$\begin{array}{l} \text{On branes in} \\ \mathcal{N}=2 \\ \text{supergravities} \end{array}$

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Branes in maximal supergravities

Supergravities with less supersymmetry

Real forms and Tits-Satake diagrams

Branes in $\mathcal{N} = 2$ supergravities

Conclusions

On branes in $\mathcal{N} = 2$ supergravities

Fabio Riccioni

Workshop on Exotic Structures of Spacetime March 10-21, 2014 YITP, Kyoto



Sezione di Roma



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Based on work with E. Bergshoeff and L. Romano

19th March 2014

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IIA/IIB string theory on $\mathcal{T}^d o O(d,d;\mathbb{Z})$ T-duality

Consider d = 1: one defines

$$p_L = \frac{m}{2R} + \frac{nR}{2\alpha'}$$
 $p_R = \frac{m}{2R} - \frac{nR}{2\alpha'}$

One finds that the spectrum contains 1/2-BPS states that are purely momentum or purely winding states

These states are charged under the graviphoton $g_{\mu9}$ with charge $p_L + p_R$, and under $B_{\mu9}$ with charge $p_L - p_R$

The invariant quantity under $O(1,1;\mathbb{Z})$ is $p_L^2 - p_R^2 \sim mn$, which vanishes for 1/2-BPS states

This generalises to any *d*

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From the low-energy supergravity view-point, one has a continuous SO(d, d) symmetry, and one finds 1/2-BPS black hole solutions with charge Q_A such that $Q^A Q_A = 0$

But maximal supergravity theories have bigger symmetries: $E_{d+1(d+1)} \supset SO(d, d) \times \mathbb{R}^+$

e.g.:

D = 7: $SL(5, \mathbb{R}) \supset SO(3, 3) \times \mathbb{R}^+$

D = 6: $SO(5,5) \supset SO(4,4) \times \mathbb{R}^+$

 $D = 5: \quad E_{6(6)} \supset SO(5,5) \times \mathbb{R}^+$

In the quantum theory these are conjectured to be broken to discrete U-duality symmetries

Hull, Townsend

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 $1/2\mbox{-BPS}$ branes in string theory correspond to $1/2\mbox{-BPS}$ brane solutions of the supergravity theory

These include D-branes, but also more non-perturbative branes

Supersymmetry and the global symmetries allow a complete classification of these branes

Such classification includes *non-standard* branes, i.e. domain walls and space-filling branes, for which one can write a supersymmetric (and κ -symmetric) effective action

Wess-Zumino term for a *p*-brane:

 $Q \int A_{p+1} + \dots$

- IIA: $Q \int A_9 + \dots$ D8-brane
- IIB: $Q \int A_{10} + \dots$ D9-brane

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10-forms of IIB:

 $A_{10}^{\alpha\beta\gamma}$ 4-plet of SL(2, \mathbb{R})

Bergshoeff, de Roo, Kerstan, FR

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Corresponding charge:

 $Q_{\alpha\beta\gamma}$

1/2-BPS constraint:

 $Q_{\alpha\beta\gamma}\epsilon^{\gamma\delta}Q_{\delta\epsilon\tau}=0$



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In any dimensions, one classifies all the fields A_{p+1} , belonging to representations of $E_{d+1(d+1)}$

FR, West Bergshoeff, De Baetselier, Nutma

The $1/2\mbox{-}BPS$ condition gives a constraint on the charge. This is always the highest-weight constraint

Ferrara, Maldacena Bergshoeff, Marrani, FR

The number of independent branes inside a representation is the number of longest weights of that representation Bergshoeff, FR, Romano

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$$D = 8$$
 Global symmetry $G = SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$

field	representation	# of branes
1-forms	(3 ,2)	6
2-forms	(3,1)	3
3-forms	(1,2)	2
4-forms	(3 ,1)	3
5-forms	(3,2)	6
6-forms	(8,1)	6
	(1,3)	2
7-forms	(6 , 2)	6
8-forms	(15, 1)	6

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Half-maximal sugra theories coupled to n_V vector multiplets in 10 - d dimensions:

 $SO(d, n_V)/[SO(d) \times SO(n_V)]$

We do not consider gaugings

String theory: $n_V = d + 16$

Real form not maximally non-compact

 $1/2\mbox{-}BPS$ branes: same constraint as before: highest-weight constraint

 $\begin{array}{l} \mbox{Counting $1/2$-BPS single branes} \rightarrow \mbox{lightcone rules} \\ \mbox{Bergshoeff, FR} \end{array}$

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Quarter-maximal sugra theories exist in six dimensions and below. They couple sugra to tensor multiplets (six dim), vector multiplets and hyper-multiplets.

