

A Reduction Formula for Fermion Determinant



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in XQCD-J collaboration

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New Type of Fermions on Lattice,

2012/02/16 at YITP



Det (D) is a key in fd LQCD



- *Fermion determinant det D(mu)*
 - *chemical potential*
 - *sign problem*
- *Techniques in fdLQCD need the evaluation of det D(mu)*
- *Determinant calculation is heavy*
 - *Time $O(N^3)$ for $N \times N$ matrix*
 - *Memory $O(N^2)$*

 – *Memory $O(N^2)$*



An idea

- *There is a formula*
 - *to perform the temporal part of $\det D(\mu)$ by hand*
 - *and reduce the rank of the determinant*
- *The rank of the determinant is reduced, and difficulty of the determinant evaluation is relaxed*
 - *CPU time and memory is suppressed $\sim 1/Nt^2$*



References

- *Derivation for Staggered type*
 - Gibbs, PLB 172, 53 ('86). Hasenfratz & Toussaint, NPB371, 539('92).
- *Derivation for Wilson type*
 - Borici, PTP. Suppl. 153, 335 ('04). Alexandru & Wenger, PRD83, 034502 ('11). KN&AN, PRD82,094027 ('10).
- *Derivation for continuum case*
 - Adams, PRL92, 162002 ('04), PRD70, 045002 ('04).
- *Application*
 - reweighting : Glasgow group e.g., Barbour&Bell, Nucl.Phys. B372 (1992) 385-402. Fodor & Katz, JHEP 0203 (2002) 014.
 - canonical formalism : Hasenfratz & Toussaint('92). de Forcrand & Kratochivila, hep-lat/0509143.
 - Hadron spectrum. Fodor, Szabo & Toth, JHEP 0708, 092('07).
 - see also refs in KN&AN, PRD82,094027 ('10).

Basic concept of the reduction formula

- *Fermion matrix as t-t matrix*
 - *spatial (diag), temporal (n.n + b.c)*
 - *temporal hop accompanies chemical potential*
- *Performing temporal determinant by hand*
 - *rank reduces to N/Nt (memory & CPU time $\sim 1/Nt^2$.)*
 - *det D is an analytic function of μ .*
- *Reduction formula*
 - *suppress CPU time and memory $1/Nt^2$*
 - *provides an analytic form of μ -dependence*
- *Limitation*
 - *direct method is heavy*
 - *limited to small lattice size*

Focus on temporal part of quark action

$$\Delta(\mu) = I - \kappa(\text{spatial hop} + \text{clover})$$

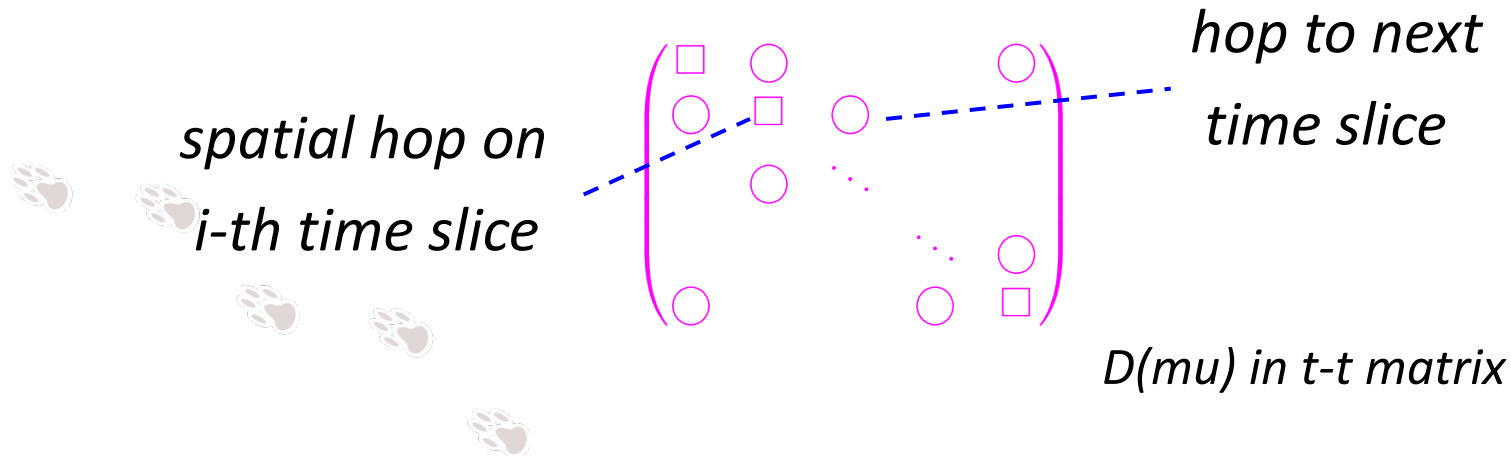
$$- \kappa [e^{\mu a} (1 - \gamma_4) U_4(x) \delta_{x', x+\hat{4}} + e^{-\mu a} (1 + \gamma_4) U_4^+(x') \delta_{x', x-\hat{4}}]$$

- *Quark action as t-t matrix*

- temporal hop accompanies chemical potential

$$\bar{\psi} \gamma^0 \psi \rightarrow \bar{\psi} \gamma^4 \psi$$

- spatial (diag), temporal (n.n + b.c)



Reduced matrix

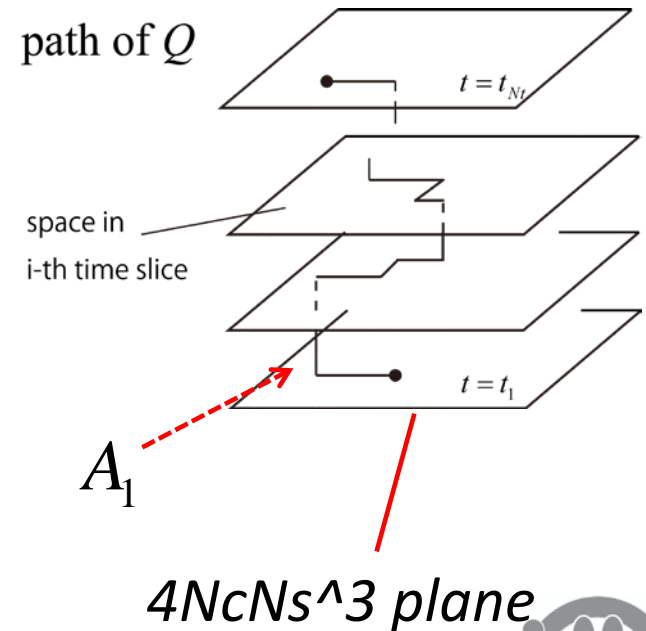
- After calculating temporal part of determinant

$$\det \Delta(\mu) = C_0 \xi^{-Nr} \det(Q + \xi),$$

$$Q \sim A_1 A_2 \cdots A_{Nt}$$

$$\xi = e^{-\mu/T} = e^{-\mu a Nt}$$

- reduced matrix Q
 - propagation from **a point at initial time** to **a point at final time**
 - rank $Nr = 4 Nc \times Ns^3$, (= full size/ Nt)
 - independent of μ
- overall factor C_0
 - spatial loops
 - independent of μ




two expressions of $\det D(\mu)$

- Calculating the eigenvalues of Q $\det(Q - \lambda I) = 0$

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \det(Q + \xi),$$

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \prod_{n=1}^{Nr} (\lambda_n + \xi)$$


$$\begin{pmatrix} \lambda_1 + \xi & & & \\ & \lambda_2 + \xi & & \\ & & \dots & \\ & & & \lambda_{Nr} + \xi \end{pmatrix}$$

