

QCD phase diagram in the Dyson-Schwinger approach

Christian S. Fischer

Justus Liebig Universität Gießen

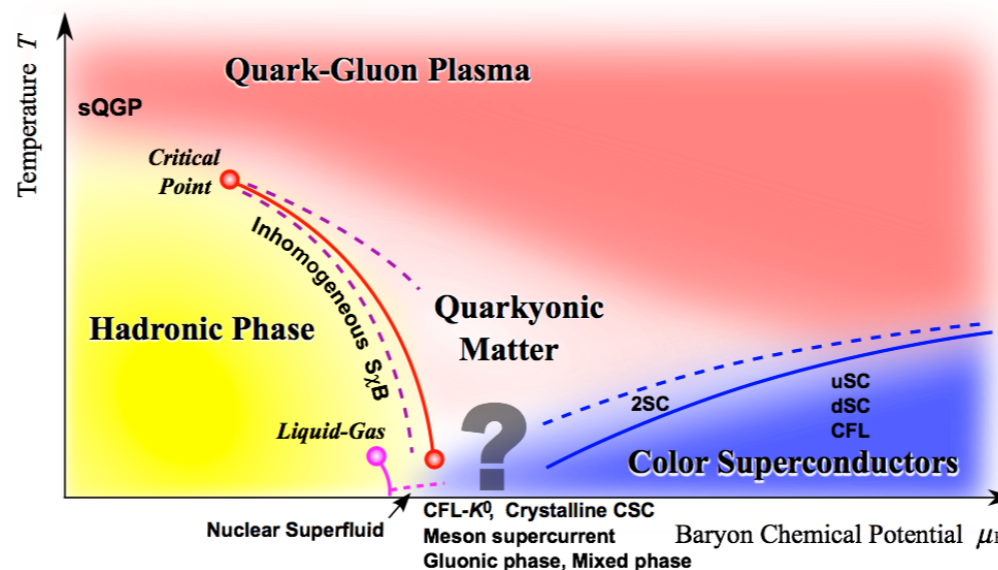
HIC | **FAIR**
for
Helmholtz International Center

