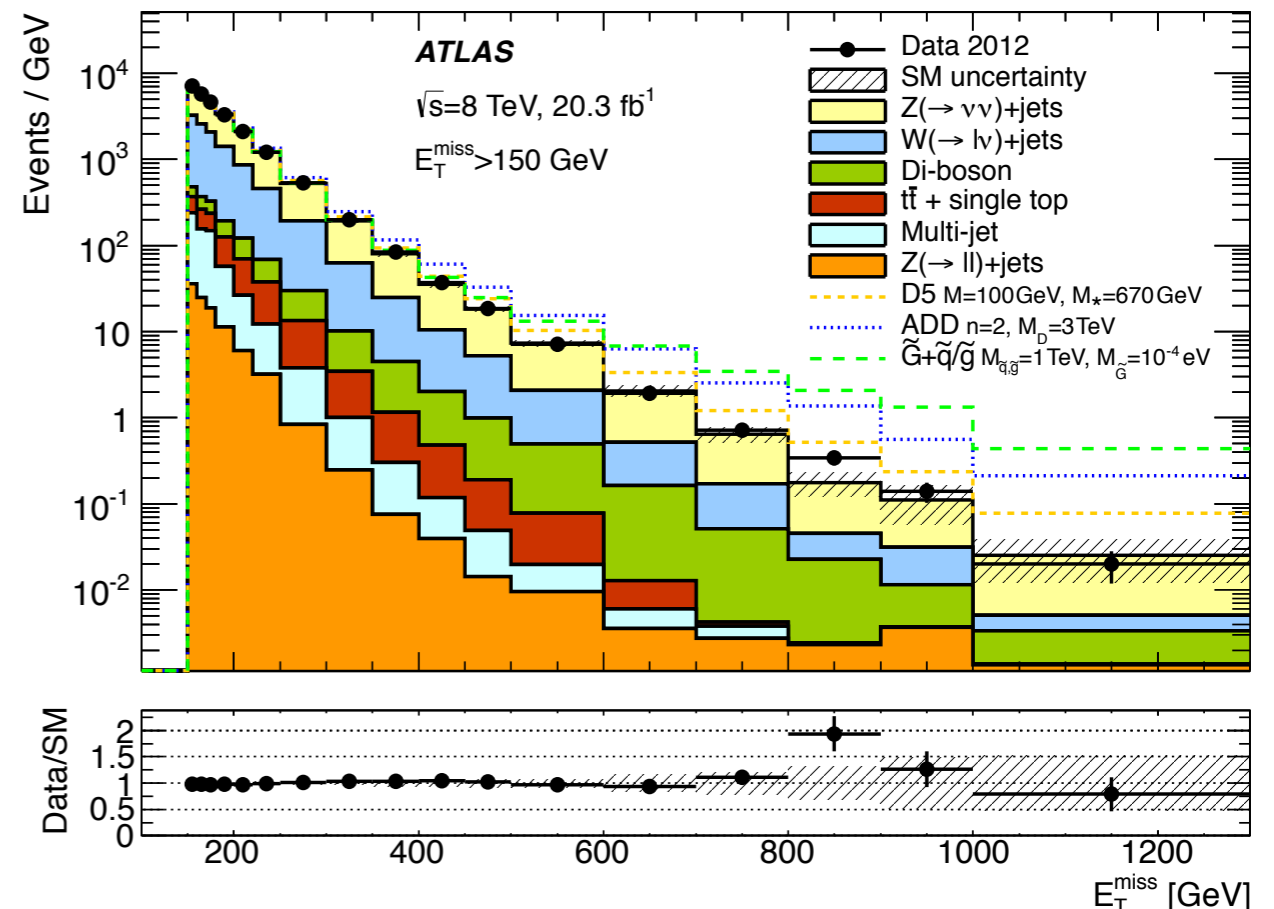
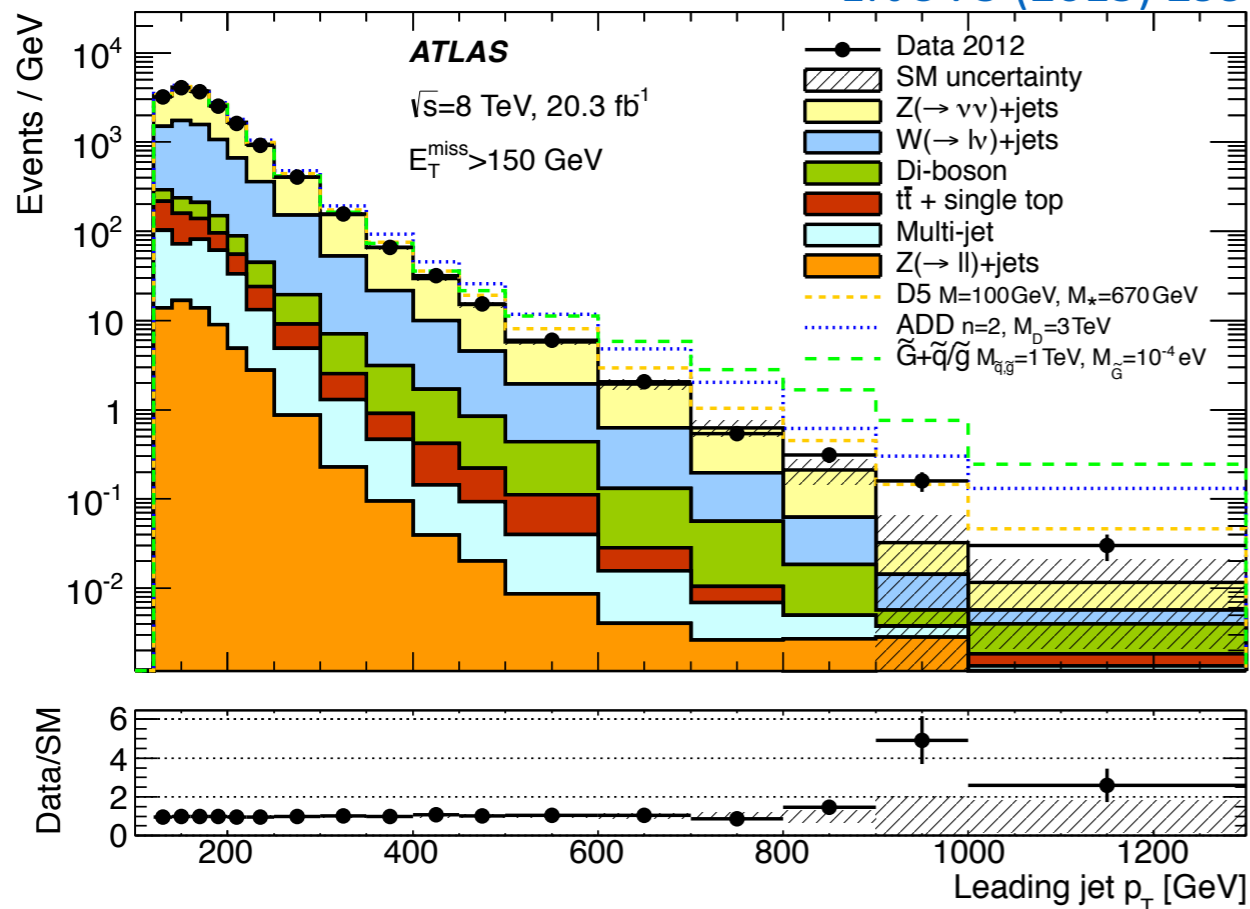
Mono-jet has the highest sensitivity to “generic” WIMP-like DM

ATLAS mono-jet search [EPJC 75:299]

- Select events with ≥ 1 jet with $p_T^{\text{jet}} > 125$ GeV and no leptons
- Require $\text{MET} > 150$ GeV and $p_T^{\text{jet}}/\text{MET} > 0.5$



EPJC 75 (2015) 299



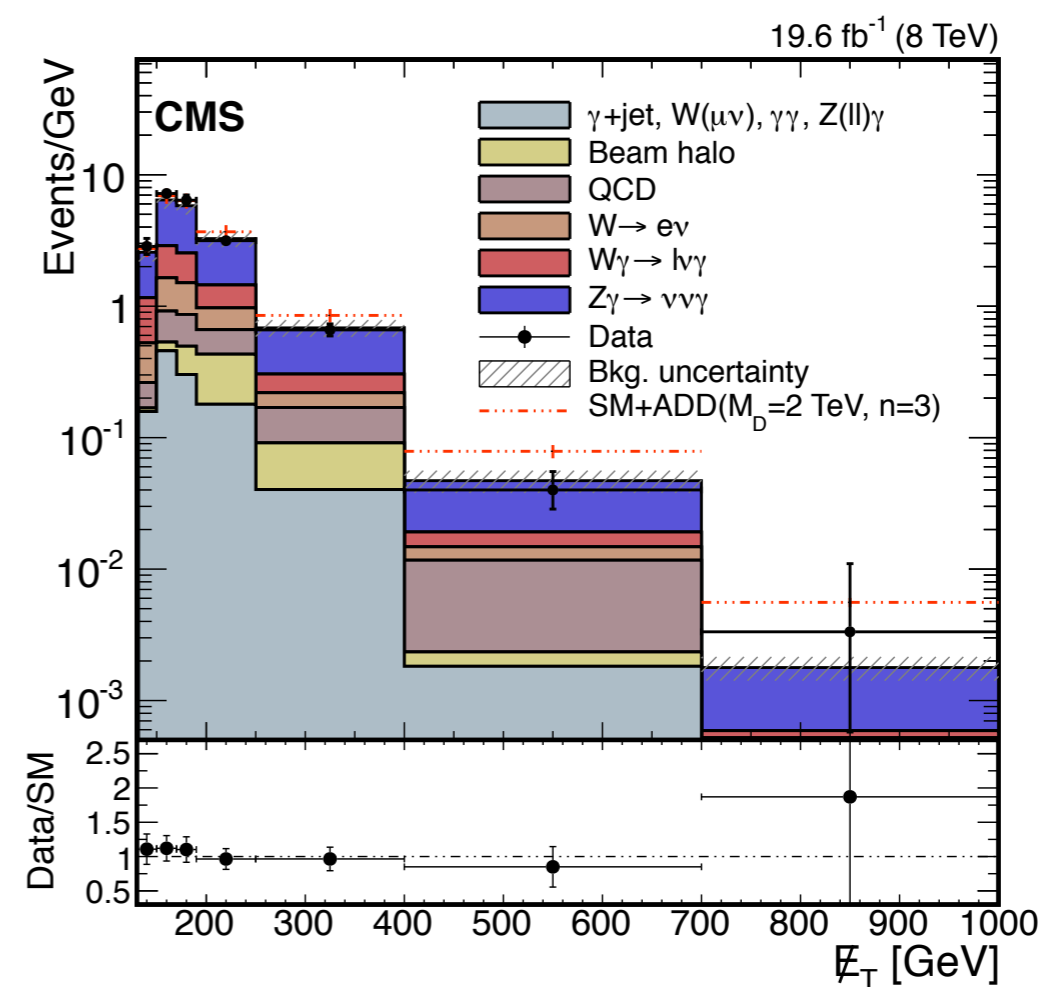
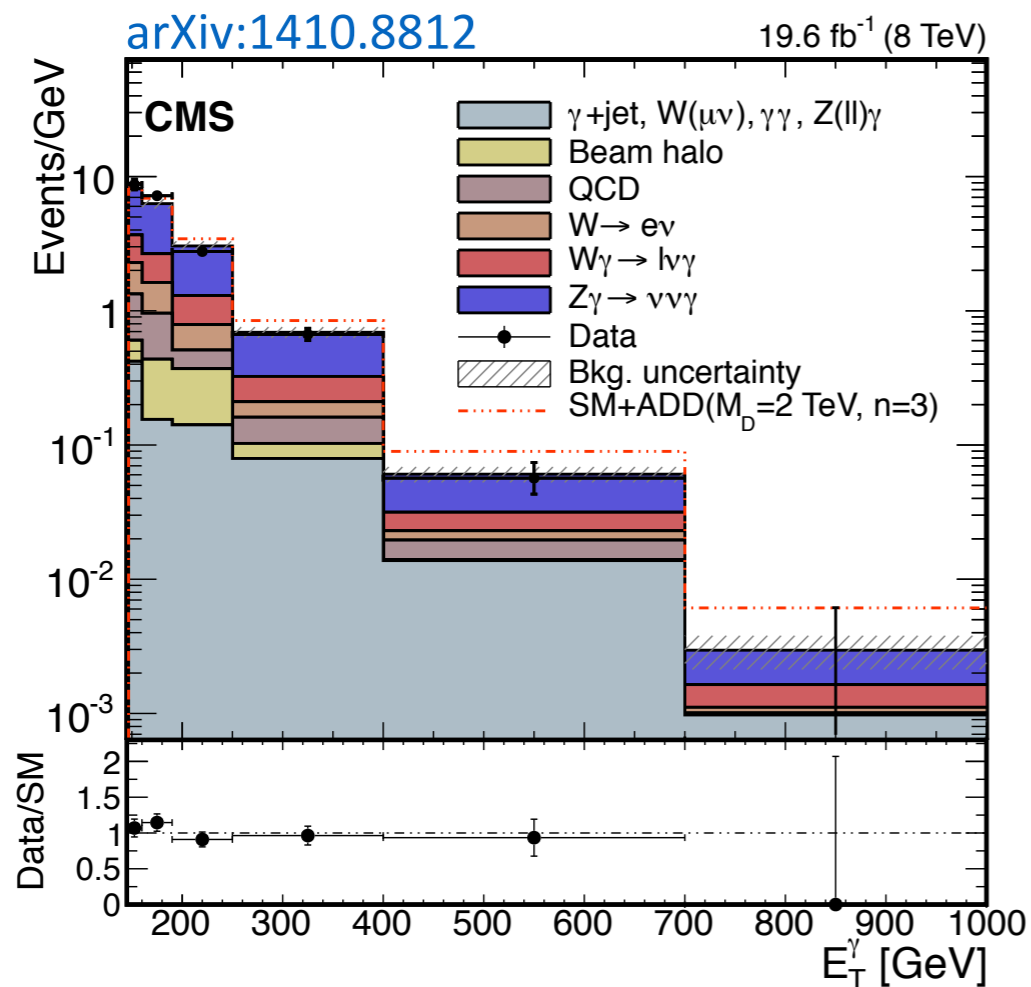
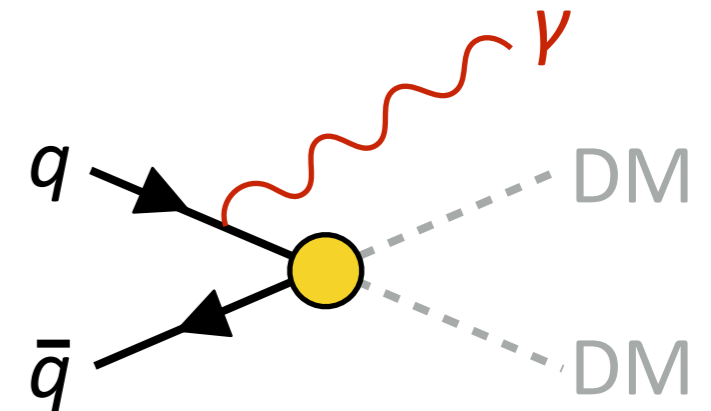
Mono-photon search

Mono-photon is similar to mono-jet

- Smaller signal cross-section, but smaller background

CMS mono-photon search [arXiv:1410.8812]

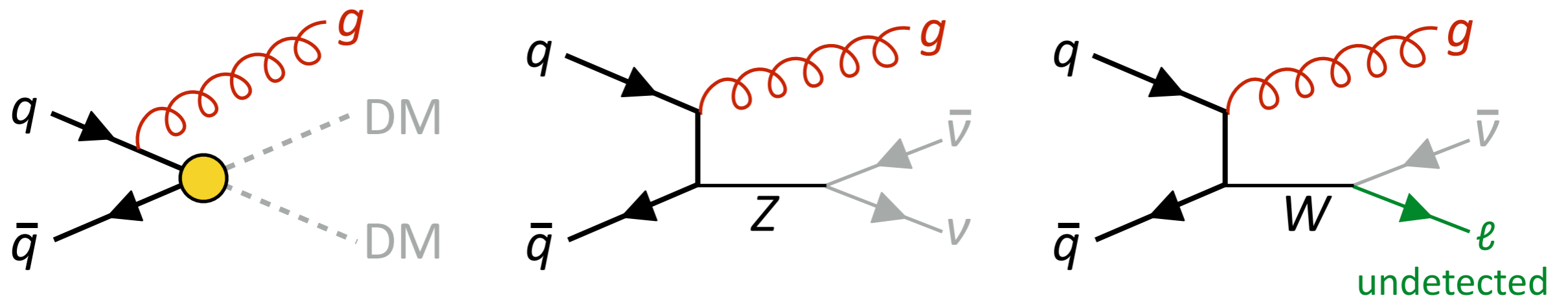
- One photon with $E_T^\gamma > 145$ GeV, and ≤ 1 jet
- Require MET > 140 GeV



Background

Background in mono-jet/photon searches are dominated by

- $Z (\rightarrow \nu\bar{\nu}) + \text{jet/photon}$ (irreducible)
- $W (\rightarrow \ell\nu) + \text{jet/photon}$ where the lepton escaped detection



Estimate them using control samples in data

- $Z (\rightarrow \ell\ell)$ and $W (\rightarrow \ell\nu) + \text{jet/photon}$ where the leptons were detected
 - ▶ NB: $\text{BR}(Z \rightarrow \ell\ell) \ll \text{BR}(Z \rightarrow \nu\bar{\nu})$
- Smaller background sources are estimated with simulation

Instrumental background has to be controlled, too

- e.g., beam-induced noise in the calorimeter that looks like a jet

Interpretation — EFT

LHC measures $\sigma(pp \rightarrow X + \text{DM})$ in bins of MET

- How does it relate to the cosmic Dark Matter?

