Alpha-Clusters in Nuclear Systems

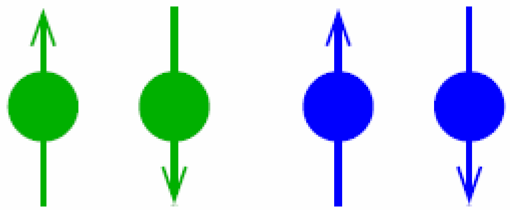
P. Schuck

Y. Funaki, H. Horiuchi, G. Röpke,
A. Tohsaki , W. von Oertzen and T. Yamada

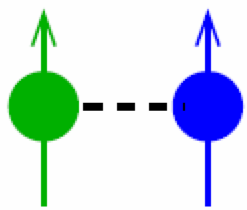
<u>**Contents:**</u>

- **General Aspects of Nuclear Clustering**
- **α－Particles and their Condensation**
- **^{8}Be and Hoyle State in ^{12}C***
- **Extensions to Heavier Nuclei: ^{16}O**
- **α － de－excitation of Compound States**
- **α’s in Compact Stars**
- **Conclusions, Outlook**

Clusters important aspect and richness of nuclear systems due to 4 Fermions :

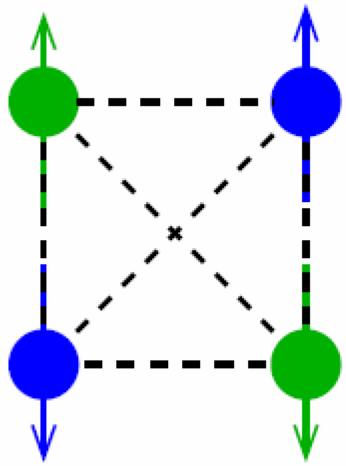


Dimer :

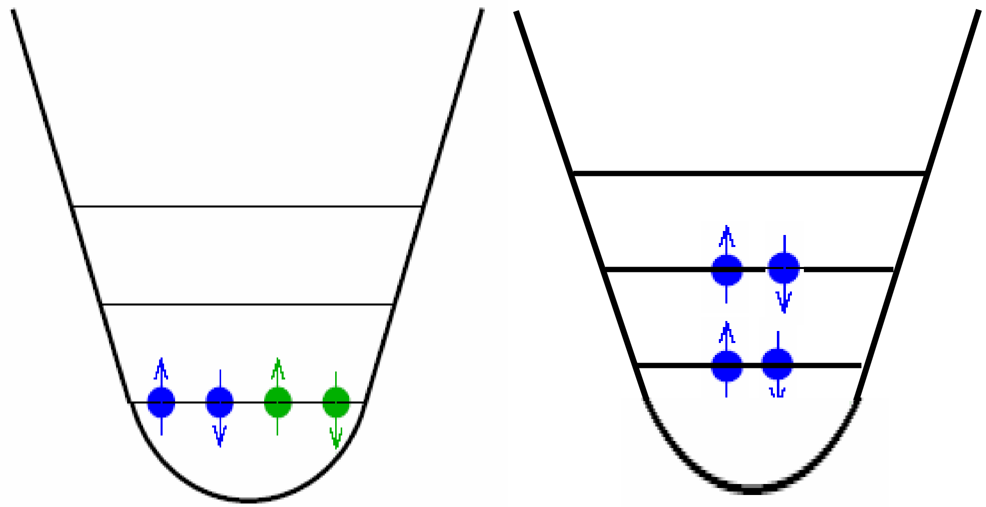


$\frac{E}{A} = 1$ MeV

Quartet :



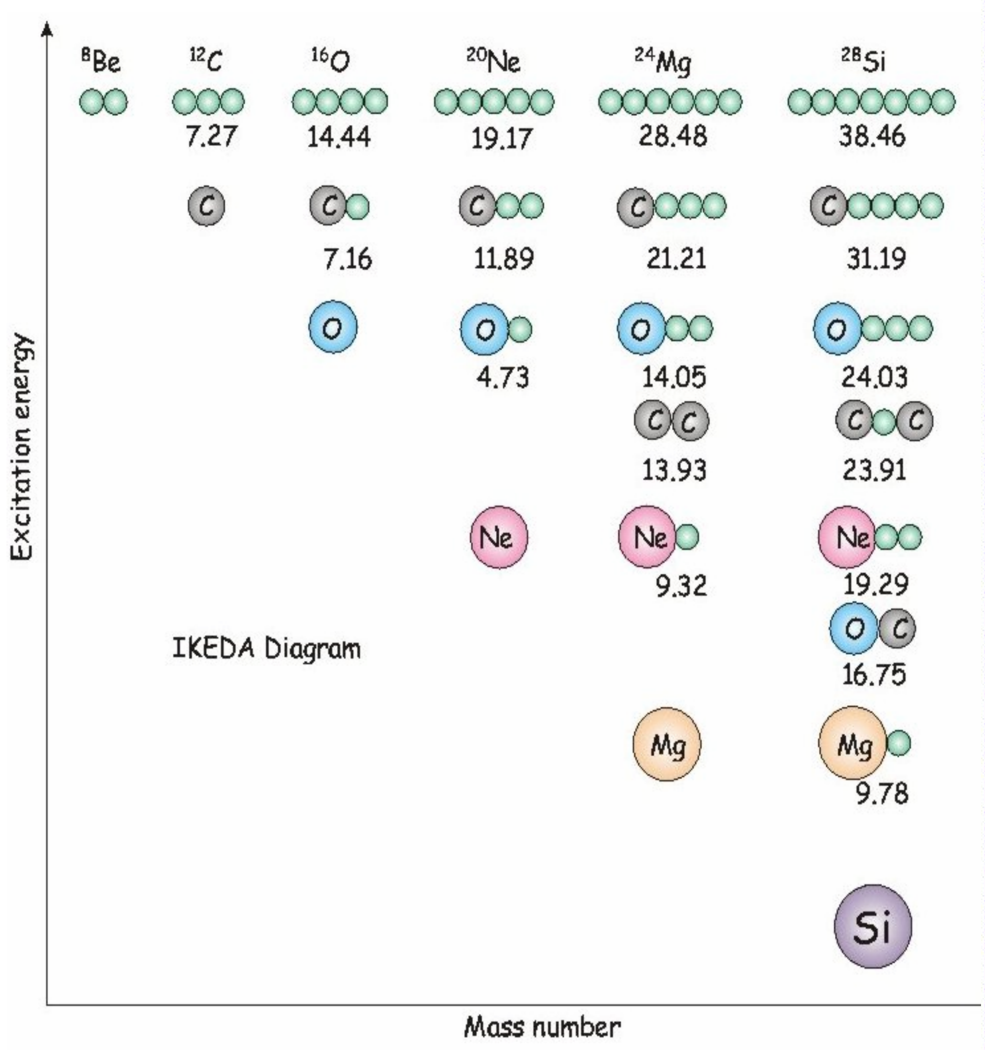
$\frac{E}{A} = 7$ MeV, $E^* = 20$ MeV



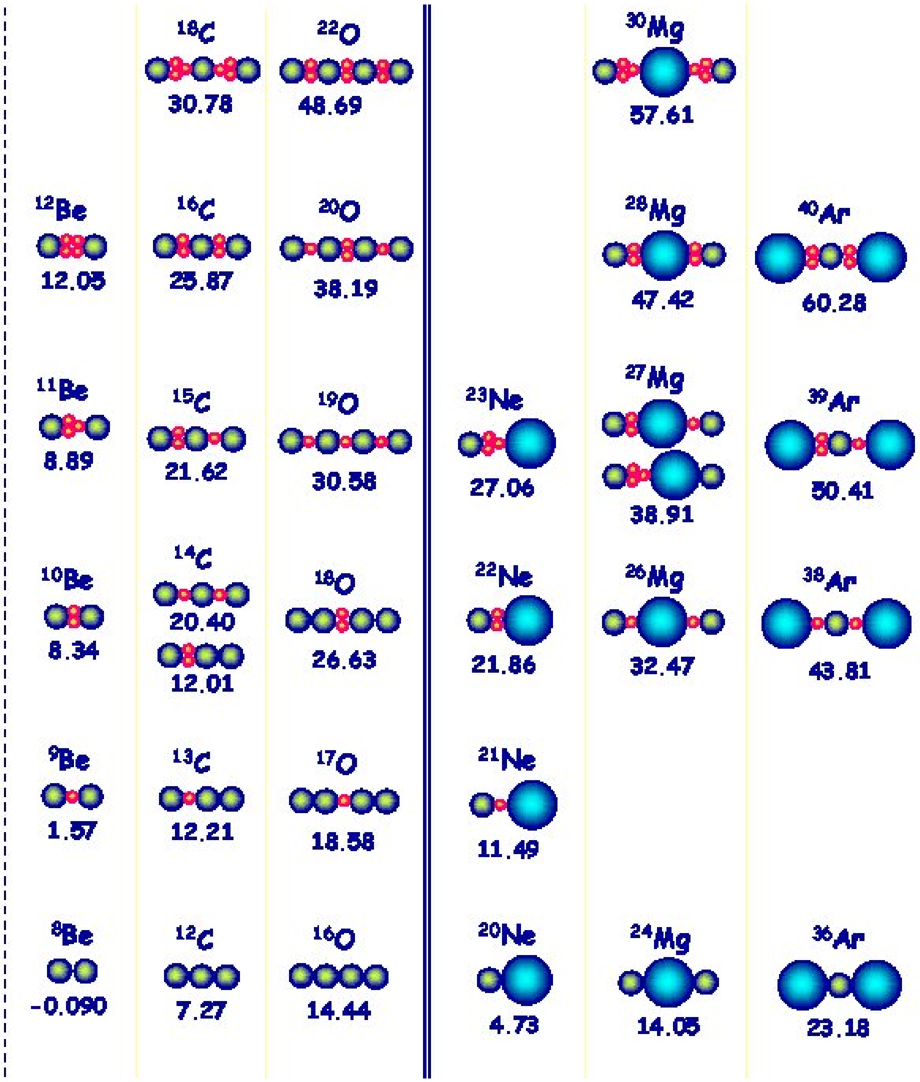
Proposal :

Trapping of 4 different species of Fermionic atoms.

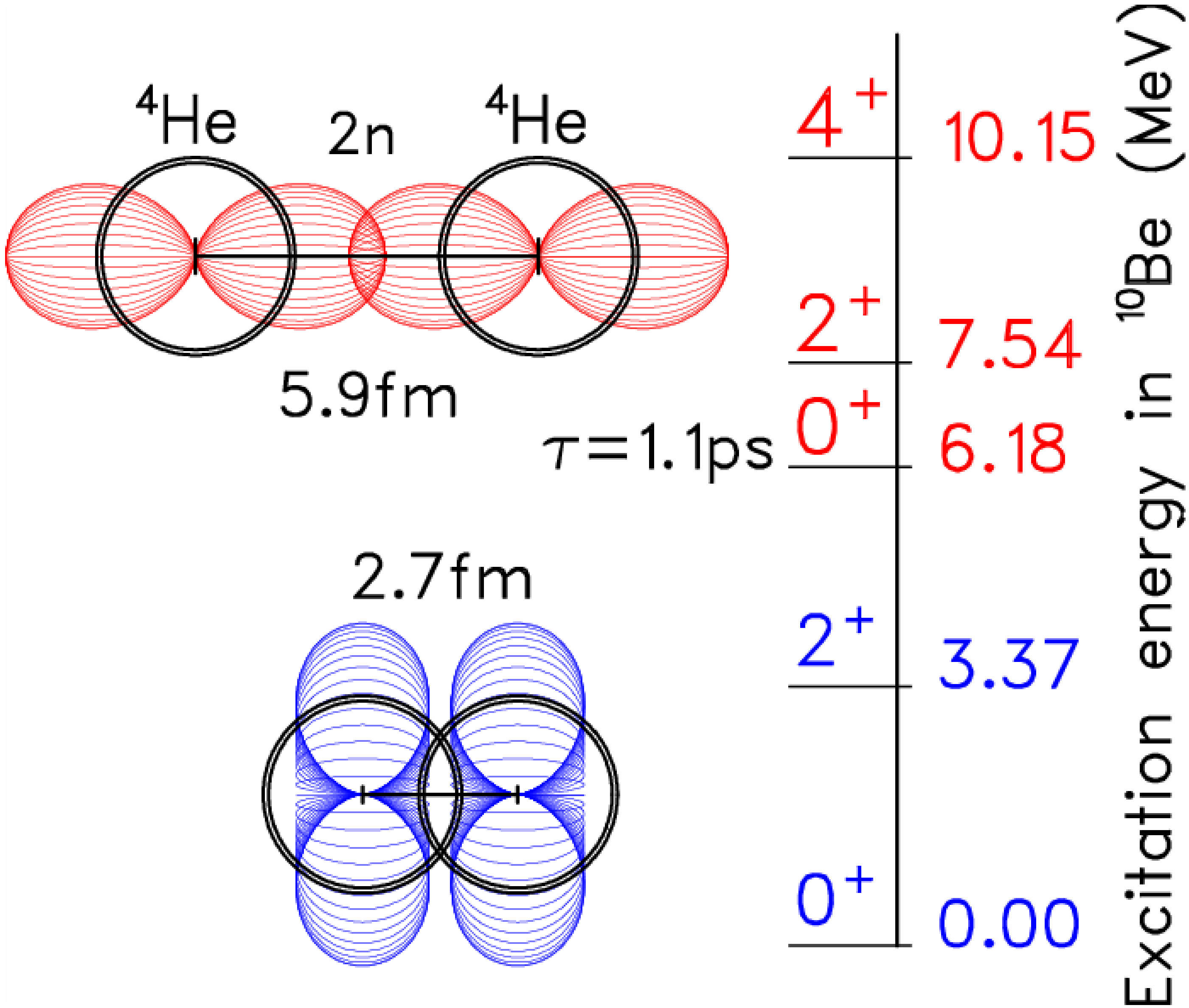
Ikeda



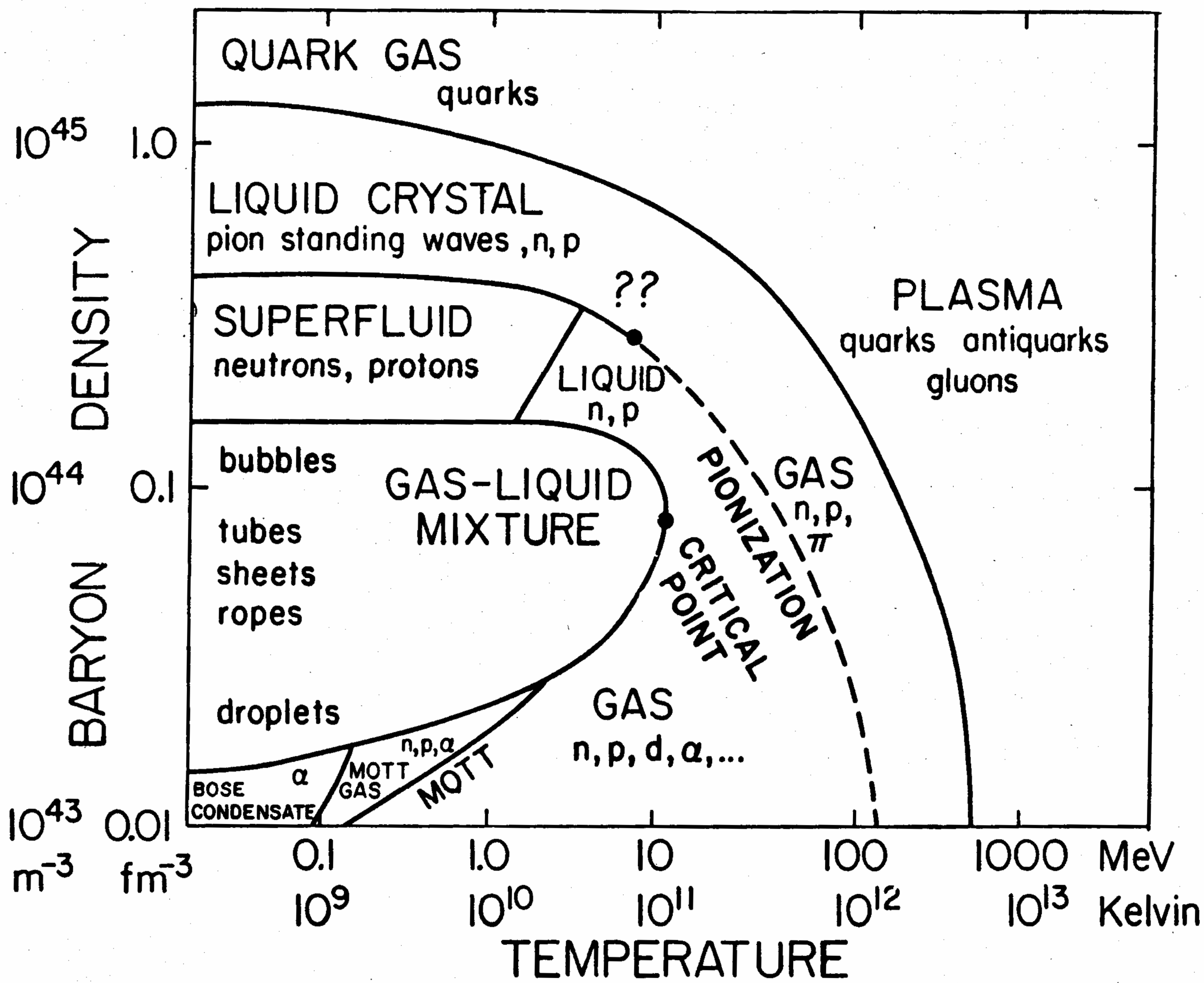
von Oertzen



H. G. Bohlen, M. Freer



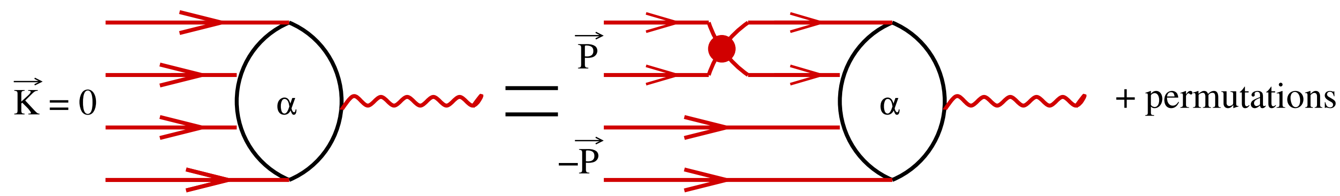
P. Siemens

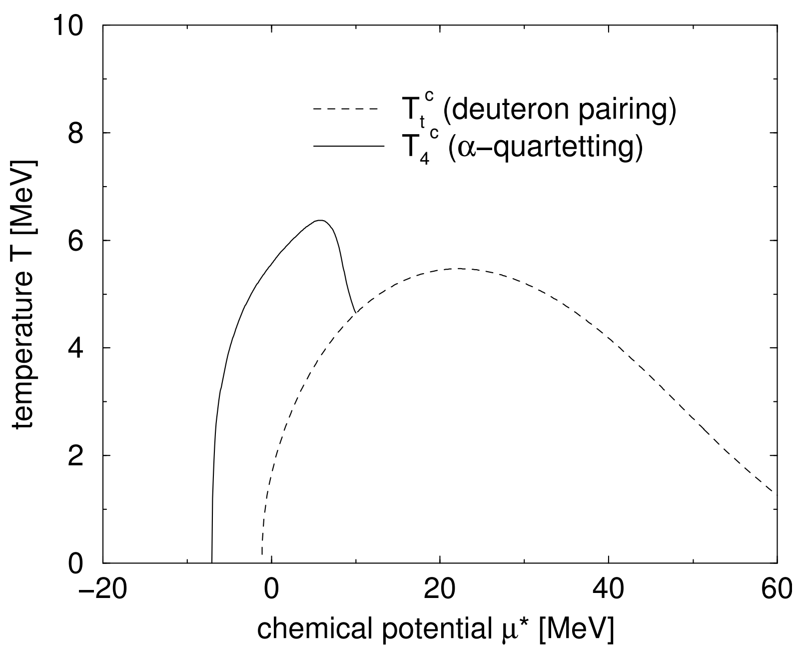


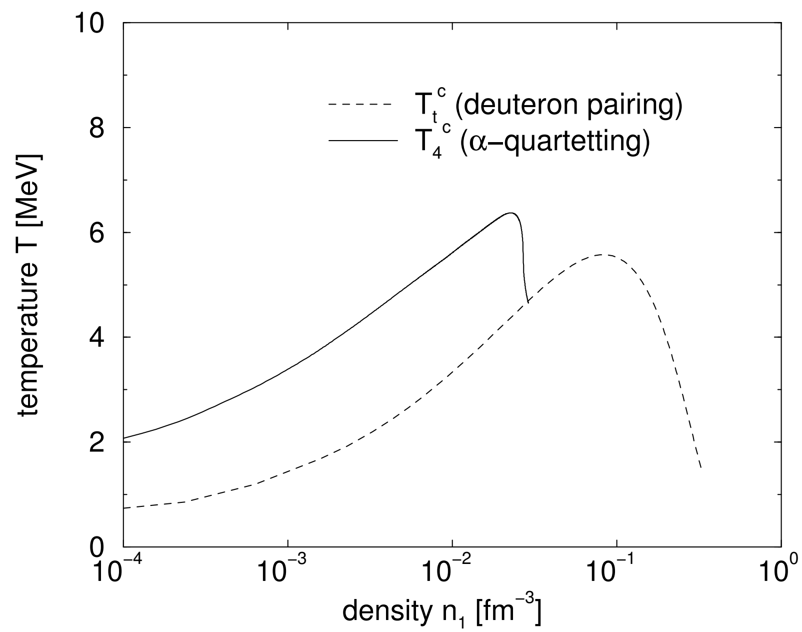
QUARK GAS
quarks
LIQUID CRYSTAL
pion standing waves , n, p
SUPERFLUID
neutrons, protons
??
LIQUID
n, p
PLASMA
quarks antiquarks
gluons
bubbles
tubes
sheets
ropes
GAS-LIQUID
MIXTURE
PIONIZATION
CRITICAL POINT
GAS
n, p, π
droplets
GAS
n, p, d, α, ...
BOSE CONDENSATE
α
MOTT GAS
n,p,α
MOTT
BARYON DENSITY
10^45 1.0
10^44 0.1
10^43 0.01
m^-3 fm^-3
0.1 1.0 10 100 1000 MeV
10^9 10^10 10^11 10^12 10^13 Kelvin
TEMPERATURE

