

# Defect Holography and the Fractional Quantum Hall Effect

**Charles M. Melby-Thompson**  
**Kavli IPMU**

Strings and Fields, YITP (26 July, 2014)

work in progress  
with M. Fujita, R. Meyer, S. Sugimoto

# Defect Holography

*and the Fractional Quantum Hall Effect*

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Today's focus:

Holography of domain walls in Yang-Mills–Chern-Simons theory.

Yang-Mills–Chern-Simons? (YMCS)

$$3\text{D: } S = -\frac{1}{4g^2} \int_{M_3} d^3x \text{Tr}(F^{\mu\nu} F_{\mu\nu}) + \frac{k}{4\pi} \int_{M_3} \omega_3(A)$$

level ↙

$$\omega_3(A) = \text{Tr}(A \wedge dA + \frac{2}{3} A \wedge A \wedge A) \leftarrow \text{Chern-Simons term}$$

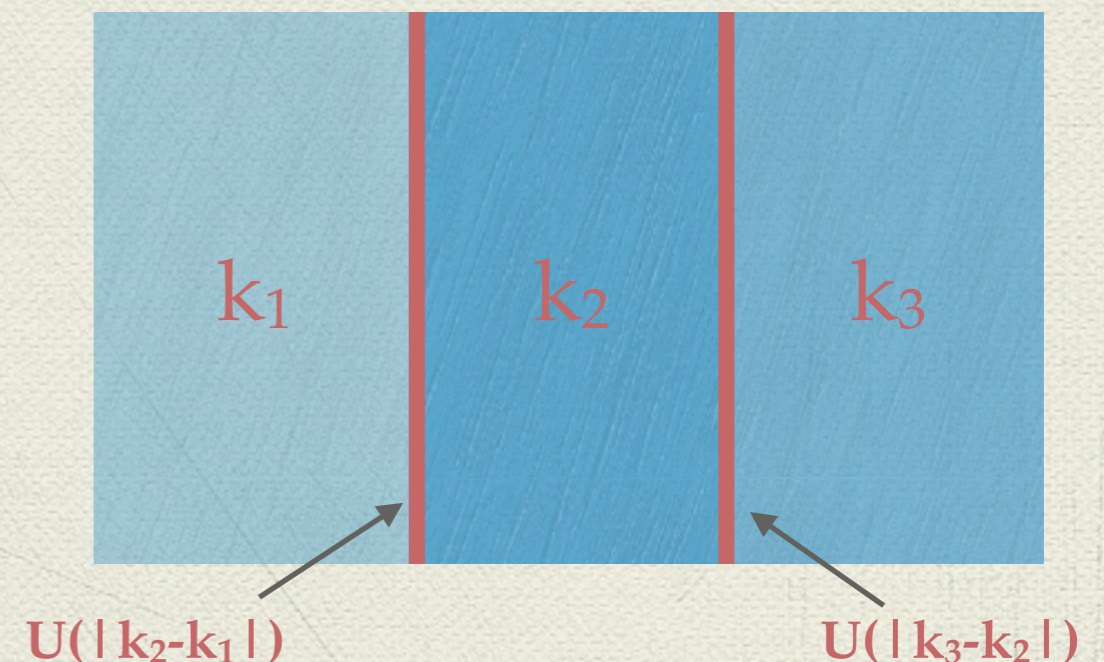
CS term only gauge invariant up to a total derivative:

$$\delta_\Lambda \omega_3(A) = d \text{Tr}(\Lambda F)$$

Consider domain walls separating phases at different levels => **not gauge invariant!**

Must be additional degrees of freedom living on the defects to compensate

New degrees transform under a  $U(|\Delta k|)$  global symmetry.



