

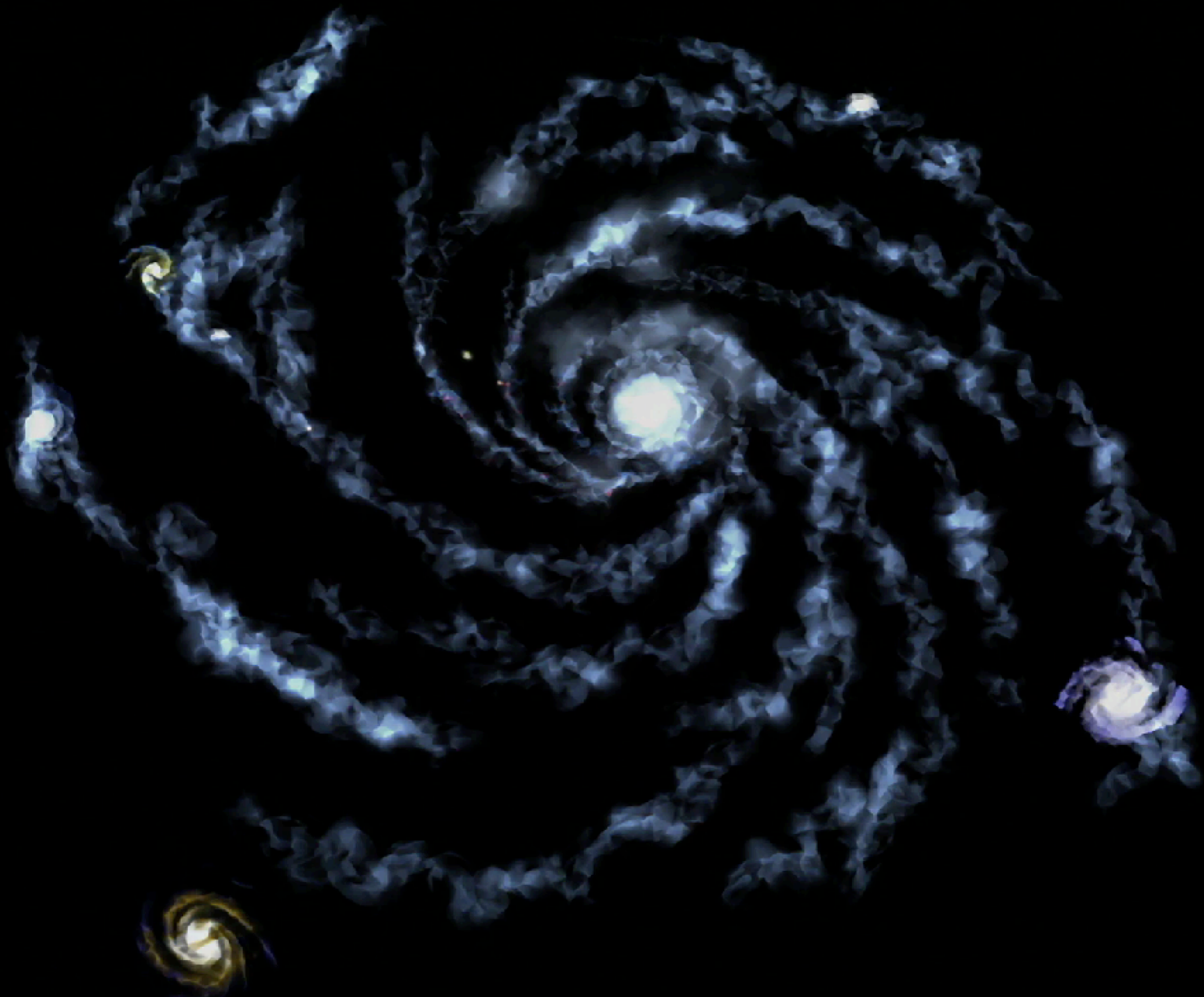
Finding Cosmic Inflation

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(MPI für Astrophysik / Kavli IPMU)

YKIS 2018 “*General Relativity: The Next Generation*”

February 19, 2018



Full-dome movie for planetarium
Director: Hiromitsu Kohsaka

日本語バージョン

「HORIZON～宇宙の果てにあるもの～」

多摩六都科学館（西東京市）

仙台市天文台（仙台市）

鹿児島市立科学館（鹿児島市）

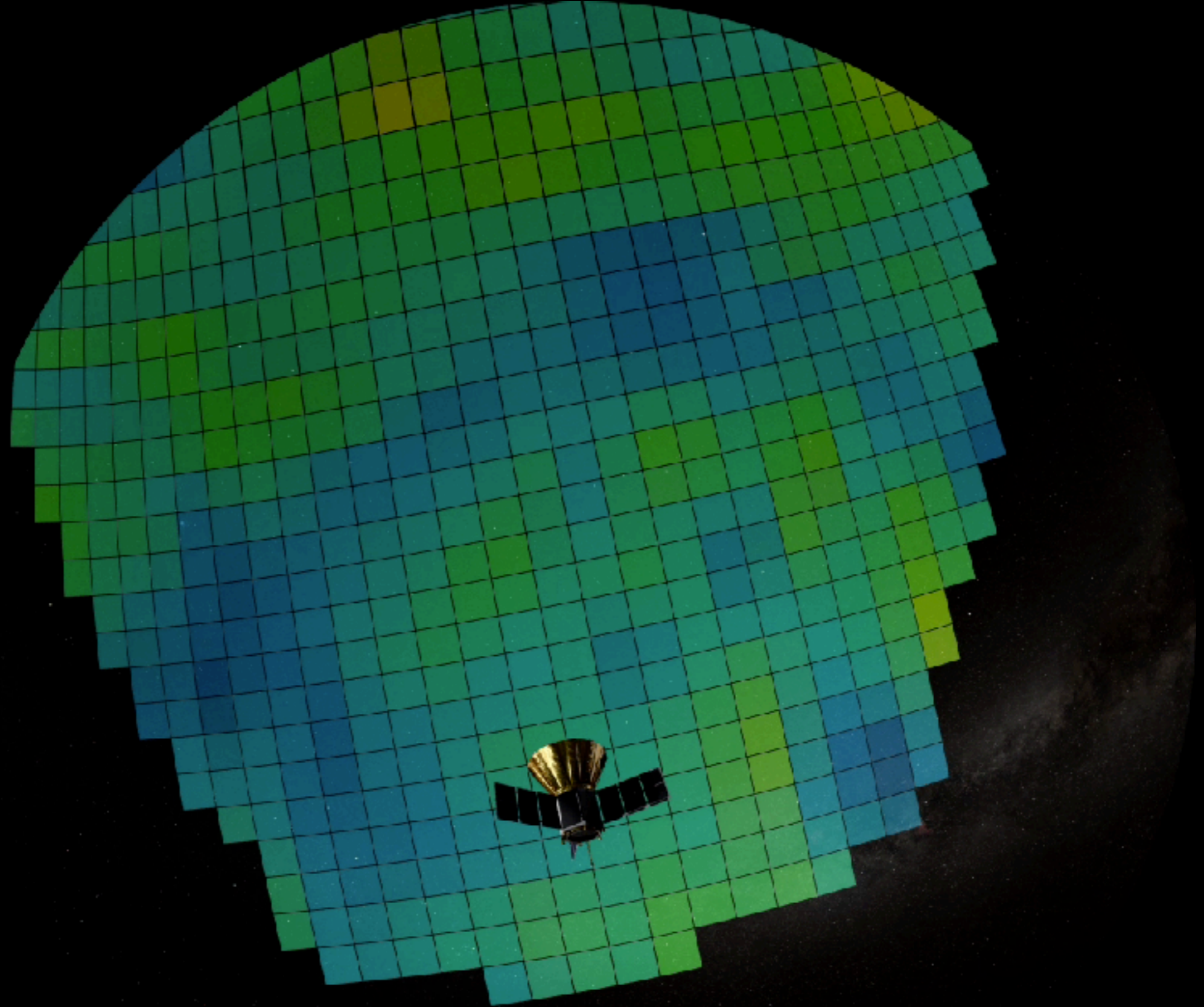
他、で上映中・上映予定

HORIZON

Beyond the Edge of the Visible Universe

▶ ⏩ 🔊 2:28 / 2:51

⚙️ HD 📺 🗉



A Remarkable Story

- Observations of the cosmic microwave background and their interpretation taught us that **galaxies, stars, planets, and ourselves originated from tiny fluctuations in the early Universe**
- *But, what generated the initial fluctuations?*

Leading Idea

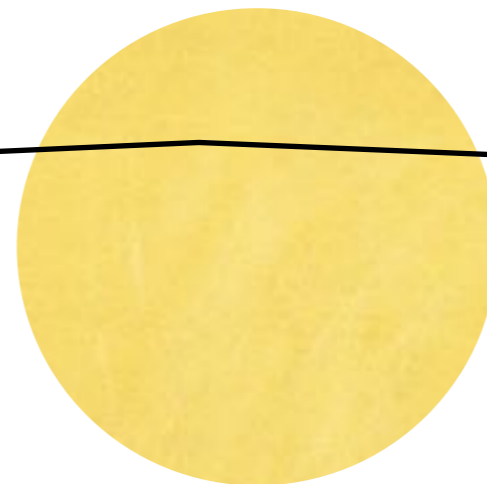
- Quantum mechanics at work in the early Universe
 - “*We all came from quantum fluctuations*”
- But, how did quantum fluctuations on the *microscopic* scales become *macroscopic* fluctuations over large distances?
 - What is the **missing link** between small and large scales?

Cosmic Inflation

Quantum fluctuations on
microscopic scales



Inflation!



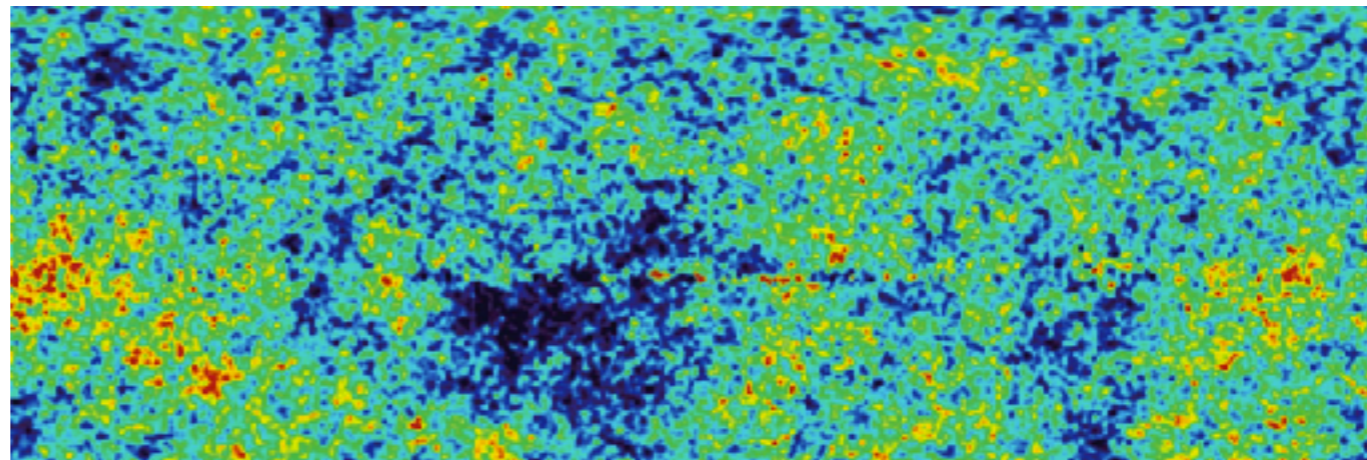
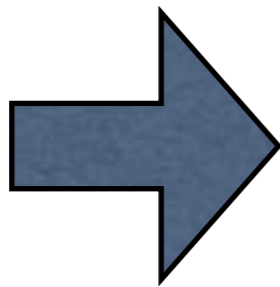
- Exponential expansion (inflation) stretches the wavelength of quantum fluctuations to cosmological scales

Key Predictions

ζ

scalar
mode

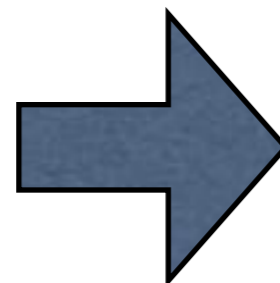
- Fluctuations we observe today in CMB and the matter distribution originate from quantum fluctuations during inflation



h_{ij}

tensor
mode

- There should also be *ultra long-wavelength* gravitational waves generated during inflation



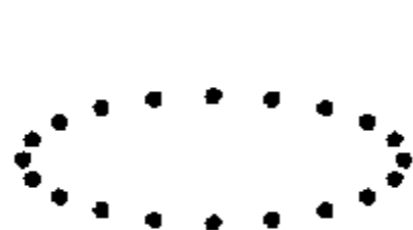
Starobinsky (1979)

We measure distortions in space

- A distance between two points in space

$$d\ell^2 = a^2(t) [1 + 2\zeta(\mathbf{x}, t)] [\delta_{ij} + h_{ij}(\mathbf{x}, t)] dx^i dx^j$$

- ζ : “curvature perturbation” (scalar mode)
 - Perturbation to the determinant of the spatial metric
- h_{ij} : “gravitational waves” (tensor mode)
 - Perturbation that does not alter the determinant



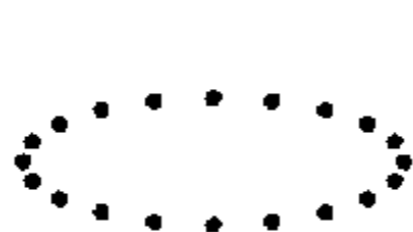
$$\sum_i h_{ii} = 0$$

We measure distortions in space

- A distance between two points in space

$$d\ell^2 = \underbrace{a^2(t)}_{\text{scale factor}} [1 + 2\zeta(\mathbf{x}, t)] [\delta_{ij} + h_{ij}(\mathbf{x}, t)] dx^i dx^j$$

- ζ : “curvature perturbation” (scalar mode)
 - Perturbation to the determinant of the spatial metric
- h_{ij} : “gravitational waves” (tensor mode)
 - Perturbation that does not alter the determinant



$$\sum_i h_{ii} = 0$$

Finding Inflation

- Inflation is the accelerated, quasi-exponential expansion. Defining the Hubble expansion rate as $H(t) = d \ln(a) / dt$, we must find

$$\frac{\ddot{a}}{a} = \dot{H} + H^2 > 0 \quad \longrightarrow \quad \epsilon \equiv -\frac{\dot{H}}{H^2} < 1$$

- For inflation to explain flatness of spatial geometry of our observable Universe, we need to have a **sustained** period of inflation. This implies $\epsilon = O(N^{-1})$ or smaller, where N is the number of e-folds of expansion counted from the end of inflation:

$$N \equiv \ln \frac{a_{\text{end}}}{a} = \int_t^{t_{\text{end}}} dt' H(t') \approx 50$$

Have we found inflation?

- *Have we found $\varepsilon \ll 1$?*

$$\varepsilon \equiv -\frac{\dot{H}}{H^2}$$

- To achieve this, we need to map out **H(t)**, and show that it does not change very much with time
 - **We need the “Hubble diagram” during inflation!**

Fluctuations are proportional to H

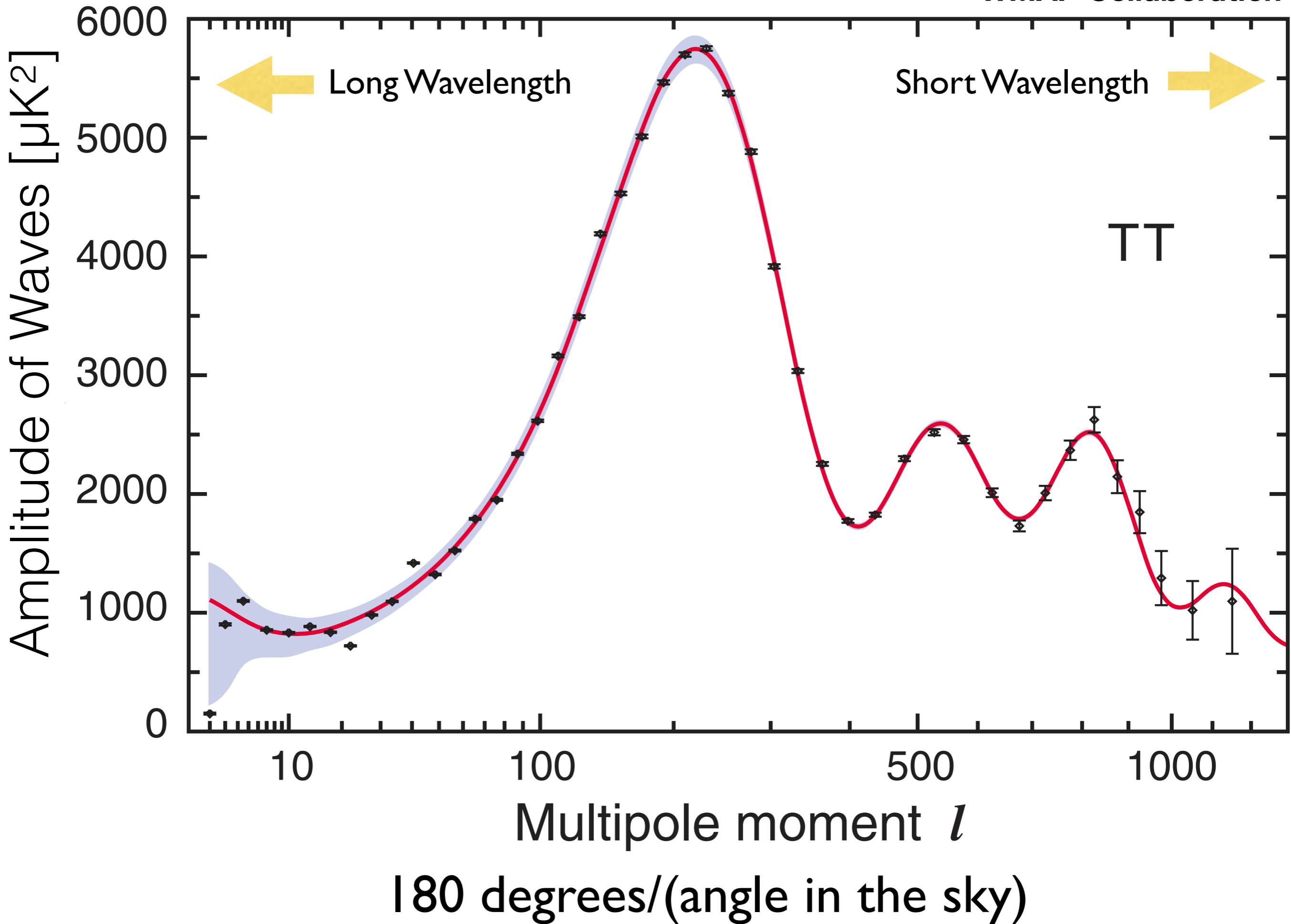
- Both scalar (ζ) and tensor (h_{ij}) perturbations are proportional to H
- Consequence of the uncertainty principle
 - [energy you can borrow] \sim [time you borrow] $^{-1} \sim H$
- **THE KEY:** The earlier the fluctuations are generated, the more its wavelength is stretched, and thus the bigger the angles they subtend in the sky. **We can map H(t) by measuring CMB fluctuations over a wide range of angles**

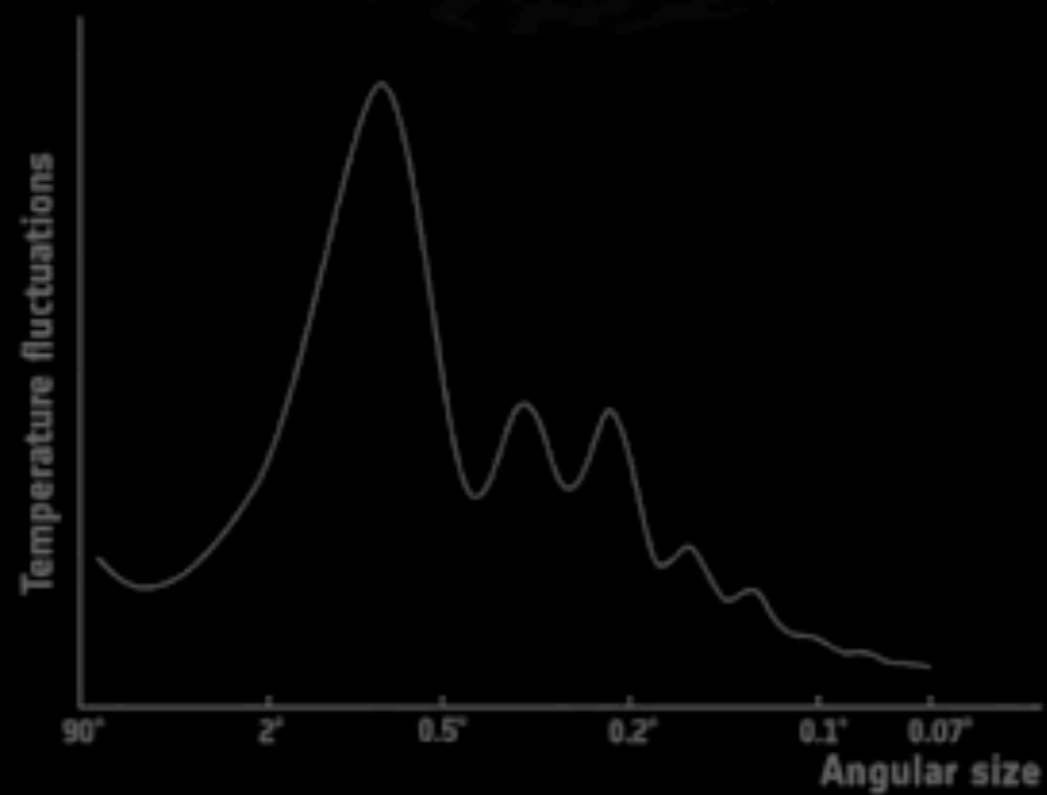
Fluctuations are proportional to H

- **We can map $H(t)$ by measuring CMB fluctuations over a wide range of angles**
 1. We want to show that the amplitude of CMB fluctuations does not depend very much on angles
 2. Moreover, since inflation must end, H would be a decreasing function of time. It would be fantastic to show that the amplitude of CMB fluctuations actually **DOES** depend on angles such that the small scale has ***slightly*** smaller power