5th of December 2013

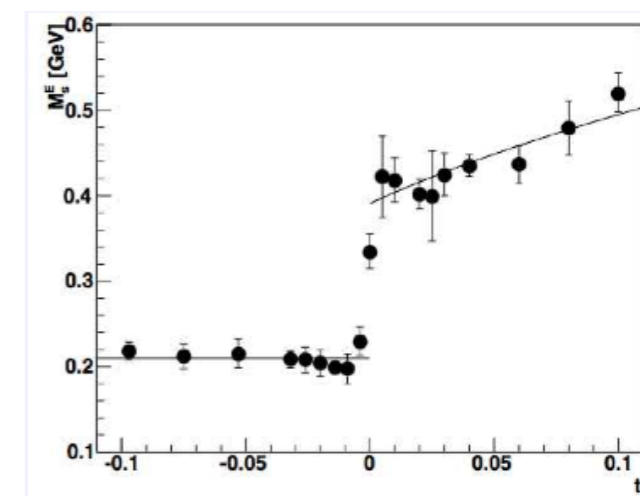
 **HELMHOLTZ**
| **GEMEINSCHAFT**

with Jens Mueller, Jan Luecker
Axel Maas, Jan Pawlowski, Leonard Fister

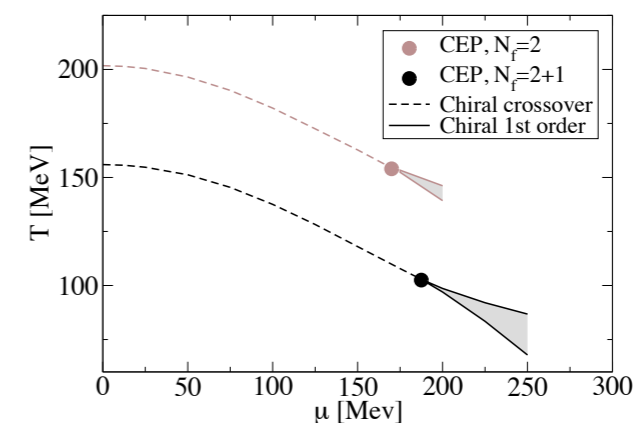
1. Introduction



2. Gluons at zero and finite temperature

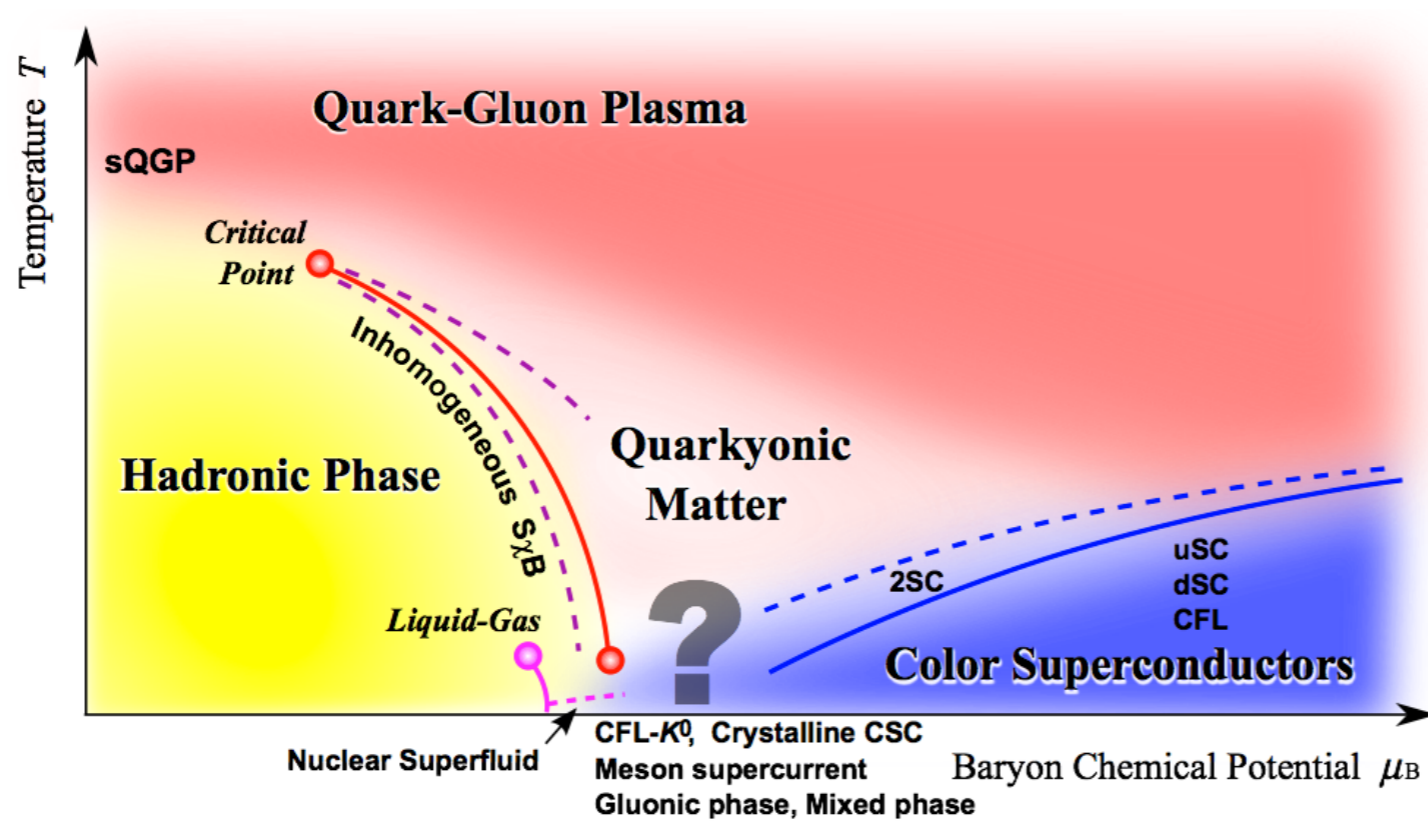


3. Quarks and the QCD phase diagram



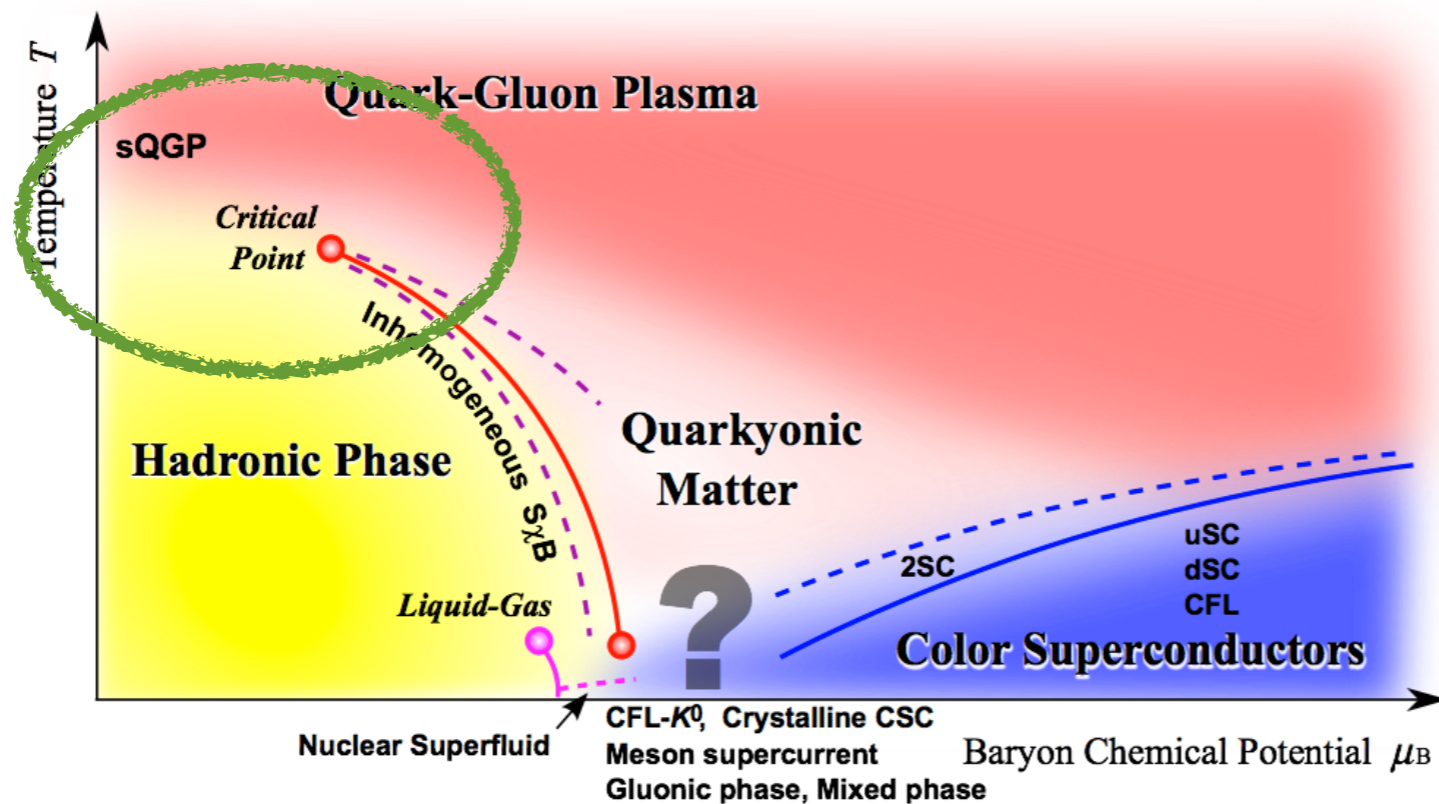
QCD phase transitions I

Fukushima, Hatsuda, Rept. Prog. Phys. 74 (2011) 014001



QCD phase transitions I

Fukushima, Hatsuda, Rept. Prog. Phys. 74 (2011) 014001



- Chiral limit ($M_{\text{weak}} \rightarrow 0$): order parameter chiral condensate

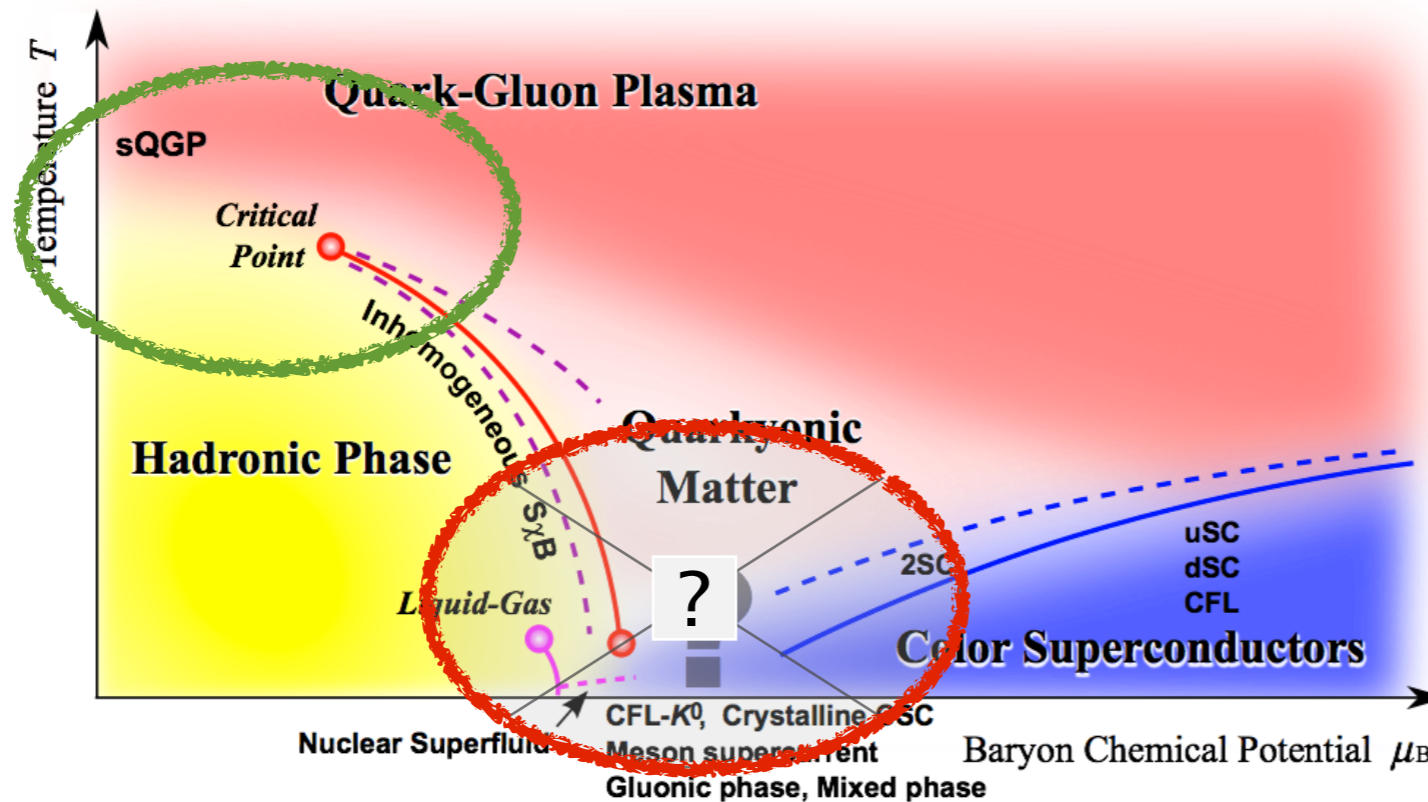
$$\langle \bar{\Psi} \Psi \rangle \sim \text{Tr} \int S$$

- Static quarks ($M_{\text{weak}} \rightarrow \infty$): order parameter Polyakov-loop

$$\langle |L| \rangle \sim e^{-F_q/T}$$

QCD phase transitions I

Fukushima, Hatsuda, Rept. Prog. Phys. 74 (2011) 014001



DSEs at low T, large μ :

