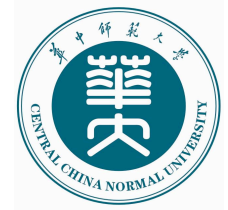
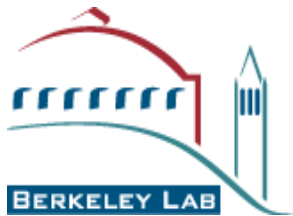


YIPQS Symposium: **Perspective in Theoretical Physics**  
-- from quark-hadron sciences to unification of theoretical physics

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# **Properties of Dense Matter in High-energy Heavy-ion Collisions**



Xin-Nian Wang  
Central China Normal University  
&  
Lawrence Berkeley National Laboratory

February 6-8, 2012

# Phases of Matter

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火

(gas)

水

(liquid)

土

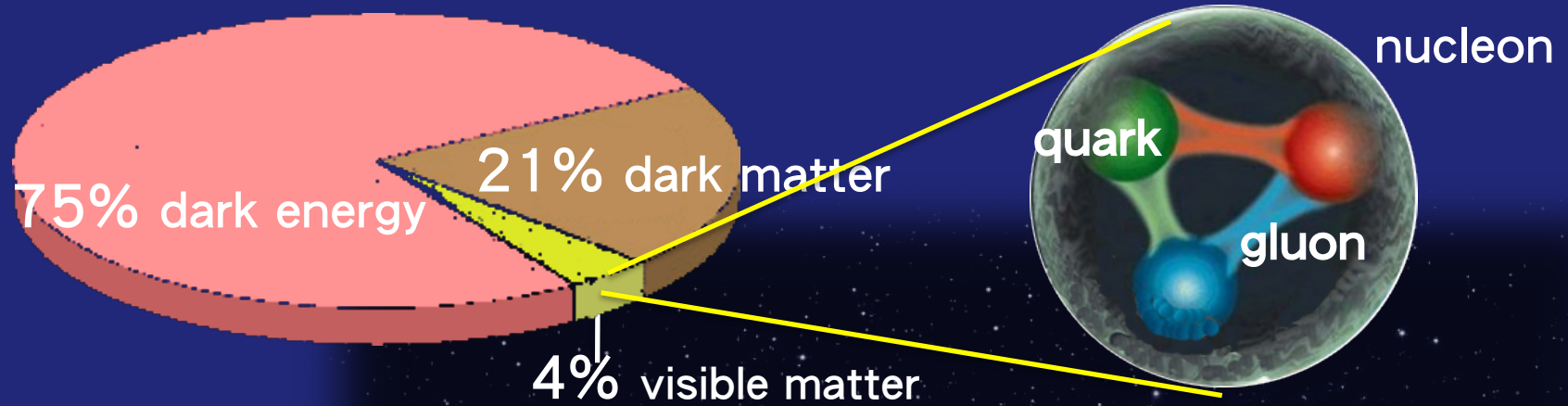
(solid)

Bose-Einstein condensate, fermionic condensate,  
superfluids, supersolids, paramagnetic,  
ferromagnetic, liquid crystals, ...

Hadron properties

Quark-gluon Plasma (QGP)

# Composition of Matter in our Universe



$$m_N \bar{u}(p)u(p) = \langle N(p) | \frac{\alpha_s}{2\pi} F_{\mu\nu}^a F_a^{\mu\nu} + \sum_q m_q \bar{\psi}_q \psi_q | N(p) \rangle$$



**95% of nucleon's mass from gluon sector**

# QCD Theory

---

$$L_{QCD} = \sum_{f=1}^{n_f} \bar{\psi} \gamma_{\mu} (i\partial^{\mu} - gA_a^{\mu} \frac{\lambda_a}{2} - m)\psi - \frac{1}{4} \sum_a F_a^{\mu\nu} F_{a,\mu\nu}$$

- SU(3) gauge symmetry (non-Abelian)

- Asymptotic freedom at short distance
- Confinement at long distance

$$\alpha_s(Q^2) = \frac{4\pi / (11 - \frac{2}{3}n_f)}{\ln(Q^2 / \Lambda_{QCD}^2)}$$

- Chiral symmetry and its spontaneous breaking

- Goldstone boson and chiral condensate

$$\langle \bar{\psi}\psi \rangle \neq 0$$

- Scale and  $U_A(1)$  anomaly
- ....

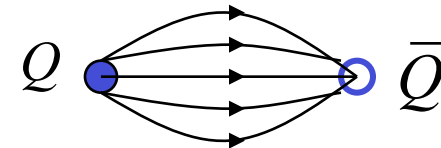
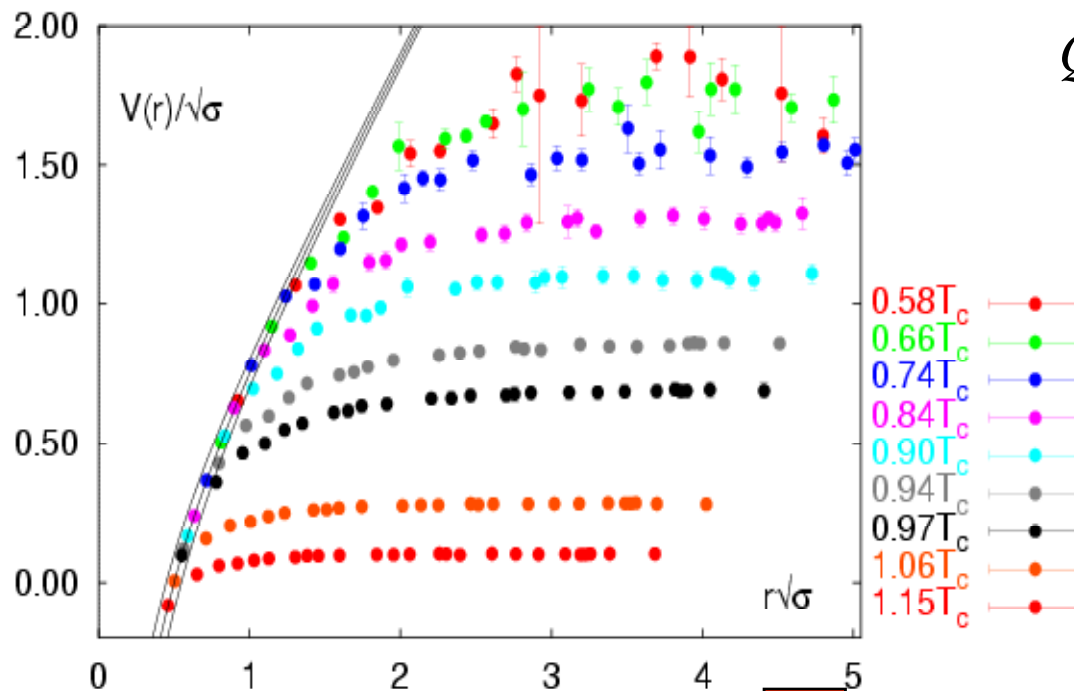


# Confinement-deconfinement

SU(3) non-Abelian gauge interaction  $\rightarrow$  confinement

$$V(r) = -\frac{\alpha}{r} + \sigma r$$

Heavy quark potential:



Karsch, Laermann  
and Peikert 2001

**Dissociation of J/ $\Psi$**

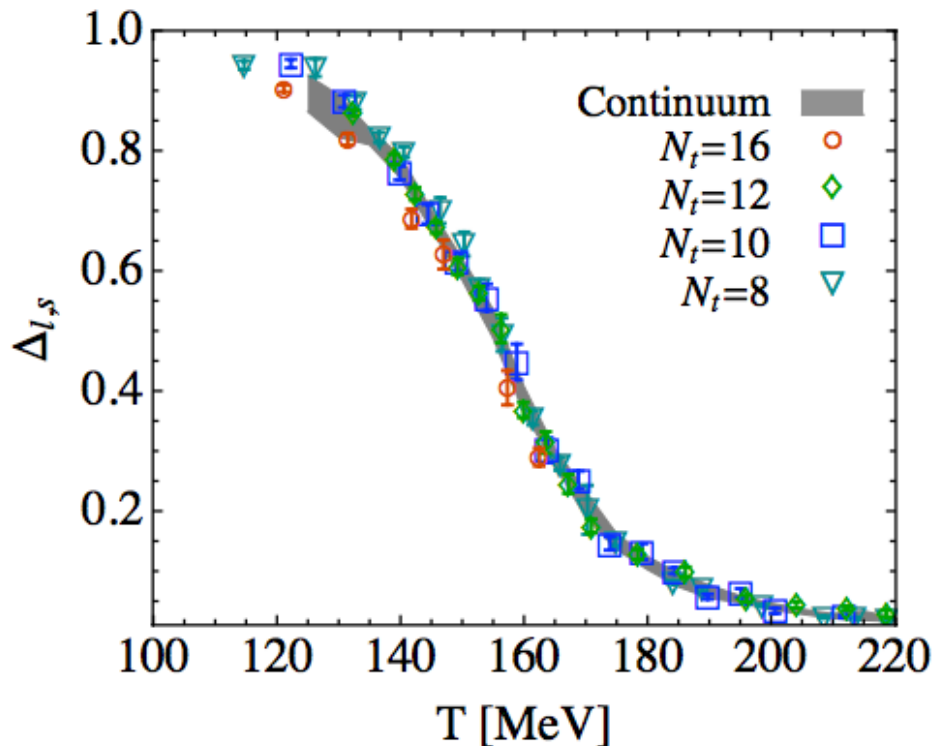
# Chiral Symmetry Restoration

$$T=0: \quad \langle \bar{\psi}\psi \rangle \neq 0$$

$$SU(3)_L \otimes SU(3)_R \rightarrow SU_{L+R}(3)$$

Spontaneously broken:

Goldstone bosons ( $\pi, K, \eta$ )



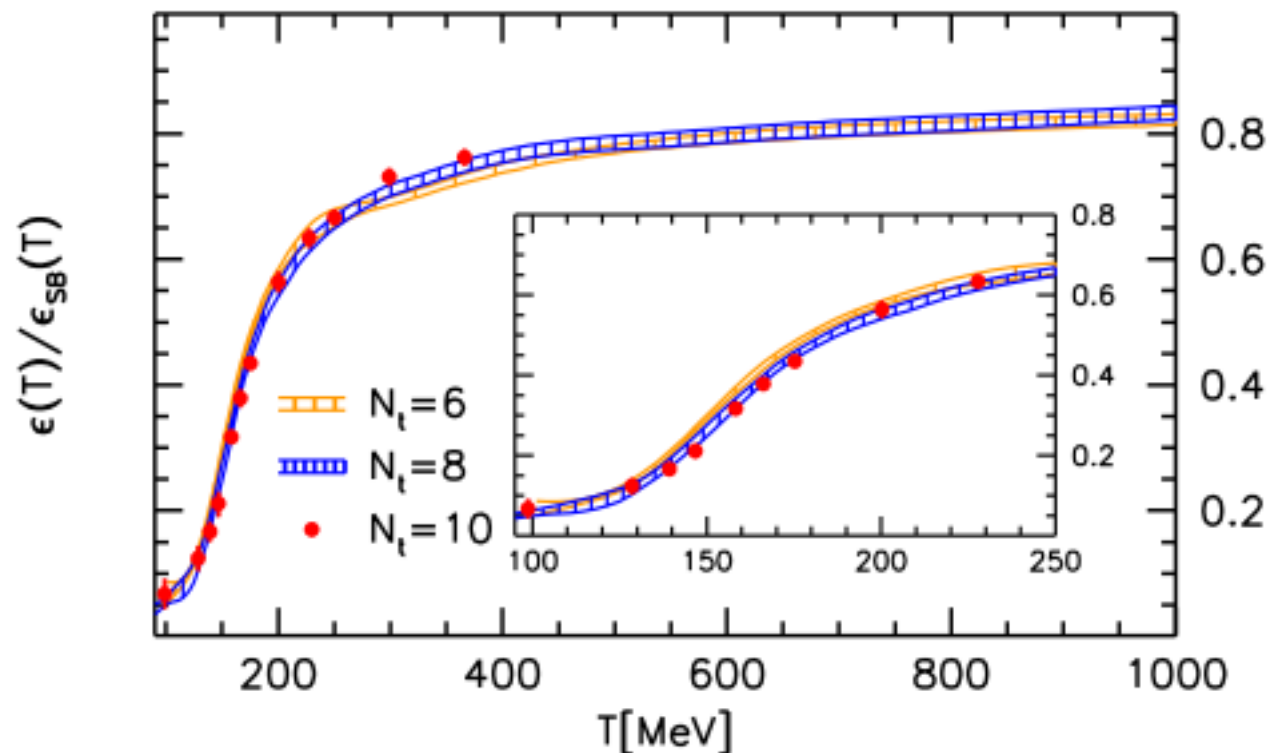
Wuppertal-Budapest,  
JHEP 1009 (2010) 073

QCD phase transition  
is a cross-over ( $\mu=0$ )