Data Analysis

- Decompose temperature fluctuations in the sky into a set of waves with various wavelengths
- Make a diagram showing the strength of each wavelength

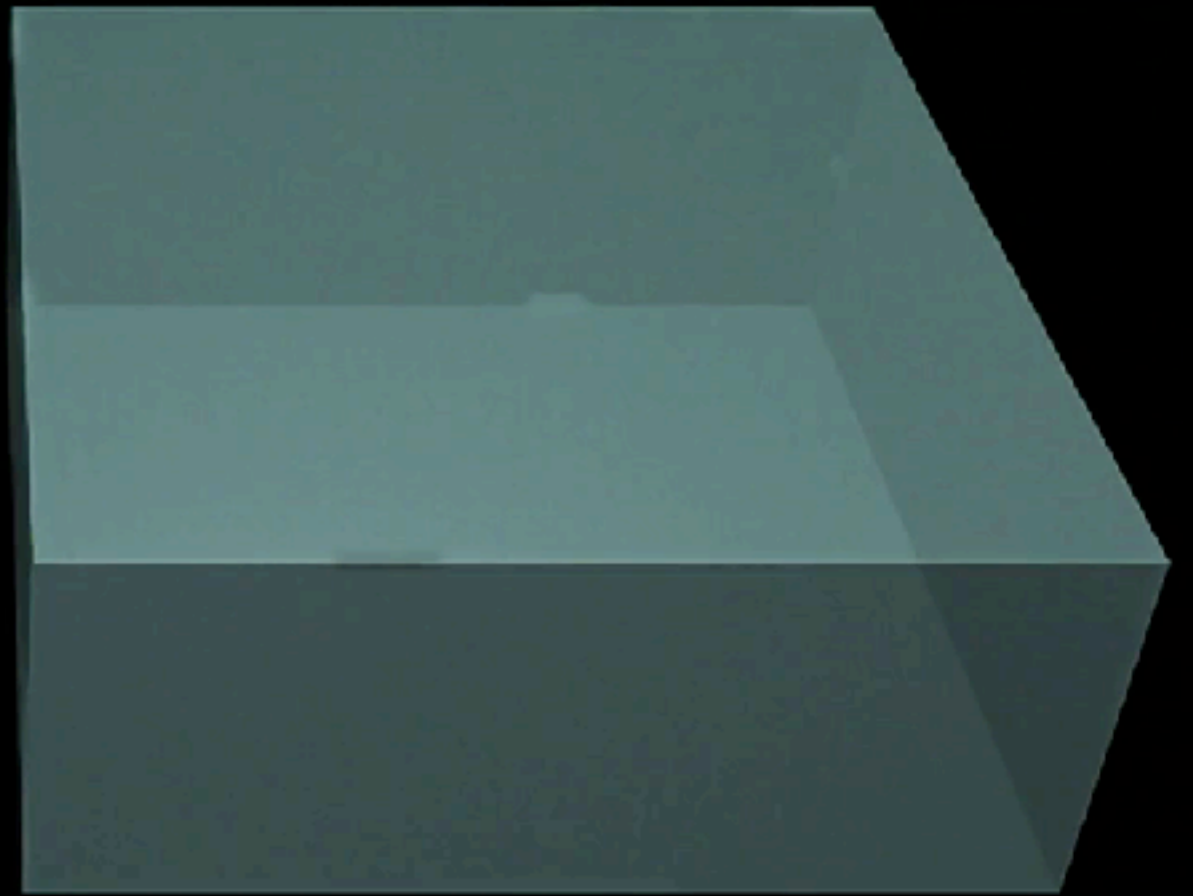
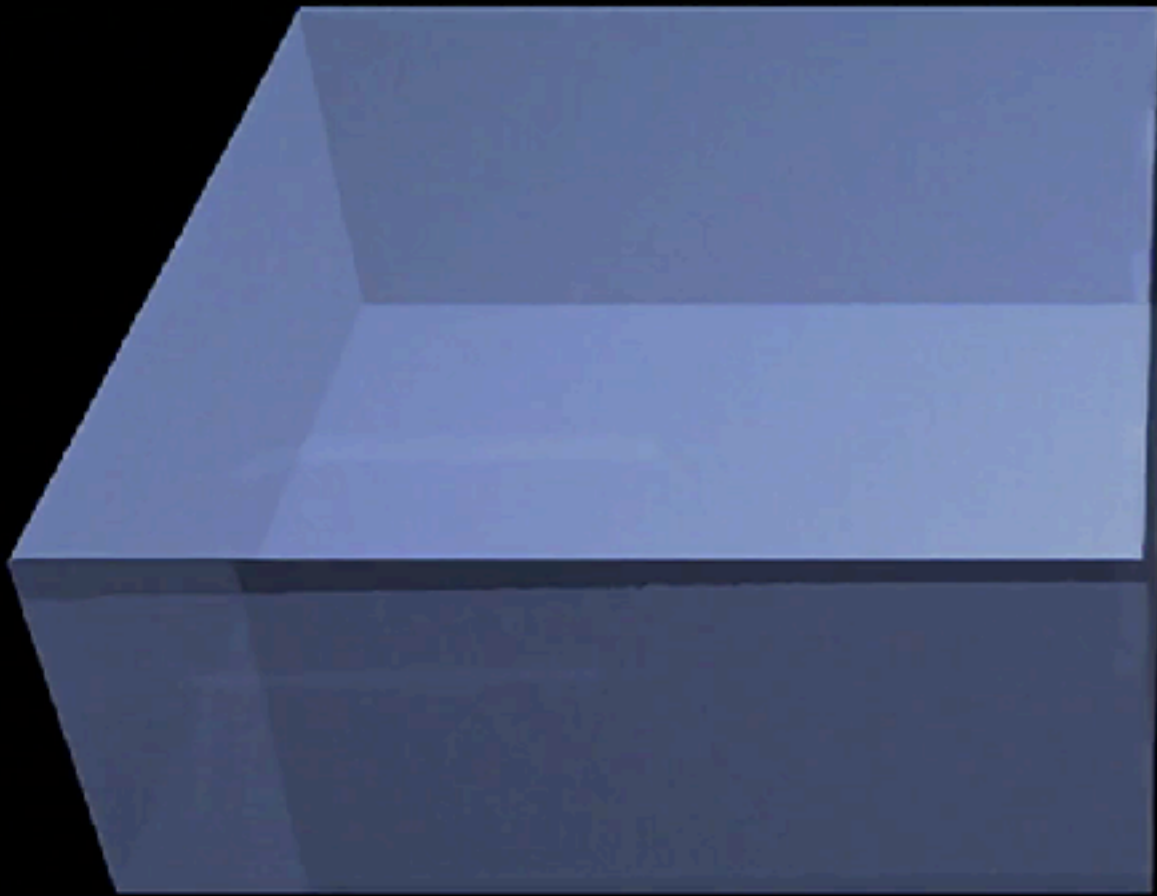




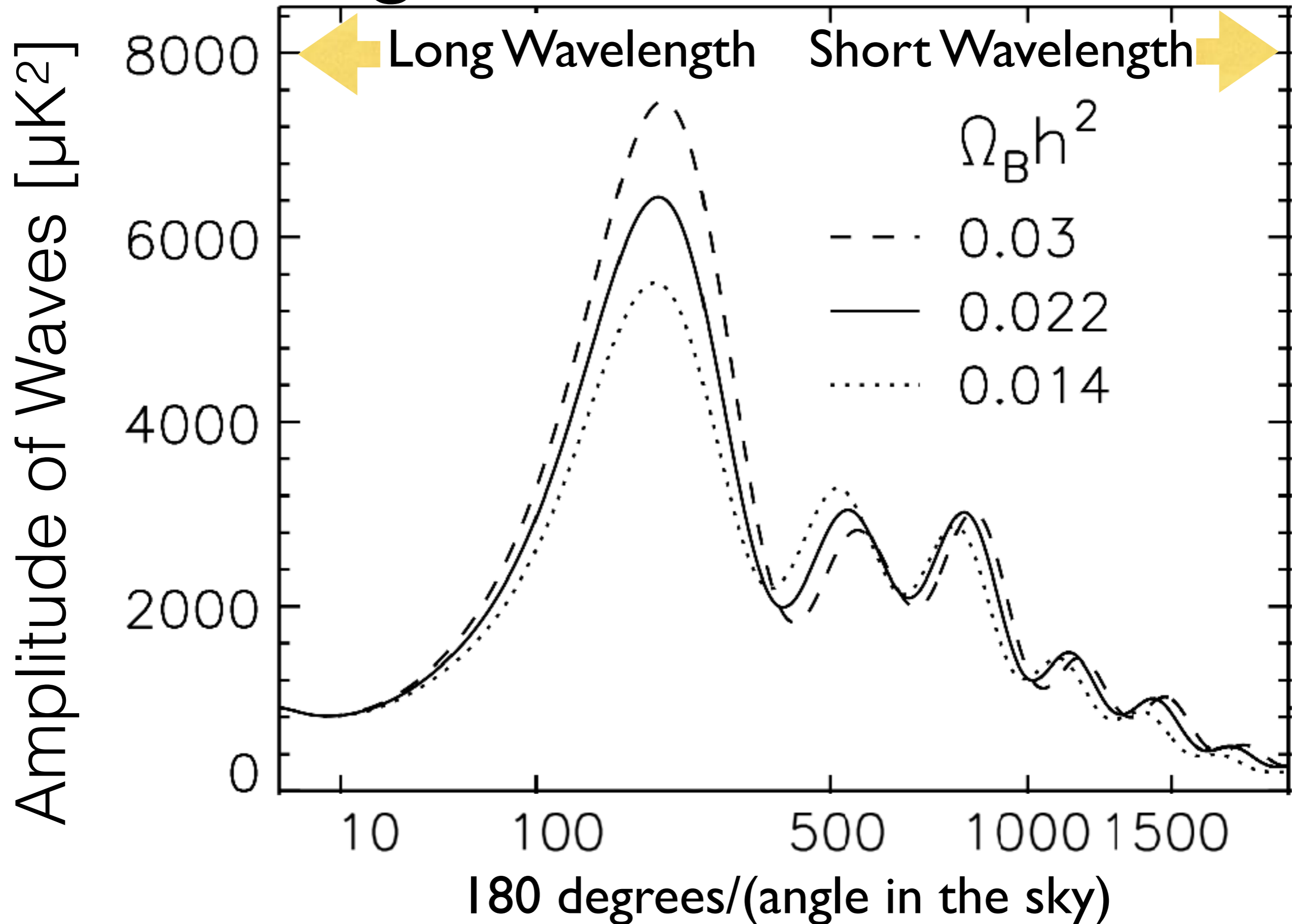


Cosmic Miso Soup

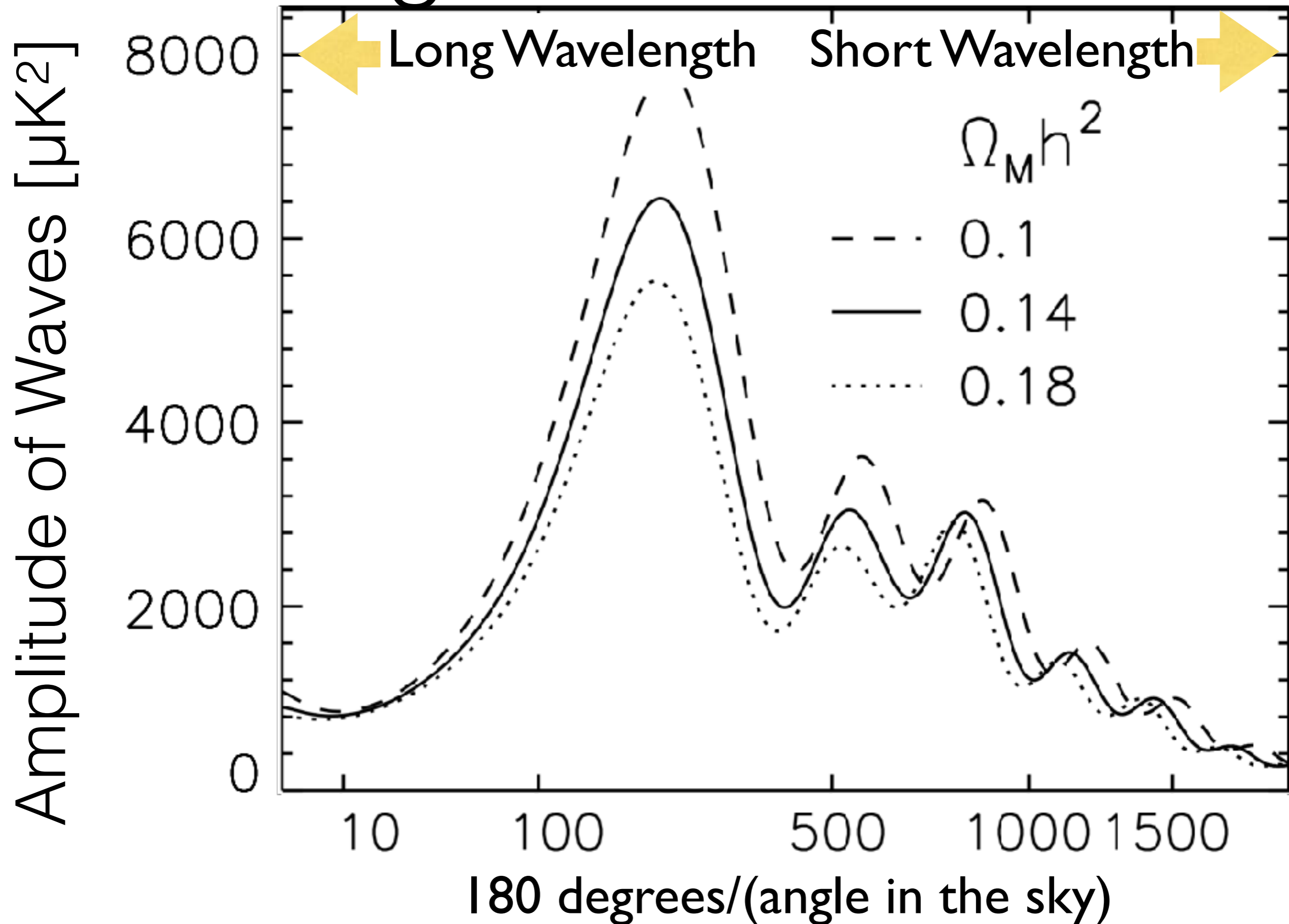
- When matter and radiation were hotter than 3000 K, matter was completely ionised. The Universe was filled with plasma, which behaves just like a soup
- Think about a Miso soup (if you know what it is). Imagine throwing Tofus into a Miso soup, while changing the density of Miso
- And imagine watching how ripples are created and propagate throughout the soup



