

# ***Avalanche of Entanglement***

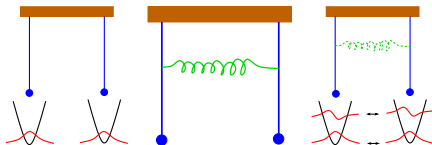
Linear fields: Gaussian (squeezed) states  $\rightarrow$  pairs of particles

$\rightarrow$  bi-partite entanglement M. Hotta, R.S., W.G. Unruh, Phys. Rev. D **91**, 124060 (2015)

E.g., Hawking radiation S. W. Hawking, Nature **248**, 30 (1974); Comm. Math. Phys. **43**, 199 (1975)

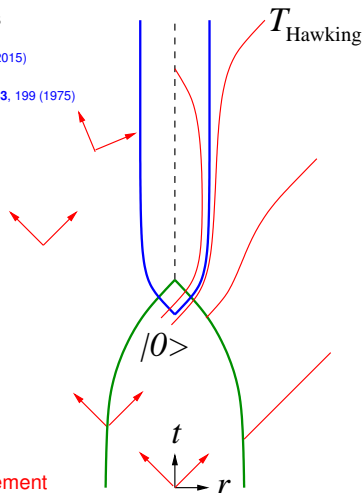
$$T_{\text{Hawking}} = \frac{1}{8\pi M} \frac{\hbar c^3}{G_N k_B}$$

Tearing apart of quantum vacuum fluctuations



But: trans-Planckian problem, information puzzle etc.

Non-linear interactions ("scrambling")  $\rightarrow$  multi-partite entanglement



Invariant under local unitary operations & non-increasing for local decoherence/dissipation etc.

- bi-partite entanglement → Bell states

$$|\text{Bell}\rangle = \frac{|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle}{\sqrt{2}},$$

- concurrence for arbitrary mixed states  $\hat{\rho}$  of two spins

$$C_2[\hat{\rho}] = f\left(\text{Eigenvalues}\left[\sqrt{\sqrt{\hat{\rho}}\hat{R}\hat{\rho}^*\hat{R}\sqrt{\hat{\rho}}}\right]\right)$$

convex roof construction: minimization over all decompositions into pure states!

- tri-partite entanglement → GHZ states

$$|\text{GHZ}\rangle_3 = \frac{|\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\rangle}{\sqrt{2}}$$

- three-tangle  $\tau_3$  of three spins in pure state

W-states  $|\text{W}\rangle_3 = (|\uparrow\downarrow\downarrow\rangle + |\downarrow\uparrow\downarrow\rangle + |\downarrow\downarrow\uparrow\rangle)/\sqrt{3}$  ???

- quadri-partite entanglement → GHZ<sub>4</sub> states

- four-tangle(s)  $\tau_4$  of four spins in pure state

- entanglement entropy ↔ entanglement between two sub-systems of pure state  $|\psi_{AB}\rangle$

$$S = -\text{Tr}\{\hat{\rho}_A \ln \hat{\rho}_A\} \quad \text{with} \quad \hat{\rho}_A = \text{Tr}_B\{|\psi_{AB}\rangle\langle\psi_{AB}|\}$$

## Hamiltonian

$$\hat{H} = -J \sum_{i=1}^N \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z - \sum_{i=1}^N \hat{\sigma}_i^x$$

$J = 0$ : paramagnetic state  $|\rightarrow\rightarrow\rightarrow\dots\rangle$

→ no entanglement

$J \rightarrow \infty$ : ferromagnetic state

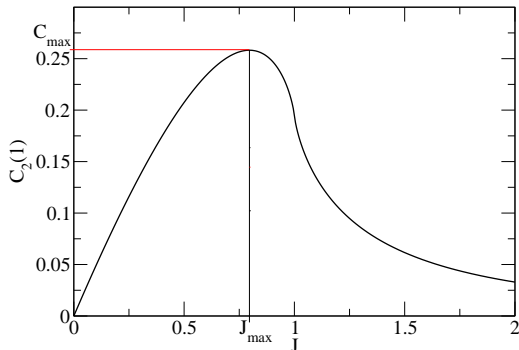
$$\frac{|\uparrow\uparrow\uparrow\dots\rangle + |\downarrow\downarrow\downarrow\dots\rangle}{\sqrt{2}}$$

→  $N$ -partite entanglement (GHZ type)

$J = 1$ : critical point (phase transition)

→ entanglement entropy  $S \sim \ln N$  (violation of “area” law...)

cannot be explained by bi-partite entanglement  $C_2$  alone → multi-partite entanglement!