No unique  $T_c$

## QCD phase transition from lattice QCD

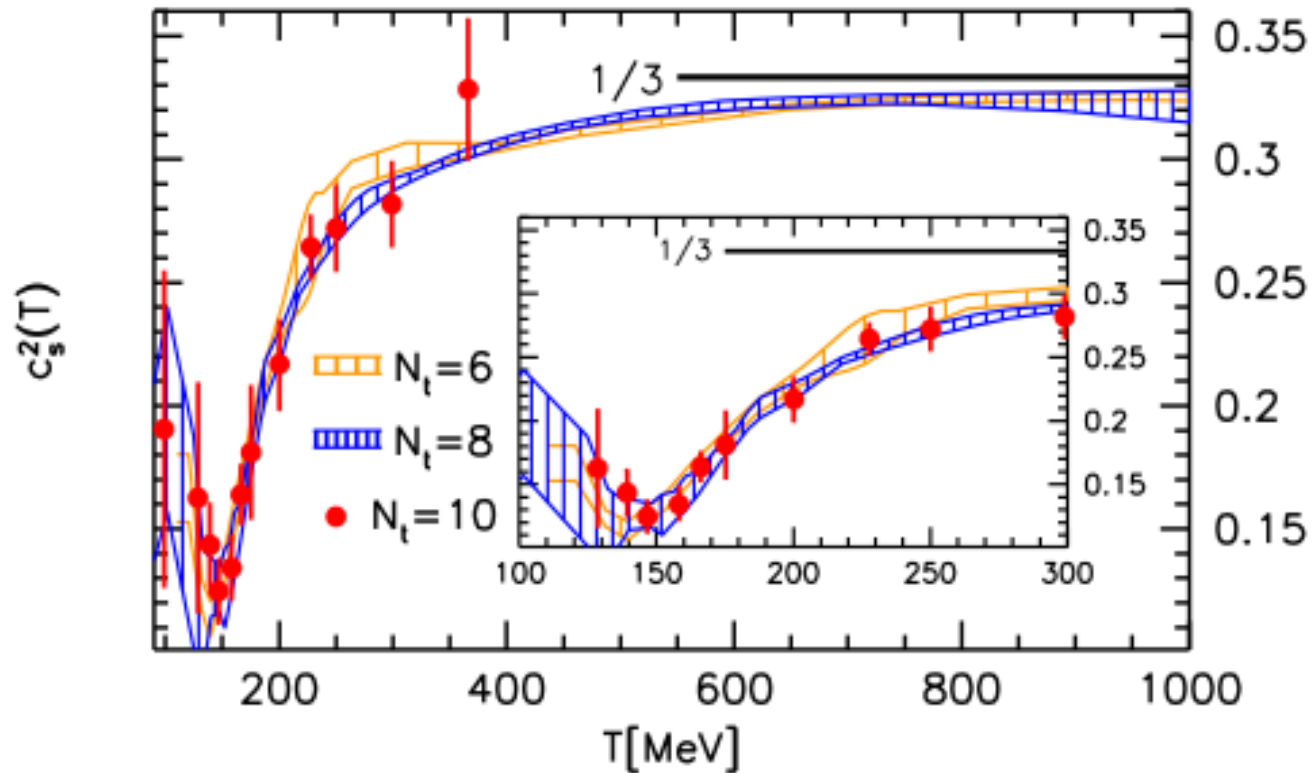
Fodor et al. 2006 (Wuppertal-Budapest)



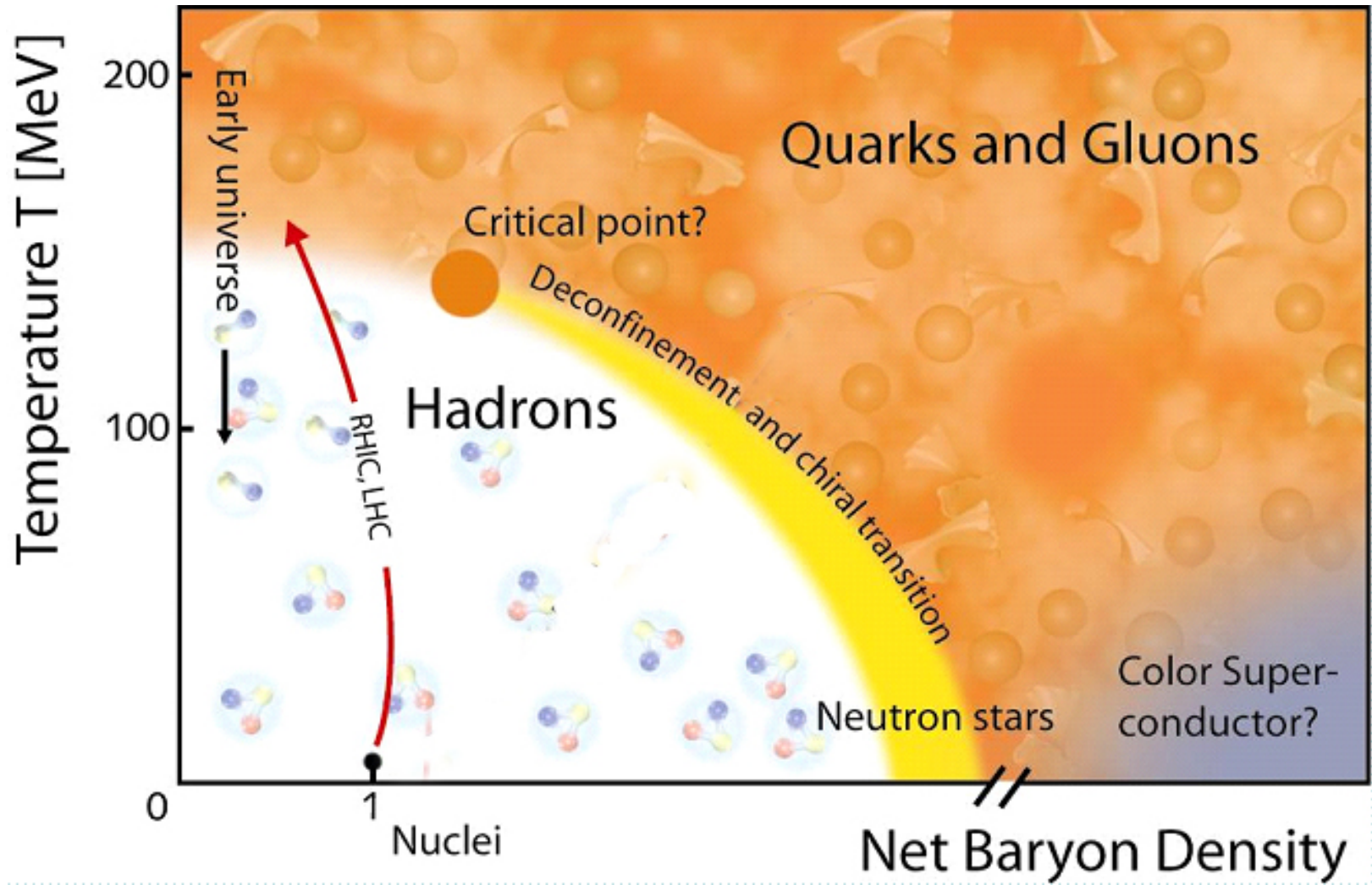
At  $T \sim 5T_c$ ,  $\epsilon$  still 80% of the Stefan-Boltzmann value:  
Quasi-particle modes at high  $T$  (Blaizot, Iancu, Rebhan '2001)

# EOS from lattice QCD

Sound velocity  $c_s^2(T) = \frac{\partial \epsilon(T)}{\partial P}$

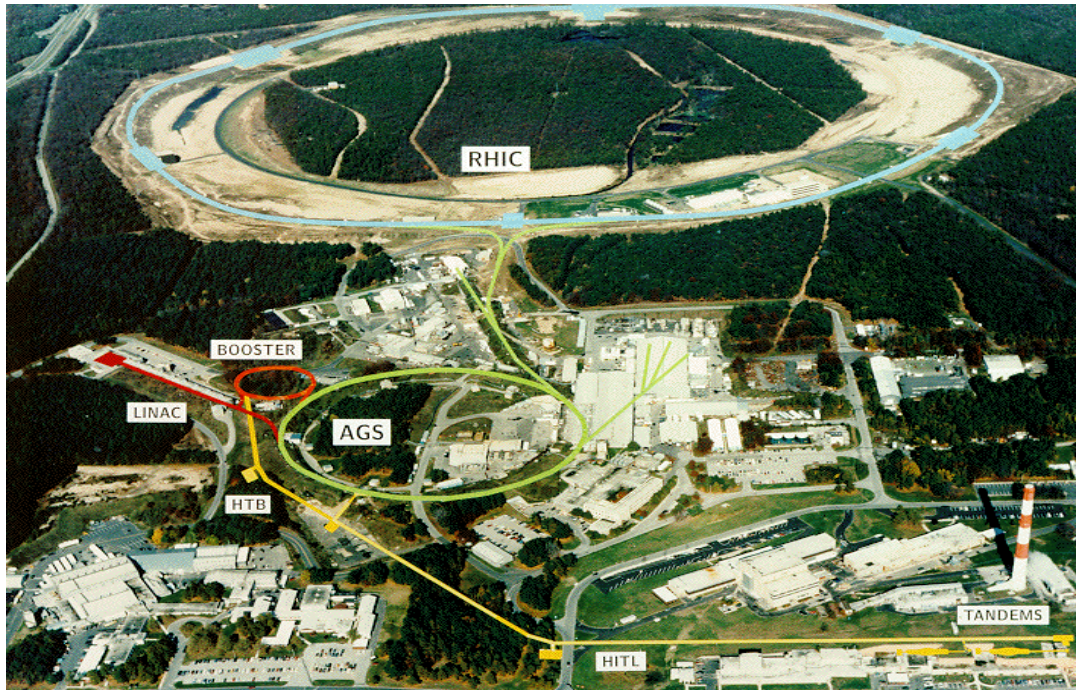


# QCD Phase Diagram





# Heavy-ion Colliders



RHIC

LHC





# Heavy Ion Collisions

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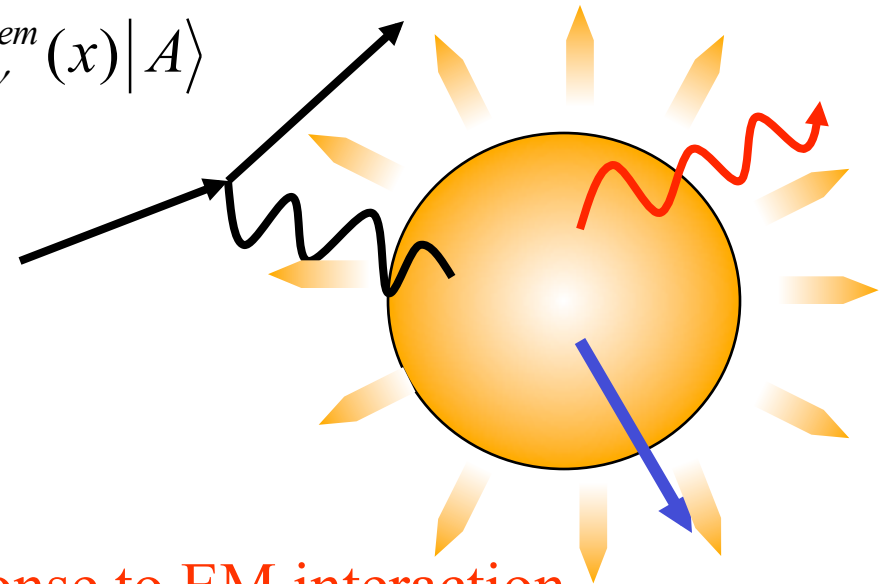


# Study properties of QGP in AA Collisions

---

$$W_{\mu\nu}(q) = \frac{1}{4\pi} \int d^4x e^{iq \cdot x} \langle A | j_\mu^{em}(0) j_\nu^{em}(x) | A \rangle$$

$$F_1(x_B) \quad x_B = -\frac{q^2}{2p \cdot q}$$

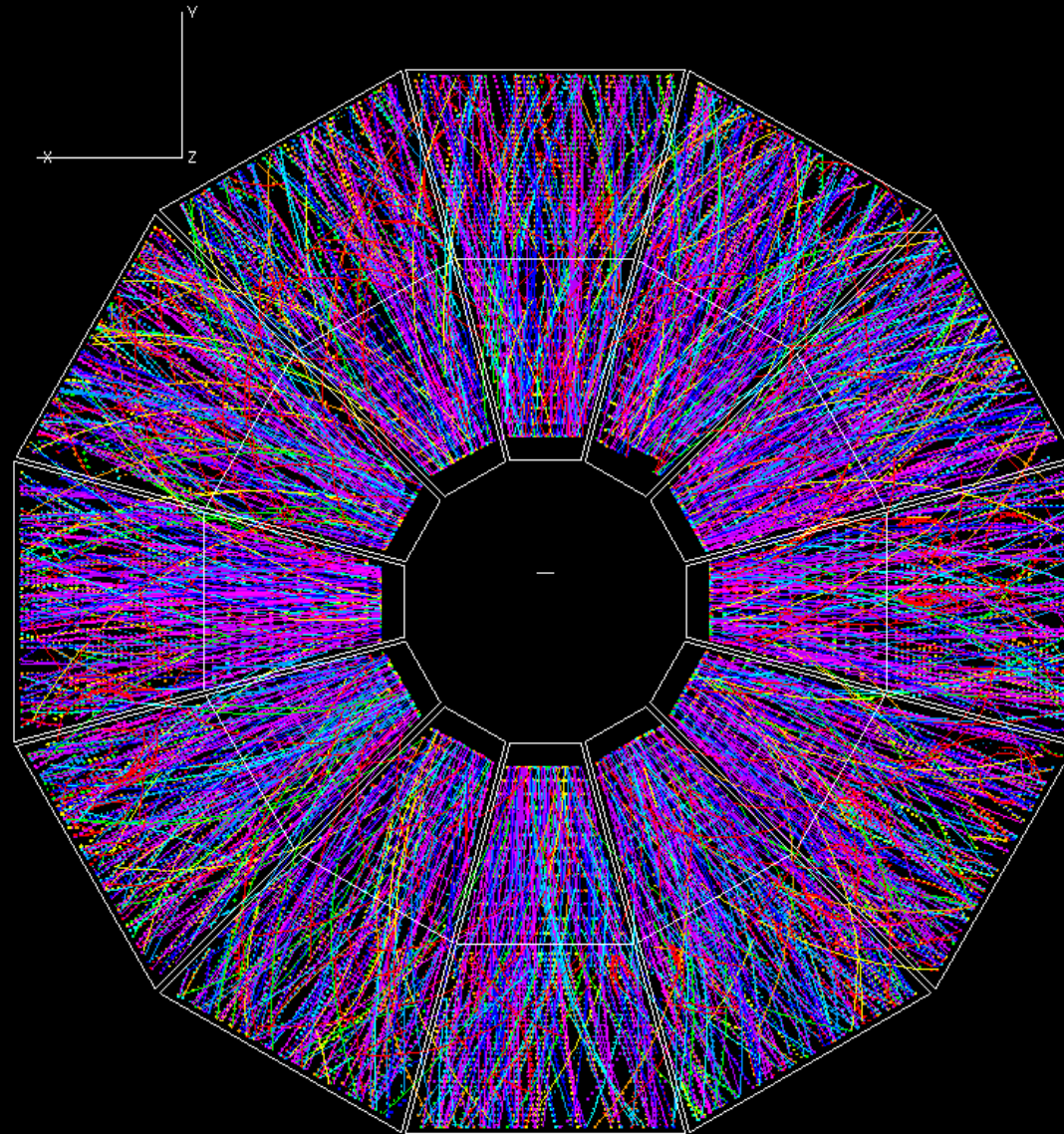


## Dynamic System:

- EM emission: Medium response to EM interaction  
 $\gamma$  production,  $J/\Psi$  suppression
- Hard probes: Medium response to strong interaction  
Jet quenching
- Soft hadrons: Bulk properties of medium, collective behavior



