

Role of non-collective excitations in fusion reaction and quasi-elastic scattering around the Coulomb barrier

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- **Introduction**
- **Coupled-channels method**
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- **Summary**

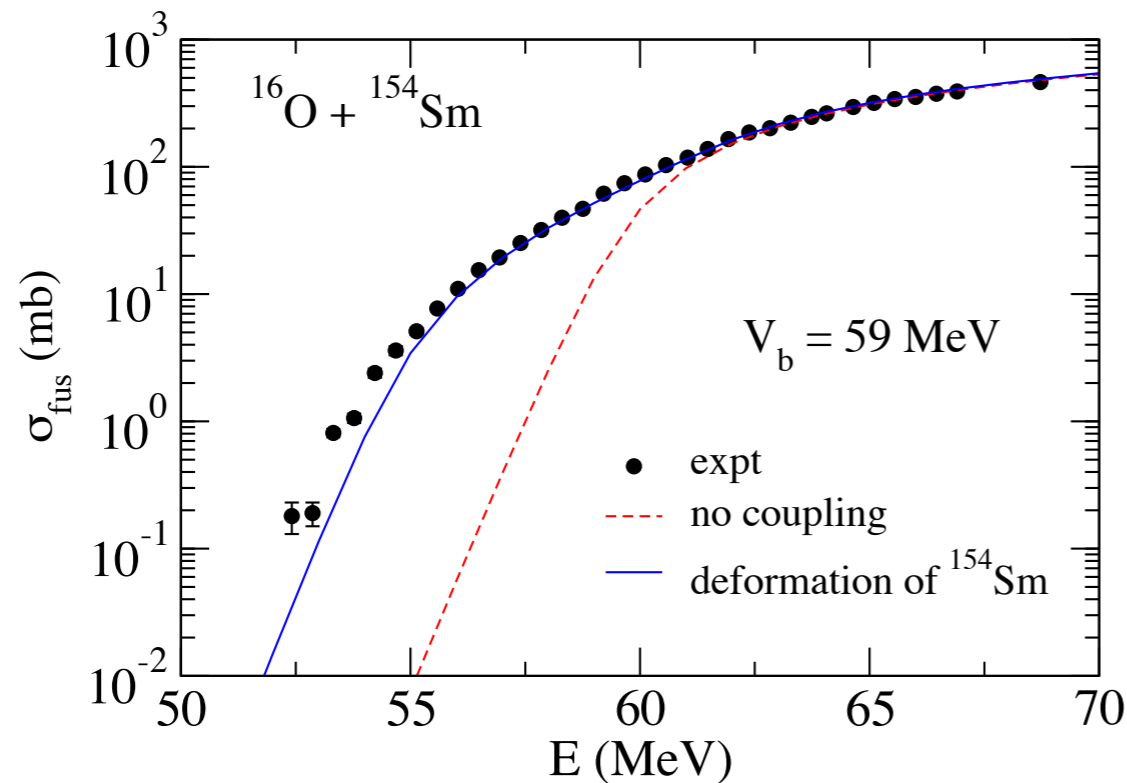
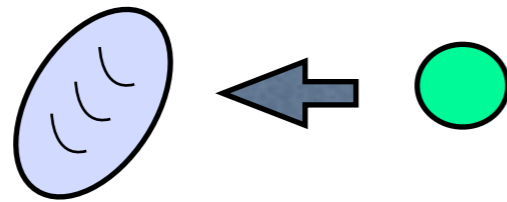
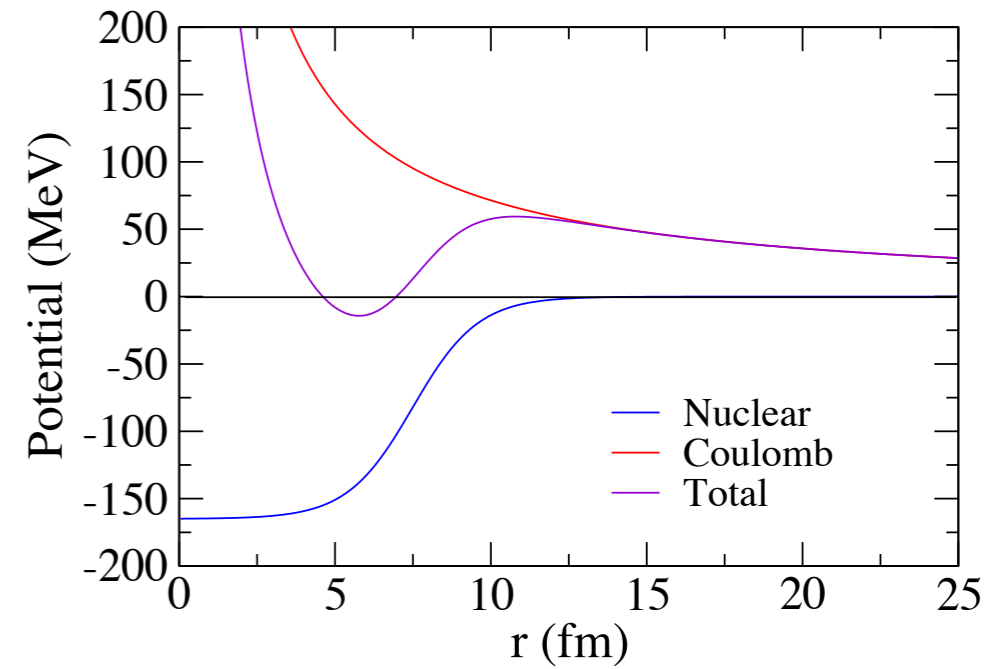
Introduction

Low energy heavy-ion reactions

- heavy-ion reactions near the Coulomb barrier

➔ **channel coupling effect**

e.g. large enhancement of subbarrier fusion cross sections

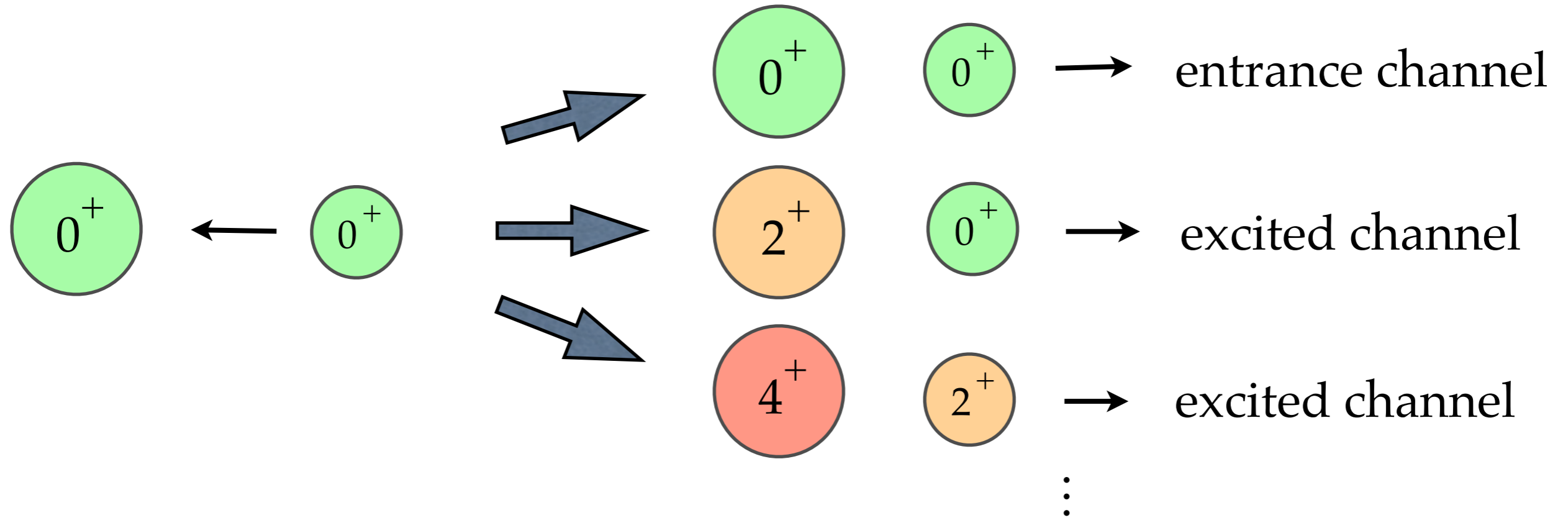


^{154}Sm

12^+	1825.7 keV
10^+	1332.8 keV
8^+	902.65 keV
6^+	543.74 keV
4^+	266.79 keV
2^+	81.976 keV
0^+	

GD band

- Coupled-Channels method



Quantum theory that takes into account excitations of colliding nuclei during the collision

➔ conventionally, **collective excitations(vibration and / or rotation)** are taken into account