$Q : (Nr, Nr)$ matrix
 $Nr = 4Nc Ns^3$

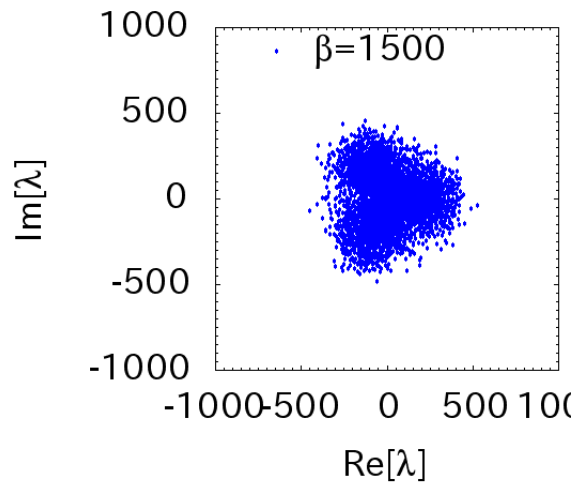
- Fugacity expansion

$$\det \Delta(\mu) \sim \sum_{n=-Nr/2}^{Nr/2} c_n \xi^n$$

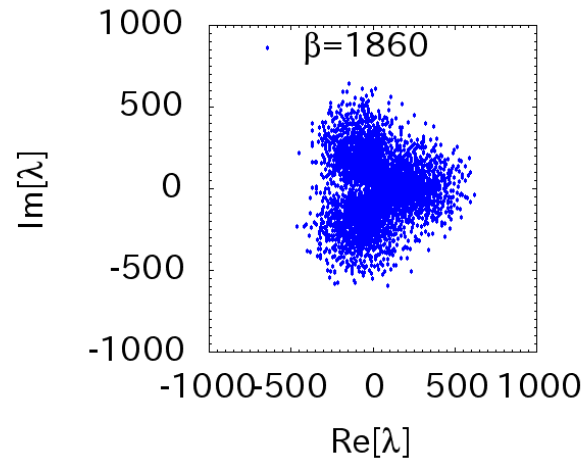
ev's or c_n determine the μ -dependence of $\det D(\mu)$

Temperature(beta) dependence 1

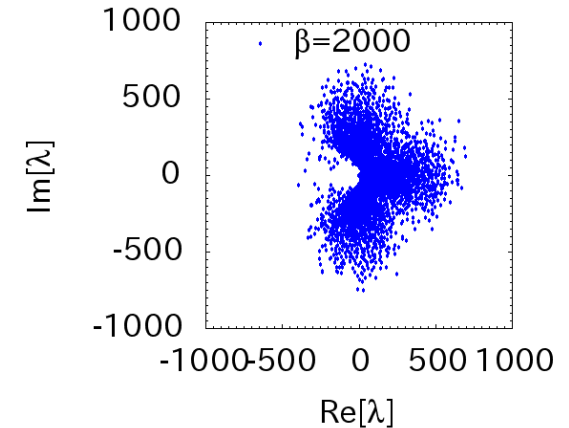
$T < T_c$



$T = T_c$



$T > T_c$



- *At low T : Z3-like distribution*
- *At high T : breaking of Z3 and broadening*
- *Blank observed at high T is related to the Polyakov loop (Alexandu&Wenger , PRD83, 034502('11)).*

Symmetry of eigenvalues

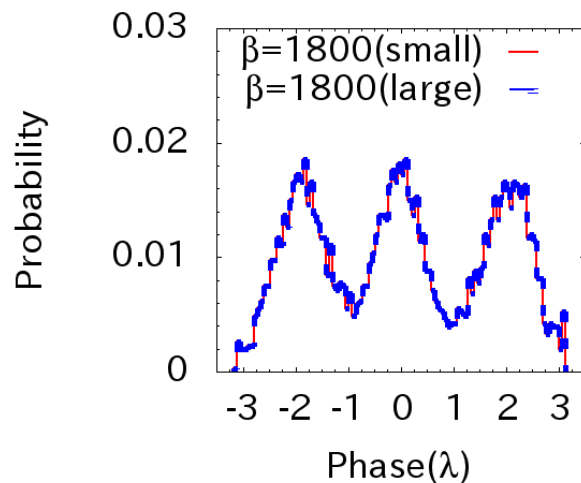
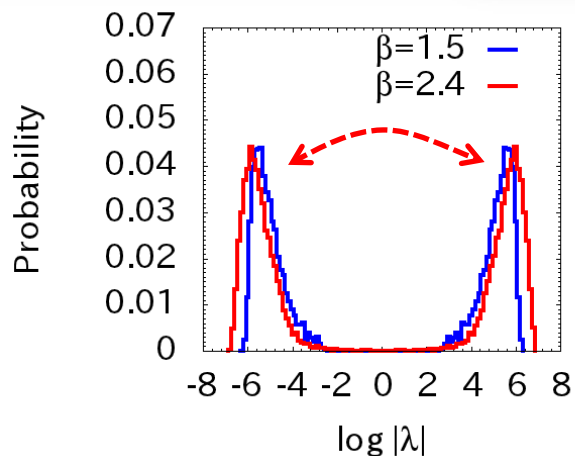
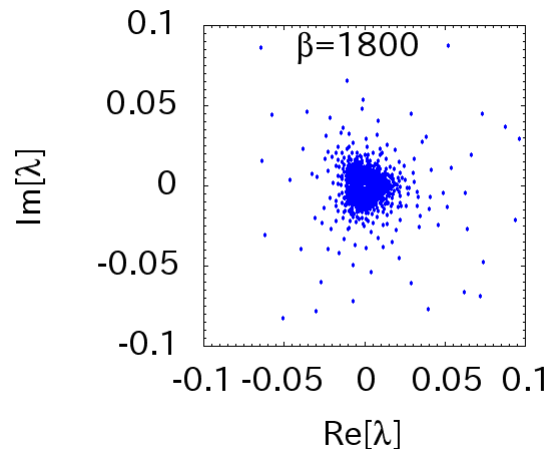
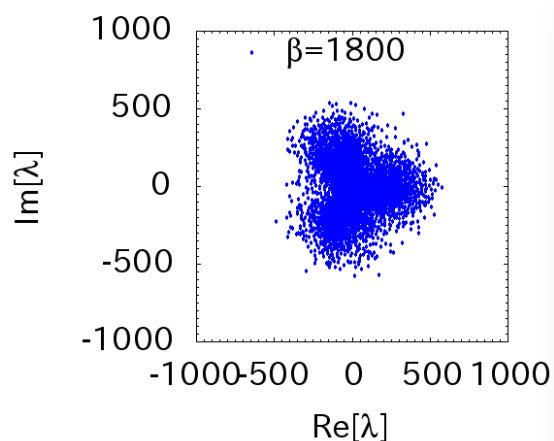
$$\lambda \leftrightarrow \frac{1}{\lambda^*}$$

clover-Wilson + RG-gauge(Nf=2)

