

Systematic study of deeply-bound pionic atom and future perspectives

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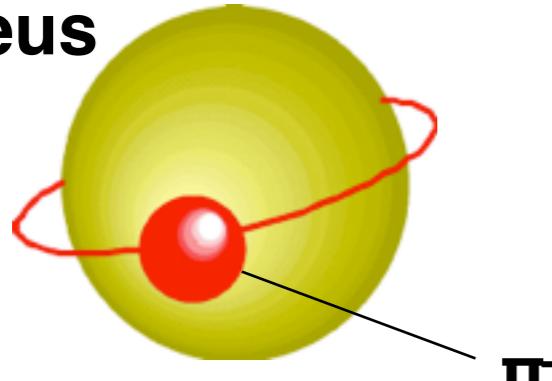
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- 2) Precise / systematic experiment in RIBF, RIKEN
- 3) Future perspectives ~ pionic unstable nuclei ~
- 4) Summary

1. Spectroscopy of deeply-bound pionic atom

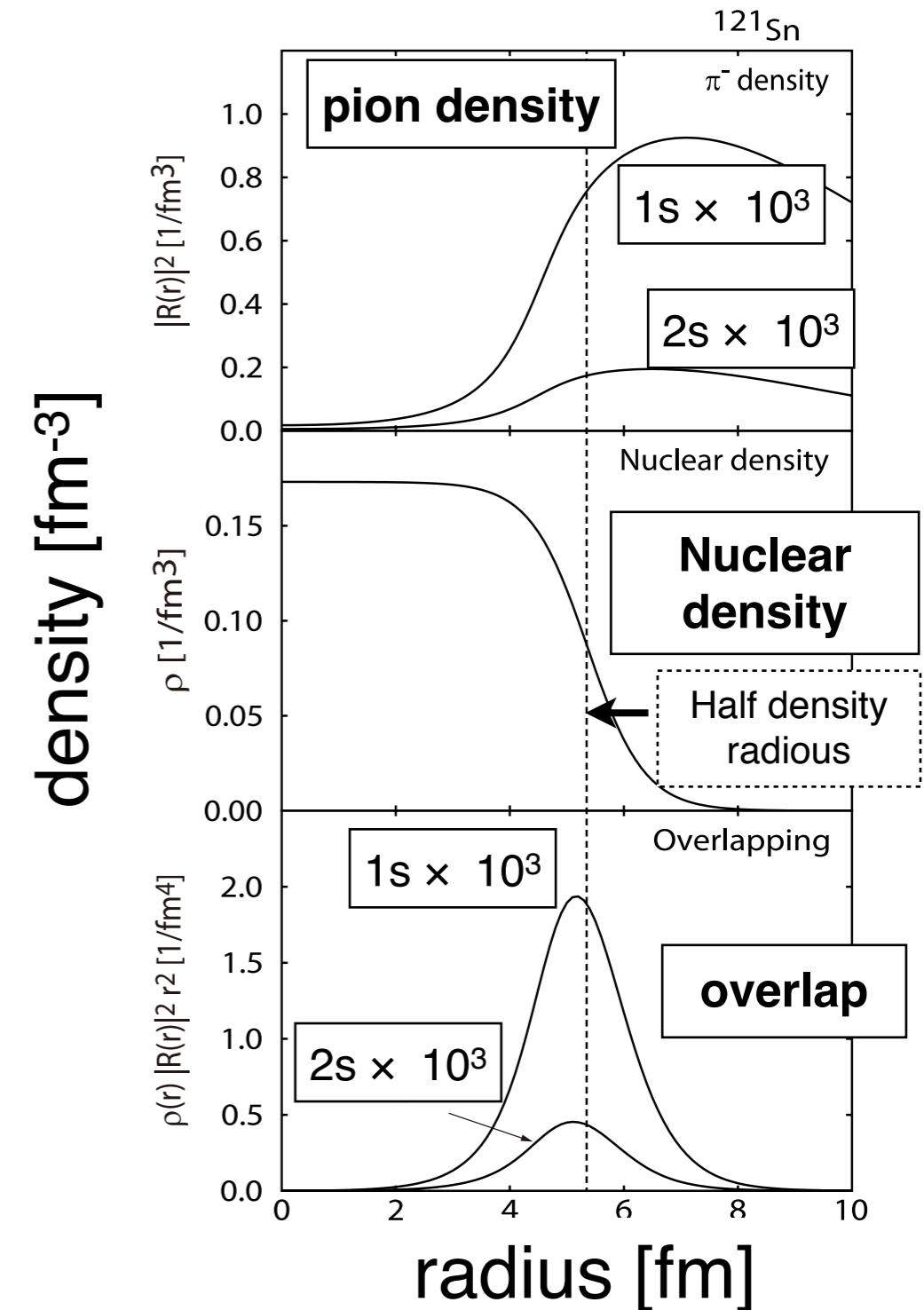
Deeply-bound pionic atom

Nucleus



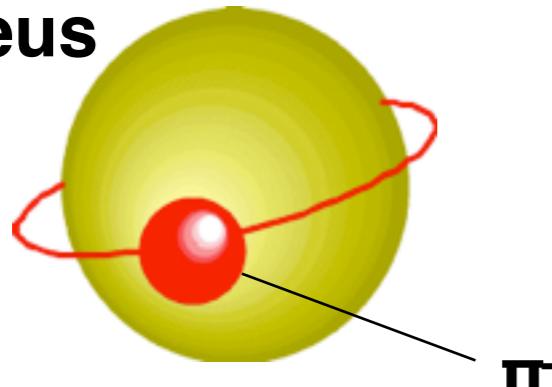
Coulomb
+
Strong

deep orbit in Heavy atom
= Large overlap between
pion and nucleus
→ probe for QCD in finite ρ



Deeply-bound pionic atom

Nucleus



Coulomb

+

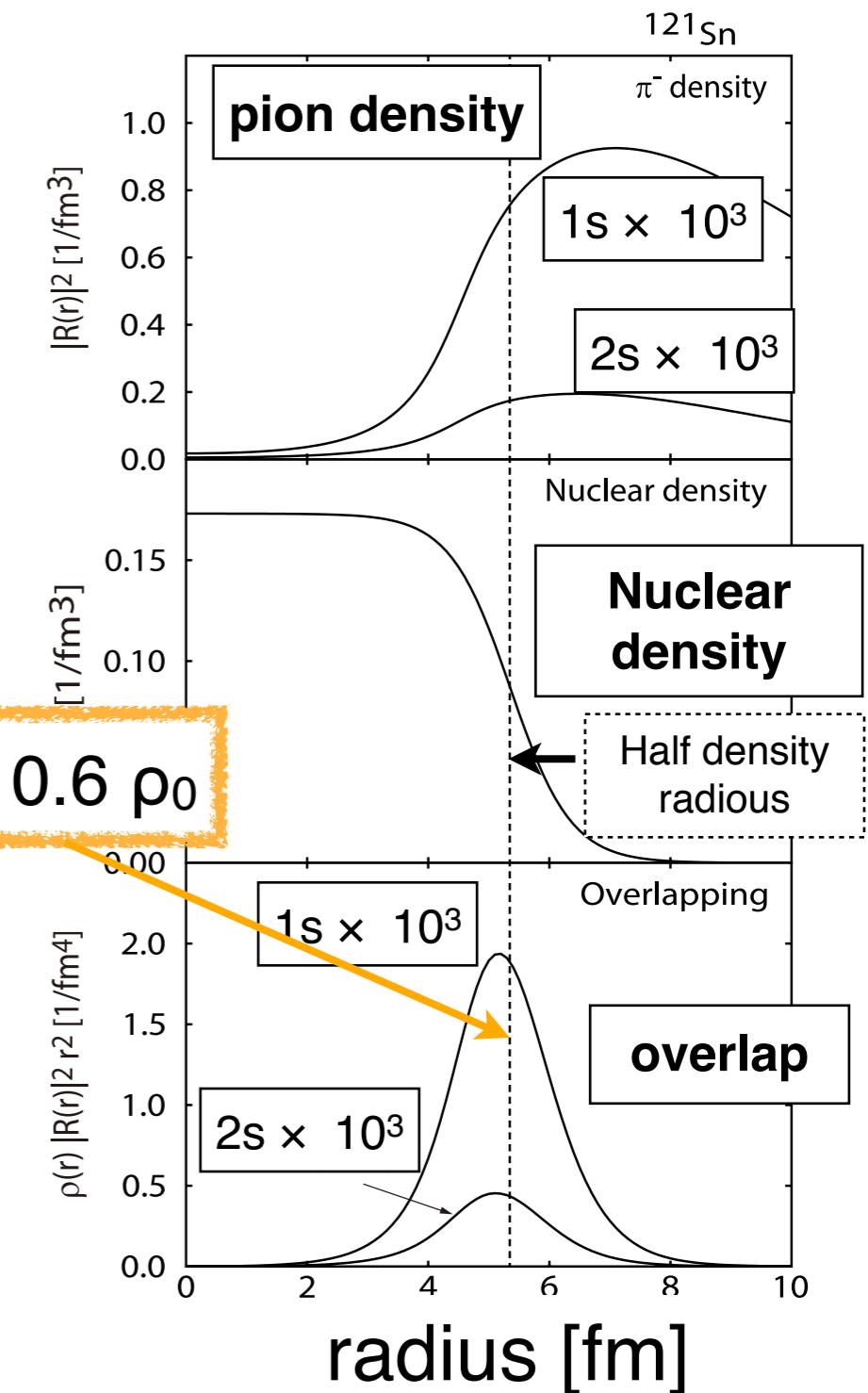
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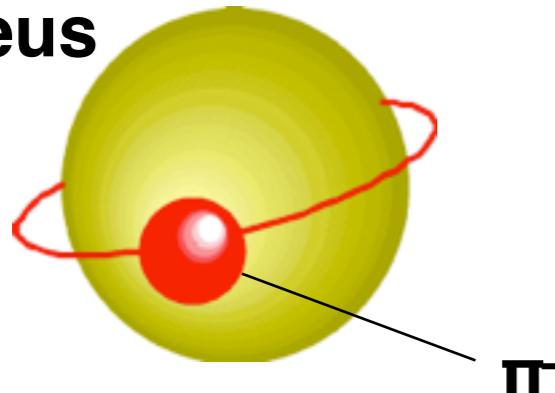
→ probe for QCD in finite ρ

\sim probing $0.6 \rho_0$



Deeply-bound pionic atom

Nucleus



Coulomb

+

Strong

deep orbit in Heavy atom

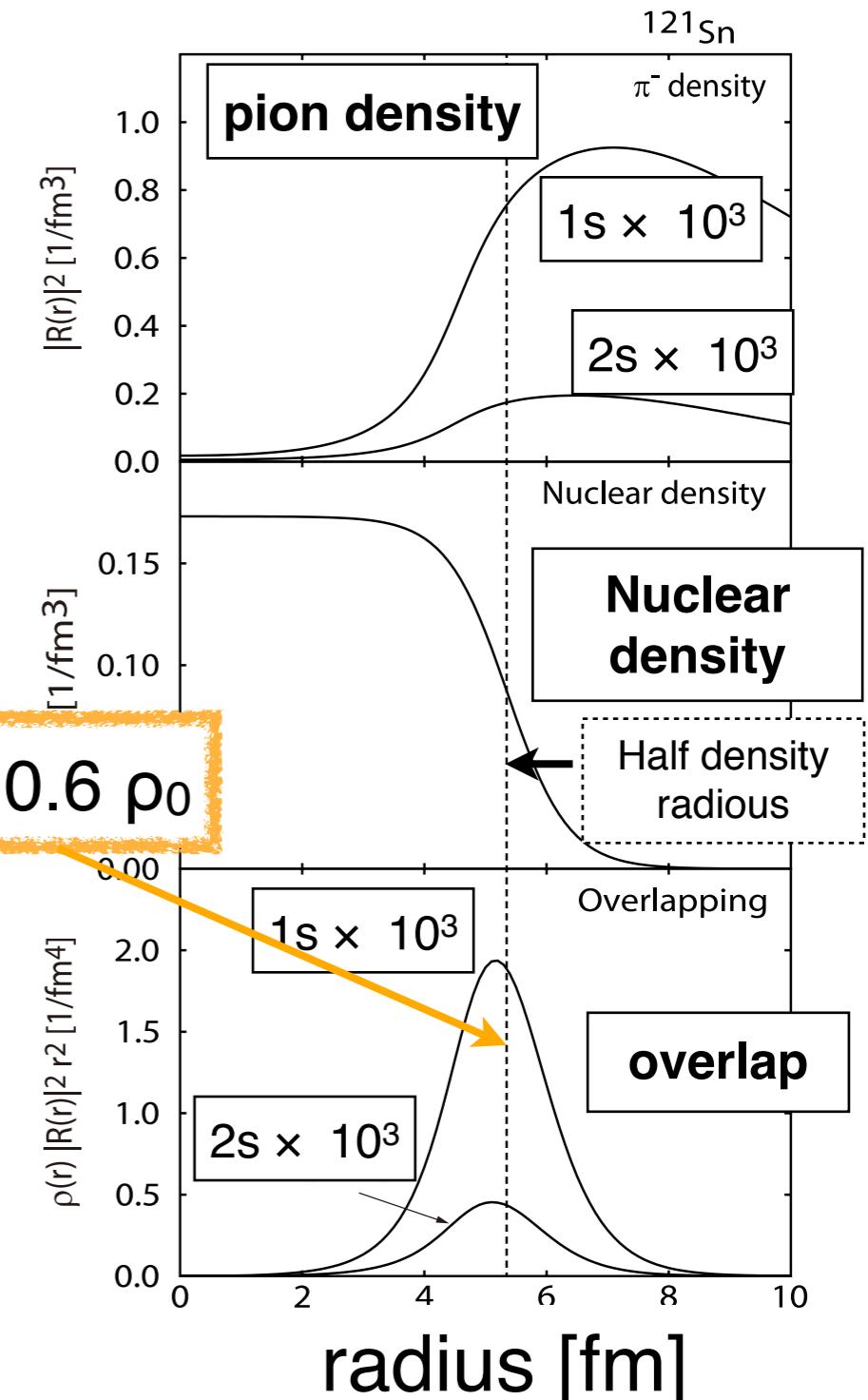
= Large overlap between
pion and nucleus

→ probe for QCD in finite ρ

E / Γ of pionic deep (1s, 2s, ...) states

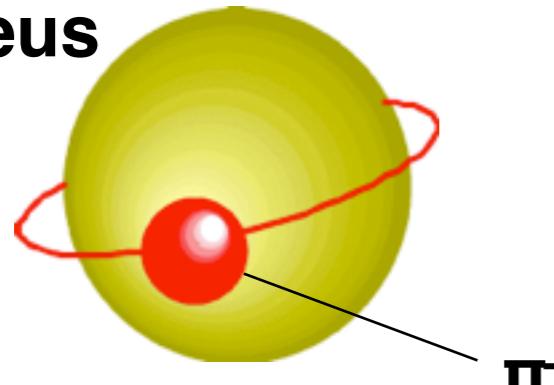
π-A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 b_0 \{\rho(r) + b_1 \delta\rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$



Deeply-bound pionic atom

Nucleus



Coulomb

+

Strong

deep orbit in Heavy atom

= Large overlap between
pion and nucleus

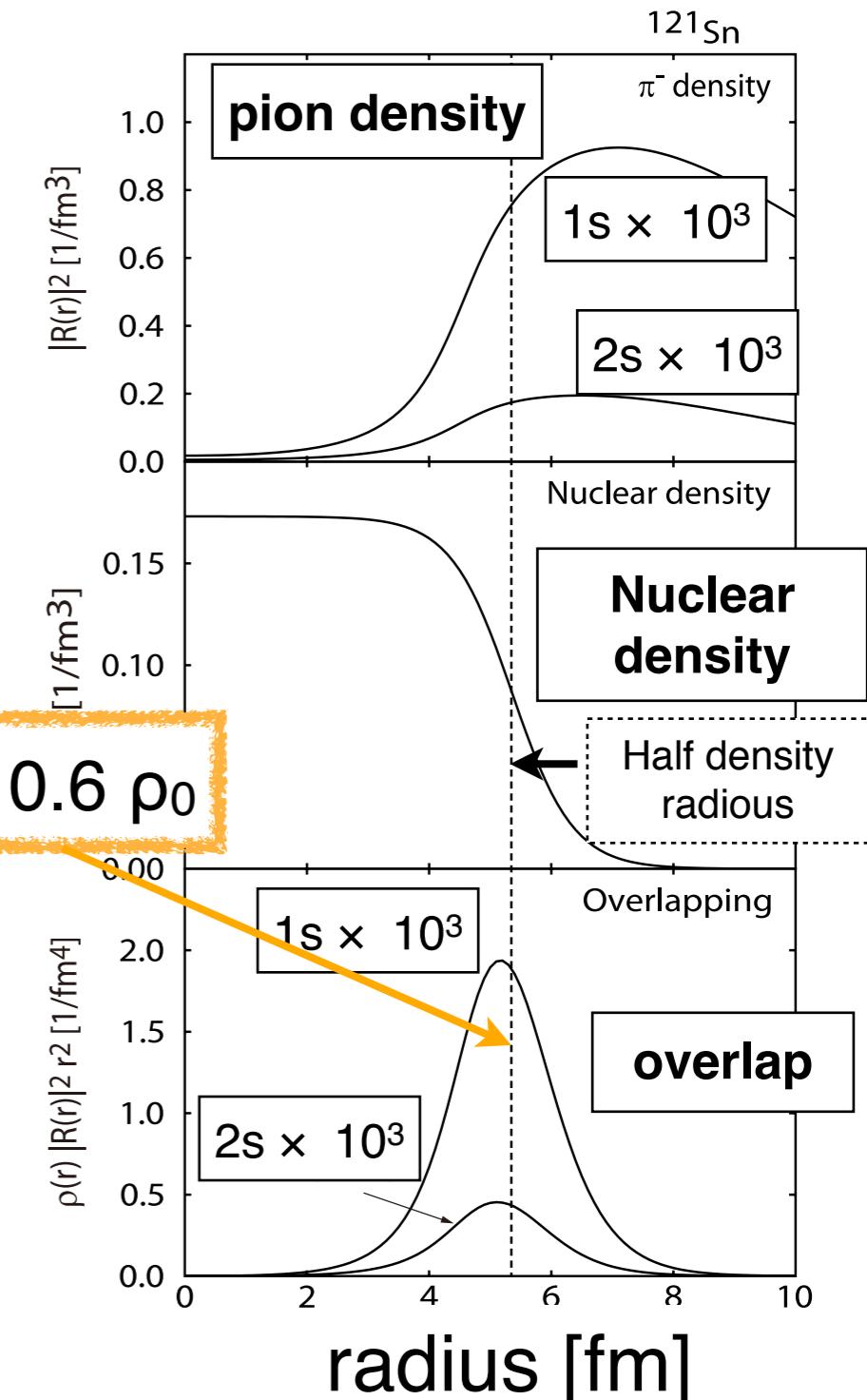
→ probe for QCD in finite ρ

E / Γ of pionic deep (1s, 2s, ...) states

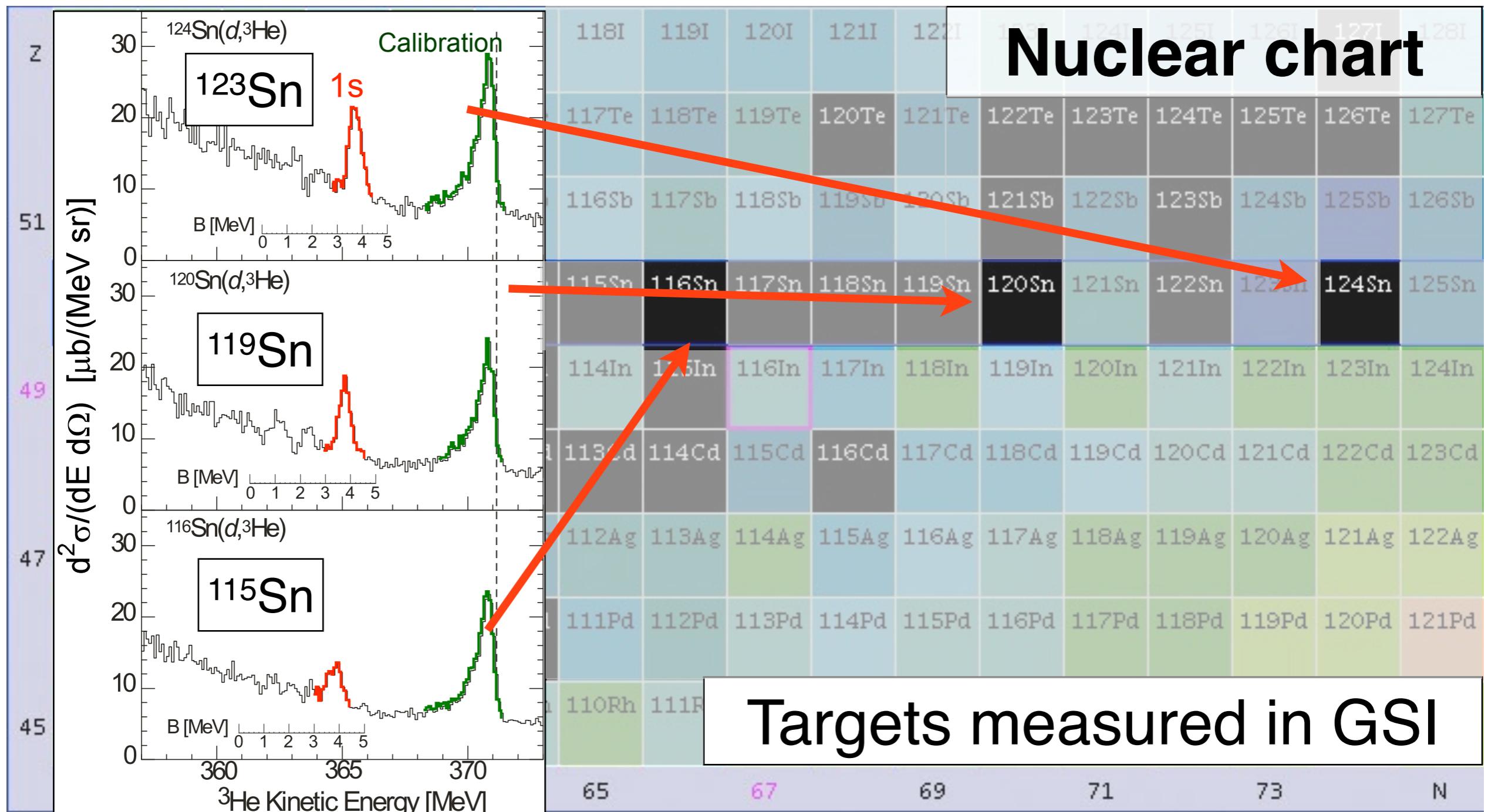
π-A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 b_0 \{\rho(r) + b_1 \delta\rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$

isovector scattering length



Experiment at GSI



Deduction of b_1 from experimental data

TABLE I. Observed binding energies (B_{1s}) and widths (Γ_{1s}) of the $1s \pi^-$ states in $^{115,119,123}\text{Sn}$ isotopes.

