

# 超伝導量子回路における開放量子系の制御

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NanoQuine  
QUantum INFORMATION Electronics

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# General picture as non-equilibrium system



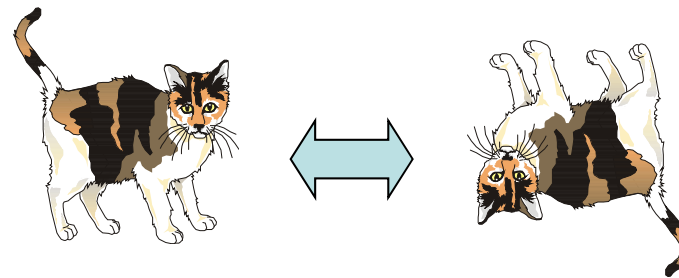
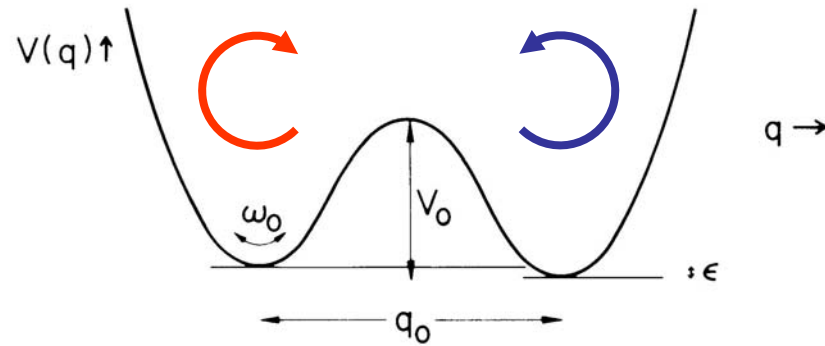
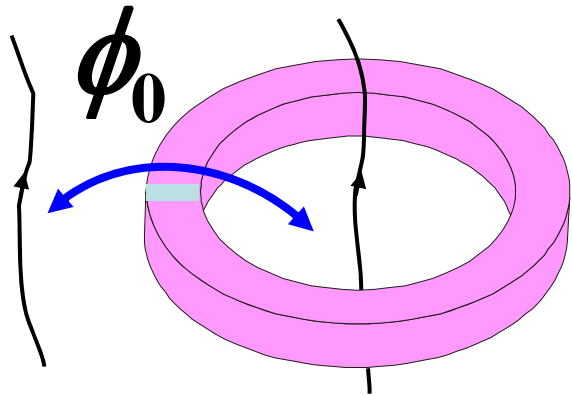
# Outline

- Superconducting qubits as an artificial atom
- Decoherence and environment
  - Qubit as quantum spectrum analyzer
- Microwave quantum optics
  - Atoms in cavity
  - Atoms in 1D waveguide

# Macroscopic quantum coherence

A.J. Leggett, Prog. Theor. Phys. Suppl. 69, 80 (1980); Phys. Scr. T102, 69 (2002).

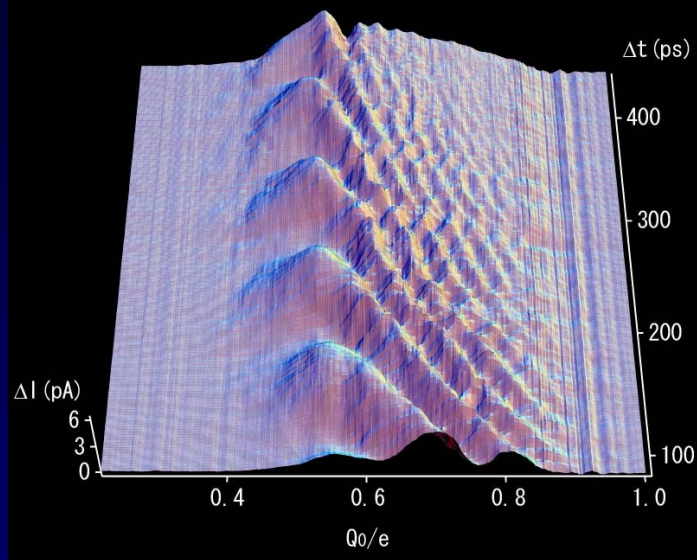
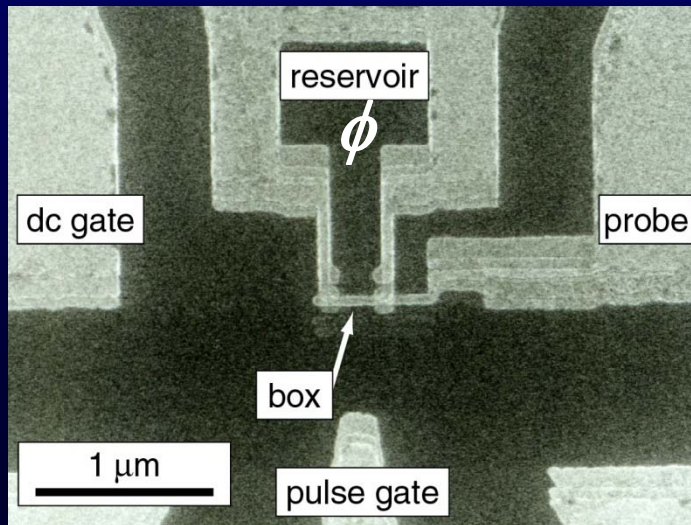
Does quantum mechanics hold in macroscopic systems?



Schrödinger 1935

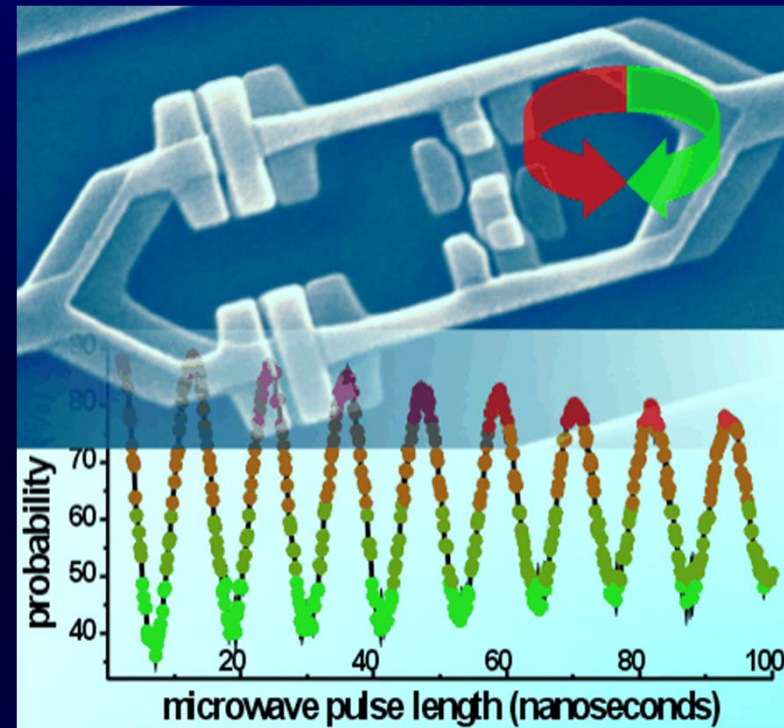
# Superconducting qubits

## Charge qubit



Y. Nakamura et al. Nature (1999)

## Flux qubit

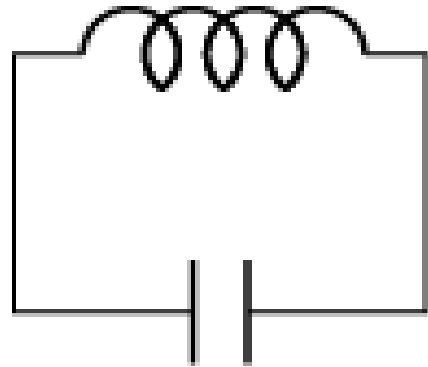


Chiorescu, Nakamura, Harmans, Mooij, Science (2003)

- Artificial two-level system in electric circuits
- Coherent control of quantum states in macroscopic systems

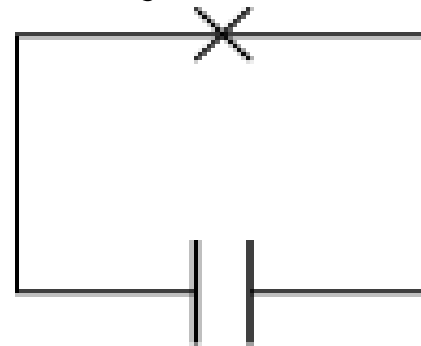
# Superconducting qubit – nonlinear resonator

LC resonator

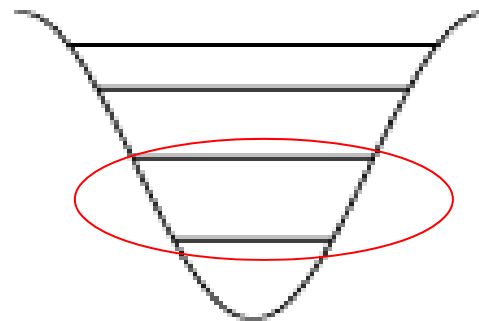
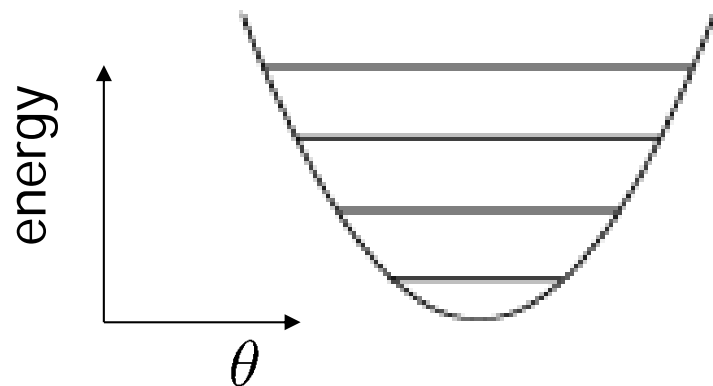


Josephson junction resonator

Josephson junction = nonlinear inductor



anharmonicity  $\Rightarrow$  effective two-level system



inductive energy = confinement potential

charging energy = kinetic energy  $\Rightarrow$  quantized states

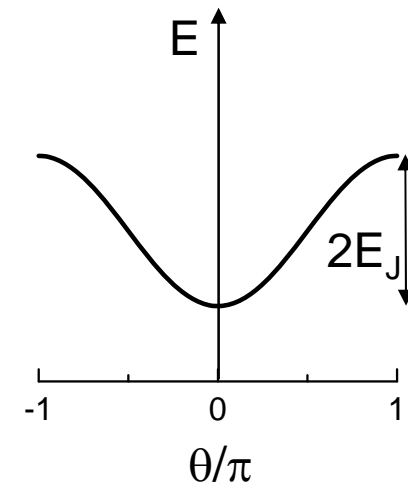
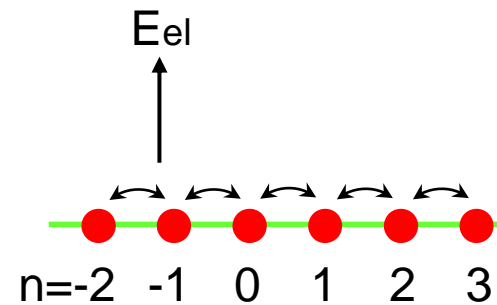
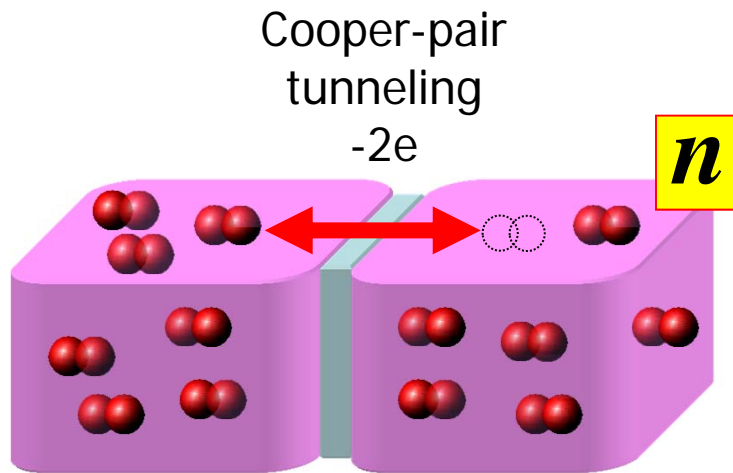


# Josephson effect

B. D. Josephson 1962

number  $n \Leftrightarrow$  phase difference  $\theta$

$$[n, \theta] = -i$$



$$H = -\frac{E_J}{2} \sum_n \{ |n\rangle \langle n+1| + |n+1\rangle \langle n| \} = - \int_0^{2\pi} d\theta E_J \cos \theta |\theta\rangle \langle \theta|$$

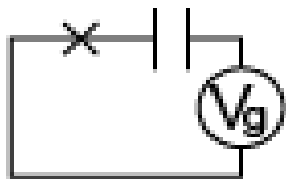
Tight-binding model in 1d lattice  $\Rightarrow$  Bloch band

$$|\theta\rangle = \sum_n e^{in\theta} |n\rangle$$

# Superconducting qubits – artificial atoms in electric circuit

small ←  $E_J/E_C$  → large

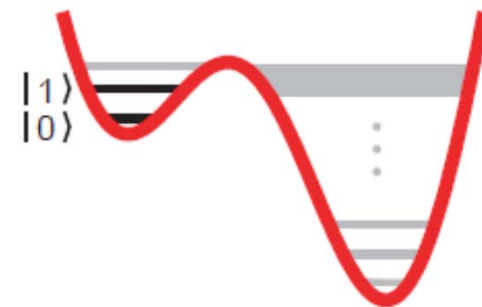
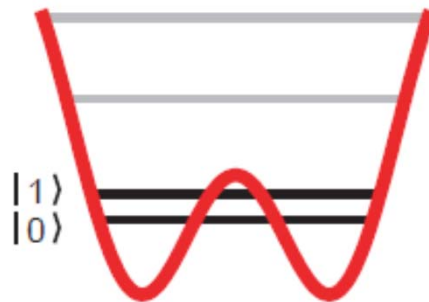
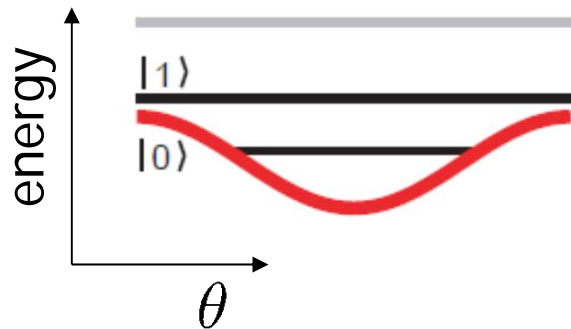
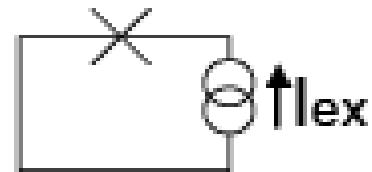
Charge qubit



