Topological susceptibility and axial symmetry at finite temperature

JICFuS

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Mini-workshop

Summary:

- ✔ Motivation
- ✔ Chiral phase transition at finite temperature and axial symmetry
- ✔ Simulating dynamical overlap fermions
- ✔ Topology fixing and friends
- ✔ Methodology and results (quenched and N_f=2)
- ✔ Discussion and conclusions

People involved in the project

JLQCD group: S. Hashimoto, S. Aoki, T. Kaneko, H. Matsufuru, J. Noaki, E. Shintani

See for example POS(Lattice2010)174 (arXiv::1011.0257), PoS(Lattice 2011)188, article in prep.

Motivation

Pattern of chiral symmetry breaking at low temperature QCD

$$SU(N_f)_V \times SU(N_f)_A \times U(1)_V \times U(1)_A \to SU(N_f)_V \times U(1)_V$$

Well known facts

- Axial anomaly – Non vanishing topological susceptibility
- Mass splitting of the η' with respect to the lighter mesons

What is the fate of the axial $U_A(1)$ symmetry at finite temperature($T \gtrsim T_c$)?

Dirac Overlap operator, retaining the maximal amount of chiral symmetry on the lattice is, theoretically, the best way to answer this question.

Goal of the current project

Check restoration of axial $U_A(1)$ symmetry by measuring (spatial) meson correlators at finite temperature in full QCD with overlap

Degeneracy of correlators is the signal that we are looking for

$$\sigma(1_4 \otimes 1_2) \xleftrightarrow{\text{Chiral sym.}} \pi(i\gamma_5 \otimes \tau^a)$$

$$\sigma(1_4 \otimes 1_2) \xleftrightarrow{U(1)_A} \eta(i\gamma_5 \otimes 1_2) \qquad \pi(i\gamma_5 \otimes \tau^a) \xleftrightarrow{U(1)_A} \delta(1_4 \otimes \tau^a)$$

$$\eta(i\gamma_5 \otimes 1_2) \xleftrightarrow{\text{Chiral sym.}} \delta(1_4 \otimes \tau^a)$$

As I will show in this talk,
there are some issues to solve before attacking the real problem...

Fixing topology: how to deal with it at T=0 (I)

Partition function at fixed topology

$$Z_Q = \frac{1}{2\pi}\int_{-\pi}^{\pi} d\theta\ \exp(-VF(\theta)) \qquad F(\theta) \equiv E(\theta) - i\theta Q/V$$

where the energy can be expanded

$$E(\theta) = \sum_{n=1}^{\infty} \frac{c_{2n}}{(2n)!}\theta^{2n} = \frac{\chi_t}{2}\theta^2 + O(\theta^4)$$

Using saddle point expansion around $\theta_c = i\frac{Q}{\chi_t V}(1 + O(\delta^2))$

one obtains the Gaussian distribution

$$Z_Q = \frac{1}{\sqrt{2\pi\chi_t V}} \exp\left[-\frac{Q^2}{2\chi_t V}\right]\left[1 - \frac{c_4}{8V\chi_t^2} + O\left(\frac{1}{V^2}, \delta^2\right)\right].$$

Fixing topology: how to deal with it at T=0 (II)

From the previous partition function we can extract the relation between correlators at fixed θ and correlators at fixed Q

In particular for the topological susceptibility and using the Axial Ward Identity we obtain a relation involving fermionic quantities:

$$\lim_{|x|\to\text{large}} \langle mP(x)mP(0)\rangle_Q^{\text{disc}} \equiv \frac{1}{V}\left(\frac{Q^2}{V} - \chi_t - \frac{c_4}{2\chi_t V}\right) + O(e^{-m_\pi |x|})$$

P(x) is the flavor singlet pseudo scalar density operator

Aoki *et al*. PRD76,054508 (2007)

What is the effect of fixing Q at finite temperature?

Results

- ✔ Simulation details
- ✔ Eigenvalues distribution
- ✔ Finite temperature quenched SU(3) at fixed topology
- ✔ Meson correlators in two flavors QCD

BG/L
Hitachi SR16K

Simulation details

Pure gauge (16^3x6, 24^3x6):

Iwasaki action + top. fixing term

β	a(fm)	T (MeV)	T/Tc
2.35	0.132	249.1	0.86
2.40	0.123	268.1	0.93
2.43	0.117	280.9	0.97
2.44	0.115	285.7	0.992
2.445	0.114	288	1.0
2.45	0.1133	290.2	1.01
2.46	0.111	295.1	1.02
2.48	0.107	305.6	1.06
2.50	0.104	316.2	1.10
2.55	0.094	347.6	1.20

Two flavors QCD (16^3x8)

