基研研究会2011 非平衡系の物理―ミクロとマクロの架け橋―

湯川記念館パナソニックホール 2011年8月19日

液晶超薄膜に微弱外場で誘起される 非平衡構造

多辺由佳

早稲田大学 先進理工学術院 物理学及応用物理学専攻



散逸構造

希薄系		凝縮系		
(要素間の直接相互作用:弱)		(要素間の直接相互作用:本質的)		
入力	出力	入力	出力	
熱	ベナール対流	電気信号	脳の情報伝達系	
電場	液体の対流ロール			
	液晶の電気流体 力学的不安定性	化学反応	生体内の代謝	
光	レーザー発振	光	→ 液晶単分子膜 の配向波	
化学反応	BZ反応	化学ポテン シャル	→ キラル液晶分子 集団歳差運動	



・構造・光学異方性 → 表示パネルへの応用 ・高い流動性(棒状低分子の非化学結合)





多彩な相転移と高次構造 次元性、トポロジー、ゆらぎ



分子構造とマクロ物性



構成分子の構造 → 液晶の相系列、配向弾性、粘性



液晶の非平衡ダイナミクス

例① Electrohydrodynamic Instability

Williams domain



DC or low frequency High frequency

j ingi noquonoj

 $http://www.uni-magdeburg.de/anp/projects/pattern_form/index.htm$

液晶の非平衡ダイナミクス 例② Lehmann effects



under temperature gradient by Lehmann (1900) under DC electric field by Madhusudana & Pratibha (1987)

液晶の非平衡ダイナミクス

例③ Thermomechanical or photo-mechanical effects in nematic elastomers

under ∇T



birubber

under illumination

By T. Ikeda

monodomain

By H. Finkelman

分子の運動 vs マクロなダイナミクス

バルクの液晶実験で、両者の関係を追うのは容易ではない



二次元液晶を対象とした分子運動とマクロダイナミクスとの相関

身の回りの二次元液晶







Two examples:

- Photo-induced orientational waves in azo LC monolayer
- Coherent molecular precession in chiral LC monolayer

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2-1. 光誘起配向波

2-2. 物質流透過によるキラル液晶の一方向回転2-3. DC電場による液晶バブルの非平衡ダイナミクス

3. まとめ

Dissipative structures in 2D LCs (monolayers) 1. Photo-induced orientational waves

Photo-reactive molecule: azobenzene



Weakly-excited LC monolayers of azobenzene derivatives



Generation and preparation of LC orientational wave



100µm

Image under polarizing microscope

Excitation polarization

ion
ation
$$I = h^2 \sin^2 \phi (\cos \phi - f)^2$$

• Oscillation of ϕ



Wave properties

- •Periodic rotation of molecular azimuth
- •Propagate in the asymmetric film
- •Wave velocity $\propto \sqrt{\text{excitation power}}$
- •90° rotation of excitation polarization →
 180° inversion of wave direction





180° wave switching by 90° rotation of excitation polarization



Polarization direction of excitation

Y. Tabe & H. Yokoyama; New J. Phys. 5, 65.1-65.11 (2003)

Theory for illuminated LC Langmuir monolayer

T. Okuzono, Y. Tabe & H. Yokoyama; Phys. Rev. E, 69 (2004)

$$F = \int d\mathbf{r} \left[\frac{1}{2} |\nabla c_i|^2 - \frac{\tau}{2} |\mathbf{c}|^2 + \frac{u}{4} |\mathbf{c}|^4 - \lambda \psi \nabla \cdot \mathbf{c} + \frac{D}{2} |\nabla \psi|^2 + \frac{\chi}{2} \psi^2 \right]$$

$$\mathbf{c}(\mathbf{r}, t) : \text{ projection of the local molecular direction}$$

$$\psi(\mathbf{r}, t) : \text{ local concentration difference between trans and c is isomers}$$

$$\frac{\partial \mathbf{c}}{\partial t} = -\frac{\delta F}{\delta \mathbf{c}} + \mathbf{f}, \quad (\mathbf{f} \equiv -\gamma_2 \frac{1 - \psi}{1 + \psi} \mathbf{c})$$

$$\frac{\partial \psi}{\partial t} = M \nabla^2 \frac{\delta F}{\delta \psi} + g,$$

$$g \equiv -(\gamma_1 + \gamma_2) \psi - (\gamma_1 - \gamma_2)$$

$$\frac{\text{anisotropic reaction in LC monolayer:}}{\gamma_1 = g_1(\mathbf{E} \cdot \mathbf{c})^2}$$

$$\gamma_1 = \psi_1(\mathbf{E} \cdot \mathbf{c})^2$$

 γ_2 : independent of electric field polarization

Simulation of wave switching

T. Okuzono, Y. Tabe & H. Yokoyama; Phys. Rev. E, 69 (2004)



$$g_1 = \gamma_2 = 0.05, \tau = u = 2, \lambda = 1$$

Possible mechanism of wave propagation



The coupling between the anisotropic absorption & splay distortion & density variation carries the orientational waves