Isotope	B_{1s} (MeV)	Stat.	ΔB_{1s} (MeV) Syst.	Total	Γ_{1s} (MeV)	Stat.	$\Delta\Gamma_{1s}$ (MeV) Syst.	Total
^{115}Sn	3.906	± 0.021	± 0.012	± 0.024	0.441	± 0.068	± 0.054	± 0.087
^{119}Sn	3.820	± 0.013	± 0.012	± 0.018	0.326	± 0.047	± 0.065	± 0.080
^{123}Sn	3.744	± 0.013	± 0.012	± 0.018	0.341	± 0.036	± 0.063	± 0.072

K. Suzuki et al., PRL92 072302 (2004)

π -A s-wave optical potential

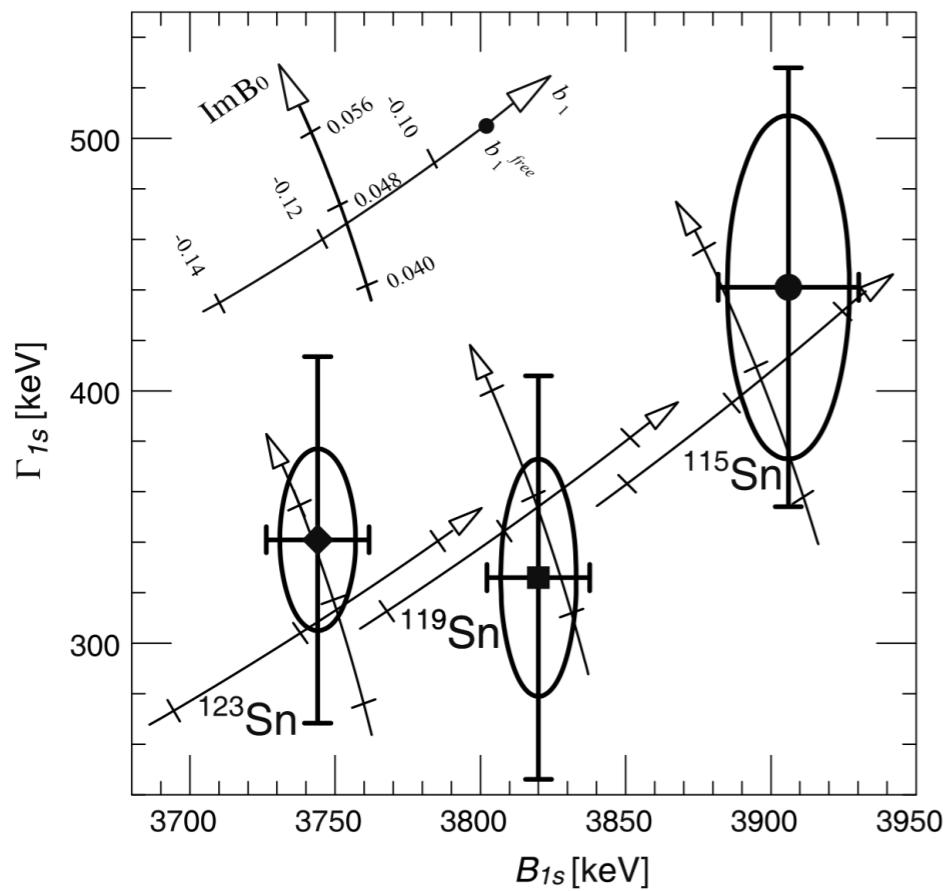
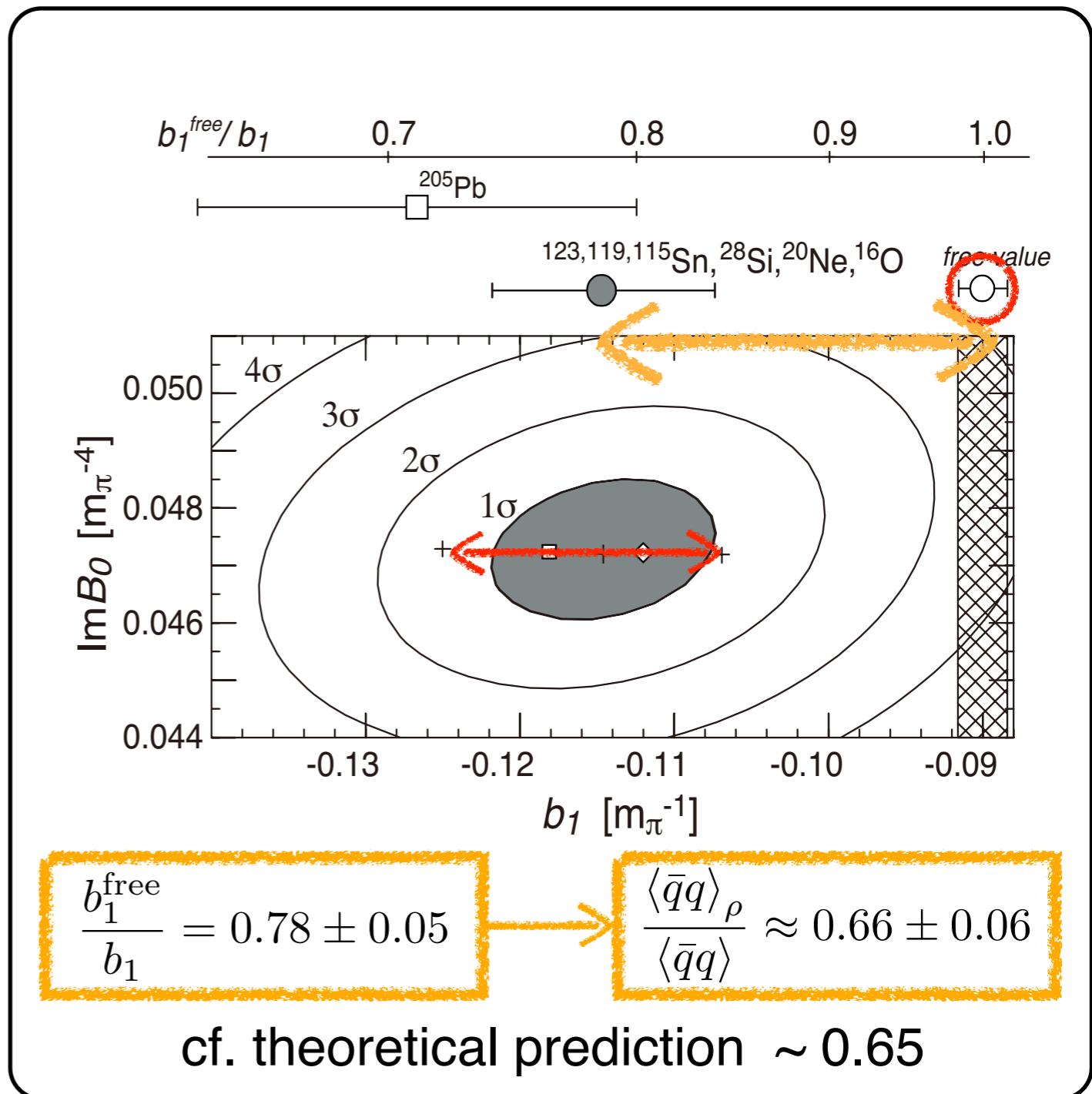
$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 b_0 \{\rho(r) + b_1 \delta\rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$

4 parameters (b_0 , b_1 , $\text{Re}B_0$, $\text{Im}B_0$) in equation
 ⇔ 2 parameters from each atom BE_{1s} , Γ_{1s}

- (i) Seki-Masutani relation (b_0 and $\text{Re}B_0$ has strong correlation)
- (ii) pionic atom of light / symmetric nuclei (^{16}O , ^{20}Ne , ^{28}Si)

$$b_0^* = b_0 + 0.025 \text{Re}B_0 = -0.0274 \pm 0.0002$$

Deduction of b_1 from experimental data

BE_{1s} vs Γ_{1s} Contour plot of χ^2 

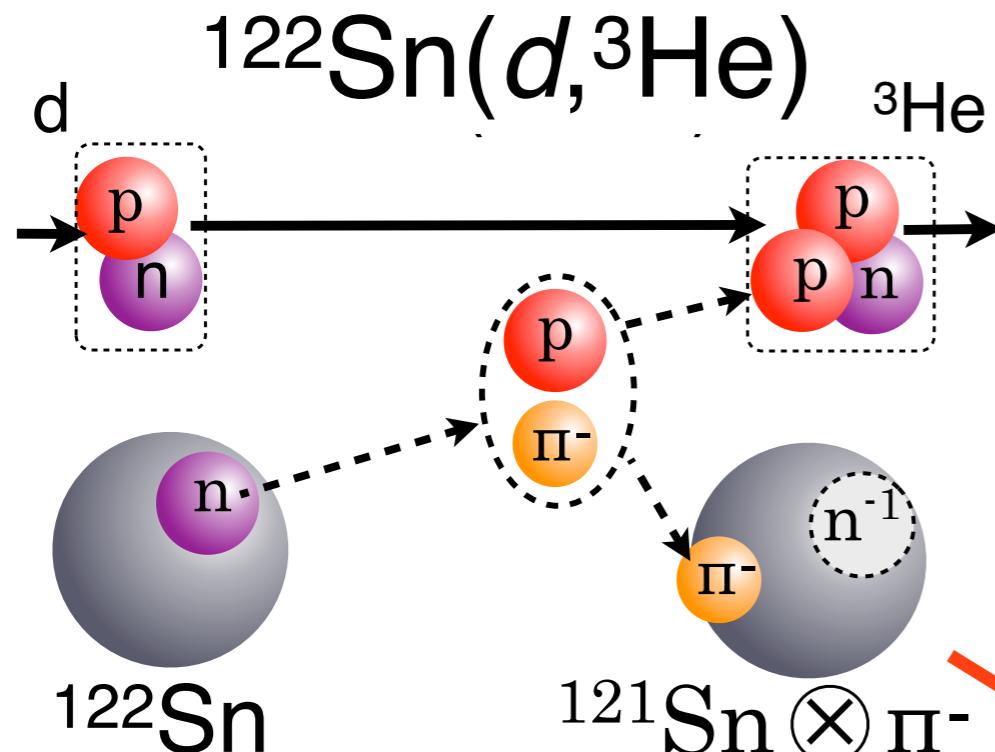
error of b_1 in medium is still large compared with that in vacuum!!

two main sources are

- **experimental error**
- **neutron distribution ambiguities**

2. Precise / systematic
experiment at RIBF, RIKEN

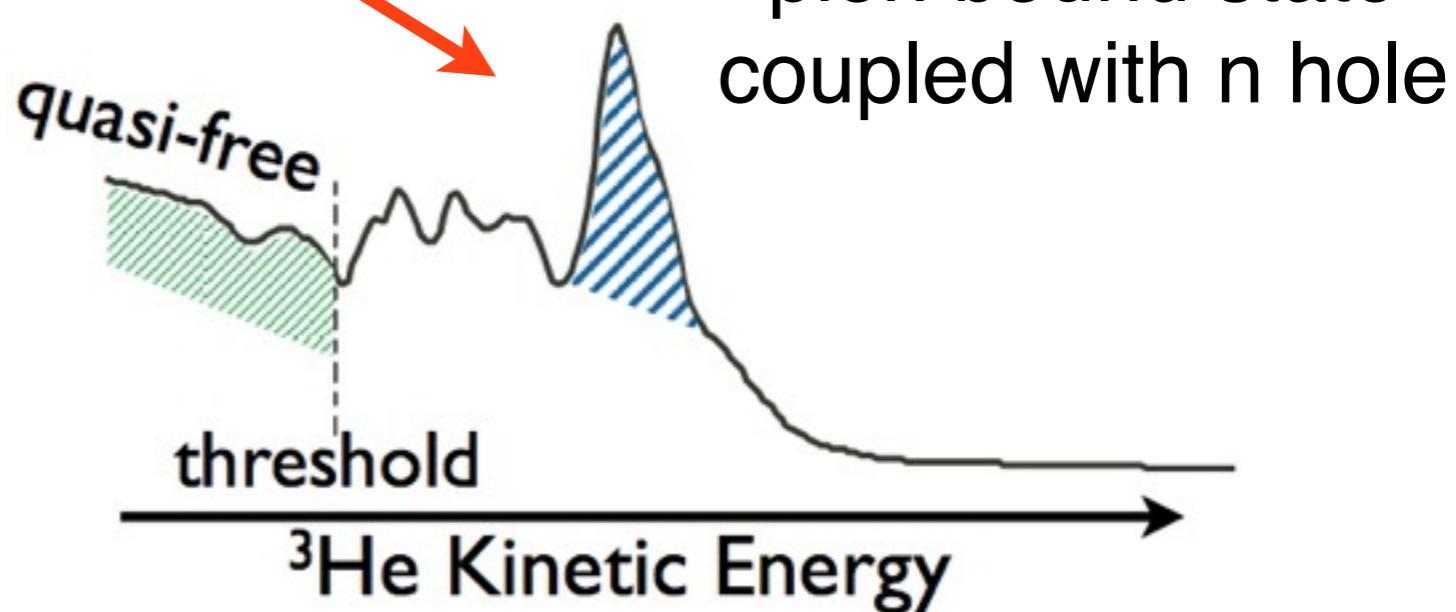
Production method; ($d, {}^3\text{He}$) reaction



2 body kinematics
 → mass of the pionic atom
 can be calculated from **Q-value**

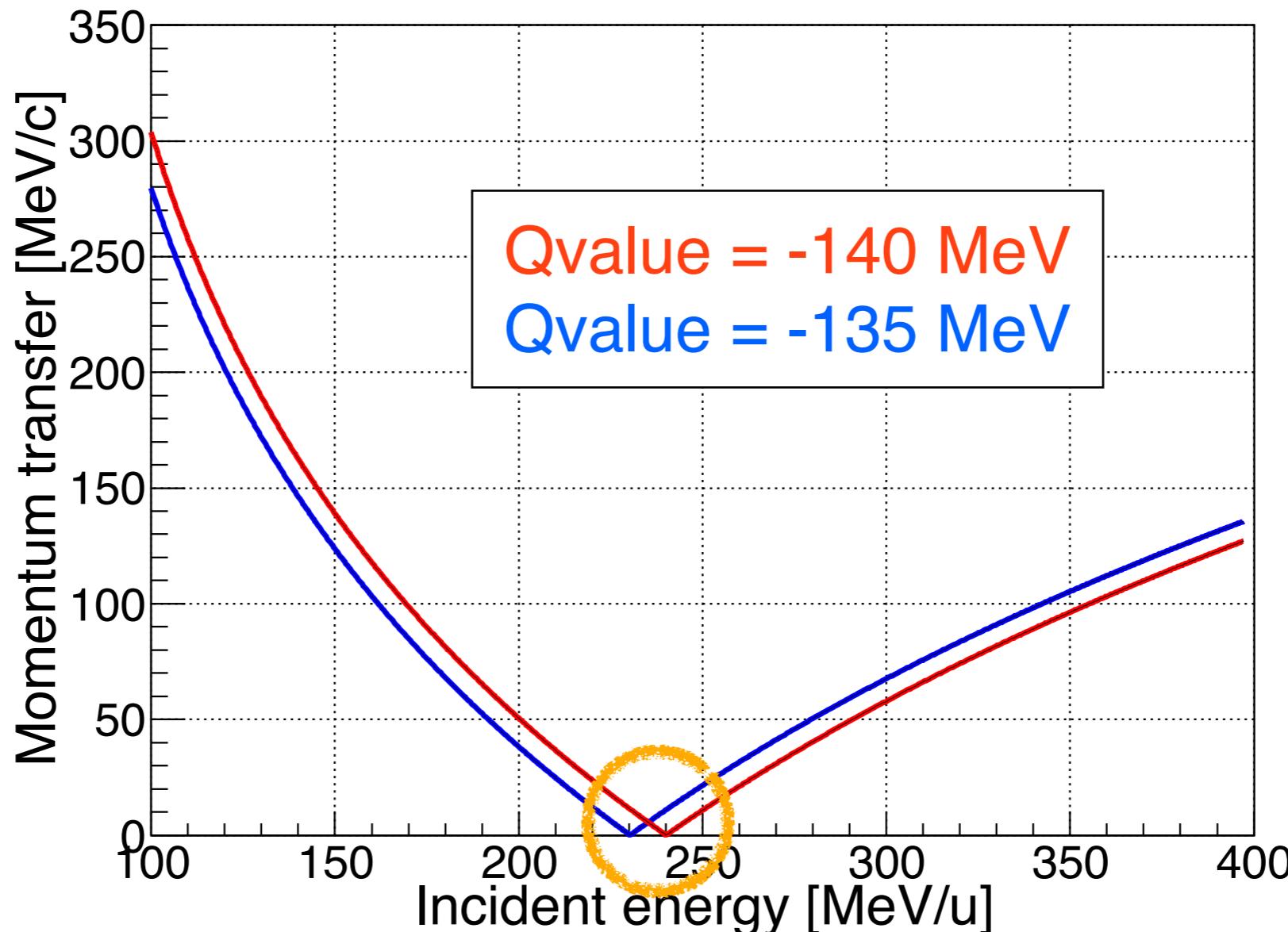
Calculated from E_d , $E_{{}^3\text{He}}$

Missing mass spectroscopy



Production method; ($d, {}^3He$) reaction

recoilless condition at $T_d \doteq 250$ MeV/u

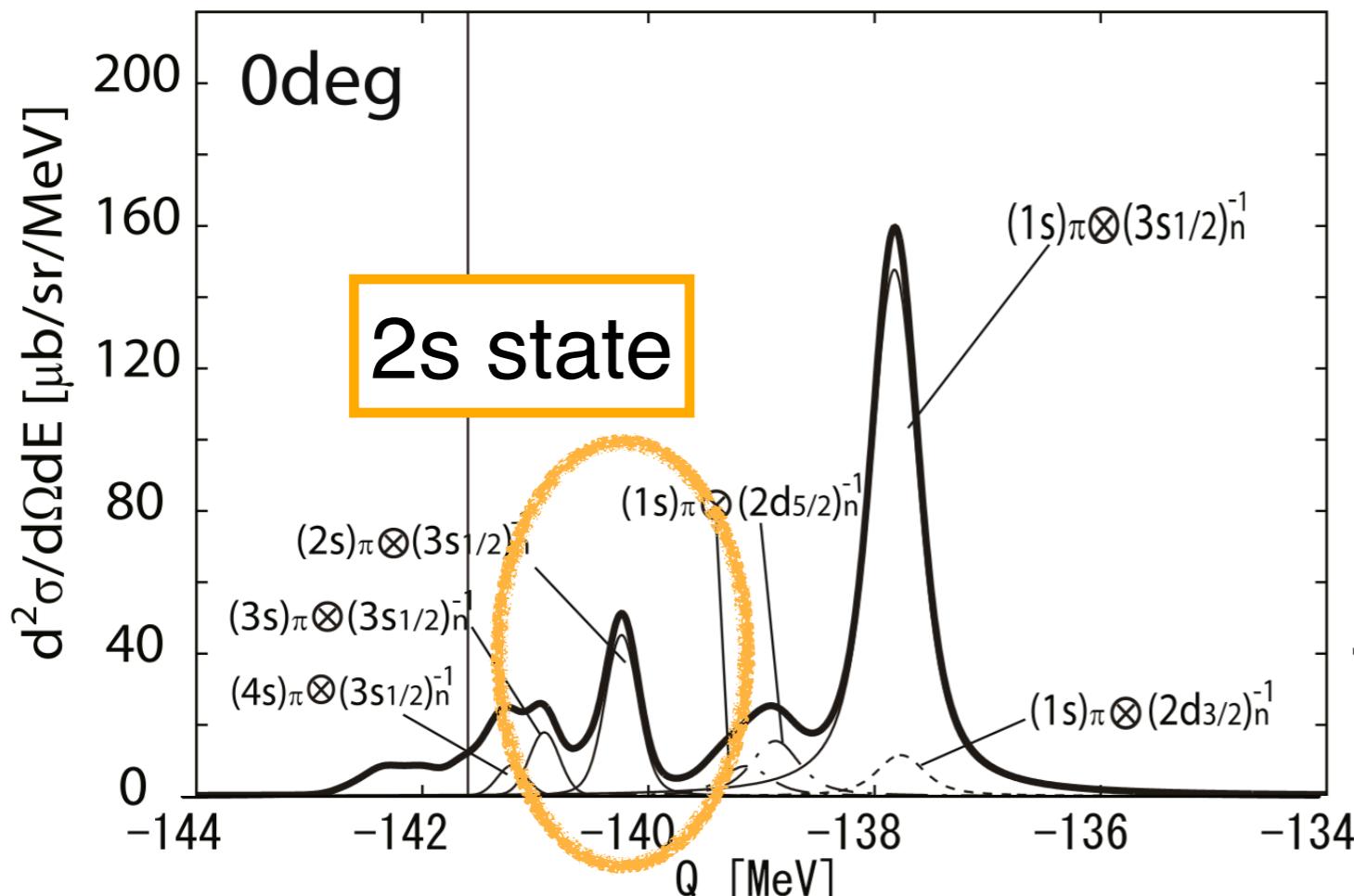


- large cross section
- Enhance $(1s)_{\pi^-} (3s)_{n-1}$ configuration

Experiment for precise / systematic study

Theoretical calculated Spectrum

$\Delta E = 300 \text{ keV}$



More statistics / resolution enable us to observe 2s state.



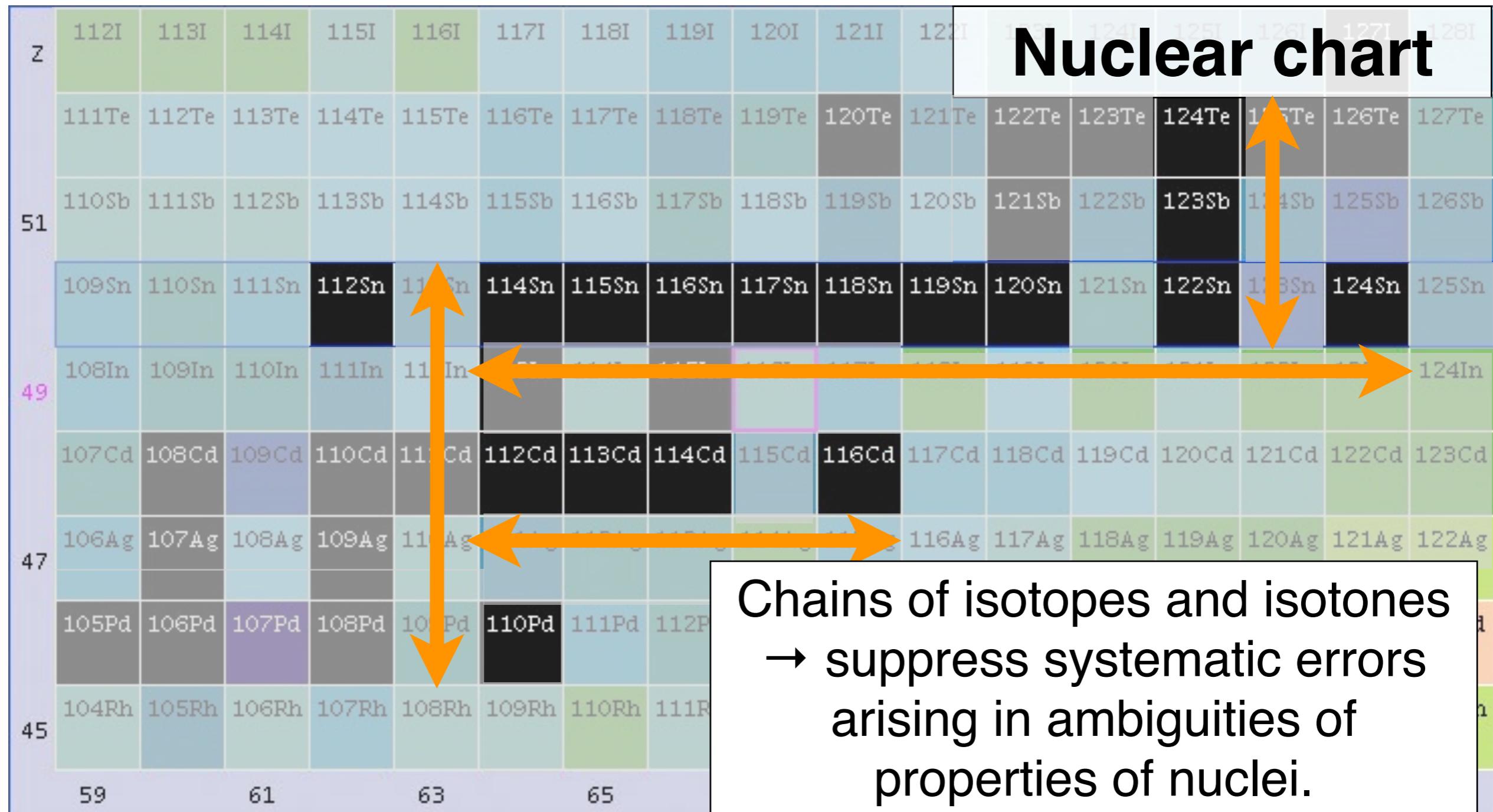
$\text{BE}_{1s} - \text{BE}_{2s} / \Gamma_{1s} - \Gamma_{2s}$
These value are less affected from the systematic errors.

