

Neutrino Spectral Property at Electroweak Scale Temperature

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References

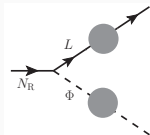
K. Miura, Y. Hidaka, D. Satow and T. Kunihiro, Phys. Rev. D **88**, 065024 (2013).

Baryon Number Asymmetry in Universe (BAU)

BAU: $\eta_B = (n_B - \bar{n}_B)/s \simeq (8.7 \pm 0.3) \times 10^{-11}$.

Sakharov Cond.	Standard Model	Leptogenesis
B	SU _w (2) Sphaleron	(N _R L decay) \rightarrow Sphaleron B
C & CP	CKM (not enough)	Loop Eff. in N _R decay
Non-Equilibrium	No	T < M _R

Leptogenesis, Particularly in Low Energy Scale



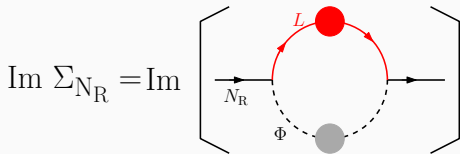
N_R = Right-Handed Neutrino

$L^T = (\nu, l)$

Φ = Higgs Doublet

If the mass difference between two N_R s is in the order of their CP-violating decay width, the CP asymmetry is dynamically enhanced (Pilaftsis ('97)), and **the leptogenesis in the EW scale** can be relevant (Pilaftsis et.al. ('05)).

Spectral Density of Leptons in Leptogenesis



N_R = Right-Handed Neutrino

$L^T = (\nu, l)$

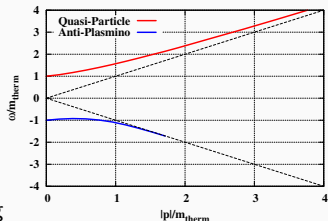
Φ = Higgs Doublet

$$\tilde{G}_{\text{Lepton}}(\mathbf{k}, i\omega_n; T) = \int_{-\infty}^{\infty} d\omega \frac{\rho_{\text{Left}}^{(\nu, l)}(\mathbf{k}, \omega; T)}{\omega - i\omega_n}. \quad (1)$$

If the standard-model leptons have non-trivial spectral properties in EW scale plasma, the lepton number creation via the N_R decay may be significantly modified (c.f. Kiessig et.al., PRD. 2010).

Spectral Property of Finite T Gauge Theory

- There is a growing interest in the collective nature of the fermions in the scenario of thermal leptogenesis (Drewes, arXiv:1303.6912).
- In QED and QCD at extremely high T , the Hard Thermal-Loop (HTL) approx. indicates that a probe fermion interacting with thermally excited gauge bosons and anti-fermions admits a **collective excitation** mode (See, The text book by LeBellac).
- In the neutrino dispersion relation in the electroweak scale plasma, the existence of a novel branch in the ultrasoft-energy region has been indicated by using the HTL and the unitary gauge (Boyanovsky, PRD. 2005).



From QGP to Particle Cosmology

Goal: We investigate Neutrino Spectral Density at $T \gtrsim M_{W,Z}$

- 1 Without restricting ourselves to the dispersion,
- 2 In R_ξ Gauge (Fujikawa et.al. PRD. 1972),
- 3 and Discuss N_R Decay Rate.

Hints in QGP Physics

- **The Spectral Density** of massless fermion coupled with the massive mesonic mode in plasma (an effective description of QGP) has been investigated and shown to have a **Three-Peak Structure with a Ultrasoft Mode**. (Kitazawa et.al. ('05-'06), Harada et.al. ('08), w.o. HTL).
- In particular, when the fermion is coupled with the massive *vectorial* meson, the **Gauge Independent Nature** of the three-peak structure has been confirmed (Satow et.al. ('10)) by using the Stueckelberg formalism.

