

Dynamics and Correlations in Exotic Nuclei (DCEN2011)
(Sept.20—Oct.28, 2011, @ Yukawa Institute for Theoretical Physics, Kyoto.)
(seminar talk on 19th October)

Microscopic interaction models for nuclear reactions with exotic beams:

- present status and future perspective -

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W. Horiuchi (RIKEN)

● ***Understanding the **interactions** between composite nuclei (AA interactions), starting from NN interaction :***

- ✓ ***one of the fundamental subject in nuclear physics***
- ✓ ***one of the key issue to understand various **nuclear reactions**:***
 - ***optical potentials: elastic scattering***
 - ***distorting potentials as doorway to various reactions (inelastic, transfer, knockout, breakup ···)***
- ✓ ***important to survey unknown nuclear structures/reaction of **unstable nuclei** far from stability lines ($N \gg Z$, $Z \gg N$), for which***
 - ***few/no elastic-scattering data & phenom. potential information is available.***

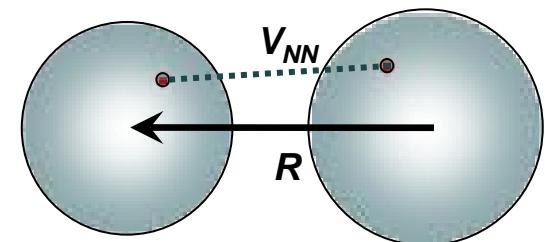
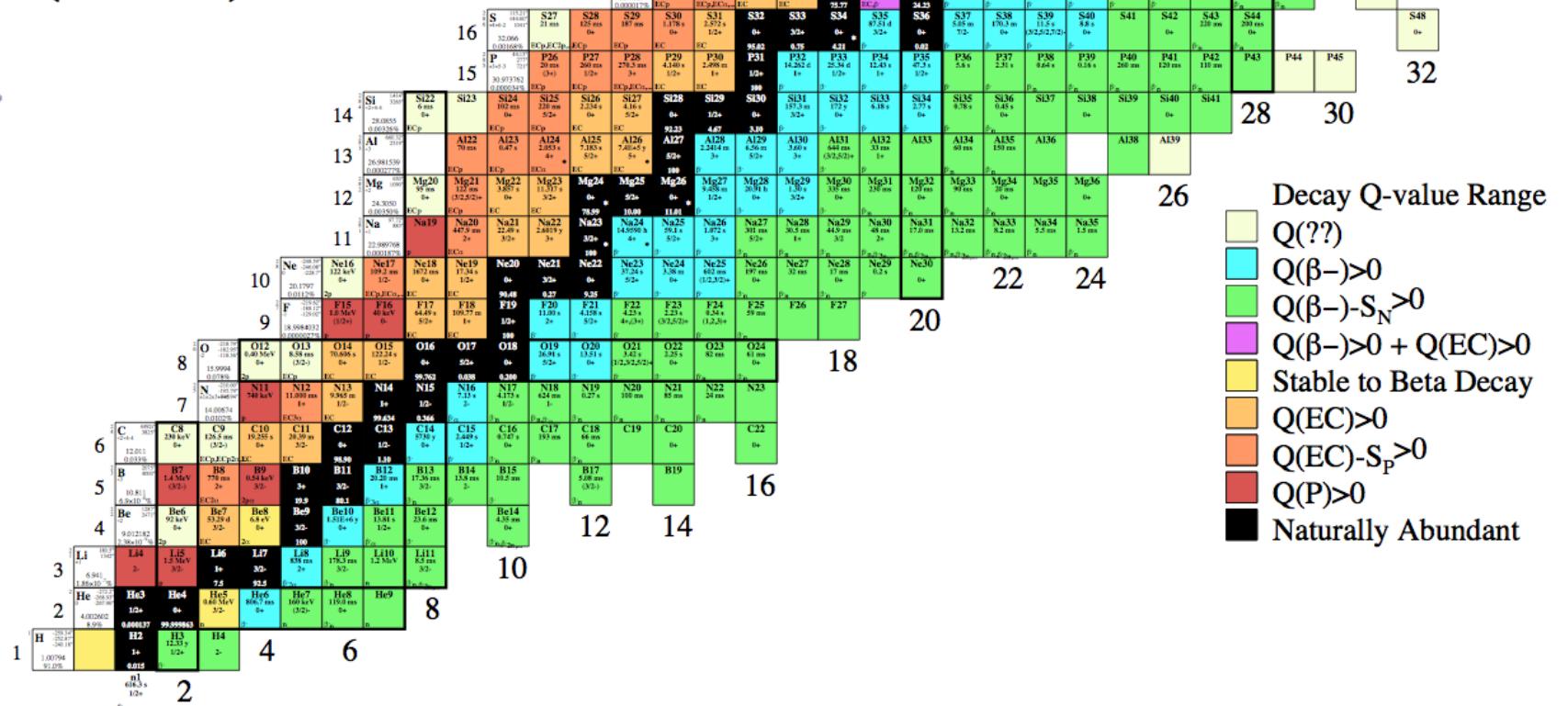


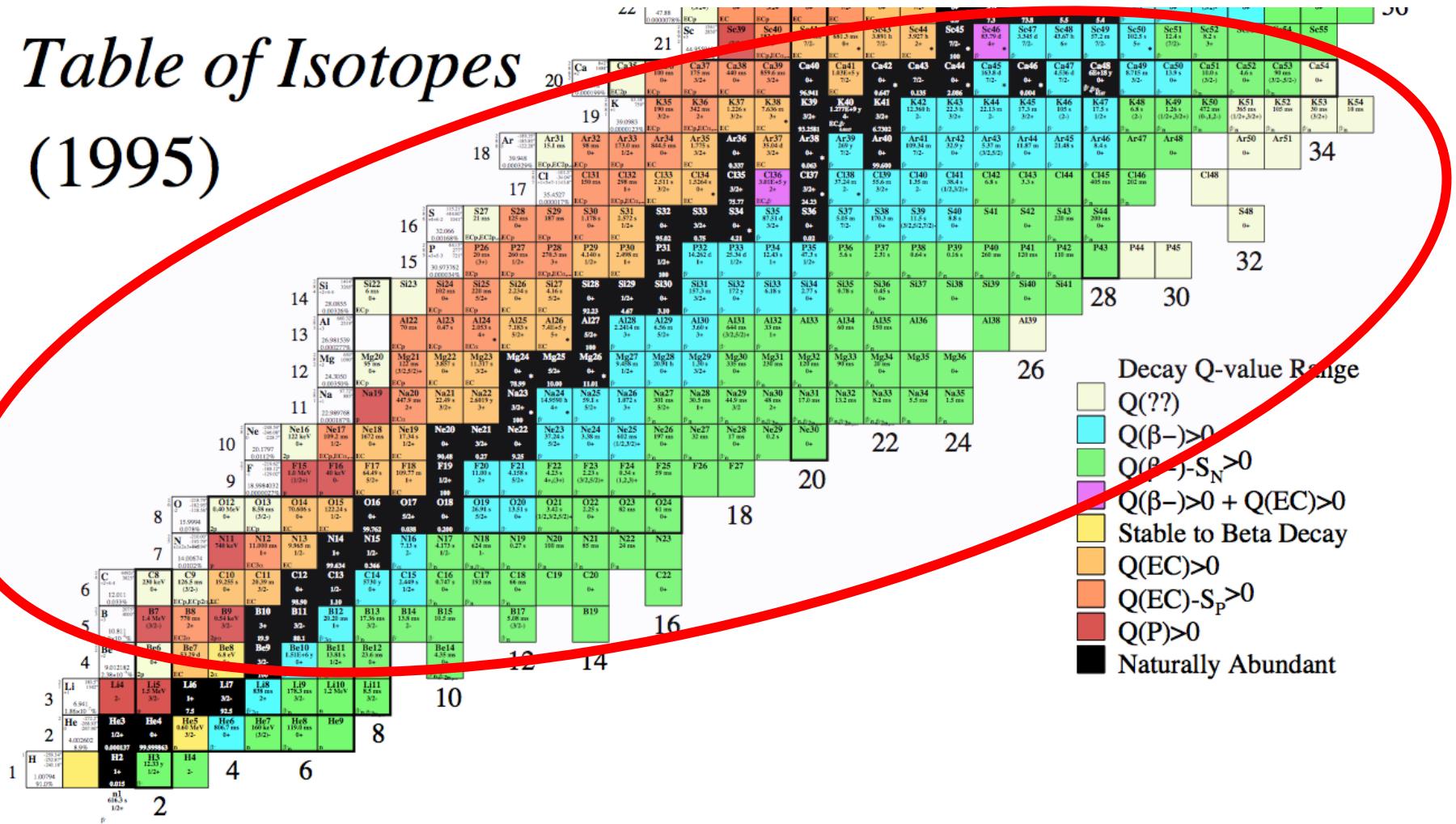
Table of Isotopes (1995)



for which

- few/no elastic-scattering data & phenom. potential information is available.

Table of Isotopes (1995)



Global potential for projectiles of unstable nuclei up to driplines
T. Furumoto, W. Horiuchi, M. Takashina, Y. Yamamoto, Y. Sakuragi (in preparation)

Global potential for projectiles of unstable nuclei up to driplines

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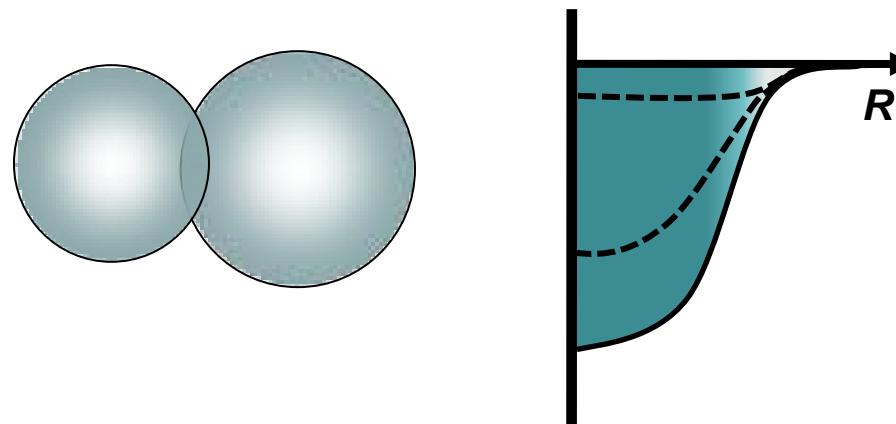
| With CEG07b | 34_Ca | 36_Ca | 38_Ca | 40_Ca | 42_Ca | 44_Ca | 46_Ca | 48_Ca | 50_Ca | 52_Ca | 54_Ca | 56_Ca | 58_Ca | 60_Ca | 62_Ca | 64_Ca | 66_Ca | 68_Ca | 70_Ca |
|-------------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|
| | 30_Ar | 32_Ar | 34_Ar | 36_Ar | 38_Ar | 40_Ar | 42_Ar | 44_Ar | 46_Ar | 48_Ar | 50_Ar | 52_Ar | 54_Ar | 56_Ar | 58_Ar | 60_Ar | 62_Ar | | |
| | 26_S | 28_S | 30_S | 32_S | 34_S | 36_S | 38_S | 40_S | 42_S | 44_S | 46_S | 48_S | 50_S | 52_S | | | | | |
| | 22_Si | 24_Si | 26_Si | 28_Si | 30_Si | 32_Si | 34_Si | 36_Si | 38_Si | 40_Si | 42_Si | 44_Si | 46_Si | 48_Si | | | | | |
| | 20_Mg | 22_Mg | 24_Mg | 26_Mg | 28_Mg | 30_Mg | 32_Mg | 34_Mg | 36_Mg | 38_Mg | 40_Mg | | | | | | | | |
| | 16_Ne | 18_Ne | 20_Ne | 22_Ne | 24_Ne | 26_Ne | 28_Ne | 30_Ne | 32_Ne | 34_Ne | 36_Ne | 38_Ne | | | | | | | |
| | 12_O | 14_O | 16_O | 18_O | 20_O | 22_O | 24_O | | | | | | | | | | | | |
| | 8_C | 10_O | 12_C | 14_C | 16_C | 18_C | 20_C | 22_C | | | | | | | | | | | |

Global parameterization of the CEG07 folding-model potentials

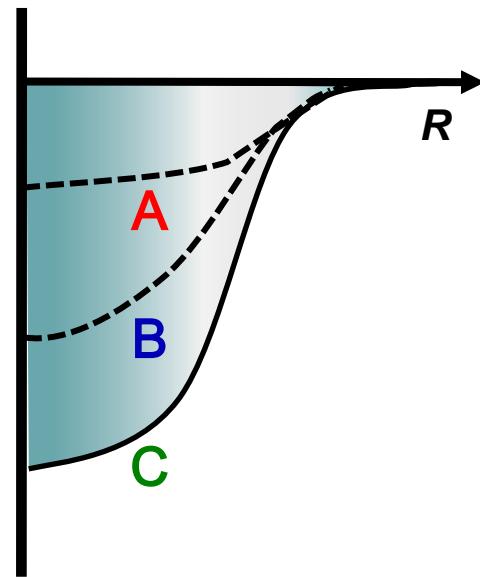
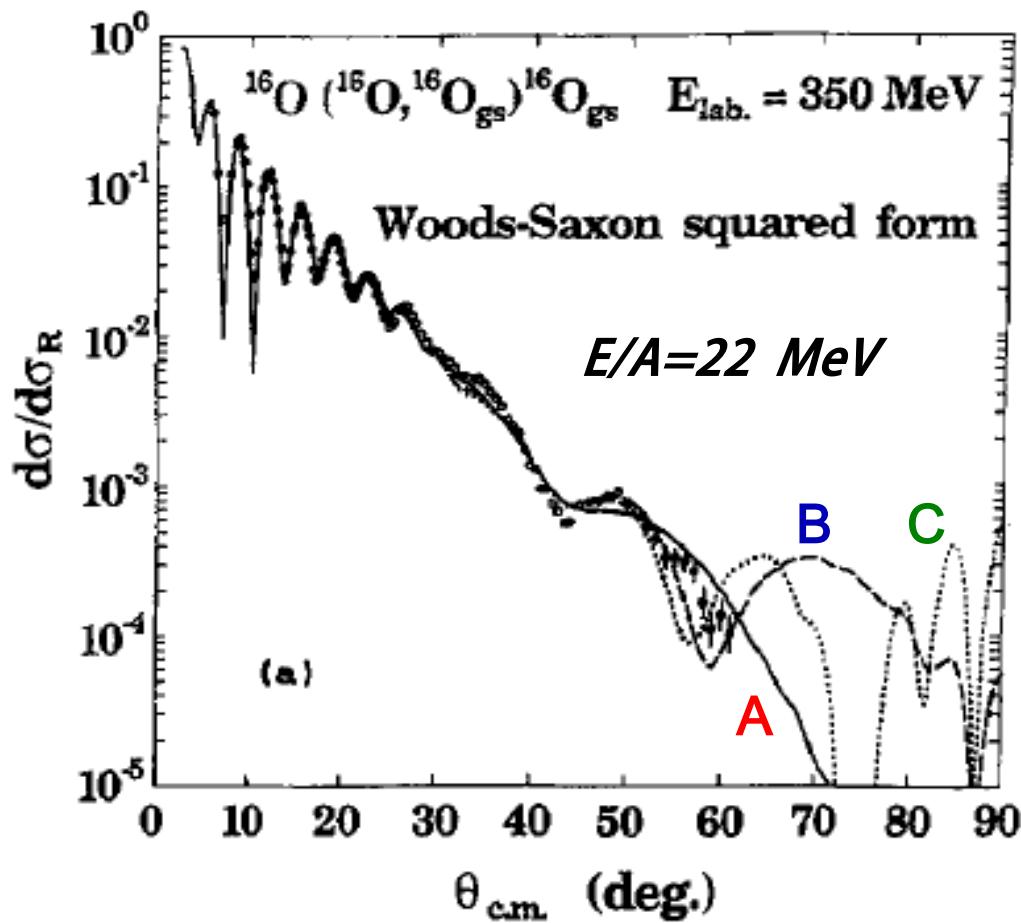
- ✓ projectiles : Z = 6 (C isotope) ~ 20 (Ca isotope) (even-even)
- ✓ targets : ^{12}C ~ ^{208}Pb (closed or sub-closed shell nuclei)
- ✓ energy range : E/A = 30 ~ 400 MeV

$$U_{opt}(R) = V_{opt}(R) + i W_{opt}(R) : \text{complex potential}$$

- *Phenomenological optical potentials:*
 - ✓ needs *Exp. Data* (*elastic scattering*)
to determine *potential parameters*
(e.g. Woods-Saxon form)
 - ✓ optical potential for heavy-ion systems (AA) has
large ambiguity in depth & shape
due to *strong absorption* (in most cases)
 - ✓ → only sensitive to potential at *nuclear surface*



discrete ambiguity of optical potential → which is correct ?



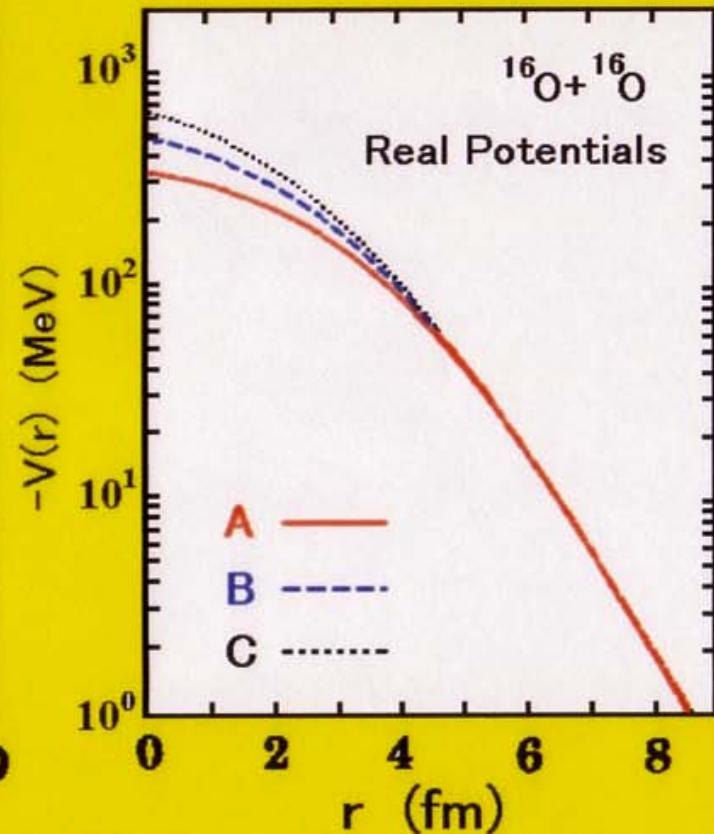
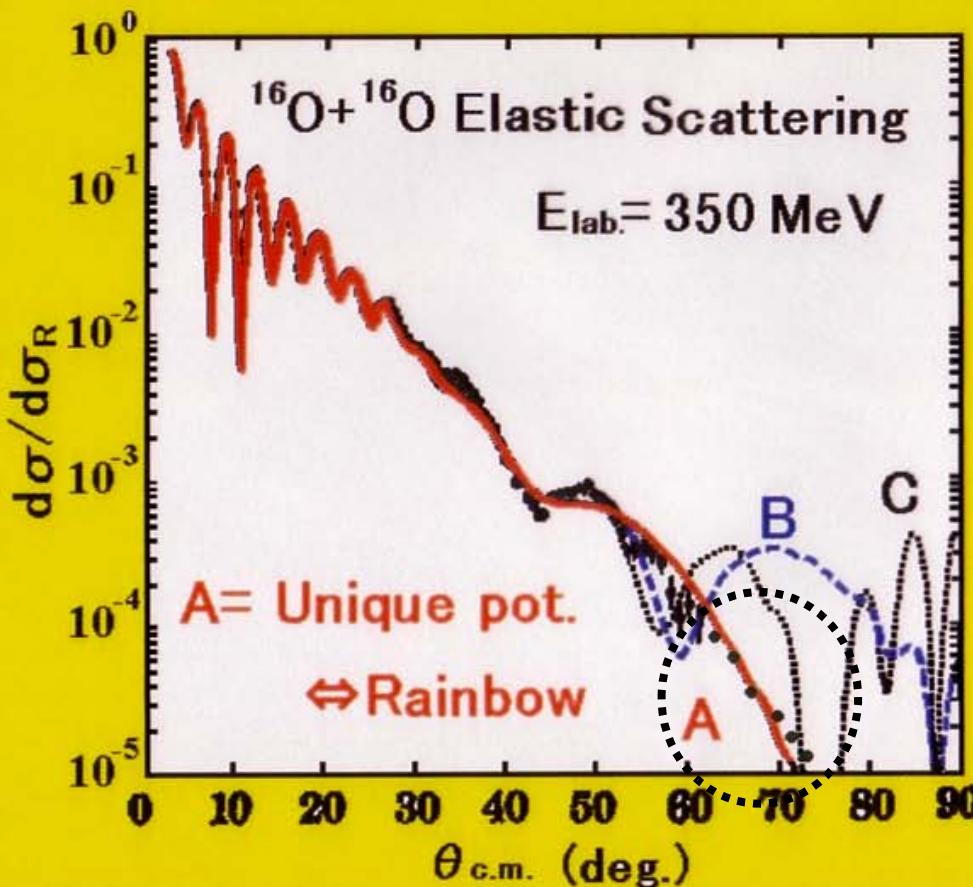
Y. Kondō, F. Michel and G. Reidemeister, Phys. Lett. B 242 (1990) 340.