- α-Particle Condensation : G. Röpke, M. Beyer

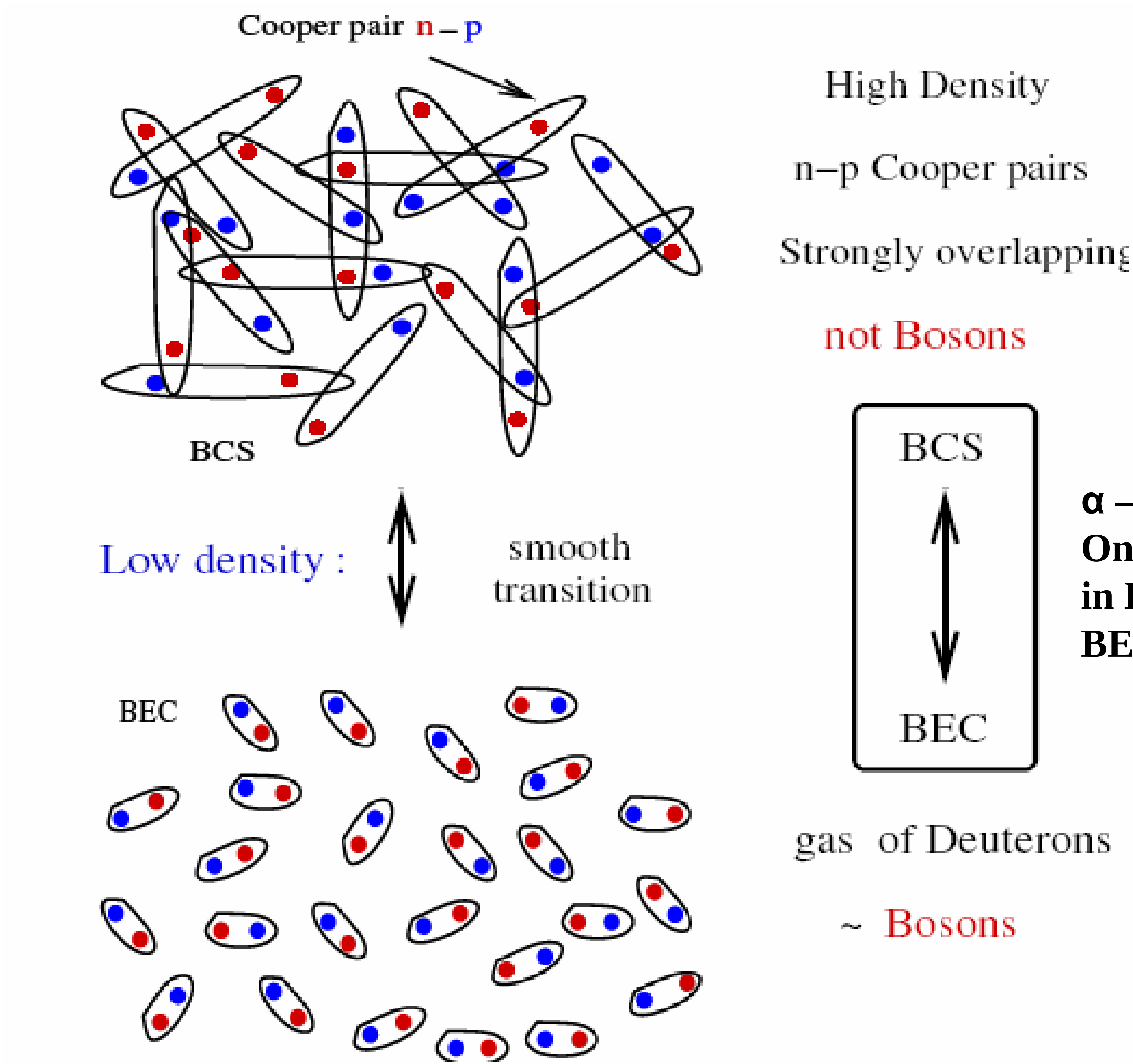
$$(4\mu - \epsilon_1 - \epsilon_2 - \epsilon_3 - \epsilon_4)\,\psi_{1234} = \quad (1 - f_1 - f_2)\sum_{1'2'} v_{121'2'}\psi_{1'2'34}$$
$$+\text{permutations}$$







α-Condensation only at very low density !

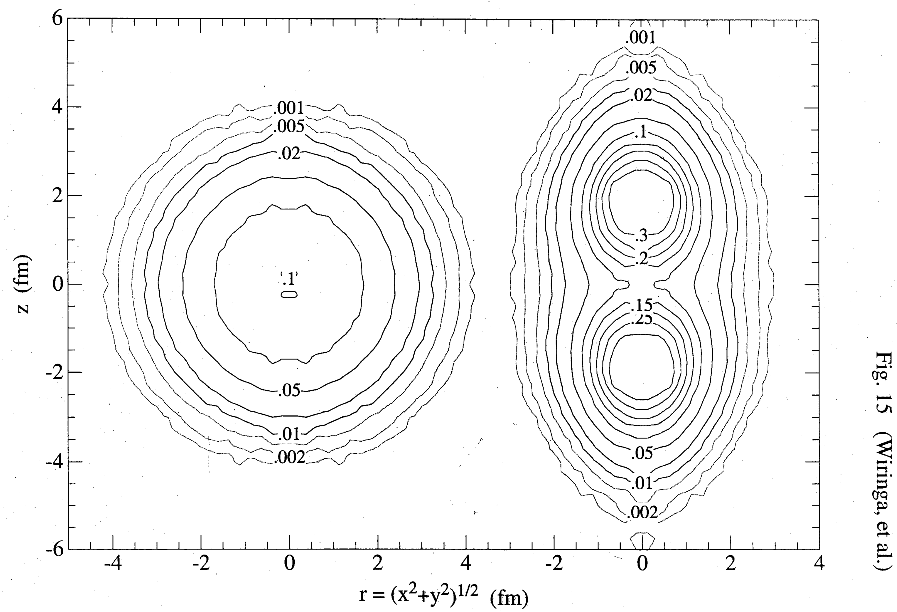


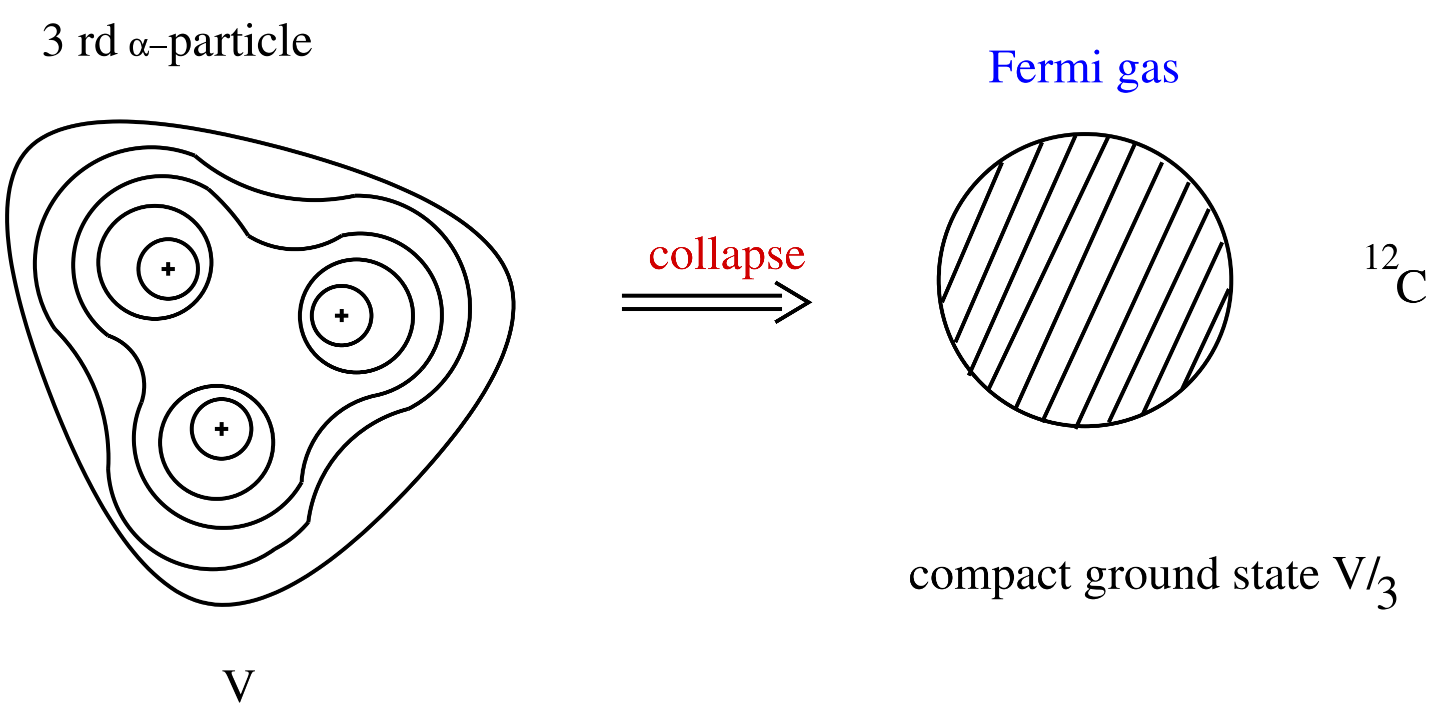
α – Particles
Only Exist
in Low Density
BEC Phase

Finite nuclei ?

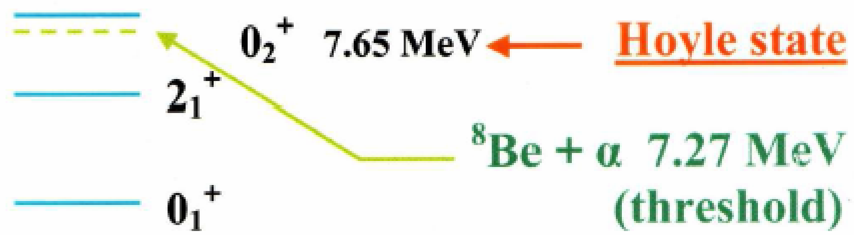
Exact ^{8}Be :

Density : $\frac{\rho_0}{3}$





Does a dilute 3α ${}^{12}C^{*}$ state exist ?
Similar to ${}^{8}Be + \alpha$?



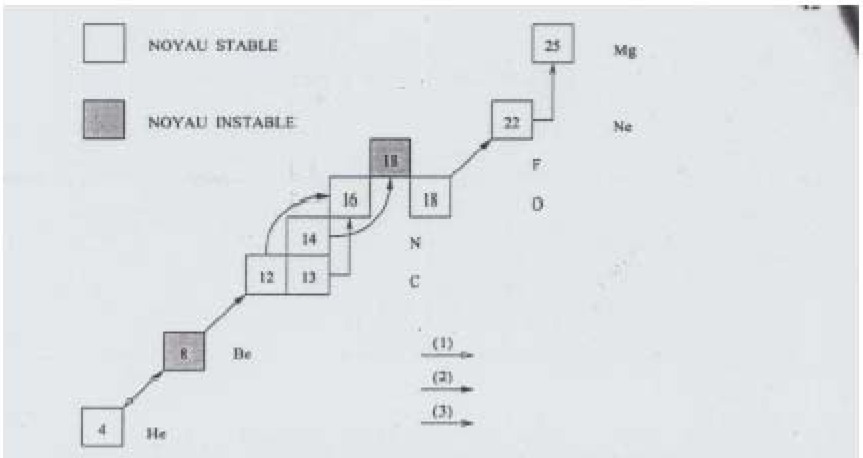


Figure 26. Réactions principales impliquées dans la combustion non explosive de He. Les flèches de type (1), (2) et (3) se rapportent respectivement à des réactions (γ, α), (α, γ), et (α, n)

At $T = 10^8 K$ helium burning
thermal equilibrium

$$\alpha + \alpha + \alpha \;\rightarrow\; {}^{8}Be \;\rightarrow\; \alpha \;\rightarrow\; {}^{12}C^{*} \quad O_2^+$$

O_2^+ : dilute 3α state hypothesis !

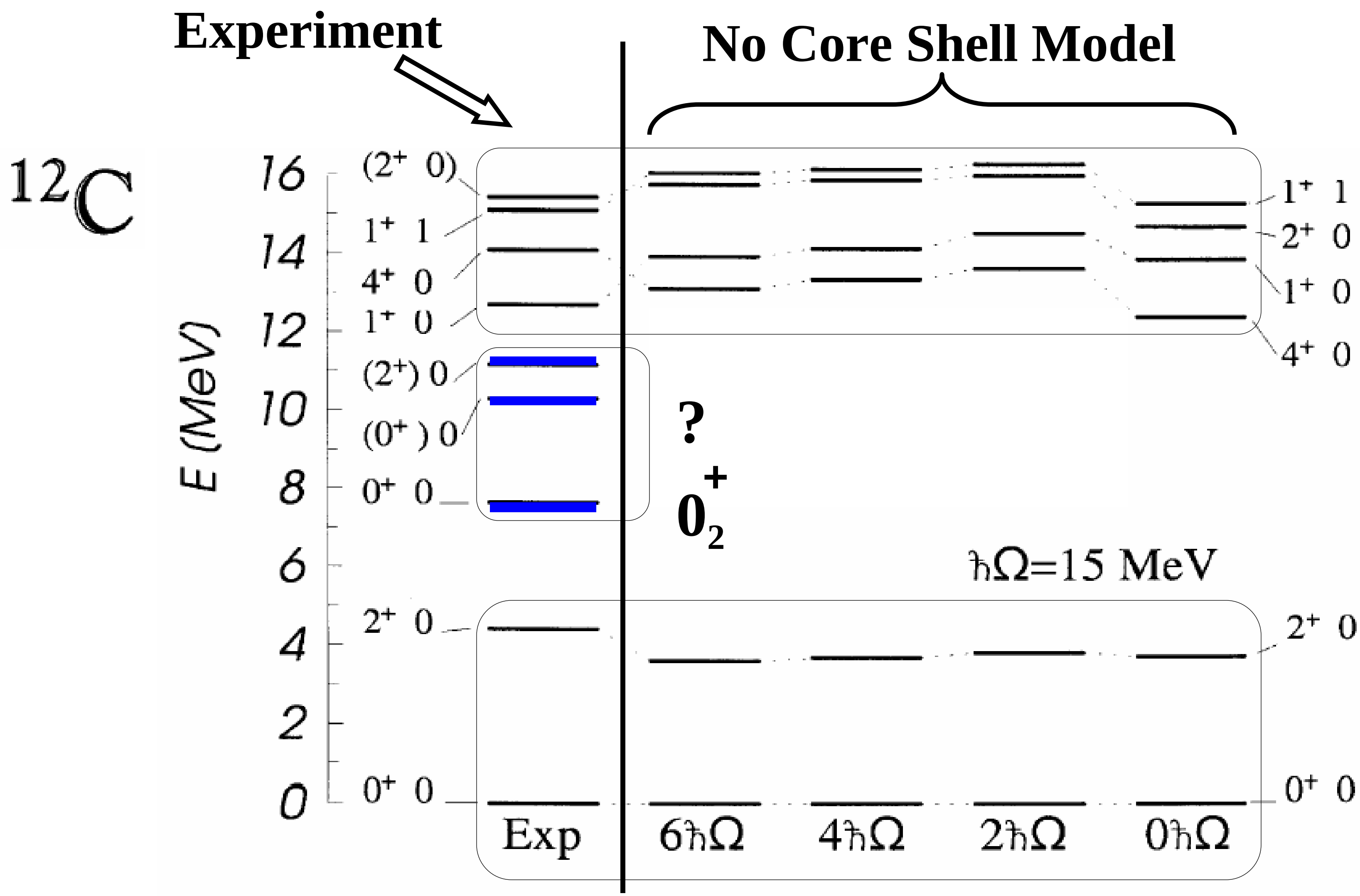
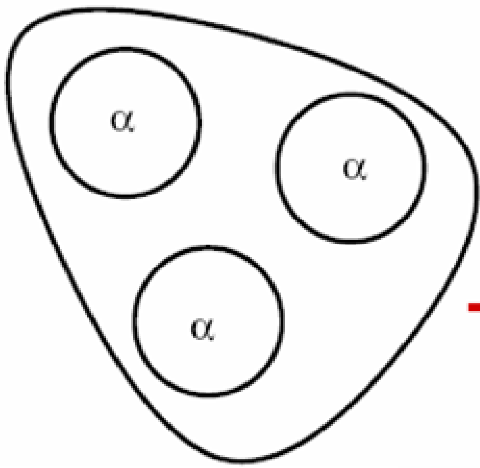


Figure 4. *Experimental and NCSM excitation spectra for ^{12}C for different model space sizes.*

it seems impossible to get Hoyle state from shell model calculation !

45 MeV B. Barret

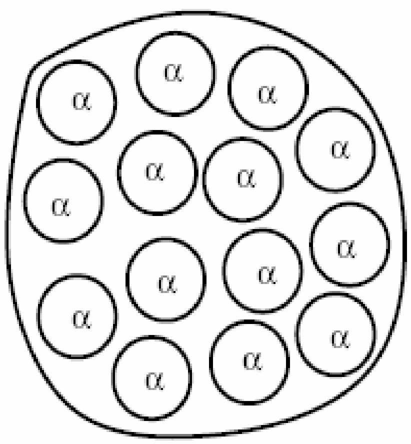
If O_2^+ in ^{12}C dilute $\alpha-$state

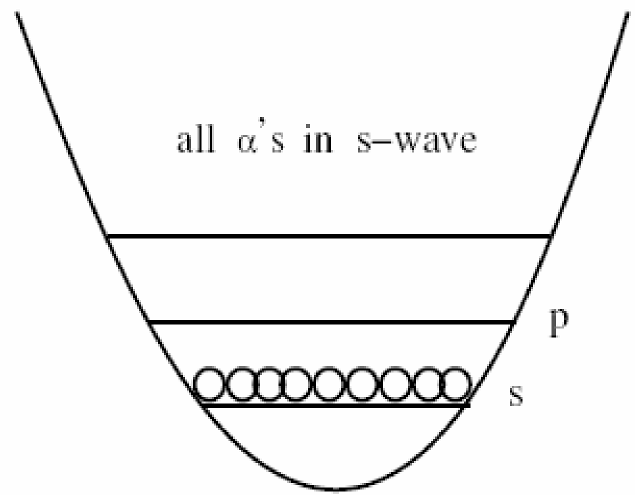


then α-condensate

infinite matter $\rho_{\text{crit}} \sim \frac{\rho_0}{3}$

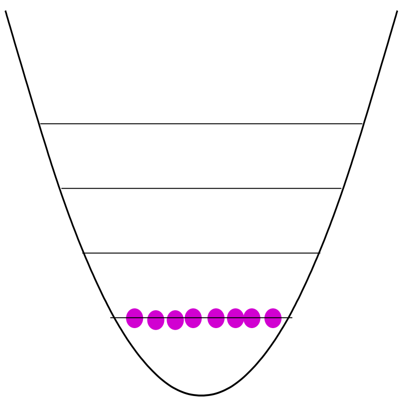
Conjecture: all n.α nuclei possess exited nα condensed state