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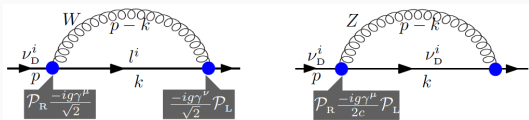
3 Results

- Neutrino Spectral Density: Overview
- Three Peak Structure: In Details
- Gauge Parameter ξ Dependence of Spectral Property

4 Discussion: Toward Leptogenesis

5 Summary

Spectral Density of Massless LH-Neutrino



$$\rho^{(\nu)}(\mathbf{p}, \omega; T) = \sum_{s=\pm} [\mathcal{P}_R \Lambda_{s, \mathbf{p}} \gamma^0 \mathcal{P}_L] \rho_s^{(\nu)}(|\mathbf{p}|, \omega; T)$$

$$\mathcal{P}_{L/R} = \frac{1 \mp \gamma_5}{2}, \quad \Lambda_{\pm, \mathbf{p}} = \frac{1 \pm \gamma^0 \boldsymbol{\gamma} \cdot \mathbf{p} / |\mathbf{p}|}{2} \quad (2)$$

$$\rho_{\pm}^{(\nu)}(|\mathbf{p}|, \omega; T) = \frac{-\text{Im} \Sigma_{\pm}^{(\nu)}(|\mathbf{p}|, \omega; T) / \pi}{\{\omega - |\mathbf{p}| \mp \text{Re} \Sigma_{\pm}^{(\nu)}(|\mathbf{p}|, \omega; T)\}^2 + \{\text{Im} \Sigma_{\pm}^{(\nu)}(|\mathbf{p}|, \omega; T)\}^2}.$$

Higgs Effective Potential (R_ξ Gauge)

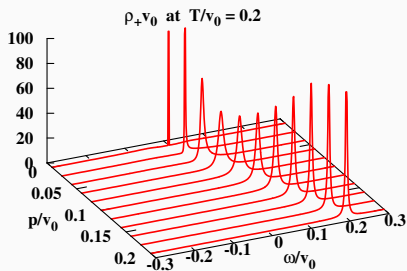
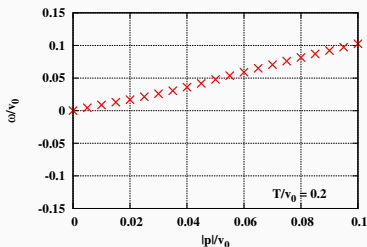
$$V_{\text{eff}} = -\frac{\mu_0^2}{2} \left[1 - \frac{T^2(2\lambda + 3g^2/4 + g'^2/4)}{4\mu_0^2} \right] v^2(T) + \frac{\lambda}{4} v^4(T), \quad (3)$$

$$M_W(T) = \frac{gv(T)}{2} + \mathcal{O}(gT), \quad M_Z(T) = \frac{\sqrt{g^2 + g'^2} v(T)}{2} + \mathcal{O}(gT) \quad (4)$$

- The ξ dependences cancel out among the Nambu-Goldstone modes, longitudinal modes of weak bosons, and the ghost (Text Book by Kapusta).
- The effective potential leads to the second-order phase transition. Note that in reality the possibility of the strong first-order transition has been ruled out within the standard model (Kajantie et.al.('96), Y.Aoki et.al.('99), Csikor et.al.('99, '00)).
- The temperature region satisfying $M_{W,Z}(T) \lesssim T \ll v(T)$ should exist, and a non-trivial spectral property is anticipated there.

Low Temperature Region

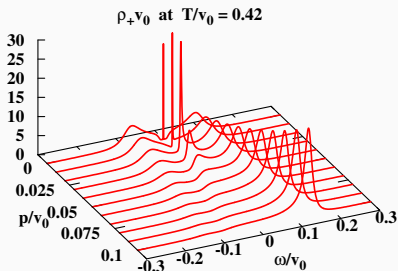
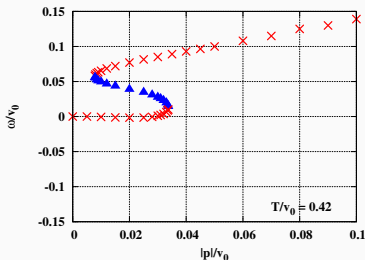
$$T/\nu_0 = 0.2, \quad T/M_W(T) \simeq 0.63, \quad (5)$$