Scalars in the vector multiplets parametrise:

- a very special real manifold in five dimensions
- a special Kähler manifold in four dimensions
- a quaternionic manifold in three dimensions

We will consider theories that give rise to symmetric spaces upon reduction to three dimensions

Such theories contain no hypers in dimension higher than three

In general the global symmetry is not maximally non-compact

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Branes in $\mathcal{N} = 2$ supergravities $SO(4, n_H)/[SO(4) \times SO(n_H)]$ $F_{4(4)}/[USp(6) \times SU(2)]$ $(n_H = 7)$ $E_{6(2)}/[SU(6) \times SU(2)]$ (*n_H* = 10) $E_{7(-5)}/[SO(12) \times SU(2)]$ $(n_H = 16)$ $E_{8(-24)}/[E_7 \times SU(2)]$ (*n_H* = 28) $G_{2(2)}/SO(4)$ (*n_H* = 2) $SU(n_H, 2)/[SU(n_H) \times SU(2) \times U(1)]$ $USp(2n_H, 2)/[USp(2n_H) \times USp(2)]$

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Symmetric quaternionic manifolds:

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D = 3	<i>D</i> = 4	D = 5	<i>D</i> = 6
$\frac{SO(4, n)}{SO(4) \times SO(n)}$ $n\mu = n$	$\frac{\mathrm{SO}(2, n-2)}{\mathrm{SO}(2) \times \mathrm{SO}(n-2)} \times \frac{\mathrm{SU}(1, 1)}{\mathrm{U}(1)}$ $n_{V} = n - 1$	$\frac{\mathrm{SO}(1, n-3)}{\mathrm{SO}(n-3)} \times \mathbb{R}^+$ $n_V = n-2$	$SO(1, 1)$ $n_T = 1, n_V = n - 4$ $SO(1, n - 3)$
 F ₄₍₄₎	Sp(6, ℝ)	SL(3.ℝ)	SO(n-3) $n_T = n-3, n_V = 0$ SO(1, 2)
$\overline{\text{USp}(6) \times \text{SU}(2)}$ $n_H = 7$	$\frac{U(3)}{U(3)}$	$\overline{SO(3)}$ $n_V = 5$	$\overline{SO(2)}$ $n_T = 2, n_V = 2$
$\frac{E_{6(2)}}{SU(6)\times SU(2)}$	$\frac{SU(3,3)}{SU(3)\timesSU(3)}$	SL(3, ℂ) SU(3)	SO(1, 3) SO(3)
$n_{H} = 10$	n _V = 9	<i>n</i> _V = 8	$n_T = 3, n_V = 4$
$\frac{E_{7(-5)}}{SO(12)\timesSU(2)}$	SO*(12) U(6)	SU*(6) USp(6)	SO(1,5) SO(5)
$n_{H} = 16$	$n_V = 15$	$n_V = 14$	$n_T = 5, n_V = 8$
$\frac{E_{8(-24)}}{E_7\timesSU(2)}$	$\frac{E_{7(-25)}}{E_6\times SO(2)}$	$\frac{E_{6(-26)}}{F_4}$	SO(1,9) SO(9)
$n_{H} = 28$	$n_V = 27$	$n_V = 26$	$n_T=9, n_V=16$
$\frac{G_{2(2)}}{SO(4)}$	$\frac{SU(1,1)}{U(1)}$	1	-
$n_H = 2$	$n_V = 1$	$n_V = 0$	
$\frac{SU(n,2)}{SU(n)\timesSU(2)\timesU(1)}$	$\frac{SU(n-1,1)}{SU(n-1)\timesU(1)}$	-	-
$n_H = n$	$n_V = n - 1$		
$\frac{\text{USp}(2n,2)}{\text{USp}(2n) \times \text{USp}(2)}$	-	-	-
$n_H = n$			

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 $1/2\text{-}\mathsf{BPS}$ branes \rightarrow highest-weight constraints

We want to count the branes

SO(4, n) theories: lightcone rule

dim.	type of brane	field	# of branes
<i>D</i> = 3	0-brane	A_{1,A_1A_2}	24
	1-brane	A _{2,AB}	8
		A _{2,A1A4}	16
<i>D</i> = 4	0-brane	A _{1,Aa}	8
	1-brane	A _{2,ab}	2
		A_{2,A_1A_2}	4
D = 5	0-brane	$A_{1,A}$ ($\alpha = 0$)	2
		$A_1 (\alpha = -2)$	1
	1-brane	$A_2 (\alpha = 0)$	1
		$A_{2,A} (\alpha = -2)$	2
<i>D</i> = 6	1-brane	A _{2,A}	2

How do we count the branes in the other theories?