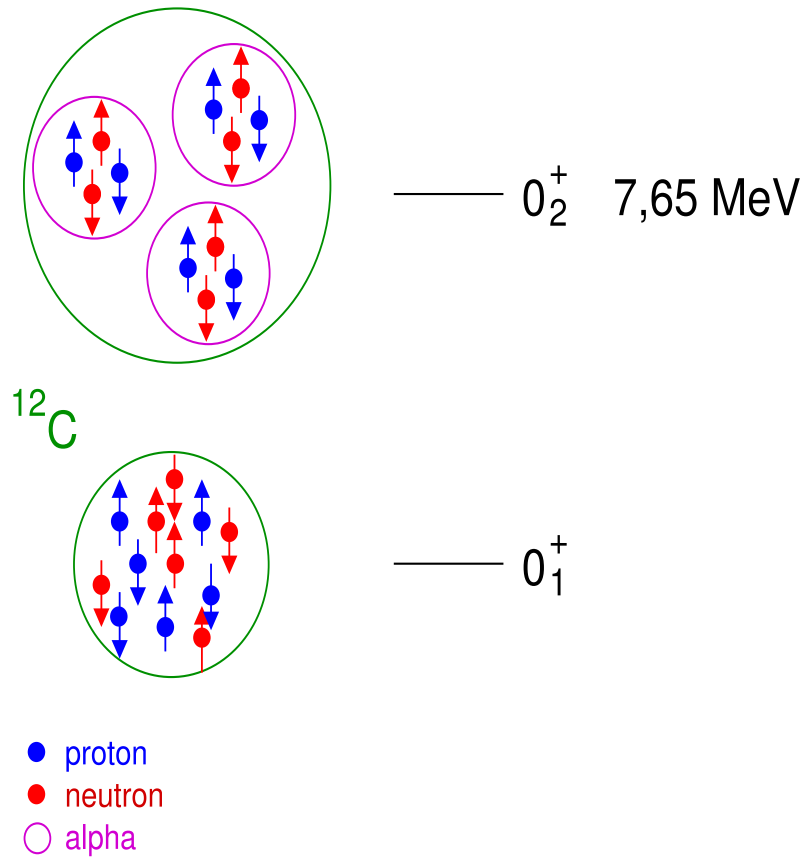


Analogy with atoms in traps ! $\rho(r) = N|\phi_0(r)|^2$ $N = 10^6$

Bosons

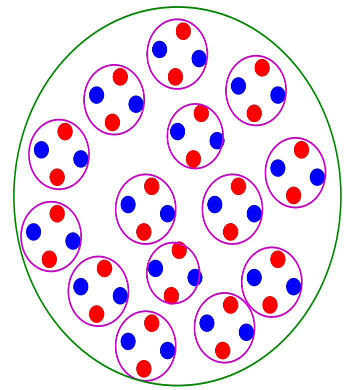


Back to nuclei

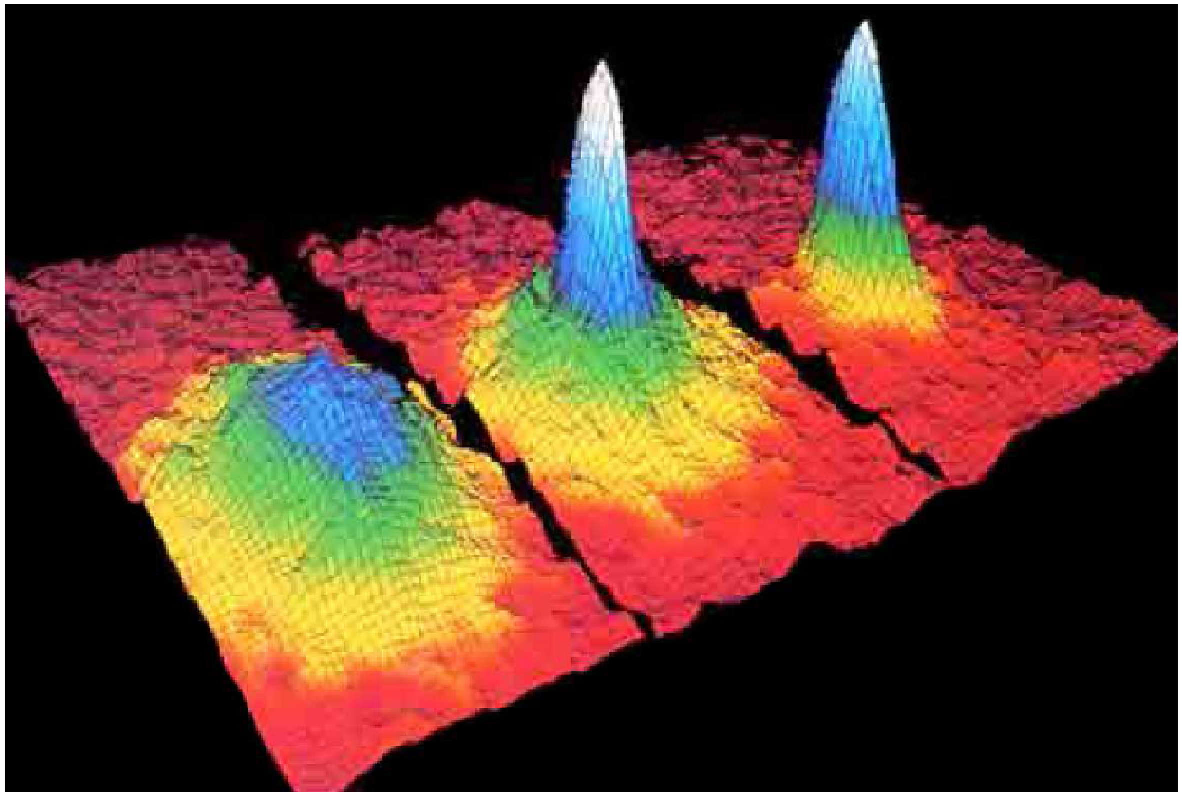


many α's

$\rightarrow$ condensate



strong cluster phenomena in lighter nuclei



Theoretical Description

Ideal Bose condensate : $|0\rangle = b_0^\dagger b_0^\dagger \cdots b_0^\dagger |vac\rangle$

α-particle condensate : $|\Phi_{\alpha C}\rangle = C_\alpha^\dagger C_\alpha^\dagger \cdots C_\alpha^\dagger |vac\rangle$

In r-space :

$$\langle \vec{r}_1, \vec{r}_2, \cdots, \vec{r}_{4n} | \Phi_{\alpha C} \rangle = \mathcal{A}\left\{ \Phi(\vec{r}_1, \vec{r}_2, \vec{r}_3, \vec{r}_4) \Phi(\vec{r}_5, \vec{r}_6, \vec{r}_7, \vec{r}_8) \cdots \Phi(\vec{r}_{4n-3}, \vec{r}_{4n-2}, \vec{r}_{4n-1}, \vec{r}_{4n}) \right\}$$

In comparison with pairing :

$$\langle \vec{r}_1, \vec{r}_2, \cdots | \mathrm{BCS} \rangle = \mathcal{A}\left\{ \Phi(\vec{r}_1, \vec{r}_2) \Phi(\vec{r}_3, \vec{r}_4) \cdots \right\}$$

Variational ansatz for $\Phi\left(\vec{r}_1,\vec{r}_2,\vec{r}_3,\vec{r}_4\right)$: $\Phi\left(\vec{r}_1,\vec{r}_2,\vec{r}_3,\vec{r}_4\right)=e^{-\frac{2}{B^2}\vec{R}^2}\phi_\alpha\left(\vec{r}_i-\vec{r}_j\right)$

Center of mass : $\vec{R}=\frac{1}{4}\left(\vec{r}_1+\vec{r}_2+\vec{r}_3+\vec{r}_4\right)$

Intrinsic α-wave function :

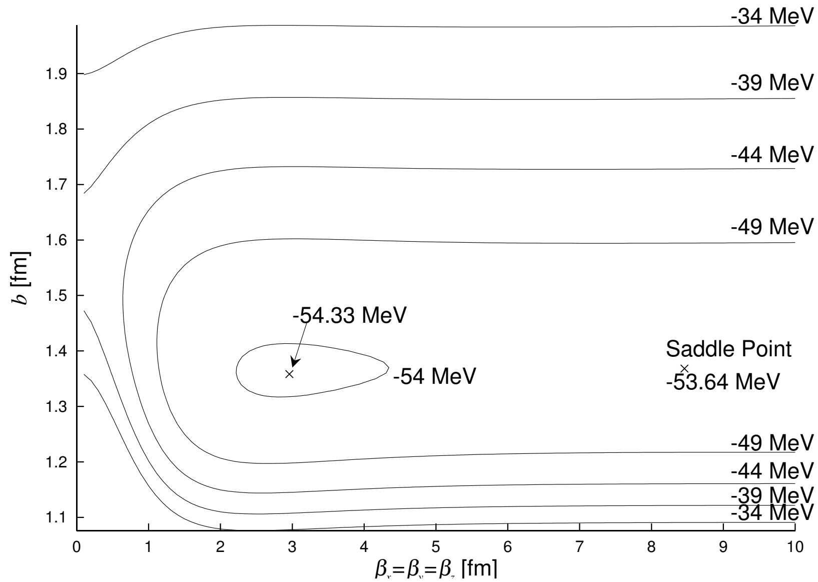
$$\phi_\alpha\left(\vec{r}_i-\vec{r}_j\right)=e^{-\frac{1}{8b^2}\left\{\left(\vec{r}_4-\vec{r}_1\right)^2+\left(\vec{r}_4-\vec{r}_2\right)^2+\left(\vec{r}_4-\vec{r}_3\right)^2+\cdots\right\}}$$

Two variational parameters : $B,\ b$

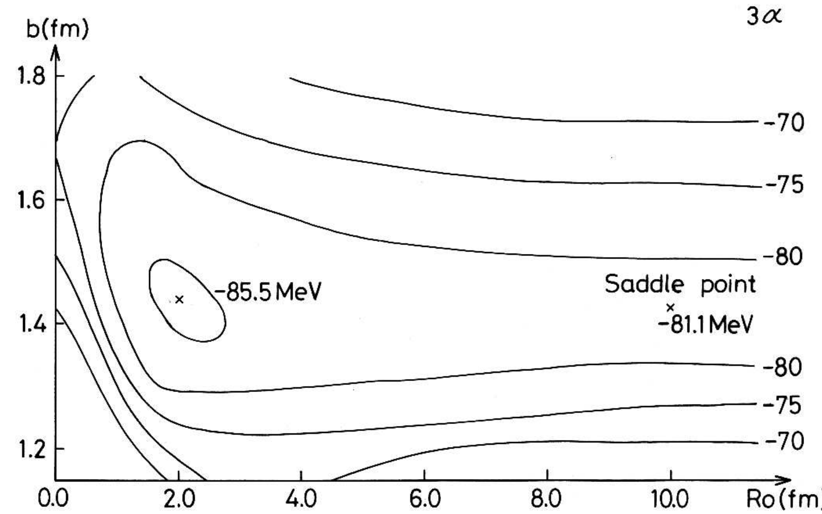
Two limits : $B=b$ $\quad\left|\Phi_{\alpha\mathrm{C}}\right\rangle=$ Slater determinant

$B\gg b$ $\quad\left|\Phi_{\alpha\mathrm{C}}\right\rangle=$ gas of independent α-particles

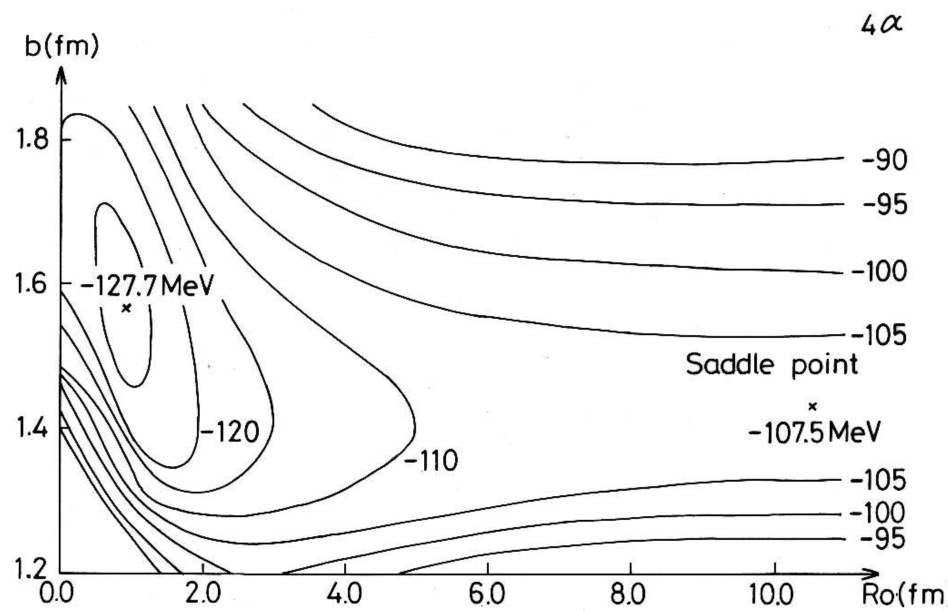
Two dimensional surface : $E\left(B,b\right)=\dfrac{\left\langle\Phi_{\alpha\mathrm{C}}\left|H\right|\Phi_{\alpha\mathrm{C}}\right\rangle}{\left\langle\Phi_{\alpha\mathrm{C}}\middle|\Phi_{\alpha\mathrm{C}}\right\rangle}$


-34 MeV
-39 MeV
-44 MeV
-49 MeV
-54.33 MeV
-54 MeV
Saddle Point
-53.64 MeV
-49 MeV
-44 MeV
-39 MeV
-34 MeV
b [fm]
1.1
1.2
1.3
1.4
1.5
1.6
1.7
1.8
1.9
0
1
2
3
4
5
6
7
8
9
10

^{8}B


3α
b(fm)
1.8
1.6
1.4
1.2
-70
-75
-80
-85.5 MeV
Saddle point
-81.1 MeV
-80
-75
-70
0.0
2.0
4.0
6.0
8.0
10.0
Ro(fm)

^{12}C


4α
b(fm)
1.8
1.6
1.4
1.2
-90
-95
-100
-105
-127.7 MeV
-120
-110
Saddle point
-107.5 MeV
-105
-100
-95
0.0
2.0
4.0
6.0
8.0
10.0
Ro(fm)

^{16}O

Hamiltonian :

$$\begin{array}{ccccccccc} H & = & T & + & V_{N-N} & + & V_C & + & V_{N-N-N} \\ & & \text{Kin. energy} & & \text{Gaussien} & & \text{Coulomb} & & \text{Gaussian} \end{array}$$

Quantization of energy surface $E(B, b)$:

Force : A. Tohsaki ~ 1990 no adjustable parameters !

Hill-Wheeler : $|\psi\rangle = \sum_B f_B |\Phi_{\alpha C}(B)\rangle$

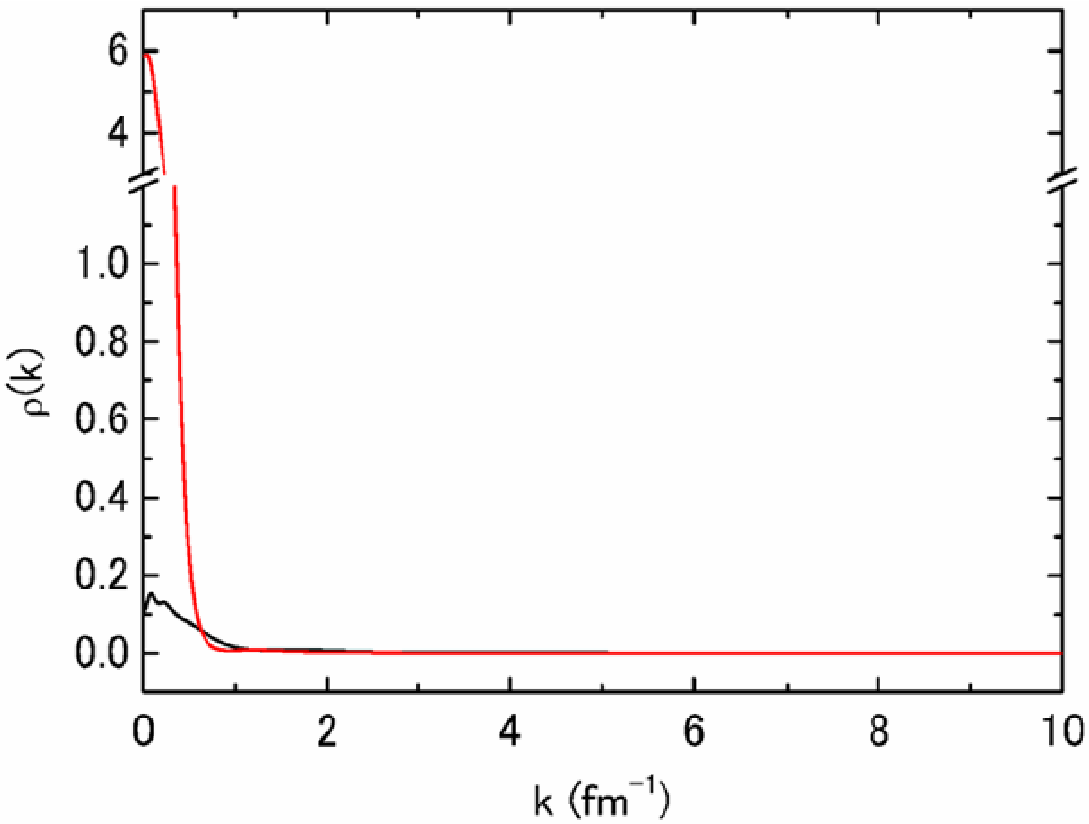
Without adjustable parameters :

$$\begin{array}{rcl} ^{12}C: \quad (E_{O_2^+} - E_{3\alpha}) & = & \text{Theory} + 0.50 \text{ MeV} \\ & & \text{Exp.} + 0.38 \text{ MeV} \\ ^{16}O: \quad (E_{O_5^+} - E_{4\alpha}) & = & \text{Theory} - 0.70 \text{ MeV} \\ & & \text{Exp.} - 0.44 \text{ MeV} \end{array}$$

r.m.s. of O_2^+ in ^{12}C 3.83 fm

ground state 2.40 fm

Momentum Distribution



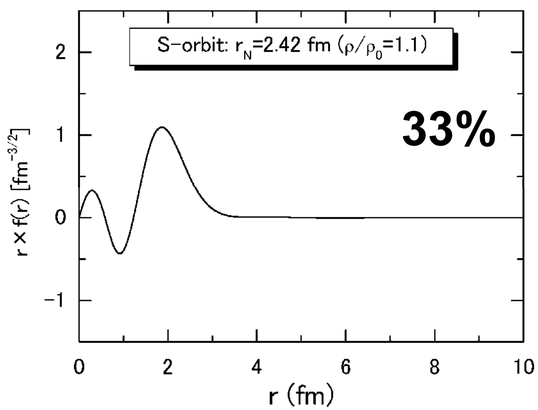
$$\frac{V_{O_2^+}}{V_{O_1^+}} \sim 3-4$$

	$Exp.$	$Cal.$
$M(O_2^+ \rightarrow O_1^+)$	$5.4 \pm 0.2 fm^2$	6.7
$R_{rms}(O_1^+)$	$2.43 fm$	2.40
$R_{rms}(O_2^+)$		3.47

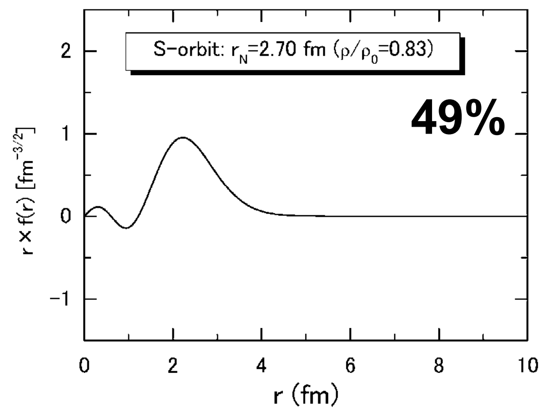
Radial behavior of *S*-wave α orbit vs. R_{rms}

R_{rms}=2.43 fm → 4.84 fm (ρ/ρ_0=**1.1** → **0.14**)

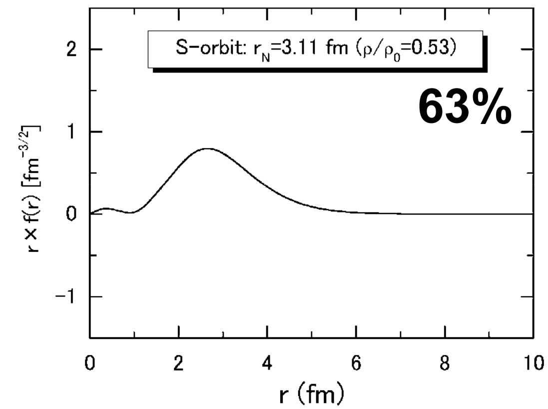
R_{rms}=2.43 fm (ρ/ρ_0=1.1)



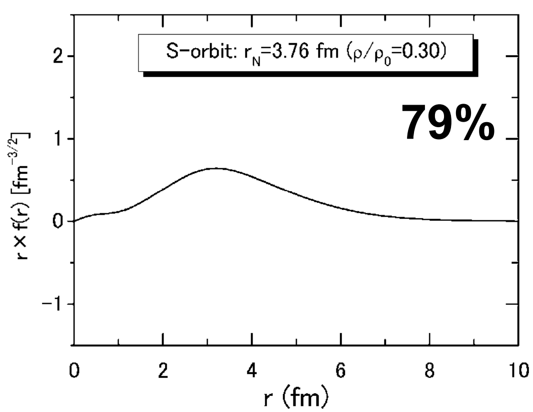
R_{rms}=2.70 fm (ρ/ρ_0=0.83)



R_{rms}=3.11 fm (ρ/ρ_0=0.53)



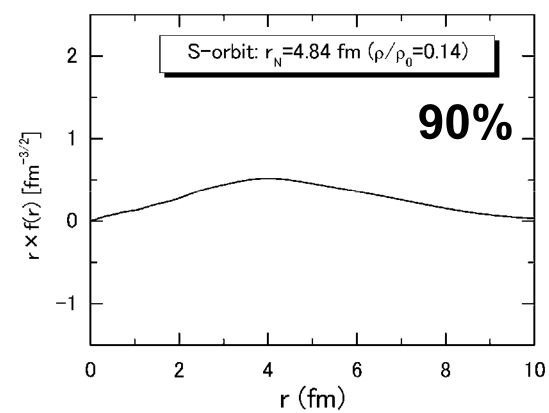
R_{rms}=3.76 fm (ρ/ρ_0=0.30)

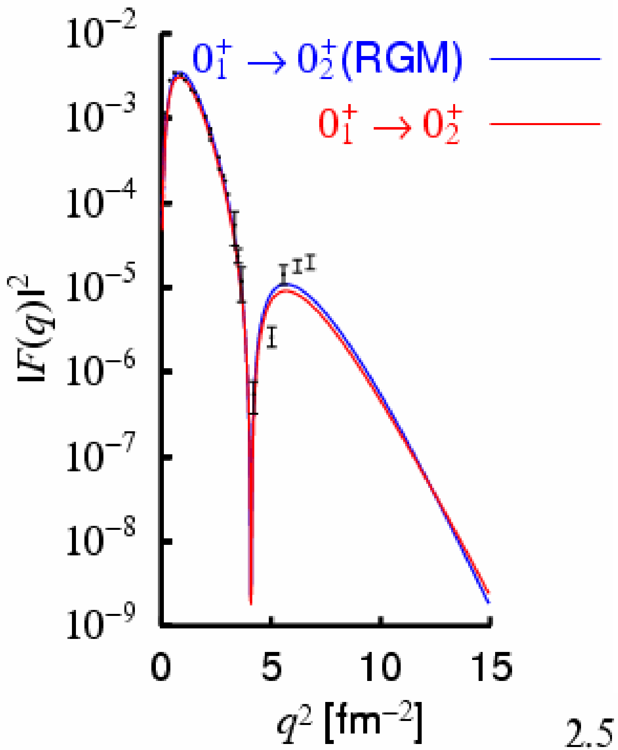


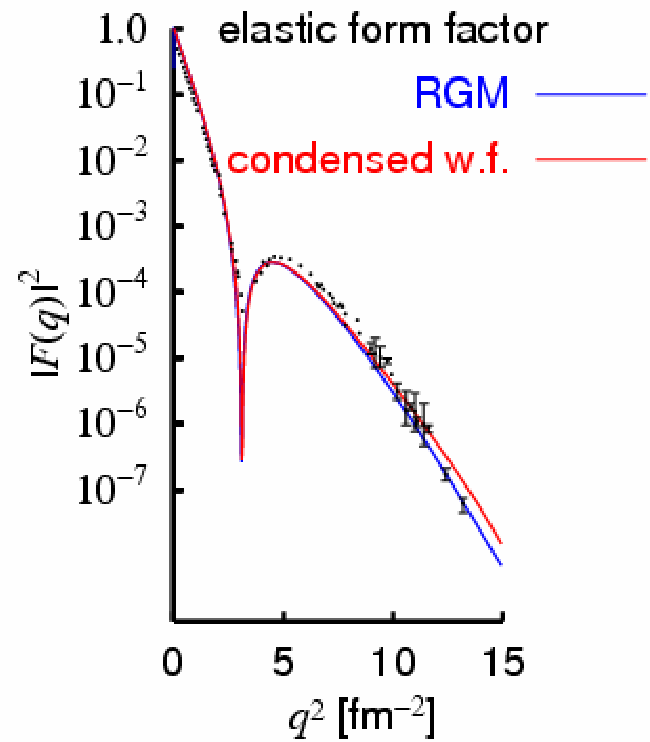
Increasing R_{rms}, we see smooth change from 2*S* to 0*S* orbit.