Flux qubit



Phase qubit



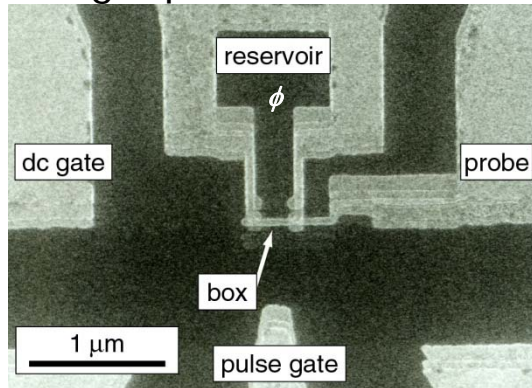
Josephson energy  $E_J$  = confinement potential  
 charging energy  $E_C$  = kinetic energy  $\Rightarrow$  quantized states

typical qubit energy  $E_{01} \sim 10 \text{ GHz} \sim 0.5 \text{ K}$

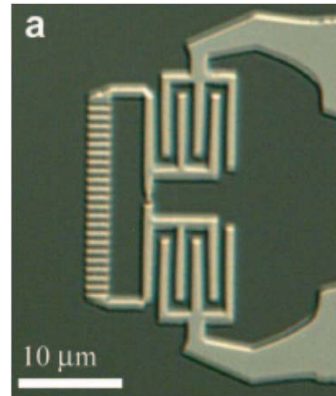
typical experimental temperature  $T \sim 0.02 \text{ K}$

# Superconducting qubits – macroscopic artificial atom in circuits

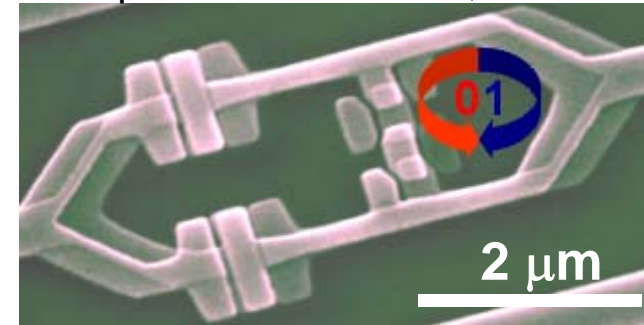
charge qubit/NEC  $E_J/E_C \sim 0.3$



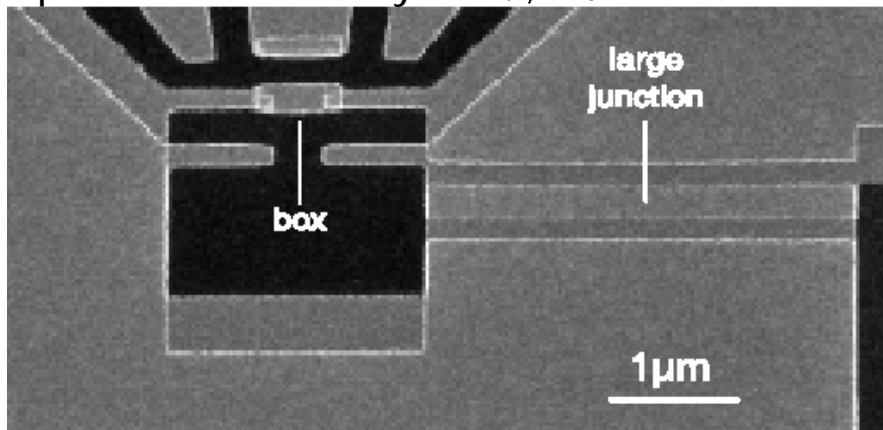
"fluxonium"/Yale



flux qubit/Delft  $E_J/E_C \sim 40$

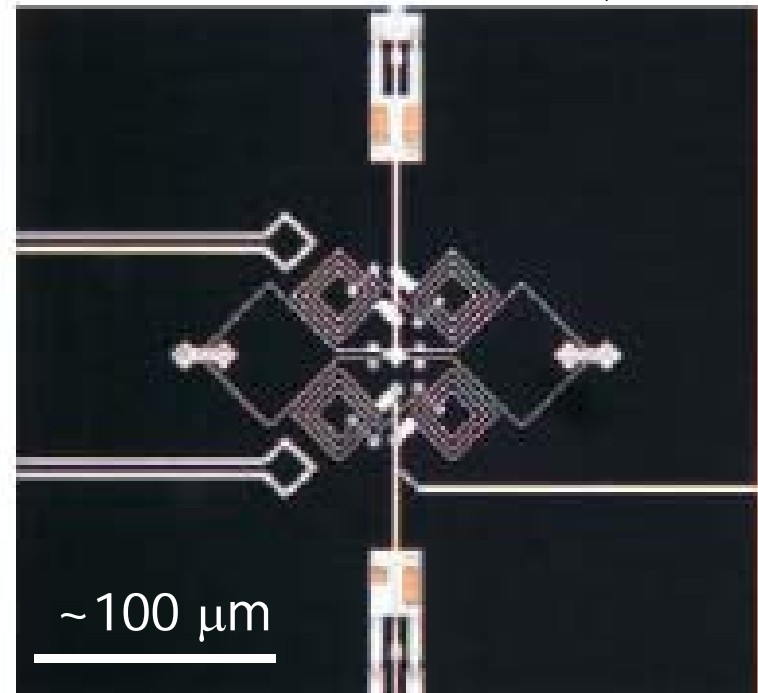


"quantronium"/Saclay  $E_J/E_C \sim 5$

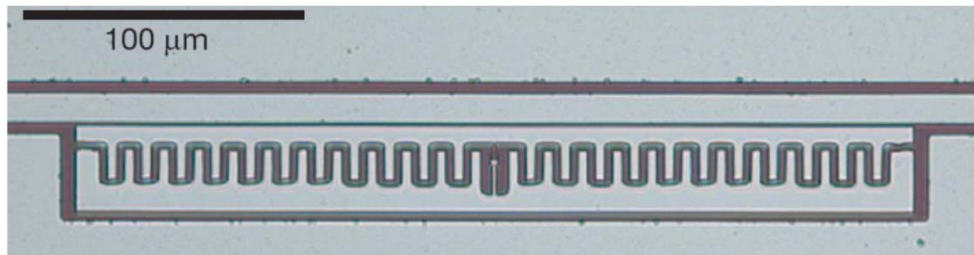


$E_J/E_C \sim 3$

phase qubit/NIST/UCSB  $E_J/E_C \sim 10^4$



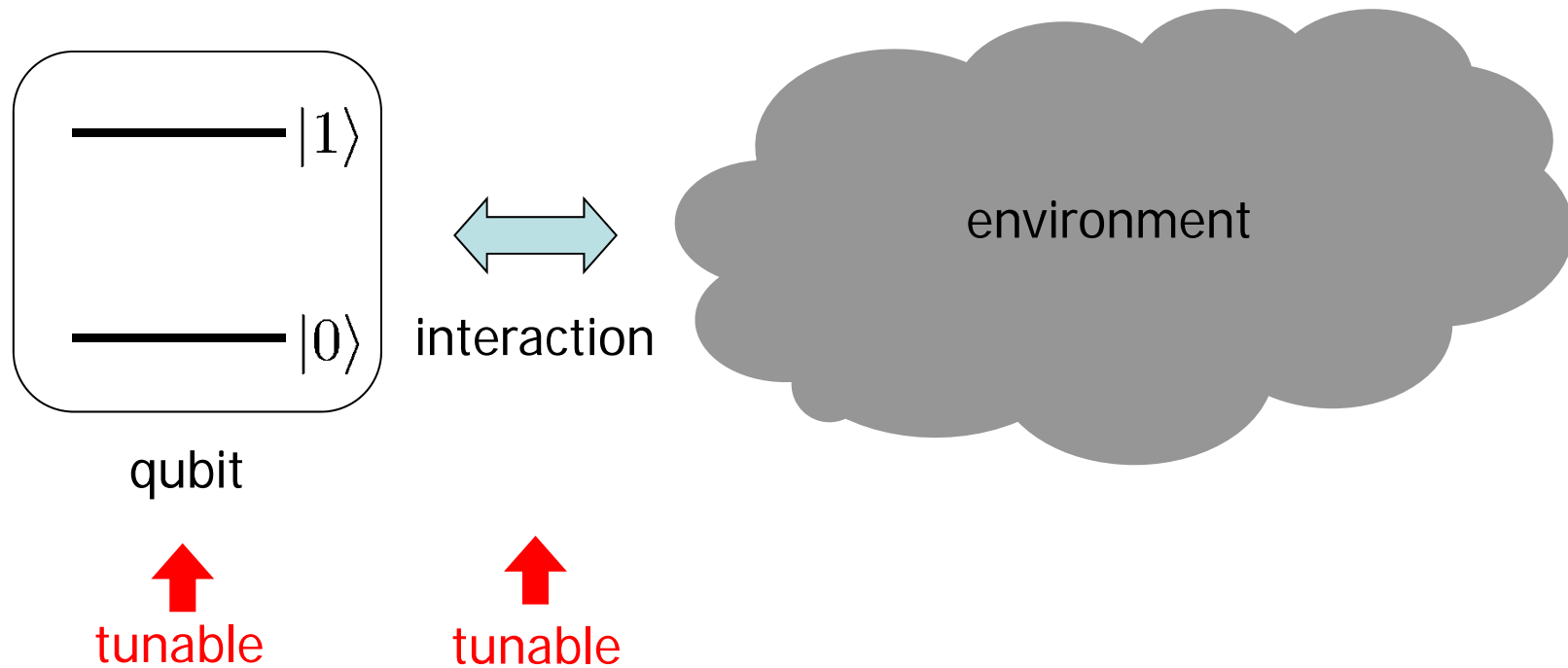
"transmon"/Yale  $E_J/E_C \sim 50$



# Decoherence

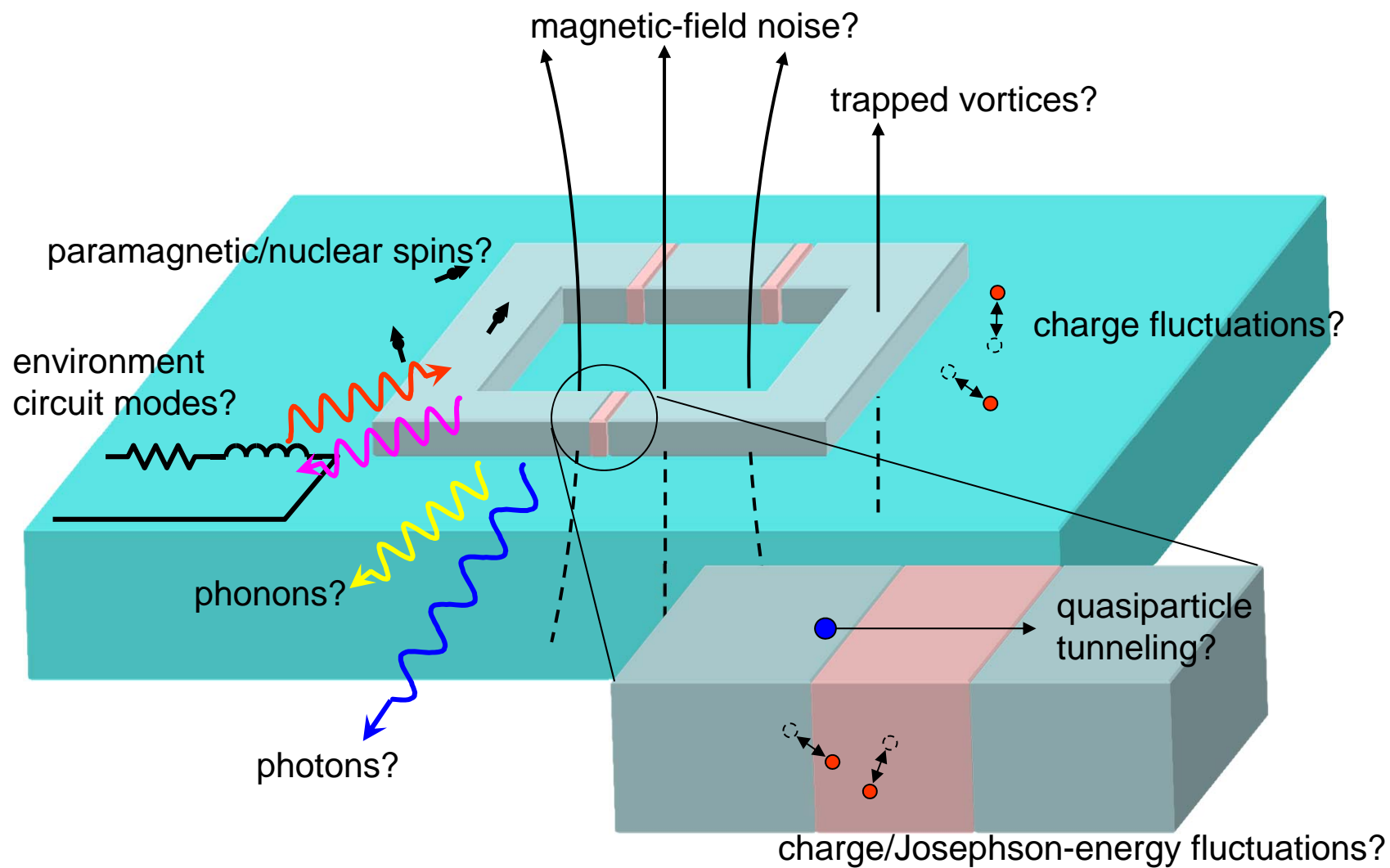
Loss of coherence due to coupling with uncontrolled environment

- dissipation
- dephasing



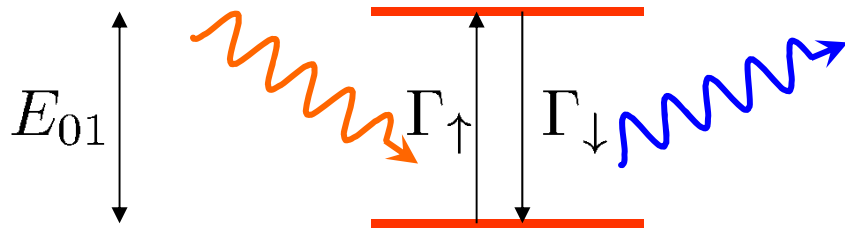
Qubit as a tool for characterization of environment

# Possible decoherence sources



# Energy relaxation

relaxation and excitation



$$\Gamma_{\downarrow} = \frac{2\pi}{\hbar^2} \left| \left\langle 0 \left| \frac{\partial H_q}{\partial \lambda} \right| 1 \right\rangle \right|^2 S_{\lambda} \left( \frac{E_{01}}{\hbar} \right)$$

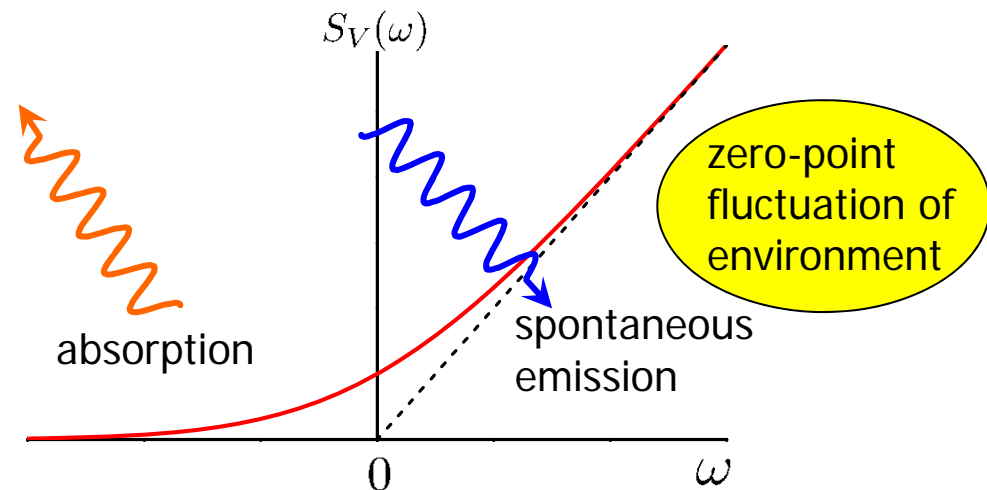
$$\Gamma_{\uparrow} = \frac{2\pi}{\hbar^2} \left| \left\langle 0 \left| \frac{\partial H_q}{\partial \lambda} \right| 1 \right\rangle \right|^2 S_{\lambda} \left( -\frac{E_{01}}{\hbar} \right)$$

for weak perturbation: Fermi's golden rule

- qubit energy  $E_{01}$  variable
  - relaxation  $\propto S(+E_{01}/\hbar)$  and excitation  $\propto S(-E_{01}/\hbar)$
- $\Rightarrow$  quantum spectrum analyzer

ex. Johnson noise in Ohmic resistor  $R$

$$S_V(\omega) = \frac{|\hbar\omega|}{2\pi} R \left( \coth \frac{|\hbar\omega|}{2k_B T} + 1 \right)$$

