Center Domains and Their Phenomenological Consequences

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Based on S. A. Bass, B. Müller, M.A., PRL 110 (2013) 202301

Some major achievements @RHIC/LHC

- i) measurement of strong elliptic flow success of relativistic viscous hydrodynamics with a very small value of η/s
- ii) measurement of very strong suppression of high energy particles rapid redistribution of the jet energy into the whole solid angle
- iii) observation of constituent quark number scaling law for the elliptic flow of identified hadrons
 - i) iii) : often attributed to the formation of QGP
 - iii) is the most direct evidence to date for quark deconfinement, but little information on dynamical properties of (s)QGP

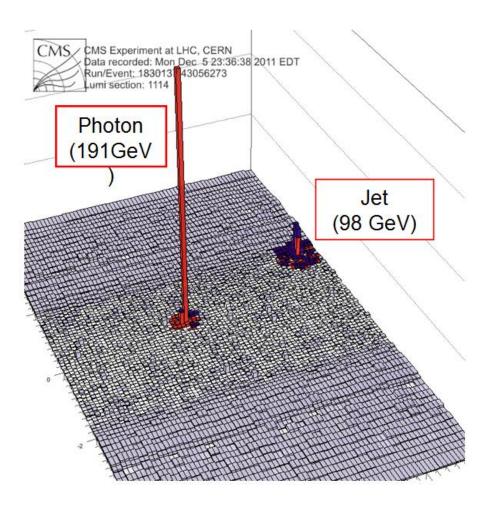
why?

Neither i) nor ii): direct evidence of (s)QGP

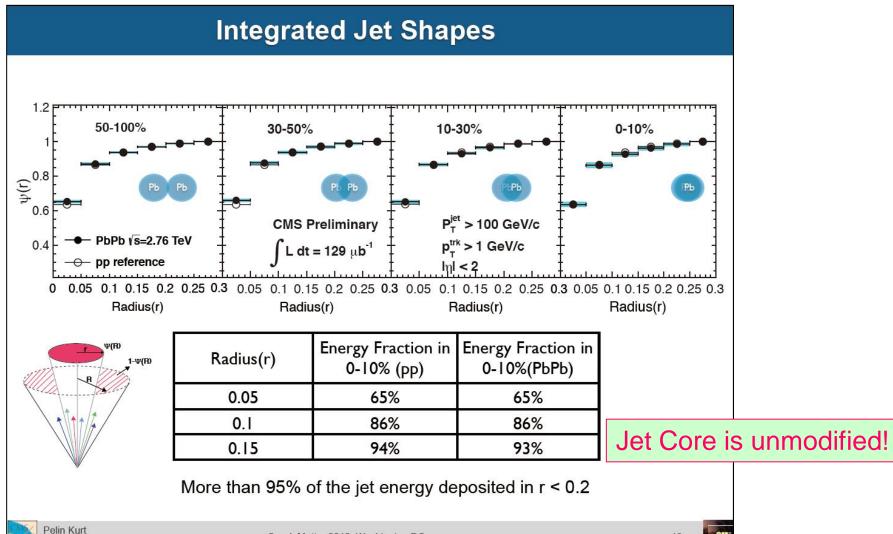
Transport properties of QGP

- i) measurement of strong elliptic flow success of relativistic viscous hydrodynamics with a very small value of η/s
- ii) measurement of very strong suppression of high energy particles rapid redistribution of the jet energy into the whole solid angle
- For i): No first principle nonperturbative calculation of η/s
 2 lattice calculations, but not completely first principle calculation
 AdS/CFT: applicability of AdS/CFT correspondence conjecture to QCD?
- For ii): Most calculations of jet energy-loss: based on pQCD
 Sensitive to the gluon content of the matter, do not distinguish phases
 No explanation so far *for the rapid redistribution of the jet energy*