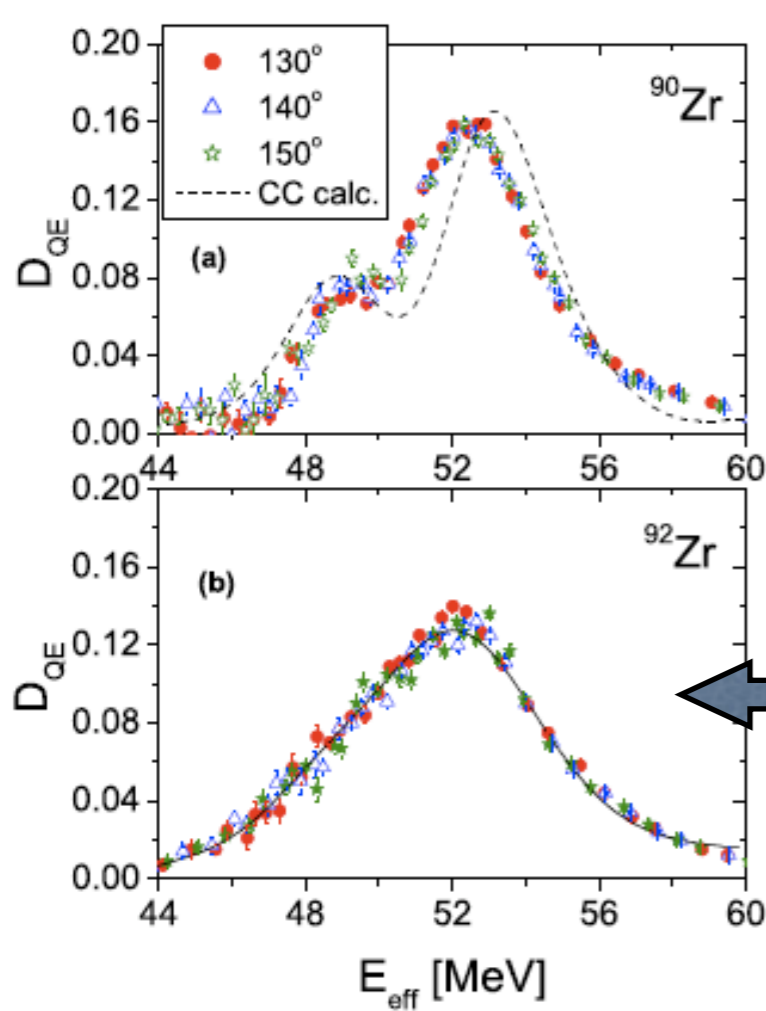
successfully accounted for heavy-ion fusion reactions and quasi-elastic scattering(elastic + inelastic + transfer)

$^{20}\text{Ne} + ^{90,92}\text{Zr}$ quasi-elastic scattering

experimental barrier distributions show **different behavior** (much more smeared distribution for $^{20}\text{Ne} + ^{92}\text{Zr}$ system)

on the other hand,

coupled-channels calculation with collective excitations results in **similar barrier distributions**



$$D_{\text{qel}} = -\frac{d}{dE} \left(\frac{\sigma_{\text{qel}}(E, \pi)}{\sigma_R(E, \pi)} \right)$$

E.Piasecki *et al.*,
PRC 80, 054613
(2009)

much more smeared

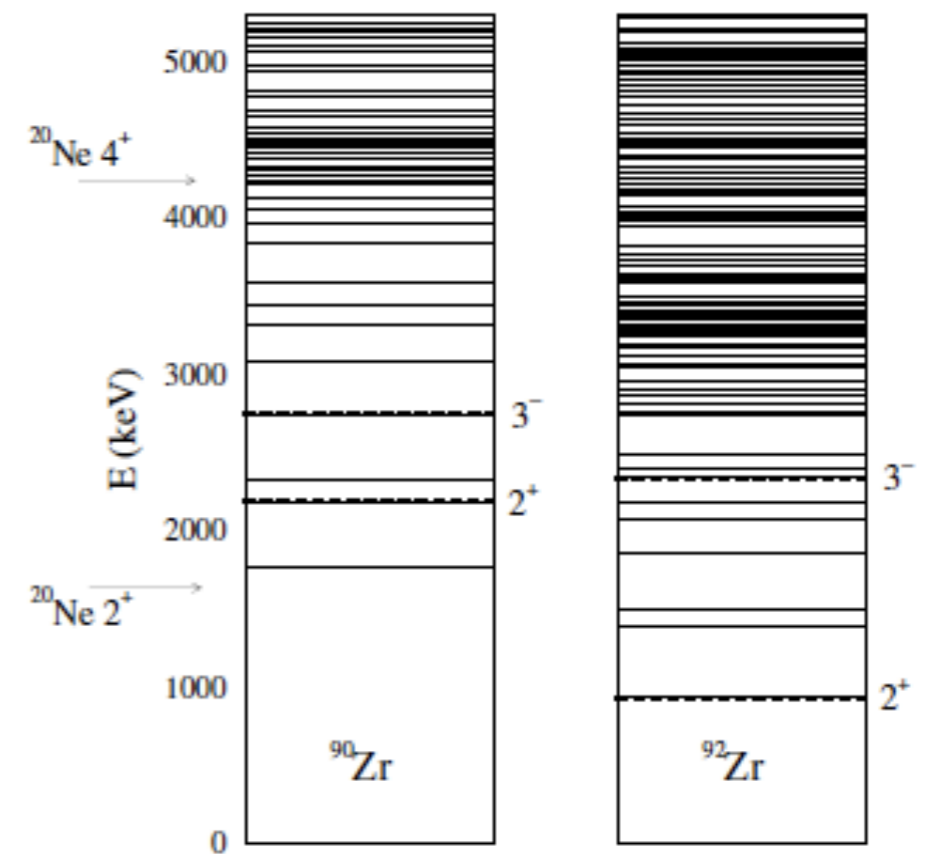
Energy spectrum for Zr isotopes

different level density

$$\begin{cases} ^{90}\text{Zr} : N = 50 \text{ (shell closure)} \\ ^{92}\text{Zr} : N = 50 + 2 \end{cases}$$