- ✓ *In general, it is rather difficult to probe the short-range part of H.I. potentials, due to **strong absorption**.*
- ✓ Can we probe H.I. potential at short distances?
 - ***Yes, we can!***
(at least for light heavy-ions)

*by the measurements of
refractive scattering at high- q region (backward),
such as nuclear-rainbow phenomena.*

■ Phenomenological Analysis of the Nuclear Rainbow

- $^{16}\text{O} + ^{16}\text{O}$ $E_{\text{lab}} = 350 \text{ MeV}$ (HMI 1989)
- The data up to 61° are reproduced by A, B and C.
- "A" pot. is found to be a unique deep potential for this system by the fits to the data up to 73° .



But, good quality of exp. data are not always in our hands.

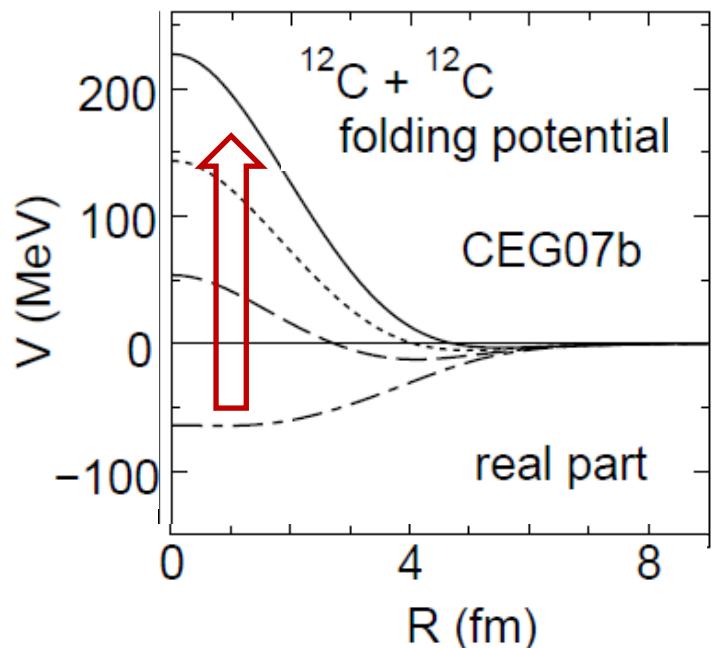
→ We need a **microscopic theory** that **explains & predicts**

- ✓ *correct depth & shape of heavy-ion optical potentials,
(hopefully, of both the real and imaginary parts)*
- ✓ *including unstable nuclei (n -rich & p -rich isotopes)*
- ✓ *correct energy dependence over the wide range of
incident energy, up to a few hundred MeV/u*

*starting from **bare NN interaction** in free space*

one of the
key word of the present talk

- ***attractive-to-repulsive transition of the optical potentials for heavy-ion systems with the increasing energy***

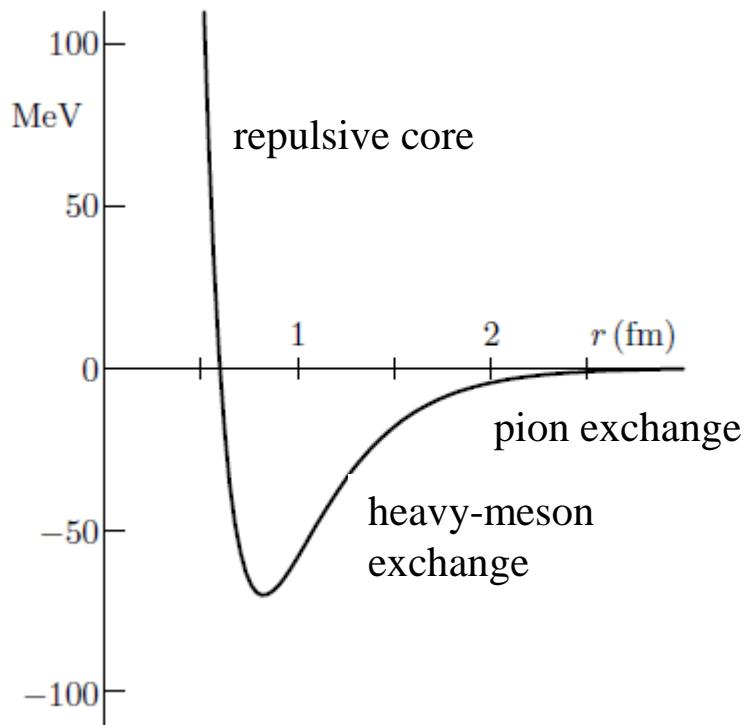


$E/A=100\sim400\text{ MeV}$

***T. Furumoto, Y. Sakuragi,
Y. Yamamoto,
PRC82 (2010) 044612***

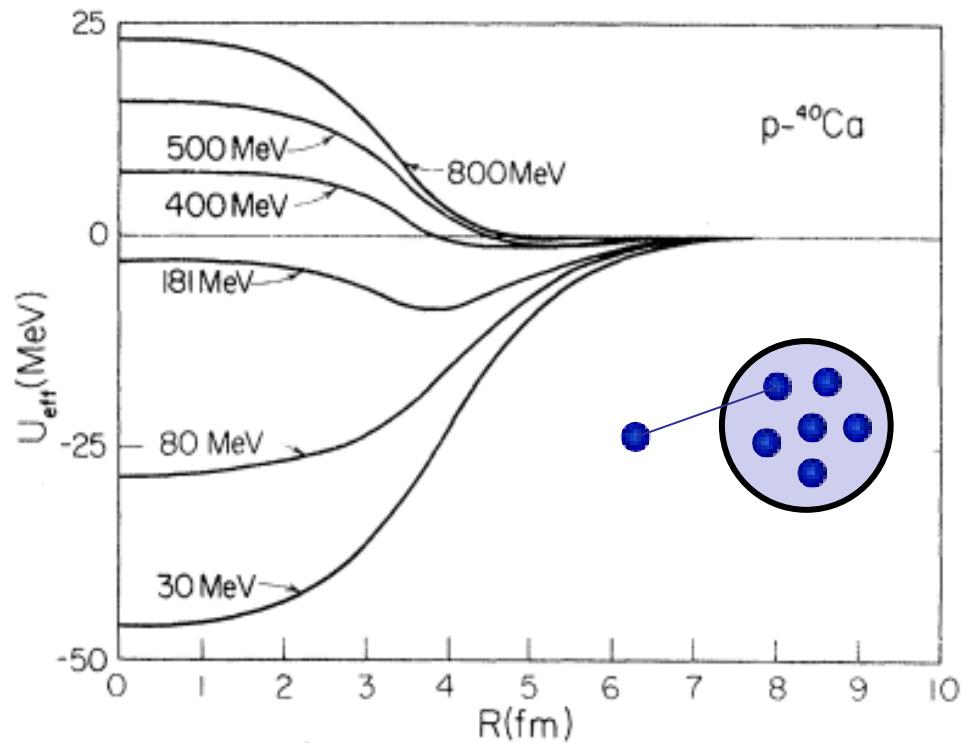
◆ NN interaction :

- long-range **attraction**
- short-range **repulsive core**



◆ nucleon-nucleus (NA) interaction :

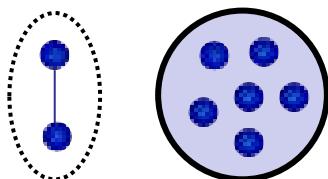
- **attractive** at low energies ($E < 200$ MeV)
- **wine-bottle-bottom (WBB)** around transitional energies
- **repulsive** at high energies ($E > 500$ MeV)



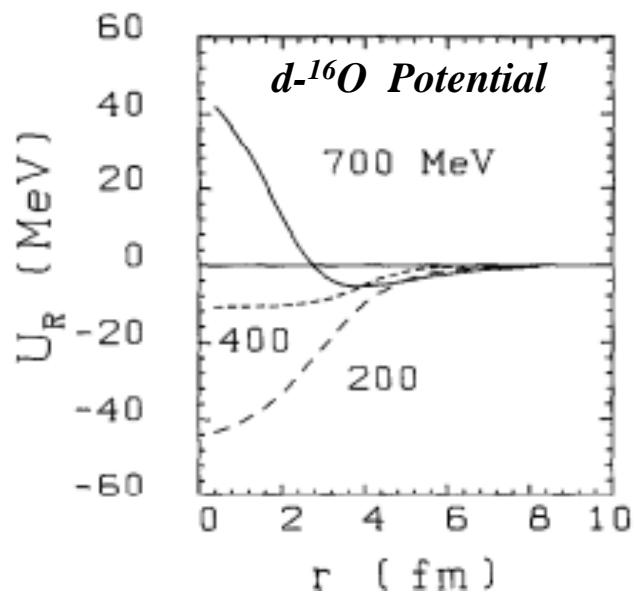
L.G.Arnold, (Phys.Rev.C25(1982)936

◆ *d-A interaction:*

- *similar behavior to NA int.*
- $f(d\text{-}A) \sim f(p\text{-}A) + f(n\text{-}A)$



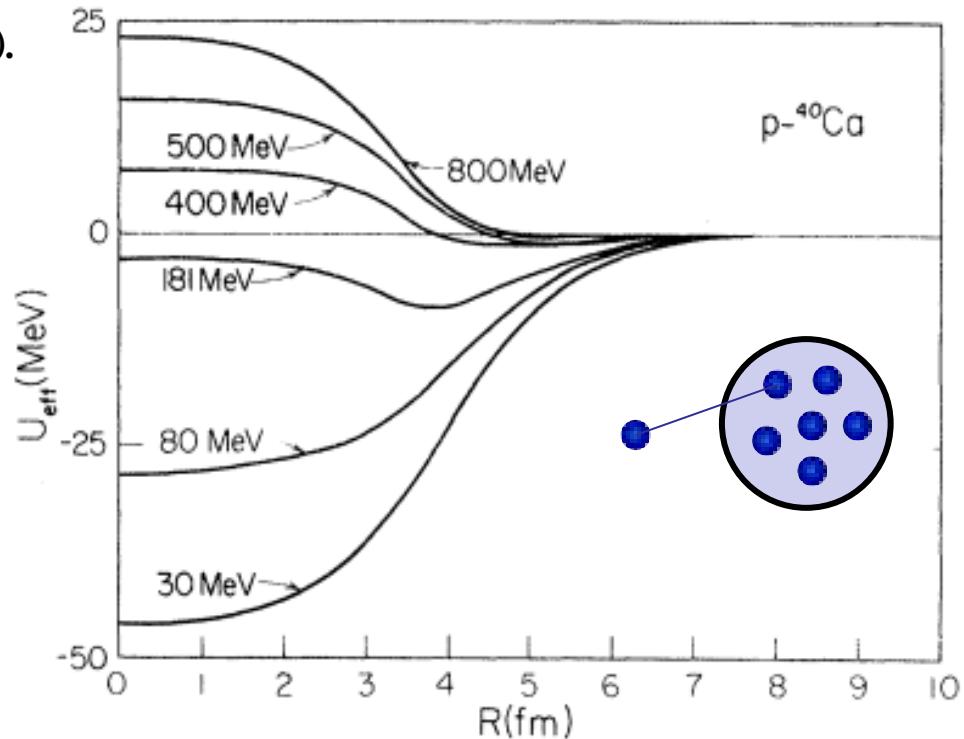
Y. Sakuragi, M. Tanifushi, NPA560, 945(1993).



N.V.Sen, NPA464 (1987) 717

◆ *nucleon-nucleus (NA) interaction:*

- *attractive at low energies ($E < 200$ MeV)*
- *wine-bottle-bottom (WBB) around transitional energies*
- *repulsive at high energies ($E > 500$ MeV)*



L.G.Arnold, (Phys.Rev.C25(1982)936

Q: How about optical potential for **heavy ions**?

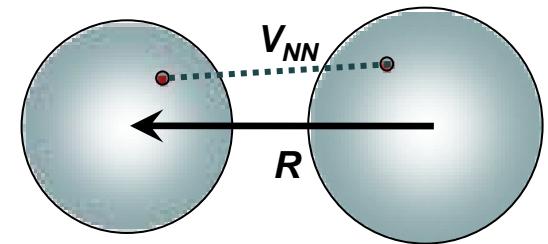
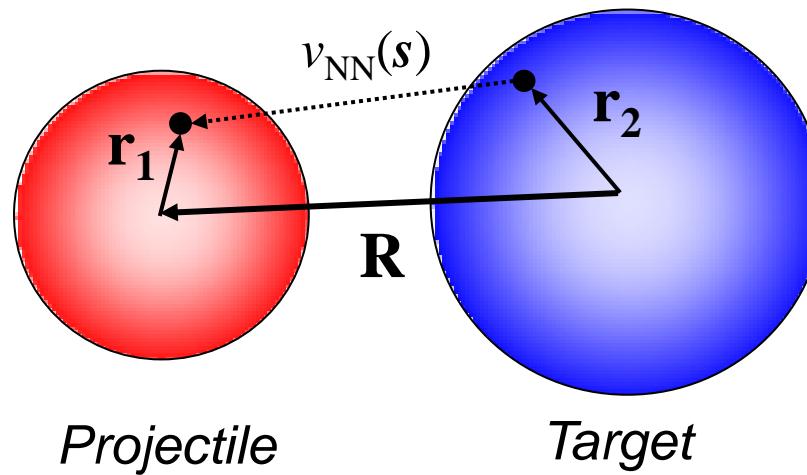
A: according to the predictions of microscopic theory,

- ✓ attractive-to-repulsive transition occurs ?
 - Yes, but thus far we have no experimental evidence.
- ✓ if so, in what energy region?
 - the transition occurs around $E/A = 300 \sim 400 \text{ MeV}$
- ✓ how can we observe the transition, if it really occurs?
 - measure the evolution of elastic scattering angular distribution with increasing energy in the energy range of $E/A = 200 \sim 400 \text{ MeV}$.
- ✓ what are the new ingredients we can learn, if we observe the transition?
 - ① repulsive three-body force (TBF) in nuclear medium & ② tensor force effects
 - besides the genuine repulsive core of NN int.

- **Microscopic / semi-microscopic models :**

- ✓ starting from **NN interactions** (V_{NN})

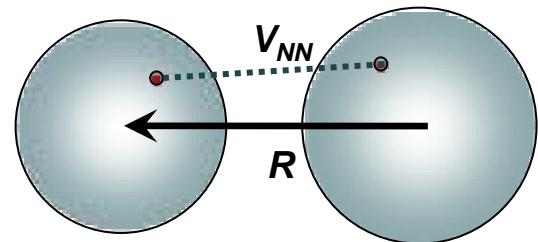
Double-Folding Model (DFM)



$$U_{DFM}(\mathbf{R}) = \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) \underline{v_{NN}(\mathbf{s}; \rho, E)} d\mathbf{r}_1 d\mathbf{r}_2$$

effective NN interaction in nuclear medium

- **Microscopic / semi-microscopic models :**
 - ✓ starting from NN interactions (V_{NN})



- ◆ G-matrix with scattering b.c.

- ✓ V_{NN} : **effective NN interaction in nuclear medium**
 - ✓ *should have proper density-dependence* (ρ -dep)
consistent with nuclear saturation properties
 - ✓ *should have proper energy-dependence* (E -dep)
 - ✓ *should be complex* (*real-part + imaginary part*)

However, no such ideal effective V_{NN} exists so far !

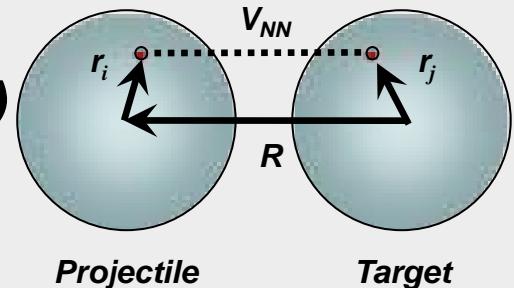


Simple M3Y (1975~1985)

- ✓ **real part only (add a phenom. imag. pot)**
- ✓ **zero-range exchange term**

$$v_{NN}(\mathbf{r}) = 7999 \frac{e^{-4r}}{4r} - 2134 \frac{e^{2.5r}}{2.5r} - \hat{J}_{00}\delta(\mathbf{r})$$

Double-Folding Model (DFM)

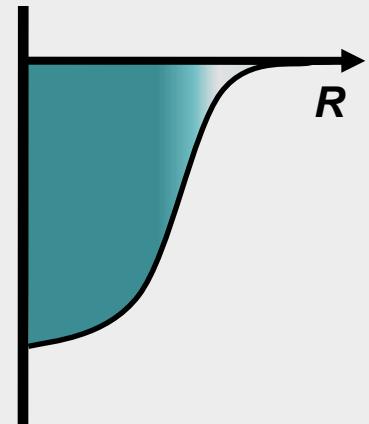


- ✓ **no density-dependence**
⇒ **too deep at short distances, but gives a reasonable strength at nuclear surface**

- ◆ **due to strong absorption for Heavy Ions (HI)**
⇒ **sensitive only to nuclear surface**

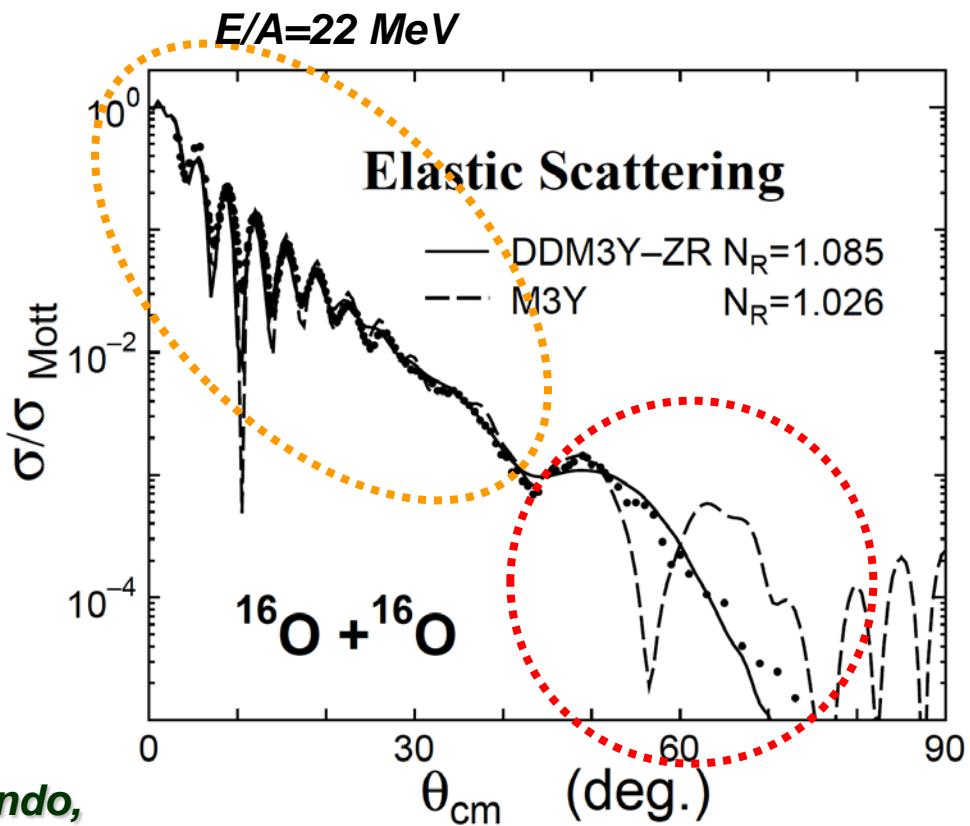
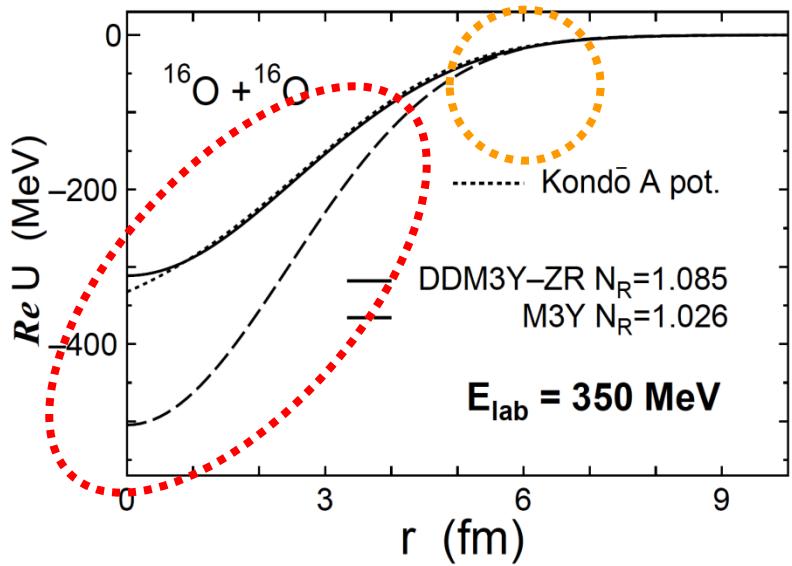
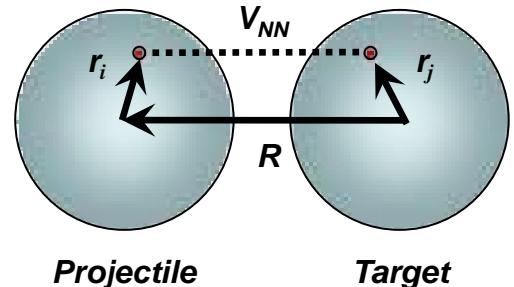
⇒ “Successful” for low-energy ($E/A < 30$ MeV) scattering of heavy-ion (HI) projectiles with $A_p < 40$