Müller, Buballa, Wambach, EPJA 49 (2013),
PLB 727 (2013) 240

see Talk of M. Buballa, prev. week

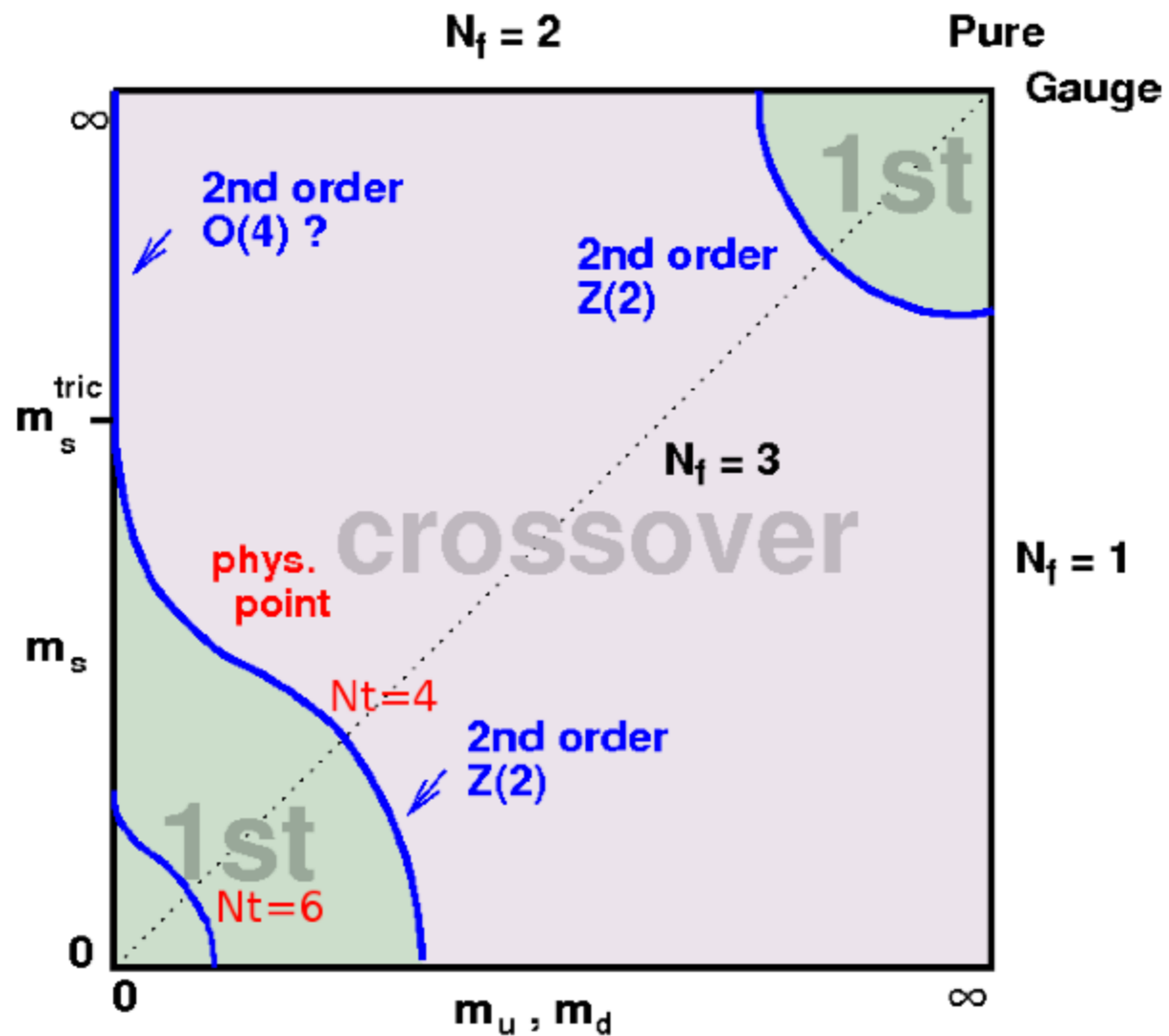
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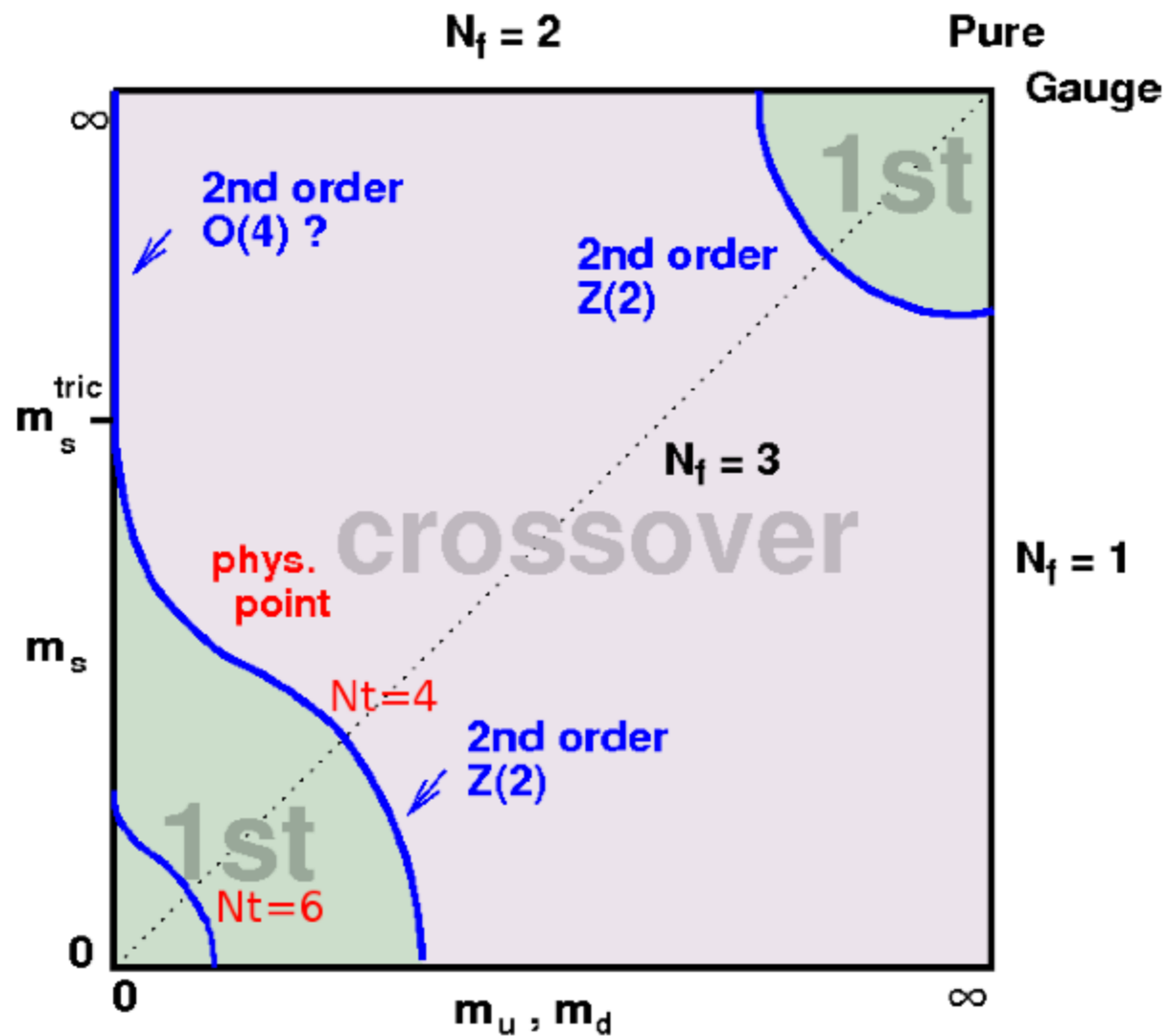
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QCD phase transitions II