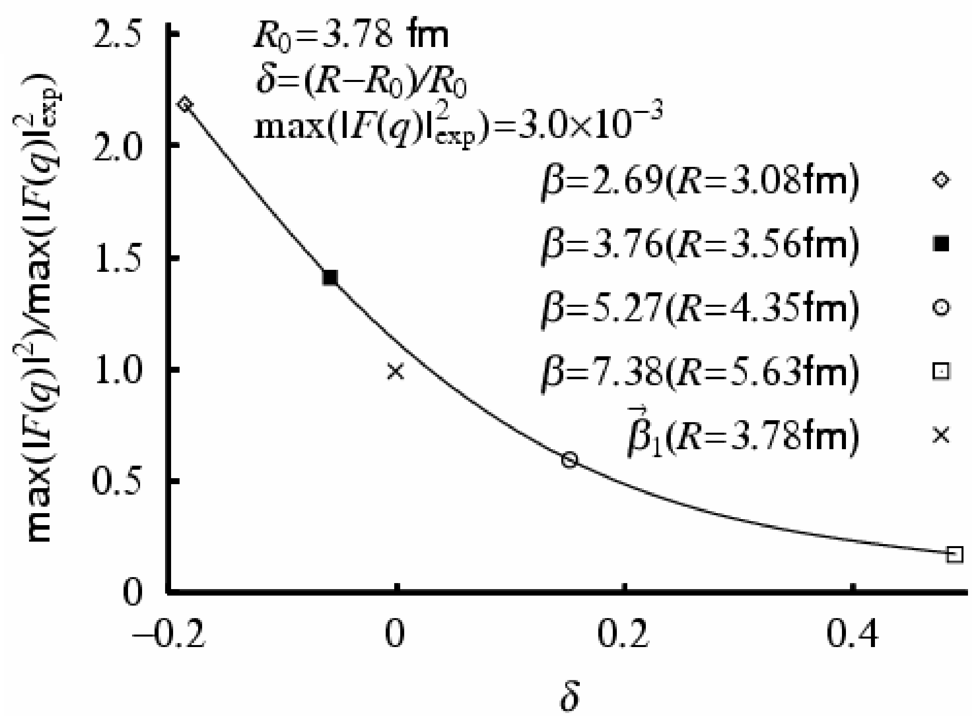
Yamada & Schuck, EPJA 2005

R_{rms}=4.84 fm (ρ/ρ_0=0.14)









R : r.m.s. radius

$\vec{\beta}_1 \coloneqq (\beta_x = \beta_y,\ \beta_z)$
$= (5.27\ \mathrm{fm}, 1.37\ \mathrm{fm})$

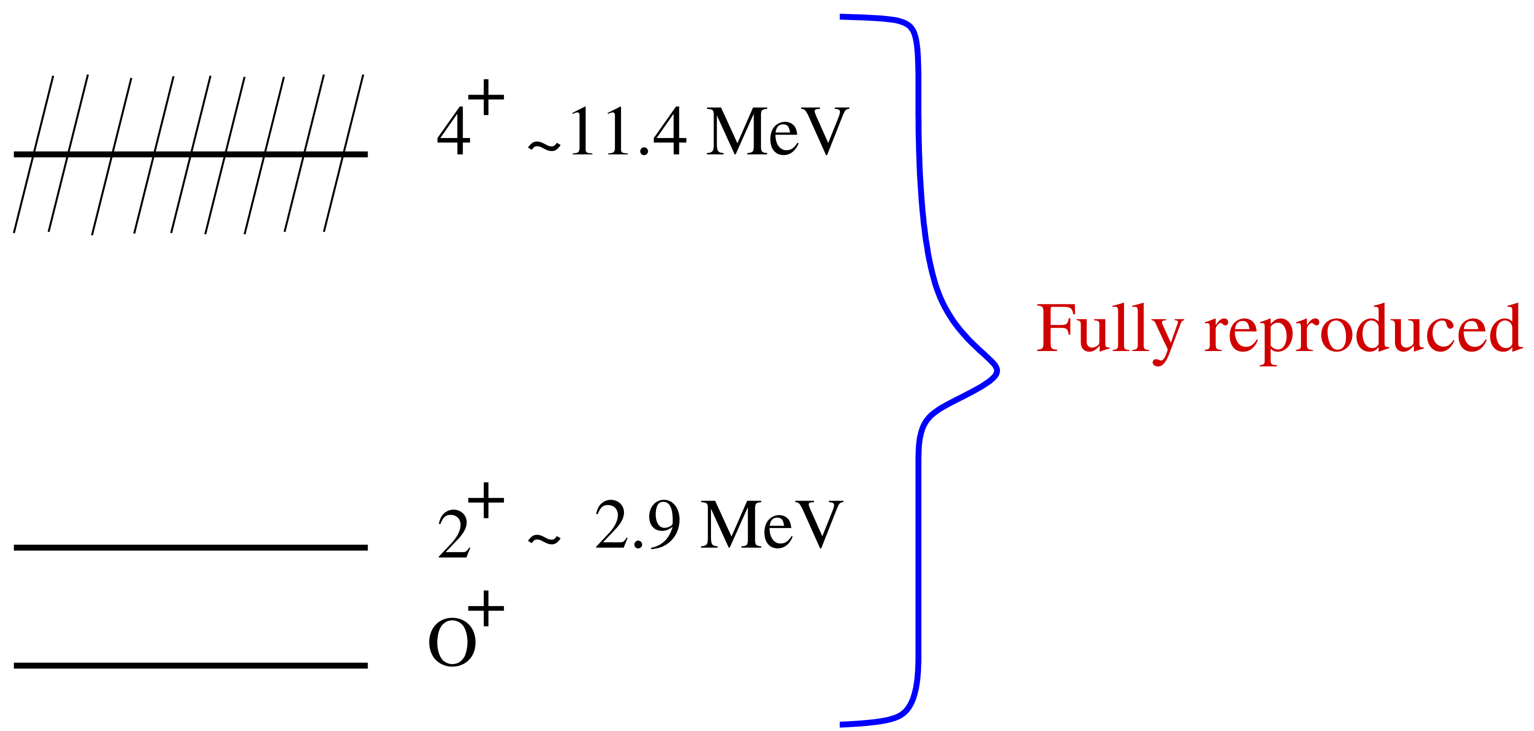
Some more numbers :

^{12}C :

	Theory	Exp.
O_1^+	-89.52	-92.16
O_2^+	-81.79	-84.51
	7.73	7.65

Threshold states $E - E_{\text{thresh.}}$	^8Be	^{12}C	^{16}O	^{20}Ne	
	2α	3α	4α	5α	
	O_1^+	O_2^+	O_3^+	O_4^+	theory
	-0.17	0.50	-0.7	1.8	theory
	O_1^+	O_2^+	O_5^+	?	experimental
	0.09	0.38	-0.44	?	experimental

Spectrum of ^{8}Be :



^{12}C : Second excited 2^+ : 2_2^+

It has been discovered recently by Itoh *et al.*
2.6 MeV above 3 α rhreshold
Width ~ 1 MeV : resonance in continuum

Theory : We start with deformed α condensate state :

$$\Phi_{n\alpha} \propto \mathcal{A} \prod_{i=1}^{n} \exp\left\{ -\frac{2\,X_{ix}^2}{B_x^2} - \frac{2\,X_{iy}^2}{B_y^2} - \frac{2\,X_{iz}^2}{B_z^2} \right\} \Phi_{\alpha_i}$$

Then projection on good angular momentum

Then Hill Wheeler or GCM
For width : ACCC method

	Position	Width
Experiment :	2.6 ± 0.3 MeV	1.0 ± 0.3 MeV
Theory :	2.1 MeV	0.64 MeV

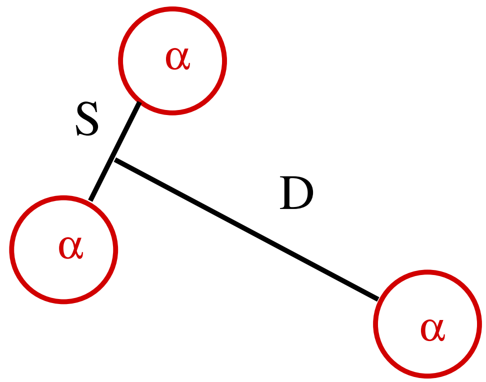
With in error bars !

RMS : $4.43\ fm$

$$\frac{V_{2_2^+}}{V_{O_1^+}} \sim 8 \quad !!$$

α – halo !

Internal structure :



Extremely dilute 3α state

Suggests a pure Boson picture $|\phi_0\rangle = b_0^+ b_0^+ b_0^+ \ldots |vac\rangle$

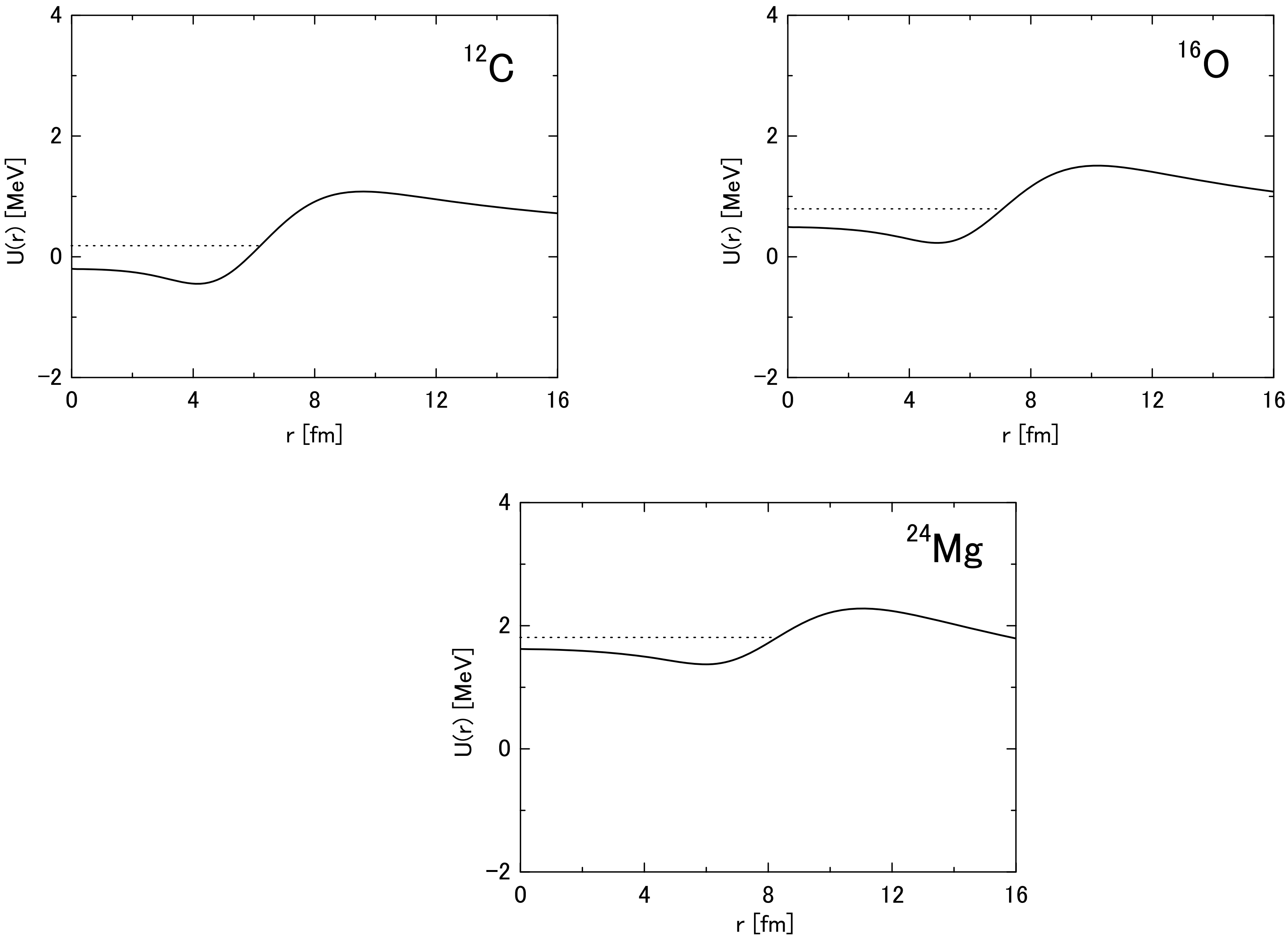
Hartree – Fock (Gross Pitaevsky eq)

for ideal bosons (α' s) :

$$\left[-\frac{\hbar}{2m_\alpha}\Delta + N \int d^3r' v(\vec{r} - \vec{r'}) |\phi_0(\vec{r'})|^2\right] \phi_0(\vec{r}) = \epsilon_0 \phi_0(\vec{r})$$

effective α – α + Coulomb

T. Yamada



Estimate for maximum number

$$\mathrm{N}^{\alpha}_{limit} \simeq 10 \qquad \Rightarrow \qquad {}^{40}\mathrm{Ca}^{**}$$

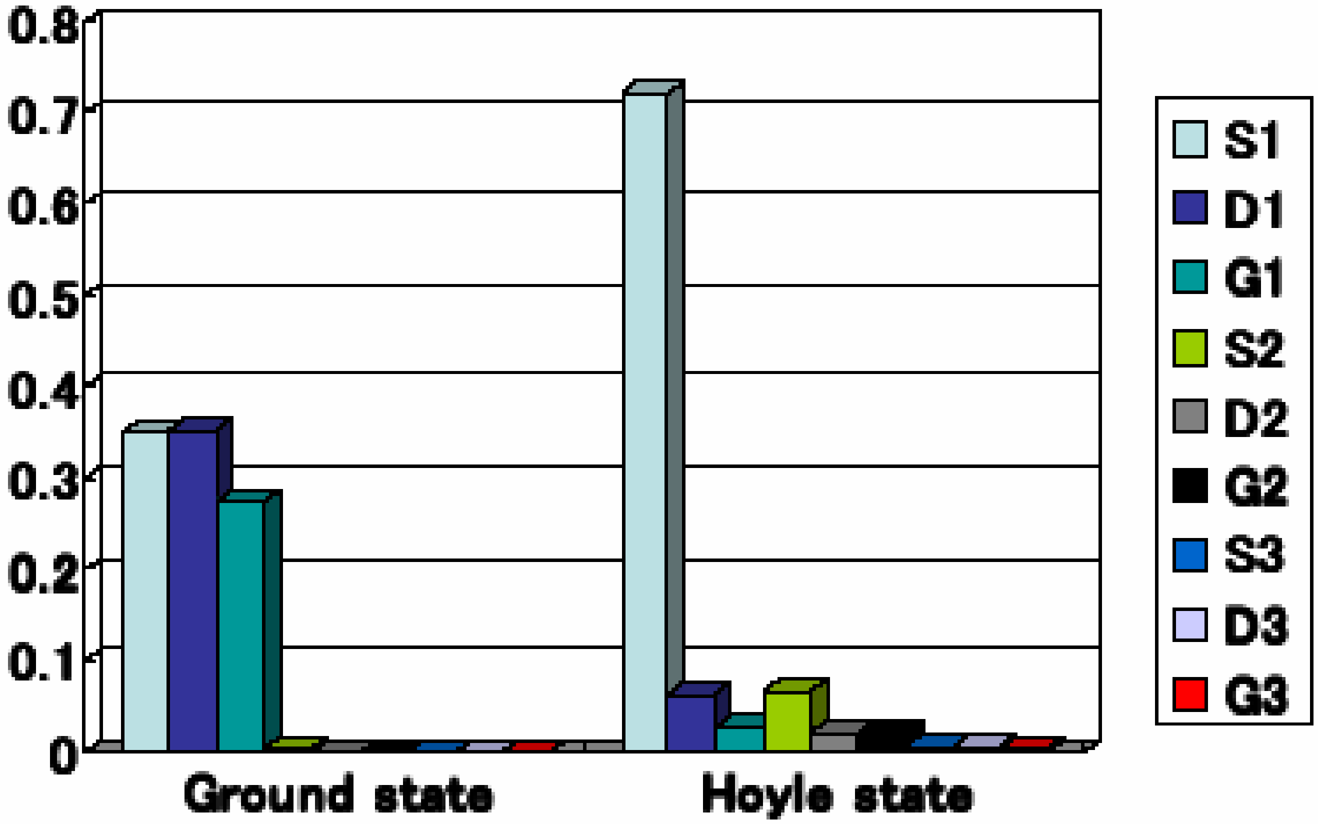
Boson occupancy :

α-particle density matrix :

$$\rho_\alpha(\vec{R}, \vec{R}'), \quad \vec{R} \text{ : c.m. of } \alpha$$

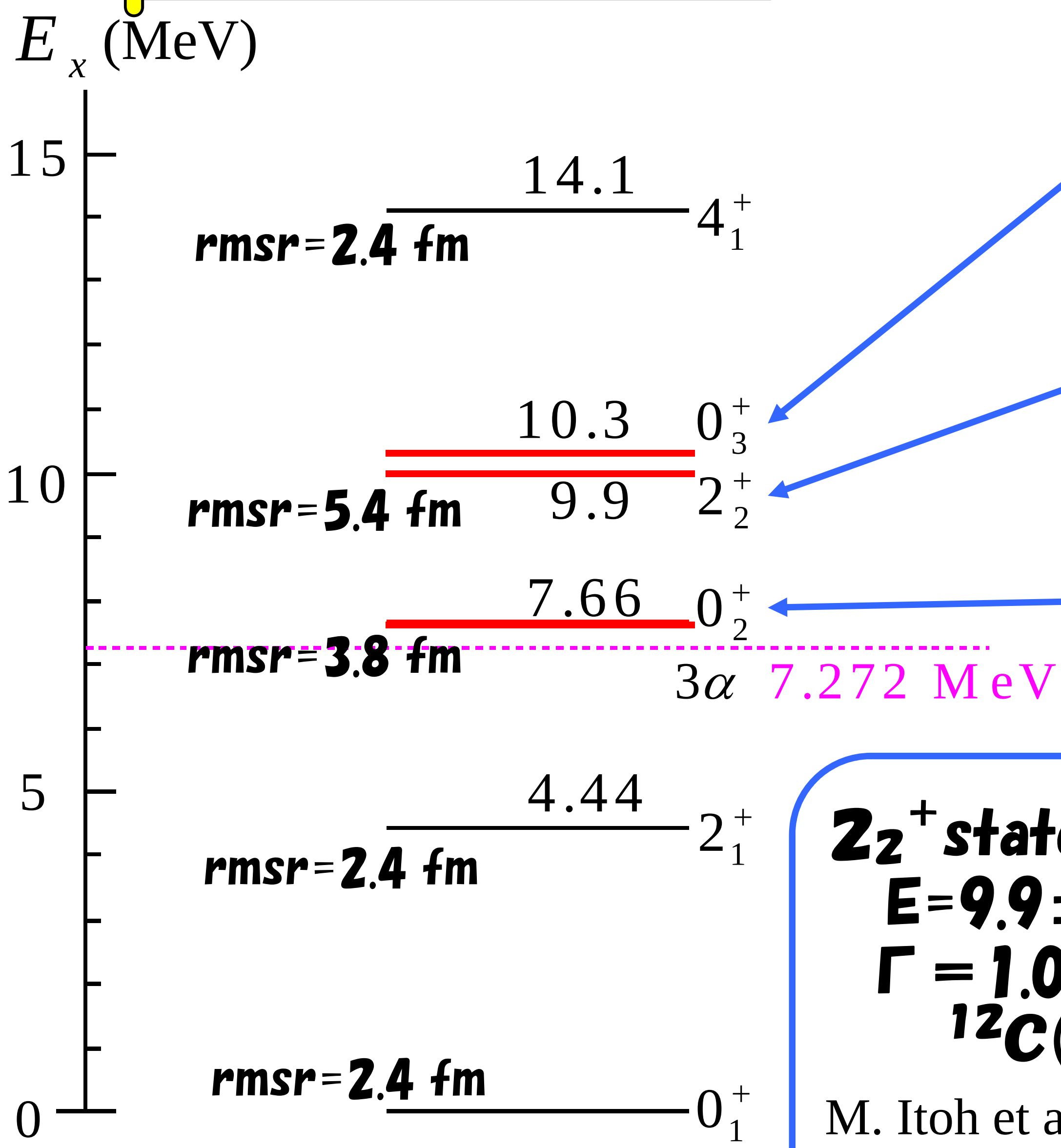
Diagonalization :

