# Self-accelerating universe from nonlinear massive gravity

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# Outline

### Introduction;

• The nonlinear massive gravity theory

• Self-accelerating solutions in open FRW universe;

- The first workable model !
- Conclusion and Discussion
  - Cosmological perturbations



#### Cosmic acceleration



• Can we give graviton a mass?

• Fierz and Pauli 1939

 $\mathcal{L}_{FP} = f^4 \left( h_{\mu
u} h_{\mu
u} - h^2 
ight)$ 

van Dam-Veltman-Zakharov discontinuity

$$T^{\mu}_{\nu}h^{
u}_{\mu} = T^{\mu}_{\nu}(\hat{h}^{
u}_{\mu} + m^2_g \delta^{
u}_{\mu}\phi) = T^{\mu}_{\nu}\hat{h}^{
u}_{\mu} + rac{1}{M_{
m Pl}}T\phi^{\mu}$$

- Vainshtein 1972 non-linear interactions
- Boulware–Deser (BD) ghost 1972

Lack of Hamiltonian constrain and momentum constrain

6th dof is BD ghost!



• Whether there exist a nonlinear model without ghost?

- N. Arkani–Hamed et al 2002 decoupling limit  $m_g \to 0$ ,  $M_p \to \infty$ ,  $\Lambda = \left(m_g^4 M_p\right)^{1/5} \to \text{finite}$ .
- C. de Rham and G. Gabadadze 2010

$$\mathcal{L} = M_{\rm Pl}^2 \sqrt{-g}R - \frac{M_{\rm Pl}^2 m^2}{4} \sqrt{-g} (U_2(g, H) + U_3(g, H) + U_4(g, H) + U_5(g, H) \cdots),$$



where  $U_i$  denotes the interaction term at *i*th order in  $H_{\mu\nu}$ ,

$$\begin{split} U_{2}(g,H) &= H_{\mu\nu}^{2} - H^{2}, \\ U_{3}(g,H) &= \underline{c_{1}}H_{\mu\nu}^{3} + \underline{c_{2}}HH_{\mu\nu}^{2} + \underline{c_{3}}H^{3}, \\ U_{4}(g,H) &= \underline{d_{1}}H_{\mu\nu}^{4} + \underline{d_{2}}HH_{\mu\nu}^{3} + \underline{d_{3}}H_{\mu\nu}^{2}H_{\alpha\beta}^{2} \\ &+ \underline{d_{4}}H^{2}H_{\mu\nu}^{2} + \underline{d_{5}}H^{4}, \end{split} \\ U_{5}(g,H) &= \underline{f_{1}}H_{\mu\nu}^{5} + \underline{f_{2}}HH_{\mu\nu}^{4} + \underline{f_{3}}H^{2}H_{\mu\nu}^{3} + \underline{f_{4}}H_{\alpha\beta}^{2}H_{\mu\nu}^{3} \\ &+ \underline{f_{5}}H(H_{\mu\nu}^{2})^{2} + \underline{f_{6}}H^{3}H_{\mu\nu}^{2} + \underline{f_{7}}H^{5}. \\ g_{\mu\nu} &= \eta_{\mu\nu} + \frac{h_{\mu\nu}}{M_{\text{Pl}}} = H_{\mu\nu} + \eta_{ab}\partial_{\mu}\phi^{a}\partial_{\nu}\phi^{b}, \\ &+ \underline{d_{4}}H^{2}H_{\mu\nu}^{2} + \underline{d_{5}}H^{4}, \\ H_{\mu\nu} &= \frac{h_{\mu\nu}}{M_{\text{Pl}}} + \partial_{\mu}\pi_{\nu} + \partial_{\nu}\pi_{\mu} - \eta_{\alpha\beta}\partial_{\mu}\pi^{\alpha}\partial_{\nu}\pi^{\beta}. \end{split}$$

• C. de Rham, G. Gabadadze and A. Tolly 2011

$$\begin{split} I_g &= M_{Pl}^2 \int d^4 x \sqrt{-g} \left[ \frac{R}{2} + m_g^2 (\mathcal{L}_2 + \alpha_3 \mathcal{L}_3 + \alpha_4 \mathcal{L}_4) \right] \\ \mathcal{L}_2 &= \frac{1}{2} \left( [\mathcal{K}]^2 - [\mathcal{K}^2] \right) , \\ \mathcal{L}_3 &= \frac{1}{6} \left( [\mathcal{K}]^3 - 3[\mathcal{K}][\mathcal{K}^2] + 2[\mathcal{K}^3] \right) , \\ \mathcal{L}_4 &= \frac{1}{24} \left( [\mathcal{K}]^4 - 6[\mathcal{K}]^2[\mathcal{K}^2] + 3[\mathcal{K}^2]^2 + 8[\mathcal{K}][\mathcal{K}^3] - 6[\mathcal{K}^4] \right) , \end{split}$$