# Holographic Embedding of YMCS

N=4 SYM  $\Leftrightarrow$  AdS5

Compactify N=4 SYM on a circle, impose anti-periodic boundary conditions on fermions  $\Rightarrow$  ~~SUSY~~

Geometry: AdS soliton

Massive fermions induce scalar masses

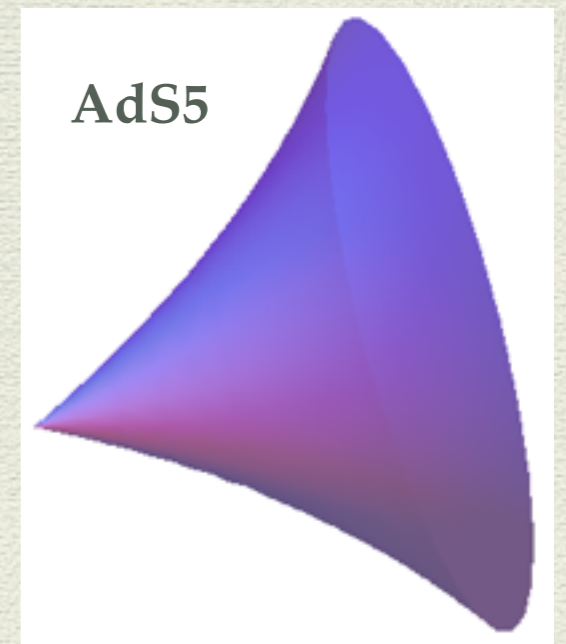
Gauge field confines

}  $\Rightarrow$  gapped theory

Dual geometry caps off at a scale determined by the compactification radius

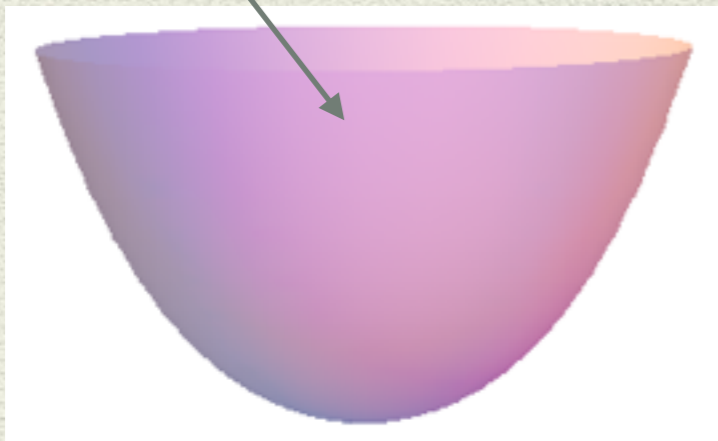
$$ds^2 = w^2 dx_3^2 + f(w)w^2 d\theta^2 + \frac{R^2}{f(w)} \frac{dw^2}{w^2}$$

cap at  $w=w_0$   $f(w) = 1 - \frac{w_0^4}{w^4}$



# Holographic Embedding of YMCS~~CS~~ -ish

string worldsheet



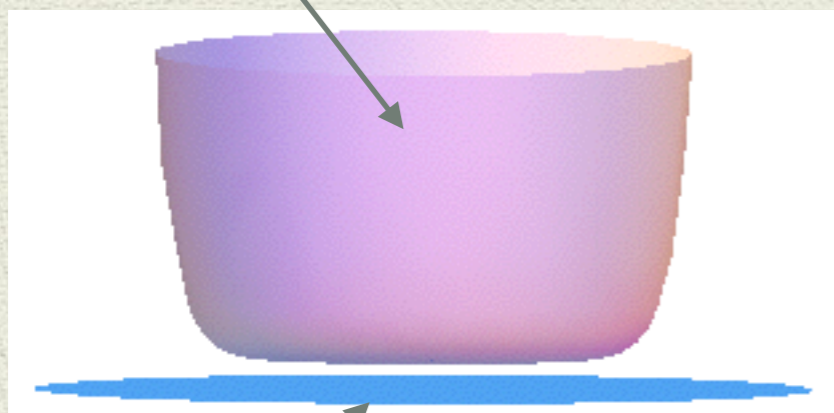
Wilson line

Computed by attaching a string.

AdS => conformal => no dependence on size  
(after renormalization)



string worldsheet



soliton tip

AdS soliton => Wilson lines larger than the  
compactification radius hit the bottom

Dominated by area of worldsheet at bottom  
=> area law \*confinement

Infrared physics = confined phase of 3D YM

# Holographic Embedding of YMCS<sup>-ish</sup>

How to get the CS in YM-CS?

Starting with D3 branes, introduce  $k$  units of  $C_0$  flux on the circle.

