

Asymmetric Neutrino Emission from Strongly Magnetized Neutron Stars

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§ 1. Introduction

High Density Matter Study \Rightarrow Exotic Phases inside Neutron Stars

Strange Matter, Ferromagnetism, Meson Condensation, Quark matter

Observable Information ···· Neutrino Emissions

S.Reddy, et al., PRD58 #013009 (1998) **Influence from Hyperons Λ , Σ**

Magnetar 10^{15} G in surface $10^{17\text{-}19}$ G inside (?) \rightarrow Large Asymmetry of ν ?

P. Arras and D. Lai, PRD60, #043001 (1999), S. Ando, P.R.D68 #063002 (2003)
(Non-Relativistic)

Our Works \Rightarrow Neutrino Reactions on High Density Matter wih Strong Mag. Fields
 in the Relativistic Mean-Field (RMF) Approach

TM et al., PRD83, 081303(R) (2011)

Application to Phenomena related with ν — Asymmetric Emission

Pulsar Kick

in Poloidal Magnetic Field

Spin Deceleration

in Toroidal Magnetic Field

§ 2. Formulation

Magnetic Field : $\vec{B} = B\hat{z}$.

Lagrangian : $\mathcal{L} = \mathcal{L}_{RMF} + \mathcal{L}_{lep.} + \mathcal{L}_{mag} + \mathcal{L}_{int}$

Baryon

Lepton

B & L – Mag.

1. Proto-Nuetron-Star (PNS) Matter without Mag. Field
2. Baryon Wave Function under Mag. Field in Perturbative Way
3. Cross-Sections for ν reactions

Weak Interaction

$$\mathcal{L}_{int} = G_F \{\bar{\psi}_l \gamma_\mu (1 - \gamma_5) \psi_l\} \{\bar{\psi}_B \gamma^\mu (c_V - c_A \gamma_5) \psi_B\}$$

$\nu_e + B \rightarrow \nu_e + B$: scattering

$\nu_e + B \rightarrow e^- + B'$: absorption

§ 2-1 EOS of Proto Neutron-Star-Matter in RMF $N, \Lambda, \sigma, \omega, \rho$

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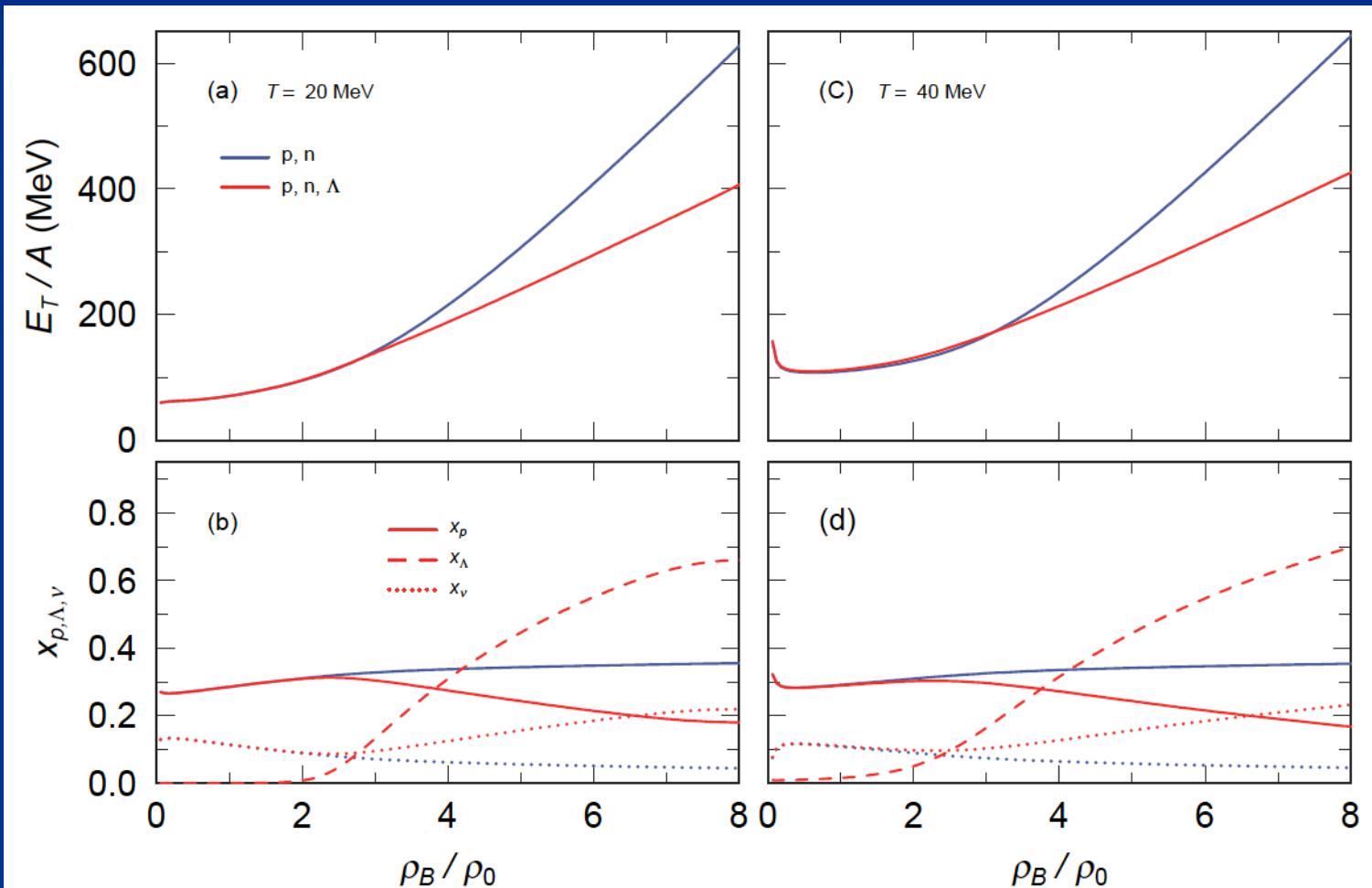
PM1-L1

