

Macroscopic collision dynamics by time-dependent energy density functional theory

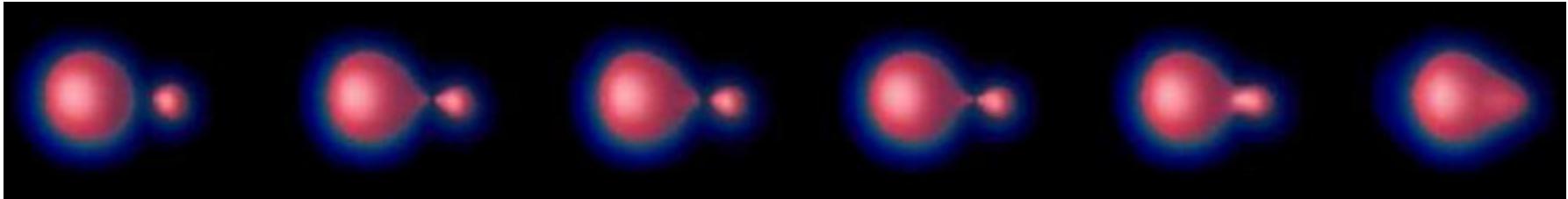
Kouhei Washiyama (Université Libre de Bruxelles)

Denis Lacroix (GANIL), Sakir Ayik (Tennessee Tech. Univ)

Mean-field dynamics

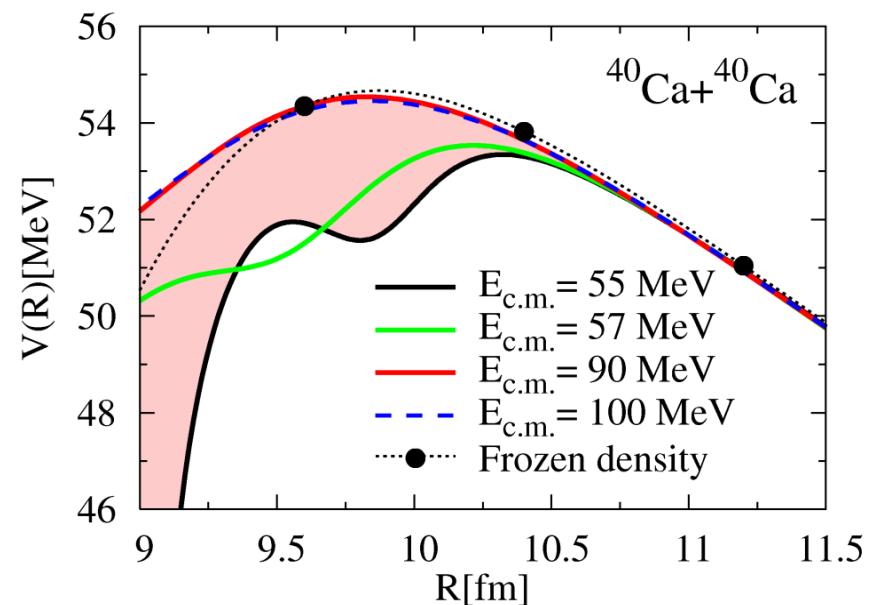
- Microscopic \longleftrightarrow Macroscopic
- Nucleus-nucleus potential
- Energy dissipation

Figure taken from Simenel, Avez, Int.J.Mod.Phys.E17(2008)31



Plan of the talk

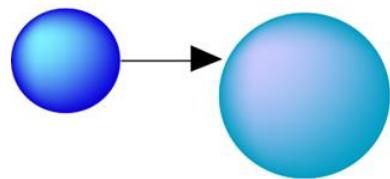
- Introduction
- Macroscopic and microscopic nuclear reaction
- Time-dependent energy density functional (TDHF)
- Nucleus-nucleus potential, energy dissipation
- Fluctuation around mean field: Width of mass distribution
- Stochastic mean-field approach
- Summary



Introduction: macroscopic, microscopic

(Energy < 10 MeV/A)

□ Macroscopic nuclear reaction



- Collective degrees of freedom
- Coupled-channels, CDCC
- Nucleus-nucleus potential

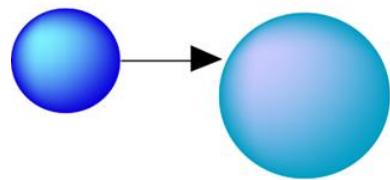
■ Nucleus-nucleus potential

- Phenomenological potential (Woods-Saxon)
- Double folding potential with realistic nuclear density

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□ Microscopic nuclear reaction: TDHF

- Microscopic degrees of freedom
(single-particle wave functions in TDHF)
- No collective degrees of freedom is explicitly considered

Introduction: TDHF

- Nuclear structure and dynamics in a unified framework
- Dynamical effect (vibration, rotation, transfer, ...) automatically included
- Single-particle motion: quantum. Collective motion: classical
- Full Skyrme forces including spin-orbit and time-odd terms
- No fitted parameter for dynamics



Predictive power of TDHF

- Bonche et al., PRC13 (1976)
- Flocard et al., PRC17 (1978)
- Kim et al., J. Phys. G23 (1997)
- Simenel et al., PRL86 (2001)
- Nakatsukasa et al., PRC71 (2005)
- Umar et al., PRC73 (2006)
- Maruhn et al., PRC74 (2006)

3D-TDHF code by P. Bonche (Kim, Otsuka, Bonche, J.Phys.G23(1997)1267)

- $i\hbar\dot{\phi}_\alpha = \hat{h}\phi_\alpha$, h : self-consistent mean field
- Three dimensional box (mesh size: 0.8 fm, time step: 0.45 fm/c)
- Skyrme energy density functional (SLy4d)

Introduction: macroscopic, microscopic

Macroscopic aspects \leftrightarrow Microscopic model



Nucleus-nucleus potential and

Microscopic potentials based on EDF

□Frozen density approximation

Denisov, Norenberg, EPJA15(2002)375

- Energy density functional of projectile and target
- Density is frozen to ground-state density at each R

$$V^{FD}(R) = \mathcal{E}[\rho_{P+T}](R) - \mathcal{E}[\rho_T] - \mathcal{E}[\rho_P]$$

\mathcal{E} : Skyrme energy density functional

□Density-constrained TDHF

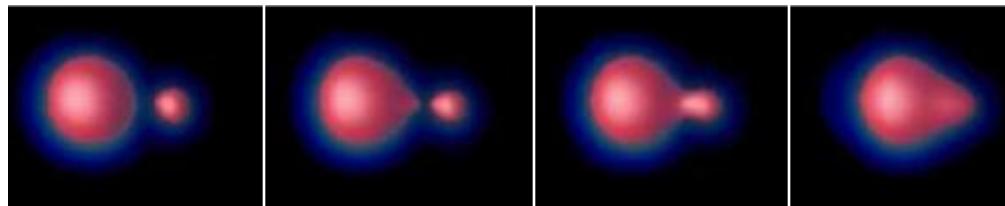
Umar, Oberacker, PRC74(2006)021601

- Density from TDHF trajectory
- Minimization of energy: constraint HF on this density

$$\hat{H} \rightarrow \hat{H} + \int d^3r \lambda(r) \hat{\rho}(r)$$

$$V(R) \rightarrow E_{DC}(R) - E_{A_1} - E_{A_2}$$

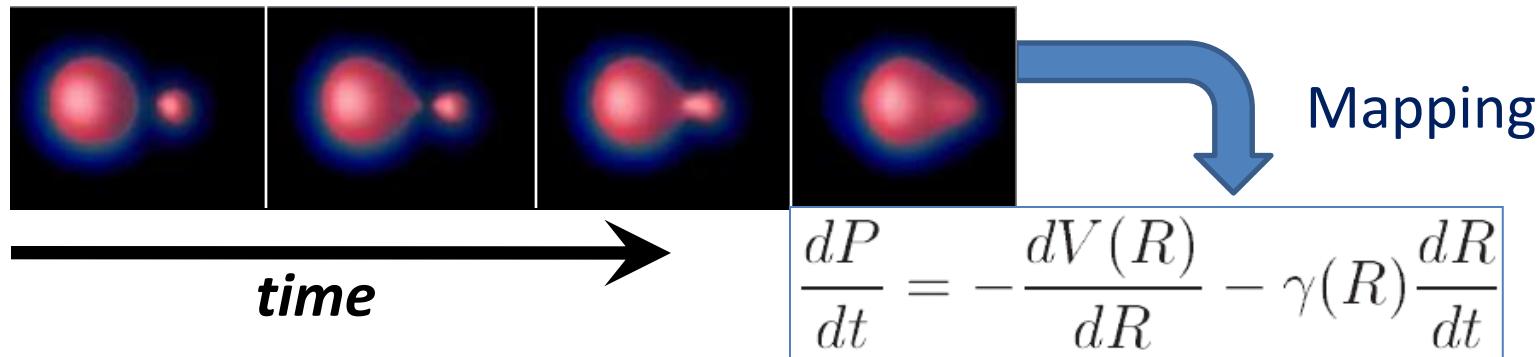
From TDHF to macroscopic dynamics (fusion)



→
time

Central collision

From TDHF to macroscopic dynamics (fusion)