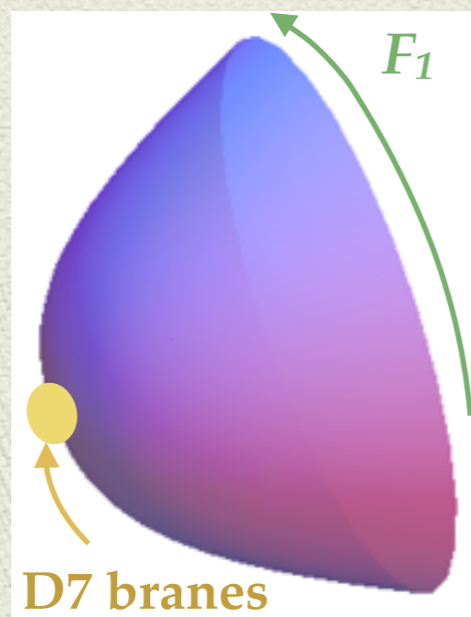
WZ part of action: 
$$\frac{1}{8\pi^2} \int_{M_4} \text{Tr}(F \wedge F) C_0 = \frac{1}{8\pi^2} \int_{M_4} \omega_3(A) \wedge (k d\theta) = \frac{k}{4\pi} \int_{M_3} \omega_3(A)$$
 Chern-Simons

=> With twisted fermion b.c., dual low-energy effective theory is YMCS, in the confined phase.

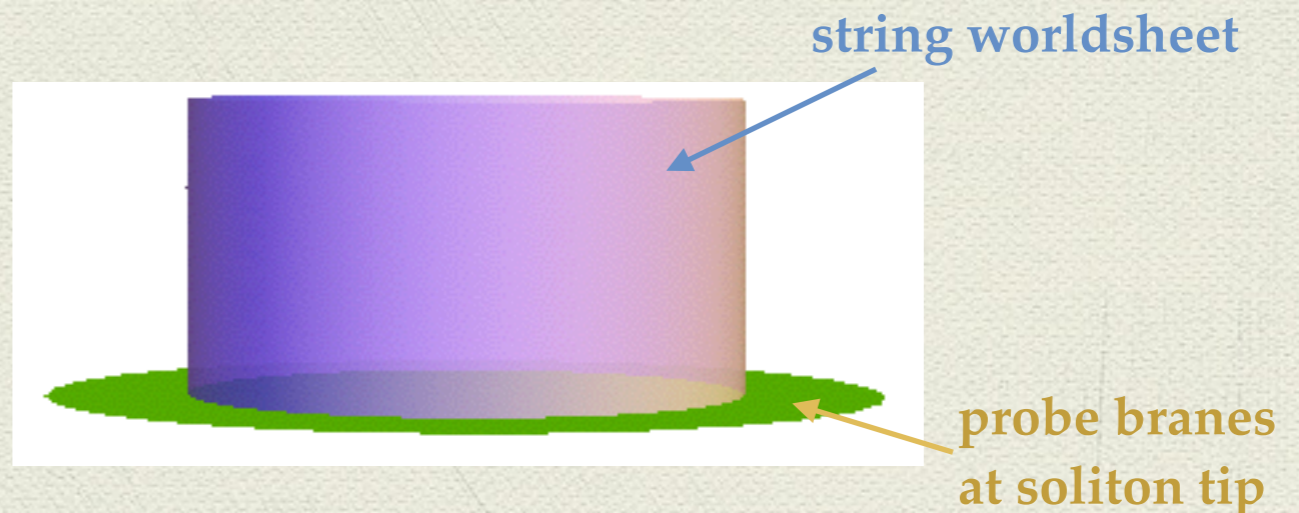
$k$  units of  $C_0$  flux on the circle  
=>  $k$  D7 branes must lie at tip