*N. Ikeno et al., Eur. Phys. J. A 47, 161 (2011)

cf. $\Delta E \sim 400 \text{ keV}$ in GSI experiment

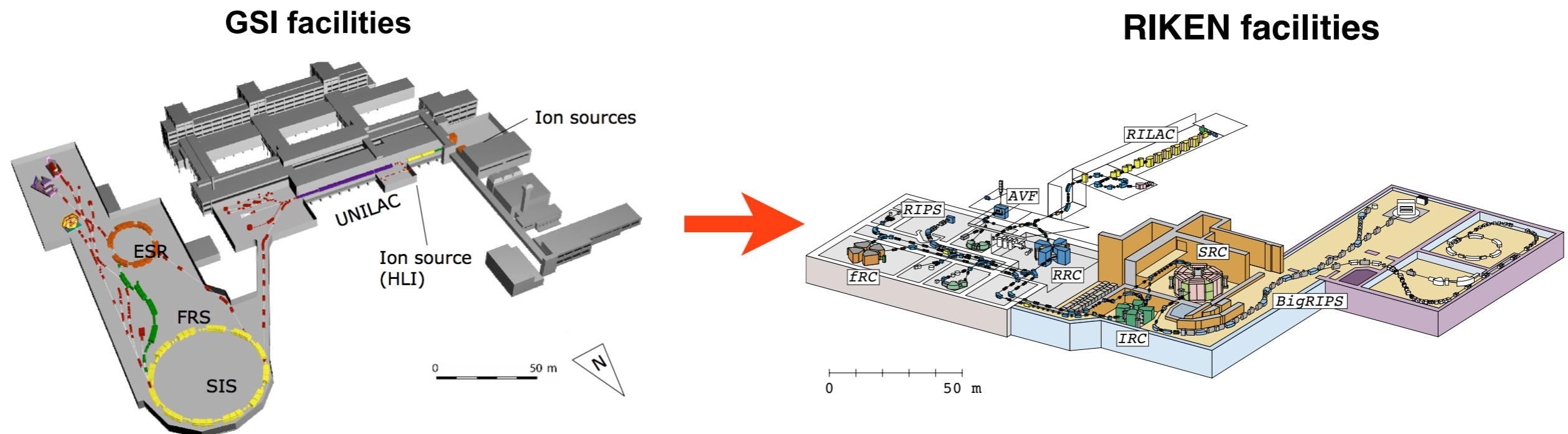
Experiment for precise / systematic study

systematic study of pionic nuclei in isotope / isotone chain



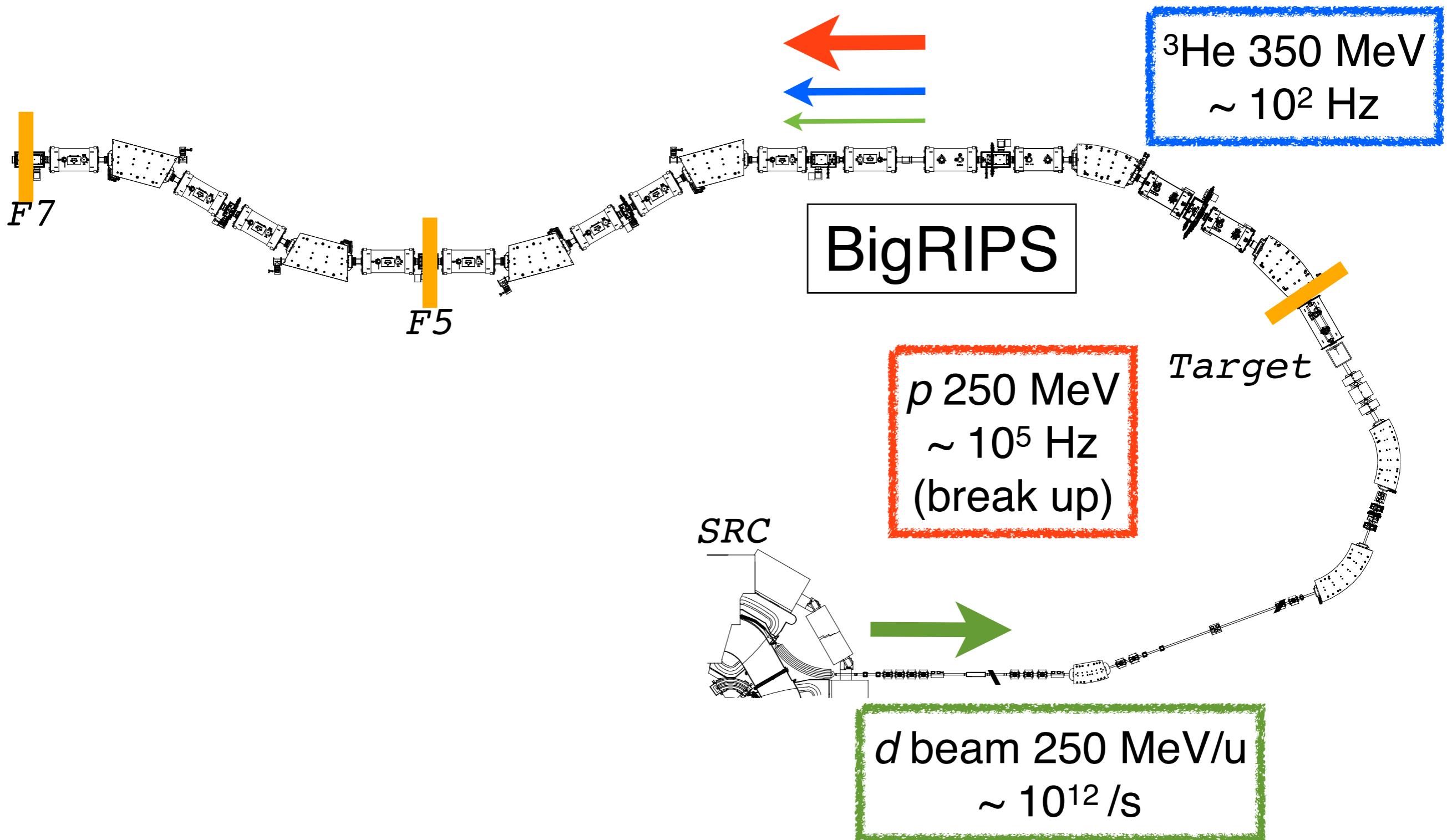
Chains of isotopes and isotones
→ suppress systematic errors
arising in ambiguities of
properties of nuclei.

Experiment at RIKEN

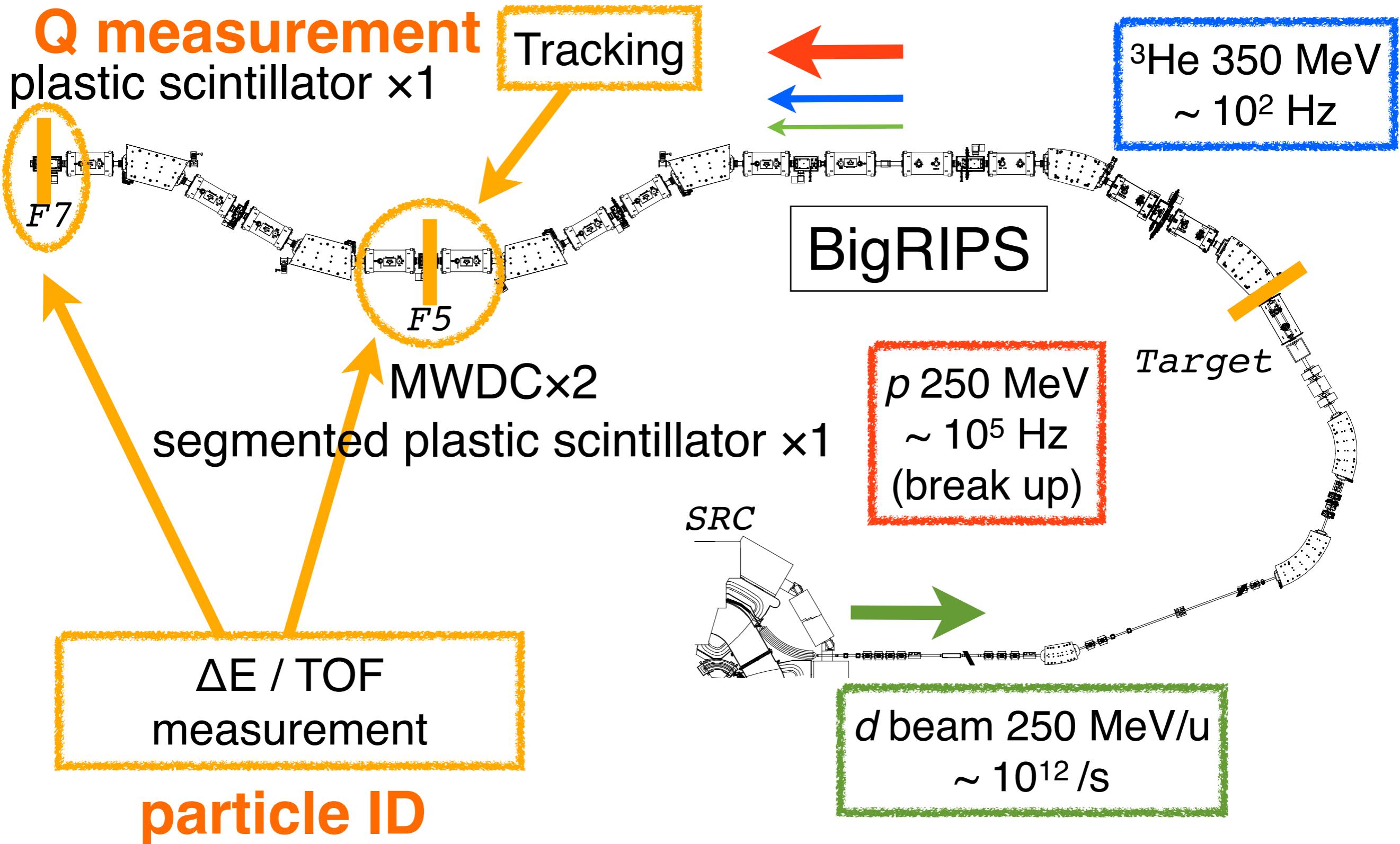


	GSI	RIBF	
intensity	$\sim 10^{11}/\text{spill}$	$\sim 10^{12}/\text{s}$	$\times 50$
angular acceptance	$\sim 10 \text{ mrad}$	$40 / 60 \text{ mrad}$	$\times 20$

Experimental setup



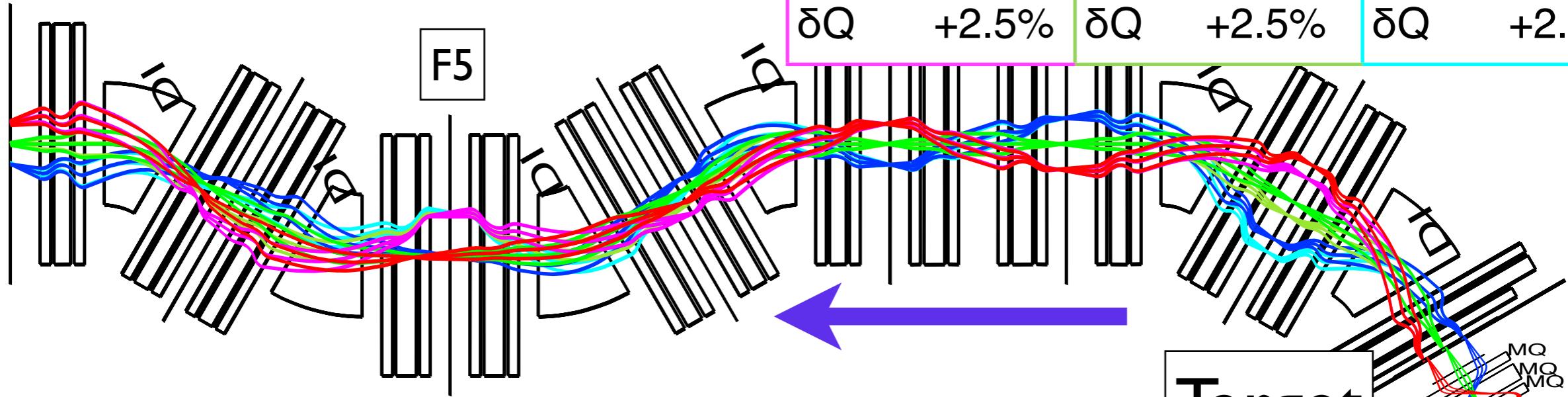
Experimental setup



Ion optics (dispersion matching)

BigRIPS

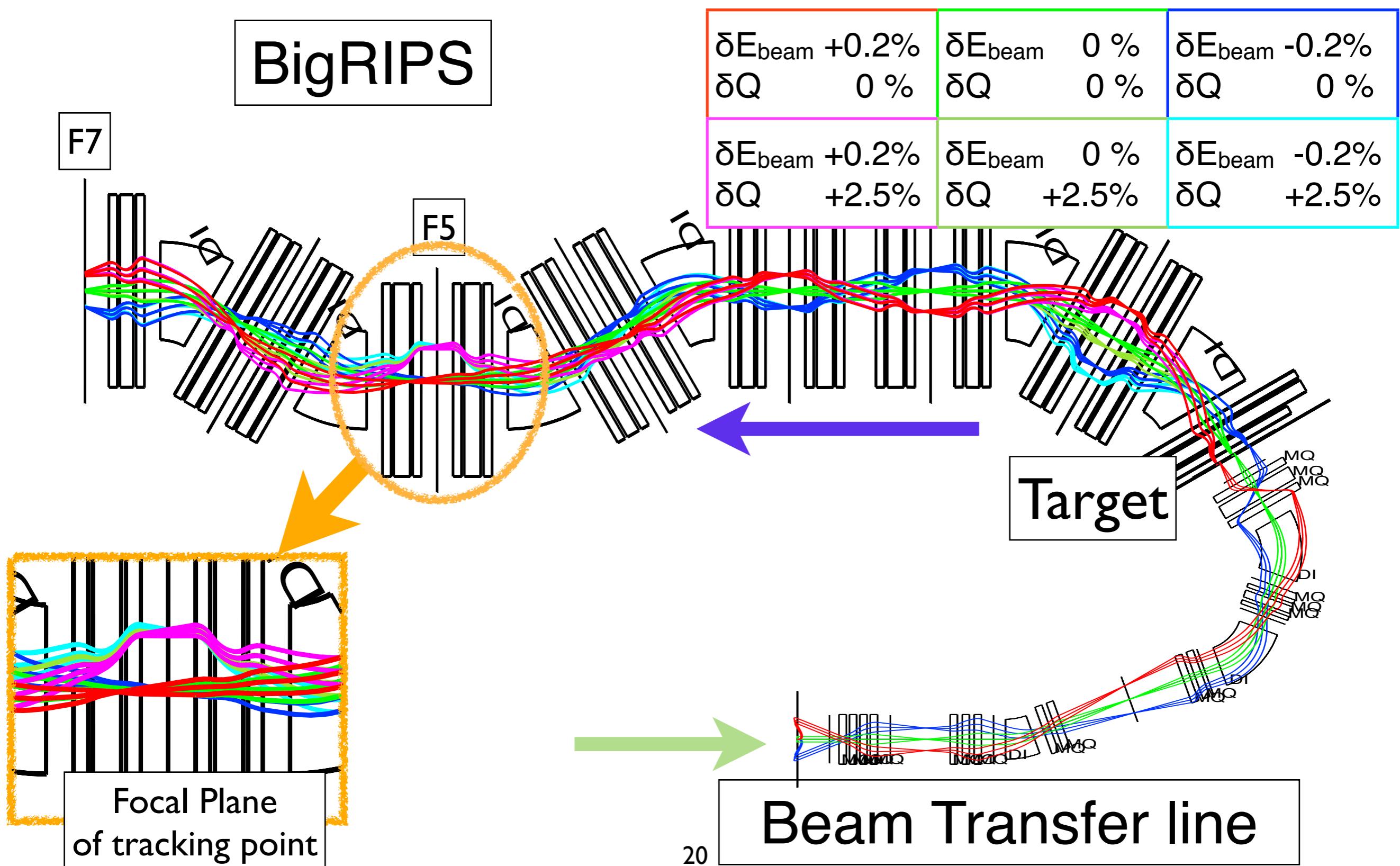
F7



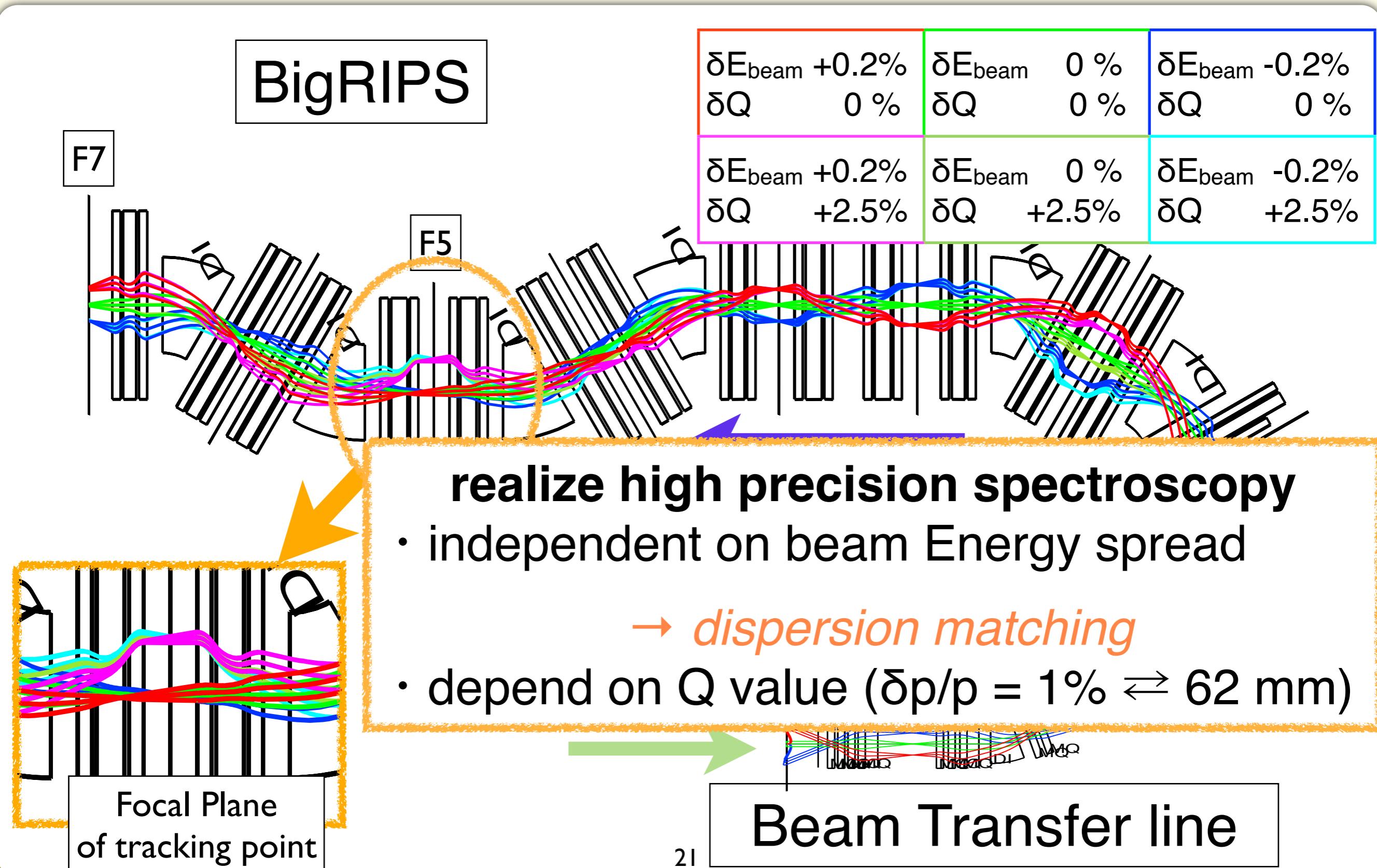
$\delta E_{beam} +0.2\%$	$\delta E_{beam} 0 \%$	$\delta E_{beam} -0.2\%$
$\delta Q 0 \%$	$\delta Q 0 \%$	$\delta Q 0 \%$
$\delta E_{beam} +0.2\%$	$\delta E_{beam} 0 \%$	$\delta E_{beam} -0.2\%$
$\delta Q +2.5\%$	$\delta Q +2.5\%$	$\delta Q +2.5\%$

Target

Ion optics (dispersion matching)

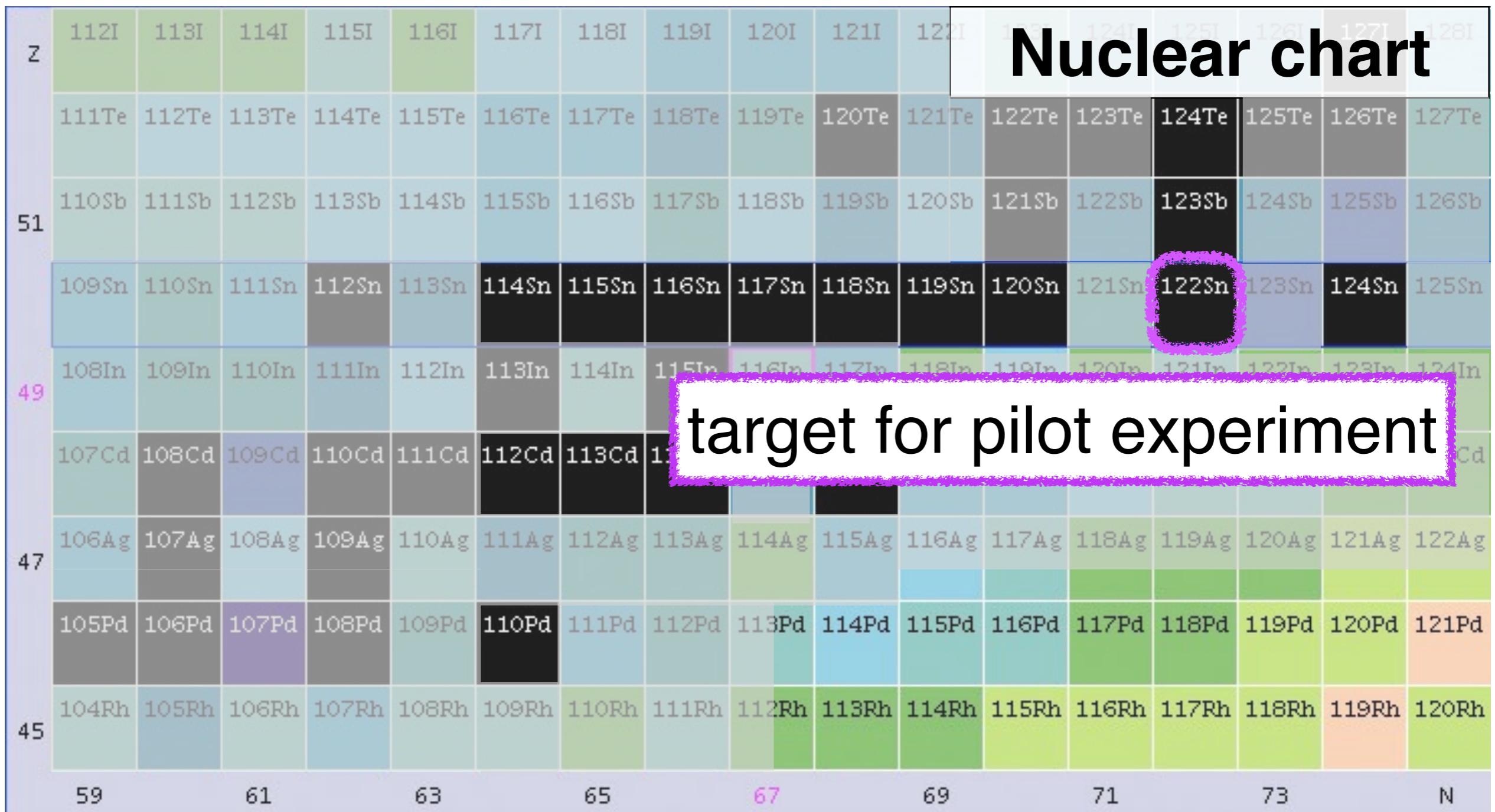


Ion optics (dispersion matching)



Pilot experiment

pilot experiment; performed in 2010



Pilot experiment

pilot experiment; performed in 2010



Target ; ^{122}Sn

Purpose ; test for

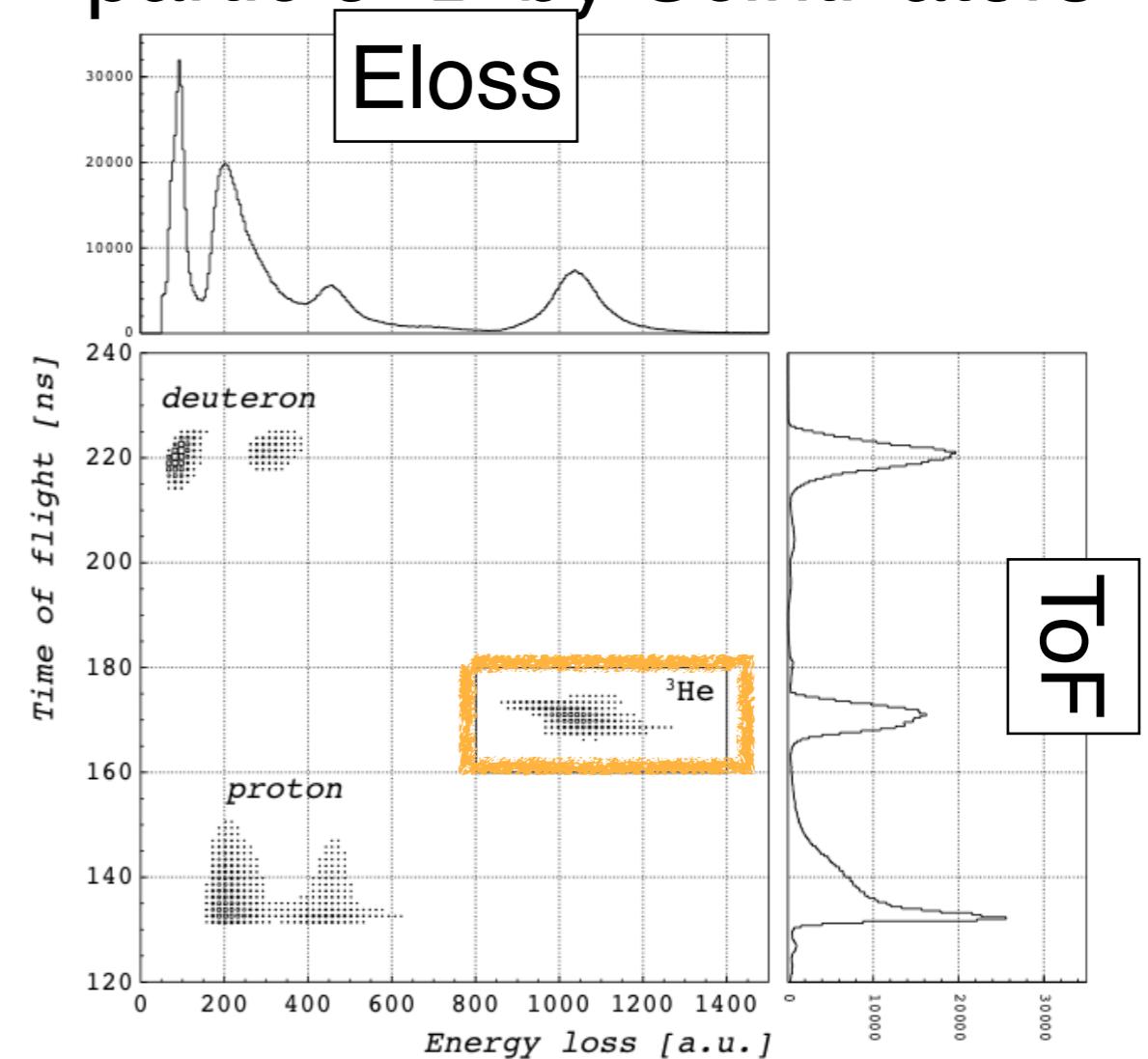
- **all detector system**
- **ion optics**