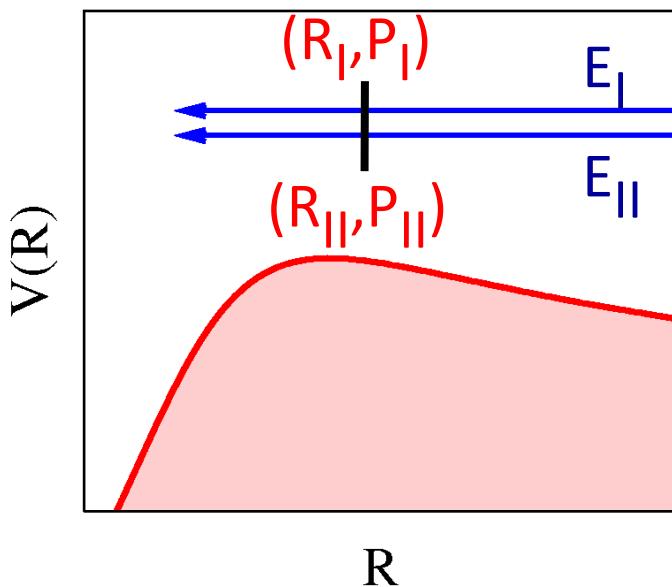


R : relative distance

P : momentum

V : potential

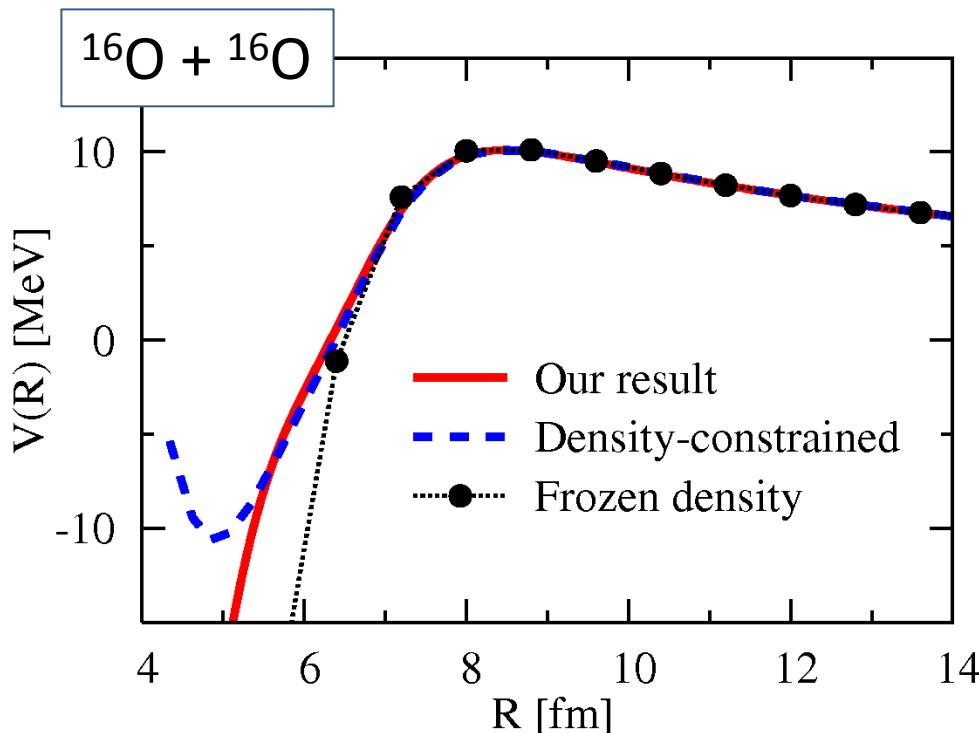
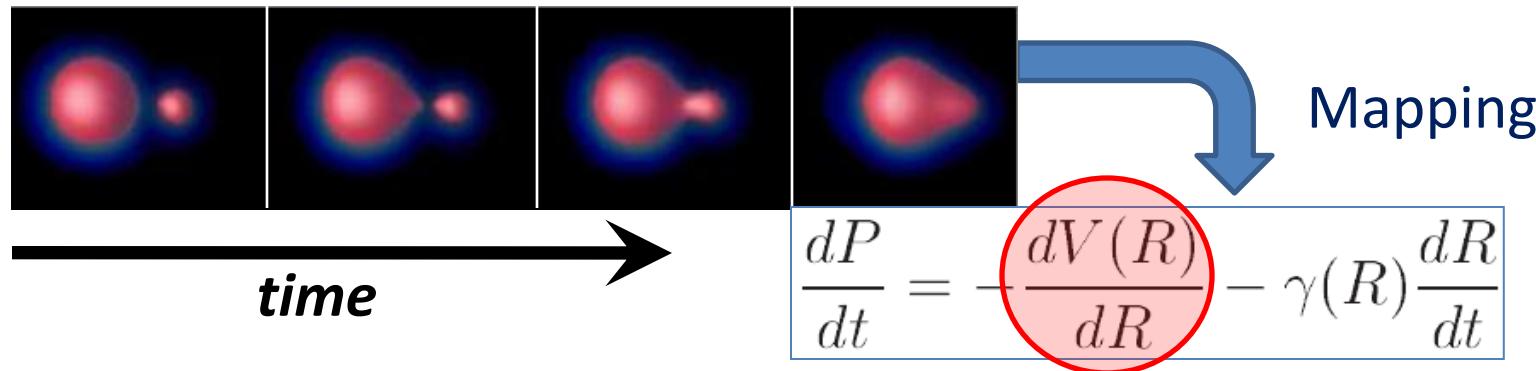
γ : friction coefficient



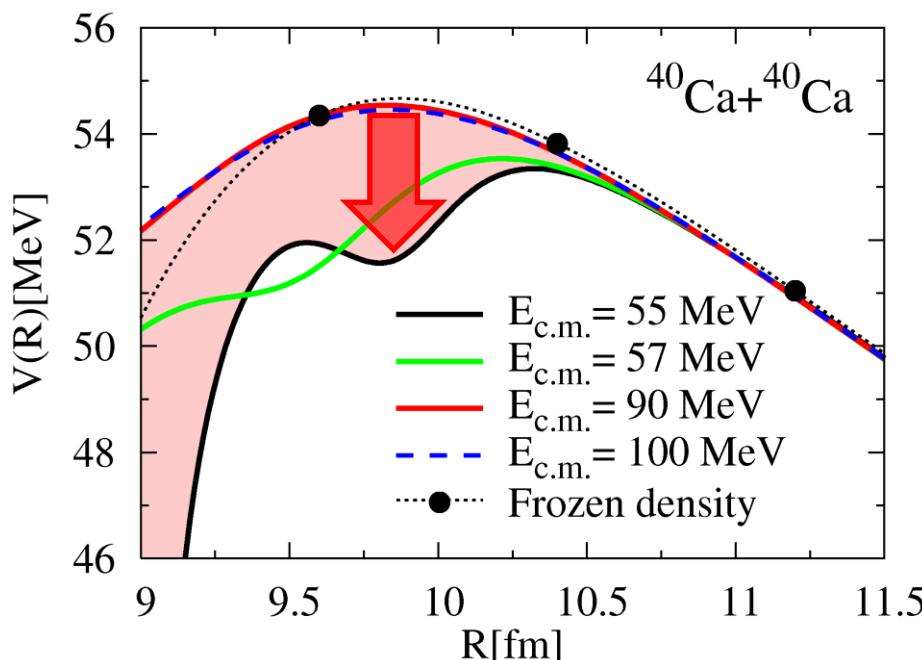
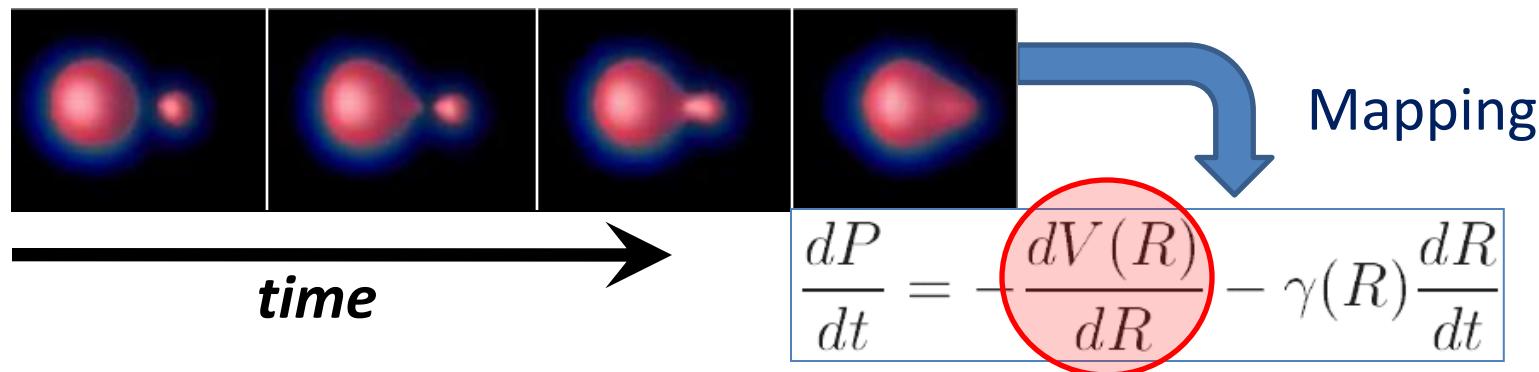
- Two trajectories $(R_I, P_I), (R_{II}, P_{II})$ at two slightly different energies (E_I, E_{II})

$$\rightarrow \gamma(R) = -\frac{dP_I/dt - dP_{II}/dt}{dR_I/dt - dR_{II}/dt}$$

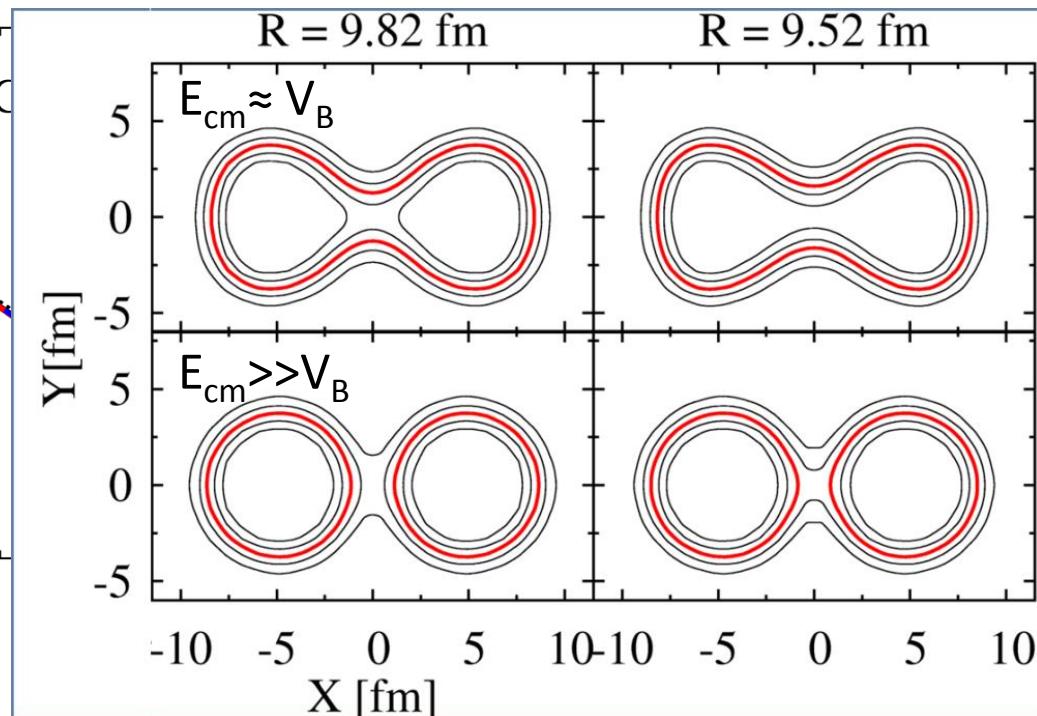
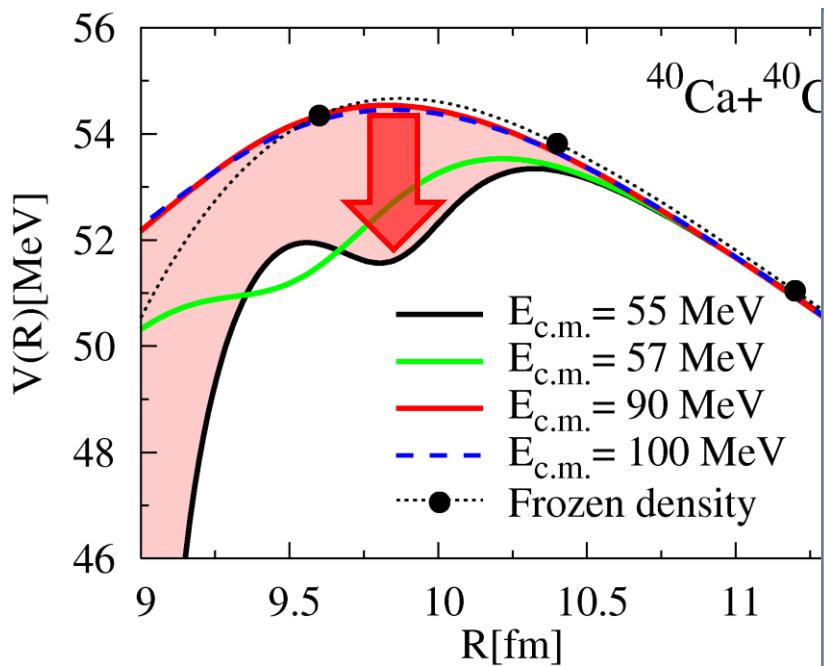
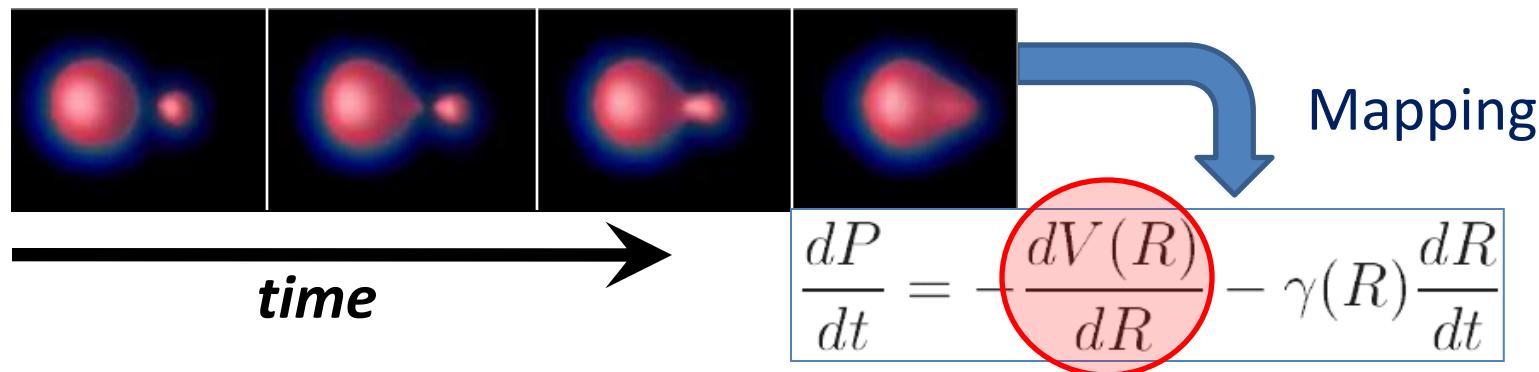
From TDHF to macroscopic dynamics (fusion)