Interpretation requires theoretical assumptions

- What kind of DM (mass, spin, parity, interactions) is it?

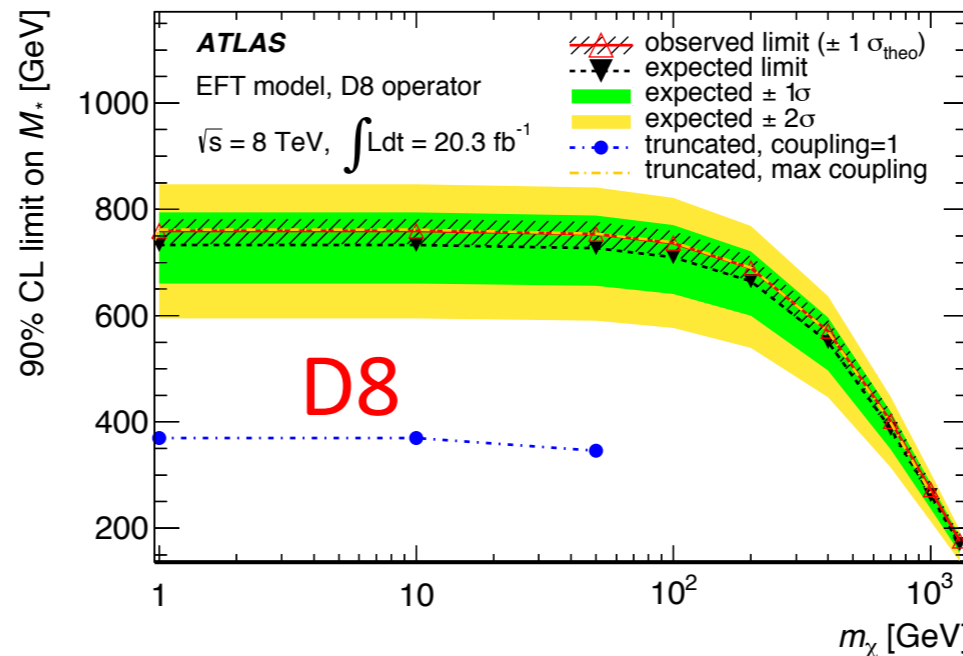
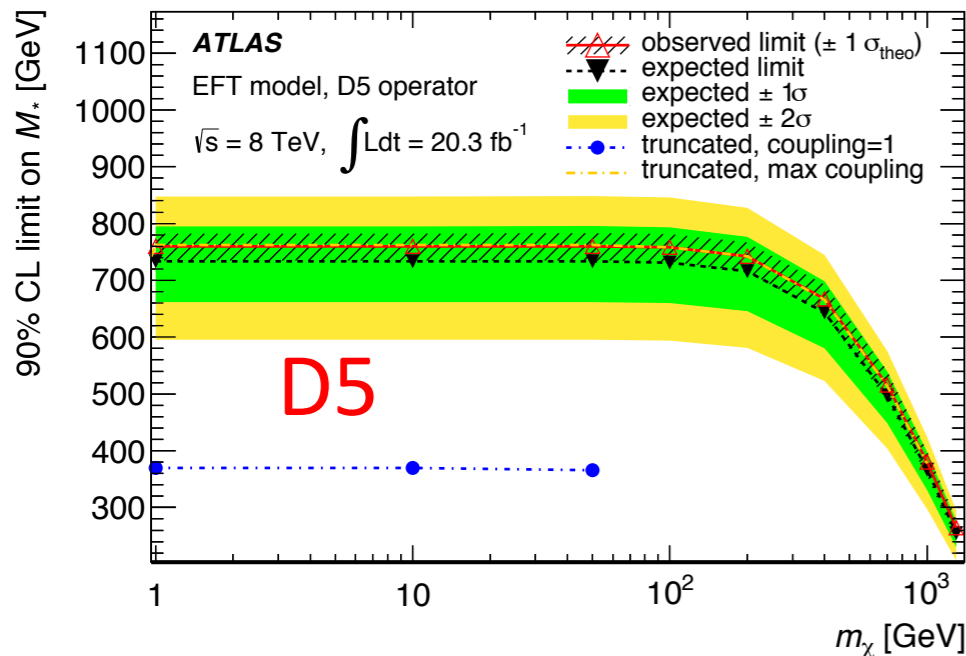
Effective Field Theory (EFT) has been our benchmark

- Constrain the scale M_* as a function of the DM mass m_χ

Goodman et al.,
PRD 82 (2010) 116010

Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

ATLAS mono-photon, PRD 91 (2015) 012008

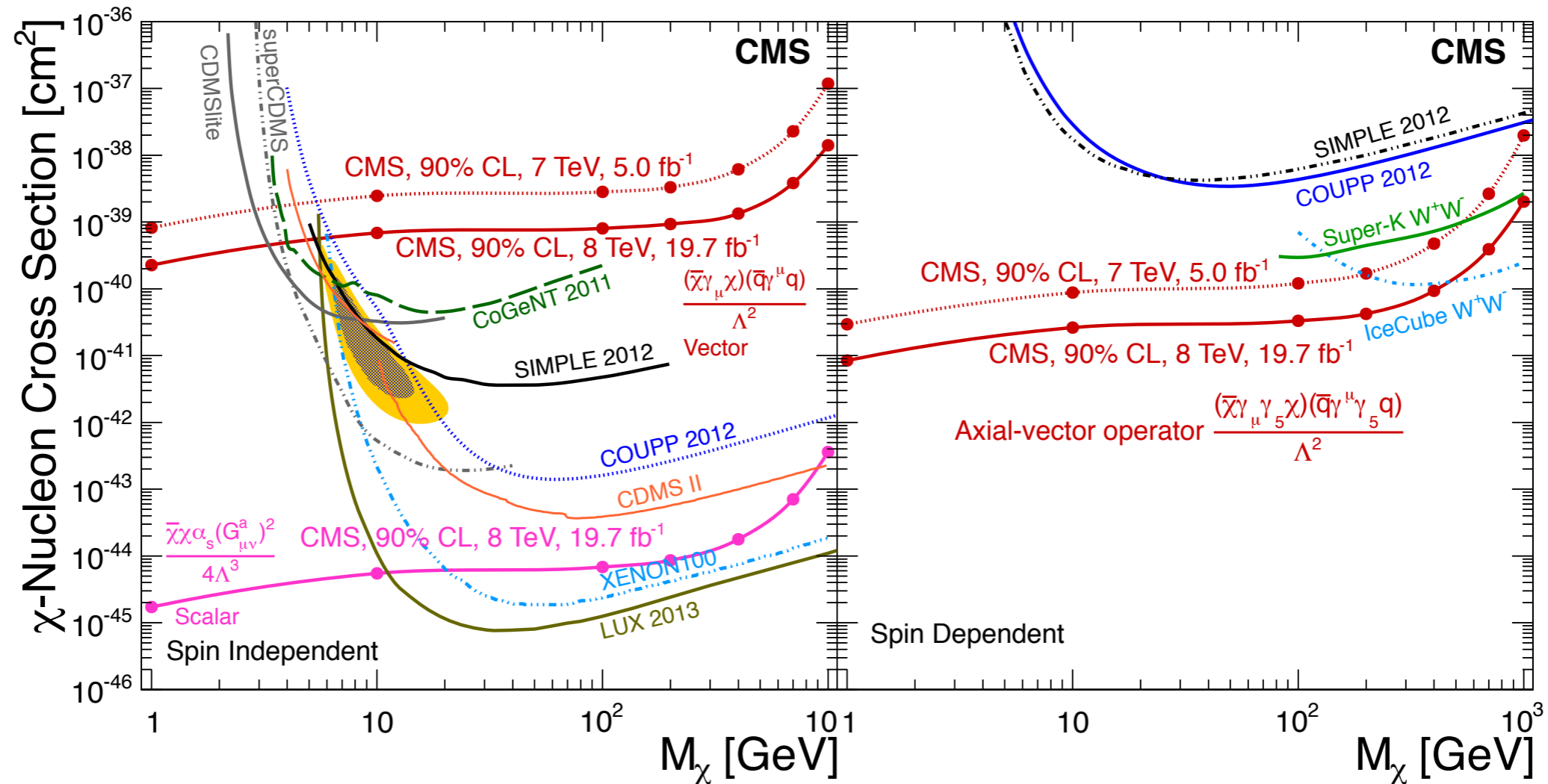


- NB: Validity of EFT is questionable at $M_* \lesssim 2m_\chi$

LHC v.s. direct-detection

Caution: comparisons are strongly model-dependent

- Ex) EFT-based results from CMS mono-jet, EPJC 75 (2015) 235



- Limits for the vector (D5) and scalar (D11) operators differ by $\sim 10^5$
- LHC is more sensitive to light DM, and to spin-dependent interactions

Interpretation — Simplified model

Validity of EFT has been a recurring question

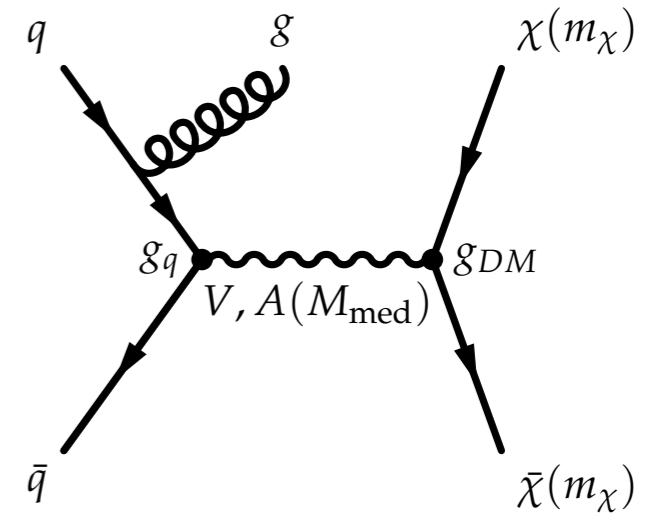
- What if the mediator is light (a few 100 GeV)?

ATLAS/CMS DM Forum report [arXiv:1507.00966]

- Recommends **simplified models** for use in new measurements from the LHC
- Mono-jet results, for example, will be presented as a bound on the coupling ($g_q \cdot g_{DM}$) on the M_{med} -v.s.- m_χ plane

Recent ATLAS/CMS papers report both EFT and simplified model results

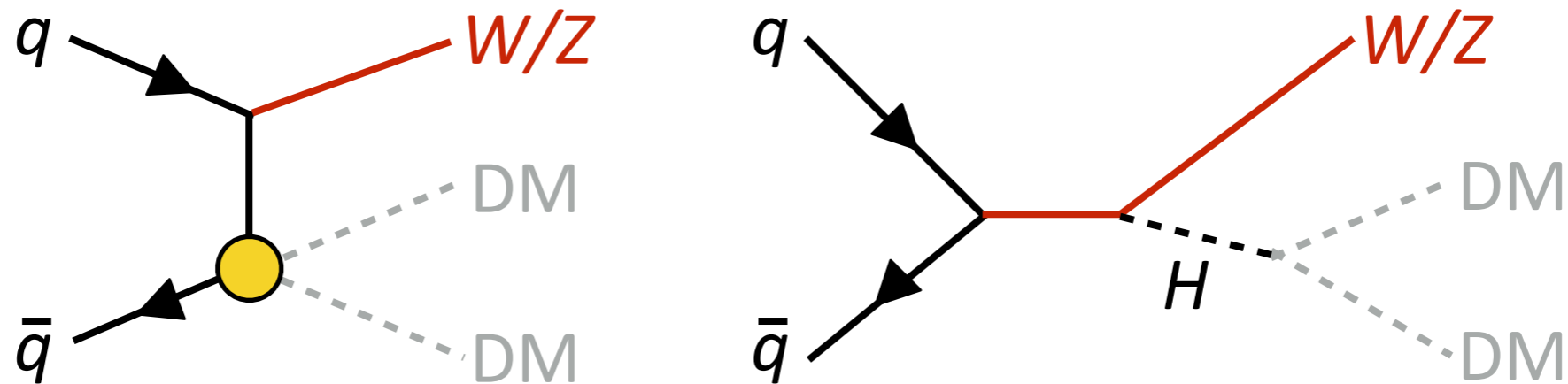
Mono-jet example



$$\mathcal{L}_{\text{vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$
$$\mathcal{L}_{\text{axial-vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma^5 q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi.$$

Mono- W/Z

Mono- W/Z can be produced in the same way as mono-jet/photon, but other possibilities exist:



- “Higgs-strahlung” (right) probes Higgs-DM coupling, or invisible Higgs decays in general