probe branes wrapping internal  $S^5$

=> 3D surface in AdS soliton



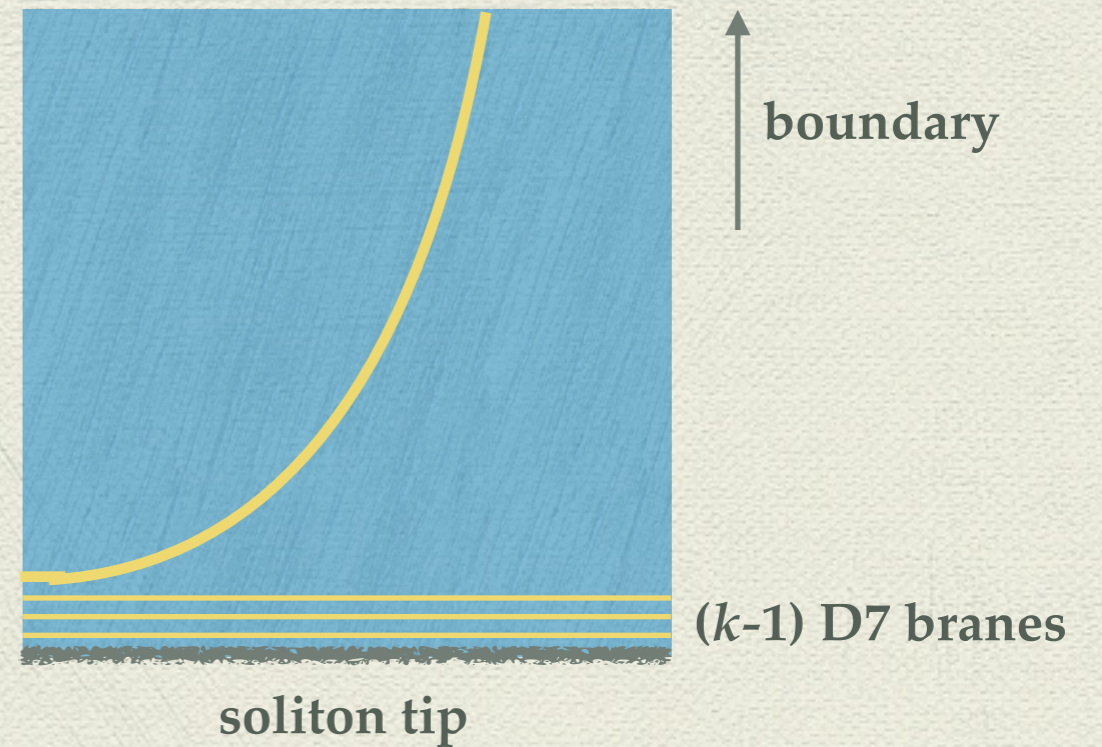
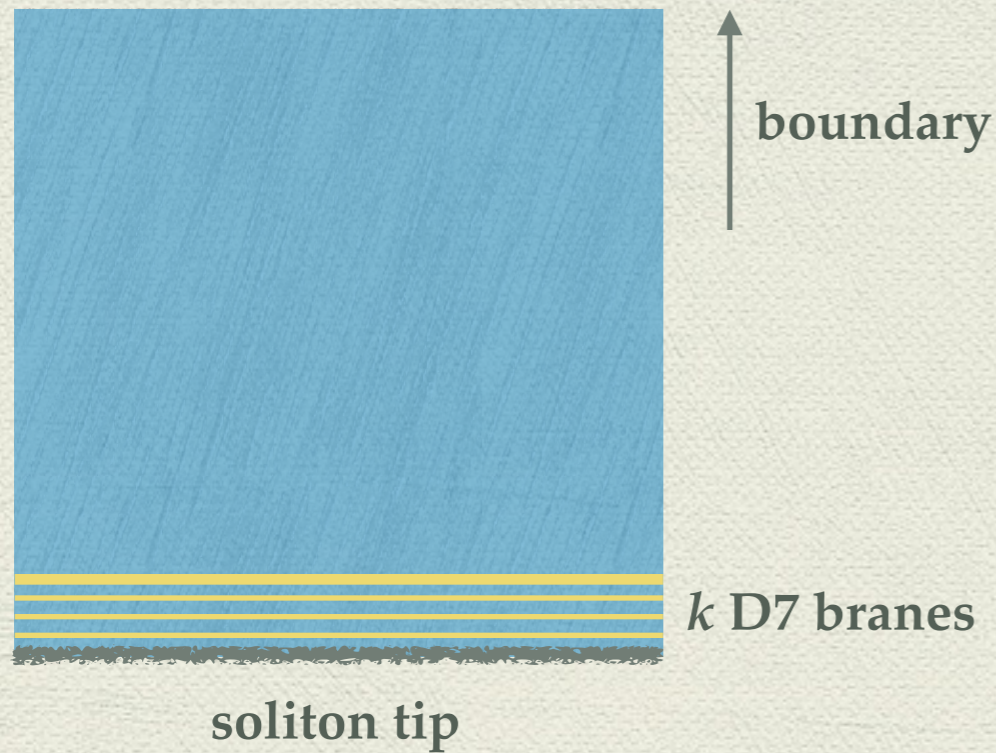
String worldsheet can attach to D7 probe branes:



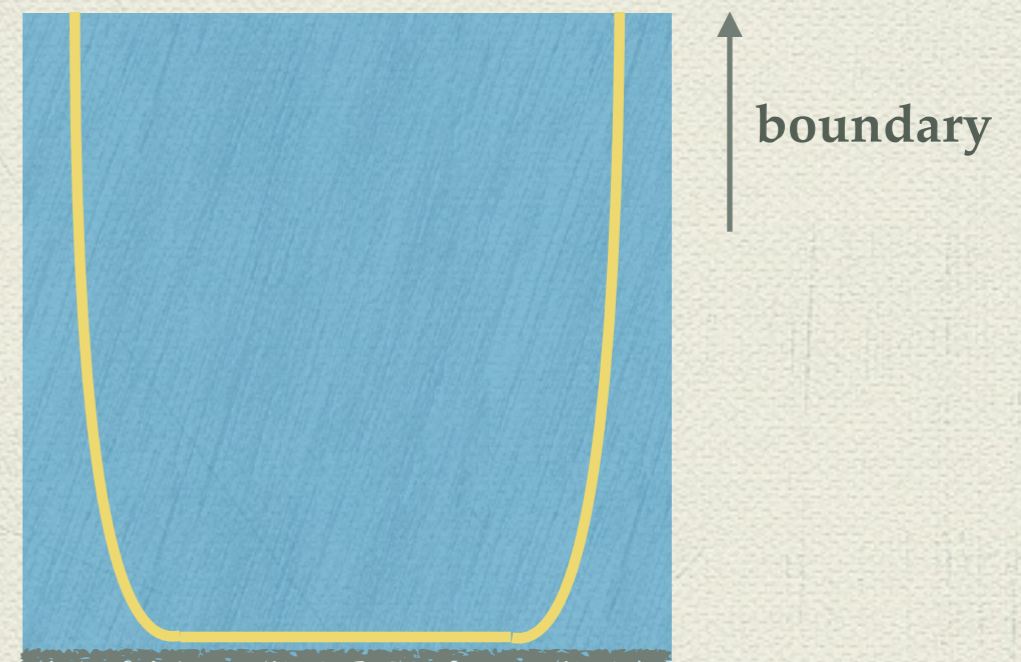
Dominant contribution to on-shell action from Wilson loop on D7 brane => level-rank duality

# Holographic Embedding of Domain Walls

Interested in defects changing the level  $k \Rightarrow$  need to change the number of branes



Will mostly consider interval of finite length with level  $k = 1$ ;  $k = 0$  elsewhere.



# Spectrum of operators

Consider only bosonic operators uncharged under  $SO(6)$  or  $SO(2)$  global symmetries.

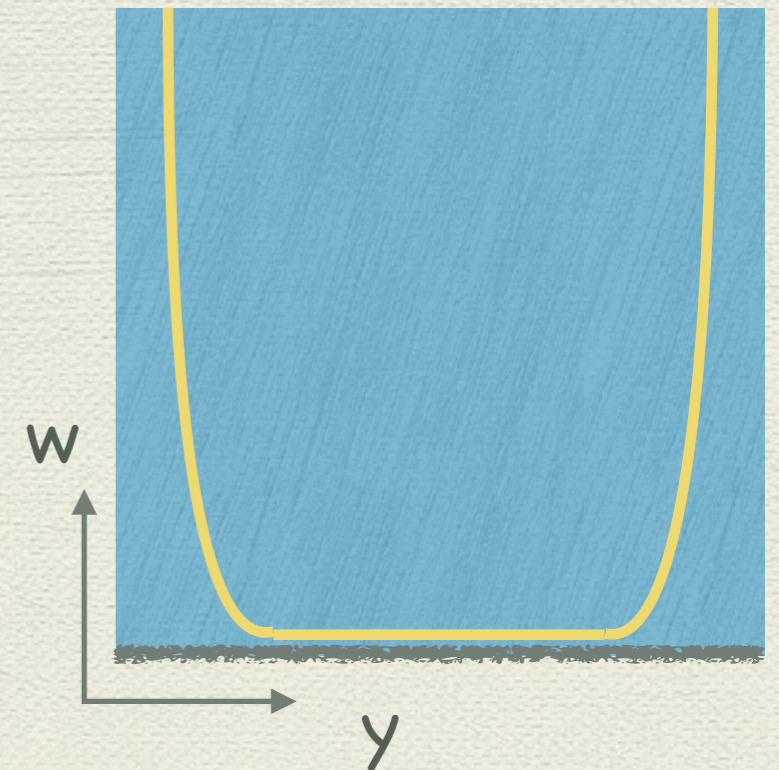
The brane is extended infinitely in the  $t, x$  directions.