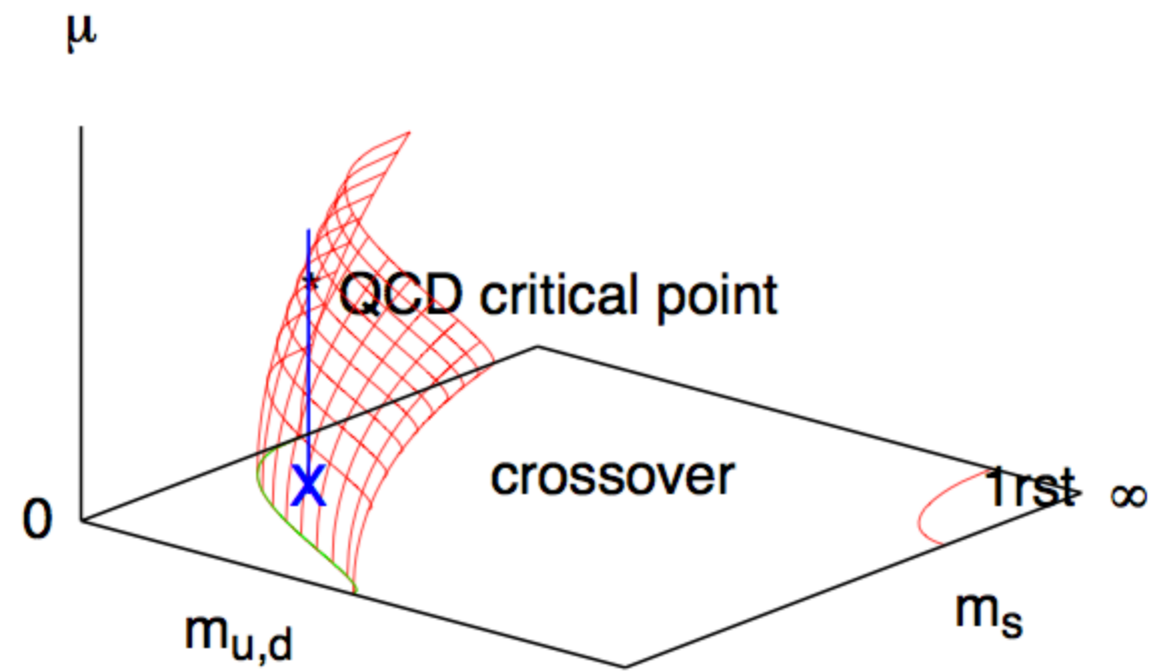


Plot: O. Philipsen

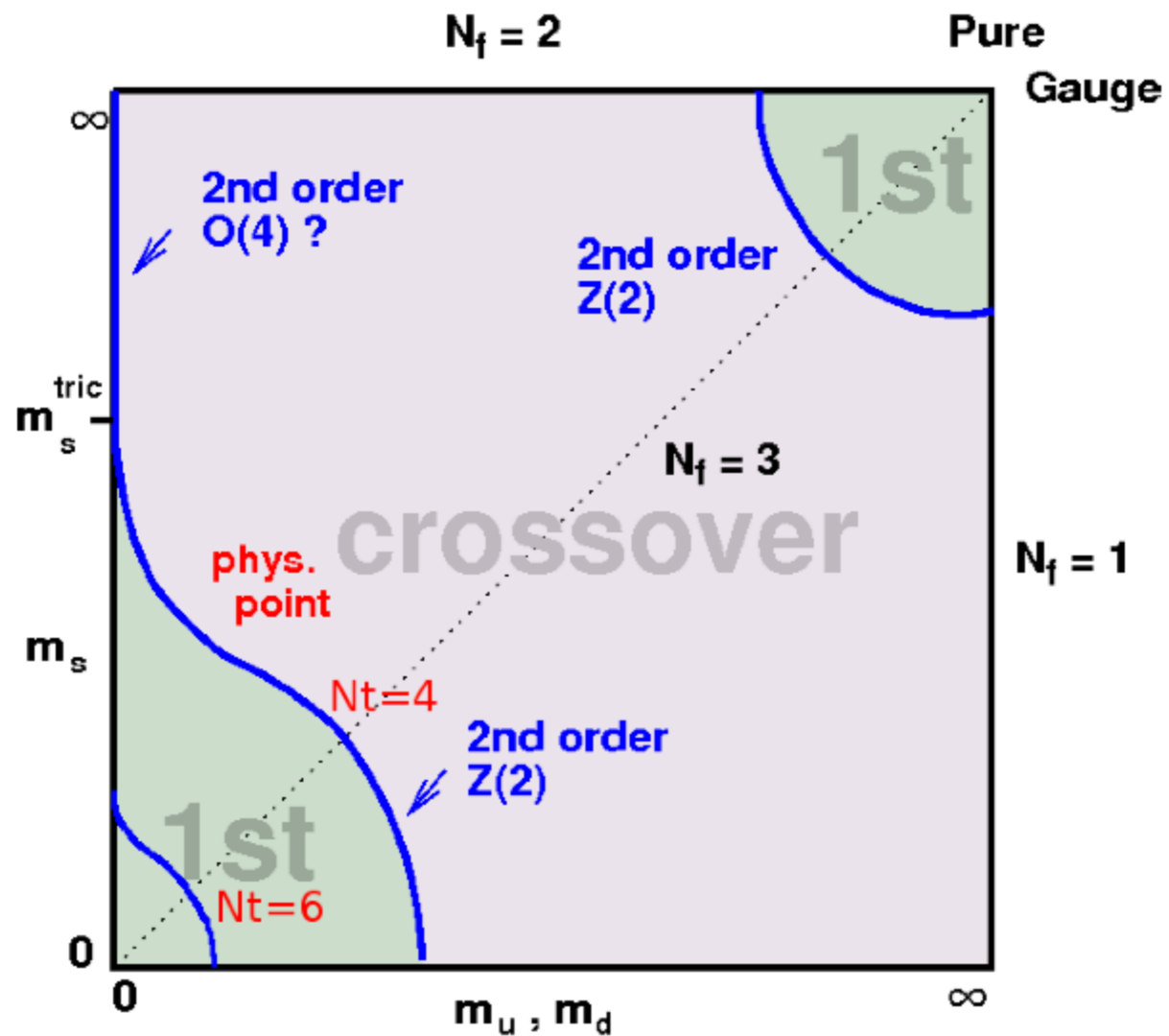
QCD phase transitions II



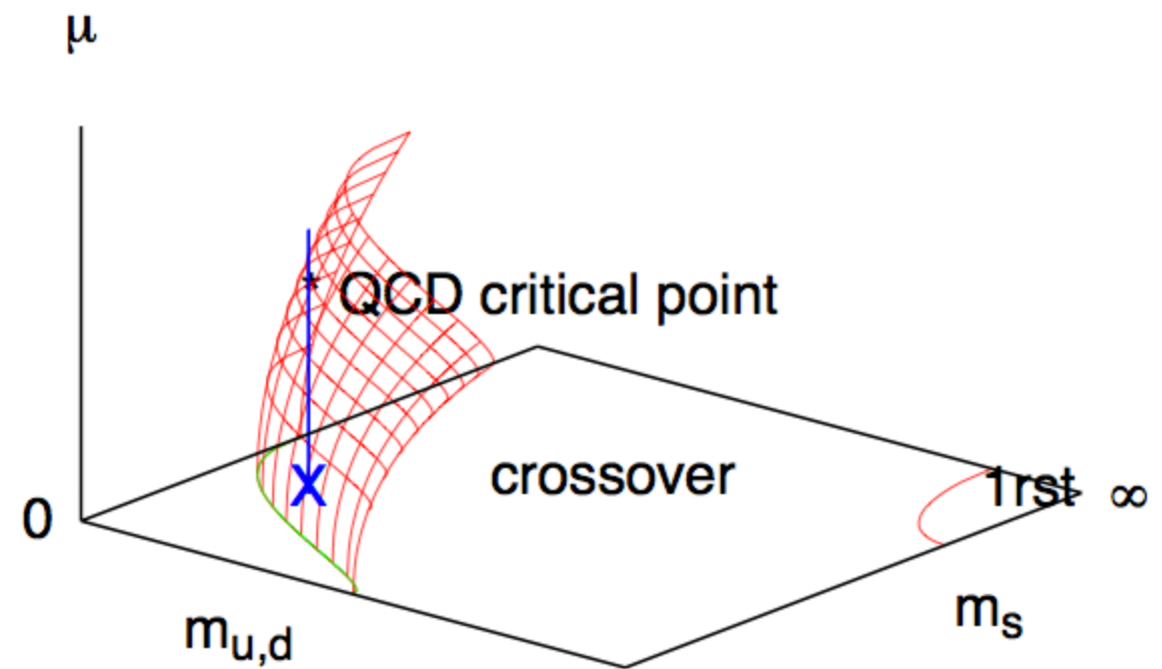
Plot: O. Philipsen



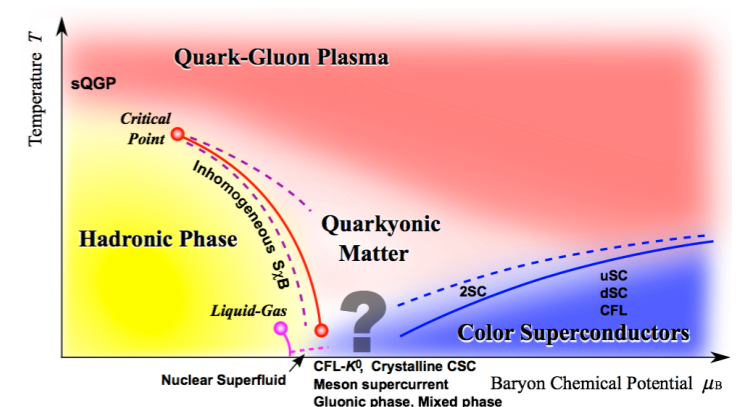
QCD phase transitions II



Plot: O. Philipsen



Is this happening ??



Nonperturbative QCD: Complementary approach

Quarks and gluons

- Lattice simulations of QCD
 - Ab initio
 - Gauge invariant
- Functional approaches to QCD (DSE, FRG, Hamilton):
 - Chiral symmetry: **physical quark masses**
 - Analytical solutions in IR
 - Infinite volume and continuum limit
 - Multi-scale problems feasible (e.g. $(g-2)_\mu$)
Goecke, CF, Williams, PRD 87 (2013) 03401
 - **Chemical potential: no sign problem**

Hadrons


- Effective theories and models (chiral \mathcal{L} , PQM, PNJL)
 - Physical degrees of freedom

→ see also Talk by J. Pawłowski

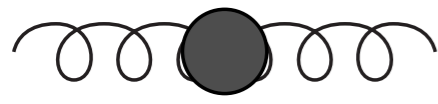
QCD in covariant gauge

Imaginary time formulation:

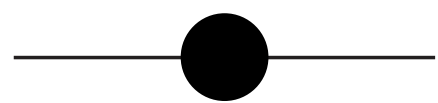
$$Z_{QCD} = \int \mathcal{D}[\Psi, A] \exp \left\{ - \int_0^{1/T} dt \int d^3x \left(\overline{\Psi} (i\not{D} + \gamma_4 \mu - m) \Psi - \frac{1}{4} (F_{\mu\nu}^a)^2 + \text{gauge fixing} \right) \right\}$$

M_{weak} 

Landau gauge propagators in momentum space, $p = (\vec{p}, \omega_p)$:



$D_{\mu\nu}^{\text{Gluon}}(p) = \frac{Z_T(p)}{p^2} P_{\mu\nu}^T(p) + \frac{Z_L(p)}{p^2} P_{\mu\nu}^L(p)$



$S^{\text{Quark}}(p) = Z_f(p) [-i \vec{\gamma} \vec{p} - i \gamma_4 \tilde{\omega}_n Z_c(p) + M(p)]^{-1}$

The Goal: gauge invariant information in a gauge fixed approach.