many non-collective excitations in ^{92}Zr

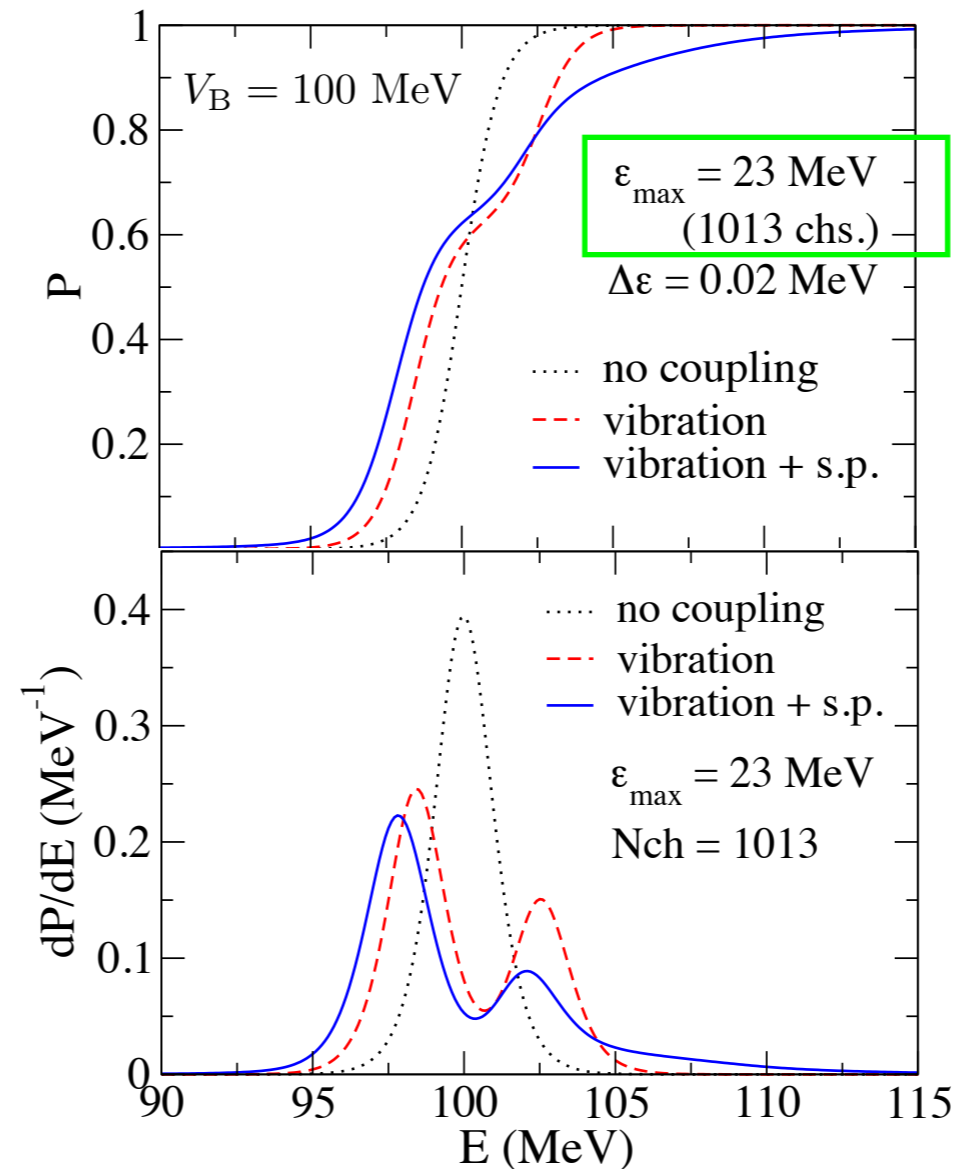
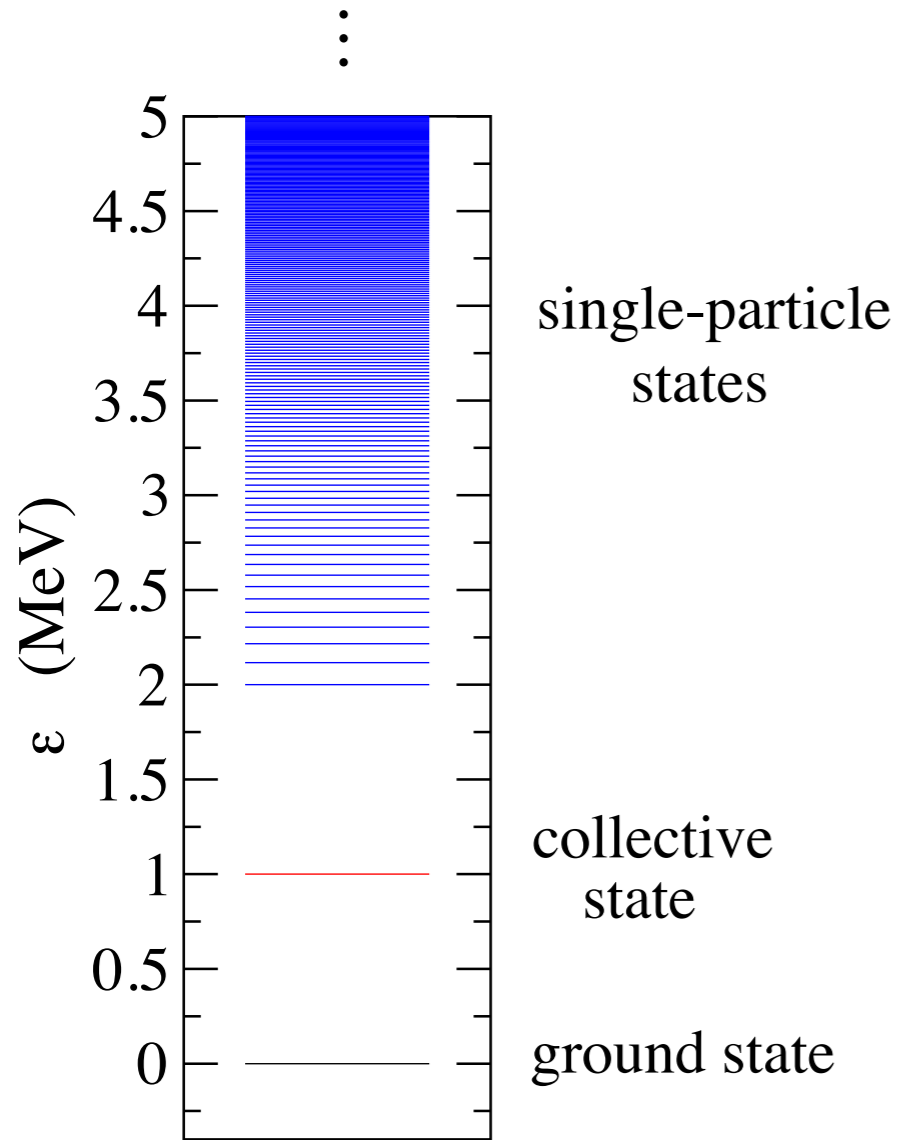
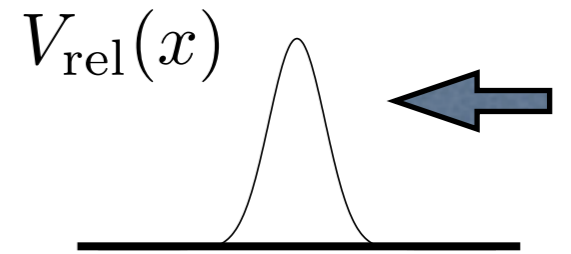


role of non-collective excitations ?

(not explicitly taken into account in the c.c. calculation)

Non-collective excitations in one-dimensional model

description of non-collective excitations based on random matrix theory



large dimensional coupled-channels calculation

S.Y., K.Hagino and N.Rowley, PRC82('10)023606

* non-coll. excitations \Rightarrow suppress the penetrability above the barrier
 \Rightarrow smear the barrier distribution

Coupled-channels method

2. Coupled-channels method

coupled-channels equations(in the isocentrifugal approximation):

$$\left[\frac{d^2}{dr^2} + k^2 - \frac{J(J+1)}{r^2} - \frac{2\mu}{\hbar^2} (V_C(r) + V_N(r) + \epsilon_n) \right] u_n^J(r) = \sum_m \frac{2\mu}{\hbar^2} V_{nm}(r) u_m^J(r)$$

$$k = \sqrt{\frac{2\mu E}{\hbar^2}}$$

$$V_N(r) = V_N^{(0)}(r) + iW_N(r) = -\frac{V_0}{1 + \exp((r - R_N)/a)} - i\frac{W_0}{1 + \exp((r - R_W)/a_W)}$$

ϵ_n : excitation energy of the n -th channel

boundary conditions

$u_n^J(r)$: regular at the origin

$$k_n = \sqrt{\frac{2\mu(E - \epsilon_n)}{\hbar^2}}$$

$$u_n^J(r) \rightarrow \delta_{n,0} H_J^{(-)}(k_n r) - \sqrt{k_0/k_n} S_{n0}^J H_J^{(+)}(k_n r) \quad (r \rightarrow \infty)$$

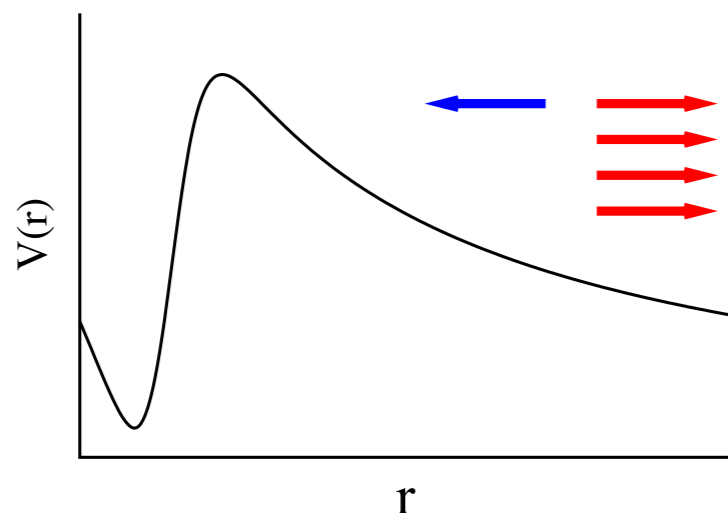
$H_J^{(\pm)}(kr)$: outgoing(imcoming) Coulomb wave functions

S_{n0}^J : S-matrix \rightarrow scattering amplitude

differential cross section for the n -th channel : $\frac{d\sigma_n}{d\Omega} = \frac{k_n}{k_0} |f_N^n(\theta) + \delta_{n,0} f_C(\theta)|^2$

quasi-elastic cross section : $\sigma_{\text{qel}}(E, \theta) = \sum_n \frac{d\sigma_n}{d\Omega}$