Remarkable features observed @CMS



Remarkable features observed @CMS

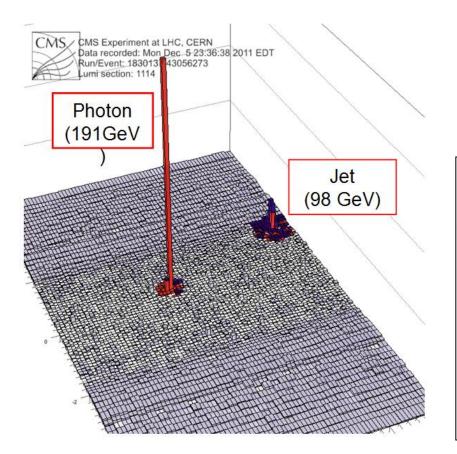


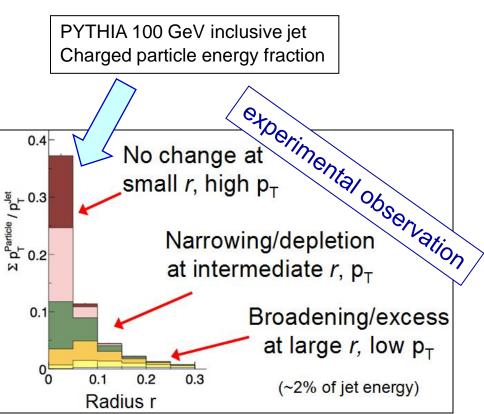
Quark Matter 2012, Washington DC

Vanderbilt University



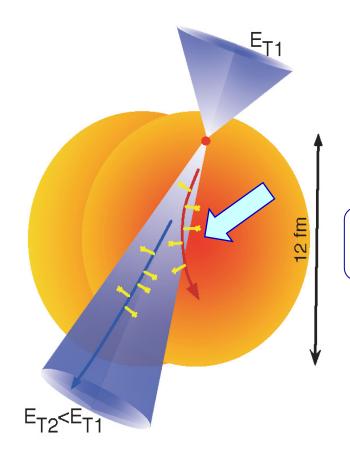
Remarkable features observed @CMS





Roland (CMS), QM2012

Idea of Frequency Collimation

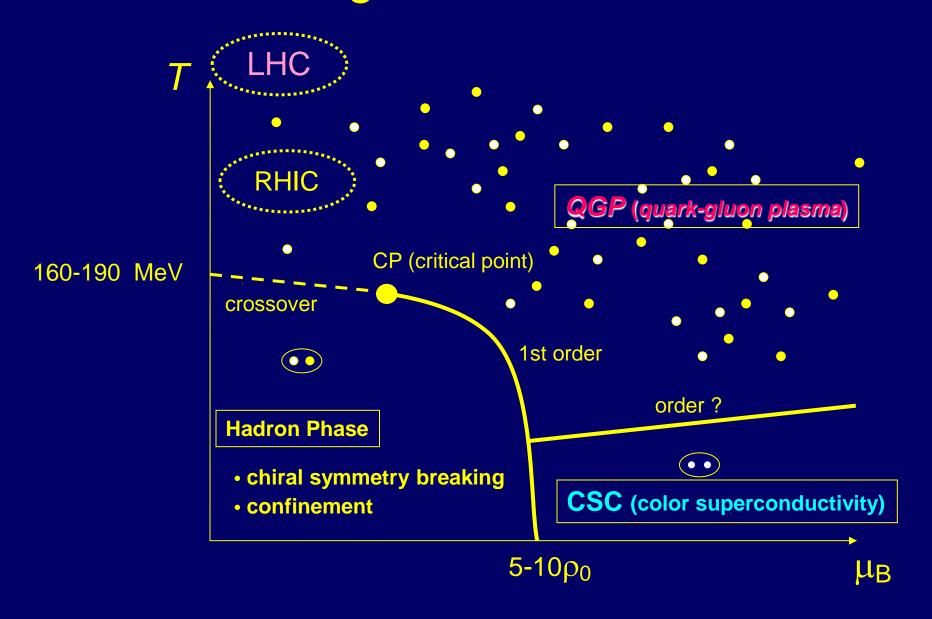


An attractive working conjecture:

Medium acts as a frequency collimator trims away the soft components of the jet

Casalderry-Solana, Milhano, and Wiedemann, J. Phys. G (2011)

QCD Phase Diagram



Order Parameters

- What characterizes QCD phase transition?
 - Chiral Symmetry : $SU(N_f)_L \times SU(N_f)_R$ in the massless limit

Quark condensate $|\langle \overline{\psi}\psi \rangle$: finite value $\rightarrow 0$

• Confinement : Z₃ in the pure gauge theory

Polyakov loop $\langle L(\vec{x}) \rangle : 0 \to 1$, $\exp\left(\frac{2}{3}\pi i\right)$, $\exp\left(\frac{4}{3}\pi i\right) \in Z_3$