# Jets in Heavy-ion Collisions

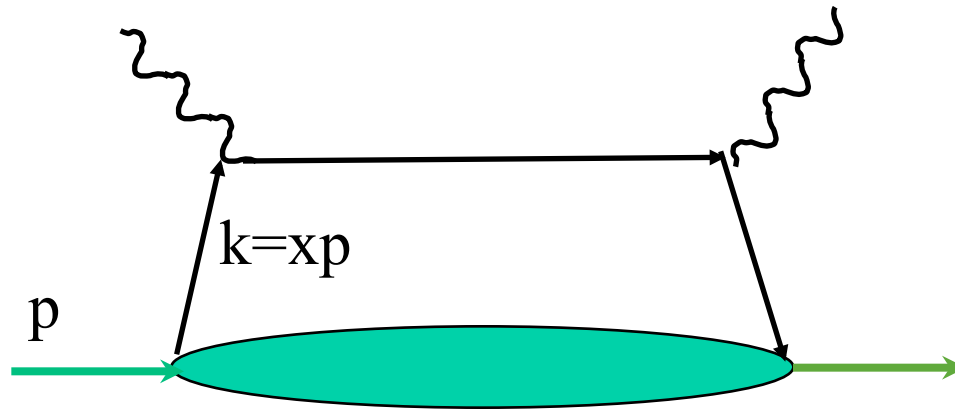






# Deeply Inelastic Scattering

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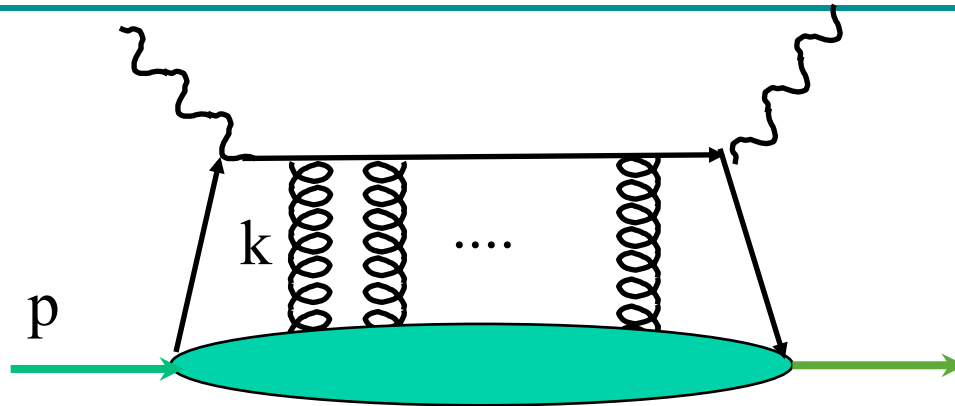
Quark distribution in collinear factorized pQCD parton model:

$$f_A^q(x) = \int \frac{dy^-}{4\pi} e^{ixp^+ y^-} \langle A | \bar{\psi}(0) \gamma^+ \psi(y^-) | A \rangle$$

quarks carrying momentum fraction  $x$  of the nucleon (nucleus)



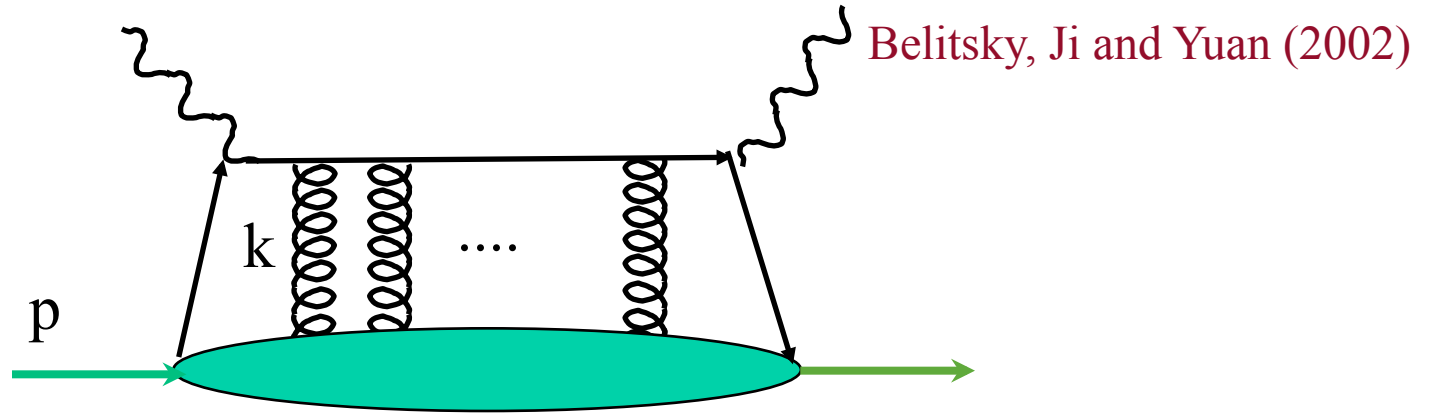
# Gauge Invariance and Multiple Interaction



$$f_A^q(x) = \int \frac{dy^-}{4\pi} e^{ixp^+ y^-} \langle A | \bar{\psi}(0) \gamma^+ \mathcal{L}_{\parallel}(0, y^-; \vec{0}_{\perp}) \psi(y^-) | A \rangle$$

$$\mathcal{L}_{\parallel}(0, y^-; \vec{0}_{\perp}) = \mathcal{P} \exp \left[ ig \int_0^{y^-} d\xi^- A_+(\xi^-, \vec{0}_{\perp}) \right]$$

# TMD parton distribution in DIS



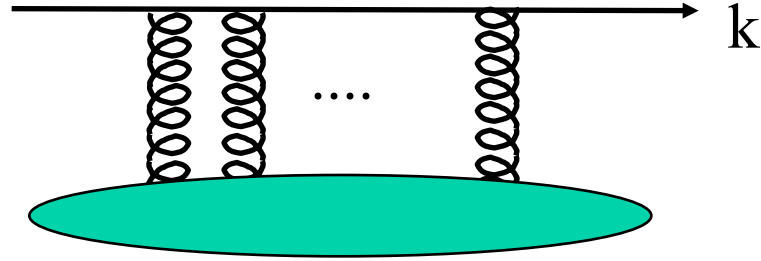
$$f_A^q(x, \vec{k}_\perp) = \int \frac{dy^-}{4\pi} \frac{d^2 y_\perp}{(2\pi)^2} e^{ixp^+ y^- - i\vec{k}_\perp \cdot \vec{y}_\perp} \langle A | \bar{\psi}(0) \gamma^+ \mathcal{L}(0, y) \psi(y) | A \rangle$$

$$\mathcal{L}(0, y) = \mathcal{L}_\parallel^\dagger(\infty, 0; \vec{0}_\perp) \mathcal{L}_\perp^\dagger(\infty; \vec{y}_\perp, \vec{0}_\perp) \mathcal{L}_\parallel(\infty, y^-; \vec{y}_\perp)$$



$$\mathcal{L}_\parallel(0, y^-; \vec{0}_\perp) = \mathcal{P} \exp \left[ ig \int_0^{y^-} d\xi^- A_+(\xi^-, \vec{0}_\perp) \right] \quad \mathcal{L}_\perp(\infty; \vec{y}_\perp, \vec{0}) = \mathcal{P} \exp \left[ -ig \int_{\vec{0}_\perp}^{\vec{y}_\perp} d\vec{\xi}_\perp \cdot \vec{A}_\perp(\infty, \vec{\xi}_\perp) \right]$$

# Jet Transport in Medium



Classical Lorentz force:  $\frac{d\vec{p}_\perp}{d\tau} = g\vec{F}_{\perp\mu}v^\mu$

$$\vec{W}_\perp(y^-, \vec{y}_\perp) \equiv i\vec{D}_\perp(y) + g \int_{-\infty}^{y^-} d\xi^- \vec{F}_{+\perp}(\xi^-, y_\perp)$$

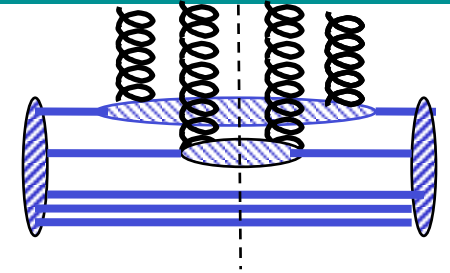
Jet Transport Operator

$$f_A^q(x, \vec{k}_\perp) = \int \frac{dy^-}{4\pi} e^{ixp^+ y^-} \langle A | \bar{\psi}(0) \gamma^+ \exp[\vec{W}_\perp(y^-) \cdot \nabla_{k_\perp}] \psi(y^-) | A \rangle \delta^{(2)}(\vec{k}_\perp)$$

Liang, XNW & Zhou (2008)

# Momentum Broadening

$$\langle\langle W_{\perp}^{2n} \rangle\rangle_A \sim \left[ \int dy \frac{\rho_A(y)}{2p^+} \langle N | F_{+\perp} F_{+\perp} | N \rangle \right]^n \sim \left[ \int dy \rho_A(y) x G_N(x) \right]^n$$



2-gluon correlation approximation

$$f_A^q(x, \vec{k}_{\perp}) \approx \frac{A}{\pi\Delta} \int d^2q_{\perp} \exp \left[ -\frac{(\vec{k}_{\perp} - \vec{q}_{\perp})^2}{\Delta} \right] f_N^q(x, \vec{q}_{\perp})$$

$$\Delta = \langle \Delta k_{\perp}^2 \rangle = \int d\xi_N^- \hat{q}(\xi_N)$$

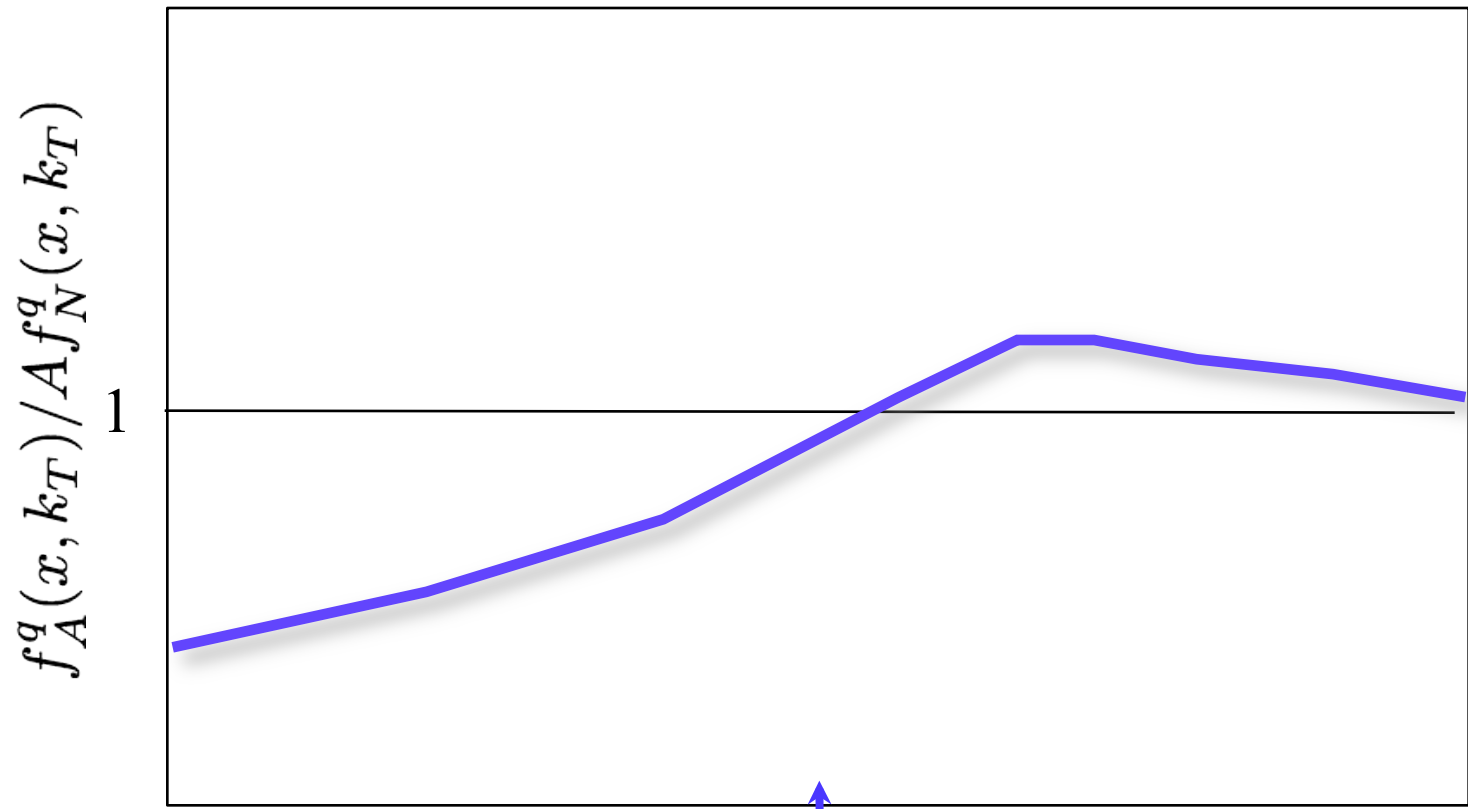
Liang, XNW & Zhou' 08  
Majumder & Muller' 07  
Kovner & Wiedemann' 01  
BDMPS' 96

$$\hat{q}(\xi_N) \equiv \frac{4\pi^2 \alpha_s C_F}{N_c^2 - 1} \rho_A(\xi_N) x G_N(x) |_{x \approx 0}$$

