

# Study of $K$ - $pp$ with an effective $K^{\text{bar}}N$ potential on coupled-channel Complex Scaling Method

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1. Introduction
2. Effective single-channel potential  
*“Feshbach method with  
coupled-channel Complex Scaling Method”*
3. Application to 2-body system:  
“Resonance of  $K^{\text{bar}}N$ - $\pi\Sigma = \Lambda^*$ ”
4. Application to 3-body system: “ $K$ - $pp$ ”
5. Summary and future plans

# *1. Introduction*

*Kaonic nuclei = Nuclear system with anti-kaon “ $K^{\text{bar}}$ ” ( $K^-$ ,  $K^{0\text{bar}}$ )*

$\Lambda(1405) \sim K\text{-}p$  quasi-bound state  $\rightarrow$  Important building block  
 $K^-pp$  ... A prototype of  $K^{\text{bar}}$  nuclei

*$K^-pp$  has been studied theoretically and experimentally*

*Experiment: FINUDA, DISTO*

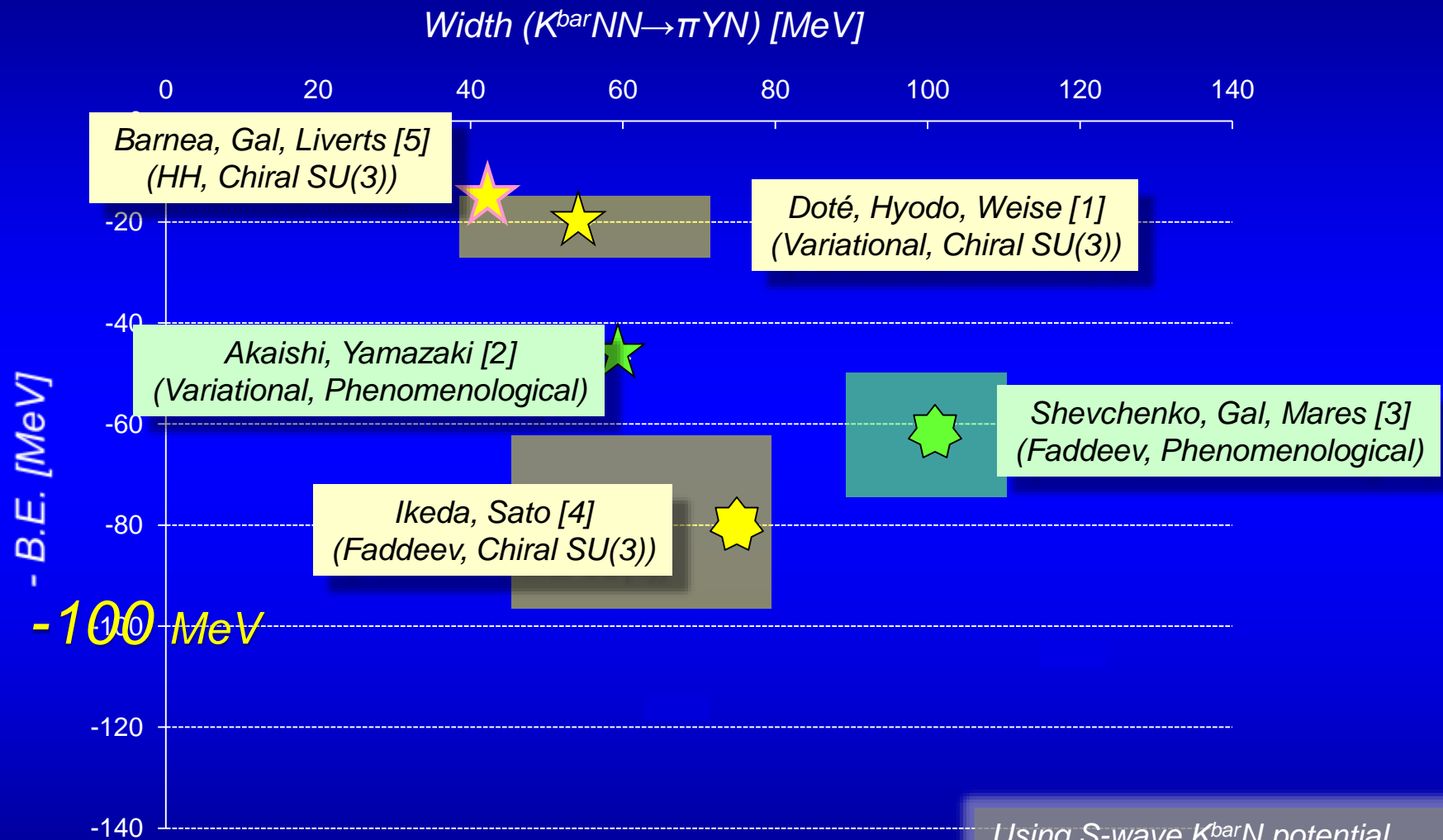
*J-PARC: E27  $d(\pi^+, K^+)$ , E15  $^3\text{He}(\text{in-flight } K^-, n)$*

*LEPS:  $d(\gamma, K^+)$ ,  $d(\gamma, K^+\pi^-)$*

*Theory: Faddeev-AGS / Variational method (Gauss base, Hyperspherical Harmonics)*