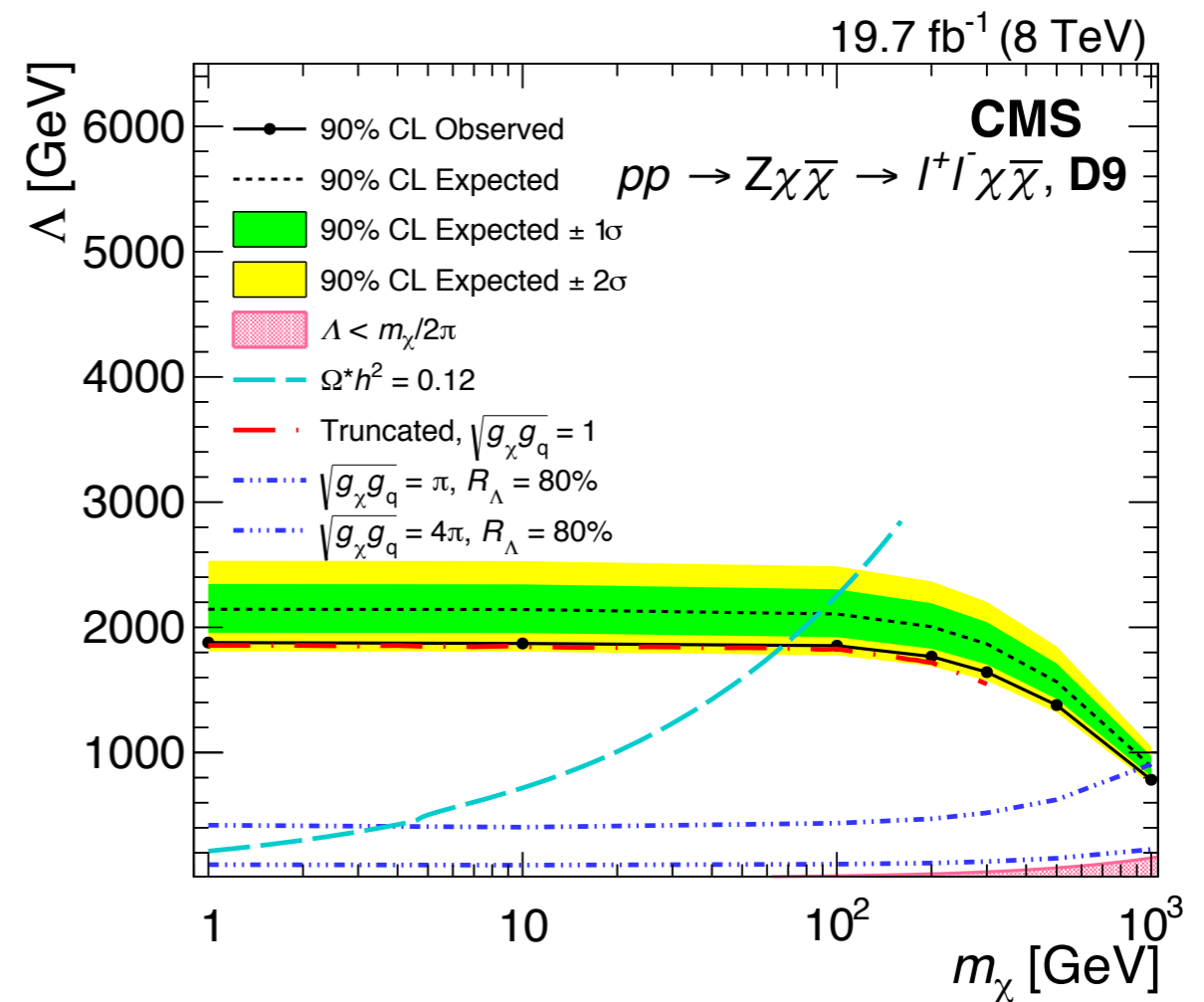
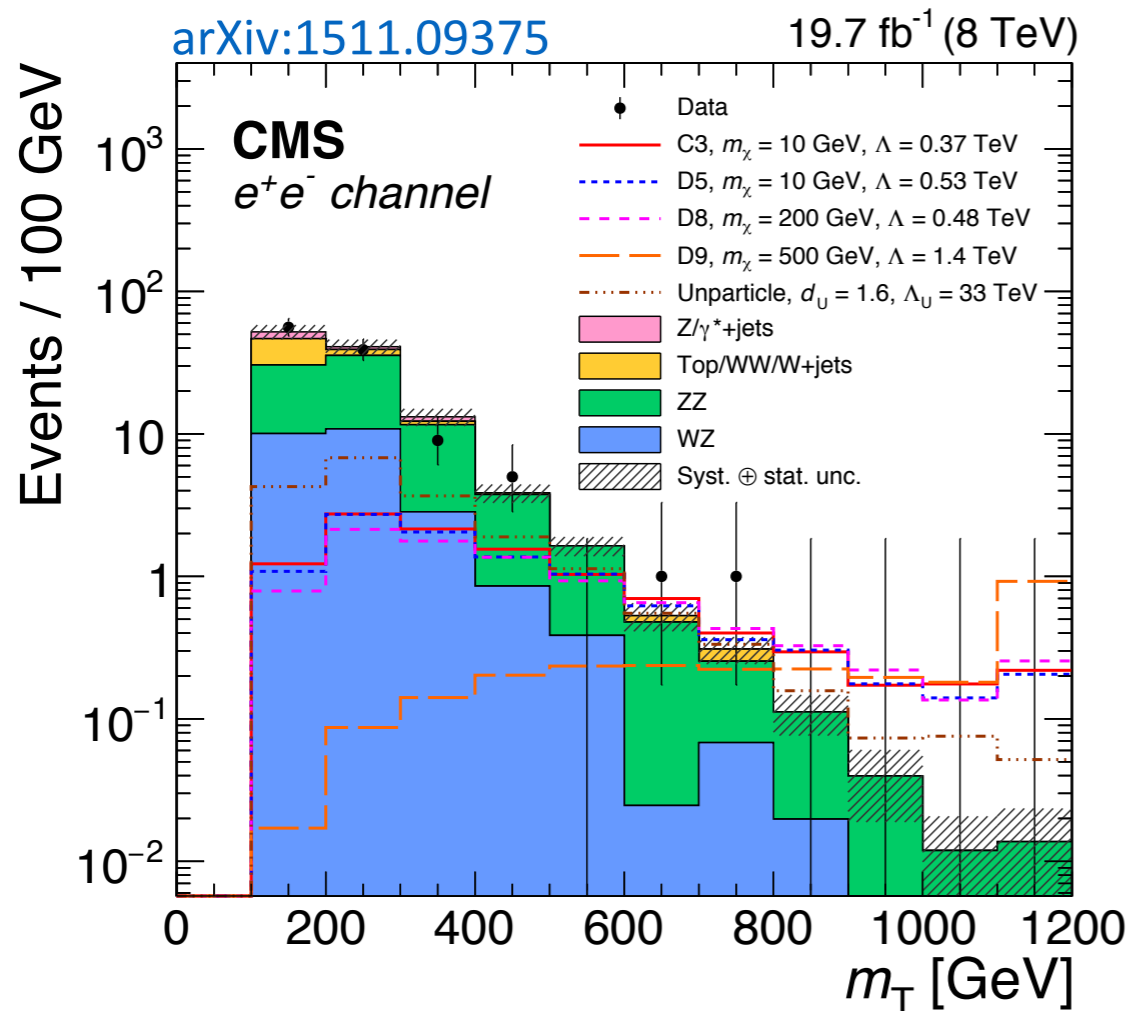
ATLAS/CMS can detect W/Z in two ways

- Leptonic final states $W \rightarrow \ell\nu$ and $Z \rightarrow \ell\ell$
 - ▶ Small background. But small branching fractions
- Hadronic final states $W/Z \rightarrow q\bar{q}$ (2 jets)
 - ▶ Large branching fractions. But large background

Leptonic mono-Z

CMS leptonic mono-Z [arXiv:1511.09375]

- $Z \rightarrow ee$ or $\mu\mu$ with $p_T^\ell > 20$ GeV, $|m_{\ell\ell} - m_Z| < 10$ GeV, and $p_T^{\ell\ell} > 50$ GeV
- MET > 80 GeV. MET is back-to-back and balanced with $p_T^{\ell\ell}$
- At most 1 jet and zero b -tagged jet (to suppress $t\bar{t}$ background)

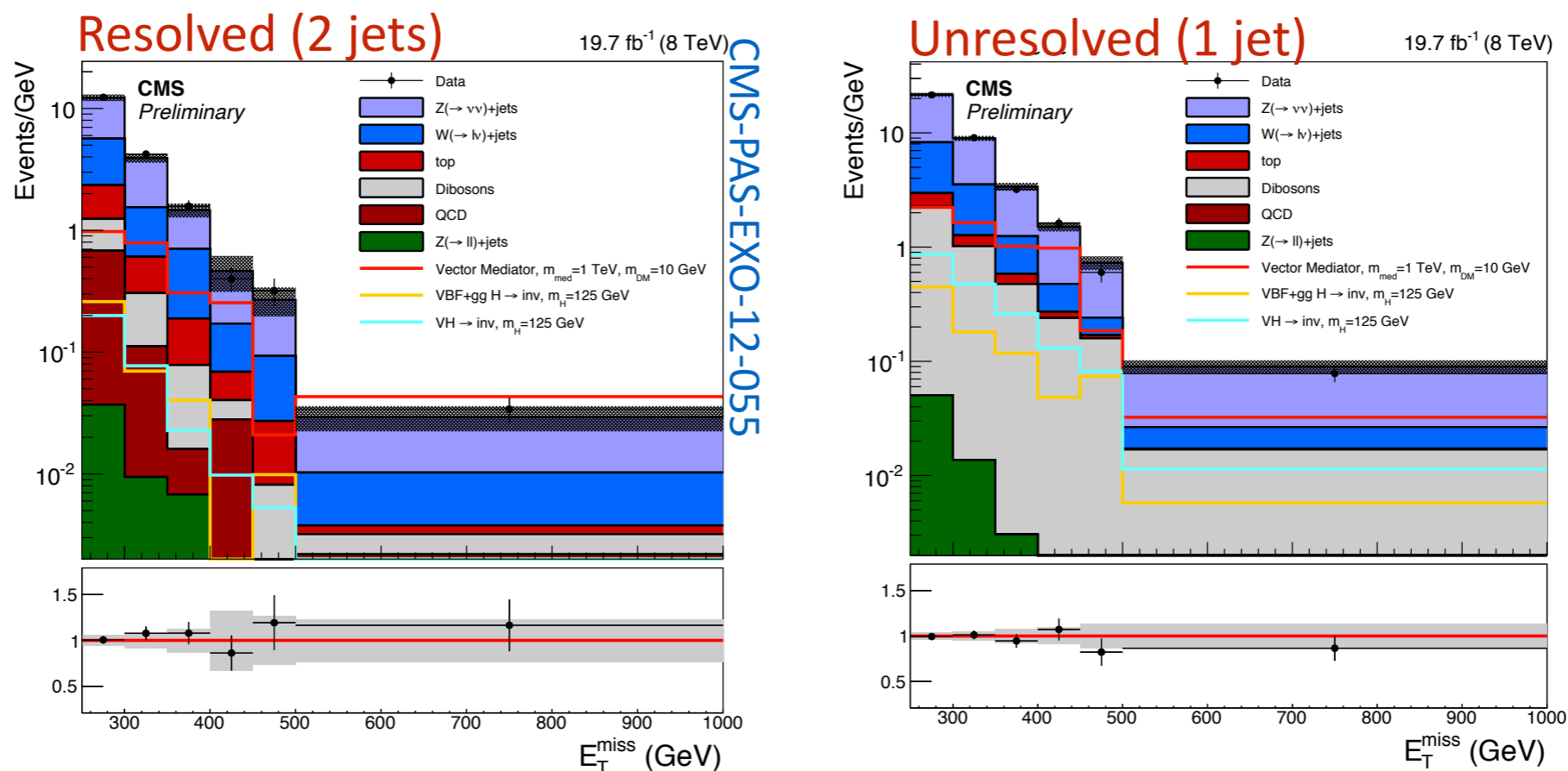


- Background is small, dominated by WZ and ZZ diboson production

Hadronic mono- W/Z

CMS hadronic mono- W/Z [CMS-PAS-EXO-12-055]

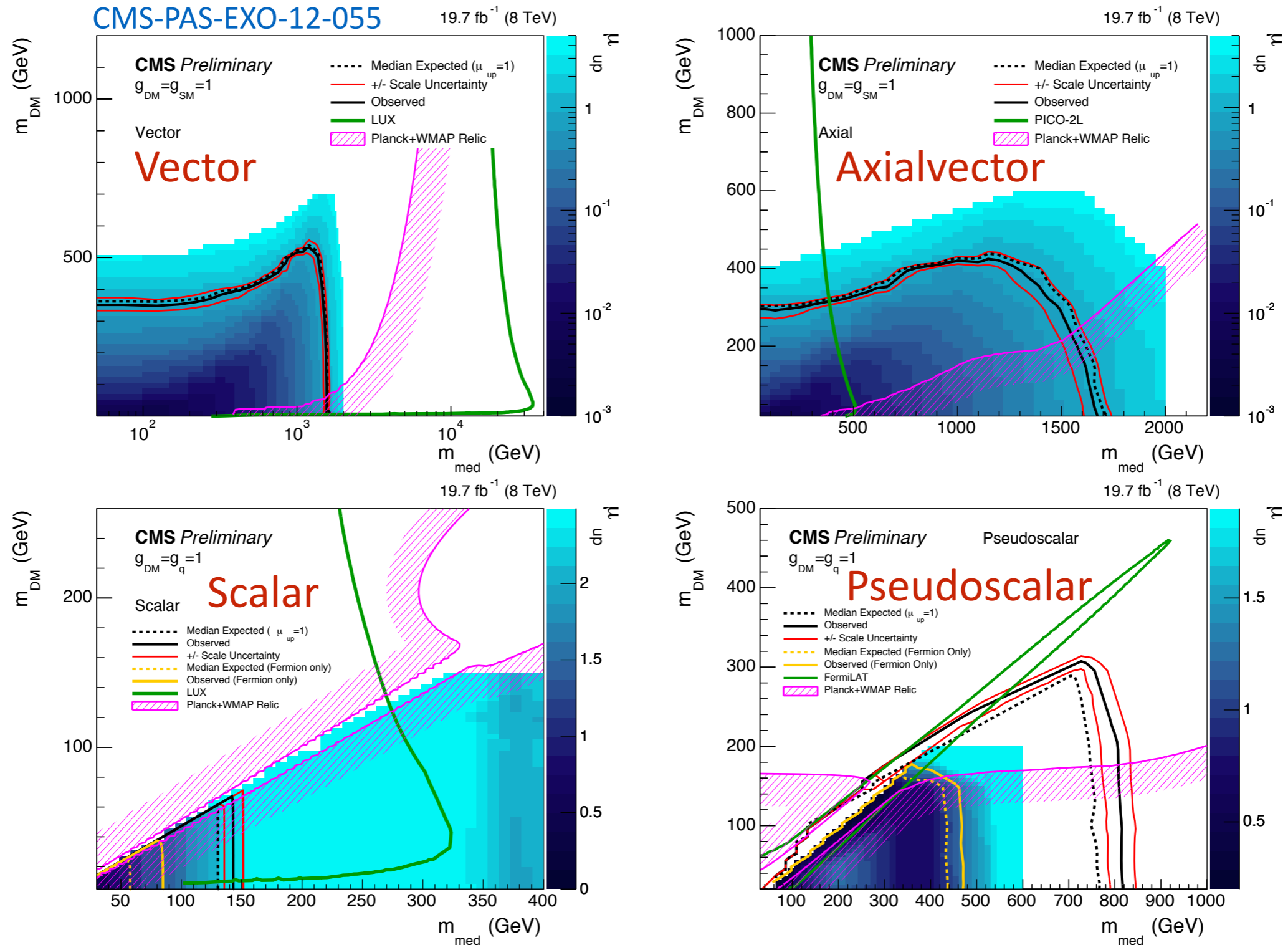
- Two jets may be detected as resolved (= 2 jets) or unresolved (= 1 jet) depending on the boost of the W/Z
 - ▶ Unresolved case requires analysis of sub-jet structure
- $60 < m_{jj} < 110$ GeV, ≤ 1 extra jet, and MET > 250 GeV



- Analysis also uses mono-jet events with $p_T^{\text{jet}} > 150$ GeV and MET > 200 GeV

Hadronic mono-W/Z

Simplified-model interpretation for different mediators



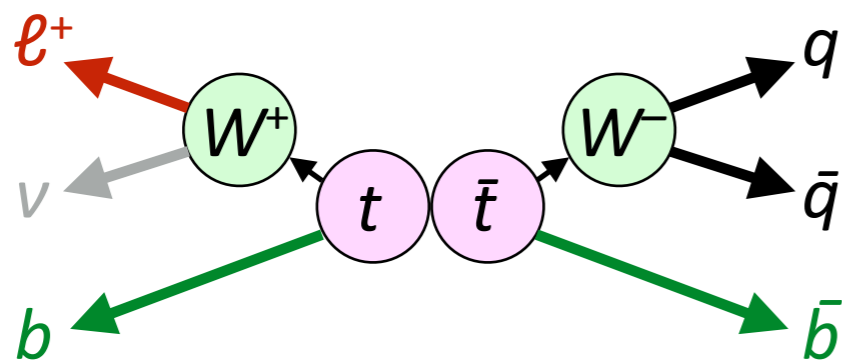
Mono(?))- $t\bar{t}$

Scalar interaction (D1) between DM and quarks is $L_{\text{int}} = \frac{m_q}{M_*^3} \bar{q}q\bar{\chi}\chi$