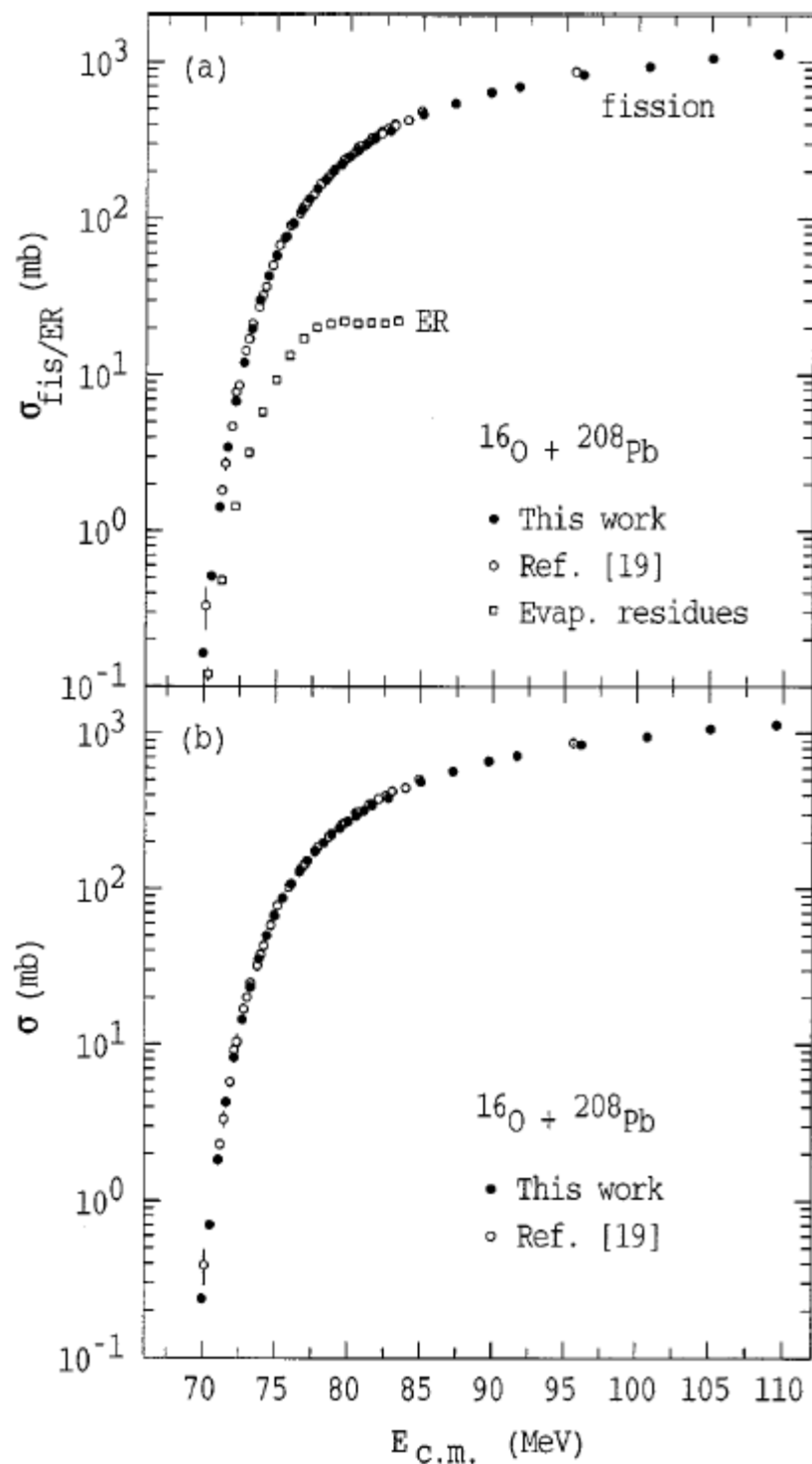
quasi-elastic barrier distribution : $D_{\text{qel}}(E, \theta) = -\frac{d}{dE} \left(\frac{\sigma_{\text{qel}}}{\sigma_R} \right)$