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A real form exists if the algebra admits a basis such that all the structure constants are real

Killing form:

 $B(X,Y) = \operatorname{Tr}(\operatorname{ad} X \operatorname{ad} Y)$

- Positive eigenvectors: non-compact generators
- Negative eigenvectors: compact generators

Cartan involution:

 $B_{\theta}(X,Y) = B(X,\theta Y)$

such that $B_{\theta}(X, Y)$ is negative-definite. Diagonalise it:

- $\theta = 1$ eigenvectors are compact
- $\theta = -1$ eigenvectors are non-compact

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Generators of SL(2, \mathbb{R}): $\sigma_1 \quad i\sigma_2 \quad \sigma_3$ Generators of SU(2): $i\sigma_1 \quad i\sigma_2 \quad i\sigma_3$ Killing forms:

 $B_{\mathfrak{sl}(2,\mathbb{R})} = \begin{pmatrix} 8 & 0 & 0 \\ 0 & -8 & 0 \\ 0 & 0 & 8 \end{pmatrix} \qquad B_{\mathfrak{su}(2)} = \begin{pmatrix} -8 & 0 & 0 \\ 0 & -8 & 0 \\ 0 & 0 & -8 \end{pmatrix}$ Define $H = \sigma_3 \quad E_+ = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \quad E_- = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$

Cartan involution of $SL(2,\mathbb{R})$ is

 $\theta H = -H \quad \theta E_+ = -E_-$

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 $\operatorname{ad}(\theta X) = -(\operatorname{ad} X)^{\dagger}$

and

 $[H, E_{\alpha}] = \alpha(H)E_{\alpha}$

one derives the reality properties of the roots $\alpha(H)$.

- real root: vanishes for compact $H \rightarrow \theta \alpha = -\alpha$
- imaginary root: vanishes for non-compact $H \rightarrow \theta \alpha = \alpha$

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complex root: otherwise

Cartan involution on the roots: $\theta H_{\alpha} = H_{\theta \alpha}$

Implies $\theta E_{\alpha} = \pm E_{\theta \alpha}$

Real forms and TS diagrams

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Real forms and TS diagrams

 $\theta \alpha = \alpha \text{ (fixed root)}$ imaginary (compact generator) $\theta \alpha = -\alpha \text{ up to fixed roots}$ $\theta \alpha_i = -\alpha_j \text{ up to fixed roots}$ complex
complex

From the TS diagram one derives the action of the Cartan involution on the generators

 $\begin{array}{l} \text{On branes in} \\ \mathcal{N}=2 \\ \text{supergravities} \end{array}$

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Example: real forms of $SL(3, \mathbb{C})$

SU(2,1): $\theta \alpha_1 = -\alpha_2$

On branes in $\mathcal{N}=2$ supergravities

and

Fabio Riccioni real form generators compact non-compact $\mathfrak{sl}(3,\mathbb{R})$ Cartan H_{α_1} H_{α_2} $E_{\alpha_1} + E_{-\alpha_1}$ $E_{\alpha_1} - E_{-\alpha_1}$ root $E_{\alpha_2} - E_{-\alpha_2}$ $E_{\alpha_2} + E_{-\alpha_2}$ generators $\frac{E_{\alpha_1+\alpha_2}-E_{-\alpha_1-\alpha_2}}{iH_{\alpha_1}}$ $E_{\alpha_1+\alpha_2} + E_{-\alpha_1-\alpha_2}$ su(3) Cartan Real forms iH_{α_2} $E_{\alpha_1} - E_{-\alpha_1}$ Tits-Satake root diagrams generators $i(E_{\alpha_1}+E_{-\alpha_1})$ $E_{\alpha_2} - E_{-\alpha_2}$ $i(E_{\alpha_2}+E_{-\alpha_2})$ $E_{\alpha_1+\alpha_2} - E_{-\alpha_1-\alpha_2}$ $i(E_{\alpha_1+\alpha_2}+E_{-\alpha_1-\alpha_2})$ $i(H_{\alpha_1} - H_{\alpha_2})$ $\frac{H_{\alpha_1} + H_{\alpha_2}}{E_{\alpha_1} + E_{\alpha_2} + E_{-\alpha_1} + E_{-\alpha_2}}$ su(2,1) Cartan $E_{\alpha_1} + E_{\alpha_2} - E_{-\alpha_1} - E_{-\alpha_2}$ root $i(E_{\alpha_1} - E_{\alpha_2} + E_{-\alpha_1} - E_{-\alpha_2}) \mid i(E_{\alpha_1} - E_{\alpha_2} - E_{-\alpha_1} + E_{-\alpha_2})$ generators

