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GRAVITATIONAL LENSING IN A DARK MATTER FREE BRANEWORLD MODEL

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Outline of presentation

- Some background information
- Our model setting
- Gravitational Lensing in braneworld models
- □ The geometry of extra dimension

Status of dark matter problem

- 1. Evidences of dark matter always inferred from motion under gravitational interaction
- 2. No simple theoretical framework for dark matter
- Supersymmetry (SUSY) offers a beautiful solution
- Signal of Higgs candidate support SUSY
- Detection of sparticle has not succeeded yet
- Dark matter direct detections also give no results yet
- Behavior of sparticle is not clear, alternative approach cannot be ignored.

Randall-Sundrum models

- 2-brane model
- 2 brane located at end point of a orbifold, AdS5 volume in between
- 1-brane model one of the brane located at infinity
- Showing a new way to hide extra dimension



Braneworld as modified 4D Einstein theory

- The matter on membrane infer discontinuity of extrinsic curvature (Israel, 1966)
- Project 5D Einstein equation to 4D (Shiromizu, Maeda, Sasaki, 2000)



The dark radiation E_{µv}

A traceless tensor from projection of 5D Weyl tensor

$$E_{\mu\nu} = {}^{(5)} C_{ABCD} n^C n^D g^A_\mu g^B_\nu$$

- Weyl tensor = component of curvature that is not governed by Einstein equation
- □ It could be "sourced" by the brane via $\nabla^{\mu}S_{\mu\nu} = k_5^4 \nabla^{\mu}E_{\mu\nu}$
- It is shown to be small so that Randall Sundrum model converge to Newtonian gravity
- Is it possible to sufficiently change the gravity? (e.g. by giving up the AdS5 geometry)

Model setting – A phenomenological approach

Assume a corrected Schwarzschild metric

+

$$ds^{2} = -\left(1 - \frac{r_{b}}{r} + B_{t}(r)\right)dt^{2} + \frac{dr^{2}}{1 - \frac{r_{b}}{r} + B_{r}(r)} + r^{2}d\Omega$$

Bt can be fixed by orbital velocity profile in rotation curves

$$v_{tg}^2 = \frac{rg_{tt}}{2g_{tt}}$$

Obtain the differential equation of Br by traceless property of dark radiation, e.g.

$$B_{t}' - \frac{2v_{tg}^{2}}{r} \left(1 - \frac{r_{b}}{r} + B_{t}\right) + \frac{r_{b}}{r^{2}} = 0^{5}$$

$$(2 + v_{tg}^{2})B_{r}' + \left(\frac{2}{r} + 4v_{tg}v_{tg}' + \frac{2v_{tg}^{2}}{r} + \frac{2v_{tg}^{4}}{r}\right)B_{r}$$

$$v_{tg}^{4} \left(\frac{2}{r} - \frac{2r_{b}}{r^{2}}\right) + v_{tg}^{2} \left(\frac{2}{r} - \frac{r_{b}}{r^{2}}\right) + 4v_{tg}v_{tg}' \left(1 - \frac{r_{b}}{r}\right) = 0$$
Distance

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Constant velocity limit

- □ We consider a limit of flat rotation curve.
- In dark matter picture, the motion is supported by a Halo with r^{-2} density.
- In braneworld, the metric correction is given by

$$B_t(r) \to -1 + \frac{r_b}{r} + C_t r^{2v_{\infty}^2}$$
$$B_r(r) \to r^{-(v_{\infty}^2+1)} \left[C_r - r^{\frac{v_{\infty}^2}{2}} \left(v_{\infty}^2 r - r_b \right) \right]$$

 Test if these metric corrections contradict what dark matter predict

Motion in weak field low velocity limit

For weak field and slow velocity, acceleration can be approximated by

 $\frac{d^2 x^{\mu}}{dt^2} = -\Gamma^{\mu}_{00}$

For radial in-falling motion

$$a = \frac{1}{2r} \left(1 - \frac{r_b}{r} + B_r \right) \left(\frac{r_b}{r} + rB'_t \right)$$
$$\approx \frac{v_\infty^2}{r} + \frac{r_b^2}{2r^3} - \frac{r_b B_r}{2r^2} \approx \frac{v_\infty^2}{r}$$

Consistent with dark matter prediction

Gravitational deflection of light?

A Simple lensing model



Compare angle between asymptotes

$$\begin{split} \delta(r_{\min}) &= 2 \int_{r_{\min}}^{\infty} \mathcal{I} - \pi \\ \mathcal{I}(r) &= \frac{1}{r} \left\{ \frac{g_{rr}(r)}{\left[g_{tt}(r_{\min})/g_{tt}(r)\right] \left(r/r_{\min}\right)^2 - 1} \right\}^{1/2} \\ & \text{K. C. Wong, YITP Kyoto, 3/3/2012} \end{split}$$

Series expansion of small parameters

 \square Small parameters $r_b/r, v_\infty^2$

$$\delta = -\int_{r_{\min}}^{\infty} \frac{r_{\min}^2 - r^2 + 2\ln(\frac{r_{\min}}{r})r^2}{(r^2 - r_{\min}^2)r} \sqrt{\frac{r_{\min}^2}{r^2 - r_{\min}^2}} v_{\infty}^2 dr$$

$$\square$$
 Integrated to
$$\delta = \frac{3\pi v_\infty^2}{2}$$

For dark matter that produce the same rotation curve

$$\delta = 4v_{\infty}^2$$

Large distance behavior beyond constant velocity

- NFW model dark matter distribution that motivated by gravitating n-body simulation
- If the rotation curve is well described by a Newtonian motion in NFW density halo
- It is general for Braneworld predict different deflection against dark matter



Energy and Pressure of Euv

- Alternative approach obtain theoretical form of dark radiation (Gergely, et al., 2011)
- Dark radiation view as energy U and pressure P of "some" fluid
- Guess equation of state based on Schwarzschild case in 2+1+1 decomposition

$$P = (a-2)U + \frac{B}{r^2}$$

LSB metric and Visible mass matching

$$-g_{tt} \approx 1 - \frac{2GM}{c^2 r} + \frac{2\gamma}{1-\alpha} \left(\frac{r}{r_c}\right)^{\alpha-1}$$
$$g_{rr} \approx \left\{1 - \frac{2G(M+M_b)}{c^2 r} + \gamma \left[\left(\frac{r}{r_c}\right)^{\alpha-1} - 1\right]\right\}^{-1}$$

 $\square \alpha$ and γ depend on a and B

Mb is a degenerate parameter in rotation curve

□ rc is obtained from matching Baryonic mass

Contribution to deflection angle

□ The numerical integration of δ BW



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Solving the geometry of the bulk

Key Question: Do we need a bulk blackhole to support the picture?

 ${}^{(5)}ds^2 = -M(r,y)^2 dt^2 + N(r,y)^2 dr^2 + Q(r,y)^2 d\Omega^2 + dy^2$

□ This metric gives dark radiation

$$\begin{split} E^t_t &= \frac{N_{,y,y}}{N} + \frac{2Q_{,y,y}}{Q} - \frac{\Lambda_5}{2} \\ E^r_r &= -\frac{N_{,y,y}}{N} + \frac{\Lambda_5}{6} \\ E^\theta_\theta &= E^\phi_\phi = -\frac{Q_{,y,y}}{Q} + \frac{\Lambda_5}{6} \end{split}$$



- The phenomenological brane metric define boundary condition
- 5D Einstein equation form a system of PDE with Neumann boundary condition, and... K. C. Wong, YITP Kyoto, 3/3/2012