[G.R.Satchler and W.G.Love, *Phys.Rep.55,183(1979)*]



Double-Folding-model potentials with M3Y (density-independent)

Double-Folding Model (DFM)



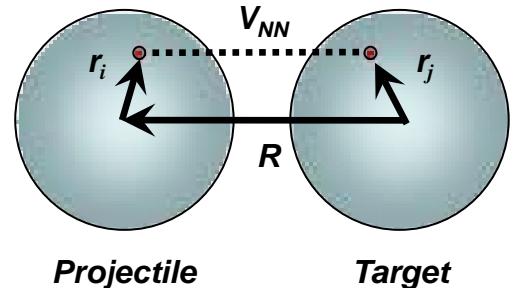
M.Katsuma, Y.Sakuragi, S.Okabe, Y.Kondo,
Prog.Theor.Phys. 107 (2002) 377

- **Introduction of density-dependence** :
DDM3Y-ZR (with zero-range exchange term)

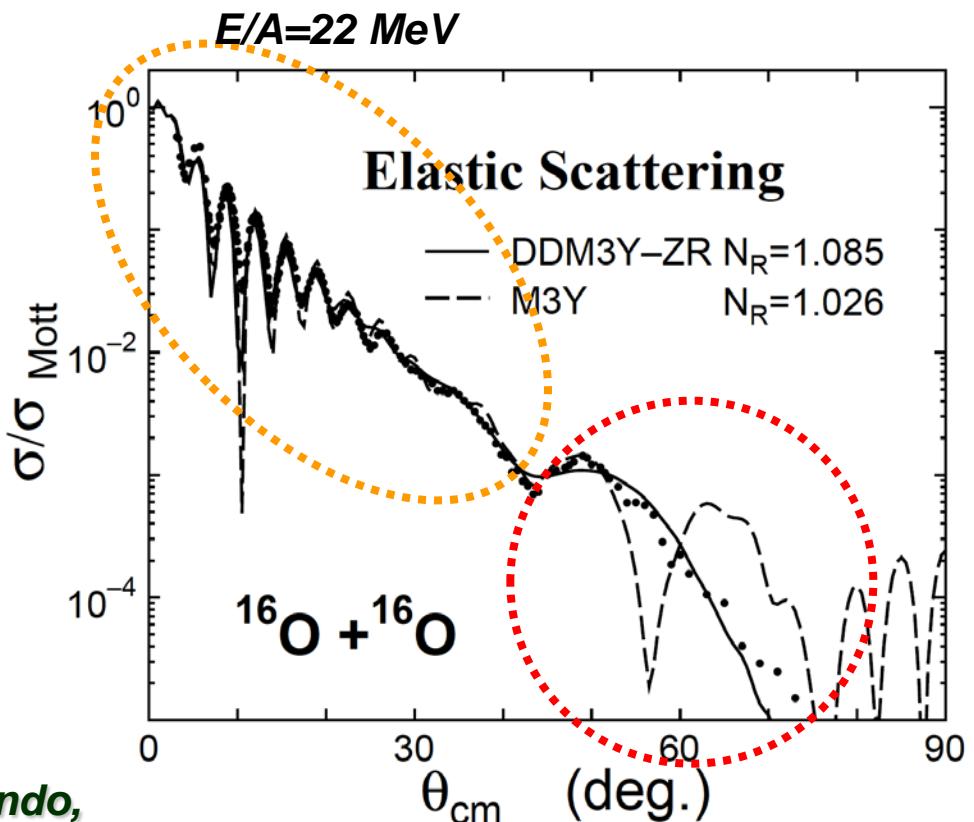
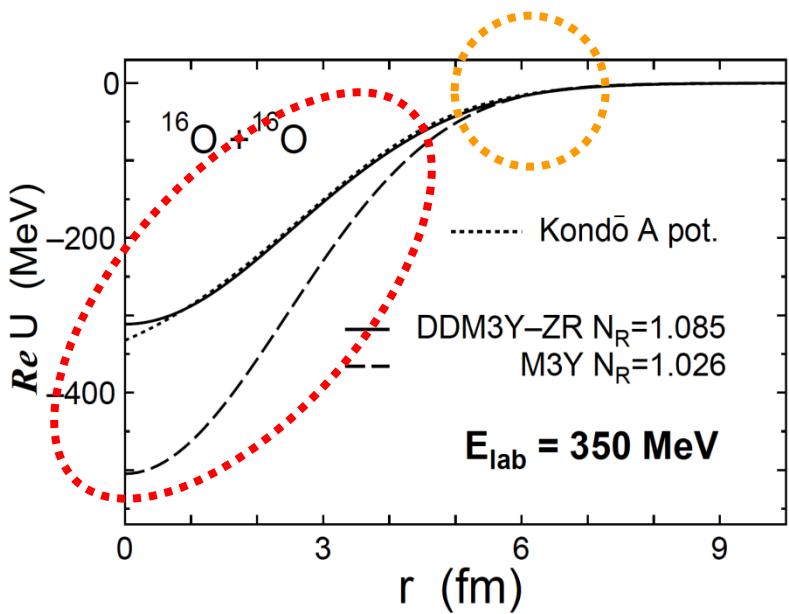
$$v_{NN}(E, \rho; \mathbf{s}) = g(E, \mathbf{s}) \underline{f(E, \rho)}$$

$$\underline{f(E, \rho)} = C(E) [1 + \alpha(E) e^{-\beta(E)\rho}]$$

- ⇒ greatly **reduce** the potential strength at **short distances**
- ⇒ reproduce refractive phenomena, such as **nuclear-rainbow** (eg. ${}^4\text{He} + \text{A}$, ${}^{16}\text{O} + {}^{16}\text{O}$)



**Double-Folding-model potentials
with M3Y (density-independent)
with DDM3Y (density-dependent)**



*M.Katsuma, Y.Sakuragi, S.Okabe, Y.Kondo,
Prog.Theor.Phys. 107 (2002) 377*

New complex G-matrix interaction (CEG07)

1. derived from ESC04

“ESC04” : the latest version of **Extended Soft-Core** force designed for **NN**, **YN** and **YY** systems

Th. Rijken, Y. Yamamoto, Phys. Rev. C 73 (2006) 044008

2. Three body force

Three-body attraction (TBA) : Fujita-Miyazawa type

Three-body repulsion (TBR) : triple-meson correl.

3. up to higher density region

for the local density prescription in the case of DFM

- ◆ *T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC78 (2008) 044610,*
- ◆ *T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC79 (2009) 011601(R),*
- ◆ *T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC80 (2009) 044614*
- ◆ *T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 029908(E)*
- ◆ *T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 044612*

Complex G-matrix interaction (CEG07)

T.Furumoto, Y. Sakuragi and Y. Yamamoto, Phys. Rev. C 78 (2008) 044610

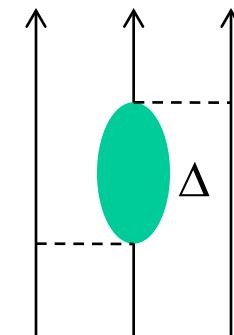
Extended Soft-Core model : “ESC04” force

designed for NN, YN and YY interactions

Th. Rijken, Y. Yamamoto, Phys.Rev.C 73 (2006) 044008

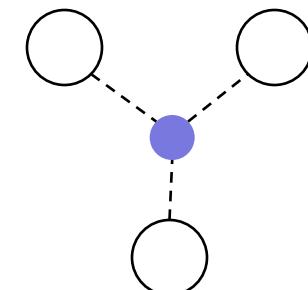
1. Three-body attractive (TBA)

- originated from Fujita-Miyazawa diagram
- important at low density region

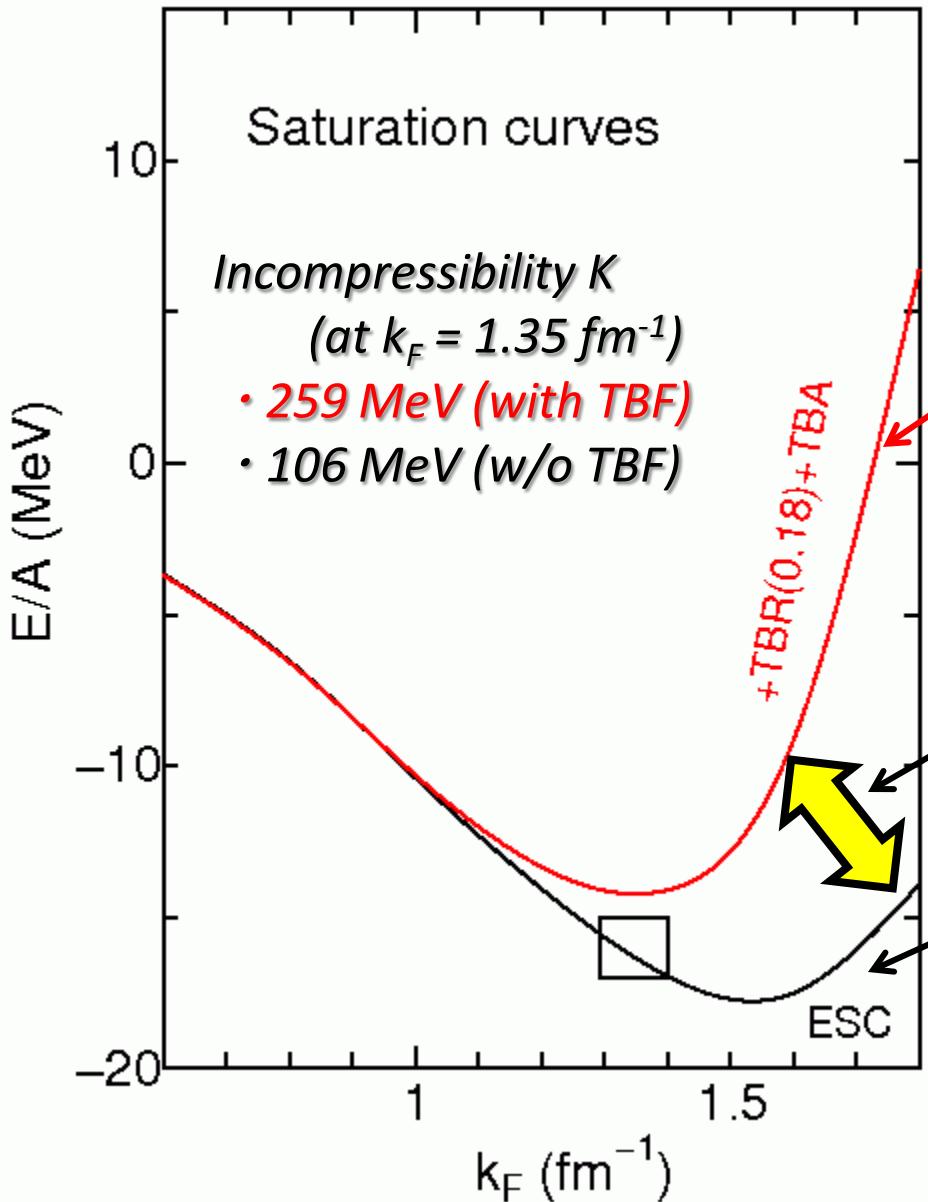


2. Three-body repulsive (TBR)

- universal three-body repulsion (NNN, NNY, NYY)
originated from triple-meson correlation
- important at high-density region



New complex G-matrix interaction (CEG07)



$$\frac{E}{A} = \frac{3}{5} \frac{k_F^2}{2m} + \frac{1}{2} \sum_{k_1, k_2 < k_F} \langle k_1 k_2 | G(\omega) | k_1 k_2 \rangle$$

CEG07b

- +Three body repulsive (TBR)
- +Three body attractive (TBA)

Decisive role to make the saturation curve realistic

CEG07a

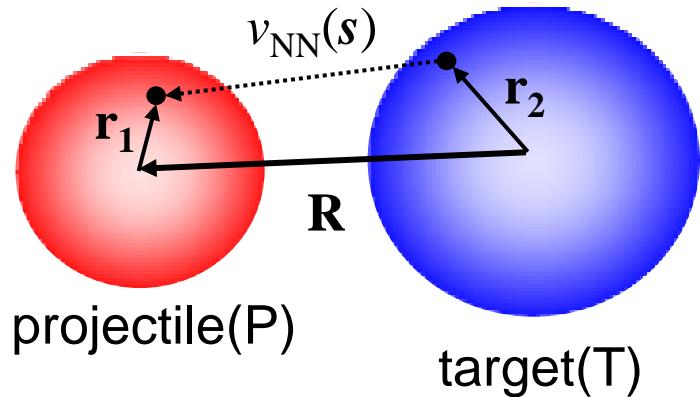
With only Two-Body Force

T.Furumoto, Y. Sakuragi , Y. Yamamoto,
Phys. Rev. C 78 (2008) 044610

- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC79 (2009) 011601(R),
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC80 (2009) 044614
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 029908(E)
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 044612

Double folding Potential with complex-G (CEG07)

$$U(\mathbf{R}) = \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) g_D(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2$$



$$+ \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) g_{EX}(\mathbf{s}; \rho, E) \exp\left[i \frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2$$

$$= V_{DFM}(\mathbf{R}) + iW_{DFM}(\mathbf{R})$$

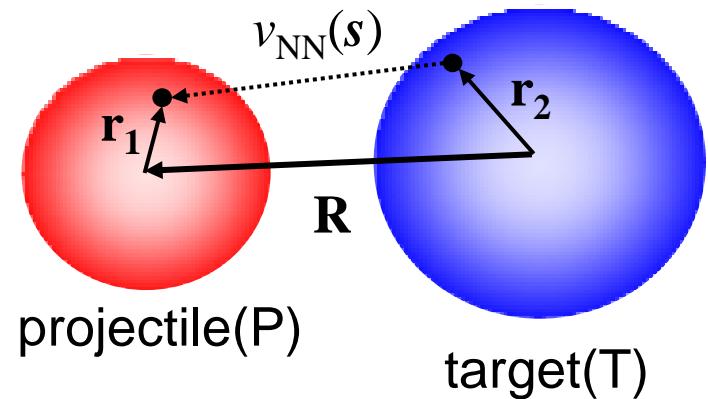
➤ Complex G-matrix interaction (**CEG07**)

$$g_{D,EX} = g_{D,EX}^{(real)} + ig_{D,EX}^{(imag)}$$

- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC79 (2009) 011601(R),
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC80 (2009) 044614
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 029908(E)
- ◆ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) (in press)

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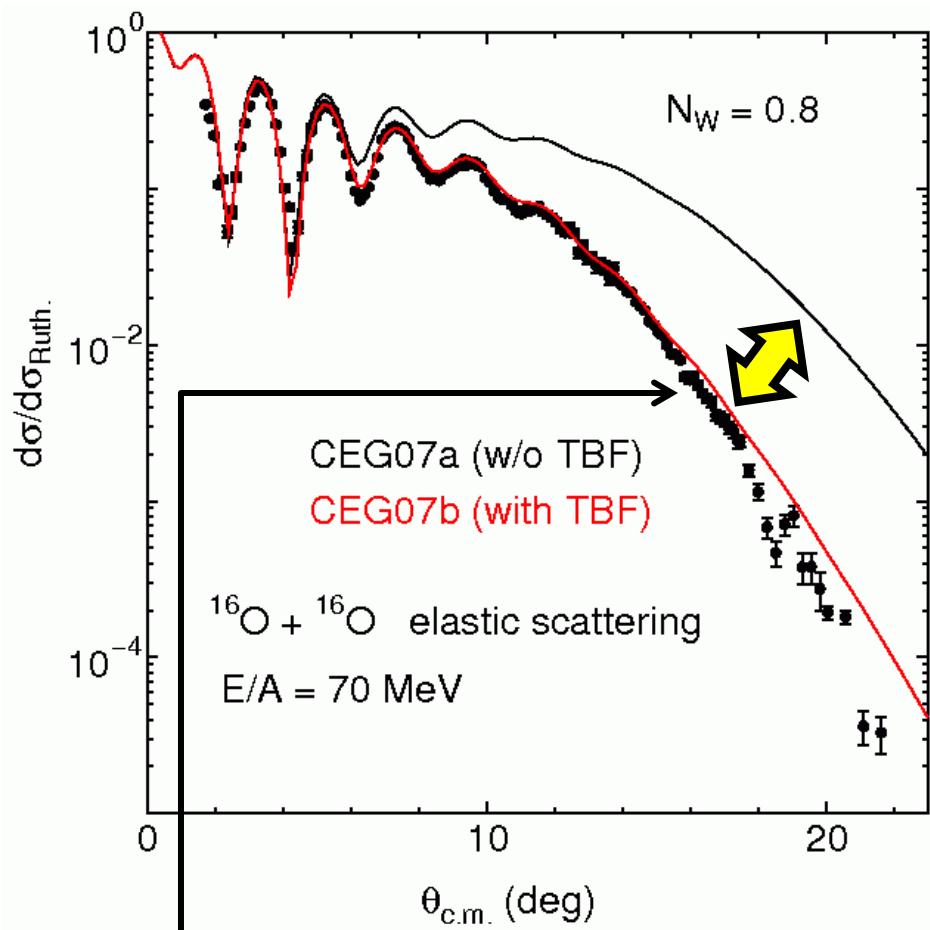
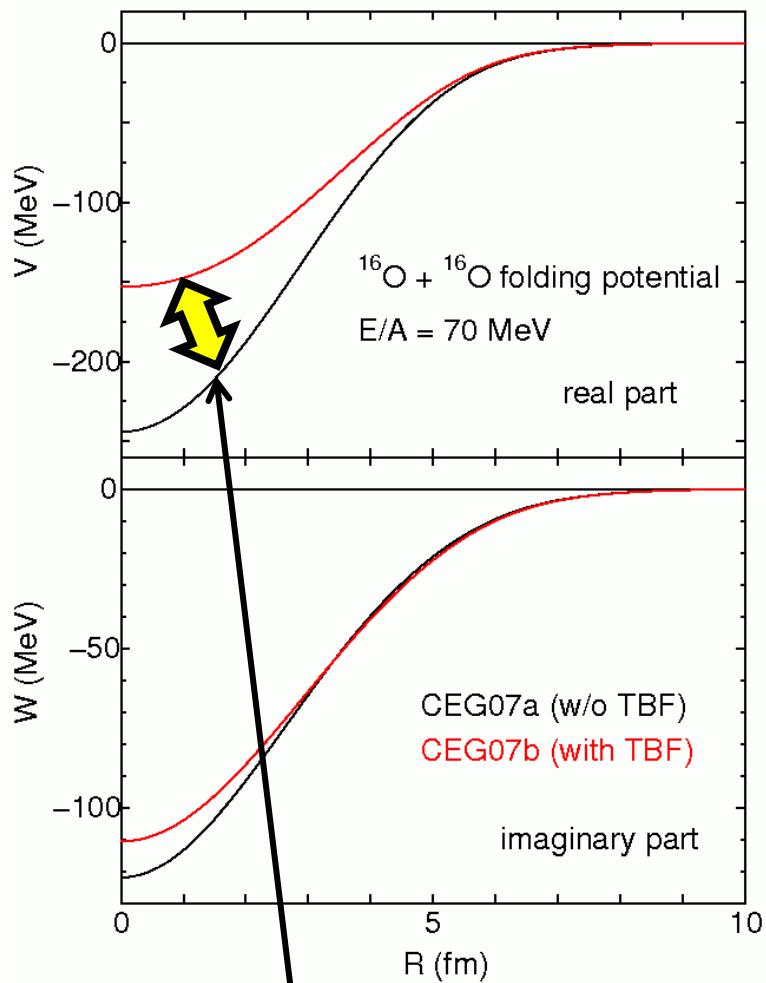


$$\begin{aligned} &+ \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) g_{EX}(\mathbf{s}; \rho, E) \exp\left[i \frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2 \\ &= V_{DFM}(\mathbf{R}) + i W_{DFM}(\mathbf{R}) \end{aligned}$$