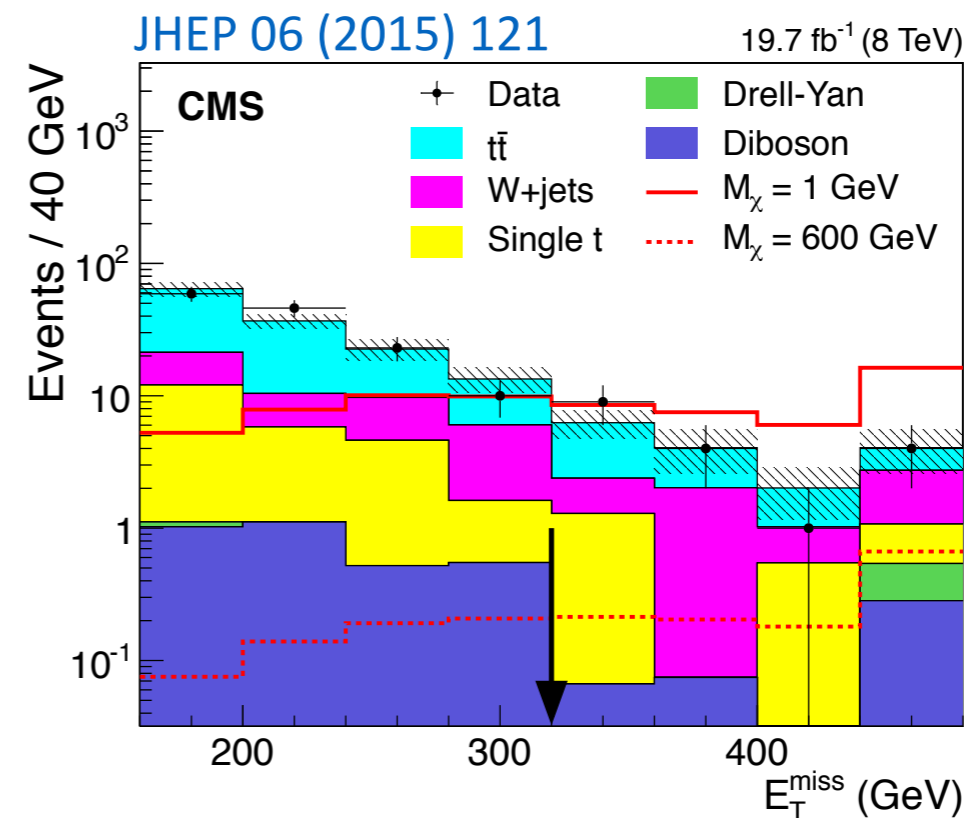
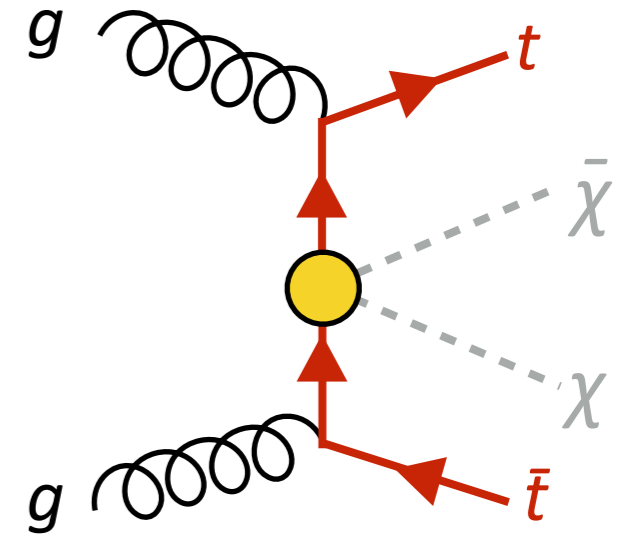
- Top quark is favored \rightarrow Search for MET + $t\bar{t}$ pair

CMS $t\bar{t}$ + MET search [JHEP 06 (2015) 121]

- Reconstruct $t(\rightarrow \ell\nu b)\bar{t}(\rightarrow q\bar{q}\bar{b})$
- Require 1 lepton, ≥ 3 jets of which ≥ 1 is b -tagged

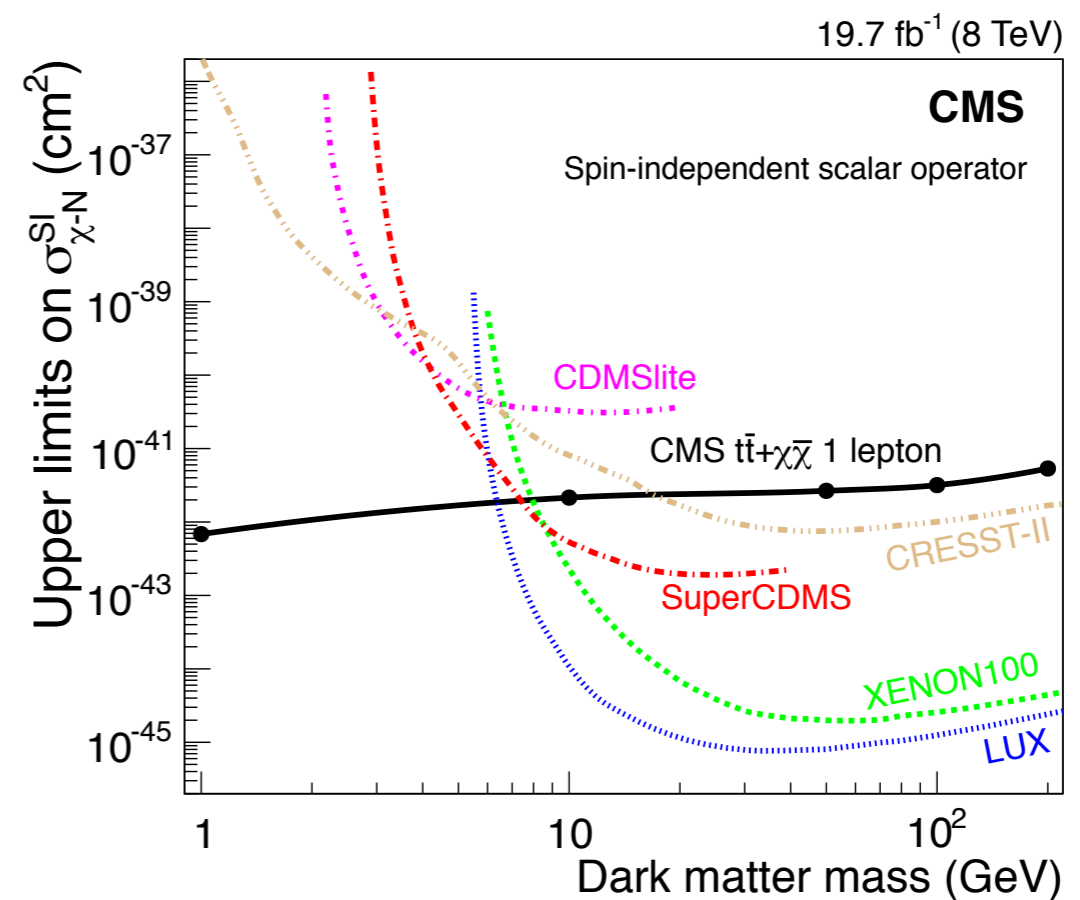
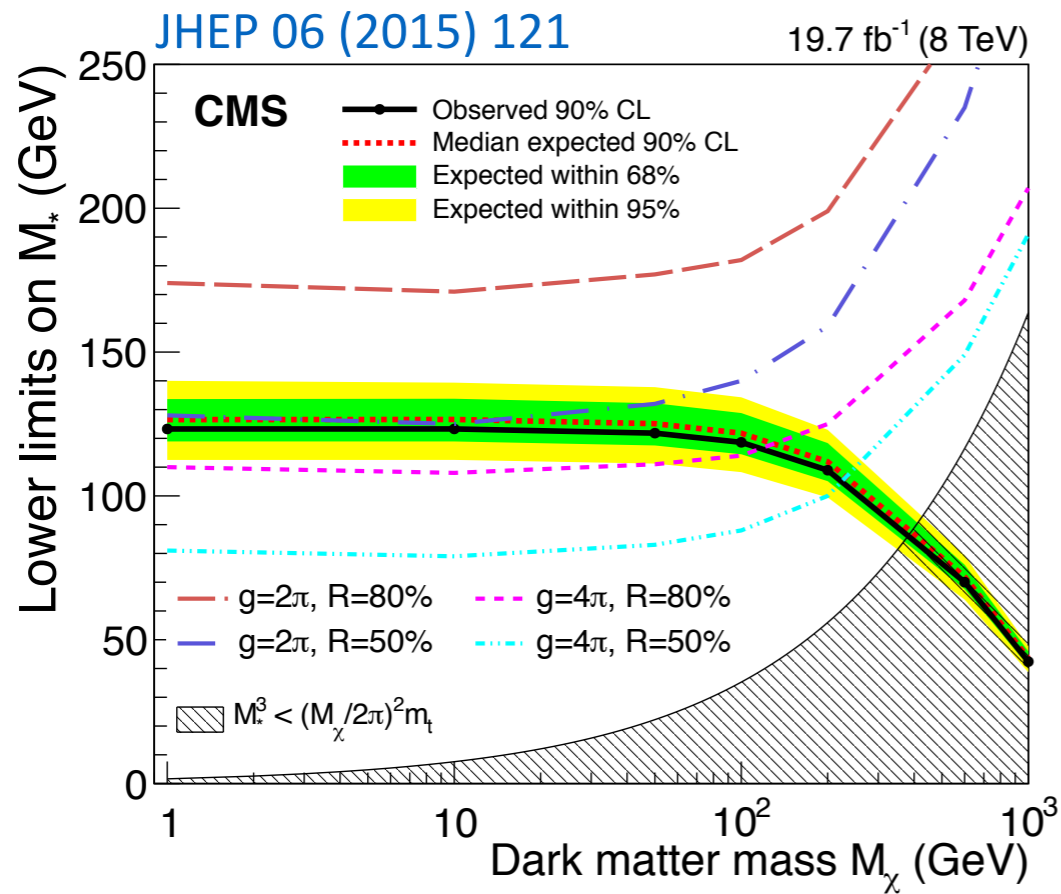


- MET > 320 GeV
- Additional kinematical cuts to reduce dominant background from SM $t\bar{t}$ events



Mono(?)- $t\bar{t}$

Interpret the result in scalar EFT (D1) $L_{\text{int}} = \frac{m_q}{M_*^3} \bar{q}q\bar{\chi}\chi$

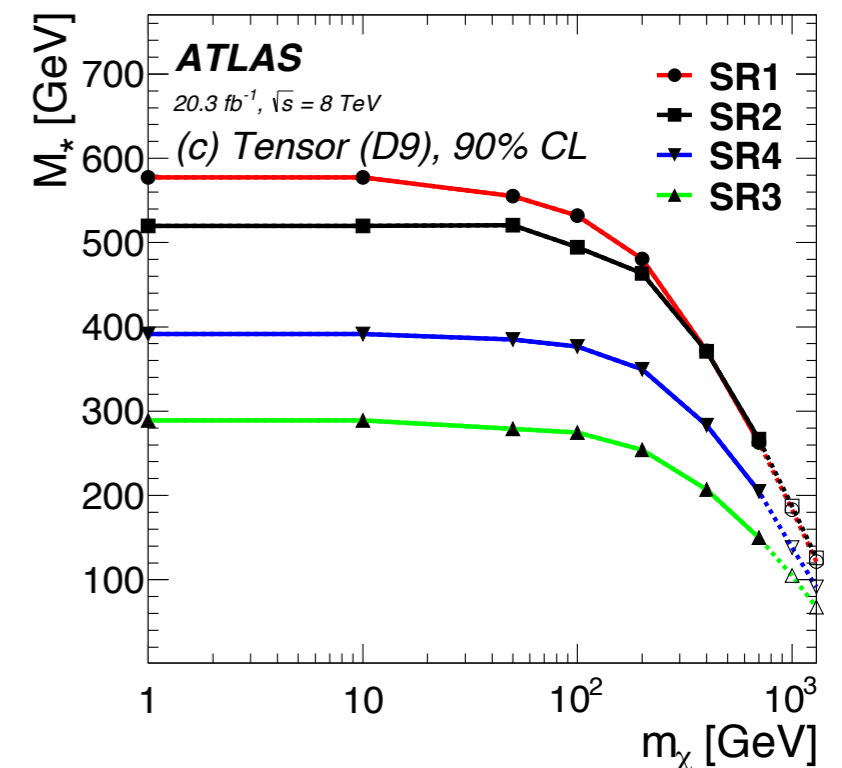
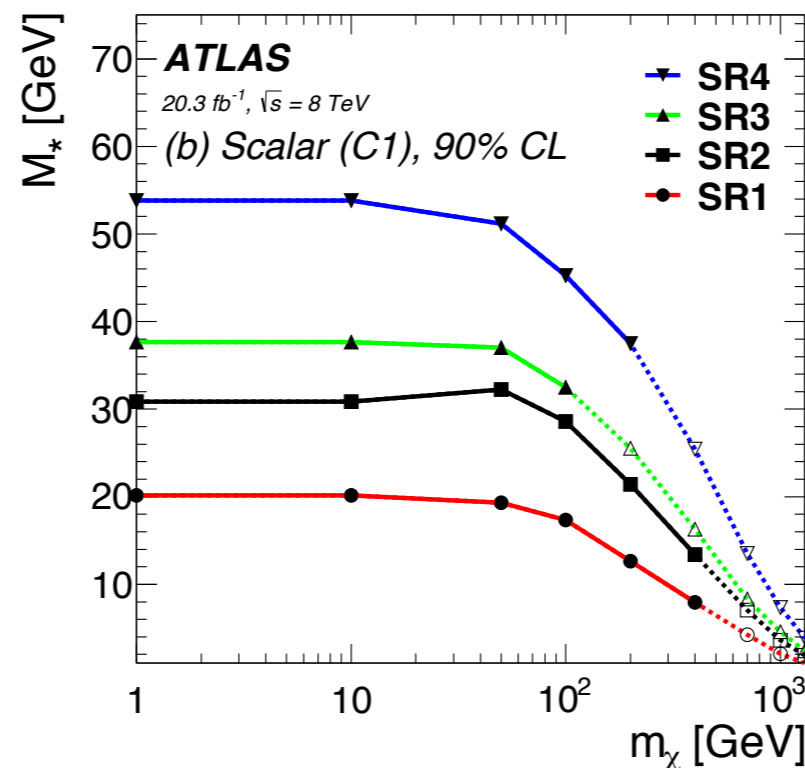
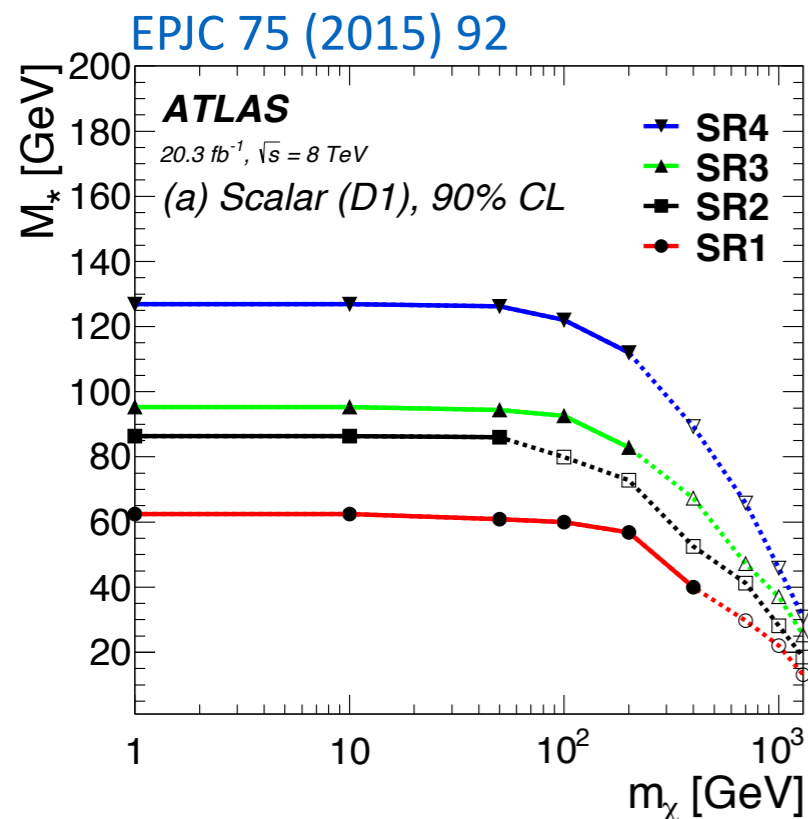
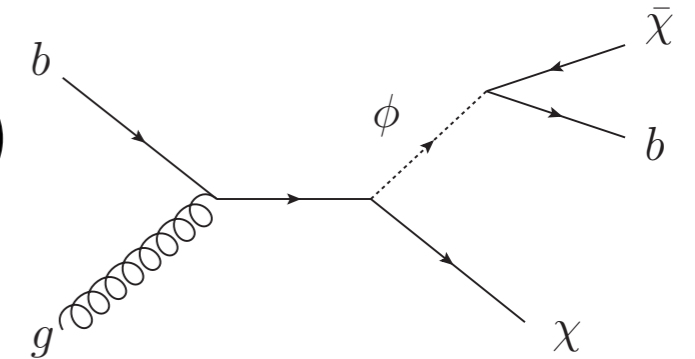