Pilot experiment

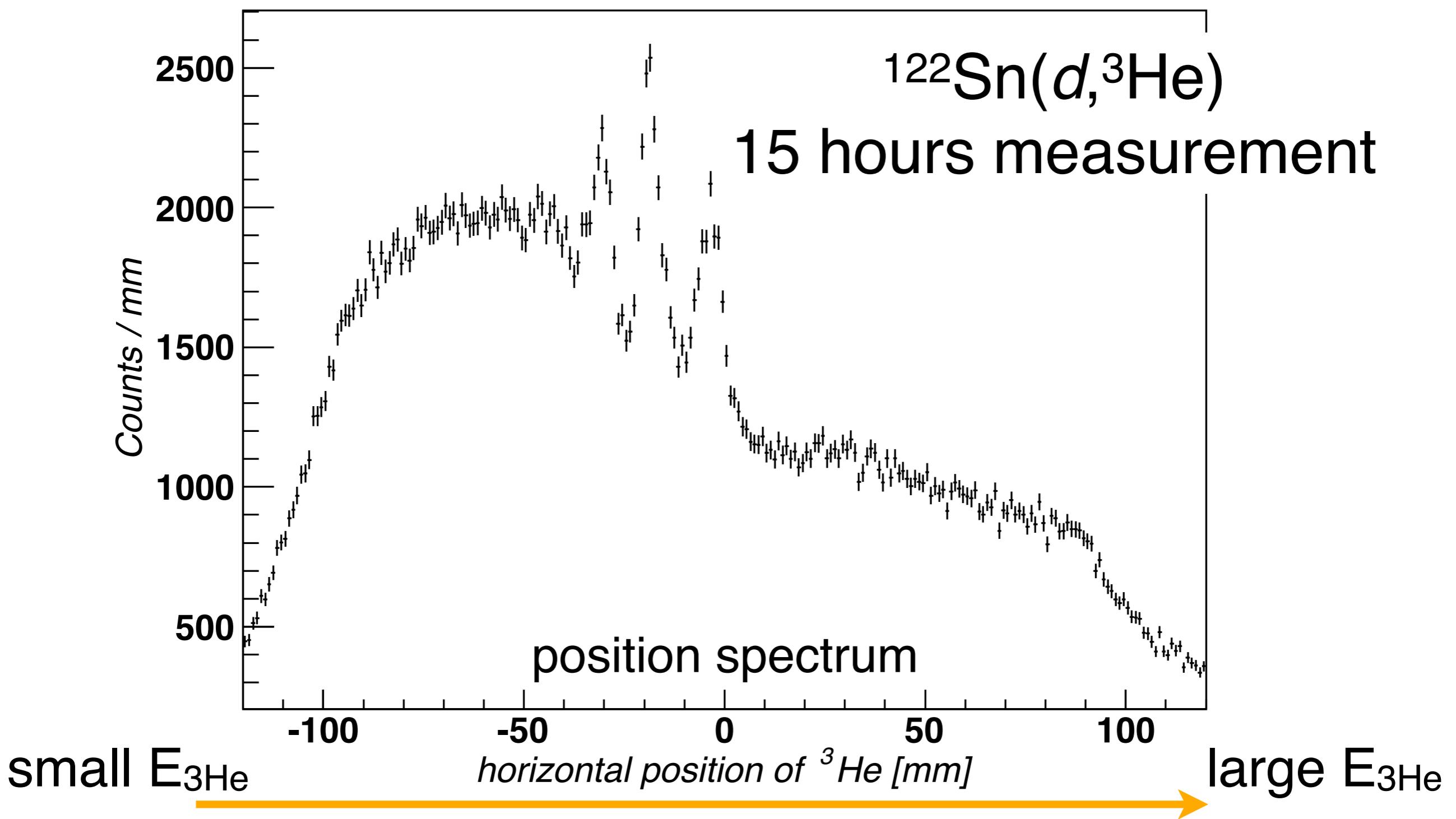
pilot experiment; performed in 2010



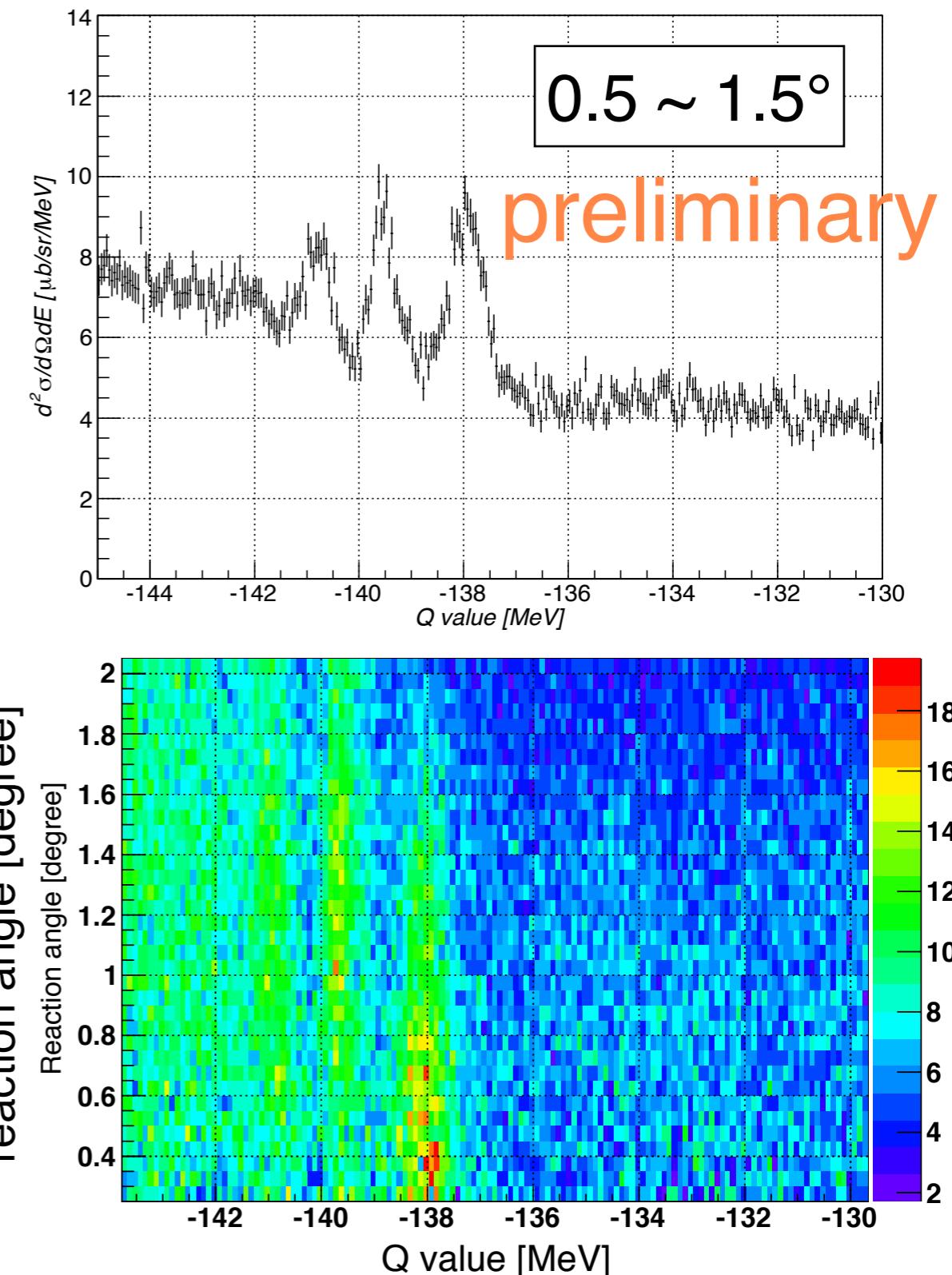
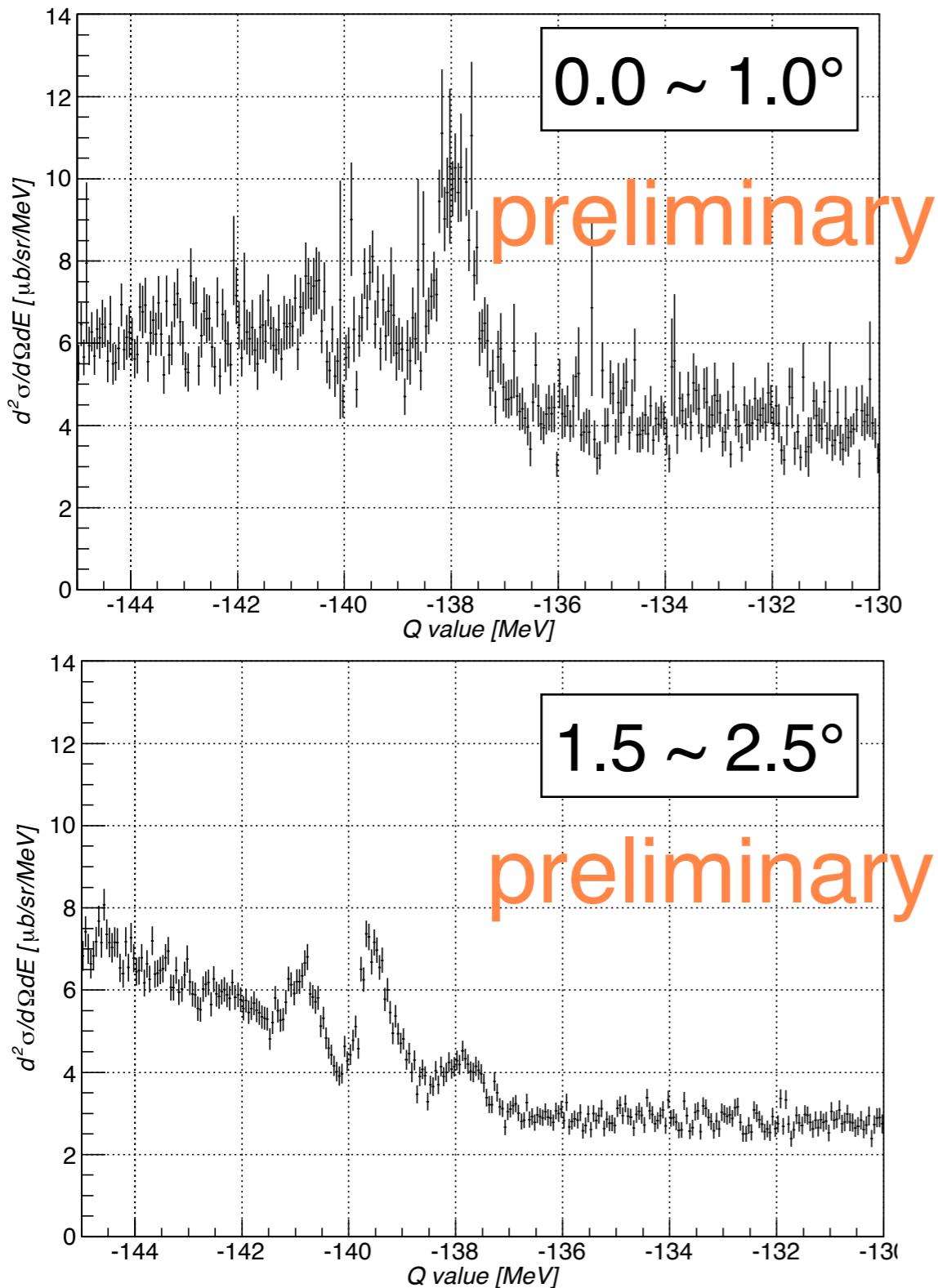
particle ID by Scintillators



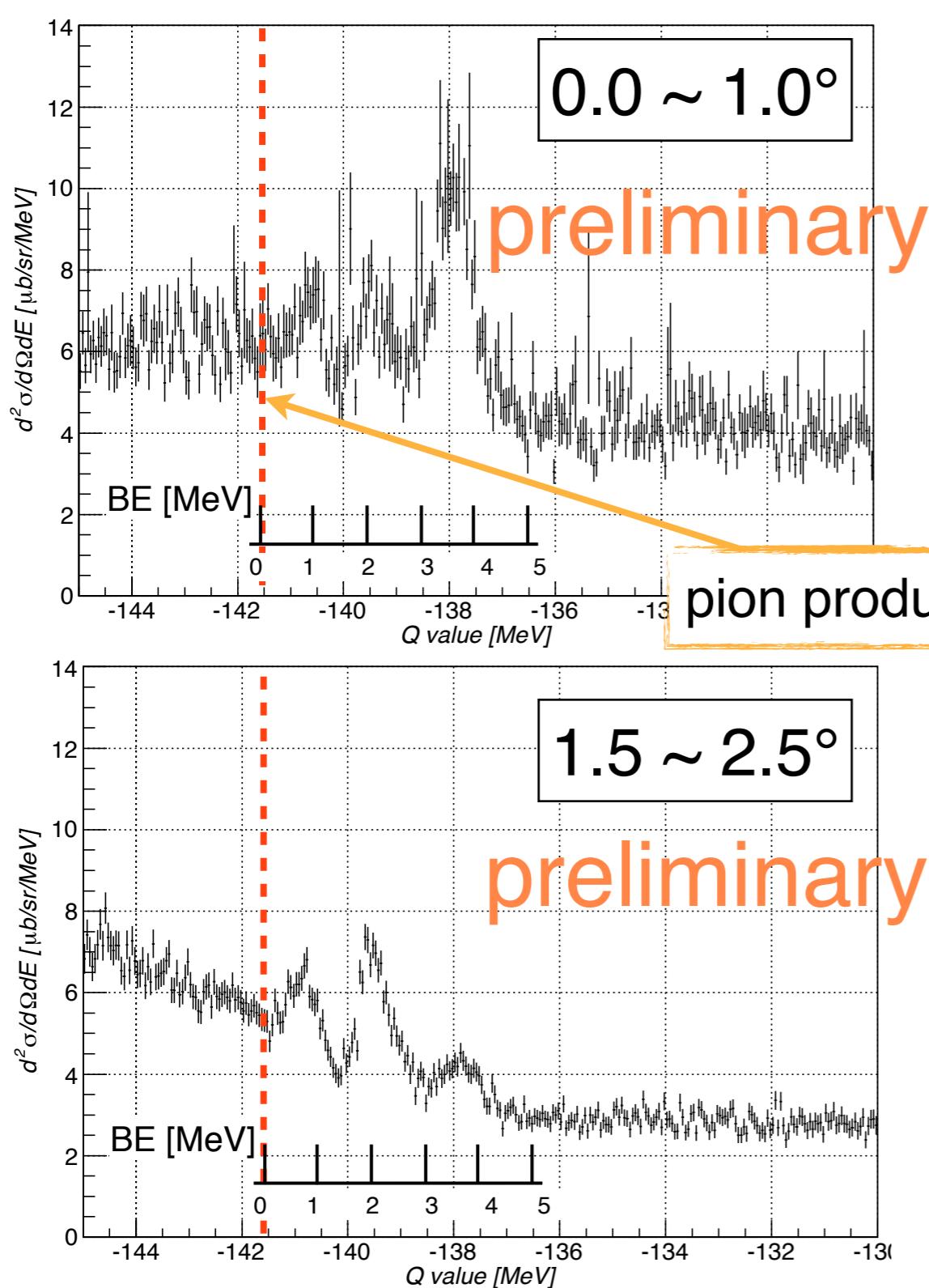
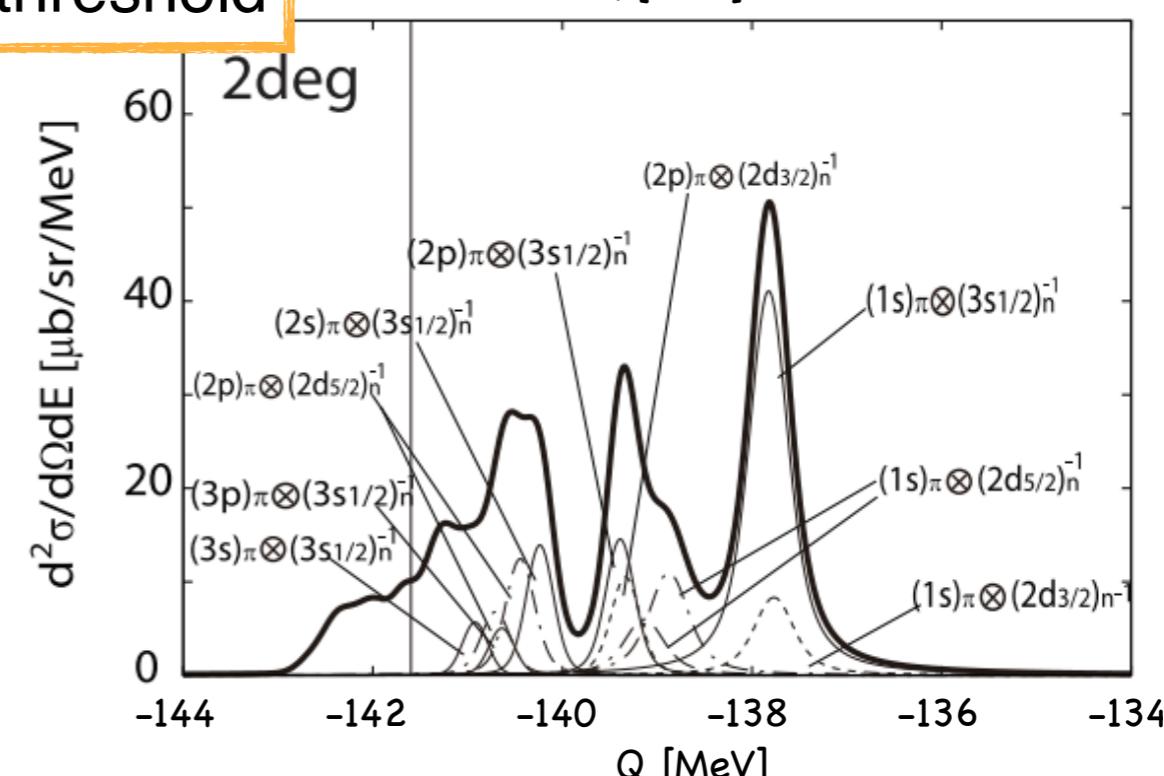
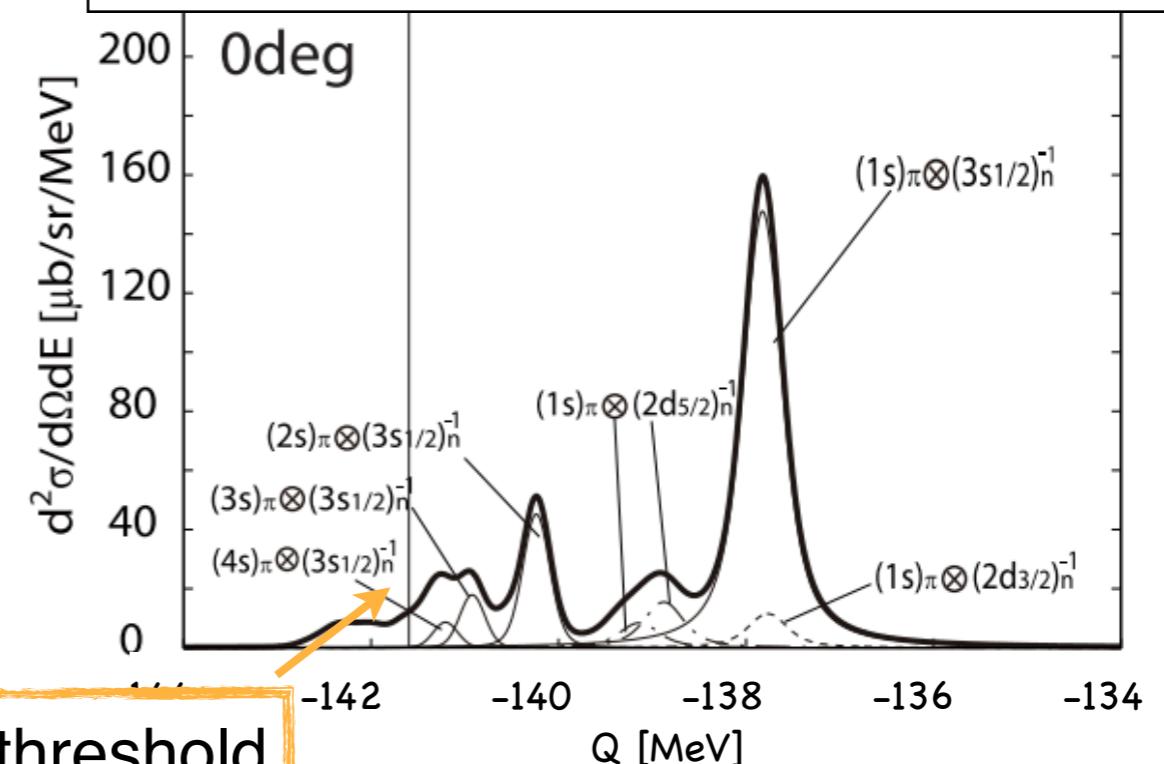
Pilot experiment



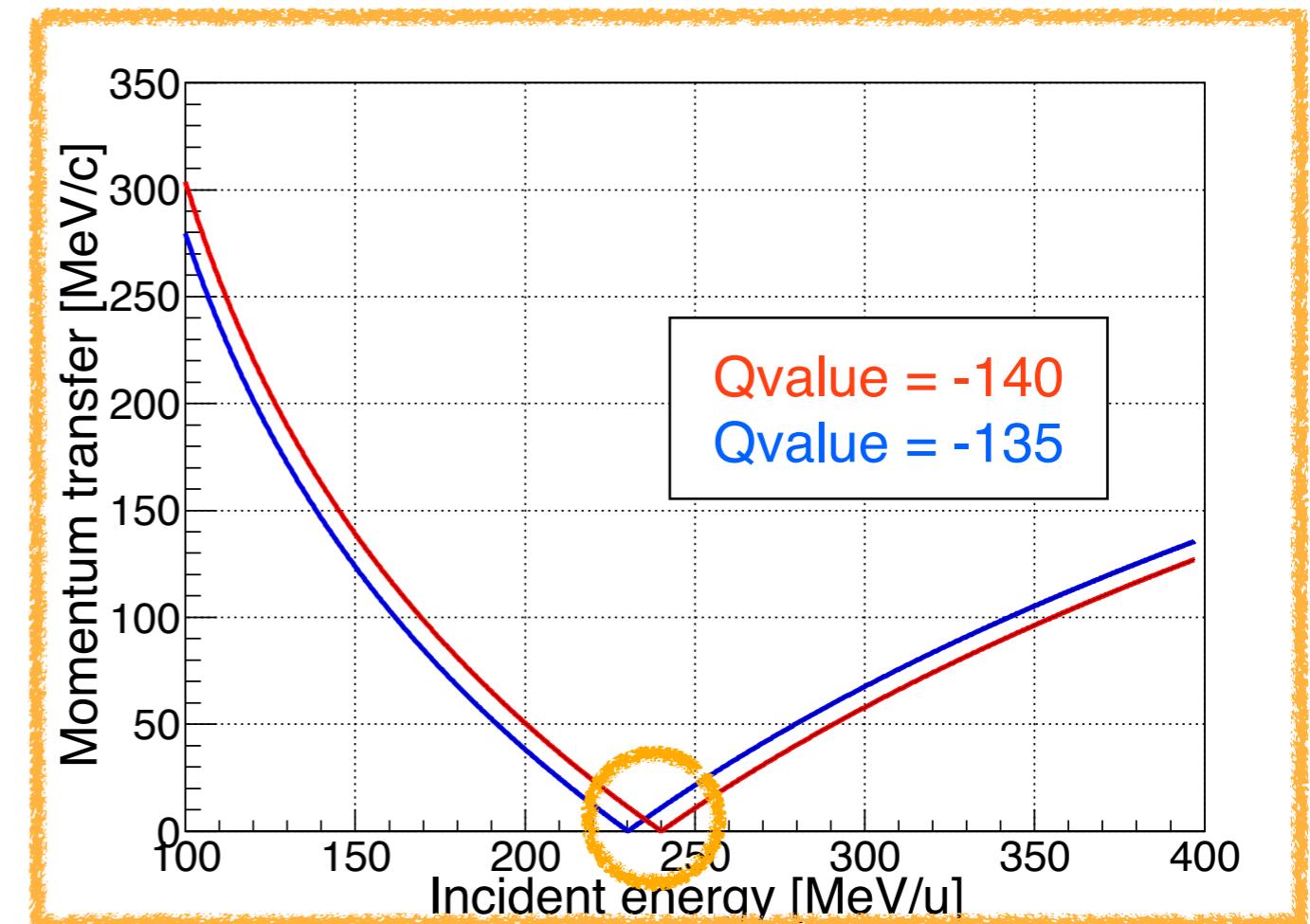
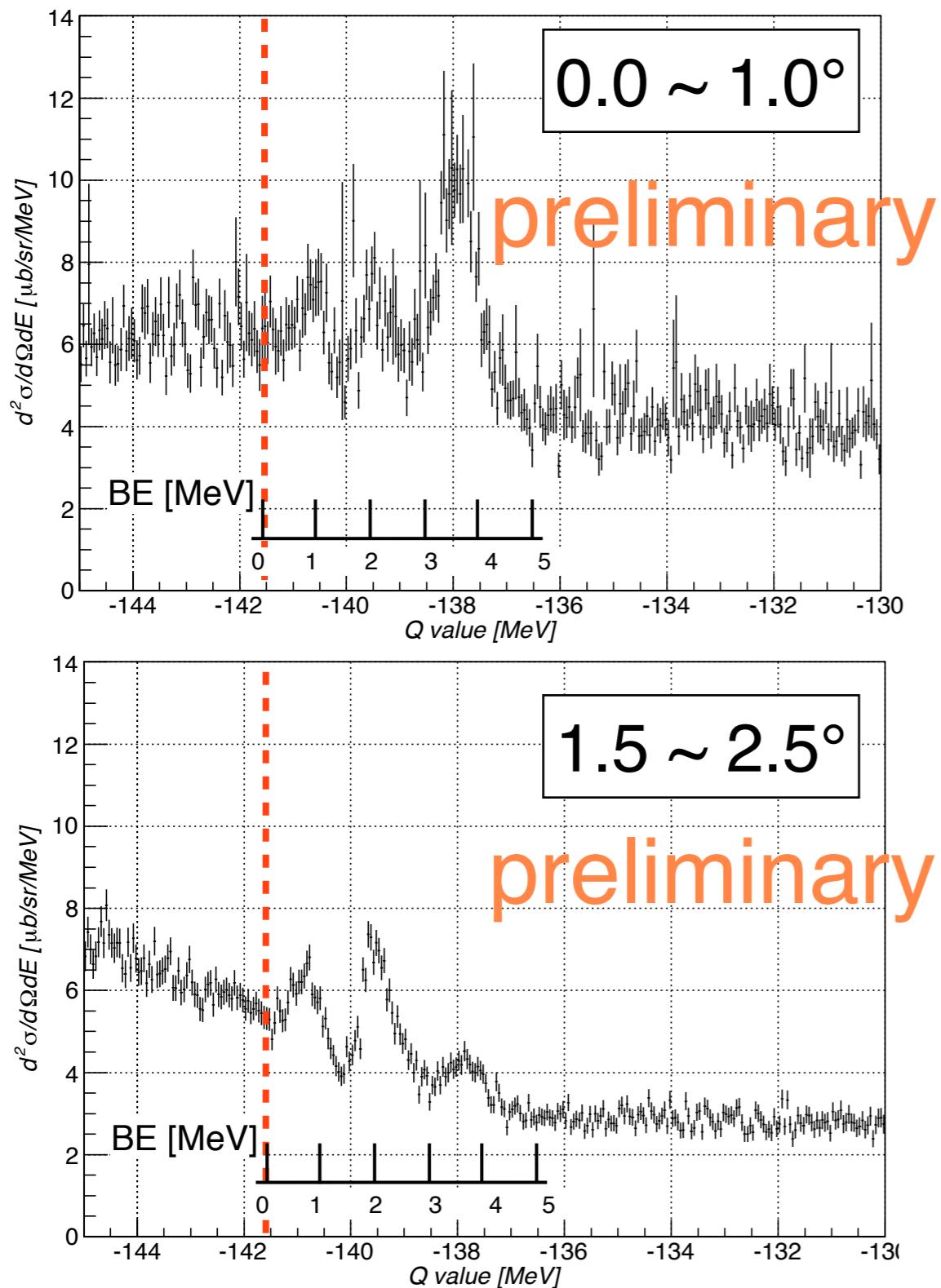
Experimental Q-value spectra