$BE = 16 \text{ MeV}$, $M_N^*/M_N = 0.7$, $K = 200 \text{ MeV}$ at $\rho_0 = 0.17 \text{ fm}^{-3}$

$$g_{\sigma,\omega}^\Lambda = \frac{2}{3} g_{\sigma,\omega} \quad \text{SU(3)}$$

T.M. et al.
PTP. 102, p809
(1999)

Charge Neutral ($\rho_p = \rho_e$) & Lepton Fraction : $Y_L = 0.4$



§ 2-2 Dirac Equation under Magnetic Fields

$\mu_N B \ll \epsilon_N$ (Chem. Pot) $\rightarrow B$ can be treated perturbatively

$$B \sim 10^{17} \text{ G}$$

Landau Level can be ignored

Lagrangian

$$\mathcal{L}_{mag} \approx - \sum_b \mu_n B \bar{\psi}_b \sigma_z \psi_b$$

Dirac Eq.

$$\hat{K}(p) u(p) = [\gamma_\mu p^\mu - M^* - \mu B \sigma_z] u(p) = 0$$

Single Part. Eng.

$$e(\mathbf{p}, s) = \left[\left(\sqrt{p_x^2 + M^{*2}} + s\mu B \right)^2 + p_z^2 \right]^{\frac{1}{2}} \approx E_p^* + s\mu B \frac{\sqrt{p_T^2 + M^{*2}}}{E_p^*}$$

Dirac Spinor

$$u(\mathbf{p}, s) \bar{u}(\mathbf{p}, s) \approx \frac{(\not{p} + M)(1 + \gamma_5 \not{d}(p)s)}{4E_p^*}$$

Spin Vector

$$a_z = \frac{E_p}{\sqrt{p_T^2 + M^2}} \quad \mathbf{a}_T = 0 \quad a_0 = \frac{p_z}{\sqrt{p_T^2 + M^2}}$$

The Cross-Section of Lepton-Baryon Scattering

$$\frac{d^2\sigma}{dk'd\Omega'_k} = \frac{G_F^2}{8\pi^2} k'^2 \sum_{s_i, s_f} \int \frac{d^3 p}{(2\pi)^3} \tilde{W}_{BL}(2\pi) \delta(|\mathbf{k}| - |\mathbf{k}'| + e_i(\mathbf{p}) - e_f(\mathbf{k} + \mathbf{p} - \mathbf{k}')) \\ \times [1 - f_l'(\mathbf{k}') n_B(e_i)] [1 - n_{B'}(e_f)]$$

$$\tilde{W}_{BL} = \text{Tr} \left\{ \frac{(\not{k}' + m_f)(1 + \gamma_5 \not{d}_{l'})}{4|\mathbf{k}'|} \gamma^\mu (1 - \gamma_5) \frac{\not{k}'}{2|\mathbf{k}|} \gamma^\nu (1 - \gamma_5) \right\} \\ \times \text{Tr} \left\{ \frac{(\not{p}' + M_f^*)(1 + \gamma_5 \not{d}_f(p'))}{4E_f^*(\mathbf{p}')} \gamma_\mu (c_V - c_A \gamma_5) \frac{(\not{p} + M_i^*)(1 + \gamma_5 \not{d}_i(p))}{4E_i^*(\mathbf{p})} \gamma_\nu (c_V - c_A \gamma_5) \right\}$$

$$m_f = 0 \quad \text{when } l_f = \nu \quad \quad m_f = m_e \quad \text{when } l_f = e$$

Fermi
Distribution

$$n(e(\mathbf{p}), s) \approx n(\varepsilon(\mathbf{p}, s)) + n'(\varepsilon(\mathbf{p}, s)) \frac{\sqrt{p_T^2 + M^{*2}}}{E_p^*} \mu B s.$$

Deformed Distribution

Perturbative
Treatment

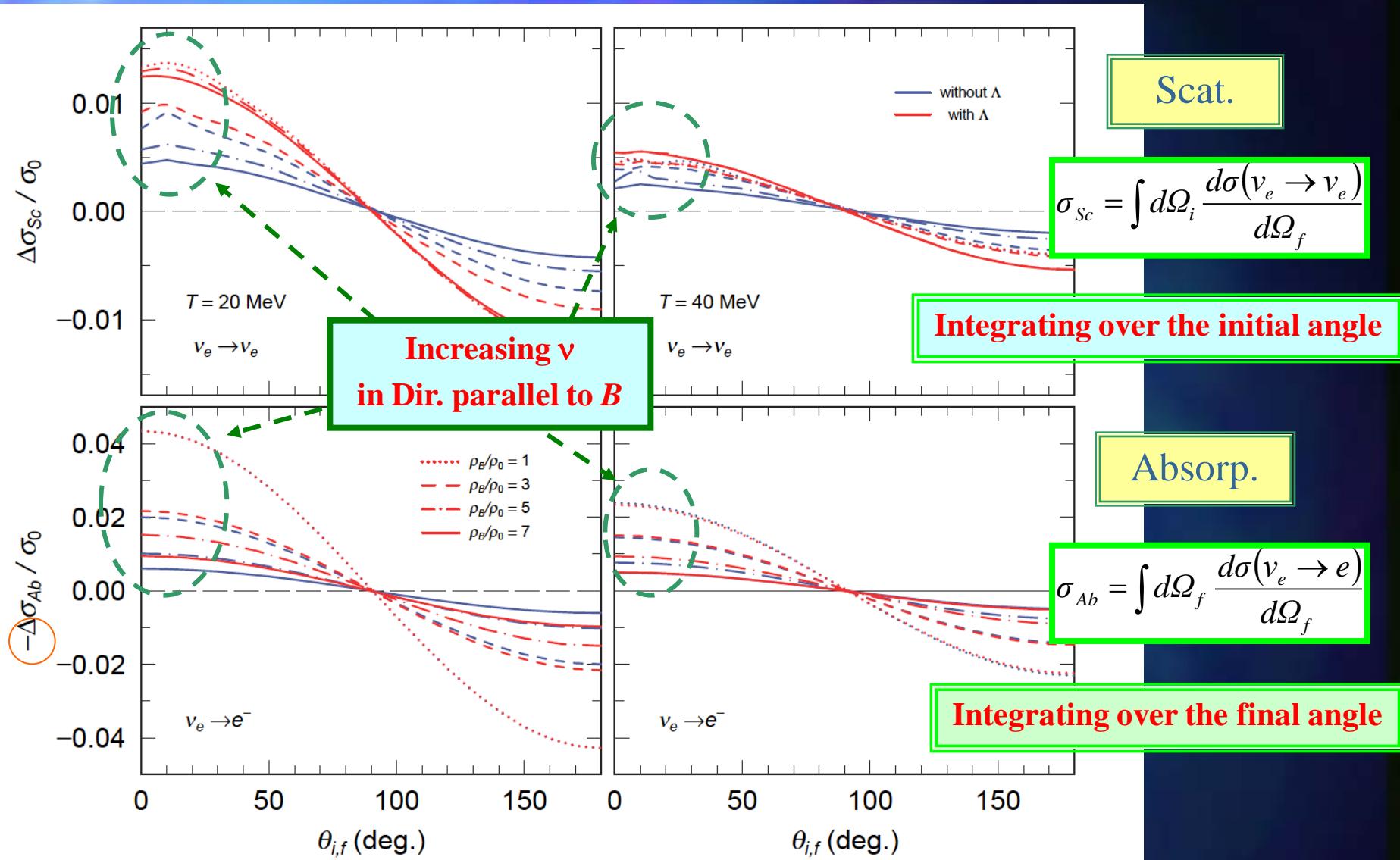
$$\sigma = \sigma_0 + \Delta\sigma \quad \Delta\sigma \propto B$$