Measuring Abundance of H&He

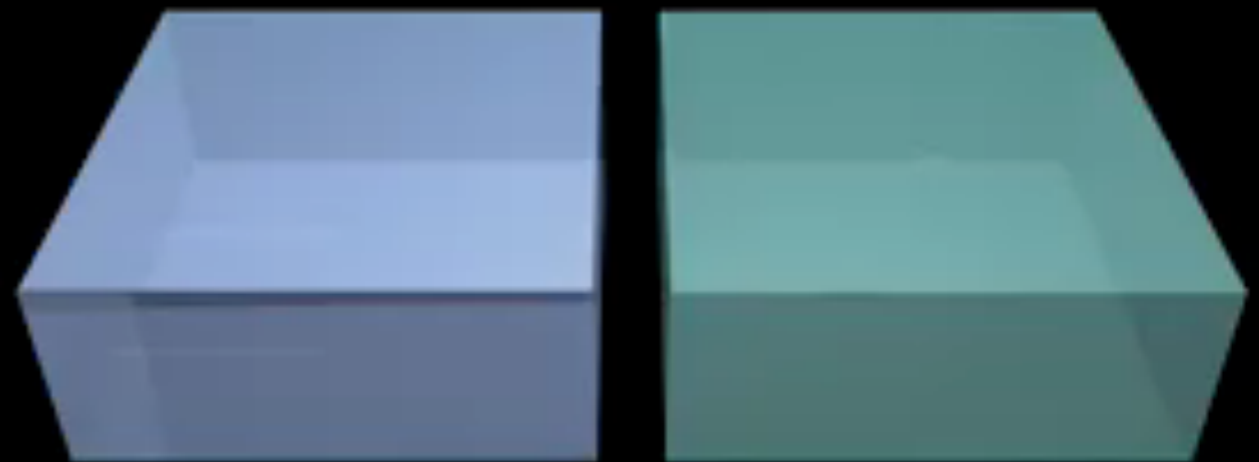


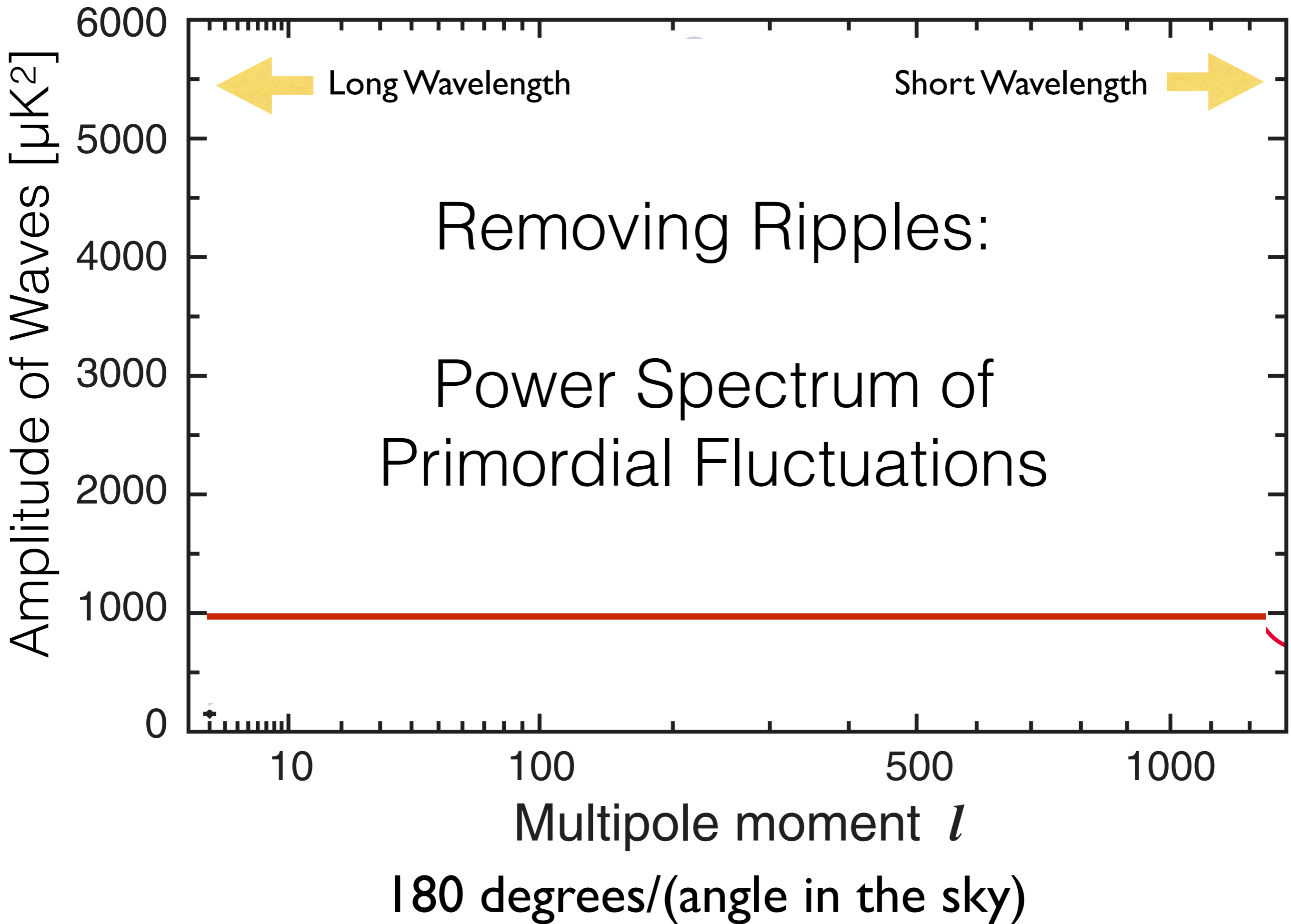
Measuring Total Matter Density

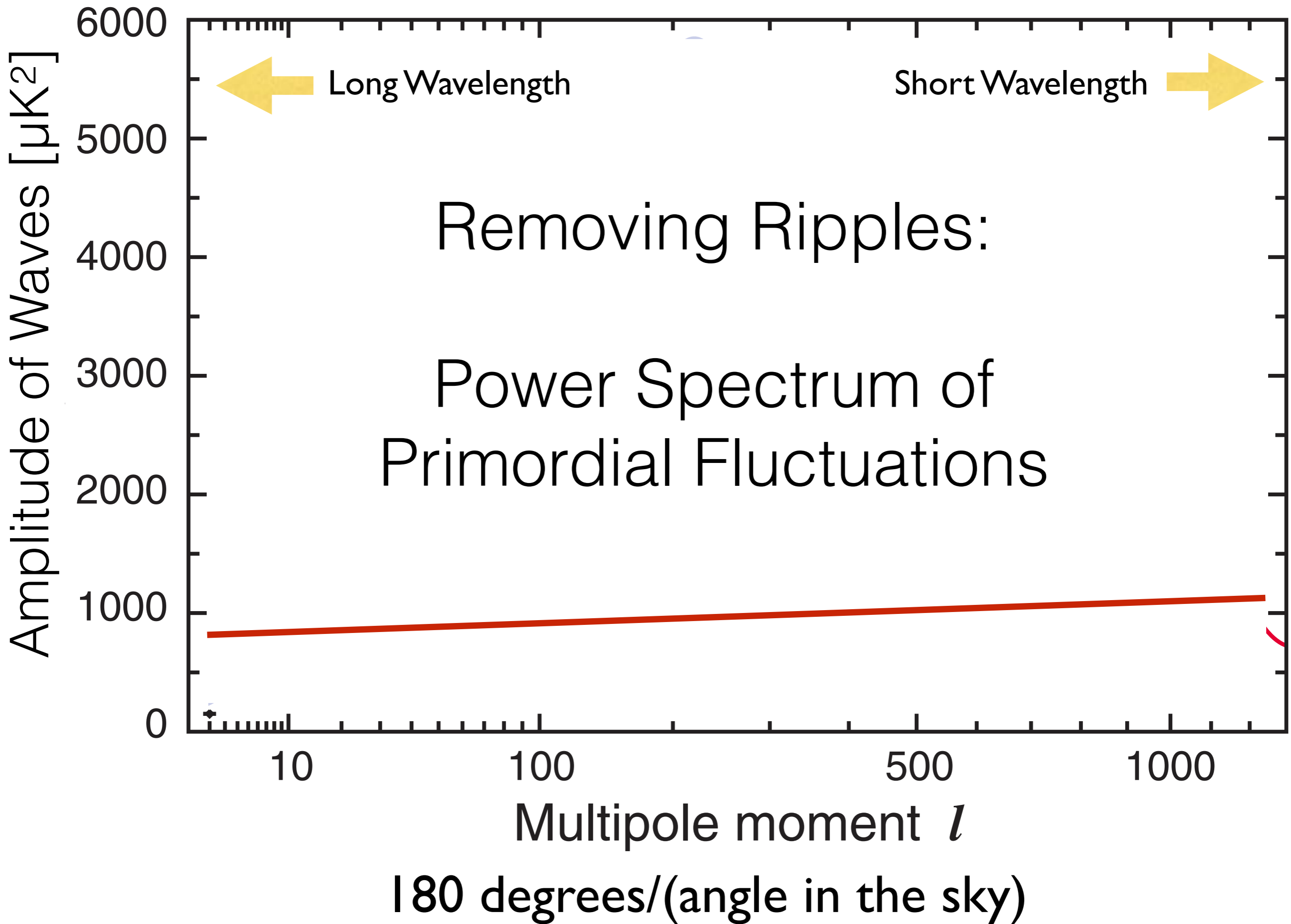


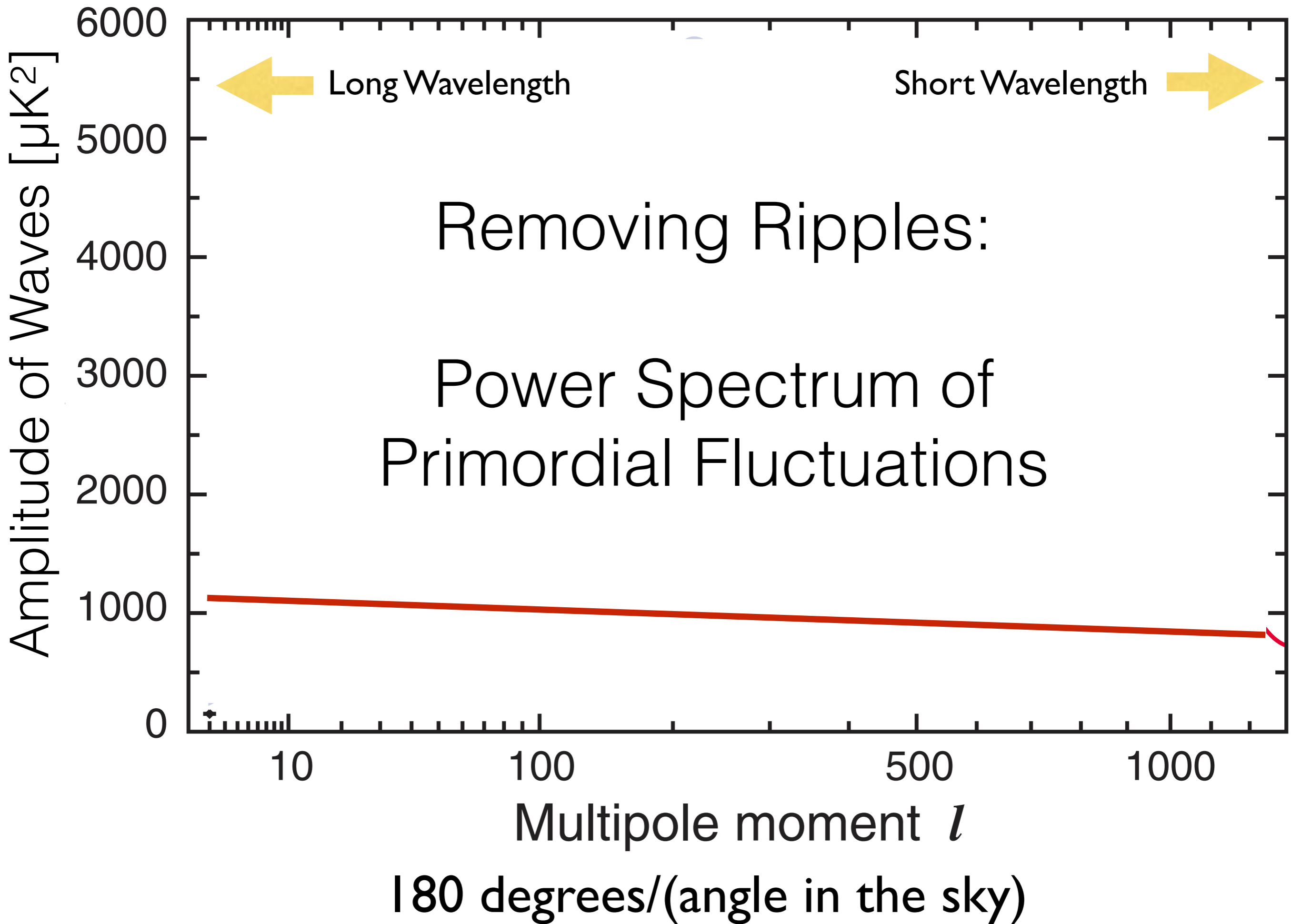
Origin of Fluctuations

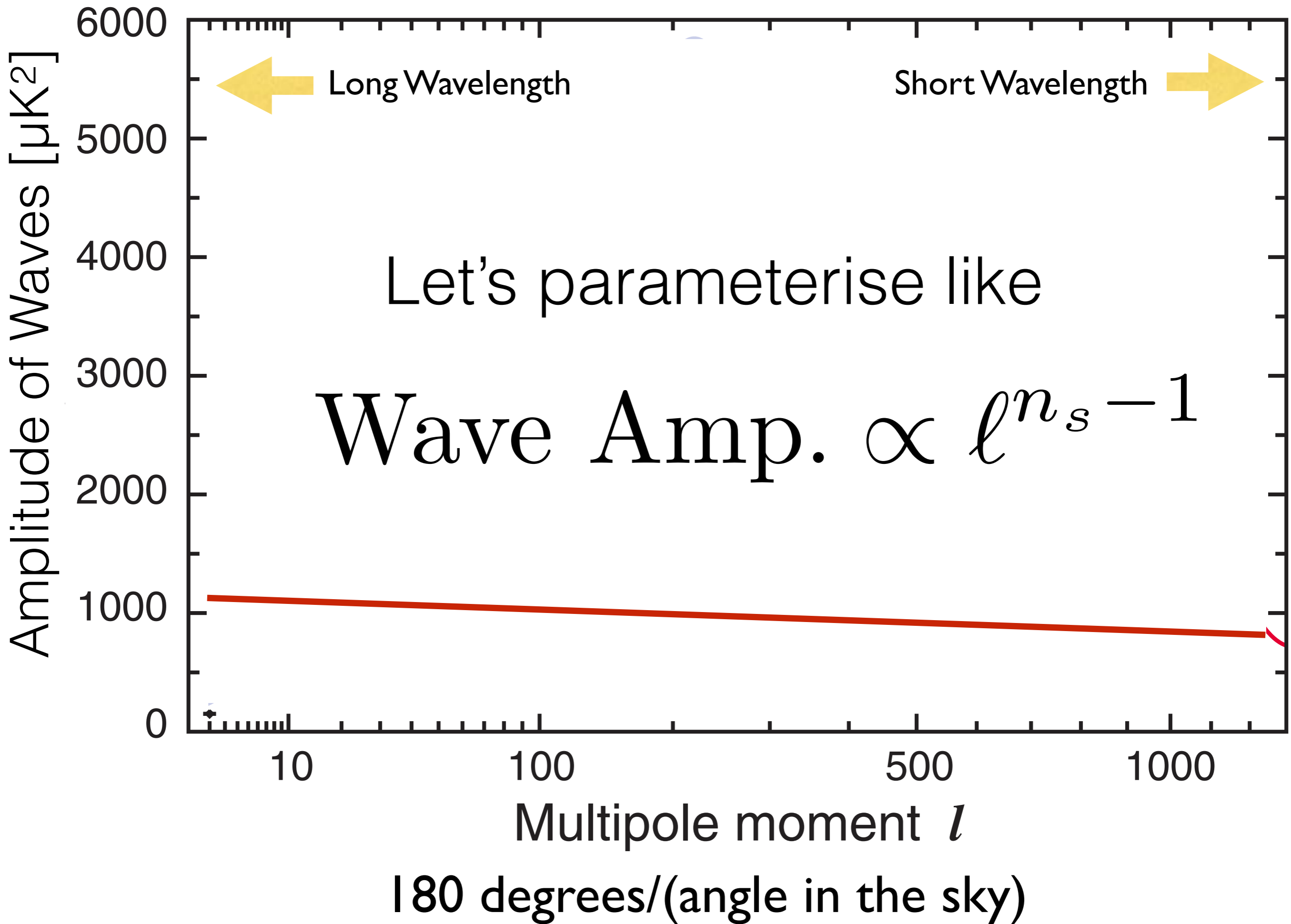
- Who dropped those Tofus into the cosmic Miso soup?

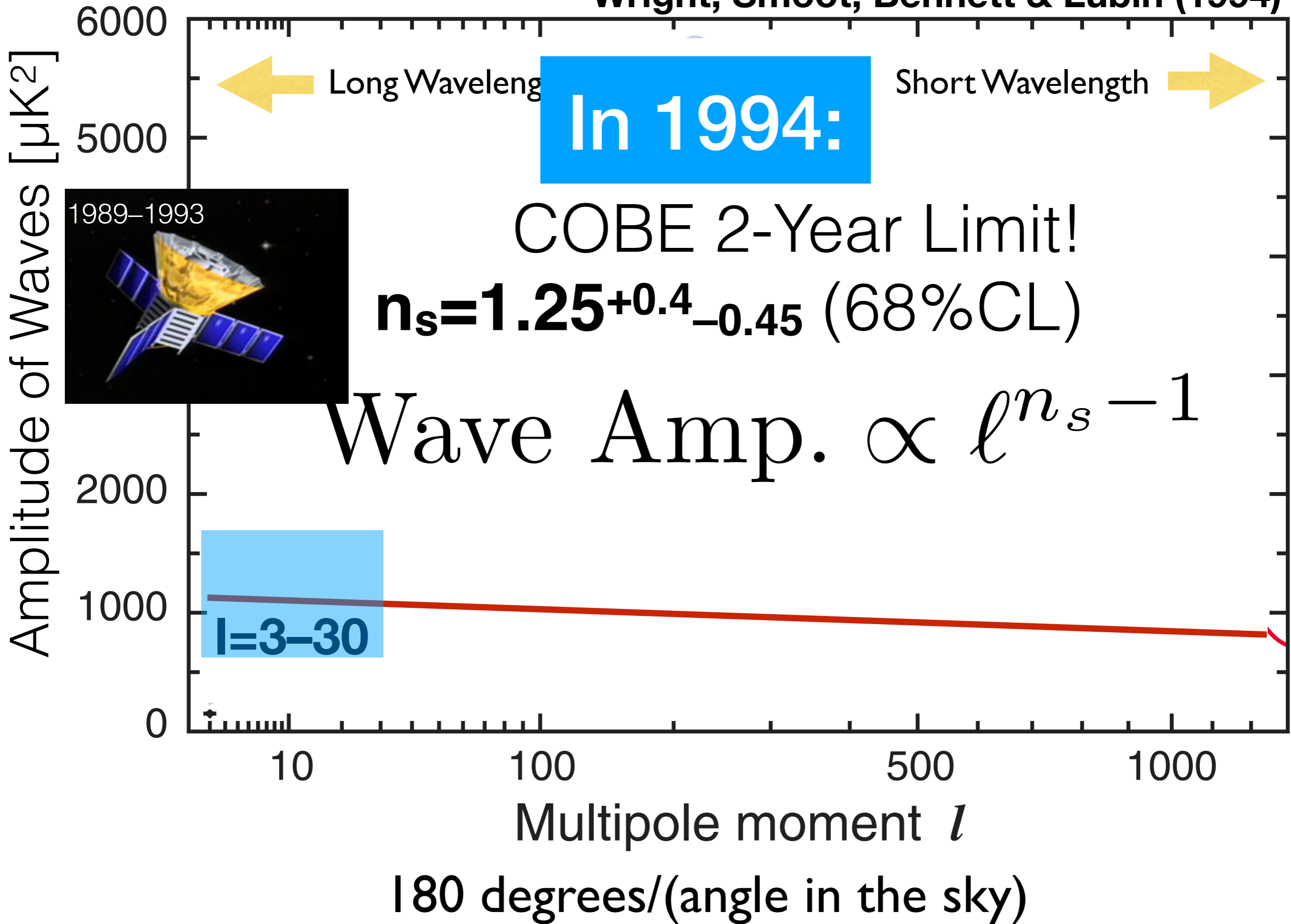


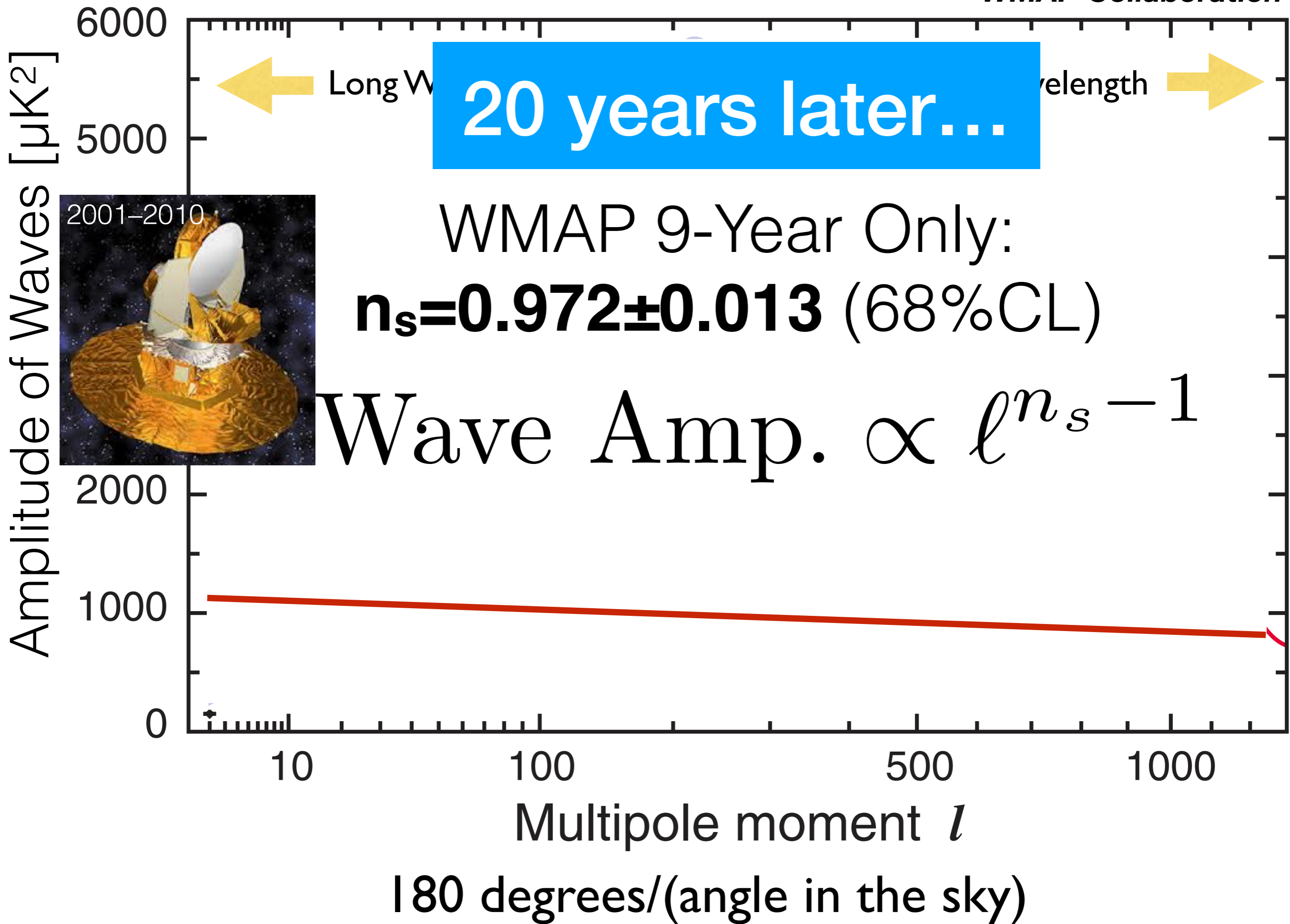












Angular scale

WMAP Collaboration

90° 2° 0.5° 0.2° 0.1°

Amplitude of ΔC_{ℓ}^2

2001–2010

South Pole Telescope
[10-m in South Pole]



$n_s = 0.965 \pm 0.010$

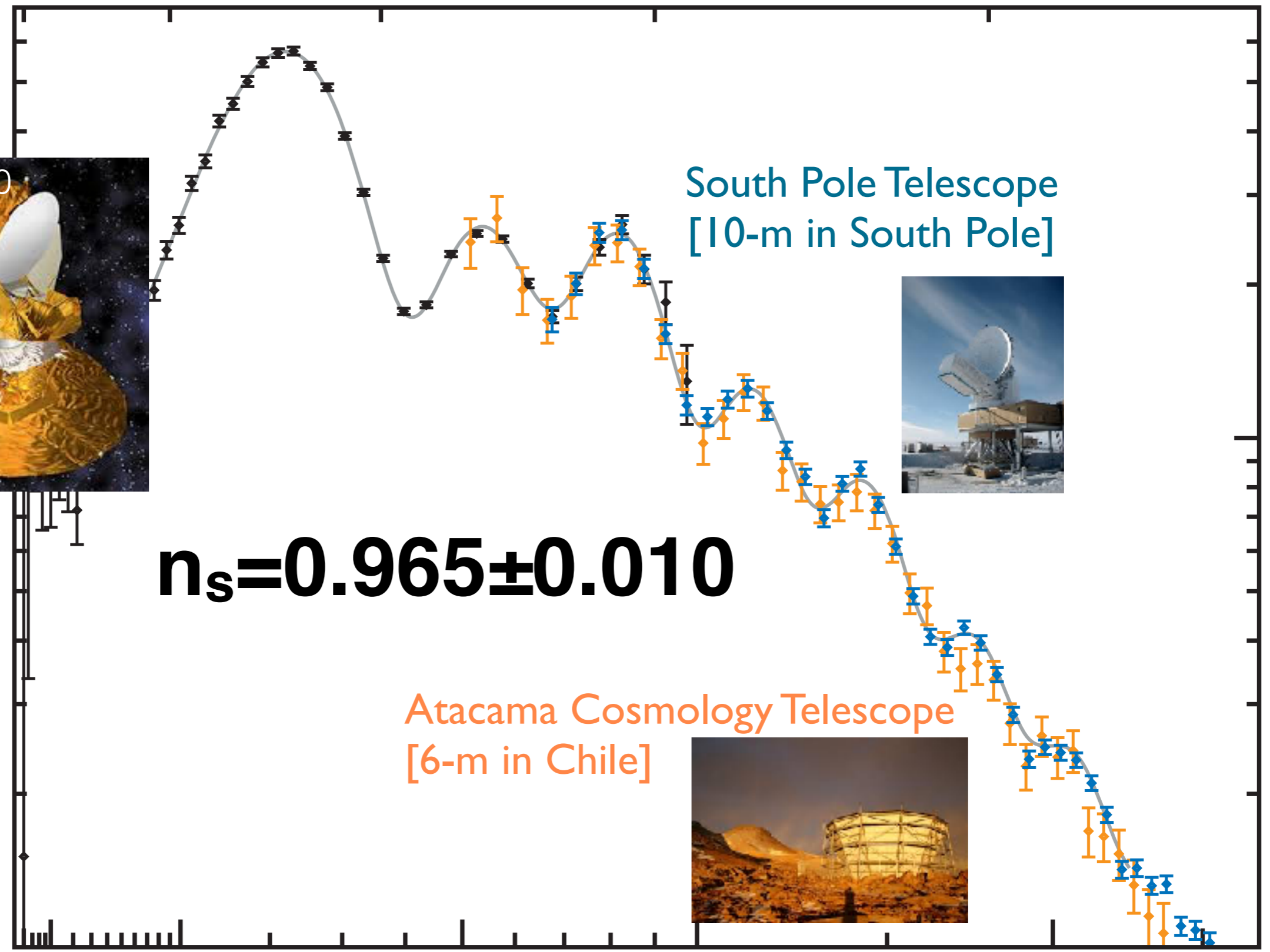
Atacama Cosmology Telescope
[6-m in Chile]



