

Gravity theory on Poisson manifold with R-flux



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References

1408.2649 [hep-th] (Int.J.Mod.Phys. A 30 (2015) 17,1550097)

1508.05706 [hep-th] (Fortsch.Phys. 63 (2015) 683-704)

Space-time Geometry probed with strings

String theory would describe quantum gravity

| | | |
|---------------------|---------------------------------------|---------------|
| | Quantum Gravity General Relativity | String Theory |
| Fundamental Objects | Point Particles | Strings |
| Space-time Geometry | Riemannian | ? ? ? ? ? |

Typical Example of T-duality

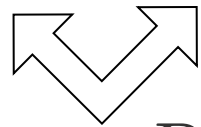
Compactifying on $X^9 \sim X^9 + 2\pi R$ (periodic)

Contribution to KK tower (Mass spectrum of reduced theory)

= KK momenta + Windings

$$P_9 = \frac{K}{R}$$

$$W_9 = \frac{1}{2\pi\alpha'} 2\pi N R = \frac{R}{\alpha'} N$$



$$R \leftrightarrow \frac{\alpha'}{R}$$

Reduced theories are physically equivalent:

T-duality

[84 Kikkawa, Yamasaki]

Why Generalized Geometry ?

T-duality implies a physical **equivalence**

between **two different background geometries**

(configurations of **space-time metric** and **NSNS B -field**)

suggesting the appearances of

- **Strange metric (T-folds)** [04 Hull], ...

- **Non-geometric fluxes** [05 Shelton, Taylor, Wecht], ...

- **Exotic branes** [10 de Boer, Shigemori], ...

➤ Machinery treating these as “geometry”

⇒ (Poisson) Generalized Geometry

Plan of Today's Talk

- Introduction & Motivations
- A Little Bit More on T-duality
- Generalized Geometry
 - Definitions & Properties
 - Generalized Riemannian Geometry
- Poisson Generalized Geometry
 - Definitions & Properties
 - Poisson Generalized Riemannian Geometry

A Little Bit More on T-duality: Buscher rule

T-duality: Physical equiv. between two backgrounds

$$(g, B) \sim (\tilde{g}, \tilde{B})$$

given by Buscher rule (“0”: isometry) [87 Buscher]

$$\tilde{g}_{00} = \frac{1}{g_{00}}, \quad \tilde{g}_{0i} = \frac{B_{0i}}{g_{00}}, \quad \tilde{g}_{ij} = g_{ij} - \frac{g_{0i}g_{0j} - B_{0i}B_{0j}}{g_{00}},$$

$$\tilde{B}_{0i} = \frac{g_{0i}}{g_{00}}, \quad \tilde{B}_{ij} = B_{ij} - \frac{g_{0i}B_{0j} - g_{0j}B_{0i}}{g_{00}}.$$

Metric and **B-field** should be on **equal footing**

Generalized Geometry

[04 Hitchin, '04 Gualtieri]

Consider tangent and cotangent bundles at the same time!

- Canonical Inner Product

$$\langle X + \xi, Y + \eta \rangle = \frac{1}{2}(i_X \eta + i_Y \xi)$$

➤ Invariant under an action of $O(D,D)$: T-duality transf.

- Operations corresponding to diffeo. + B-field gauge transf.

$$[X + \xi, Y + \eta] = [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2}d(i_X \eta - i_Y \xi)$$

Riemannian Geom. based on Gen. Geom.

- $O(D,D)$ -invariant inner product

⇒ Decompose in Positive-/Negative-definite subbundles

$$C_{\pm} = \{X + (\pm g + B)(X) | X \in \Gamma(TM)\}$$

- Define a connection ∇ on positive-def. subbundle C_+ :

Coefficients = **Christoffel Symbol** + **H-flux** (B -field's field strength)

$$\nabla_{\partial_i} (\partial_j)^+ = g^{lk} (2\underline{\Gamma_{kij}} + \underline{H_{kij}}) (\partial_l)^+$$

Gravity based on Gen. Geom.

- Curvature tensor:

$$R(X, Y)u := (\nabla_X \nabla_Y - \nabla_Y \nabla_X - \nabla_{[X, Y]_C})u$$

- Ricci scalar:

$$\mathcal{R} - \frac{1}{4} H^{ijk} H_{ijk}$$

- Einstein-Hilbert-like action :

$$S = \int d^D x \sqrt{g} \left(\mathcal{R} - \frac{1}{4} H^2 \right)$$

➤ This is the same as NSNS-sector of SUGRA

G.G. would be a good tool to “geometrize” string theory

Variant of Generalized Geometry

Slightly modifying structures of GG would be interesting :