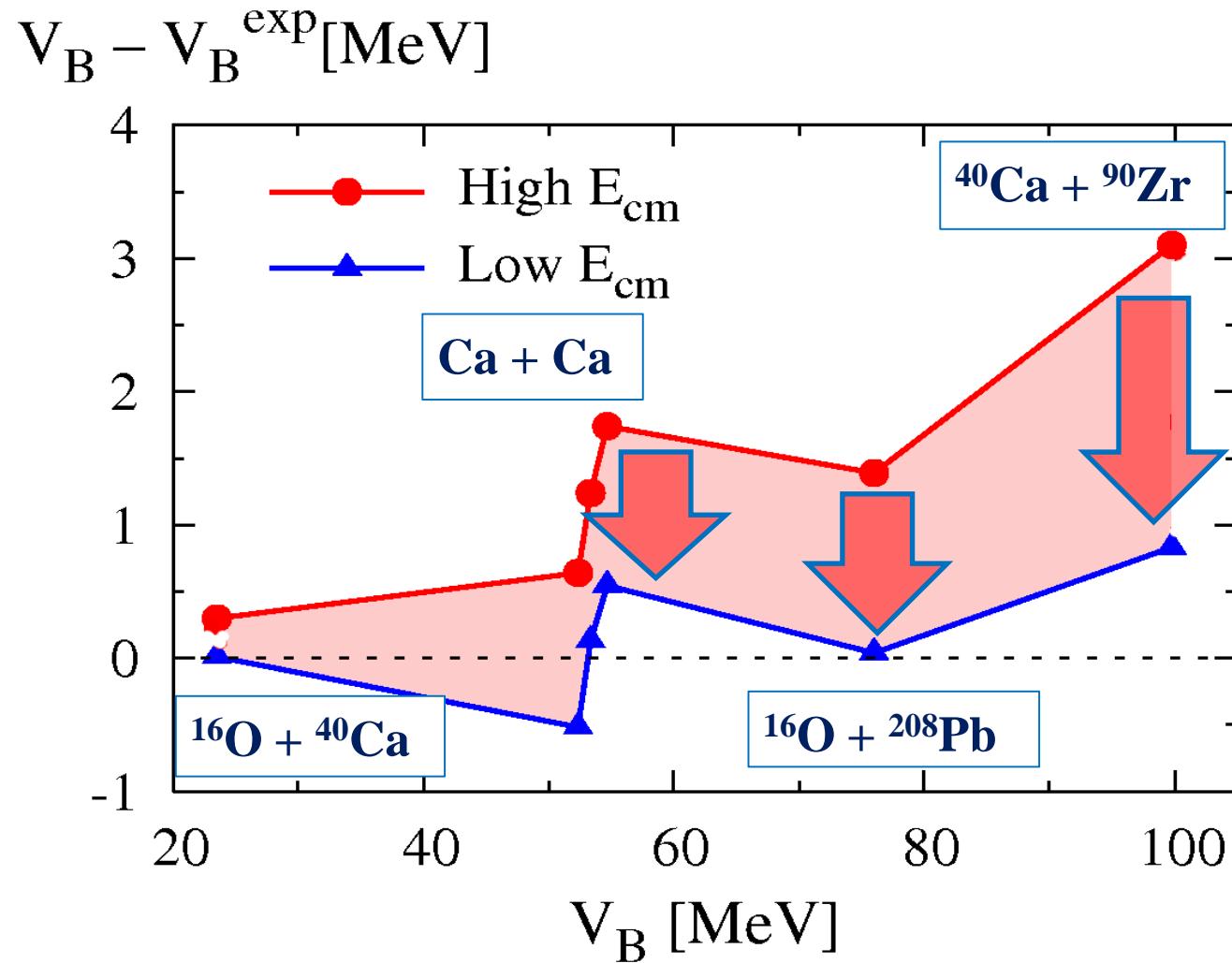
Energy dependence of potential



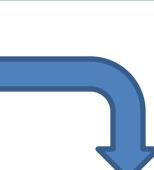
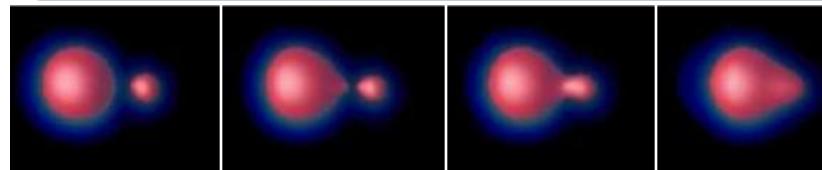
Energy dependence of potential



Energy dependence of potential



Nuclear friction: energy dissipation

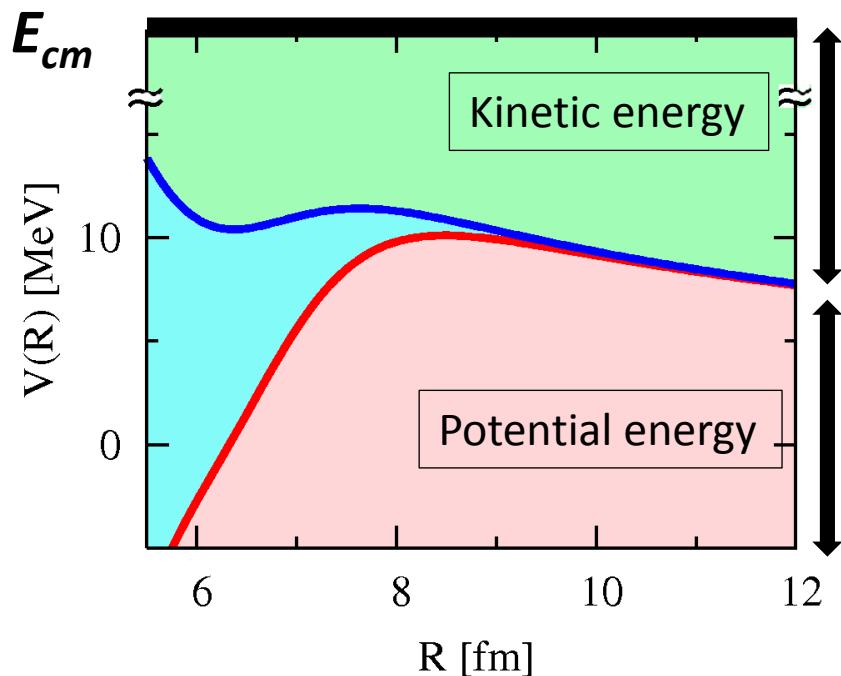


Mapping

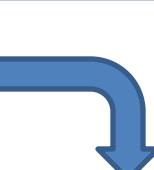
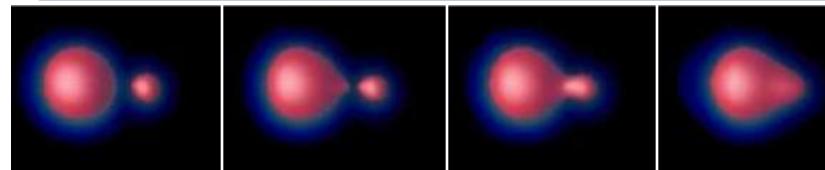
$$\frac{dP}{dt} = -\frac{dV(R)}{dR} - \gamma(R) \frac{dR}{dt}$$