 $i(E_{\alpha_1+\alpha_2}+E_{-\alpha_1-\alpha_2})$

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 $i(E_{\alpha_1+\alpha_2}-E_{-\alpha_1-\alpha_2})$

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G/H	Tits-Satake diagram of G	Cartan involution
$\frac{SO(4,n)}{SO(4)\timesSO(n)}$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{l} \theta(\alpha_1)=-\alpha_1\\ \theta(\alpha_2)=-\alpha_2\\ \theta(\alpha_3)=-\alpha_3\\ \theta(\alpha_4)=-\alpha_4-2\sum_{i=5}^k\beta_i\\ -2^a\bmod 2(\beta_{k+1}+\beta_{k+2})\\ \theta(\beta_i)=\beta_i i=5,,k+2 \end{array}$
$\frac{F_{4(4)}}{USp(6) \times SU(2)}$		$\theta(\alpha_i) = -\alpha_i$ $i = 1, 2, 3, 4$
$\frac{E_{6(2)}}{SU(6)\times SU(2)}$		$\theta(\alpha_1) = -\alpha_5$ $\theta(\alpha_2) = -\alpha_4$ $\theta(\alpha_3) = -\alpha_3$ $\theta(\alpha_4) = -\alpha_2$ $\theta(\alpha_5) = -\alpha_1$ $\theta(\alpha_6) = -\alpha_6$
$\frac{E_{7(-5)}}{SO(12)\timesSU(2)}$		$\begin{aligned} \theta(\alpha_2) &= -\alpha_2 - \beta_1 - \beta_3 \\ \theta(\alpha_4) &= -\alpha_4 - \beta_3 - \beta_7 \\ \theta(\alpha_5) &= -\alpha_5 \\ \theta(\alpha_6) &= -\alpha_6 \\ \theta(\beta_i) &= \beta_i i = 1, 3, 7 \end{aligned}$
$\frac{E_{8(-24)}}{E_7\times SU(2)}$		$\begin{array}{l} \theta(\alpha_1) = -\alpha_1 \\ \theta(\alpha_2) = -\alpha_2 \\ \theta(\alpha_3) = -\alpha_3 - 2\beta_4 - 2\beta_5 - \beta_6 - \beta_8 \\ \theta(\alpha_7) = -\alpha_7 - \beta_4 - 2\beta_5 - 2\beta_6 - \beta_8 \\ \theta(\beta_i) = \beta_i \qquad i = 4, 5, 6, 8 \end{array}$
G2(2) SO(4)		$\theta(\alpha_i) = -\alpha_i$ $i = 1, 2$
$\frac{SU(n,2)}{SU(n) \times SU(2) \times U(1)}$	n=2 1 2 3 $n>2$ $n>2$ $n>2$ $n>2$ $n>2$	$\begin{array}{l} \theta(\alpha_1) = -\alpha_3\\ \theta(\alpha_2) = -\alpha_2\\ \theta(\alpha_3) = -\alpha_1\\ \theta(\alpha_1) = -\alpha_{n+1}\\ \theta(\alpha_2) = -\alpha_n - \sum_{i=3}^{n-1} \beta_i\\ \theta(\beta_i) = \beta_i i = 3,, n-1 \end{array}$
115-(2-2)		$\begin{aligned} \theta(\alpha_n) &= -\alpha_2 - \sum_{i=3}^{n-1} \beta_i \\ \theta(\alpha_{n+1}) &= -\alpha_1 \\ \theta(\beta_1) &= \beta_1 \\ \theta(\alpha_2) &= -\alpha_2 - 2\beta_1 \end{aligned}$
$\frac{\text{USp}(2n, 2)}{\text{USp}(2n) \times \text{USp}(2)}$	n > 1 n > 1 1 2 3 n n + 1	$\theta(\beta_i) = \beta_i$ $i \neq 2$ $\theta(\alpha_2) = -\alpha_2 - \beta_1$ $-2\sum_{i=3}^n \beta_i - \beta_{n+1}$

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On branes in $\mathcal{N} = 2$ supergravities

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Conclusions

Branes in $\mathcal{N} = 2$ sugras

Project roots under $P = \frac{1}{2}(1-\theta)$

Projected roots define the restricted-root algebra

Take the highest weight of the representation of the field

If highest weight is not fixed under P: no branes

If it is fixed: write it as highest weight of a representation of the restricted-root algebra