Volume : $8^3 \times 4$

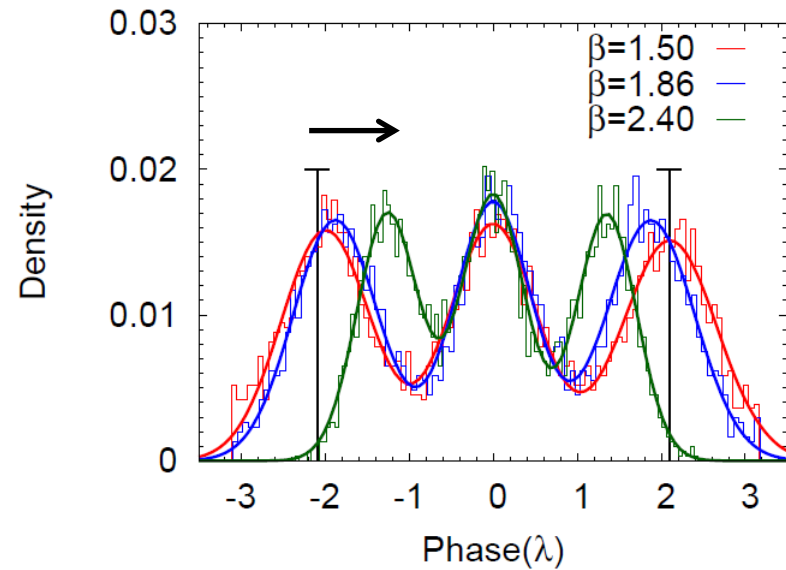
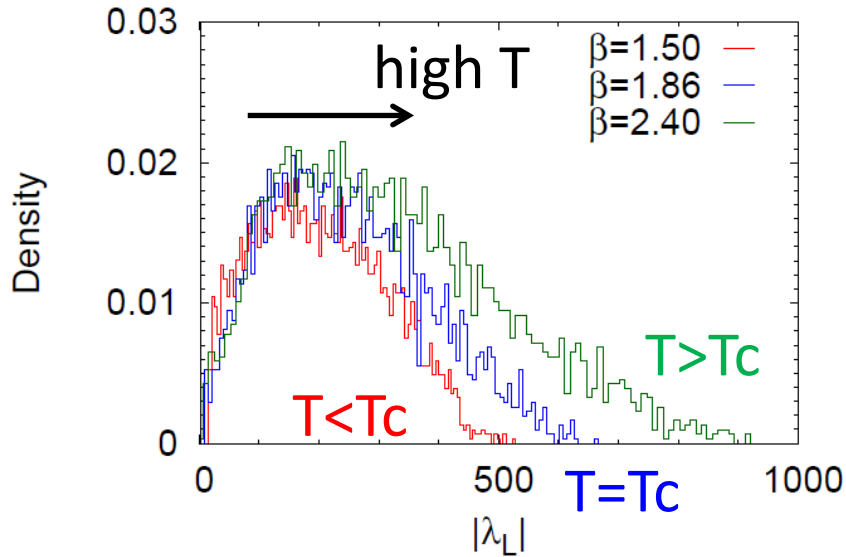
quark mass : $m_{ps}/m_V \sim 0.8$

Configurations : HMC at $\mu=0$

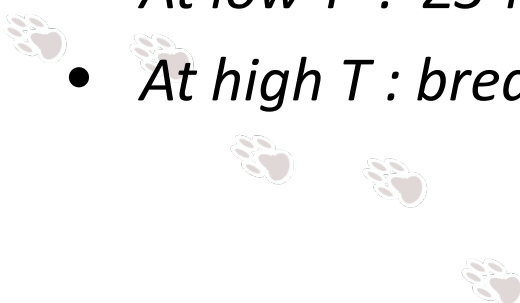


Temperature(beta) dependence 2

Histogram : Absolute (left) and phase (right)



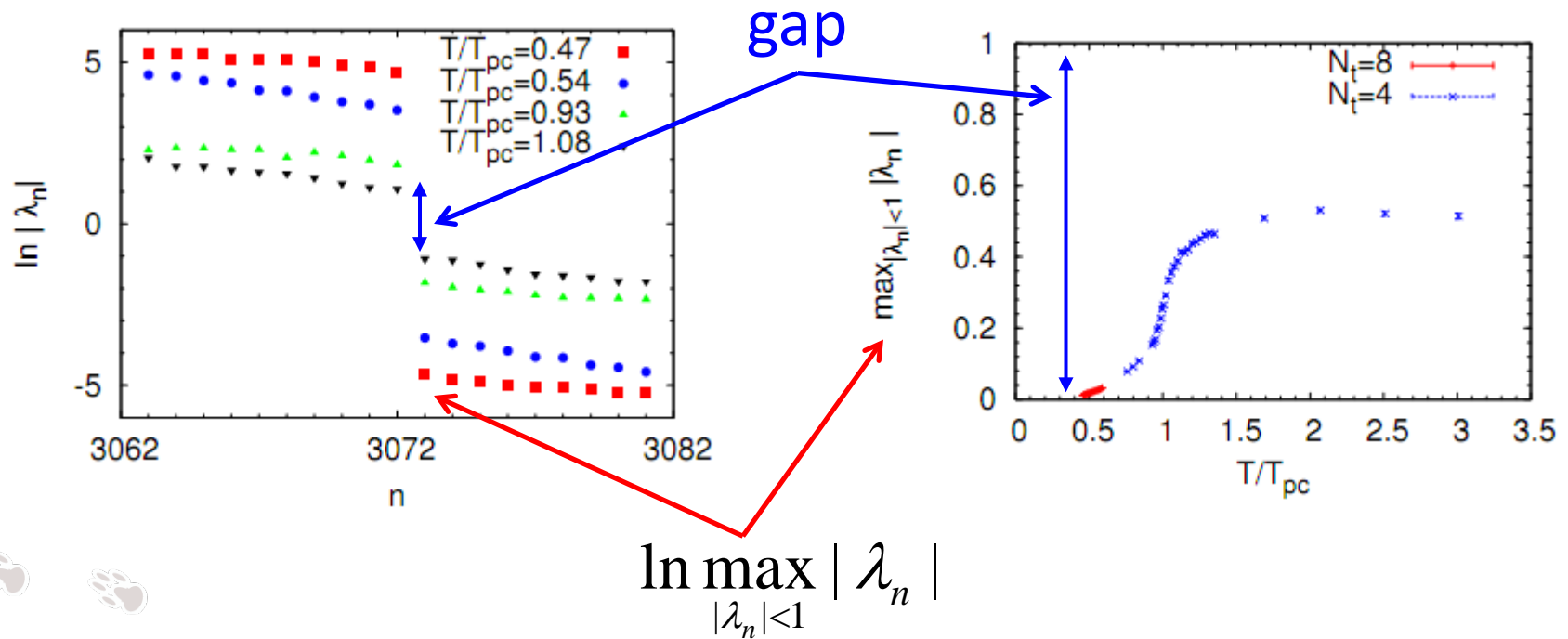
- At low T : Z3-like distribution
- At high T : breaking of Z3 and broadening



Gap of ev's near unit circle

Left : eigenvalues close to unity

Right : max. ev among the smaller half



The gap decreases as T increases.