$$\omega - |\mathbf{p}| - \text{Re } \Sigma_+(\omega, |\mathbf{p}|, T) = 0$$

Intermediate Temperature Region I

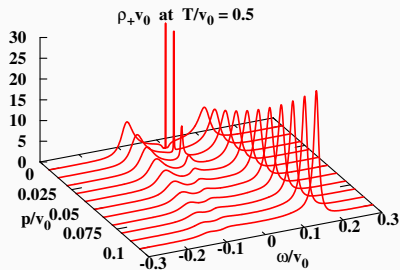
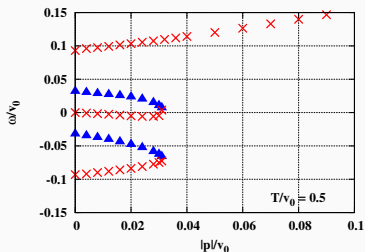
$$T/v_0 = 0.42, \quad T/M_W(T) \simeq 1.45, \quad (6)$$



$$\omega - |\mathbf{p}| - \text{Re } \Sigma_+(\omega, |\mathbf{p}|, T) = 0$$

Intermediate Temperature Region II

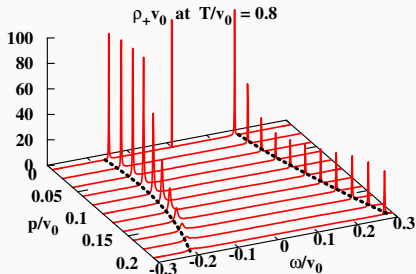
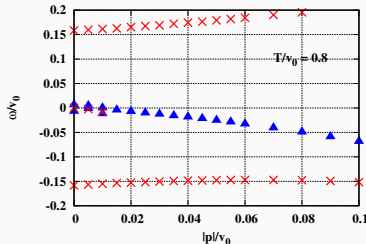
$$T/v_0 = 0.5, \quad T/M_W(T) \simeq 1.83, \quad (7)$$



$$\omega - |\mathbf{p}| - \text{Re } \Sigma_+(\omega, |\mathbf{p}|, T) = 0$$

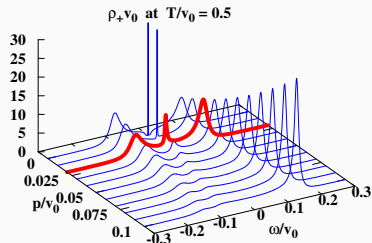
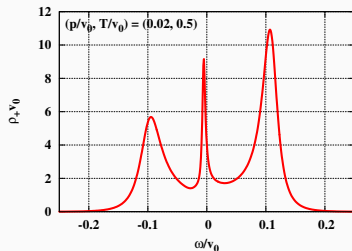
High Temperature Region

$$T/\nu_0 = 0.8, \quad T/M_W(T) \simeq 4.9.$$

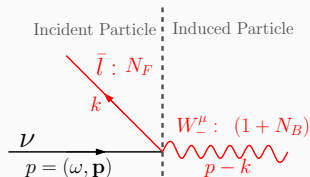


- The spectral property becomes closer to the HTL result.
- $T/\nu(T) \simeq 1.59 > 1$: The additional thermal-loop corrections may modify the spectral property (Hidaka-Satow-Kunihiro, Nucl.Phys.A, 2012).

Spectral Density at $(|p|/v_0, T/v_0) = (0.02, 0.5)$



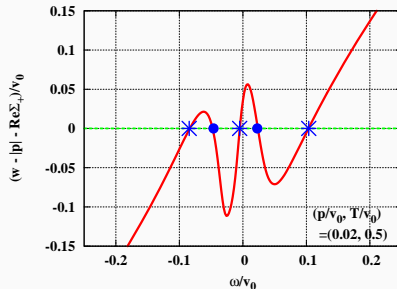
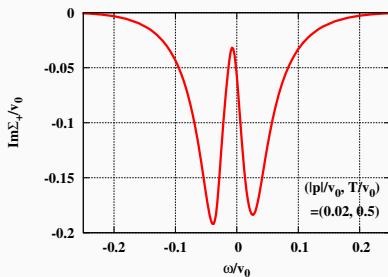
Landau Damping