Jet transport parameter

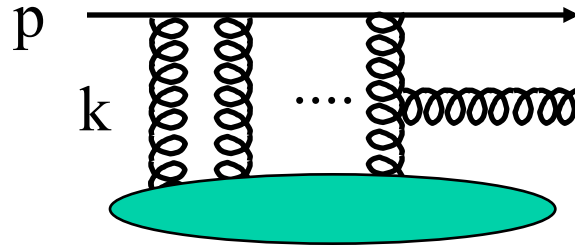
# P<sub>T</sub> Broadening

$$f_N^q(x, k_T) \sim 1/(k_T^2 + p_0^2)^\alpha$$



$$\hat{q}(\xi_N) \equiv \frac{4\pi^2 \alpha_s C_F}{N_c^2 - 1} \rho_A(\xi_N) x G_N(x) |_{x \approx 0}$$

# Parton Energy Loss



Splitting functions in medium

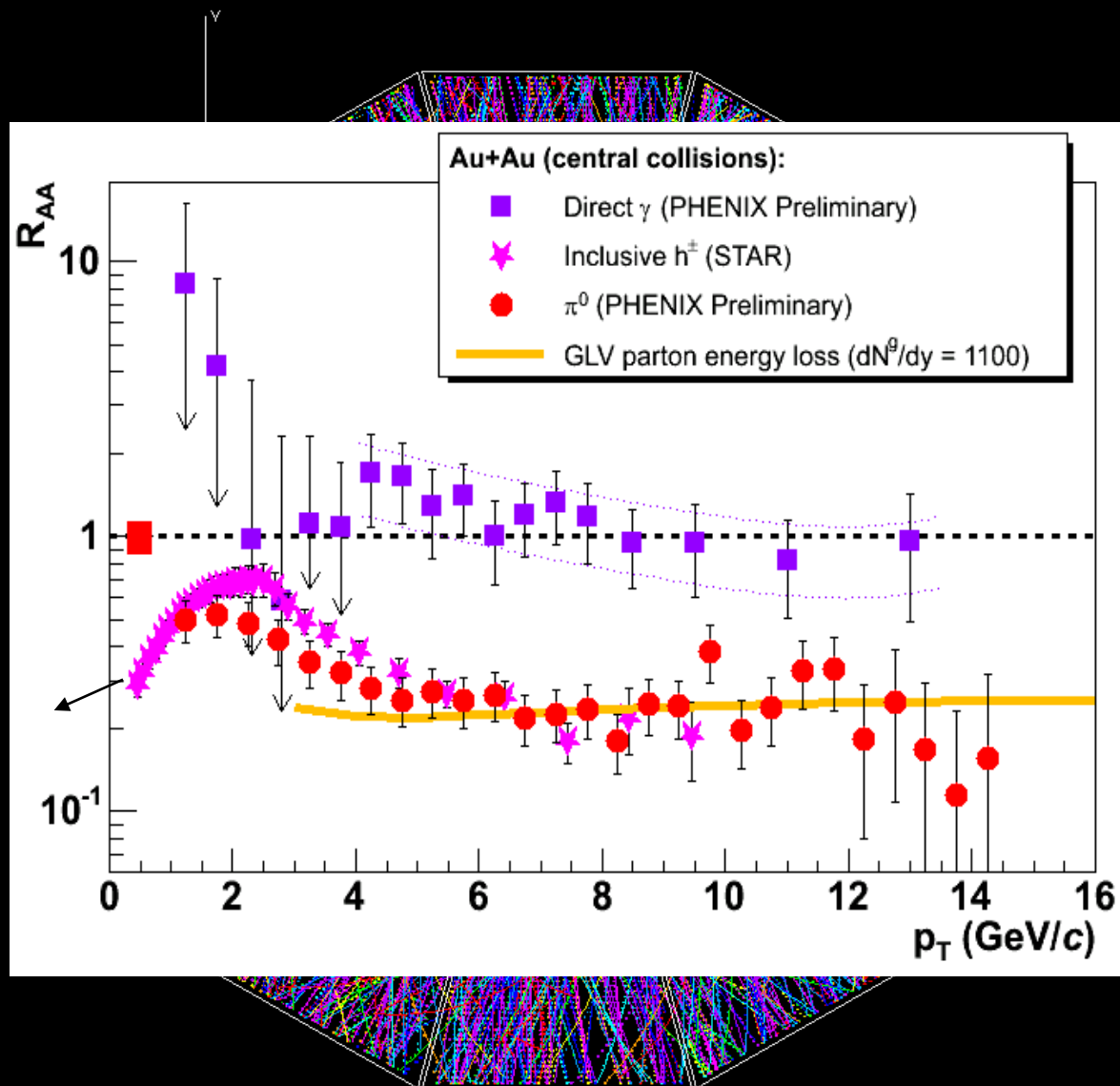
$$\Delta\gamma(z, \ell_{\perp}^2) = C_A \frac{1+z^2}{(1-z)_+} \frac{2}{\ell_{\perp}^4} \int d\xi^- \hat{q}(\xi) [1 - \cos(x_L p^+ \xi^-)]$$

Parton Energy Loss

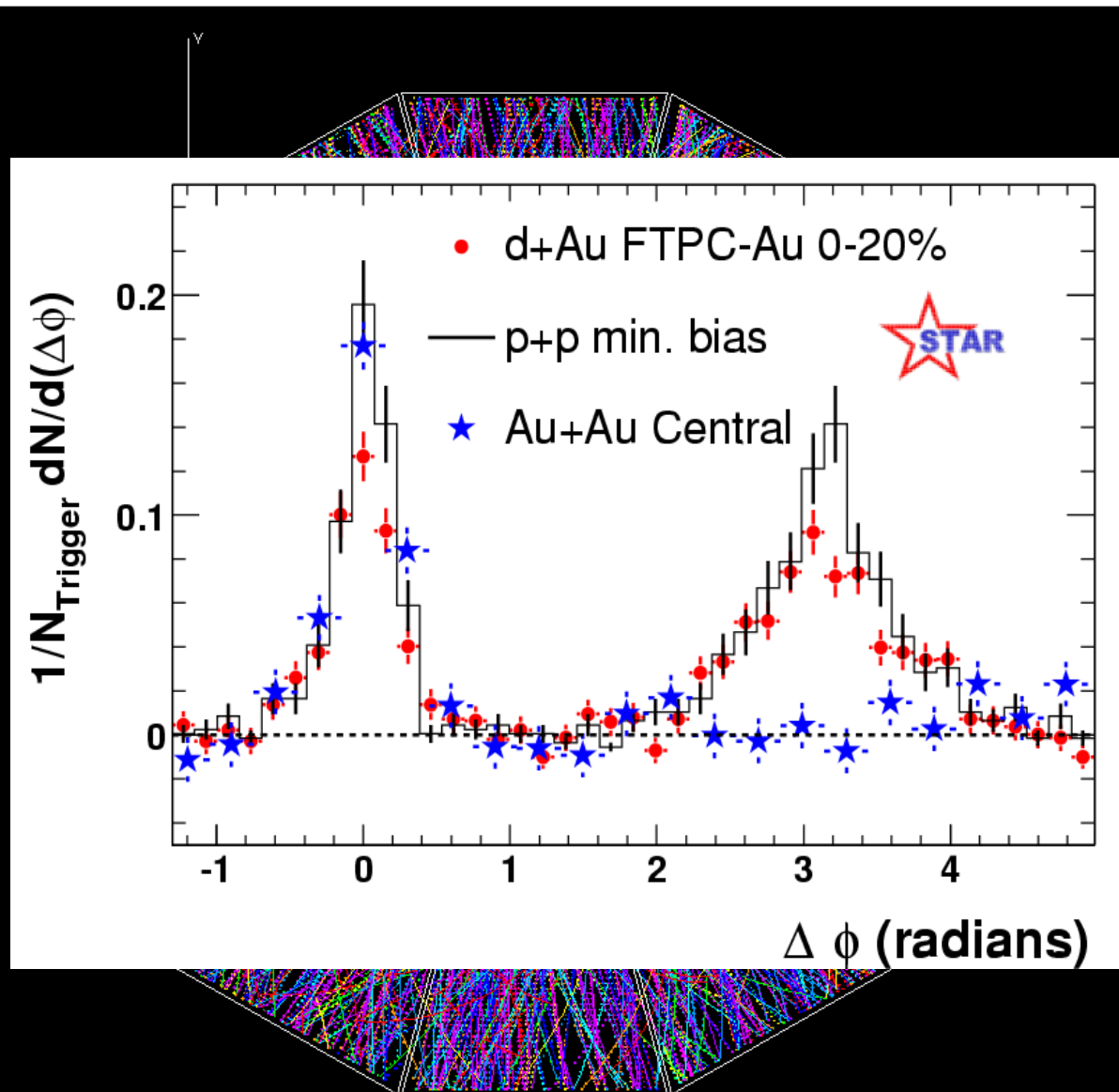
$$\frac{\Delta E}{E} = C_A \frac{\alpha_s}{2\pi} \int \frac{dl_T^2}{l_T^4} \int dz [1 + (1-z)^2] \int d\xi^- \hat{q}(\xi) 4 \sin^2(x_L p^+ \xi^- / 2)$$