Non-Magnetic Part

Magnetic Part

§ 2-3 Magnetic parts of Cross-Sections

$$\sigma = \sigma_0 + \Delta\sigma \quad \Delta\sigma \propto B$$



$$k_i = \varepsilon_\nu \text{ (neutrino chem. pot.)}, \quad B = 2 \times 10^{17} \text{ G} \quad \text{and} \quad \theta_i = 0^\circ$$

§ 2-4 Neutrino Transportation

Neutrino Phase Space Distribution Function

$$f(\mathbf{p}, \mathbf{r}) \approx f_0(\mathbf{p}, \mathbf{r}) + \Delta f(\mathbf{p}, \mathbf{r}), \quad f_0(\mathbf{p}, \mathbf{r}) = 1 / \{1 + \exp[(-|\mathbf{p}| - \varepsilon_\nu) / T]\}$$

Equib. Part

Non-Equib. Part

Neutrino Propagation \Rightarrow Boltzmann Eq.

$$c \frac{\partial}{\partial \mathbf{r}} f_0(\mathbf{p}, \mathbf{r}) \approx c \frac{\partial}{\partial \mathbf{r}} f_0(\mathbf{p}, \mathbf{r}) + c \frac{\partial}{\partial \mathbf{r}} \Delta f(\mathbf{p}, \mathbf{r}) = I_{coll} \approx -cb_\nu \Delta f(\mathbf{p}, \mathbf{r}), \quad b_\nu = \frac{\sigma_{ab}}{V}$$

Neutrinos Propagate on

Strait Line

only absorption

Solution \Rightarrow

$$\Delta f(\mathbf{p}, \mathbf{r}_T, z) = \int_0^z dx \left[-\frac{\partial}{\partial x} f_0(\mathbf{p}, \mathbf{r}_T, x) \right] \exp \left[- \int_x^z dy b_\nu(y) \right],$$

$$z = \mathbf{r} \cdot \hat{\mathbf{p}}, \quad \frac{\partial}{\partial z} f_0(\mathbf{p}, \mathbf{r}_T, z) = \frac{d\varepsilon_\nu}{dz} \frac{\partial}{\partial \varepsilon_\nu} f_0(\mathbf{p}, \mathbf{r}_T, z)$$

§ 3 Estimating Pulsar Kick Velocities of Proto-Neutron Star

A.G.Lyne, D.R.Lomier, Nature 369, 127 (94)

Asymmetry of Supernova Explosion

kick and translate Pulsar with

Kick Velocity: Average ... 400km/s,

Highest ... 1500km/s

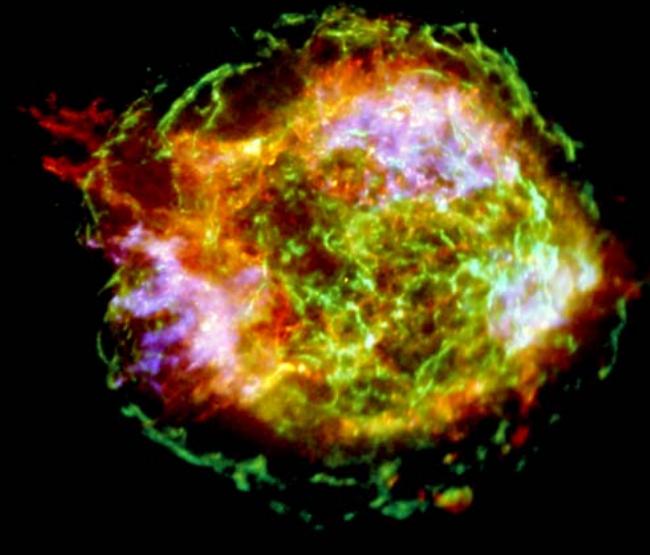
Explosion Energy $\sim 10^{53}$ erg

(almost Neutrino Emissions)