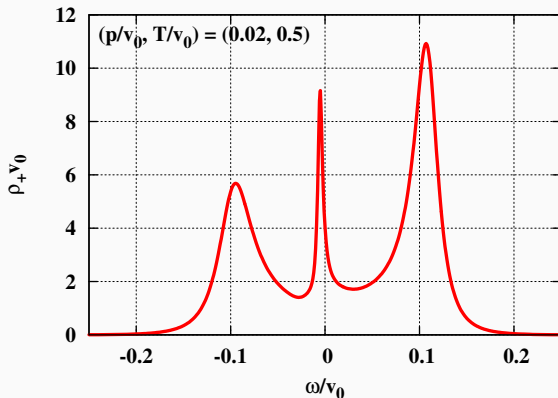
$$\begin{aligned}
 \text{Im}\Sigma_+^{(\nu)} &\ni \int_k \delta\left[\omega + |\mathbf{k}| - \sqrt{|\mathbf{p} - \mathbf{k}|^2 + M_{W,Z}^2}\right] \\
 &\quad \times [N_F(1 + N_B) + N_B(1 - N_F)] \cdot [\dots] \\
 &\ni \int_{x_0}^{\infty} dx \, x N_B(x) = \sum_{n=1}^{\infty} \frac{e^{-nx_0}}{n^2} [1 + nx_0] \\
 &\rightarrow \zeta(2) \quad (T \gg M_{W,Z}, \omega, |\mathbf{p}|), \quad x_0 = \frac{\omega^2 - |\mathbf{p}|^2 - M_{W,Z}^2}{2T(\omega - |\mathbf{p}|)} > 0.
 \end{aligned}$$

Self-Energy in Three-Peak Region

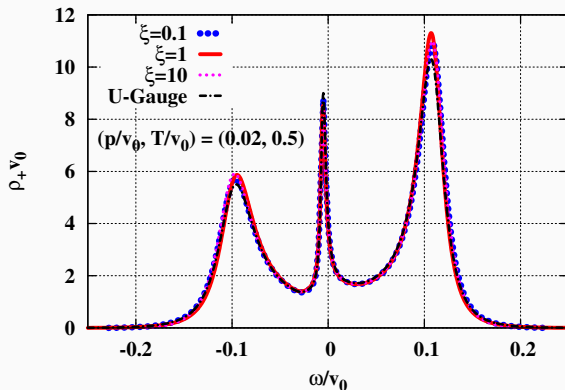
$$\frac{T}{v_0} = 0.5, \quad \frac{|\mathbf{p}|}{v_0} = 0.02. \quad (8)$$



(c.f. Kitazawa-Kunihiro-Nemoto, PTP 2007).

Spectral Density at $(|p|/v_0, T/v_0) = (0.02, 0.5)$ 

ξ Dependence of Three-Peak Spectral Density



Sphaleron Freeze-out Temperature

The net baryon number N_b is produced in the sphaleron process when the changing rate of N_b is larger than the expanding rate of the universe,

$$\left| \frac{1}{N_b} \frac{dN_b}{dt} \right| \geq H(T), \quad (9)$$

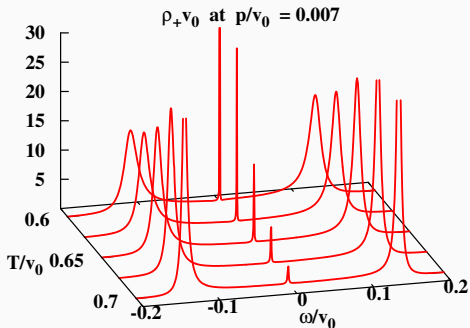
where,

$$H(T) = 1.66 \sqrt{N_{\text{dof}}} \frac{T^2}{M_{\text{PL}}} \simeq T^2 \times 1.41 \times 10^{-18} \text{ (GeV)}, \quad (10)$$

$$\frac{1}{N_b} \frac{dN_b}{dt} = -1023 \cdot g^7 v(T) \exp \left[-1.89 \frac{4\pi v(T)}{gT} \right], \quad (11)$$

and we obtain

$$T \geq T_* \simeq 160 \text{ GeV}, \quad T_*/v_0 \simeq 0.65. \quad (12)$$

Neutrino Spectral Density around $T = T_*$ 

$$\frac{T_*}{\nu_0} \simeq 0.65, \quad \frac{T_*}{\nu(T)} \sim 1. \quad (13)$$

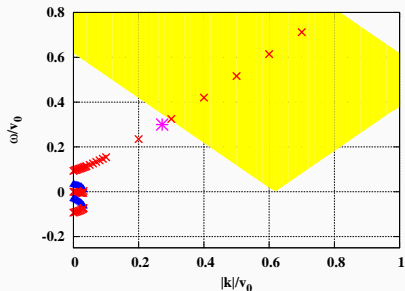
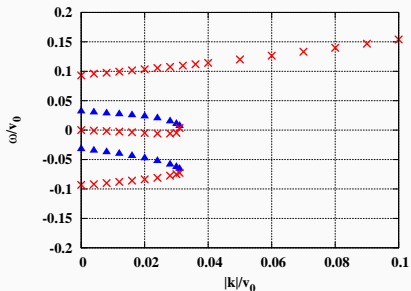
N_R Decay Rate