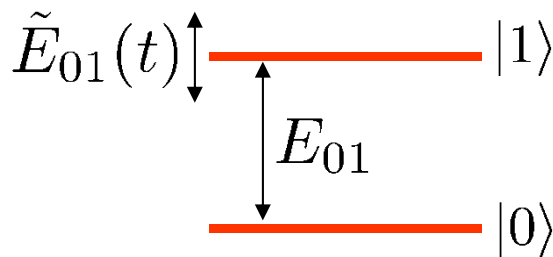


# Dephasing

free evolution of the qubit phase

$$\alpha|0\rangle + \beta|1\rangle \longrightarrow \alpha|0\rangle + \beta e^{i\varphi(t)}|1\rangle$$

$$\varphi(t) = \bar{\varphi}(t) + \underbrace{\tilde{\varphi}(t)}_{\text{dephasing}} \quad \bar{\varphi}(t) = \frac{\overline{E_{01}}}{\hbar} t$$



$$\langle \exp i\tilde{\varphi}(t) \rangle = \left\langle \exp \left( \frac{i}{\hbar} \int_0^\tau dt \tilde{E}_{01}(t) \right) \right\rangle$$

$$= \exp \left[ -\frac{1}{2\hbar^2} \left( \frac{\partial E_{01}}{\partial \lambda} \right)^2 \int_{-\infty}^{\infty} d\omega S_\lambda(\omega) \left( \frac{\sin(\omega t/2)}{\omega/2} \right)^2 \right]$$

for Gaussian fluctuations

sensitivity of qubit energy to the fluctuations of external parameter

tunable

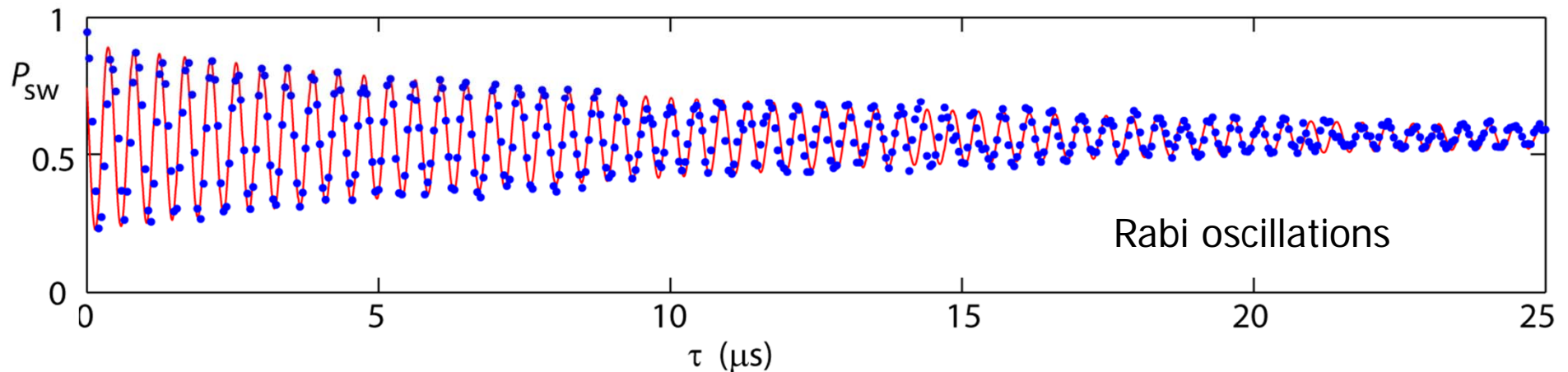
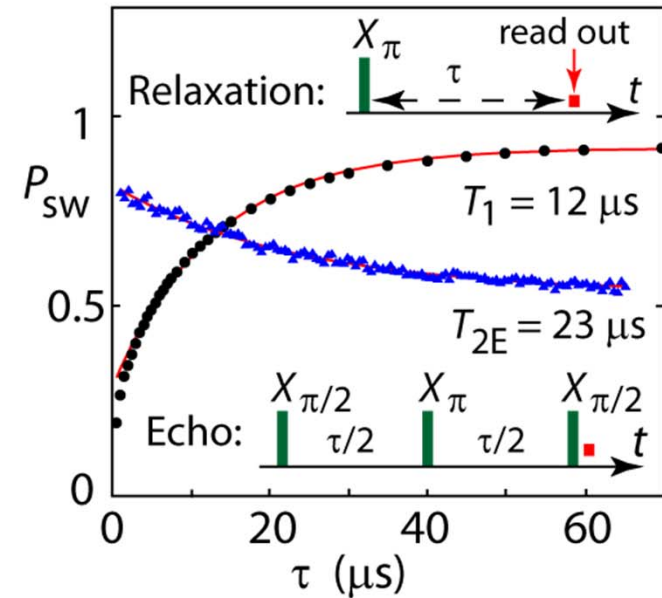
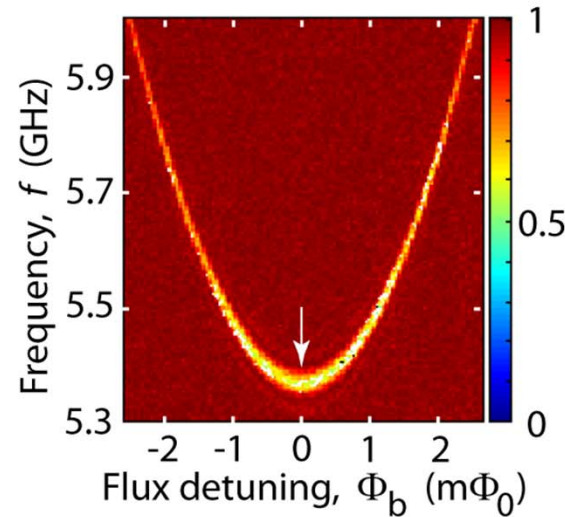
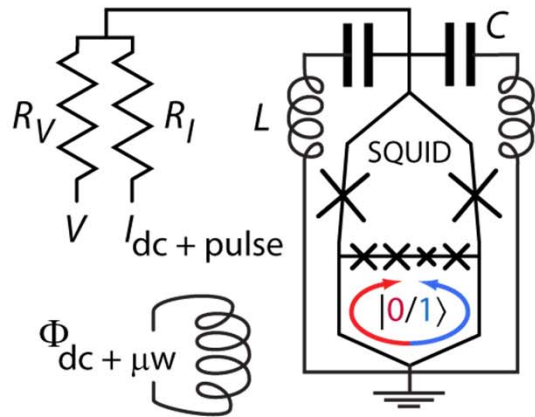
tunable

information of  $S(\omega)$  at low frequencies

# A long-lived flux qubit

Aluminum 4JJ flux qubit on SiO<sub>2</sub>/Si

$$Q = \omega_{01} T_1 \approx 400,000$$

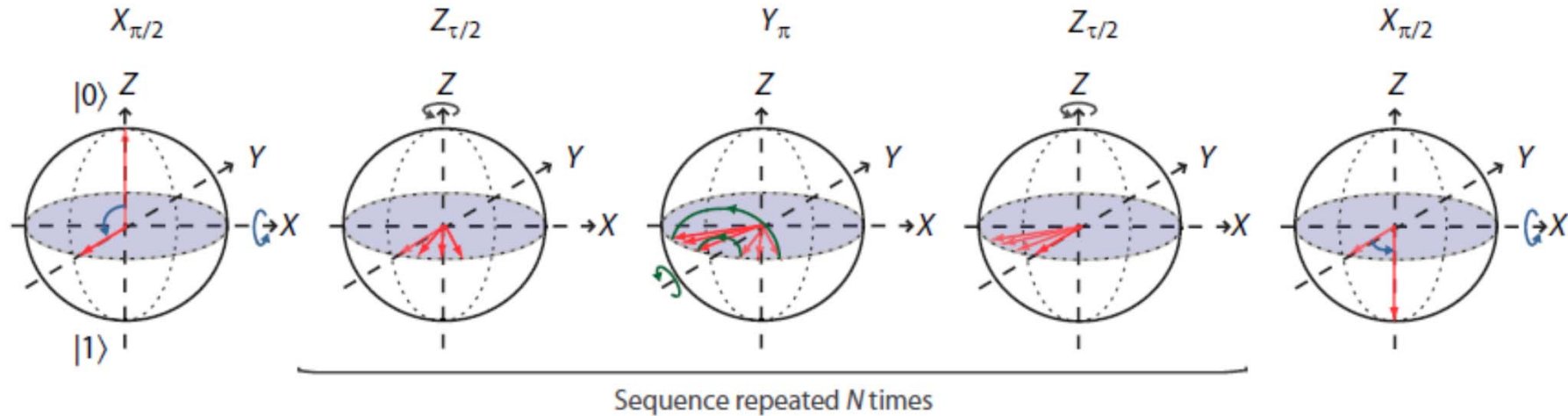


In collaboration with Jonas Bylander, Simon Gustavsson, Fei Yan, Will Oliver (MIT)