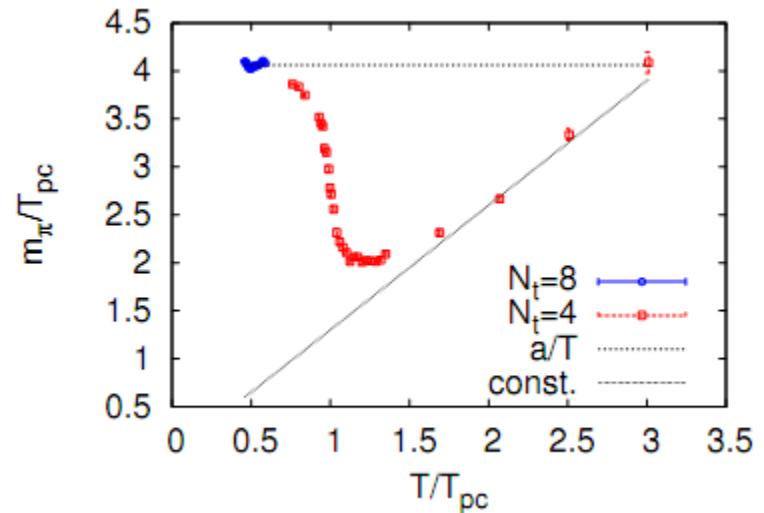
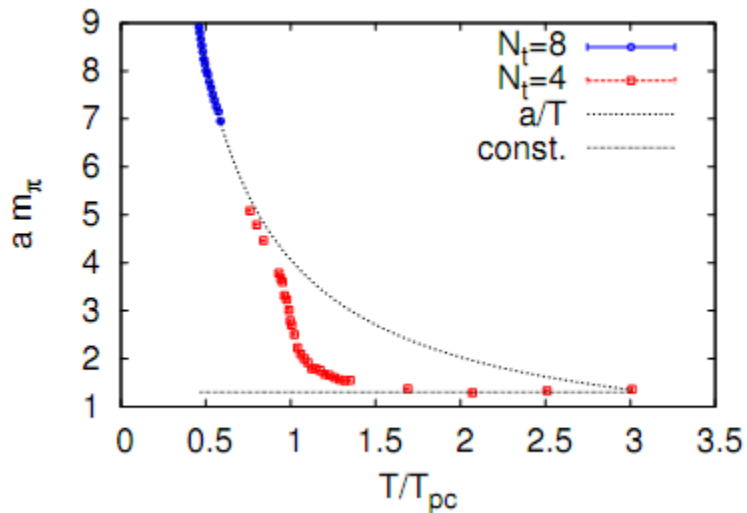
Gap is related to pion mass

$$am_\pi = -\frac{1}{Nt} \ln \max_{|\lambda_n| < 1} |\lambda_n|^2$$

Gibbs('86). Eigenvalues and m_π

See also, Fodor, Szabo, Toth ('06). Eigenvalues and hadron spectrum

Left : m_π/T . Right : m_π/T_{pc}



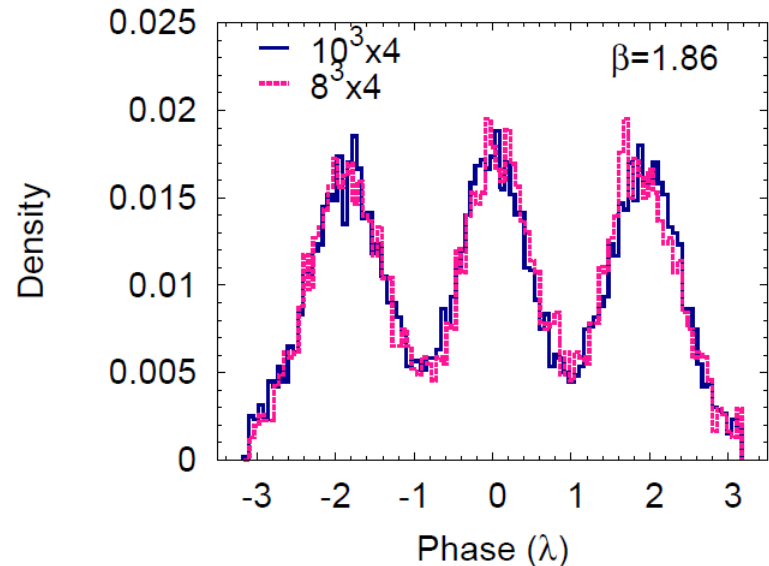
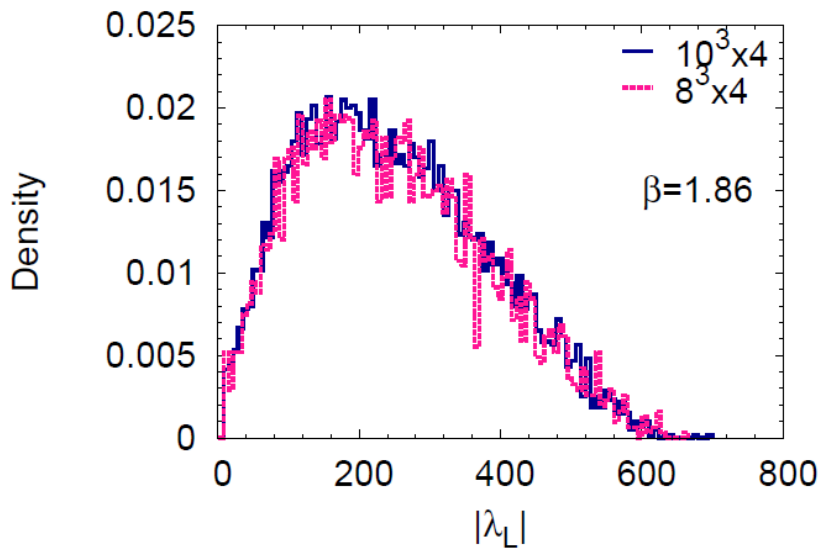
- At low T , m_π/T is well fitted with a/T , $a = 4 T_{pc}$ (m_q heavy)
- At high T , m_π approaches to a constant



Check 1. Volume dependence

Histograms : $|ev|$ (Left), $\arg(ev)$ (Right)

blue : $10^3 \times 4$, red : $8^3 \times 4$



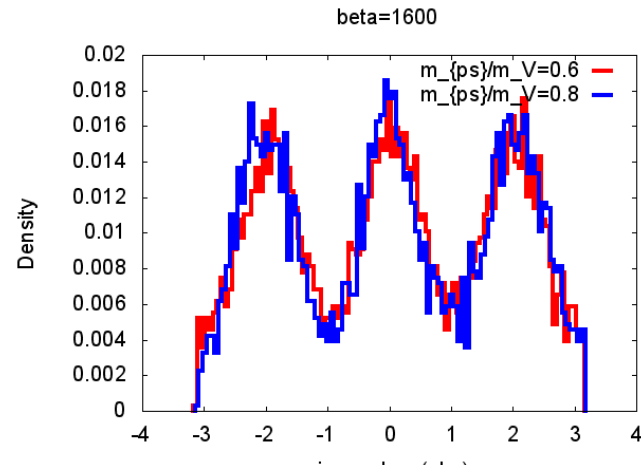
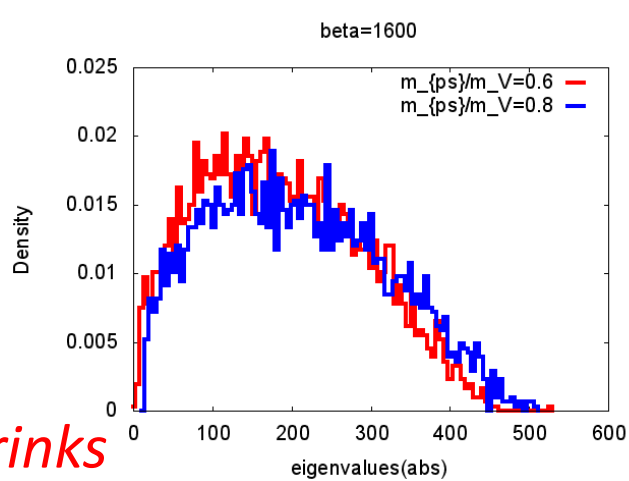
- *The small V -dependence suggests long distance propagations are small*
- *Further investigation for small quark mass is important*