100

10 100 500 1000 2000

Multipole moment l



Angular scale

WMAP Collaboration

90°

2°

0.5°

0.2°

0.1°

Amplitude of ΔC_{ℓ}^2

2001–2010

South Pole Telescope
[10-m in South Pole]

$n_s = 0.961 \pm 0.008$

~5 σ discovery of $n_s < 1$ from the
CMB data combined with the
distribution of galaxies

Atacama Cosmology Telescope
[6-m in Chile]

100

10

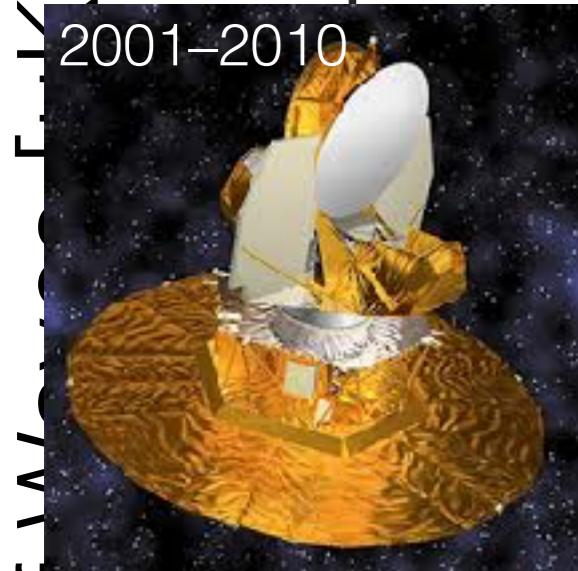
100

500

1000

2000

Multipole moment l



Residual Amplitude of Waves [μK^2]

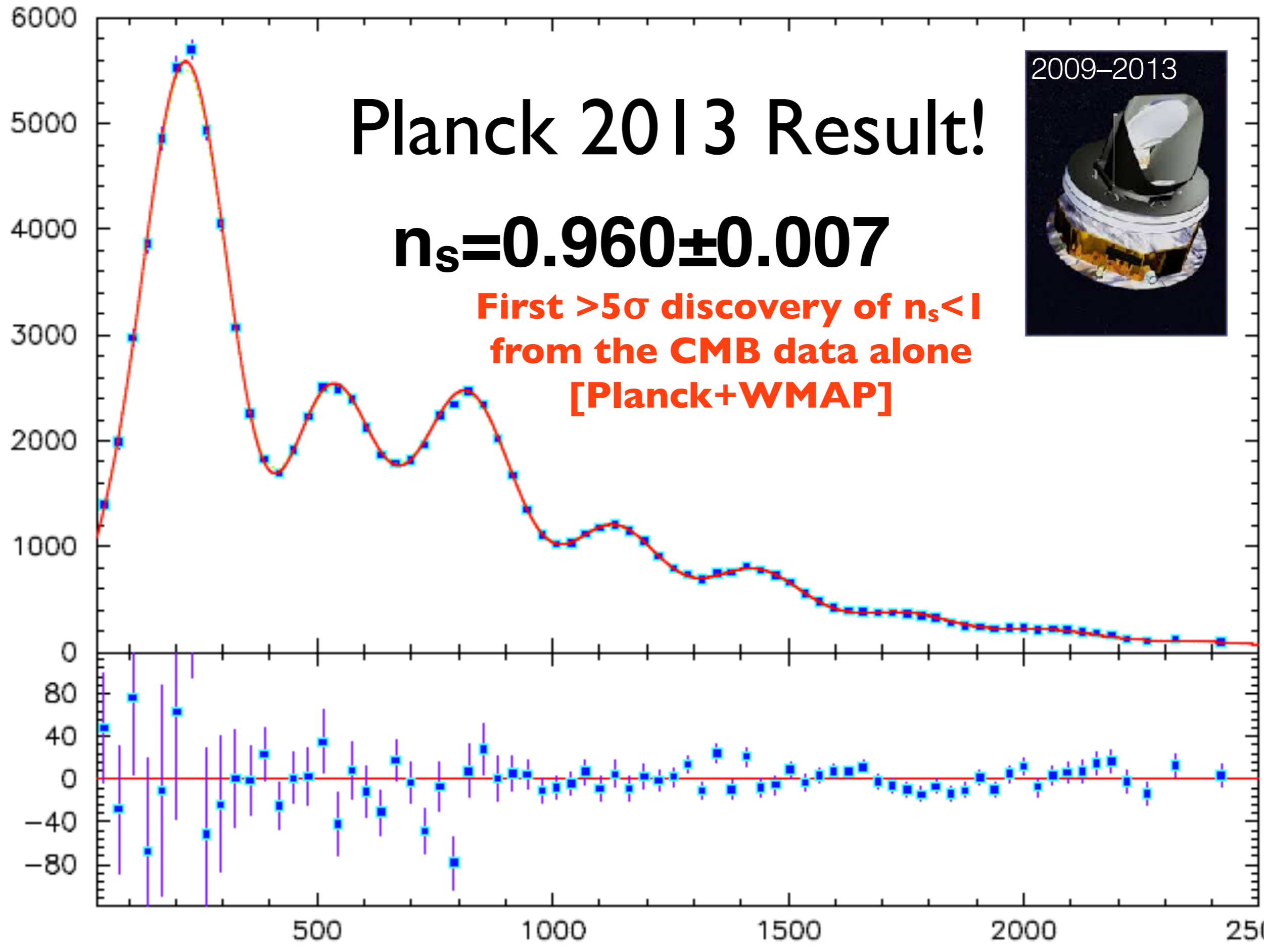
2009–2013



Planck 2013 Result!

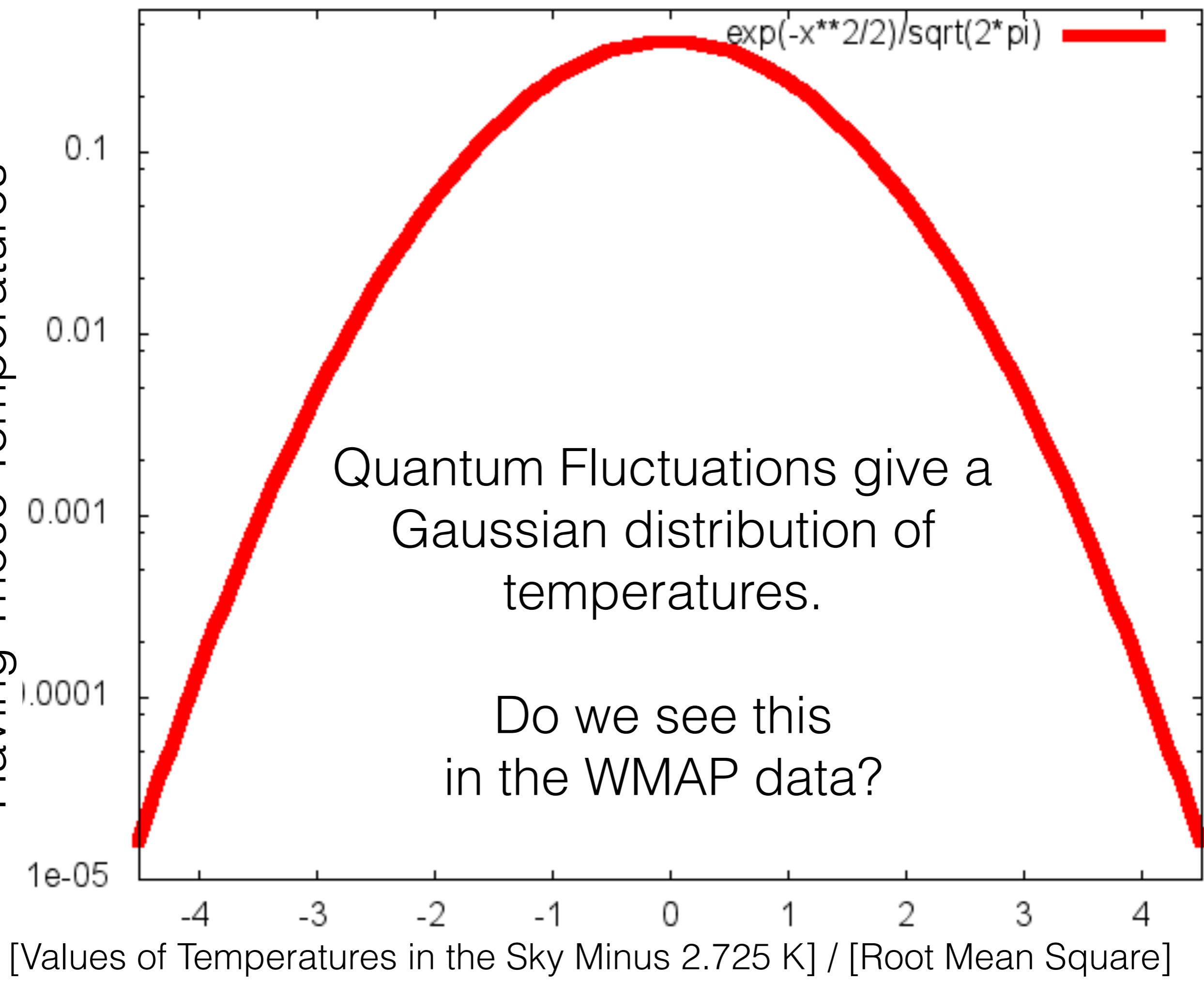
$$n_s = 0.960 \pm 0.007$$

First $>5\sigma$ discovery of $n_s < 1$
from the CMB data alone
[Planck+WMAP]



l 80 degrees/(angle in the sky)

Fraction of the Number of Pixels
Having Those Temperatures



Fraction of the Number of Pixels
Having Those Temperatures

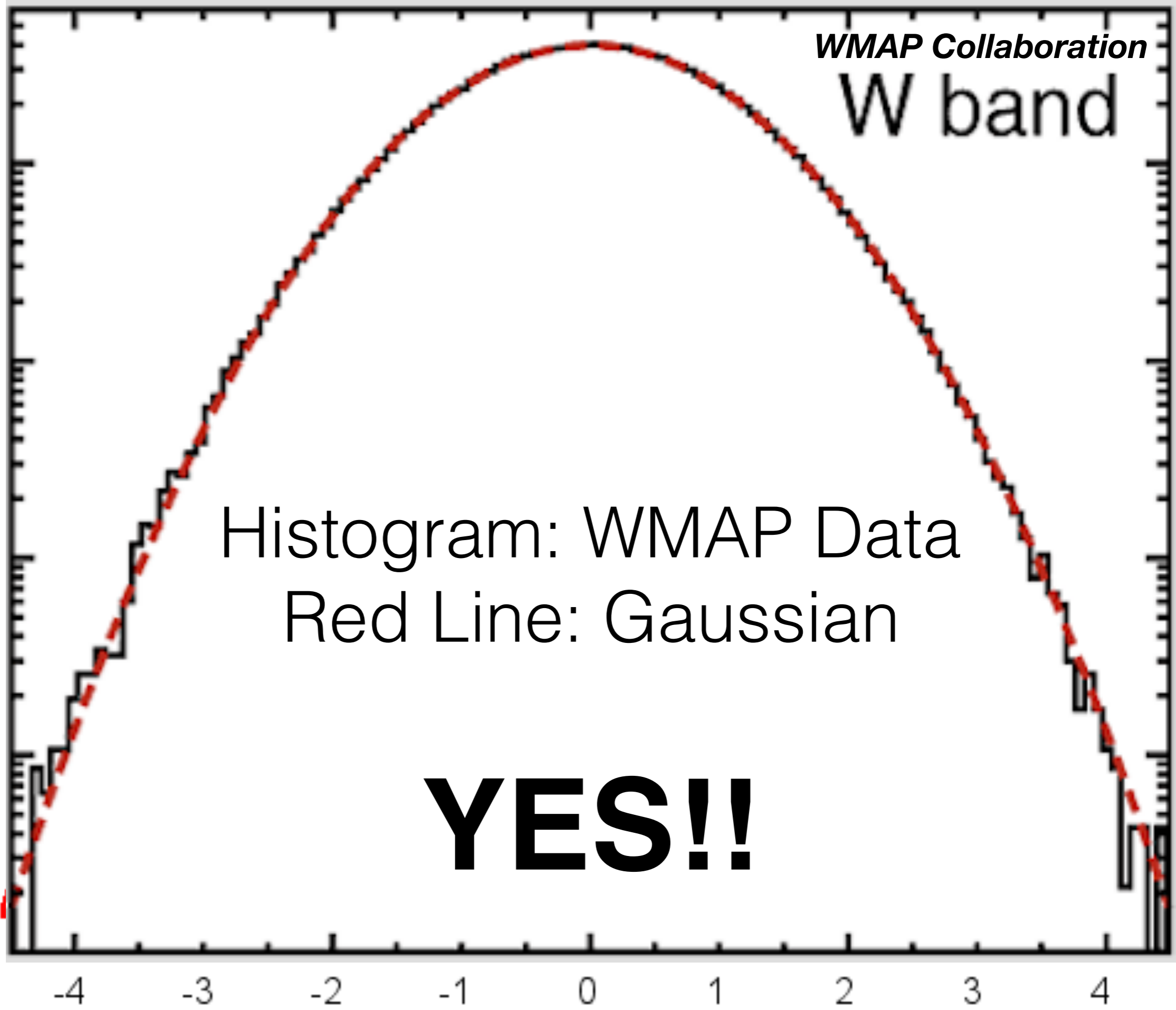
WMAP Collaboration
W band

0.1
0.01
0.001
0.0001
1e-05

Histogram: WMAP Data
Red Line: Gaussian

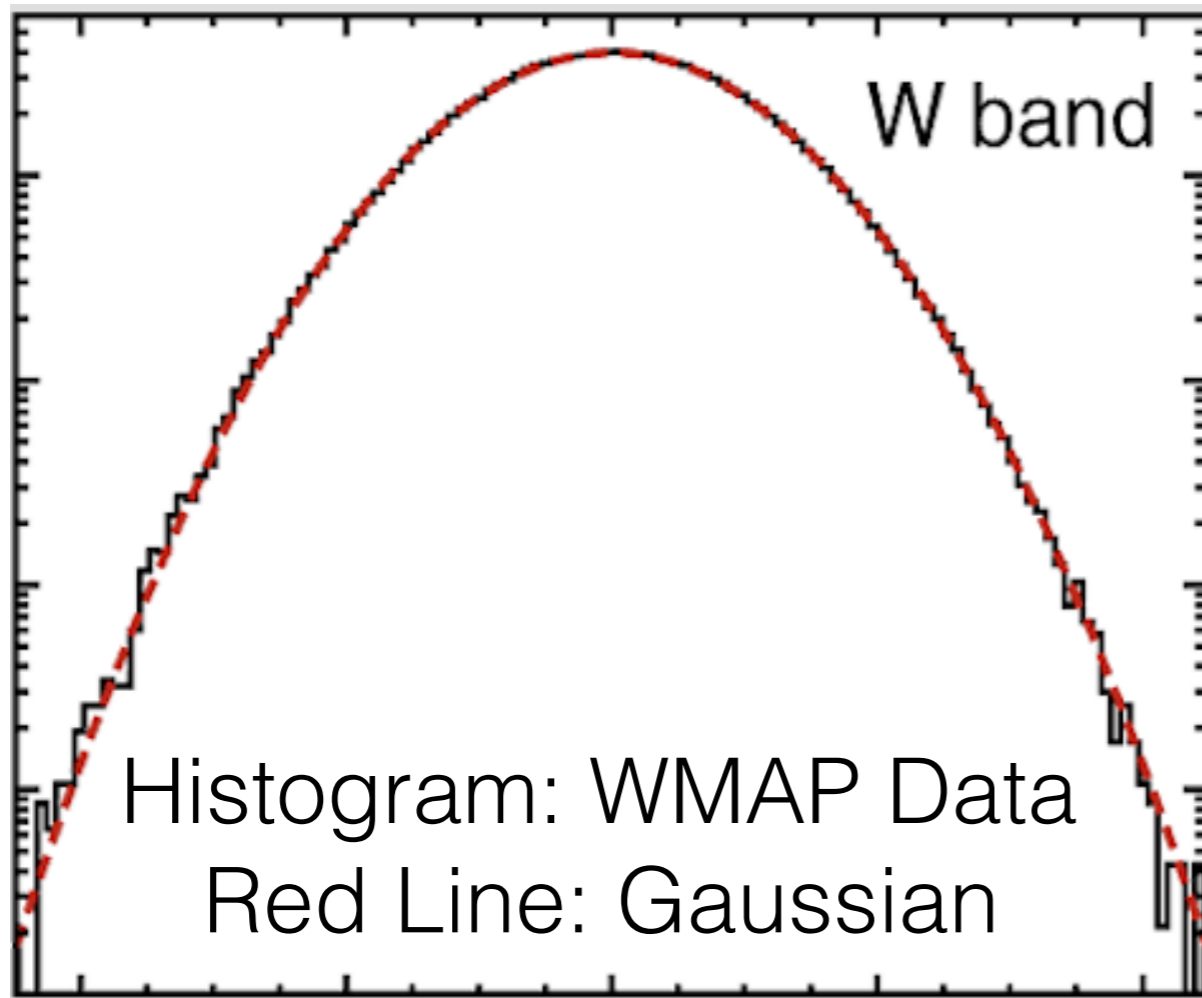
YES!!

[Values of Temperatures in the Sky Minus 2.725 K] / [Root Mean Square]



Testing Gaussianity

Fraction of the Number of Pixels
Having Those Temperatures



[Values of Temperatures in the Sky Minus
2.725 K]/ [Root Mean Square]

- Since a Gauss distribution is symmetric, it must yield a vanishing **3-point function**

$$\langle \delta T^3 \rangle \equiv \int_{-\infty}^{\infty} d\delta T P(\delta T) \delta T^3$$

- More specifically, we measure this by averaging the product of temperatures at three different locations in the sky

$$\langle \delta T(\hat{n}_1) \delta T(\hat{n}_2) \delta T(\hat{n}_3) \rangle$$

Lack of non-Gaussianity

- The WMAP data show that the distribution of temperature fluctuations of CMB is very precisely Gaussian
 - with an upper bound on a deviation of **0.2%** (95%CL)

$$\zeta(\mathbf{x}) = \zeta_{\text{gaus}}(\mathbf{x}) + \frac{3}{5} f_{\text{NL}} \zeta_{\text{gaus}}^2(\mathbf{x}) \text{ with } f_{\text{NL}} = 37 \pm 20 \text{ (68\% CL)}$$

WMAP 9-year Result

- The Planck data improved the upper bound by an order of magnitude: deviation is **<0.03%** (95%CL)

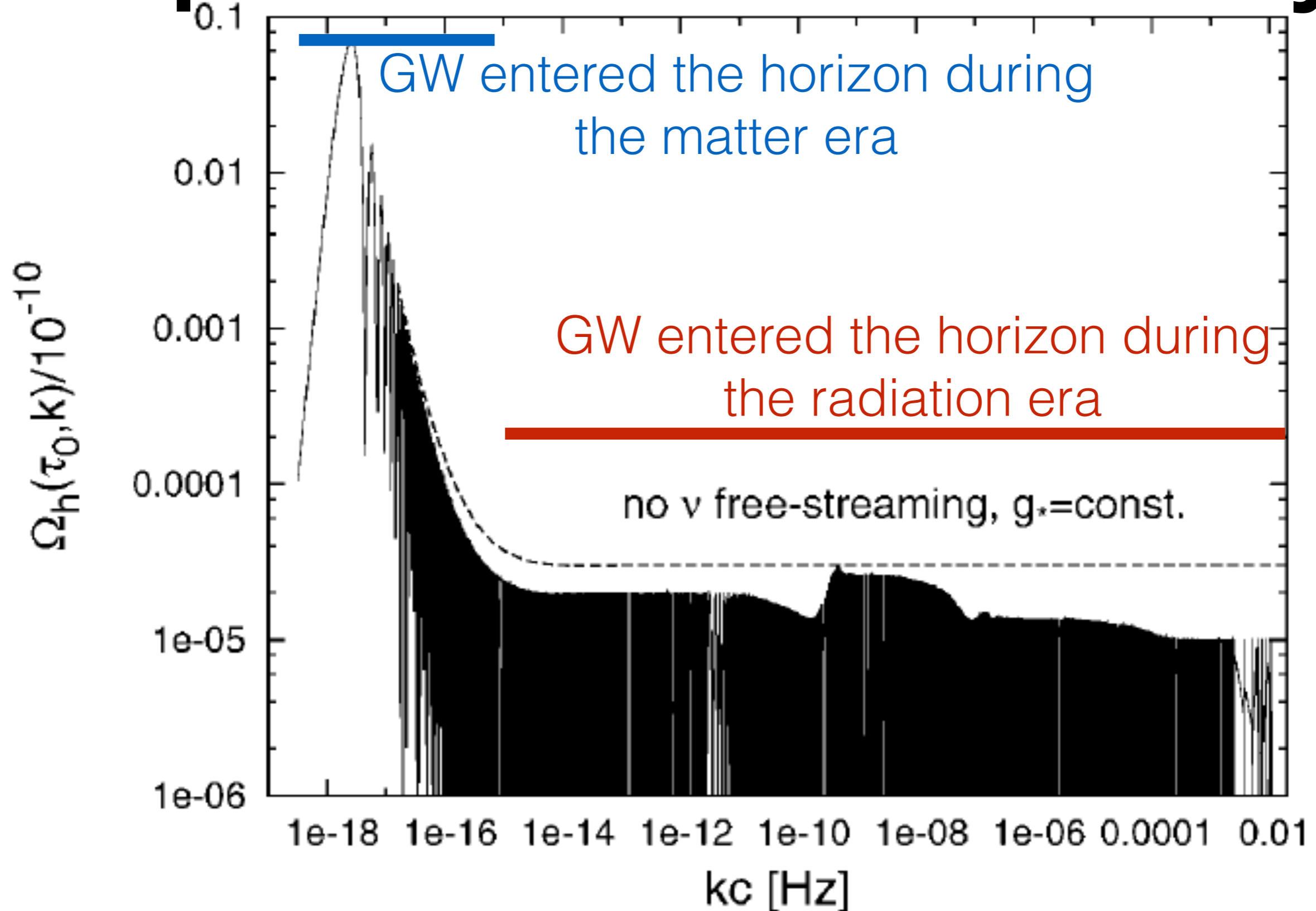
$$f_{\text{NL}} = 0.8 \pm 5.0 \text{ (68\% CL)}$$

Planck 2015 Result

So, have we found inflation?