Experimental Q-value spectra

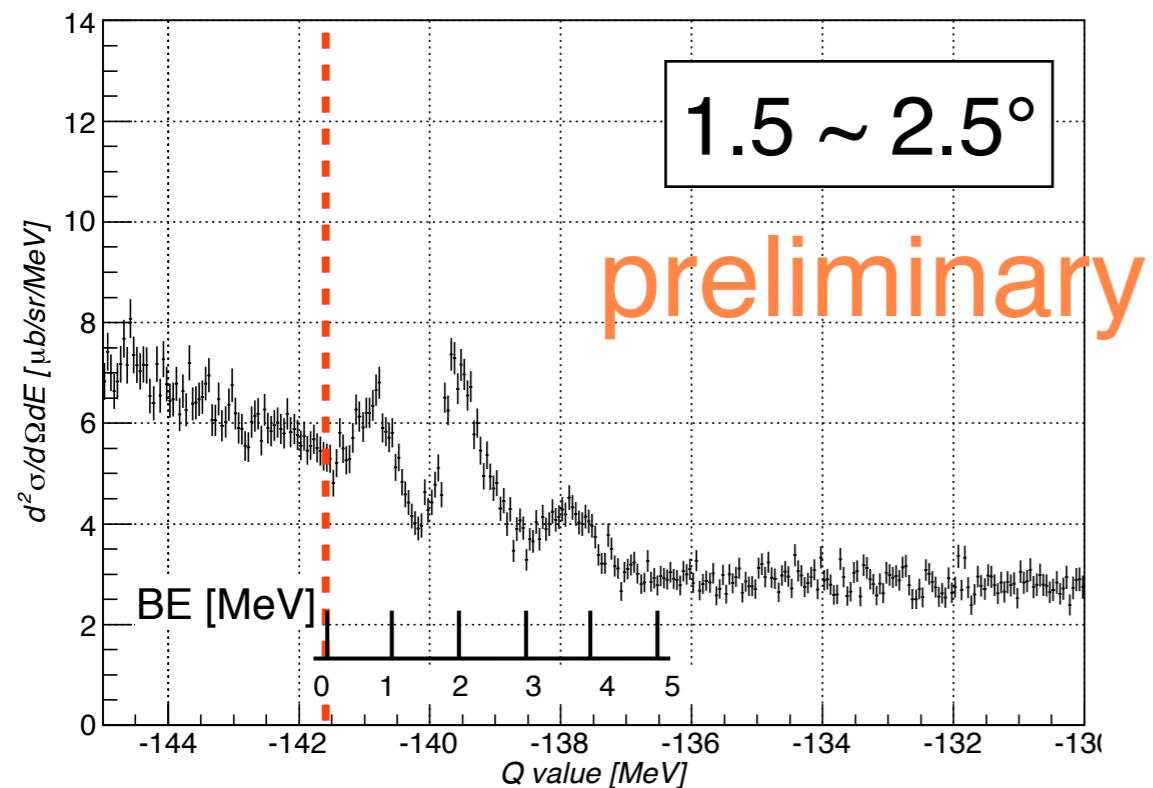
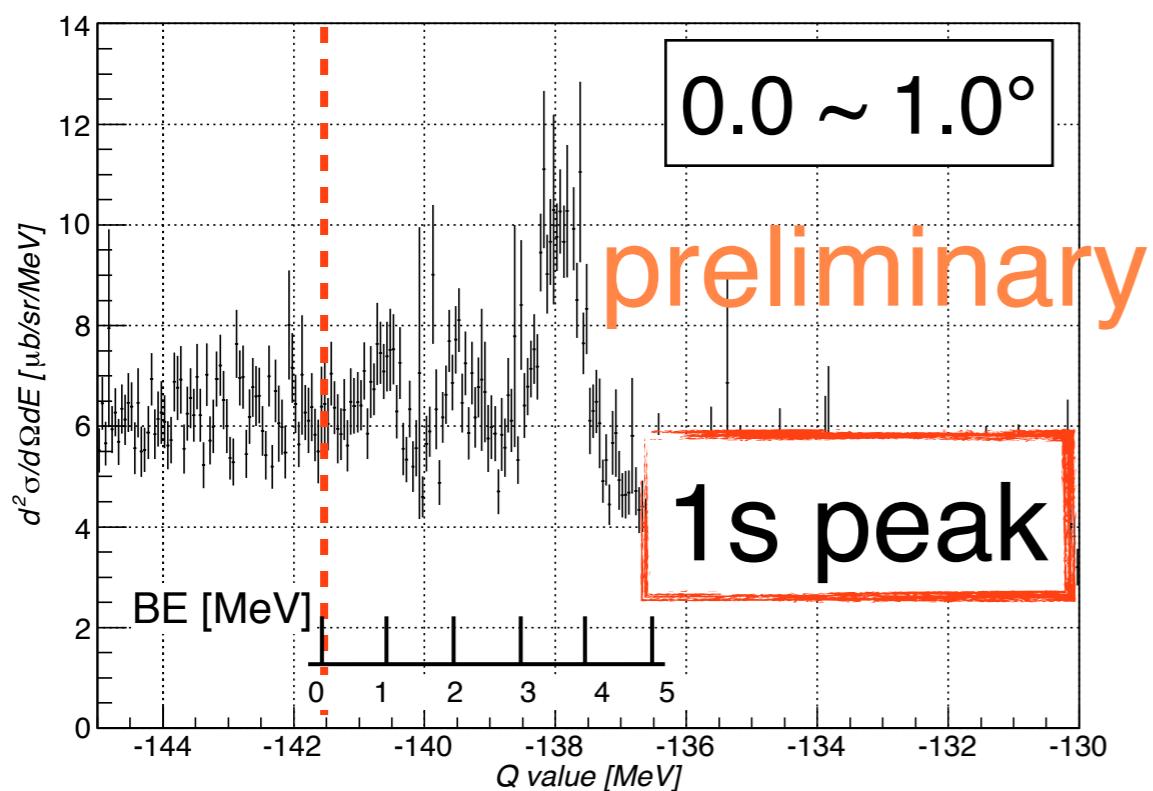
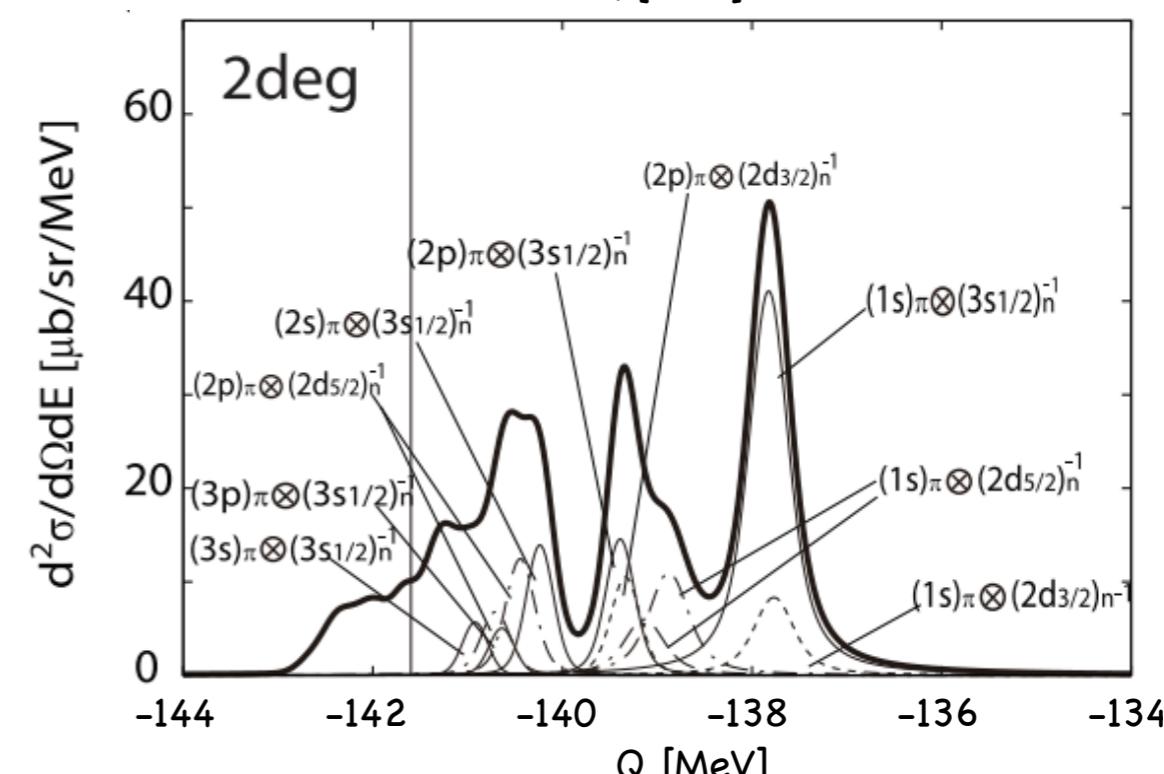
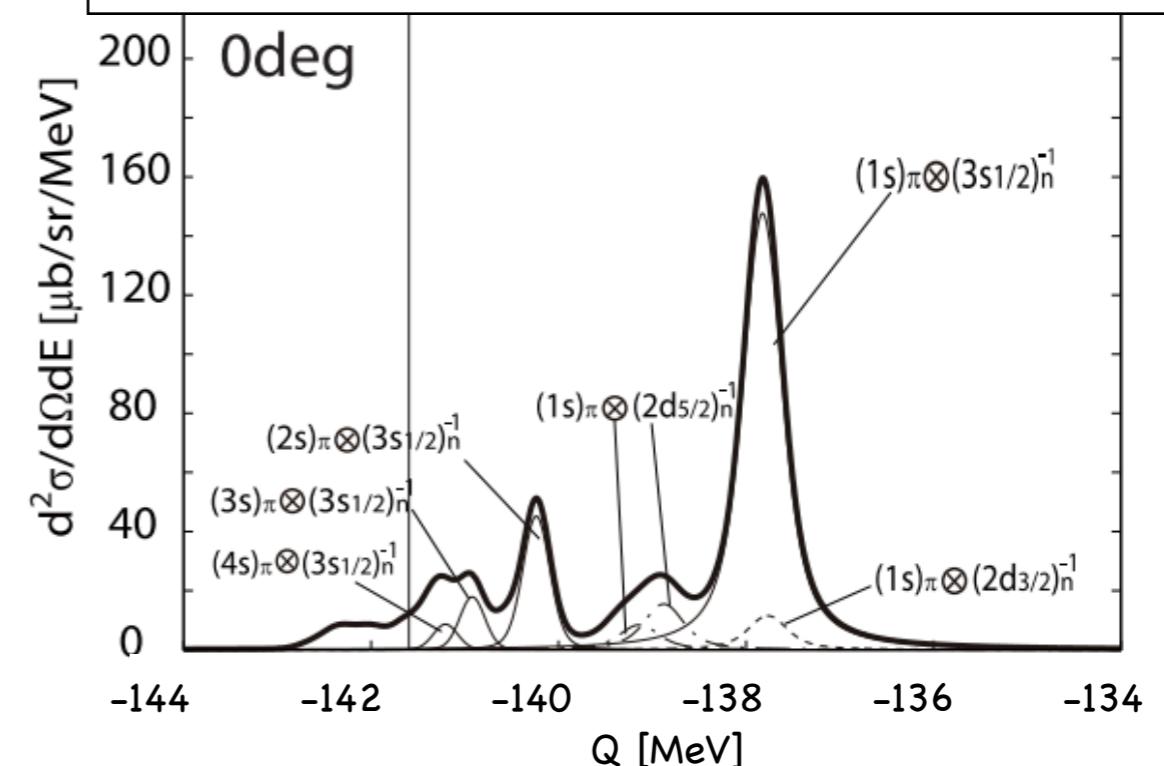
Theoretical spectrum ($\Delta E = 300$ keV)

Experimental Q-value spectra

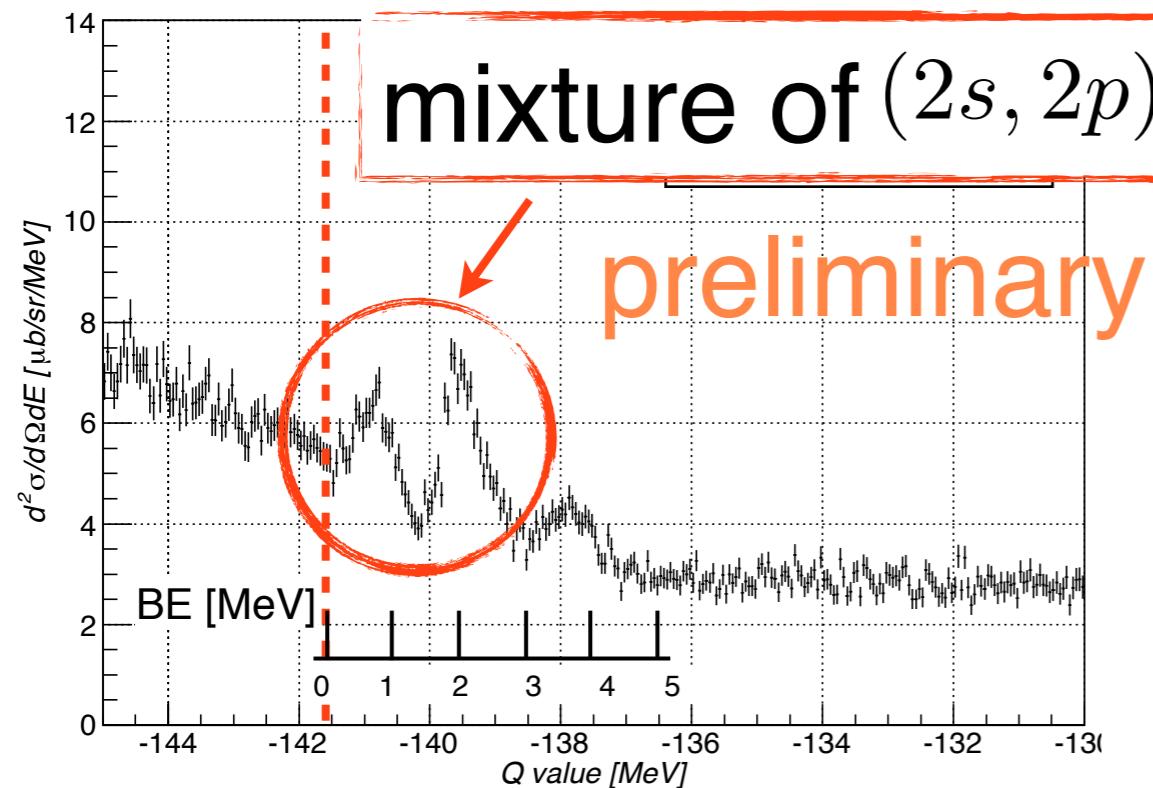
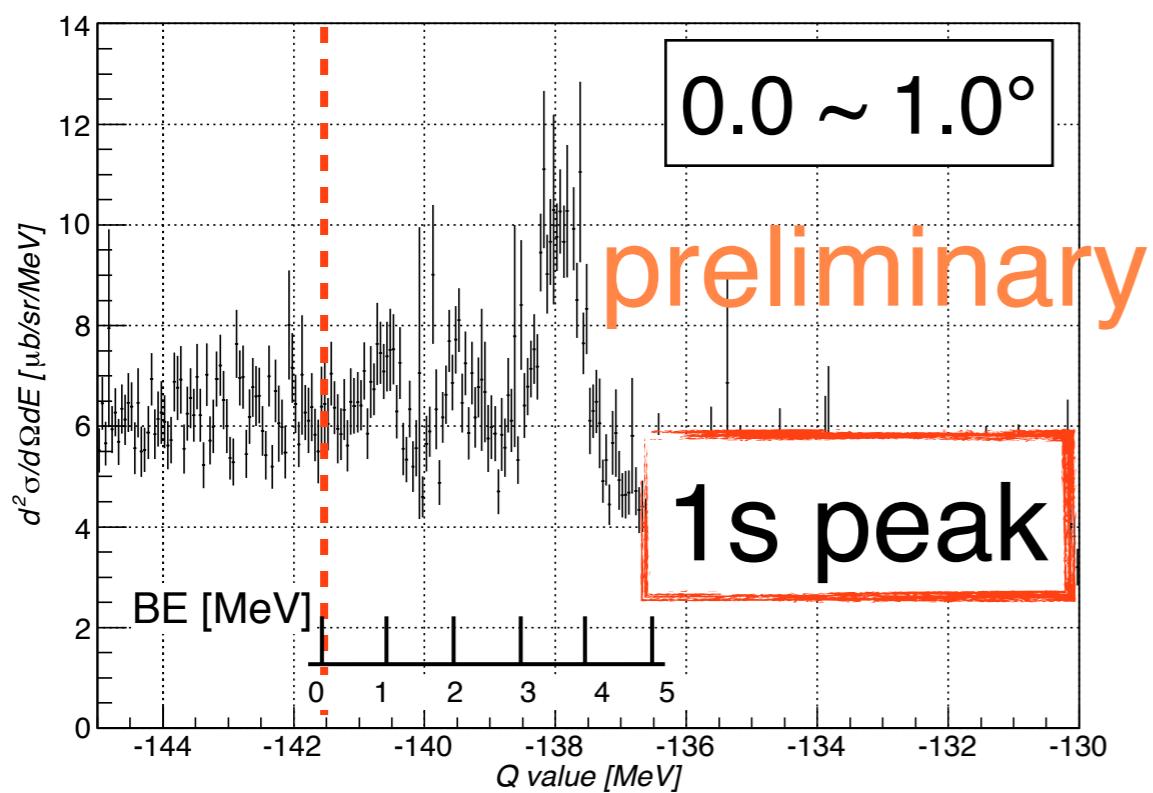


finite angle
 → recoilless condition is broken
 → configuration with finite Δl grow up

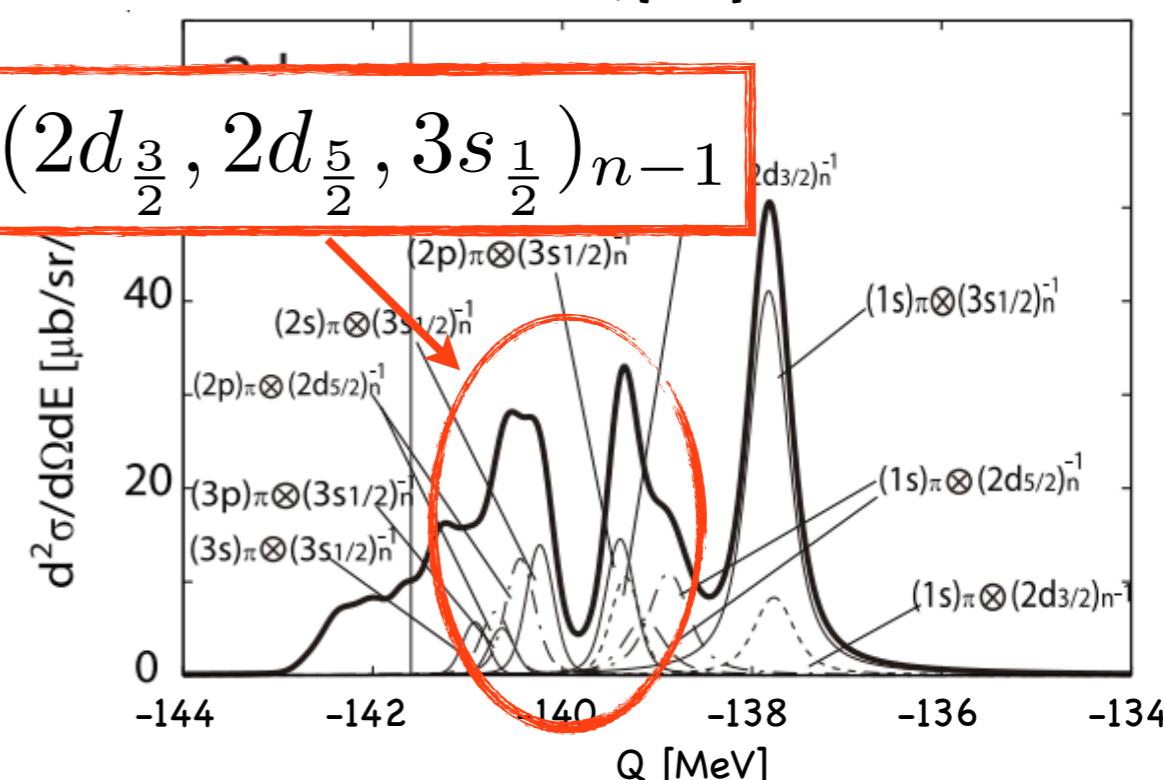
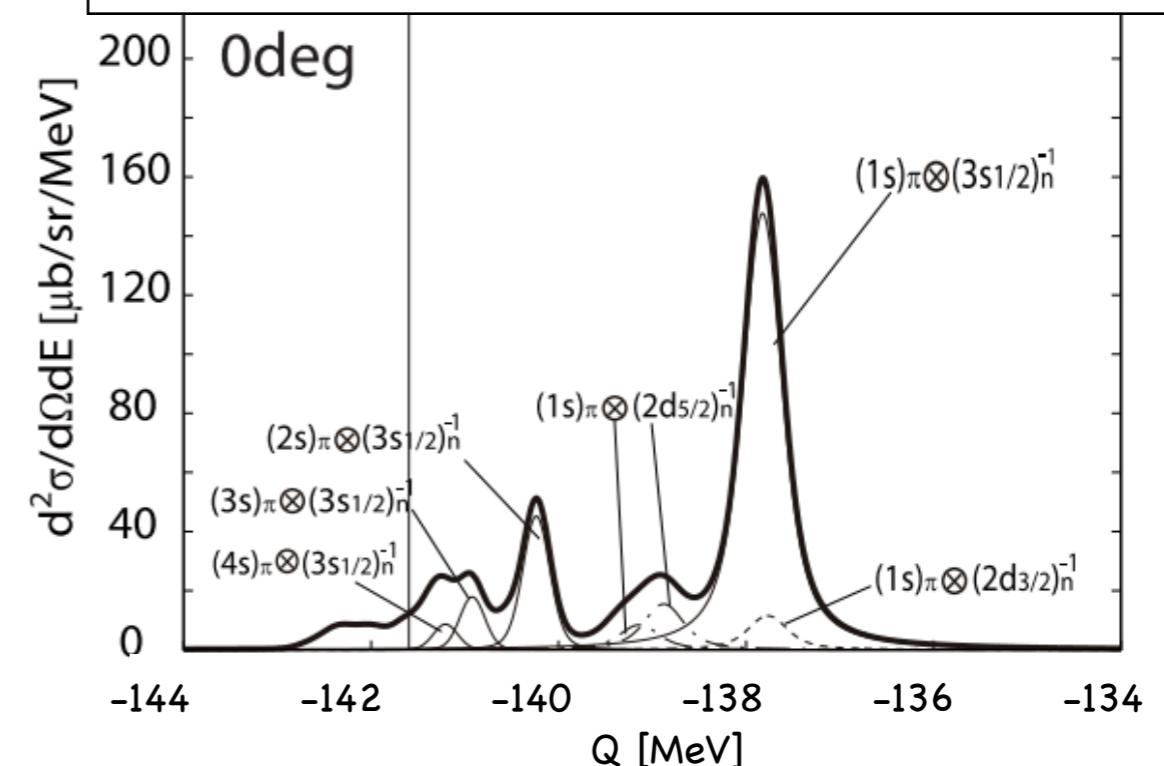
Experimental Q-value spectra