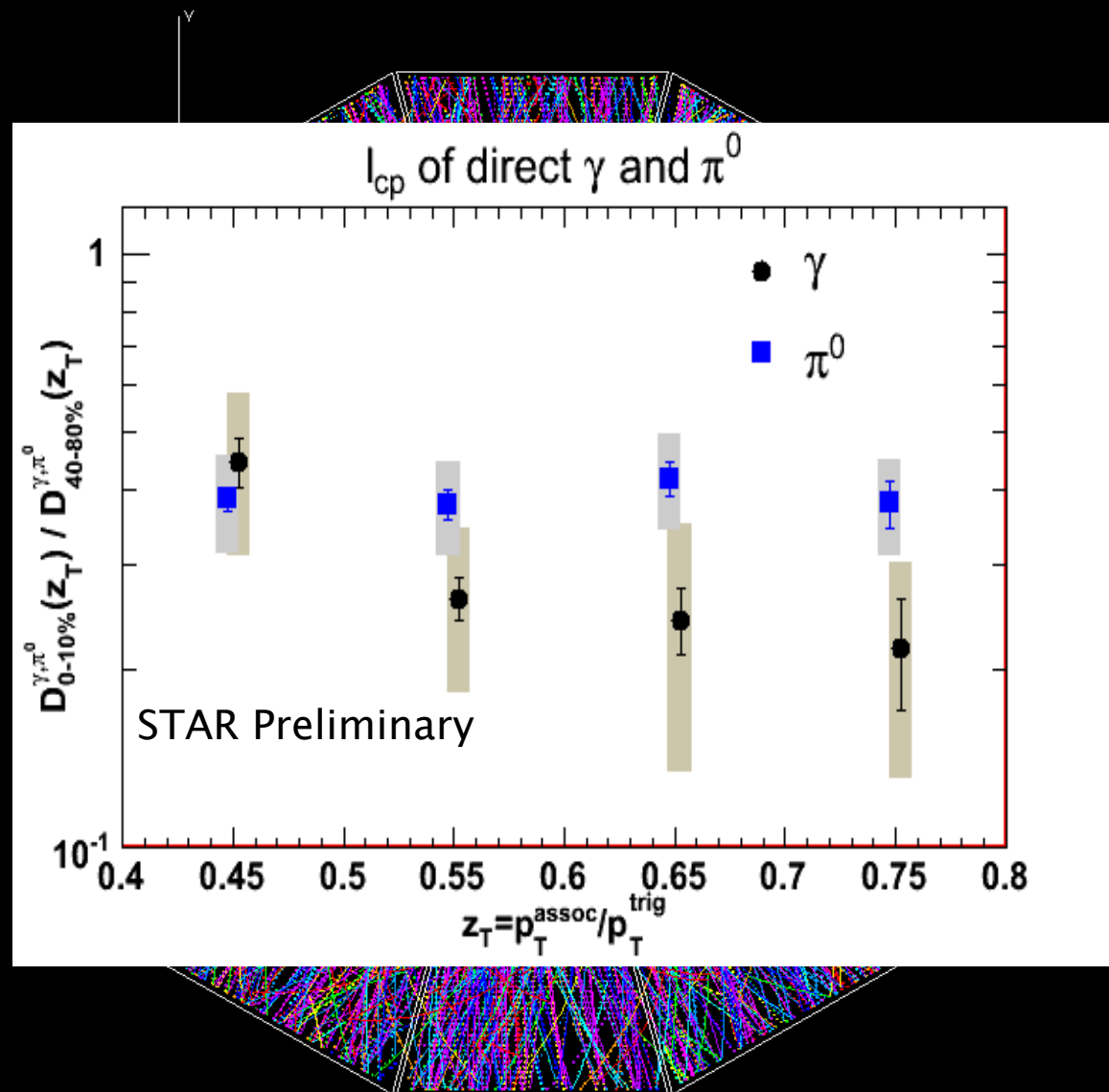
# Jet Quenching phenomena at RHIC



# Jet Quenching phenomena at RHIC

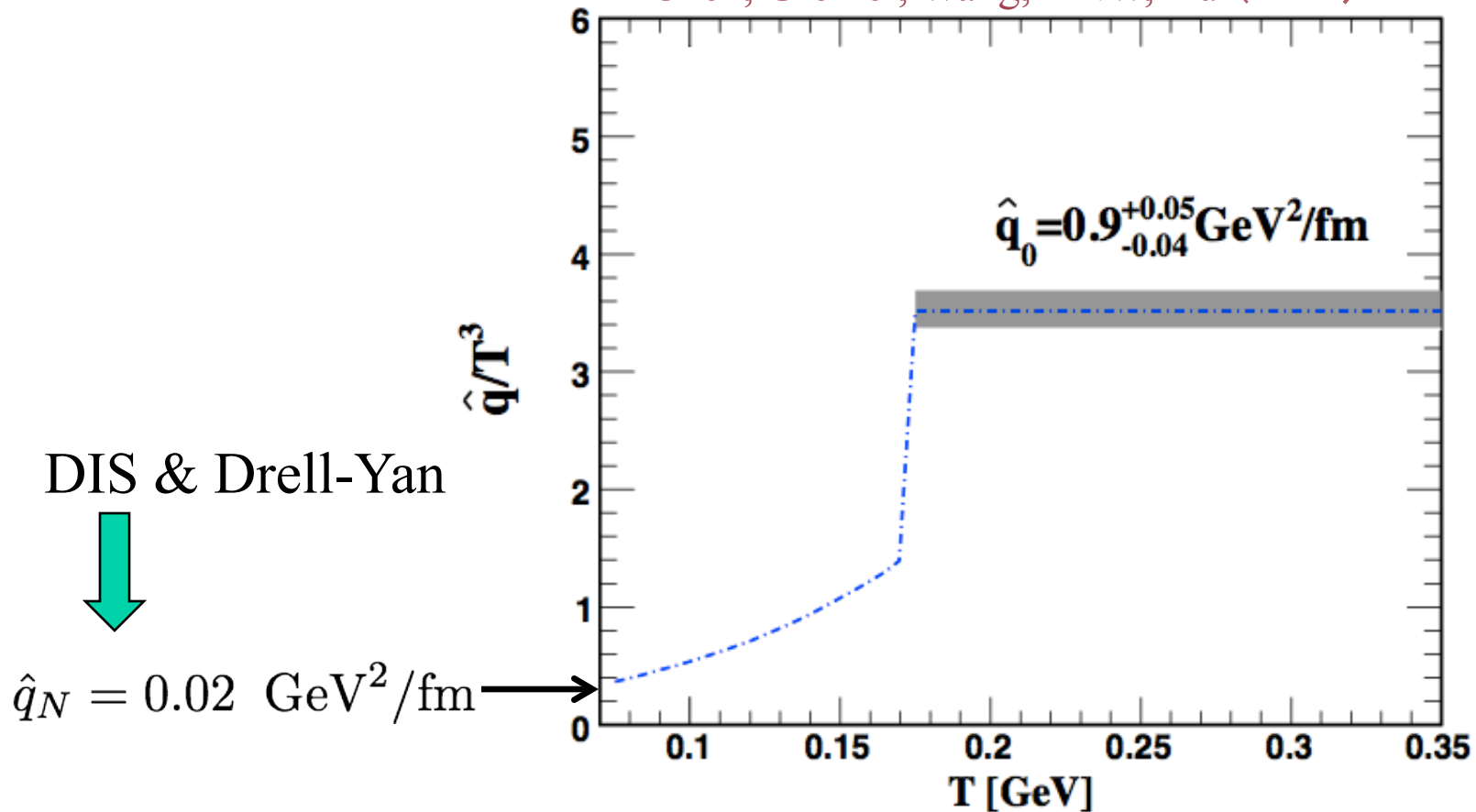


# Jet Quenching phenomena at RHIC



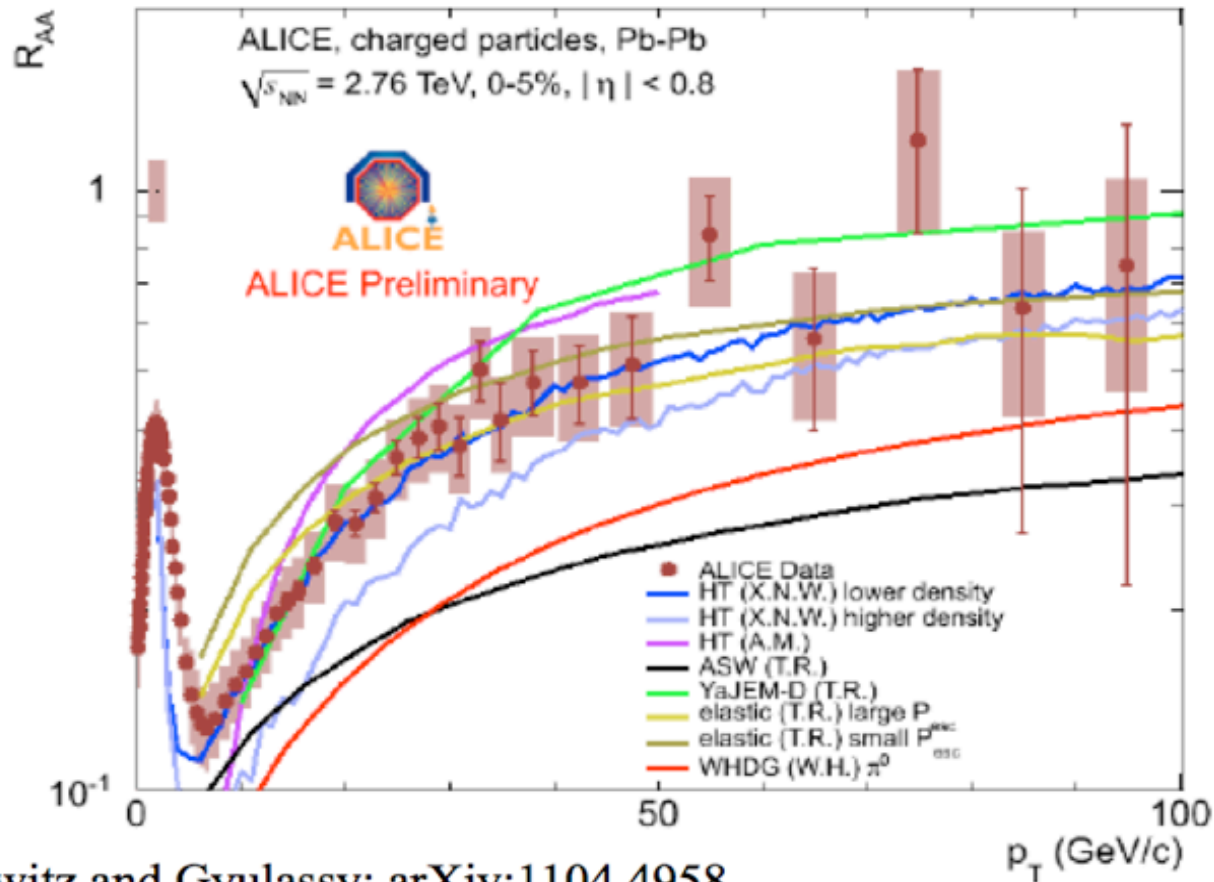
# Jet quenching in QGP & hadronic phase

Chen, Greiner, Wang, XNW, Xu (2010)



30% quenching from hadronic phase

# Jet quenching at LHC



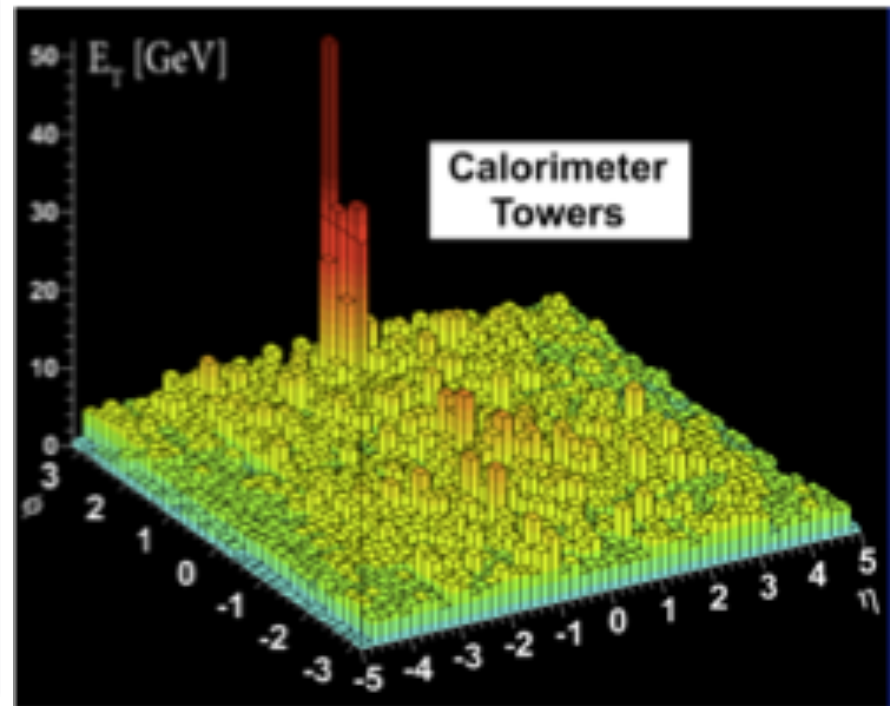
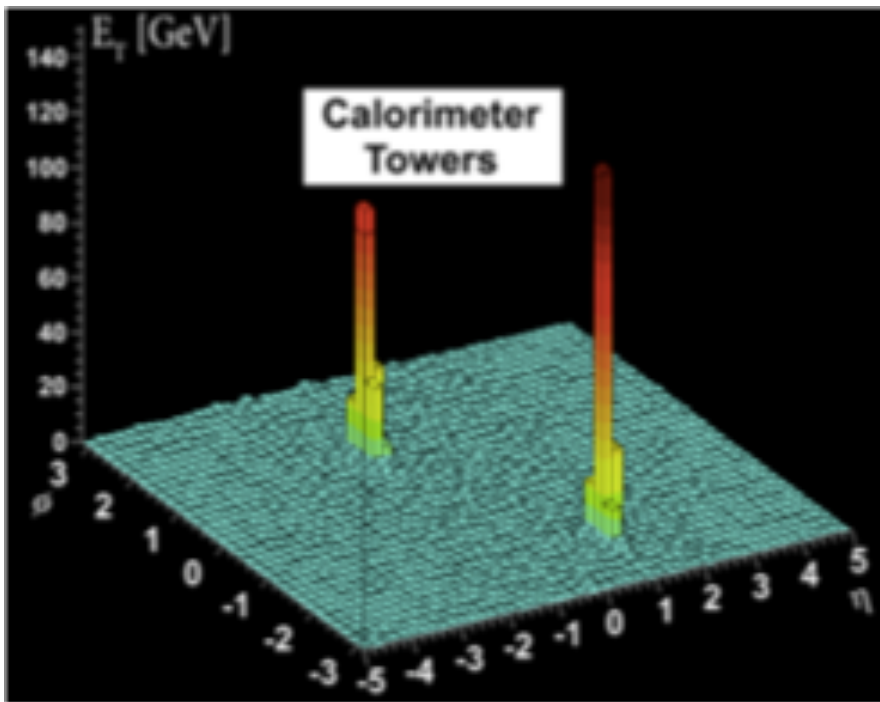
$$\hat{q} \propto \frac{1}{\tau_0 \pi R_A^2} \frac{dN_{ch}}{d\eta}$$

Horowitz and Gyulassy: [arXiv:1104.4958](https://arxiv.org/abs/1104.4958)

Majumder and Shen: [arXiv:1103.0809](https://arxiv.org/abs/1103.0809)