1 % Asymmetry is sufficient to explain the Pulsar Kick

D.Lai & Y.Z.Qian, Astrophys.J. 495 (1998) L103

CasA

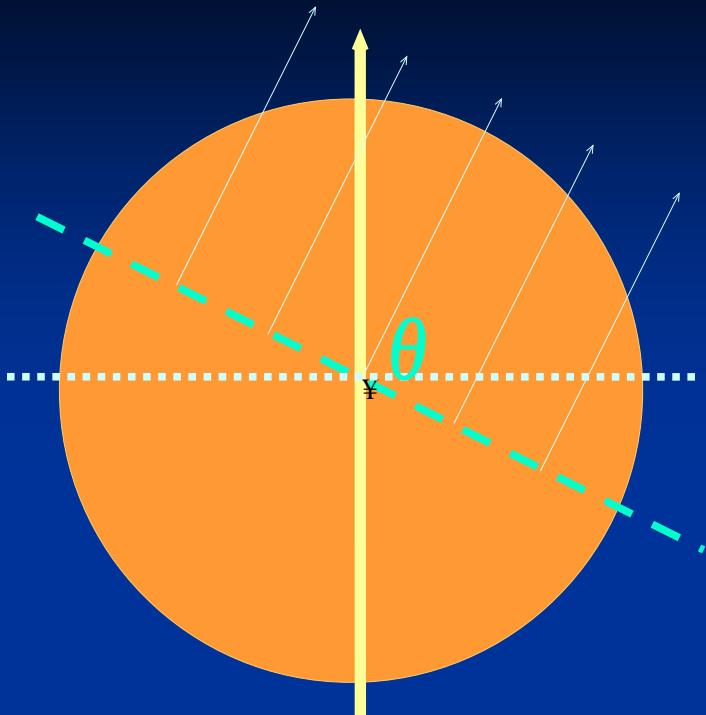


[http://chandra.harvard.edu/photo/
2004/casa_xray.jpg](http://chandra.harvard.edu/photo/2004/casa_xray.jpg)

Estimating Kick Velocity of PNS with $T = 20$ MeV and $B = 2 \times 10^{17}$ G

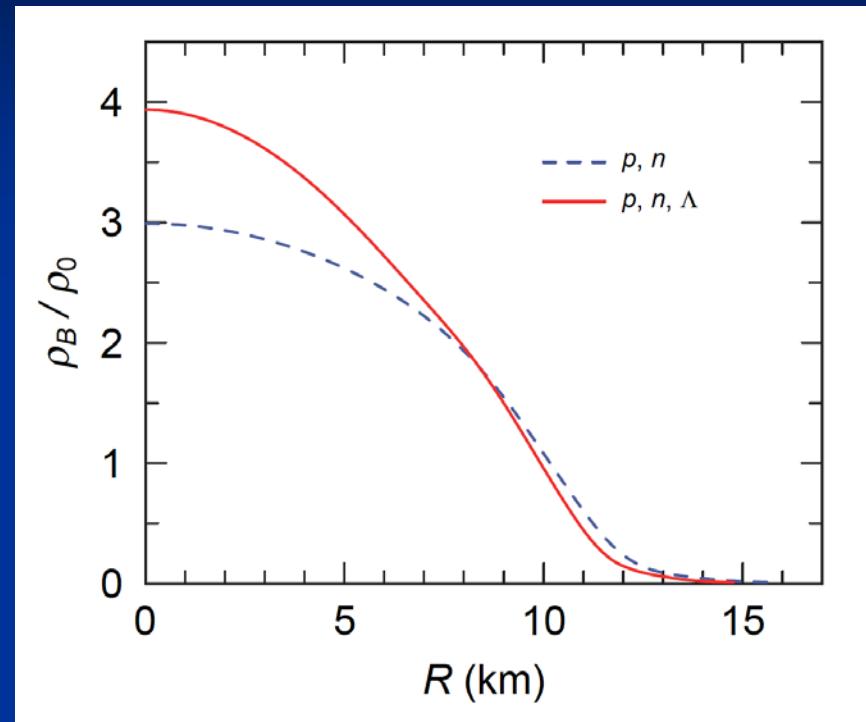
Poloidal Magnetic Field 2 - 3% Asymmetry in Absorption

Baryon density in PNS



- 1) Neutrinos propagate on the straight lines
- 2) Neutrino are produced and absorbed at all positions on the lines

