

CDCCC法による3体・4体分解反応の 系統的解析

**Systematic Analyses of Three- and Four-Body Breakup
Reactions in the CDCC Method**

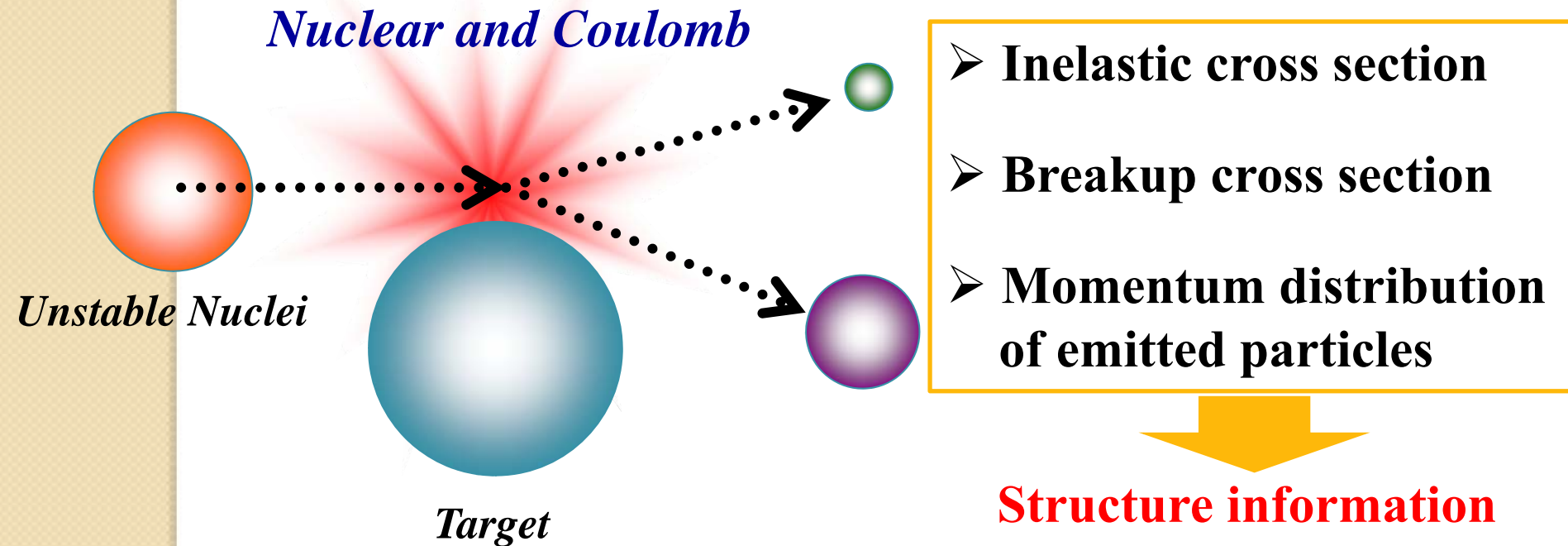
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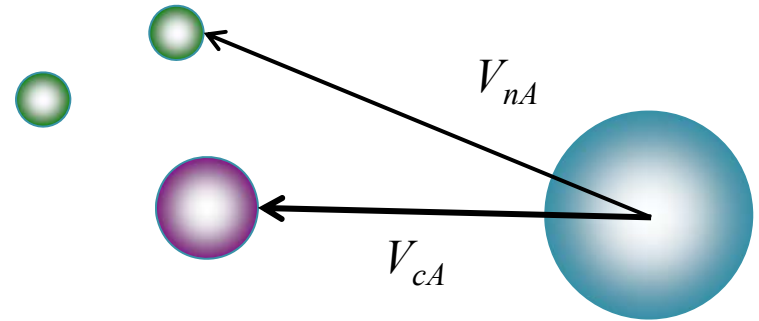
Background

- The unstable nuclear structure can be efficiently investigated via **the breakup reactions**.



- **An accurate method** is needed to analyze breakup reactions **systematically**.

Introduction



□ Systematic analyses

- Reactions for various nuclei
 - Projectile (**light neutron rich nuclei**)
 - ✓ one-nucleon halo → **Three-body breakup**
 - ✓ two-nucleon halo → **Four-body breakup**
 - **Continuum-Discretized Coupled-Channel**
 - Target → proton, light and heavy nuclei
 - Wide range of incident energy
- Potentials for nucleon + target, core + target (Subsystem)
 - nucleus target
 - **Microscopic Optical Potential (K. Minomo)**
 - nucleon target (or neutron induced)
 - **Effective interaction (JLM)**

Continuum-Discretized Coupled-Channels

□ The Continuum-Discretized Coupled-Channels method (CDCC)

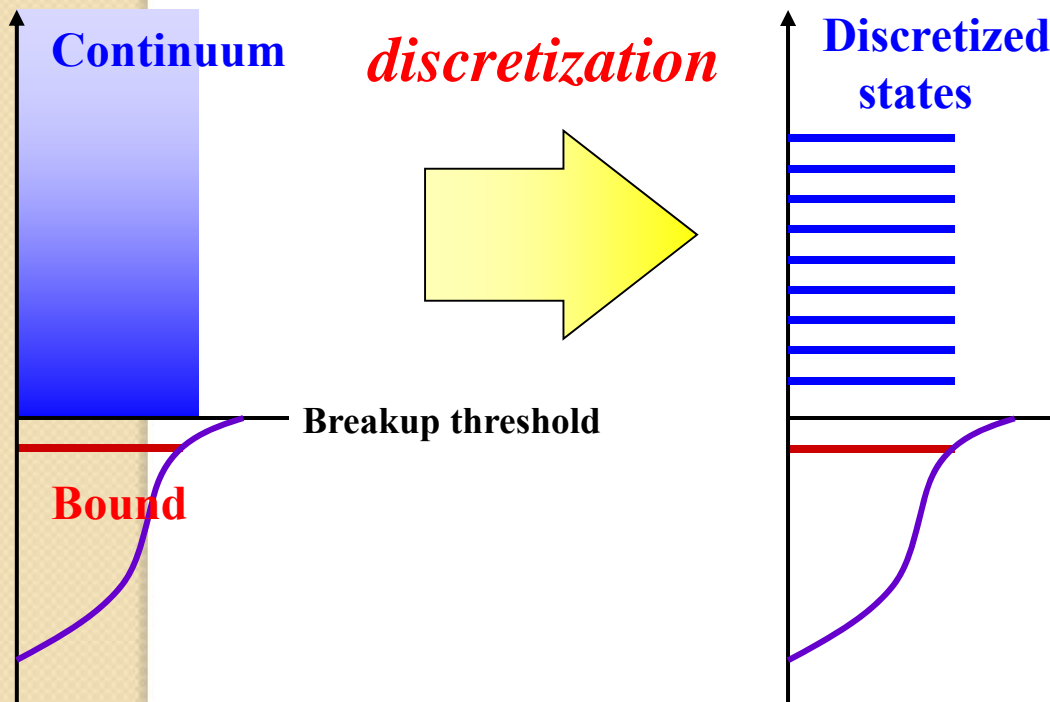
- Developed by Kyushu group about 20 years ago

M.Kamimura, M.Yahiro, Y.Iseri, Y.Sakuragi, H.Kameyama and M.Kawai, PTP Suppl. 89, 1 (1986)

- Fully-quantum mechanical method
- Successful for analyses of **nuclear and Coulomb** breakup reactions
- Extended to describing **four-body reaction system**

T.M., E. Hiyama, K. Ogata, Y. Iseri, M. Kamimura, S. Chiba, M. Yahiro, PRC 70, 061601 (2004).

□ Essence of CDCC



- Breakup continuum states of the projectile are described by a finite number of **discretized states**

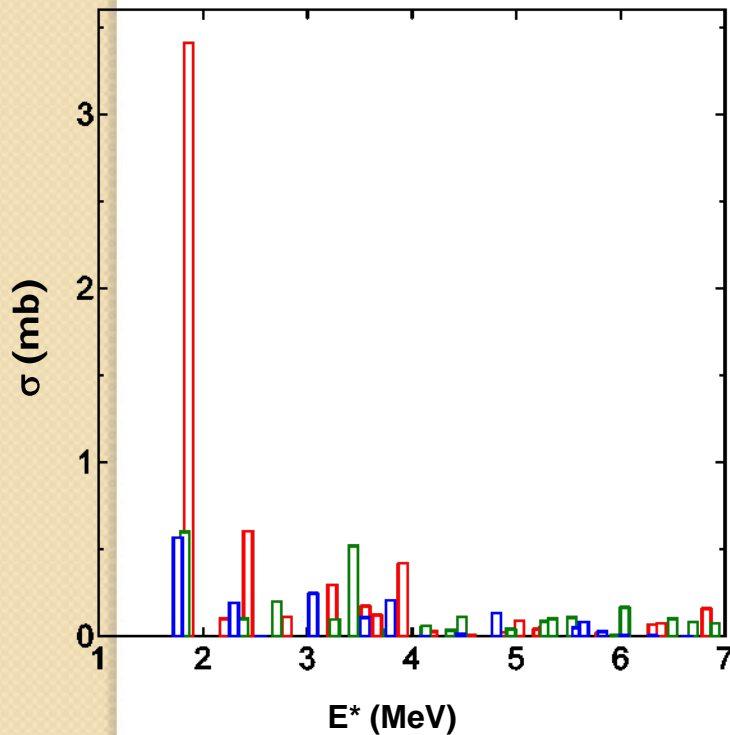
- A set of eigenstates forms a **complete set within a finite model space** that is important for breakup processes