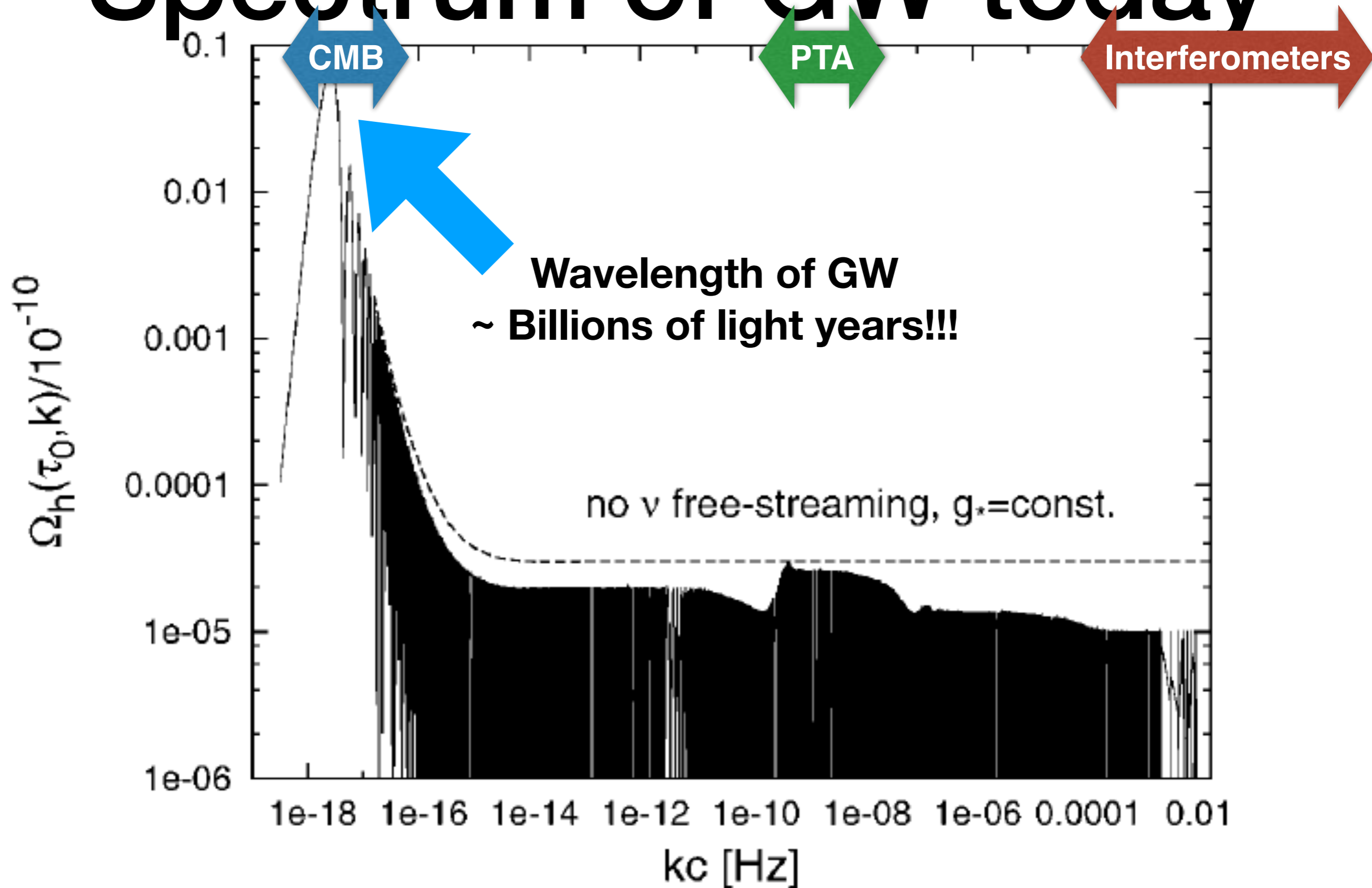
- Single-field slow-roll inflation looks remarkably good:
 - **Super-horizon fluctuation**
 - **Adiabaticity**
 - **Gaussianity**
 - **$n_s < 1$**
- What more do we want? **Gravitational waves**. Why?
 - Because the “*extraordinary claim requires extraordinary evidence*”

Theoretical energy density Spectrum of GW today



Theoretical energy density

Spectrum of GW today



Finding Signatures of Gravitational Waves in the CMB

- **Next frontier in the CMB research**
 1. Find evidence for nearly scale-invariant gravitational waves
 2. Once found, test Gaussianity to make sure (or not!) that the signal comes from vacuum fluctuation
 3. Constrain inflation models

New
Research
Area!

Measuring GW

- GW changes distances between two points

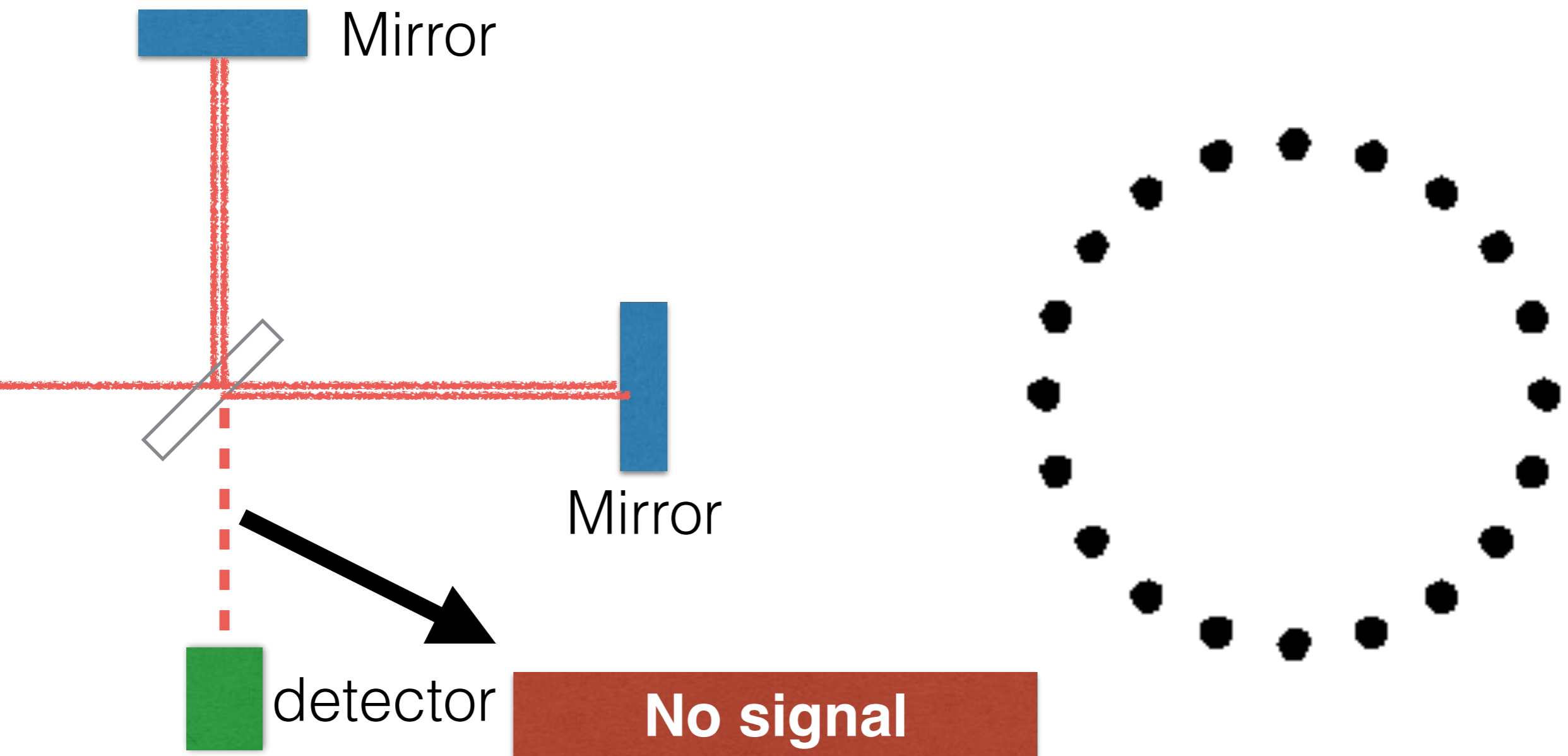
$$d\ell^2 = d\mathbf{x}^2 = \sum_{ij} \delta_{ij} dx^i dx^j$$



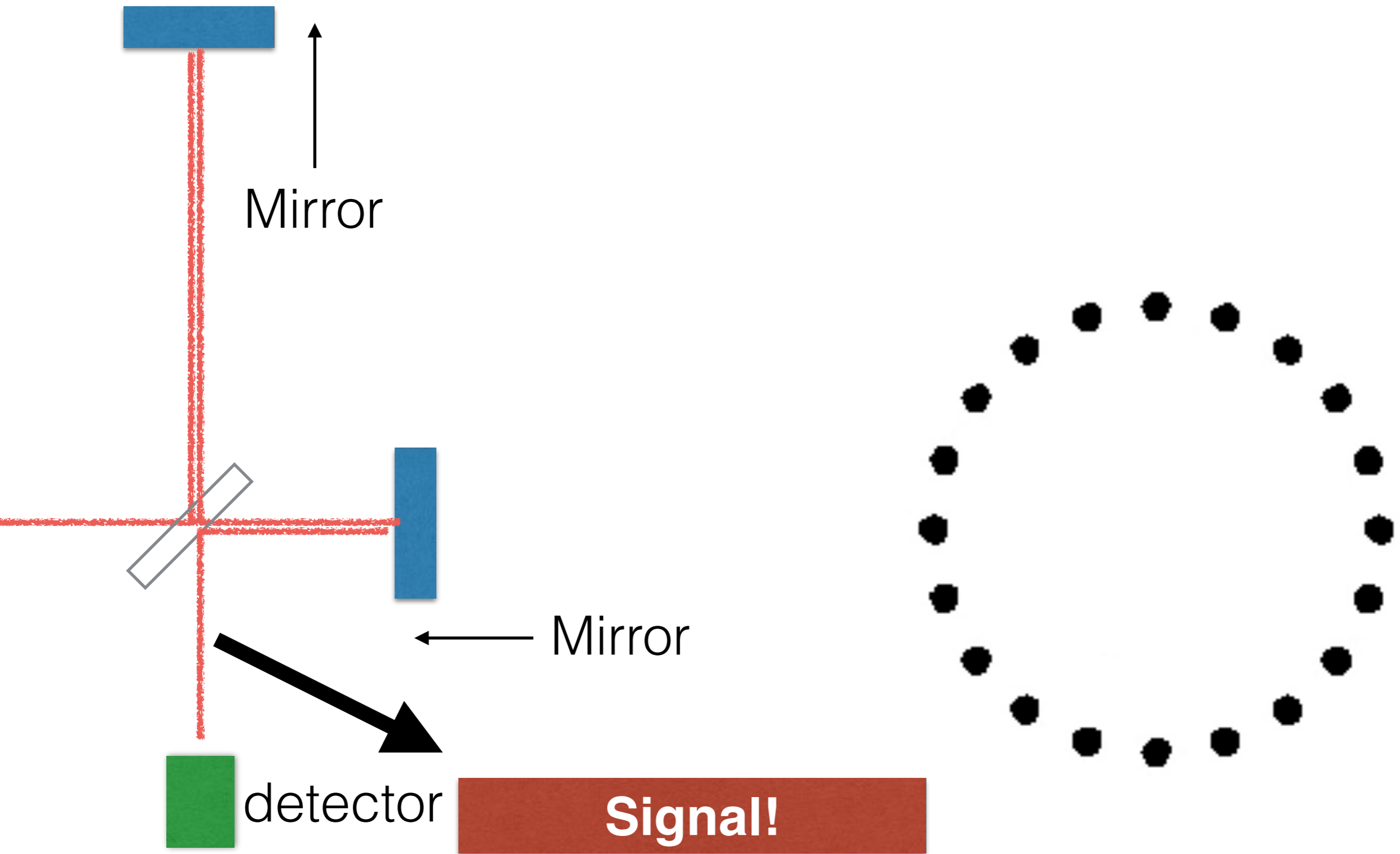
$$d\ell^2 = \sum_{ij} (\delta_{ij} + \underline{h_{ij}}) dx^i dx^j$$



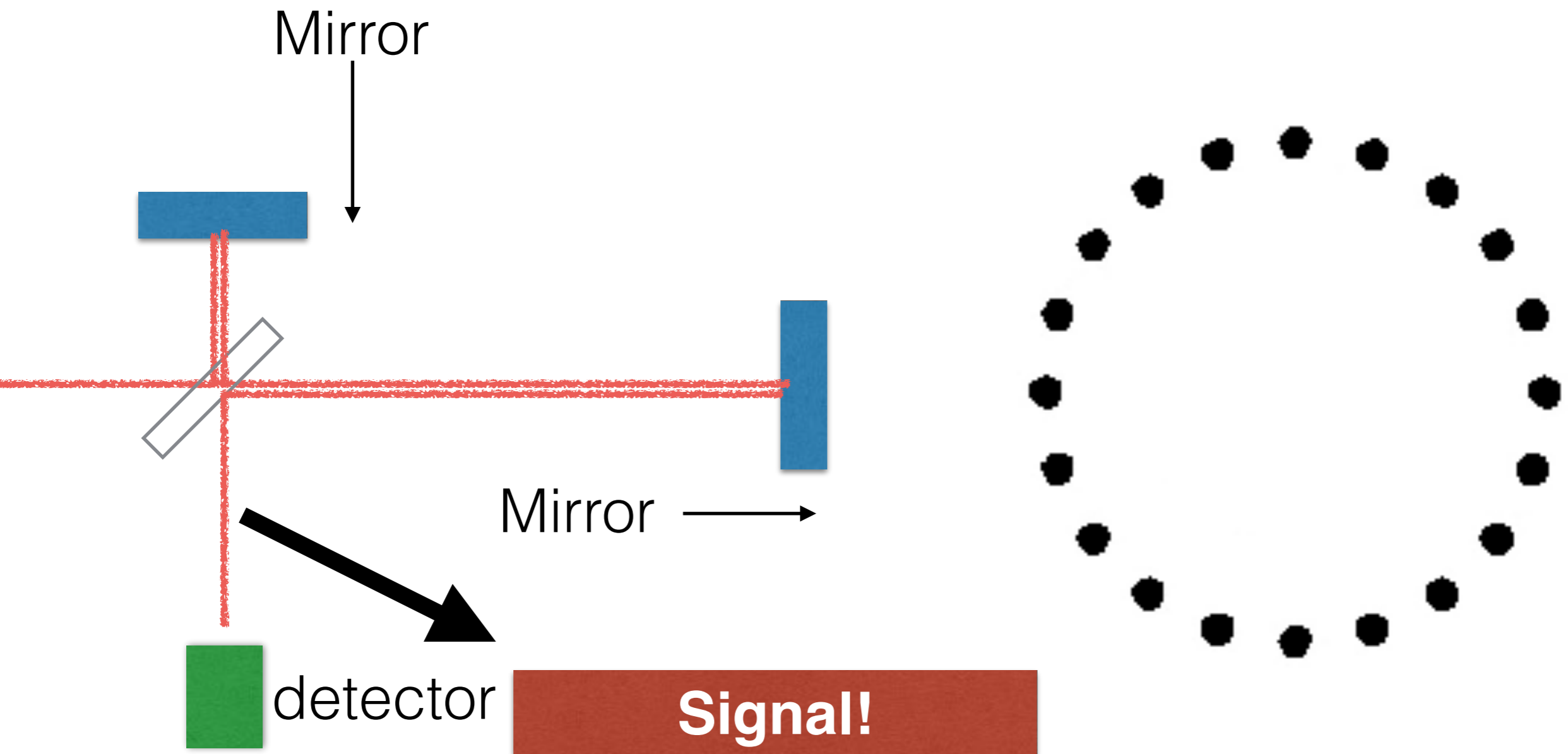
Laser Interferometer



Laser Interferometer



Laser Interferometer



LIGO detected GW from a binary blackholes, with the wavelength of thousands of kilometres

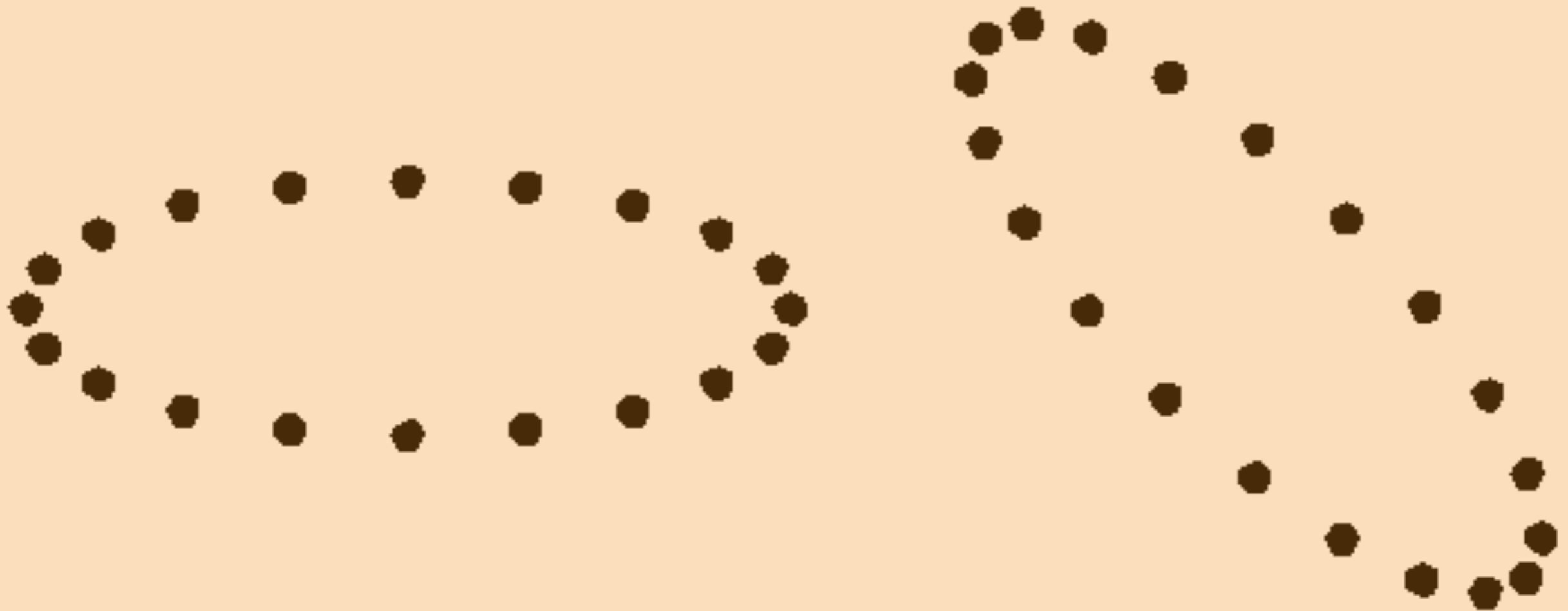
But, the primordial GW affecting the CMB has a wavelength of **billions of light-years!!** How do we find it?