Mean-Free path : V/σ_{ab}

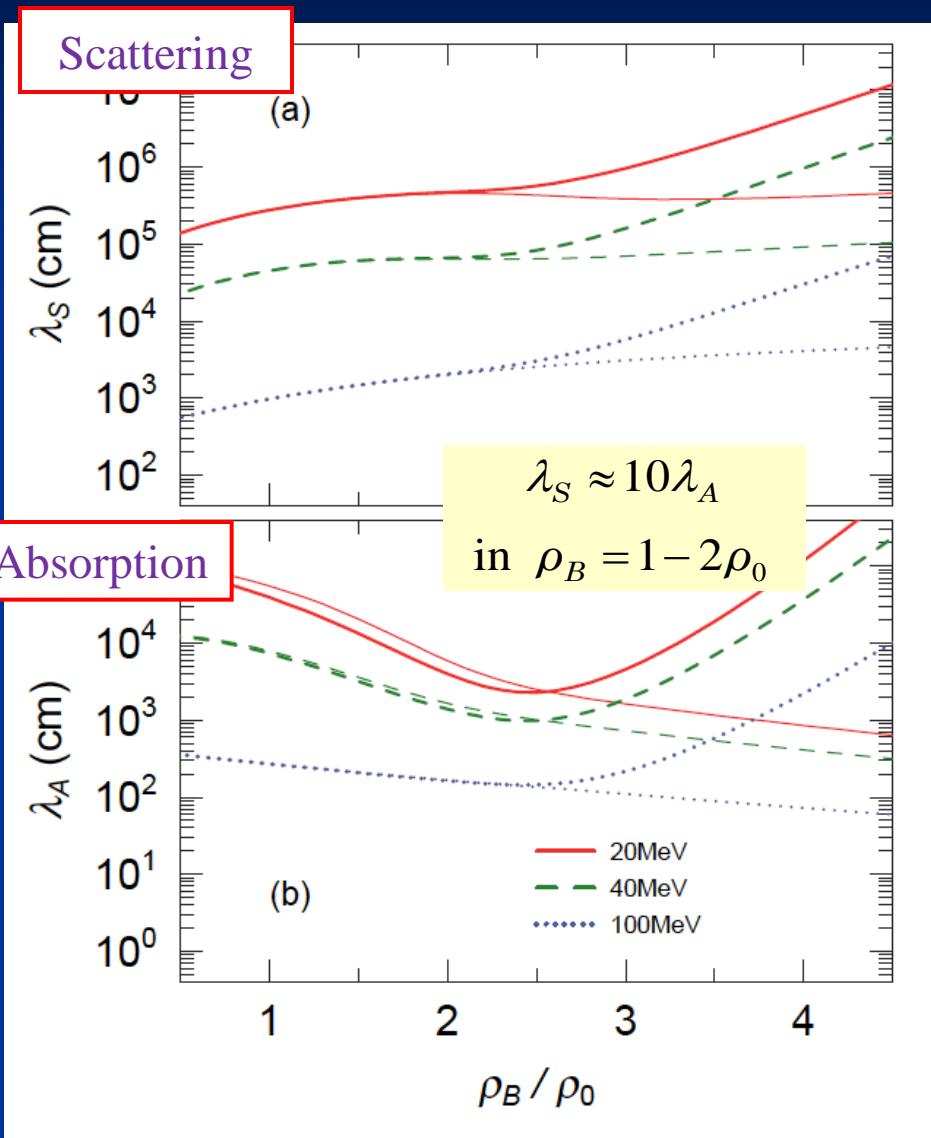


$$M = 1.68M_{solar}, Y_L = 0.4, T = 20 \text{ MeV}$$

Calculating Neutrino Propagation above $\rho_B = \rho_0 \rho_0$

Mean-Free Paths

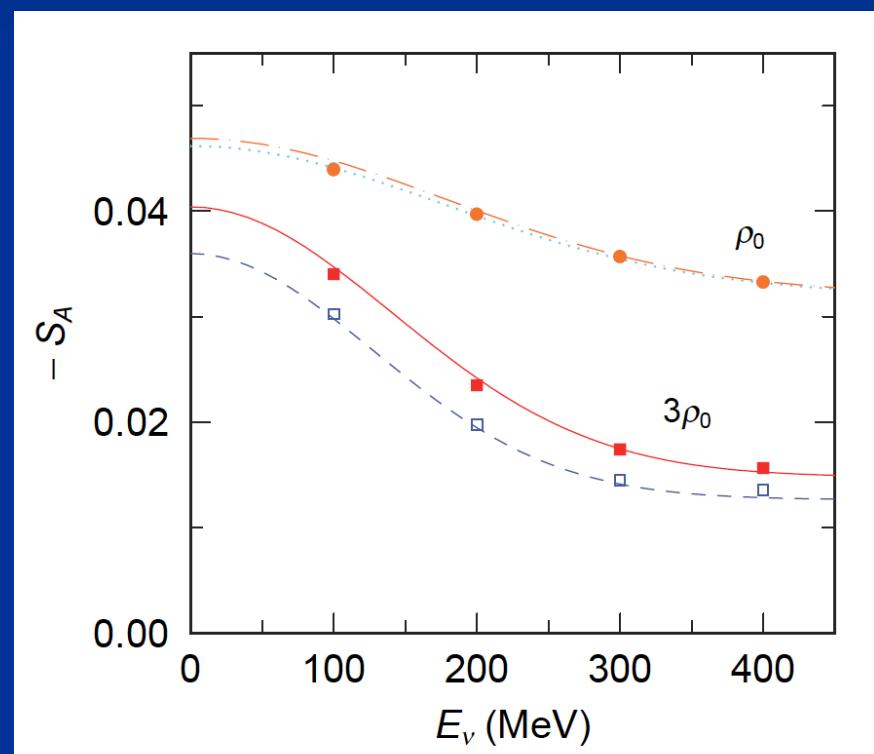
$$\lambda_{S,A} = V / \sigma_{S,A}^0$$



Magnetic Parts

$$\sigma_A = \sigma_A^0 (1 + S_A \cos \theta_{i,f})$$

S_A : fitting function



Angular Dep. of Emitted Neutrinos in Uniform Poloidal Mag. Field

$$M_{NS} = 1.68 M_{solar} \text{ [g]},$$

$$E_T = 3 \times 10^{53} \text{ [erg]}$$

$$\left\langle p_z \right\rangle = \frac{P_1}{E_T} = \frac{P_1}{3P_0} = 2.0 \times 10^{-2} \quad T = 20 \text{ MeV}$$

$$\langle p(n) \rangle \approx P_0 + P_1 \cos \theta$$

$$B = 2 \times 10^{17} \text{ [G]}$$

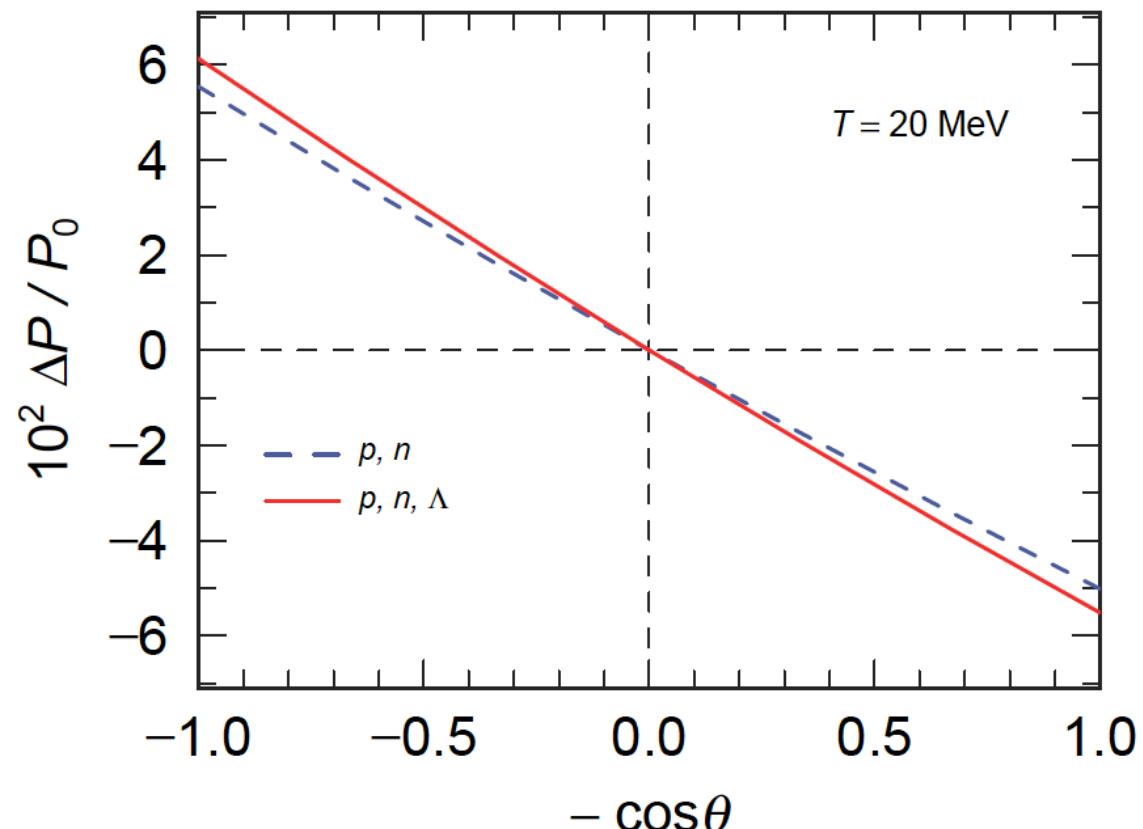
$$v_{kick} = \frac{\langle p_z \rangle}{M} \approx 600 \text{ [km/s]}$$