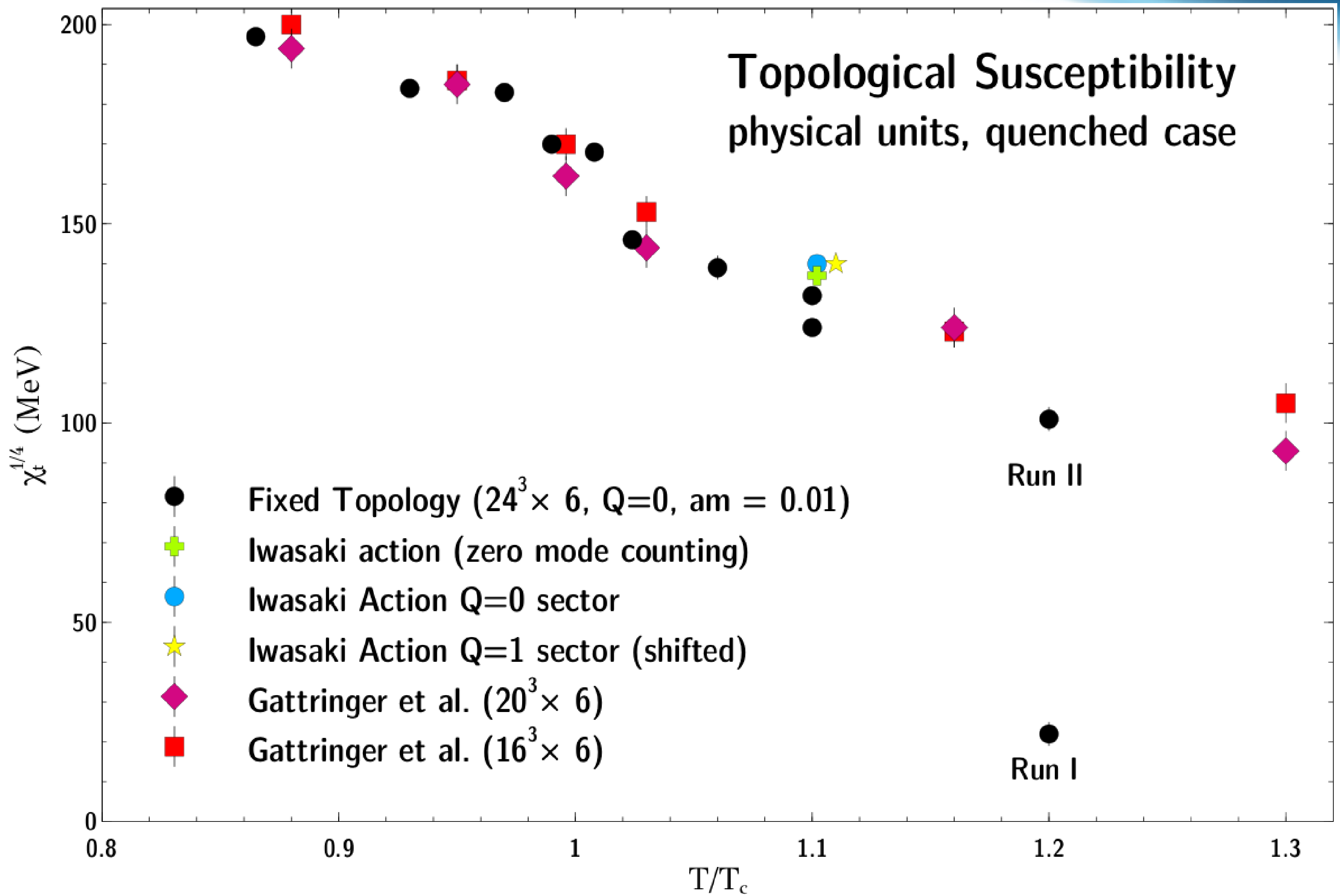
Iwasaki + Overlap + top. Fix

O(300) trajectories per T

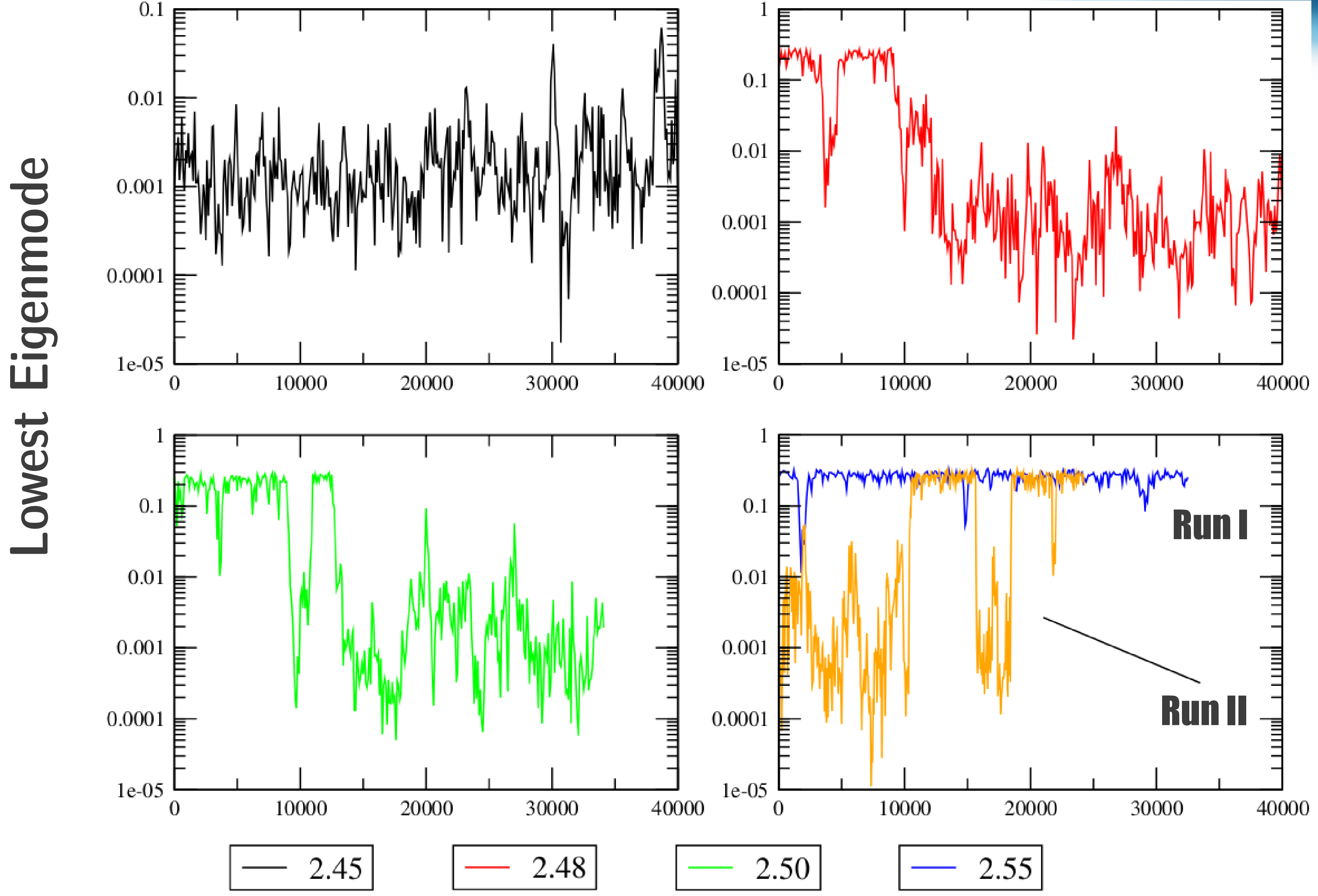
am=0.05, 0.025, 0.01

β	a(fm)	T (MeV)	T/Tc
2.18	0.1438	171.5	0.95
2.20	0.1391	177.3	0.985
2.25	0.12818	192.2	1.06
2.30	0.1183	208.5	1.15
2.40	0.1013	243.5	1.35
2.45	0.0940	262.4	1.45

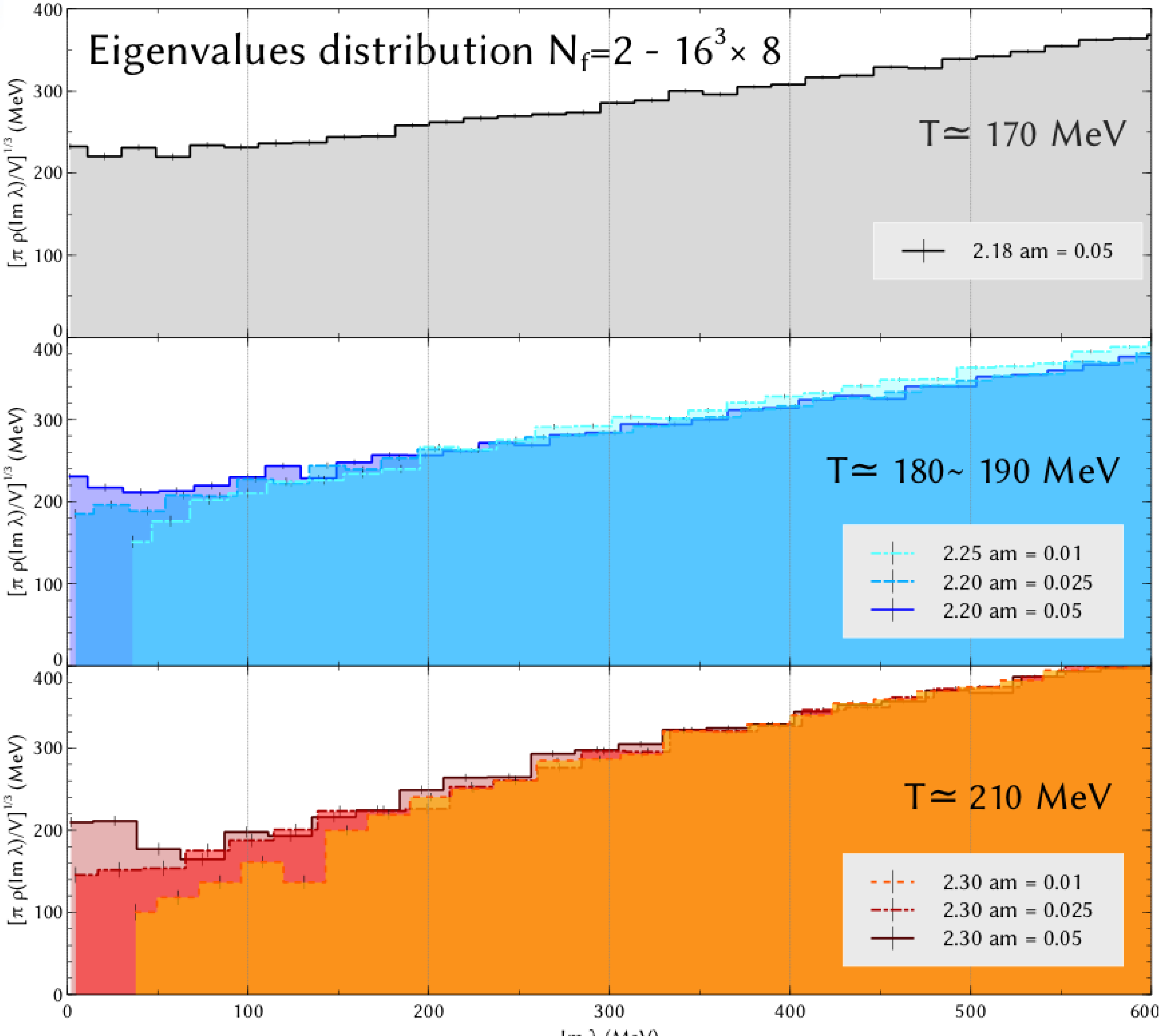
Topological susceptibility in pure gauge theory - I



Topological susceptibility in pure gauge theory - III



Full QCD – Eigenvalues



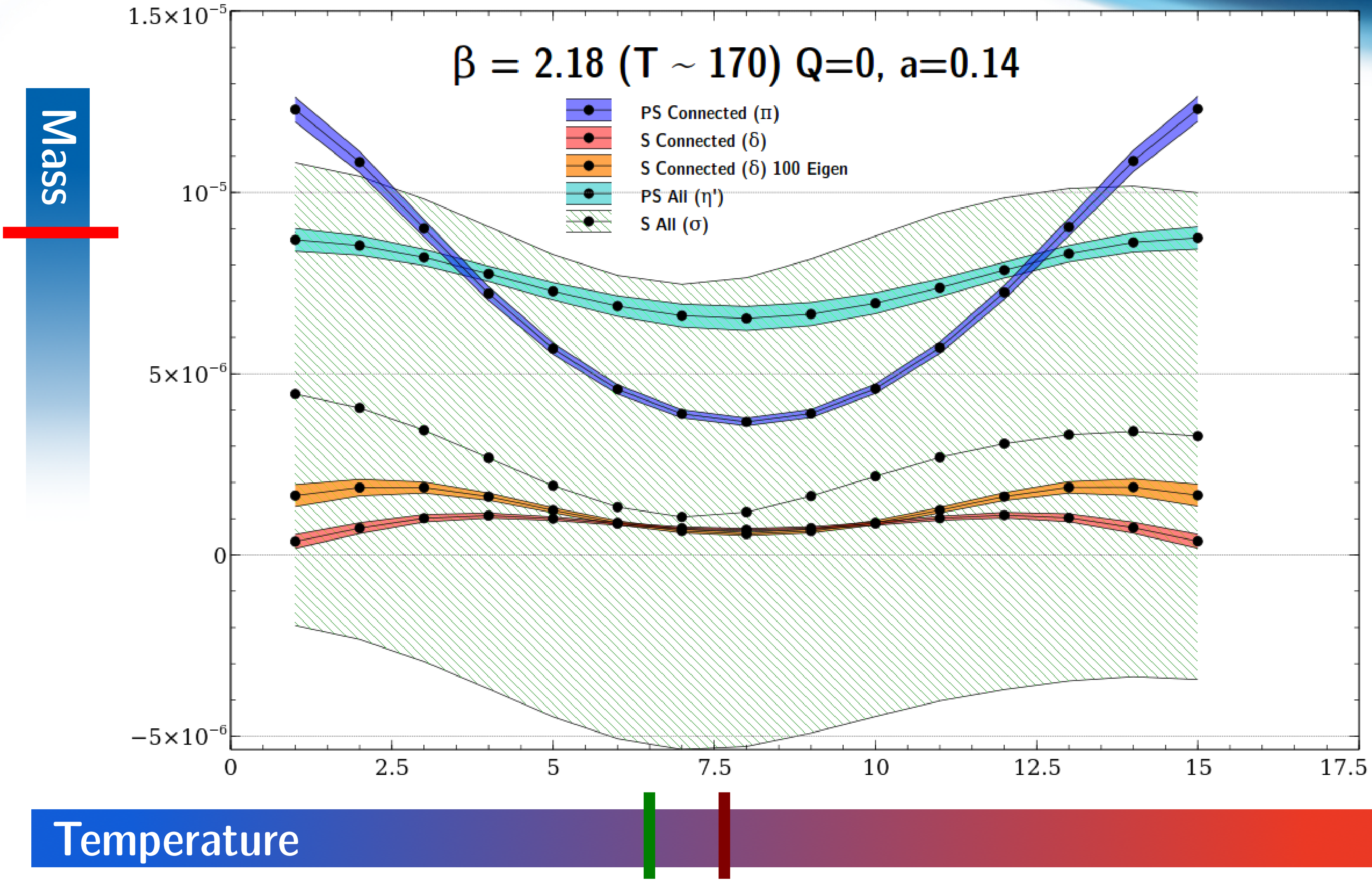
Effect of axial symmetry on the Dirac sperctrum

$$\chi^{\pi-\delta} = \int d\lambda \rho_m(\lambda) \frac{4m^2}{\lambda^2 + m^2}$$