■ c.f. Lower limit on D1 from mono-jet ([ATLAS EPJC 75:299](#)) is ~ 38 GeV

Mono(?))- $b/b\bar{b}/t\bar{t}$

ATLAS [EPJC 75 (2015) 92] searched more broadly for events with MET plus b or t quarks

- Four signal regions targeting 1, 2 b -quarks (SR1, 2) and $t\bar{t}$ decaying hadronically and semileptonically (SR3, 4)
- Additional sensitivity to scalar (D1, C1) and tensor (D9), as well as a bottom-flavored DM model (b -FDM)



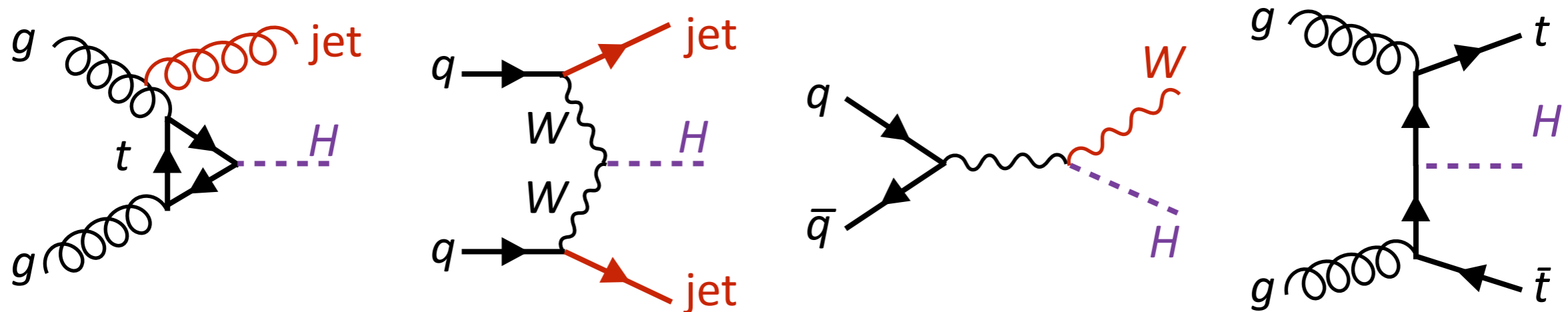
Higgs portal models

What if the Dark Matter couples only to the Higgs?

- Both $\sigma(pp \rightarrow \chi\bar{\chi})$ and $\sigma(N\chi \rightarrow N\chi)$ are strongly suppressed
 - ▶ Experimental constraints are very weak

One obvious test: invisible decay of Higgs \rightarrow DM

- Similar to mono- X searches, look for events in which an invisible Higgs is produced with something else



- Combine all measurements to set upper limit on $\text{BR}(H \rightarrow \text{invisible})$

Expt.	Reference	BR limit (95%)
ATLAS	arXiv:1509.00672, submitted to JHEP	< 23%
CMS	CMS-PAS-HIG-15-012	< 36%

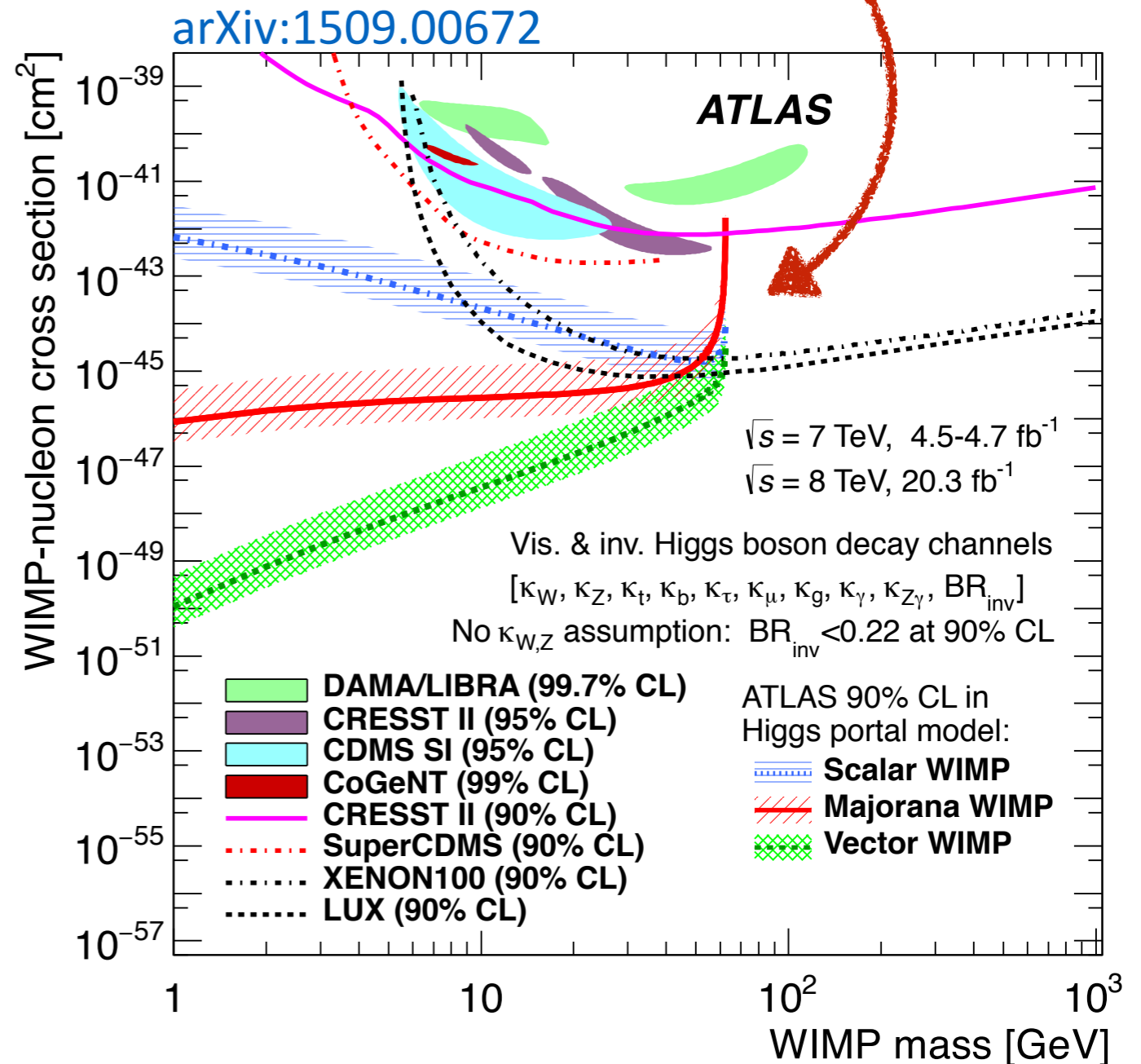
Higgs portal models

Expt.	Reference	BR limit (95%)
ATLAS	arXiv:1509.00672, submitted to JHEP	< 23%
CMS	CMS-PAS-HIG-15-012	< 36%

Translate from invisible Higgs decay to N - χ scattering cross-section

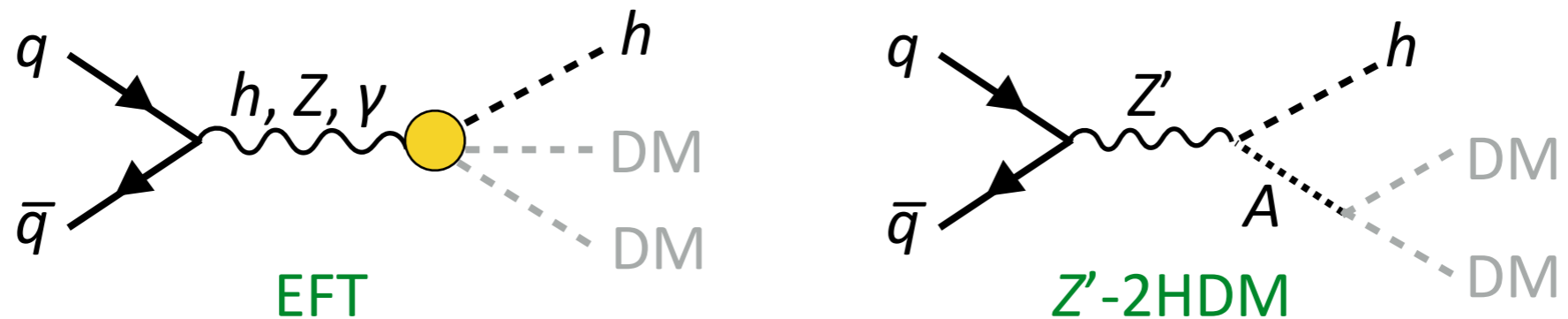
- Strong limit on Higgs Portal models for $m_\chi < m_H/2$
- NB: valid only if the DM does not interact with other SM particles

What if $m_\chi > m_H/2$?



Mono-Higgs

Mono-Higgs is another probe to DM with special (or unique) couplings to the Higgs



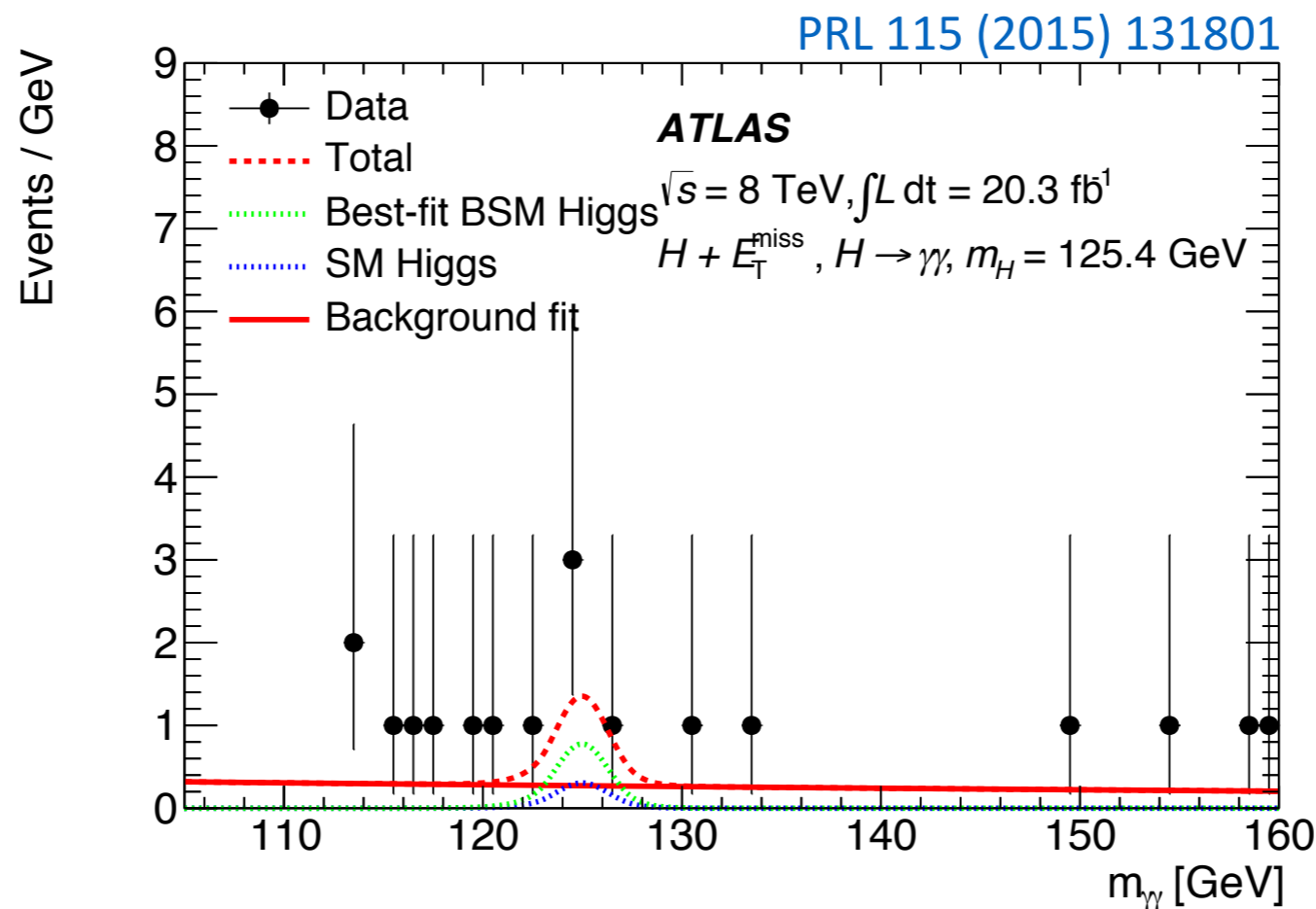
The Higgs can be detected in several decay channels

- $H \rightarrow \gamma\gamma$ has a small BR (0.23%), with small background
- $H \rightarrow b\bar{b}$ has a large BR (58%), with large background
- Other decays are too small to measure

Mono-Higgs $\rightarrow \gamma\gamma$

ATLAS mono- $H \rightarrow \gamma\gamma$ [PRL 115 (2015) 131801]

- Select events with 2 photons with $105 < m_{\gamma\gamma} < 160$ GeV
- Require MET > 90 GeV and $p_T^{\gamma\gamma} > 90$ GeV