WZ term +  $F_5$  flux  $\Rightarrow$  CS term on brane worldvolume

Level =  $N$ , number of D3 branes  $\Rightarrow$

Probe brane theory = DBI + CS at level  $N$

Fields (near right-hand boundary; left-hand boundary fields have opposite chirality):



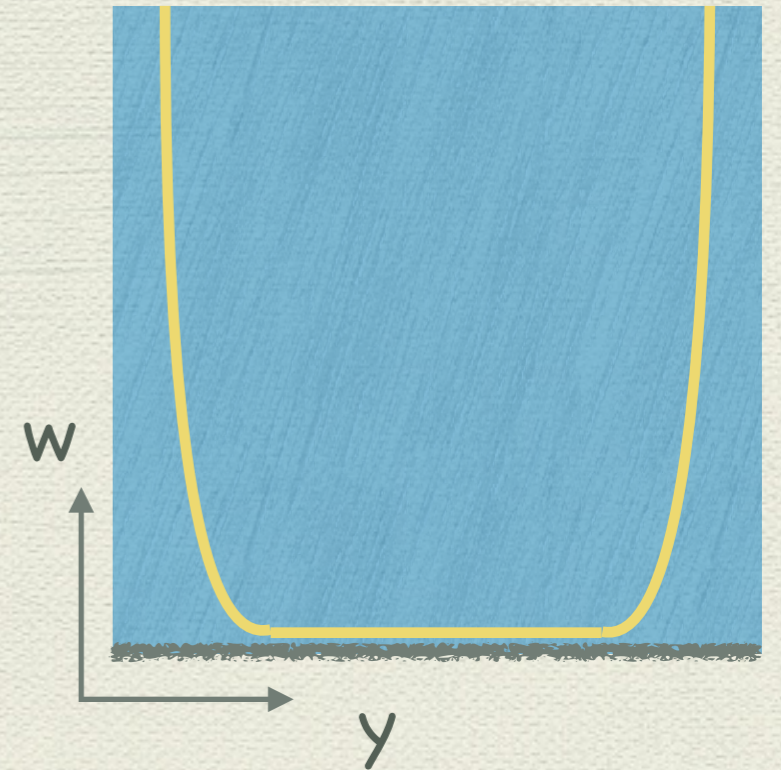
Fields	Asymptotic form
$y$	$y_0 - \frac{Rw_*^3}{4w^4} + \dots$
$A$	$a_-^{(4)} w^4 + \dots + (\text{flat})$
$A$	$(a_-^{(4)} \text{ terms}) + (\text{flat}) + a_+^{(-4)} w^{-4} + \dots$



# Spectrum of operators

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Source	Response	Dimension
$y$		3
$a$	$a$	5
$a$	$a$	1



← irrelevant chiral operator

← chiral current =>  $U(k)$  symmetry

# Spectrum of operators

Harvey, Roysten  
0804.2854

Source	Response	Dimension
$\gamma$		3
$\mathbf{a}$	$\mathbf{a}$	5
$\mathbf{a}$	$\mathbf{a}$	1

*c.f.*

Operator
$\Phi^{(3)} \sim \bar{q}\gamma_+ F_{-y} q$
$V^{(5)} \sim \bar{q}\gamma_+ F_{-y} F_{-y} q$
$J_+ \sim \bar{q}\gamma_+ q$