Breakup Cross Section

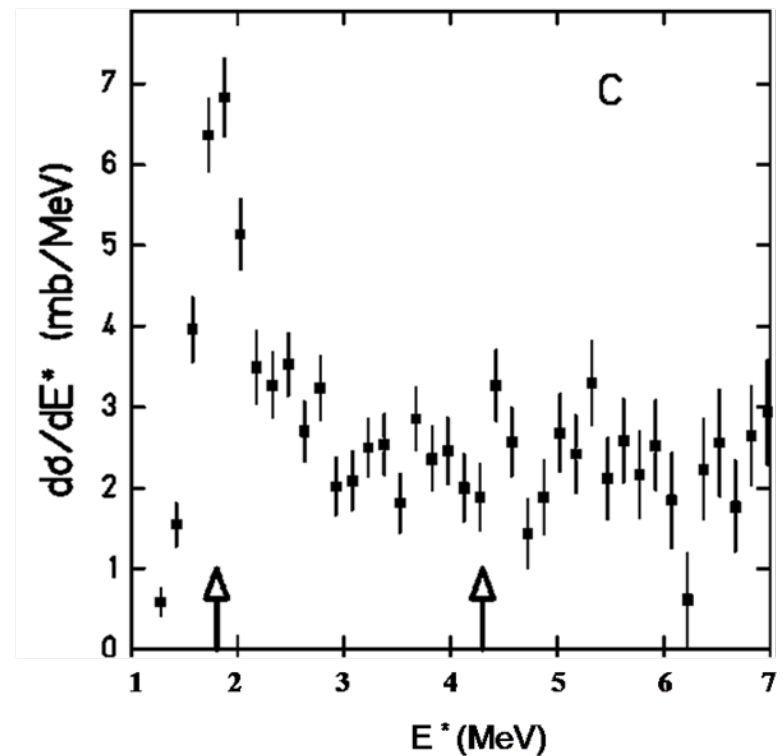
Breakup cross sections calculated by CDCC are **discrete** in the internal energy of the projectile.

${}^6\text{He}+{}^{12}\text{C}$ scattering at 240 MeV/nucl.

4-body CDCC calc.



PRC59, 1252(1999), T. Aumann *et al.*



*How to calculate the **continuum** breakup cross section*

New Smoothing Procedure with *CSM*

T.M., K. Kato, and M. Yahiro, PRC82, 051602 (2010).

$$\frac{d\sigma}{dE} = \int T^\dagger(E')T(E')\delta(E - E')dE' = \frac{1}{\pi}\text{Im}\mathcal{R}(E)$$

$$T(E) = \psi^{(-)}(E, \xi) \chi_C^{(-)}(\mathbf{R}) |V| \Psi^{(+)}(\xi, \mathbf{R})$$

Response function

$$\mathcal{R}(E) = \int d\xi d\xi' \langle \Psi^{(+)}(\xi, \mathbf{R}) | V^* | \chi_C^{(-)}(\mathbf{R}) \rangle_{\mathbf{R}} \mathcal{G}^{(-)}(E, \xi, \xi') \langle \chi_C^{(-)}(\mathbf{R}) | V | \Psi^{(+)}(\xi, \mathbf{R}) \rangle_{\mathbf{R}}$$

Final state of the projectile

Green's function with Complex-Scaling Method (CDCS Green's function)

$$\mathcal{G}^{(-)}(E, \xi, \xi') = U^{-\theta} \frac{1}{E - H^\theta - i\epsilon} U^\theta \approx \sum_\nu U^{-\theta} \frac{|\Phi_\nu^\theta\rangle \langle \tilde{\Phi}_\nu^\theta|}{E - E_\nu^\theta} U^\theta$$

$$\mathcal{G}^{(-)}(E, \xi, \xi') \approx \sum_\nu \sum_{i,j} |\Phi_i\rangle \frac{\langle \Phi_i | U^{-\theta} | \Phi_\nu^\theta \rangle \langle \tilde{\Phi}_\nu^\theta | U^\theta | \Phi_j \rangle}{E - E_\nu^\theta} \langle \Phi_j |$$

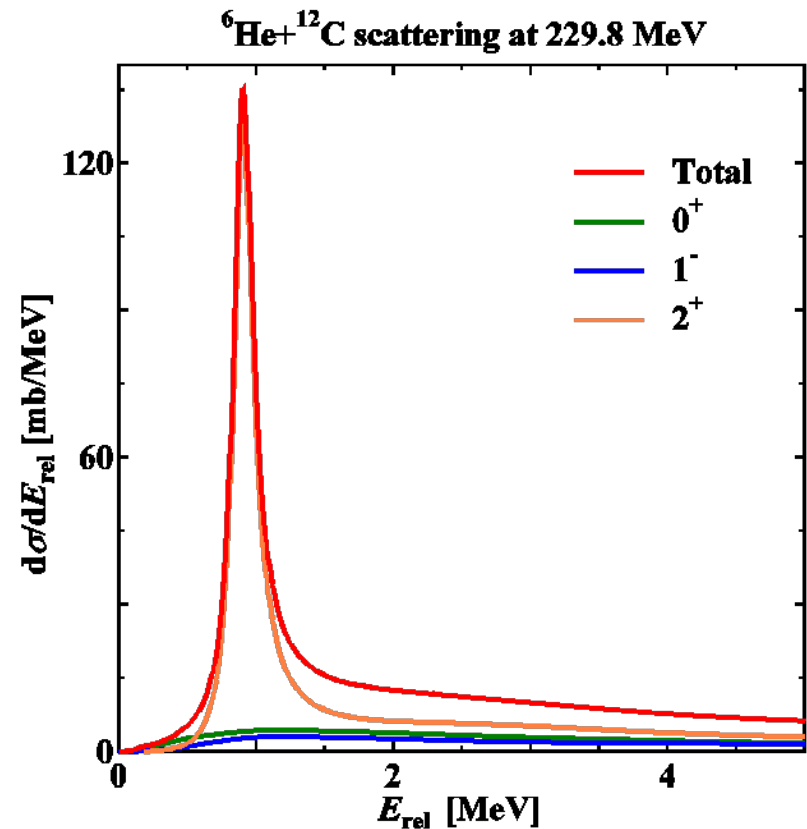
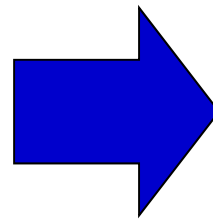
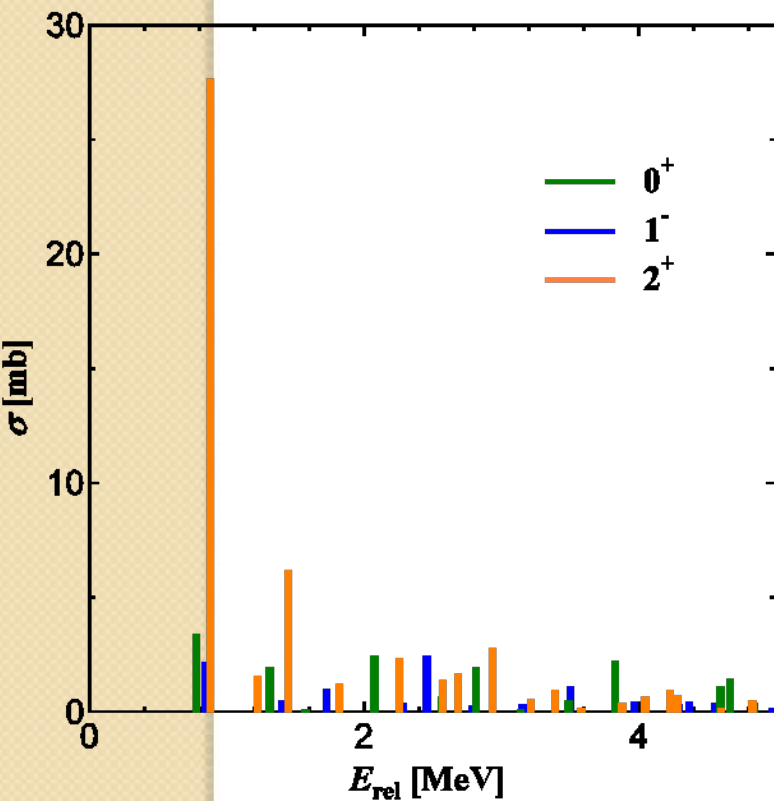
$$\mathcal{R}(E) = \sum_\nu \sum_{i,j} \langle \Psi^{(+)} | V^* | \chi_C^{(-)} \Phi_i \rangle \frac{\langle \Phi_i | U^{-\theta} | \Phi_\nu^\theta \rangle \langle \tilde{\Phi}_\nu^\theta | U^\theta | \Phi_j \rangle}{E - E_\nu^\theta} \langle \Phi_j \chi_C^{(-)} | V | \Psi^{(+)} \rangle$$

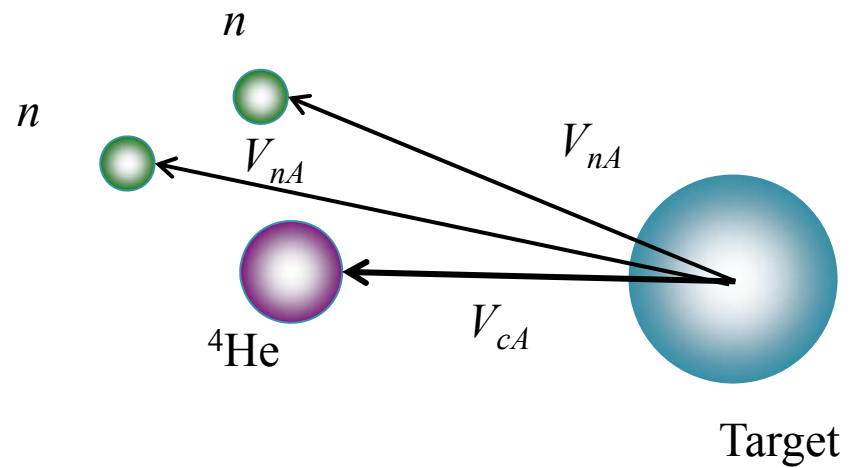
T-matrix calculated by CDCC

Differential Breakup Cross Section

New description of differential breakup cross section

$$\frac{d\sigma}{dE} = \frac{1}{\pi} \text{Im} \sum_{\nu} \sum_{i,j} T_i^{\text{CDCC}\dagger} \frac{\langle \Phi_i | U^{-\theta} | \Phi_{\nu}^{\theta} \rangle \langle \tilde{\Phi}_{\nu}^{\theta} | U^{\theta} | \Phi_j \rangle}{E - E_{\nu}^{\theta}} T_j^{\text{CDCC}}$$