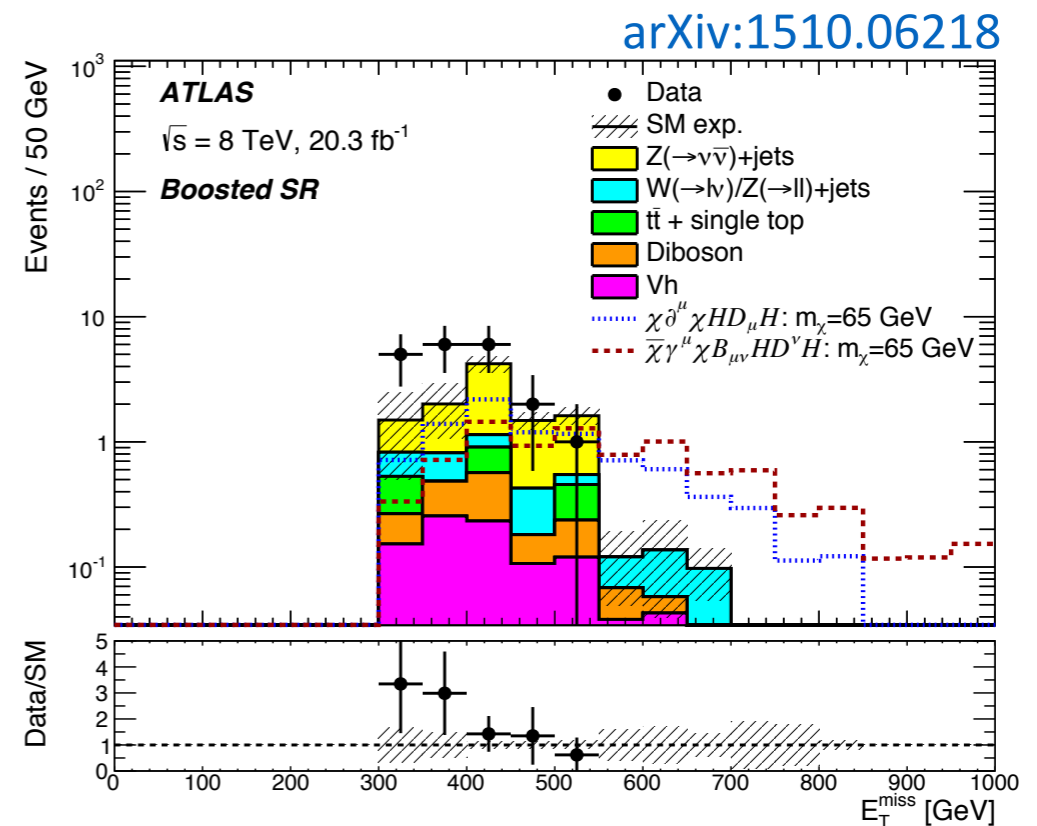
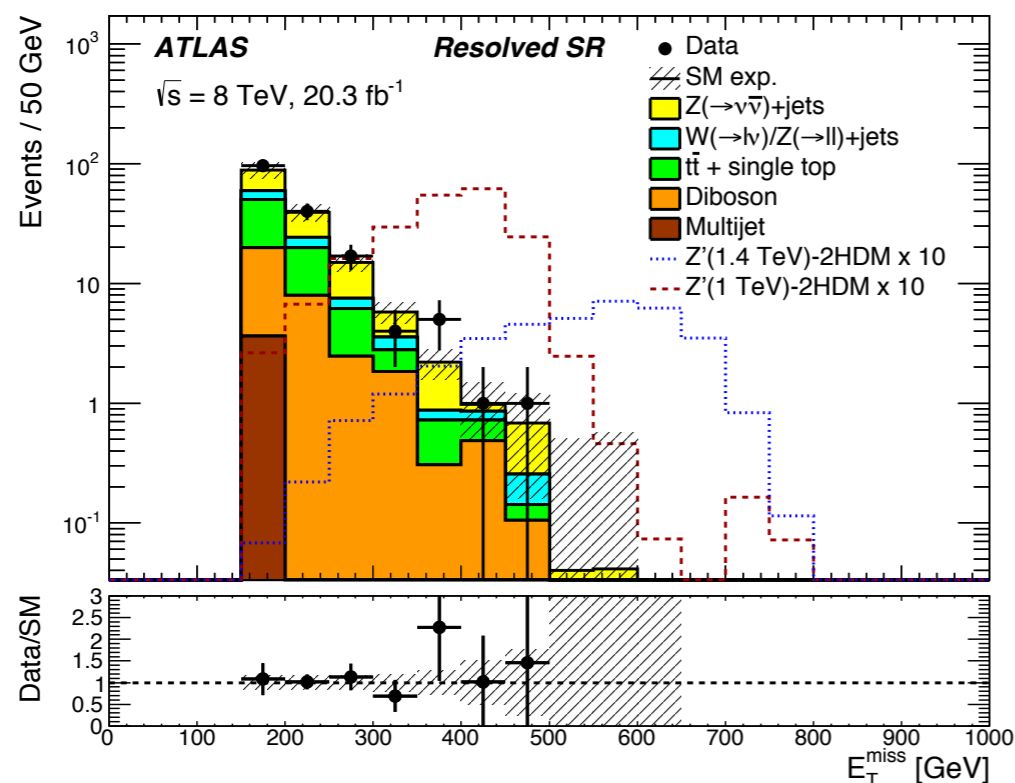
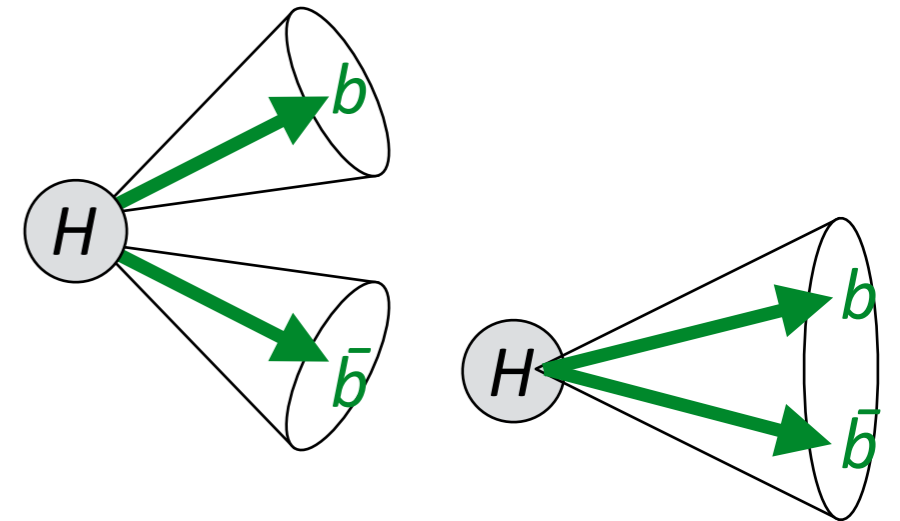


- Look for a peak at $m_{\gamma\gamma} = 125$ GeV. Practically zero background
- Limited statistics makes $H \rightarrow \gamma\gamma$ less powerful than $H \rightarrow bb$ for now

Mono-Higgs $\rightarrow b\bar{b}$

ATLAS mono- $H \rightarrow b\bar{b}$ [arXiv:1510.06218]

- Two b -jets may be resolved or merged
- Resolved: 2 b -tagged jets with $p_T > 100/60$ GeV. $90 < m_{bb} < 150$ GeV. MET > 150 GeV
- Merged: 1 fat jet with $p_T > 350$ GeV that contains 2 b -tagged sub-jets. $90 < m_J < 150$ GeV. MET > 300 GeV



- Background is dominated by $Z(\rightarrow v\bar{v}) + \text{jets}$

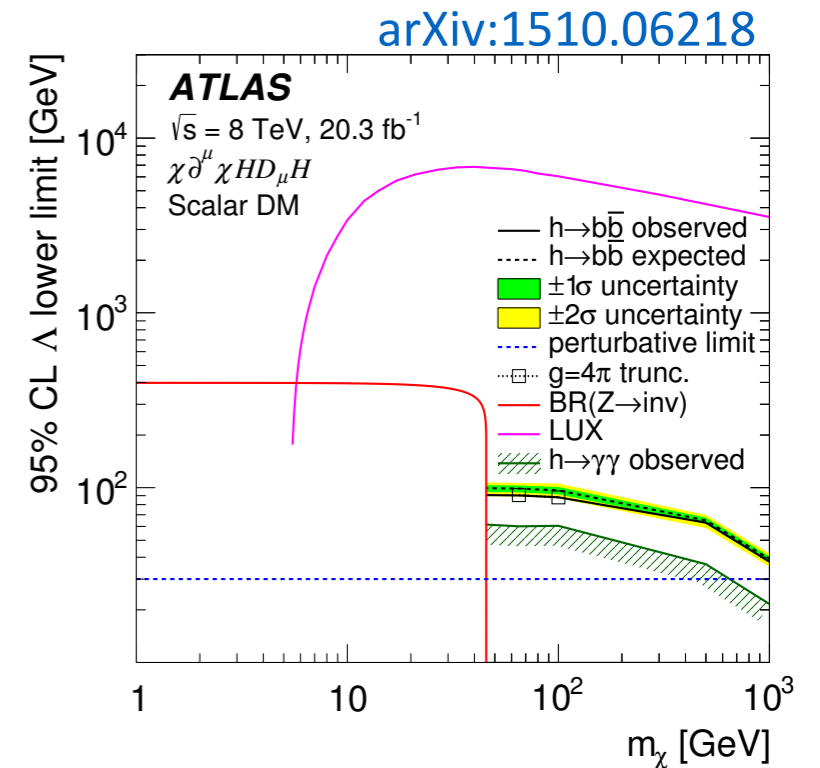
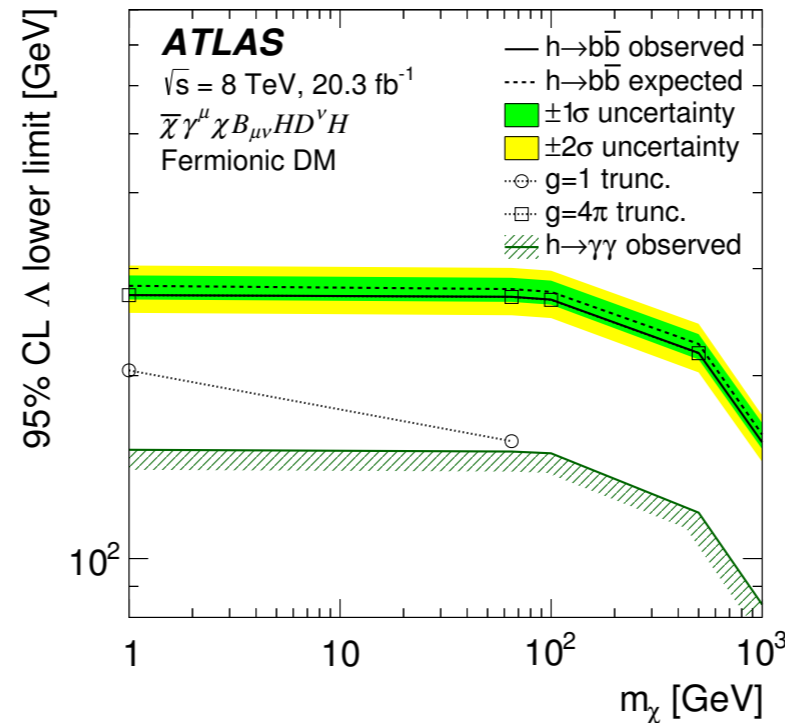
Mono-Higgs

ATLAS mono- $H \rightarrow b\bar{b}$ results interpreted

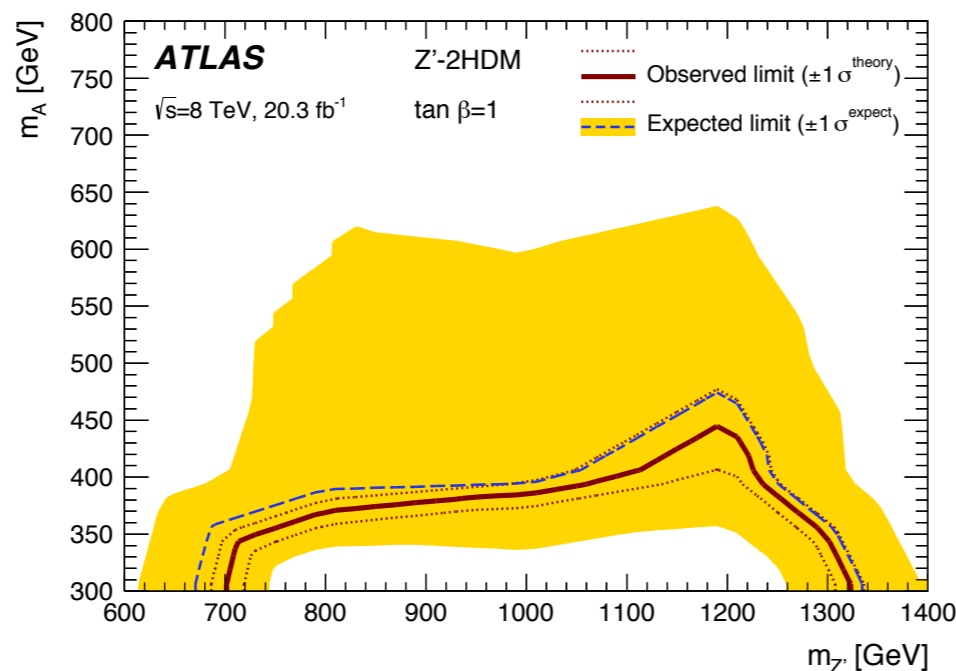
EFT scale limits for

▶ $\frac{1}{\Lambda^4} \bar{\chi} \gamma^\mu \chi B_{\mu\nu} H^\dagger D^\nu H$ (left)

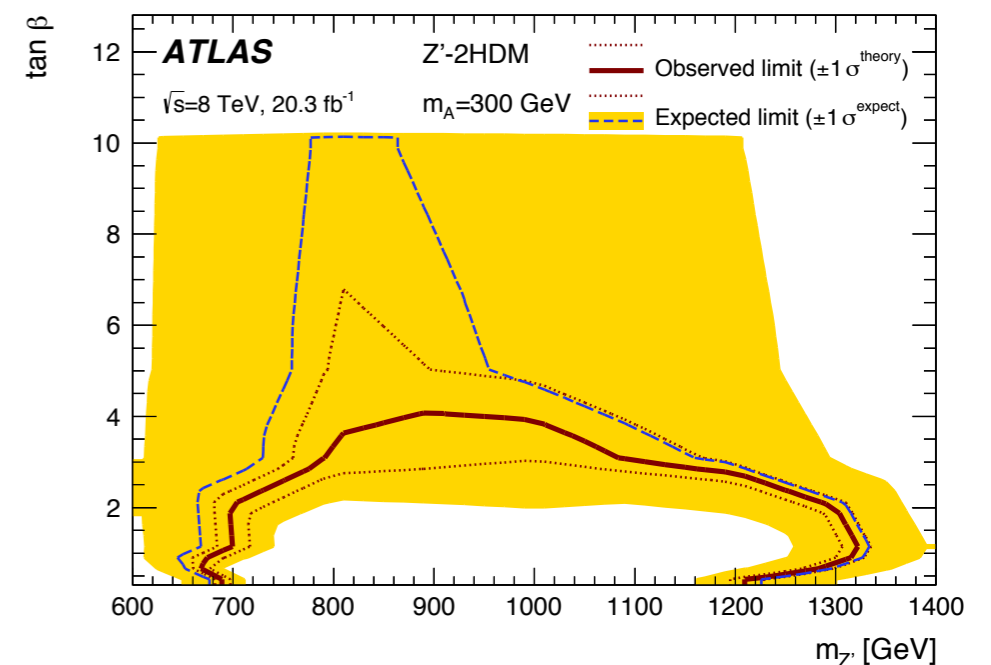
▶ $\frac{1}{\Lambda^2} \chi^\dagger \partial^\mu \chi H^\dagger D_\mu H$ (right)