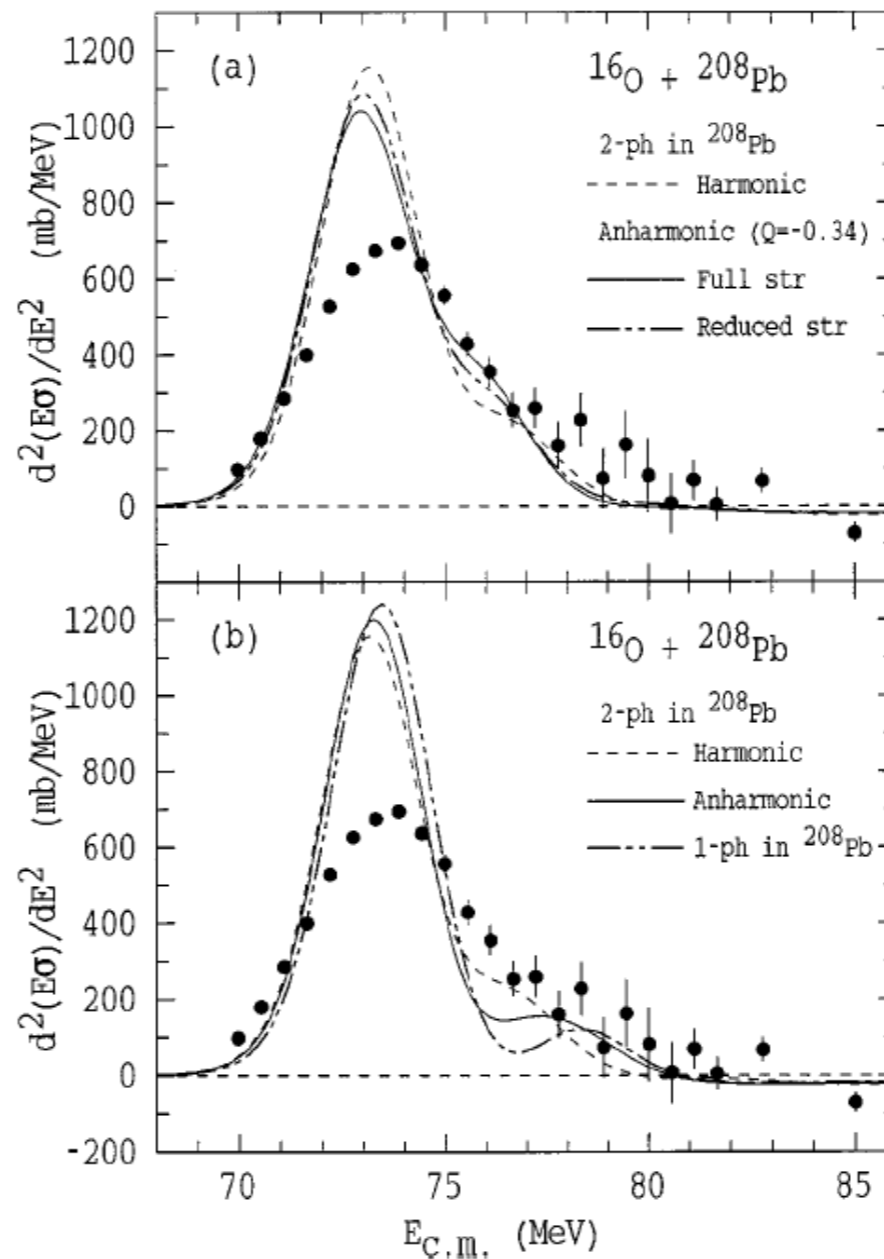
$^{16}\text{O} + ^{208}\text{Pb}$

fusion reaction and quasi-elastic scattering with non-coll. excitations

precise fusion data



fusion barrier distribution



C. R. Morton *et al.*,
PRC60, 044608(1999)

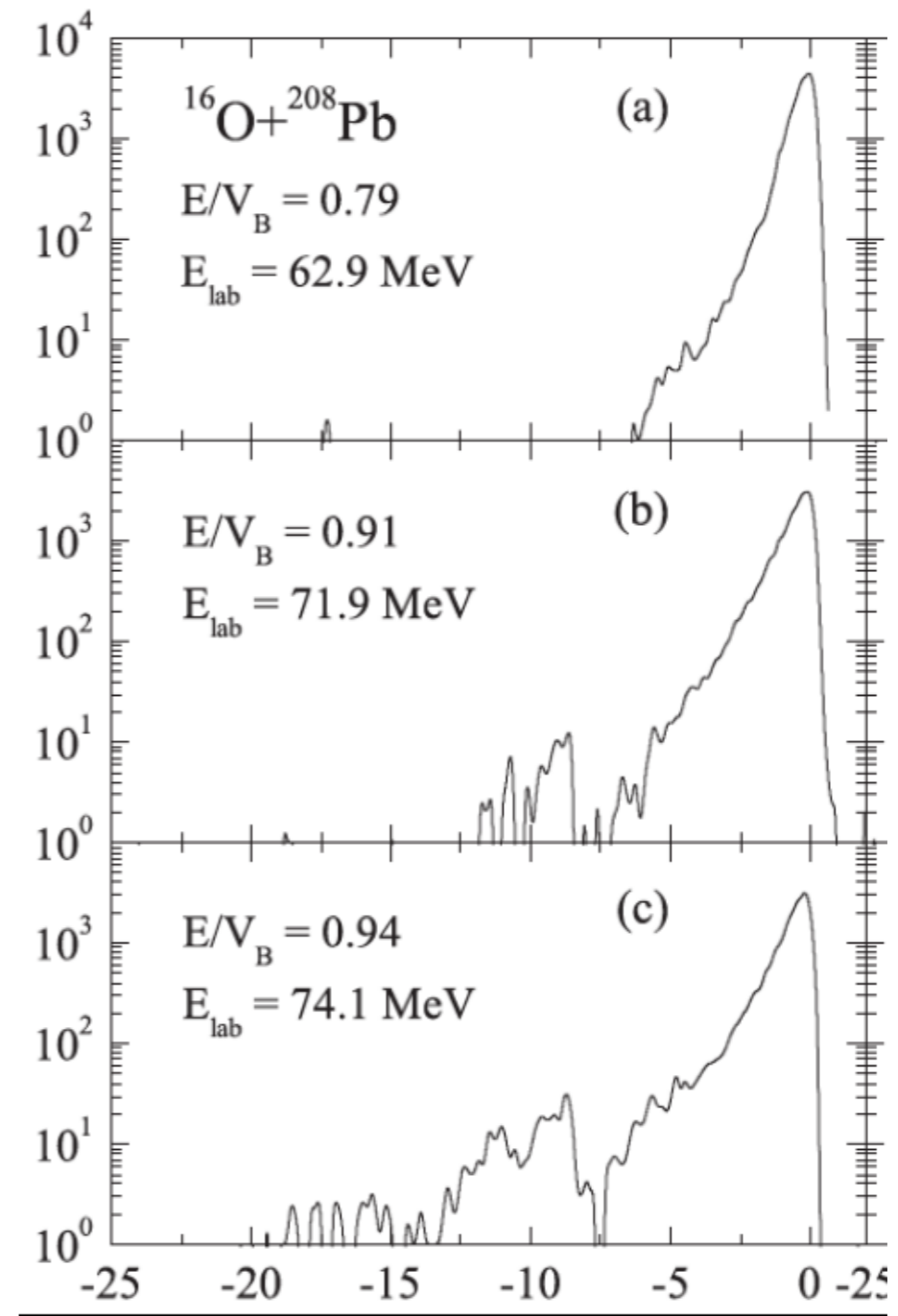
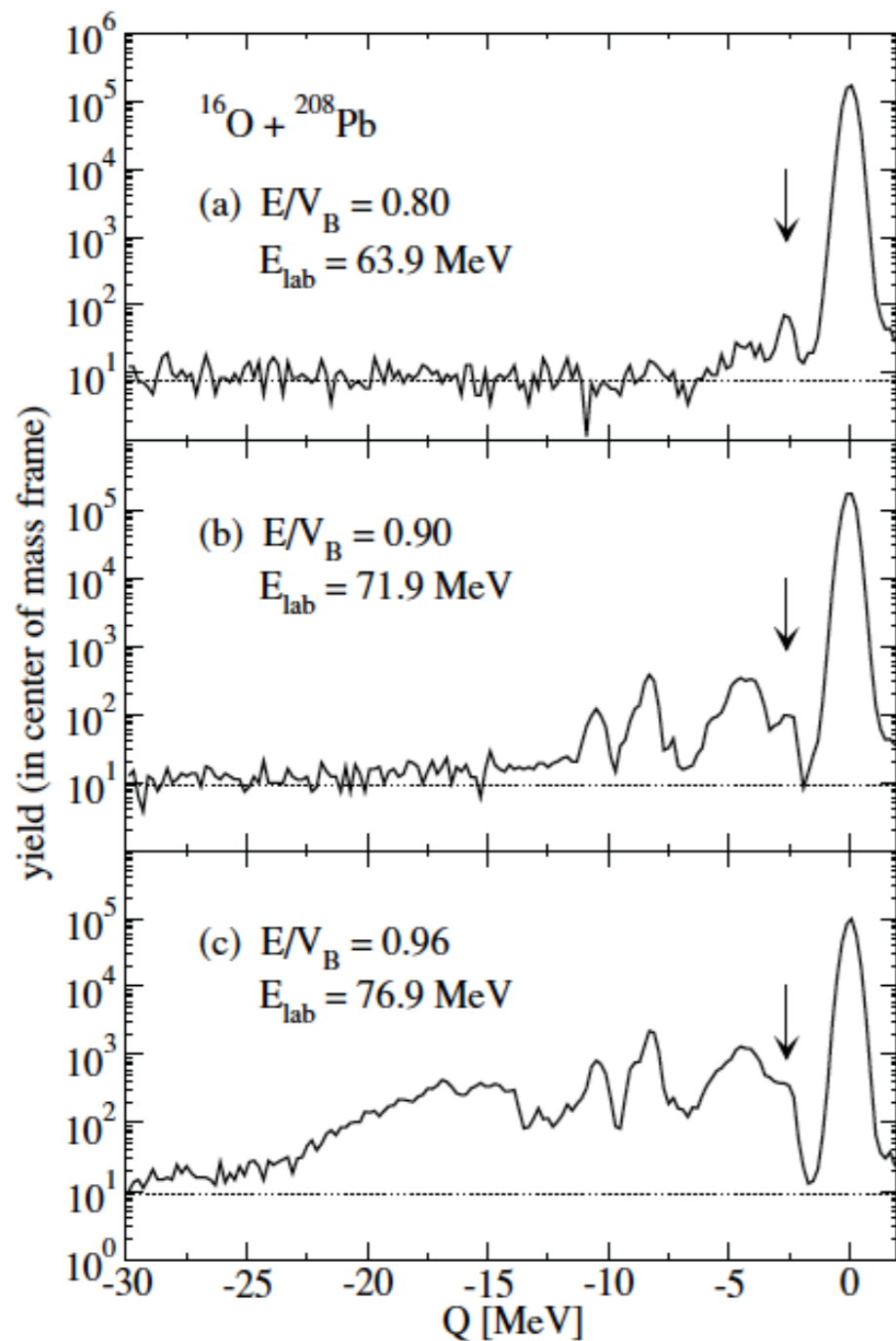
analysis with collective excitations

→ no satisfactory description has obtained

$^{16}\text{O} + ^{208}\text{Pb}$

fusion reaction and quasi-elastic scattering with
non-coll. excitations

Energy dependence of Q-value distribution



$^{16}\text{O} + ^{208}\text{Pb}$

fusion reaction and quasi-elastic scattering with non-coll. excitations

high precision proton inelastic scattering experiment of ^{208}Pb

W.T.Wagner, *et al.* PRC $\underline{12}$ 757(1975)