*Phenomenological / Chiral SU(3)-based  $K^{\text{bar}}N$  potential*

# Typical results of theoretical studies of $K^-pp$



[1] PRC79, 014003 (2009)

[5] PLB94, 712 (2012)

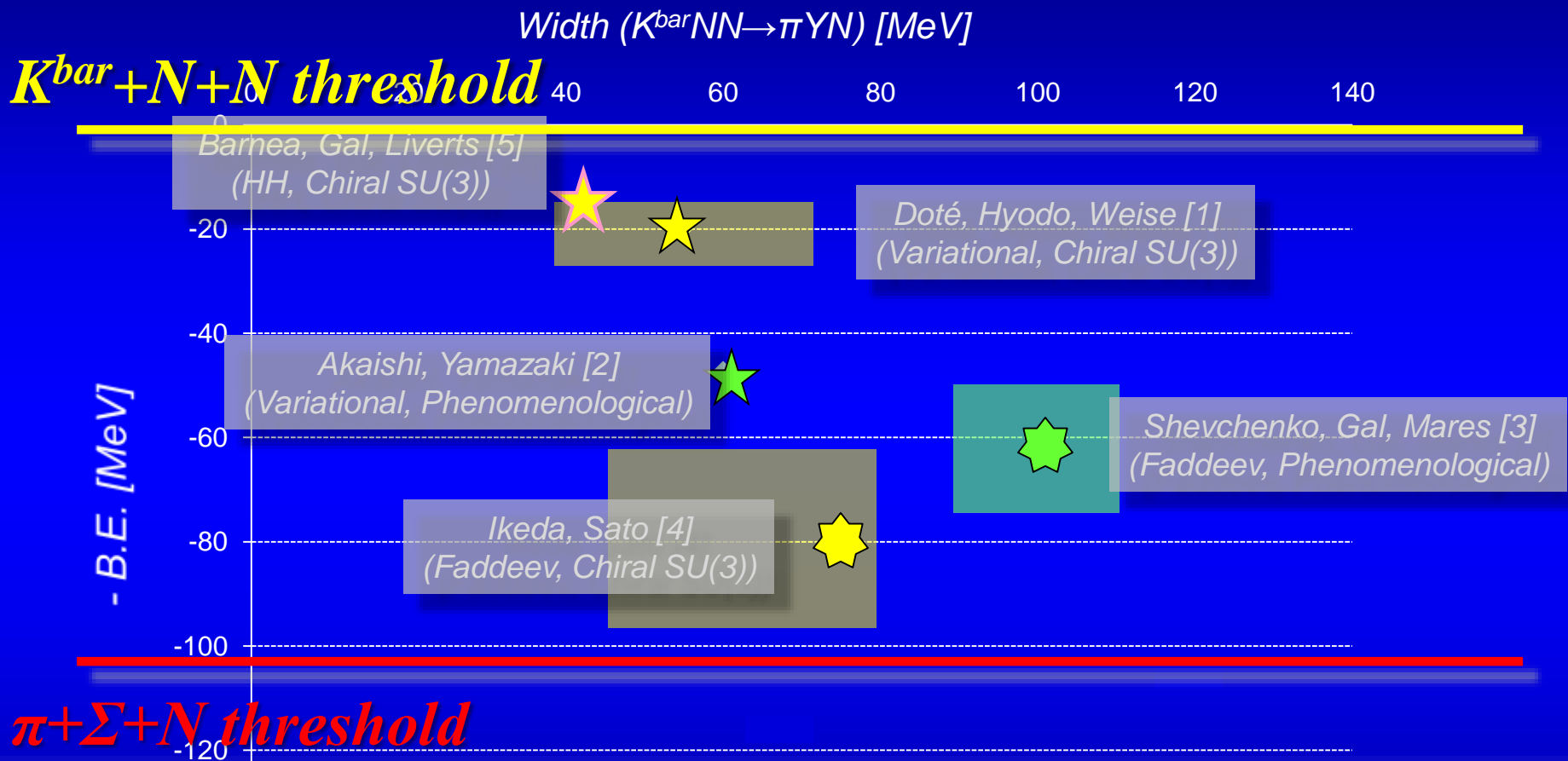
[2] PRC76, 045201 (2007)

[3] PRC76, 044004 (2007)

[4] PRC76, 035203 (2007)

Using S-wave  $K^{\text{bar}}N$  potential constrained by experimental data. ...  $K^{\text{bar}}N$  scattering data, Kaonic hydrogen atom data, " $\Lambda(1405)$ " etc.

# Typical results of theoretical studies of $K^-pp$



From theoretical viewpoint,  
 $K^-pp$  exists between  $K^{\text{bar}}-N-N$  and  $\pi-\Sigma-N$  thresholds!

[4] PRC76, 035203 (2007)

Atomic hydrogen atom data,  
 “ $\Lambda(1405)$ ” etc.