Check 2. quark mass dependence

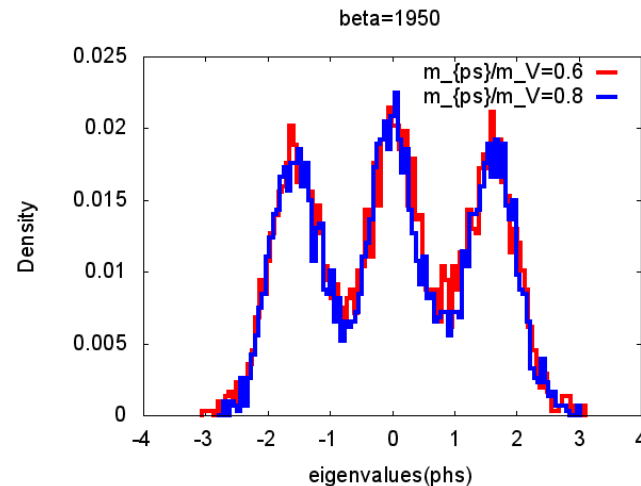
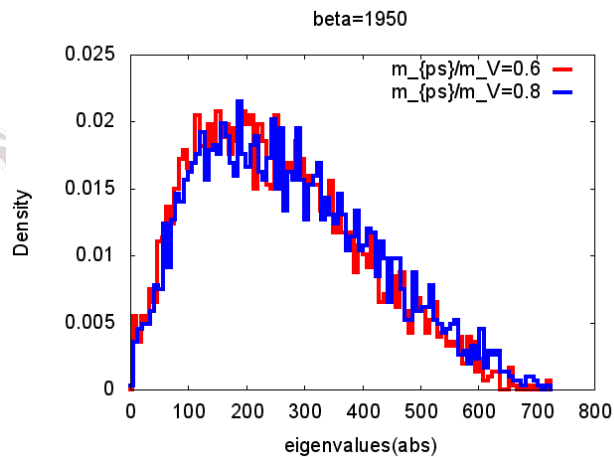
$m_{ps} / m_V = 0.6$ (red), 0.8 (blue)

Histograms : $|ev|$ (Left), $arg(ev)$ (Right)

confinement (top), deconfinement(bottom)

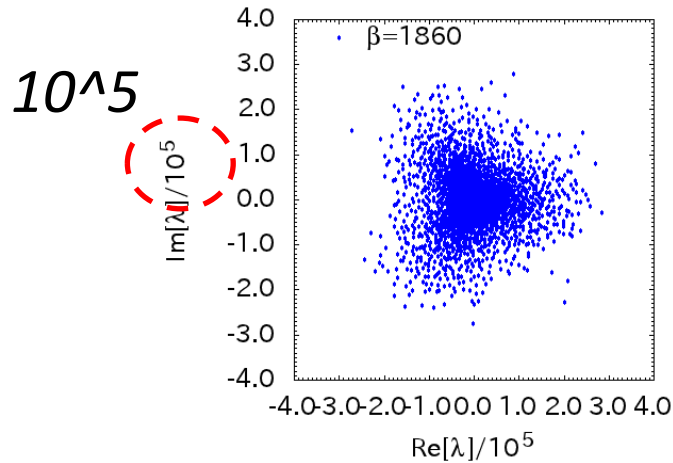


Gap shrinks

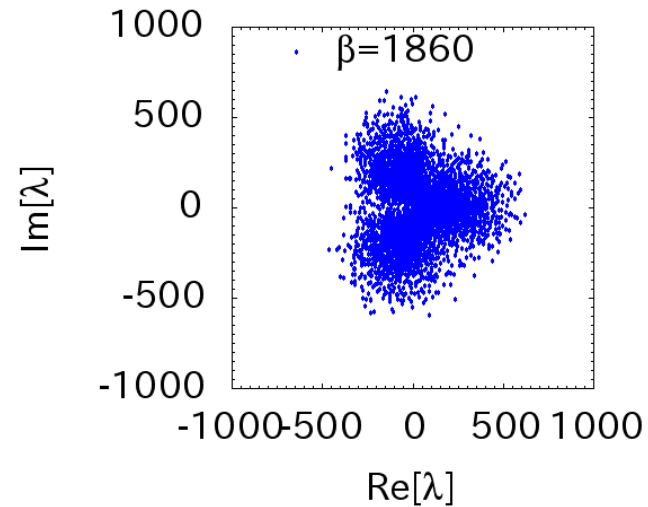


Temperature(Nt) dependence

$Nt=8 (T/T_c=0.5)$



$Nt=4 (T/T_c=1)$



- As Nt increases (T decrease)



— larger half of ev's become larger

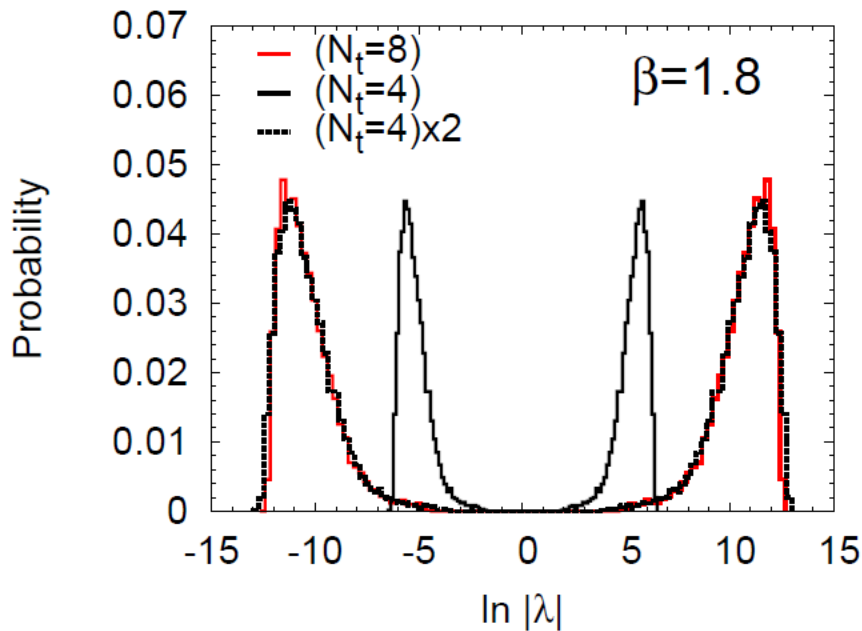
— smaller half of ev's become smaller

$\lambda \leftrightarrow \frac{1}{\lambda^*}$



Temperature(Nt) dependence 2

Power-law of ev 's with an exponent Nt



$$|\lambda_{N_t=4}|^2 = |\lambda_{N_t=8}|$$

$$|\lambda| = l^{N_t}$$

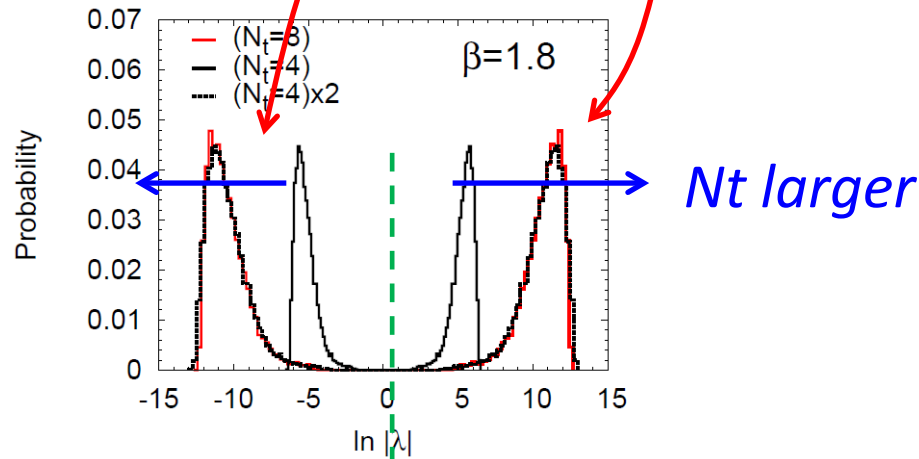
- Q is a product of Nt -matrices $Q \sim A_1 A_2 \cdots A_{N_t}$
- In equilibrium, it is expected $A_i = \bar{A} + \delta A_i$
- Power law $Q \sim \bar{A}^{N_t}$