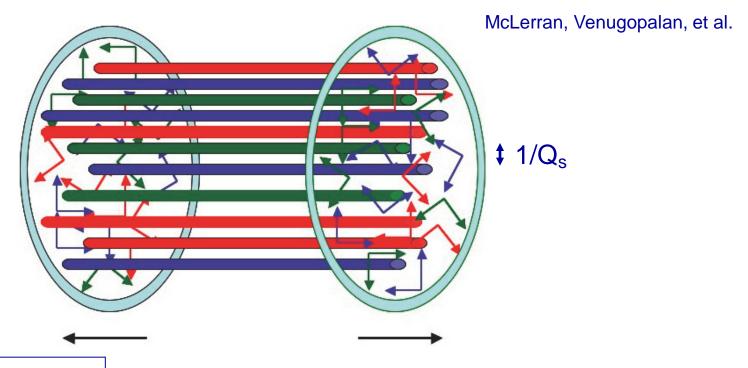
$$L(\vec{x}) = \frac{1}{3} \operatorname{tr} \operatorname{P} \exp \left(ig \int_0^{1/T} A_4(\tau, \vec{x}) d\tau \right)$$

Z₃ symmetry is *spontaneously broken* in QGP

Center symmetry is *spontaneously broken* in QGP

How is QGP created in HI collisions?

Glasma [glæzmə]: longitudinally slowly varying coherent color field



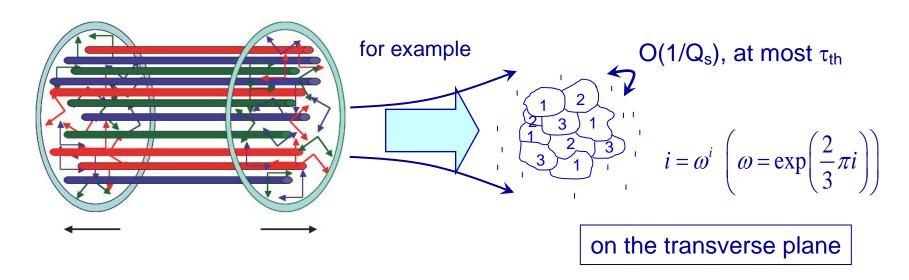
Standard scenario

Glasma decays into QGP

Typical transverse correlation scale: 1/Q_s @LHC Q_s ~ 2GeV

Z₃ valued QGP

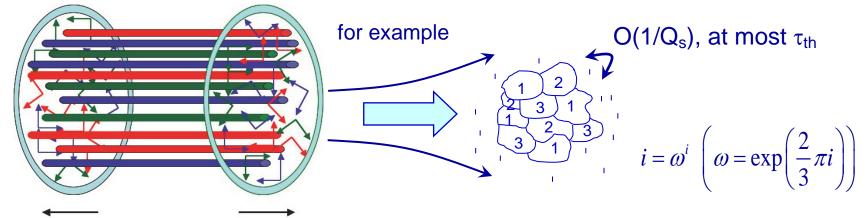
Since the typical transverse scale of the gauge field $\sim 1/Q_s$, in the transverse direction in the QGP created in HI collisions, changes from one of the three Z_3 elements to another with scale $1/Q_s$



Note: Z₃ is a discrete group

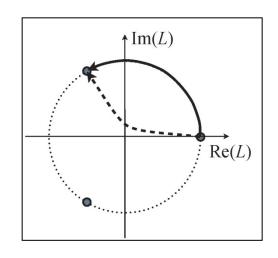
Effective potential for L has three distinct minima above T_c

L between Center Domains



on the transverse plane

How does L change between domains?



- > two of possibilities
 - modulus ~ const. (solid)
 - modulus changes (dashed)

Empirical form of Polyakov Loop potential

Commonly used form in PNJL model

$$U(\langle L \rangle) = -bT \left[54e^{-a/T} \left| \langle L \rangle \right|^2 + \log P(\langle L \rangle, \langle L^{\dagger} \rangle) \right]$$

$$P(z, \overline{z}) = 1 - 6 \left| z \right|^2 + 3 \left| z \right|^4 + 4(z^3 + \overline{z}^3)$$

$$a = 0.664 \text{ GeV}, b = 0.0075 \text{ GeV}^3$$

Form of *P* : from SU(3) Haar measure

- From lattice data
 - dominated by information around the potential minima
 - a lot of ambiguity

Note: Z_3 domain-like distribution of L is observed also in Full QCD calculation @ physical quark mass