GEM+MOP+CDCC

Gaussian Expansion Method

E. Hiyama, Y. Kino, M. Kamimura, Prog. Part Nucl. Phys. 51, 223.

Microscopic Optical Potential

K. Minomo.

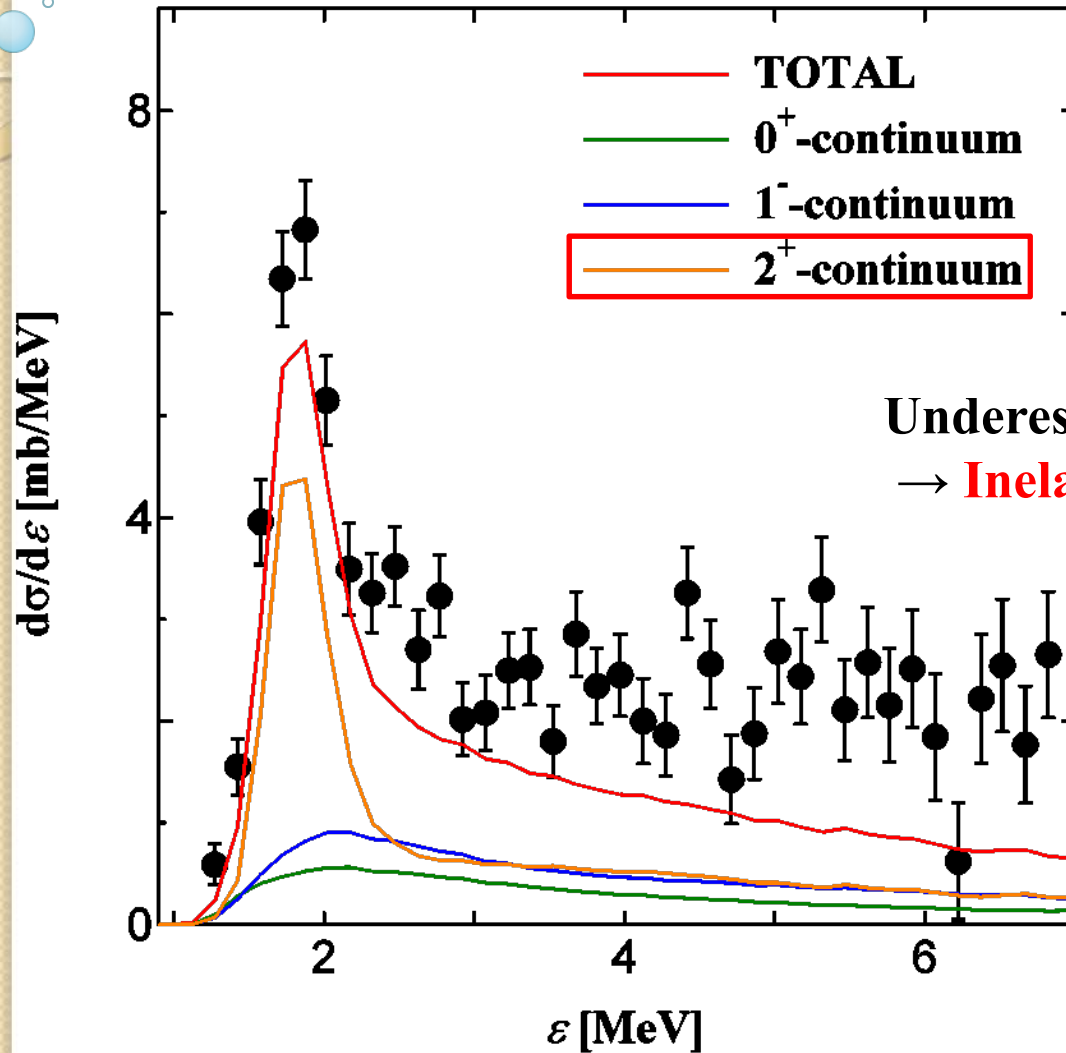
${}^6\text{He} + {}^{12}\text{C}$ and ${}^{208}\text{Pb}$ scattering at 240 MeV/A

Nucleus targets

n - ${}^{12}\text{C}$ and ${}^4\text{He} - {}^{12}\text{C}$ potentials

n - ${}^{208}\text{Pb}$ and ${}^4\text{He} - {}^{208}\text{Pb}$ potentials

${}^6\text{He}+{}^{12}\text{C}$ scattering @ 240 MeV/nucleon.



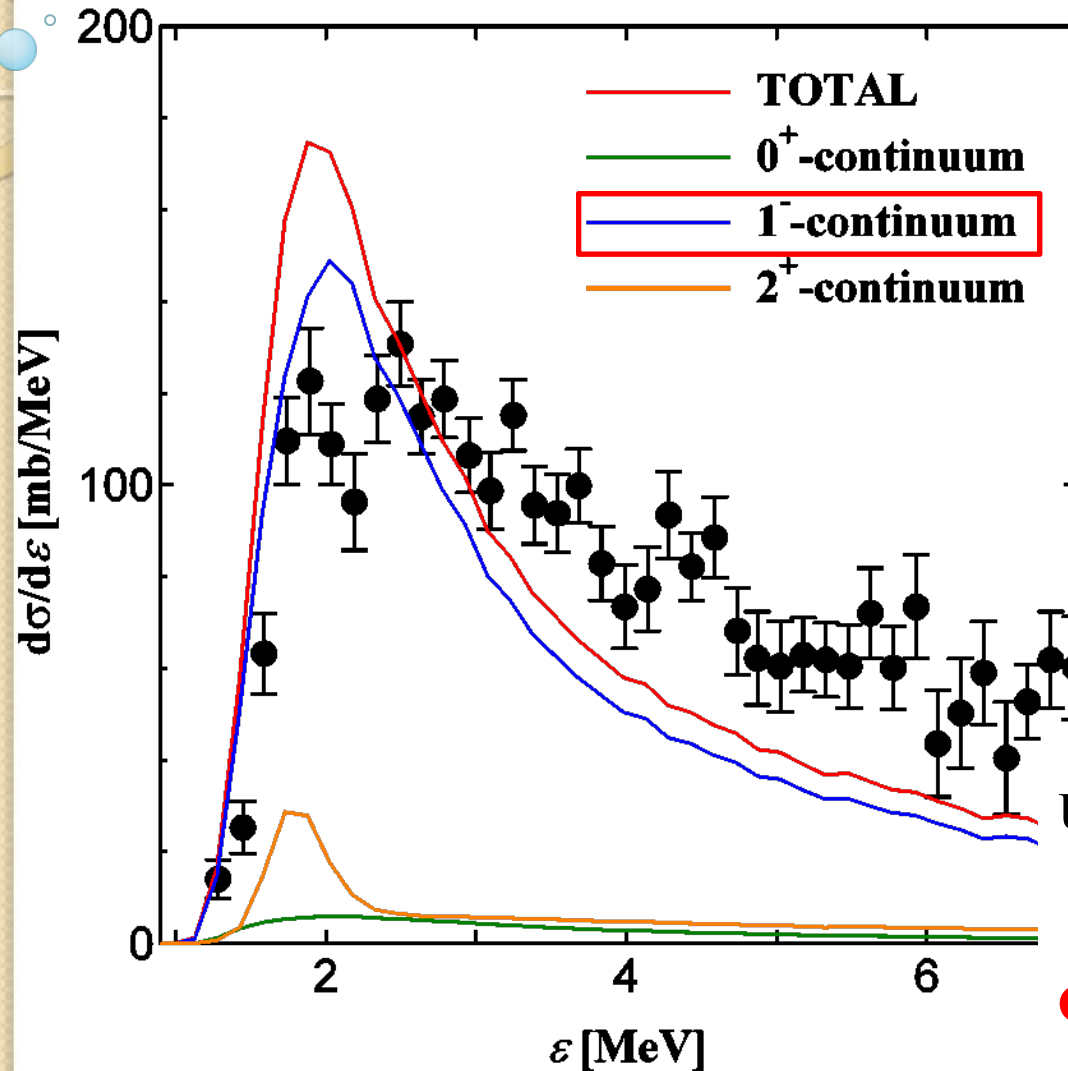
Nuclear Breakup is dominant

Underestimation

→ Inelastic breakup effect ~ 20%

Exp. data from PRC59, 1252 (1999), T. Aumann *et al.*

${}^6\text{He}+{}^{208}\text{Pb}$ scattering @ 240 MeV/nucl.



Coulomb Breakup is dominant

Underestimation

→ Inelastic breakup effect

Overestimation ???

Exp. data from PRC59, 1252 (1999), T. Aumann *et al.*

GEM+JLM+CDCC

Gaussian Expansion Method

E. Hiyama, Y. Kino, M. Kamimura, Prog. Part Nucl. Phys. 51, 223.

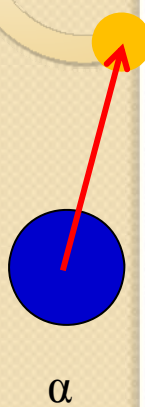
JLM effective interaction

J.-P. Jeukenne, A. Lejeune, and C. Mahaux, Phys. Rev. C16, 80.