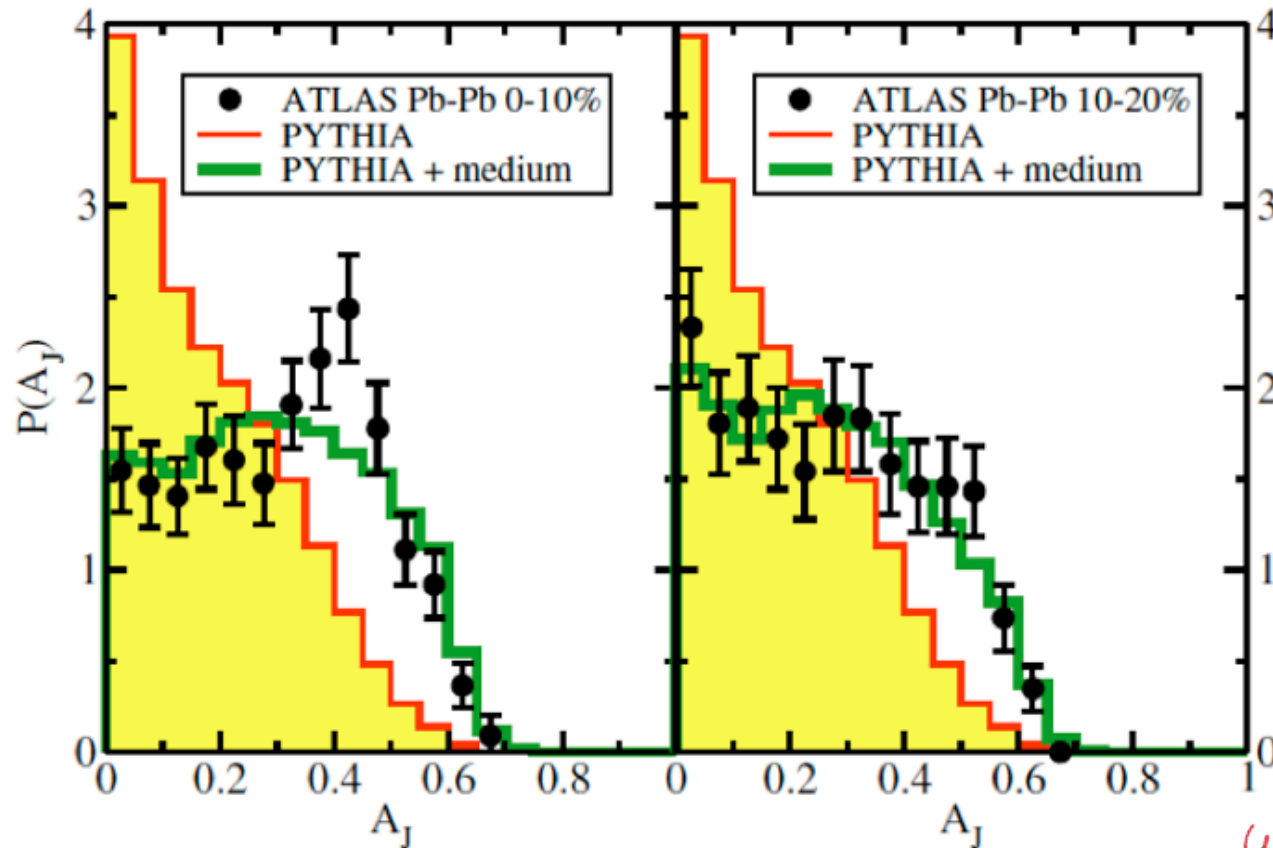
Chen et. al. [arXiv:1102.5614](https://arxiv.org/abs/1102.5614)

# Jet Quenching at LHC





# Di-Jet Asymmetry



$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}$$

$$\omega > \omega_c$$

$$\omega_c = 1 - 2 \text{ GeV}$$

Qin and Muller Phys.Rev.Lett.106:162302,2011

B. Schenke, C. Gale, S. Jeon, Phys.Rev.C80:054913 (2009)

Y. He, Vite, B.W. Zhang, arXiv:1105.2566

# Relativistic Hydrodynamics

---

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$T^{\mu\nu} = (\epsilon + P)u^{\mu}u^{\nu} - Pg^{\mu\nu} + \Delta T^{\mu\nu}$$

$$\Delta T^{\mu\nu} = \eta(\Delta^{\mu}u^{\nu} + \Delta^{\nu}u^{\mu}) + \left(\frac{2}{3}\eta - \zeta\right)H^{\mu\nu}\partial_{\rho}u^{\rho}$$

Inputs:

$\eta$  shear viscosity

$\zeta$  bulk viscosity

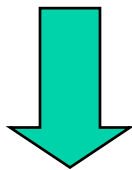
EOS

Initial conditions for  $\epsilon$  &  $u$

# Elliptic Flow

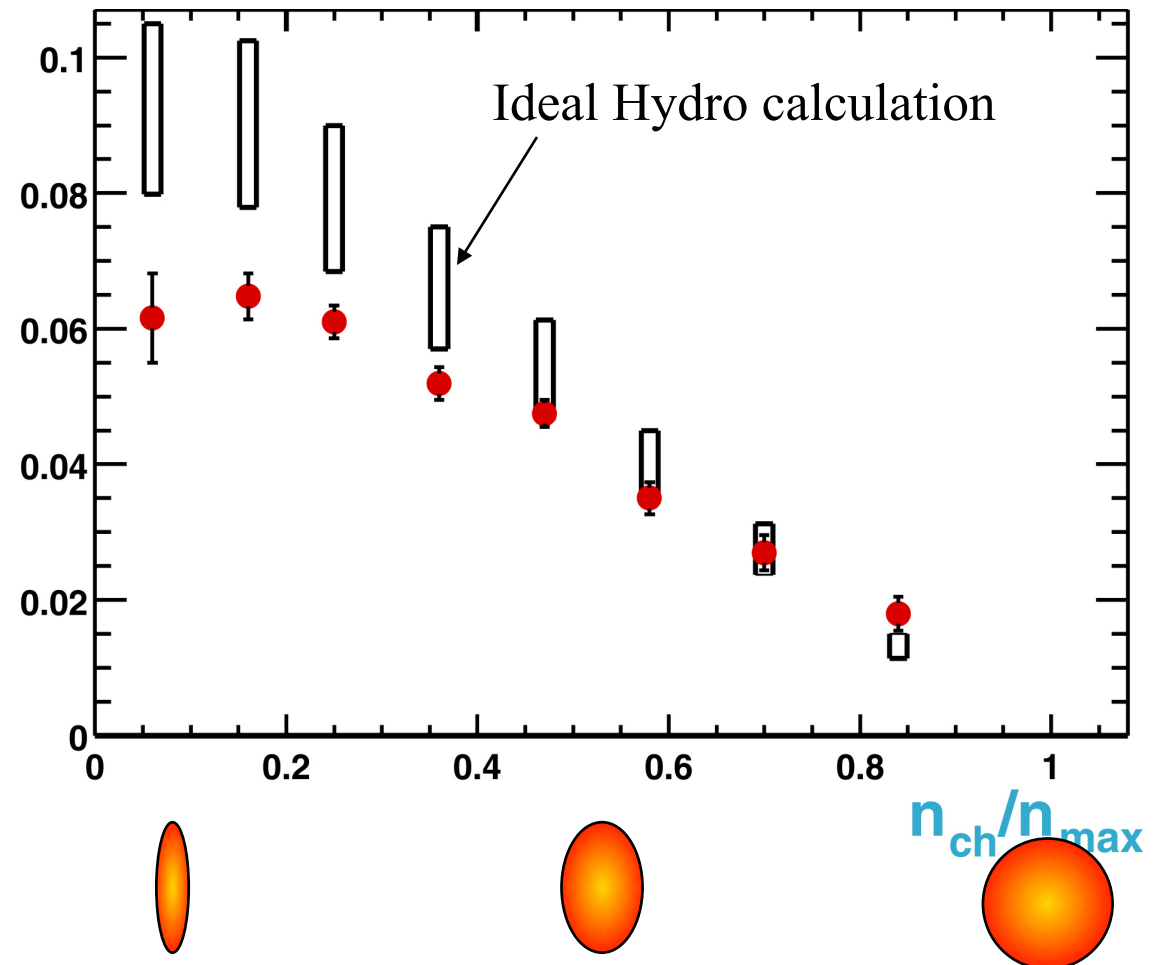


Pressure gradient  
anisotropy

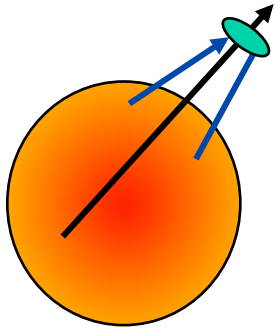


$$v_2 = \langle \cos 2\varphi \rangle$$

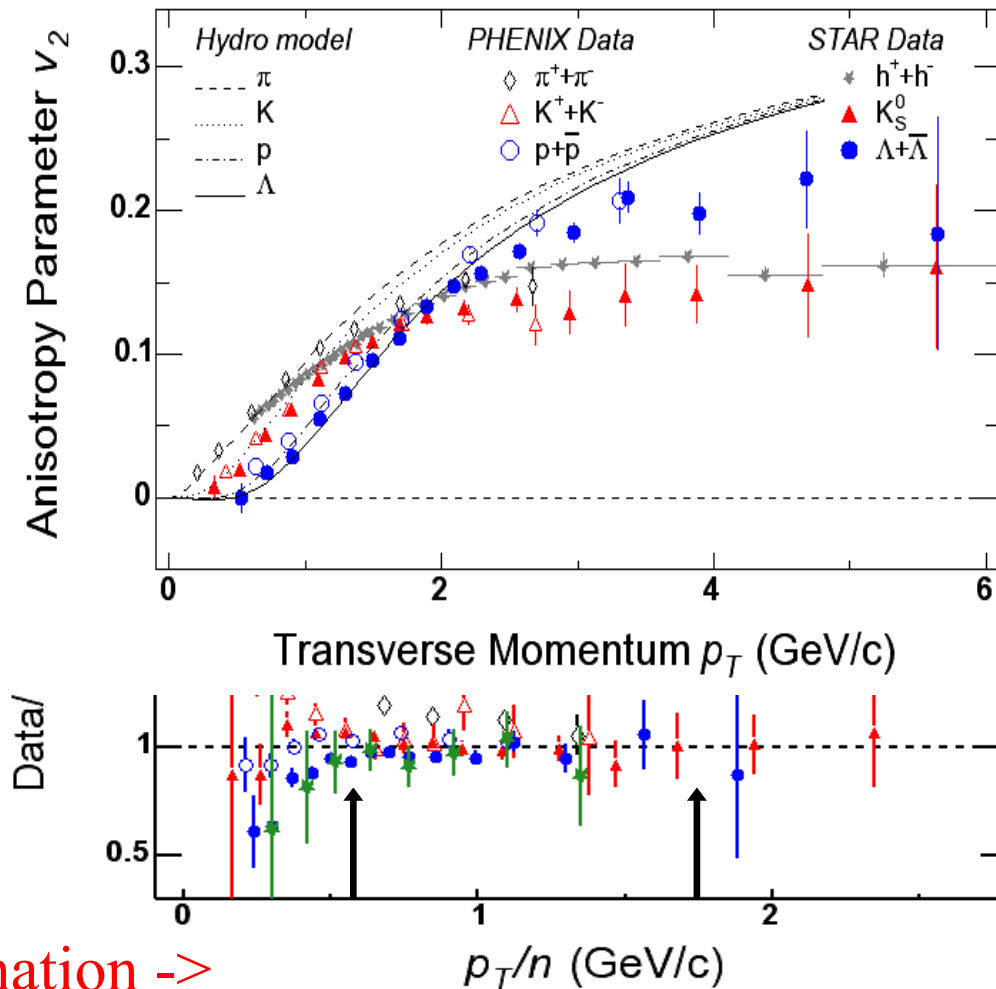
$$\frac{dN_{ch}}{d\varphi} = N_0 (1 + v_1 \cos \varphi + 2v_2 \cos 2\varphi + \dots)$$