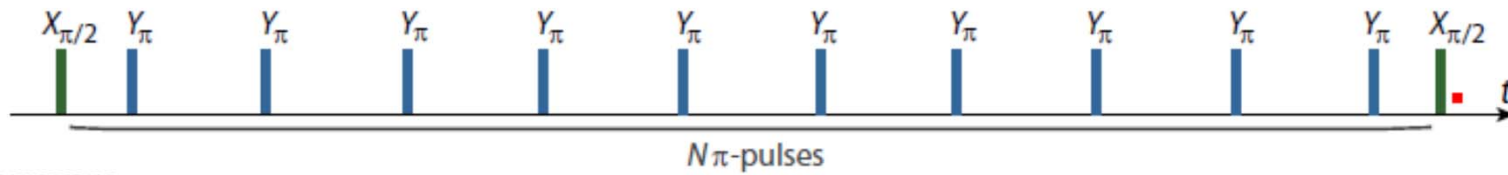


# Dynamical decoupling pulse sequences

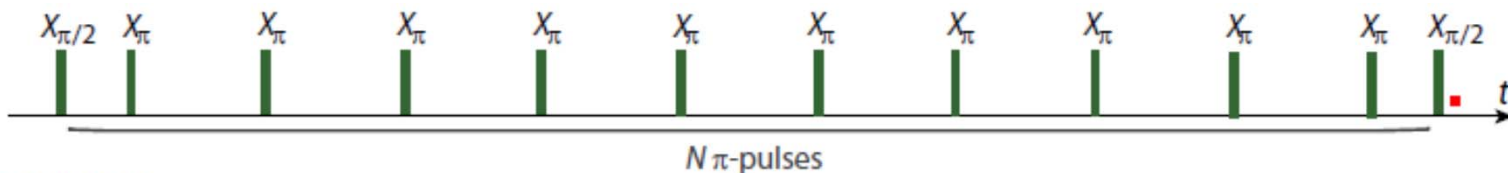
CPMG rotations



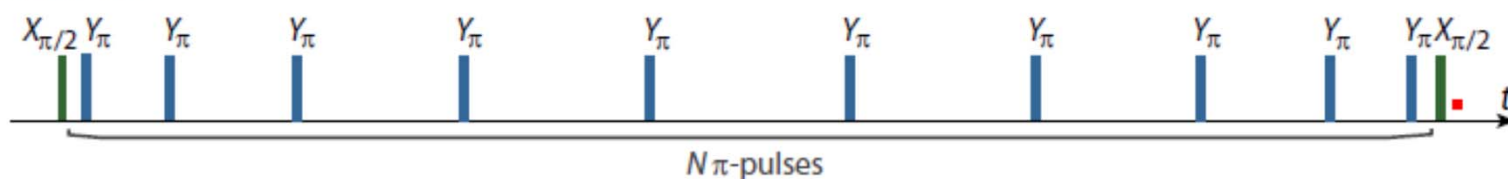
CPMG sequence



CP sequence

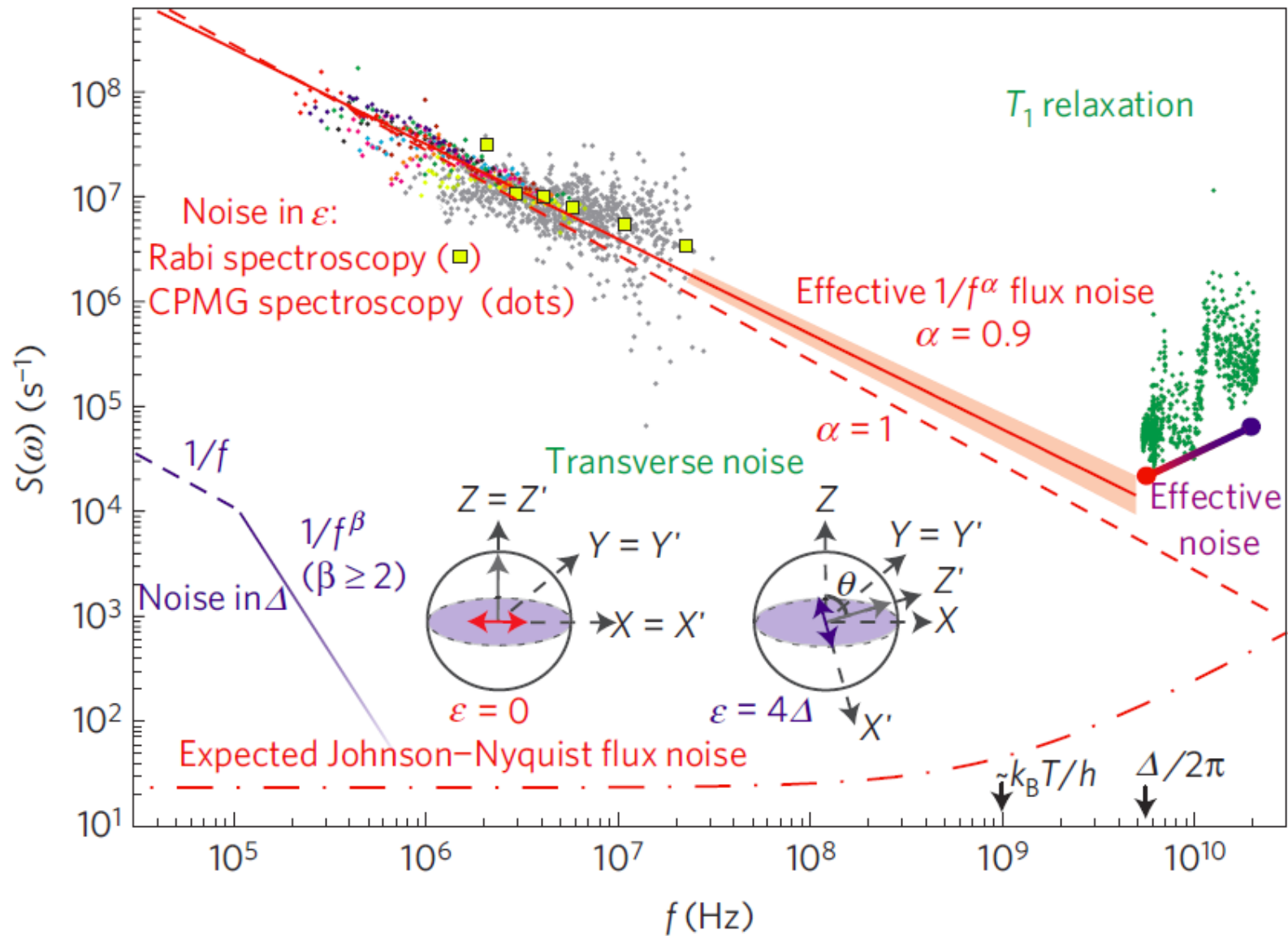


UDD sequence

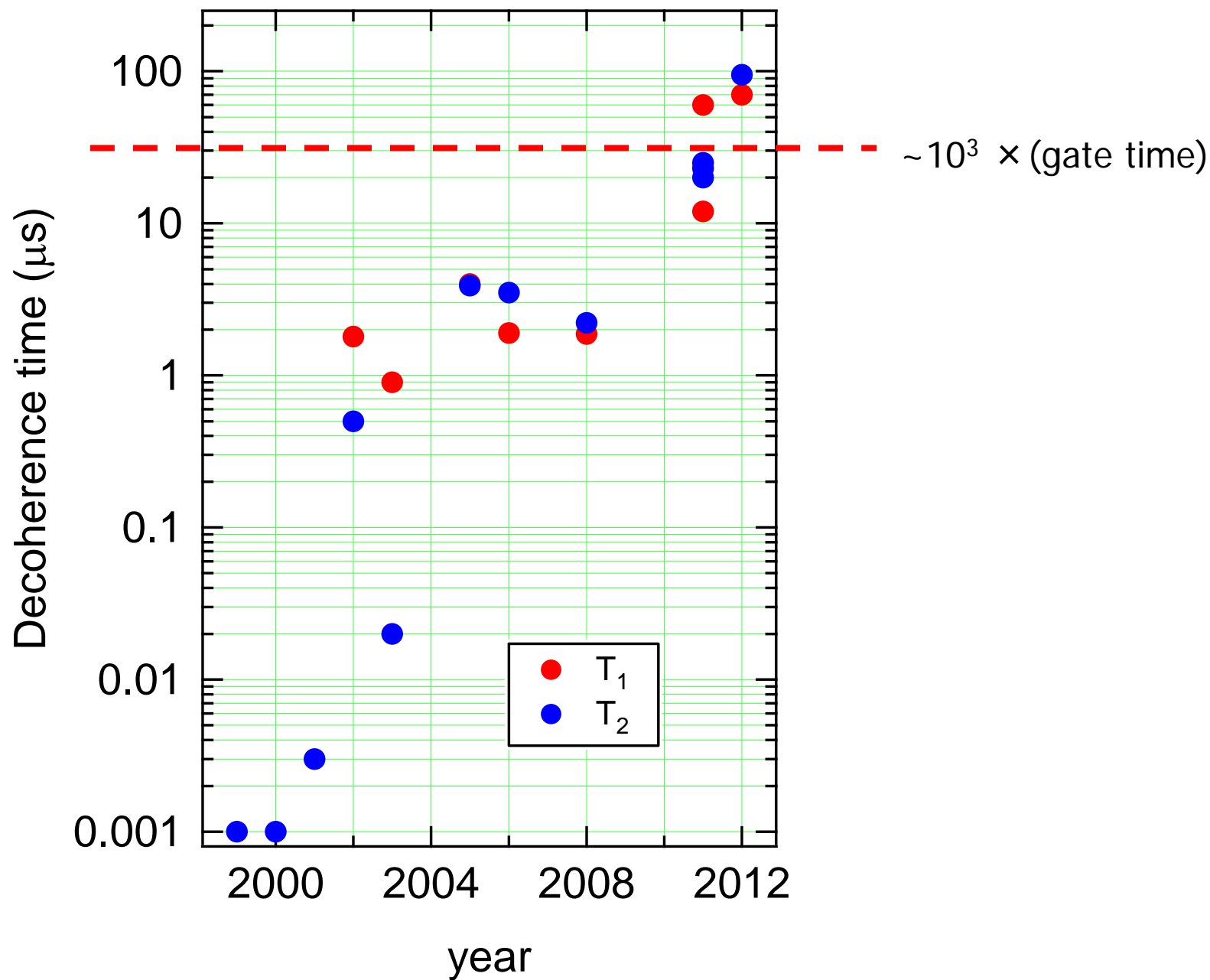




# Noise spectrum



# Decoherence time of superconducting qubits



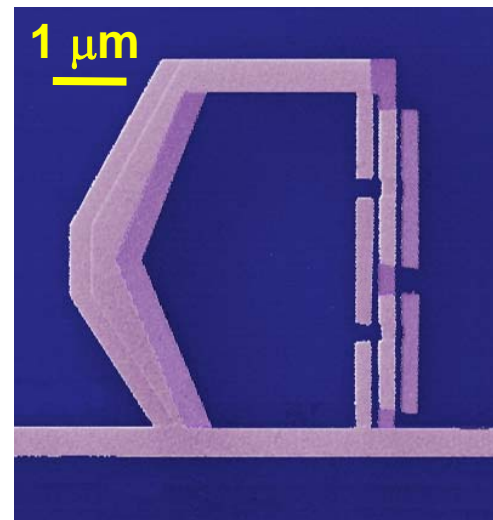
# Atoms

(a stereotype of) atom

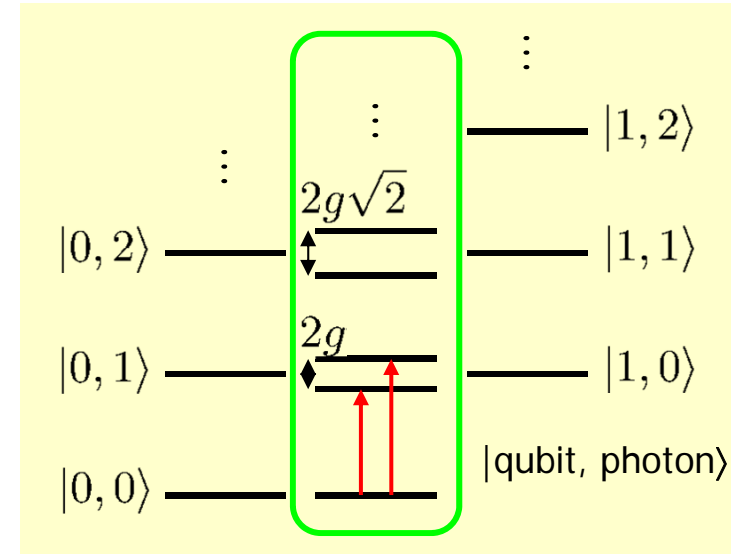
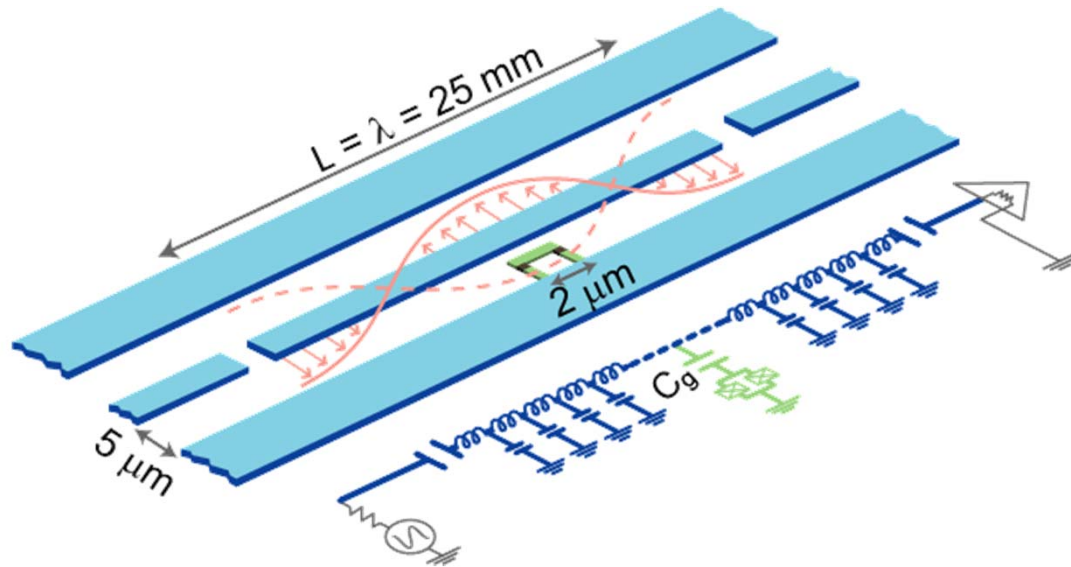


$\sim \text{\AA}$

Our artificial atom



# Circuit quantum electrodynamics (circuit QED)



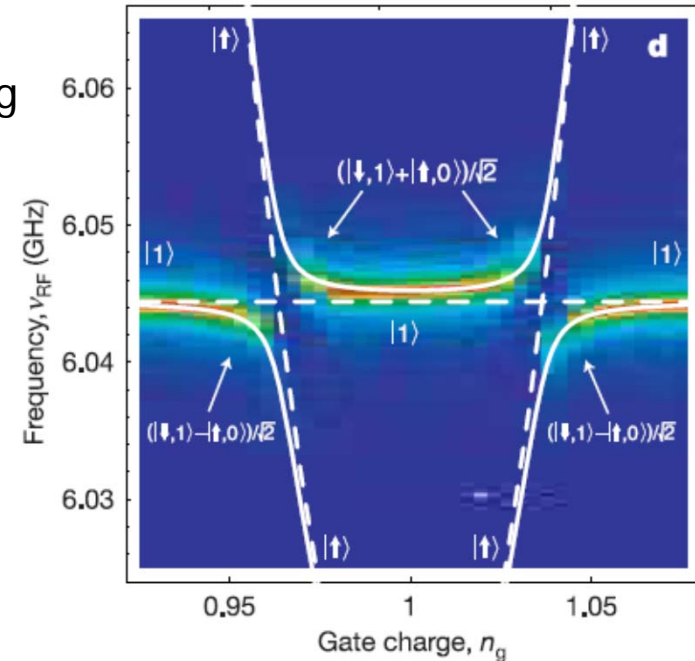
vacuum Rabi splitting

Jaynes-Cummings Hamiltonian:

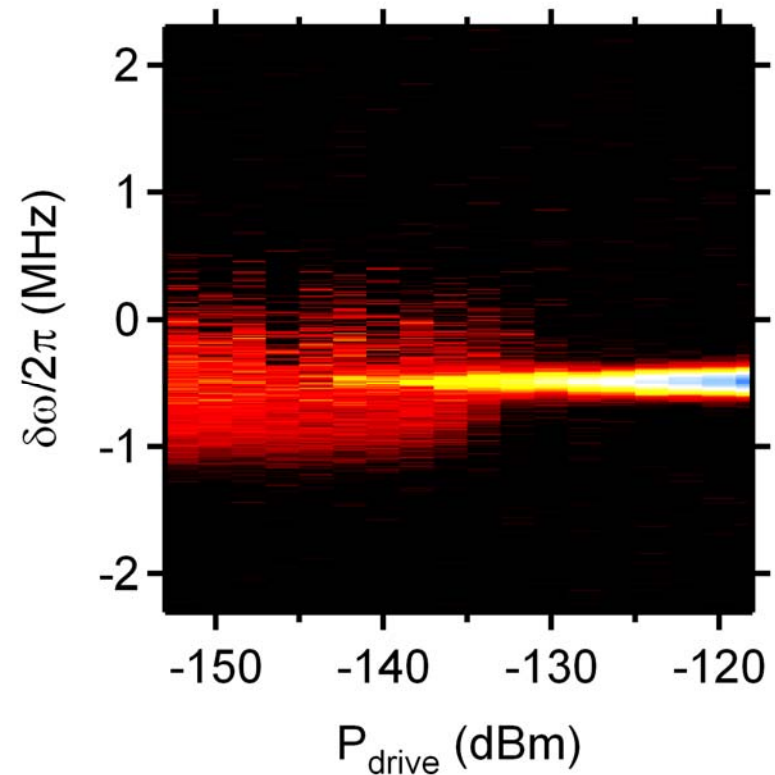
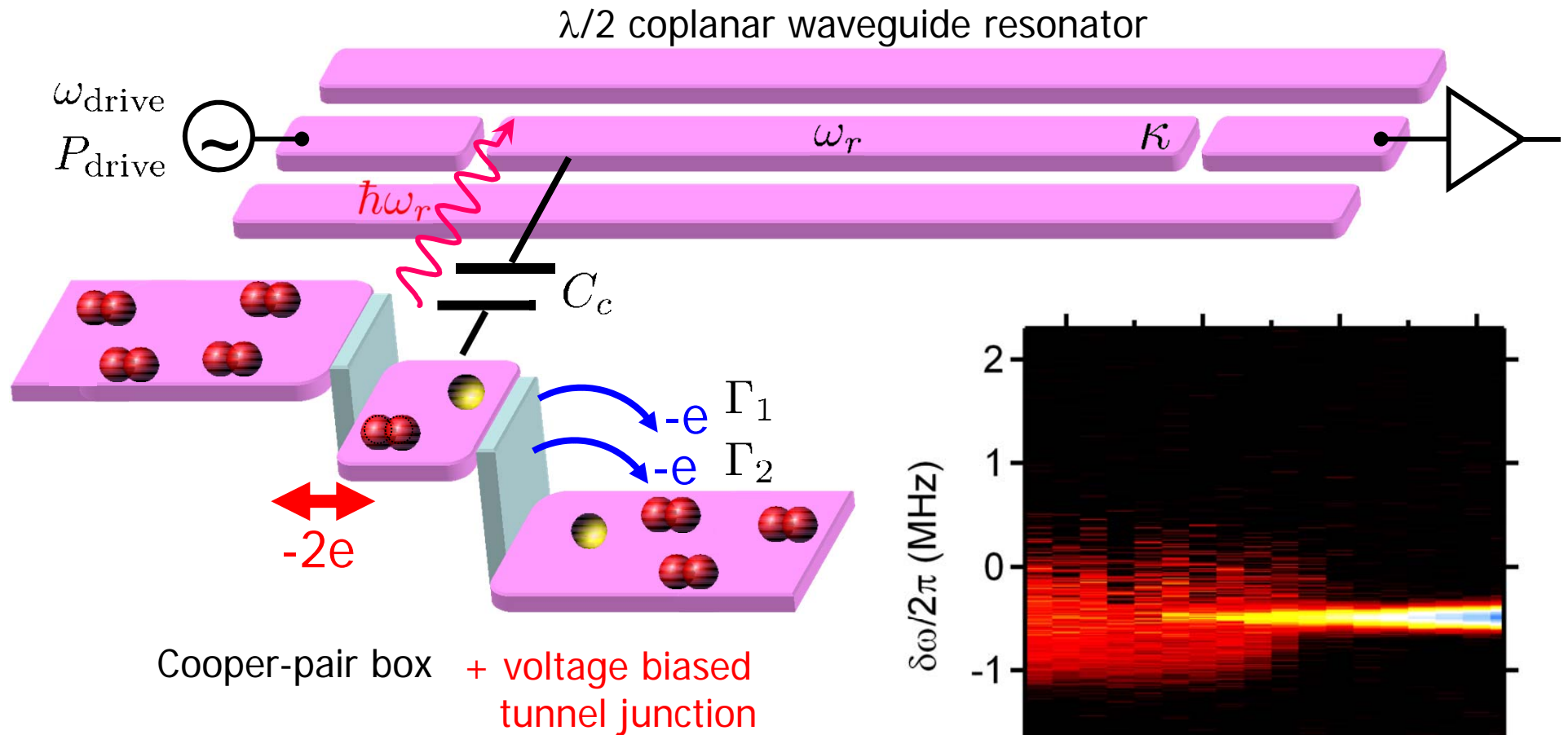
$$H = \hbar\omega_r \left( a^\dagger a + \frac{1}{2} \right) - \frac{\hbar\omega_q}{2} \sigma_z + \hbar g (a^\dagger \sigma_- + a \sigma_+)$$