Proton Target (neutron induced reaction)



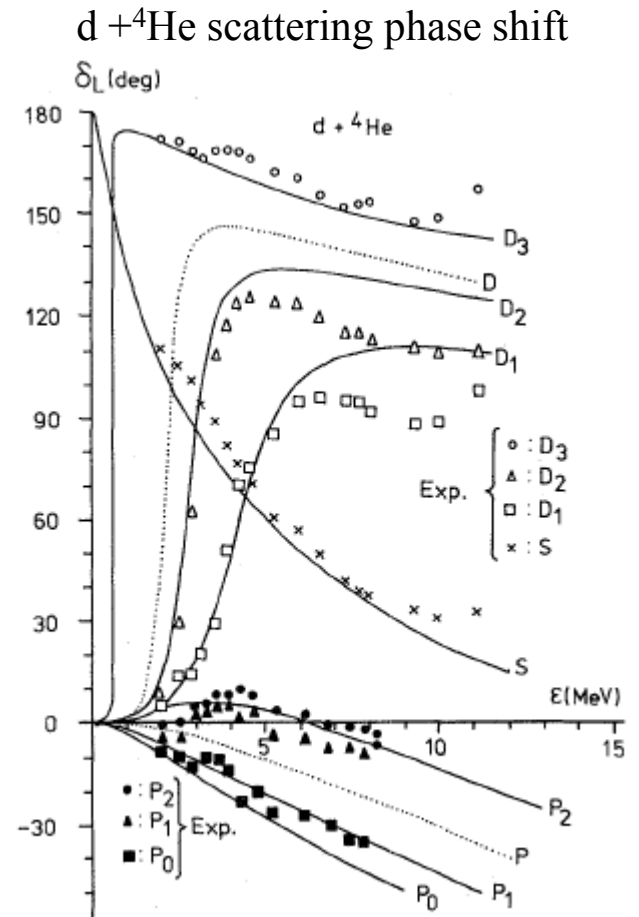
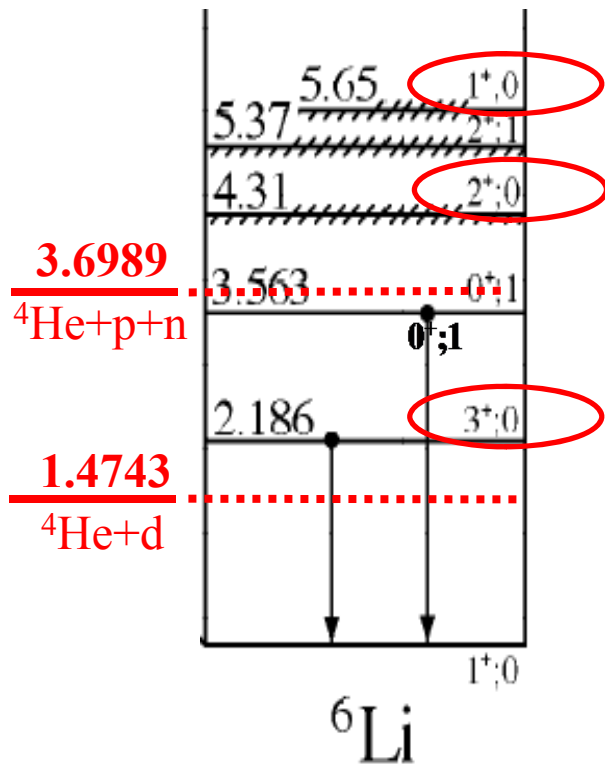
${}^6\text{Li}: d + \alpha$ two-body model



$d, s_d = 1$

r, l

$$V_{d\alpha}(r) = V^C(r) + V^{LS}(r)$$



Y. Sakuragi, M. Yahiro and M. Kamimura, Prog. Theor. Phys. 89, 136 (1986)

JLM effective interaction

JLM interaction (J.-P. Jeukenne, A. Lejeune, and C. Mahaux, Phys. Rev. C16, 80 (1977))

$$v_{j0}(R_{j0}; \rho, E) = \lambda_v V(\rho, E) \exp(-R_{j0}/t_R^2) + i\lambda_w W(\rho, E) \exp(-R_{j0}/t_I^2)$$

In generally, $t_R = t_I = 1.2$, $\lambda_v = 1.0$ and $\lambda_w = 0.8$ (single channel calculation)

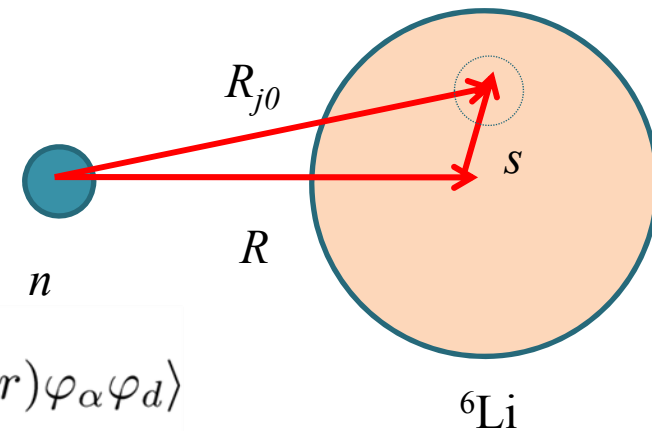
→ λ_w is optimized

Coupling potential

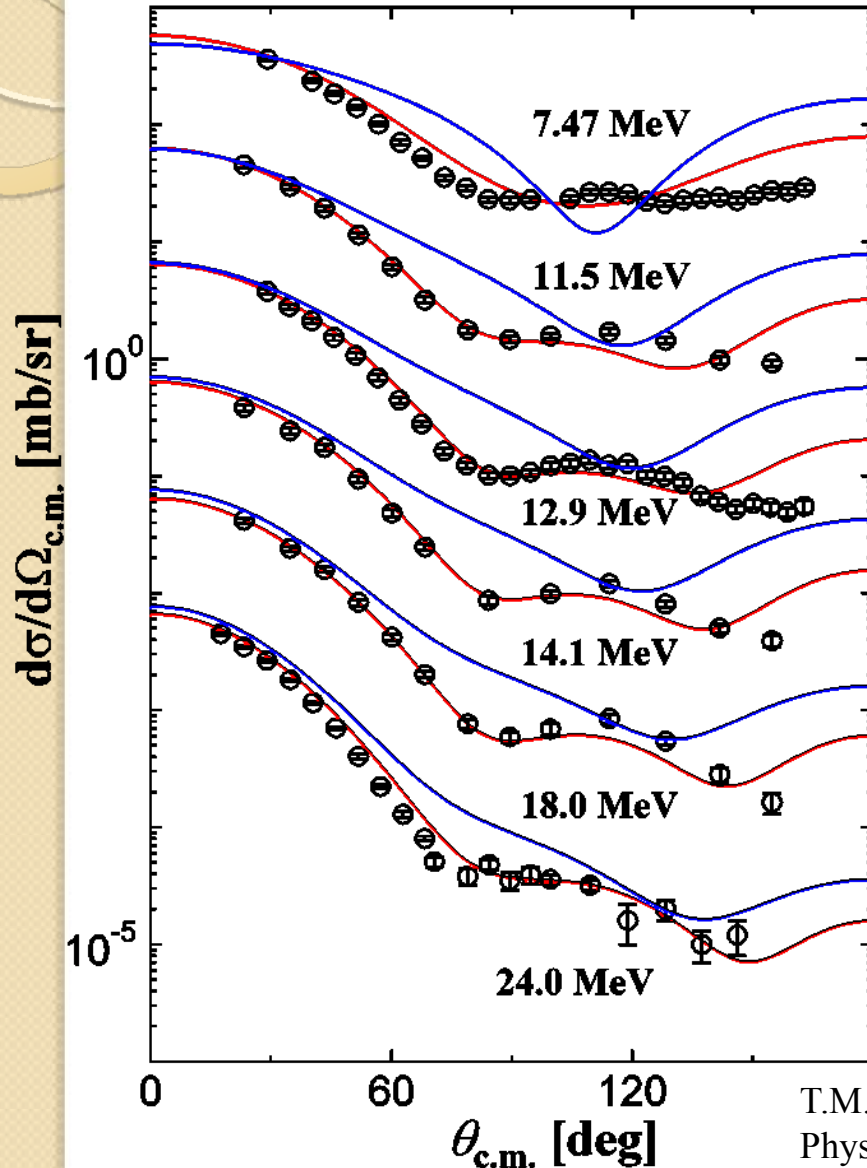
$$V_{\gamma'\gamma}(R) = \int ds \rho_{\gamma'\gamma}(s) v_{j0}(R_{j0}; \rho, E)$$

Transition density of ${}^6\text{Li}$

$$\rho_{\gamma'\gamma}(s) = \langle \phi_{i'l'}^{I'}(r) \varphi_\alpha \varphi_d | \sum_{k=1}^6 \delta(s - r_k) | \phi_{il}^I(r) \varphi_\alpha \varphi_d \rangle$$