Count the longest weights of that representation: this is the number of real longest weights, i.e. the number of branes

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Conclusions

dim.	type of brane	repr.	highest weight	# of branes
<i>D</i> = 3	0-brane	52 (1 0 0 0)	$2\alpha_1 + 3\alpha_2 + 4\alpha_3 + 2\alpha_4$	24
	1-brane	324 (0 0 0 2)	$2\alpha_1 + 4\alpha_2 + 6\alpha_3 + 4\alpha_4$	24
<i>D</i> = 4	0-brane	14 (1 0 0)	$\frac{3}{2}\alpha_2 + 2\alpha_3 + \alpha_4$	8
	1-brane	21 (0 0 2)	$\alpha_2 + 2\alpha_3 + 2\alpha_4$	6
D = 5	0-brane	6 (2 0)	$\frac{4}{3}\alpha_{3} + \frac{2}{3}\alpha_{4}$	3
	1-brane	6 (0 2)	$\frac{2}{3}\alpha_3 + \frac{4}{3}\alpha_4$	3
<i>D</i> = 6	1-brane	3 (2)	α_4	2

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$\begin{array}{l} \text{On branes in} \\ \mathcal{N} = 2 \\ \text{supergravities} \end{array}$

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Conclusions

Consider $E_{6(2)}$:



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Restricted-root algebra:

 $P\alpha_3 = \alpha_3$

 $P\alpha_6 = \alpha_6$

 $P\alpha_2 = \frac{1}{2}(\alpha_2 + \alpha_4)$

 $P\alpha_1 = \frac{1}{2}(\alpha_1 + \alpha_5)$

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Conclusions

dim.	brane	repr.	highest weight	restricted repr.	#
D = 3	0-brane	78 (0 0 0 0 0 1)	$\alpha_1+2\alpha_2+3\alpha_3+2\alpha_4+\alpha_5+2\alpha_6$	52 (1 0 0 0)	24
	1-brane	650 (1 0 0 0 1 0)	$2\alpha_1+3\alpha_2+4\alpha_3+3\alpha_4+2\alpha_5+2\alpha_6$	324 (0 0 0 2)	24
	2-brane	5824 (1 1 0 0 0 0)	$3\alpha_1+5\alpha_2+6\alpha_3+4\alpha_4+2\alpha_5+3\alpha_6$		
		5824 (0 0 0 1 1 0)	$2\alpha_1+4\alpha_2+6\alpha_3+5\alpha_4+3\alpha_5+3\alpha_6$	-	
<i>D</i> = 4	0-brane	20 (0 0 1 0 0)	$\frac{1}{2}\alpha_1 + \alpha_2 + \frac{3}{2}\alpha_3 + \alpha_4 + \frac{1}{2}\alpha_5$	14 (1 0 0)	8
	1-brane	35 (1 0 0 0 1)	$\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4 + \alpha_5$	21 (0 0 2)	6
	2-brane	70 (1 1 0 0 0)	$\frac{3}{2}\alpha_1 + 2\alpha_2 + \frac{3}{2}\alpha_3 + \alpha_4 + \frac{1}{2}\alpha_5$	-	-
		70 (0 0 0 1 1)	$\frac{1}{2}\alpha_1 + \alpha_2 + \frac{3}{2}\alpha_3 + 2\alpha_4 + \frac{3}{2}\alpha_5$	-	
	3-brane	280 (2 0 0 1 0)	$2\alpha_1 + 2\alpha_2 + 2\alpha_3 + 2\alpha_4 + \alpha_5$	-	-
		280 (0 1 0 0 2)	$\alpha_1+2\alpha_2+2\alpha_3+2\alpha_4+2\alpha_5$	-	-
D = 5	0-brane	(3,3) (0 1 1 0)	$\frac{1}{3}\alpha_1 + \frac{2}{3}\alpha_2 + \frac{2}{3}\alpha_4 + \frac{1}{3}\alpha_5$	6 (2 0)	3
	1-brane	(3,3) (1 0 0 1)	$\frac{2}{3}\alpha_1 + \frac{1}{3}\alpha_2 + \frac{1}{3}\alpha_4 + \frac{2}{3}\alpha_5$	6 (0 2)	3
	2-brane	(8,1) (1100)	$\alpha_1 + \alpha_2$	-	-
		(1,8) (0 0 1 1)	$\alpha_4 + \alpha_5$	-	-
	3-brane	(6,3) (2010)	$\frac{4}{3}\alpha_1 + \frac{2}{3}\alpha_2 + \frac{2}{3}\alpha_4 + \frac{1}{3}\alpha_5$	-	-
		(3,6) (0 1 0 2)	$\frac{1}{3}\alpha_1 + \frac{2}{3}\alpha_2 + \frac{2}{3}\alpha_4 + \frac{4}{3}\alpha_5$	-	-
	4-brane	(15,3) (2 1 0 1)	$\frac{5}{3}\alpha_1 + \frac{4}{3}\alpha_2 + \frac{1}{3}\alpha_4 + \frac{2}{3}\alpha_5$	-	-
		(3, 15) (1 0 1 2)	$\frac{2}{3}\alpha_1 + \frac{1}{3}\alpha_2 + \frac{4}{3}\alpha_4 + \frac{5}{3}\alpha_5$	-	-
D = 6	0-brane	(2,1) (1 0)	$\frac{1}{2}\alpha_1$	-	-
		(1,2) (01)	$\frac{1}{2}\alpha_5$	-	
	1-brane	(2,2) (1 1)	$\frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_5$	3 (2)	2
	2-brane	(2,1) (1 0)	$\frac{1}{2}\alpha_1$	-	-
		(1,2) (01)	$\frac{1}{2}\alpha_5$	-	
	3-brane	(1,3) (0 2)	Ω5	-	
		(3,1) (20)	α ₁	-	
	4-brane	(2,3) (1 2)	$\frac{1}{2}\alpha_1 + \alpha_5$	-	-
		(3,2) (21)	$\alpha_1 + \frac{1}{2}\alpha_5$	-	-
	5-brane	(4,2) (31)	$\frac{3}{2}\alpha_1 + \frac{1}{2}\alpha_5$	-	-
		(2,4) (1 3)	$\frac{1}{2}\alpha_1 + \frac{3}{2}\alpha_5$	-	-