✓ Renormalization factor for the **imaginary** part

$$\rightarrow U_{DFM} = V_{DFM} + i N_W W_{DFM}$$

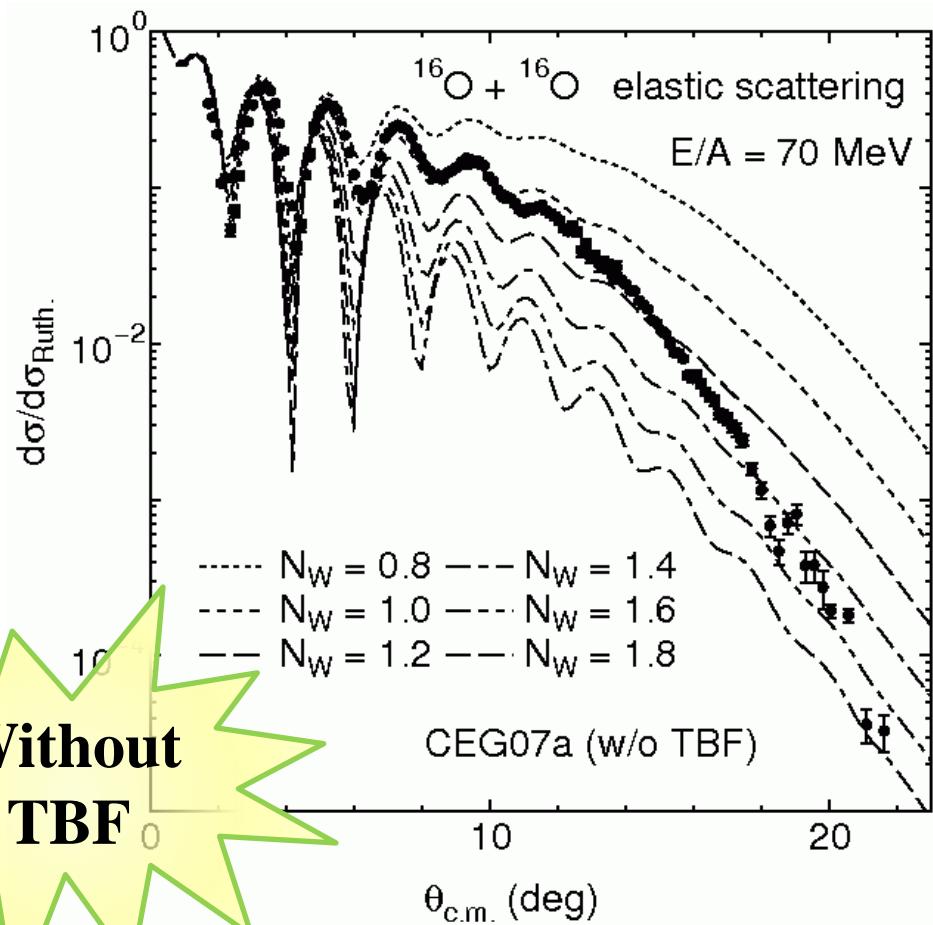
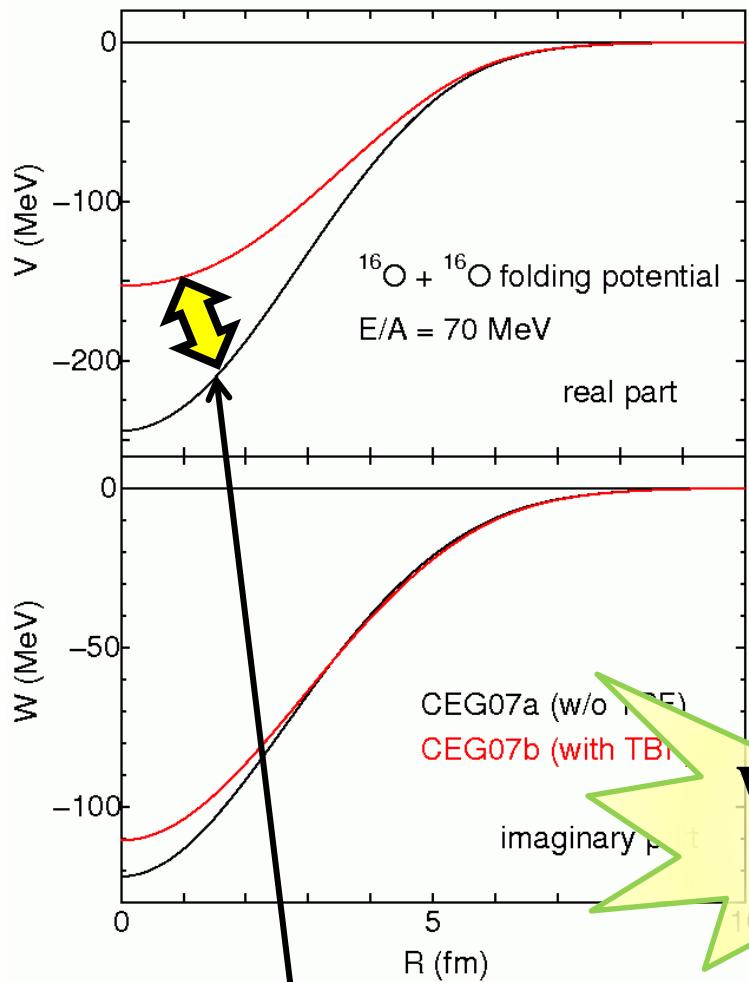
$^{16}\text{O} + ^{16}\text{O}$ elastic scattering $E/A = 70$ MeV



important effect of three-body force

$$U_{DFM} = V_{DFM} + iN_W W_{DFM}$$

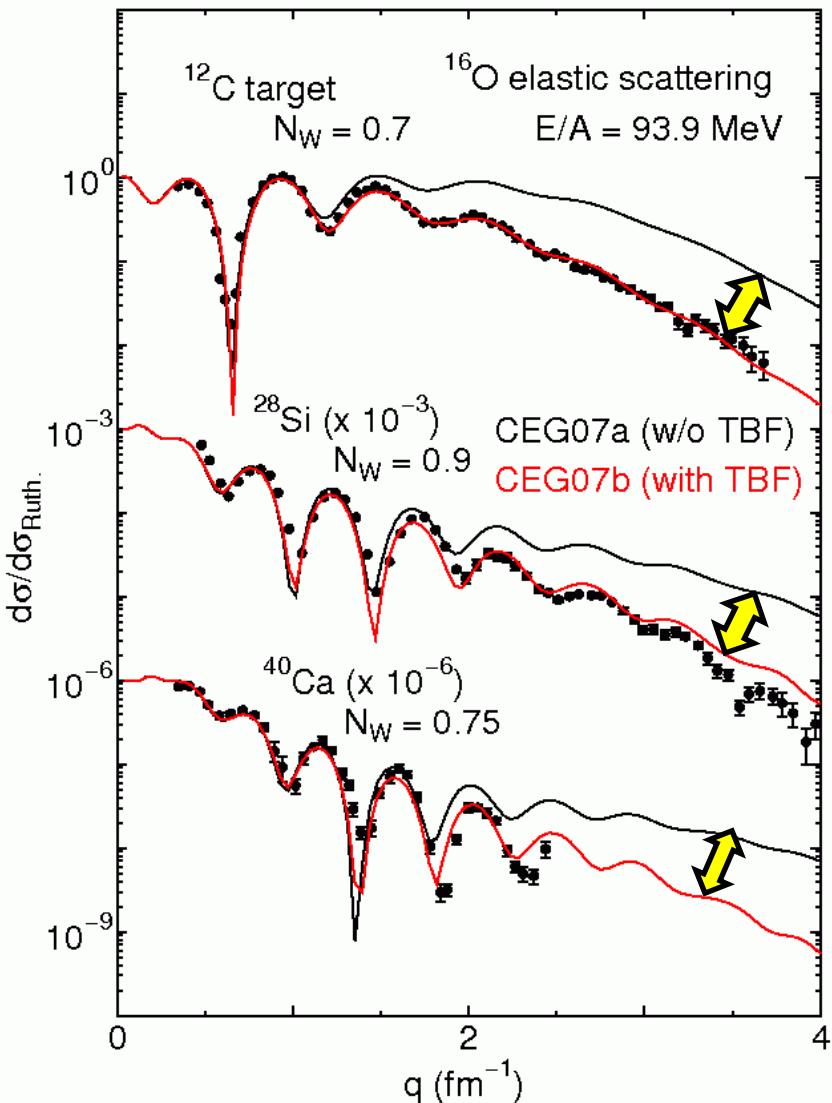
$^{16}\text{O} + ^{16}\text{O}$ elastic scattering $E/A = 70$ MeV



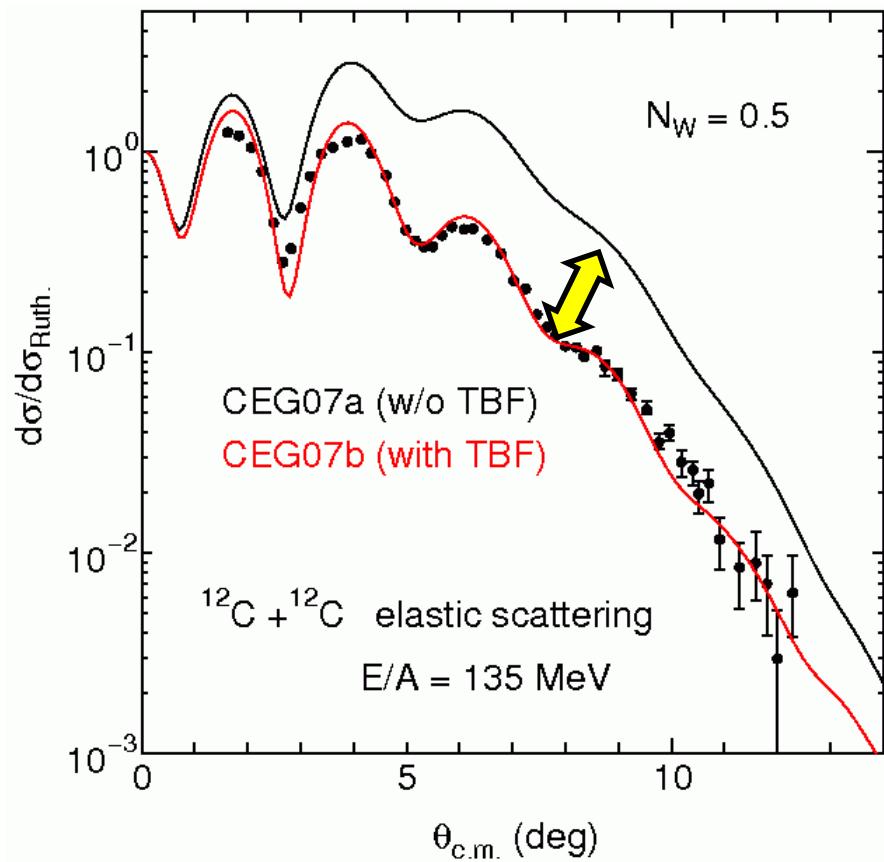
important effect of three-body force

$$U_{DFM} = V_{DFM} + iN_W W_{DFM}$$

$^{16}\text{O} + ^{12}\text{C}, ^{28}\text{Si}, ^{40}\text{Ca}$



$^{12}\text{C} + ^{12}\text{C}$ elastic scattering

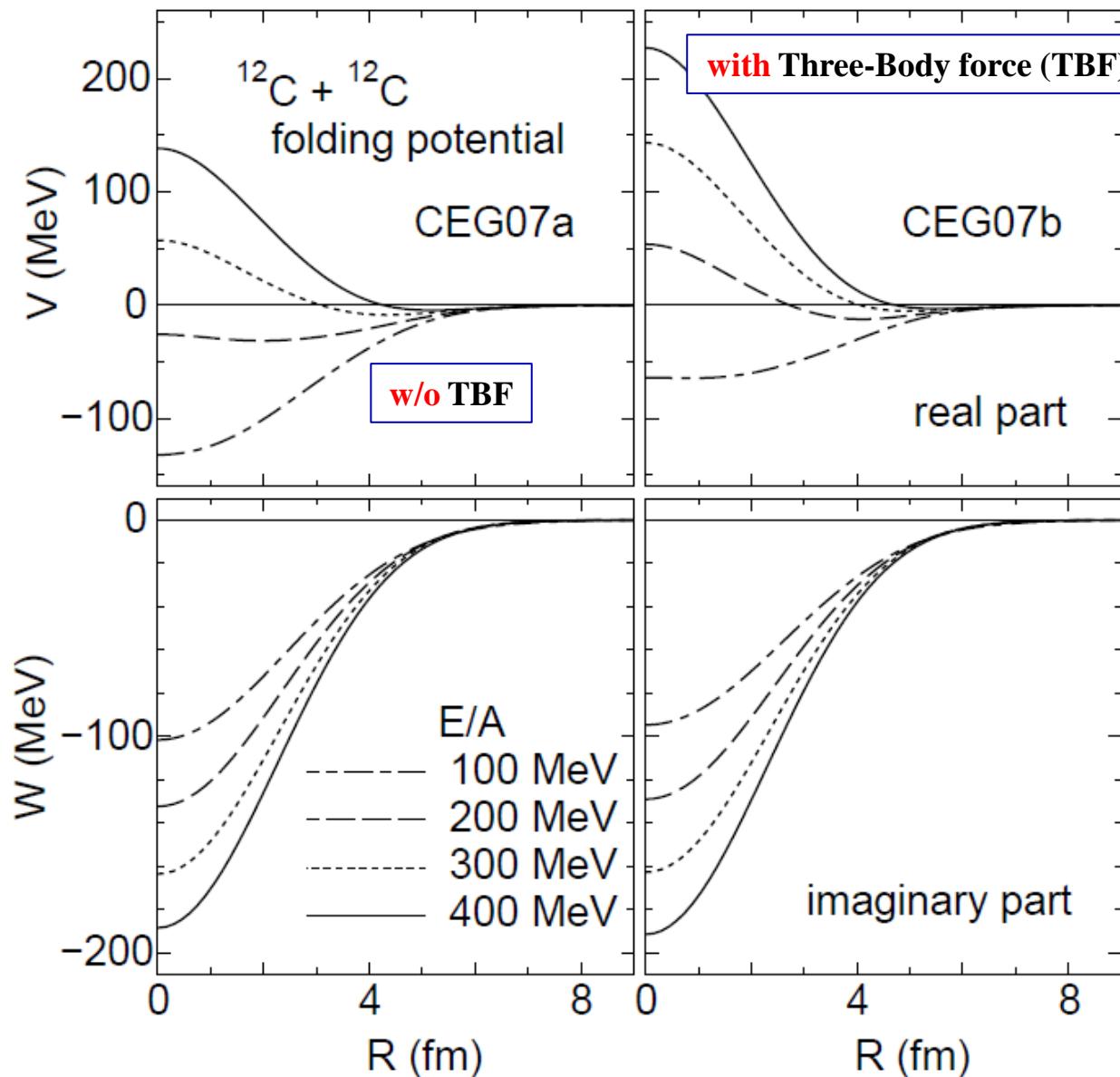


$$U_{DFM} = V_{DFM} + iN_W W_{DFM}$$

important effect of **three-body force**

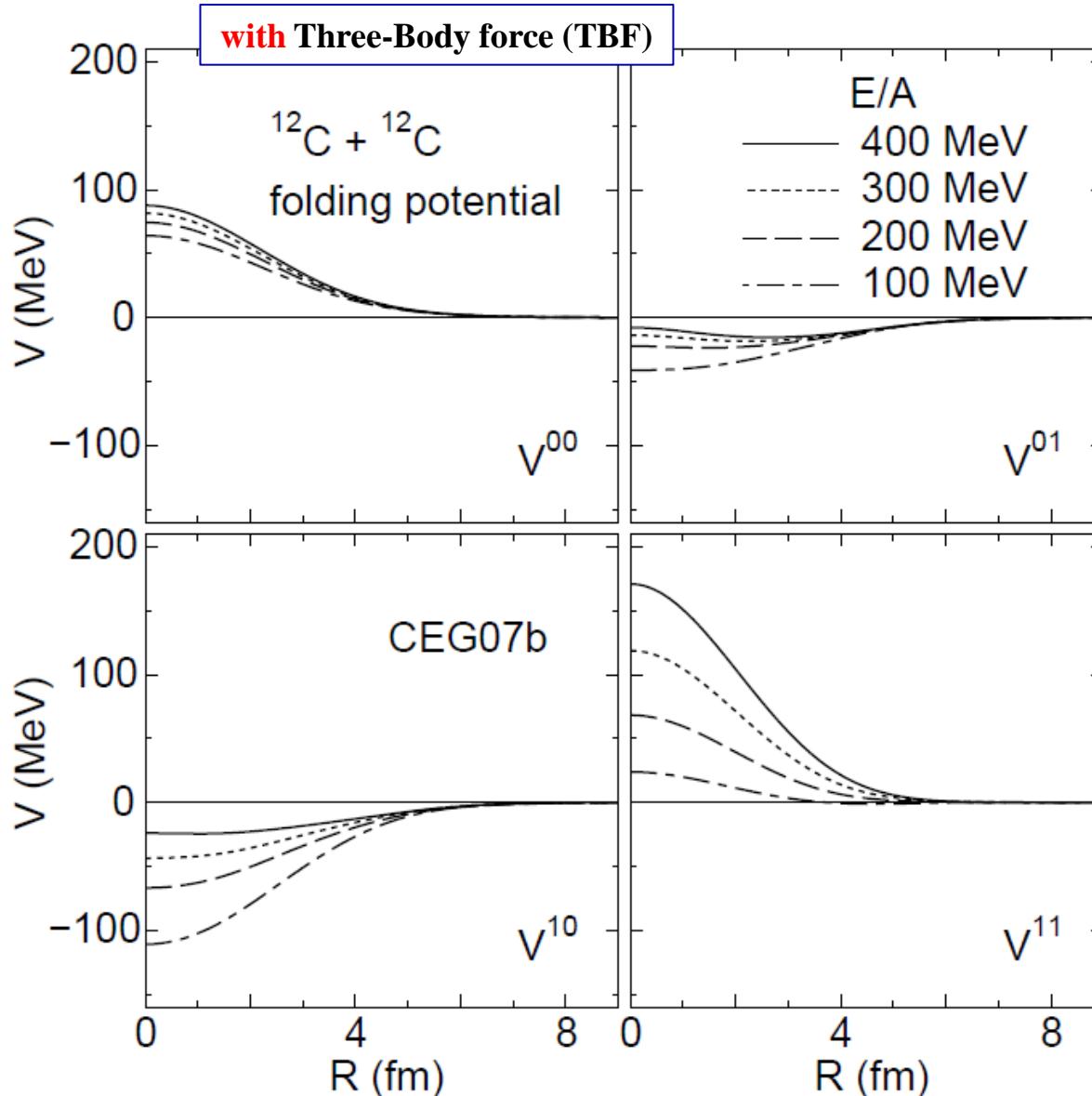
$^{12}\text{C} + ^{12}\text{C}$ elastic scattering at $E/A = 100 \sim 400$ MeV

➤ real potential becomes **repulsive** around $E/A = 300 \sim 400$ MeV



*T. Furumoto, Y. Sakuragi,
Y. Yamamoto,
PRC82 (2010) 044612*

NN tensor force plays an essential role in the **attractive-to-repulsive transition** of the **A-A potentials**

