

Shear viscosity in a large- N_c NJL model from Kubo formalism

- Advertisement -

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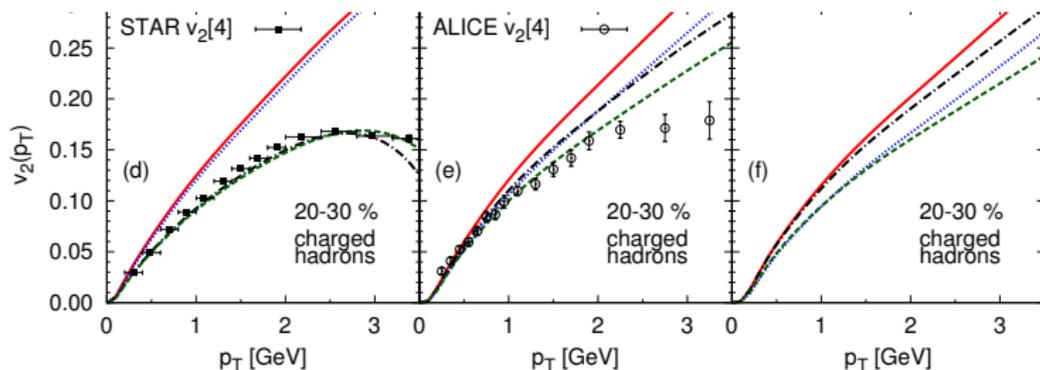
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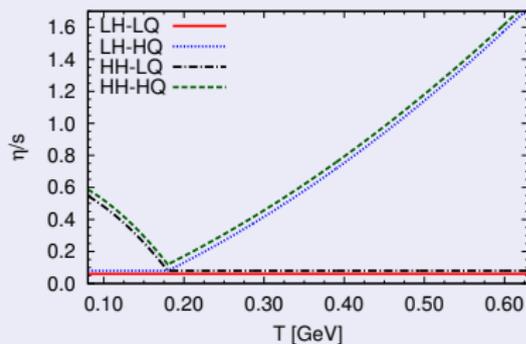


Temperature dependence of shear viscosity



[Niemi, Denicol, Molnár, Rischke: PRL 106 212302 (2011)]

T -dependent ratios η/s



Aim of my work:

Microscopic calculation of the (T, μ) -dependence of η/s within a large- N_c NJL model

NJL model: large- N_c analysis



QCD building blocks: A **four-gluon vertices**,
 B **three-gluon vertices**, and $C \geq N_f$ **quark-gluon vertices**

\Rightarrow Any connected diagram obeys: $A + \frac{1}{2}(B + C) = N_f - 1 + L$

$$N_c \text{ scaling of a } 2N\text{-vertex : } \frac{1}{N_c^{N-1+L}} \quad \text{e.g. } G \sim \frac{1}{N_c}$$

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Application: rediscover standard approaches in many-body physics

DSE $\mathcal{O}(1)$:

BSE $\mathcal{O}(N_c^{-1})$:

Shear viscosity from Kubo formula

Shear viscosity η relates to a four-point correlation function $\langle T_{\mu\nu}(t, \vec{x}) T^{\mu\nu}(0) \rangle_{\text{ret}}$:

$$\Pi(\omega_n) = \gamma_2 \text{---} \text{---} \text{---} \text{---} \gamma_2 = \mathcal{O}(N_c^1) + \mathcal{O}(N_c^0) + \dots$$

$$\eta[\Gamma(p)] = \frac{64N_c N_f}{15\pi T} \int_{-\infty}^{\infty} \frac{d\epsilon}{2\pi} \int_0^{\infty} \frac{dp}{2\pi} \frac{p^6 m^2 \Gamma^2(p) n_F(\epsilon) [1 - n_F(\epsilon)]}{[(\epsilon^2 - p^2 - m^2 + \Gamma^2(p))^2 + 4m^2 \Gamma^2(p)]^2}$$

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perturbative aspects: for “small” Γ we might perform a Laurent-series expansion of $\eta[\Gamma(p)]$:

$$\eta[\Gamma] = \frac{A_{-1}}{\Gamma} + A_0 + A_1 \Gamma + A_2 \Gamma^2 + \dots$$

$$\Rightarrow \eta \cdot \Gamma = A_{-1} + A_0 \Gamma + A_1 \Gamma^2 + A_2 \Gamma^3 + \dots$$