Reduced density matrices  $\hat{\rho}_i = \text{Tr}_{\text{lattice} \setminus \{i\}} \{ |\Psi\rangle \langle \Psi| \}$ ,  $\hat{\rho}_{ij} = \text{Tr}_{\text{lattice} \setminus \{ij\}} \{ |\Psi\rangle \langle \Psi| \}$ , etc.

Diagonalization (exact)

$$\hat{\rho}_{ij\dots} = \sum_{\alpha} \rho_{\alpha} \left| \Psi_{ij\dots}^{\alpha} \right\rangle \left\langle \Psi_{ij\dots}^{\alpha} \right|$$

First two eigenvalues are dominant

$$\hat{\rho}_{ij\dots} \approx \sum_{\alpha=1}^2 \rho_{\alpha} \left| \Psi_{ij\dots}^{\alpha} \right\rangle \left\langle \Psi_{ij\dots}^{\alpha} \right|$$

→ rank-two matrices

→ calculation of tangles

first bi-partite  $C_2$

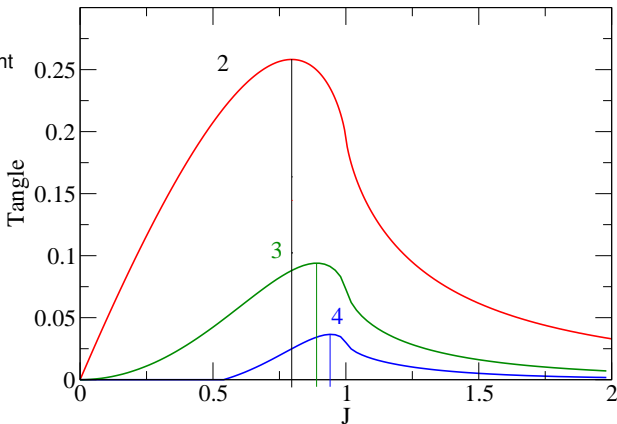
then tri-partite  $\sqrt{\tau_3}$

later quadri-partite  $\tau_4$

...

finally  $N$ -partite

→ avalanche of entanglement



Correlations, e.g.,

$$\langle \hat{\sigma}_i^z \hat{\sigma}_j^z \rangle^{\text{corr}} = \langle \hat{\sigma}_i^z \hat{\sigma}_j^z \rangle - \langle \hat{\sigma}_i^z \rangle \langle \hat{\sigma}_j^z \rangle$$

Correlated density matrices

$$\hat{\rho}_{ij}^{\text{corr}} = \hat{\rho}_{ij} - \hat{\rho}_i \hat{\rho}_j \quad \text{etc.}$$

Upper bound (spectral norm)

$$\langle \hat{\sigma}_i^z \hat{\sigma}_j^z \rangle^{\text{corr}} = \text{Tr}\{\hat{\sigma}_i^z \hat{\sigma}_j^z \hat{\rho}_{ij}^{\text{corr}}\} \leq \|\hat{\rho}_{ij}^{\text{corr}}\|_1$$

Concurrence for two spins

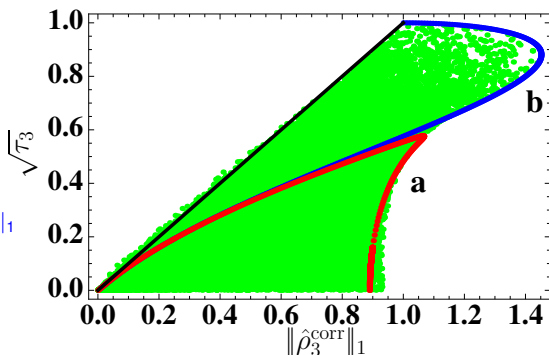
$$C_2[\hat{\rho}_{ij}] \leq \|\hat{\rho}_{ij}^{\text{corr}}\|_1$$

→ entanglement ↔ correlations

Question: three (or more) spins?

→ three-tangle  $\sqrt{\tau_3[\hat{\rho}_{ijk}]}$  as approximate lower bound for maximum correlation  $\|\hat{\rho}_{ijk}^{\text{corr}}\|_1$

Questions: outlier states? four spins? ...



Small  $J$ : hierarchy of correlations

two-point  $\gg$  three-point  $\gg$  four-point

$$\|\hat{\rho}_{ij}^{\text{corr}}\|_1 \gg \|\hat{\rho}_{ijk}^{\text{corr}}\|_1 \gg \|\hat{\rho}_{ijkl}^{\text{corr}}\|_1$$