Elastic cross section of ${}^6\text{Li}(n, n)$



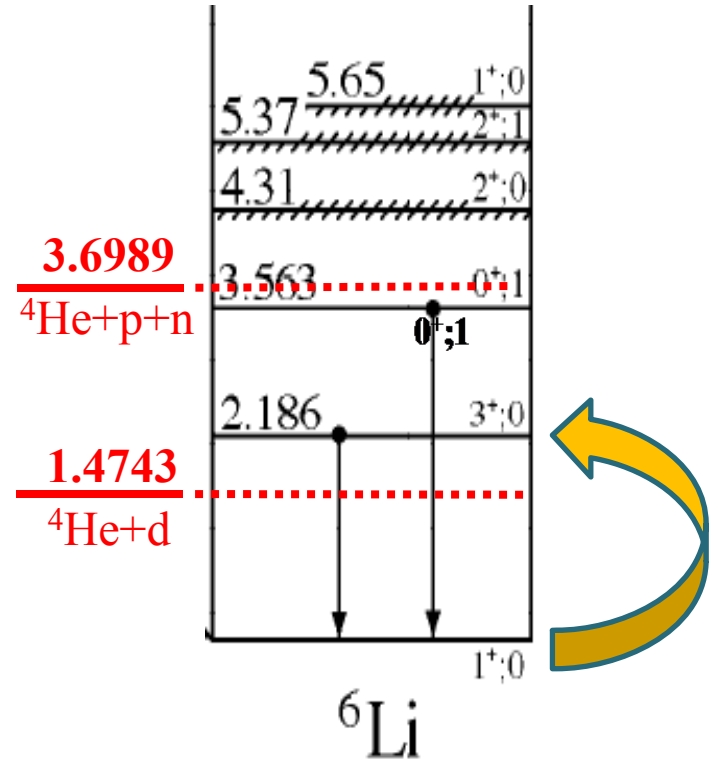
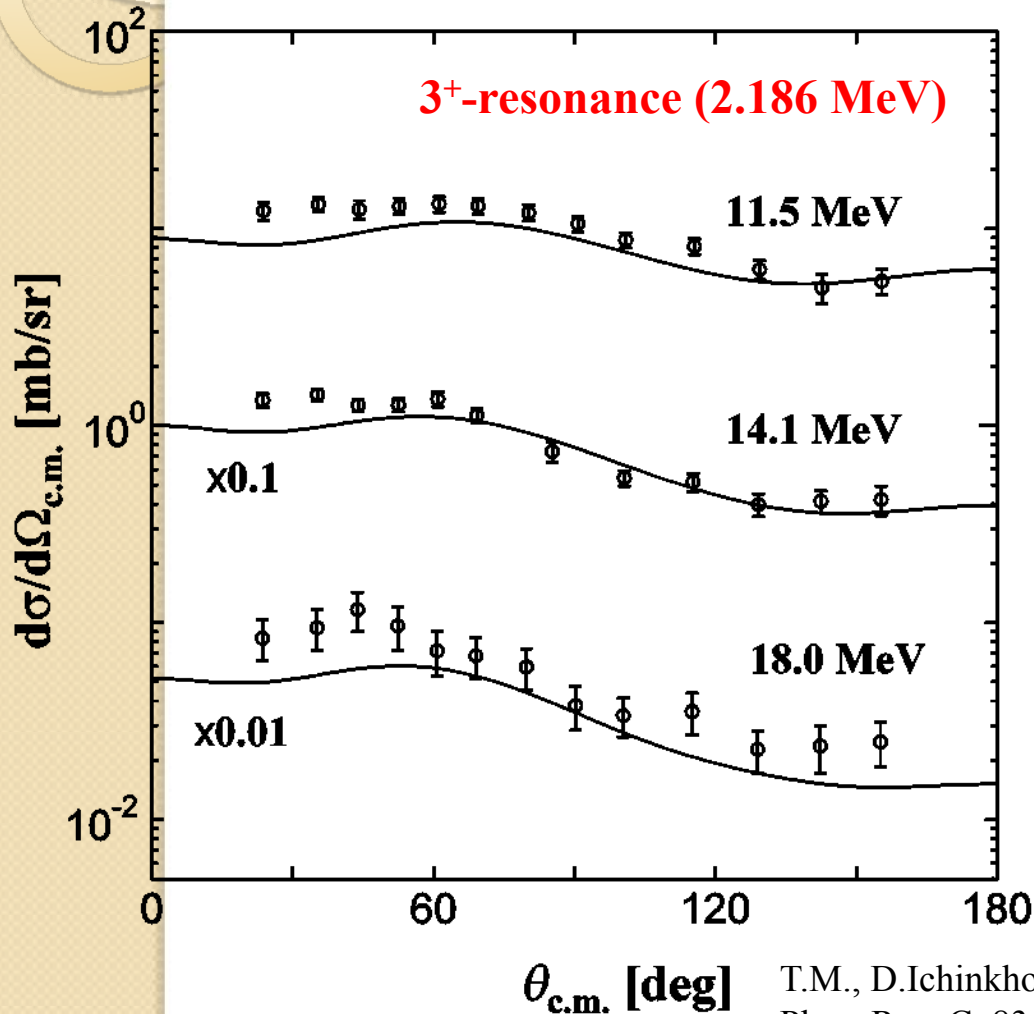
Blue:
Single channel calc.
(without BU effects)

Red: Full CC

➤ The optimized λ_w is **0.1**.

➤ **Breakup effect is significant**

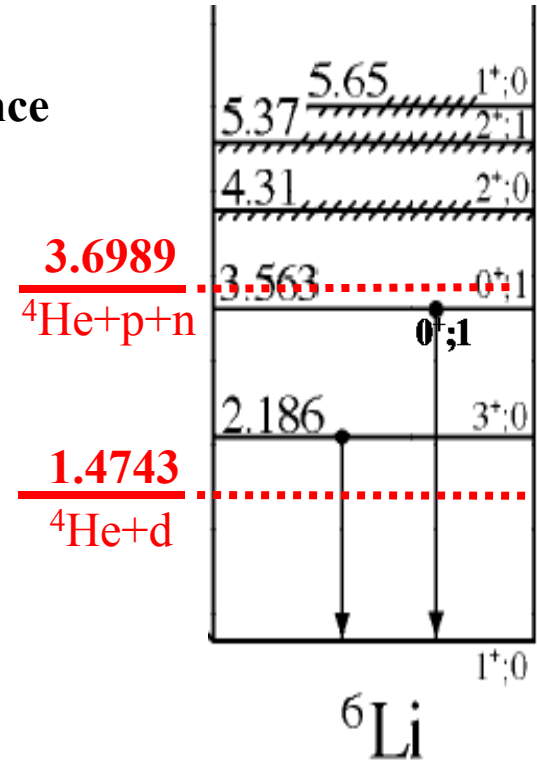
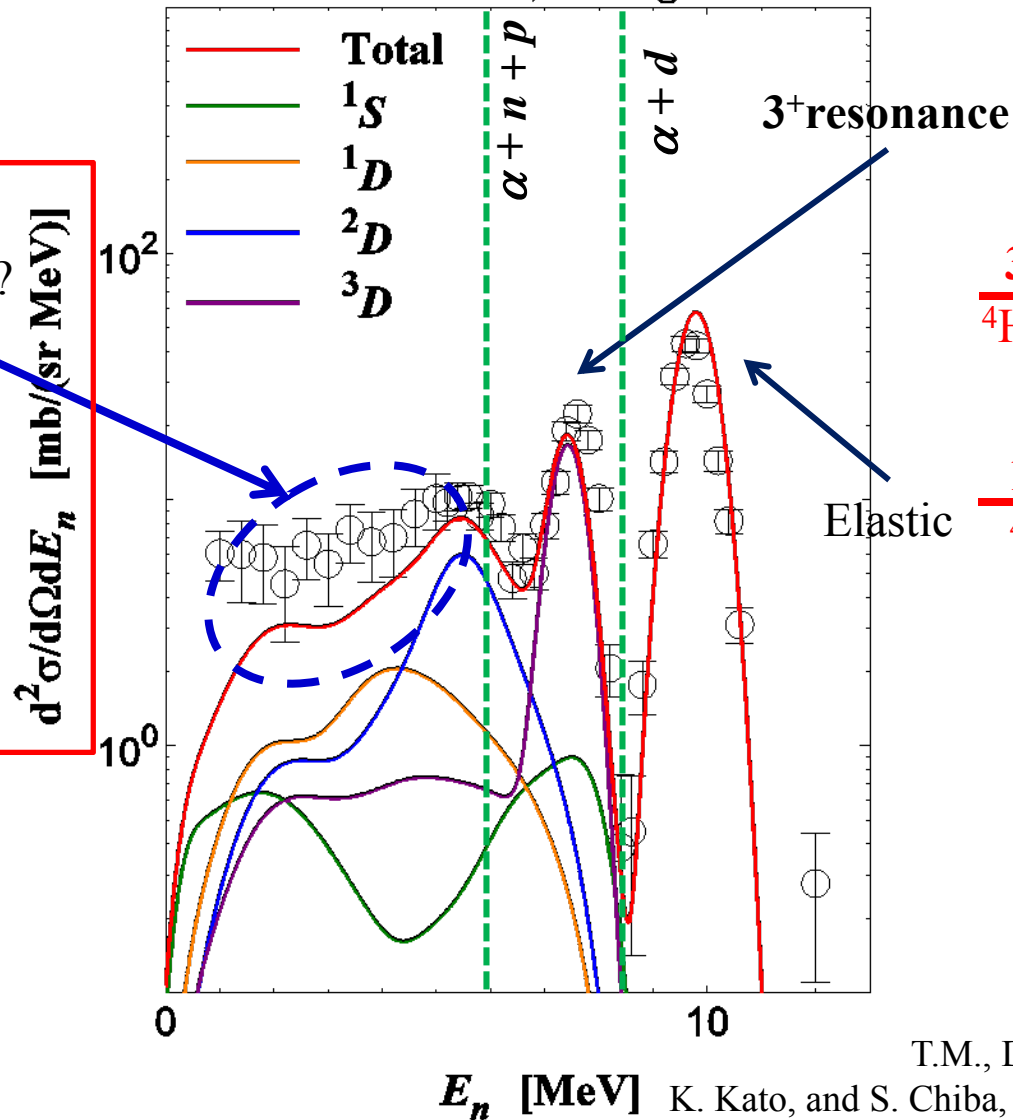
Inelastic cross section



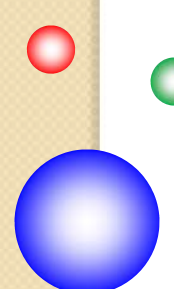
T.M., D. Ichinkhoroloo, Y. Hirabayashi, K. Kato, and S. Chiba,
 Phys. Rev. C. 83. 064611 (2011)

Neutron spectrum of ${}^6\text{Li}(n, n')$

11.5-MeV, 60-deg



$n+p+{}^4\text{He}$ Breakup?