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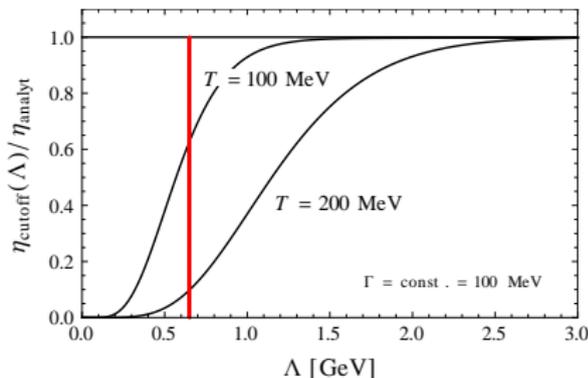
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Hope to see and discuss with you at my poster...



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Motivation

Heavy-ion collisions at RHIC and LHC are one possible approach to explore the QCD phase diagram [1]. In such collisions one probes the phase transition between confined hadronic matter and deconfined quarkonic matter. At the same time this is a transition between phases where chiral symmetry is spontaneously broken and where it is restored.

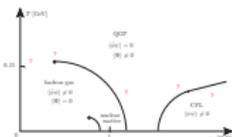


Figure: Sketch of the QCD phase diagram

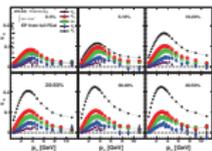


Figure: Flow coefficients in heavy-ion collisions at ATLAS for different centrality classes [2]

Comparisons between hydrodynamic simulations and the measured elliptic flow v_2 indicate that the quark-gluon plasma produced in such collisions behaves as an almost-perfect fluid. The ratio η/s of the quark-gluon plasma is the lowest value found so far in nature. It is known that a T -dependent shear viscosity is crucial to understand the flow coefficients measured at LHC [3].

Aim of this work: Calculate the thermal dependencies of the shear viscosity $\eta(T, \mu)$ within the NJL model.

Shear viscosity from Kubo formalism: a non-perturbative approach

In the Kubo formalism [4] the shear viscosity is expressed as imaginary part of the retarded four-point quark correlator $\Pi^R(p)$. At leading order [5] only ring diagrams contribute to $\Pi(\omega_n, \vec{p})$:

$$\text{---} \text{---} \text{---} = \Pi(\omega_n, \vec{p}) = \text{---} \text{---} \text{---} + \mathcal{O}(N_c^0)$$

Shear viscosity and strong cutoff effects

From a purely mathematical point of view, the shear viscosity $\eta[\Gamma(p)]$ does not require necessarily a momentum-cutoff:

In order for the shear viscosity $\eta[\Gamma]$ as functional of $\Gamma(p)$ to be convergent, the asymptotic $\Gamma(p)$ should not converge too rapidly to zero:

$$\eta[\Gamma(p)] < \infty \Leftrightarrow p^3 e^{-\beta p/2} \in o(\Gamma(p)),$$

$$\begin{aligned} \Gamma_{\text{const}} &= 100 \text{ MeV}, \\ \Gamma_{\text{exp}}(p) &= \Gamma_{\text{const}} e^{-\beta p/8}, \\ \Gamma_{\text{lin}}(p) &= \Gamma_{\text{const}} \frac{\beta p}{1 + (\beta p)^2}, \\ \Gamma_{\text{dis}}(p) &= \Gamma_{\text{const}} \sqrt{\beta p}. \end{aligned}$$

We find within the NJL model a strong cutoff dependence for these four parameterizations:

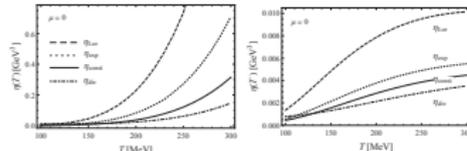


Figure: Shear viscosity WITHOUT applying a momentum cutoff (left) and WITH a physical NJL cutoff $\Lambda = 651 \text{ MeV}$ (right).

The perturbative sector in the NJL model and thermal effects on η/s

Only if the spectral width Γ is small enough, a perturbative treatment of $\eta[\Gamma(p)]$ is possible. Then it can be expanded in a Laurent series:

$$\eta[\Gamma] = \frac{A_{-1}}{\Gamma} + A_0 + A_1 \Gamma + A_2 \Gamma^2 + \dots$$

Only in the perturbative sector ladder-diagram resummation [8, 9, 10] becomes a relevant contribution to the shear viscosity. The full non-perturbative $\eta[\Gamma(p)]$ takes the whole Laurent series into account.