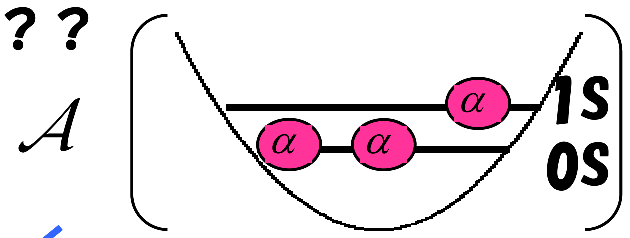
$$^{12}C : \quad O_2^+ \qquad 70\% \quad \text{S-wave occupancy}$$



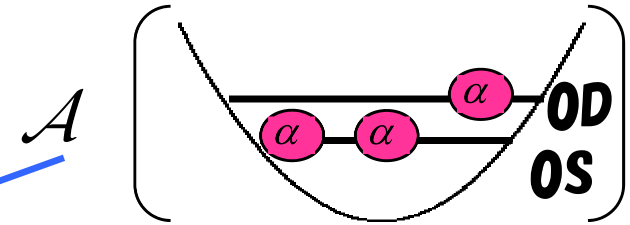
"BEC" in ^{12}C

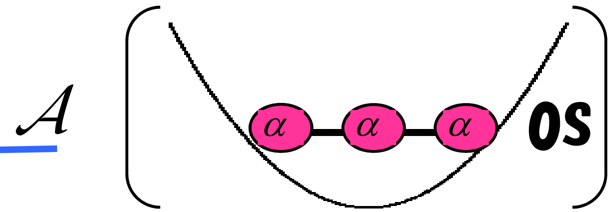
Observed levels of ^{12}C





C. Kurokawa and K. Katō, PRC 71, 021301 (2005).





2_2^+ state :
E=9.9±0.3 MeV
Γ=1.0±0.3 MeV
$^{12}C(\alpha, \alpha')$

M. Itoh et al., Nucl. Phys. A 738 (2004) 268-272

α cond. +ACCC
E=9.38 MeV
Γ=0.64 MeV
Volkov No. 1 force is adopted

Y. F. et al., EPJA 24, 321 (2005).

Hamiltonian of 4α OCM

$$H = T + \sum_{i<j}\left[V_{2\alpha}(r_{ij}) + V_{2\alpha}^{Coul}(r_{ij})\right] + V_{3\alpha} + V_{4\alpha} + V_{Pauli}$$

Equation of motion

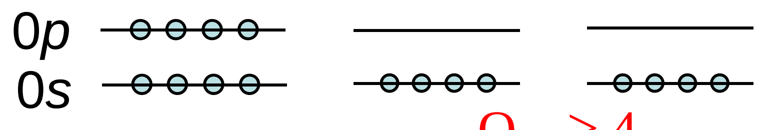
$$\delta\left[\left\langle \Phi_L(^{16}\mathrm{O}) \middle| H - E \middle| \Phi_L(^{16}\mathrm{O}) \right\rangle\right] = 0$$

Pauli blocking operator on α－α motions

$$V_{Pauli} = \lim_{\lambda\to\infty} \lambda \sum_{2n+\ell<4} \sum_{ij} \left| u_{n\ell}(r_{ij}) \right\rangle \left\langle u_{n\ell}(r_{ij}) \right|$$

Pauli forbidden state: h.o.w.f.

$$^{8}\mathrm{Be}(\mathrm{Q} \geq 4) \rightarrow \alpha(\mathrm{Q{=}0}) + \alpha(\mathrm{Q{=}0})$$



$\mathrm{Q_{rel}} \geq 4$

α－α motion with $Q_{rel} < 4$ should be eliminated

Q_{rel} : H.O. quantum number regarding α－α motion

2-body force (folding MHN force)

$$V_{2\alpha}(r) = \sum_n V_n^{(2)} \exp\left(-\beta_n^{(2)} r^2\right)$$

Coulomb force

$$V_{2\alpha}^{Coul}(r) = \frac{4e^2}{r} \operatorname{erf}(ar)$$

Phenomenological 3-body force (repulsive)

$$V_{3\alpha} = V^{(3)} \sum_{i<j<k} \exp\left[-\beta(r_{ij}^2 + r_{jk}^2 + r_{ki}^2)\right]$$

$$V^{(3)} = 87.5\ \mathrm{MeV}, \quad \beta = 0.15\ \mathrm{fm}^{-2}$$

Phenomenological 4-body force (repulsive)

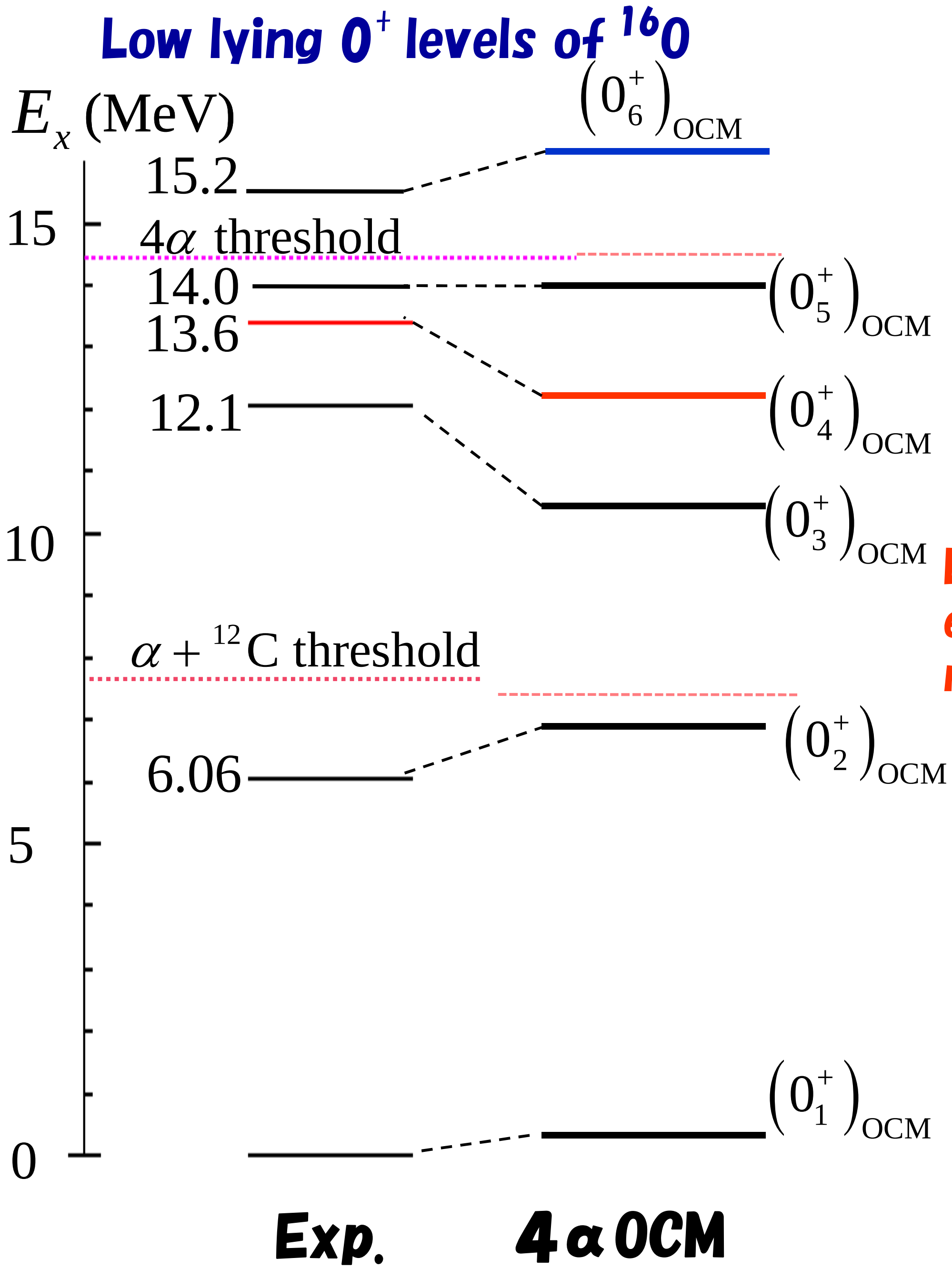
$$V_{4\alpha} = V^{(4)} \exp\left[-\beta(r_{12}^2 + r_{13}^2 + r_{14}^2 + r_{23}^2 + r_{24}^2 + r_{34}^2)\right]$$

$$V^{(4)} = 12000\ \mathrm{MeV}, \quad \beta = 0.15\ \mathrm{fm}^{-2}$$

Energies from 4α threshold

	Cal. (MeV)	Exp. (MeV)
^{12}C(g.s.)	−7.32	−7.28
^{12}C(2_1^+)	−4.88	−2.84
^{12}C(4_1^+)	2.06	6.43
^{12}C(0_2^+)	0.70	0.38
^{16}O(g.s.)	−14.2	−14.44

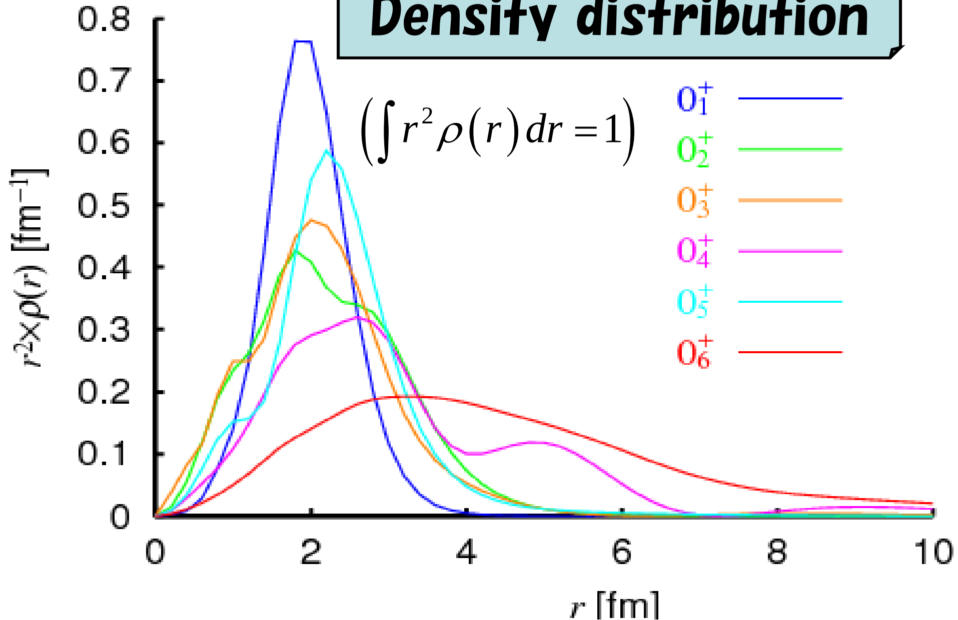
Energy levels, rms radii, monopole matrix elements and density distribution.



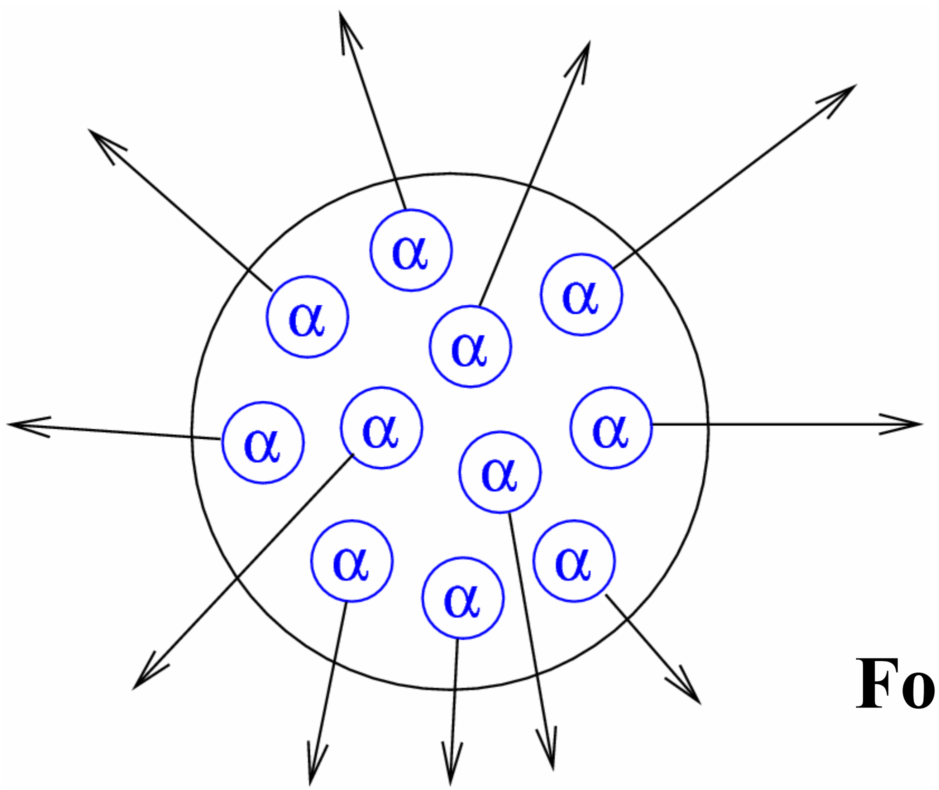
	R_{rms} (fm)	M(E0) (fm^2)	M(E0) (fm^2) Exp.
$(0_1^+)_{OCM}$	2.7		
$(0_2^+)_{OCM}$	3.1	4.2	0_2^+: 3.55
$(0_3^+)_{OCM}$	2.9	4.1	0_3^+: 4.03
$(0_4^+)_{OCM}$	3.9	2.4	0_4^+: no data
$(0_5^+)_{OCM}$	3.1	2.0	0_5^+: 3.3
$(0_6^+)_{OCM}$	5.4	1.4	0_6^+: no data

Large monopole matrix element can be the evidence of cluster states (Yamada, Y.F. et el., nucl-th/0703045)

Density distribution



Expanding Many α – Particle Coherent State:



All α's can be detected in coincidence

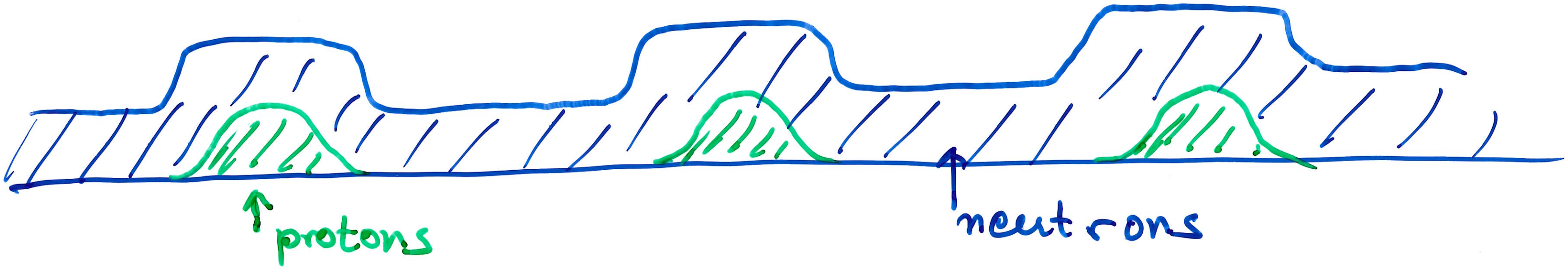
→Coherent State Established

For example, $^{40}Ca^{*} \rightarrow 10\ \alpha$'s

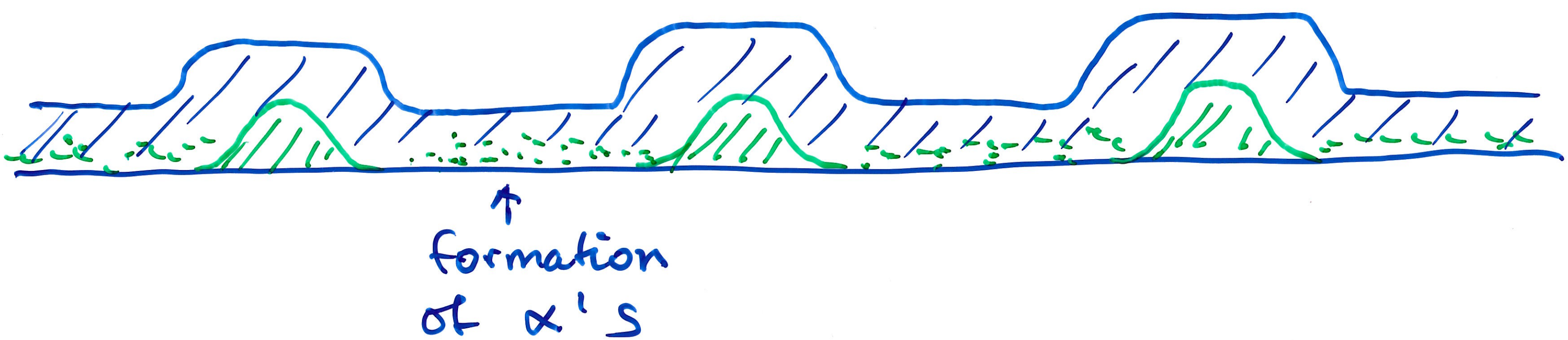
Search at Orsay

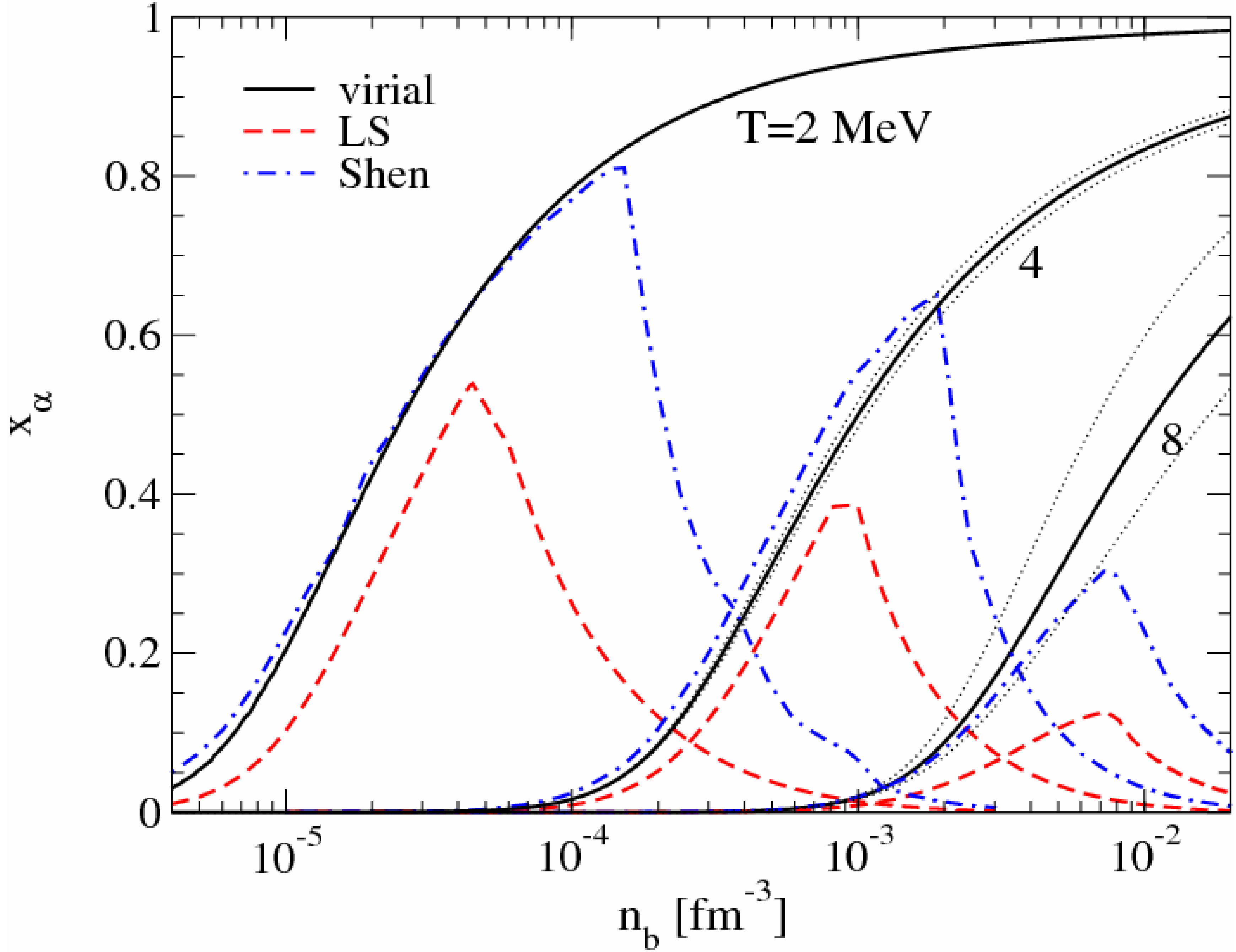
Alpha particle condensation in crust of neutron stars?