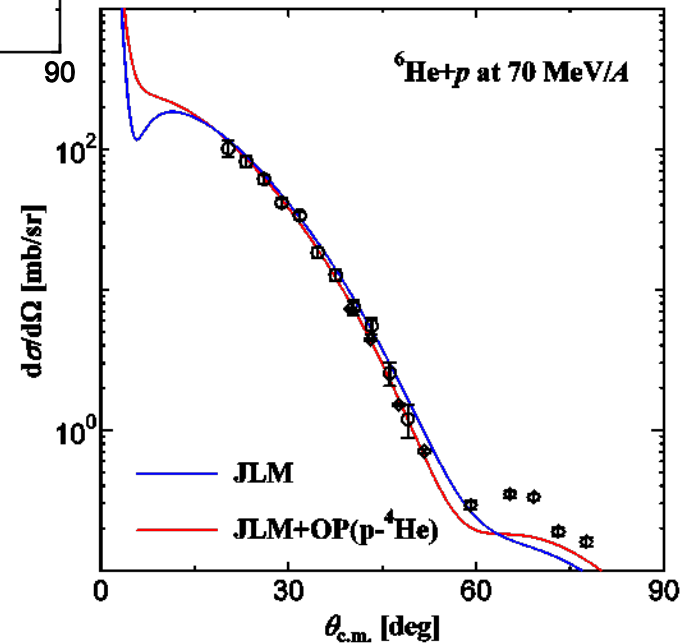
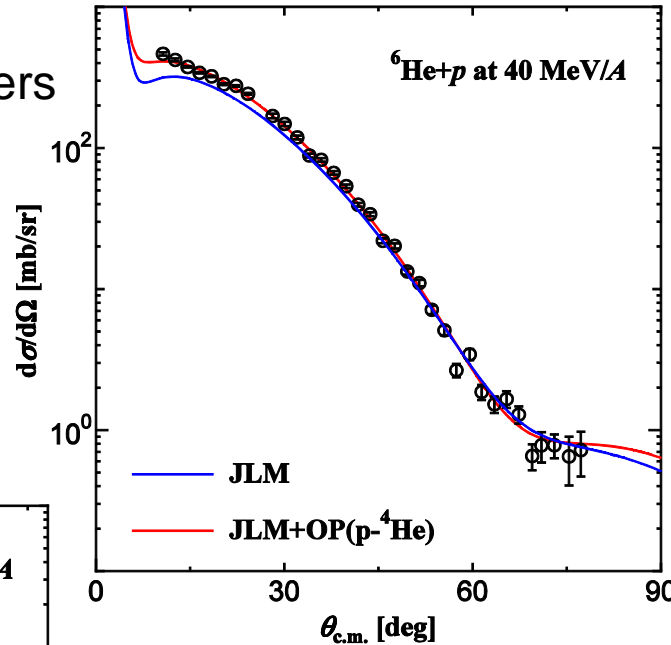
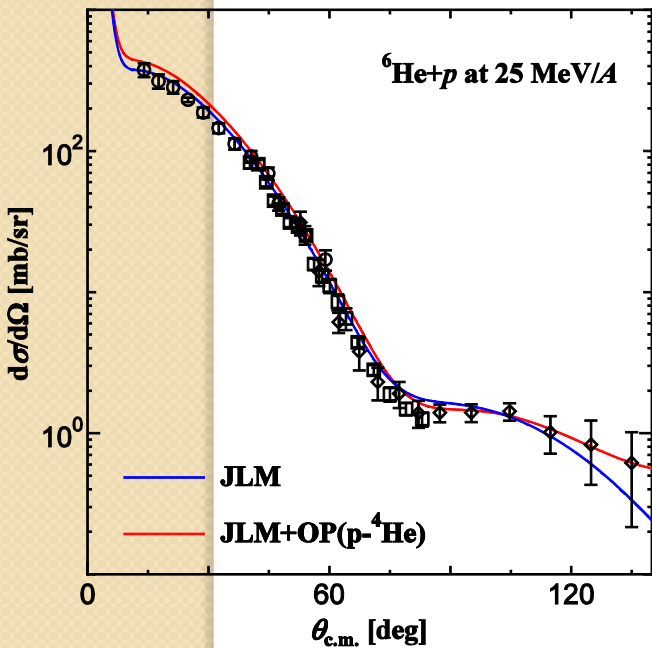


New results

${}^6\text{He} + p$ elastic cross section

OP(p - ${}^4\text{He}$)

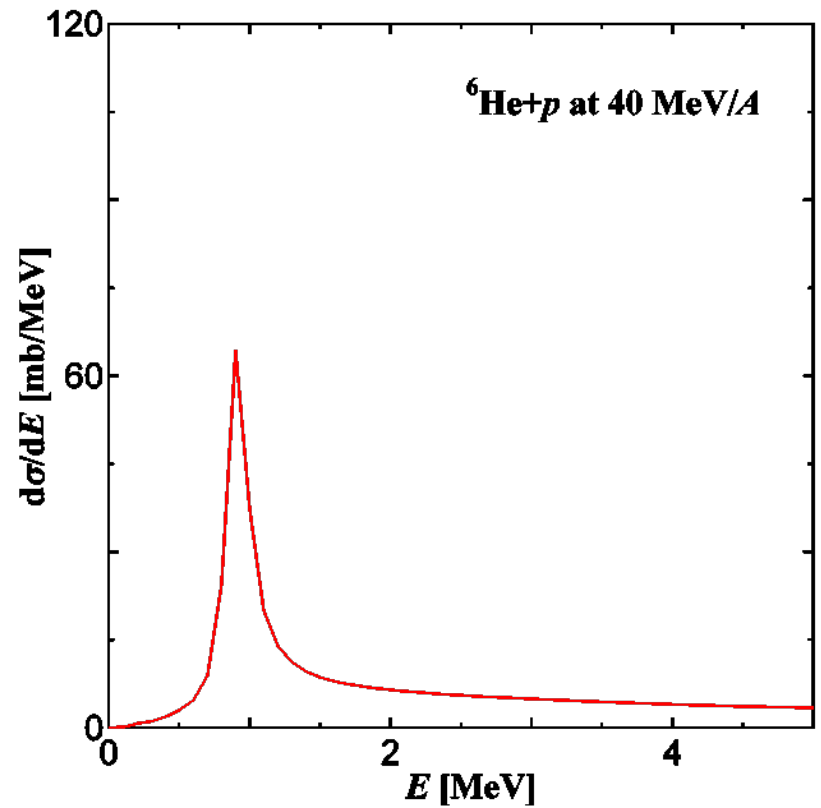
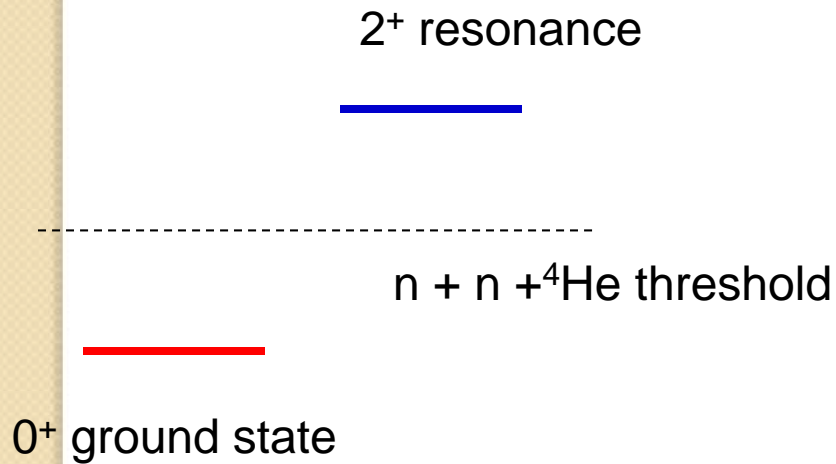
:Optical Potential parameters
taken from **NPA 674, 77**
(2000) D. Gupta *et al.*