Kaonic nuclei = Nuclear system with anti-kaon “ $K^{\text{bar}}$ ” ( $K^-$ ,  $K^{0\text{bar}}$ )

$\Lambda(1405)$  = Important building block

$K^-pp$  = a prototype of  $K^{\text{bar}}$  nuclei

$K^-pp$  has been studied theoretically and experimentally

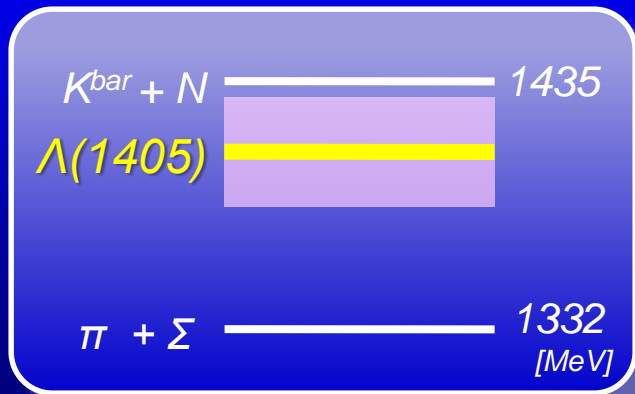
Experiment: FINUDA, DISTO

J-PARC: E27  $d(\pi^+, K^+)$ , E15  $^3\text{He}(\text{in-flight } K^-, n)$

LEPS:  $d(\gamma, K^+)$ ,  $d(\gamma, K^+\pi)$

Theory: Faddeev-AGS / Variational method (Gauss base, Hyperspherical Harmonics)

Phenomenological / Chiral SU(3)-based  $K^{\text{bar}}N$  potential



Resonant state

and

a coupled-channel system

(  $K^{\text{bar}}N(N) - \pi Y(N)$  )

$K^{bar} + N + N \dots$

$K^{bar}$  nucleus

$\pi + \Sigma + N \dots$

Resonant state

and

a coupled-channel system

(  $K^{bar}N(N) - \pi Y(N)$  )

*“coupled-channel Complex Scaling Method”*

- ✓ Consider a coupled-channel problem.
- ✓ Treat resonant states adequately.
- ✓ Obtain the wave function to help the analysis of the state.
- ✓ Applicable to many-body systems.

## *coupled-channel Complex Scaling Method*

- ✓ Applied to  $K^{\text{bar}}N\text{-}\pi Y$  system
- ✓ Treat resonant and scattering states with Gaussian base like bound-states study  
“CSWF” method for scattering problem

A. T. Kruppa, R. Suzuki and K. Katō, PRC 75, 044602 (2007)

## Chiral SU(3)-based potential (KSW-type potential)

N. Kaiser, P. B. Siegel and W. Weise,  
NPA 594, 325 (1995)

- Based on Chiral SU(3) theory  
→ **Energy dependence**
- WT term,  $r$ -space, Gaussian form
- Semi-rela. / Non-rela.

$$V_{ij}^{(I=0,1)}(r) = -\frac{C_{ij}^{(I=0,1)}}{8f_\pi^2} (\omega_i + \omega_j) \sqrt{\frac{M_i M_j}{s \omega_i \omega_j}} g_{ij}(r)$$

$$g_{ij}(r) = \frac{1}{\pi^{3/2} d_{ij}^3} \exp\left[-(r/d_{ij})^2\right]$$

## Constrained by $K^{\text{bar}}N$ scattering length

$$a_{KN(I=0)} = -1.70 + i0.67 \text{ fm}, \quad a_{KN(I=1)} = 0.37 + i0.60 \text{ fm}$$

A.D.Martin, NPB179, 33(1979)



# Complex Scaling Method for Resonance

## Complex rotation of coordinate (Complex scaling)

$$U(\theta): \mathbf{r} \rightarrow \mathbf{r} e^{i\theta}, \quad \mathbf{k} \rightarrow \mathbf{k} e^{-i\theta}$$

$$H_\theta \equiv U(\theta) H U^{-1}(\theta), \quad |\Phi_\theta\rangle \equiv U(\theta) |\Phi\rangle$$

## By Complex scaling,

➤ Resonance wave function: divergent function  $\Rightarrow$  damping function



*Boundary condition is the same as that for a bound state.*

➤ *The pole position of resonance doesn't change.*

### ABC theorem

*“The energy of bound and resonant states is independent of scaling angle  $\theta$ .”*

† J. Aguilar and J. M. Combes, Commun. Math. Phys. 22 (1971),269.