DWBA analysis

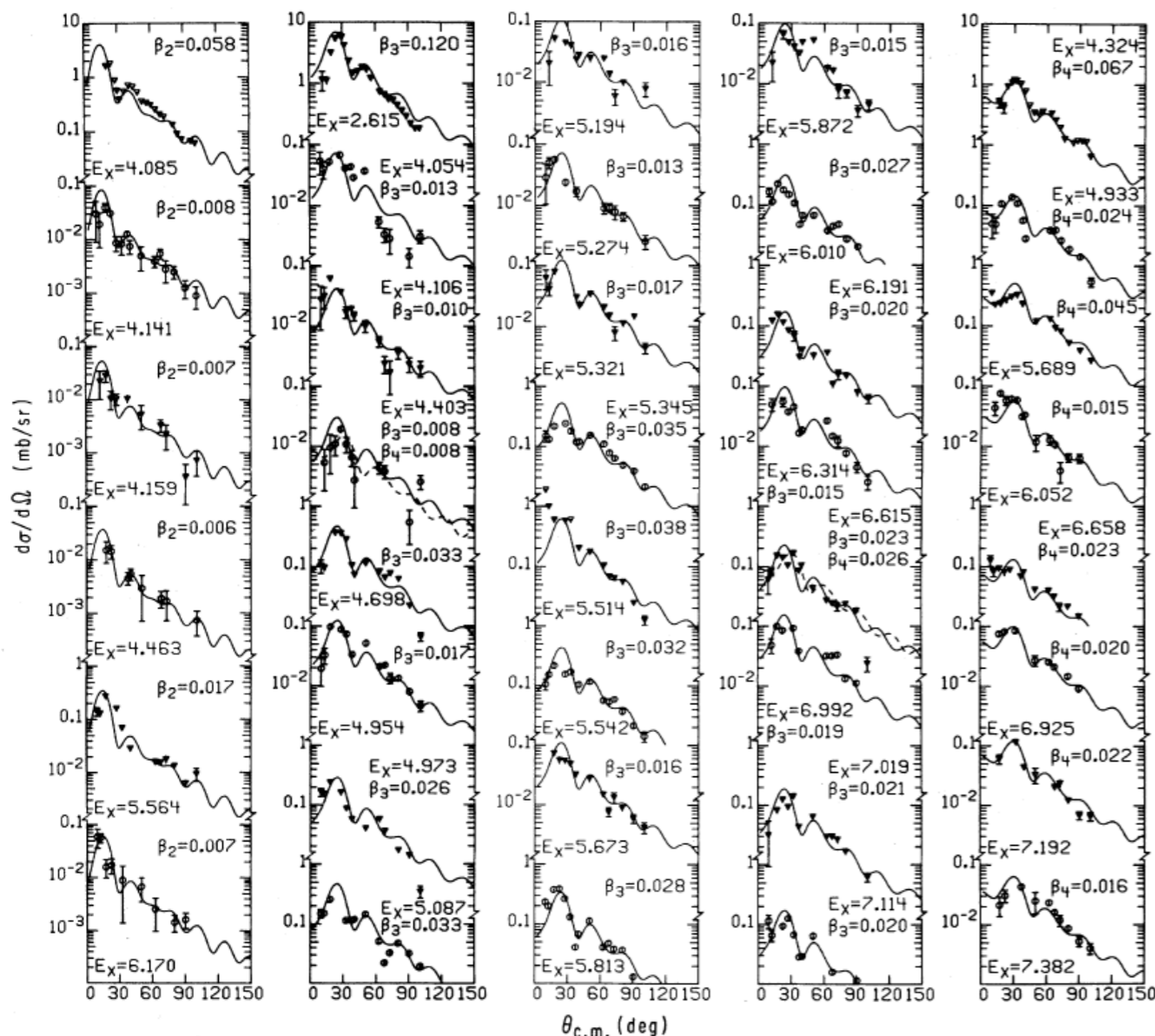


energy E^* and deformation parameter β_λ of excited states up to about 7 MeV



couplings to these non-collective excitations

(with coupling form factor for vibrational coupling)



Results

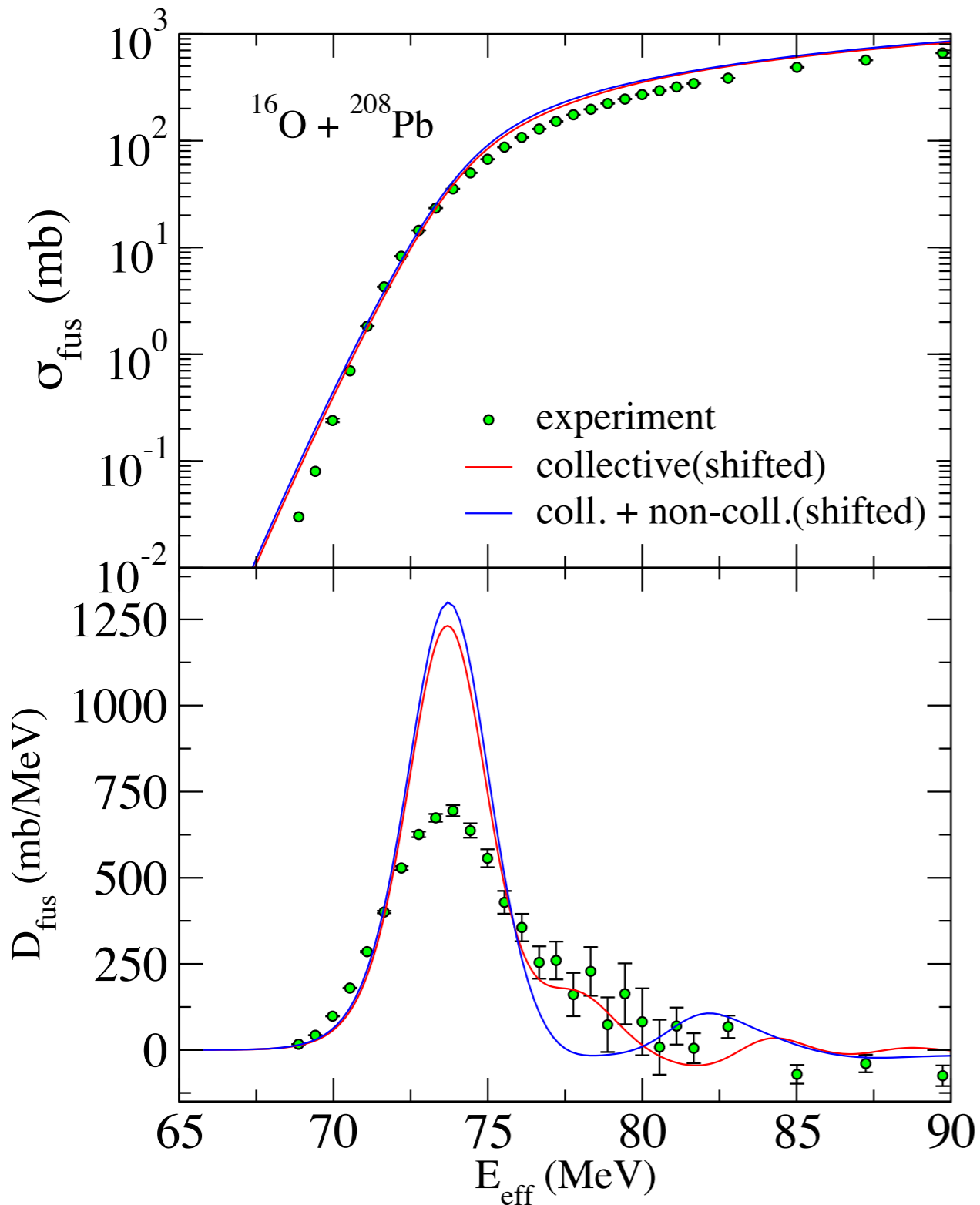
$^{16}\text{O} + ^{208}\text{Pb}$ fusion reaction
with 70 non-coll. states

※ collective states

$^{208}\text{Pb} : 3^- : 2.615 \text{ MeV}, \beta_3 = 0.122$

$5^- : 3.198 \text{ MeV}, \beta_5 = 0.058$

$^{16}\text{O} : 3^- : 6.13 \text{ MeV}, \beta_3 = 0.733$

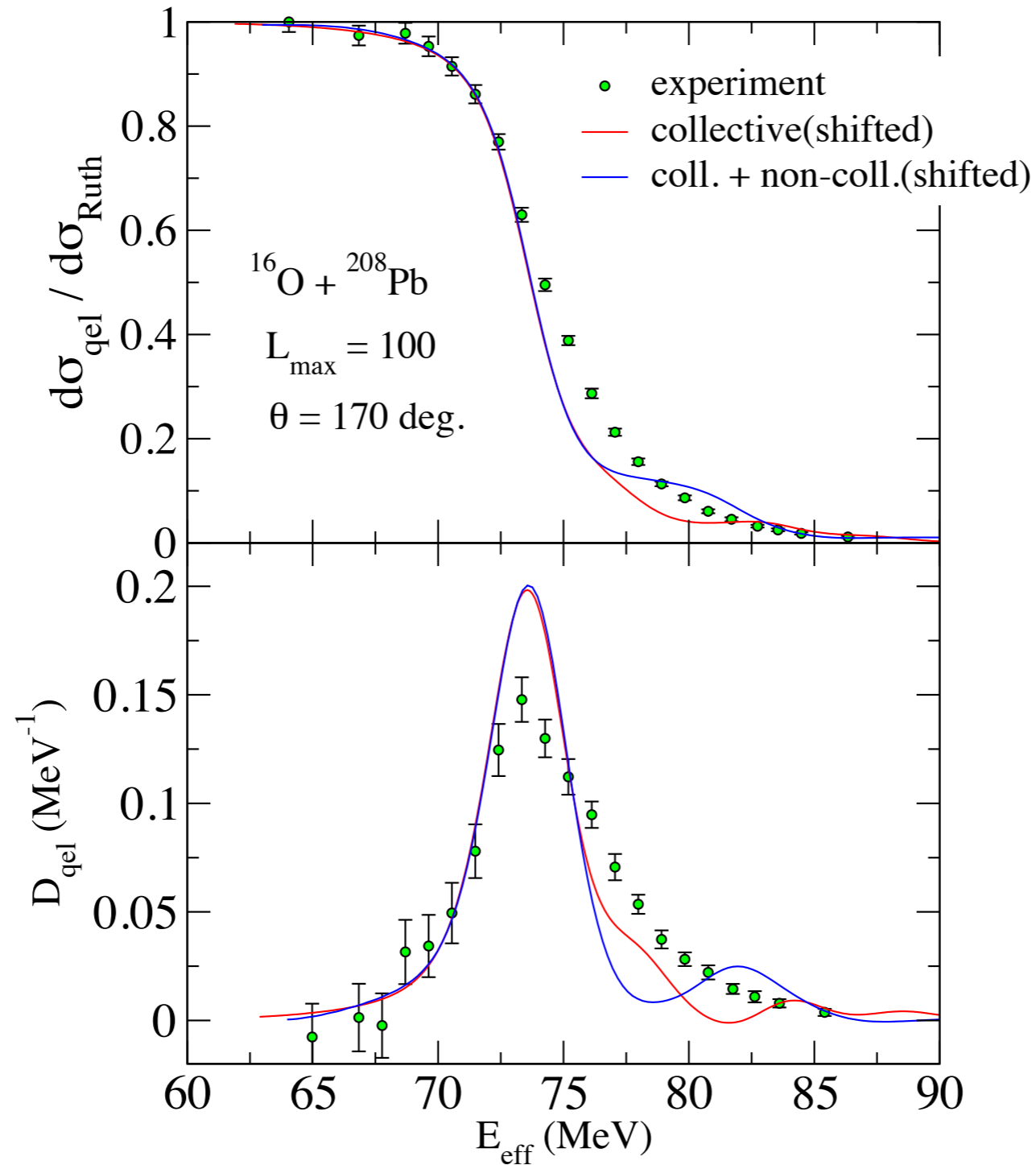


✓ barrier distribution becomes single peak structure due to the non-collective excitations(smearing effect)