Strong coupling:  $g \gg \gamma, \kappa$

$$g/2\pi \sim 10 \text{ MHz} \quad (\gamma, \kappa)/2\pi \sim 1 \text{ MHz}$$



# Single artificial-atom maser



- Population inversion generated by current injection
- Capacitive coupling with cavity mode

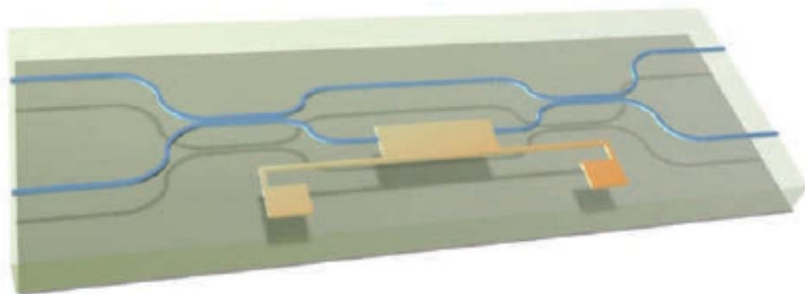
# Photons in transmission line

## Optics

Frequency 100-1000 THz  
Wavelength 3  $\mu\text{m}$  – 300 nm

Optical fiber, low loss  $\sim 0.2$  dB/km  
Photonic on-chip circuits,  $\sim 0.1$  dB/cm

Weak nonlinearity

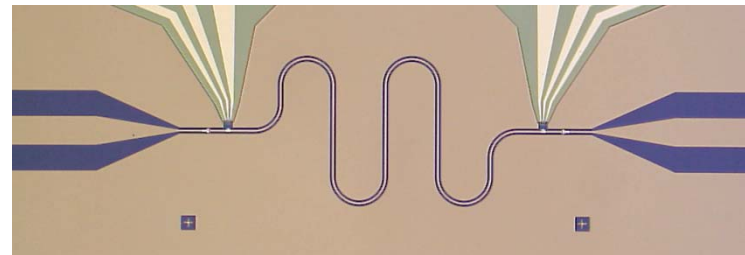


## Microwave

Frequency 1-10 GHz  
Wavelength 30 cm – 3 cm

Microwave on-chip circuits,  $\sim 0.3$  dB/km(?)

Strong nonlinearity available



coplanar waveguide





# Superconducting transmission line

$$Z = \sqrt{l/c}$$

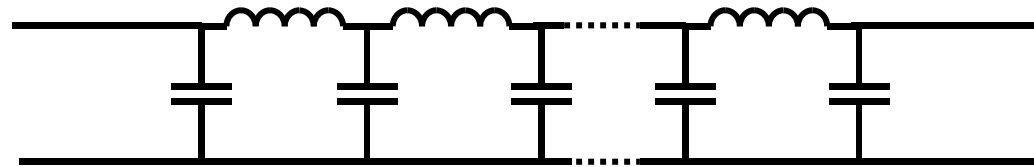
$$v = 1/\sqrt{lc}$$

$$\omega = vk$$



distributed element

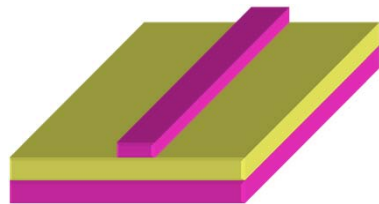
↕  
lumped element



- small dissipation for  $\hbar\omega_r, k_B T \ll \Delta$
- 1D transmission mode
- Photon life time  $\sim 100 \mu\text{s}$ , 10 km (?)

Variety of transmission lines:

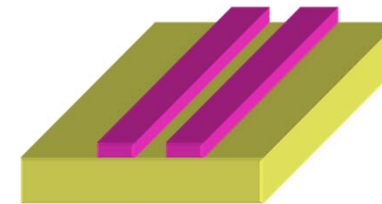
microstrip line



coplanar waveguide



slot line



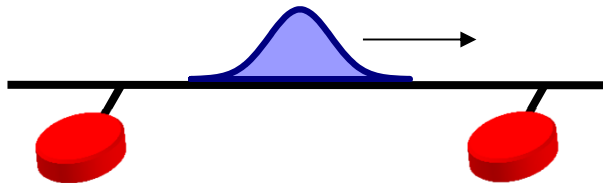
# Confined photon and flying photon

- in resonator (confined “0D” photon; single-mode)



$$H = \hbar\omega_r a^\dagger a$$

- through transmission line (flying photon; multi-mode; continuum)

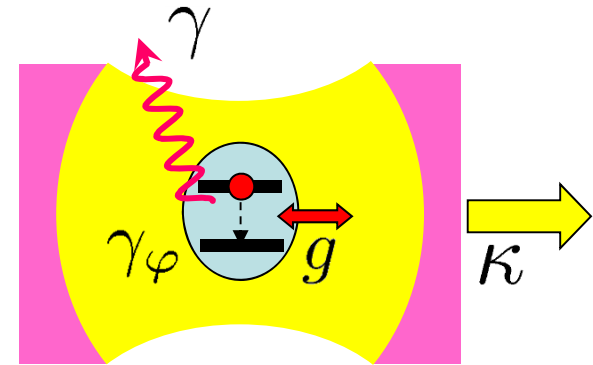


$$H = \hbar \int_0^\infty d\omega \omega a^\dagger(\omega) a(\omega)$$

# Atom-photon strong coupling

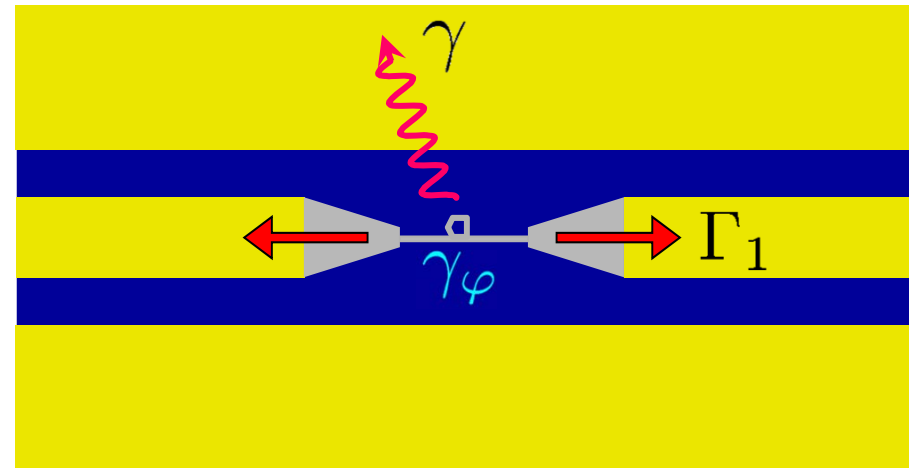
Strong coupling in cavity QED

$$g \gg \kappa, \gamma, \gamma_\varphi$$



“Strong coupling” in 1D waveguide

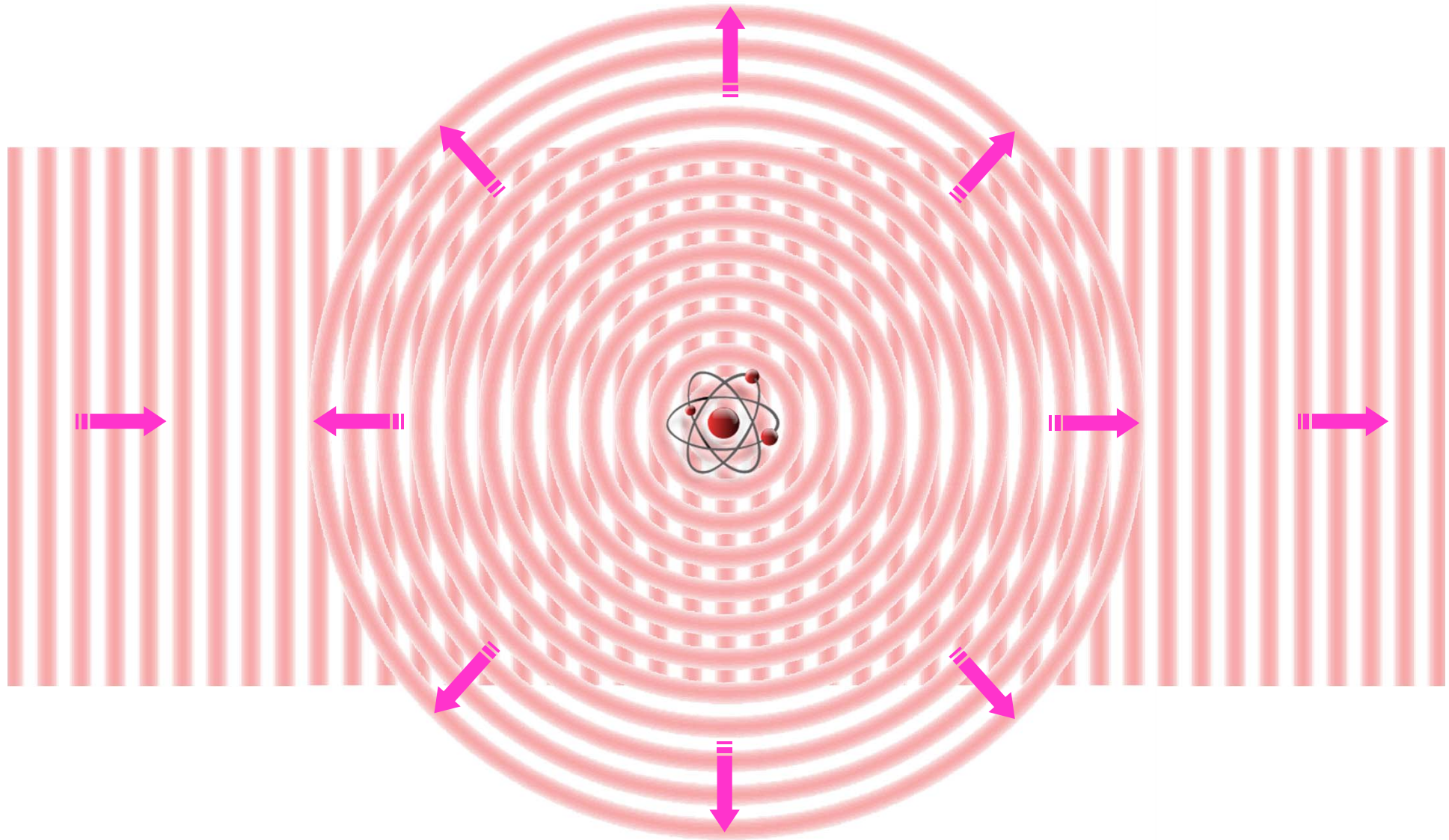
$$\Gamma_1 \gg \gamma, \gamma_\varphi$$



# Superconducting qubits coupled to a transmission line

- Superconducting qubits as artificial atoms
  - Fixed on chip
  - Strong coupling
  - Multi levels, selection rules
- Beauty of 1D
  - Microwave transmission line as 1D channel
  - Perfect spatial mode matching
- Use of interference
  - Importance of temporal modes
  - Limitation with bandwidth
- Spontaneous emission – coherent process

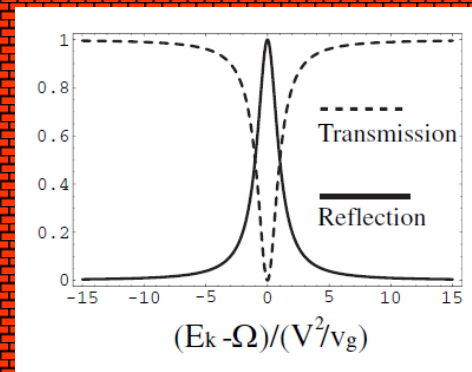
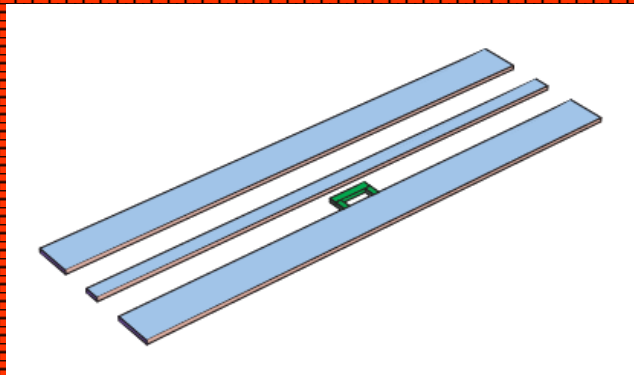
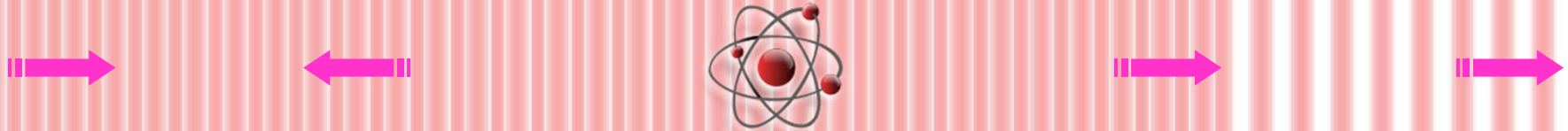
# Resonant scattering in 3D space



- Small scattering cross section
- Spatial mode mismatch between incident and radiated waves