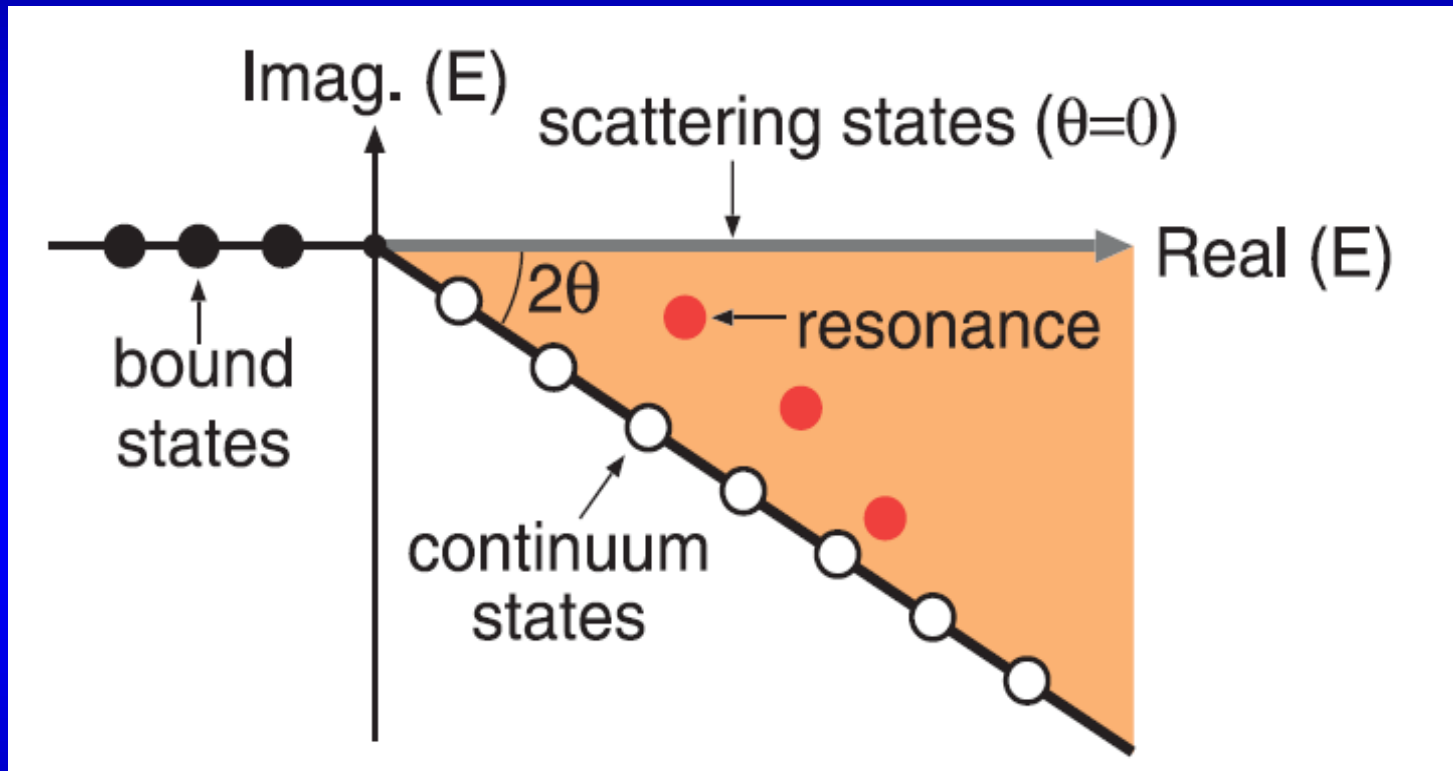
E. Balslev and J. M. Combes, Commun. Math. Phys. 22 (1971),280

## Diagonalize $H_\theta$ with Gaussian base,

*we can obtain resonant states, in the same way as bound states!*

# Complex Scaling Method for Resonance

*Eigenvalue distribution after complex scaling*



- *Continuum state appears on  $2\vartheta$  line.*
- *Resonance pole is off from  $2\vartheta$  line, and independent of  $\vartheta$ .*

***2. Effective potential  
for single-channel calculation***

*“Feshbach method with ccCSM”*

# Feshbach method

$H \Phi = E \Phi$  : Schrödinger eq. in  $P$  and  $Q$  spaces

$$\begin{pmatrix} T_P + v_P & V_{PQ} \\ V_{QP} & T_Q + v_Q \end{pmatrix} \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix} = E \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix}$$

$P$ : model space

$Q$ : out of model space



Schrödinger eq. in  $P$  space :

$$\left( T_P + U_P^{\text{Eff}}(E) \right) \Phi_P = E \Phi_P$$

Effective potential

$$\begin{aligned} U_P^{\text{Eff}}(E) &= v_P + V_{PQ} G_Q(E) V_{QP} \\ &= v_P + V_{PQ} \frac{1}{E - H_{QQ}} V_{QP} \end{aligned}$$

# Green function expressed with ECR<sup>†</sup>

$$G_Q(E) = \frac{1}{E - H_{QQ}}$$

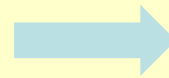


† Y. Kikuchi, T. Myo, M. Takashina, K. Kato and K. Ikeda, PTP122, 499 (2009); PRC81, 044308(2010)

Complex scaling  $U(\theta): r \rightarrow r e^{i\theta}, k \rightarrow k e^{-i\theta}$

$$H_{QQ}^\theta = U(\theta) H_{QQ} U^{-1}(\theta)$$

$$G_Q^\theta(E) = \frac{1}{E - H_{QQ}^\theta}$$



$$G_Q(E) = U^{-1}(\theta) G_Q^\theta(E) U(\theta)$$

## Extended Completeness Relation

$$H_{QQ}^\theta |\chi_n^\theta\rangle = \varepsilon_n^\theta |\chi_n^\theta\rangle$$

$$\int_C \sum_{R+B} |\chi_n^\theta\rangle \langle \chi_n^\theta| = 1$$

Diagonalize  $H_{QQ}^\theta$  with Gaussian base,

$$\sum_n |\chi_n^\theta\rangle \langle \chi_n^\theta| \approx 1$$

Well approximated

T. Myo, A. Ohnishi and K. Kato, PTP99, 801 (1998)

$$G_Q^\theta(E) \approx \sum_n |\chi_n^\theta\rangle \frac{1}{E - \varepsilon_n^\theta} \langle \chi_n^\theta|$$



$$U_P^{Eff}(E) = v_P + V_P(E)$$

✓ Non-local  
✓ Energy dependent

$$V_P(E) = V_{PQ} G_Q(E) V_{QP}$$

$$\approx \sum_n V_{PQ} U^{-1}(\theta) |\chi_n^\theta\rangle \frac{1}{E - \varepsilon_n^\theta} \langle \chi_n^\theta| U(\theta) V_{QP}$$