R : relative distance
 P : momentum
 V : potential
 γ : friction coefficient

$$E_{cm} = P^2/2\mu + V$$



Nuclear friction: energy dissipation

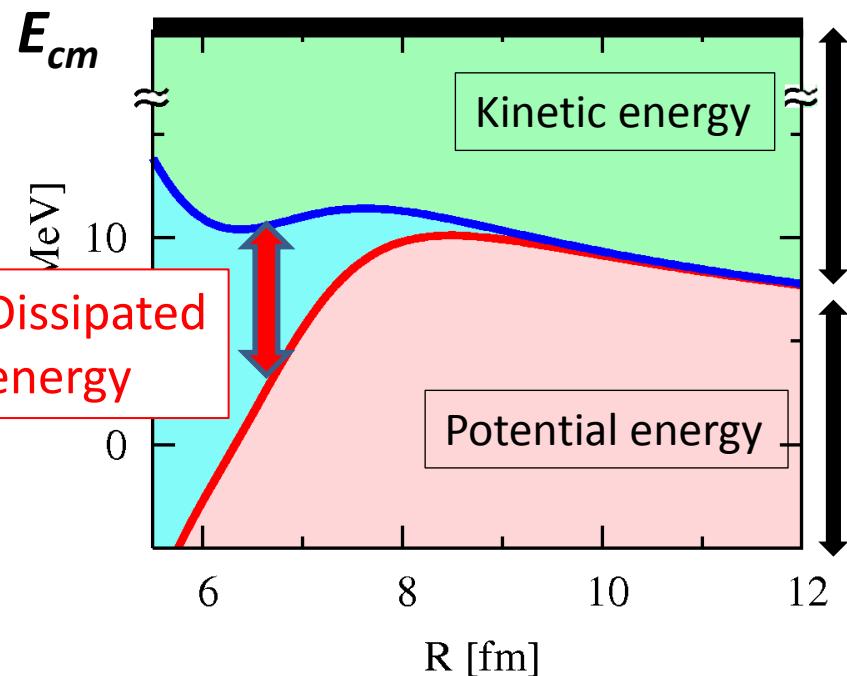


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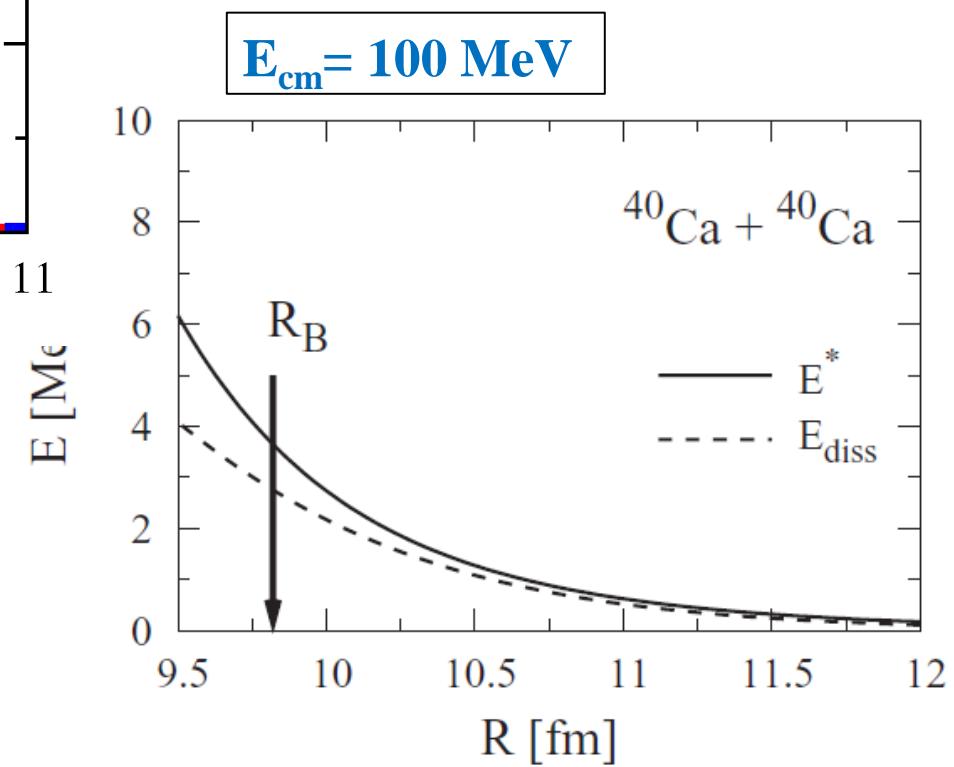
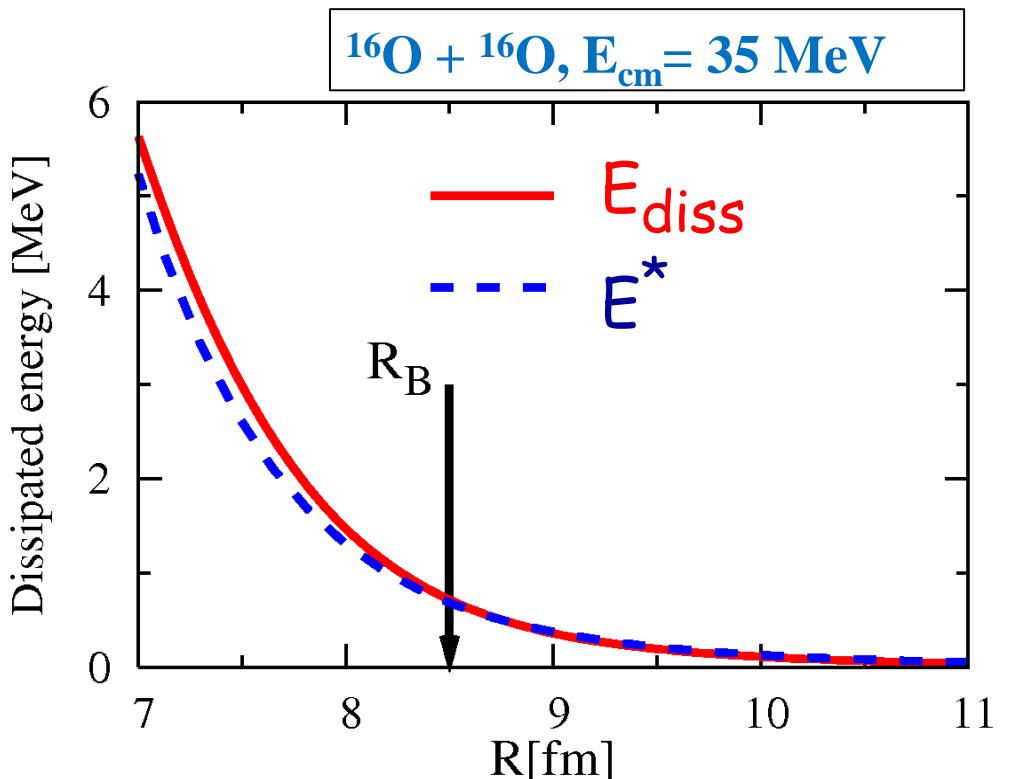
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$$E_{cm} = P^2/2\mu + V + E_{dissipation}$$

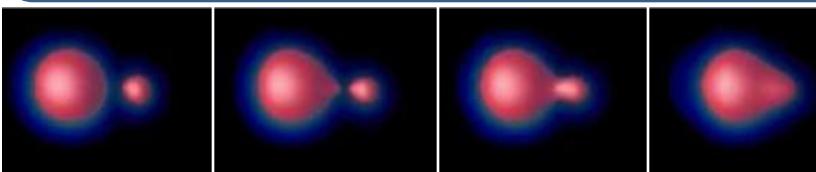


Nuclear friction: dissipated energy



Nuclear friction: energy dissipation

Four sequential images showing the evolution of a nuclear system.

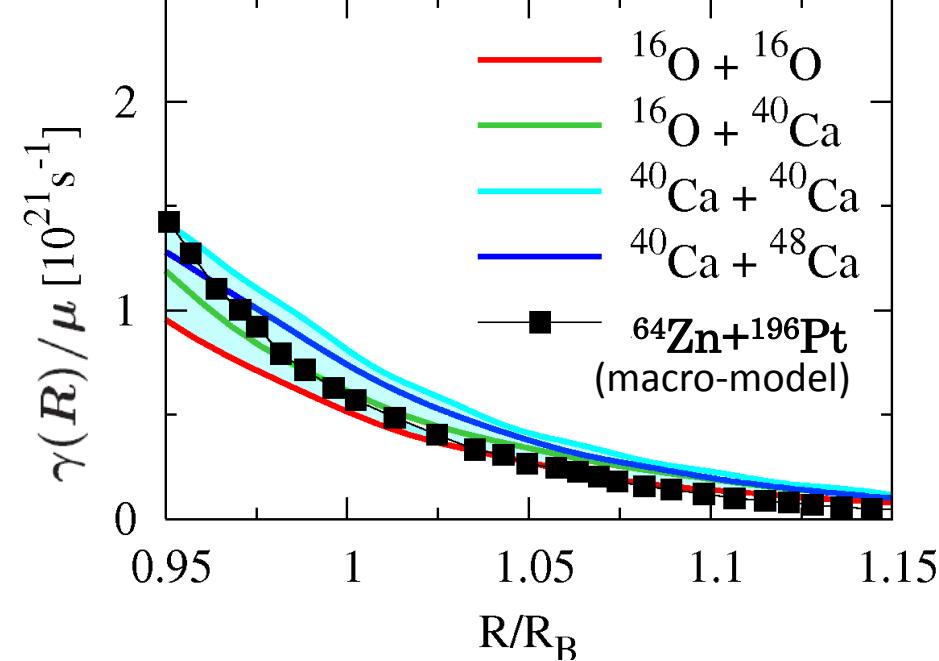
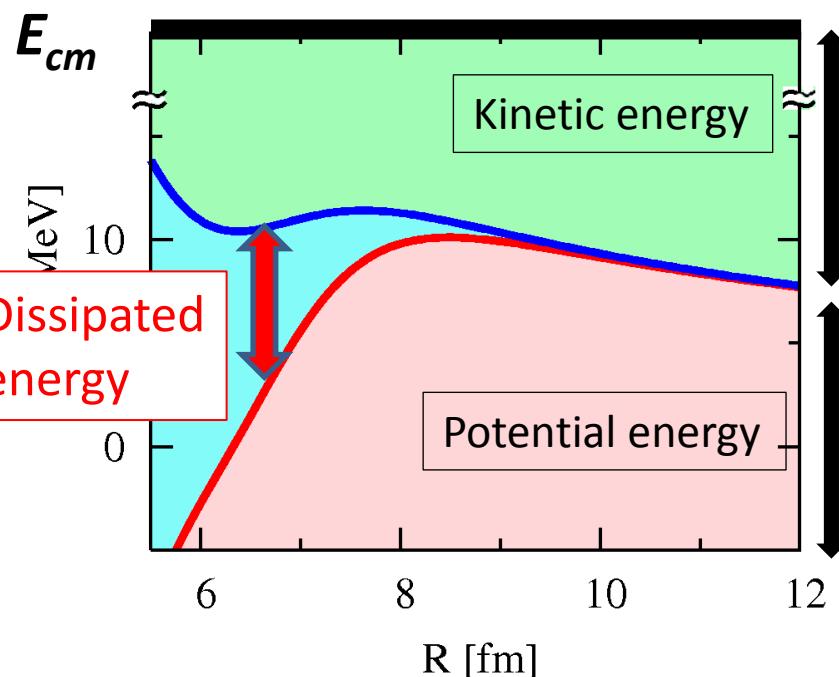


