

# Symmetry protected topological phase in magnetization plateaux

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Collaboration with

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# Introduction

- Gapped systems

Long range entangled state • • • anyonic excitation  
(Intrinsic topological order)

Short Long range entangled state

Local unitary  
transformation



Without any symmetry

Trivial (direct product) state

- Short Long range entangled state can be **nontrivial** if some symmetry is imposed.

- Local order parameter → Ginzburg-Landau theory

- No Local order parameter → **Symmetry protected topological (SPT) phase**

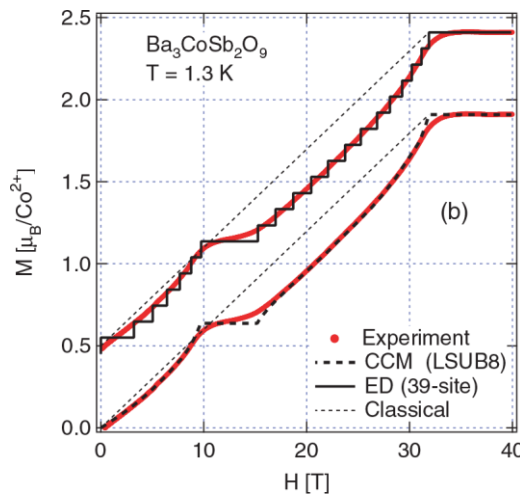
# Results

Magnetization plateau states in 1D antiferromagnets are in an SPT phase protected by **link-inversion** symmetry if  $S-m = \text{odd}$  integer.

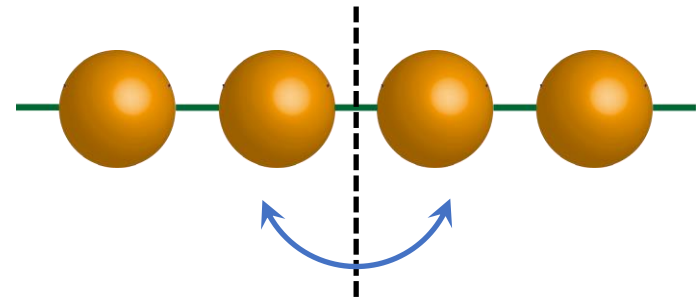
Size of spin      Magnetization per site

## Magnetization plateau

- Region where  $M$  is unchanged with increasing  $H$  in magnetization curves



## Link-inversion



Shirata et al., 2012

# Methods

## 1. Field theories

- Nonlinear sigma model,  
dual boson-vortex theory (sine-Gordon model)

## 2. Numerical calculations

- Infinite time-evolving block decimation (iTEBD),  
entanglement spectra

## 3. Matrix product state (MPS) representation

# Field theories

## Oshikawa-Yamanaka-Affleck condition

Oshikawa-Yamanaka-Affleck, 1997

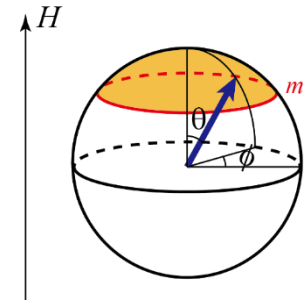
$$r(S - m) \in \mathbb{Z}$$

Number of sites in a unit cell      Size of spin      Magnetization per site

## Canted spin configuration

Tanaka-Totsuka-Hu, 2009

$$\mathbf{S}_j(\tau) = S \begin{pmatrix} (-1)^j \cos \phi_j(\tau) \sin \theta_0 \\ (-1)^j \sin \phi_j(\tau) \sin \theta_0 \\ \cos \theta_0 \end{pmatrix}$$