Scale separation

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \prod_{n=1}^{Nr/2} \underline{(1/\lambda_n^* + \xi)} \prod_{n=1}^{Nr/2} \underline{(\lambda_n + \xi)}$$

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \prod_{n=1}^{Nr/2} \underline{(l_n^{*-Nt} + \xi)} \prod_{n=1}^{Nr/2} \underline{(l_n^{Nt} + \xi)}$$

$$\lambda = |l|^{Nt}$$



Nt larger

$$\xi = e^{-\mu/T} = e^{-\mu a Nt}$$



low T limit

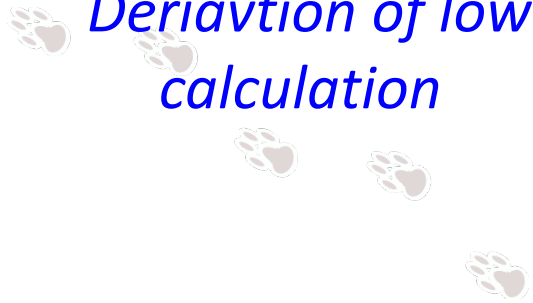
Eigenvalues at low T $T \rightarrow 0, (Nt \rightarrow \infty)$

$$\lambda_L \sim l^{Nt} \rightarrow \infty, \lambda_S \sim l^{-Nt} \rightarrow 0$$

Fugacity depends on T $\xi = e^{-\mu/T} = e^{-\mu a Nt}$

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \prod_{n=1}^{Nr/2} (l_n^{*-Nt} + \xi) \prod_{n=1}^{Nr/2} (l_n^{Nt} + \xi)$$

Derivation of low- T limit of $\det D(\mu)$ needs a careful calculation



Symmetry of Reduced matrix

$$\det \Delta(\mu) = C_0 \xi^{\xi^{-Nr/2}} \prod_{n=1}^{Nr/2} (l_n^{*-Nt} + \xi) \prod_{n=1}^{Nr/2} (l_n^{Nt} + \xi)$$

1. $\mu/T = \text{fixed}$ ($x_i = \text{const.}$)

$$\det \Delta(\mu) = C_0 \prod_{n=1}^{Nr/2} \lambda_n^L \in R$$

2. $\mu = \text{fixed}$ ($x_i \rightarrow \infty$, $T \rightarrow 0$)

a. small μ $e^{\mu a} < \bar{l}$

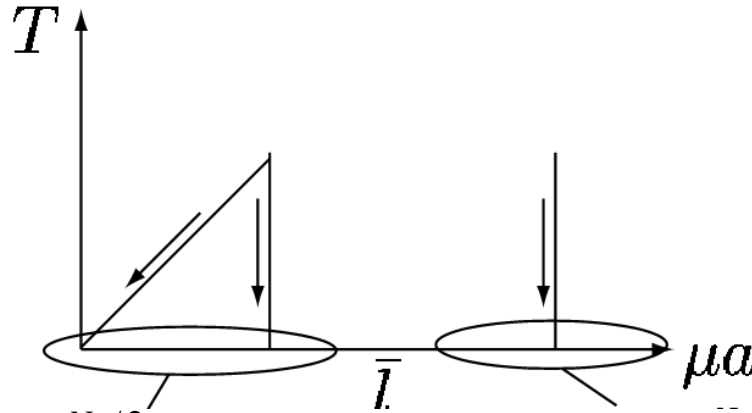
$$\det \Delta(\mu) = C_0 \prod_{n=1}^{Nr/2} \lambda_n^L \in R$$

b. large μ $e^{\mu a} > \bar{l}$

$$\det \Delta(\mu) = C_0 \xi^{\xi^{-Nr/2}} \det Q \in R$$



Interpretation



$$\xi^{-Nr/2} = \exp(2N_c N_f N_s^3)$$

$$\det \Delta(\mu) = C_0 \prod_{n=1}^{Nr/2} \lambda_n^L \in R$$

$$\det \Delta(\mu) = C_0 \xi^{-Nr/2} \det Q \in R$$

$\langle n \rangle = 0$

$\langle n \rangle = 2 N_c N_f$ [lattice unit]
(all states are occupied)

"Silver Blaze problem"

Independence of $\det D$ on μ at low T : Cohen,
PRL91,222001 ('03)Adams, PRD70, 045002 ('04).



Lee-Yang zero

- *Lee-Yang theorem : origin of thermodynamical singularities, i.e., phase transition*

$$F = -T \ln Z(\xi), \quad \frac{\partial F}{\partial \mu} \infty - T \frac{1}{Z} \frac{\partial Z}{\partial \mu}, \quad \xi = \exp(-\mu/T)$$

$$Z(\xi) = 0 \Rightarrow F' \rightarrow \infty$$



Canonical formalism of QCD

- *Grand partition function is expanded in powers of fugacity*

$$Z(\mu) = \sum_{n=-N_q}^{N_q} Z_n \xi^n$$

- *N_q (= Max. # of the quark) is given by (d.o.f x V)*
- *Negative power terms come from the anti-quark*



Canonical formalism of QCD

Preliminary

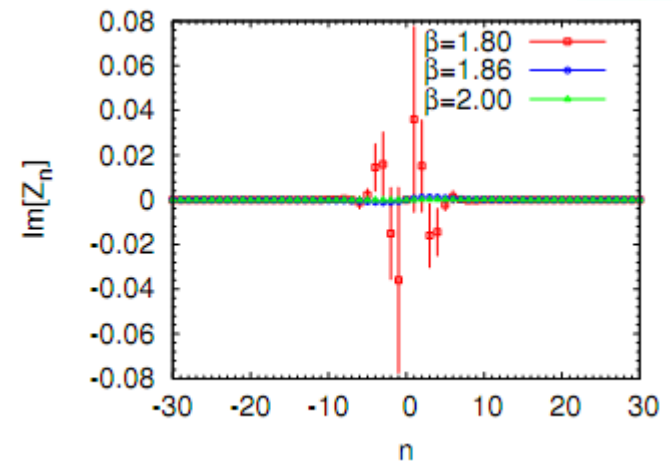
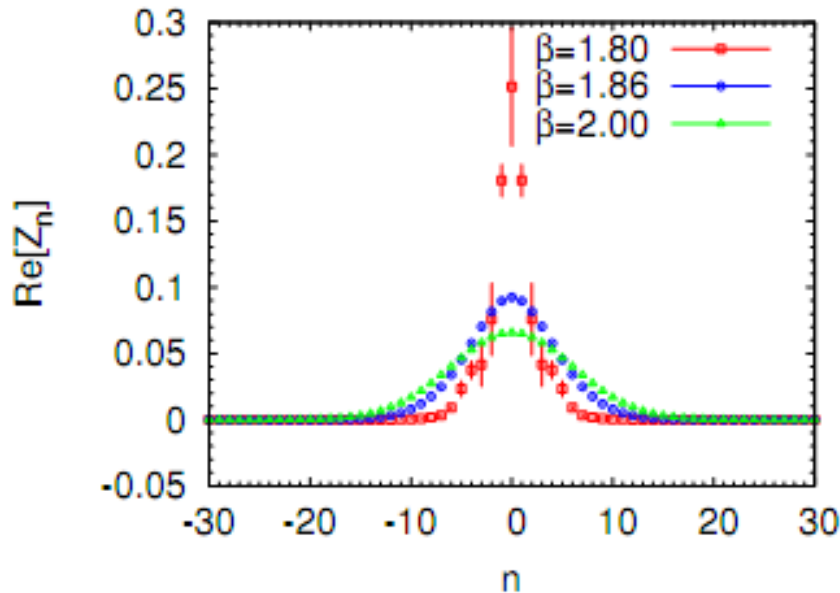
- We employ the Glasgow method
 - Reduction formula + reweighting

$$Z_n = \left\langle \frac{c_n}{(\det \Delta(0))^2} \right\rangle_0 \quad (\det \Delta(\mu))^2 = \sum_{n=-Nr}^{Nr} c_n \xi^n$$