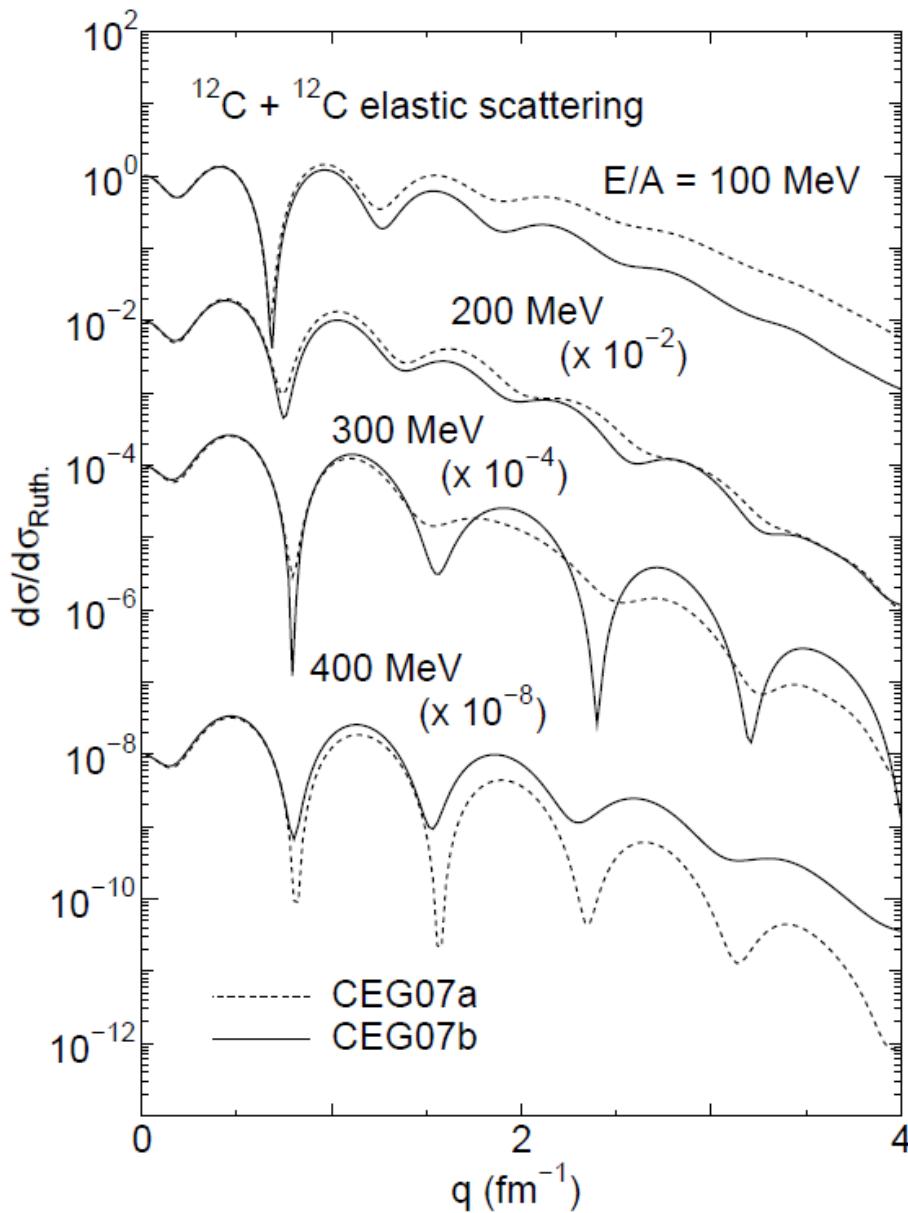


spin(S) and isospin(T)
components V^{ST}
of folding potential

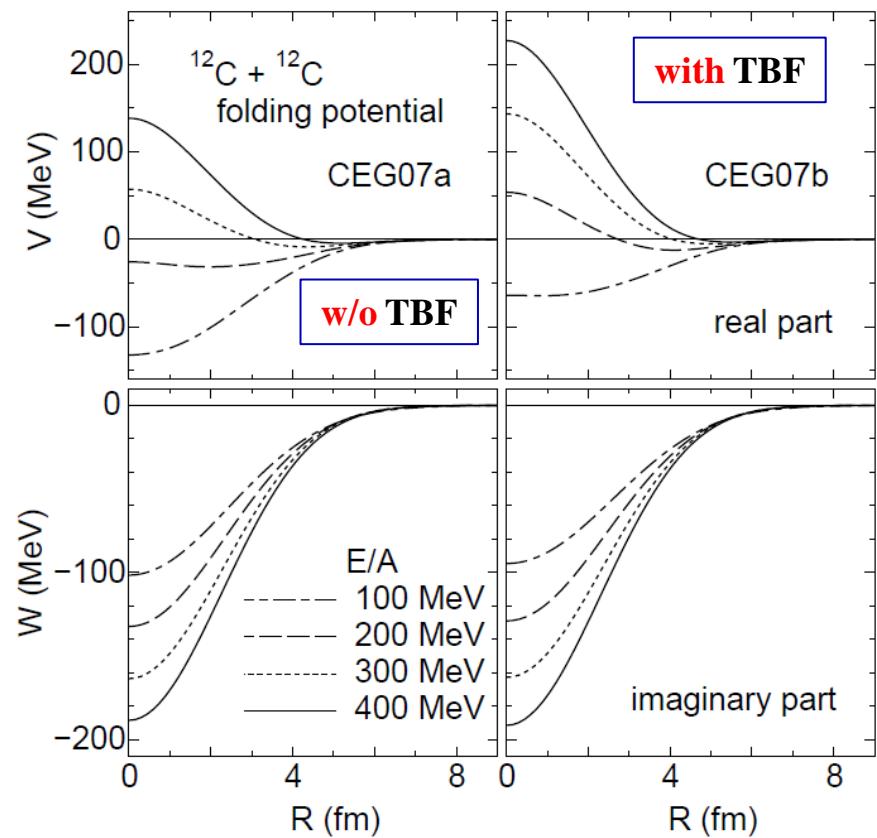
★ $(S,T) = (0,0)$ and $(0,1)$
do not include the tensor
force.

★ $(S,T) = (1,0)$ and $(1,1)$
components include the
tensor force,

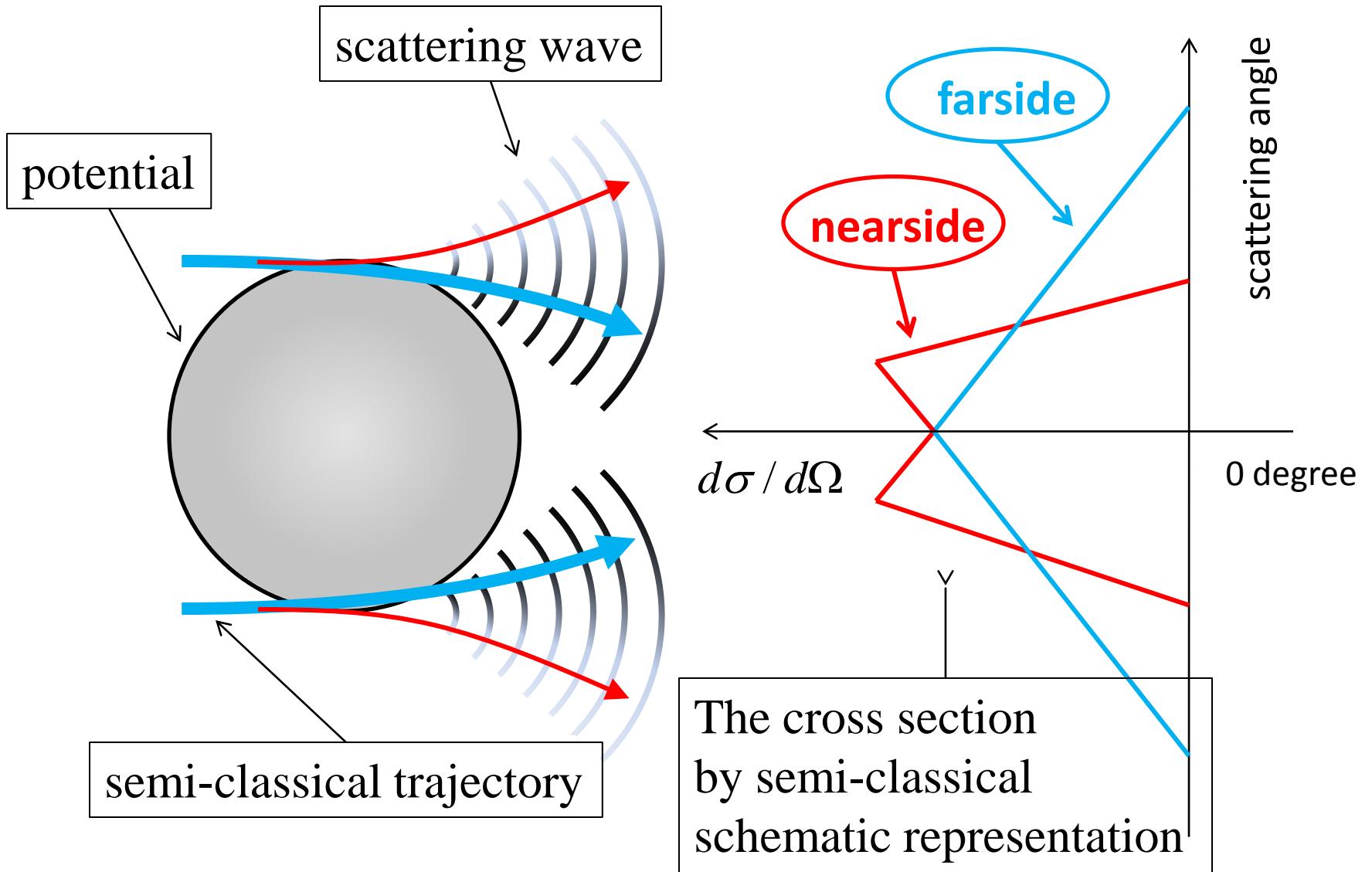
$^{12}\text{C} + ^{12}\text{C}$ elastic scattering at $E/A = 100 \sim 400$ MeV



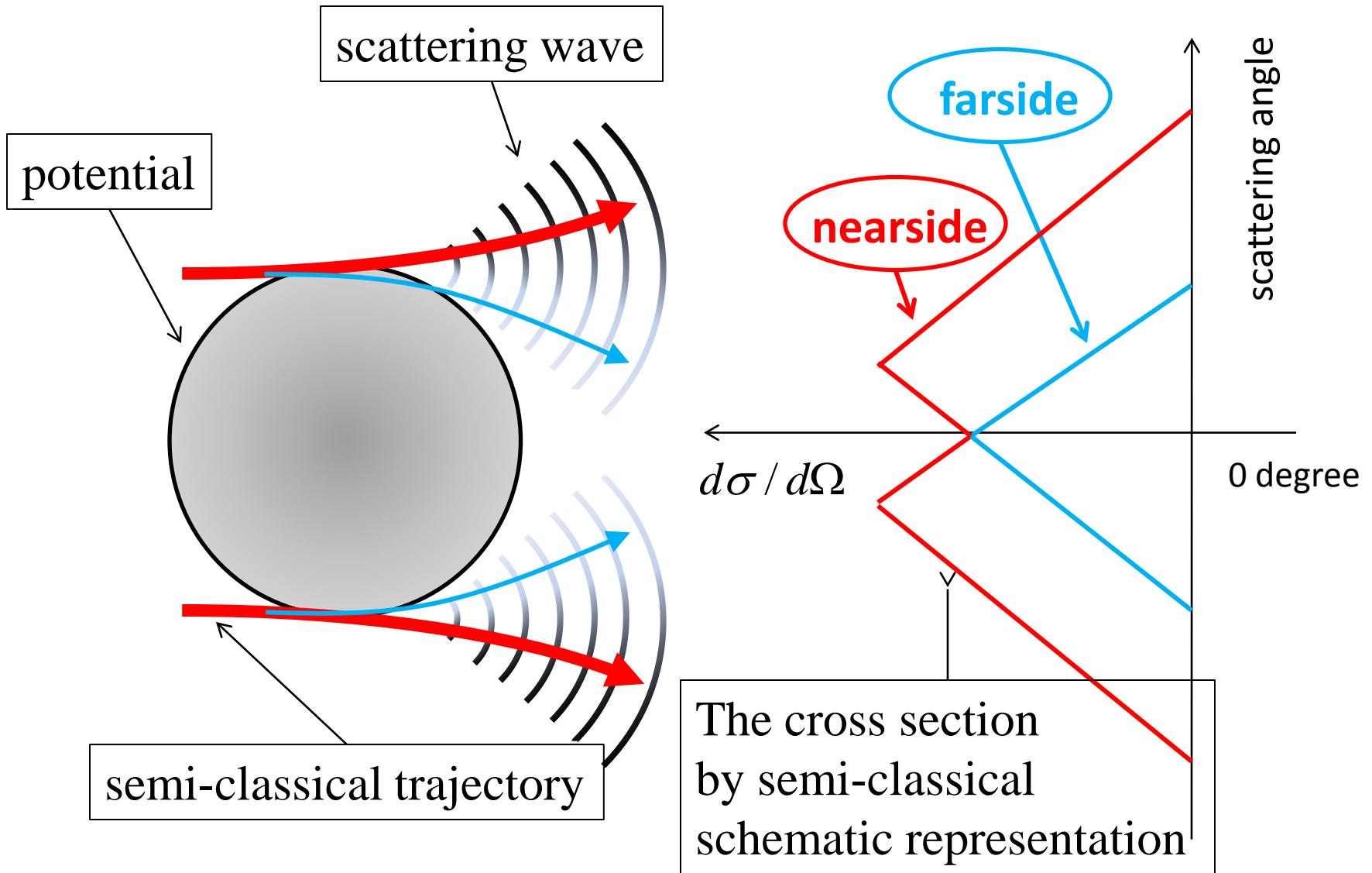
➤ real potential : **repulsive**
around $E/A = 300 \sim 400$ MeV



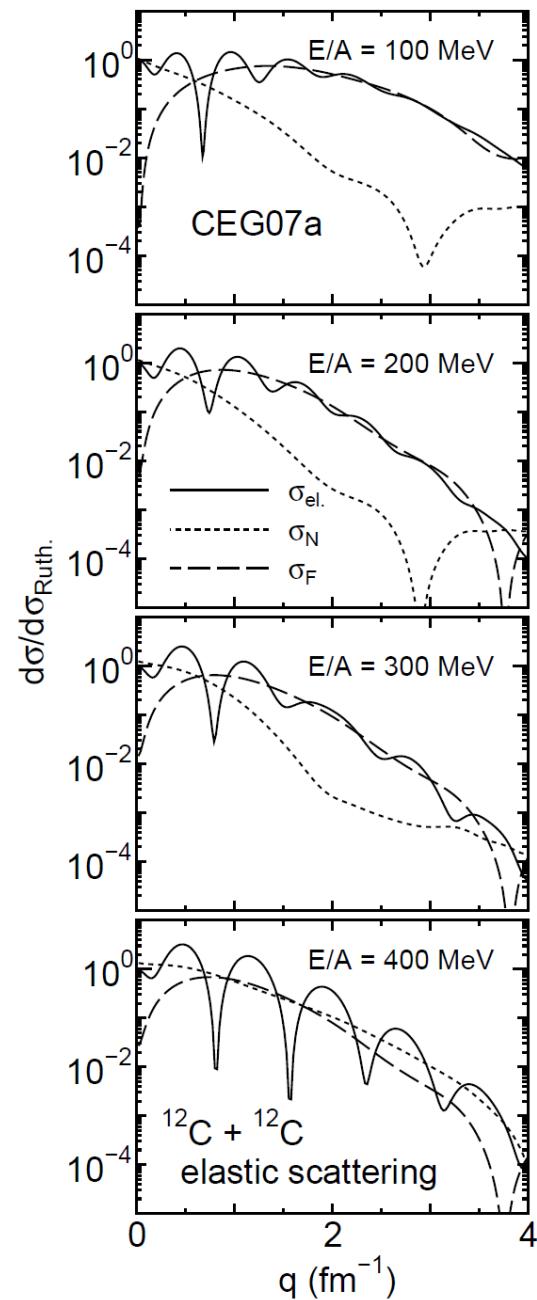
(a) Attractive potential ($V < 0$)



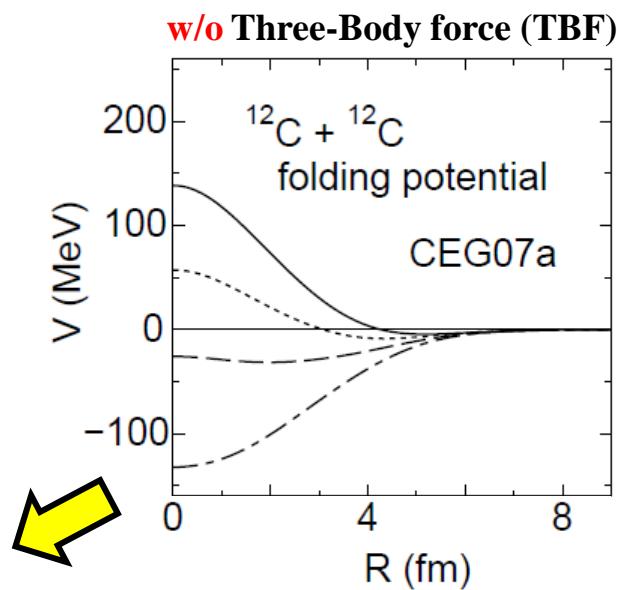
(b) Repulsive potential ($V > 0$)



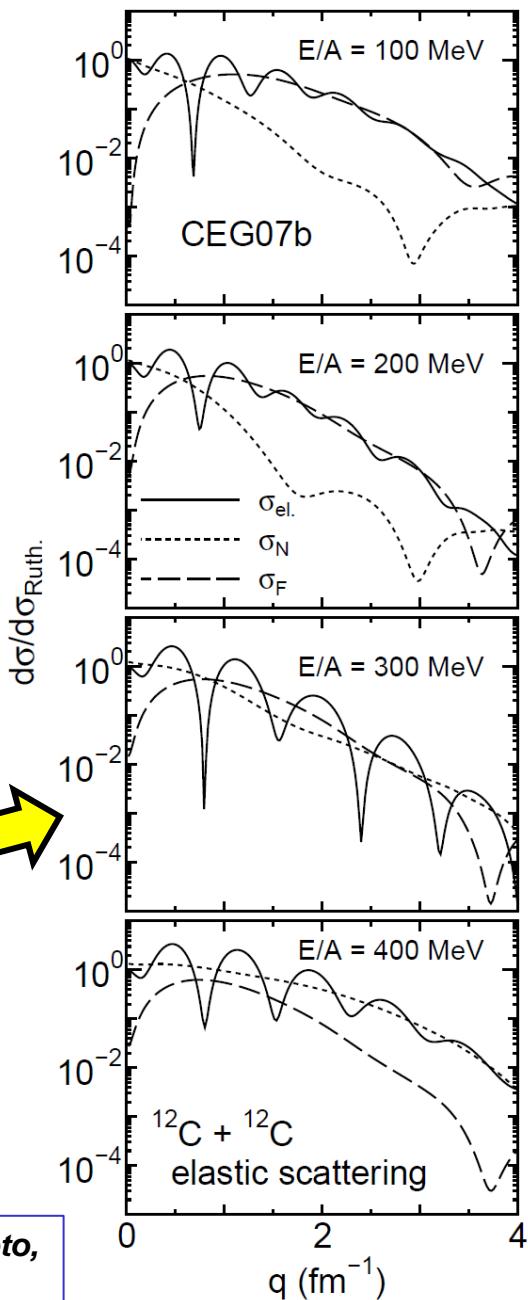
w/o Three-Body force (TBF)



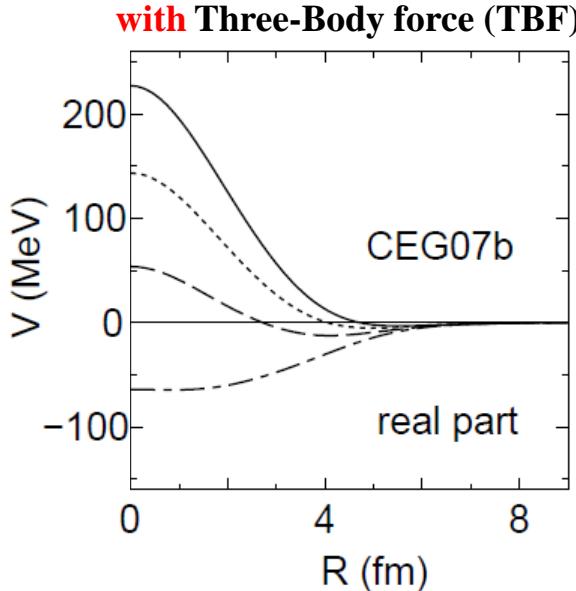
w/o Three-Body force (TBF)



with Three-Body force (TBF)