- ❖ Vector + 1-form \Rightarrow unchanged
- ❖ Inner product (T-dual) \Rightarrow unchanged
- ❖ diffeo. + B-field gauge trnsf. \Rightarrow **changed!**

$$[X + \xi, Y + \eta] = [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2} d(i_X \eta - i_Y \xi)$$

\Rightarrow **Then how do we change it ???**

Hint: A chain of T-duality

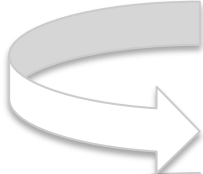
String action $\mathcal{L} \sim g(\dot{X}^2 + X'^2) + B\dot{X}X'$


⇒ conjugate momenta : $\Pi_i \sim g_{ij}\dot{X}^j + B_{ij}X'^j$

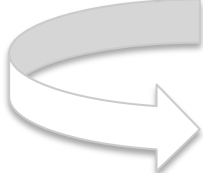
$[p_i, p_j] \sim H_{ijk}w^k$

momentum winding

T-duality : $p_i \leftrightarrow w^i = X'^i$

T  $[p_i, w^k] \sim F_{ij}^k w^j$

T  $[p_i, w^k] \sim Q_i^{kj} p_j$

T  $[w^i, w^j] \sim R^{ijk} p_k$

Fluxes associated with non-geom. background

Fluxes in Generalized Geometry

$$H_{ijk} \longleftrightarrow F_{ij}^k \quad \Big| \quad Q_k^{ij} \longleftrightarrow R^{ijk}$$



$$[\partial_i, \partial_j]$$



$$[\partial_i, dx^k]$$

$$\sim H_{ijk} dx^k$$

$$\sim F_{ij}^k dx^j$$

Understood in terms of GG

$$[X, Y] \neq 0 \quad [\xi, \eta] = 0$$



$$[\partial_i, dx^j]$$



$$[dx^i, dx^j]$$

$$\sim Q_{ik}^j \partial_k ??$$

$$\sim R^{ijk} \partial_k ??$$

MISSING in GG!!

$$[\xi, \eta] \neq 0 \quad [X, Y] = 0$$



Variant of GG interchanging roles of Vector and 1-form

Poisson Geometry

Poisson bi-vector $\theta = \frac{1}{2}\theta^{ij}\partial_i \wedge \partial_j$

-Poisson bracket $\{f, g\} = \theta^{ij}\partial_i f \partial_j g$

-Jacobi id. $\{\{f, g\}, h\} + \{\{g, h\}, f\} + \{\{h, f\}, g\} = 0$

$$\Leftrightarrow [\theta, \theta]_S = 0 \quad \text{Poisson cond. } \theta^{l[i}\partial_l\theta^{jk]} = 0$$

Schouten bracket: **e.g.** $[X \wedge Y, Z]_S = [X, Z] \wedge Y - [Y, Z] \wedge X$
Extension of Lie br. to multi-vector

Lie algebra on 1-forms

$$\text{Lie bracket: } [\xi, \eta]_\theta = \mathcal{L}_{\theta(\xi)}\eta - \mathcal{L}_{\theta(\eta)}\xi + d(\theta(\eta, \xi))$$

Cartan Algebra on Poly-Vector Fields

“Interior product” \bar{i}_ξ $\bar{i}_\xi(X \wedge Y) = (i_X \xi)Y - (i_Y \xi)X$

“Exterior derivative” $d_\theta = [\theta, \cdot]_S$

-Nilpotent $d_\theta^2 = 0 \iff [\theta, \theta]_S = 0$

“Lie derivative” $\bar{\mathcal{L}}_\zeta f := \bar{i}_\zeta d_\theta f$

$\bar{\mathcal{L}}_\zeta \xi := [\zeta, \xi]_\theta$ $([\xi, \eta]_\theta = \mathcal{L}_{\theta(\xi)}\eta - i_{\theta(\eta)}d\xi)$

$\bar{\mathcal{L}}_\zeta X := (d_\theta \bar{i}_\zeta + \bar{i}_\zeta d_\theta)X$

“Cartan algebra”: enables diff. calculus induced by 1-form

$$\{\bar{i}_\xi, \bar{i}_\eta\} = 0, \quad \{d_\theta, \bar{i}_\xi\} = \bar{\mathcal{L}}_\xi, \quad [\bar{\mathcal{L}}_\xi, \bar{i}_\eta] = \bar{i}_{[\xi, \eta]_\theta},$$

$$[\bar{\mathcal{L}}_\xi, \bar{\mathcal{L}}_\eta] = \bar{\mathcal{L}}_{[\xi, \eta]_\theta}, \quad [d_\theta, \bar{\mathcal{L}}_\xi] = 0.$$

Poisson Generalized Geometry 1408.2649

- Vector field and 1-form (same as GG)
- $O(D,D)$ -invariant inner product (same as GG)