$$T = 20 \text{ MeV}$$

Observable

Average ... 400 km/s,

Highest ... 1500 km/s



§ 4 Angular Deceleration in Toroidal Magnetic Field

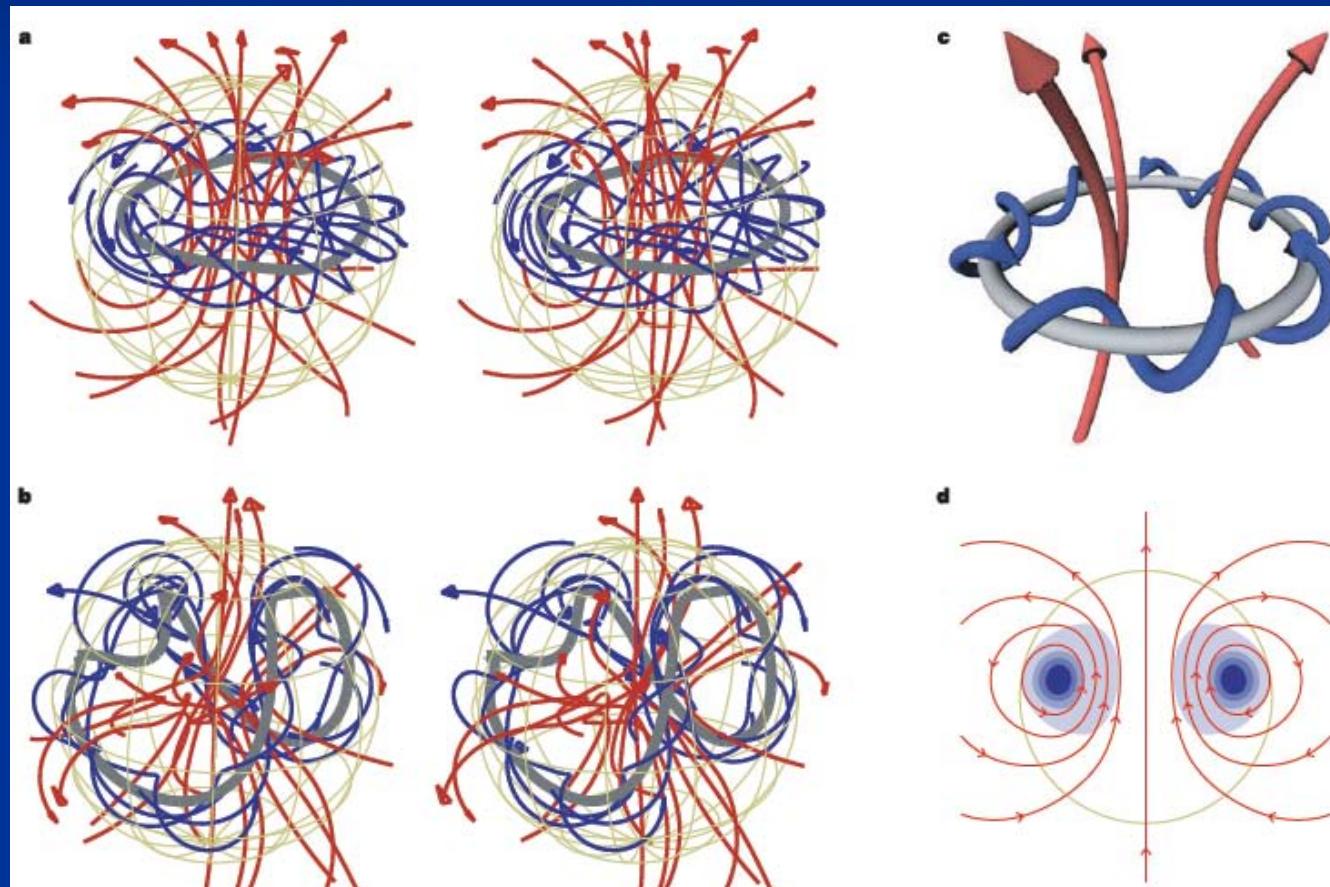
Stability of Magnetic Field in Compact Objects

(Braithwaite & Spruit 2004)