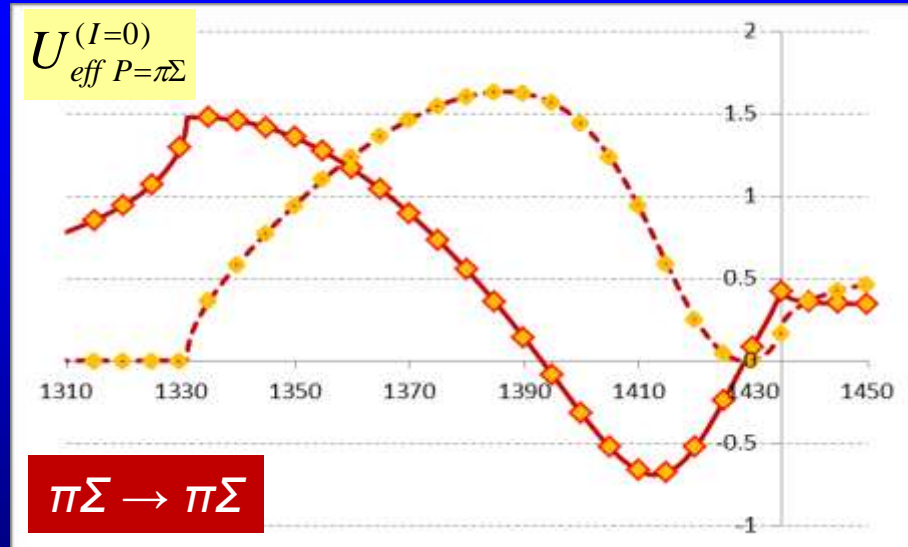
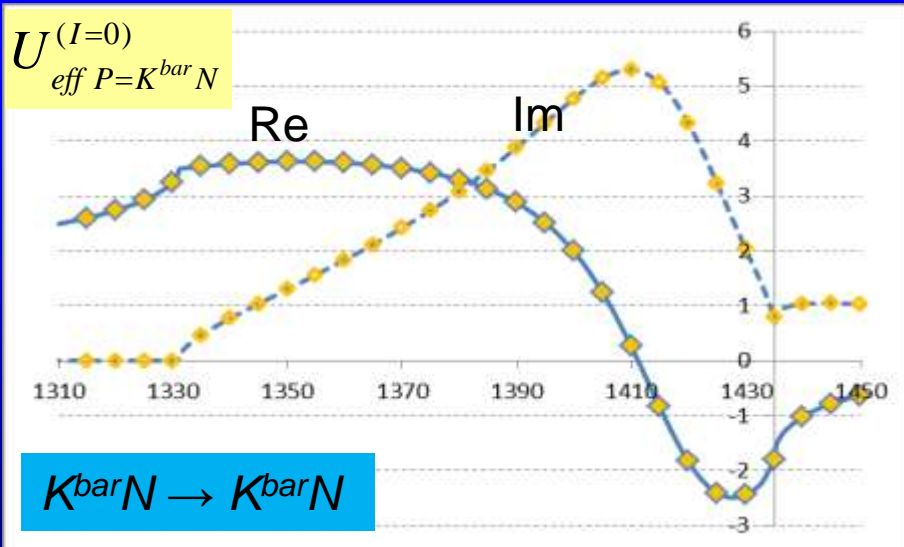
# 3. Application to 2-body system

“Resonance of  $K^{\text{bar}}N\text{-}\pi\Sigma(I=0) = \Lambda^*$ ”

Confirmed that

“Feshbach with ccCSM” works well for scattering states,  
But how is resonance case?

- Scattering amplitude



KSW - NRv2 ( $I=0$ ),  $f_\pi=110$  MeV

# AY potential<sup>†</sup> (Non-rela. / E-indep.)

<sup>†</sup>Y. Akaishi and T. Yamazaki,  
PRC 52 (2002) 044005

Schrödinger eq. in P space :  $(T_P + U_P^{Eff}(Z))\Phi_P = Z\Phi_P$

Resonance → Self-consistency for complex energy “Z”

## Feshbach+ccCSM

P-space	KN + π Σ	KN	π Σ
B(KN)	28.1698	28.1698	28.1698
Γ / 2	20.0288	20.0288	20.0289
[MeV]			
Mean distance			
KN	1.31 - i0.35	1.25 - i0.27	*
π Σ	0.31 - i0.21	*	0.21 + i0.91
Total	1.34 - i0.39	1.25 - i0.27	0.21 + i0.91
[fm]			

# KSW-type potentials (E-dep.)

$$V_{ij}^{(I=0)}(r) \sim -\frac{C_{ij}^{(I=0)}}{8f_\pi^2} (\omega_i + \omega_j) \times g_{ij}(r)$$

- NRv2  
(Non-rela. / E-dep.)

- SR-A  
(Semi-rela. / E-dep.)