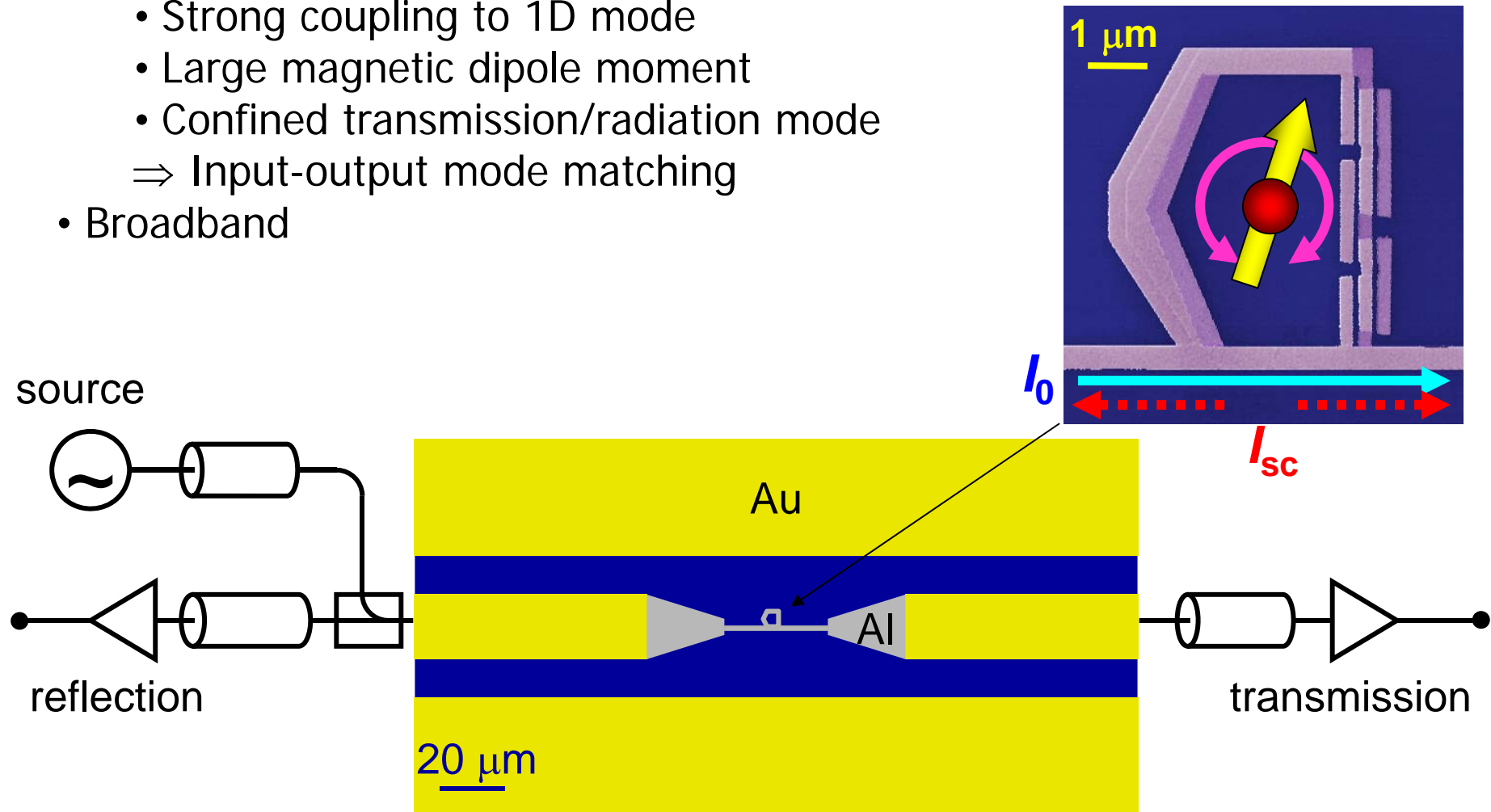
# Resonant scattering in 1D waveguide

Destructive interference of transmitted wave  
⇒ Extinction of transmittance  
⇒ Perfect reflection



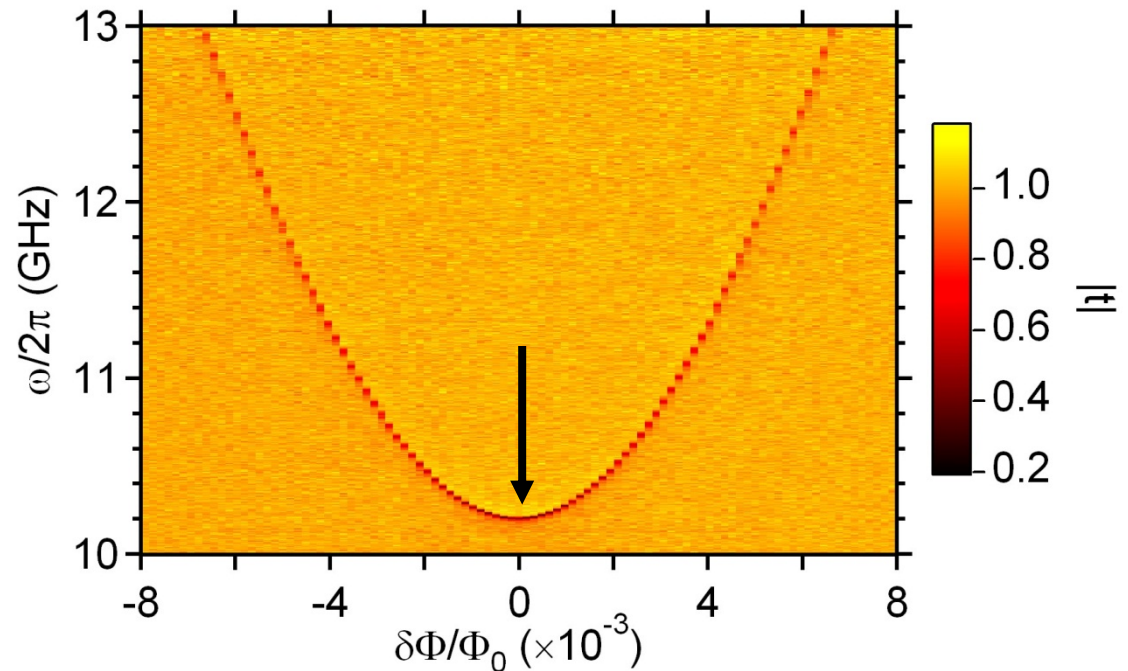
# Artificial atom in 1D open space

- Flux qubit coupled to transmission line via kinetic inductance
  - Strong coupling to 1D mode
  - Large magnetic dipole moment
  - Confined transmission/radiation mode $\Rightarrow$  Input-output mode matching
- Broadband



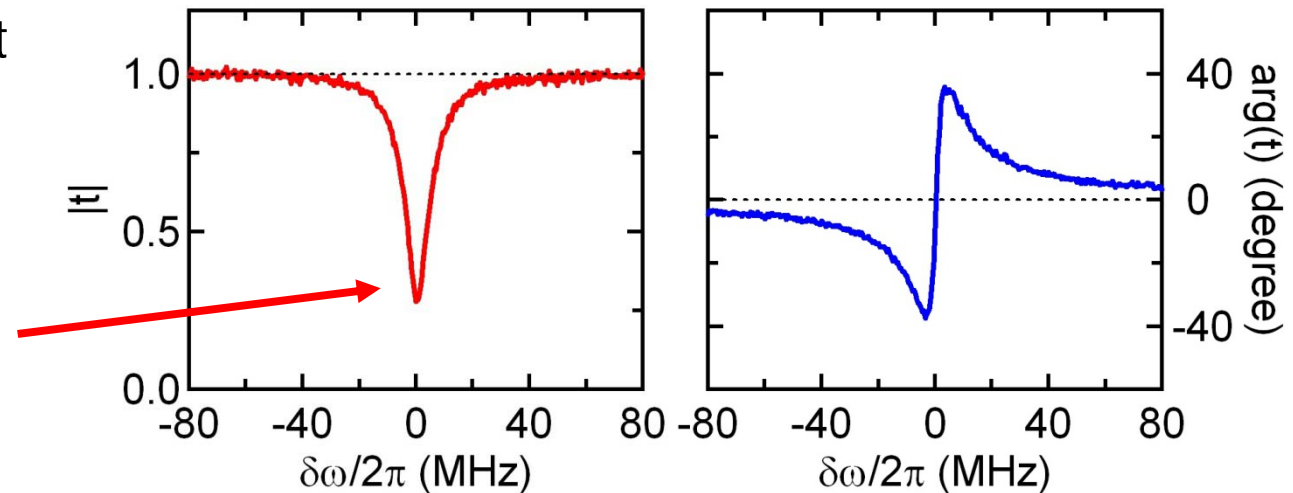
# Transmission spectroscopy — elastic scattering

$$\Delta/h = 10.20 \text{ GHz}$$
$$I_p = 195 \text{ nA}$$



At degeneracy point  
 $\delta\Phi/\Phi_0 = 0$

Strong extinction  
of transmitted wave

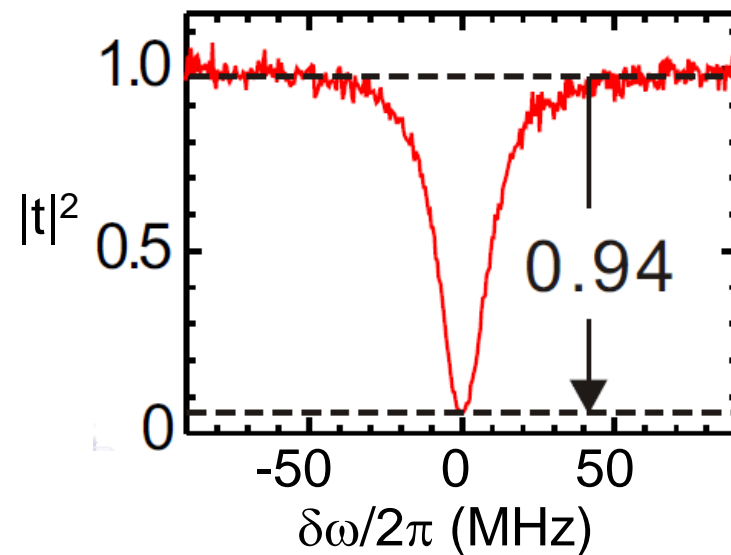
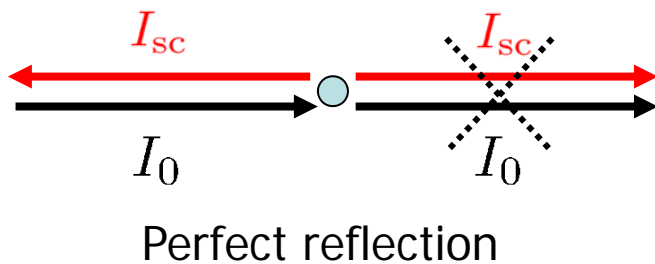
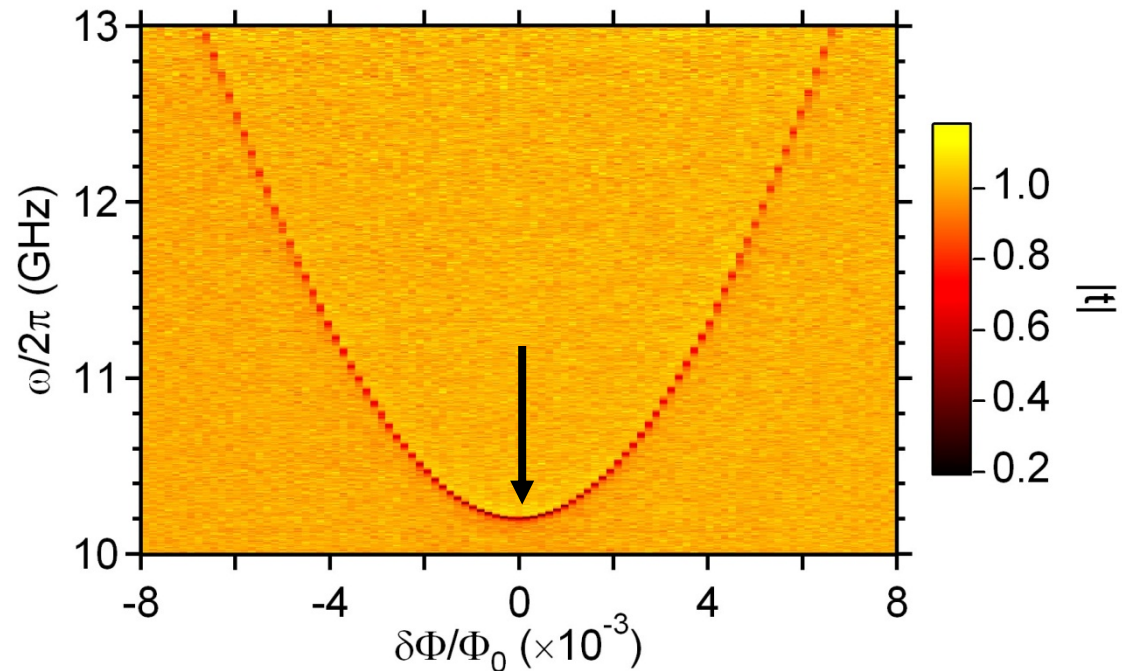




# Transmission spectroscopy — elastic scattering

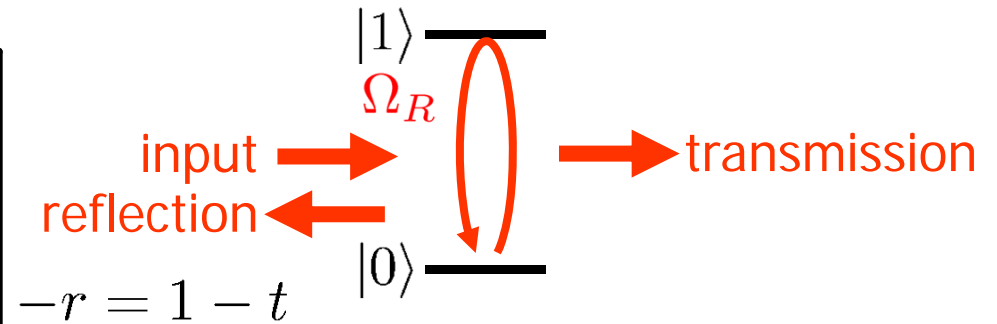
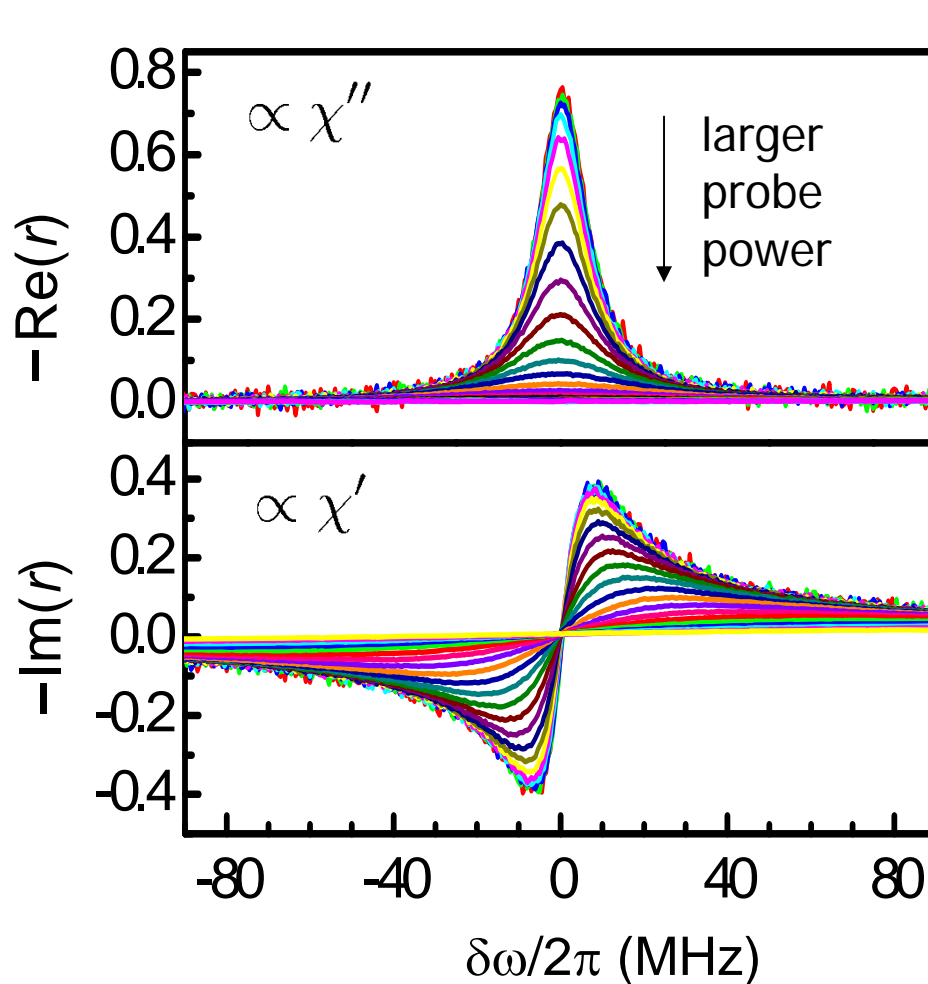
$$\Delta/h = 10.20 \text{ GHz}$$

$$I_p = 195 \text{ nA}$$



$$|r_0|^2 = \left( \frac{\Gamma_1}{2\Gamma_2} \right)^2$$

# Power dependence — saturation of atom



$$r = -r_0 \frac{1 + i\delta\omega/\Gamma_2}{1 + (\delta\omega/\Gamma_2)^2 + \Omega_R^2/\Gamma_1\Gamma_2}$$

$$r_0 = \frac{\Gamma_1}{2\Gamma_2} = \frac{\Gamma_1}{\Gamma_1 + 2\Gamma_\varphi}$$