D3-D7 brane intersection

# System Properties

**Metric:**  $ds^2 = w^2 dx_3^2 + f(w)w^2 d\theta^2 + \frac{R^2}{f(w)} \frac{dw^2}{w^2}$

**Embedding:**  $y(w) = \int_{w_*}^w \frac{R w^4 dw}{\sqrt{(w^6 - w_*^6)(w^4 - w_0^4)}}$

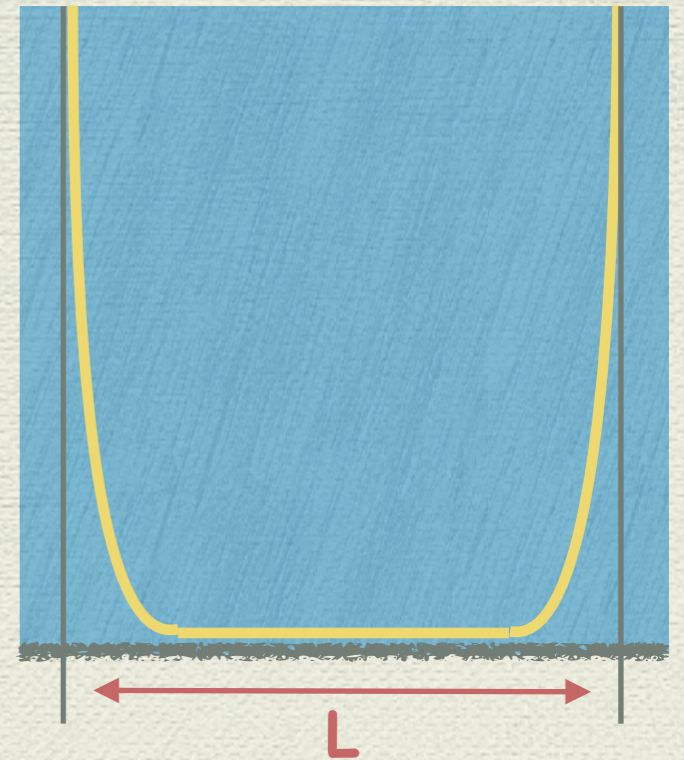
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**Free energy:**  $F = 2R \left( \int_{w_*}^{\Lambda} \frac{w^6 dw}{\sqrt{(w^6 - w_*^6)(w^4 - w_0^4)}} - \frac{1}{2} \Lambda^2 \right)$

For  $L$  large,  $F$  grows linearly with length.

Behavior due to the bulk geometry ending; if the effective background were  $AdS_4$ , there would be no  $L$  dependence.

Interpretation as order parameter for confinement?



# System Properties

c.f.: Yee, Zahed 1103.6286

## Edge mode [right-hand side chiralities]

Conserved current sector given by boundary values of flat connections:  $A = d\varphi$

Source:  $A_+^{(\text{flat})} = \text{external gauge field } \mathcal{A}_+$       VeV:  $A_-^{(\text{flat})} = \text{current } j_-$

Flatness condition:  $\partial_+ A_- = \partial_- A_+$

When source is turned off, flatness  $\Rightarrow \partial_+ j_- = 0$  so current is chiral.

## Anomaly [right-hand side chiralities]

c.f.: Jensen 1012.4831  
Yee, Zahed 1103.6286

Flatness condition:  $\partial_+ A_- = \partial_- A_+ \Rightarrow \partial_\mu j^\mu = -\partial_+ j_- \propto \partial_- \mathcal{A}_+ = E_{\text{external}}$

anomaly in 2 dimensions

Consistent with expectation from brane construction.