## Feshbach+ccCSM

P-space	KN+π Σ	KN	π Σ
B(KN)	17.1605	17.1606	17.1728
Γ / 2	16.6176	16.6178	16.613
[MeV]			
Mean distance			
KN	1.37 - i0.37	1.28 - i0.40	*
π Σ	0.37 + i0.04	*	0.23 + i0.93
Total	1.42 - i0.34	1.28 - i0.40	0.23 + i0.93
[fm]			

## Feshbach+ccCSM

P-space	KN+π Σ	KN	π Σ
B(KN)	15.5336	15.5336	15.524
Γ / 2	25.0158	25.015	24.9997
[MeV]			
Mean distance			
KN	1.22 - i0.47	1.07 - i0.46	*
π Σ	0.13 + i0.05	*	0.05 - i0.27
Total	1.22 - i0.47	1.07 - i0.46	0.05 - i0.27
[fm]			

$$(T_P + U_P^{Eff}(Z)) \Phi_P = Z \Phi_P$$



$$v_{ij}(Z)$$

*Even when the original interaction has energy dependence, a self-consistent solution with complex energy can be obtained.*



# 4. Application to 3-body system

“ $K^-pp$ ”

...  $K^{\text{bar}}NN - \pi\Sigma N - \pi\Lambda N$  ( $J^\pi=0^-, T=1/2$ )

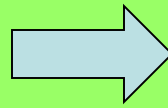
# 4. Application to 3-body system

“ $K^-pp$ ”

...  $K^{bar}NN - \pi\Sigma N - \pi\Lambda N$  ( $J^\pi=0^-, T=1/2$ )

Feshbach + ccCSM

$$\begin{array}{l} V(K^{bar}N - \pi Y; I=0,1) \\ V(\pi Y - \pi Y'; I=0,1) \end{array}$$



$$U_{K^{bar}N(I=0,1)}^{Eff}(E)$$

Schrödinger eq. for  $K^{bar}NN$  channel :

$$\left( T_{K^{bar}NN} + V_{NN} + \sum_{i=1,2} U_{K^{bar}N_i(I)}^{Eff}(E_{K^{bar}N}) \right) \Phi_{K^{bar}NN} = E \Phi_{K^{bar}NN}$$

# How to solve ...

- Wave function

$$| "K^- pp" \rangle = \sum_a C_a^{(KNN,1)} \left\{ G_a^{(KNN,1)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) + G_a^{(KNN,1)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN} = 0\rangle \left[ [K[NN]_1]_{T=1/2} \right]$$

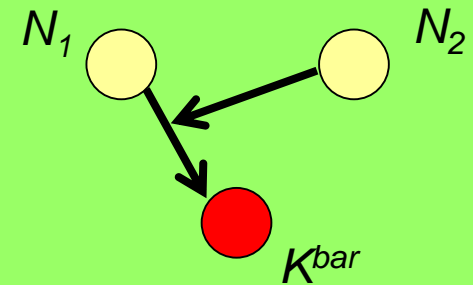
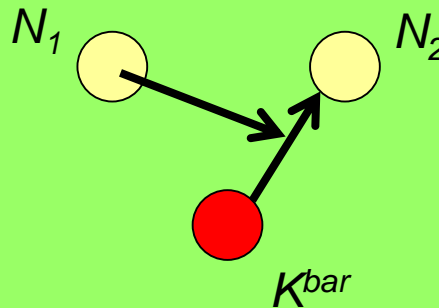
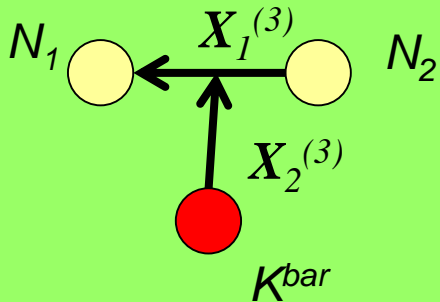
Ch. 1:  $K^{bar}NN$ ,  $NN:1E$

$$+ \sum_a C_a^{(KNN,2)} \left\{ G_a^{(KNN,2)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) - G_a^{(KNN,2)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN} = 0\rangle \left[ [K[NN]_0]_{T=1/2} \right]$$

Ch. 2:  $K^{bar}NN$ ,  $NN:1O$

- Basis function = Correlated Gaussian

$$G_a^{(KNN,i)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) = N_a^{(KNN,i)} \exp \left[ -(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) A_a^{(KNN,i)} \begin{pmatrix} \mathbf{x}_1^{(3)} \\ \mathbf{x}_2^{(3)} \end{pmatrix} \right]$$