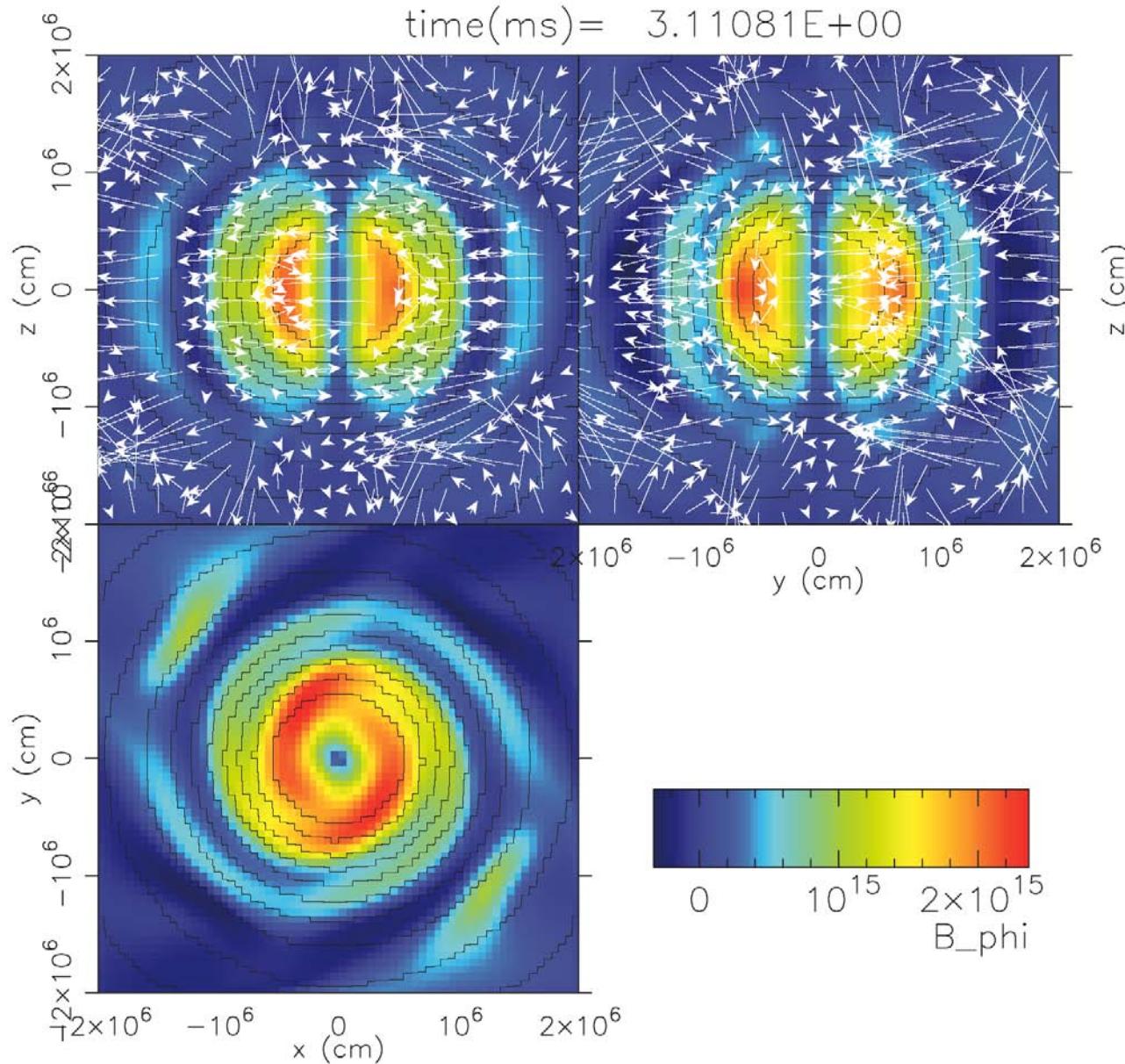
Toroidal Magnetic
Field is stable !!

Mag. Field
Parallel
to Baryonic Flow

Assym. of ν -Emit.
must decelerate
PNS Spin



No poloidal Magnetic Field at the beginning



Single
Toroidal

by T. Kuroda

Toroidal Magnetic Field

$$\mathbf{B}(r_T, z) = B_0 G_T(r_T) G_L(z) \hat{e}_\phi$$

$$G_T(r_T) = \frac{16 \exp[-(r_T - R_0)/\Delta r]}{\{1 + \exp[-(r_T - R_0)/\Delta r]\}^2}$$