It is often called fiducial metric

- Automatically produce the "appropriate coefficients" to eliminate BD ghost at any order in decoupling limit !
- Free of BD ghost away from the decoupling limit, at fully nolinear level

Hassan, Rosen '11

Stukelberg fields

A.Emir Gumrukcuoglu, Chunshan Lin, Shinji Mukohyama

arXiv:1109.3845

- No go result for flat FRW solution (G. D'Amico et al 2011 Aug.)
- O However... (A.Gumrukcuoglu, C. Lin, S. Mukohyama: 1109.3845)

It does not extend to open FRW universe

$$g_{\mu\nu}dx^{\mu}dx^{\nu} = -N(t)^{2}dt^{2} + a(t)^{2}\Omega_{ij}dx^{i}dx^{j},$$
  
$$\Omega_{ij}dx^{i}dx^{j} = dx^{2} + dy^{2} + dz^{2} - \frac{|K|(xdx + ydy + zdz)^{2}}{1 + |K|(x^{2} + y^{2} + z^{2})},$$

The 4 Stukelberg scalars

motivated by...

$$\phi^{0} = f(t)\sqrt{1 + |K|(x^{2} + y^{2} + z^{2})},$$
  

$$\phi^{1} = \sqrt{|K|}f(t)x,$$
  

$$\phi^{2} = \sqrt{|K|}f(t)y,$$
  

$$\phi^{3} = \sqrt{|K|}f(t)z.$$
  
Minkowski  
metric  

$$\phi$$
  
Open FRW  
chart

• Fiducial metric respect FRW symmetry

$$\begin{split} \phi^{0} &= f(t)\sqrt{1+|K|(x^{2}+y^{2}+z^{2})}, \\ \phi^{1} &= \sqrt{|K|}f(t)x, \\ \phi^{2} &= \sqrt{|K|}f(t)y, \\ \phi^{3} &= \sqrt{|K|}f(t)z. \end{split} \qquad f_{\mu\nu} \equiv \eta_{ab}\partial_{\mu}\phi^{a}\partial_{\nu}\phi^{b} \\ = -(\dot{f}(t))^{2}\delta^{0}_{\mu}\delta^{0}_{\nu} + |K|f(t)^{2}\Omega_{ij}\delta^{i}_{\mu}\delta^{j}_{\nu}, \end{split}$$

- (0i) components of the equation of motion for  $g_{\mu\nu}$  are trivially satisfied;
- Evolution equations for cosmic perturbations fully respect homogeneity and isotropy at any order.

• Constraint from Stuckelberg scalars:

$$\begin{aligned} (\dot{a} - \sqrt{|K|}N) \left[ \left( 3 - \frac{2\sqrt{|K|}f}{a} \right) + \alpha_3 \left( 3 - \frac{\sqrt{|K|}f}{a} \right) \left( 1 - \frac{\sqrt{|K|}f}{a} \right) \\ + \alpha_4 \left( 1 - \frac{\sqrt{|K|}f}{a} \right)^2 \right] &= 0. \end{aligned}$$

• Branch I

$$\dot{a} = \sqrt{|K|}N,$$

• Branch  $||_{\pm}$ 

$$f = \frac{a}{\sqrt{|K|}} X_{\pm}, \quad X_{\pm} \equiv \frac{1 + 2\alpha_3 + \alpha_4 \pm \sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4}}{\alpha_3 + \alpha_4}$$

Please notice that these 2 solutions do not exist when K=0.

• Freedmann equation

$$3H^2 - \frac{3|K|}{a^2} = \rho_m + c_{\pm}m_g^2$$
$$-\frac{2\dot{H}}{N} - \frac{2|K|}{a^2} = \rho_m + p_m,$$

where

$$c_{\pm} \equiv -\frac{1}{(\alpha_3 + \alpha_4)^2} \left[ 1 + \alpha_3 \pm \sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4} \right] \\ \times \left[ 1 + \alpha_3^2 - 2\alpha_4 \pm (1 + \alpha_3)\sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4} \right].$$

The effective cosmological constant

$$\Lambda_{\pm} = c_{\pm} m_g^2.$$



$$\Lambda_{\pm} \equiv -\frac{m_g^2}{(\alpha_3 + \alpha_4)^2} \left[ (1 + \alpha_3) \left( 2 + \alpha_3 + 2 \alpha_3^2 - 3 \alpha_4 \right) \pm 2 \left( 1 + \alpha_3 + \alpha_3^2 - \alpha_4 \right)^{3/2} \right]$$

### Conclusion and discussion

- The nonlinear massive gravity theory
- For Minkowski fiducial metric, only K<0 FRW solution exists</li>
   Extensions of the theory with generic fiducial metric Hassan, Rosen, 11
   Cosmological perturbations
  - Scalar sector & vector sector ... ?
  - Tensor sector ... !