$$N_R \longrightarrow \mathcal{P}_L \begin{bmatrix} \nu_D \\ l \end{bmatrix} \mathcal{P}_R + \begin{bmatrix} \phi_2 + i\phi_1 \\ h - i\phi_3 \end{bmatrix}, \quad (14)$$

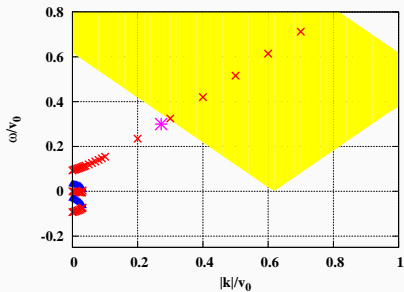
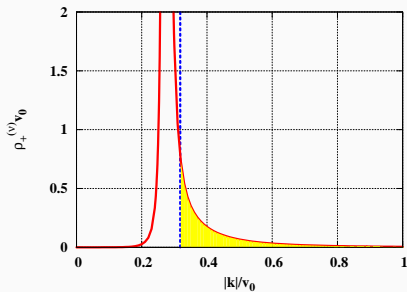
$$\Gamma_{N_R \rightarrow \nu + \phi_{1,2}}(p) = \frac{-(Y^\dagger Y)_{11}}{4p_0 |\mathbf{p}|} \int_{\omega} \int_{k_-(\omega)}^{k_+(\omega)} d|\mathbf{k}| \rho_{\pm}^{(\nu)}(\omega, |\mathbf{k}|) \\ \times \left[\frac{1}{e^{\omega/T} + 1} + \frac{1}{e^{(\omega - p^0)/T} - 1} \right] F_{\text{quad}}(\omega, |\mathbf{k}|, |\mathbf{p}|, M_R). \quad (15)$$

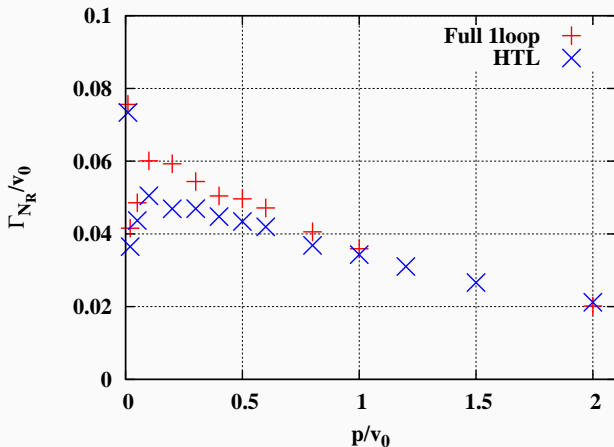
N_R Decay Rate: Support vs 3-Peak Region

Preliminary Result for $(M_R/v_0, |\mathbf{p}|/v_0, T/v_0) = (1.0, 0.5, 0.5)$



N_R Decay Rate: Support vs Width



N_R Decay Rate: Full 1loop vs HTLPreliminary Result for $(M_R/v_0, T/v_0) = (1.0, 0.5)$ 

Summary

- We have investigated the spectral properties of standard-model left-handed neutrinos at finite T around the electroweak scale in R_ξ gauge.
- The spectral density of SM neutrino has the three-peak structure with the ultrasoft mode with a physical significance when $T/M_{W,Z} \gtrsim 1$.
- The three-peak appears at temperature comparable to T_* within the one-loop approximation. However, it may not have a significant contribution to N_R the decay rate...
- The decay rate including the full one-loop spectral density of ν gets larger than the HTL result. This would be because of the width effect.
- Future Work: Fermion mass effect, Resummation scheme at high temperature, Spectral property of N_R , and many others!

Thanks for Your Attention!