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Conclusions

dim.	brane	repr.	highest weight	restricted repr.	#
D = 3	0-brane	133 (0 0 0 0 0 1 0)	$\beta_1+2\alpha_2+3\beta_3+4\alpha_4+3\alpha_5+2\alpha_6+2\beta_7$	52 (1 0 0 0)	24
	1-brane	1539 (0 1 0 0 0 0 0)	$2\beta_1 + 4\alpha_2 + 5\beta_3 + 6\alpha_4 + 4\alpha_5 + 2\alpha_6 + 3\beta_7$	324 (0 0 0 2)	24
	2-brane	40755 (1 0 0 0 0 0 1)	$3\beta_1+5\alpha_2+7\beta_3+9\alpha_4+6\alpha_5+3\alpha_6+5\beta_7$	-	-
<i>D</i> = 4	0-brane	32 (0 0 0 0 1 0)	$\frac{1}{2}\beta_1 + \alpha_2 + \frac{3}{2}\beta_3 + 2\alpha_4 + \frac{3}{2}\alpha_5 + \beta_7$	14 (1 0 0)	8
	1-brane	66 (0 1 0 0 0 0)	$\beta_1+2\alpha_2+2\beta_3+2\alpha_4+\alpha_5+\beta_7$	21 (0 0 2)	6
	2-brane	352 (1 0 0 0 0 1)	$\frac{3}{2}\beta_1 + 2\alpha_2 + \frac{5}{2}\beta_3 + 3\alpha_4 + \frac{3}{2}\alpha_5 + 2\beta_7$	-	-
	3-brane	462 (0 0 0 0 0 0 2)	$\beta_1+2\alpha_2+3\beta_3+4\alpha_4+2\alpha_5+3\beta_7$	-	-
		2079 (1 0 1 0 0 0)	$2\beta_1+3\alpha_2+4\beta_3+4\alpha_4+2\alpha_5+2\beta_7$	-	-
D = 5	0-brane	15 (0 0 0 1 0)	$\frac{1}{3}\beta_1 + \frac{2}{3}\alpha_2 + \beta_3 + \frac{4}{3}\alpha_4 + \frac{2}{3}\beta_7$	6 (2 0)	3
	1-brane	15 (0 1 0 0 0)	$\frac{2}{3}\beta_1 + \frac{4}{3}\alpha_2 + \beta_3 + \frac{2}{3}\alpha_4 + \frac{1}{3}\beta_7$	6 (0 2)	3
	2-brane	35 (1 0 0 0 1)	$\beta_1 + \alpha_2 + \beta_3 + \alpha_4 + \beta_7$	-	-
	3-brane	21 (0 0 0 0 2)	$\frac{1}{3}\beta_1 + \frac{2}{3}\alpha_2 + \beta_3 + \frac{4}{3}\alpha_4 + \frac{5}{3}\beta_7$	-	-
		105 (1 0 1 0 0)	$\frac{4}{3}\beta_1 + \frac{5}{3}\alpha_2 + 2\beta_3 + \frac{4}{3}\alpha_4 + \frac{2}{3}\beta_7$	-	-
	4-brane	384 (1 1 0 0 1)	$\frac{5}{3}\beta_1 + \frac{7}{3}\alpha_2 + 2\beta_3 + \frac{5}{3}\alpha_4 + \frac{4}{3}\beta_7$	-	-
<i>D</i> = 6	0-brane	(4 , 2) (0 0 1 1)	$\frac{1}{4}\beta_1 + \frac{1}{2}\alpha_2 + \frac{3}{4}\beta_3 + \frac{1}{2}\beta_7$	-	-
	1-brane	(6,1) (0 1 0 0)	$\frac{1}{2}\beta_1 + \alpha_2 + \frac{1}{2}\beta_3$	3 (2)	2
	2-brane	(4,2) (1 0 0 1)	$\frac{3}{4}\beta_1 + \frac{1}{2}\alpha_2 + \frac{1}{4}\beta_3 + \frac{1}{2}\beta_7$	-	-
	3-brane	(1,3) (0 0 0 2)	β ₇	-	-
		(15,1) (1 0 1 0)	$\beta_1 + \alpha_2 + \beta_3$	-	-
	4-brane	(20,2) (1 1 0 1)	$\frac{5}{4}\beta_1 + \frac{3}{2}\alpha_2 + \frac{3}{4}\beta_3 + \frac{1}{2}\beta_7$	-	-
	5-brane	(10,3) (2 0 0 2)	$\frac{3}{2}\beta_1 + \alpha_2 + \frac{1}{2}\beta_3 + \beta_7$	-	-
		(64 , 1) (1 1 1 0)	$\frac{3}{2}\beta_1 + 2\alpha_2 + \frac{3}{2}\beta_3$	-	-