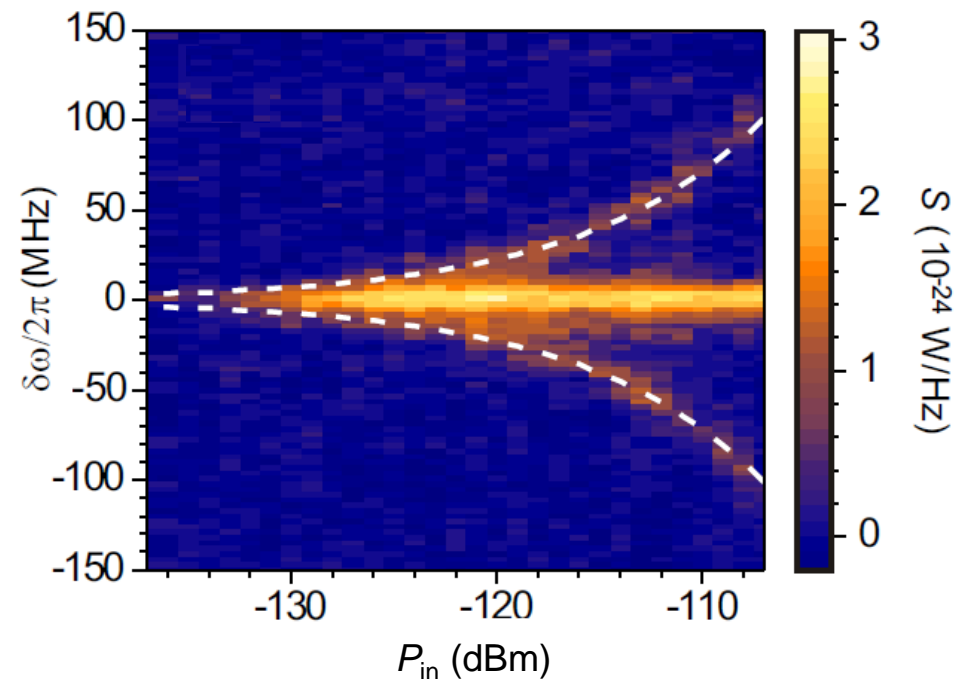
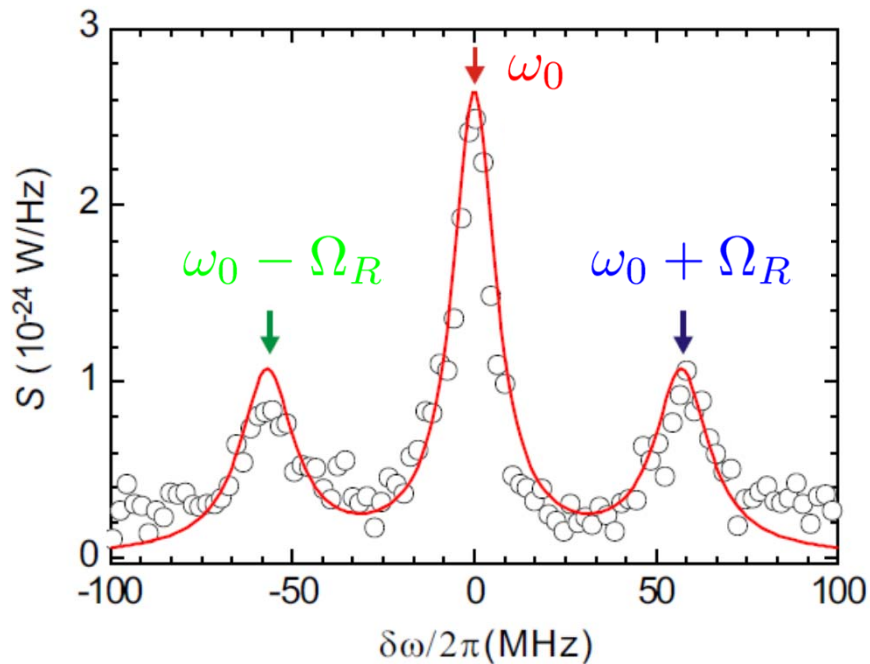
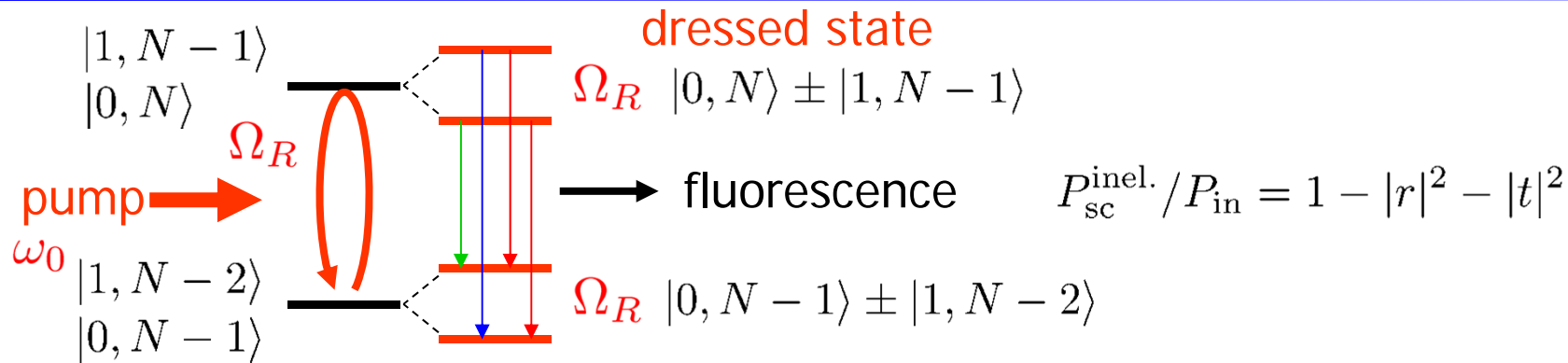
$$\Gamma_1/2\pi = 11.0 \text{ MHz}$$

$$\Gamma_\varphi/2\pi = 1.7 \text{ MHz}$$

$$M = 12.2 \text{ pH}$$

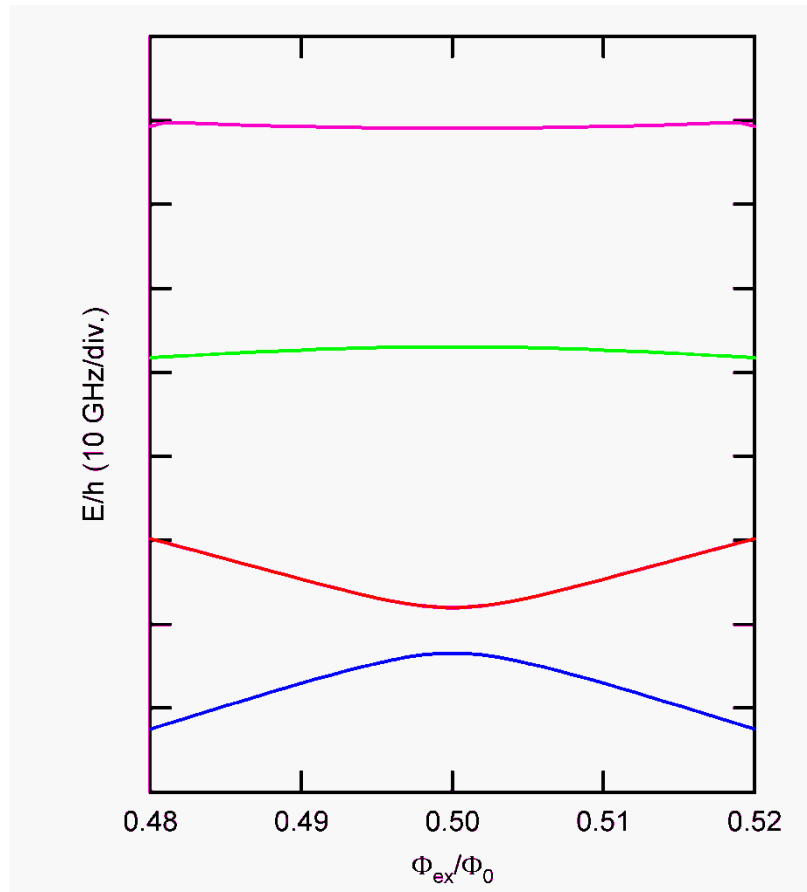
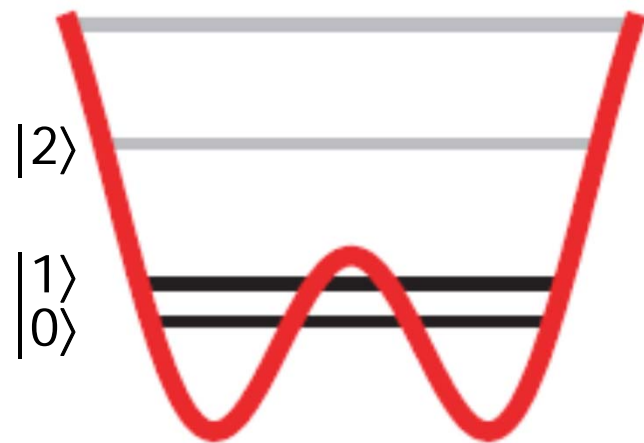
Inherent nonlinearity of the two-level atom

# Resonance fluorescence: inelastic scattering



Mollow triplet: 
$$S(\omega) \approx \frac{1}{2\pi} \frac{\hbar\omega\Gamma_1}{8} \left( \frac{\gamma_s}{(\delta\omega + \Omega)^2 + \gamma_s^2} + \frac{2\gamma_c}{\delta\omega^2 + \gamma_c^2} + \frac{\gamma_s}{(\delta\omega - \Omega)^2 + \gamma_s^2} \right)$$

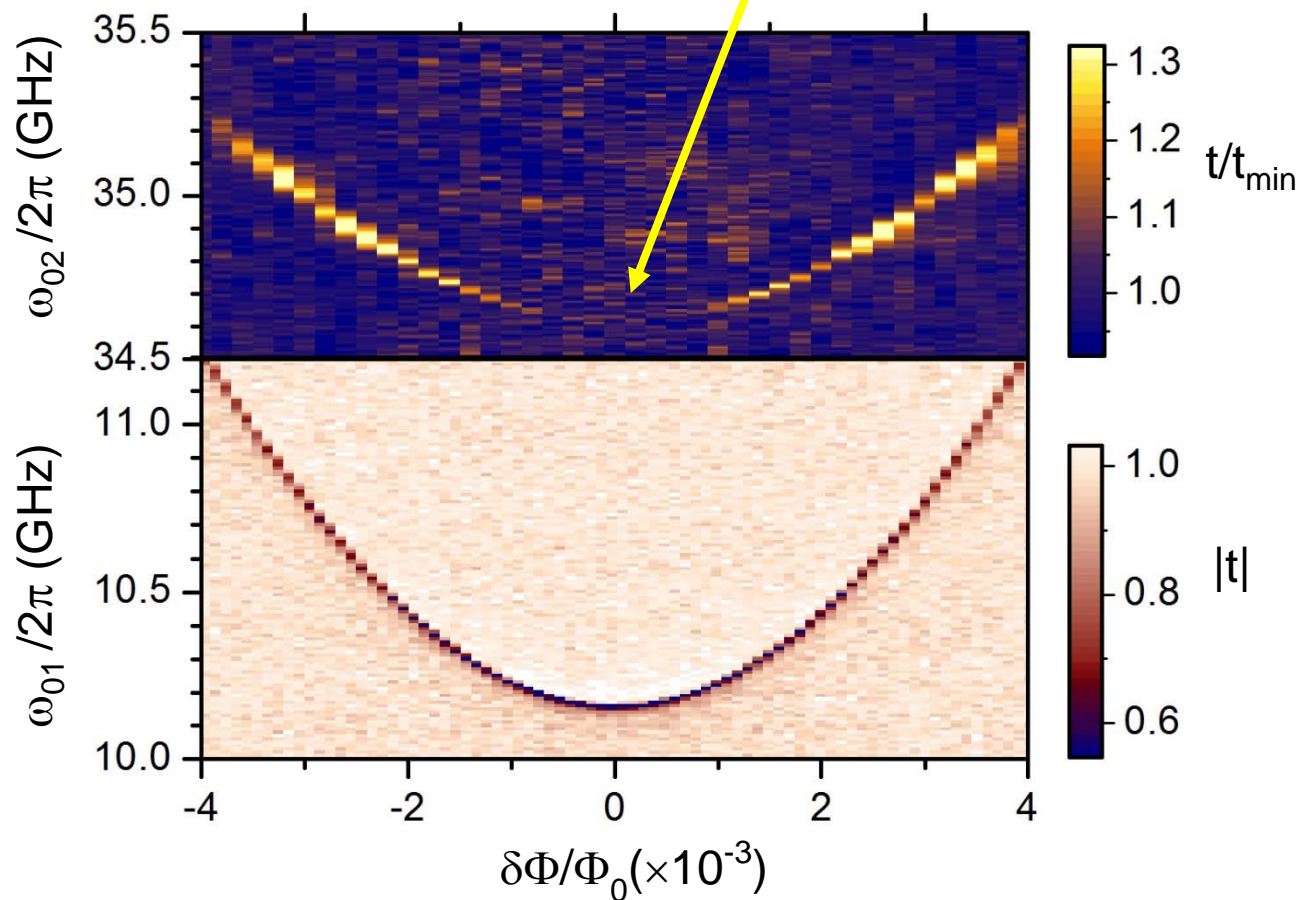
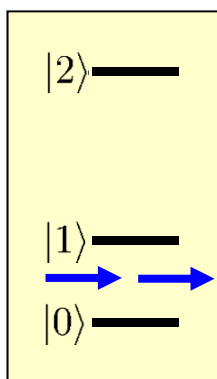
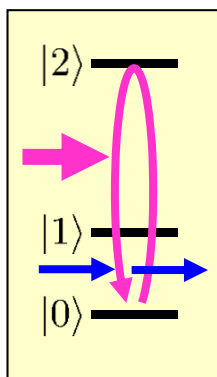
# Flux qubit as a three-level artificial atom



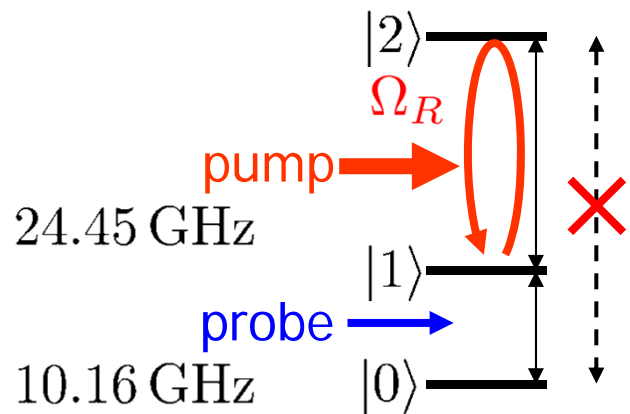
- Josephson junction qubits = **effective** two-level system
- presence of auxiliary states
- large anharmonicity/nonlinearity
- selection rule due to symmetry when flux bias  $\delta\Phi=0$