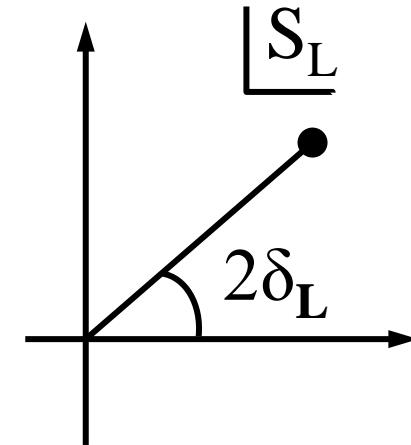
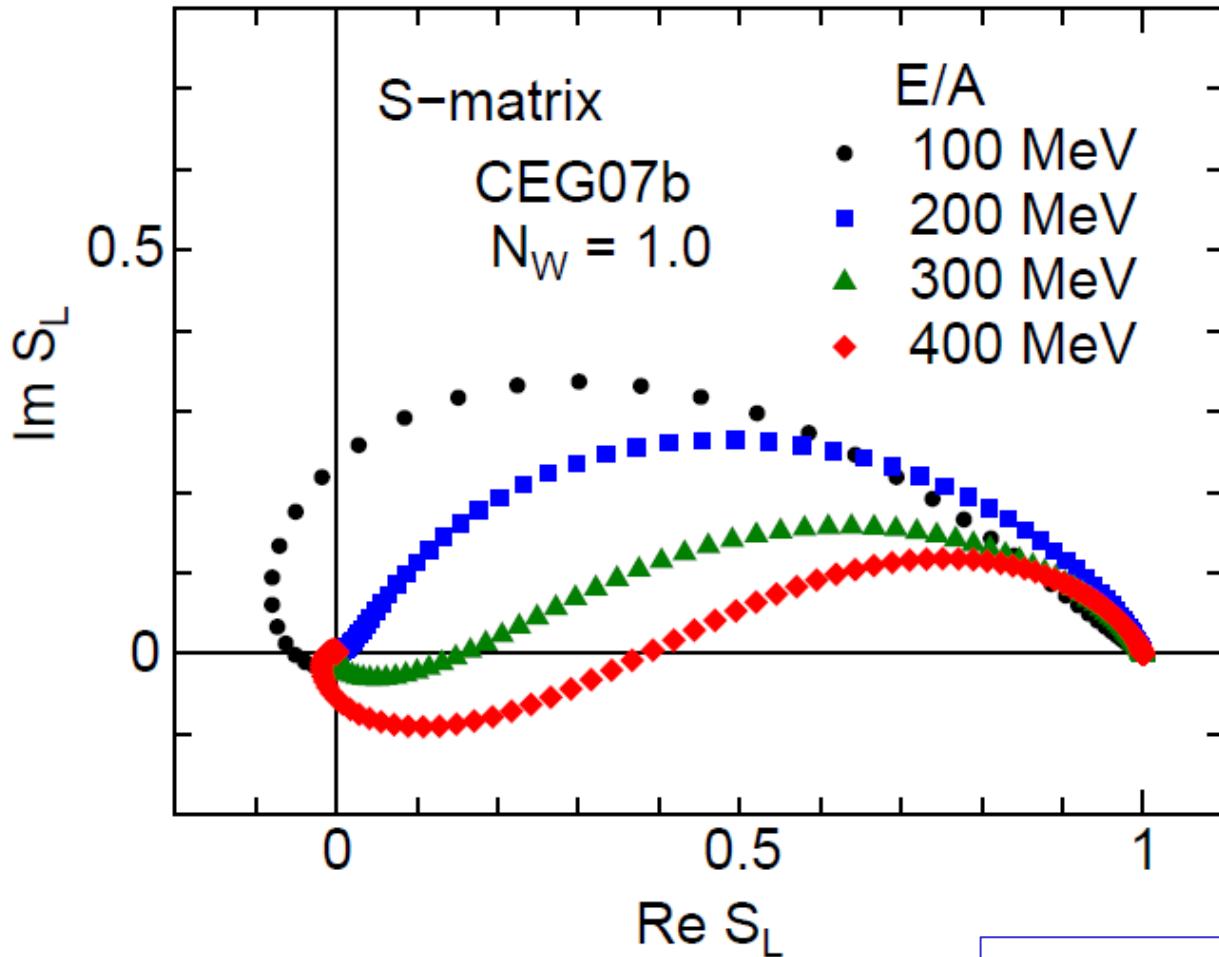


with Three-Body force (TBF)



S-matrix elements of the $^{12}\text{C} + ^{12}\text{C}$ elastic scattering

at $E/A = 100 \sim 400$ MeV with CEG07b (with TBF effects)



$$S_L = \exp(2i\delta_L)$$

$\delta_L < 0$: repulsive
 $\delta_L > 0$: attractive

Summary & Conclusion of part 1

complex G-matrix folding model with a new G-matrix
CEG07 predicts that

- ✓ attractive-to-repulsive transition occurs also in
heavy-ion optical potentials *around $E/A = 300 \sim 400 \text{ MeV}$*
→ *but, no experimental evidence* → **BIG CHALLENGE!**
- ✓ can be observed by *measuring the energy-evolution of elastic scattering angular distribution* in the energy range of $E/A = 200 \sim 400 \text{ MeV}$.
- ✓ new ingredients we have learnt are the important roles of
 - ① *repulsive three-body force (TBF) in nuclear medium*
 - ② *tensor force effects*

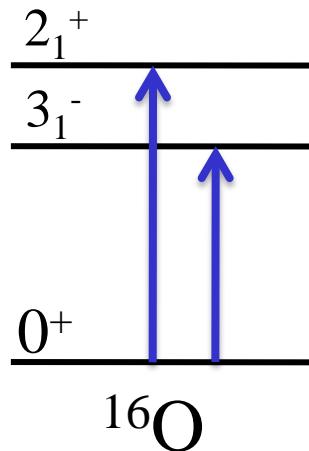
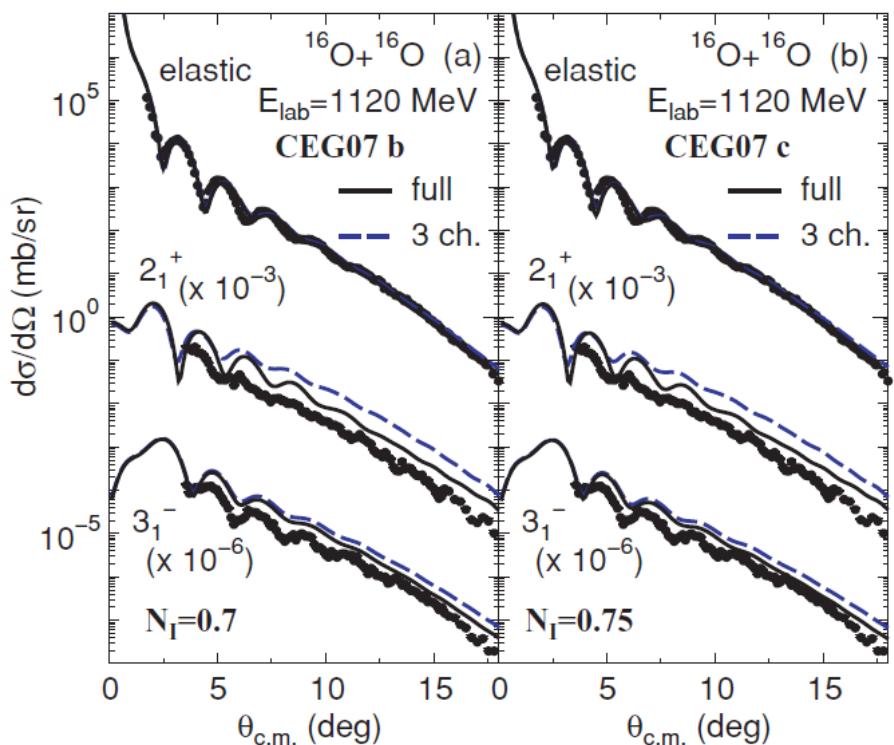
Part 2:

Applications of microscopic FMP to

- 1. reaction calculations (CC, CDCC etc.)*
- 2. scattering of unstable nuclei*
- 3. global parameterization*

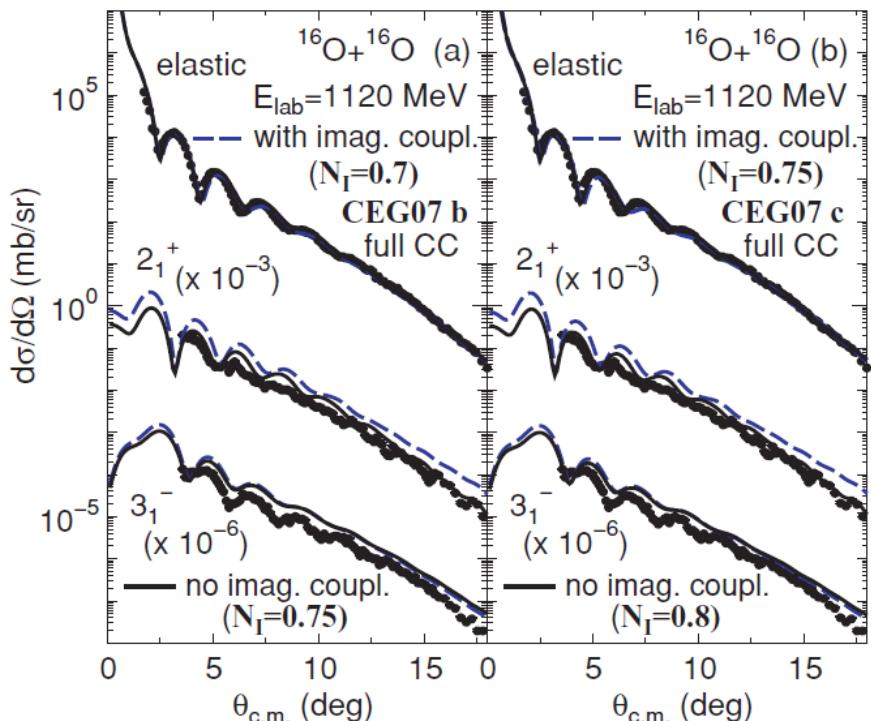
**$^{16}\text{O} + ^{16}\text{O}$ inelastic scattering studied
by a complex *G*-matrix interaction
(@ $E/A=70\text{ MeV}$)**

by M.Takashina, T. Furumoto, Y.Sakuragi
PRC 81, 047605 (2010)



CC cal. with complex-*G* (CEG07)

$$\mathbf{U}_{ij}^{DFM}(\mathbf{R}) = \mathbf{V}_{ij}^{DFM}(\mathbf{R}) + iN_W \mathbf{W}_{ij}^{DFM}(\mathbf{R})$$



$^{9,11}\text{Li} + ^{12}\text{C}$ “quasi-elastic” scattering

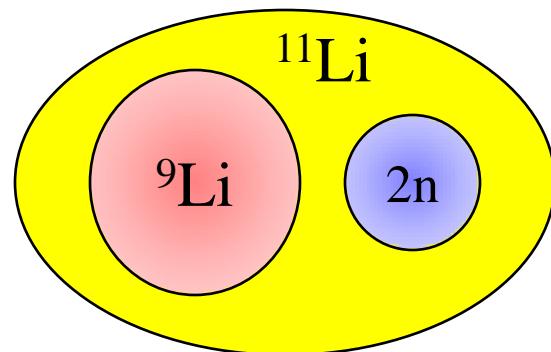
◆ ^9Li density : proton, neutron \Rightarrow single Gaussian form

$$\begin{cases} R_{\text{r.m.s}}^p = 2.18 \text{ (fm)} \\ R_{\text{r.m.s}}^n = 2.39 \text{ (fm)} \end{cases}$$

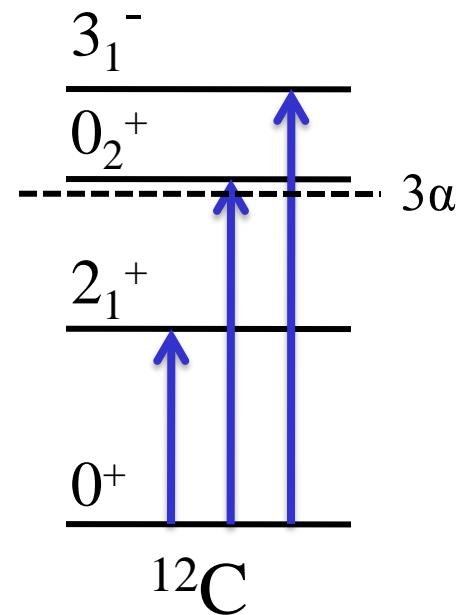
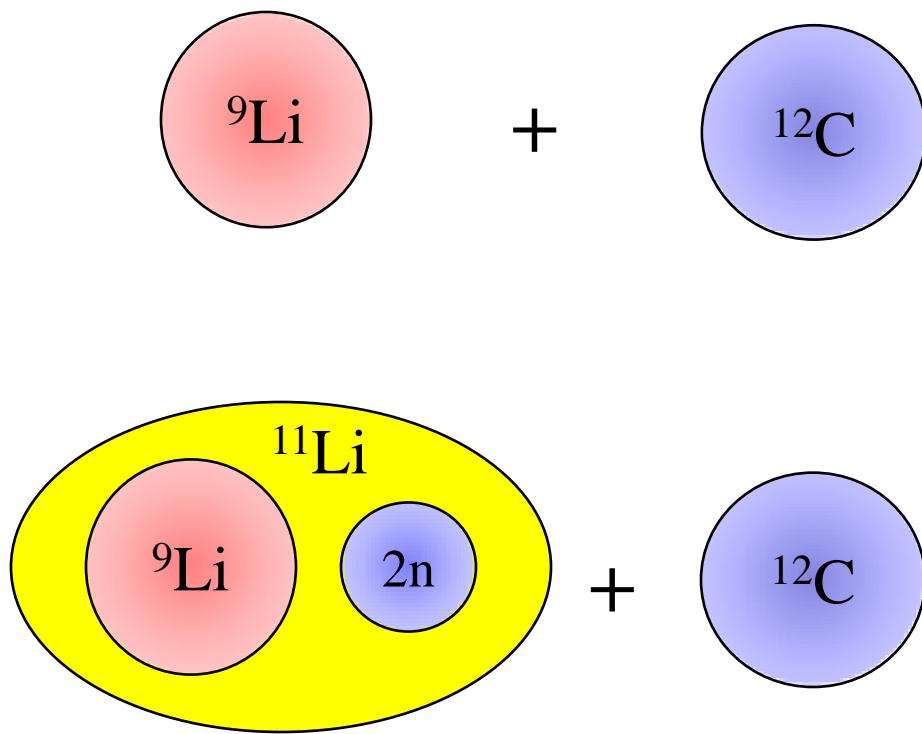
◆ ^{11}Li density : $^9\text{Li} + \text{di-neutron}$ model

$$\rho^{(11\text{Li})}(r) = \left\langle \psi_0(R) \mid \rho^{(9\text{Li})}(\vec{r} - \frac{2}{11} \vec{R}) + \rho^{(2\text{n})}(\vec{r} + \frac{9}{11} \vec{R}) \mid \psi_0(R) \right\rangle$$

$$R_{\text{r.m.s}} = 3.16 \text{ (fm)}$$



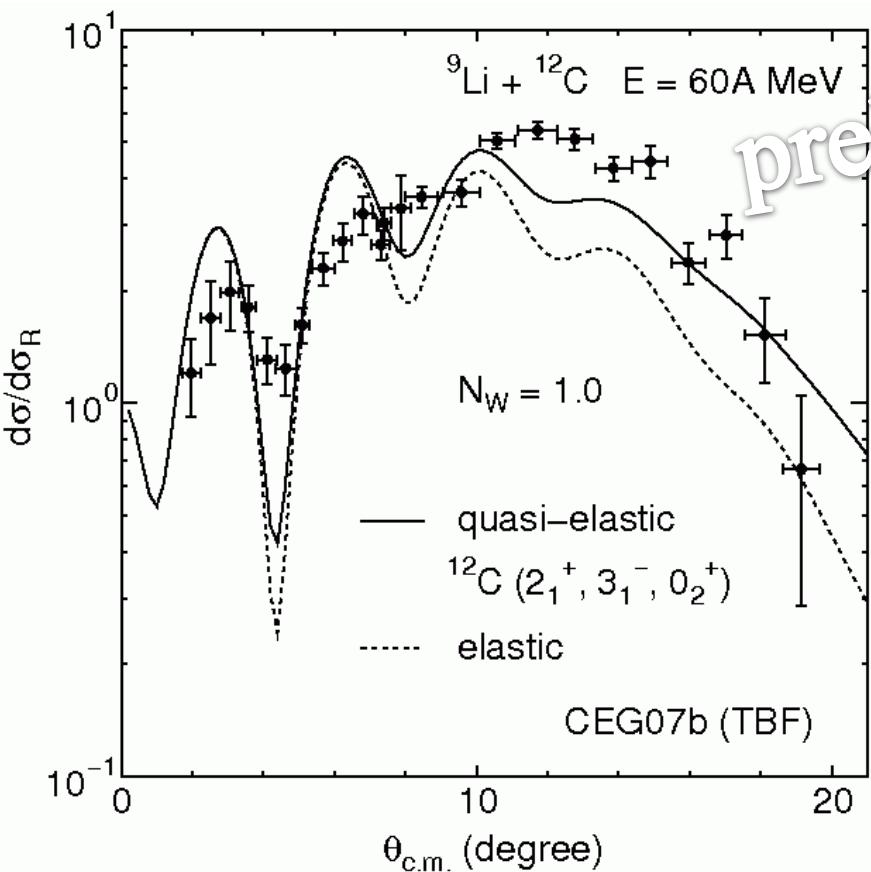
$^{9,11}\text{Li} + ^{12}\text{C}$ “quasi-elastic” scattering



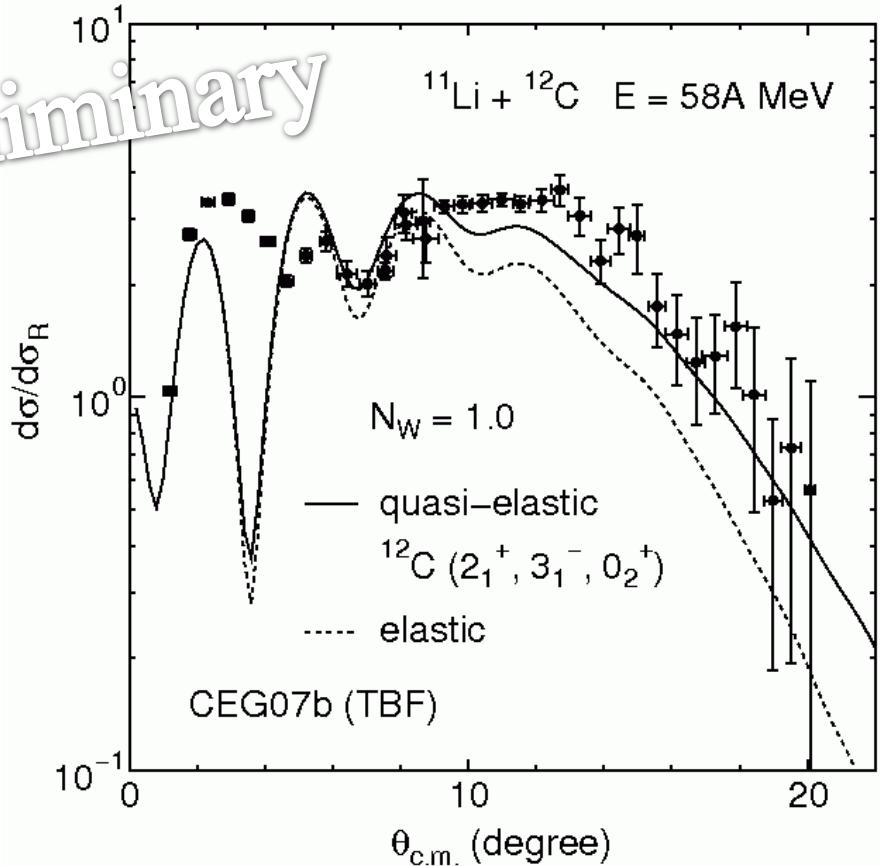
$$\mathbf{U}_{ij}^{DFM}(\mathbf{R}) = \mathbf{V}_{ij}^{DFM}(\mathbf{R}) + i N_w \mathbf{W}_{ij}^{DFM}(\mathbf{R})$$