If chiral symmetry is restored we can conclude that

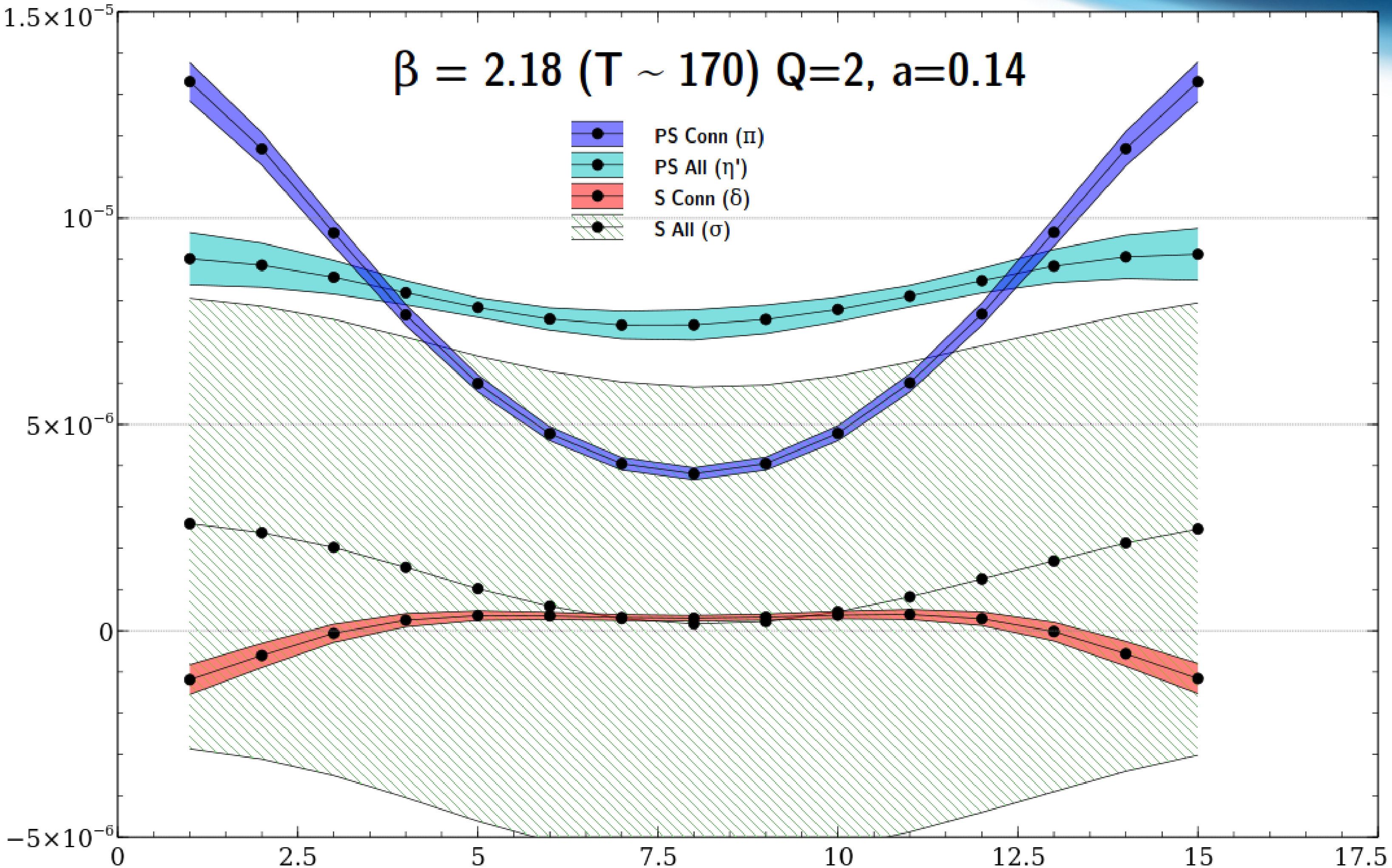
$$\lim_{\lambda\to 0} \lim_{m\to 0} \frac{\rho_m(\lambda)}{\lambda} = 0$$