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Conclusions

dim.	brane	repr.	highest weight	restricted repr.	#
<i>D</i> = 3	0-brane	248 (1 0 0 0 0 0 0 0 0)	$2\alpha_1 + 3\alpha_2 + 4\alpha_3 + 5\beta_4 + 6\beta_5 + 4\beta_6 + 2\alpha_7 + 3\beta_8$	52 (1 0 0 0)	24
	1-brane	3875 (0 0 0 0 0 0 1 0)	$2\alpha_1 + 4\alpha_2 + 6\alpha_3 + 8\beta_4 + 10\beta_5 + 7\beta_6 + 4\alpha_7 + 5\beta_8$	324 (0 0 0 2)	24
	2-brane	147250 (0 0 0 0 0 0 0 1)	$3\alpha_1 + 6\alpha_2 + 9\alpha_3 + 12\beta_4 + 15\beta_5 + 10\beta_6 + 5\alpha_7 + 8\beta_8$	-	-
<i>D</i> = 4	0-brane	56 (1 0 0 0 0 0 0)	$\frac{3}{2}\alpha_2 + 2\alpha_3 + \frac{5}{2}\beta_4 + 3\beta_5 + 2\beta_6 + \alpha_7 + \frac{3}{2}\beta_8$	14 (1 0 0)	8
	1-brane	133 (0 0 0 0 0 1 0)	$\alpha_2 + 2\alpha_3 + 3\beta_4 + 4\beta_5 + 3\beta_6 + 2\alpha_7 + 2\beta_8$	21 (0 0 2)	6
	2-brane	912 (0 0 0 0 0 0 1)	$\frac{3}{2}\alpha_2 + 3\alpha_3 + \frac{9}{2}\beta_4 + 6\beta_5 + 4\beta_6 + 2\alpha_7 + \frac{7}{2}\beta_8$	-	-
	3-brane	8645 (0 0 0 0 1 0 0)	$2\alpha_2+4\alpha_3+6\beta_4+8\beta_5+6\beta_6+3\alpha_7+4\beta_8$	-	-
D = 5	0-brane	27 (1 0 0 0 0 0)	$\frac{4}{3}\alpha_3 + \frac{5}{3}\beta_4 + 2\beta_5 + \frac{4}{3}\beta_6 + \frac{2}{3}\alpha_7 + \beta_8$	6 (2 0)	3
	1-brane	27 (0 0 0 0 1 0)	$\frac{2}{3}\alpha_3 + \frac{4}{3}\beta_4 + 2\beta_5 + \frac{5}{3}\beta_6 + \frac{4}{3}\alpha_7 + \beta_8$	6 (0 2)	3
	2-brane	78 (0 0 0 0 0 1)	$\alpha_3+2\beta_4+3\beta_5+2\beta_6+\alpha_7+2\beta_8$	-	-
	3-brane	351 (0 0 0 1 0 0)	$\frac{4}{3}\alpha_3 + \frac{8}{3}\beta_4 + 4\beta_5 + \frac{10}{3}\beta_6 + \frac{5}{3}\alpha_7 + 2\beta_8$	-	-
	4-brane	1728 (0 0 0 0 1 1)	$\frac{5}{3}\alpha_3 + \frac{10}{3}\beta_4 + 5\beta_5 + \frac{11}{3}\beta_6 + \frac{7}{3}\alpha_7 + 3\beta_8$	-	-
<i>D</i> = 6	0-brane	16 (1 0 0 0 0)	$\frac{5}{4}\beta_4 + \frac{3}{2}\beta_5 + \beta_6 + \frac{1}{2}\alpha_7 + \frac{3}{4}\beta_8$	-	-
	1-brane	10 (0 0 0 1 0)	$\frac{1}{2}\beta_4 + \beta_5 + \beta_6 + \alpha_7 + \frac{1}{2}\beta_8$	3 (2)	2
	2-brane	16 (0 0 0 0 1)	$\frac{3}{4}\beta_4 + \frac{3}{2}\beta_5 + \beta_6 + \frac{1}{2}\alpha_7 + \frac{5}{4}\beta_8$	-	-
	3-brane	45 (0 0 1 0 0)	$\beta_4 + 2\beta_5 + 2\beta_6 + \alpha_7 + \beta_8$	-	-
	4-brane	144 (0 0 0 1 1)	$\frac{5}{4}\beta_4 + \frac{5}{2}\beta_5 + 2\beta_6 + \frac{3}{2}\alpha_7 + \frac{7}{4}\beta_8$	-	-
	5-brane	320 (0 0 1 1 0)	$\frac{3}{2}\beta_4 + 3\beta_5 + 3\beta_6 + 2\alpha_7 + \frac{3}{2}\beta_8$	-	-