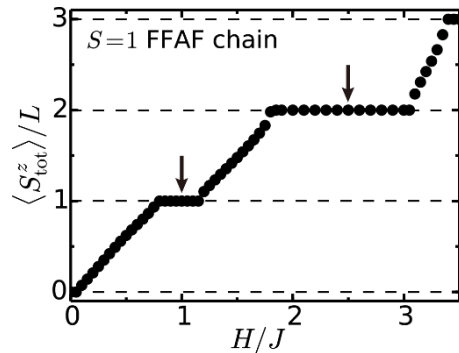
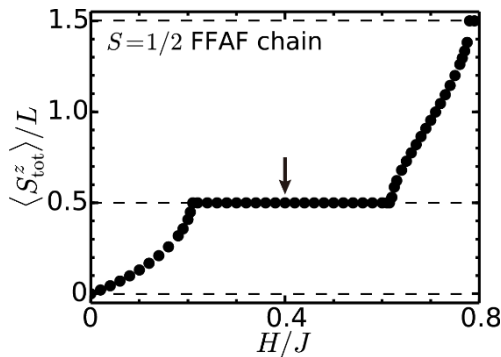
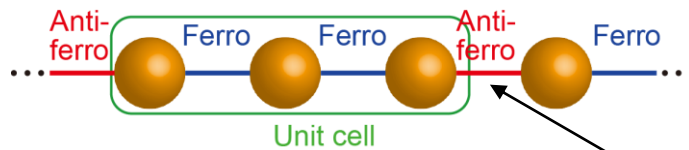
$$\mathcal{S} = \mathcal{S}_{\text{kin}} + \mathcal{S}_{\text{BP}}^{\text{tot}} \quad \mathcal{S}_{\text{kin}} \rightarrow \int dx d\tau \frac{\zeta}{2} \left\{ \frac{1}{v^2} (\partial_\tau \phi)^2 + (\partial_x \phi)^2 \right\}$$

O(3) nonlinear sigma model

Also dual vortex theory approach  
 $\rightarrow$  sine-Gordon theory

# Numerical calculations

- Ferro-ferro-antiferro model
- Infinite-time evolving block decimation (iTEBD).
- Magnetization curves and entanglement spectra (ES).

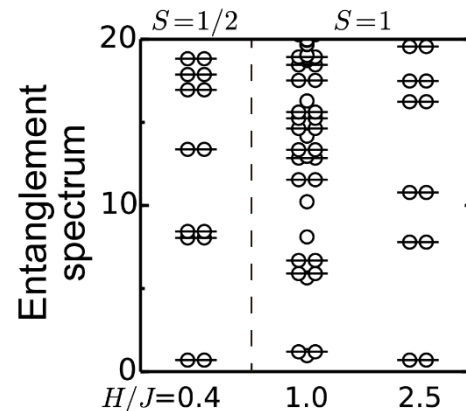


Schmidt decomposition

Bipartition at the AF bond

$$|\Psi\rangle = \sum_{\alpha} \lambda_{\alpha} |\Psi_A\rangle_{\alpha} \otimes |\Psi_B\rangle_{\alpha}$$

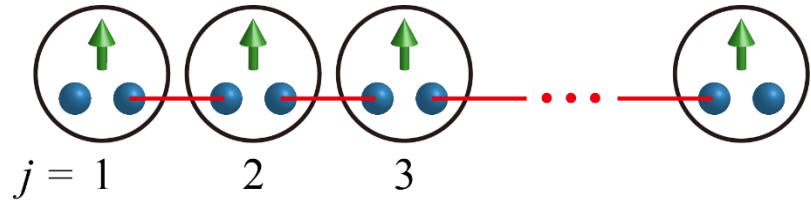
$$\text{ES } \{-\ln(\lambda_{\alpha}^2)\} \quad (\alpha = 1, \dots, \chi)$$



$$S - m = \text{odd} \quad \text{even} \quad \text{odd}$$

# MPS representation

VBS picture for  $m=1/2$  plateau in  $S=3/2$



## Schwinger boson representation

$$|\Psi\rangle = \prod_j P_j a_j^\dagger (a_j^\dagger b_{j+1}^\dagger - a_j^\dagger b_{j+1}^\dagger) \otimes_j |0\rangle_j$$

↑
↑  
 up                  down

## Canonical representation

$$|\Psi\rangle = \sum_{S_j^z = -3/2}^{3/2} \dots \Lambda \Gamma[S_{j-1}^z] \Lambda \Gamma[S_j^z] \Lambda \Gamma[S_{j+1}^z] \Lambda \dots \otimes_j |S_j^z\rangle$$

Link- Inversion  $\mathcal{I}$  acts on MPS as  $\Gamma^T = e^{i\theta_{\mathcal{I}}} U_{\mathcal{I}}^\dagger \Gamma U_{\mathcal{I}}$

We can find that  $U_{\mathcal{I}} = i\sigma_y = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$

$U_{\mathcal{I}}^2 = -1 : \text{Nontrivial}$

# Conclusion

Magnetization plateau states in 1D antiferromagnets are in an SPT phase protected by **link-inversion** symmetry if  $S-m = \text{odd}$  integer.

1. Field theories
2. Numerical calculations
3. Matrix product state (MPS) representation

- ✓ Please visit PS-D3.
- ✓ Discussions in free time are also welcome.