Full QCD – Meson correlators



Full QCD – Meson correlators

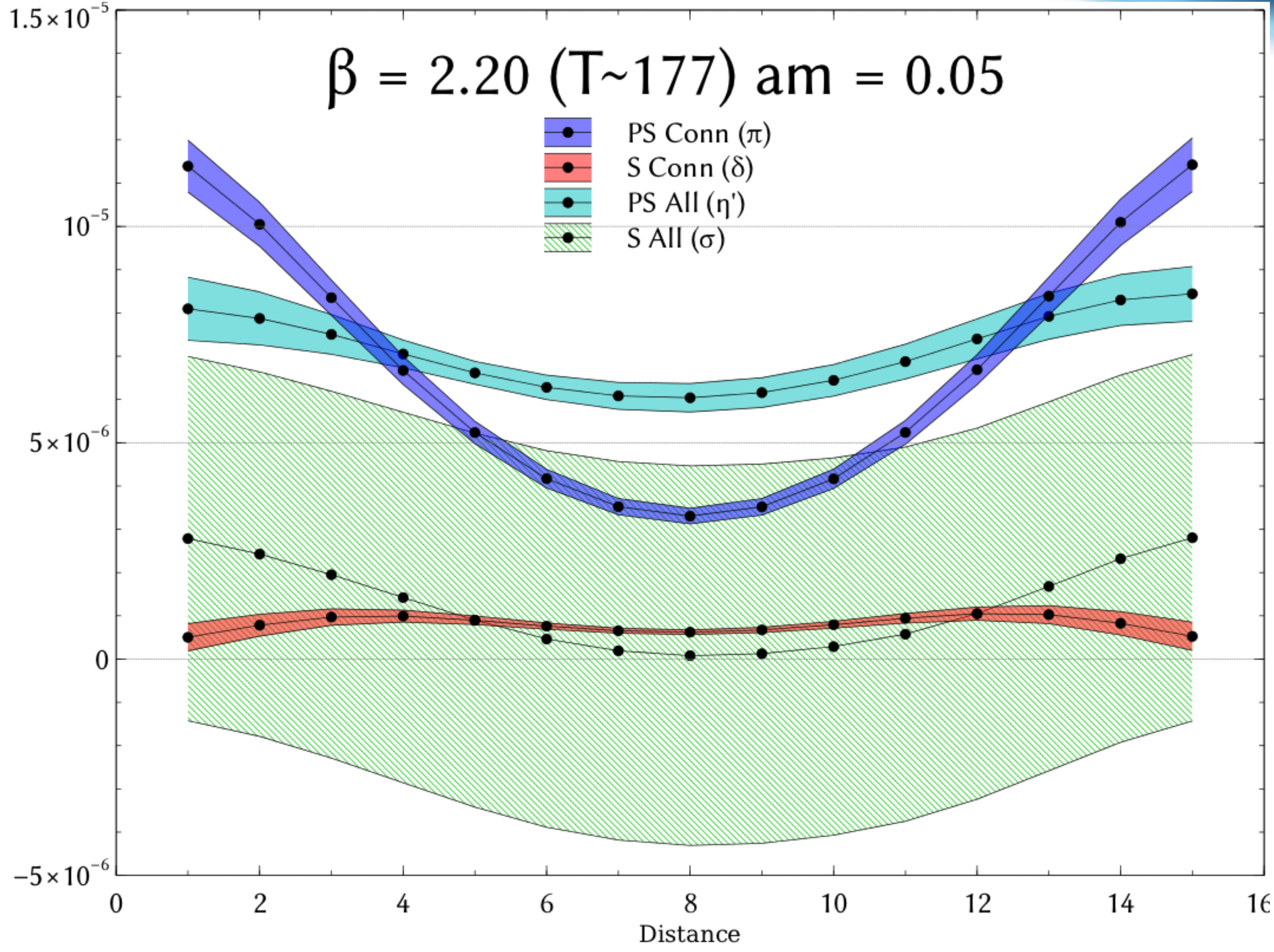
Mass



Temperature

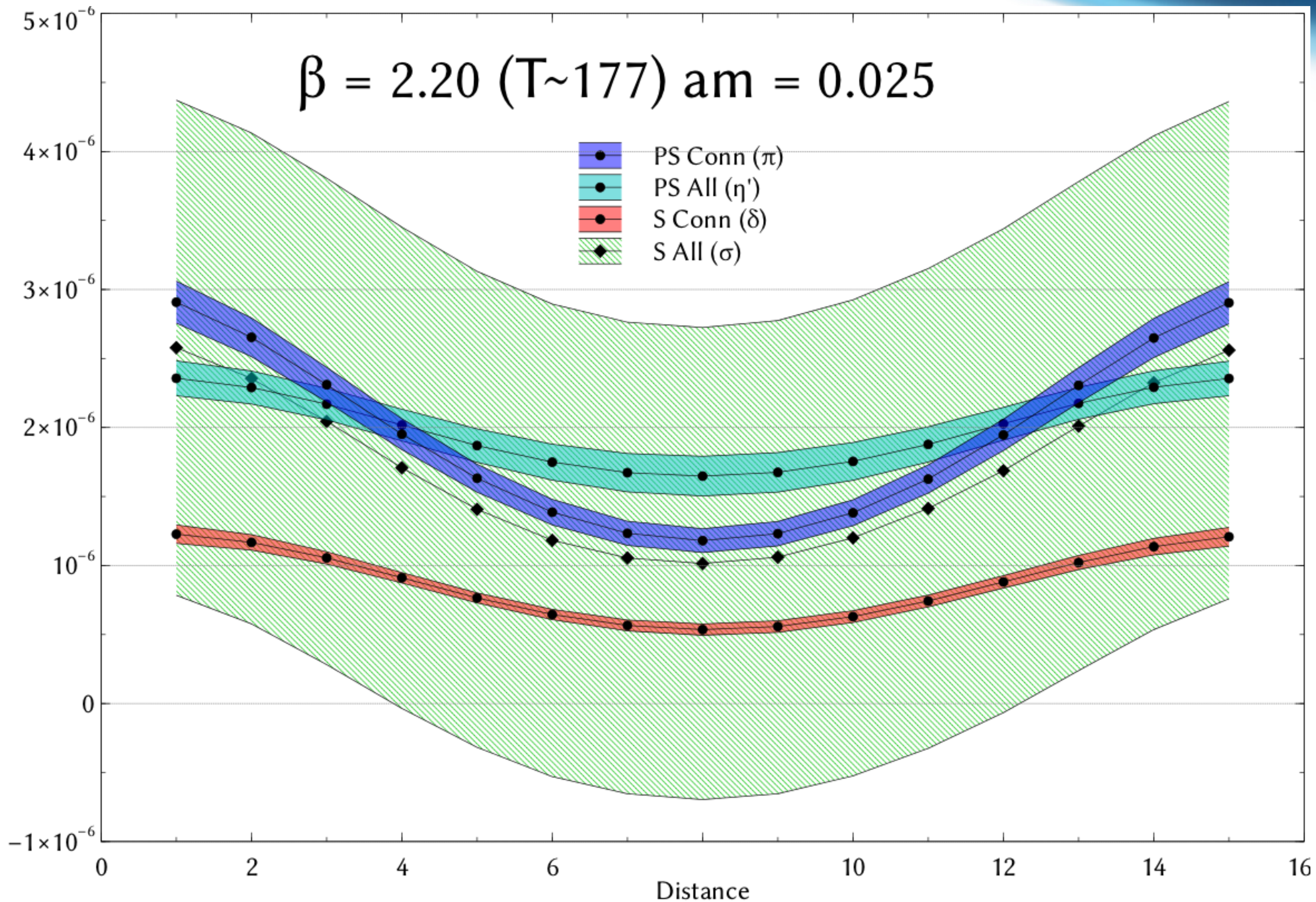
Full QCD – Meson correlators

Mass



Temperature

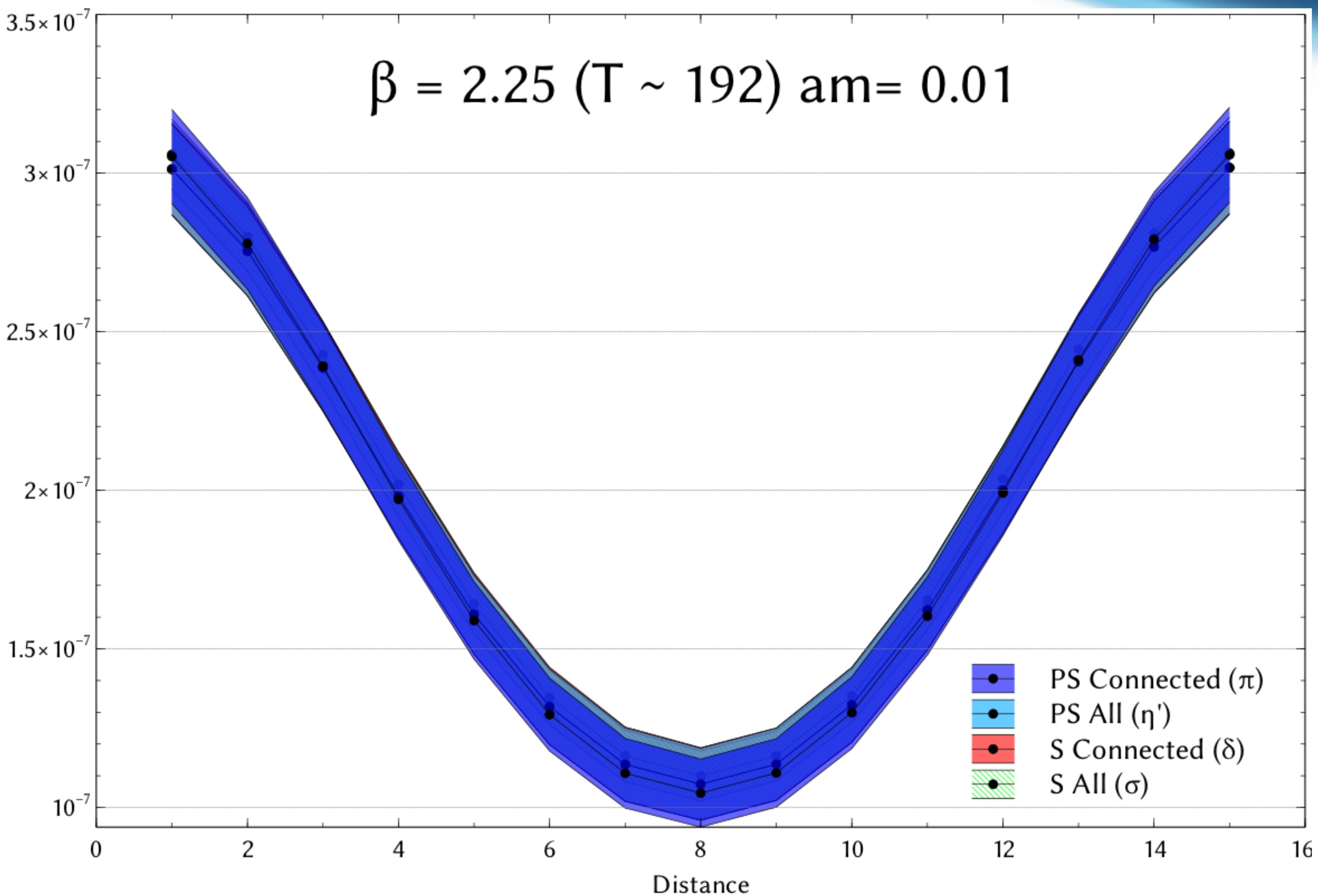
Full QCD – Meson correlators



Mass

Temperature

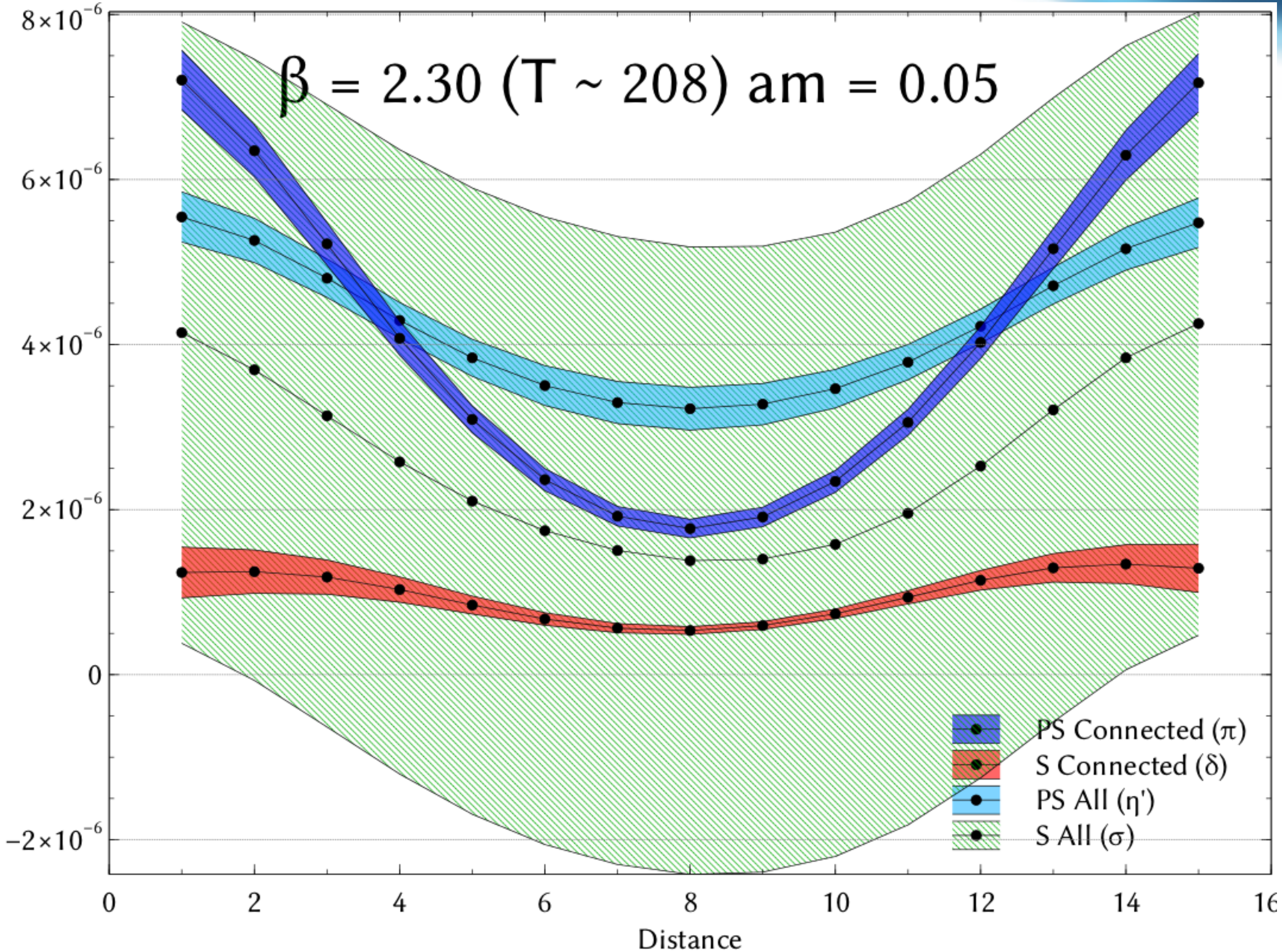
Full QCD – Meson correlators

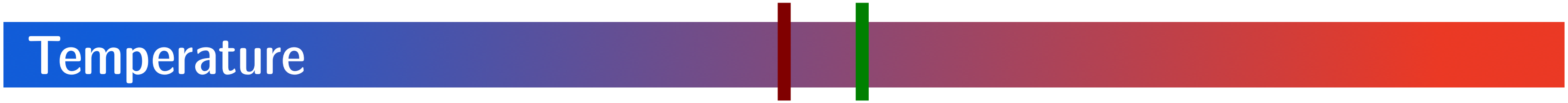