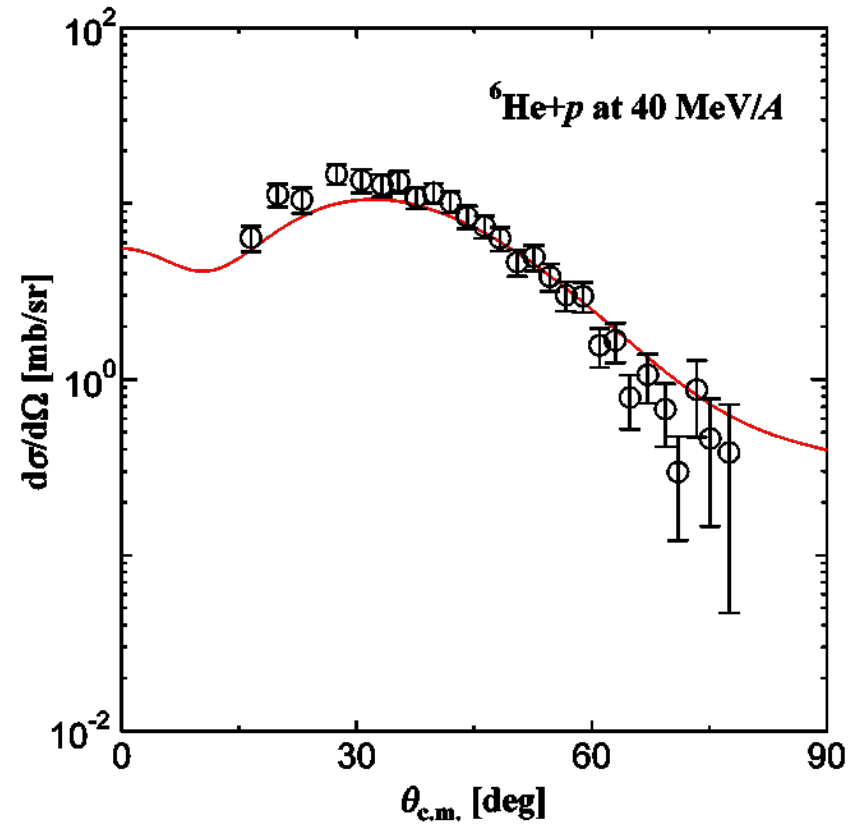
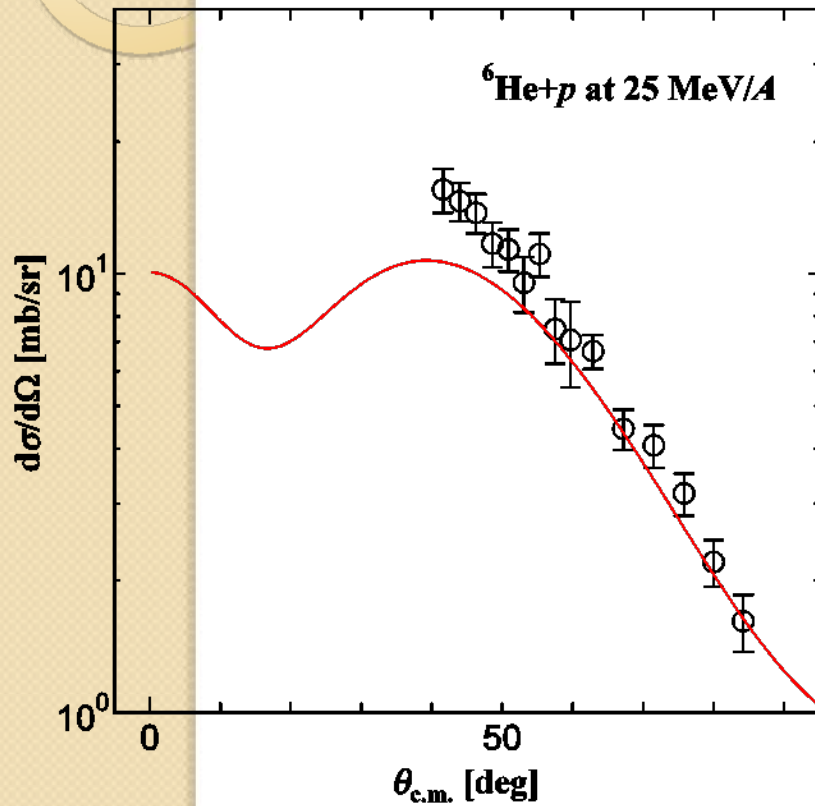
Normalization of JLM
interaction is 0.5 for
the imaginary part.

Inelastic Cross Section to 2^+

Energy spectral of ${}^6\text{He}$



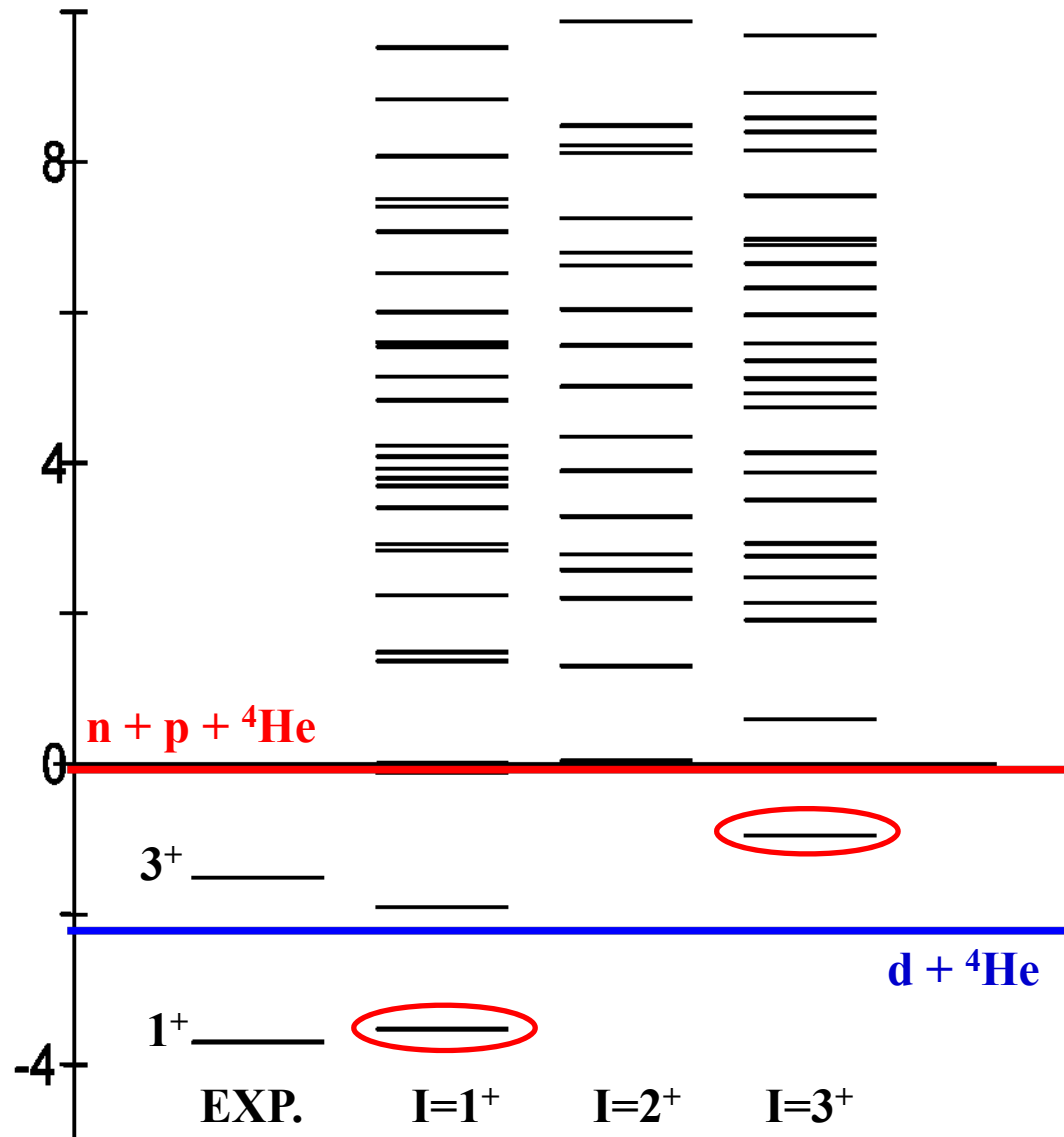
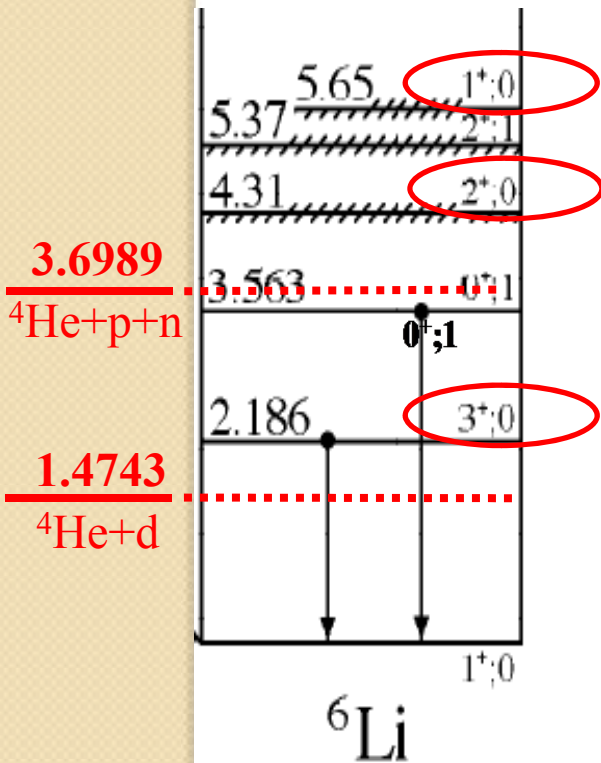
Inelastic Cross Section



