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Close connection to Axiverse project:

Project outline and Targets: Kodama's talk
 Latest results: Yoshino's talk

 black hole superradiance instabilities Helvi's, Vitor's and Oscar's talks

• Very welcome to join the Axiverse sphere!

Towards singularity theorems in asymptotically anti-de Sitter spaces

Akihiro Ishibashi 5 April 2012 at YITP Work in progress w/ Kengo Maeda

Instability in AdS spaces

• Numerical (in part, analytical) results

Bizon – Rostworowski 2011 (massless scalar field, spherical system) Dias – Horowitz – Santos 2011 (vacuum, gravitational d.o.f. excited)

- Non-linear effects
- Oscar's talk : "AdS is linearly stable but non-linearly unstable. Generic small perturbations of AdS become large and eventually form black holes"

So far only a few pieces of supporting evidence and still speculative

Need further study by both numerical and analytical methods to proceed

- One of their motivations: Hawking-Penrose's singularity theorem in closed universe
 - "This theorem does not apply to AdS directly (because spacelike surfaces are not compact) but morally speaking, the negative cosmological constant acts like a confining box for fields inside. So one expects that generic solutions will be singular."

Dias-Horowitz-Santos '11

- Outline:
 - 1. AdS instability
 - 2. Singularity theorems
 - 3. What we wish to show
 - 4. A singularity theorem in AdS
 - 5. Towards more general theorems

Singularity theorems

 "Singularity" in singularity theorem is defined as "incomplete causal" geodesic curve

- Show the existence of such an incomplete curve under
 - 1. Generic, and energy (convergence) conditions
 - 2. Global conditions (Causality)
 - 3. Strong-gravity (trapped region) condition

Theorem 1 - (null geodesic incompleteness)
(1) null convergence
(2) non-compact Cauchy surface
(3) closed trapped surface

- Theorem 2 (timelike or null incompleteness)
 (1) convergence for causal curves
 - (1) convergence for causal curves
 - (2) generic condition
 - (3) chronology condition
 - (4) trapped set
 - (i) compact, achronal set without edge
 - (ii) closed trapped surface
 - (iii) point p s.t. the null geodesics from p are focussed and start to reconverge

- Thoerem 3 (past incomplete causal geodesic)
 (1) convergence for every causal curves
 (2) strong causality
 - (3) some past-dir unit timelike vector W at p s.t.

- Theorem 4 (timelike incomplete geodesic)
- (1) convergence for every causal curve
- (2) compact spacelike hypersurface S (without edge)
- (3) unit normals to S are everywhere converging on S

(2) Can be replaced w/ (2') S is a Cauchy surface

Key notion I @ Global Hyperbolicity_ U (open set): globally hyperbolic ► <u>Def</u>. $: \Leftrightarrow \int \mathcal{D} \forall pair of p. g \in \mathcal{U}$ $I^{(p)} \cap I^{(g)} has compact clourure.$ $\mathcal{U}.$ 3 strong causality. ► importance in physics -P = Cauchy surface in U.

Key notion I @ <u>Conjugate</u> points ▶ <u>Def</u>. Point "r" is conjugate to "p" : (=) = infinitesimally neighboring geodesic from "p" which intersects "r" in (p. 8] attractive. ► Gravity $\frac{d\theta}{ds} = \Theta Rab k^{a} k^{b} \Theta Tab 0^{ab} \Omega \frac{1}{n}$

© <u>Significance</u> of (I) Global Hyperbolicity <u>L</u>(I) Conjugate points. (I) ∀ causally related p. g ∈ U. ∃ timelike (null) geodesic by tween p and g which maximizes the length of causal curves p→8.

(II) If Y contains a pair of conj. p.t. (p.r) => Y can be varied to a longer causal curve

• Hawking-Penrose Theorem shows:

There must be a globally hyperbolic region in which there should be a pair of conjugate points on every causal geodesic curve.

S: closed 3 ochrond surface Pros Ht(s) is NON-compact or show Hts I inextendible causal curve DIS hyperboli globally S H(S)

Let {xn3. {yn3 be sequences of points on V. as in Fig. ∃ timelike geodesic un : xn → yn: maximal length. I fimelike geodesic 11 as a limit curve of 1 µm}. ⊕ If µ: in-complete _ > we are done! ② If µ: complete → µ contains pair conj. pts. (p.r) For sufficiently large n, un would also contain pair conj. p.ts (p.r.) L p contradiction!



What we wish to show

• Starting from No assumption of the trapped (strong-gravity) condition,

Show the occurrence of a trapped set

$$\begin{split} & \bigotimes System. \qquad \{ \text{ spherical} \\ \text{ is perfect fluid.} \\ & \text{ds}^2 = - \int dt^2 + h \, dr^2 + R^2 \, d\Omega^2. \qquad \text{f. h. R: (t.r).} \\ & \text{Tob} = \mu \, \text{VaVb} + P \, (\mathcal{J}_{ab} + \text{VaVb}) \\ & \text{t-coord. along fluid (ines.} \\ & \text{W} \\ & \text{Raychaudhuri equation} \\ & \frac{d\Theta}{ds} = -4\pi(\mu+3P) - |\Lambda| - \sigma^2 - \frac{1}{3}\Theta^2 \\ & + \nabla_a \dot{v}^a \qquad \text{Negative } \Lambda \text{ enhances} \\ & \text{the contraction} \\ & \dot{v}^a \equiv V^c \nabla_c V^a. \end{split}$$

Blackboard talk