$^{9,11}\text{Li} + ^{12}\text{C}$ “quasi-elastic” scattering E/A \sim 60 MeV

**coupled-channel (CC)
calculation with
complex-G (CEG07)**



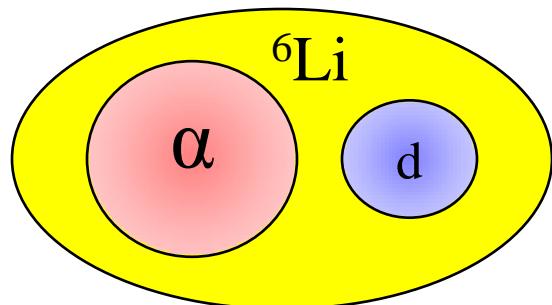
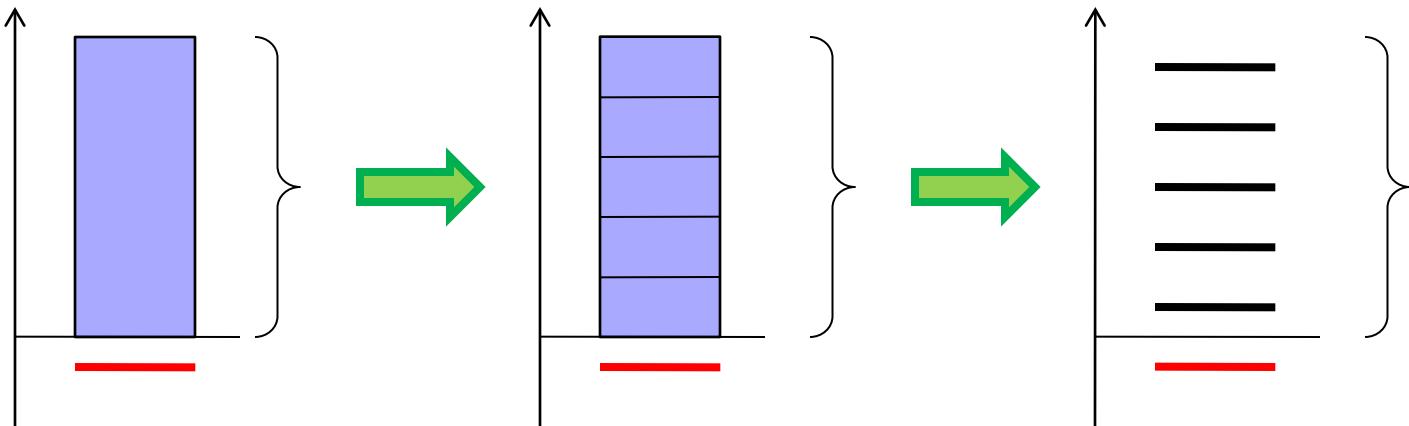
$$U_{DFM}(\mathbf{R}) = V_{DFM}(\mathbf{R}) + iN_W W_{DFM}(\mathbf{R})$$



Exp. data : J. J. Kolata et al., (Phys. Rev. Lett. 69 (1993) 2631)

${}^6\text{Li}$ elastic scattering with ${}^6\text{Li} \rightarrow \alpha + d$ break-up

⇒ Continuum-Discretized Coupled-Channels (CDCC) method



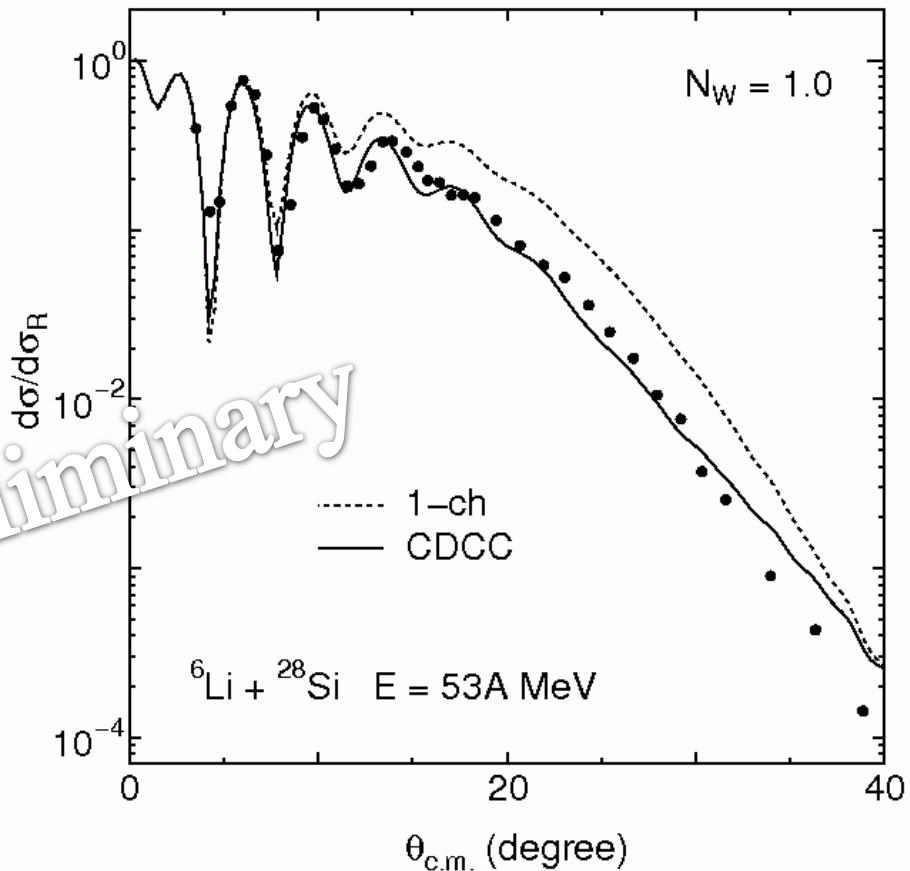
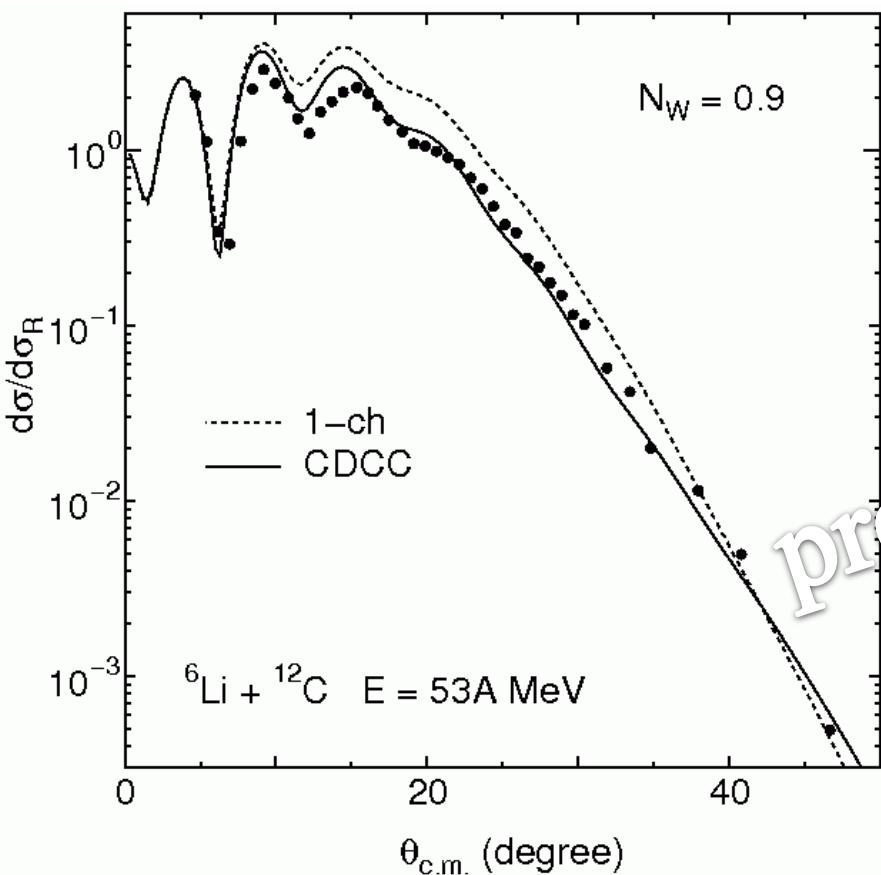
$$U_{ij}^{DFM}(\mathbf{R}) = V_{ij}^{DFM}(\mathbf{R}) + iN_w W_{ij}^{DFM}(\mathbf{R})$$

*Y. Sakuragi, M. Ito, Y. Hirabayashi, C. Samanta
(Prog. Theor. Phys. 98 (1997) 521)*

elastic scattering of ${}^6\text{Li}$ by ${}^{12}\text{C}$, ${}^{28}\text{Si}$ at $E/A = 53$ MeV

➤ **CDCC cal. with complex-G (CEG07) folding model**

$$U_{ij}^{DFM}(\mathbf{R}) = V_{ij}^{DFM}(\mathbf{R}) + iN_W W_{ij}^{DFM}(\mathbf{R})$$



Exp. data : A. Nadasen et al., (Phys. Rev. C. 47 (1993) 674)

Global potential for projectiles of unstable nuclei up to driplines

T. Furumoto, W. Horiuchi, M. Takashina, Y. Yamamoto, Y. Sakuragi (in preparation)

| With CEG07b | 34_Ca | 36_Ca | 38_Ca | 40_Ca | 42_Ca | 44_Ca | 46_Ca | 48_Ca | 50_Ca | 52_Ca | 54_Ca | 56_Ca | 58_Ca | 60_Ca | 62_Ca | 64_Ca | 66_Ca | 68_Ca | 70_Ca |
|-------------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|-------|
| | 30_Ar | 32_Ar | 34_Ar | 36_Ar | 38_Ar | 40_Ar | 42_Ar | 44_Ar | 46_Ar | 48_Ar | 50_Ar | 52_Ar | 54_Ar | 56_Ar | 58_Ar | 60_Ar | 62_Ar | | |
| | 26_S | 28_S | 30_S | 32_S | 34_S | 36_S | 38_S | 40_S | 42_S | 44_S | 46_S | 48_S | 50_S | 52_S | | | | | |
| | 22_Si | 24_Si | 26_Si | 28_Si | 30_Si | 32_Si | 34_Si | 36_Si | 38_Si | 40_Si | 42_Si | 44_Si | 46_Si | 48_Si | | | | | |
| | 20_Mg | 22_Mg | 24_Mg | 26_Mg | 28_Mg | 30_Mg | 32_Mg | 34_Mg | 36_Mg | 38_Mg | 40_Mg | | | | | | | | |
| | 16_Ne | 18_Ne | 20_Ne | 22_Ne | 24_Ne | 26_Ne | 28_Ne | 30_Ne | 32_Ne | 34_Ne | 36_Ne | 38_Ne | | | | | | | |
| | 12_O | 14_O | 16_O | 18_O | 20_O | 22_O | 24_O | | | | | | | | | | | | |
| | 8_C | 10_O | 12_C | 14_C | 16_C | 18_C | 20_C | 22_C | | | | | | | | | | | |

Global parameterization of the CEG07 folding-model potentials

- ✓ projectiles : Z = 6 (C isotope) ~ 20 (Ca isotope) (even-even)
- ✓ targets : ^{12}C ~ ^{208}Pb (closed or sub-closed shell nuclei)
- ✓ energy range : E/A = 30 ~ 400 MeV

◆ Folding-model potential with CEG07a, CEG07b

$$\begin{aligned}
 U_{\text{D}}(R) &= \int \rho_1(\mathbf{r}_1)\rho_2(\mathbf{r}_2)v_{\text{D}}(s; \rho, E/A)d\mathbf{r}_1d\mathbf{r}_2 \\
 &= \int \{\rho_1^{(\text{p})}(\mathbf{r}_1)\rho_2^{(\text{p})}(\mathbf{r}_2)v_{\text{D}}^{(\text{pp})}(s; \rho, E/A) + \rho_1^{(\text{p})}(\mathbf{r}_1)\rho_2^{(\text{n})}(\mathbf{r}_2)v_{\text{D}}^{(\text{pn})}(s; \rho, E/A) \\
 &\quad + \rho_1^{(\text{n})}(\mathbf{r}_1)\rho_2^{(\text{p})}(\mathbf{r}_2)v_{\text{D}}^{(\text{np})}(s; \rho, E/A) + \rho_1^{(\text{n})}(\mathbf{r}_1)\rho_2^{(\text{n})}(\mathbf{r}_2)v_{\text{D}}^{(\text{nn})}(s; \rho, E/A)\}d\mathbf{r}_1d\mathbf{r}_2,
 \end{aligned}$$

$$\begin{aligned}
 U_{\text{EX}}(R) &= \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 + \mathbf{s})\rho_2(\mathbf{r}_2, \mathbf{r}_2 - \mathbf{s})v_{\text{EX}}(s; \rho, E/A) \exp\left[\frac{i\mathbf{k}(R) \cdot \mathbf{s}}{M}\right]d\mathbf{r}_1d\mathbf{r}_2 \\
 &= \int \{\rho_1^{(\text{p})}(\mathbf{r}_1, \mathbf{r}_1 + \mathbf{s})\rho_2^{(\text{p})}(\mathbf{r}_2, \mathbf{r}_2 - \mathbf{s})v_{\text{EX}}^{(\text{pp})}(s; \rho, E/A) + \rho_1^{(\text{p})}(\mathbf{r}_1, \mathbf{r}_1 + \mathbf{s})\rho_2^{(\text{n})}(\mathbf{r}_2, \mathbf{r}_2 - \mathbf{s})v_{\text{EX}}^{(\text{pn})}(s; \rho, E/A) \\
 &\quad + \rho_1^{(\text{n})}(\mathbf{r}_1, \mathbf{r}_1 + \mathbf{s})\rho_2^{(\text{p})}(\mathbf{r}_2, \mathbf{r}_2 - \mathbf{s})v_{\text{EX}}^{(\text{np})}(s; \rho, E/A) + \rho_1^{(\text{n})}(\mathbf{r}_1, \mathbf{r}_1 + \mathbf{s})\rho_2^{(\text{n})}(\mathbf{r}_2, \mathbf{r}_2 - \mathbf{s})v_{\text{EX}}^{(\text{nn})}(s; \rho, E/A)\} \\
 &\quad \times \exp\left[\frac{i\mathbf{k}(R) \cdot \mathbf{s}}{M}\right]d\mathbf{r}_1d\mathbf{r}_2,
 \end{aligned}$$

◆ Globally-parameterized density (“Sao Paolo density”)

L. C. Chamon, B. V. Carlson, L. R. Gasques, D. Pereira, C. D. Conti, M. A. Alvarez, M. S. Hussein, M. A. C. Ribeiro, E. S. Rossi, Jr., et al., Phys. Rev. C **66**, 014610 (2001).

◆ Globally-parameterized density (“Sao Paolo density”)

L.C.Chamon et al., PRC66, 014601 (2001)

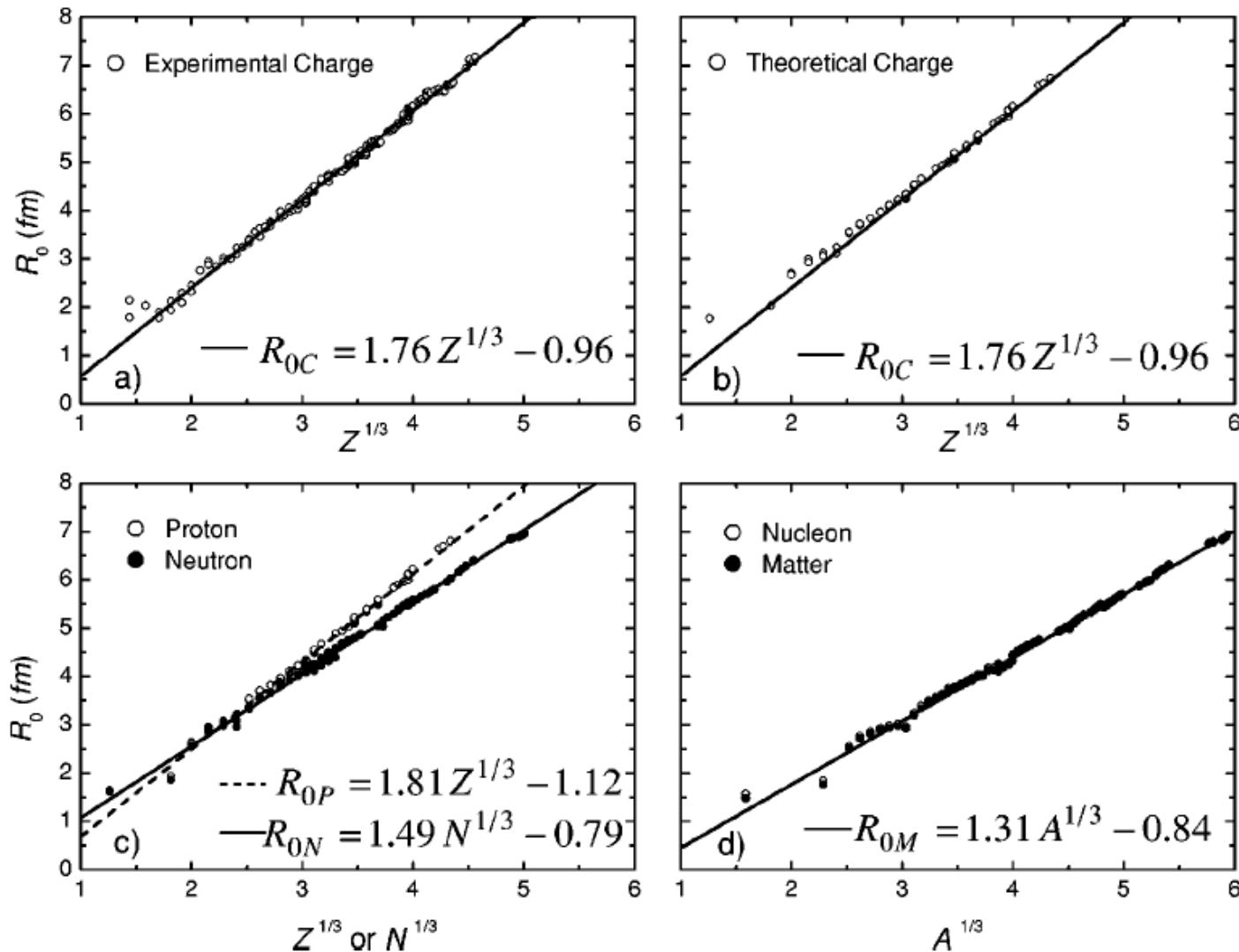


FIG. 3. The R_0 parameter obtained for charge distributions extracted from electron scattering experiments and for theoretical densities obtained from Dirac-Hartree-Bogoliubov calculations.

◆ Globally-parameterized density (“Sao Paolo density”)

L.C.Chamon et al., PRC66, 014601 (2001)

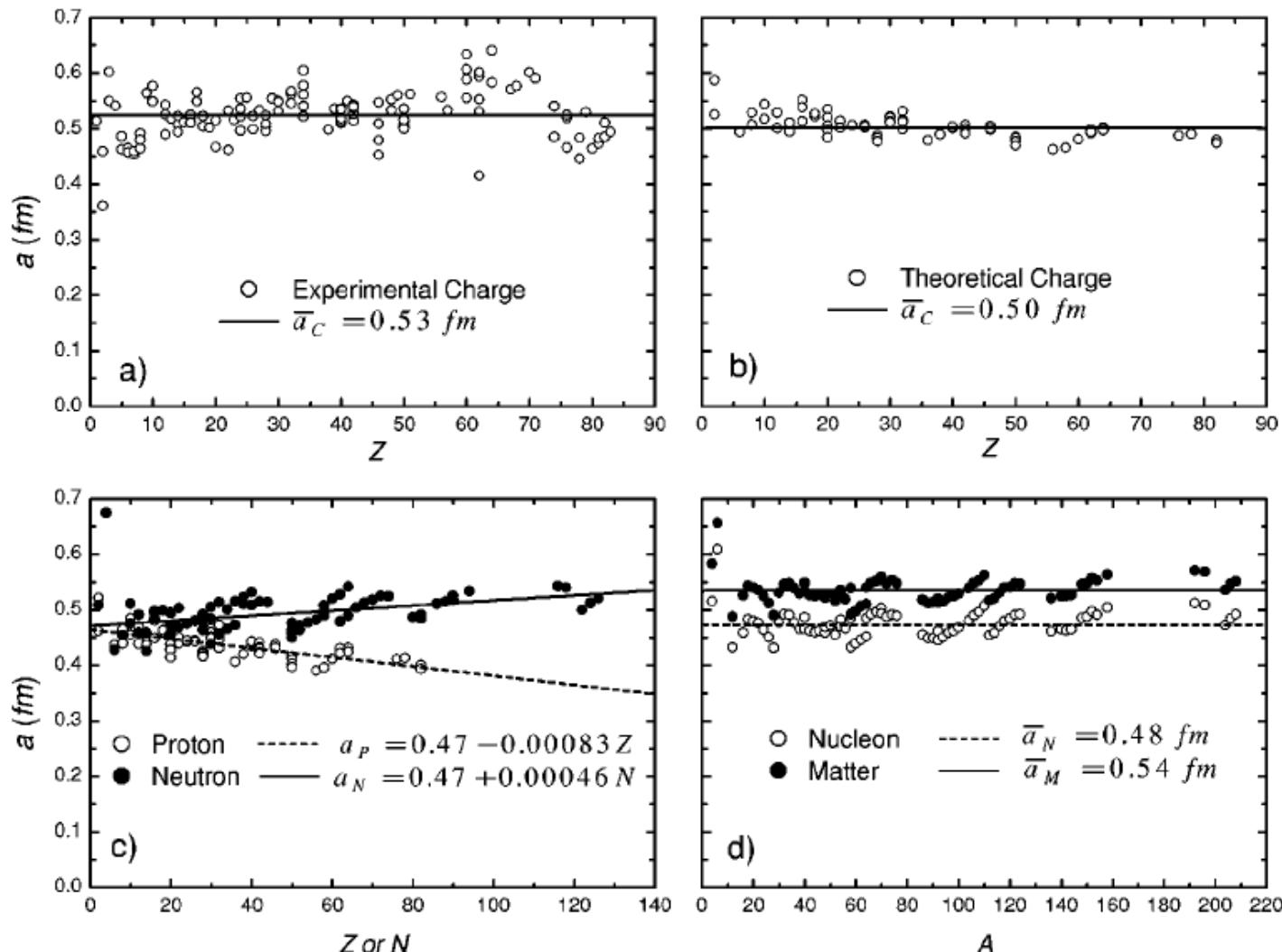


FIG. 2. Equivalent diffuseness values obtained for charge distributions extracted from electron scattering experiments and for theoretical densities obtained from Dirac-Hartree-Bogoliubov calculations.