Mass

Temperature

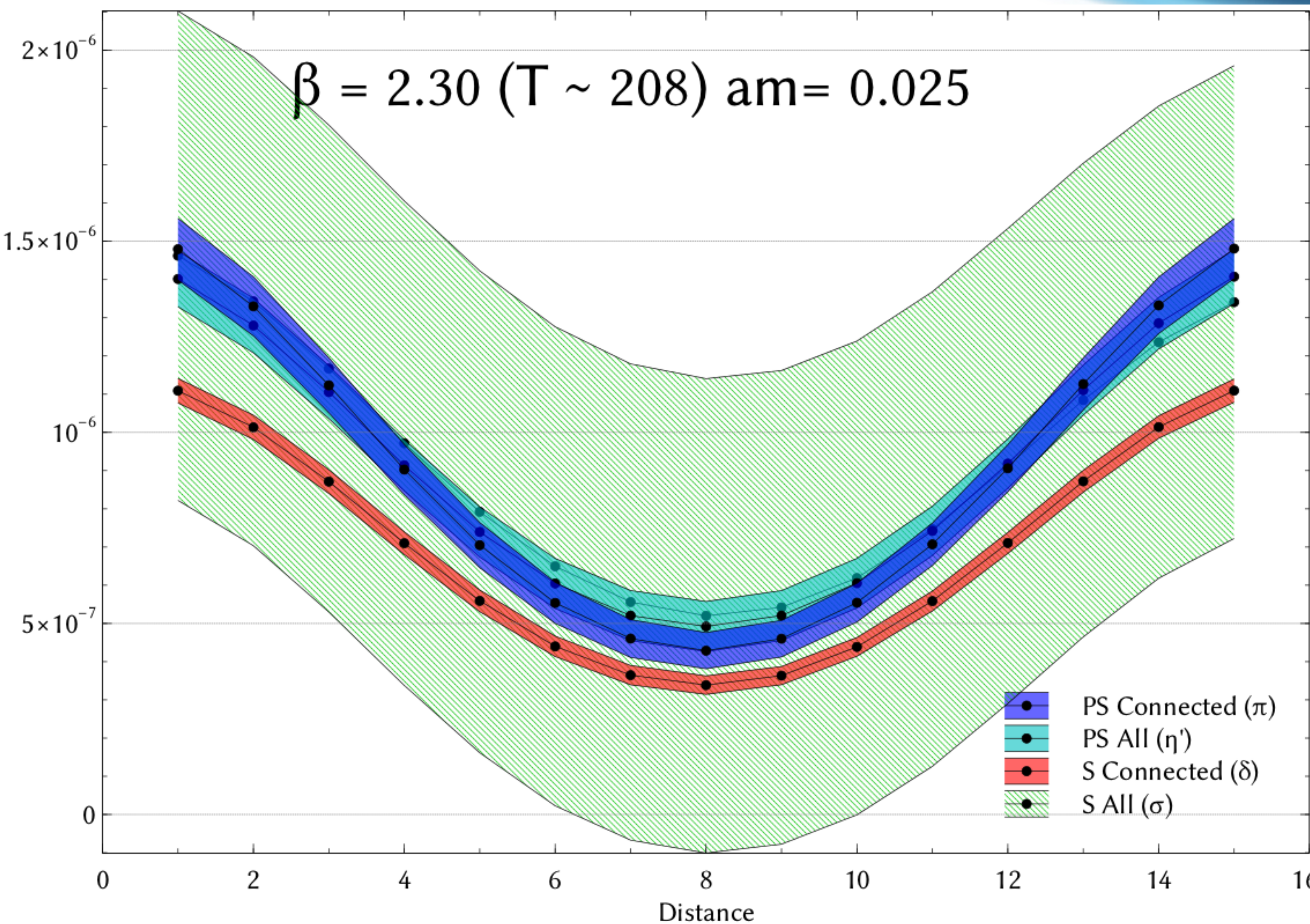
Full QCD – Meson correlators





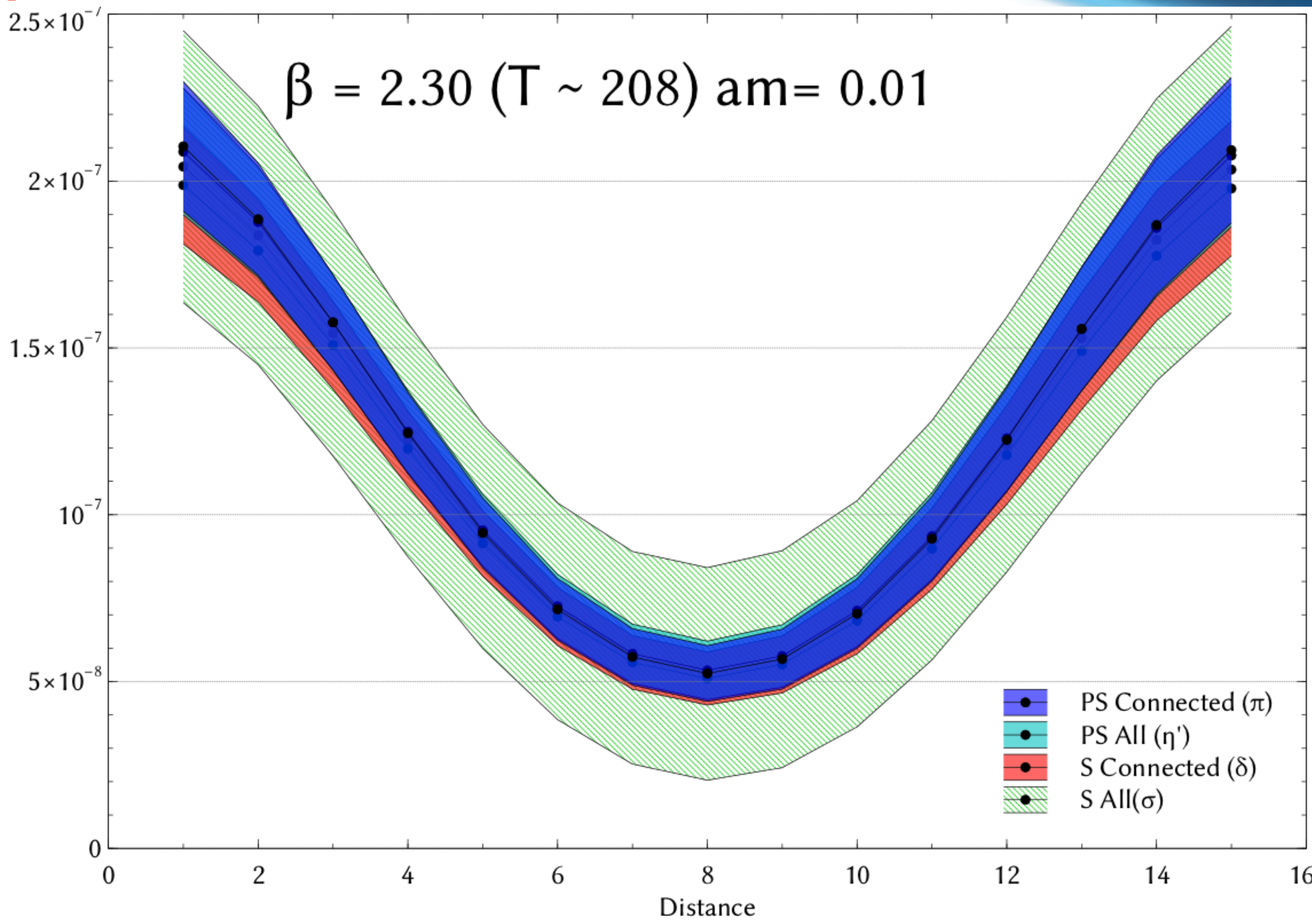
Full QCD – Meson correlators

Mass



Temperature

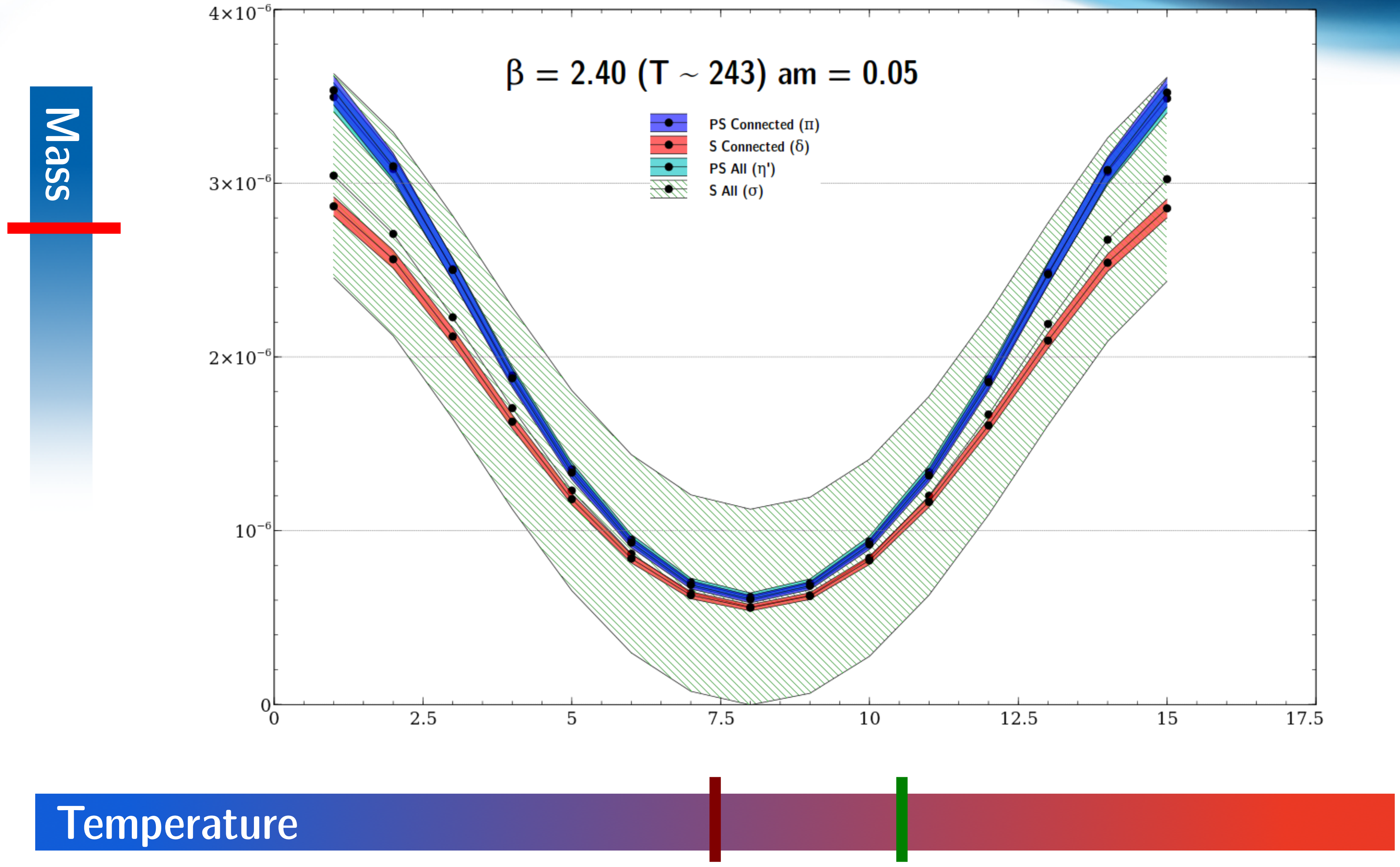
Full QCD – Meson correlators



Mass

Temperature

Full QCD – Meson correlators



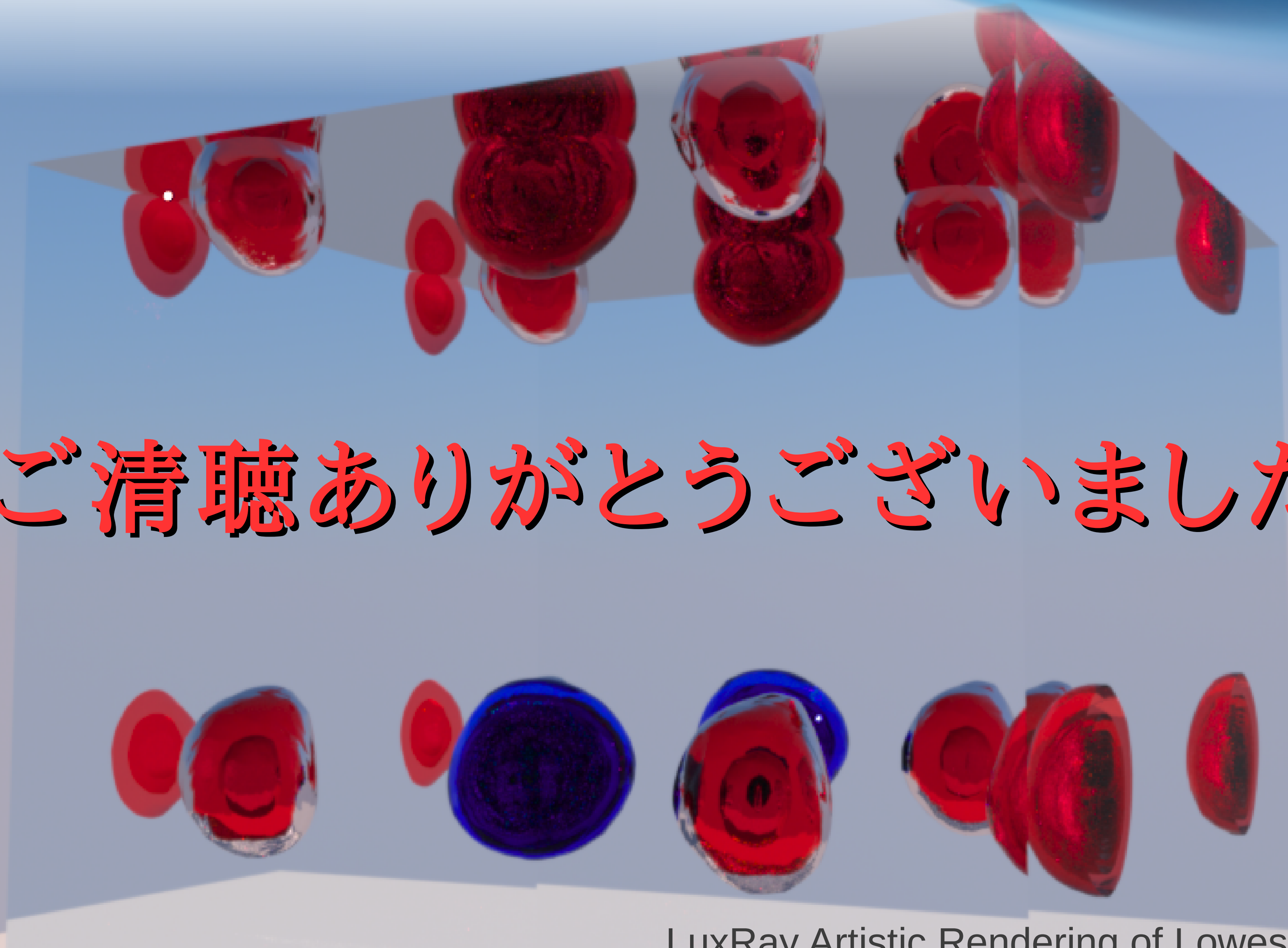
Summary

- Overlap fermions are the best choice to check axial anomaly at finite temperature
- Current machine and algorithms permit now realistic simulations...
- ...at the cost of fixing topology
- We checked feasibility by test runs in pure gauge theory
- In pure gauge systematic errors are under control.

- Results in Full QCD show signal of effective "restoration" of axial $U_A(1)$ symmetry
- We need chiral limit and check approach to T_c (ongoing) (data still to be analyzed)