QCD order parameters from propagators

$$\text{---} \bullet \text{---}^{-1} = \text{---}^{-1} - \text{---} \bullet \text{---} \text{---} \text{---}$$

Chiral order parameter:

$$\langle \bar{\Psi} \Psi \rangle = Z_2 N_c \text{Tr}_D \frac{1}{T} \sum_{\omega} \int \frac{d^3 p}{(2\pi)^3} S(\vec{p}, \omega)$$

Deconfinement:

- dressed Polyakov loop

$$\Sigma = - \int_0^{2\pi} \frac{d\varphi}{2\pi} e^{-i\varphi} \langle \bar{\Psi} \Psi \rangle_{\varphi}$$

Synatschke, Wipf, Wozar, PRD 75, 114003 (2007)
 Bilgici, Bruckmann, Gattringer, Hagen, PRD 77 094007 (2008)
 CF, PRL 103 052003 (2009)

- Polyakov loop potential

$$L = \frac{1}{N_c} \text{Tr} e^{ig\beta A_0}$$

$$\frac{\delta(\Gamma - S)}{\delta A_0} = \frac{1}{2} \left(\text{---} \bullet \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} \text{---} \right) - \left(\text{---} \text{---} \text{---} \text{---} \right) - \frac{1}{6} \left(\text{---} \bullet \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} \right)$$

Braun, Gies, Pawłowski, PLB 684, 262 (2010)
 Braun, Haas, Marhauser, Pawłowski, PRL 106 (2011)
 Fister, Pawłowski, PRD 88 045010 (2013)
 CF, Fister, Luecker, Pawłowski, arXiv:1306.6022

The DSE for the quark propagator

$$\text{---}\bullet\text{---}^{-1} = \text{---}^{-1} - \text{---}\bullet\text{---}^{-1}$$

$$[S(p)]^{-1} = [-i\not{p} + M(p^2)]/Z_f(p^2)$$

Input:

- dressed Gluon propagator
- dressed Quark-Gluon-Vertex

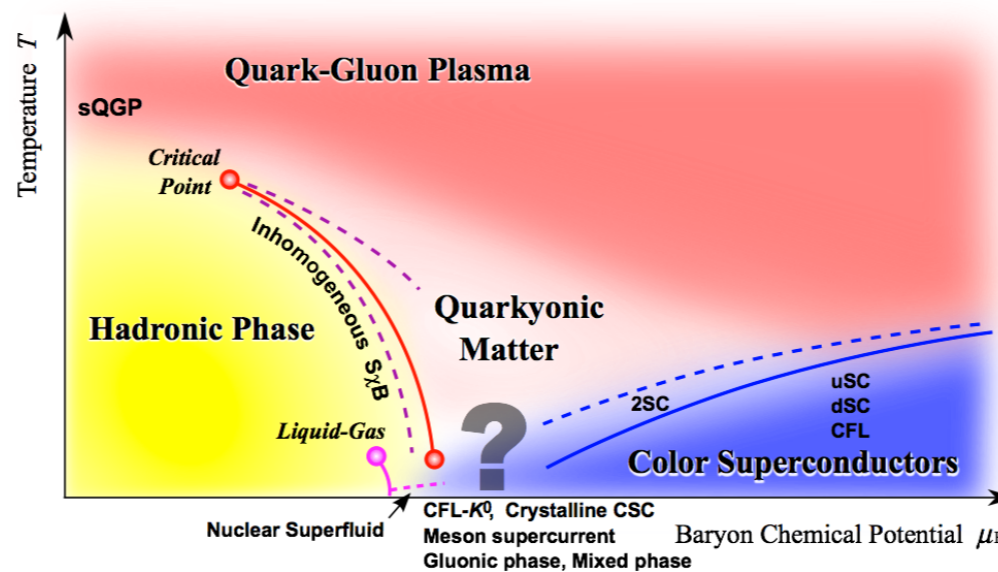
Two strategies: I. use **model** for gluon and vertex

Qin, Chang, Chen, Liu and Roberts, PRL 106 (2011) 172301
Gutierrez, Ahmad, Ayala, Bashir and Raya, arXiv:1304.8065.

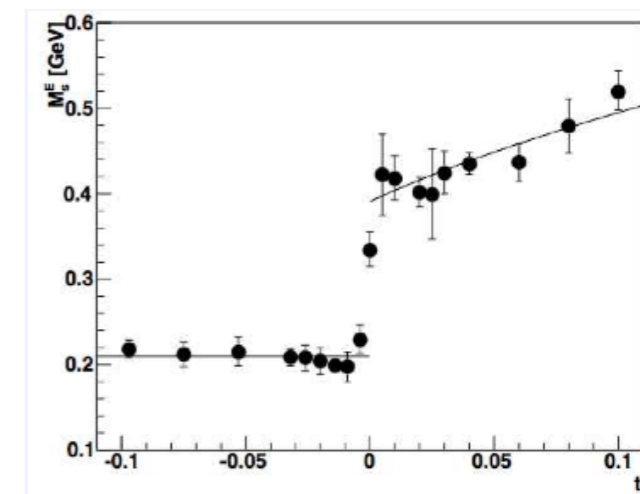
- ok for first insights
- not good enough for systematic study

II. **determine gluon and vertex** explicitly

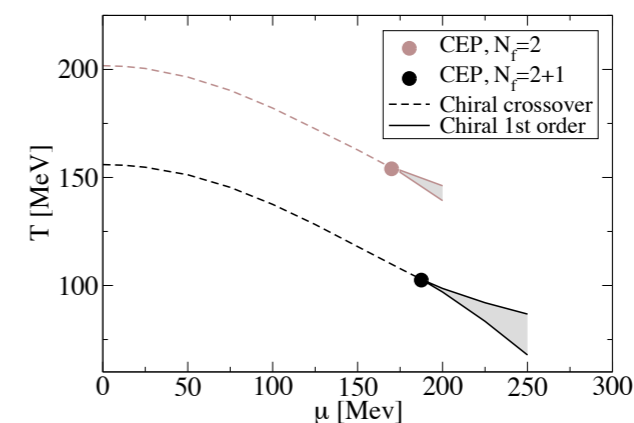
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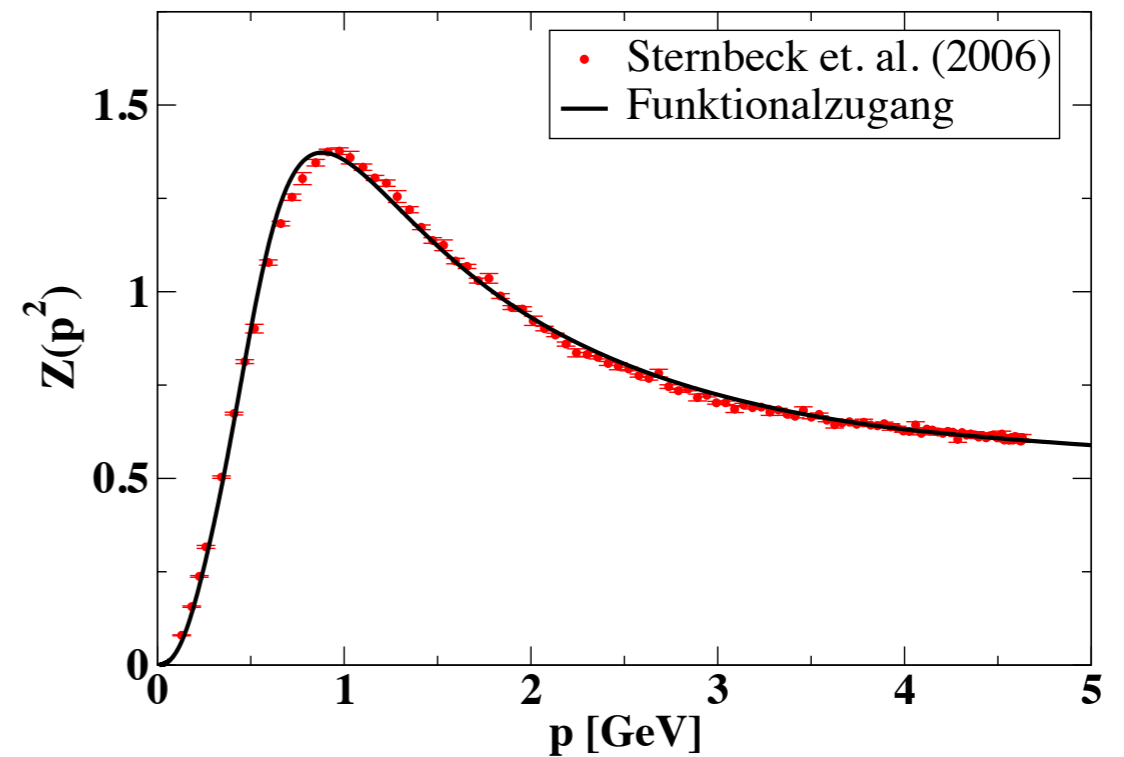
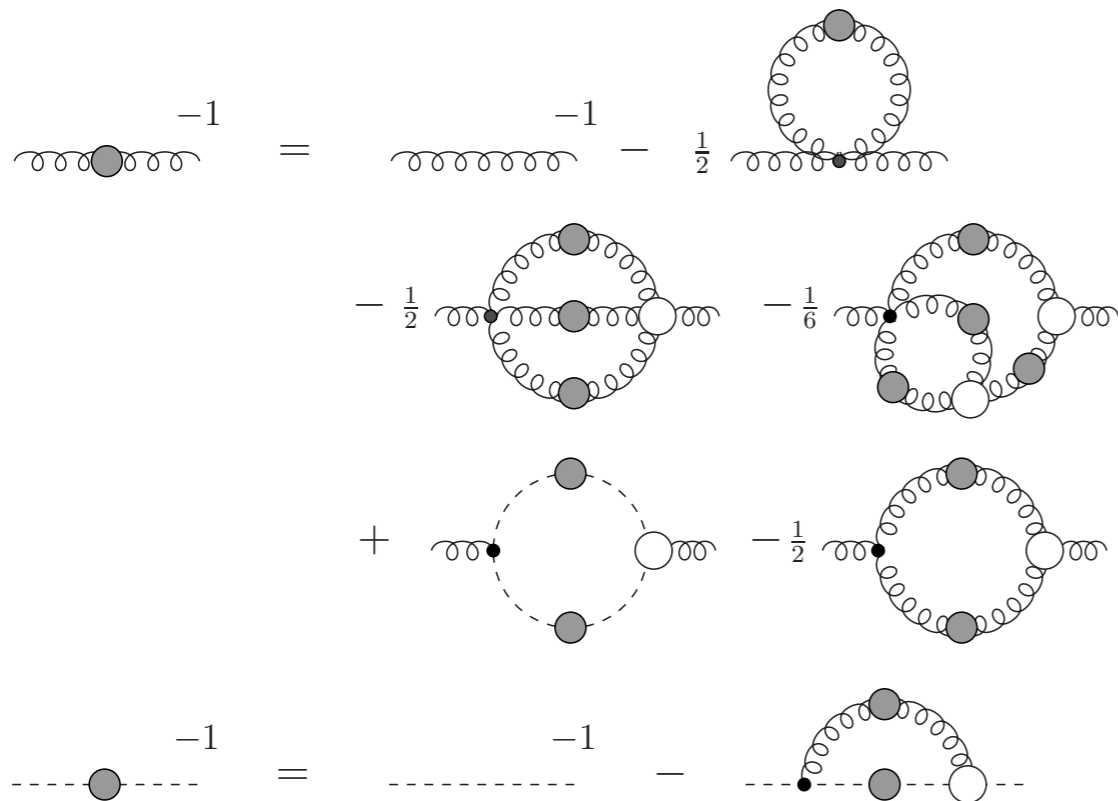
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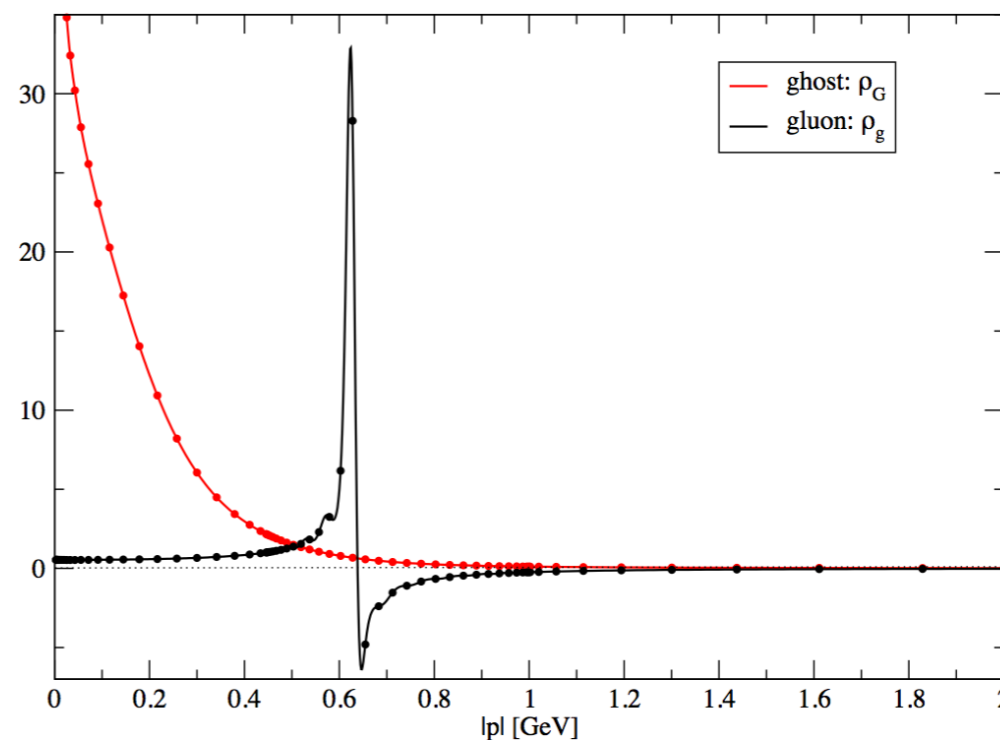