Global parameterization of the CEG07 folding-model potentials

T. Furumoto, W. Horiuchi, M. Takashina, Y. Yamamoto, Y. Sakuragi (in preparation)

- ✓ projectiles : Z = 6 (C isotope) ~ 20 (Ca isotope) (even-even)
- ✓ targets : ^{12}C ~ ^{208}Pb (closed or sub-closed shell nuclei)
- ✓ energy range : E/A = 30 ~ 400 MeV

$$V_F(R) = \sum_{n=1}^{10} \left\{ \alpha_n \exp\left(-\frac{R^2}{\gamma_n^2}\right) \right\} ,$$

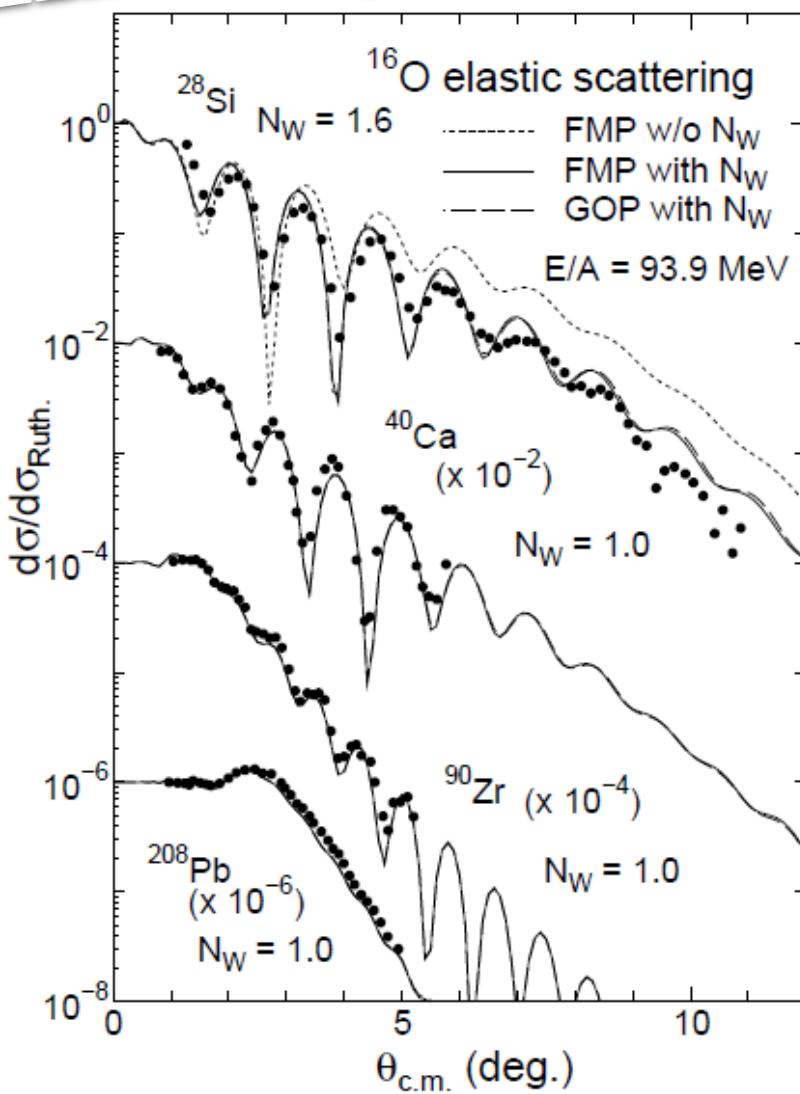
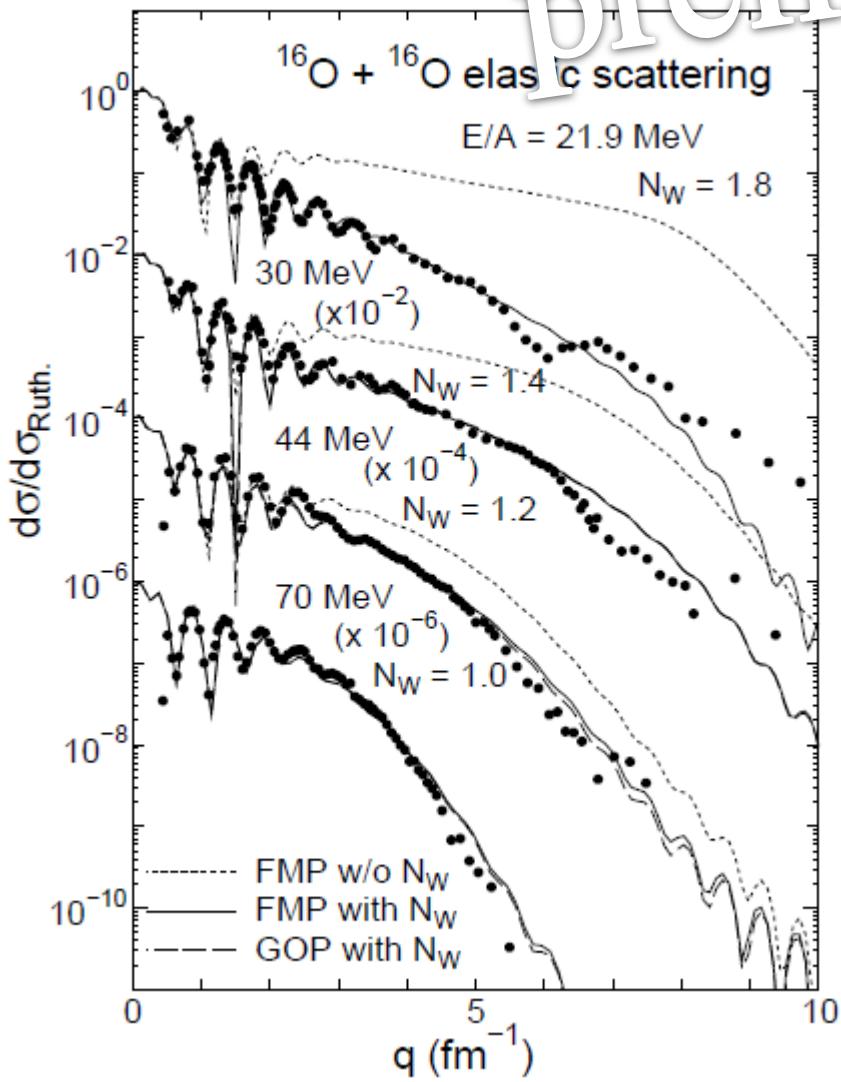
$$W_F(R) = \sum_{n=1}^{10} \left\{ \beta_n \exp\left(-\frac{R^2}{\gamma_n^2}\right) \right\} ,$$

$$\alpha_n = \alpha_n(A_p, Z_p, A_t, E/A) ,$$

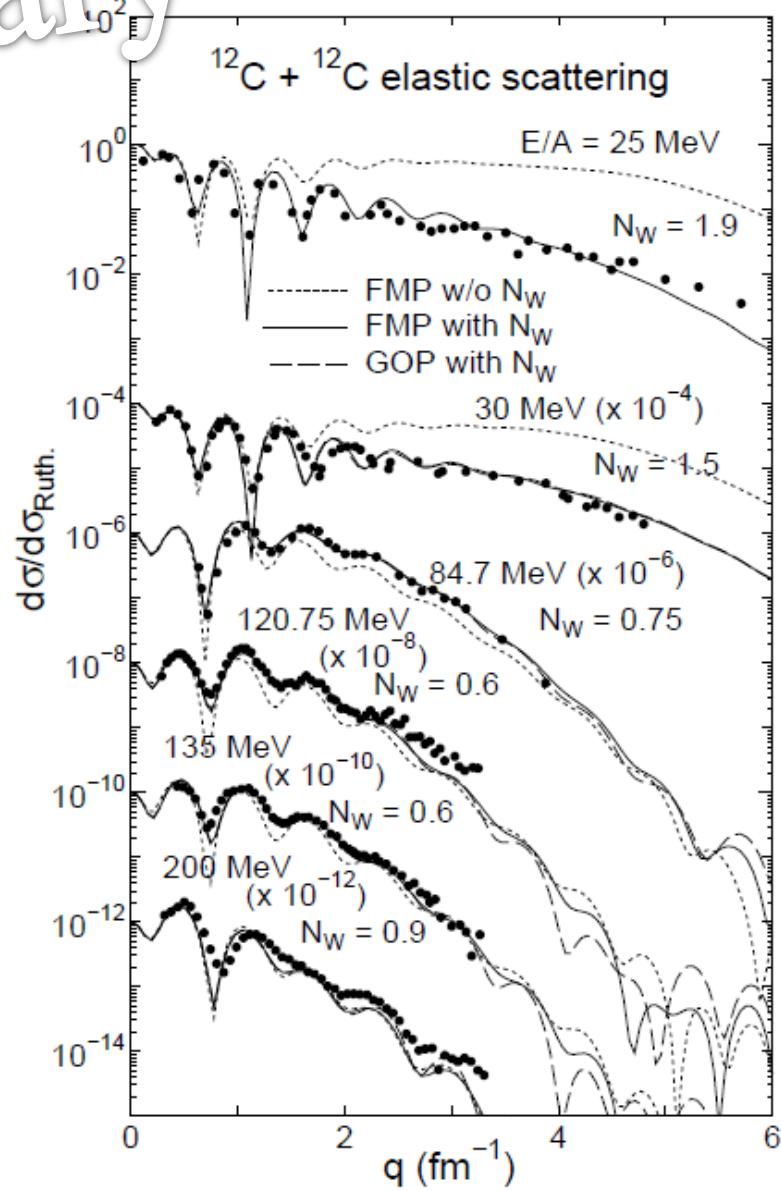
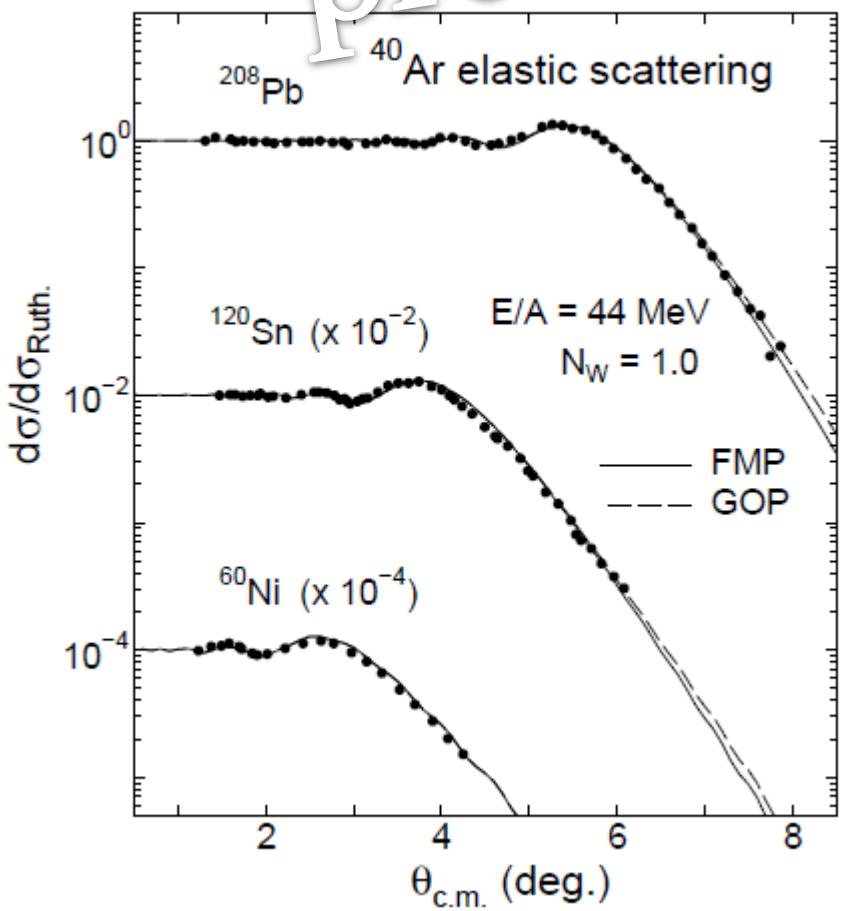
$$\beta_n = \beta_n(A_p, Z_p, A_t, E/A) ,$$

$$\gamma_n = 0.45 \left(\frac{n+8}{18} \right) (A_p^{1/3} + A_t^{1/3} + 1)$$

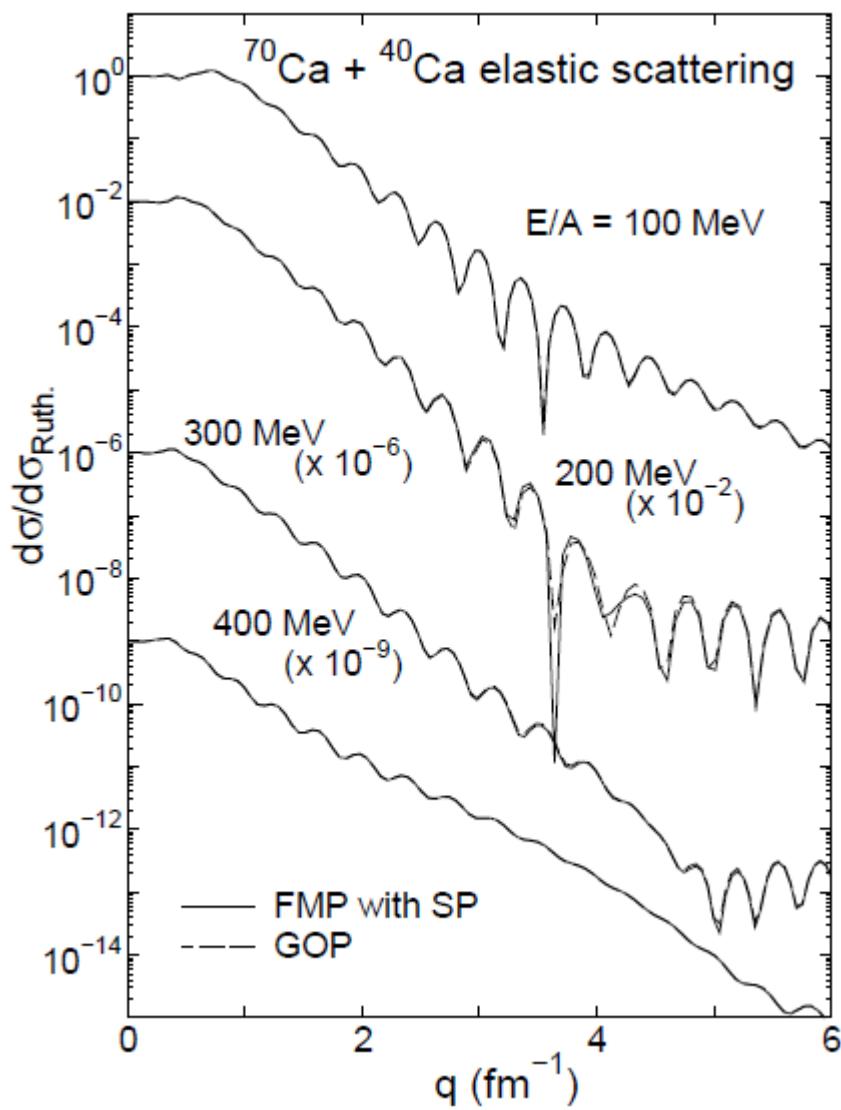
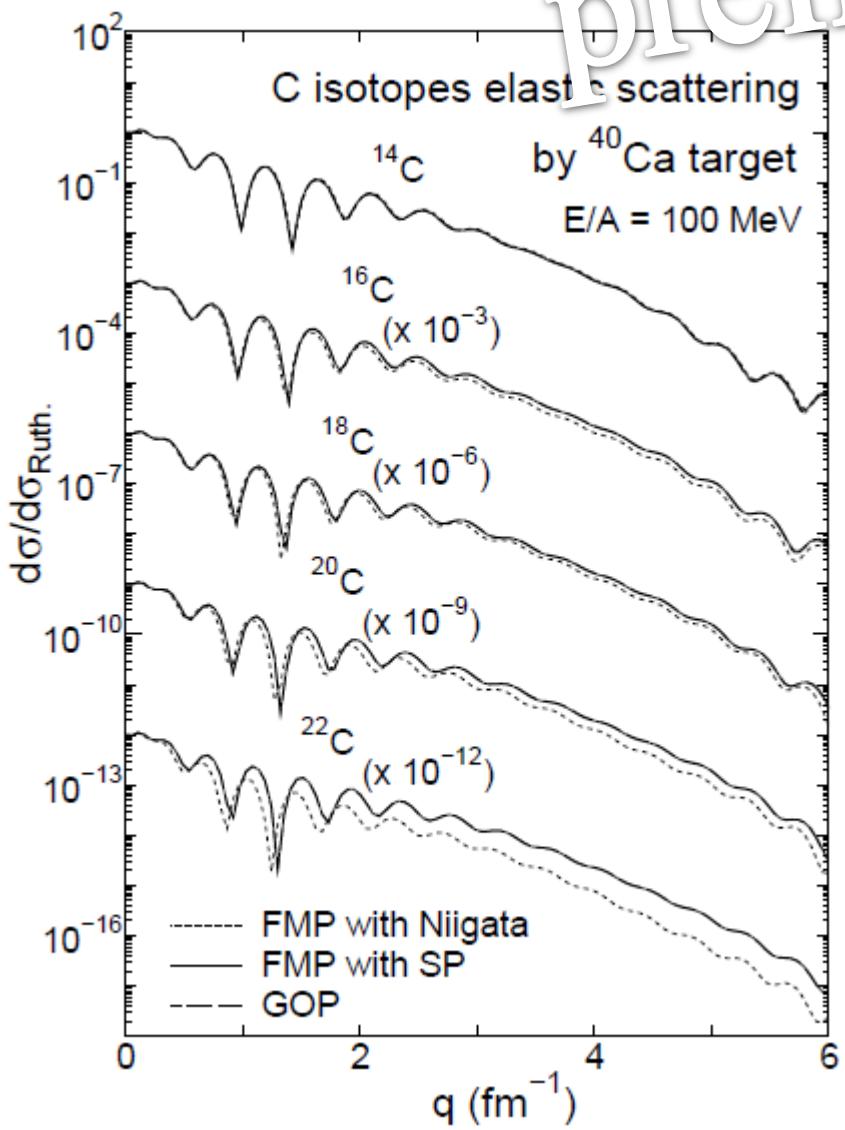
preliminary



preliminary



preliminary



Summary

- We have proposed a **new complex G-matrix** (“**CEG07**”),
 - derived from **ESC04**(extended soft-core) **NN force**
 - include **three-body force (TBF)** effect
 - calculated up to higher density (about twice the normal density)
- We have applied DFM with **new complex G-matrix** (“**CEG07**”) to **nucleus-nucleus (AA)** elastic/inelastic scattering & breakup
- **CEG07** is successful for **nucleus-nucleus** elastic scattering
 - reproduce **cross section** data for ^{12}C , ^{16}O elastic scattering by ^{12}C , ^{16}O , ^{28}Si , ^{40}Ca targets at various energies.
- We have found a decisive role of **Three-body repulsive** force effect
- We also demonstrated possible applications to **nuclear reactions** (inelastic/breakup) including those with **unstable nuclei**
- We constructed **Global potentials** for projectiles of **unstable nuclei** up to **driplines**, based on the microscopic **CEG07 folding potentials**.