- Approaches to improve overlap for canonical and LY
 - Fodor & Katz, *Phys.Lett. B534* (2002) 87-92
 - de Forcrand & Kratochivila, *hep-lat/0509143*
 - Ejiri & Yoneyama, *arXiv:0911.2257 [hep-lat]*
 - Li, Alexandru & Liu, *Phys.Rev. D84* (2011) 071503



Canonical partition function



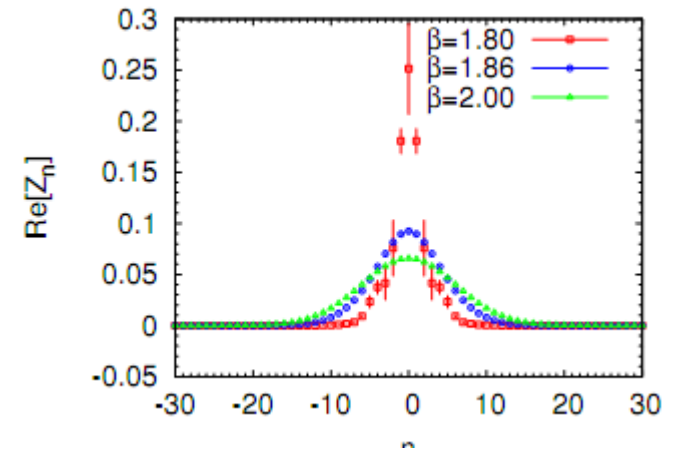
*large phase fluctuation
at low T*

- *Zn decreases exponentially*
 - *broadening at high T : increase of effective d.o.f*
- *Zn (mod(n,3) !=3) does not vanish even in confinement phase*
 - *because of importance sampling at mu=0*
 - *including all Z3 sectors by hand leads to cancelation*

Lee-Yang Zero 1

- In the vicinity of $\mu = 0$, $Z(\mu)$ is dominated by terms near $n=0$*

$$Z(\mu) = \sum_{|n| \leq N'} Z_n \xi^n + \sum_{|n| > N'} Z_n \xi^n$$



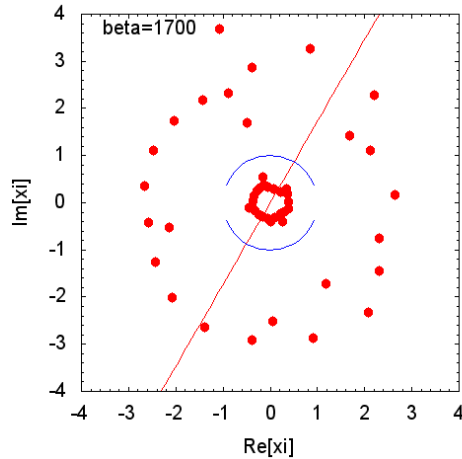
- The behavior of the root distribution near $x_i=1$ can be studied by the first terms.*
- (Density depends on N' ($f(x) = 1 + x + x^2 + \dots + x^N$))*
- This approximation should not be used for the study of the order of the phase.)*



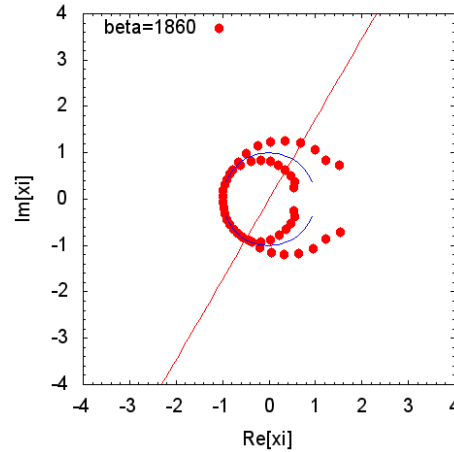
Lee-Yang Zero

$$Z(\mu) = \sum_{|n| \leq N'} Z_n \xi^n = (\xi - r_1)(\xi - r_2) \cdots$$

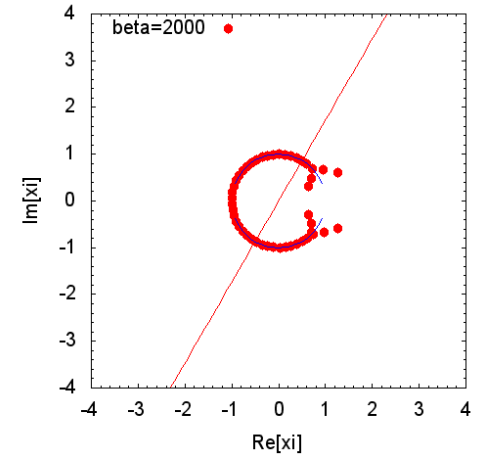
$T < T_c$



$T = T_c$



$T > T_c$

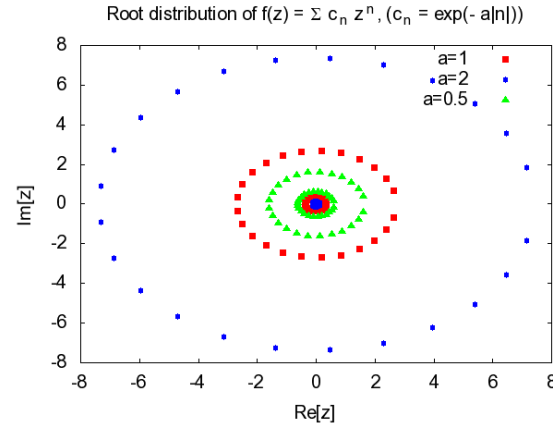
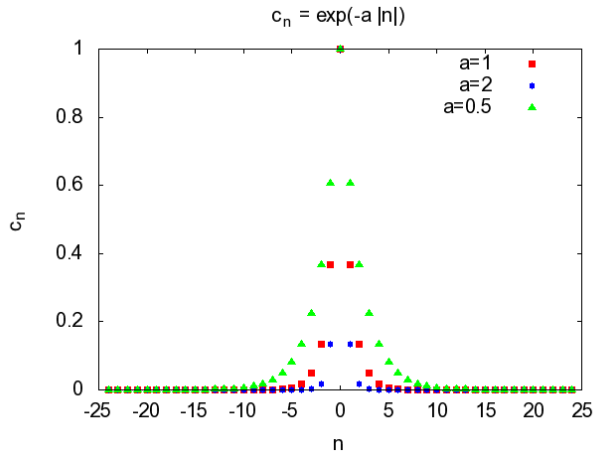


- *Root distribution shows two different behavior*
 - *at low T : two circle ($Z_n = \exp(-a |n|)$)*
 - *at high T : unit circle + 3 branches ($Z_n = \exp(-a n^2)$)*
 - *branch point is located at near RW endpoint*