T=0

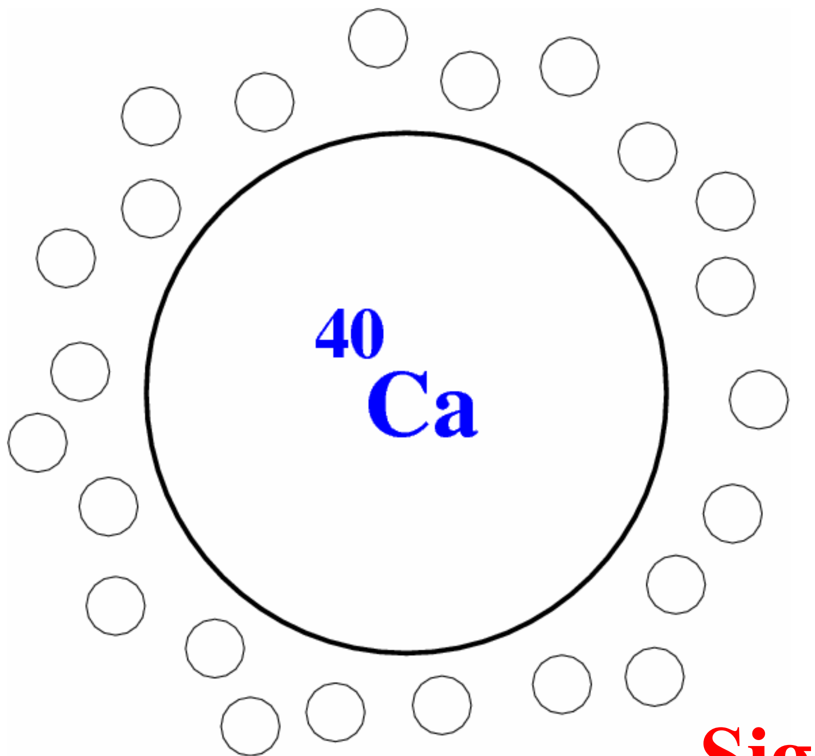


T>0




virial
LS
Shen
T=2 MeV
4
8
x_α
n_b [fm^{-3}]
0
0.2
0.4
0.6
0.8
1
10^{-5}
10^{-4}
10^{-3}
10^{-2}

Compound Nuclei with α – Gas:

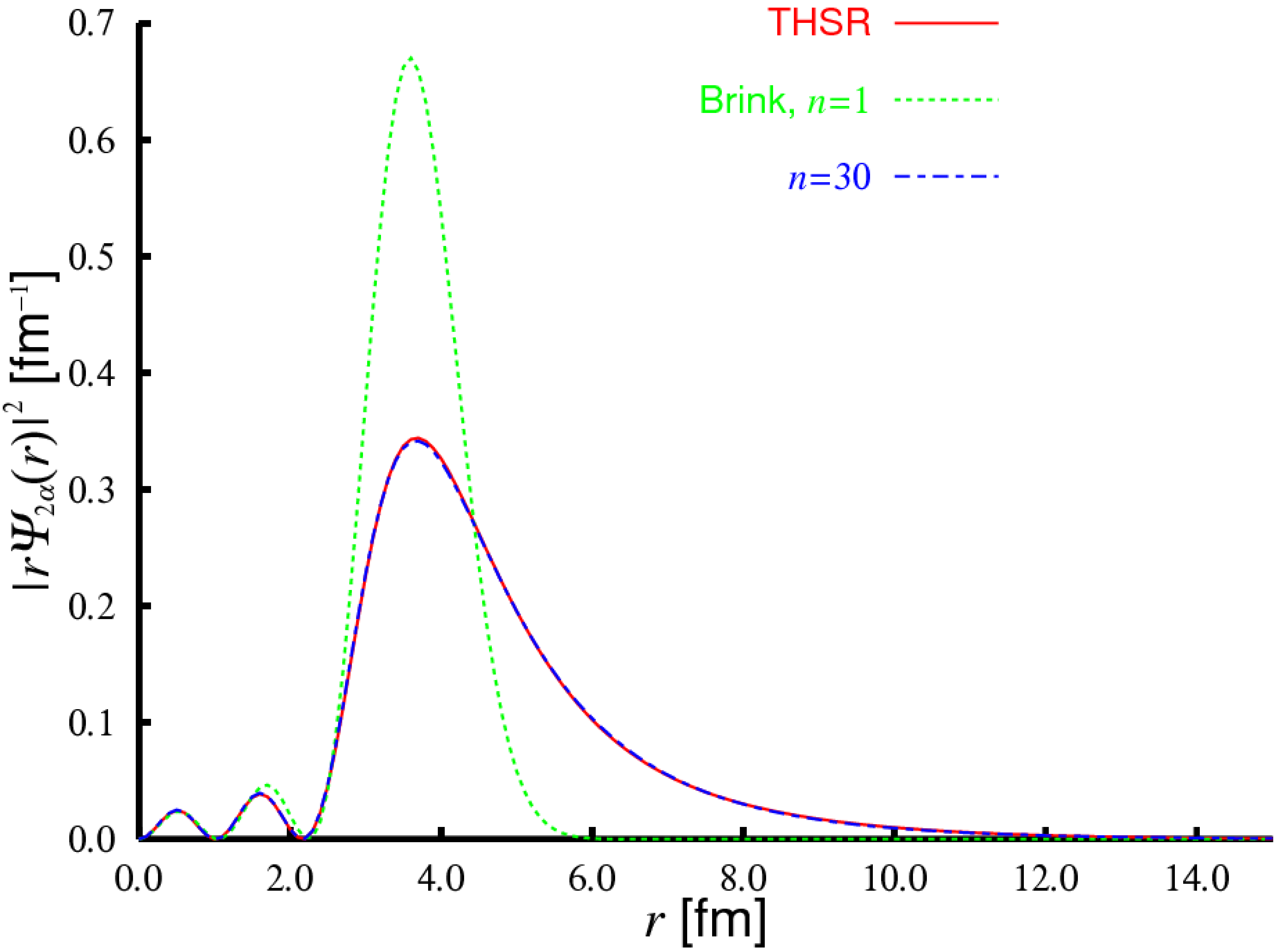


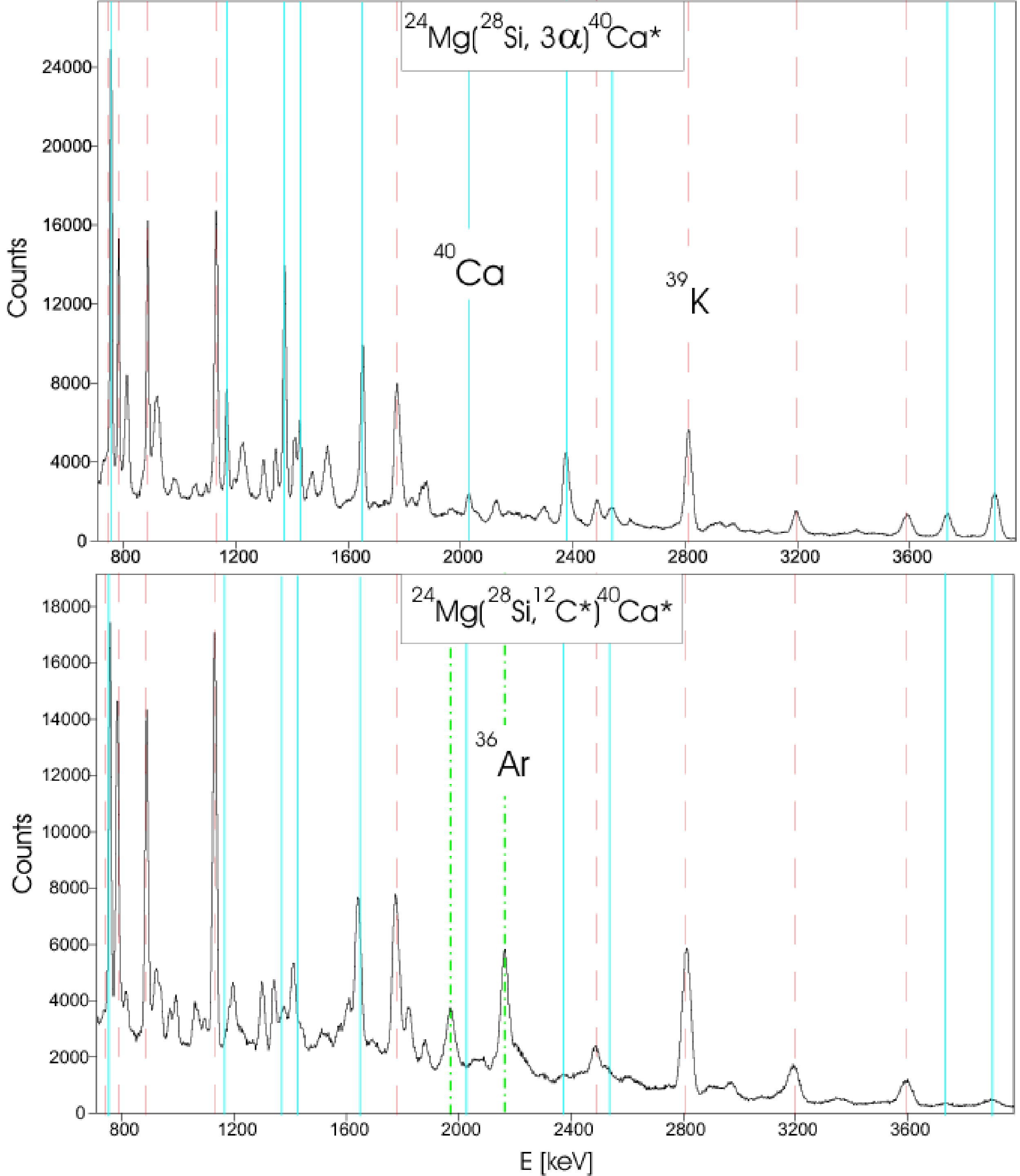
Disintegration into 2 α's

Enhanced ➔

Sign of α – Particle Condensate

(W. V. O.)


THSR
Brink, n=1
n=30
$|r\Psi_{2\alpha}(r)|^2$ [fm^{-1}]
r [fm]
0.0
0.1
0.2
0.3
0.4
0.5
0.6
0.7
0.0
2.0
4.0
6.0
8.0
10.0
12.0
14.0



$^{24}Mg(^{28}Si, 3\alpha)^{40}Ca^*$
^{40}Ca
^{39}K
Counts
$^{24}Mg(^{28}Si, ^{12}C^*)^{40}Ca^*$
^{36}Ar
Counts
E [keV]

Critical Temperature for α-particle Condensation in symmetric and asymmetric nuclear matter

T. Sogo, P. Schuck, G. Röpke

Well known: T_c for pairing

Gap eq. $$\Delta_p = \sum_{p'} V_{pp'} \frac{\Delta_{p'}}{2E_{p'}} (1 - 2f(E_{p'}))$$

$$\boxed{E_p = \sqrt{(\varepsilon_p - \mu)^2 + \Delta_p^2}} \qquad f(\omega) = \frac{1}{e^{\omega/T} + 1}$$

Critical temperature $\Delta_p \to 0$

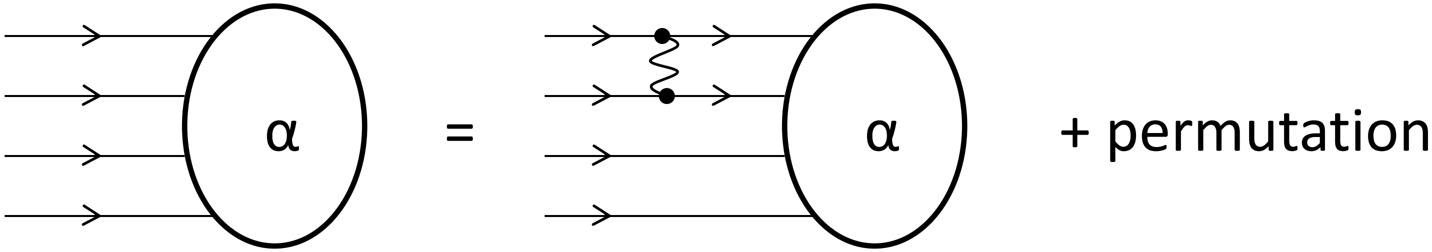
$$2\varepsilon_p \kappa_p - (1 - 2f(\varepsilon_p - \mu)) \sum_{p'} V_{pp'} \kappa_{p'} = 2\mu\kappa_p$$

<u>Thouless criterion!</u> two body equation

Generalization to four body (α-particle) condensation

<u>eq. for T_c</u>

$$\boxed{\begin{aligned}&(4\mu-\varepsilon_1-\varepsilon_2-\varepsilon_3-\varepsilon_4)\psi_{1234}\\&=(1-f_1-f_2)v_{121'2'}\psi_{1'2'34}+(1-f_1-f_3)v_{131'3'}\psi_{1'23'4}\\&+(1-f_1-f_4)v_{141'4'}\psi_{1'234'}+(1-f_2-f_3)v_{232'3'}\psi_{12'3'4}\\&+(1-f_2-f_4)v_{242'4'}\psi_{12'34'}+(1-f_3-f_4)v_{343'4'}\psi_{123'4'}\end{aligned}}$$

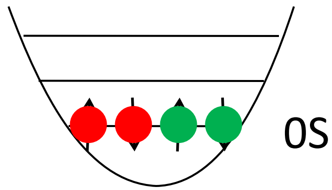


Solution with realistic NN force: R. Lazauskas, PRC**79**,1051801

Heavy numerical calculation! Approximation demanded!

Trial: momentum popjected mean field

$$\psi_{1234} = \delta(\vec{K} - \vec{k}_1 - \vec{k}_2 - \vec{k}_3 - \vec{k}_4)\varphi(k_1)\varphi(k_2)\varphi(k_3)\varphi(k_4)$$



For condensation $\underline{\vec{K} = 0}$!

Only one unknown single particle wave function $\varphi(k)$

+ Effective interaction:

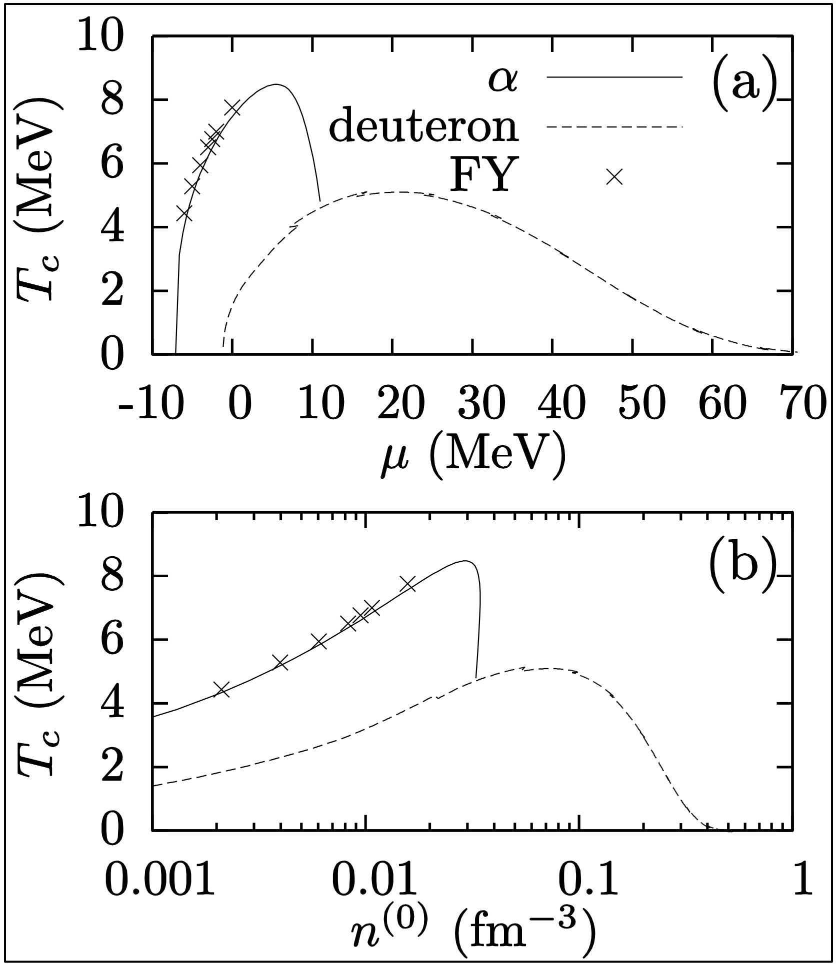
$$v_{1234} = \lambda w(\vec{k}_1 - \vec{k}_2) w(\vec{k}_3 - \vec{k}_4)\delta(\vec{k}_1 + \vec{k}_2 - \vec{k}_3 - \vec{k}_4) \qquad w(\vec{k}) = e^{-k^2/(4k_0^2)}$$

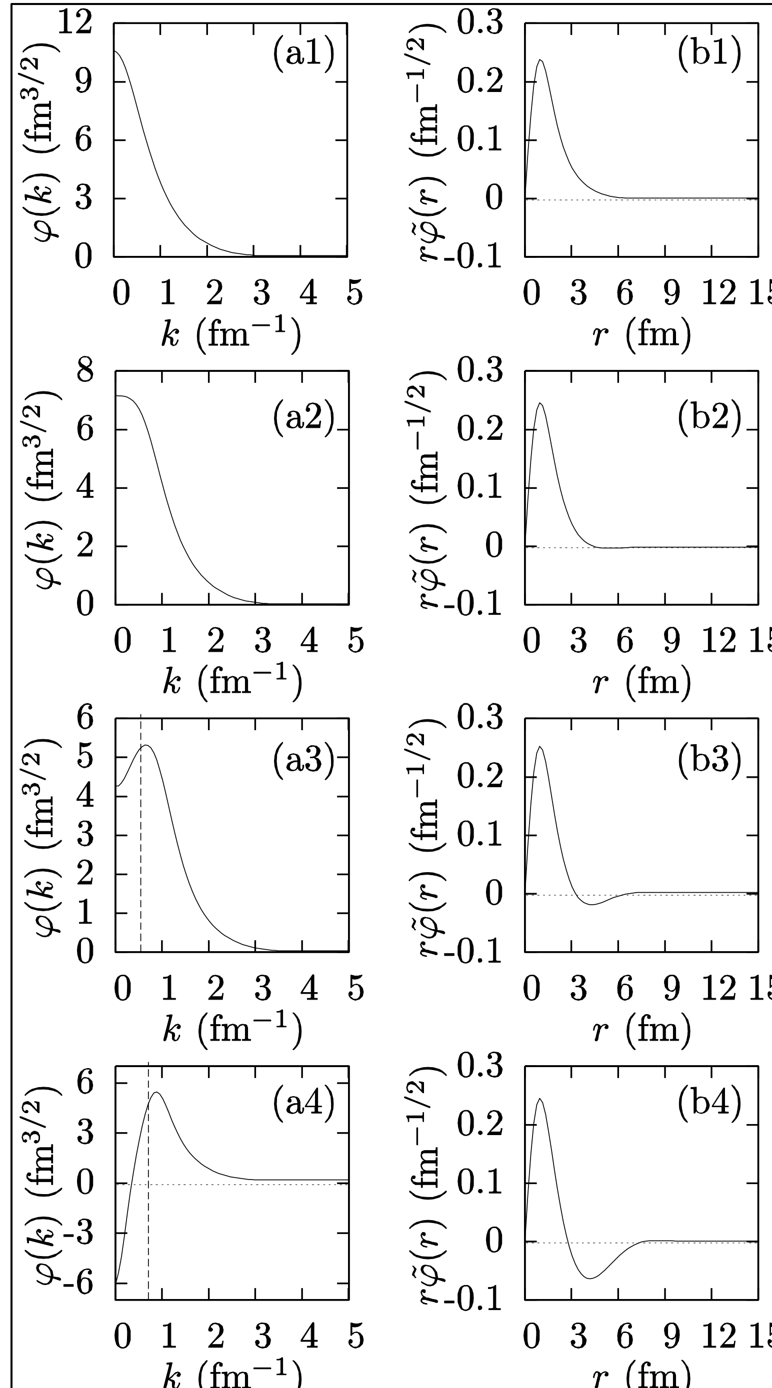
$\lambda,\ k_0$: such that binding and radius of α-particle correct!

HF type of eq. for $\varphi(k)$

$$\boxed{\varphi(k) = -\frac{3B[\varphi, T, \mu]}{A[\varphi, T, \mu] + 3C[\varphi, T, \mu]}}$$

Results: Symmetric nuclear matter



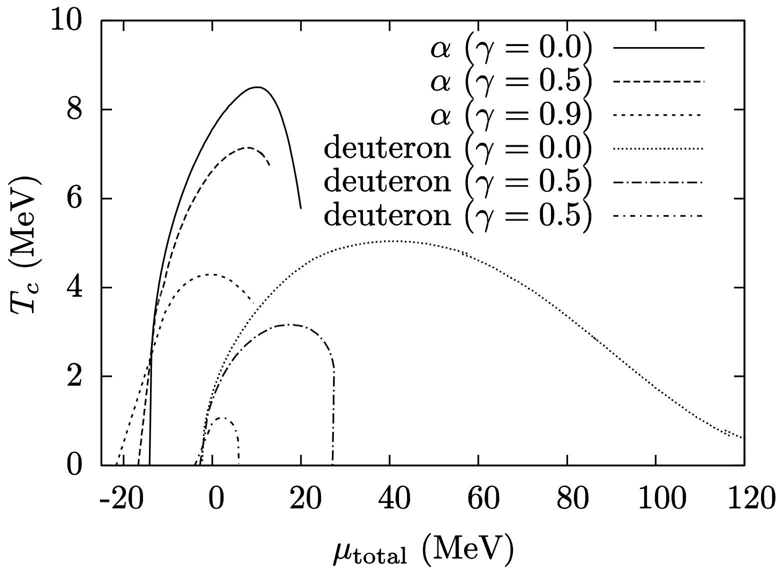


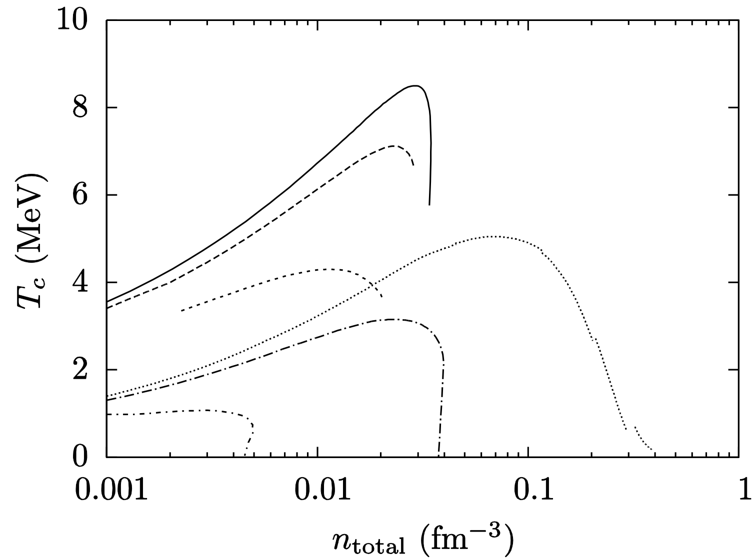
[T. Sogo, et. al., PRC **79**, 051301 (2009)]

We generalize to asymmetric matter, like in compact stars.