Simplified $Z' + 2\text{HDM}$



(a) $m_{Z'} - m_A$

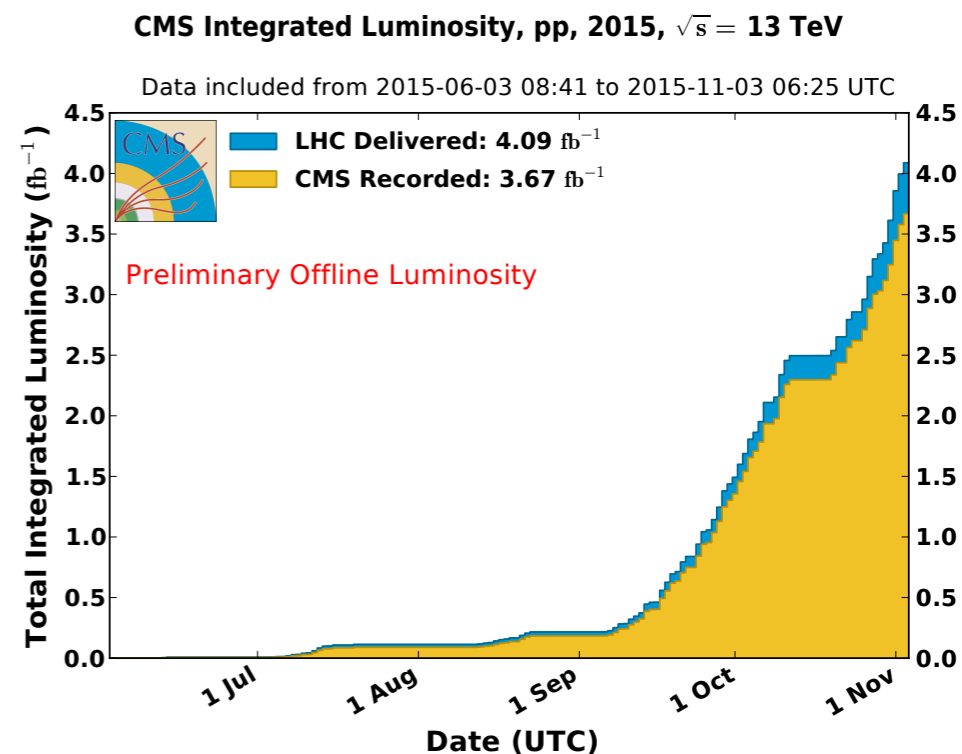
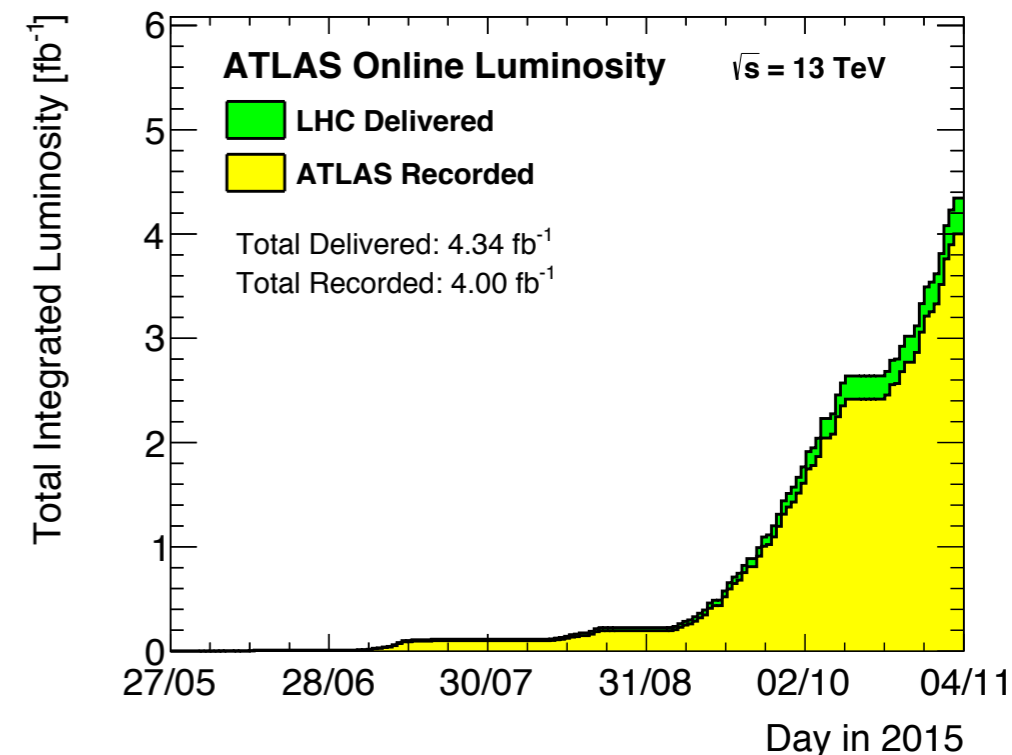
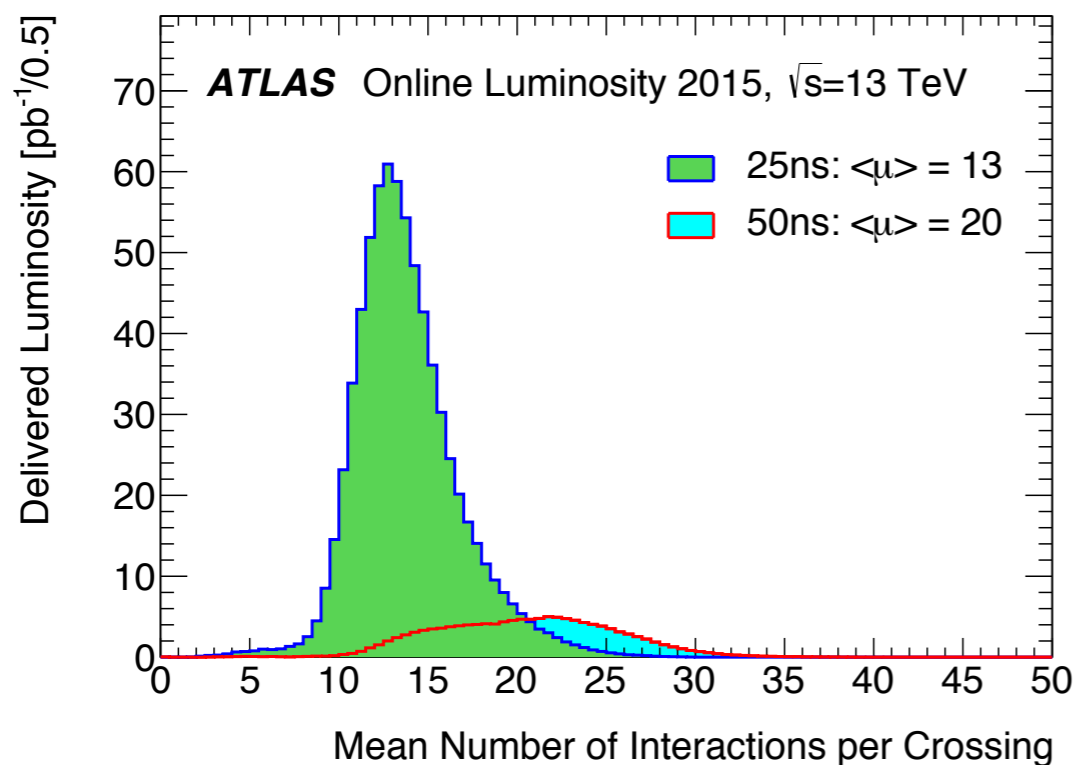


(b) $m_{Z'} - \tan \beta$

LHC Run 2

LHC re-started in June 2015 with 13 TeV

- About 4 fb⁻¹/experiment recorded in 2015
 - ▶ Will accumulate 100 fb⁻¹ by 2018
- 8 TeV → 13 TeV increases parton luminosity
 - ▶ 2015 data can improve over Run 1
- LHC bunch spacing changed 50 ns → 25 ns
 - ▶ Reduced pile-up for the same luminosity



ATLAS mono-W/Z ($\rightarrow q\bar{q}$) in Run 2

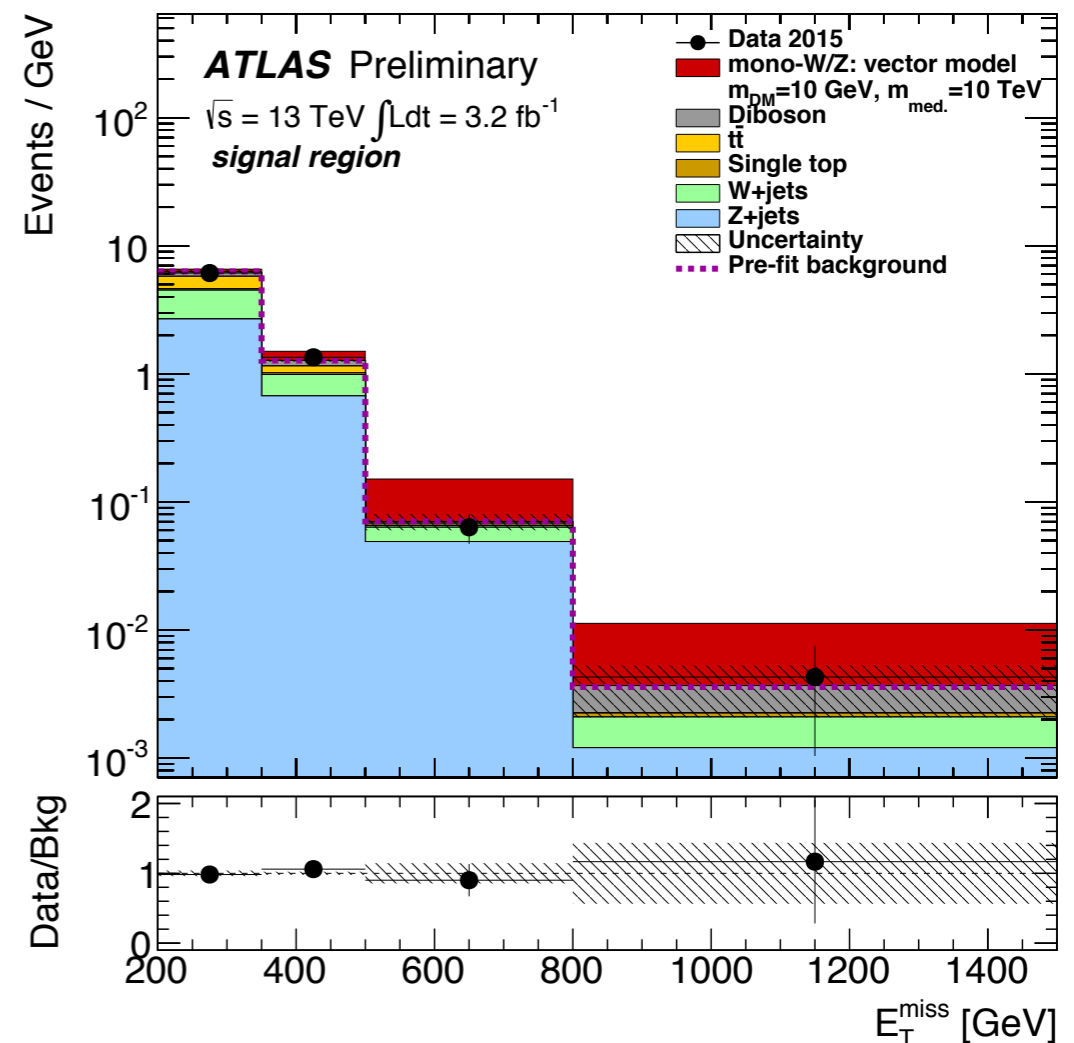
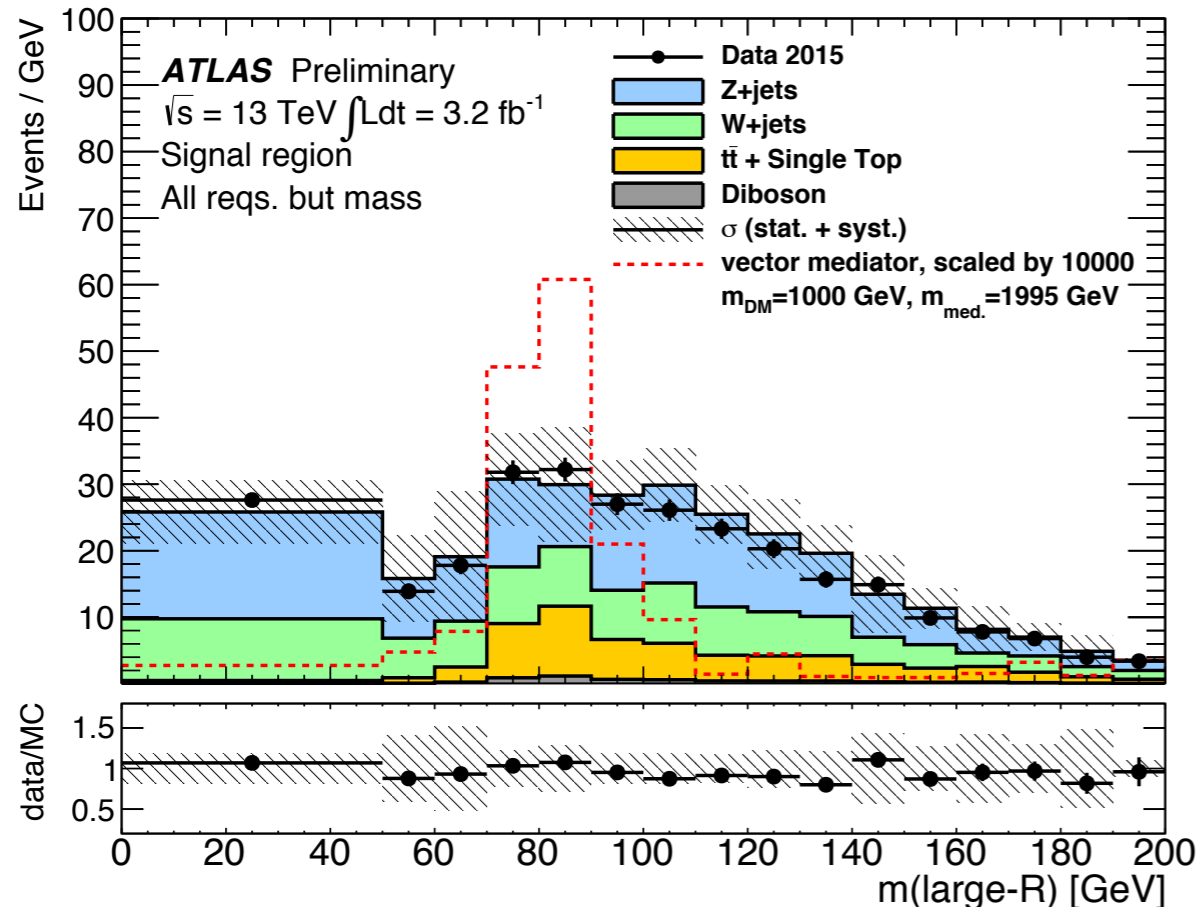
ATLAS-CONF-2015-080

ATLAS has just released a Run 2 search for mono-W/Z

- 3.3 fb⁻¹ of p - p collision data at 13 TeV
- Analysis is similar to the Run 1 search [PRL 112 (2014) 041802]

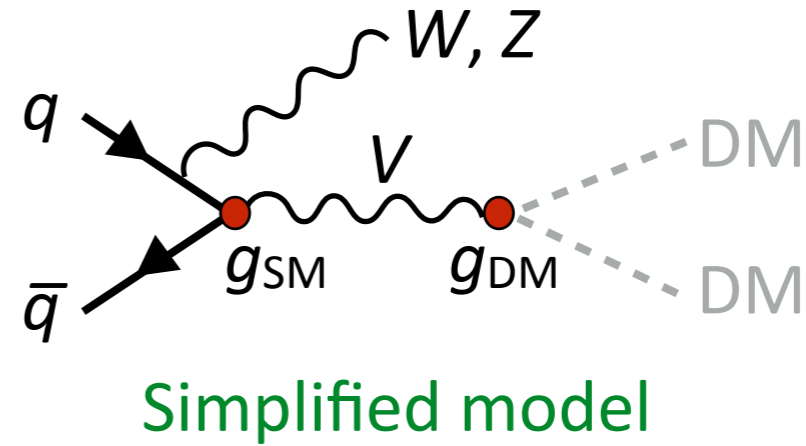
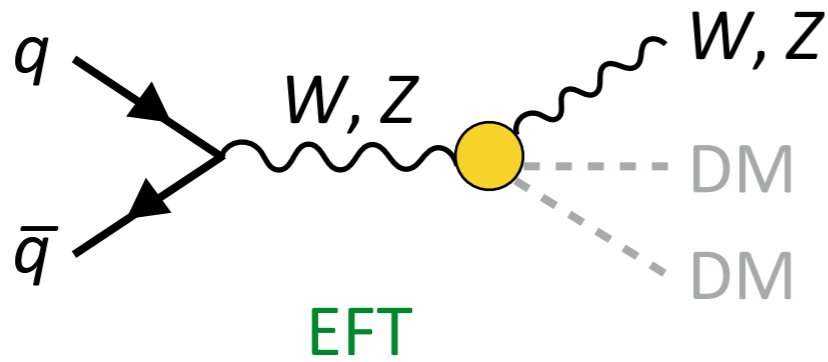
Select events with a fat jet (= boosted $W/Z \rightarrow q\bar{q}$) + MET

- The fat jet has $p_T > 200$ GeV and sub-jets consistent with $W/Z \rightarrow q\bar{q}$
- MET > 250 GeV. No leptons