Theoretical spectrum ($\Delta E = 300$ keV)

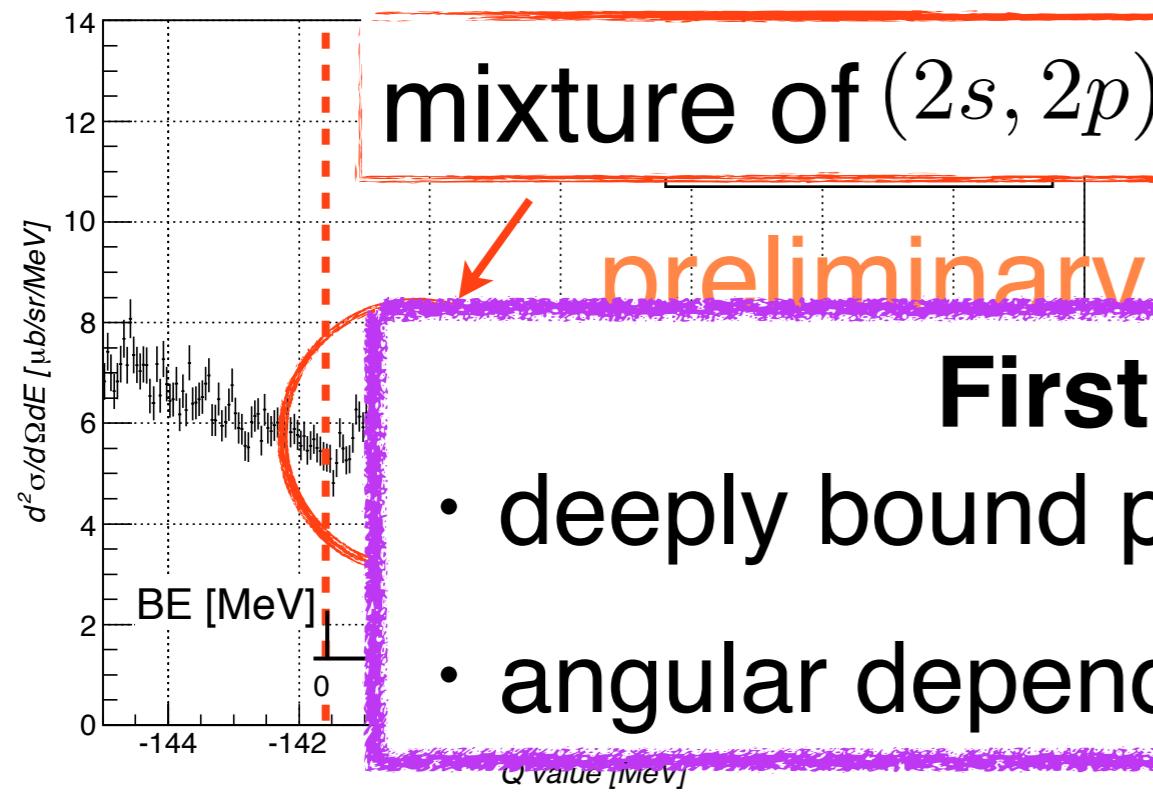
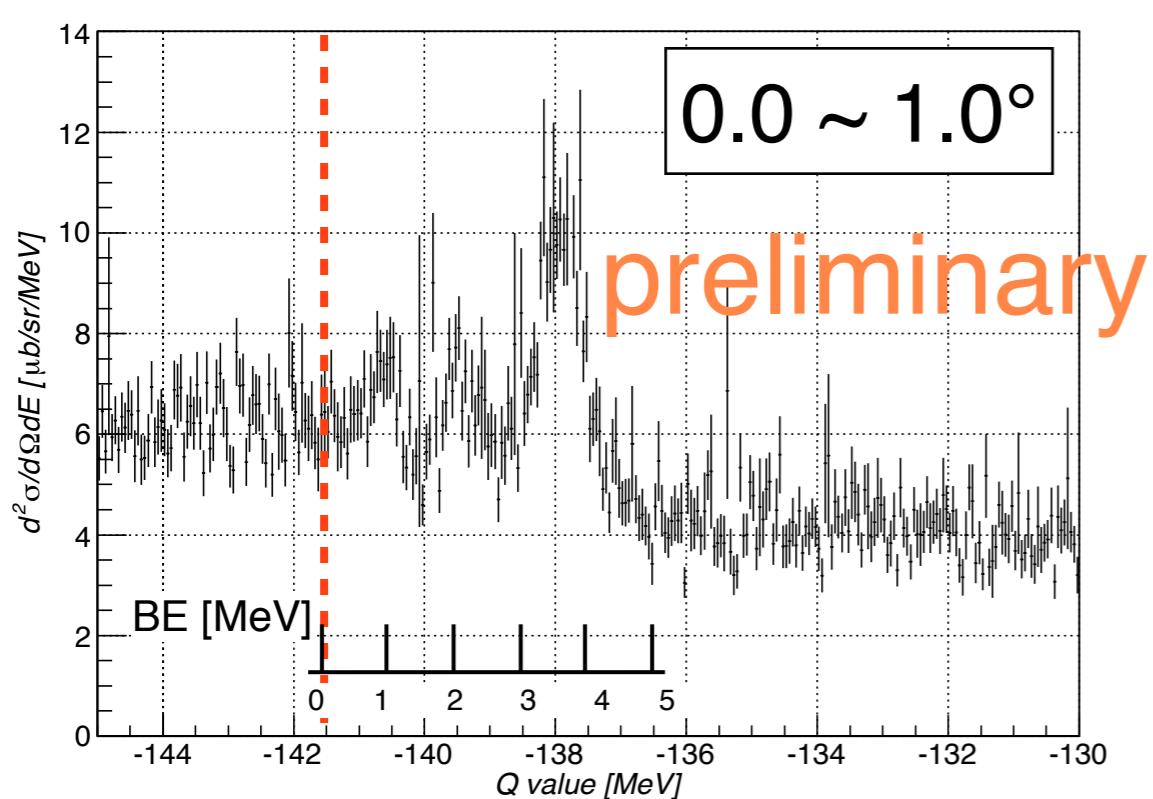
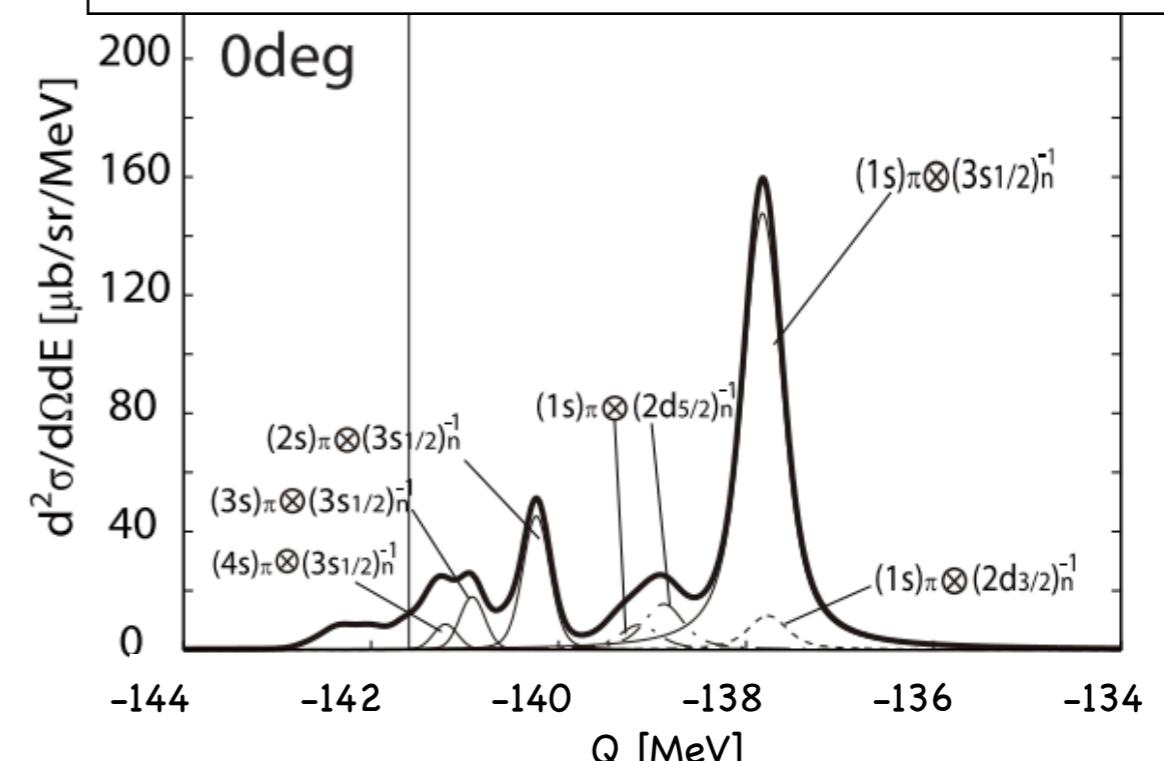
Experimental Q-value spectra



Theoretical spectrum ($\Delta E = 300$ keV)



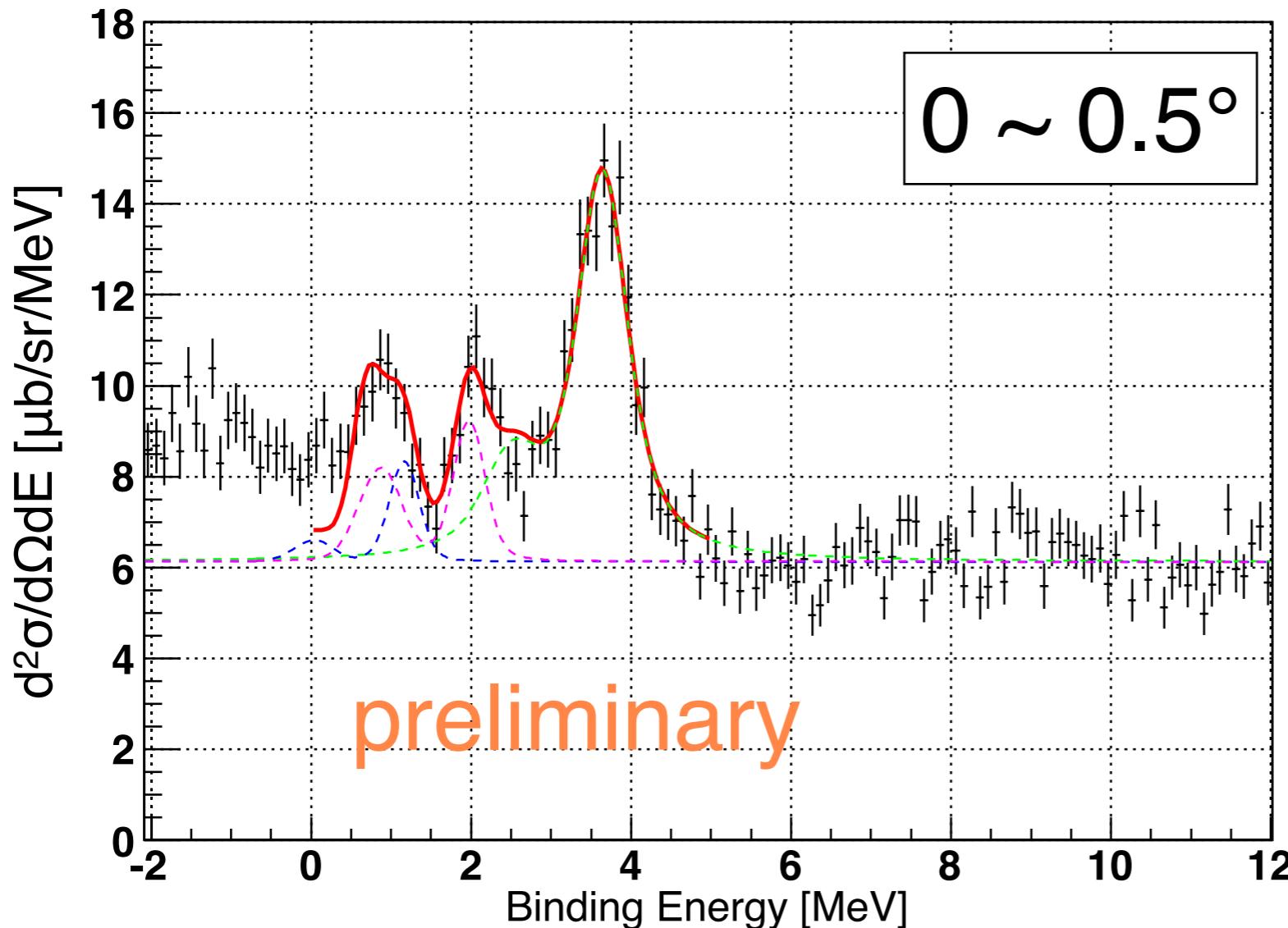
Experimental Q-value spectra

Theoretical spectrum ($\Delta E = 300 \text{ keV}$)

First observation of

- deeply bound pionic states in ^{121}Sn
- angular dependence of $(d, {}^3\text{He})$ reaction

Decomposition of Spectra



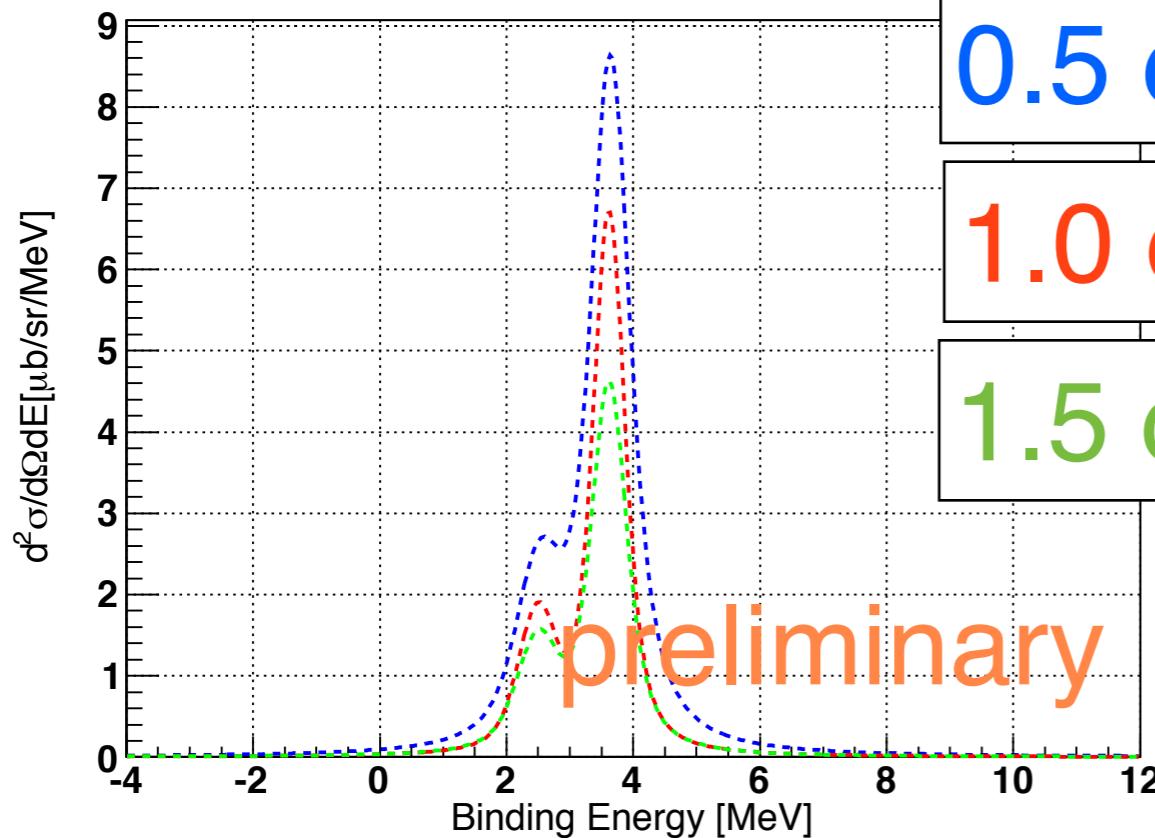
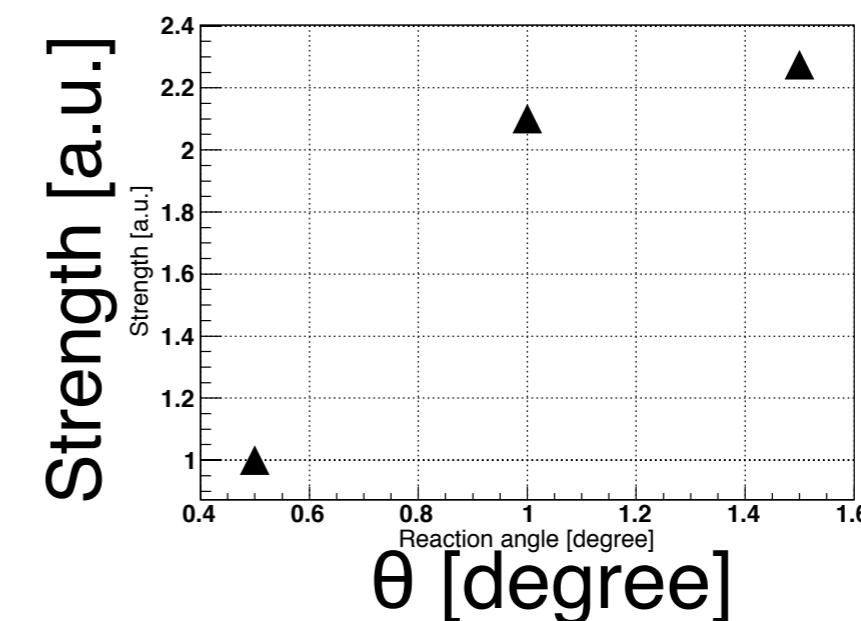
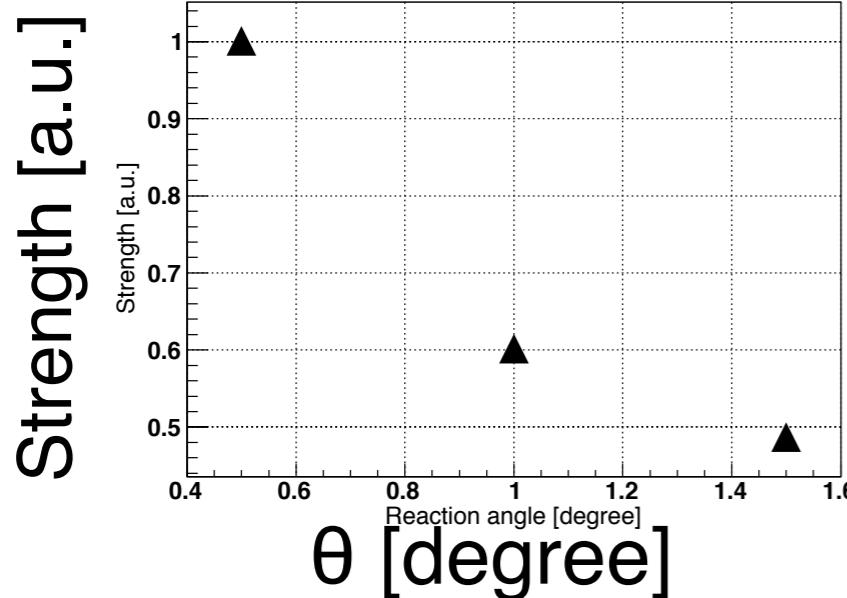
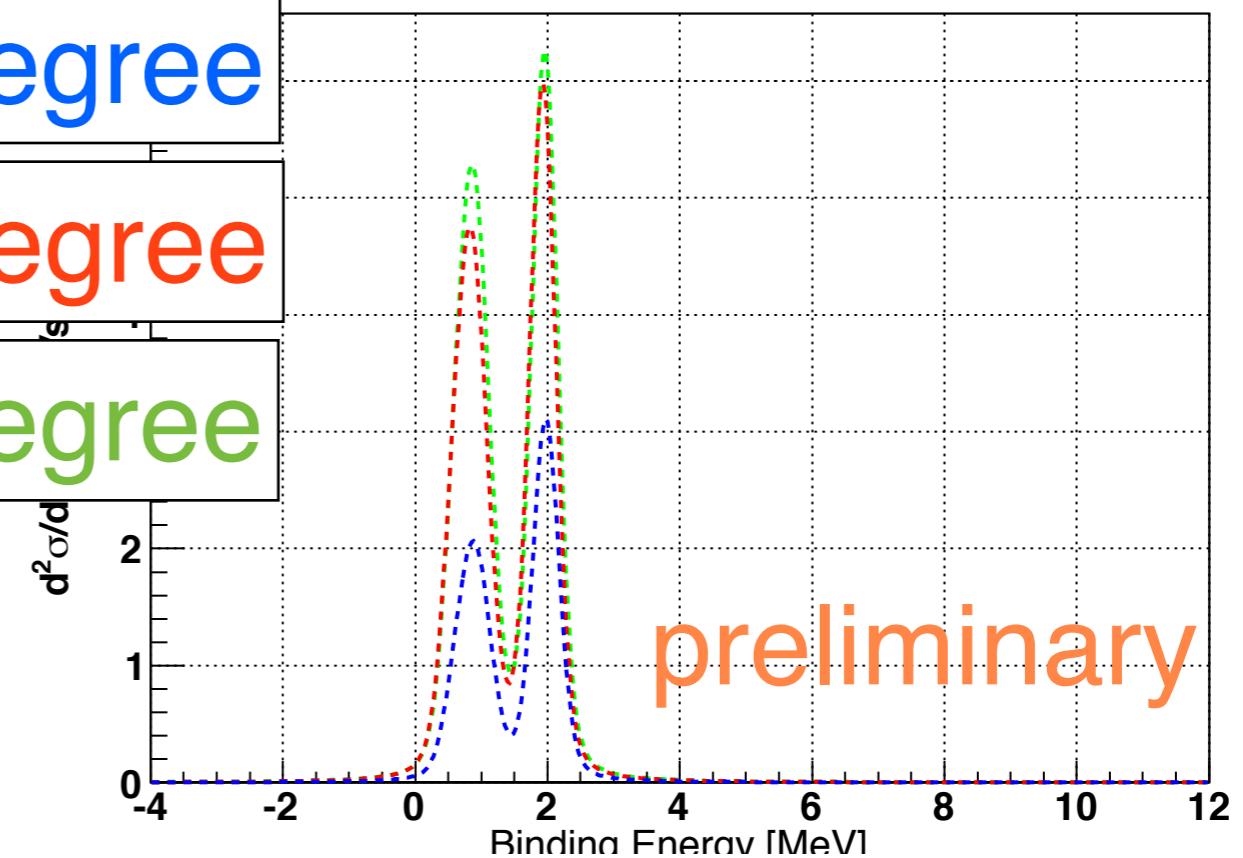
fit function; $F(E) = F_{bg} + \sum_{nl} F_{nl}$