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Conclusions

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Hypermultiplet sector: we consider in particular the $SO(4, n_H)$ case

We obtain the brane content (including domain walls and space-filling branes) by requiring that the gauge algebra is a truncation of the gauge algebra of the half-maximal theory

e.g. in three dimensions: $SO(8, n + m) \supset SO(4, n) \times SO(4, m)$

By analogy with the previous analysis, we expect the number of branes to be the same as for the $F_{4(4)}$, $E_{6(2)}$, $E_{7(-5)}$ and $E_{8(-24)}$

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Conclusions

dim.	type of brane	field	# of branes
<i>D</i> = 3	0-brane	A_{1,M_1M_2}	24
	1-brane	A _{2,MN}	8
		A_{2,M_1M_4}	16
		$A_{2,A_1A_2M_1M_2}$	576
	2-brane	$A_{3,MN_1N_2N_3A_1A_2}$	2304
		$A_{3,AB_1B_2B_3M_1M_2}$	2304
<i>D</i> = 4	1-brane	A_{2,M_1M_2}	24
	2-brane	A_{3,M_1M_2Aa}	192
	3-brane	$A_{4,M_1M_2A_1A_2ab}$	192
		$A_{4,MN_1N_2N_3}$	96
		A_{4,ABM_1M_2}	96
<i>D</i> = 5	2-brane	A_{3,M_1M_2} ($\alpha = -2$)	24
	3-brane	$A_{4,AM_1M_2} (\alpha = -2)$	48
		$A_{4,M_1M_2} \ (\alpha = -4)$	24
<i>D</i> = 6	3-brane	A_{4,M_1M_2} ($\alpha = -2$)	24

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Conclusions

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On branes in $\mathcal{N} = 2$ supergravities

Fabio Riccioni

- Branes in maximal supergravities
- Supergravities with less supersymmetry
- Real forms and Tits-Satake diagrams
- Branes in $\mathcal{N} = 2$ supergravities
- Conclusions

- Our analysis reveals a universal structure underlying the various theories with eight supercharges that can be uplifted to six dimensions
- Low-energy of heterotic theory: wrapping rules are satisfied as in the maximal and half-maximal theories
- Central charges and degeneracies
- We find vector 3-branes in four dimensions
- Many of the branes we find are exotic!