Same formulation with projected mean field

Now $\varphi_{\mathrm{p}} \neq \varphi_{\mathrm{n}} \rightarrow$ two coupled eqs.

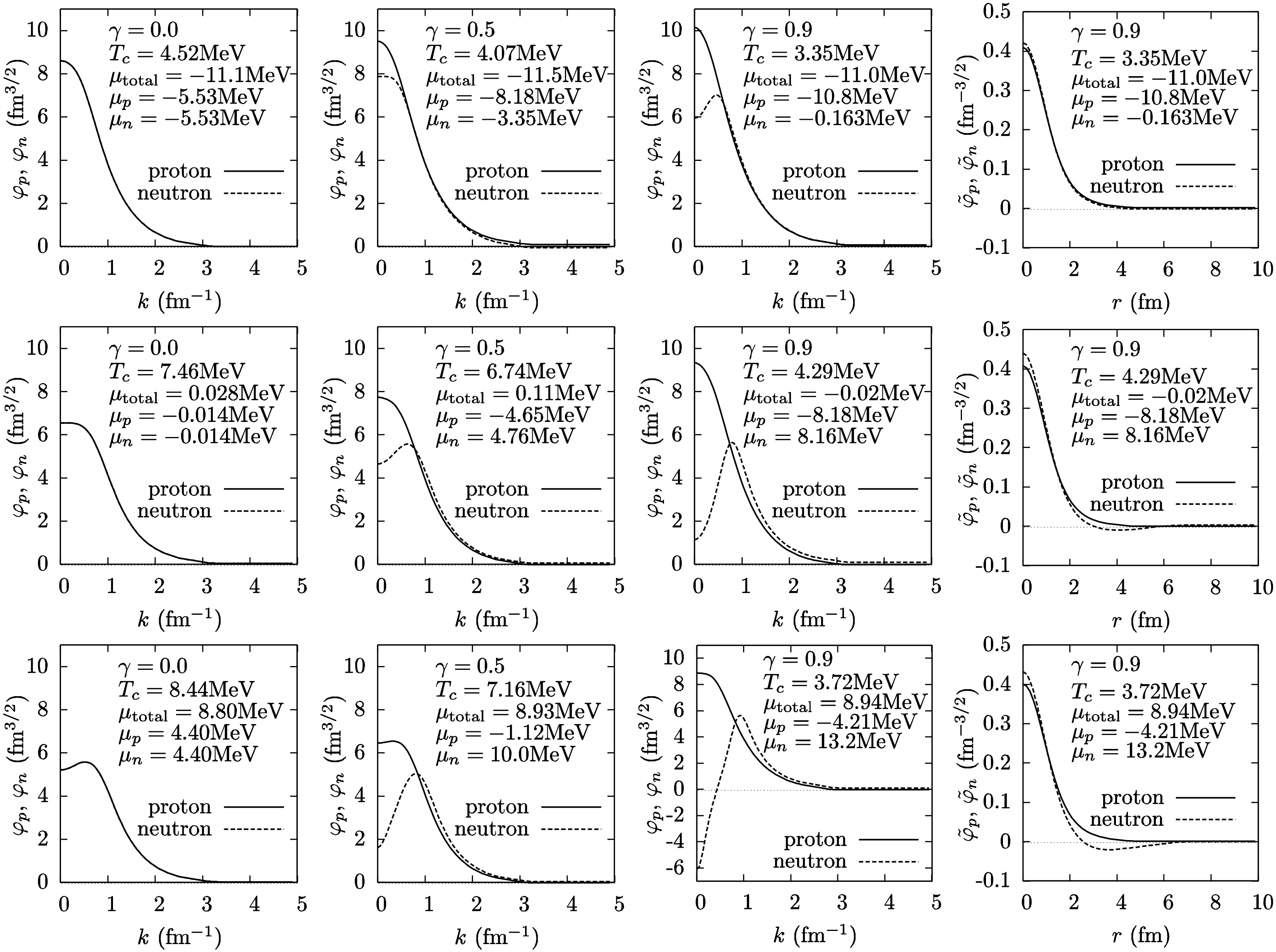




$$\gamma = \frac{n_{\mathrm{n}} - n_{\mathrm{p}}}{n_{\mathrm{n}} + n_{\mathrm{p}}}$$

Critical temperature quite high
even for very strong asymmetry (γ=0.9)!

Wave functions (Asymmetric nuclear matter)



In BCS we have the following two coupled eqs.:

$$hU + \Delta V = EU$$

$$\Delta U - hV = EV$$

In first eq we can eliminate V component using second eq.:

$$V = \frac{1}{E+h}\Delta U$$

insert into first eq.:

$$(E-h)U + \frac{\Delta\Delta}{E+h}U = 0$$

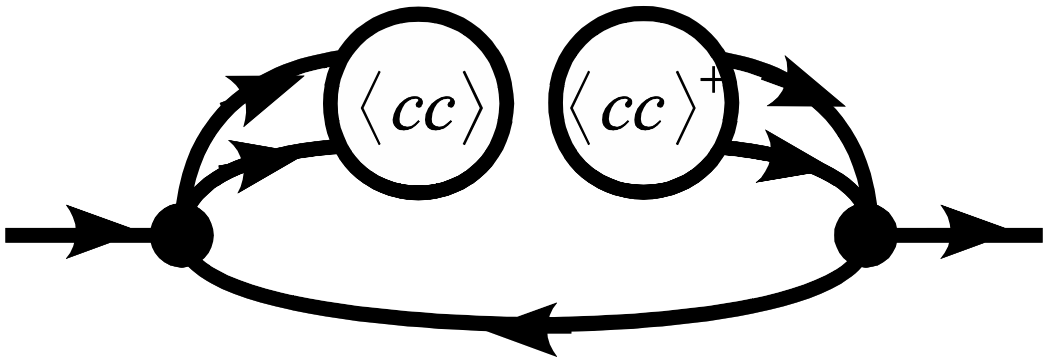
The effective single particle potential

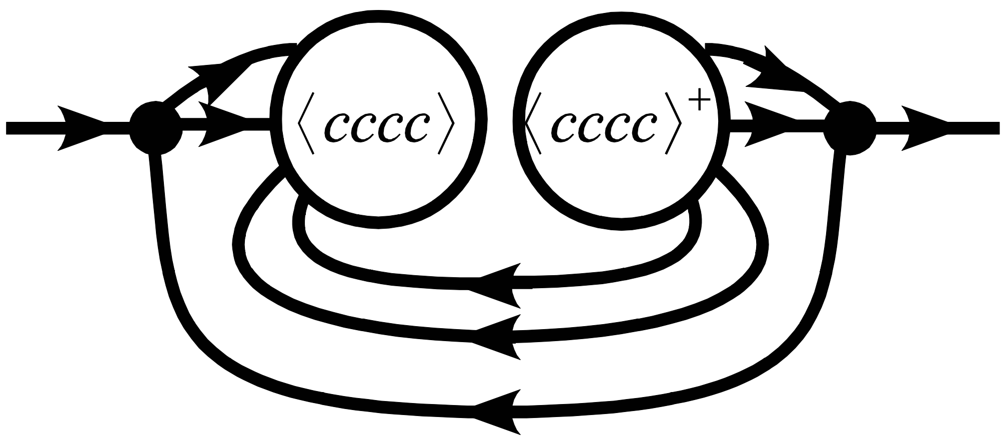
$$M^{BCS} = \frac{|\Delta|^2}{E+h}$$

is called the **single particle mass operator in BCS approximation**

Full nonlinear quartet order parameter equation + solution

Pairing and quartet mass operators:

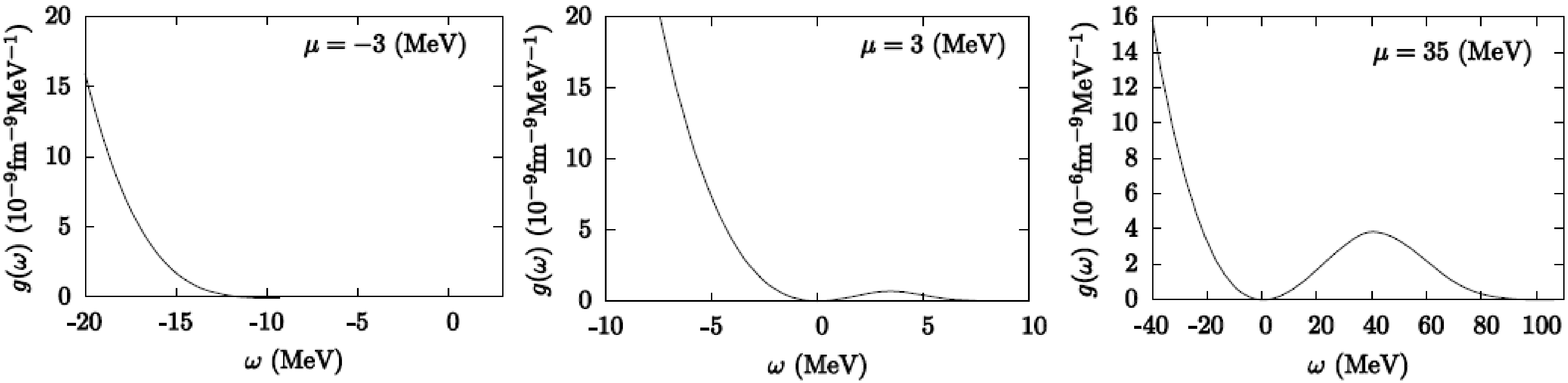




Three hole level density

$$\bar{f} = 1 - f$$

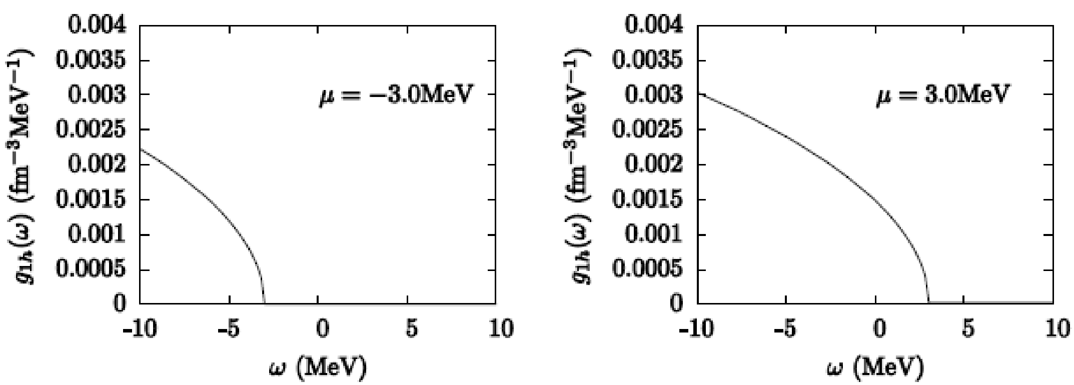
$$g(\omega) = \int \frac{d^3k_1}{(2\pi)^3} \frac{d^3k_2}{(2\pi)^3} \frac{d^3k_3}{(2\pi)^3} \left(\bar{f}_1 \bar{f}_2 \bar{f}_3 + f_1 f_2 f_3 \right) \delta(\omega + \varepsilon_1 + \varepsilon_2 + \varepsilon_3)$$



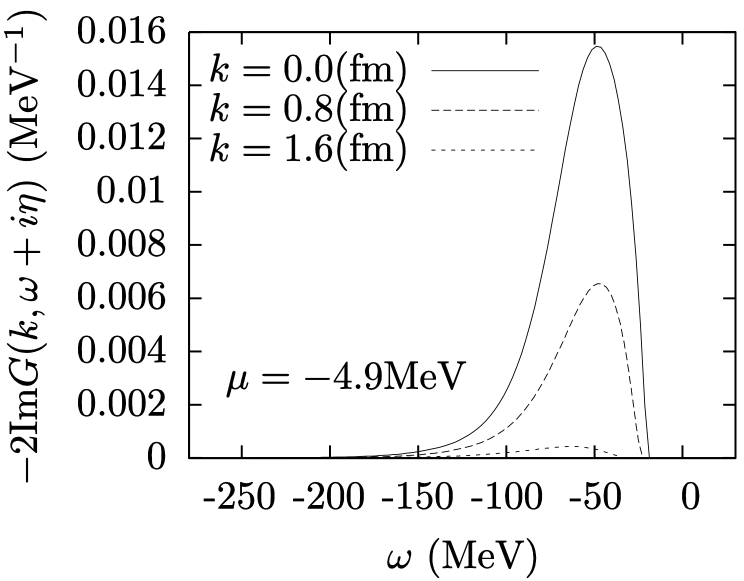
c.f. one hole level density

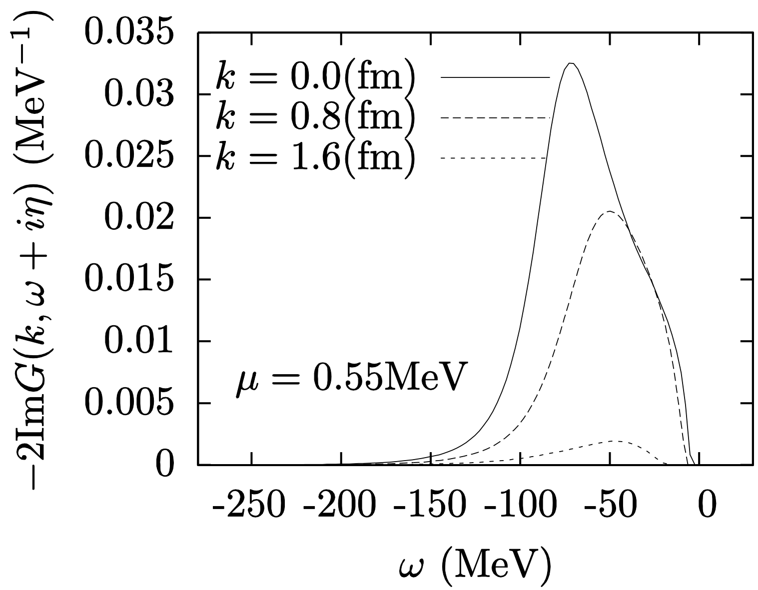
$$g_{1h}(\omega) = \int \frac{d^3k}{(2\pi)^3} \pi \left[f_k + \bar{f}_k \right] \delta(\omega + \varepsilon)$$

$$= \begin{cases} \dfrac{m}{2\pi} \sqrt{2m(\mu - \omega)} & (\mu > \omega) \\ 0 & (\mu < \omega) \end{cases}$$



Spectral function of single-particle GF





No sharp quasiparticle pole!

[T. Sogo, et. al., published in PRC (arXiv:1004.0359)]

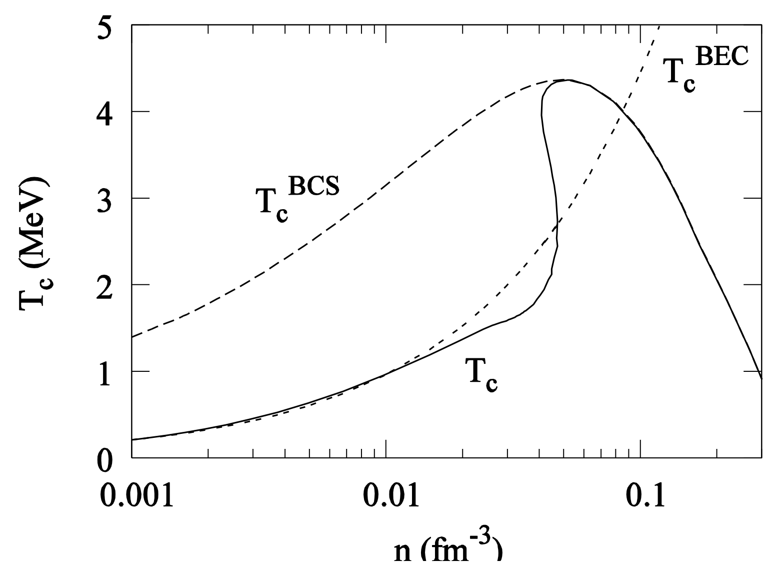
Conclusion

Projected product ansatz allows to tackle with the difficult quartet condensation process.

Still missing :

Theoretical description of the transition weakly bound quartets (present) to strongly bound ones (BEC)

For pairing, this has been achieved by Nozières and Schmitt-Rink. Transition between weakly bound pairs and strongly bound molecules seems to be smooth. However, BEC transition temperature much lower than BCS one.



What will it be for α-particle? Future task to be solved!

Conclusions, Outlook:

- **Loosely Bound 3 – α State in $^{12}C^{*}$ Finally Established**
- **More α – Particle Condensates Very Likely to Exist**
- **Nuclei Unique Fermi Systems for Cluster Effects. Eventually also in Future Cold Atom Experiments**
- **May be α – Condensation Important in Proto – Neutron stars**

Collaborators:

Y. Funaki, H. Horiuchi, G. Röpke, A. Tohsaki,

T. Yamada, W. von Oertzen

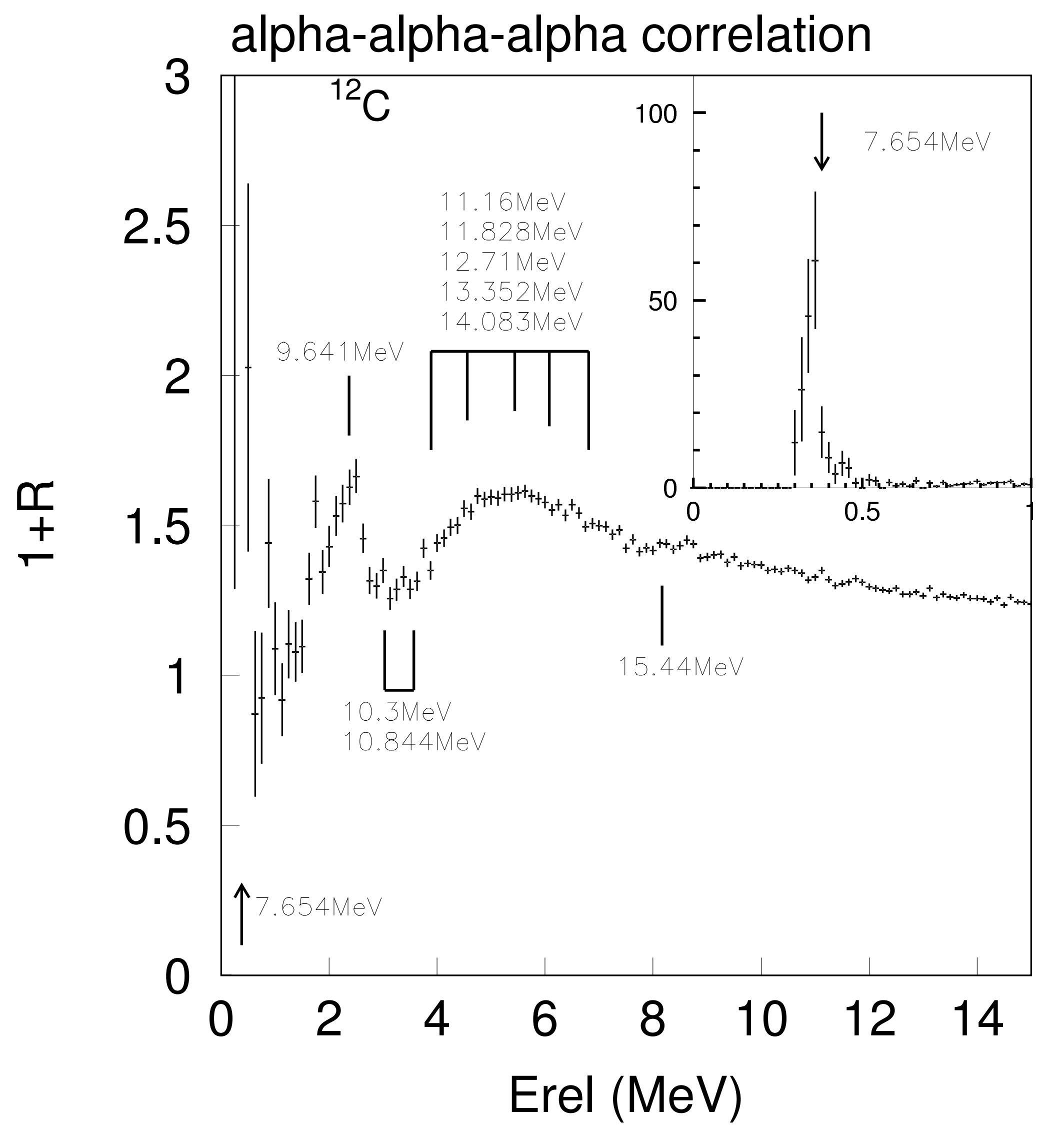
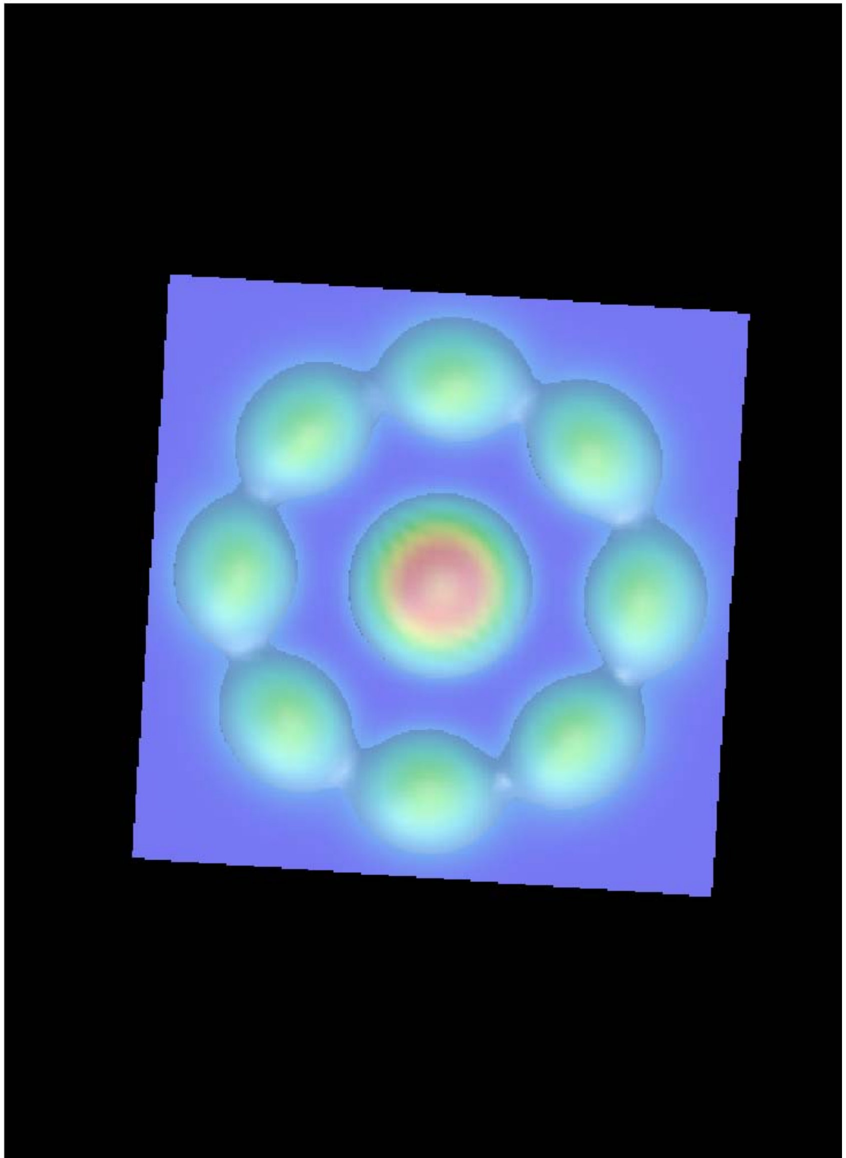
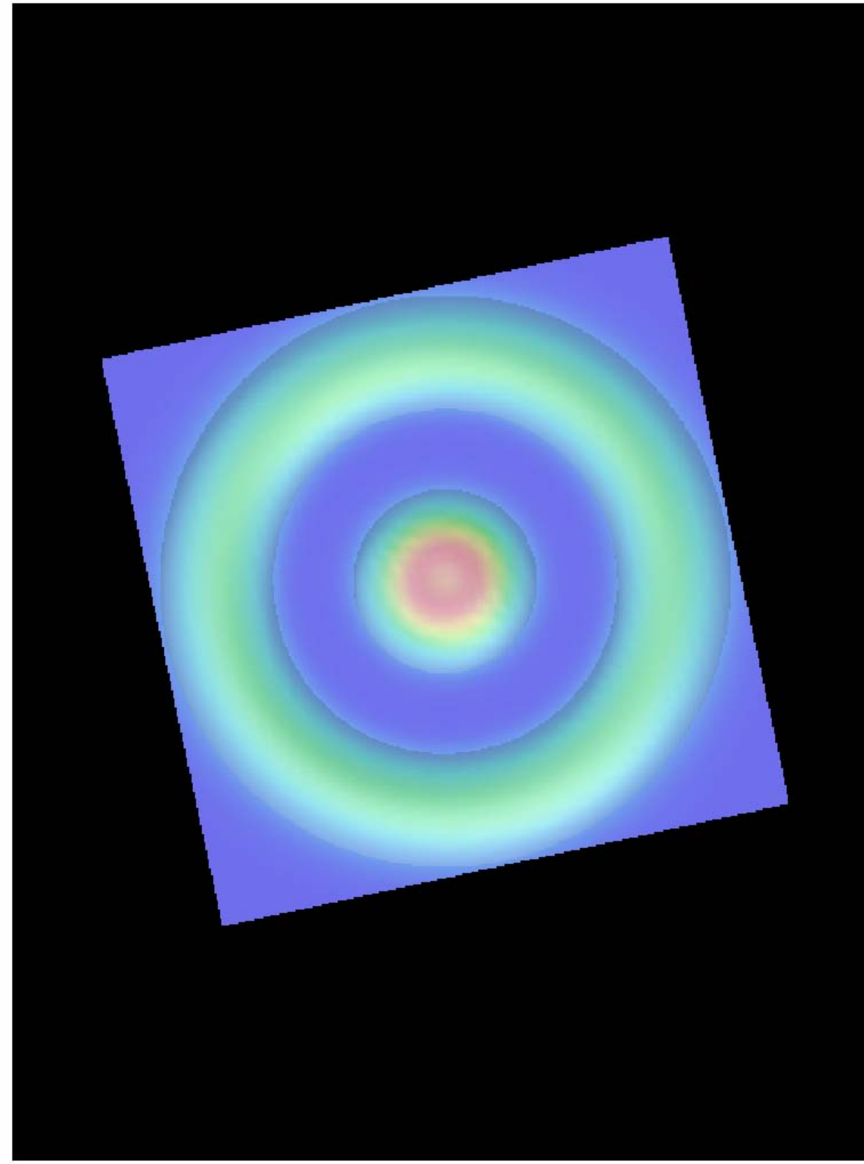
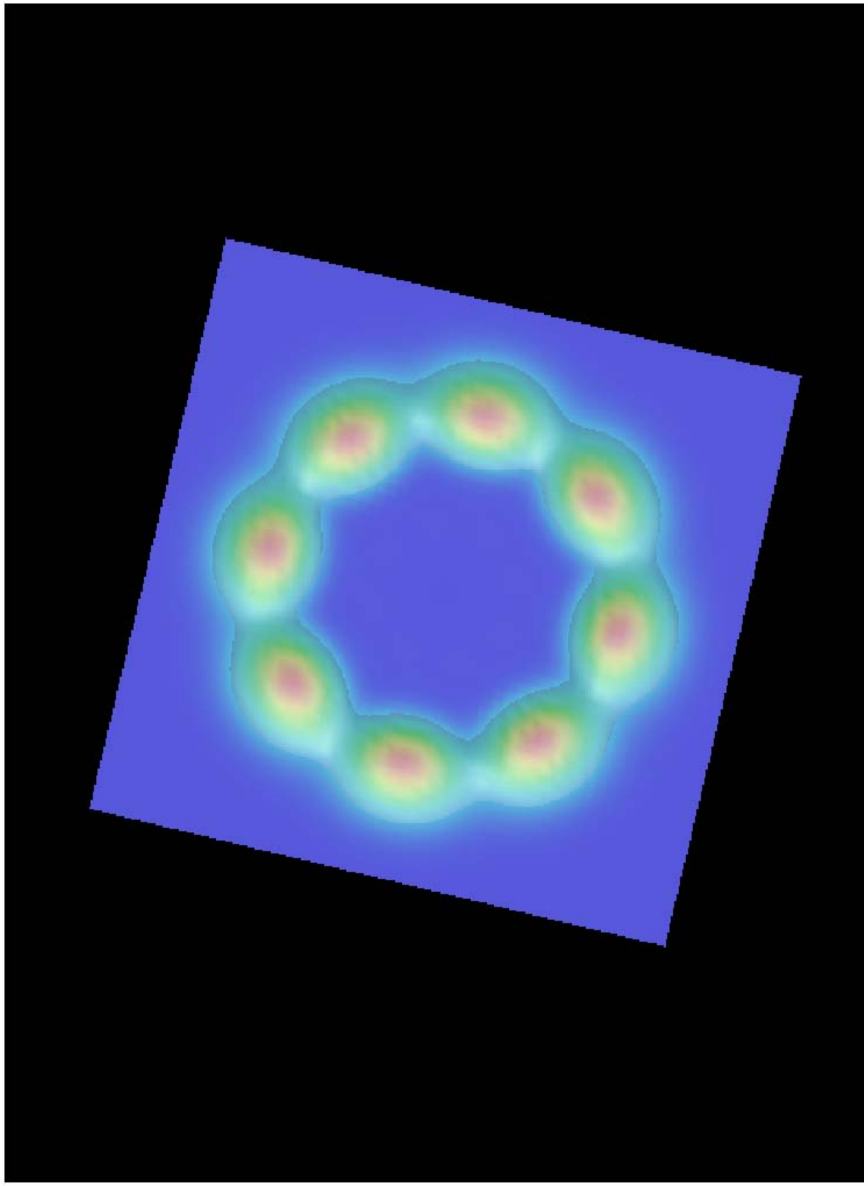
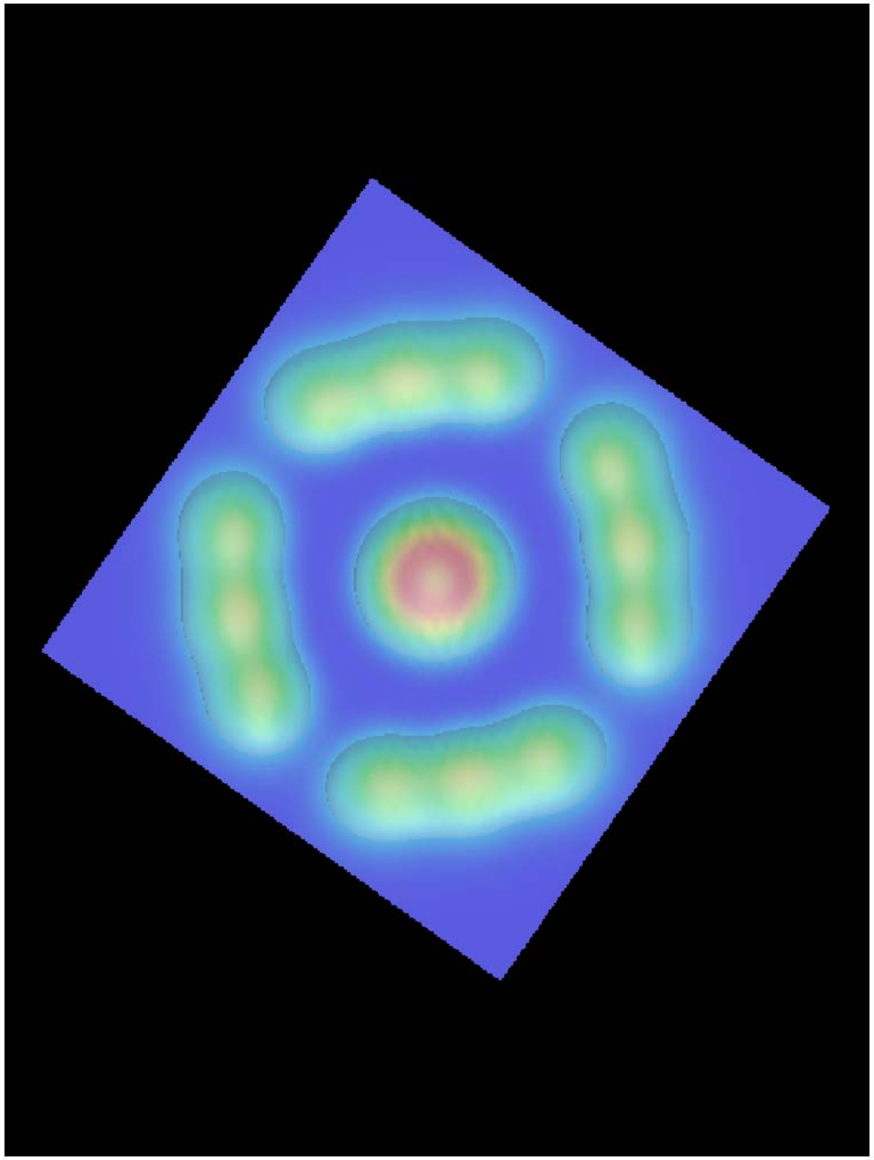


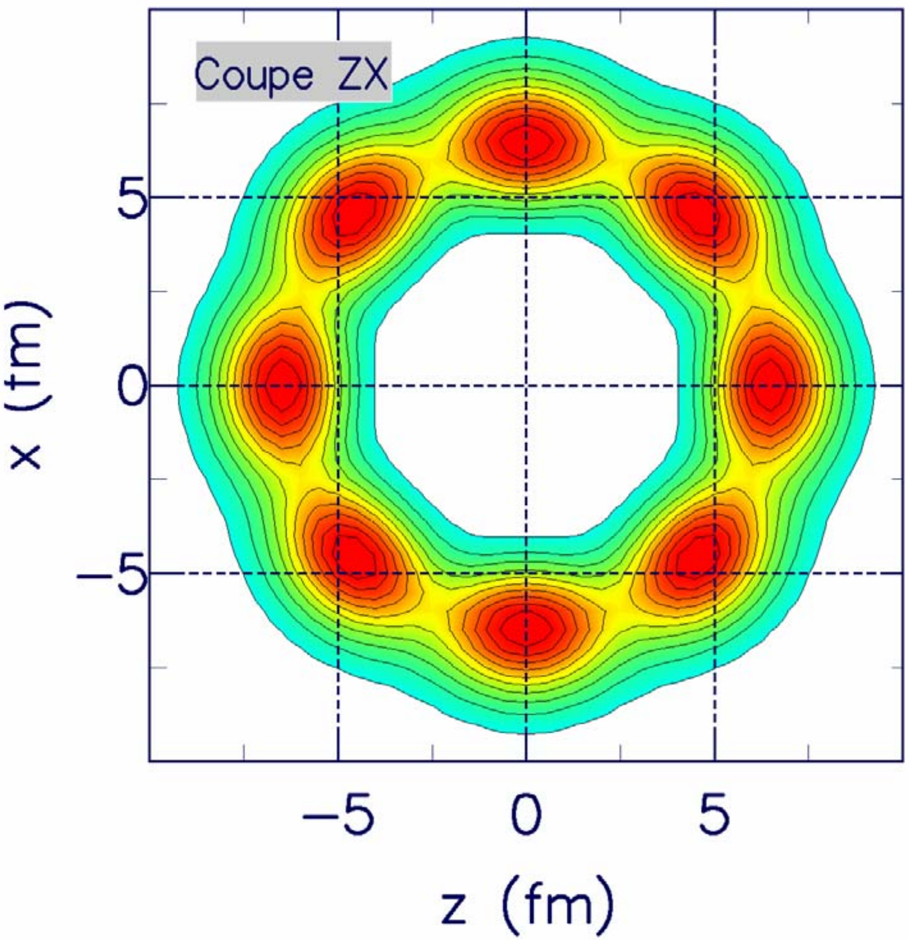
Figure 5.27: The α-α-α correlation function is shown. Resonances from the excited states of ^{12}C are labelled with the first peak seen more clearly in the inner upright panel.



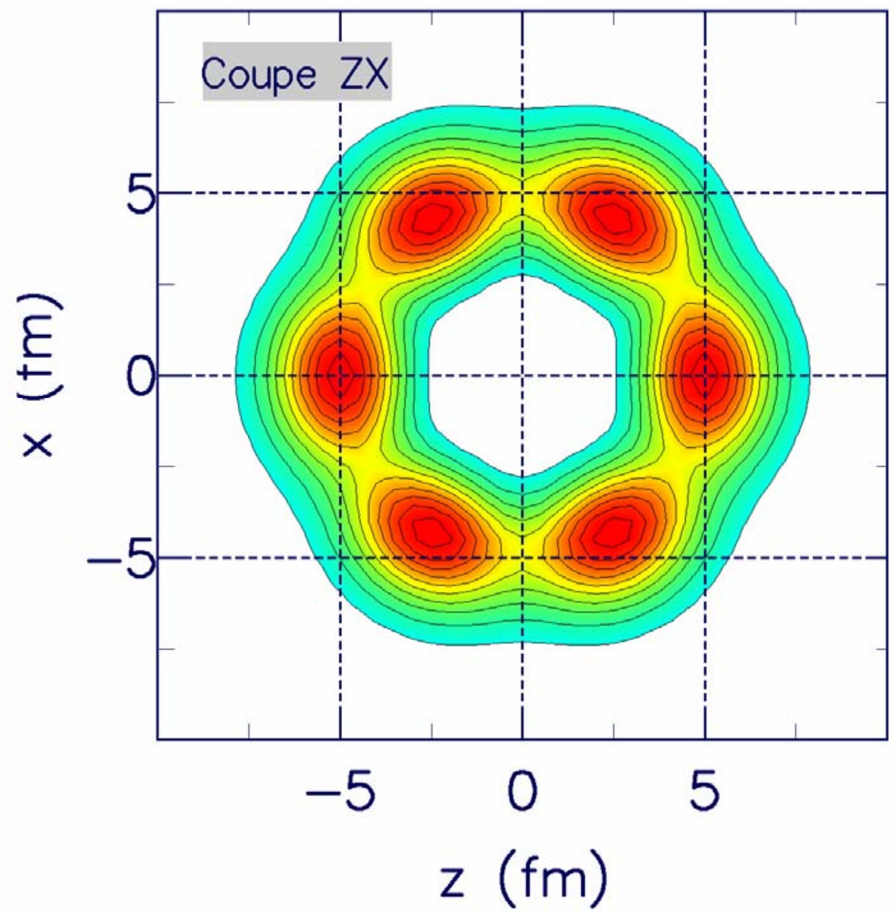


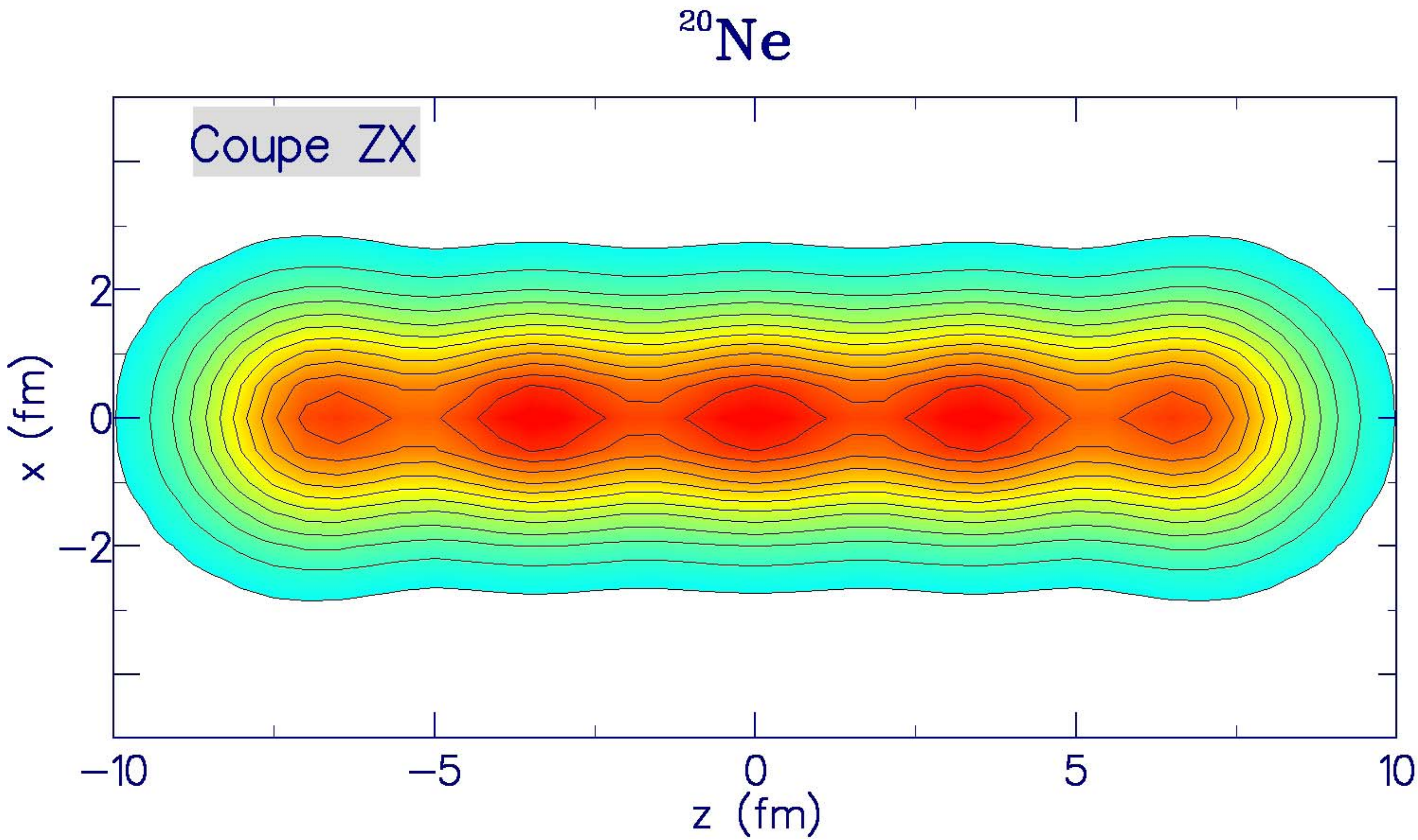



^{32}S



^{24}Mg





^{20}Ne
Coupe ZX
x (fm)
z (fm)
2
0
−2
−10
−5
0
5
10