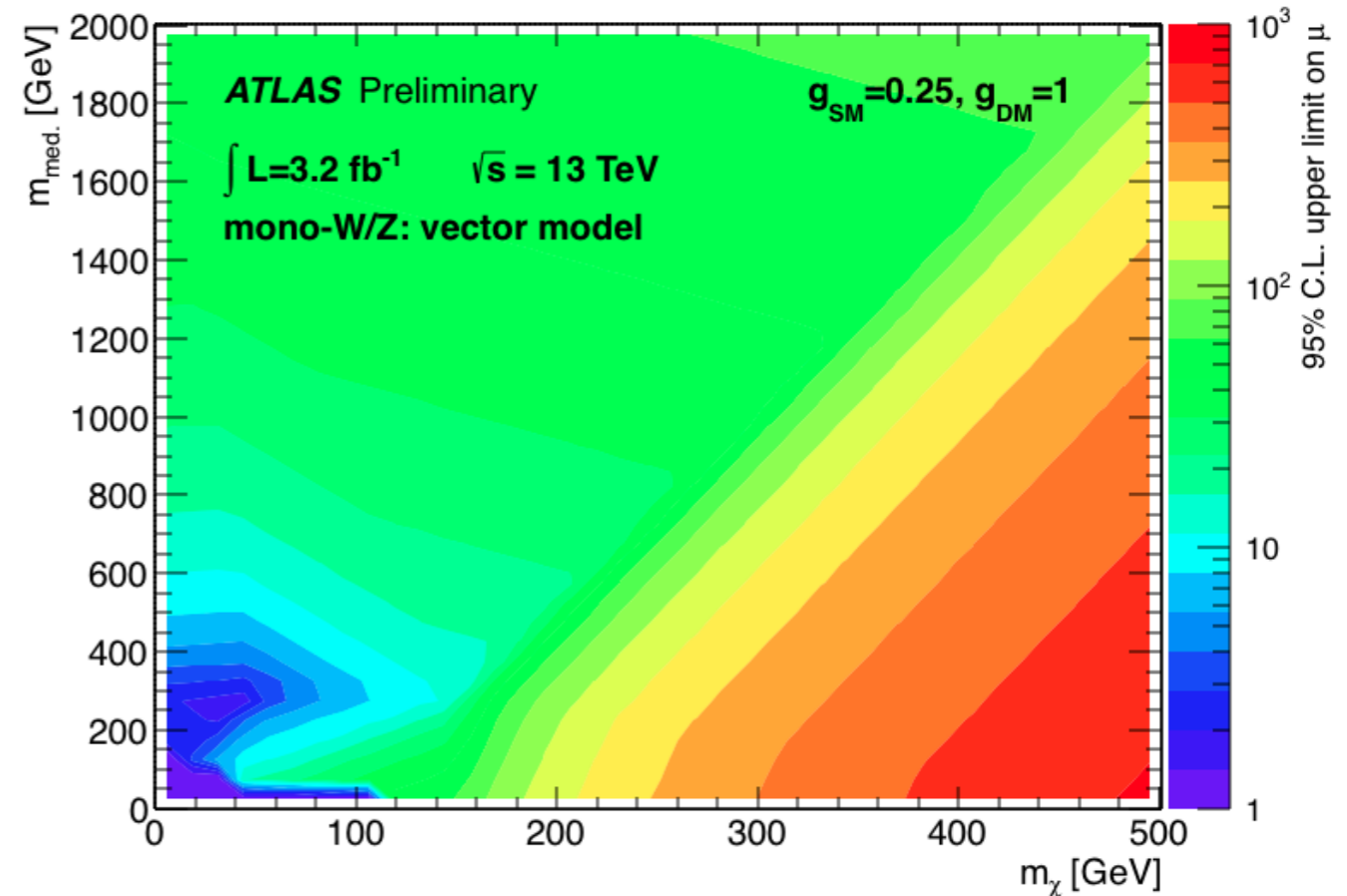
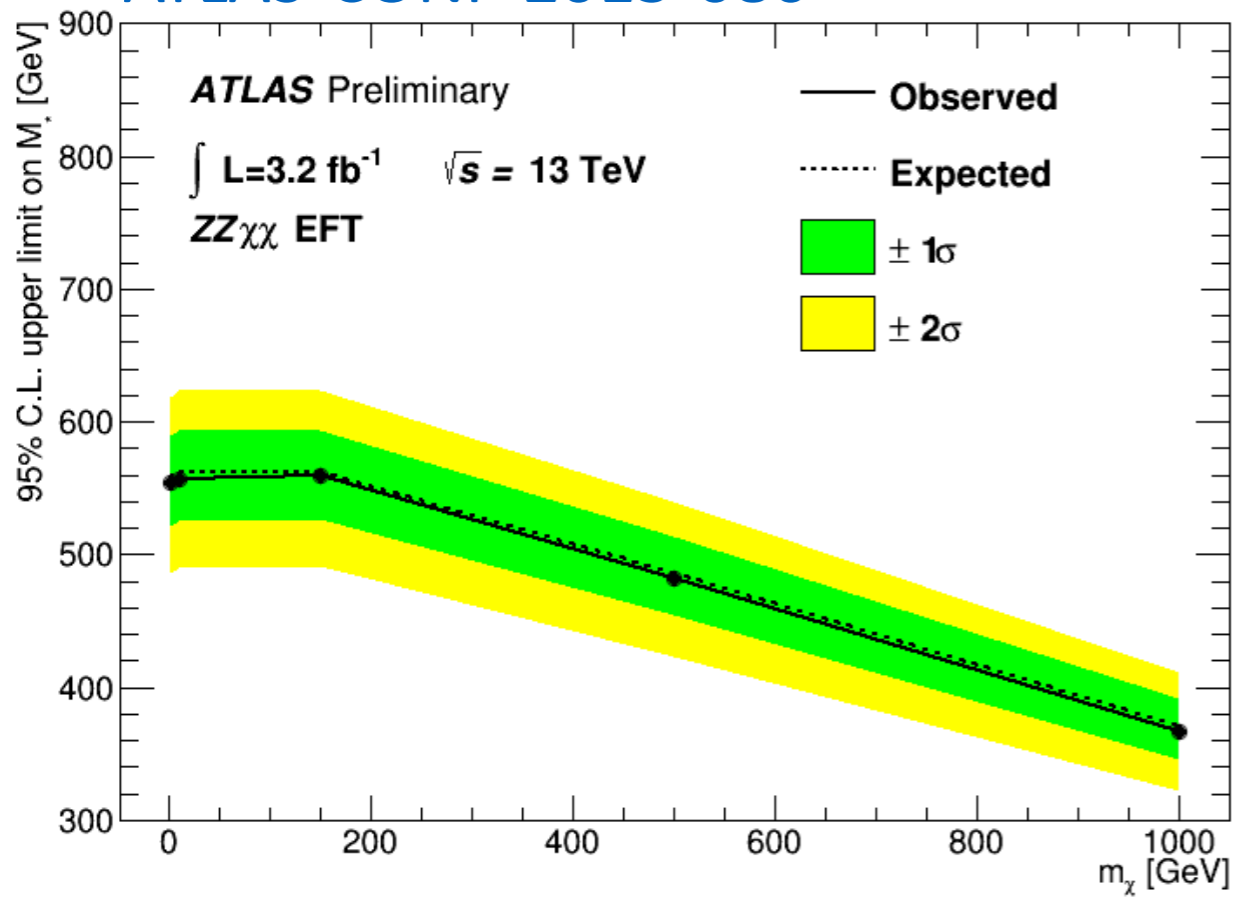


ATLAS mono-W/Z ($\rightarrow q\bar{q}$) in Run 2

Interpretation in EFT and simplified models



ATLAS-CONF-2015-080



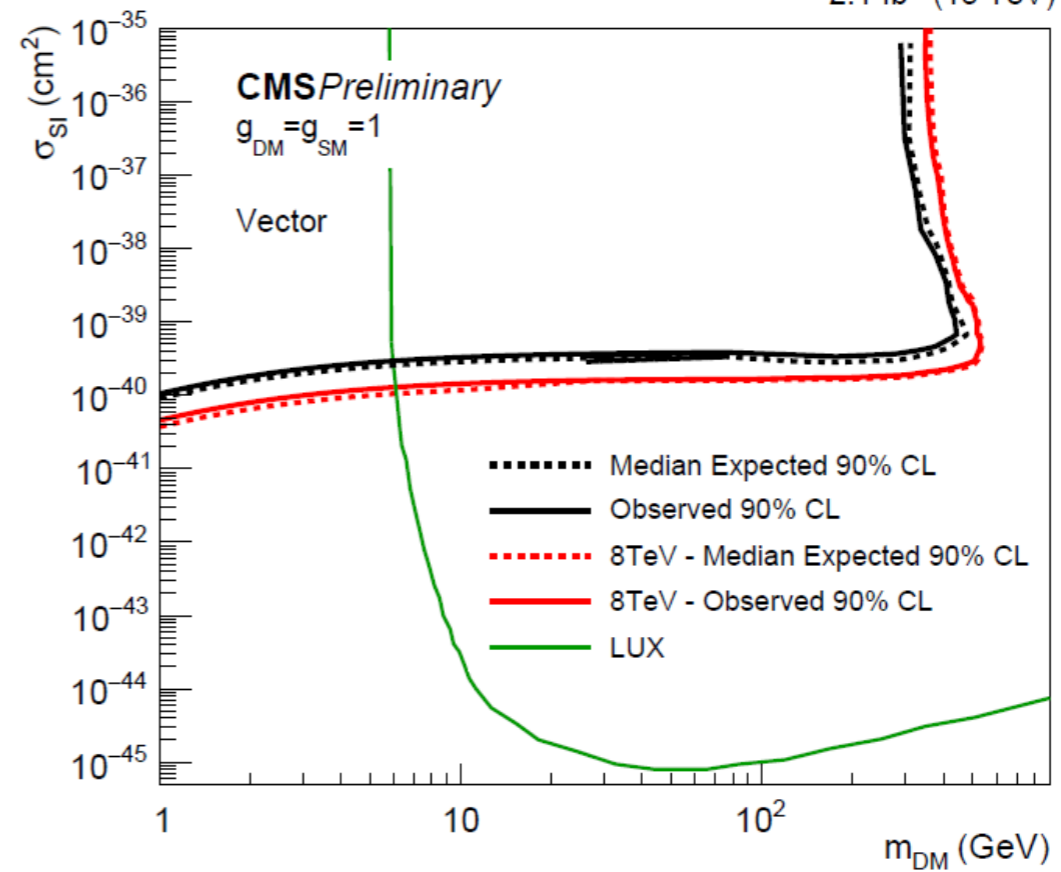
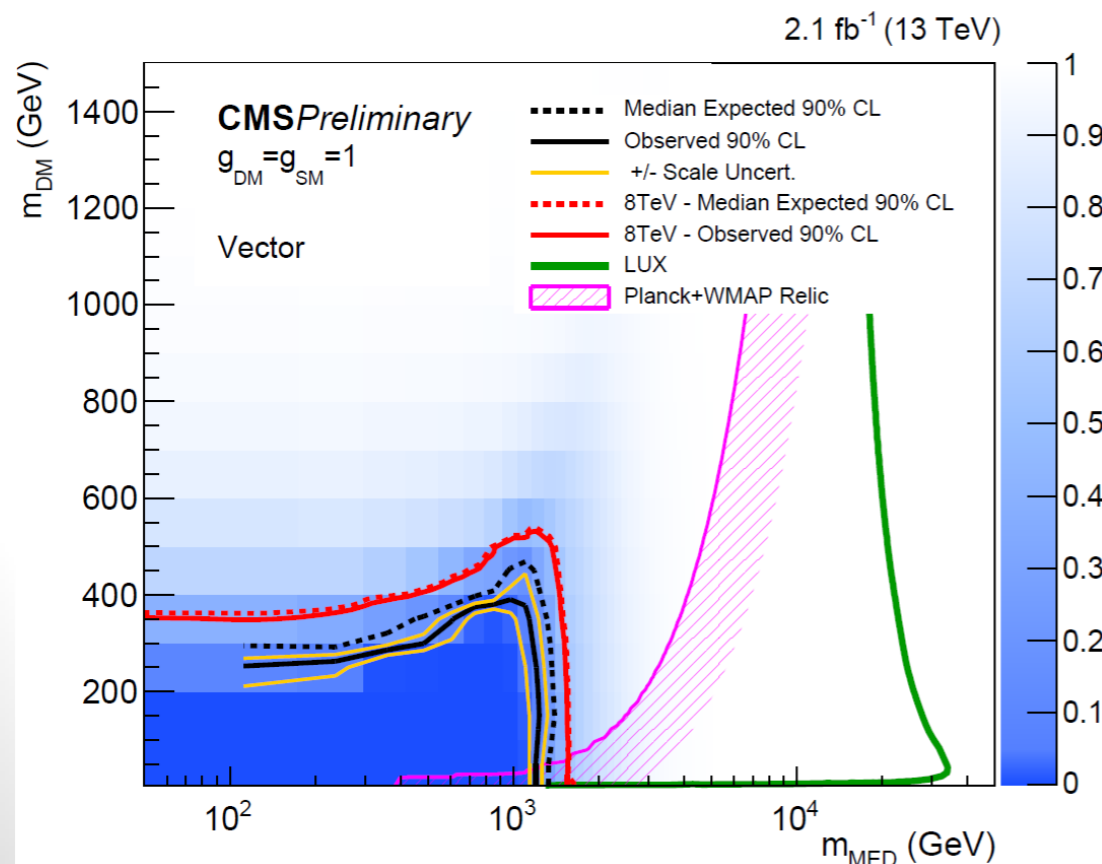
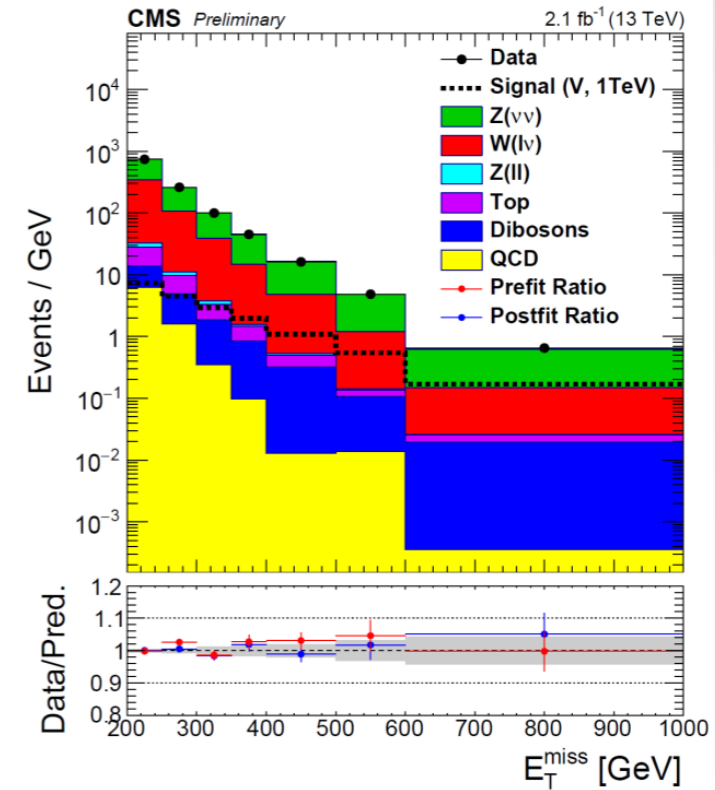
CMS mono-jet in Run 2

Slide by Jim Olsen

EXO-15-003

Search for dark matter

- Search for generic dark matter in **final states with jets and large missing transverse energy**
- Traditional monojet search extended to multijet final states, searching for DM pairs produced via a **vector mediator**
- Limits comparable to those set in Run 1



Summary

LHC hunts for Dark Matter in $pp \rightarrow X + \text{invisible}$ searches

- Mono-jet/photon/ $W/Z/t\bar{t}$ searches cover a wide variety of DM interactions, for DM masses between 0 and several 100 GeV
 - ▶ Sensitive to light (as well as heavy) DM, spin-dependent (as well as spin-independent) interactions
 - ▶ Complementary to direct-detection experiments
- Higgs portal models are tested by invisible Higgs and mono-Higgs searches

Interpretation of results are often model-dependent

- Careful with plots that are valid only for particular forms of interaction
- Expect to see more simplified-model interpretations

LHC Run 2 has started with 13 TeV

- Results from 2015 data are coming out just now
- Next 3 years will vastly increase the data and improve our sensitivity