Detecting GW by CMB

Isotropic electro-magnetic fields

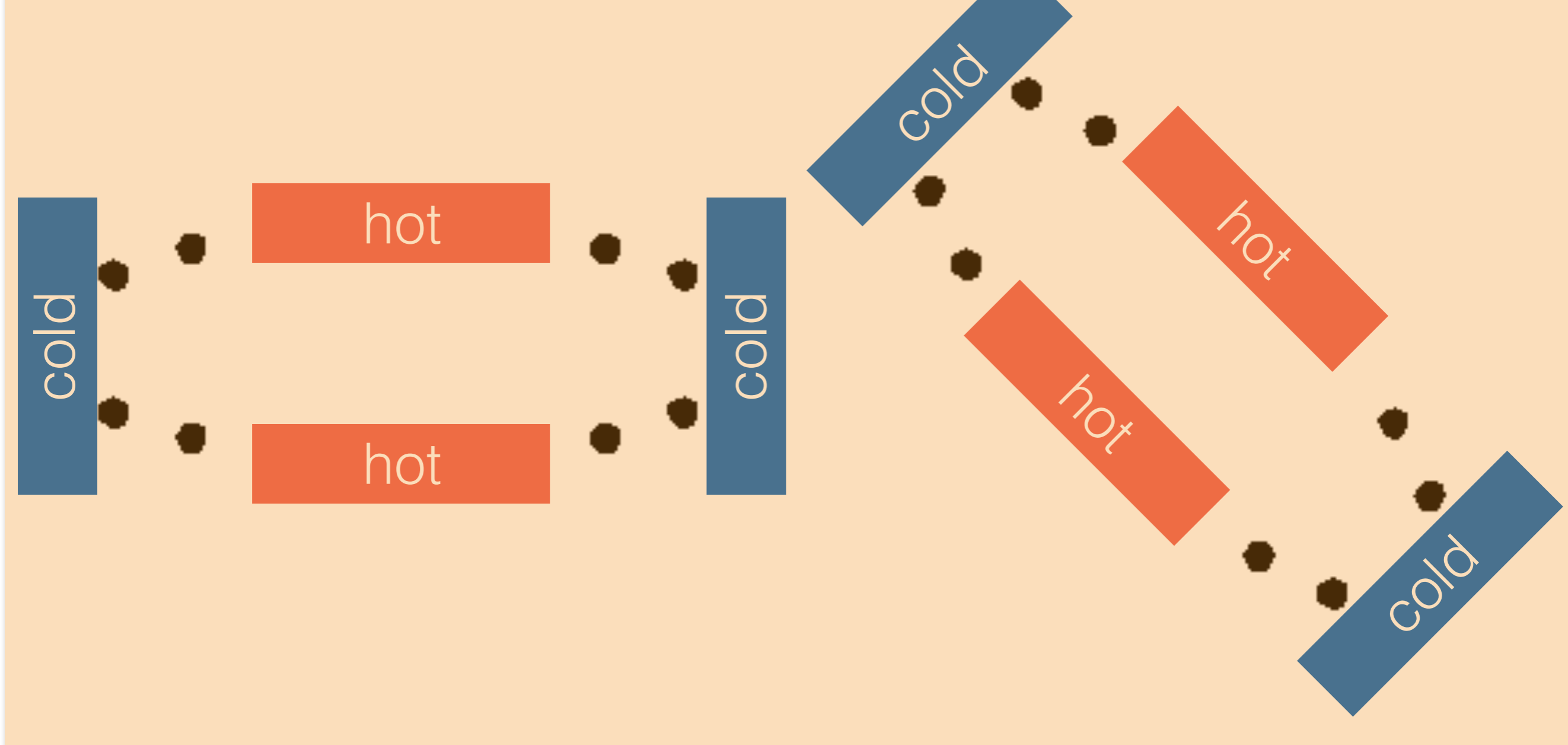
Detecting GW by CMB

GW propagating in isotropic electro-magnetic fields



Detecting GW by CMB

Space is stretched => Wavelength of light is also stretched



Detecting GW by CMB Polarisation

Space is stretched => Wavelength of light is also stretched

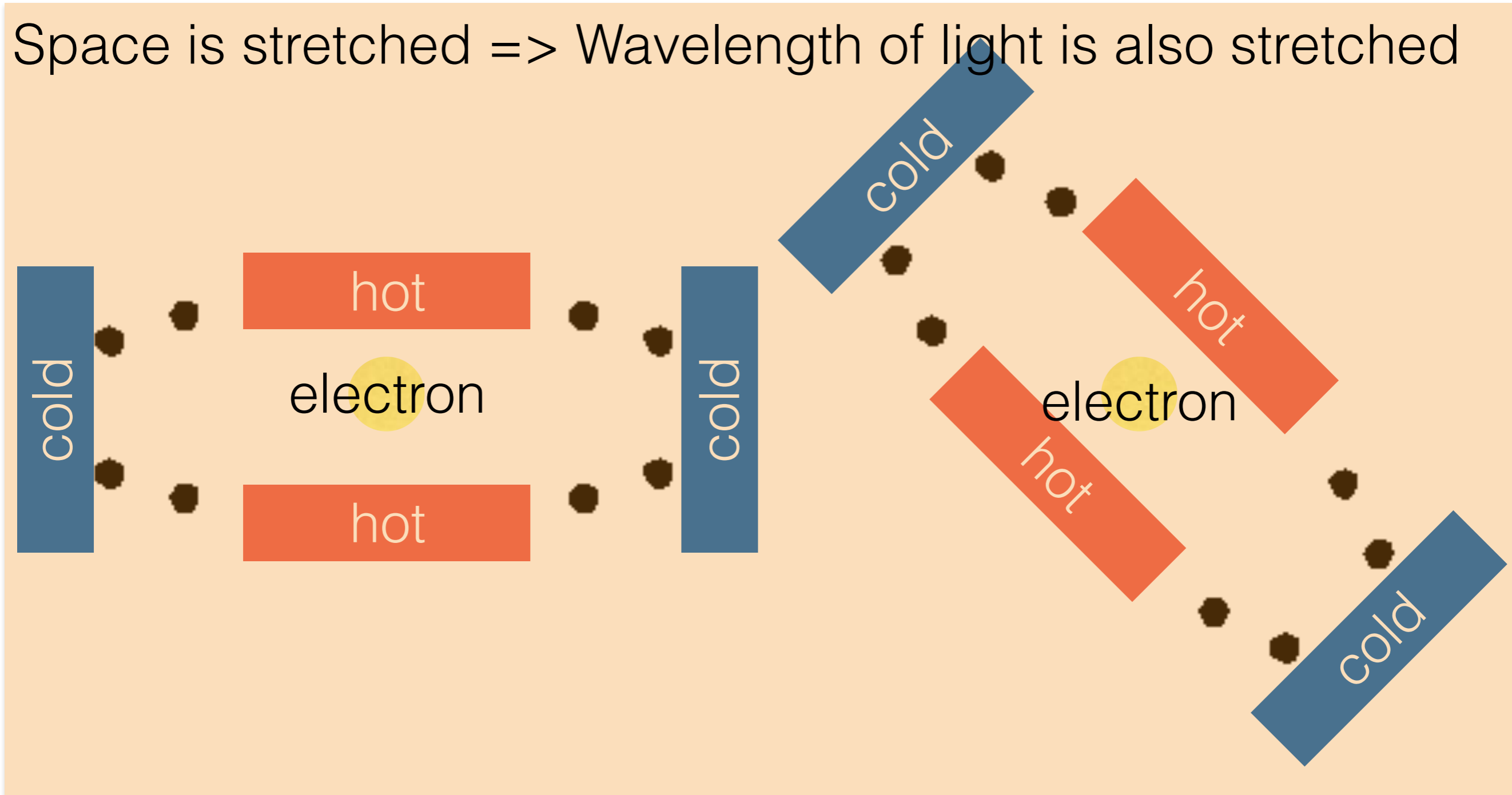


Photo Credit: TALEX



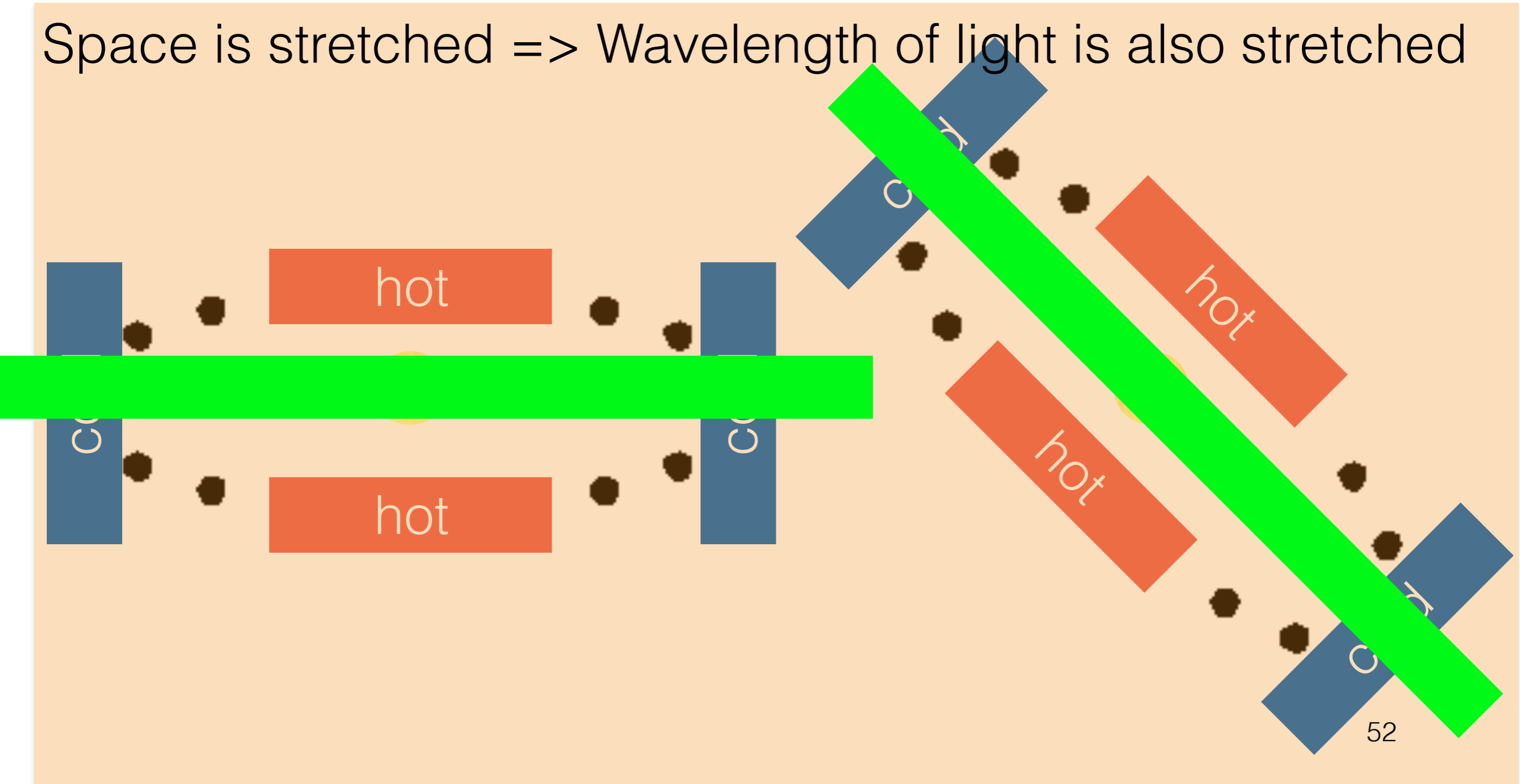
horizontally polarised

Photo Credit: TALEX



Detecting GW by CMB Polarisation

Space is stretched => Wavelength of light is also stretched

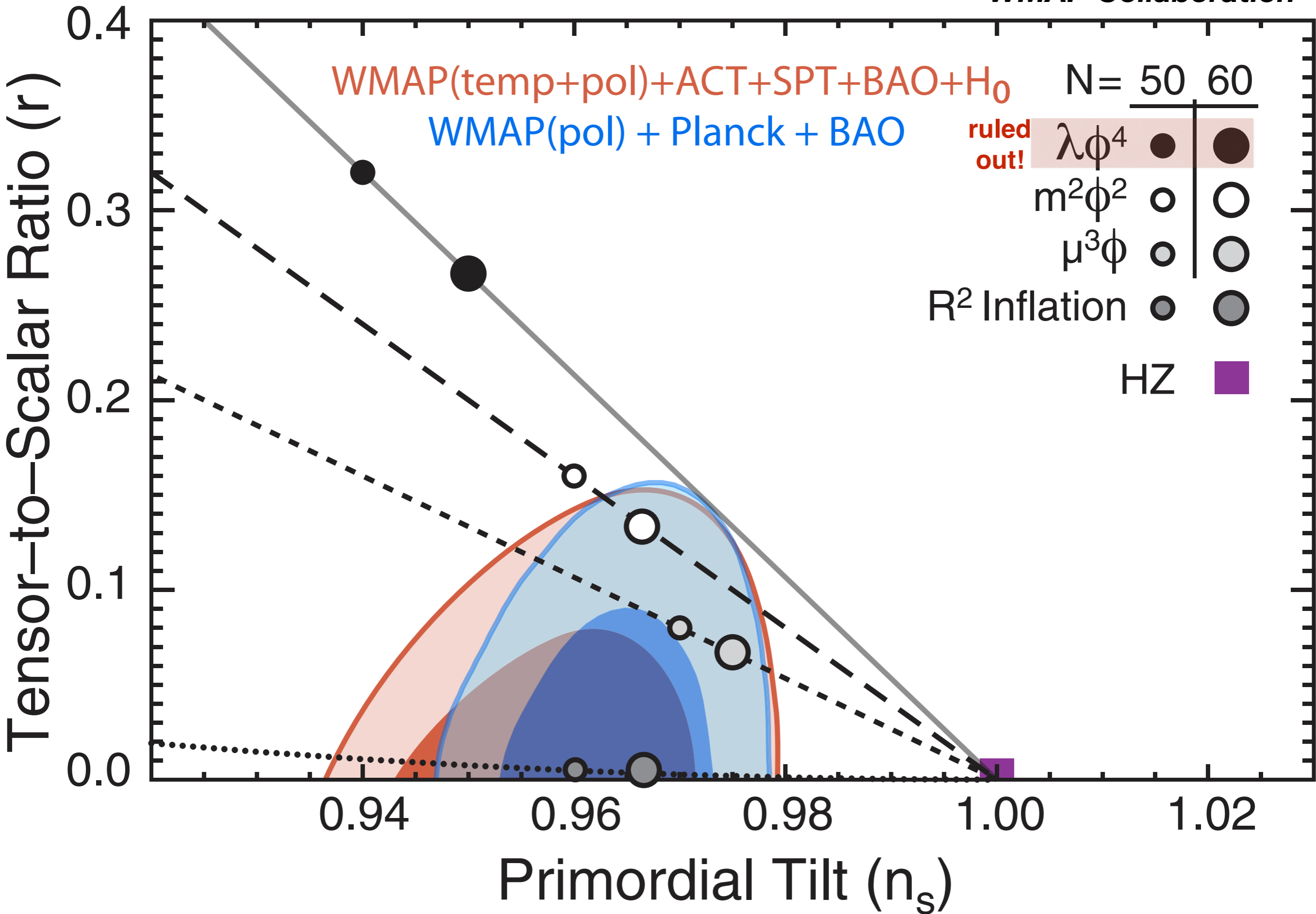


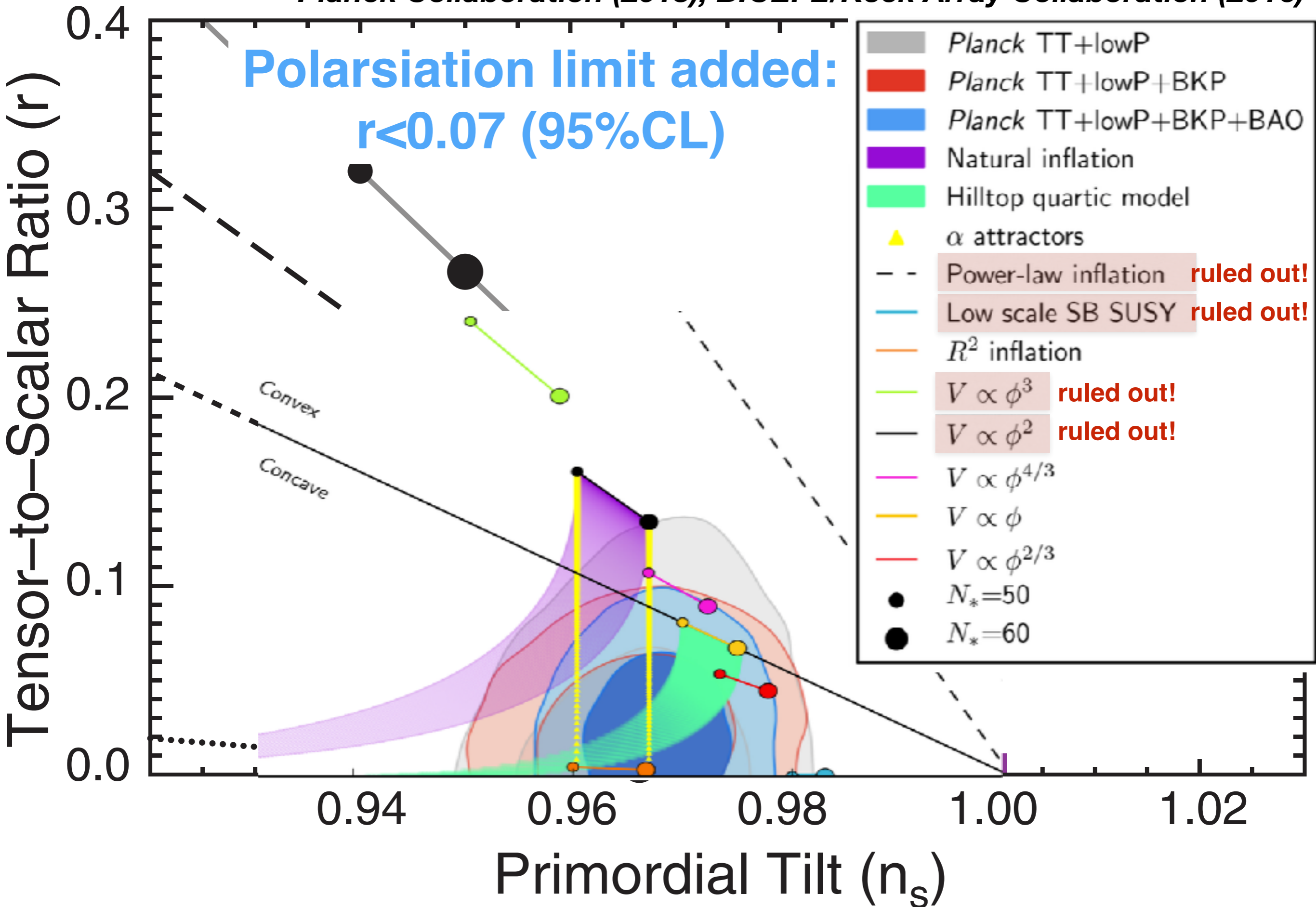
Tensor-to-scalar Ratio

$$r \equiv \frac{\langle h_{ij} h^{ij} \rangle}{\langle \zeta^2 \rangle}$$

- We really want to find this! The current upper bound is **$r < 0.07$** (95%CL)

BICEP2/Keck Array Collaboration (2016)





But, wait a minute...

Are GWs from vacuum fluctuation in spacetime, or from sources?

$$\square h_{ij} = -16\pi G \pi_{ij}$$

- **Homogeneous solution:** “GWs from vacuum fluctuation”
- **Inhomogeneous solution:** “GWs from sources”
 - Scalar and vector fields cannot source tensor fluctuations at linear order (possible at non-linear level)
 - SU(2) gauge field can!

Maleknejad & Sheikh-Jabbari (2013); Dimastrogiovanni & Peloso (2013);
Adshead, Martinec & Wyman (2013); Obata & Soda (2016); ...

Important Message

$$\square h_{ij} = -16\pi G \pi_{ij}$$

- Do not take it for granted if someone told you that detection of the B-mode polarisation would be a signature of “quantum gravity”!
- Only the homogeneous solution corresponds to the vacuum tensor metric perturbation. **There is no *a priori* reason to neglect an inhomogeneous solution!**
- Contrary, we have several examples in which detectable B-modes are generated by **sources** [U(1) and SU(2)]

A New Paradigm

- We must **not** assume that detection of gravitational waves (GWs) from inflation immediately implies that GWs are from the vacuum fluctuation in tensor metric perturbation
- The homogeneous solution is related to $H(t)$ (or the inflaton field excursion; “Lyth bound”) during inflation, but the inhomogeneous solution is **not**.
- **Detection of B-mode polarisation \neq Vacuum fluctuation in metric**



One does not simply read off H from r

From Matteo Fasiello

Experimental Strategy

Commonly Assumed So Far

1. Detect B-mode polarisation in multiple frequencies, to make sure that it is the B-mode of the CMB
2. Check for scale invariance: Consistent with a scale invariant spectrum?
 - Yes => Announce discovery of the vacuum fluctuation in spacetime
 - No => WTF?

New Experimental Strategy: New Standard!

1. Detect B-mode polarisation in multiple frequencies, to make sure that it is the B-mode of the CMB
 2. Consistent with a scale invariant spectrum?
 3. Parity violating correlations (TB and EB) consistent with zero?
 4. Consistent with Gaussianity?
- If, and **ONLY IF** Yes to **all** => Announce discovery of the vacuum fluctuation in spacetime

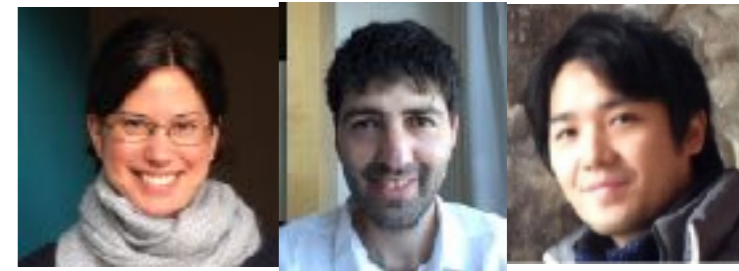
If not, you may have just discovered new physics during inflation!