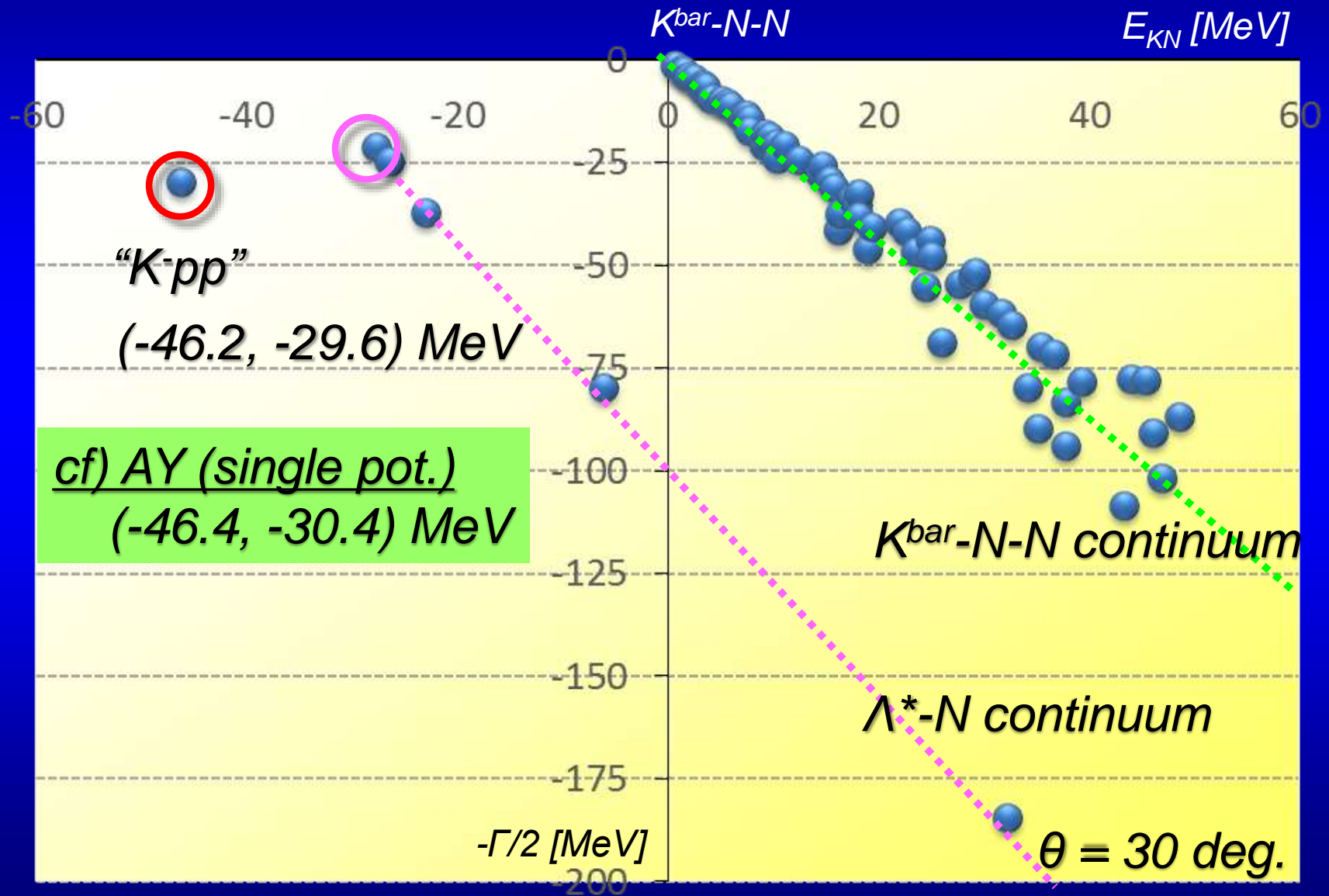


- Diagonalize the complex Hamiltonian with base functions

$$H_{K^{bar}NN} = T_{K^{bar}NN} + V_{NN} + \sum_{i=1,2} U_{K^{bar}N_i(I)}^{Eff}(E_{K^{bar}N})$$

# Test calculation with AY potential

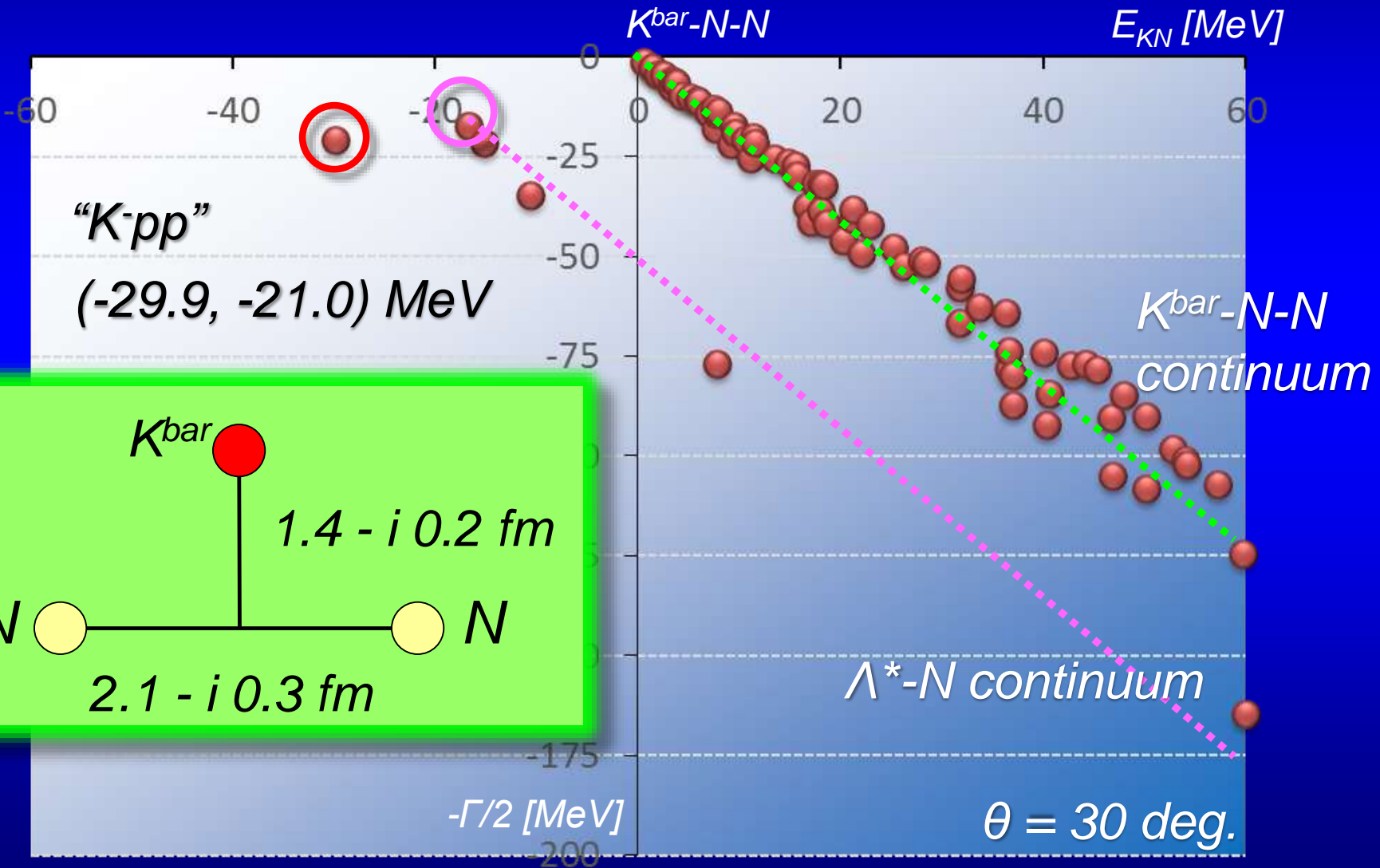
NN potential : Tamagaki potential (G3RS-case1)