Poisson structure

- Operations different from GG : Vector field \longleftrightarrow 1-form

$$[X + \xi, Y + \eta] = [\xi, \eta]_{\theta} + \bar{\mathcal{L}}_{\xi}Y - \bar{\mathcal{L}}_{\eta}X - \frac{1}{2}d_{\theta}(\bar{i}_{\xi}Y - \bar{i}_{\eta}X)$$

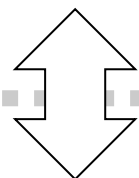
cf. $[X + \xi, Y + \eta]_C = [X, Y] + \mathcal{L}_X\eta - \mathcal{L}_Y\xi - \frac{1}{2}d(i_X\eta - i_Y\xi)$

Physical intuition of **New Operation**

- Operation in GG :

$$[X + \xi, Y + \eta]_C = [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2} d(i_X \eta - i_Y \xi)$$

$$\Pi_i = p_i + B_{ij} w^j \quad \uparrow$$



T-dual (*cf.* [Marc's talk](#))

$$\tilde{\Pi}^i = w^i + \beta^{ij} p_j \quad \Downarrow \text{Poisson structure } \theta$$

- Operation in PGG :

$$[X + \xi, Y + \eta] = [\xi, \eta]_\theta + \mathcal{L}_\xi Y - \mathcal{L}_\eta X - \frac{1}{2} d_\theta(i_\xi Y - i_\eta X)$$

- $O(D,D)$ -inv inner product
 \Rightarrow Decompose in Positive-/Negative-definite subbundles

$$C_{\pm} = \{ \xi + (\pm G + \beta)(\xi) \mid \xi \in \Gamma(T^*M) \}$$

- Define a connection $\bar{\nabla}$ on positive-def. subbundle C_+ :

Coefficients = **Contravariant Levi-Civita** + **R-flux**

$$\bar{\nabla}_{dx^i} (dx^j)^+ = \left(\underline{2\bar{\Gamma}_k^{ij}} + \underline{G_{lk} R^{kij}} \right) (dx^l)^+$$

$$R = d_{\theta} \beta = [\theta, \beta]_S$$

$$= (\theta^{l[i} \partial_l \beta^{jk]} + \beta^{l[i} \partial_l \theta^{jk]}) \partial_i \wedge \partial_j \wedge \partial_k$$

Contravariant Levi-Civita conn. [00 Fernandes]...

$$\left[\begin{aligned} \bar{\Gamma}_k^{\{ij\}} &= \frac{1}{2} [\theta^{mn} (\partial_m G^{ji}) - \theta^{mi} (\partial_m G^{jn}) \\ &\quad - \theta^{mj} (\partial_m G^{in}) - G^{jl} (\partial_l \theta^{in}) - G^{il} (\partial_l \theta^{jn})] G_{nk} \\ \bar{\Gamma}_k^{[ij]} &= \frac{1}{2} (\partial_k \theta^{ij}) \end{aligned} \right.$$

Two characteristic tensors G^{ij} & θ^{ij}

$$\Rightarrow \left[\begin{aligned} \bar{\nabla}_{dx^k} G^{ij} &= 0 && : \text{Compatible with metric} \\ \bar{\nabla}_{dx^i} \theta^{jk} + (\text{cyclic}) &= 0 && : \text{Respecting Poisson struc.} \end{aligned} \right.$$

In particular, θ^{ij} is covariantly constant $\nabla_i \theta^{jk} = 0 \Rightarrow$ LC conn.


Extension of the Levi-Civita respecting Poisson structure !

Gravity based on PGG

- Curvature tensor :

$$\bar{R}(\xi, \eta)u := (\bar{\nabla}_\xi \bar{\nabla}_\eta - \bar{\nabla}_\eta \bar{\nabla}_\xi - \bar{\nabla}_{[\xi, \eta]})u$$

- Ricci scalar:

$$\bar{\mathcal{R}} - \frac{1}{4} R^{ijk} R_{ijk}$$


- Einstein-Hilbert-like action :

$$S = \int d^D x \sqrt{G} \left(\bar{\mathcal{R}} - \frac{1}{4} R^2 \right)$$

Summary

We gave

- a **new geometric framework** based on **Poisson structure**
 - T-dual counterpart of Generalized Geometry
- a well-defined **formulation of R -flux**
 - defined as a field strength of local bi-vector gauge potentials
- an **Riemann geometry** compatible with **Poisson structure**

Future directions

We established Riemann Geometry compatible with Poisson

- Well describable **Non-geometric background** ?
- Extension to **quasi-Poisson**, **Nambu-Poisson** structures
- Poisson is semi-classical limit of **Non-commutativity**
⇒ “Riemann Geometry” on Non-commutative space
- etc.....