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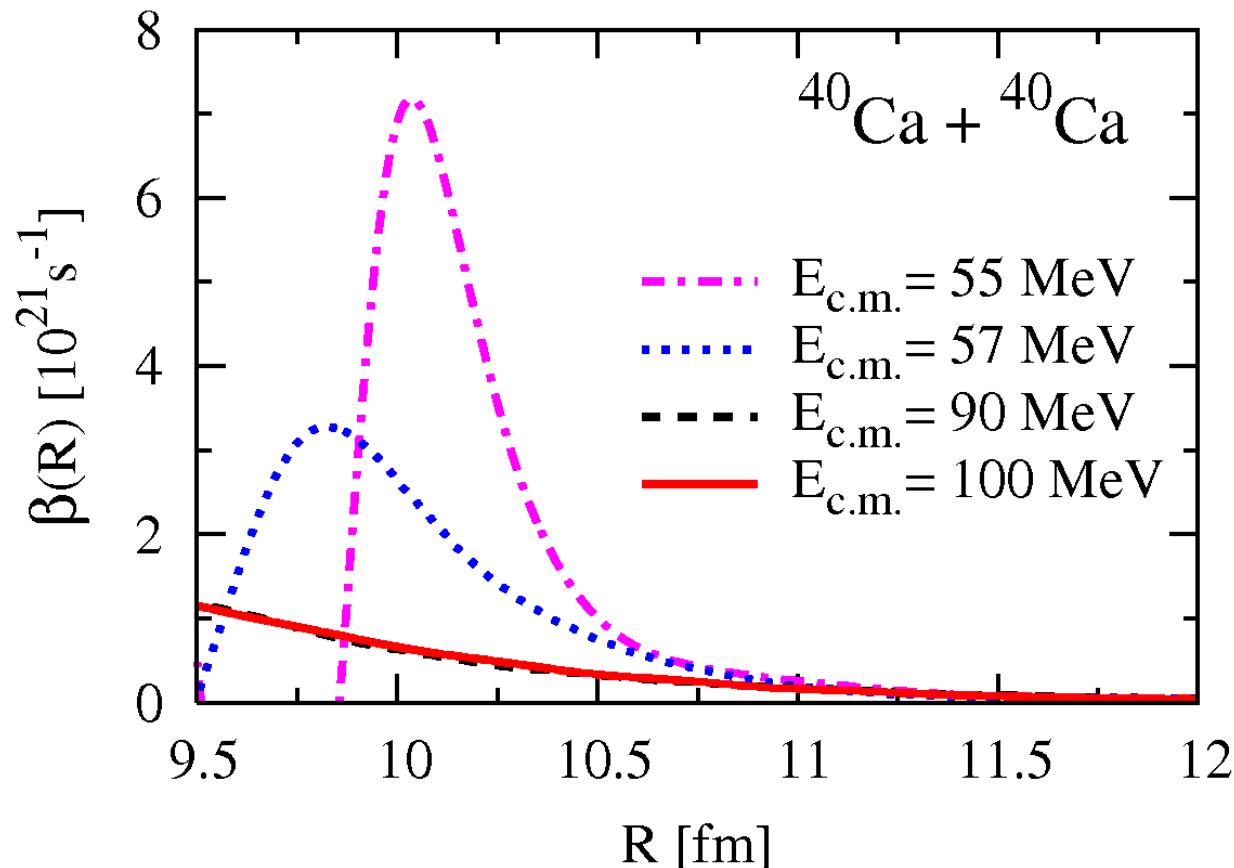
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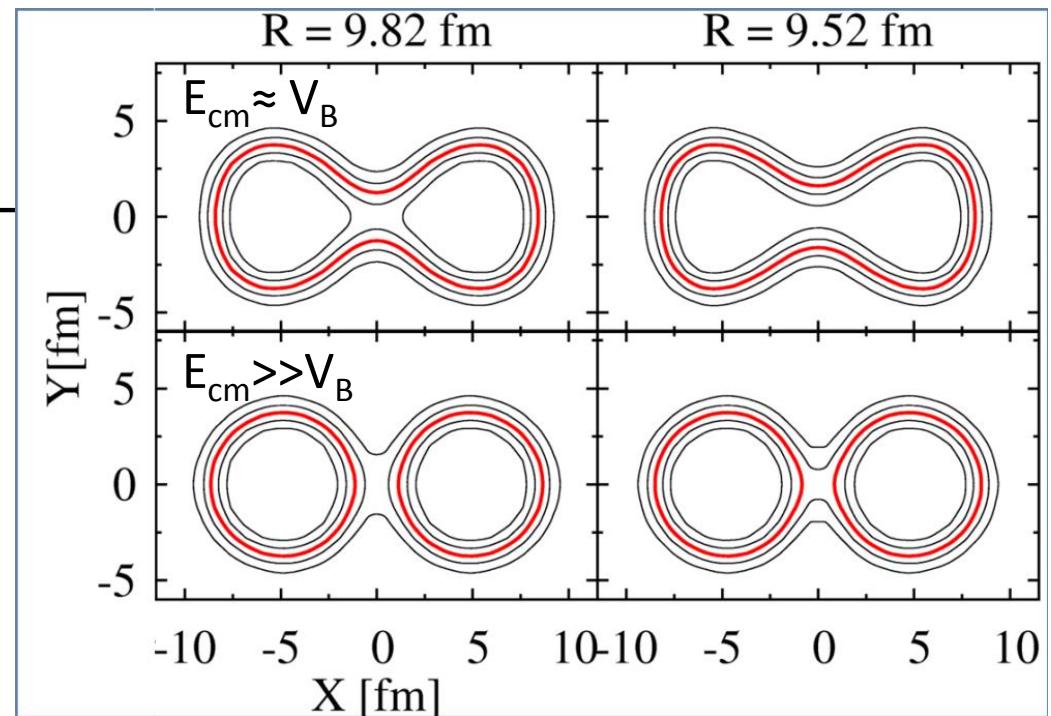
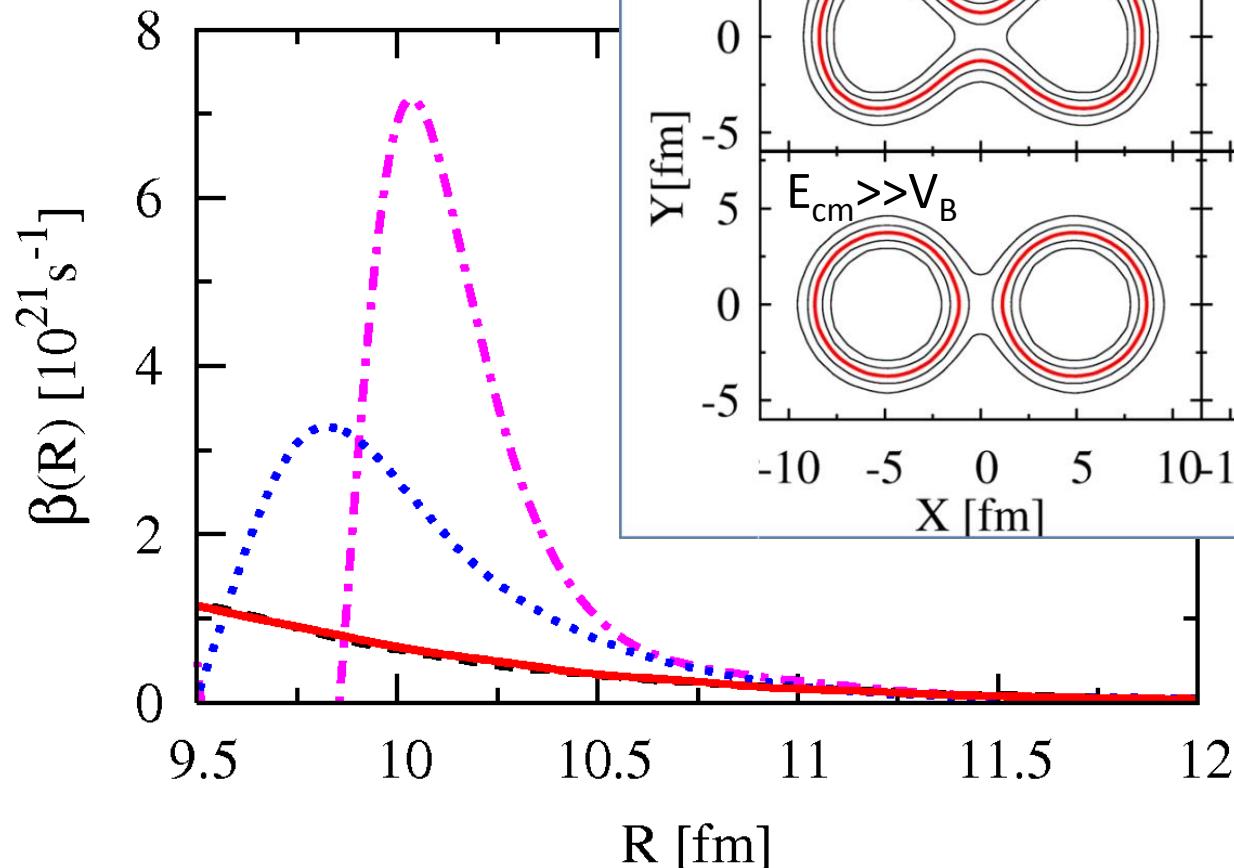
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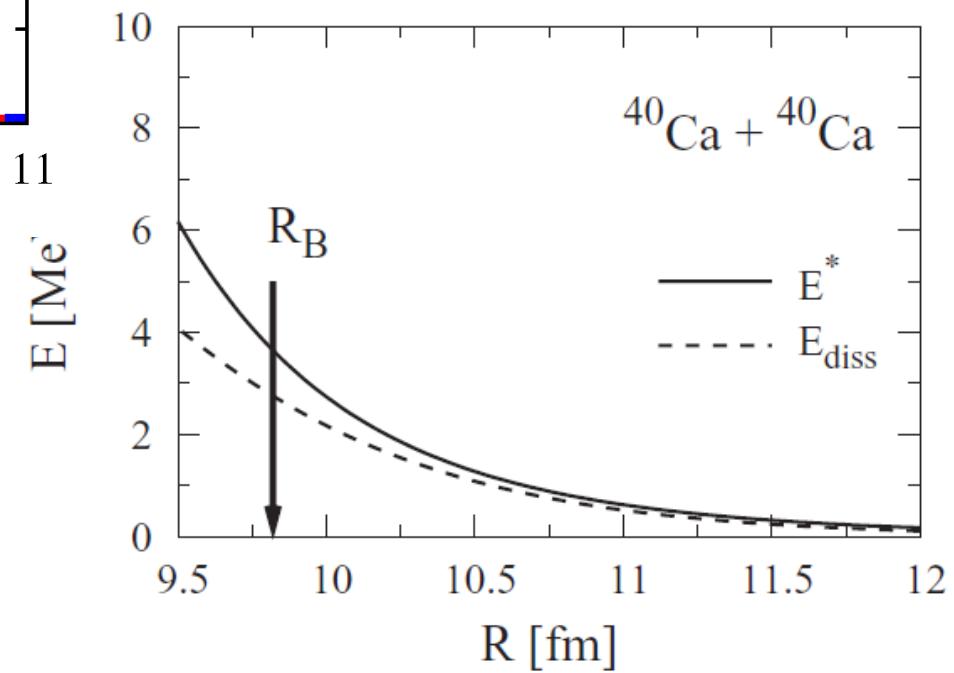
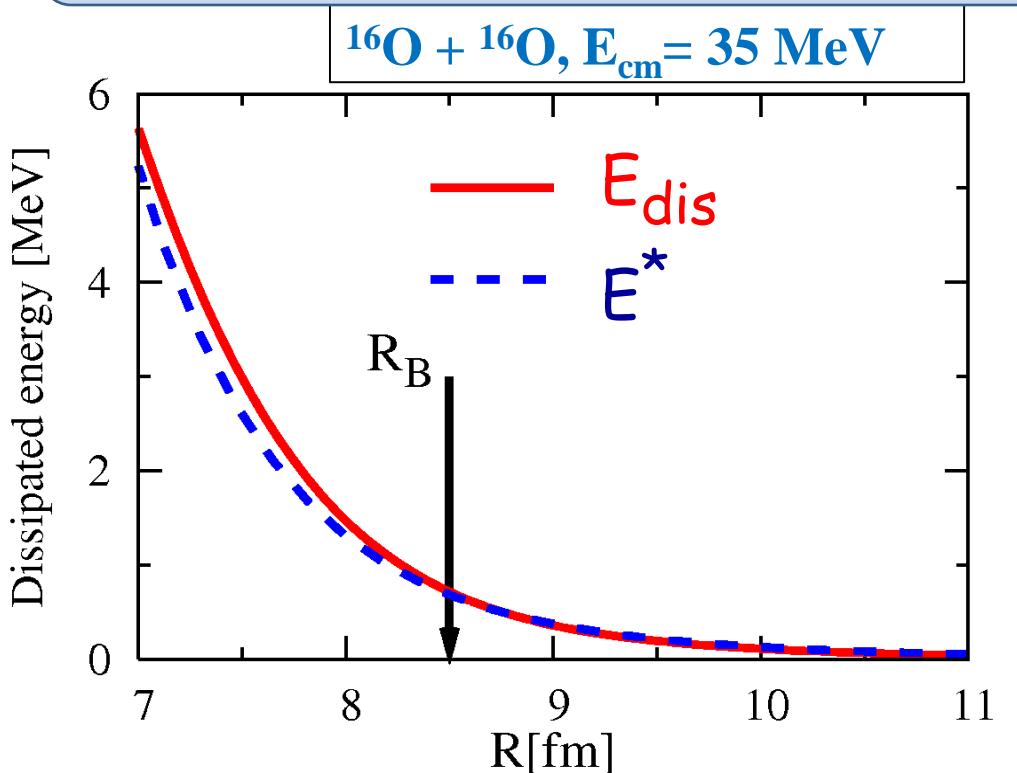
Energy dependence of nuclear friction