# Partonic flows



Parton recombination ->  
Partonic degrees of freedom



$n$  = number  
of constituent  
quarks

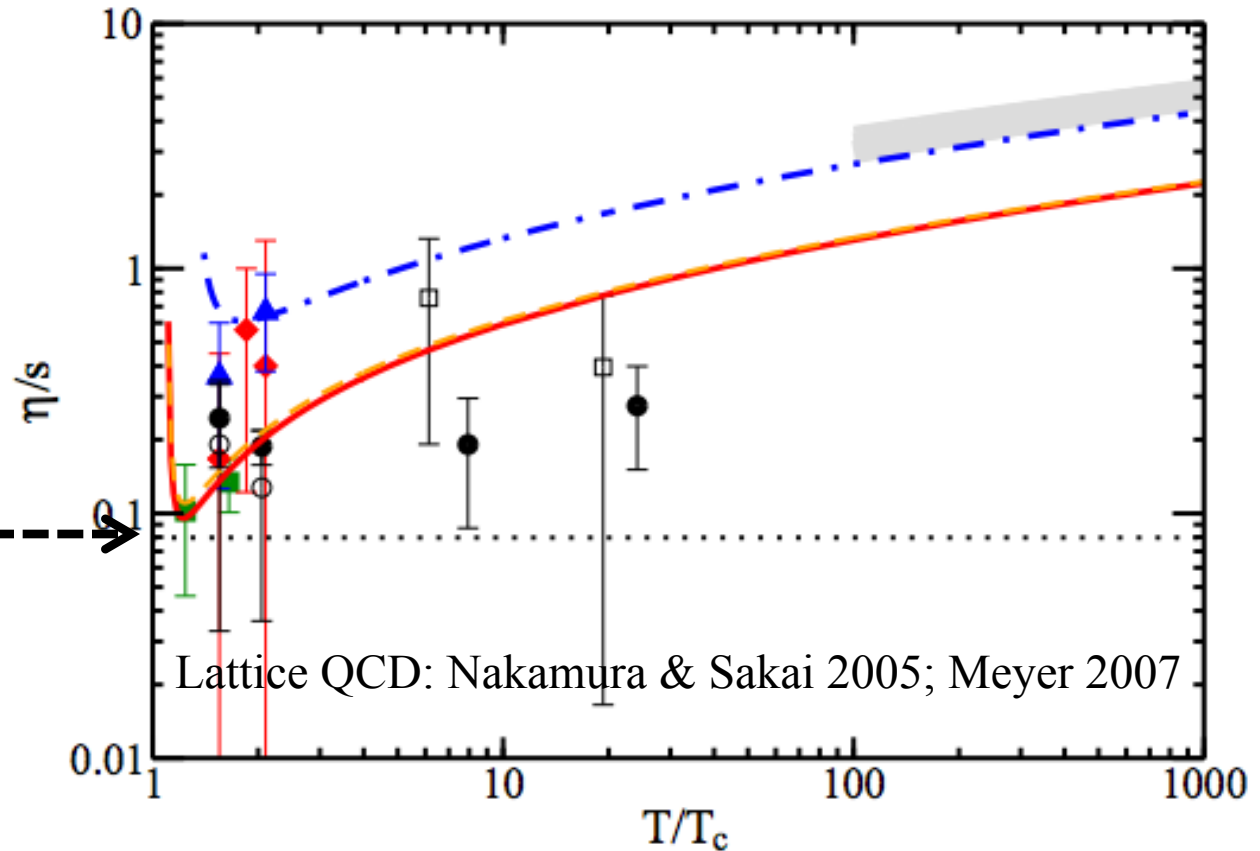
# Transport properties of QGP

Bluhm, Kampfer & Redlich 2010

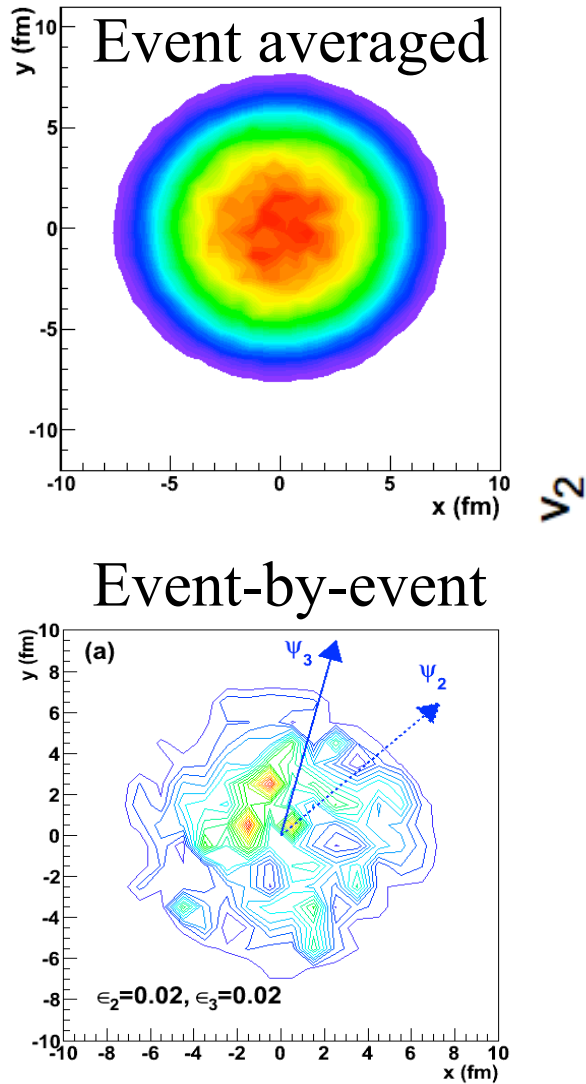
$$\frac{\eta}{s} = \frac{15}{16\pi} \frac{f_\pi^4}{T^4}$$

For pion gas

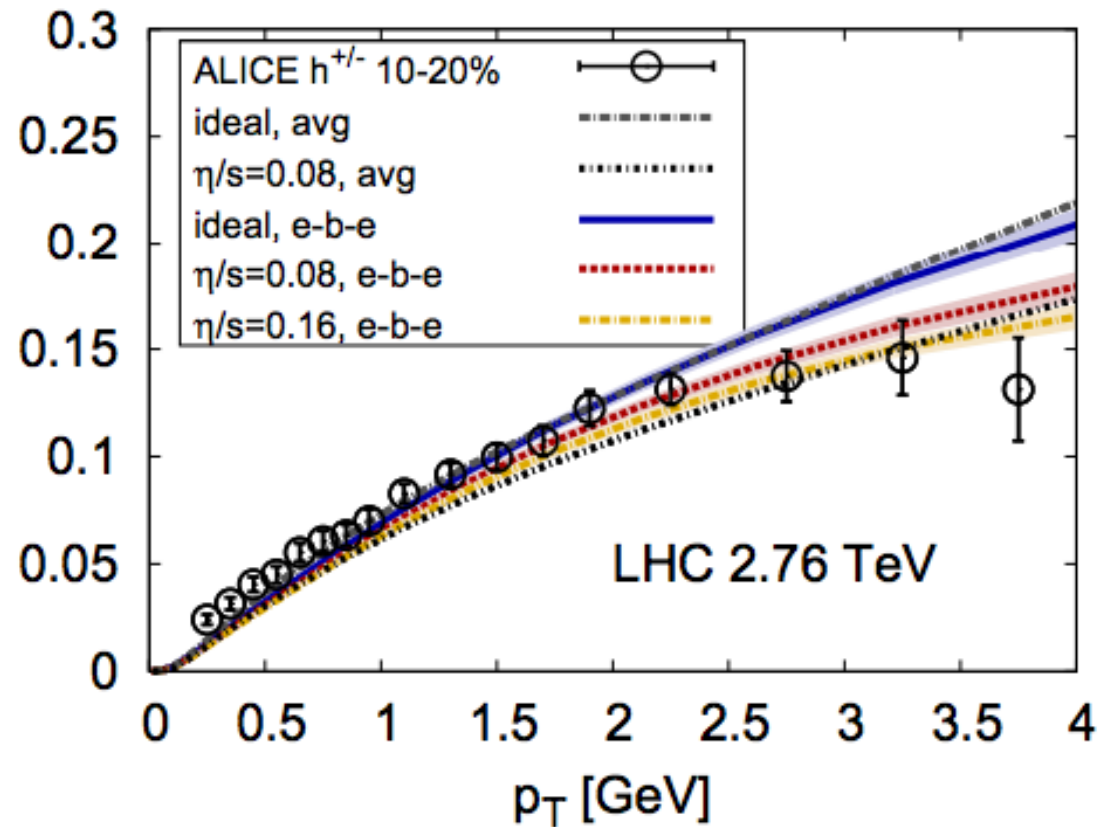
AdS/CFT



# Event-by-event viscous hydro

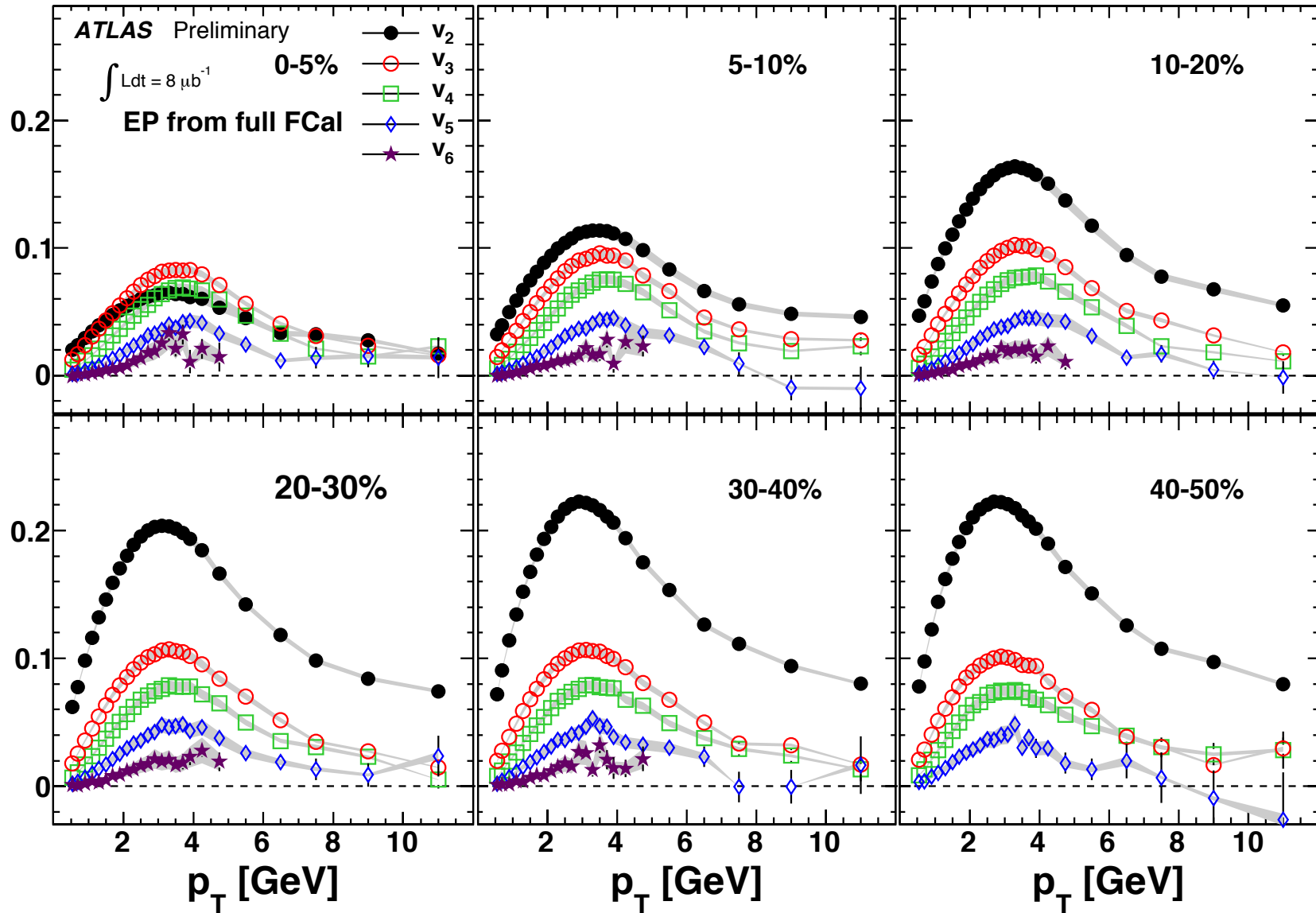


Event-by-event (3+1)D viscous hydro



Schenke, Jeon & Gale 2011

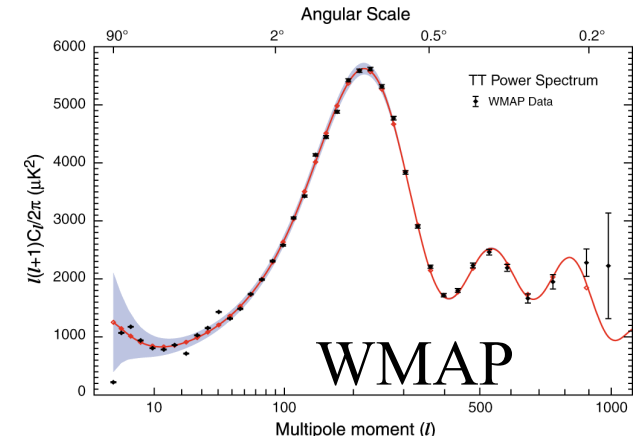
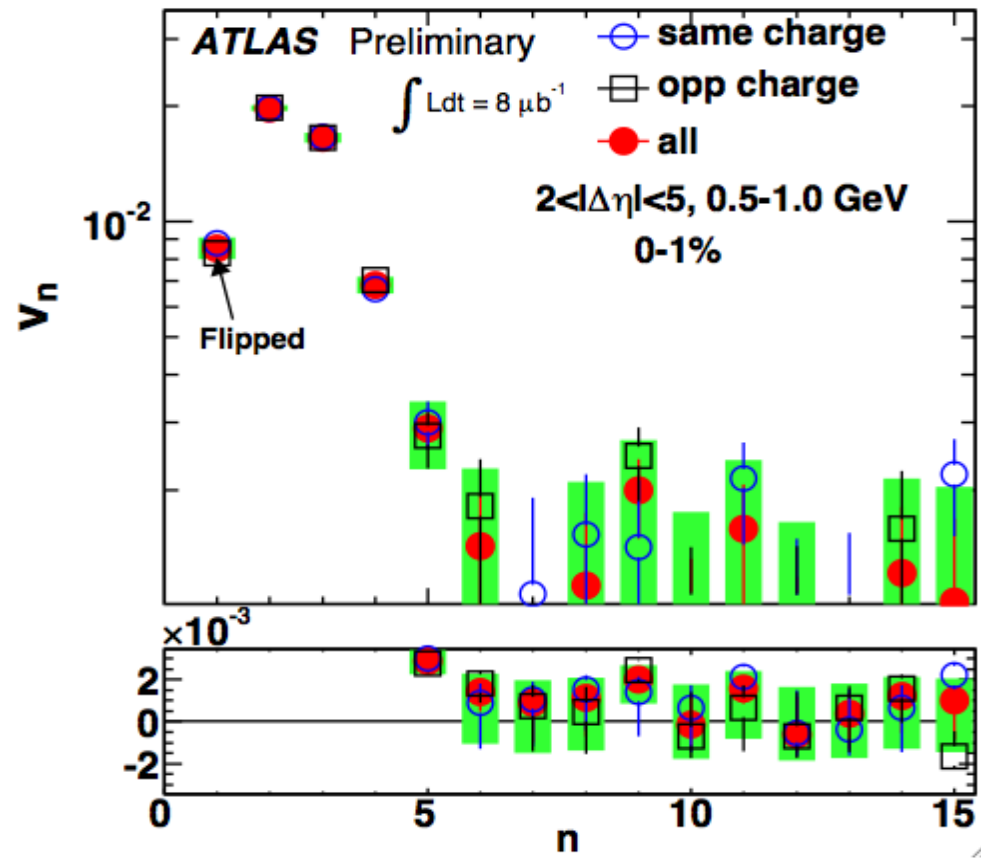
# Anisotropic flow