# Spectroscopy of a three-level atom

suppressed excitation due to selection rule



# Ladder system at degeneracy point: induced transparency

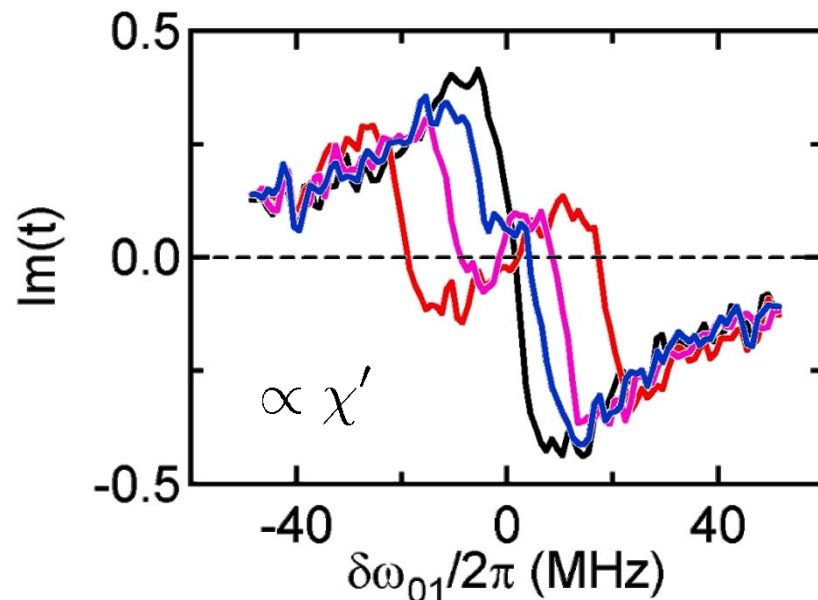
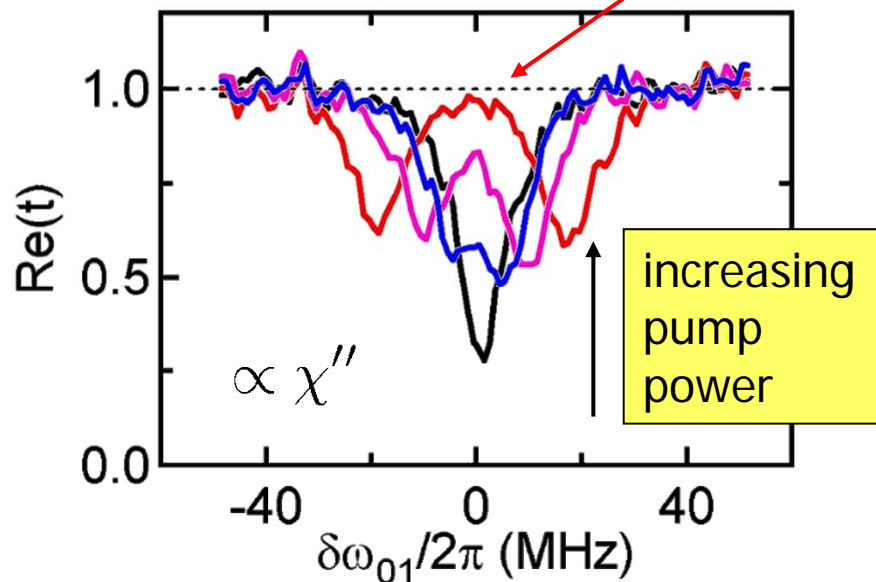


Biased at degeneracy point  
Transition  $0 \leftrightarrow 2$  not allowed

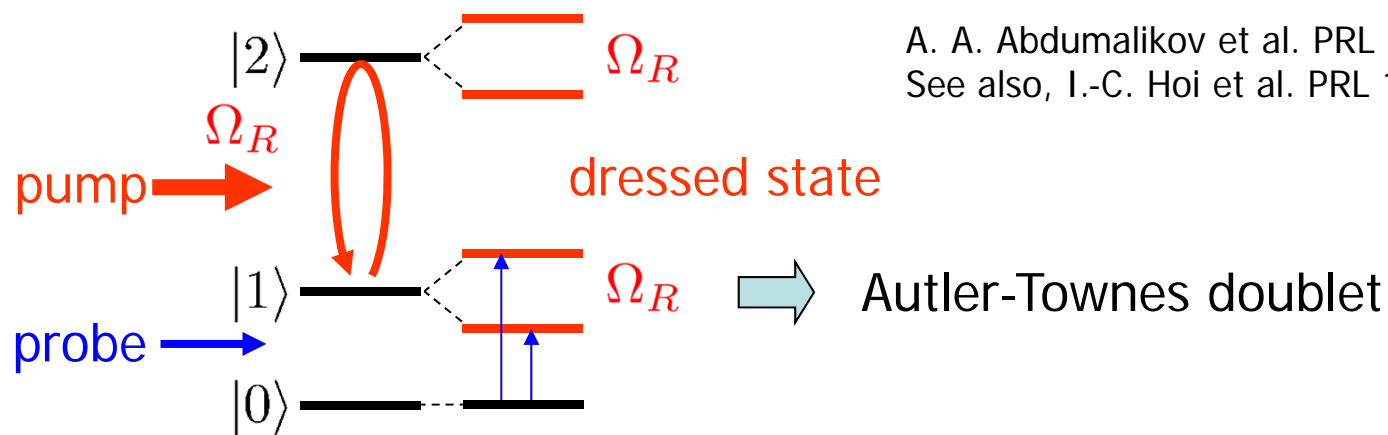
ladder-type

“Electromagnetically-induced transparency”

$$\Gamma_{10} < \Gamma_{21} < \Omega_R$$

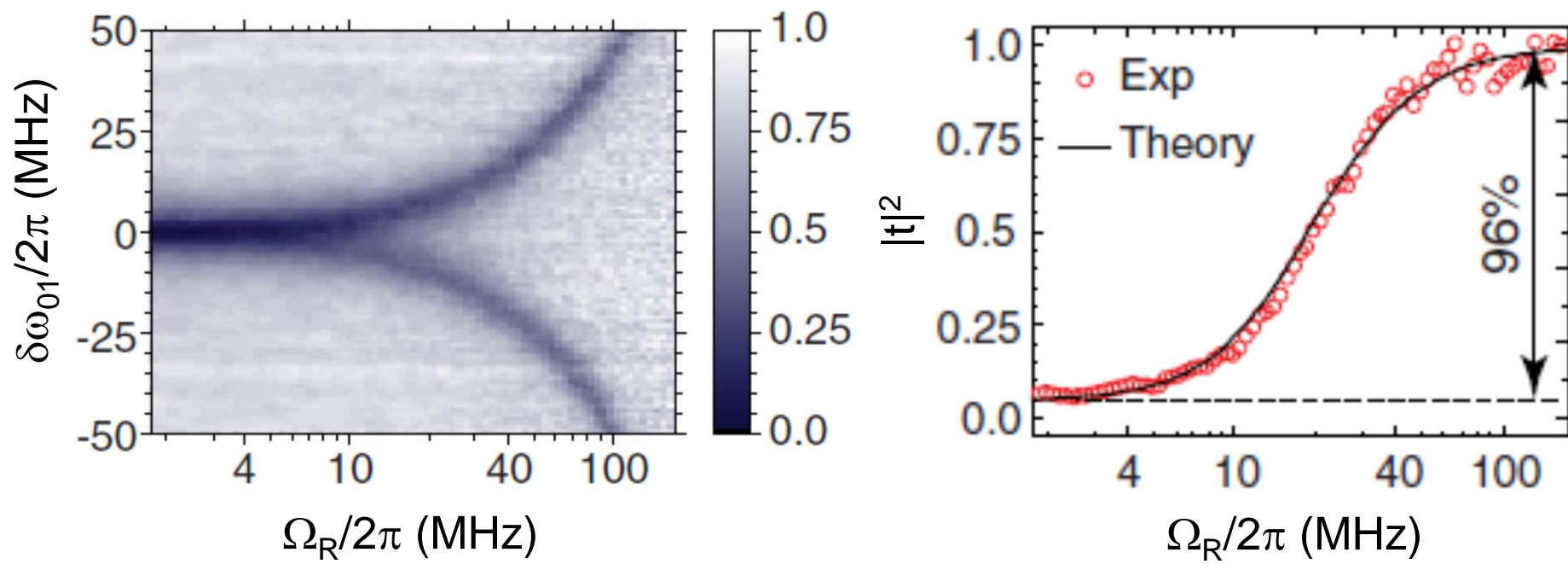


# Ladder system at degeneracy point: induced transparency



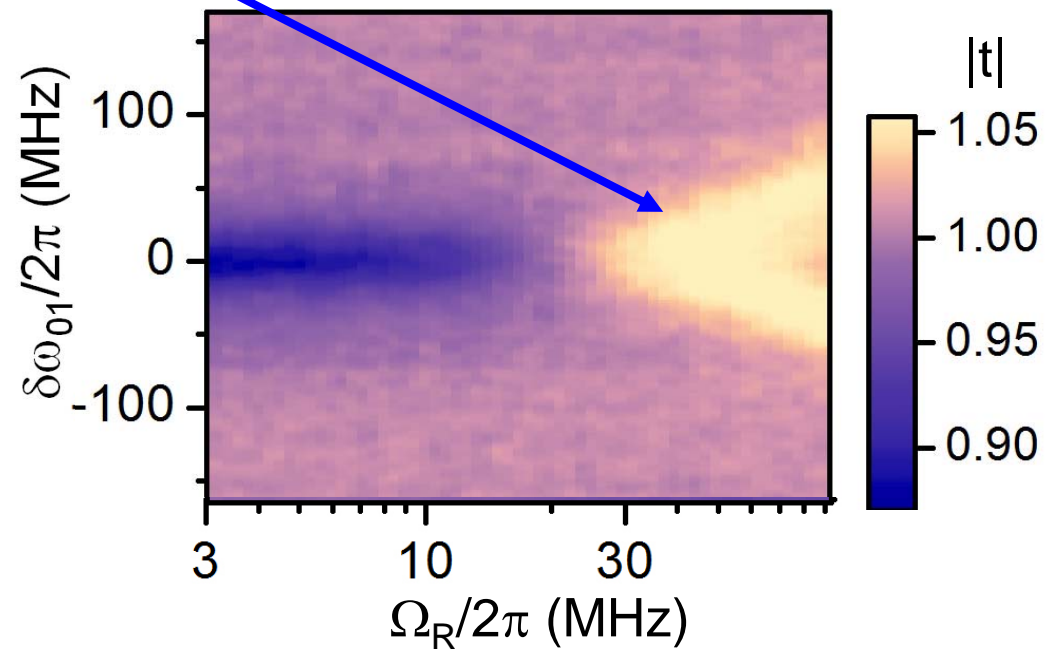
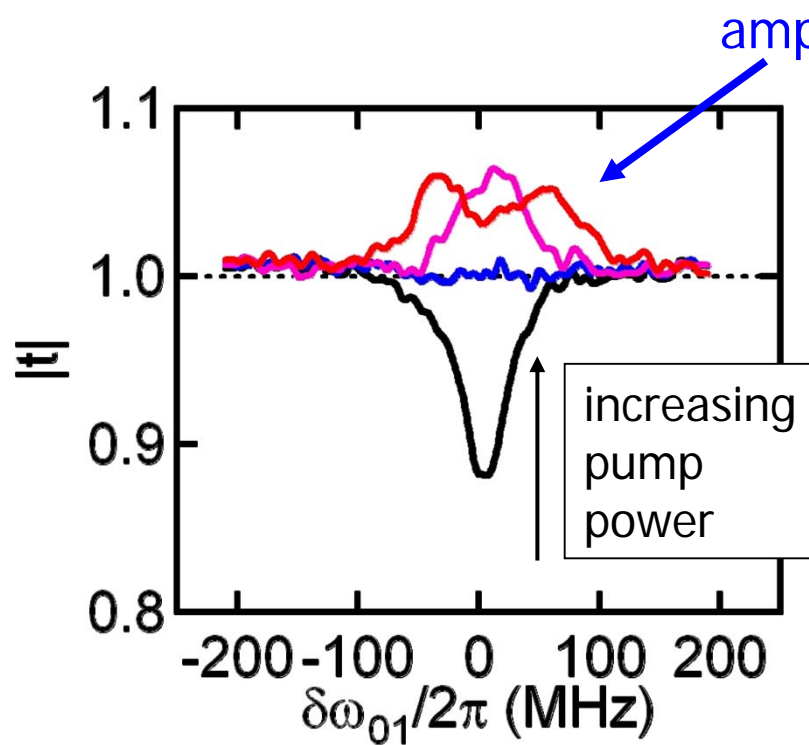
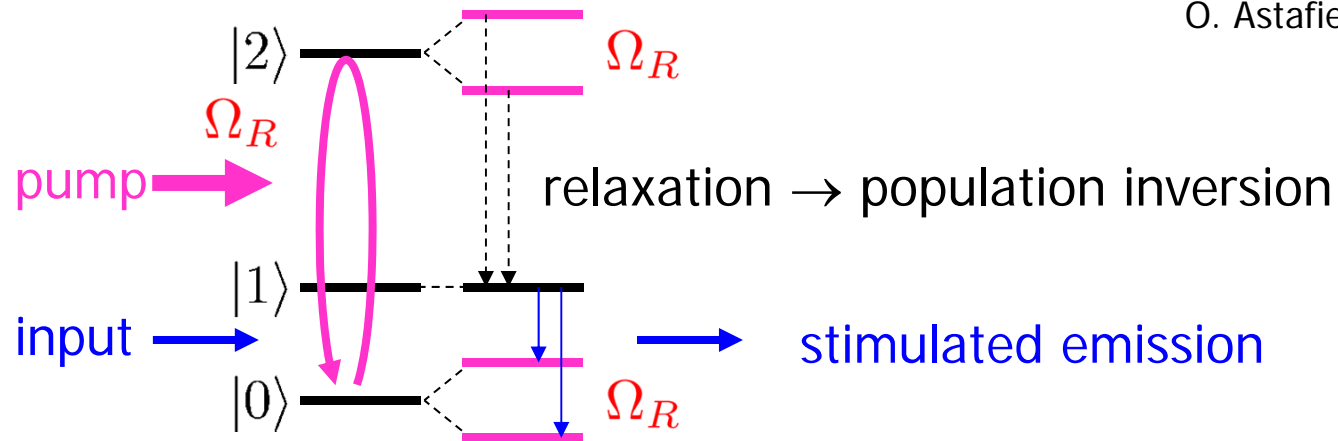
A. A. Abdumalikov et al. PRL 104, 193601 (2010)  
 See also, I.-C. Hoi et al. PRL 107, 073601 (2011) (Chalmers)

Transmission of probe signal



# Stimulated emission and amplification

O. Astafiev et al. PRL 104, 183603 (2010)





# Summary

- Superconducting qubits as artificial atoms
  - Electrical circuits fixed on chip
  - Gigantic dipole, strong coupling with EM modes
  - Multiple levels, selection rules
  - Qubit as quantum spectrum analyzer
- Coupling to 1D channel
  - Microwave transmission line as 1D channel
  - Perfect spatial mode matching
  - Interference between transmitted and scattered fields
  - Design and control of modes
- Future: quantum-optics tools in microwave domain
  - Single photon source/detectors
  - Squeezed state generators
  - Parametric amplifiers