Strategy I: Landau gauge gluon propagator



CF, Maas, Pawłowski, *Annals Phys.* 324 (2009) 2408.
 Huber and von Smekal, *JHEP* 1304 (2013) 149

- spacelike momenta: excellent agreement with lattice
 - spectral function: positivity violations
- Gluon cannot appear in detector!**

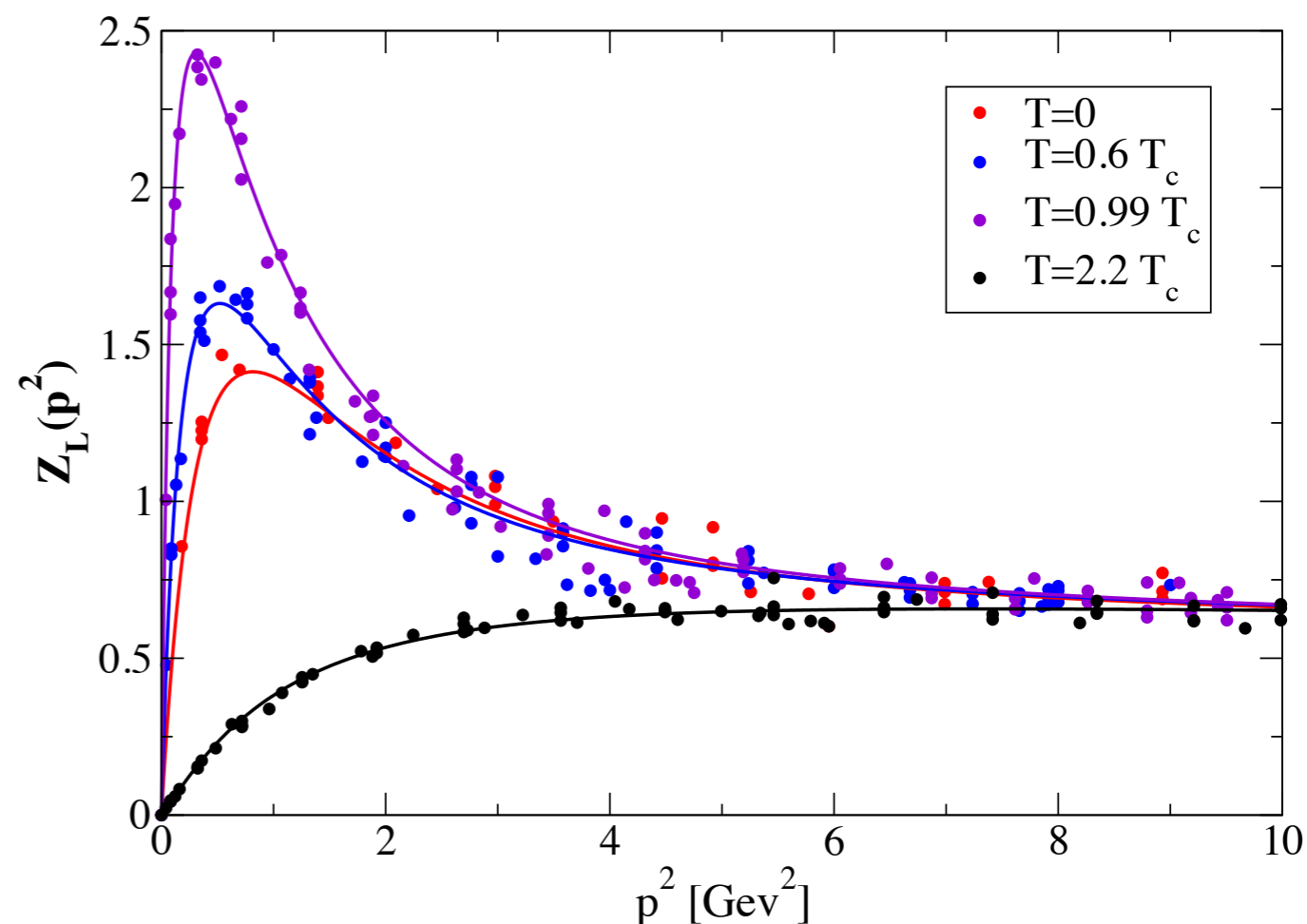
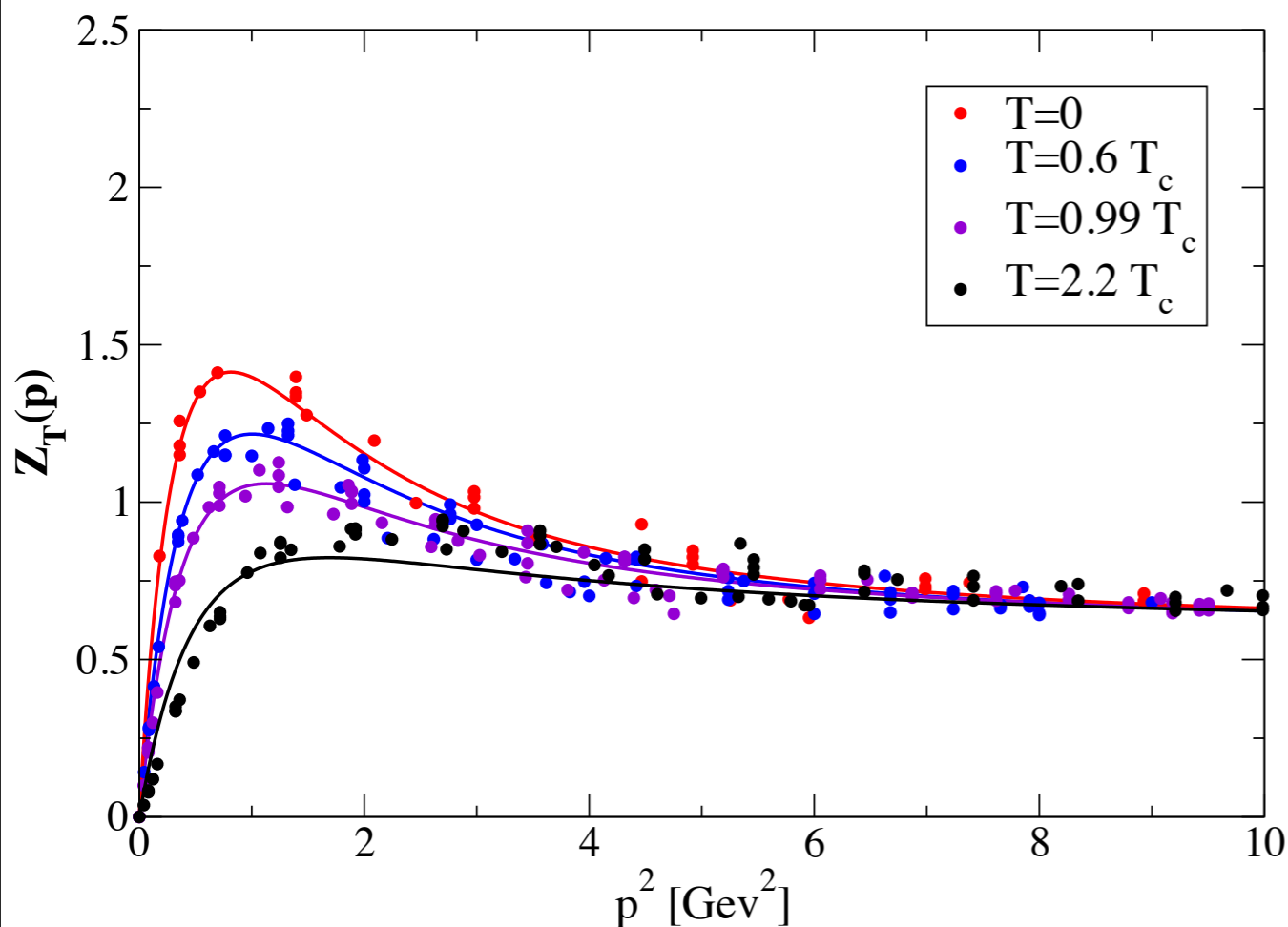