2. Consistent with a scale invariant spectrum?
 3. Parity violating correlations (TB and EB) consistent with zero?
 4. Consistent with Gaussianity?
- If, and **ONLY IF** Yes to **all** => Announce discovery of the vacuum fluctuation in spacetime

Further Remarks

- “*Guys, you are complicating things too much!*”
- **No.** These sources (eg., gauge fields) should be ubiquitous in a high-energy universe. They have every right to produce GWs if they are around
- Sourced GWs with $r \gg 0.001$ can be phenomenologically more attractive than the vacuum GW from the large-field inflation [requiring super-Planckian field excursion]. Better radiative stability, etc
- Rich[er] phenomenology: Better integration with the Standard Model; reheating; baryon synthesis via leptogenesis, etc. **Testable using many more probes!**

GW from Axion-SU(2) Dynamics

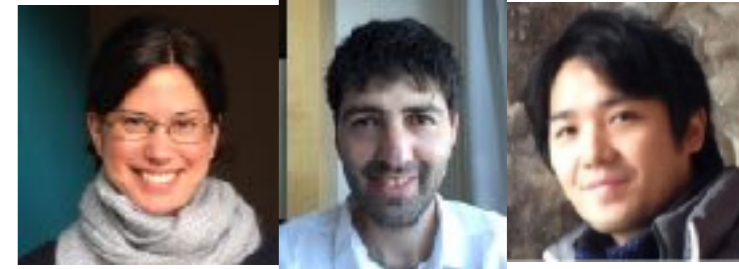


$$\mathcal{L} = \mathcal{L}_{GR} + \mathcal{L}_\phi + \mathcal{L}_\chi - \frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + \frac{\lambda \chi}{4f} F_{\mu\nu}^a \tilde{F}^{a\mu\nu}$$

- ϕ : inflaton field => Just provides quasi-de Sitter background
- χ : pseudo-scalar “axion” field. Spectator field (i.e., negligible energy density compared to the inflaton)
- Field strength of an SU(2) field A_ν^a :

$$F_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - g\epsilon^{abc} A_\mu^b A_\nu^c$$

Background and Perturbation



- In an inflating background, the SU(2) field has a background solution:

$$A_i^a = [\text{scale factor}] \times Q \times \delta_i^a$$

$$Q \equiv (-f \partial_\chi U / 3g\lambda H)^{1/3}$$

U: axion potential

- Perturbations contain a tensor mode (as well as S&V)

$$\delta A_i^a = t_{ai} + \dots$$

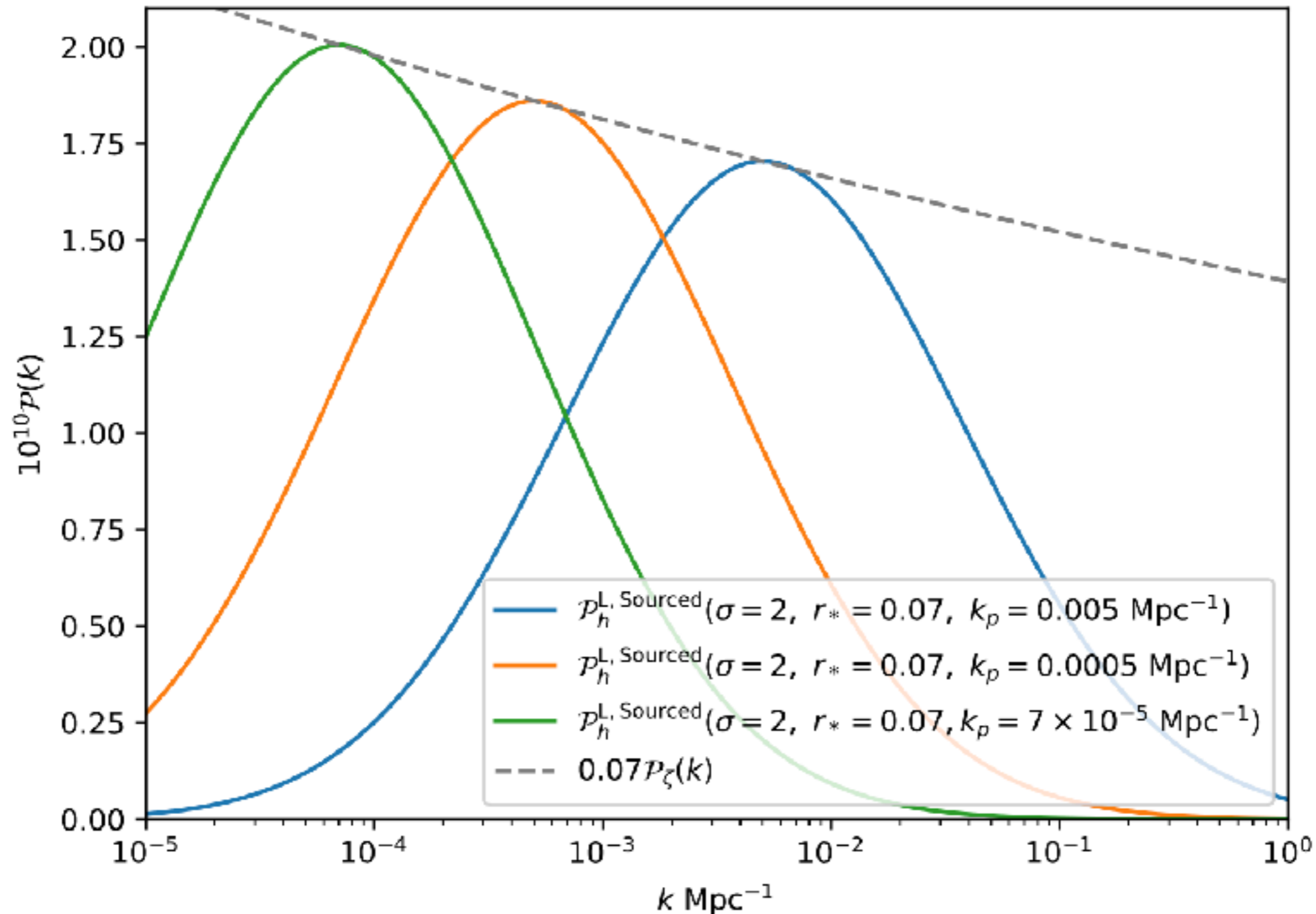
$$t_{ii} = \partial_a t_{ai} = \partial_i t_{ai} = 0$$

Scenario

- The SU(2) field contains tensor, vector, and scalar components
- The tensor components are amplified strongly by a coupling to the axion field
 - But, only one helicity is amplified \Rightarrow GW is **chiral**
(well-known result)
- Brand-new result: **GWs sourced by this mechanism are strongly non-Gaussian!**

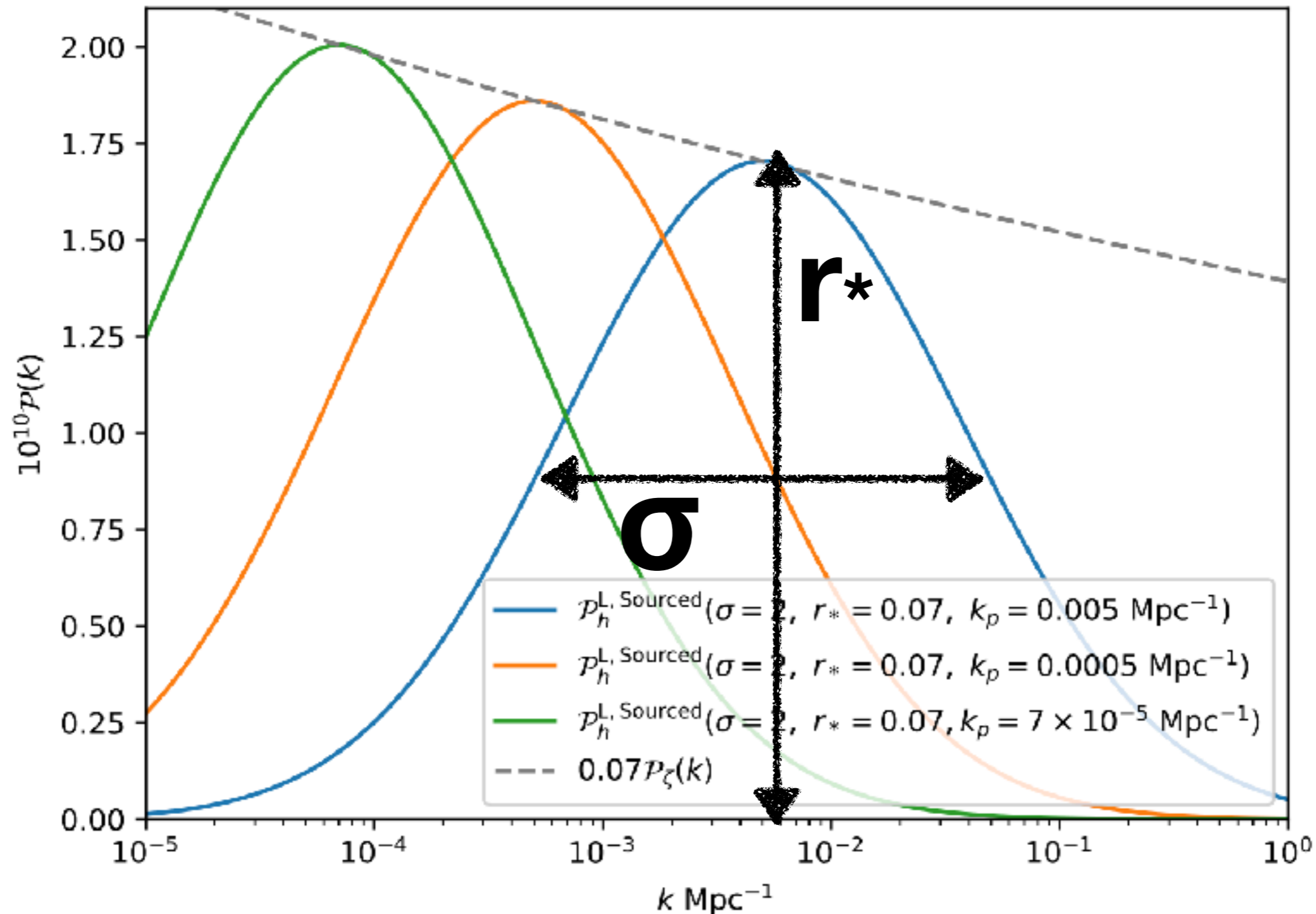
Agrawal, Fujita & EK (2017)

Example Tensor Spectra



- Sourced tensor spectrum can be close to scale invariant, but can also be bumpy

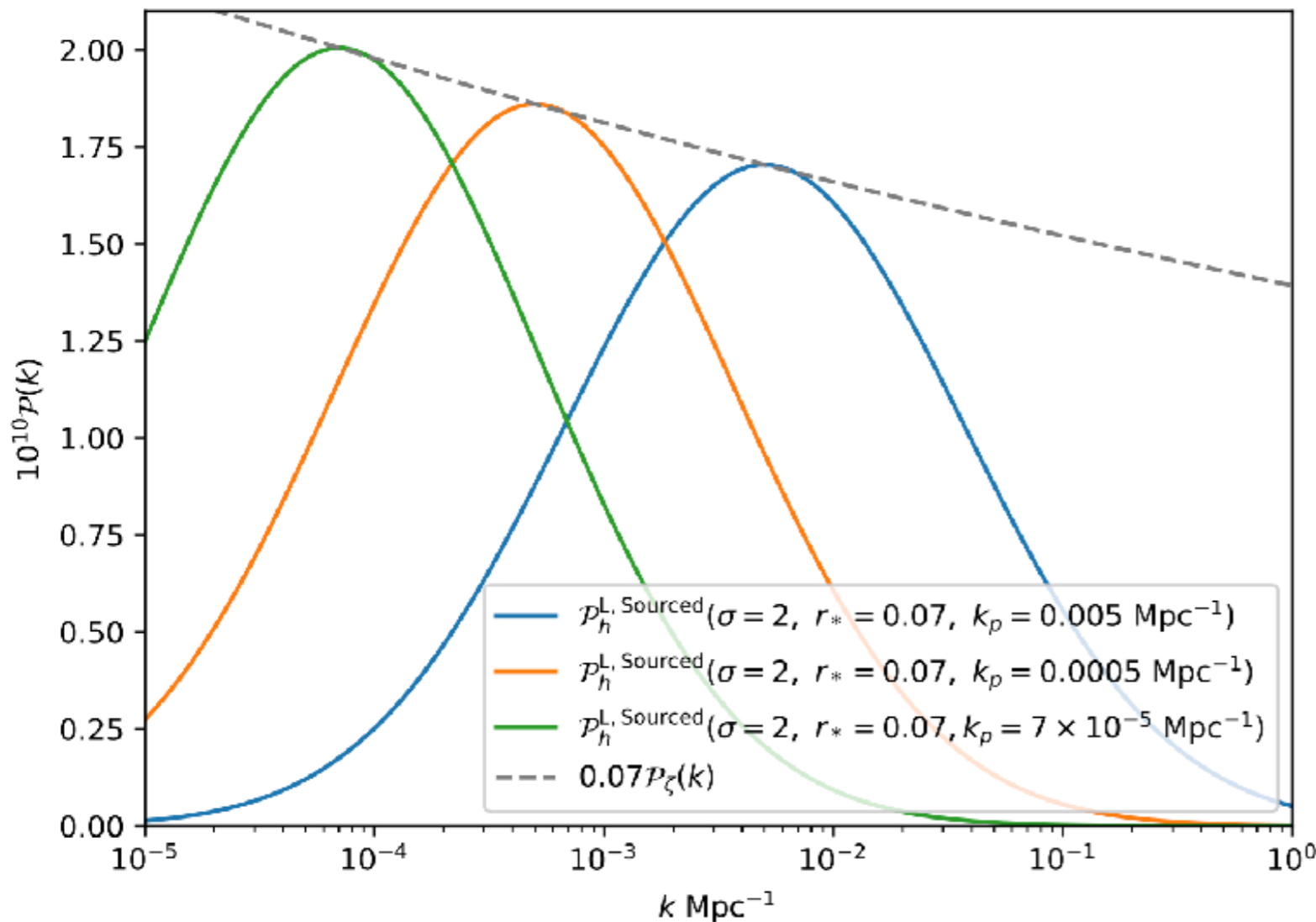
Example Tensor Spectra



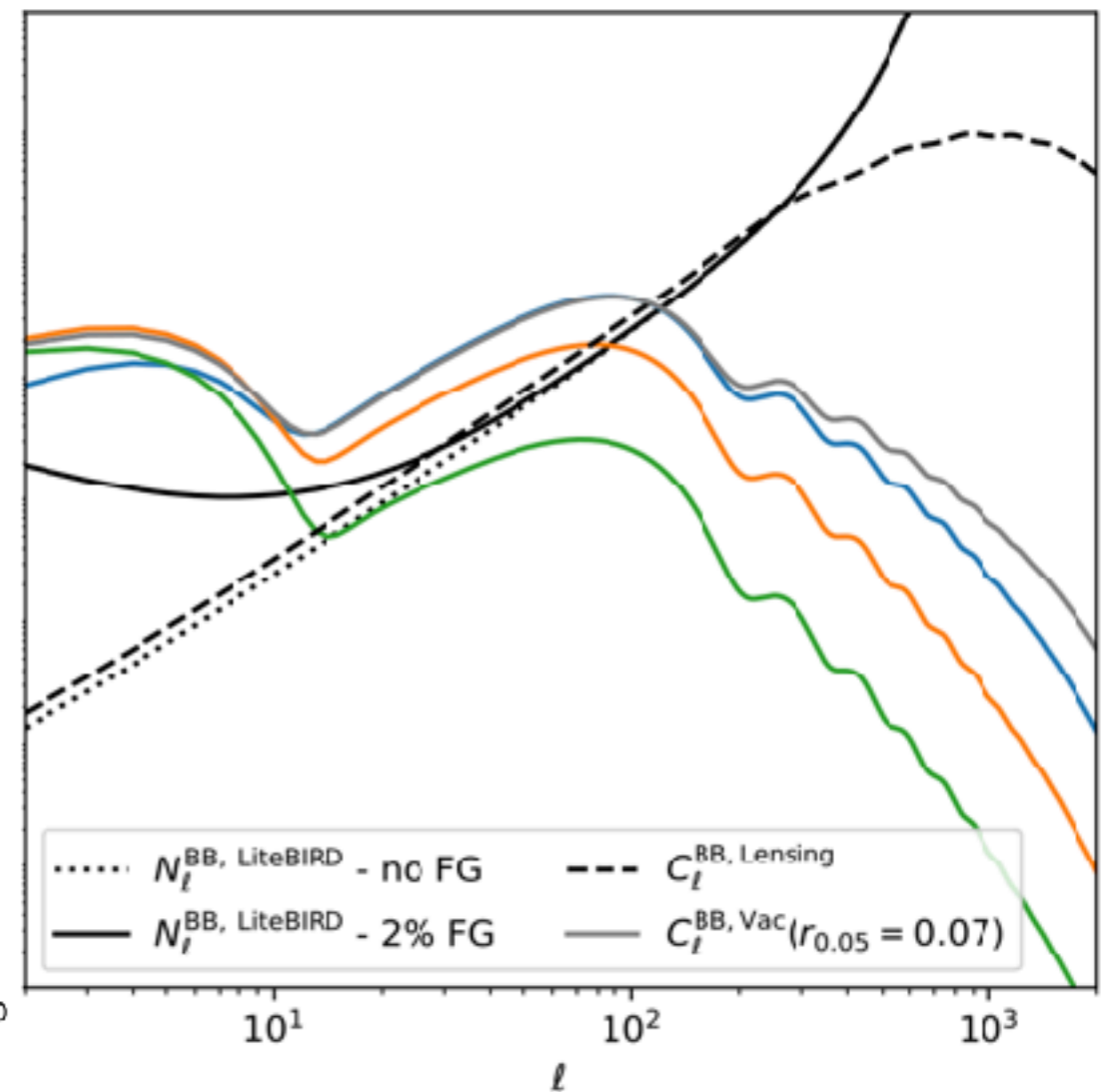
- Sourced tensor spectrum can be close to scale invariant, but can also be bumpy

Example Tensor Spectra

Tensor Power Spectrum, $P(k)$

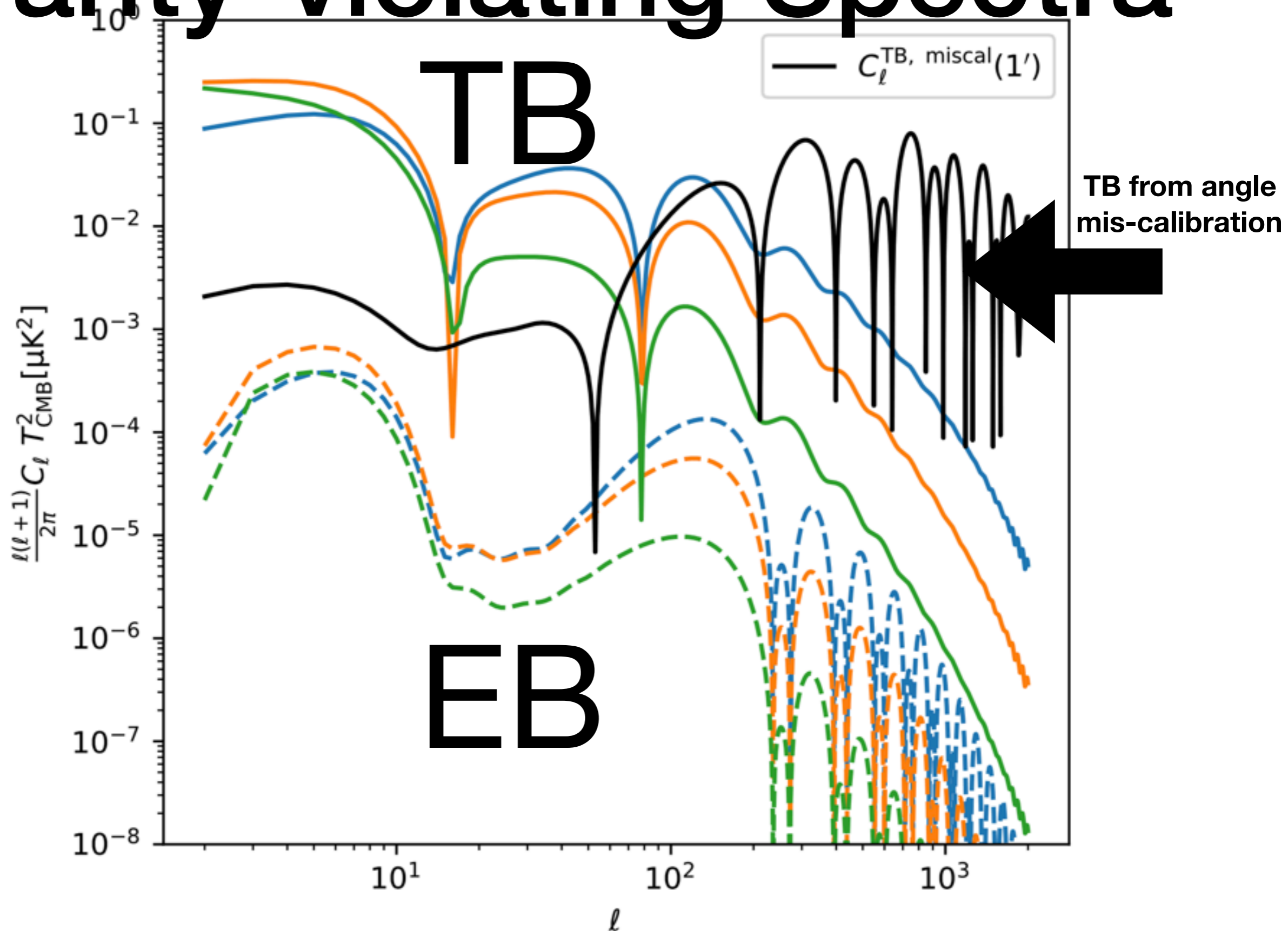


B-mode CMB spectrum, C_l^{BB}



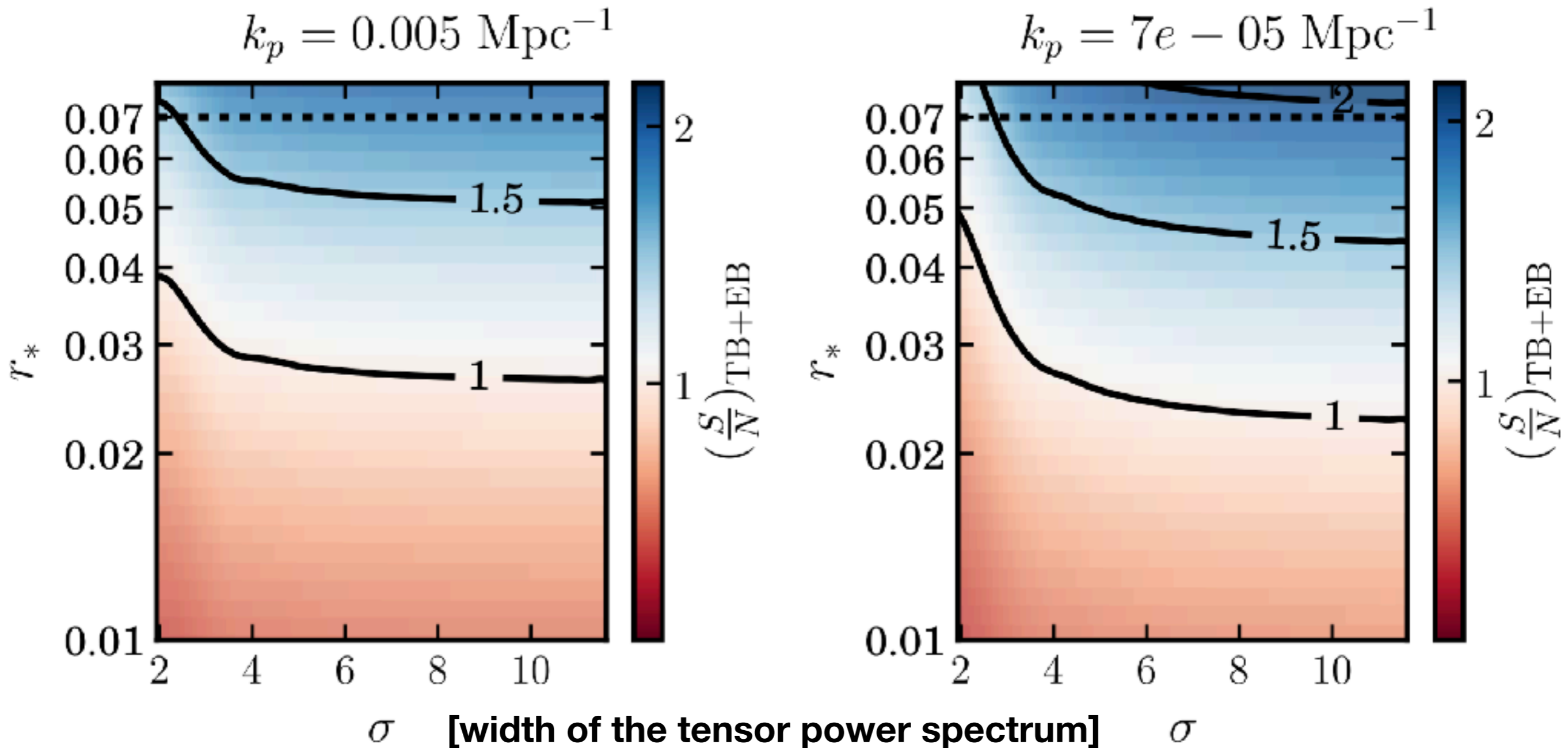
- Sourced tensor spectrum can be close to scale invariant, but can also be bumpy