$F_{nl}(E) = \underline{S_{nl}} \sum_{j_m} N_{eff}^{nl, j_m} \text{Voigt}(E - \underline{B.E.}_{\pi^-}^{nl} + E_{hole}^{j_m}, \underline{\Gamma}_{\pi^-}^{nl}, \sigma_{exp})$

nl ; index for pionic states / j_m ; index for neutron hole states

free parameter

Decomposition of Spectra

(1s)_{π-} × n hole(2p)_{π-} × n hole

Current status

Finalizing

- optical aberration
- acceptance including angular dependence
- energy calibration by $p(d,^3\text{He})\pi^0$ reaction

main experiment

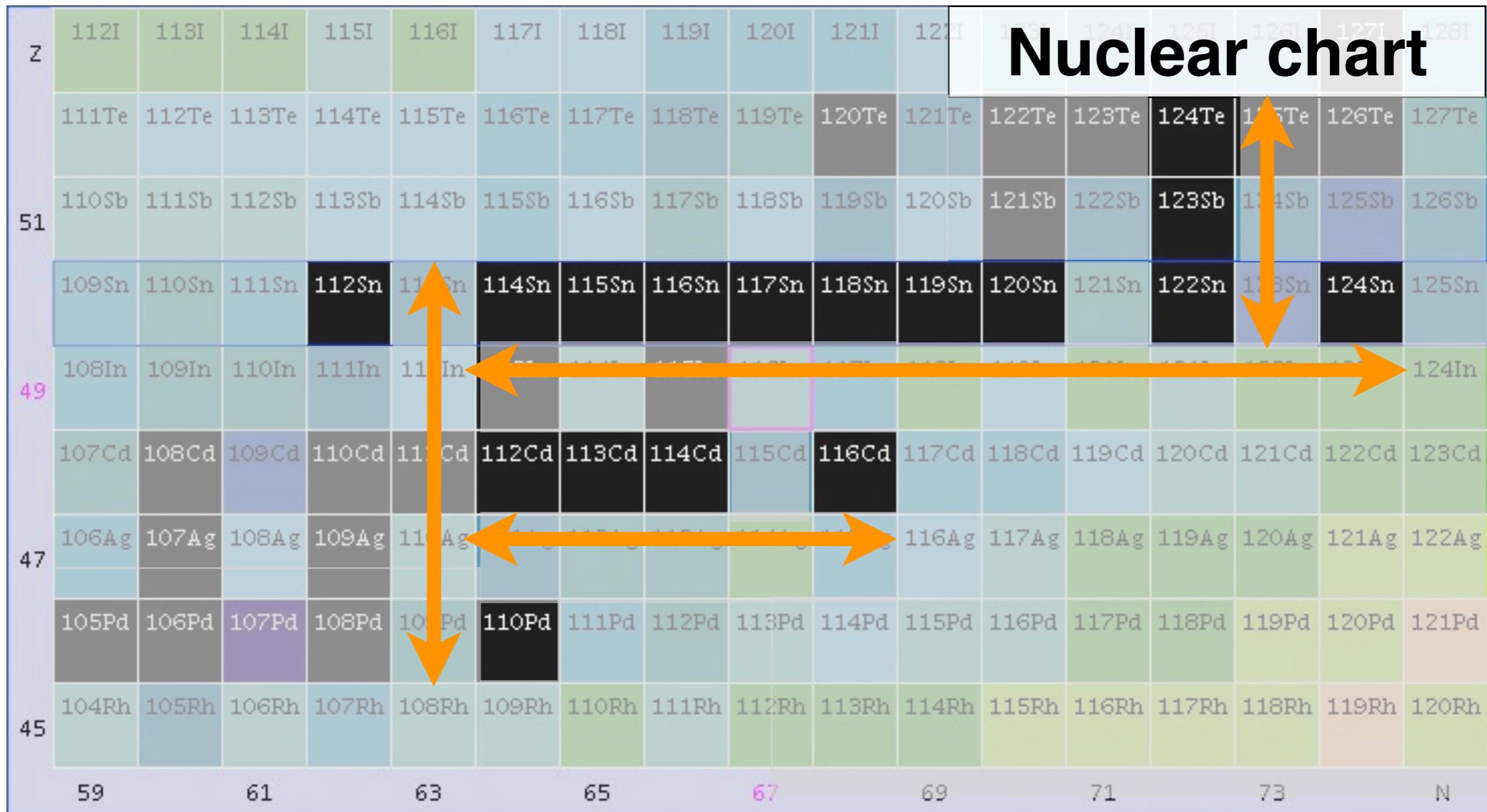
- some improvements are suggested
- preparing for the main experiment

3. Future perspectives

~ pionic unstable nuclei ~

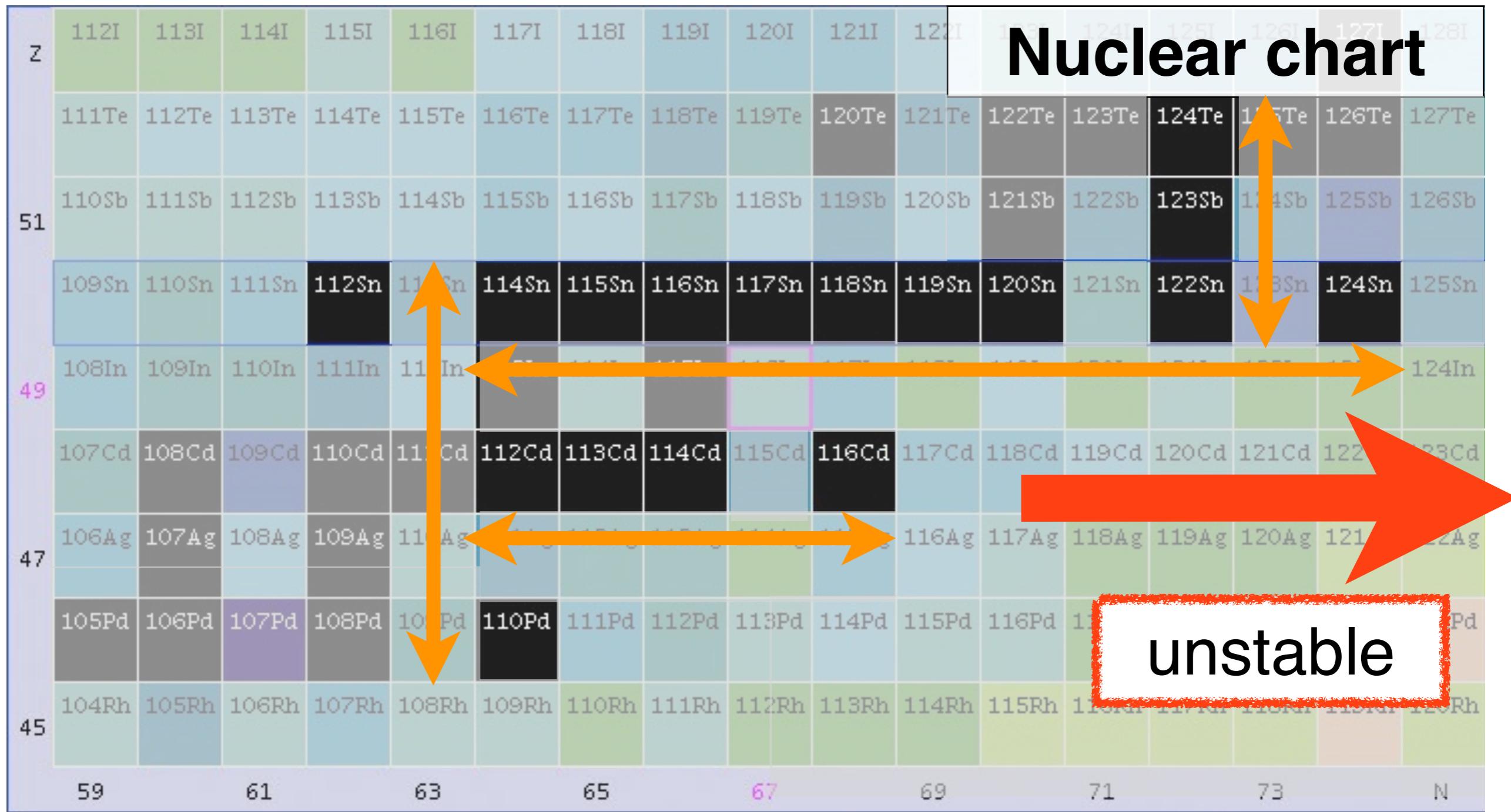
Pionic atom of unstable nuclei

systematic study of stable nuclei



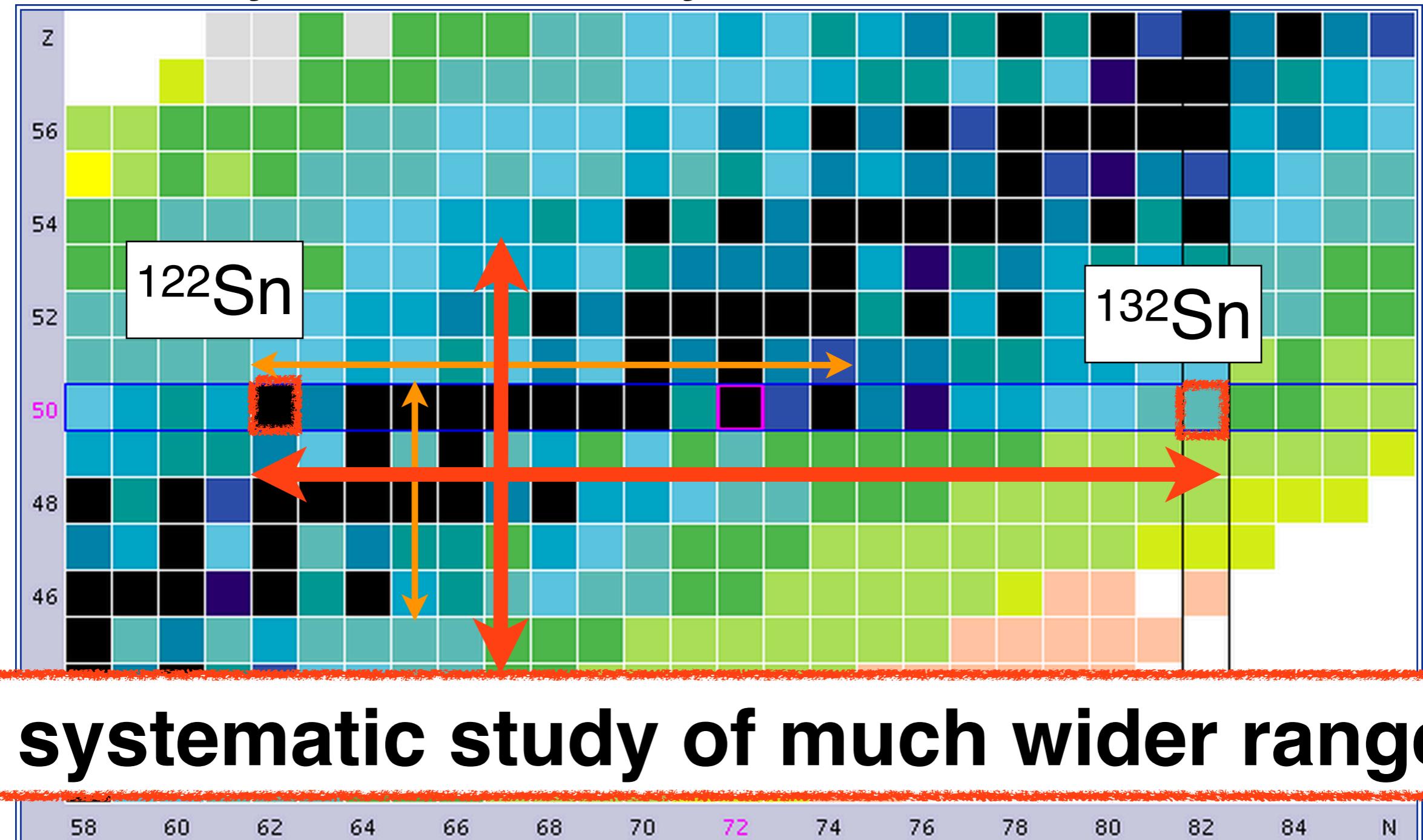
Pionic atom of unstable nuclei

systematic study of stable nuclei

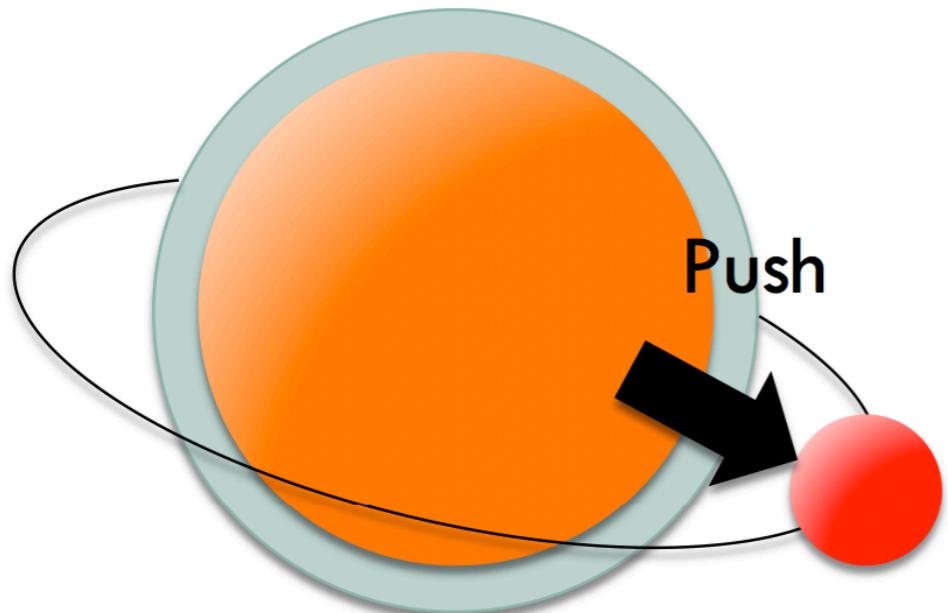


Pionic atom of unstable nuclei

systematic study of unstable nuclei



Pionic atom of neutron rich nuclei



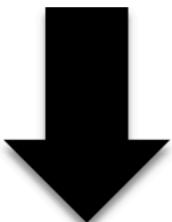
Pion-Neucleus Interaction

$$V_{s-wave} = b_0 \rho + \frac{b_1 (\rho_n - \rho_p)}{\rho^2} + B_0 \rho^2$$

Neutron-Pion: Repulsive!

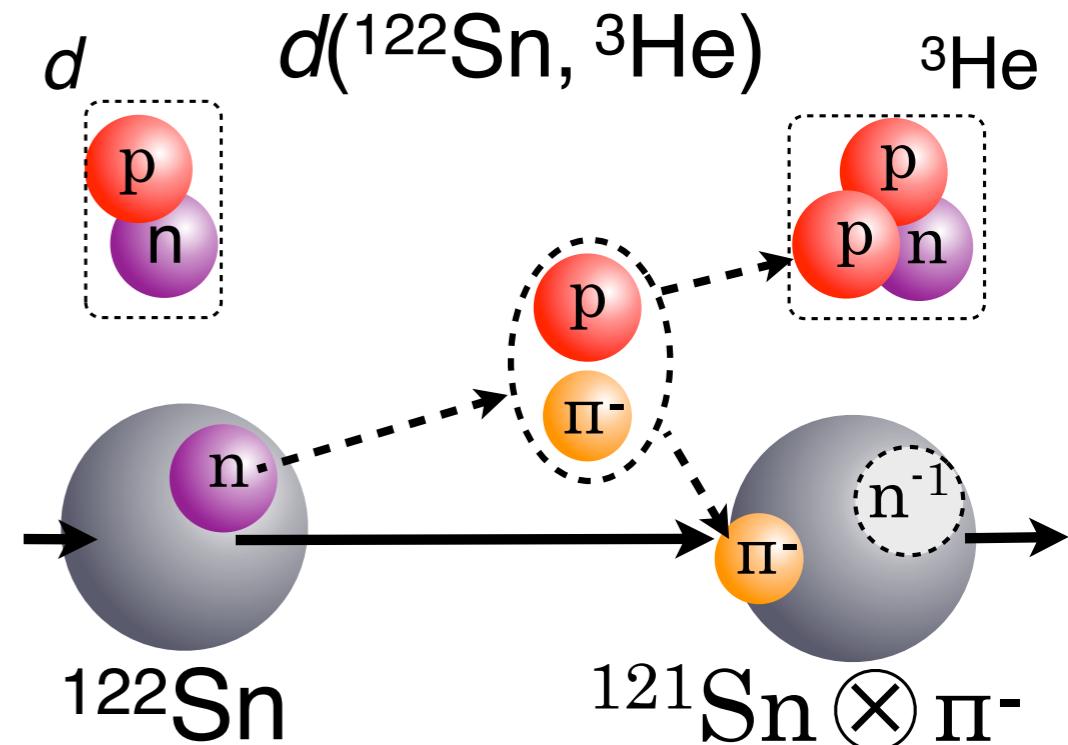
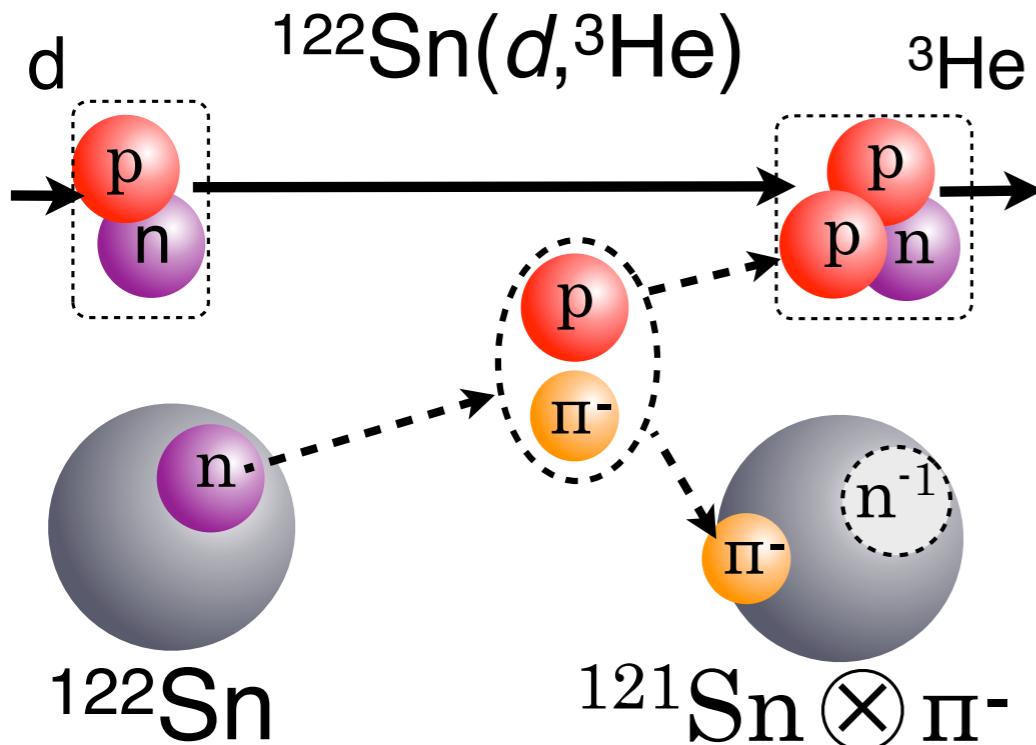
Pion is pushed outward

Different density lower than $0.6 \rho_0$ is probed



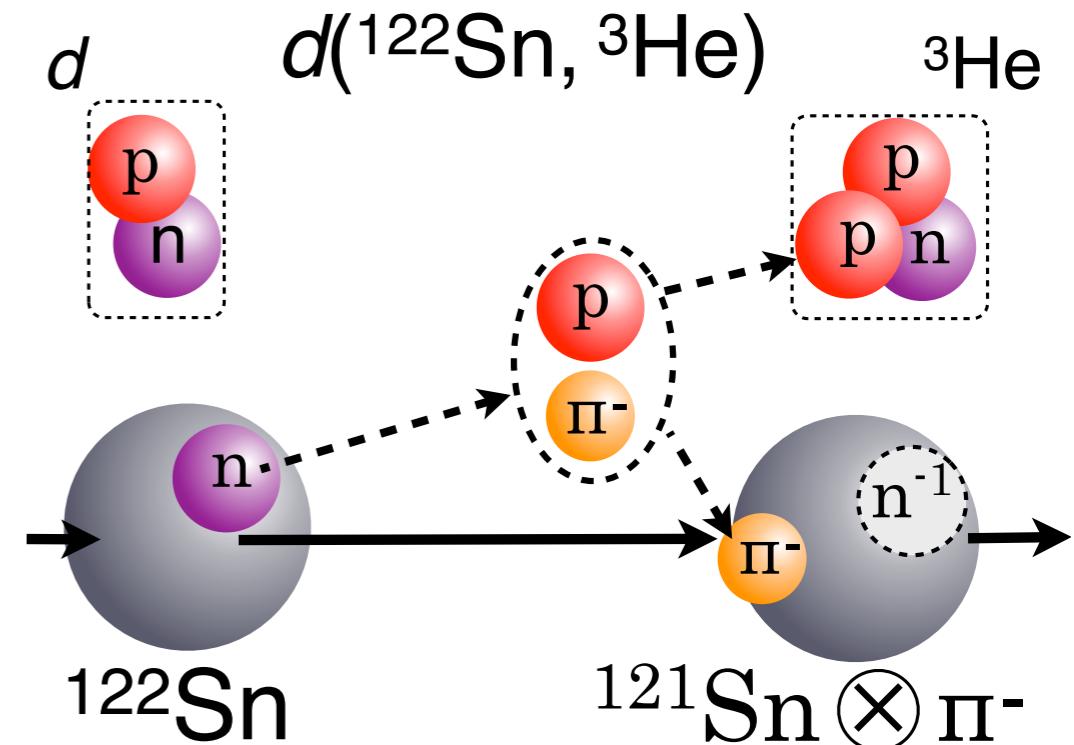
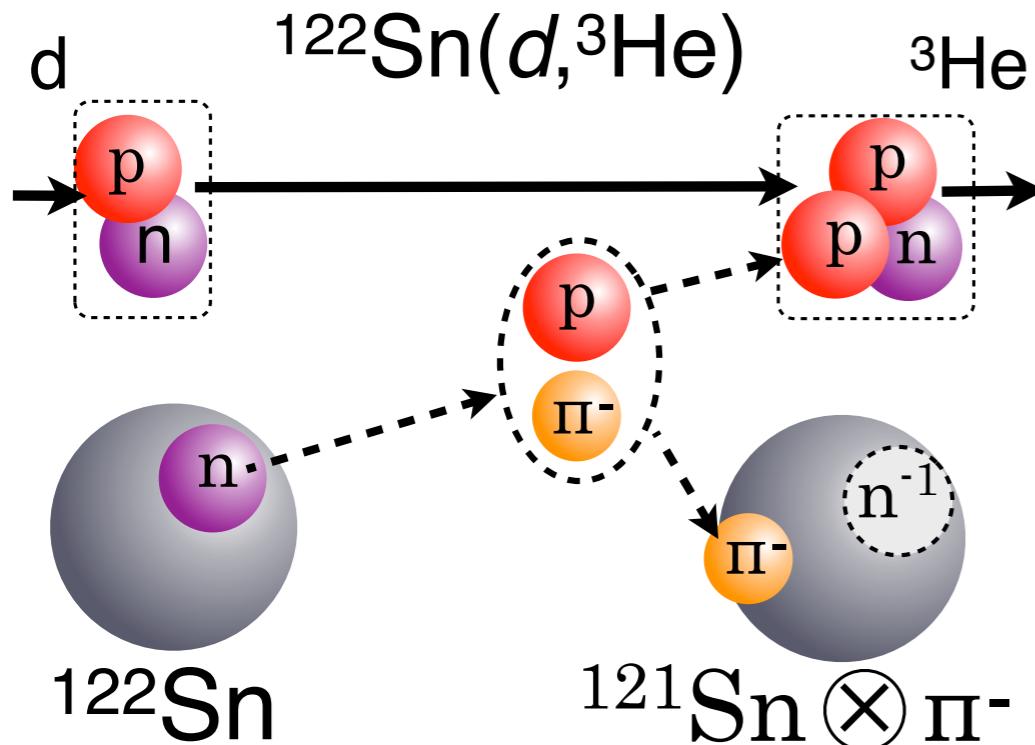
Density dependence of $|\langle \bar{q}q \rangle|$ will be deduced