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Dissipative structures in 2D LCs

2. Chiral molecular precession driven by transport currents





Chiral LC monolayer on pure glycerol

Example 1

(S)-OPOB on pure glycerol

C8H17-(

Clariant Japan Ltd.







Depolarized polarizing light microscopy: CW or CCW ?



<u>Crossed polarizers (P&S):</u> $I(\phi) = h \sin^{2} \phi(\cos \phi - f)^{2}$ <u>Slightly depolarized (P+2° & S+Δ°):</u> $I(\phi) = h \left[\sin^{2} \phi(\cos \phi - f)^{2} - g \sin \phi(\cos \phi - f) \right]$ $I(\phi) \neq I(-\phi)$ $\theta_{i} = 30^{\circ}, \Delta P = 2^{\circ}, n_{g} = 1.4746; \beta = 10^{\circ}, \varepsilon_{\perp} = 2, \varepsilon_{\parallel} = 2.1$ $\longrightarrow f = 2.5, g = 1.1$







azimuthal angle ϕ



LC motor is driven by water transfer



molecular rotational speed ∞ vapor pressure difference



Phenomenological equations



Phenomenological equation:

$$\frac{d\phi}{dt} = \frac{1}{\gamma}\tau + b \ \upsilon_m \ \Delta P$$

$$J = b \tau + \frac{1}{\eta} \upsilon_m \Delta P$$

D_m: molecular volume of water in vapor
γ: rotational viscosity of LC director
η: water mobility through the film
b: cross-coupling coefficient (depend on chirality)



when
$$\tau = 0: \frac{d\phi}{dt} = b \ v_m \ \Delta P \quad b = 10^{24} \ [\text{N}^{-1}\text{m}^{-1}\text{s}^{-1}]$$

 $\Delta P_w/P_0 = P_v - P_s$

P_v: actual water vapor pressure

P_s: saturated water vapor pressure



Molecular Dynamics Simulation Microscopic interaction between LC molecules & current molecules:

Chiral LC Molecules (detailed atomic model)









Torque on an isolated LC molecule along the molecular axis per single water collision



	OP	D 608		
	R-isomer	S-isomer	1-000	
Mean value	3.23E-19	-4.97E-19	-1.56E-20	
Standard deviation	2.91E-20	3.56E-19	1.67E-20	

Axial rotation and precession of chiral molecules in monolayer state



M. Yoneya, Y. Tabe and H. Yokoyama J. Phys. Chem.B *114* (25), 8320 (2010)

- hydrogens are treated as united atoms CH_n
- Gaussian03 HF/6-31g* RESP fitted atomic charges
- GROMACS force field (with extension)
 - Δt =4fs (with fixed bonds), cut-off = 0.9/1.8nm
- z-position restraints on dummy atoms at the mol. center

$$V(z_i) = k_{pr} (z_i - z_0)^4$$

Directions of axial rotation under Ar flow

Dependence on molecular chirality and flow direction



<u>Possible mechanism of flow-driven molecular</u> <u>precession in chiral monolayers:</u>

Collision of flow molecules with the chiral propeller



axial rotation of each molecule

develops into the collective precession

Molecular chiral propeller is essential in monolayers.

What happens in thick films with macroscopic helix?