✓ non-collective excitations do not improve the agreement

$^{16}\text{O} + ^{208}\text{Pb}$

quasi-elastic scattering



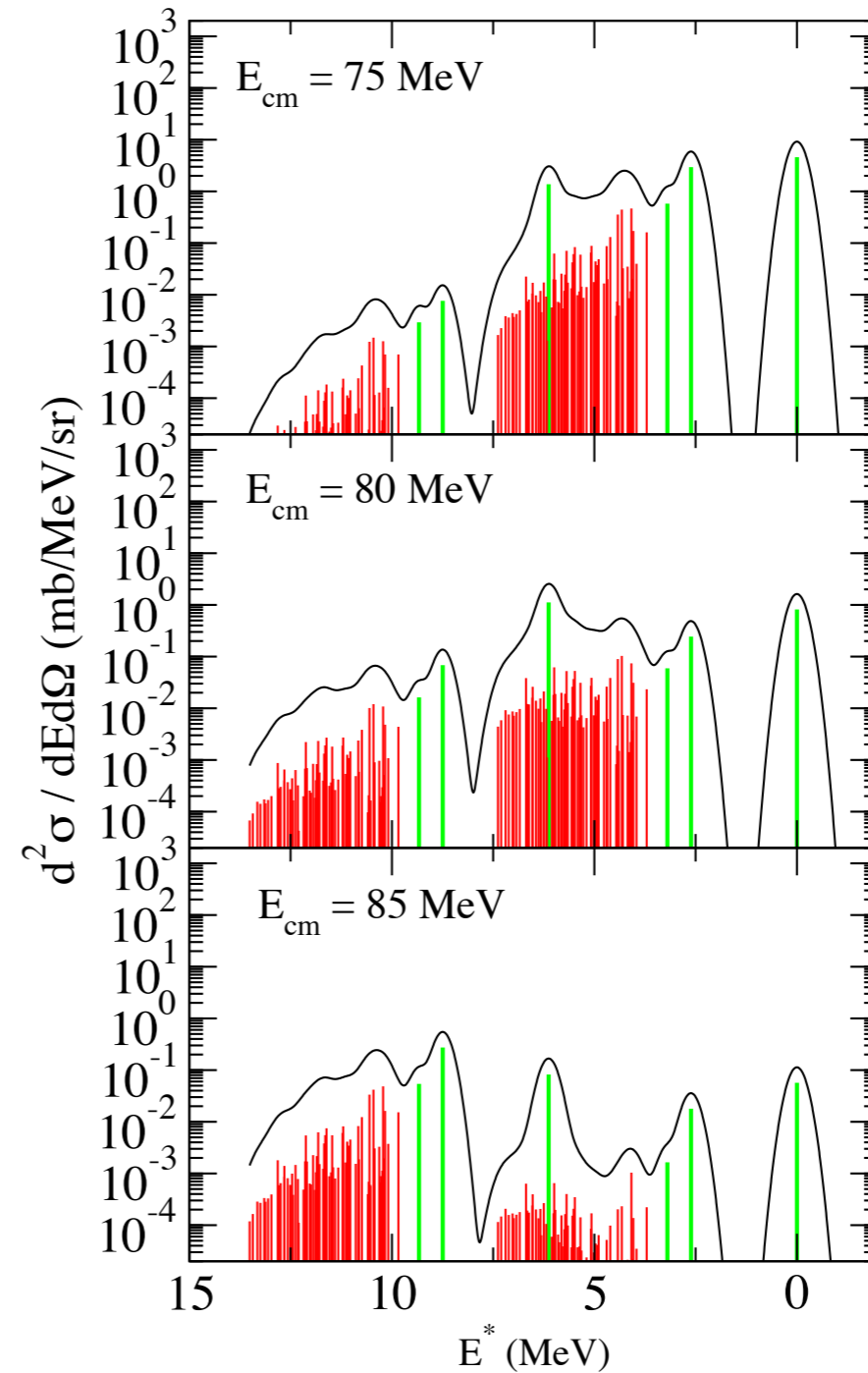
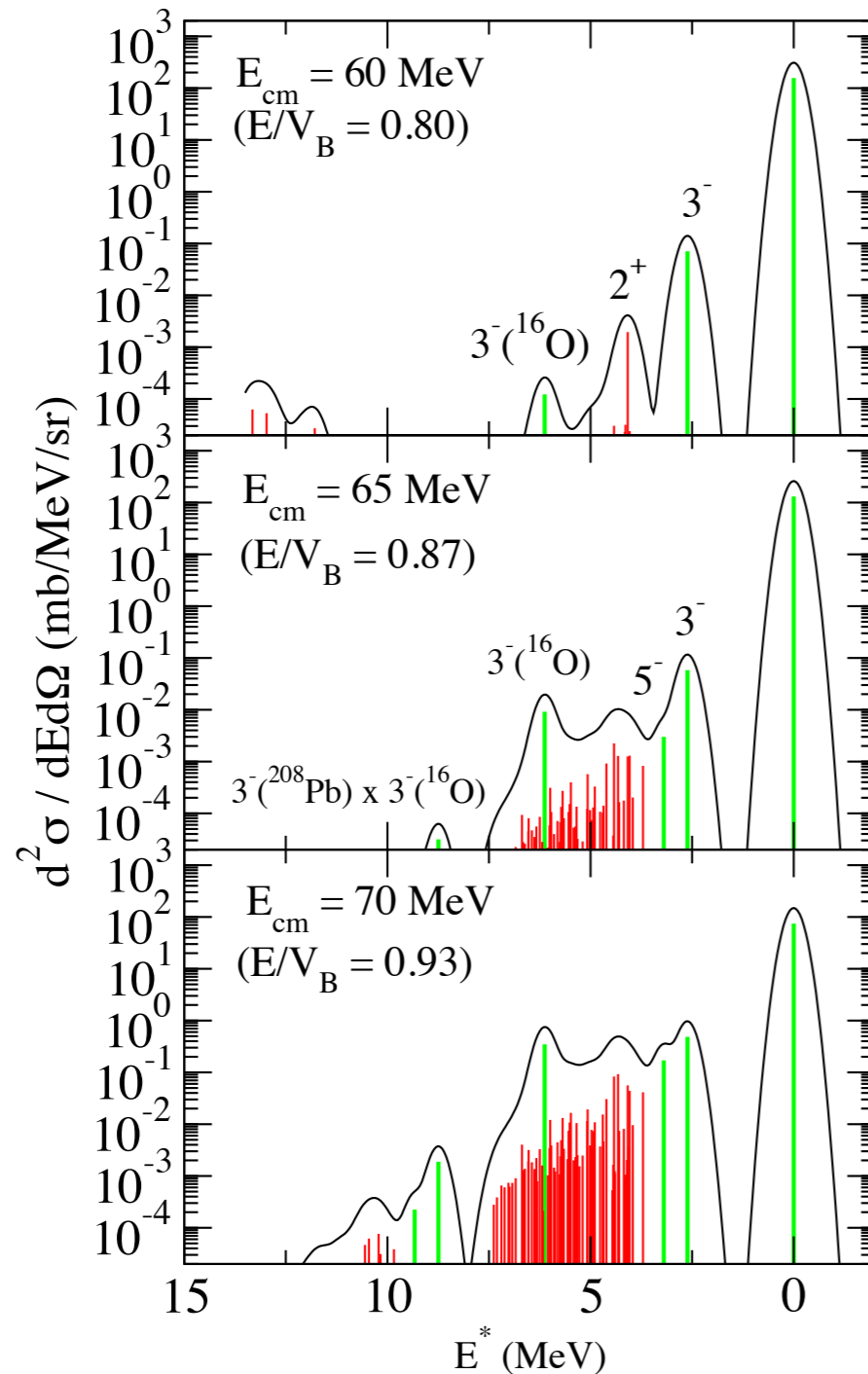
- smearing effect as in the case for fusion
- non-collective excitations do not improve the agreement between experimental barrier distributions also in this case

$^{16}\text{O} + ^{208}\text{Pb}$

$$F(E^*) = \sum_n f_n \frac{1}{\sqrt{2\pi\sigma}} e^{-\frac{(E^* - \epsilon_n)^2}{2\sigma^2}}$$