Energy dependence of nuclear friction

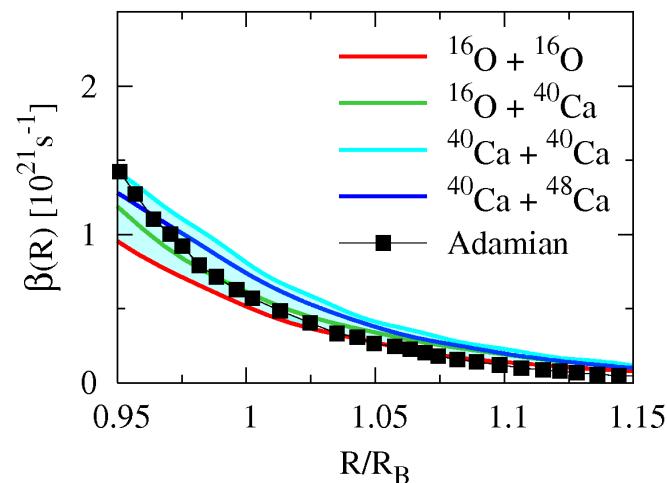
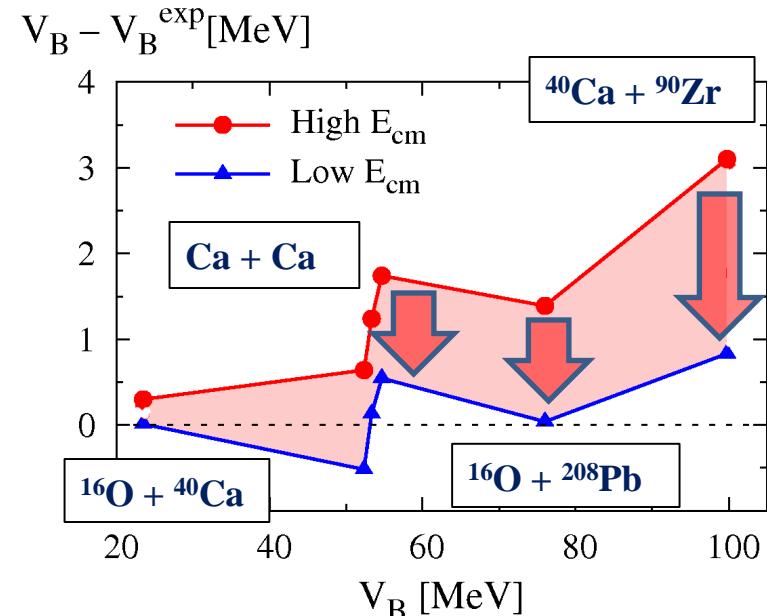


Nuclear friction: dissipated energy



Summary

- Nucleus-nucleus potentials are systematically extracted from TDHF.
- Dynamical reduction of potential energy close to the Coulomb barrier energy
- Extracted potentials agree with exp. data
- One-body energy dissipation are extracted from microscopic theory
- Universal property of extracted friction
- Energy dependence of friction



Mean-field (one-body) fluctuations

Mean-field dynamics

- good description for average evolution
- fluctuations are underestimated

- Bonche et al., PRC13 (1976)
- Flocard et al., PRC17 (1978)
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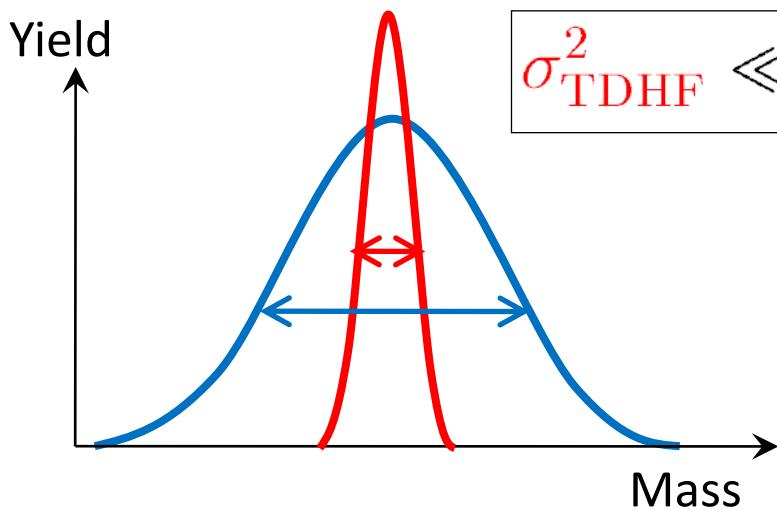
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Width of mass distribution in DIC



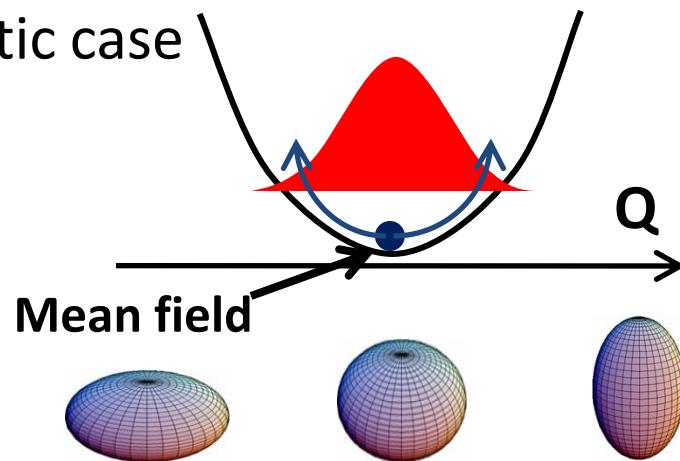
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Beyond mean field with a stochastic method

Stochastic Mean Field

- Static case

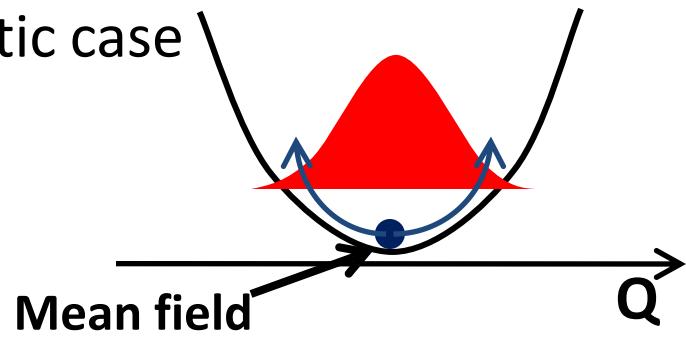


Configuration mixing calculation
(Generator coordinate method)

$$|\Psi\rangle = \int dQ f(Q) |\Phi(Q)\rangle$$

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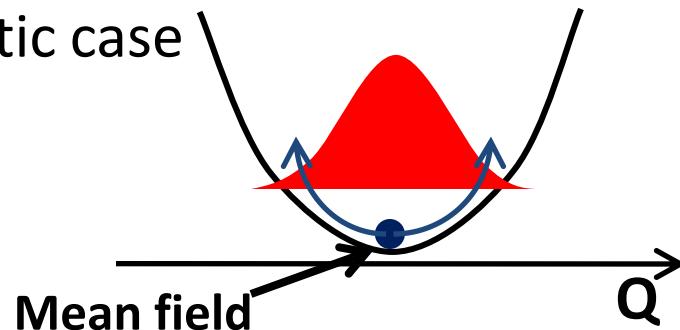
Mean Field



$$i\hbar \frac{\partial}{\partial t} \rho = [h(\rho), \rho]$$

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Mean Field $\xrightarrow{\hspace{1cm}}$ Stochastic Mean Field



$$i\hbar \frac{\partial}{\partial t} \rho = [h(\rho), \rho]$$

$$|\Psi^{\lambda=1}(t_0)\rangle$$

$$|\Psi^{\lambda=2}(t_0)\rangle$$

:

$$|\Psi^{\lambda=N}(t_0)\rangle$$



$$|\Psi^{\lambda=1}(t_1)\rangle$$

$$|\Psi^{\lambda=2}(t_1)\rangle$$

:

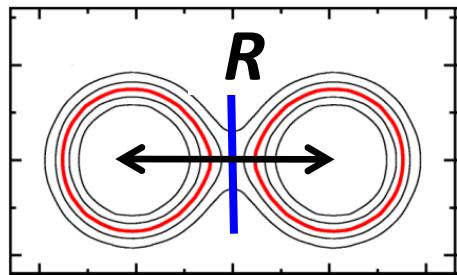
$$|\Psi^{\lambda=N}(t_1)\rangle$$

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Ayik, PLB658, 174(2008)

Ayik, Washiyama, Lacroix, PRC79, 054906(2009)