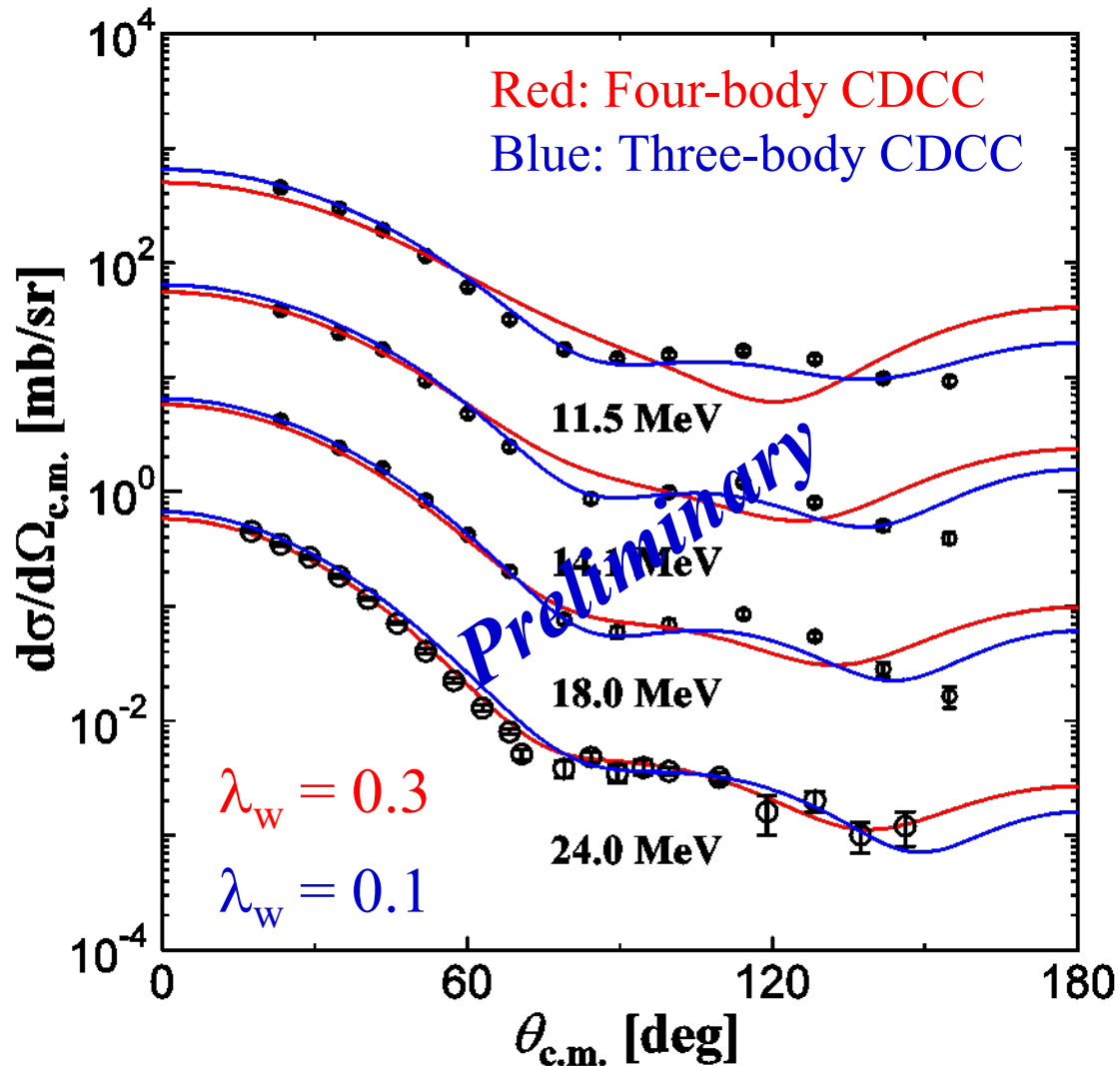
GEM + JLM + Four-body CDCC well reproduces the elastic and inelastic data

${}^6\text{Li}: n + p + \alpha$ three-body model

${}^6\text{He} (T=1) \rightarrow {}^6\text{Li} (T=0)$



Elastic cross section ${}^6\text{Li}(n, n)$



Summary

□ Systematic analyses for three-body and four-body breakup reactions

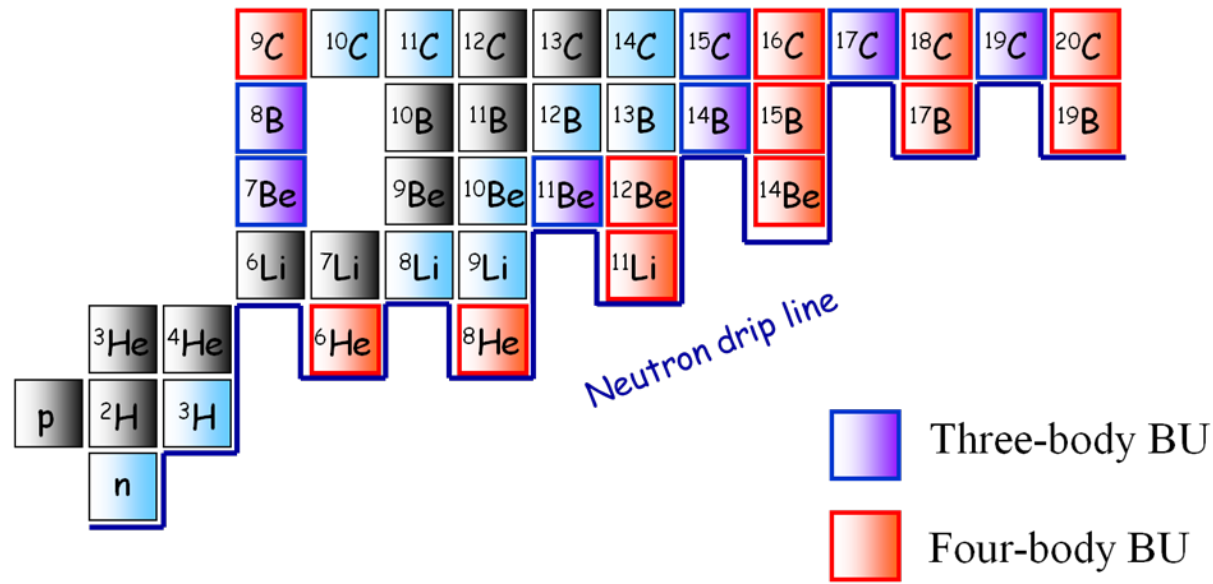
➤ Development of the CDCC method

✓ Four-body breakup system

✓ Calculation of breakup cross section with CSM

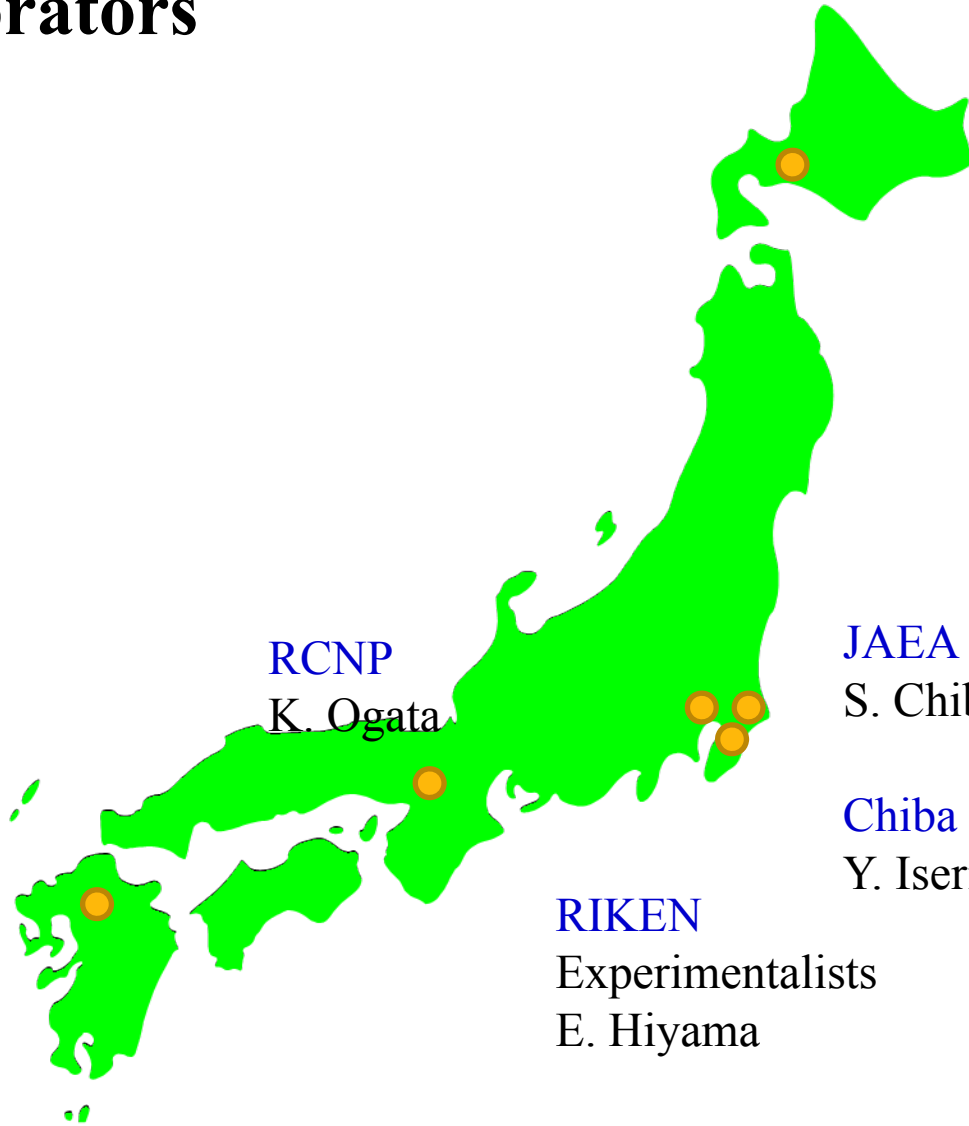
➤ Using microscopic optical potentials for nucleus targets and JLM effective interactions for nucleon targets

□ Future Plan



Collaborators

Kyushu Univ.
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Experimentalists
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D. Ichinkhorloo