# Result (preliminary): $K^{\bar{S}}-NRv2$ potential

NN potential : Av18 (Central + spin-spin)

$K^{\bar{S}}N$  energy is fixed at  $\Lambda^*$ . (Not self-consistent!)

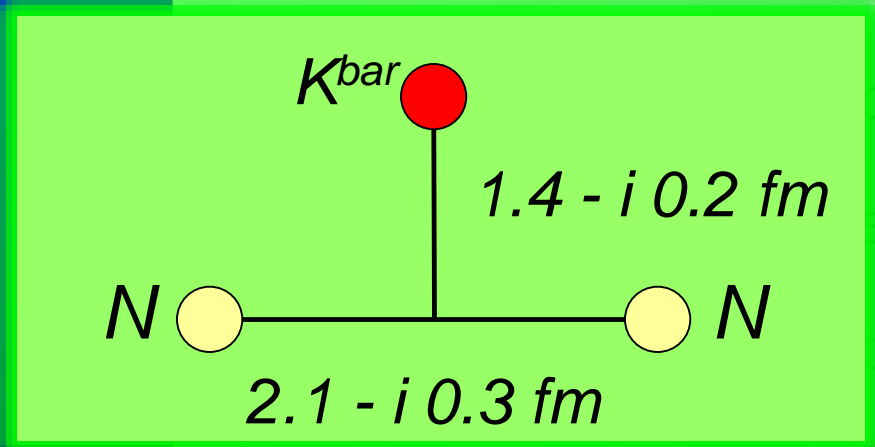


"K-pp"  
(-29.9, -21.0) MeV

$K^{\bar{S}}-N-N$   
continuum

$\Lambda^*-N$  continuum

$\theta = 30$  deg.



$-\Gamma/2$  [MeV]  
-175  
-200

***5. Summary  
and future plans***

# 5. Summary and future plans

The simplest  $K^{\text{bar}}$  nucleus “ $K\text{-}pp$ ” = Resonance state of  $K^{\text{bar}}\text{NN-}\pi\text{YN}$  coupled system

“Feshbach method + coupled-channel Complex Scaling Method”

... Represent the Green function in  $Q$ -space with the Extended Complete Set

⇒ Effective single-channel potential (for  $K^{\text{bar}}\text{N}$ , eliminate  $\pi\text{Y}$  channels)

Two-body case “ $K^{\text{bar}}\text{N-}\pi\Sigma = \Lambda^*$ ”

Reproduce perfectly the result of the full coupled-channel calculation (AY, KSW-NRv2)

Three-body case “ $K\text{-}pp$ ”

← Use **Correlated Gaussian base**, diagonalize the complex Hamiltonian

• AY potential: agree with Akaishi’s result

(B. E.,  $\Gamma$ ) = (46.4, 60.8) MeV with Tamagaki NN potential

• KSW-NRv2 potential:

Based on chiral SU(3) theory, Energy dependent, Non-relativistic version

Constraint by the  $K^{\text{bar}}\text{N}$  scattering length (A.D.Martin)

\*  $K^{\text{bar}}\text{N}$  energy is fixed to that of  $\Lambda^*$        $\Lambda^*$ : (B. E.,  $\Gamma$ ) = (17.2, 33.2) MeV

⇒ (B. E.,  $\Gamma$ ) ~ (30, 42) MeV with Av18 NN potential

1. Self consistency of  $K^{\text{bar}}\text{N}$  energy in the effective potential

2. Full coupled-channel calculation with the coupled-channel Complex Scaling Method