spectral function

Strauss, CF, Kellermann, *Phys. Rev. Lett.* 109, (2012) 252001

Glue at finite temperature ($T \neq 0$)

T-dependent gluon propagator from quenched lattice simulations:



- Crucial difference between magnetic and electric gluon
- Maximum of electric gluon near T_c

Cucchieri, Maas, Mendes, PRD 75 (2007)

CF, Maas, Mueller, EPJC 68 (2010)

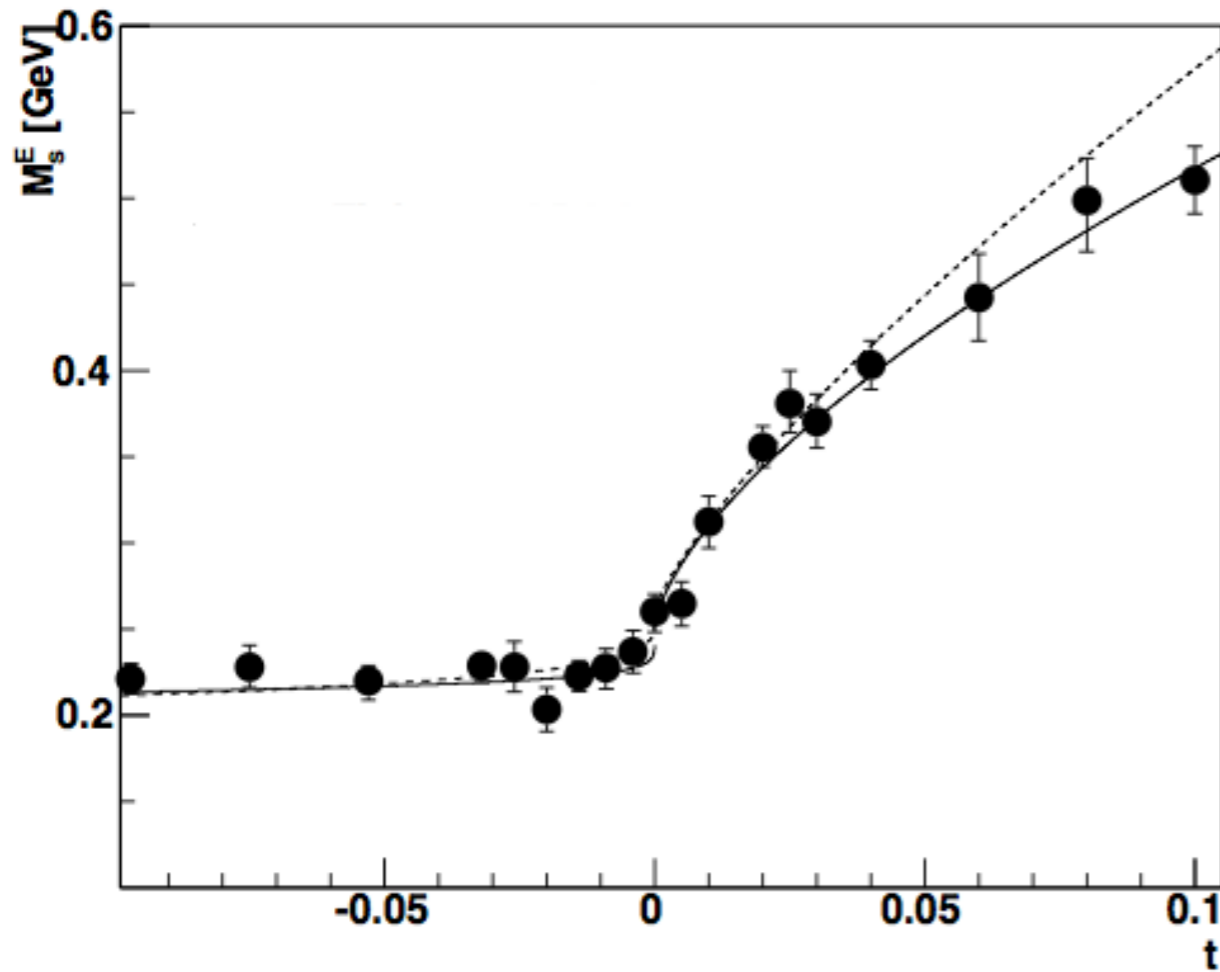
Cucchieri, Mendes, PoS FACESQCD 007 (2010)

Aouane, Bornyakov, Ilgenfritz, Mitryushkin, Muller-Preussker and Sternbeck, PRD 85 (2012) 034501