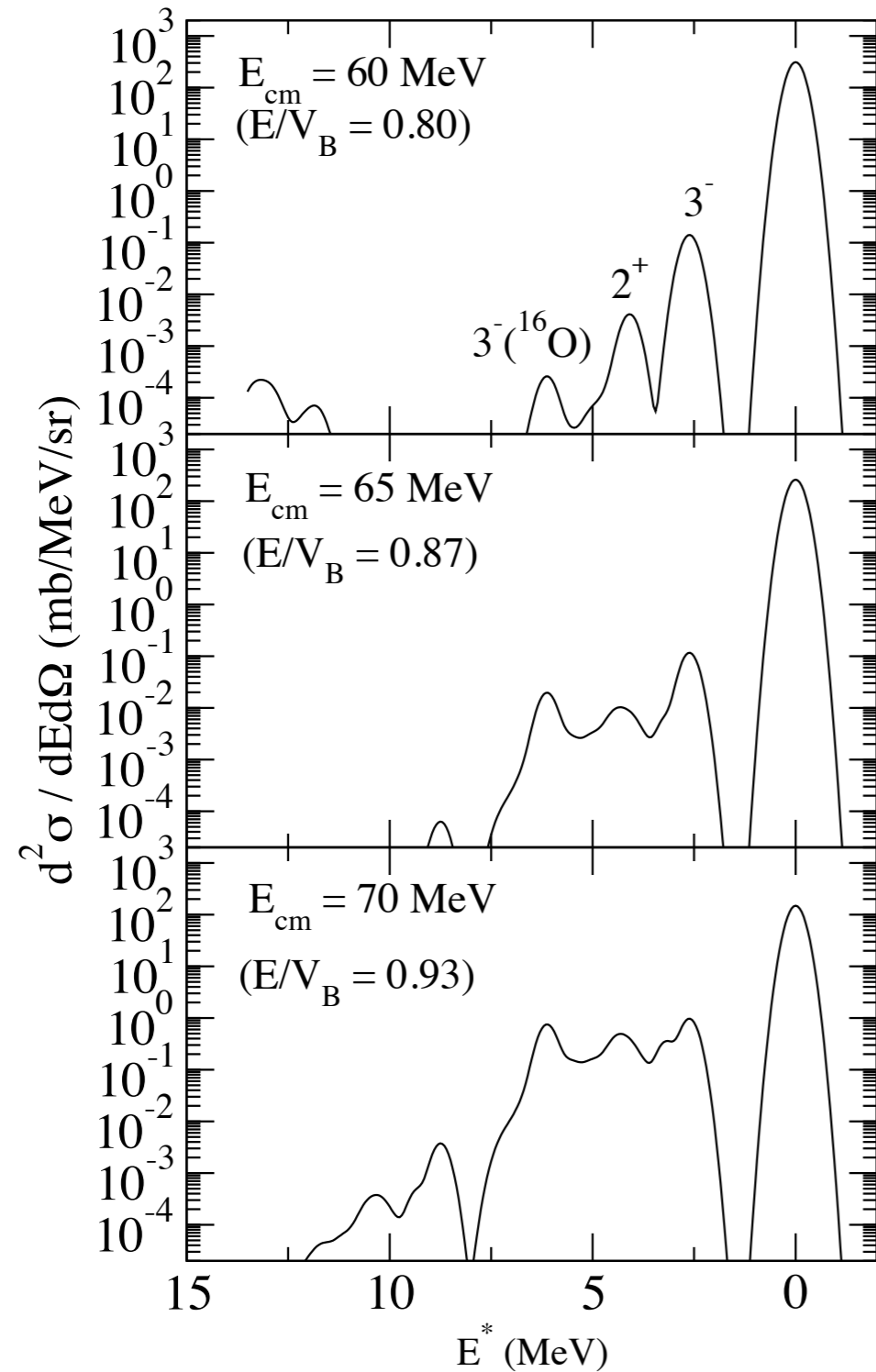
Energy dependence of Q-value distribution ($\theta = 170^\circ$)

($\sigma = 0.2$ MeV)

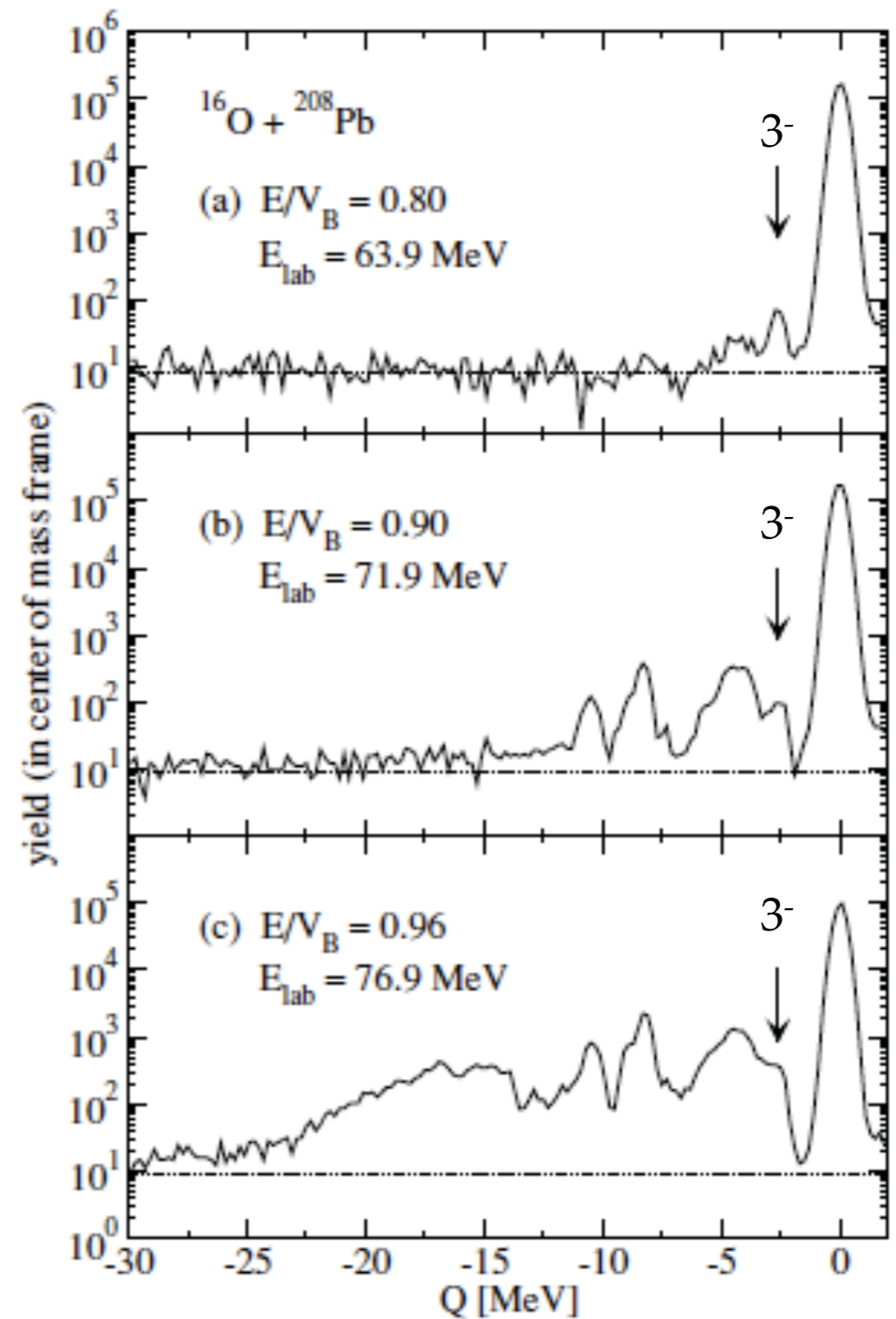


non-collective excitations becomes more and more important as the incident energy increases

$^{16}\text{O} + ^{208}\text{Pb}$



experiment



M. Ever et al. Phys.Rev.C 78('08)034614

obtained Q -value distributions are consistent with the experimental data

Summary

Low-energy heavy-ion reaction (fusion reaction and quasi-elastic scattering) taking into account non-collective excitations

- ✓ Conventional coupled-channels calculation failed to reproduce quasi-elastic barrier distributions for $^{20}\text{Ne} + ^{90,92}\text{Zr}$
- ✓ Coupled-channels calculation with **non-collective excitations** is applied to $^{16}\text{O} + ^{208}\text{Pb}$ system for fusion reaction and quasi-elastic scattering
- ✓ **Non-collective excitations in this calculation do not improve the fusion and quasi-elastic barrier distribution**
- ✓ **Energy dependence of the Q-value distribution is consistent with experimental data**

Future perspectives

- Application to $^{20}\text{Ne} + ^{90,92}\text{Zr}$ systems for which nuclear information about the non-collective excitations has not been obtained
- Description of non-collective excitations based on random matrix theory