Application to fusion reaction

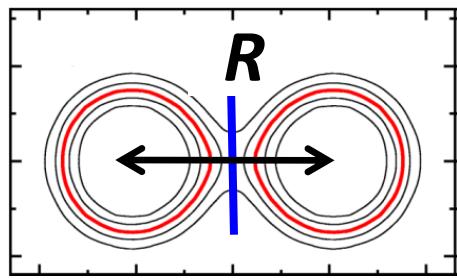


Mean-field



$$\frac{d}{dt}P = -\frac{d}{dR}U(R) - \gamma(R)\dot{R}$$

Application to fusion reaction



Mean-field

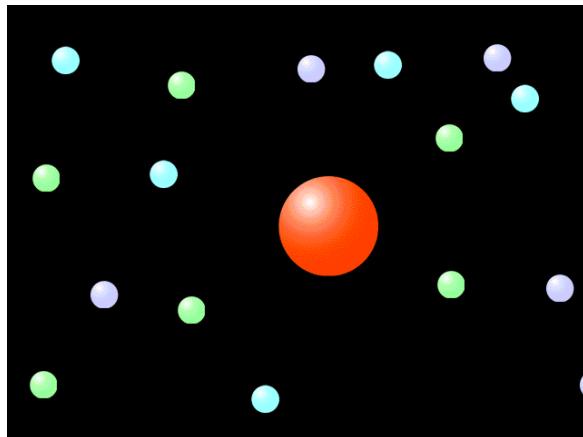


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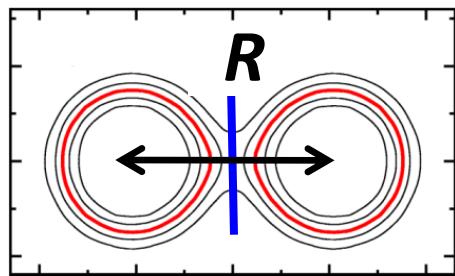
Stochastic mean-field

$$\frac{d}{dt}P^\lambda = -\frac{d}{dR^\lambda}U(R^\lambda) - \gamma(R^\lambda)\dot{R}^\lambda + \xi_P^\lambda(t)$$

cf. Brownian motion



Application to fusion reaction



Mean-field



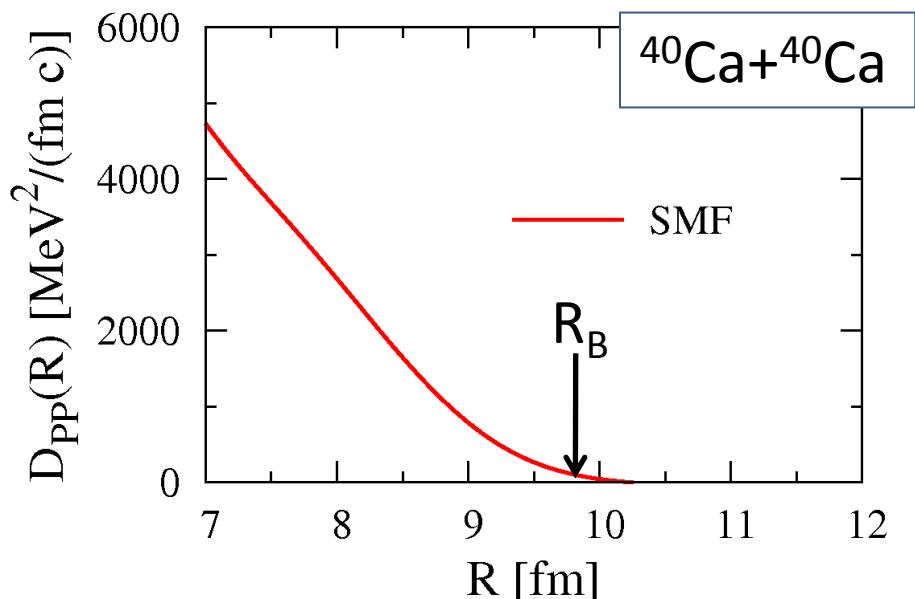
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$$\overline{\xi_P^\lambda(t)\xi_P^\lambda(t')} = 2\delta(t-t')D_{PP}(R)$$



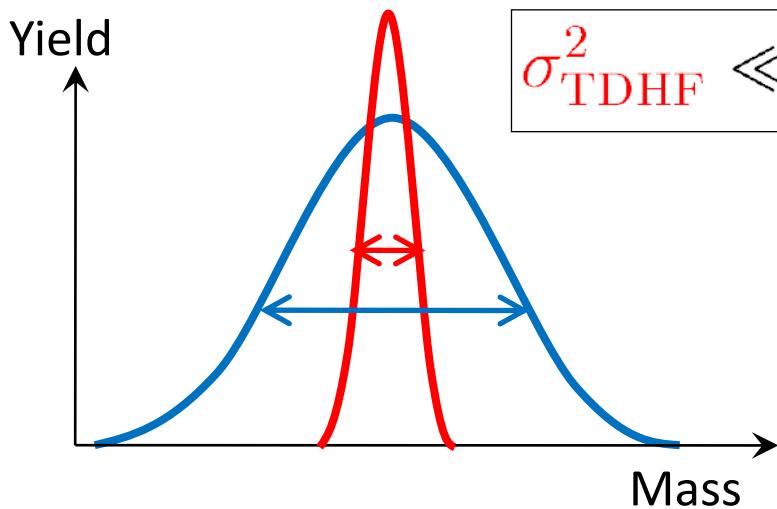
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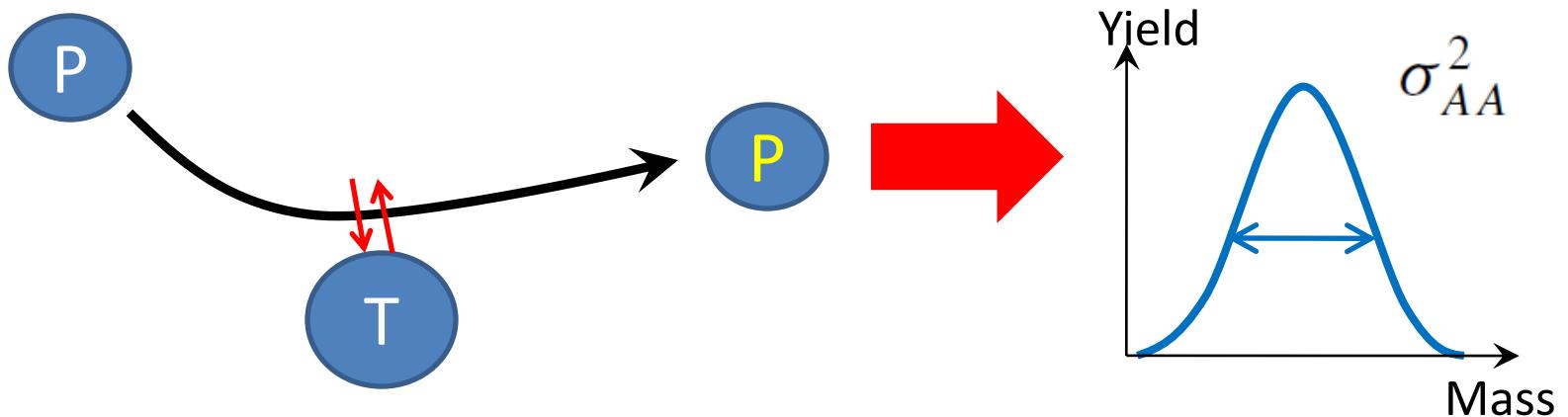


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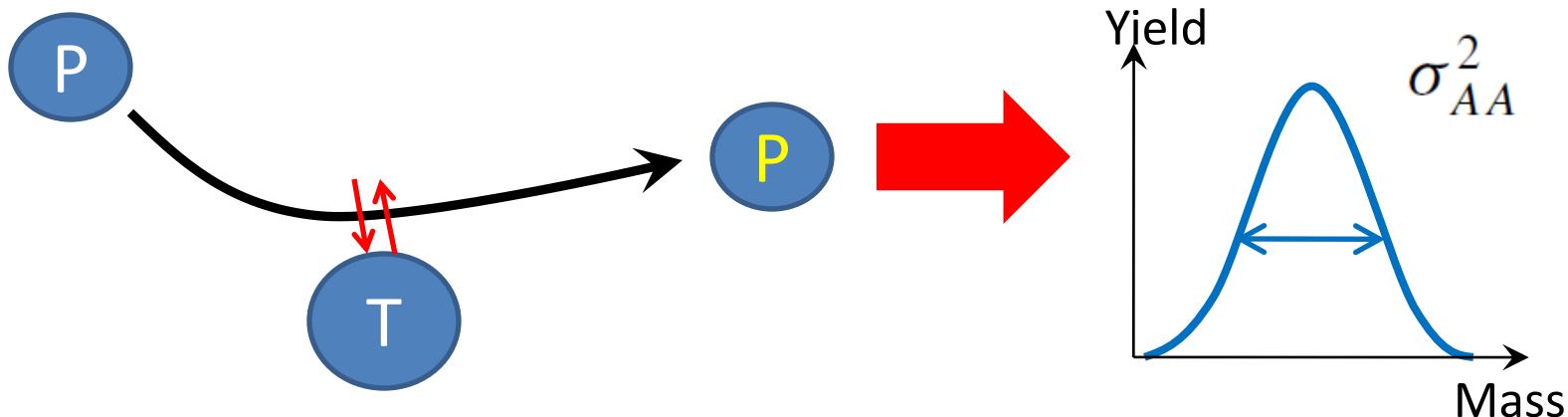


Beyond mean field with a stochastic method

Failure of TDHF for mass variance



Failure of TDHF for mass variance



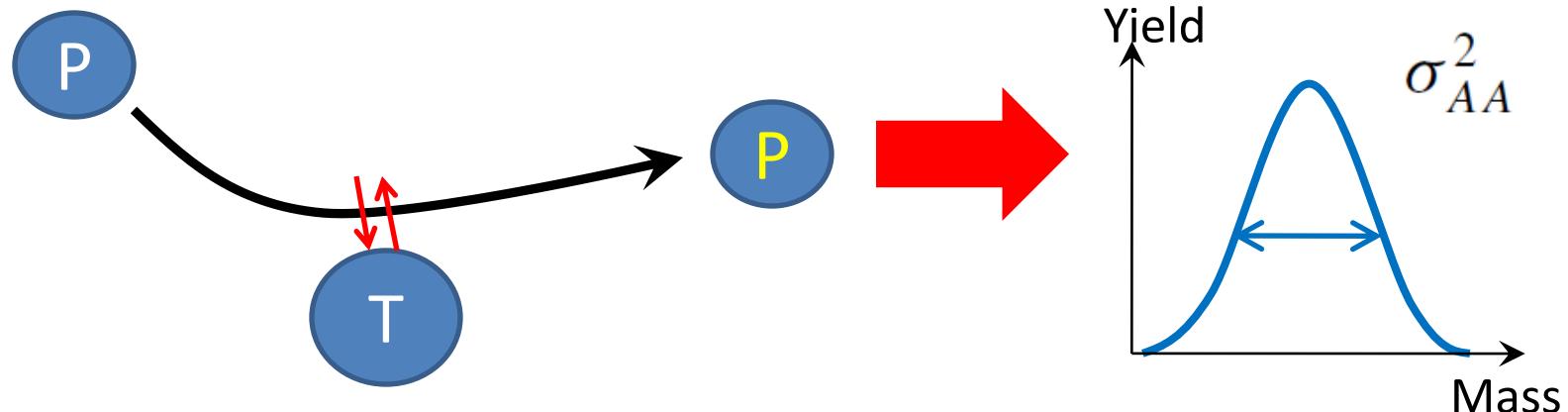
- Analysis based on phenomenological classical model

$$\sigma_{AA}^2(t) \simeq N_{\text{exc}}(t)$$

(N_{exc} : the number of exchanged nucleons)

Freiesleben, Kratz,
Phys. Rep. 106(1984)1

Failure of TDHF for mass variance



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$V_B = 53.4 \text{ MeV}$	$^{40}\text{Ca} + ^{40}\text{Ca}$	$E_{\text{c.m.}} (\text{MeV})$	σ_{TDHF}^2	N_{exc}
		51.0	0.004	0.432
		52.5	0.008	1.441
		53.0	0.008	3.634

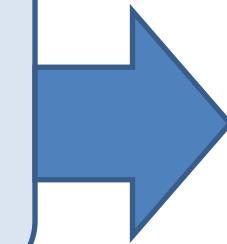
Below the table, there is a black line plot showing a peak, with a blue sphere representing a nucleus positioned near the peak.