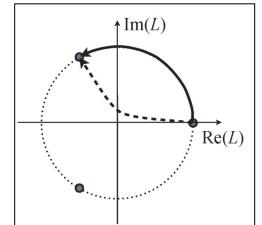
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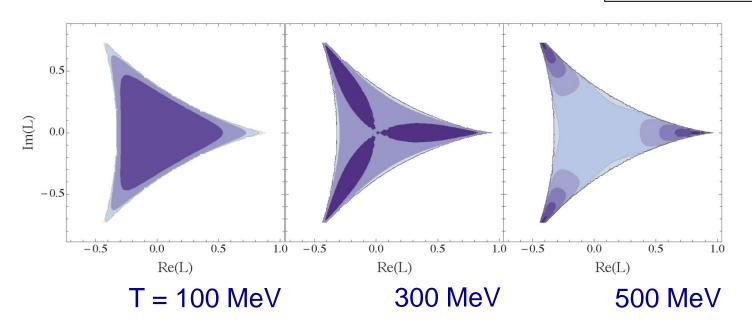
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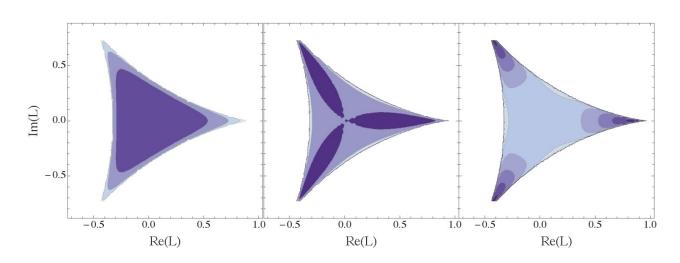
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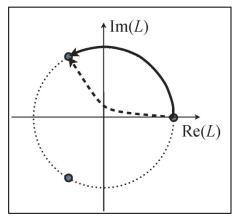
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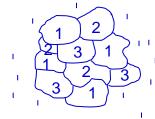
Gauge configuration on domain boundaries





Relation between heavy static quark free energy and L in the pure gauge theory

$$F_{Q}(\vec{x},T) = -T \log \left| \left\langle L(\vec{x},T) \right\rangle \right|$$





In between center domains, gauge field configuration: similar to those in the confined phase

Consequence I: shear viscosity

> Partons with thermal momenta cannot penetrate the walls but are reflected on them

Domain size R_d , and mean free path λ_f

$$\lambda_{\rm f} pprox rac{R_{
m d}}{2}$$
 and $\eta pprox rac{1}{3} n \overline{p} \lambda_{
m f}$

$$\frac{\eta}{s}$$

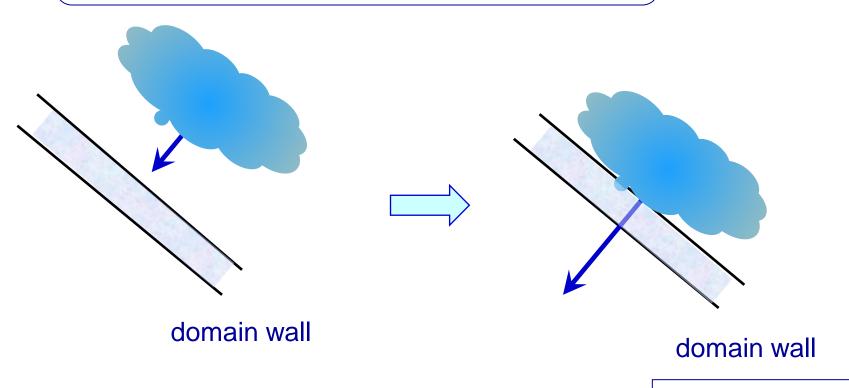
$$\boxed{\frac{\eta}{s} \approx \frac{1}{8} T R_d \quad \text{for} \quad p \approx 3T, \quad \frac{s}{n} \approx 4}$$

in the right ball park!

Consequence II: Frequency Collimation

> For partons with momenta higher than the thermal scale:

domain walls act like the combination of a frequency collimator and an irregular undulator



frequency collimation!