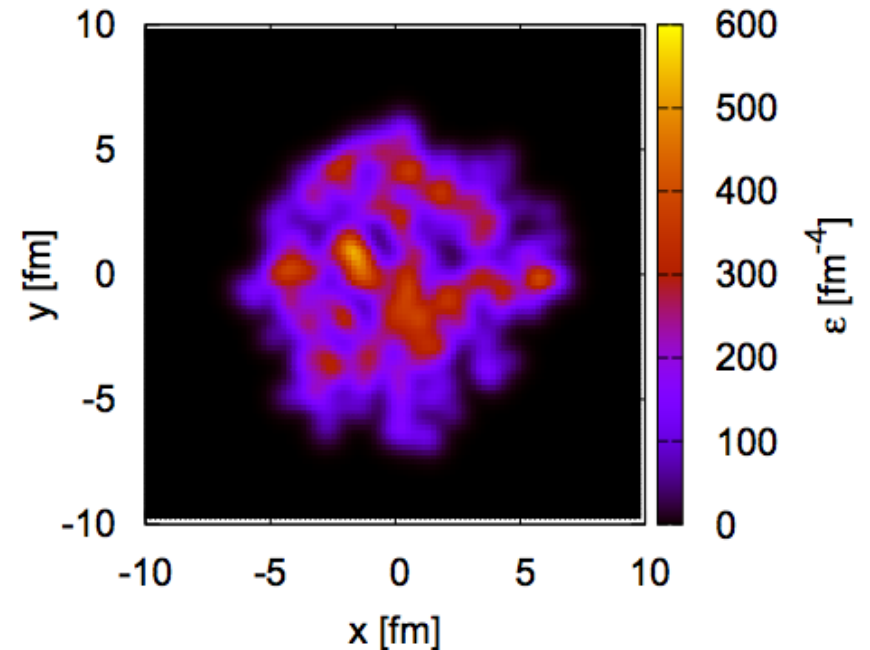


# Hadron Anisotropy Probe of the Little Bang

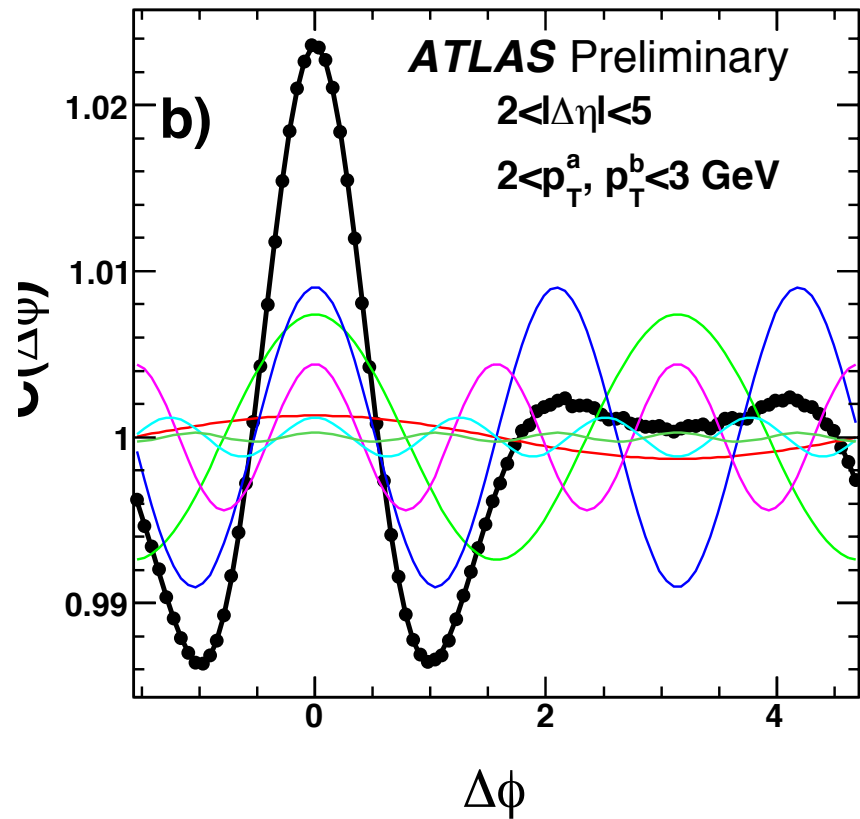
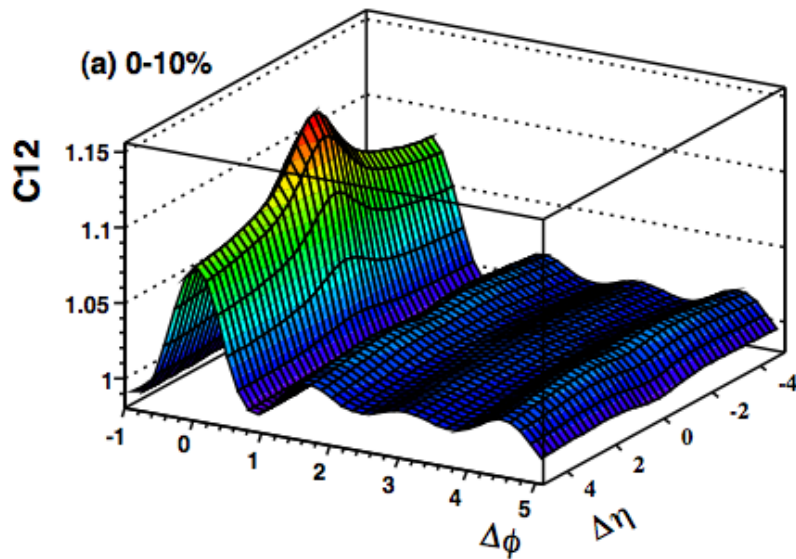
## Central Pb+Pb collisions



$\tau = 0.4 \text{ fm}/c$

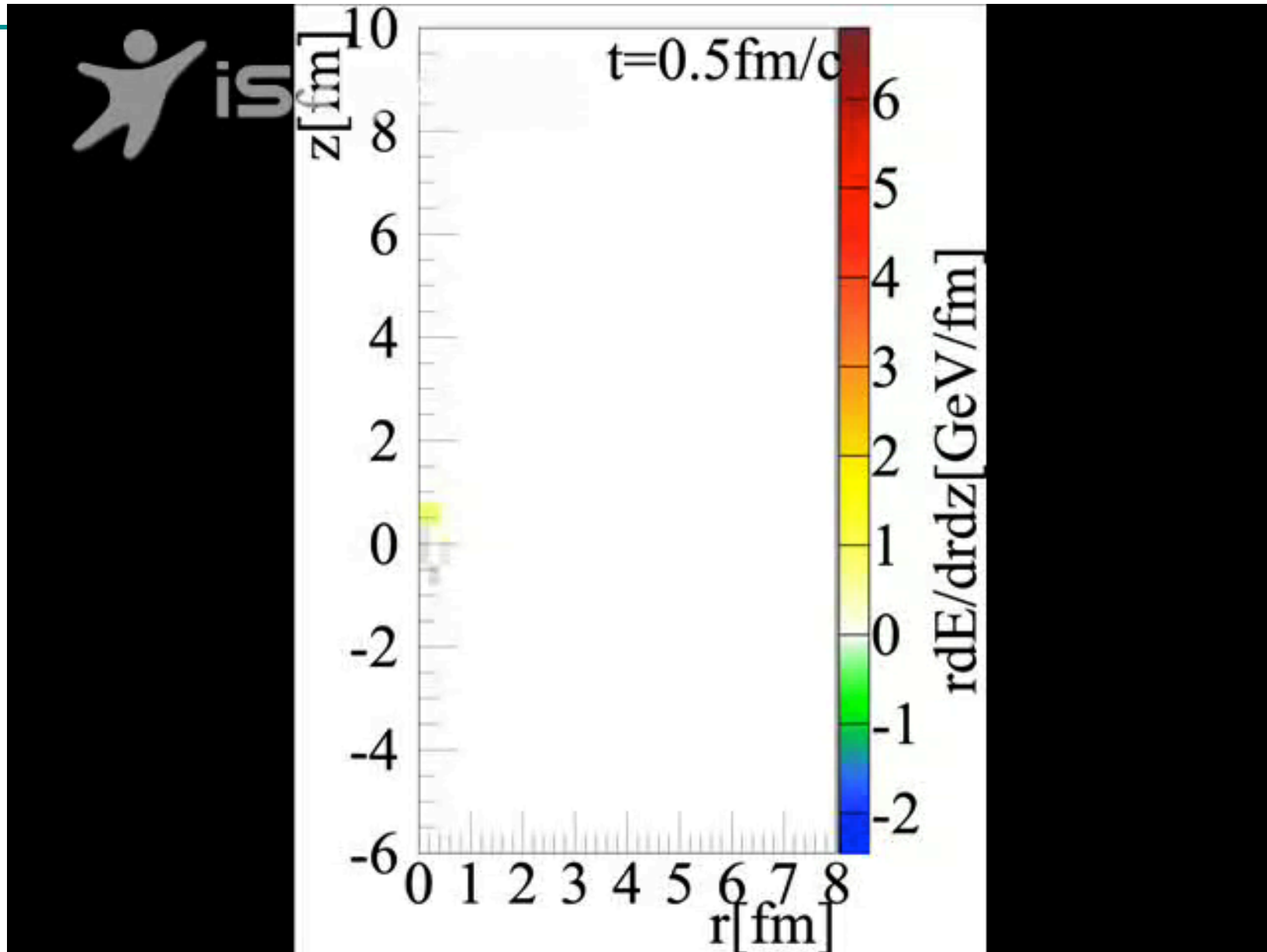


# Long range initial longitudinal correlation



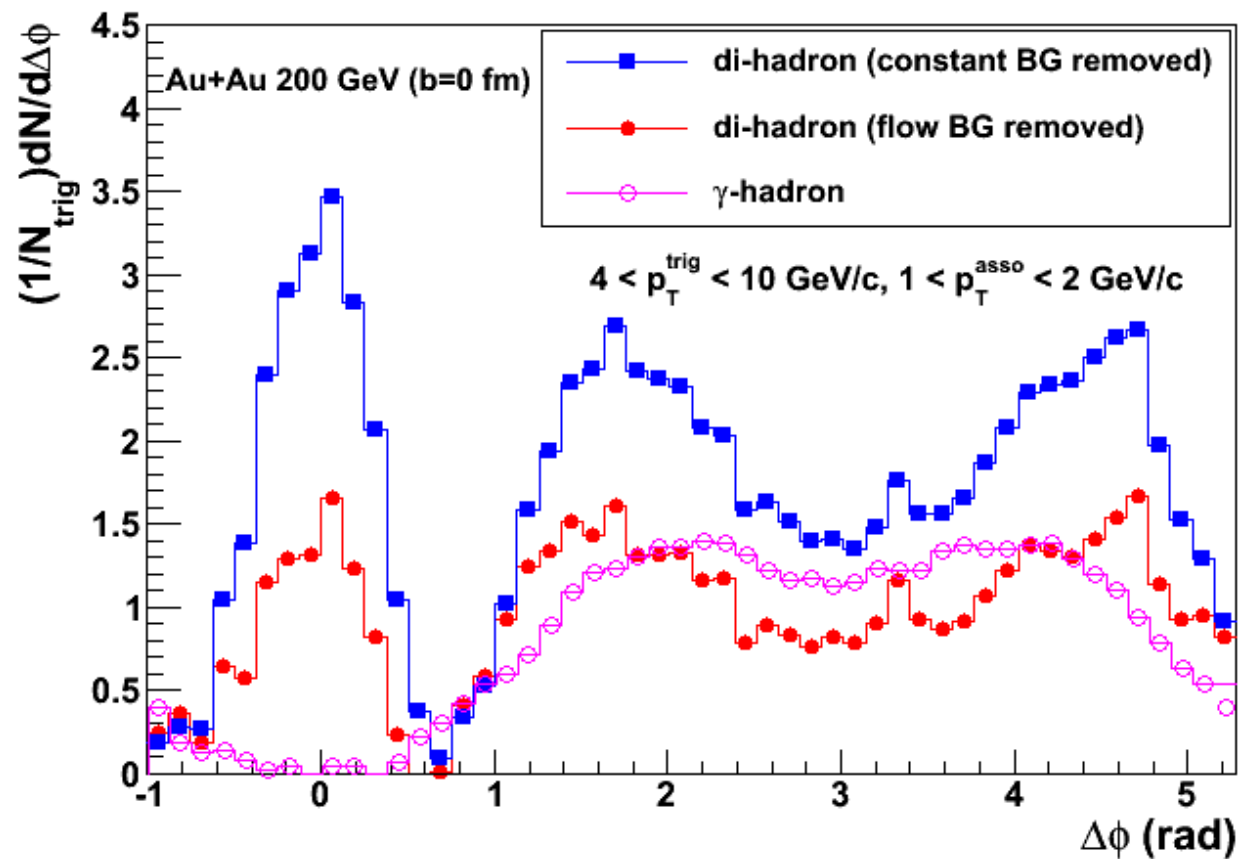
Hydro calculation

# Mach-cone-like excitation



# gamma-hadron correlation

Guo-liang Ma & XNW (2011)



# What I have not talked about

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- Initial gluon structure of nuclei: gluon saturation
- Quarkonium suppression
- Chiral symmetry restoration reflected in properties of vector-mesons in dilepton channel
- Search for critical points through measurement of fluctuations of conserved charges: baryon, electric charges, strangeness
- AdS/CFT: viscosity, jet transport, heavy  $Q$  diffusion

# Summary

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- Heavy-ion collisions can test many properties of QCD
  - Deconfinement phase transition
  - Chiral symmetry restoration
- Current RHIC data indicate formation of strongly interacting QGP
  - High energy density 20 GeV/fm ( $t_0=1$  fm/c) from jet quenching,  $dN/dy$ , radial flow
  - Elliptic flow  $\rightarrow$  early thermalization, low viscosity
  - Parton recombination  $\rightarrow$  partonic matter
  - $J/\Psi$  suppression  $\rightarrow$  deconfinement
- Quantitative study at LHC
  - New discoveries

# New opportunities

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# Back up slides

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