$$G_L(z) = \frac{\exp[-z/\Delta r]}{\{1 + \exp[-z/\Delta r]\}^2}$$

$$\hat{e}_\phi = (-\sin \varphi, \cos \varphi, 0)$$

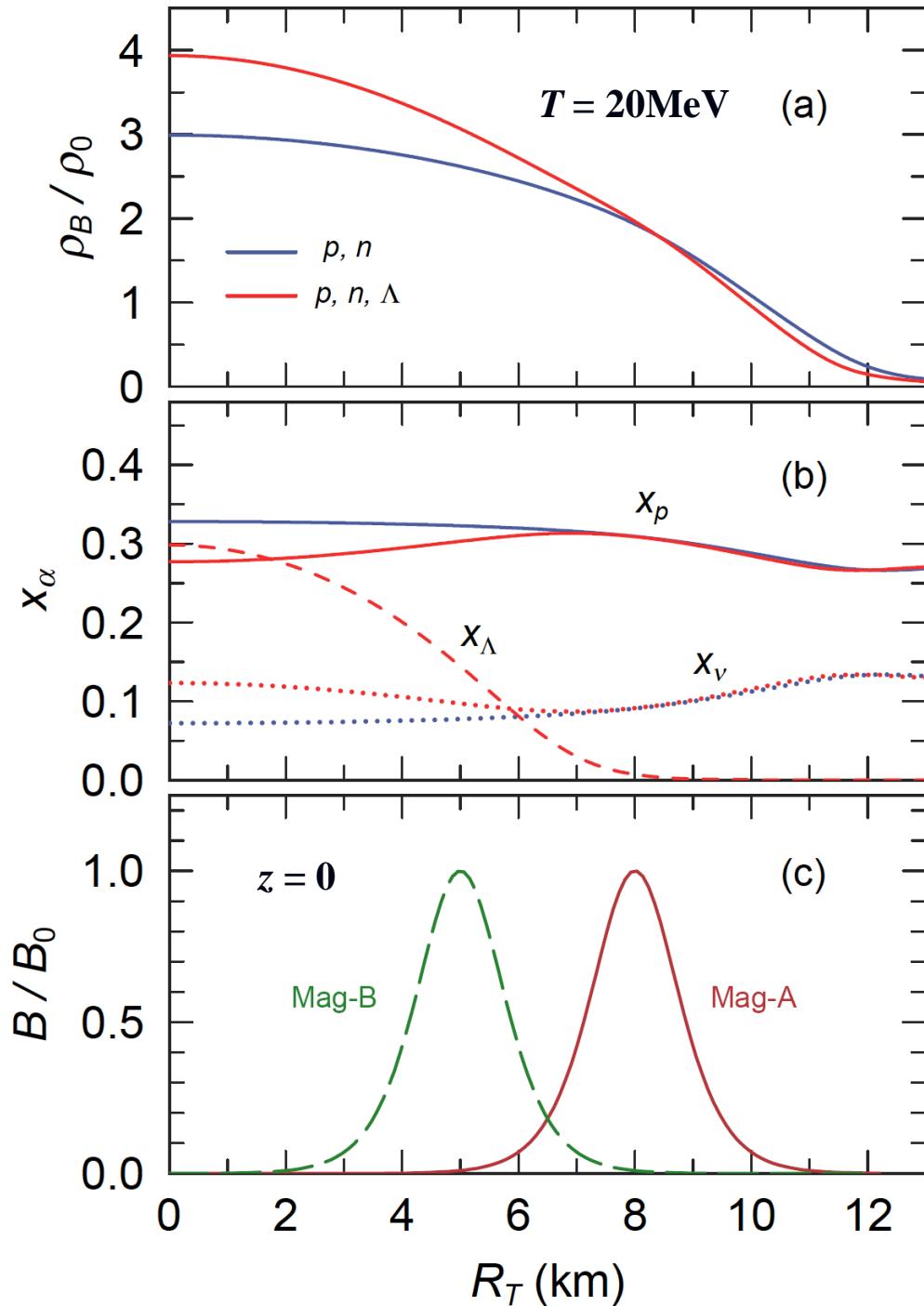
$$\Delta r = 0.5 \text{ (km)}$$

$$R_0 = 8 \text{ (km)}$$

Mag. - A

$$R_0 = 5 \text{ (km)}$$

Mag. - B



Angular Deceleration

$$\frac{dL_z}{dt} = c \int_{S_N} d\mathbf{r} \int d\mathbf{n} \int dp_L \Delta f(\mathbf{r}, p_L, \mathbf{n}) (\mathbf{r} \times \mathbf{p})_z$$

$$\dot{\omega} = \frac{d\omega}{dt} = \frac{1}{I_{NS}} \frac{dL_z}{dt} = \frac{1}{I_{NS}} \left(\frac{dE_T}{dt} \right)_v \frac{cdL_z / dt}{dE_T / dt}$$

Neutrino Luminosity

$$(dE_T/dt)_v \sim 3 \times 10^{52} \text{ erg/s}$$

$$M_{NS} = 1.68 M_{solar}$$

Period $P = 10\text{s}$

**Magnetic
Dipole Rad.**

$$\frac{\dot{P}}{P} = \frac{P}{2\pi c I_{NS}} \left(\frac{cdL_z / dt}{dE_T / dt} \right) \mathcal{L}_\nu.$$

Mag Distr.	Bary.	$-\frac{cdL_z / dt}{dE_T / dt}$ (cm)	$-\dot{\omega}$ (rad/s)	\dot{P}/P		
				ρ_0	$0.1\rho_0$	(Mg.Dpl-rad.)
Mag-A	p,n	66.9	4.94×10^{-2}	7.82×10^{-2}	1.03×10^{-2}	5.32×10^{-12}
	p,n, Λ	109	7.07×10^{-2}	11.26×10^{-2}	1.11×10^{-2}	6.77×10^{-11}
Mag-B	p,n	9.64	7.09×10^{-3}	1.13×10^{-3}	4.50×10^{-5}	5.32×10^{-12}
	p,n, Λ	7.81	5.07×10^{-3}	8.07×10^{-4}	2.29×10^{-5}	6.77×10^{-11}

In Early Stage (~ 10 s) ν Asymmetric Emission must affect PNS Spin
More Significantly than Magnetic Dipole γ -Radiation

§ 5 Summary

- EOS of Neutron-Star-Matter with p, n, Λ in RMF Approach
- Exactly Solving Dirac Eq. with Magnetic Field in Perturbative Way
- Cross-Sections of Neutrino Scattering and Absorption under Strong Magnetic Field, calculated in Perturbative Way
- Neutrinos are More Scattered and Less Absorbed
 - in Direction Parallel to Magnetic Field
 - ⇒ More Neutrinos are Emitted in Arctic Area
 - Scattering 1.7 %
 - Absorption 2.2 % at $\rho_B = 3\rho_0$ and $T = 20 \text{ MeV}$
 - ⇒ Convection, Pulsar-Kick

Pulsar Kick in Poroidal Mag. Field $B = 2 \times 10^{17} \text{G}$ \Rightarrow Perturbative Calc. Dir.

$$\begin{aligned} v_{\text{kick}} &= 580 \text{ km/s (p,n)}, \quad 610 \text{ km/s (p,n,}\Lambda\text{)} \quad \text{at } T = 20 \text{ MeV} \\ &= 230 \text{ km/s (p,n)}, \quad 270 \text{ km/s (p,n,}\Lambda\text{)} \quad \text{at } T = 30 \text{ MeV} \end{aligned}$$

400 km/s (Average of Observed Values)

Spin Deceleration in Toroidal Magnetic Field

$$\begin{array}{lll} \dot{P}/P [\text{s}^{-1}] & \text{Asymmetric Neutrino Emit.} & 10^{-2} \sim 10^{-5} \text{ when } B \sim 10^{17} \text{G} \\ & \text{Magnetic Dipole Radiation} & 10^{-11} \sim 10^{-12} \end{array}$$

Asymmetric ν -Emission plays an important role in PNS Spin in Early Time

Future Plans

ν -Scattering

Iso-Temp. \Rightarrow Iso-Entropy

Exact Solution of Dirac Eq. in Non-Perturbative Cal.
→ Landau Level at least for Electron

Neutrino Propagation in Low Density



Making Data Table and
Applying it to Supernovae Simulations