Violation of hierarchy at  $J \approx 0.7$

→ approximation schemes

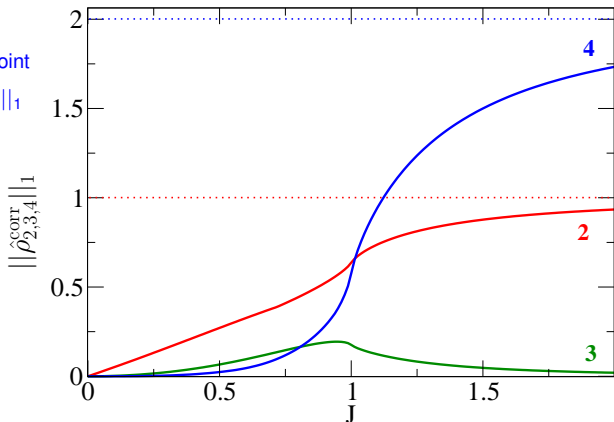
$$\langle \sigma_i^x \sigma_j^x \sigma_k^x \sigma_l^x \rangle =$$

$$\langle \sigma_i^x \rangle \langle \sigma_j^x \rangle \langle \sigma_k^x \rangle \langle \sigma_l^x \rangle +$$

$$\langle \sigma_i^x \sigma_j^x \rangle^{\text{corr}} \langle \sigma_k^x \rangle \langle \sigma_l^x \rangle + \dots +$$

$$\langle \sigma_i^x \sigma_j^x \sigma_k^x \rangle^{\text{corr}} \langle \sigma_l^x \rangle + \dots +$$

$$\langle \sigma_i^x \sigma_j^x \sigma_k^x \sigma_l^x \rangle^{\text{corr}}$$



Hamiltonian

$$\hat{H} = -J \sum_{i=1}^N \left( \hat{b}_i^\dagger \hat{b}_{i+1} + \hat{b}_{i+1}^\dagger \hat{b}_i \right) + \frac{1}{2} \sum_{i=1}^N \hat{b}_i^\dagger \hat{b}_i^\dagger \hat{b}_i \hat{b}_i$$

Small  $J$ : paramagnetic  $\rightarrow$  Mott insulator

Large  $J$ : ferromagnetic  $\rightarrow$  superfluid

Small  $J$ : hierarchy of correlations

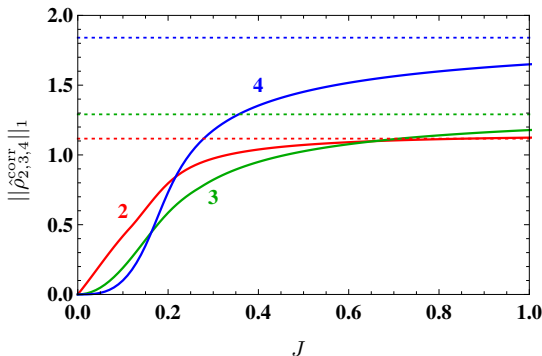
two-point  $\gg$  three-point  $\gg$  four-point

$$\|\hat{\rho}_{ij}^{\text{corr}}\|_1 \gg \|\hat{\rho}_{ijk}^{\text{corr}}\|_1 \gg \|\hat{\rho}_{ijkl}^{\text{corr}}\|_1$$

Violation of hierarchy at  $J \approx 0.16$

well before the critical point

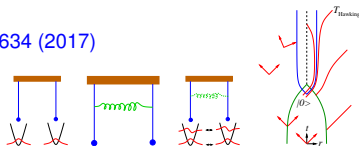
(here  $J_{\text{crit}} \approx 0.3$ )





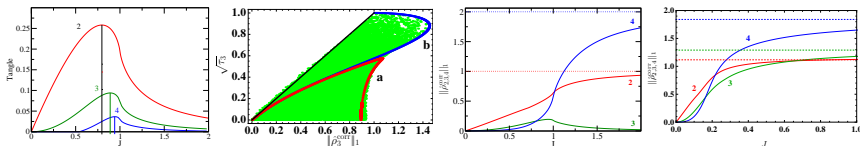
K.V. Krutitsky, A. Osterloh, R.S., Nature Scientific Reports 7, 3634 (2017)

- Motivation: bi-partite entanglement  $\leftrightarrow$  pairs (partners)  
multi-partite entanglement  $\leftrightarrow$  ???



- Ising model: avalanche of entanglement  $\rightarrow$  correlations

Hubbard model



approximation 1)

$$\hat{\rho}_{ij} \approx \rho_1 |\Psi_{ij}^1\rangle \langle \Psi_{ij}^1| + (1 - \rho_1) |\Psi_{ij}^2\rangle \langle \Psi_{ij}^2|$$

approximation 2)

$$\hat{\rho}_{ij} \approx \frac{\rho_1 |\Psi_{ij}^1\rangle \langle \Psi_{ij}^1| + \rho_2 |\Psi_{ij}^2\rangle \langle \Psi_{ij}^2|}{\rho_1 + \rho_2}$$

note that  $\rho_3 < 0.5\%$