Topological susceptibility and axial symm. at FT

ご清聴ありがとうございました

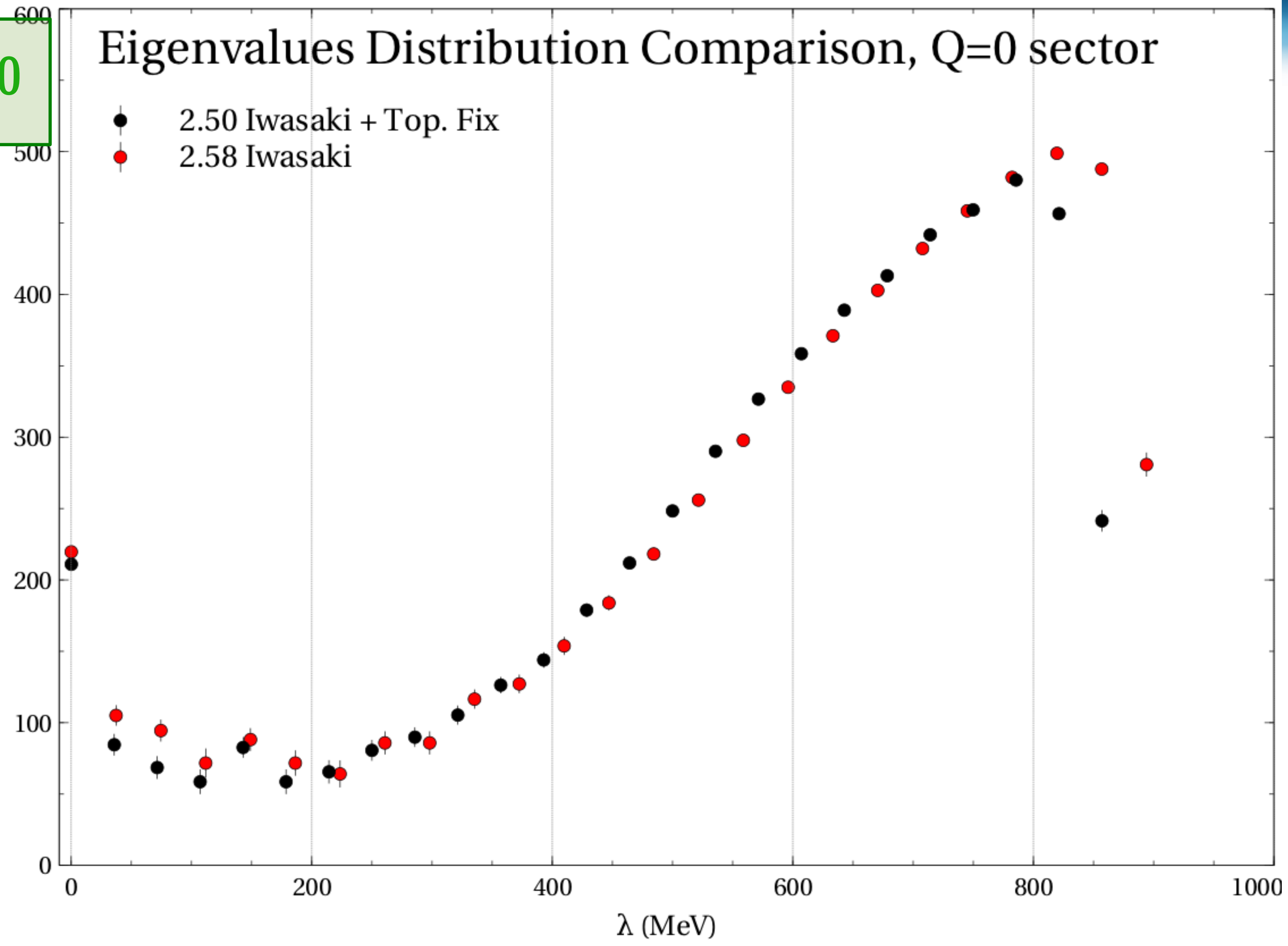


LuxRay Artistic Rendering of Lowest Eigenmode

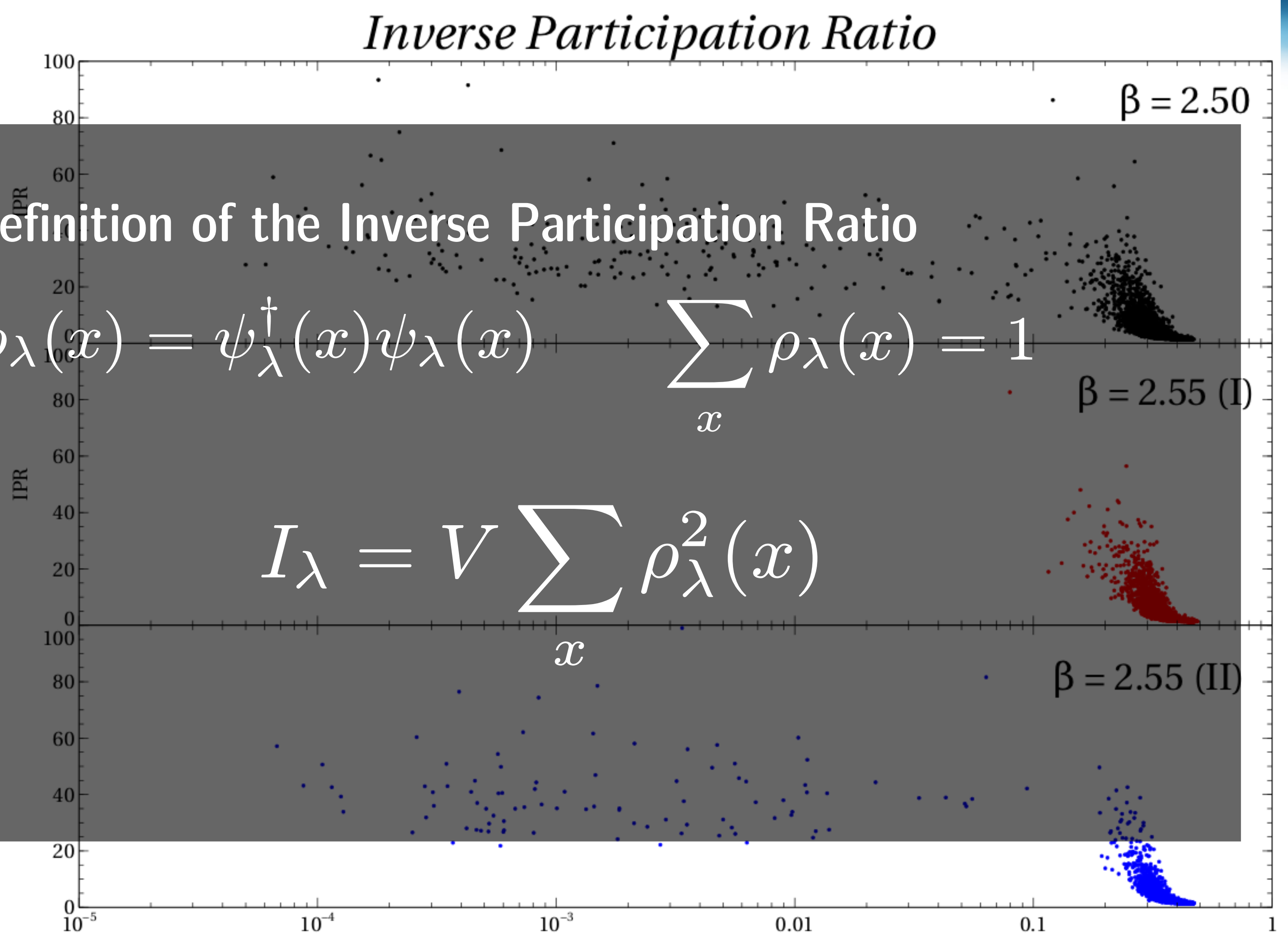
Backup slides

Topological susceptibility in pure gauge theory

$\beta = 2.50$



Topological susceptibility in pure gauge theory - VII

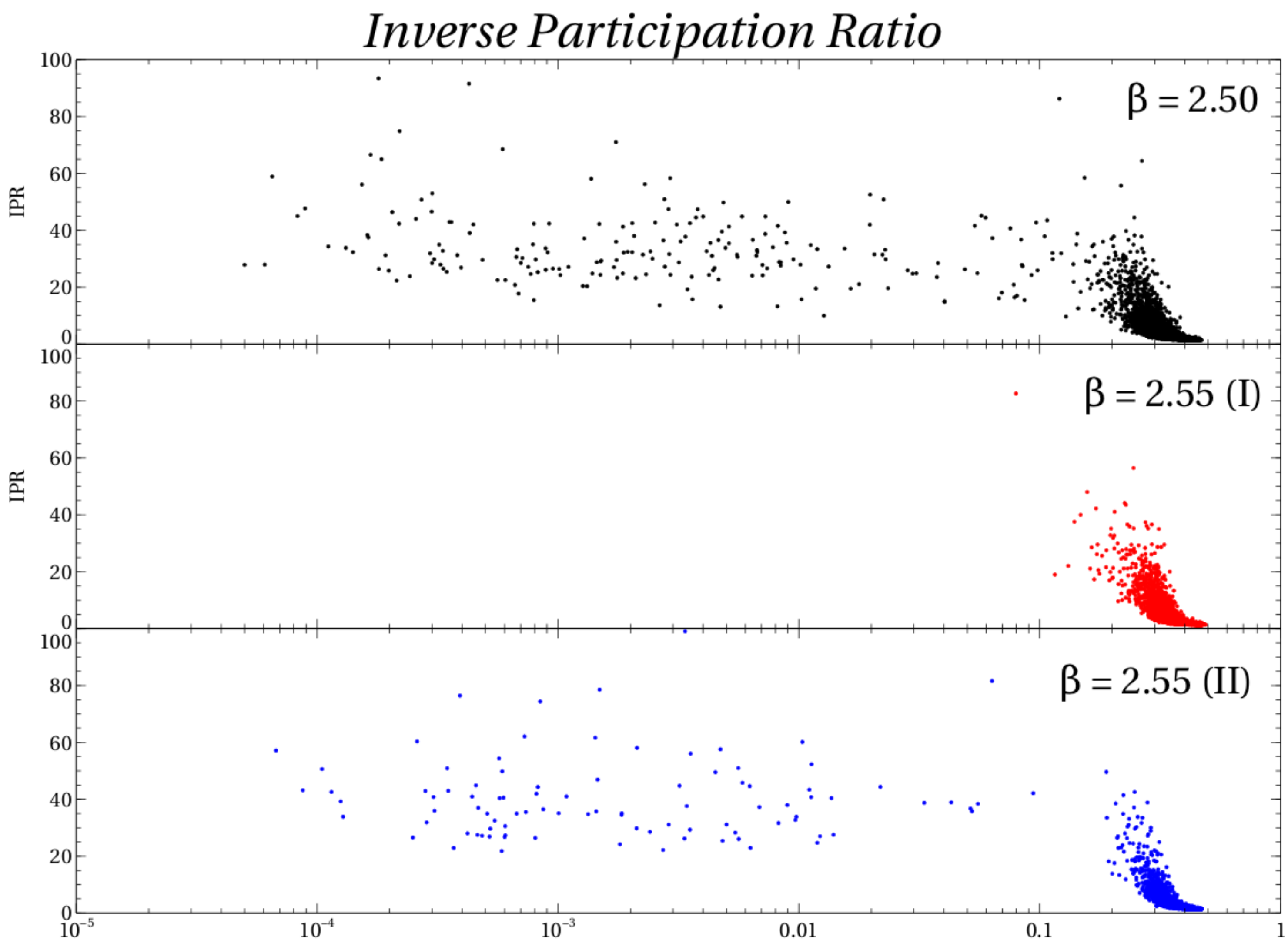


Definition of the Inverse Participation Ratio

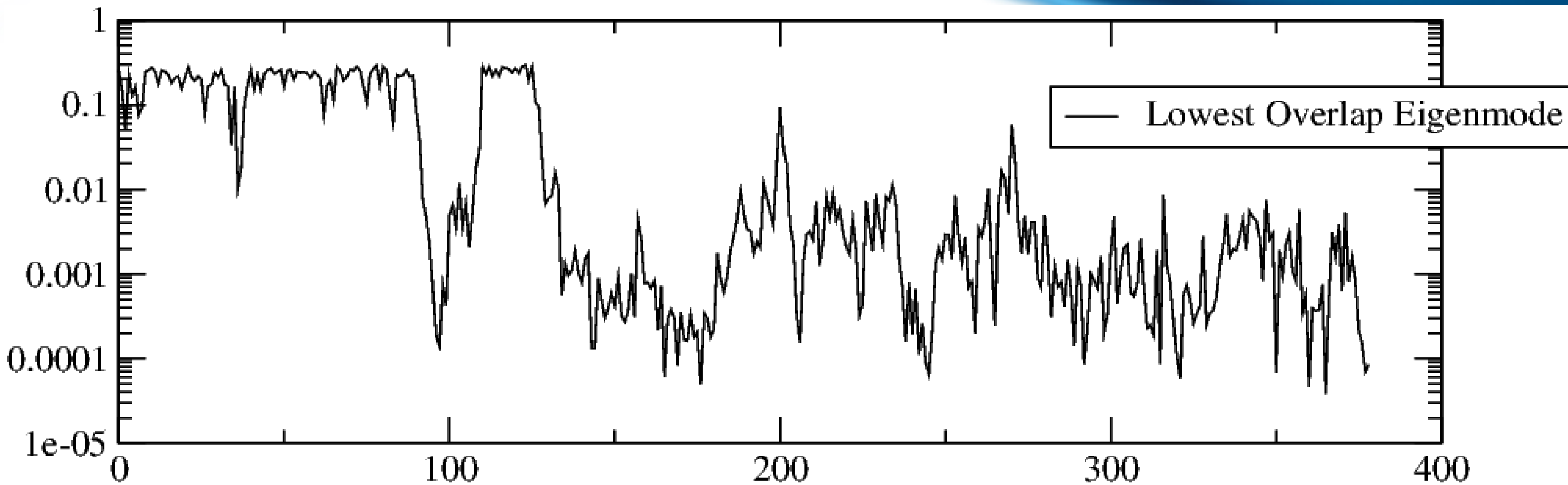
$$\rho_\lambda(x) = \psi_\lambda^\dagger(x)\psi_\lambda(x) \qquad \sum_x \rho_\lambda(x) = 1$$