Mass fluctuations and exchanged nucleons

Stochastic mean-field

$$\frac{d}{dt} A_P^\lambda = v(A_P^\lambda, t) + \xi_A^\lambda(t)$$

$$\overline{\xi_A^\lambda(t)\xi_A^\lambda(t')} = 2\delta(t-t')D_{AA}$$



$$\sigma_{AA}^2(t) \simeq 2 \int_0^t D_{AA}(s)ds$$

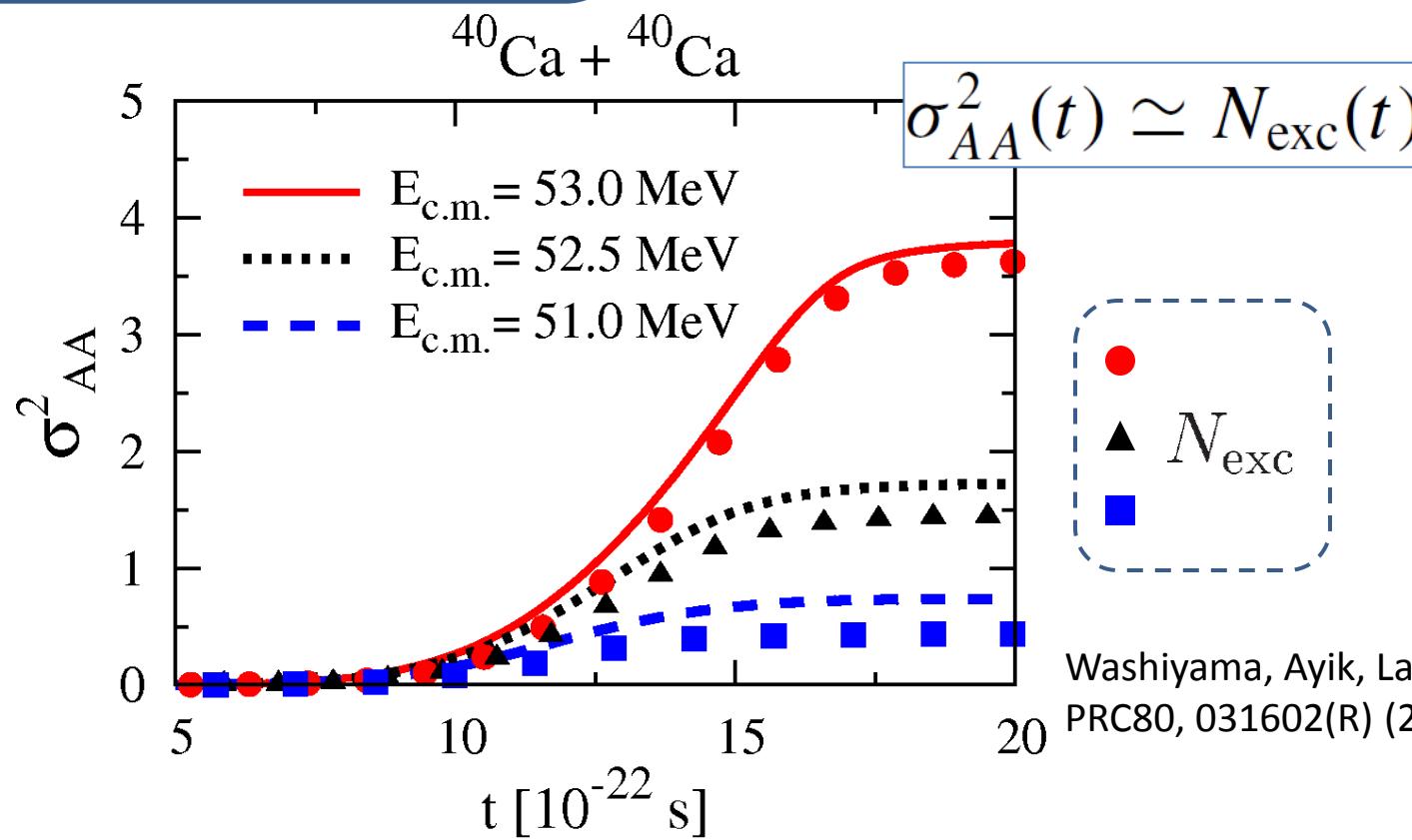
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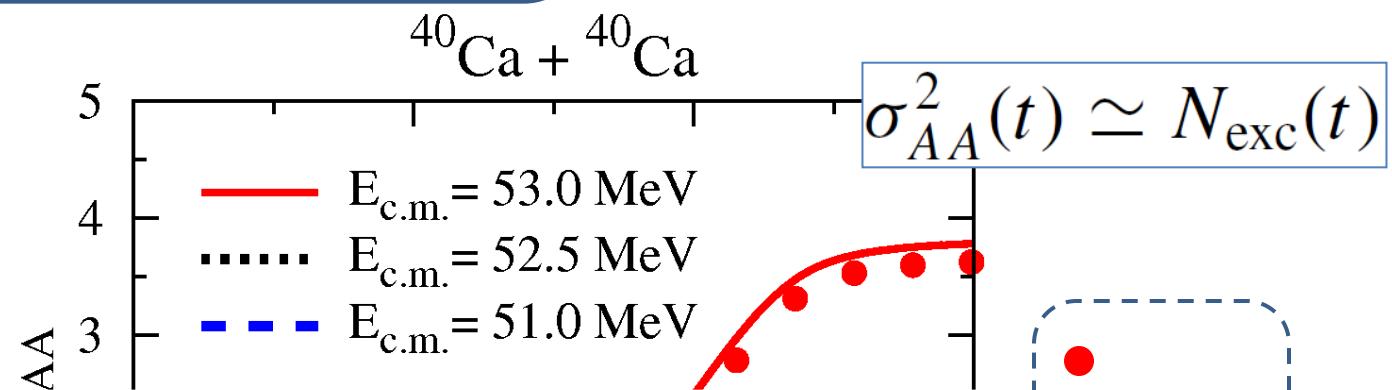
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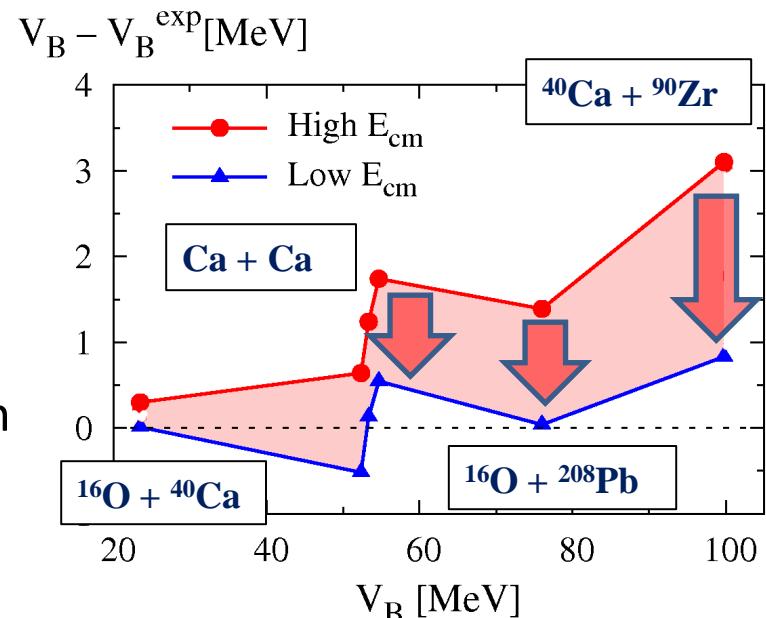
$E_{\text{c.m.}}(\text{MeV})$	σ_{TDHF}^2	N_{exc}	σ_{AA}^2
51.0	0.004	0.432	0.730
52.5	0.008	1.441	1.718
53.0	0.008	3.634	3.790

Washiyama, Ayik, Lacroix,
PRC80, 031602(R) (2009)

Summary

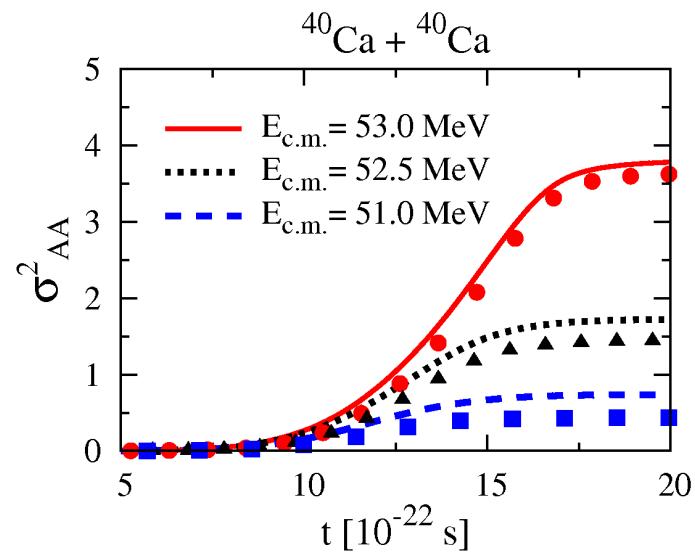
● Mean-field dynamics

- Dynamical reduction in potential
- Good agreement with experiments
- Universal behavior of Energy dissipation



● Beyond mean field

- Stochastic mean field
- Good description for mass variance σ_{AA}^2
- Fluctuation for other collective variables



Introduction: Low energy nuclear dynamics

(Energy < 10 MeV/A)

□ Goal: Unified microscopic theory for nuclear structure and dynamics