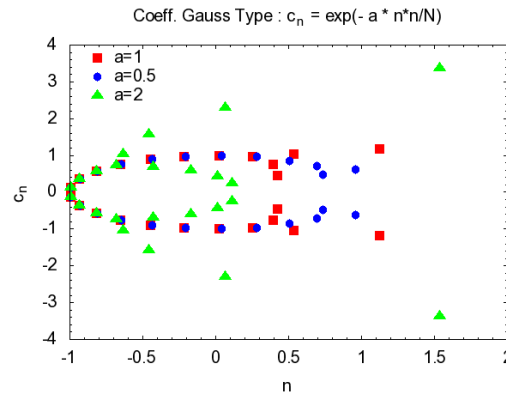
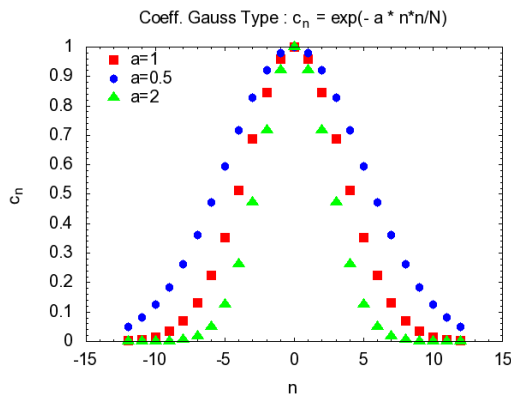


Examples for Polynomial's Root

$$c_n = \exp(-a|n|)$$

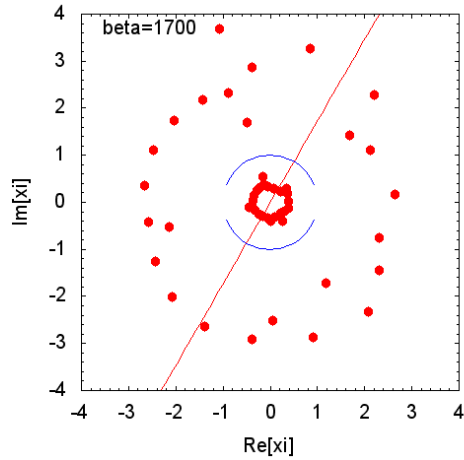


$$c_n = \exp(-an^2)$$

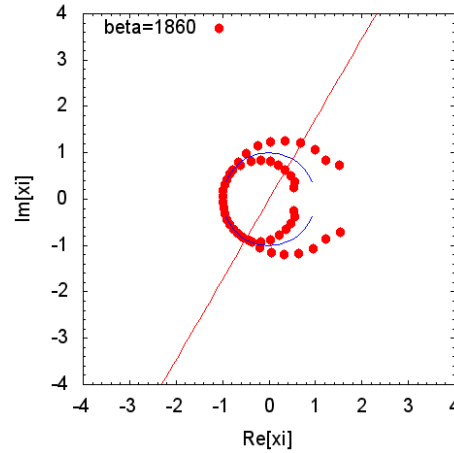


Lee-Yang Zero 2

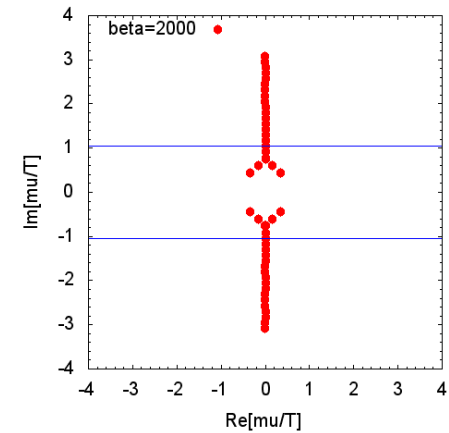
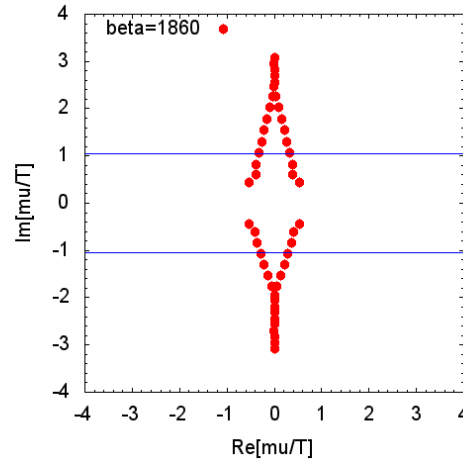
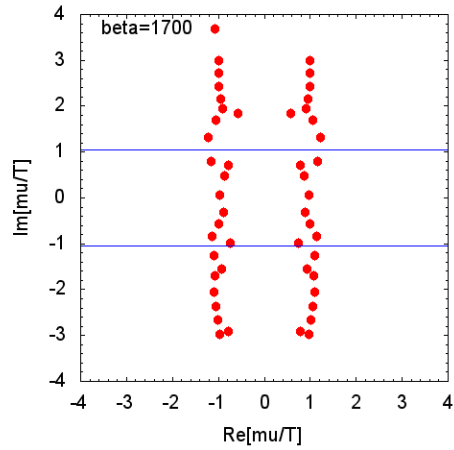
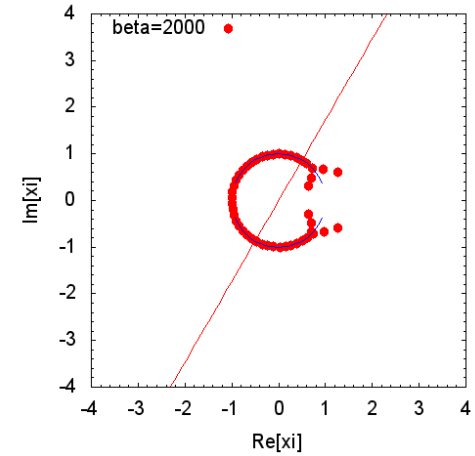
$T < T_c$

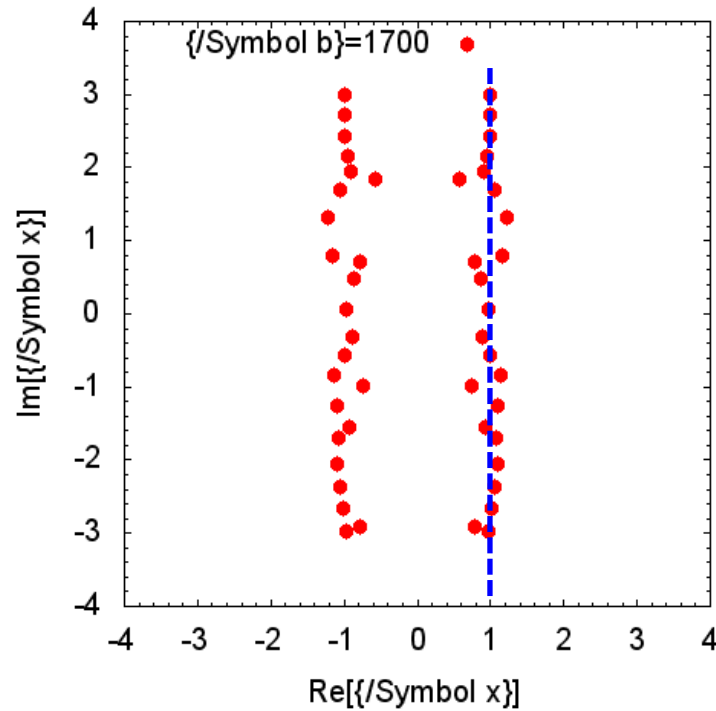


$T = T_c$



$T > T_c$





- *LYzeros seems to approach to real axis near $\mu/T = 1$*
- *Further investigations are needed for*
 - *overlap problem & sign problem*
 - *thermodynamic limit*

Summary

- *Reduction formula*
 - *relax the difficulty to calculate fermion determinant*
 - *provides the analytic expression of $\det D(\mu)$*
 - *although it is still heavy, the formula is sometimes useful to study finite density systems.*
- *Properties of reduced matrix*
 - *angular distribution is related to Z_3*
 - *gap is related to chiral symmetry breaking*
 - *the whole eigenvalues are important in overlap & sign problem*
- *In this talk, we have studied*
 - *numerical properties of transfer matrix*
 - *low- T limit of $\det D(\mu)$*
 - *canonical partition function & Lee-Yang zero*