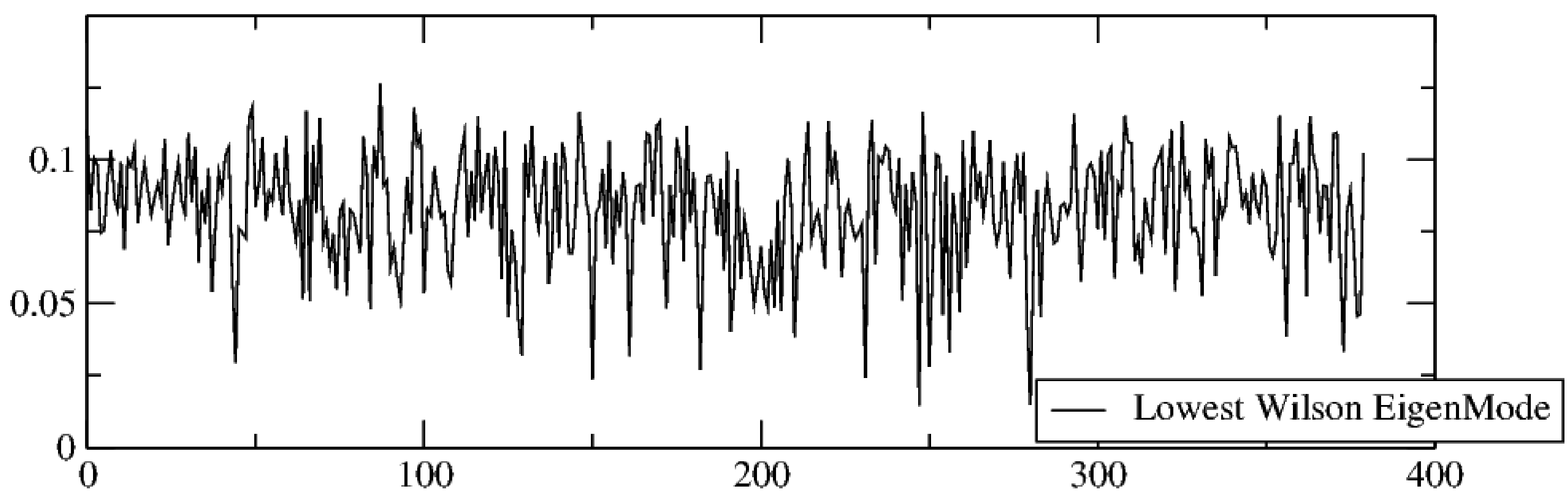
$$I_\lambda = V \sum_x \rho_\lambda^2(x)$$

Topological susceptibility in pure gauge theory - VII

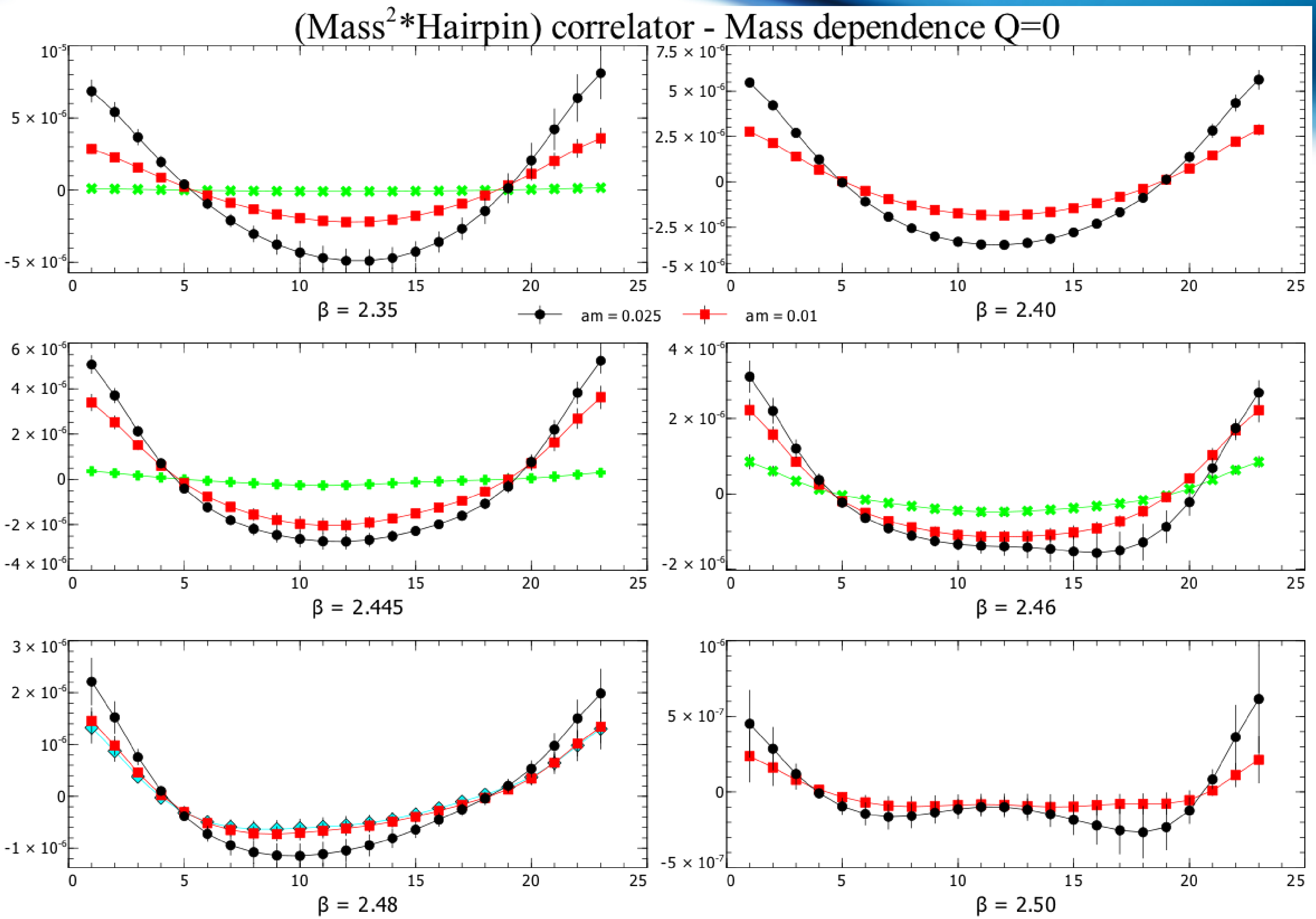


Topological susceptibility and axial symm. at FT



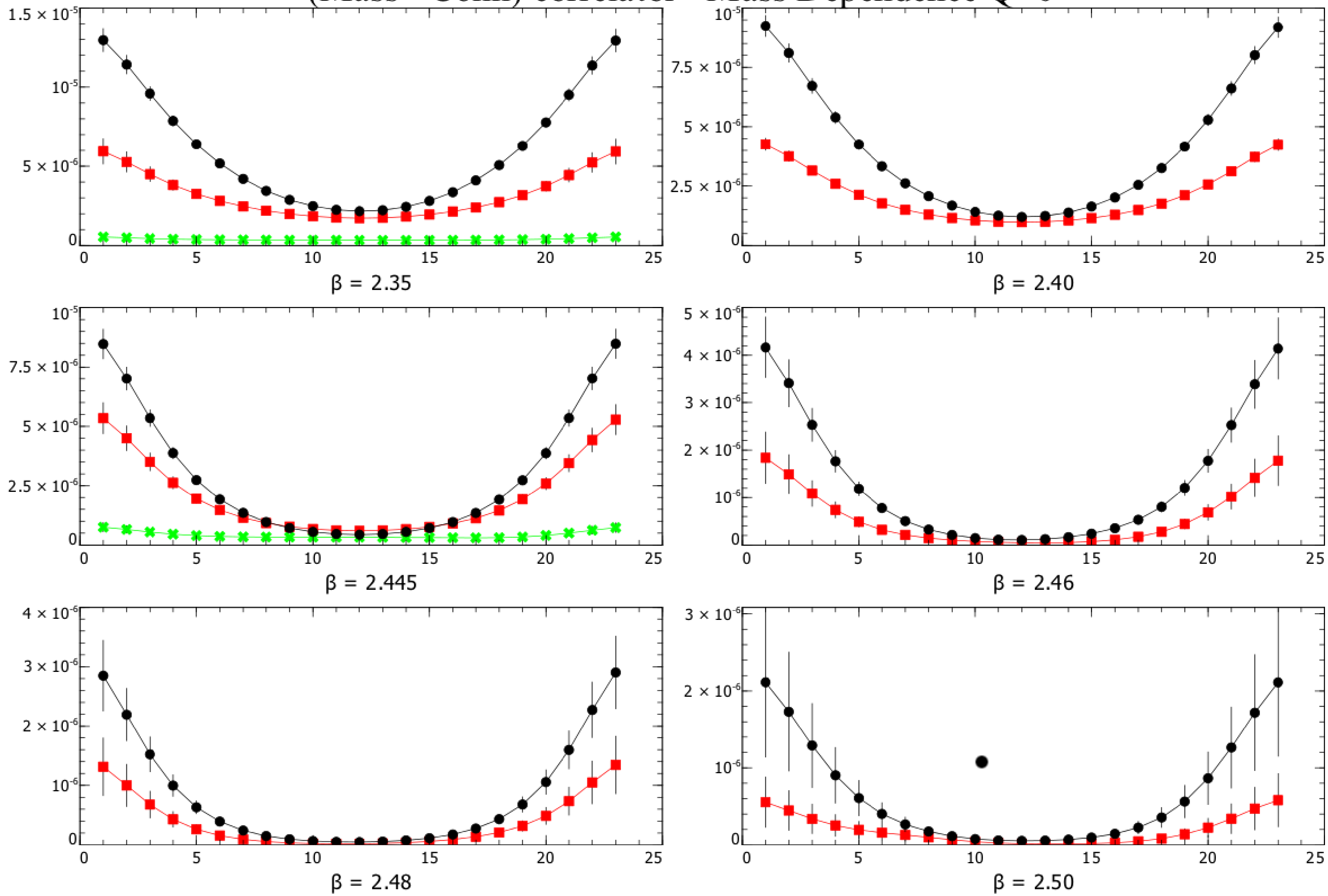


Topological susceptibility and axial symm. at FT



Topological susceptibility and axial symm. at FT

(Mass^2*Conn) correlator - Mass Dependence Q=0



Topological susceptibility and axial symm. at FT