Inverse kinematics of ($d, {}^3\text{He}$) reaction



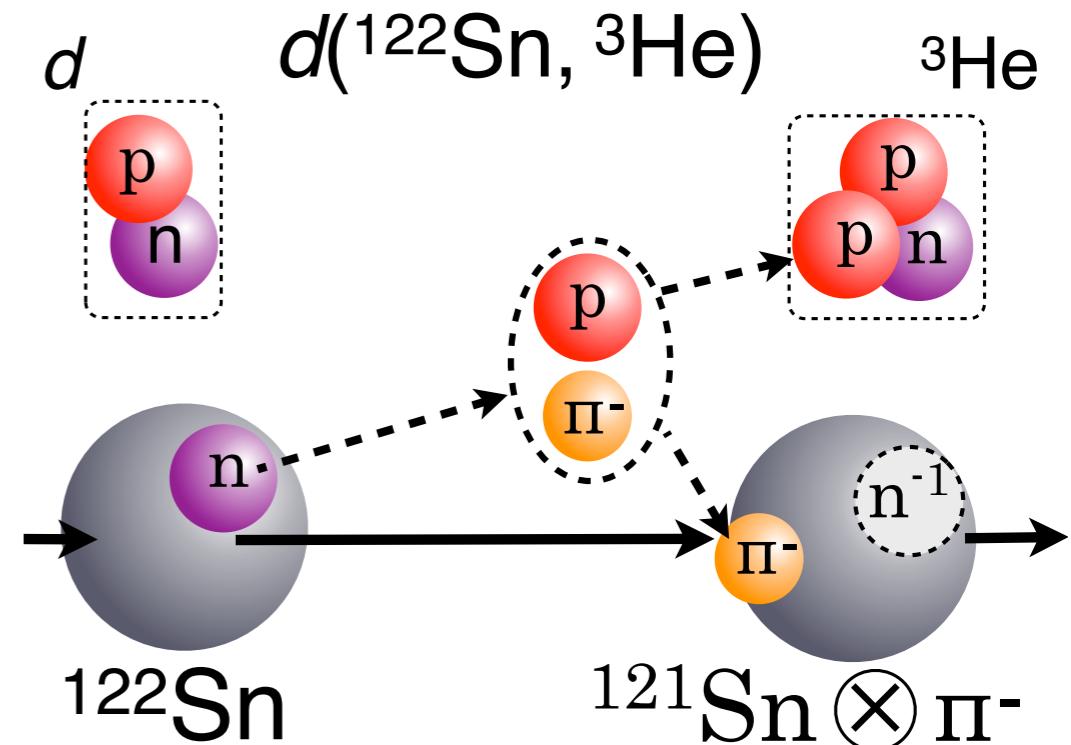
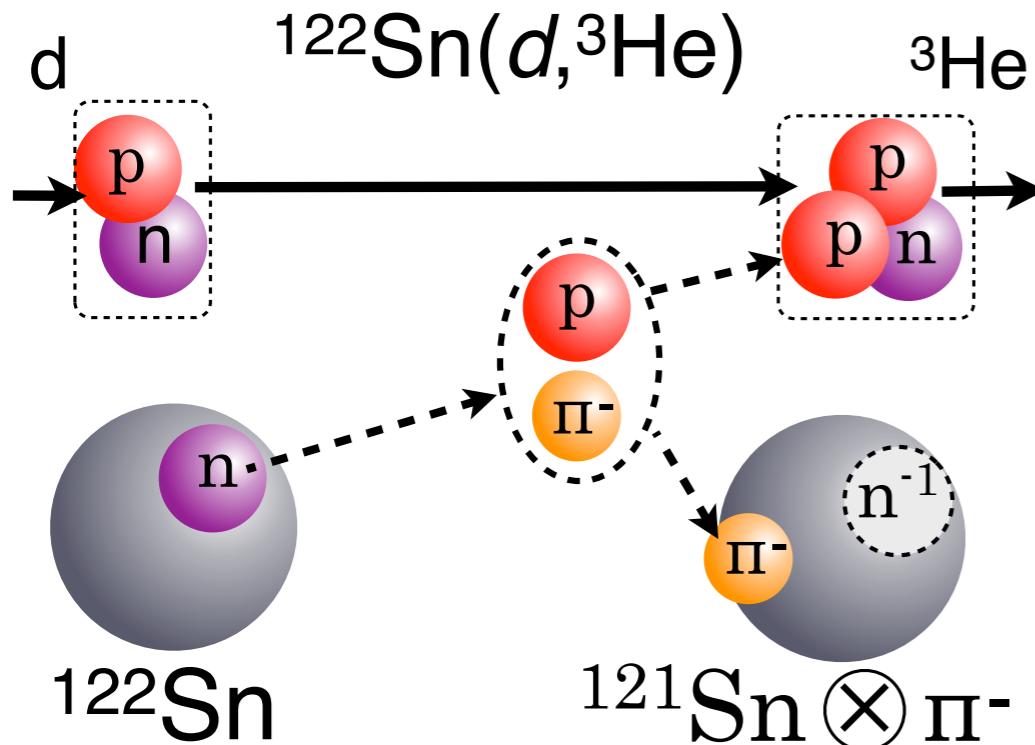
	normal	inverse
Beam	deuteron	(unstable) nuclei
target	stable nuclei	deuteron
Beam energy	250 MeV/u	250 MeV/u
Detecting particle	${}^3\text{He}$	${}^3\text{He}$

Inverse kinematics of ($d, {}^3\text{He}$) reaction



	normal	inverse
Beam	deuteron	(unstable) D_2 gas
target	stable nuclei	deuteron
Beam energy	250 MeV/u	250 MeV/u
Detecting particle	${}^3\text{He}$	${}^3\text{He}$

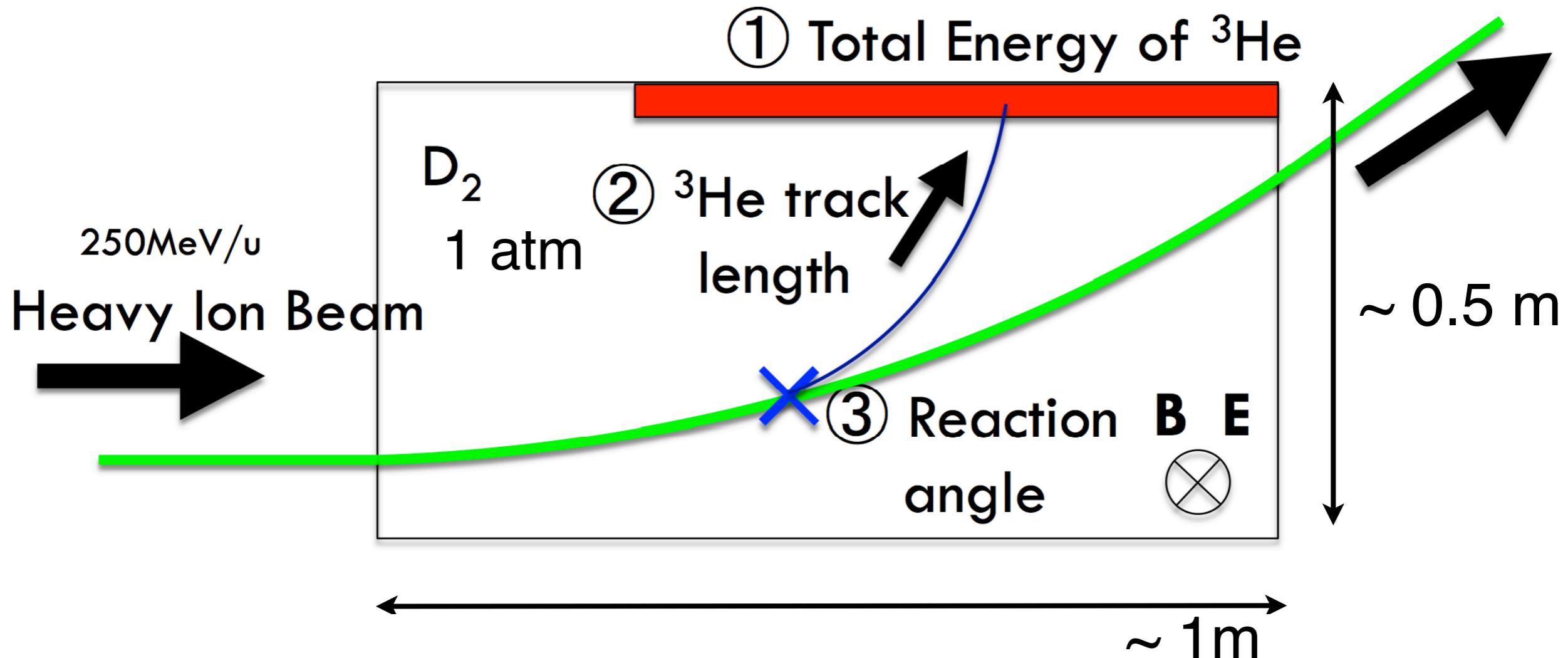
Inverse kinematics of ($d, {}^3\text{He}$) reaction



	normal	inverse
Beam	deuteron	(unstable) D_2 gas
target	stable nuclei	deuteron
Beam energy	$\sim 365 \text{ MeV/u}$	250 MeV/u $\sim 60 \text{ MeV}$
Detecting particle	${}^3\text{He}$	${}^3\text{He}$

Experimental design for inverse kinematics

Detectors ; Active target TPC & Si detector



- ① ${}^3\text{He}$ energy is measured at Si
- ② ΔE deduced from track length
- ③ θ deduced from track



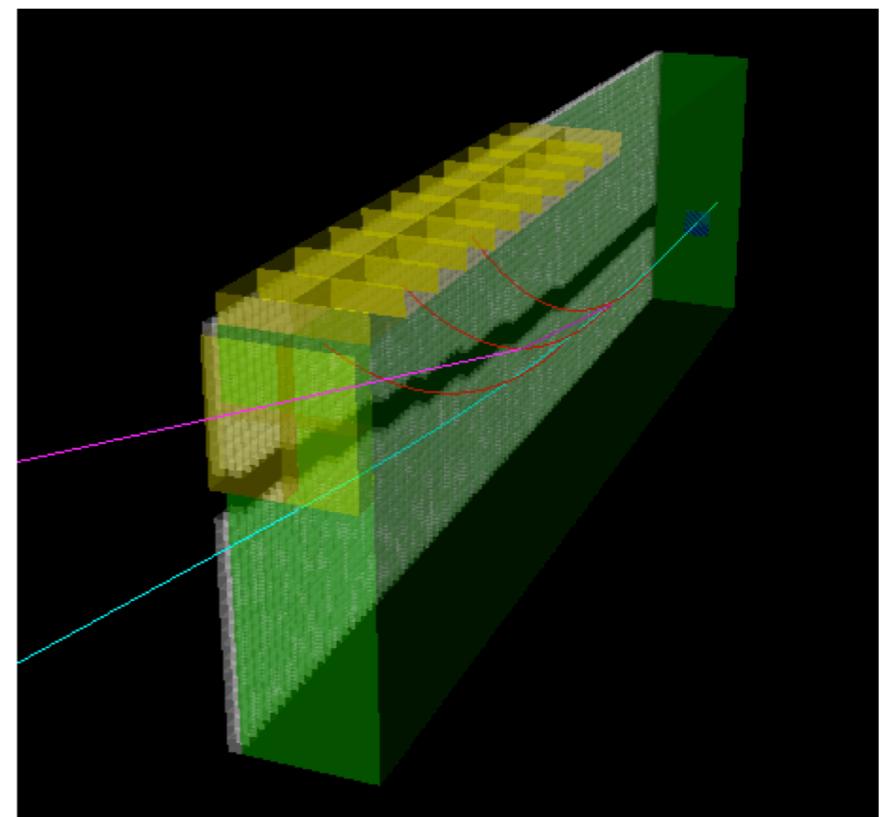
$p_{{}^3\text{He}}$ at reaction point

Yield and resolution estimation (simulation)

	requirement
yield	$10^2 / \text{day}$ (for $1s$ state)
resolution	$\Delta Q < 1 \text{ MeV}$ (FWHM)

input parameter

intensity	$\sim 10^6/\text{s}$ (^{124}Sn)
TPC resolution	$\sim 500 \mu\text{m}$
Si resolution	0.1%@60MeV



GEANT4 simulation

Yield and resolution estimation (simulation)

yield	
Target	D ₂ gas (1 atm, 293.15 K) 100 cm
Beam Intensity [/s]	1 × 10 ⁶
Cross Section [μb]	2.54 × 10 ⁻¹
Yield [/day]	1.1 × 10 ²

Yield and resolution estimation (simulation)

resolution

Cause	ΔQ (FWHM) [keV]
Energy Resolution of Si at $T_{He} \sim 60$ MeV $\sigma_{Si} = 0.1\%$	~ 350
Energy Straggling of 3He in TPC	~ 350
Vertex Reconstruction With Incident Beam $\sigma_{TPC} = 500$ μm	~ 130
Total	~ 500

Yield and resolution estimation (simulation)

requirement	
yield	$10^2 / \text{day}$ (for $1s$ state)
resolution	$\Delta Q < 1 \text{ MeV}$ (FWHM)
simulation result	
yield	$1.1 \times 10^2 / \text{day}$ (for $1s$ state)
resolution	$\Delta Q \sim 500 \text{ keV}$ (FWHM)

This design satisfy the requirement of the experiment

Current status

We are evaluating the resolution of Si detector with a source.



Summary

- ❖ Deeply-bound pionic atom is good probe for quark condensate in finite density.
- ❖ We are planning the precise / systematic measurement of deeply-bound pionic atom and performed a pilot experiment.
- ❖ We succeed to observe the deeply bound pionic states in ^{121}Sn and angular dependence of the (d , ^3He) reaction cross section for the first time with in 15-hour measurement.
- ❖ Now we are finalizing the result of the pilot experiment to extract binding energy and width of deeply bound pionic states.
- ❖ We are also preparing for main experiment with ^{122}Sn target.
- ❖ Experiment with inverse kinematics enable us to produce unstable pionic atom and to approach the density dependence of quark condensate.
- ❖ The feasibility was confirmed by GEANT4 simulation and we are starting the study of detectors.

Deduction of b_1 from experimental data

TABLE I. Observed binding energies (B_{1s}) and widths (Γ_{1s}) of the $1s \pi^-$ states in $^{115,119,123}\text{Sn}$ isotopes.

Isotope	B_{1s} (MeV)	Stat.	ΔB_{1s} (MeV) Syst.	Total	Γ_{1s} (MeV)	Stat.	$\Delta\Gamma_{1s}$ (MeV) Syst.	Total
^{115}Sn	3.906	± 0.021	± 0.012	± 0.024	0.441	± 0.068	± 0.054	± 0.087
^{119}Sn	3.820	± 0.013	± 0.012	± 0.018	0.326	± 0.047	± 0.065	± 0.080
^{123}Sn	3.744	± 0.013	± 0.012	± 0.018	0.341	± 0.036	± 0.063	± 0.072

K. Suzuki et al., PRL92 072302 (2004)

π -A s-wave optical potential

$$V_s(r) = -\frac{2\pi}{\mu} [\epsilon_1 b_0 \{\rho(r) + b_1 \delta\rho(r)\} + \epsilon_2 B_0 \rho(r)^2].$$

4 parameters (b_0 , b_1 , $\text{Re}B_0$, $\text{Im}B_0$) in equation
 \Leftrightarrow 2 parameters from experiment BE_{1s} , Γ_{1s}

Physics motivation; quark condensate

Measurement of quark condensate in finite density

Gell-Mann-Oakes-Renner relation

M. Gell-Mann *et al.*, PR175(1968)2195.

$$f_\pi^2 m_{\pi^\pm}^2 = -m_q \langle \bar{q}q \rangle + O(m_q^2)$$

f_π : pion decay constant

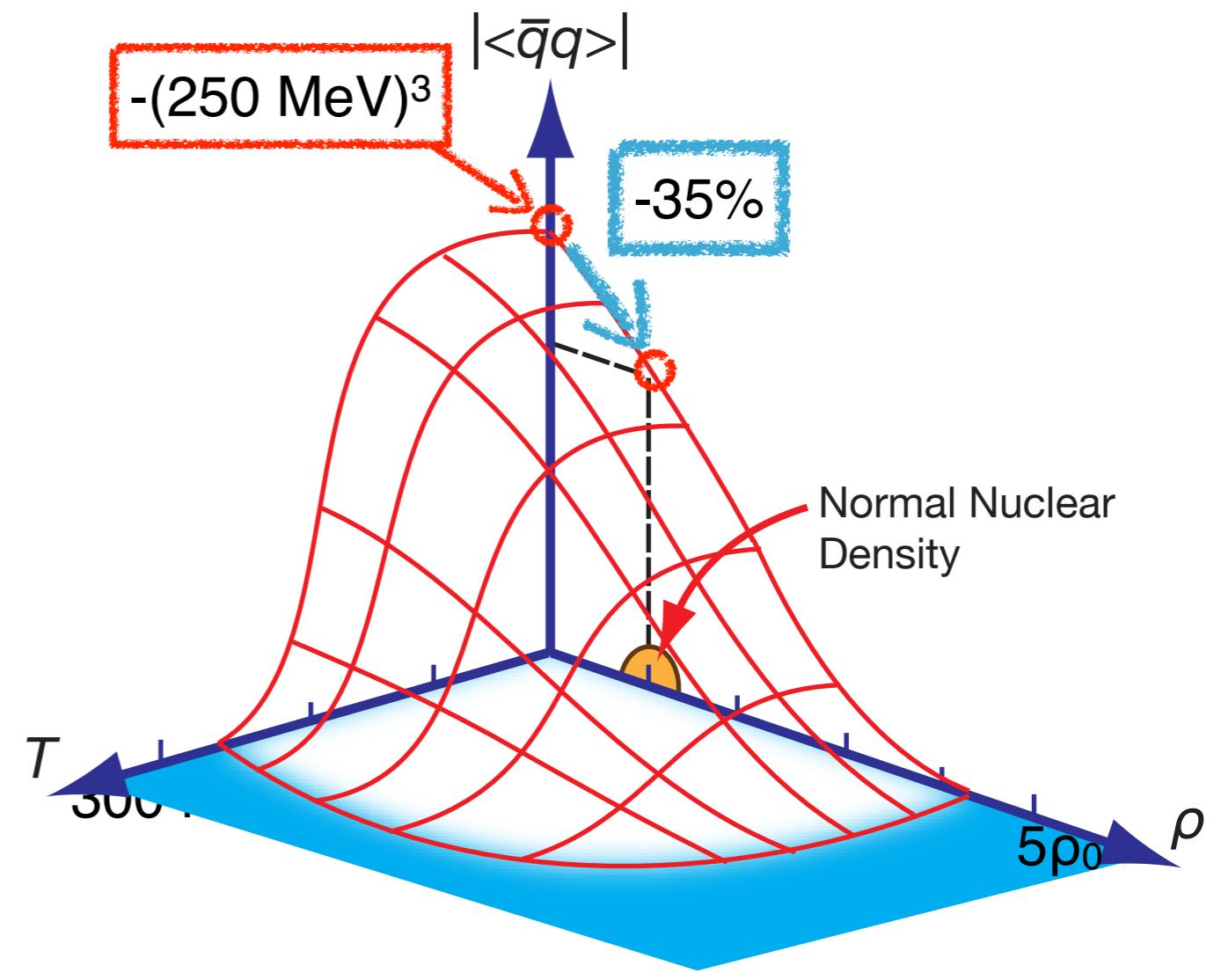
Tomozawa-Weinberg relation

Y.Tomozawa, NuovoCimA46(1966)707.

S.Weinberg, PRL17(1966)616.

$$b_1 \simeq \frac{1}{4(1 + m_\pi/M_{Nucleon})} \frac{m_\pi}{2\pi f_\pi^2}$$

b_1 : isovector scattering length



W. Weise, NPA553(93)59.

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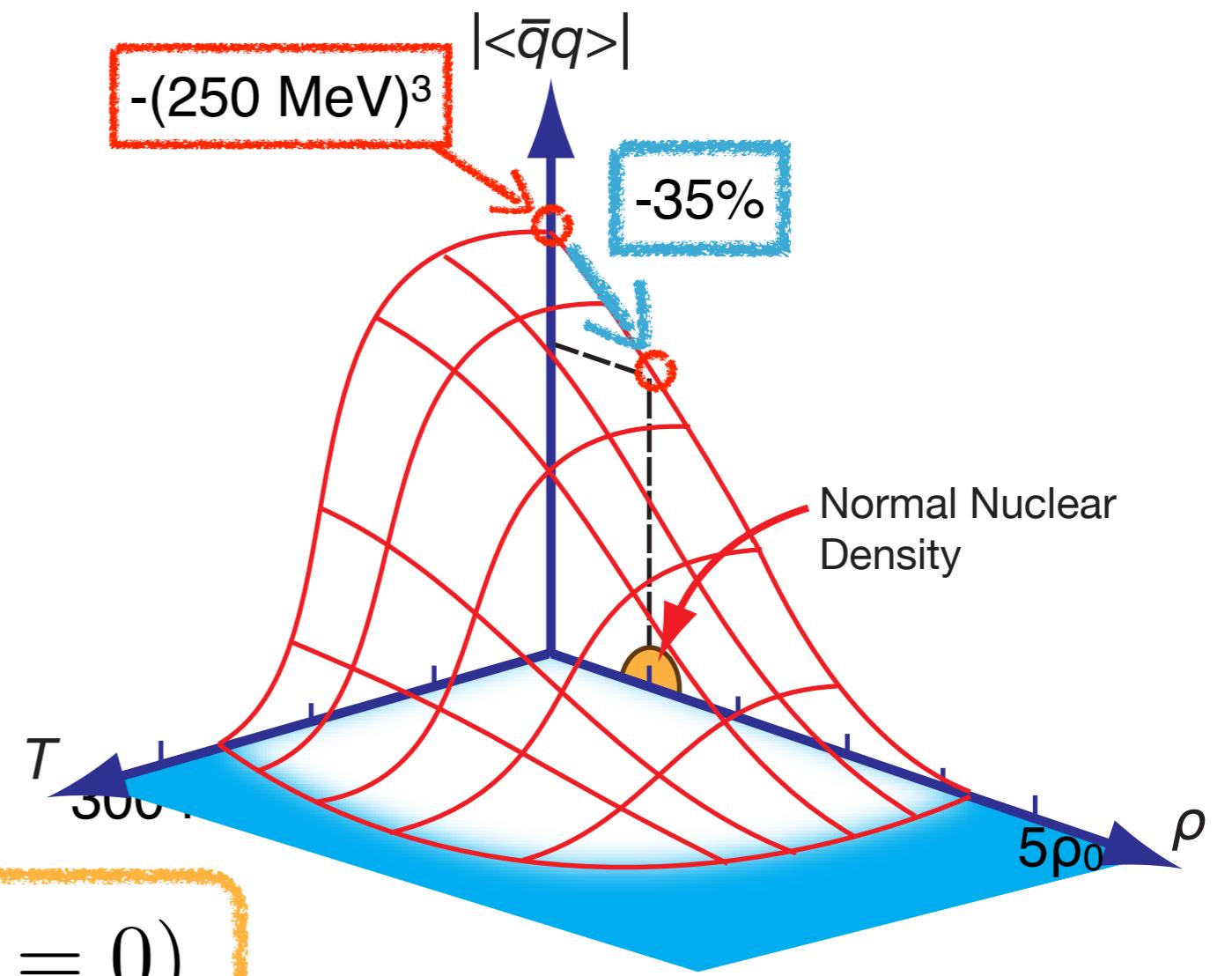
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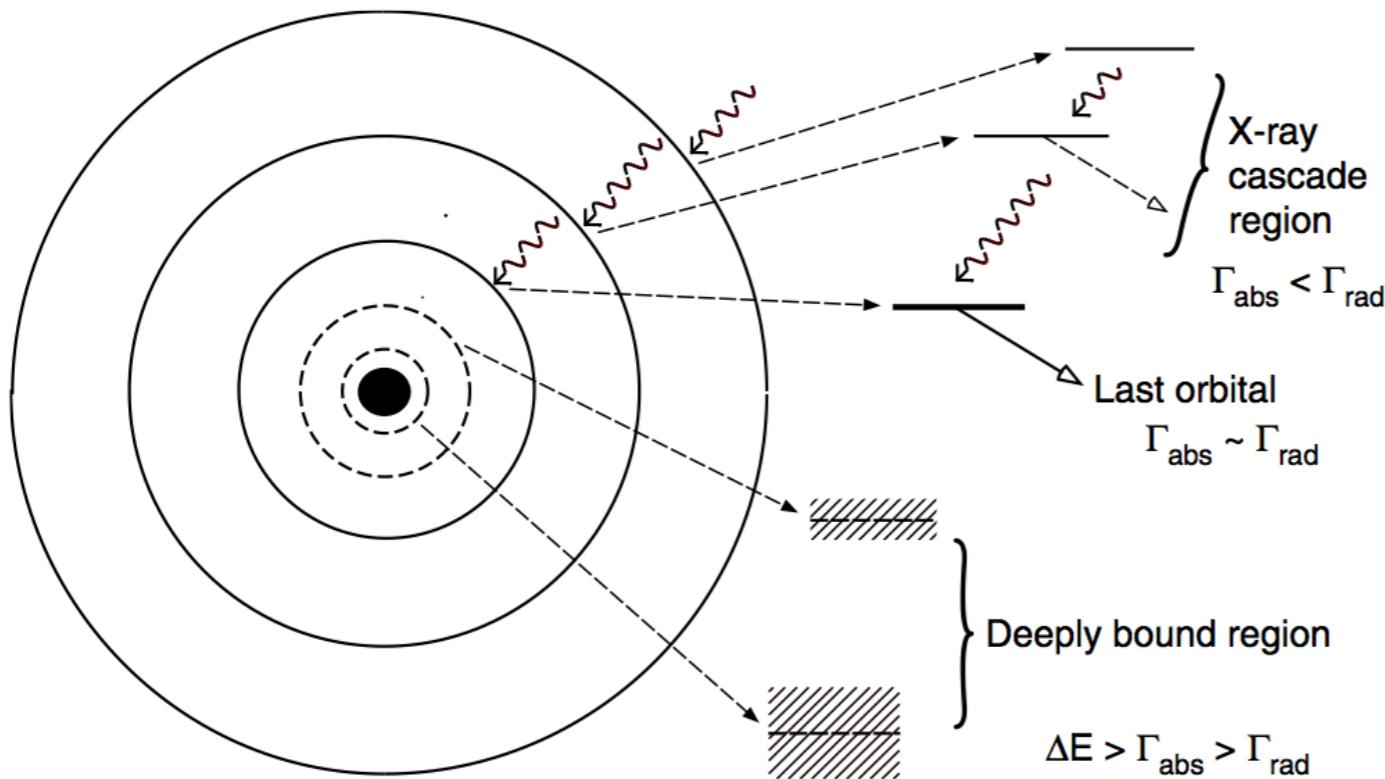
$$\frac{\langle \bar{q}q \rangle_{\rho=\rho_0}}{\langle \bar{q}q \rangle_{\rho=0}} \simeq \left(\frac{m_\pi^*}{m_\pi} \right)^2 \frac{b_1(\rho=0)}{b_1(\rho=\rho_0)}$$



W. Weise, NPA553(93)59.

Conventional method; use π^- beam

slow pion beam is captured in target nuclei
and cascade emitting X-ray



pionic 1s state in H
 $\rightarrow b_1$ in vacuum

This method is applied only for

- higher orbits
- light nuclei ($\sim {}^{24}\text{Mg}$ for 1s)

Yamazaki *et al*, Phys. Rep. 514, 1(2012)

This method cannot produce “deeply-bound” pionic atom...