Parity-violating Spectra



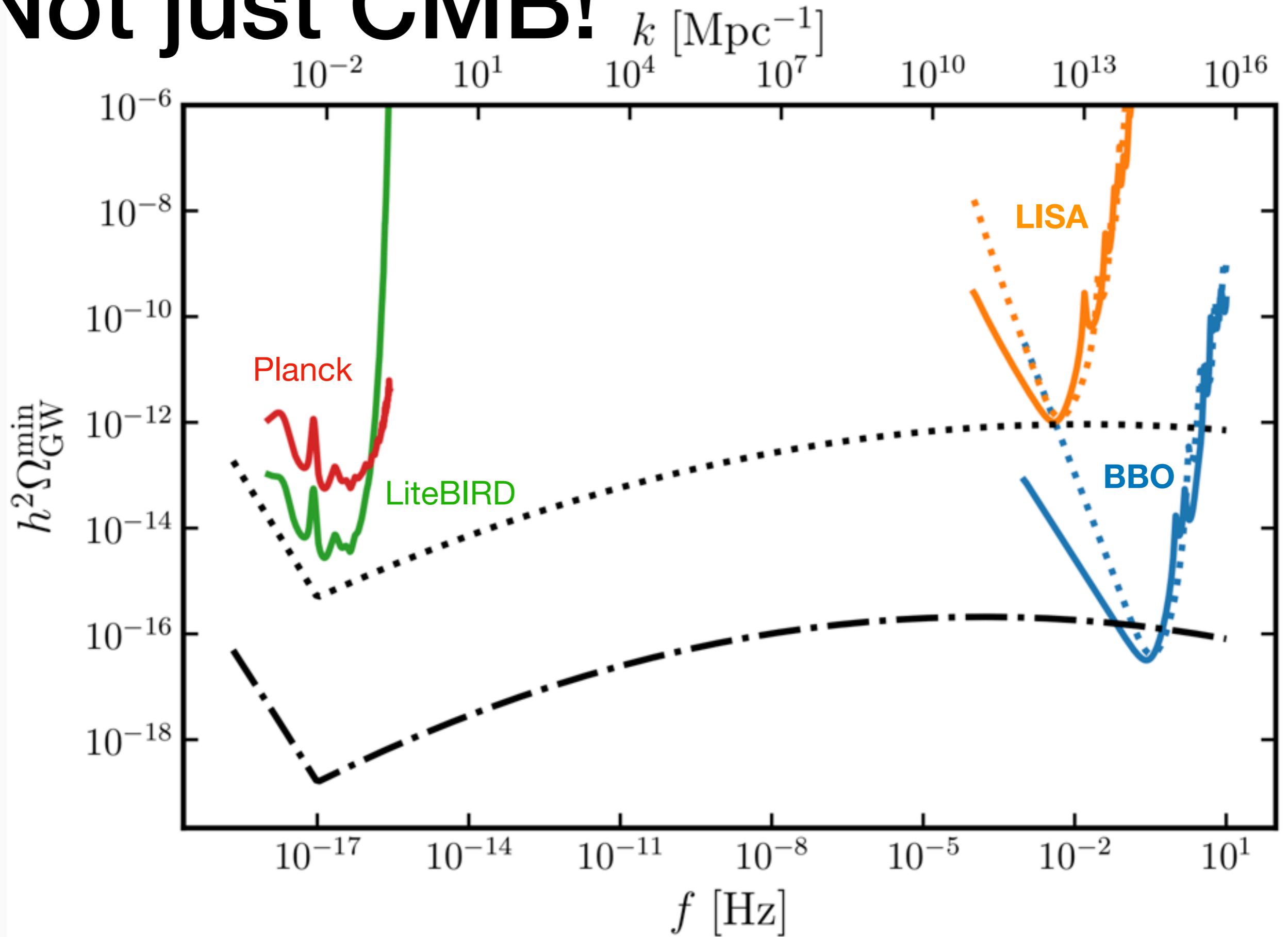
- Angle mis-calibration can be distinguished easily!

Signal-to-noise [LiteBIRD]



- S/N ~ a couple for the peak r_* of 0.07. It's something!

Not just CMB!



Large bispectrum in GW from SU(2) fields



Aniket Agrawal
(MPA)

$$\frac{B_h^{RRR}(k, k, k)}{P_h^2(k)} \approx \frac{25}{\Omega_A}$$



Tomo Fujita
(Kyoto)

$$\langle \hat{h}_R(\mathbf{k}_1) \hat{h}_R(\mathbf{k}_2) \hat{h}_R(\mathbf{k}_3) \rangle = (2\pi)^3 \delta \left(\sum_{i=1}^3 \mathbf{k}_i \right) B_h^{RRR}(k_1, k_2, k_3)$$

- $\Omega_A \ll 1$ is the energy density fraction of the gauge field
- B_h/P_h^2 is of order unity for the vacuum contribution
[Maldacena (2003); Maldacena & Pimentel (2011)]
- *Gaussianity offers a powerful test of whether the detected GW comes from the vacuum or sources*

NG generated at the tree level

$$L_3^{(i)} = c^{(i)} \left[\epsilon^{abc} t_{ai} t_{bj} \left(\partial_i t_{cj} - \frac{m_Q^2 + 1}{3m_Q \tau} \epsilon^{ijk} t_{ck} \right) \right.$$

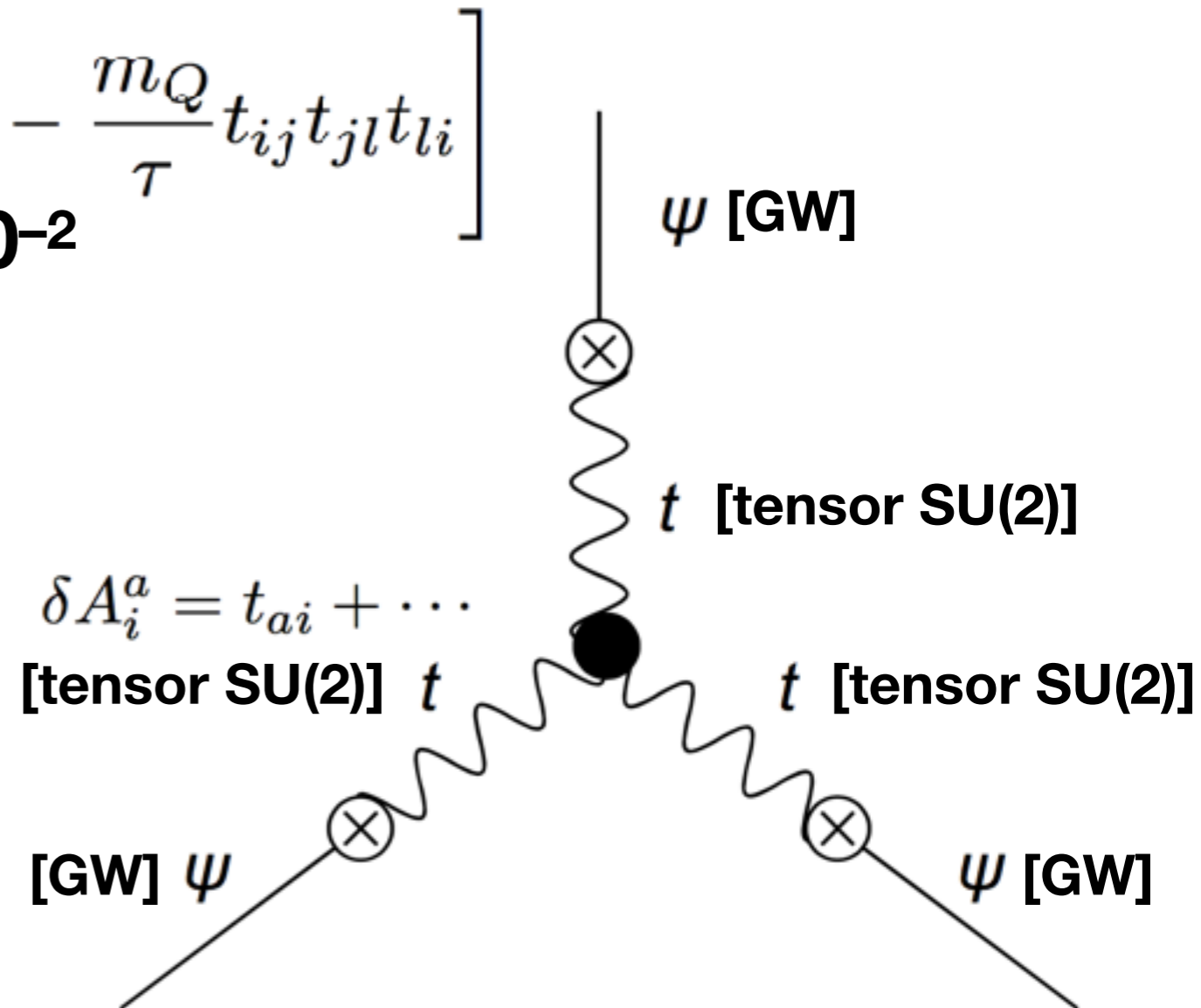
$$\left. - \frac{m_Q}{\tau} t_{ij} t_{jl} t_{li} \right]$$

$$c^{(i)} = g = m_Q^2 H / \sqrt{\epsilon_B} M_{\text{Pl}} \sim 10^{-2}$$

$$\epsilon_B \equiv \frac{g^2 Q^4}{H^2 M_{\text{Pl}}^2} \simeq \frac{2\Omega_A}{1 + m_Q^{-2}} \ll 1$$

$$m_Q \equiv gQ/H \quad [m_Q \sim \text{a few}]$$

- This diagram generates second-order equation of motion for GW



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$$\left. - \frac{m_Q}{\tau} t_{ij} t_{jl} t_{li} \right] \quad \psi \text{ [GW]}$$

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$$\delta A_i^a = t_{ai} + \dots$$

[tensor SU(2)]

t

[tensor SU(2)]

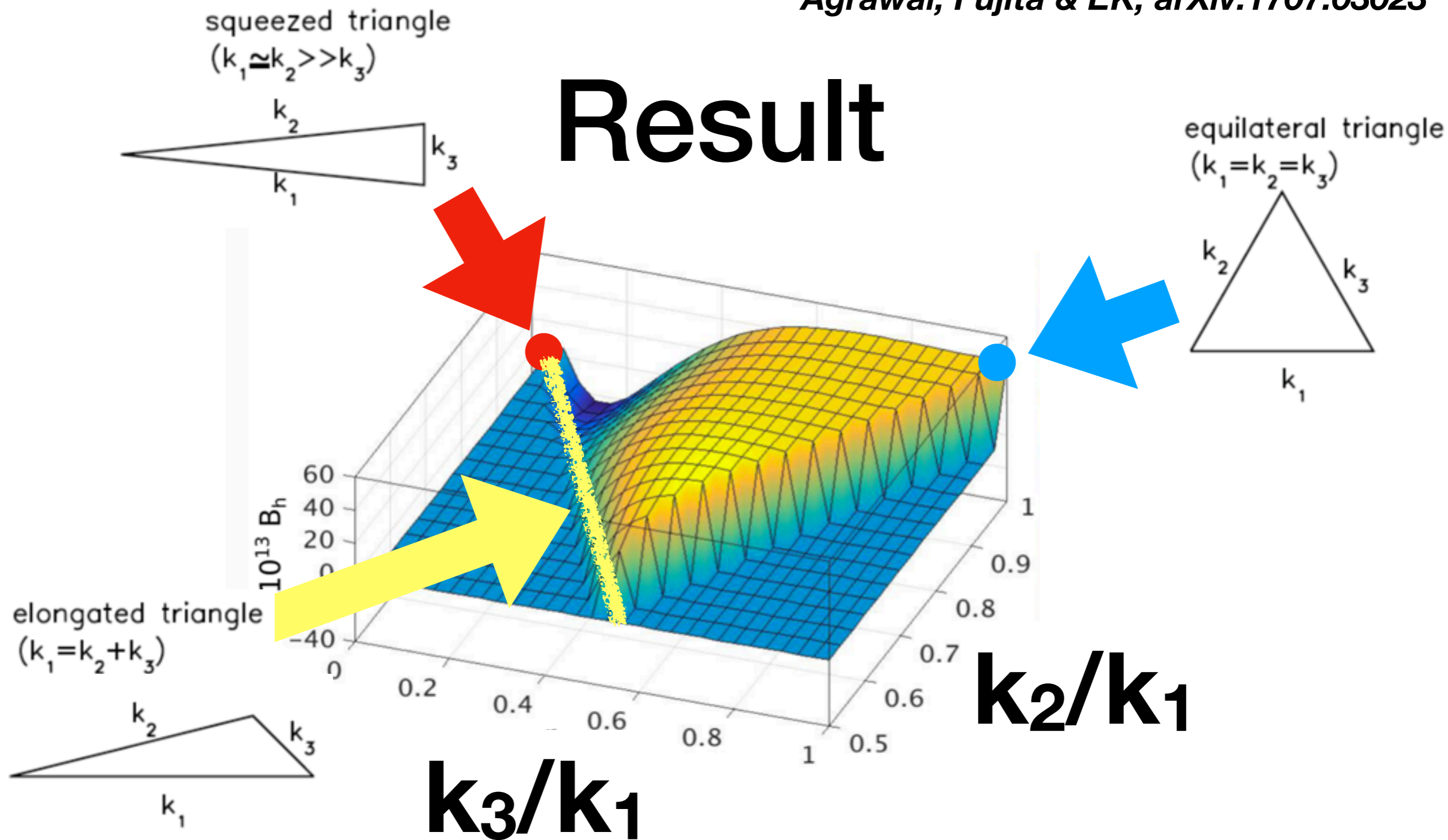
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BISPECTRUM

$$\langle \hat{\psi}_1(\tau, \mathbf{k}_1) \hat{\psi}_1(\tau, \mathbf{k}_2) \hat{\psi}_2(\tau, \mathbf{k}_3) \rangle$$

+perm.

Result



- This shape is similar to, but not exactly the same as, what was used by the Planck team to look for tensor bispectrum

Current Limit on Tensor NG

- The Planck team reported a limit on the tensor bispectrum in the following form:

$$f_{\text{NL}}^{\text{tens}} \equiv \frac{B_h^{+++}(k, k, k)}{F_{\text{scalar}}^{\text{equil.}}(k, k, k)}$$

- The denominator is the **scalar** equilateral bispectrum template, giving $F_{\text{scalar}}^{\text{equil.}}(k, k, k) = (18/5)P_{\text{scalar}}^2(k)$
- The current 68%CL constraint is $f_{\text{NL}}^{\text{tens}} = 400 \pm 1500$

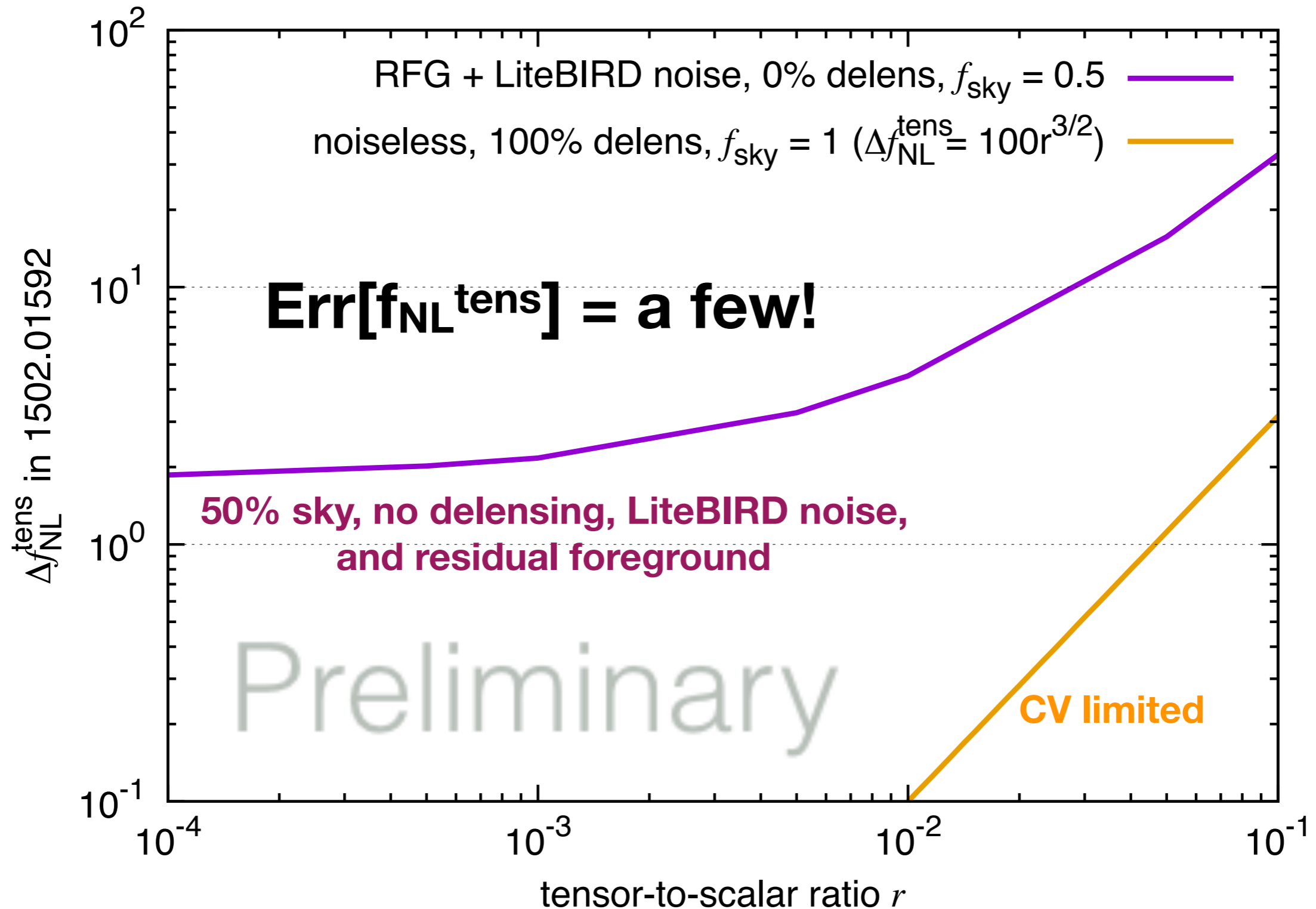
SU(2), confronted

- The SU(2) model of Dimastrogiovanni et al. predicts:

$$f_{\text{NL}}^{\text{tens}} \approx \frac{125}{18\sqrt{2}} \frac{r^2}{\epsilon_B} \approx 2.5 \frac{r^2}{\Omega_A}$$

- The current 68%CL constraint is $f_{\text{NL}}^{\text{tens}} = 400 \pm 1500$
 - This is already constraining!

LiteBIRD would nail it!



JAXA

See Masashi Hazumi's talk on Thursday

+ possible participations from USA, Canada, Europe

LiteBIRD

2025– [proposed]



Target: $\delta r < 0.001$

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LiteBIRD

2025– [proposed]



**Polarisation satellite dedicated to
measure CMB polarisation from
primordial GW, with a few thousand
super-conducting detectors in space**

JAXA

See Masashi Hazumi's
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+ possible participations
from USA, Canada,
Europe

LiteBIRD

2025– [proposed]



Down-selected by JAXA as
one of the two missions
competing for a launch in mid 2020's

Summary

- Single-field inflation looks good: all the CMB data support it
- **Next frontier**: Using CMB polarisation to find GWs from inflation. **Definitive evidence for inflation!**
- With LiteBIRD we plan to reach $r \sim 10^{-3}$, i.e., 100 times better than the current bound
- GW from vacuum or sources? An exciting window to new physics