Microscopic torque vs macroscopic torque





under temperature gradient by Lehmann (1900)





Characteristics of transferred materials

	water	ammonia	methanol	ethanol	toluene
Vapor pressure $(25^{\circ}C)$	3.2 kPa	800 kPa	16.5 kPa	7.4 kPa	2.9 kPa
Molecular weight	18	17	32	46	92
Solubility into LC (calculated, a.u.)	~0.006	~0.007	0.86	0.93	0.85
Dielectric constant	81	22	32	24	2.3

Critical thickness for torque inversion under water transfer

Pitch (um)	Spontaneous polarization (nC/cm ²)	Critical thickness (layer)
6	58	3~4
20	39	7~8
110	42	NA
20	9	NA

Competition between torque caused by helix and by molecular propeller

✓ Chemical interaction between the molecules is small
✓ Helix has short pitch

Microscopic torque < Macroscopic torque

FELIX013: director changes its in-plane direction by 0.06 degree each layer

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Bubble deformation under DC field

DC

Electrode



Smectic bubble

- sample
 - 8CB
- condition
 - r_0 (initial bubble radius): 1.30-3.27mm
 - *d* (electrode distance):2.0-4.4mm
 - DC voltage:0-2000V

Static deformation induced by low electric field



Initial bubble radius $r_0 = 2.22$ mm Distance between electrodes d = 3.00mm

Bubble is pulled by the electrostatic force

Competition between electrostatic force and surface tension

• Force per unit area of bubble surface

$$f_r = \left(p_{in} + \frac{1}{2} \rho E_r \right) - \left(p_0 + 2\sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) \right)$$

expand shrink

 $\rho(z)$: induced charge density $E_r(z)$: electric field normal to the bubble surface σ : surface tension $R_1(z)$, $R_2(z)$: curvature radius

- Assumptions:
 - Bubble shape is ellipsoidal
 - $E_{r_top} \approx V/(d-h)$
 - Constant molar number inside the bubble



Static deformation in the equilibrium

•
$$f_r = 0 \rightarrow \frac{4\sigma}{R_{top}} = \frac{1}{2} \varepsilon_0 E_r^2 + 2\sigma \left(\frac{1}{R_{10}} + \frac{1}{R_{20}}\right)$$

Normalization
 $x(x_d - x)^2 = \frac{\varepsilon_0 V^2}{24\sigma r_0} \quad x = r_V/r_0 - 1$
Slope gives:
 $\sigma = 0.025 \text{ [N/m]}$
agree with the reported value $\sigma = 0.024 \text{ [N/m]}$
By R. Stannarius and C. Cramer (1997)
 $x = r_V/r_0 - 1$
 $x_d = d/r_0 - 1$
 $x_d = d/r_0 - 1$
 $x_d = d/r_0 - 1$
 $x_d = 0.024 \text{ [N/m]}$

Derive the surface tension and the induced charge

Under high voltage: Non-equilibrium states

• In the equilibrium state, force must be balanced:

$$\frac{x(x_{d} - x)^{2}}{f(x)} = \frac{\varepsilon_{0}V^{2}}{\frac{24\sigma r_{0}}{E(V)}} \qquad \begin{array}{l} x = r_{V}/r_{0}-1 \\ x_{d} = d/r_{0}-1 \end{array}$$

• When $V > V_{\text{th}}$, *x* has no solution except for $x > x_d$

 \Rightarrow must touch the electrode

Threshold radius & voltage: $x_{\text{th}} = x_d/3 \iff r_{\text{th}} = (2r_0 + d)/3$ $E_{\text{th}} = \frac{4}{27} x_d^3 \iff V_{\text{th}} \approx \sqrt{\frac{96\sigma r_0}{27\varepsilon_0} x_d^3}$

DC電場下での液晶バブルに誘起される不安定性 Oscillation under high DC field

Mass (charge) transport -> interface instability • Oscillation period : ~60ms



Dynamics of bubble deformation

• Equation of motion of bubble top:

$$\rho_{LC} \frac{d^2 r_{\nu}}{dt^2} + \frac{v}{r_0} \frac{dr_V}{dt} = \frac{1}{2} \varepsilon_0 E_r^2 - 2\sigma \left(\frac{2}{R_{top}} - \frac{1}{R_{10}} - \frac{1}{R_{20}}\right)$$



• Normalized equation of motion :

$$A\frac{dx}{d\tau} = \frac{E(1 - e^{-\tau})}{(x_d - x)^2} - x \qquad A = \frac{v r_0}{12\sigma RC} <<1 \qquad \tau = \frac{t}{RC} \\ E = \frac{\varepsilon_0 V^2}{24\sigma r_0}$$



ニ次元液晶(単分子膜 or 分子数層からなる自己保持膜)に光・ ポテンシャル流・電場により誘起される非平衡構造

✓液晶の強い分子間相互作用が主役 ✓二次元系を対象とすることで、分子の運動の偏りがそのままマ クロダイナミクスに反映

✓現象論では説明できるが、分子の運動とマクロダイナミクスの間には大きな時空間的隔たりがあり、これを埋めることが課題