# System Properties

**Irrelevant vector operator:**

$A_-|_{\text{right}} \sim a_- w^4 + \dots$  at the right-hand boundary  
leads to  $A_-|_{\text{left}} \sim c a_- w^{-4} + \dots$  at the left boundary,  
and the opposite for  $A_+$

Thus,  $V_+|_{\text{right}} \sim A_+|_{\text{left}}$  and  $V_-|_{\text{left}} \sim A_-|_{\text{right}}$ ;

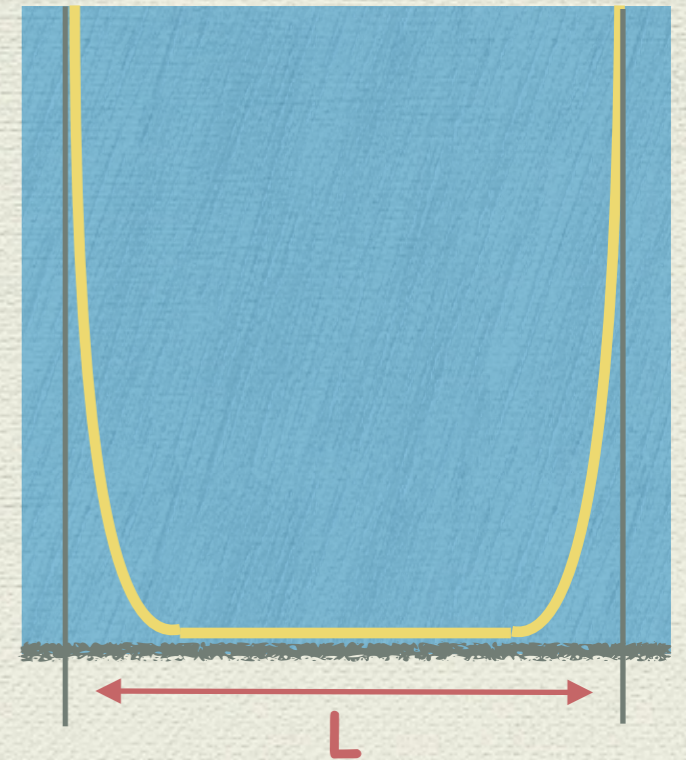
the field varied to obtain the VeV does not determine the VeV.

$\langle V_-|_{\text{left}} V_+|_{\text{right}} \rangle \sim e^{-2\pi L/l}$  with  $l$  the circle compactification length.

Implies mediation by field with mass on the KK scale.

**Meson spectrum, response to bulk fields, general correlation functions,  
multiples types of phase transitions, ...**

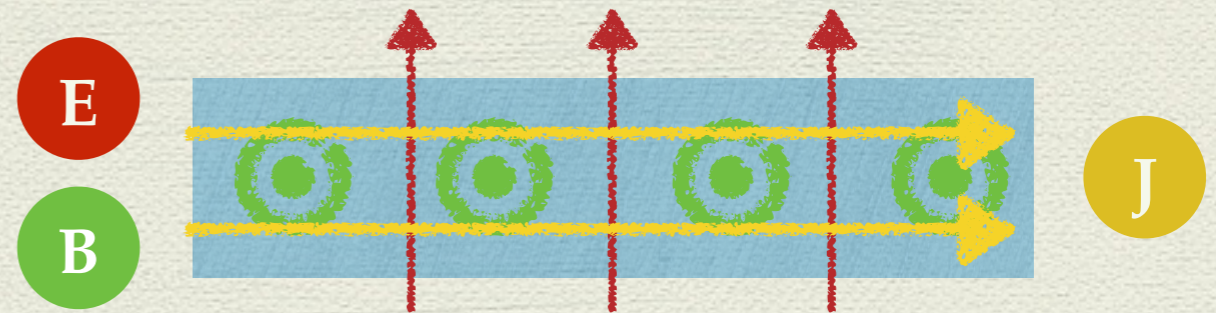
**Surprisingly rich.**



# Fractional Quantum Hall Effect

Quantum Hall Effect:

Apply a strong magnetic field normal to a conducting film



When an electric field is applied, a perpendicular current flows:  $\mathbf{E} \perp \mathbf{J}$

For strong magnetic fields, this conductivity becomes quantized.

Effective field theory description by Chern-Simons action.

Gapped, exhibits edge currents, relies crucially on interactions (unlike IQHE)

Response to external EM field (encoded in  $B$  field) reproduces response in edge currents (probably)

Some peculiar subtleties in holographic renormalization...

Examples of Hall behavior in holographic systems may shed light on FQHE.

# Summary

Holography for level-changing defects in YM-CS

Few degrees of freedom

Surprisingly rich structure:

anomalies, affine Lie symmetries, unusual correlation properties, relationship with confinement, (large  $N$ ) symmetry breaking, phase transitions...

Relationship to FQH

終わり