Energy loss by frequency collimation

Energy loss in a single wall crossing

$$\Delta E = \int_0^{\omega_c} \omega \frac{dN_g}{d\omega} d\omega \qquad \qquad \omega_c: \text{ critical frequency}$$

Weizsäcker-Williams approximation for gluon spectrum

$$\frac{dN_g}{d\omega} \approx \frac{C_2 \alpha_s}{\pi \omega} \log \left(\frac{\omega}{\omega_0}\right) \theta(\omega - \omega_0) \qquad \omega_0: \text{ infrared cutoff}$$

Energy loss per unit length

$$\frac{dE}{dx} \approx \frac{C_2 \alpha_s}{\pi R_d} \left\{ \omega_c \log \left(\frac{\omega_c}{\omega_0} \right) - (\omega_c - \omega_0) \right\}$$

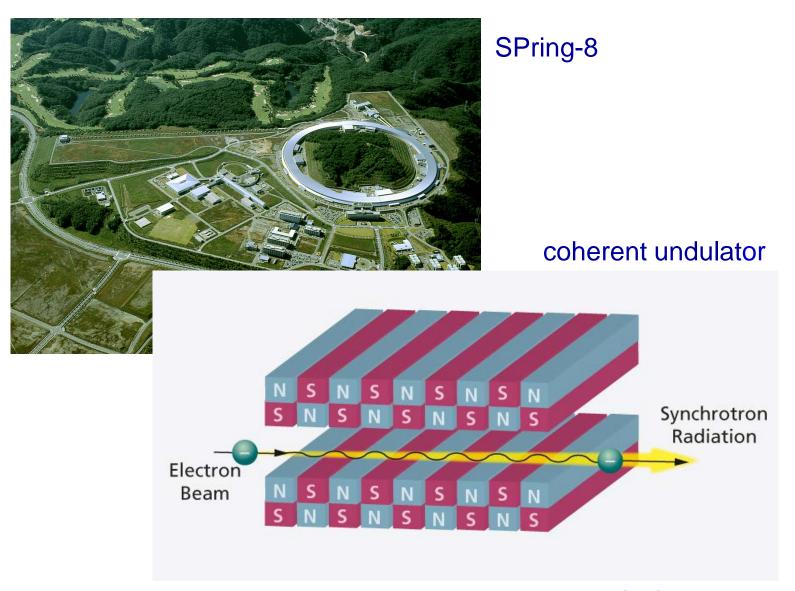
For $\omega_c \approx 1 - 2 \,\text{GeV}$, $\omega_0 \approx 0.4 \,\text{GeV}$, and $R_d \approx 0.5 \,\text{GeV}$,

$$\frac{dE}{dx} \approx (0.2 - 1) C_2 \alpha_s \text{ GeV/fm}$$

in the right ball park!

Consequence III: Irregular Undulation

Undulator @ photon factories



from SLS homepage

Consequence III: Irregular Undulation

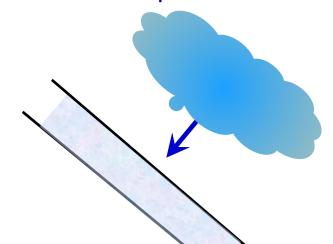
$$\left\langle L(\vec{x}) \right\rangle = \frac{1}{3} \left\langle \text{tr Pexp} \left[ig \int_{0}^{1/T} A_{0}(\tau, \vec{x}) d\tau \right] \right\rangle$$

$$\approx \exp \left[-\frac{g^{2}}{2T^{2}} \operatorname{tr} \left\langle A^{0}(\vec{x})^{2} \right\rangle \right]$$

✓ Inside domain walls (L=0), A⁰ fluctuates with large amplitude

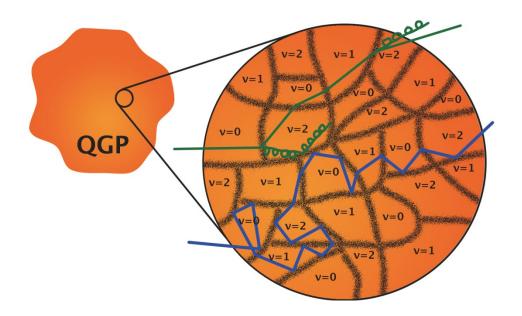


- Strong and uncorrelated radiation of gluons
- Prompt restoration of Weizsäcker-Williams field



Comments

- ✓ This energy loss mechanism works only in Z₃ valued QGP.
- ✓ This mechanism is nonperturbative.
- ✓ This mechanism distinguishes between the QCD phases
- ✓ This picture naturally leads to the immediate randomization of soft gluons



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- ✓ This mechanism is nonperturbative.
- This mechanism distinguishes between the QCD phases
- ✓ This picture naturally leads to the immediate randomization of soft gluons
- Full QCD Lattice calculation observes domain structure Early stage of QGP is dominated by gluons: close to pure gauge So, at least in this context, not much sense to fuss over, e.g., perturbative correction to the pure gauge Polyakov loop potential due to finite quark mass