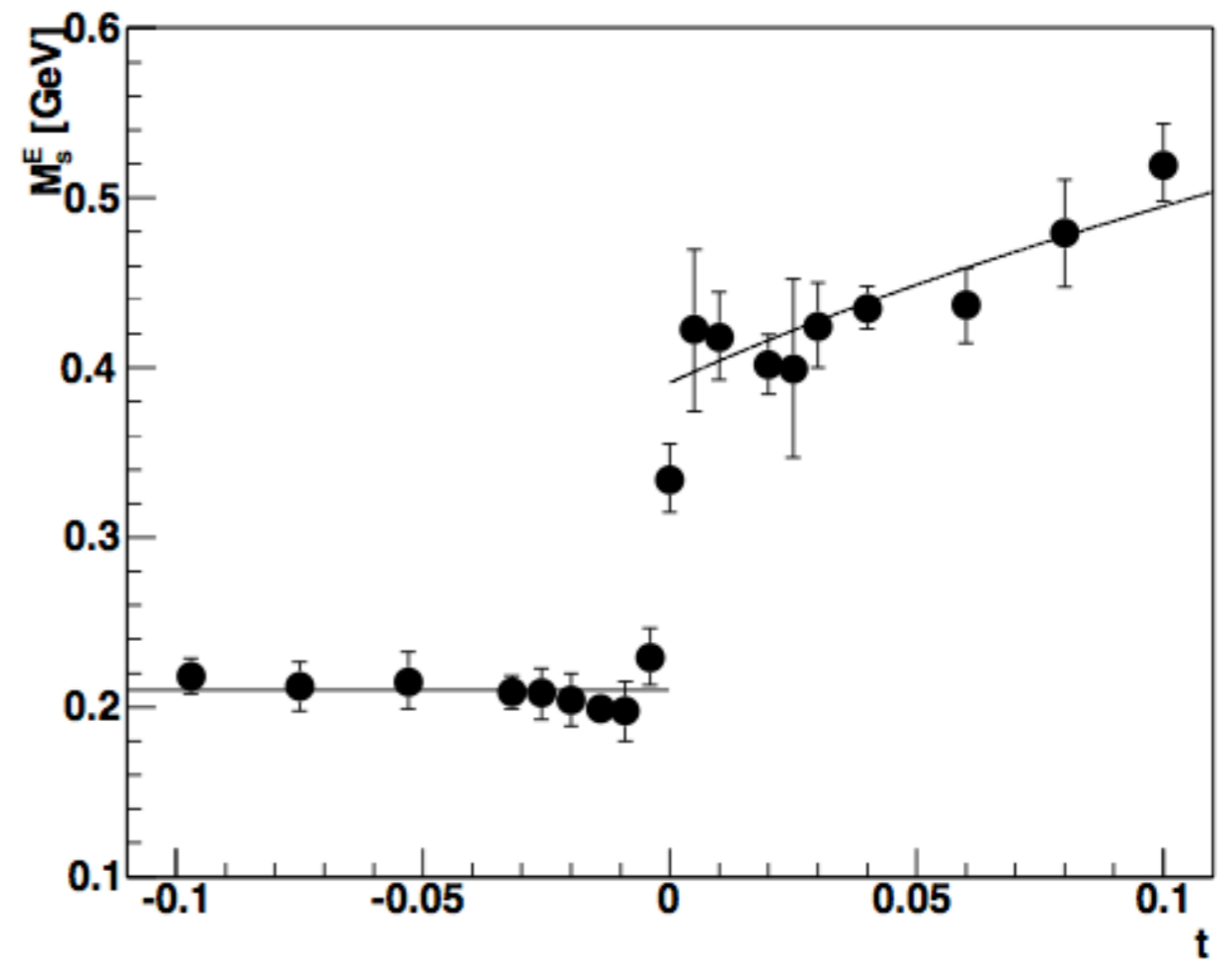
FRG: Fister, Pawlowski, arXiv:1112.5440

Glueon electric screening mass: SU(2) vs. SU(3)

SU(2)



SU(3)

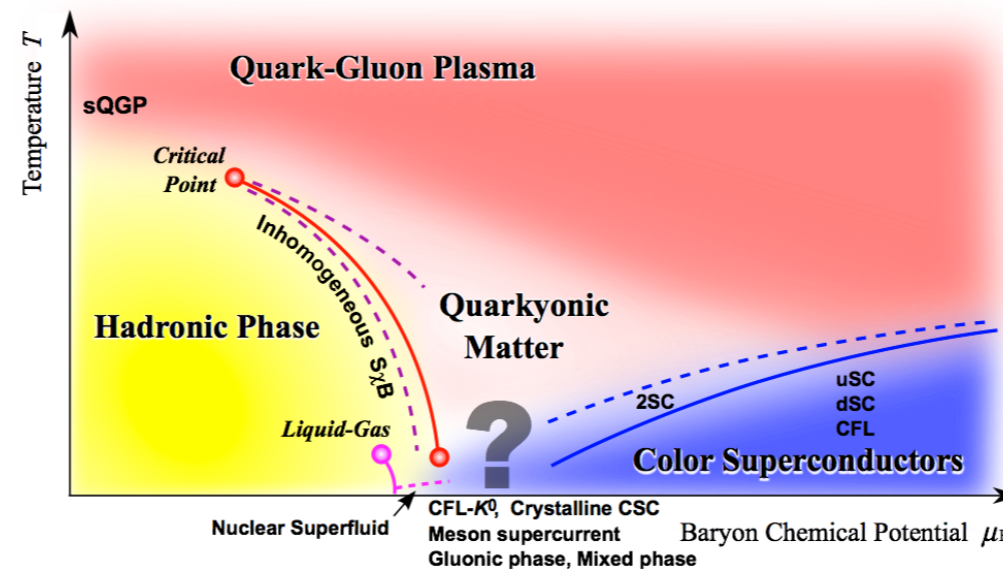


Maas, Pawłowski, Smekal, Spielmann, PRD 85 (2012) 034037
CF, Maas, Mueller, EPJC 68 (2010)

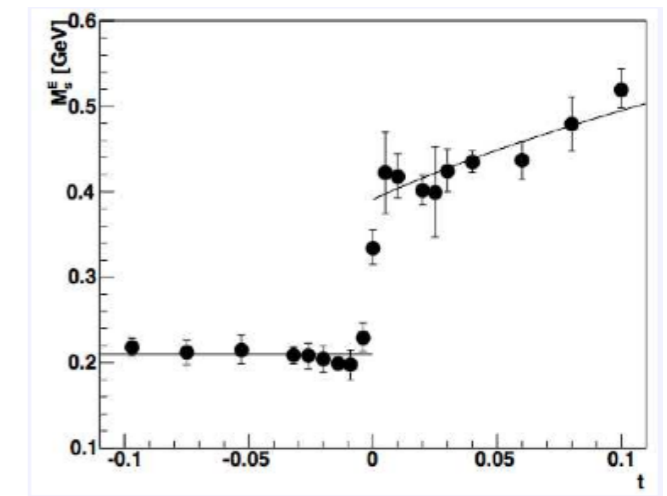
$$t = (T - T_c) / T_c$$

- phase transition of **second** and **first** order visible in electric screening mass

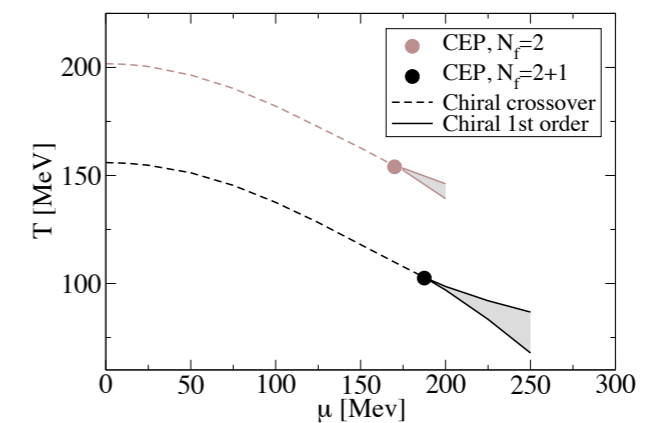
1. Introduction



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DSEs of QCD

$$\begin{aligned}
 \text{gluon propagator}^{-1} &= \text{gluon propagator}^{-1} - \frac{1}{2} \text{ghost loop} \\
 &\quad - \frac{1}{2} \text{quark loop} - \frac{1}{6} \text{quark loop with ghost} \\
 &\quad + \text{quark loop with ghost} - \frac{1}{2} \text{quark loop with ghost}
 \end{aligned}$$

quenched, T-dependent
lattice propagator

$$\text{quark-gluon vertex}^{-1} = \text{quark-gluon vertex}^{-1} - \text{quark-gluon vertex}^{-1}$$

quark gluon vertex

$$\text{quark propagator}^{-1} = \text{quark propagator}^{-1} - \text{quark self-energy}$$

- under study at T=0
Skullerud, Kizilersu, JHEP 0209 (2002) 013
Alkofer, CF, Llanes-Estrada, Schwenzer, Annals Phys. 324 (2009) 106
CF, Williams PRL 103 (2009) 122001
- T ≠ 0: ansatz,
T, m, μ dependent

Approximation for Quark-Gluon interaction

- T,μ,m-dependent vertex:

Abelian WTI

$$\Gamma_\nu(q, k, p) = \tilde{Z}_3 \left(\delta_{4\nu} \gamma_4 \frac{C(k) + C(p)}{2} + \delta_{j\nu} \gamma_j \frac{A(k) + A(p)}{2} \right) \times$$

$$\times \left(\frac{d_1}{d_2 + q^2} + \frac{q^2}{\Lambda^2 + q^2} \left(\frac{\beta_0 \alpha(\mu) \ln[q^2 / \Lambda^2 + 1]}{4\pi} \right)^{2\delta} \right)$$

perturbation theory

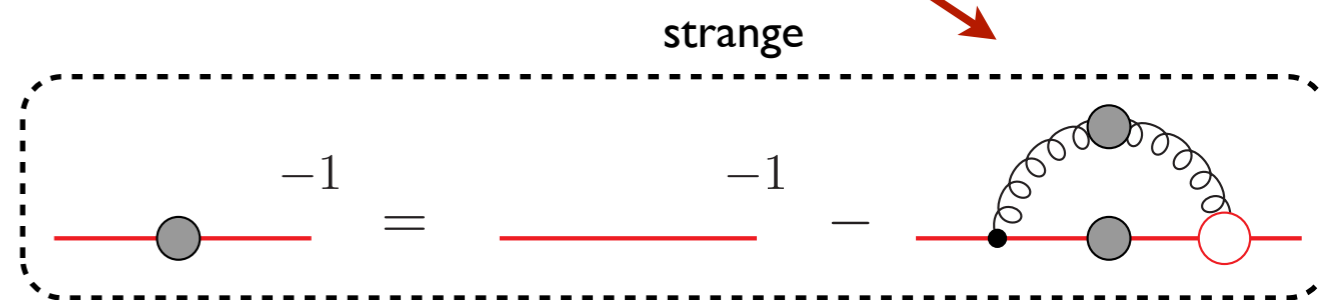
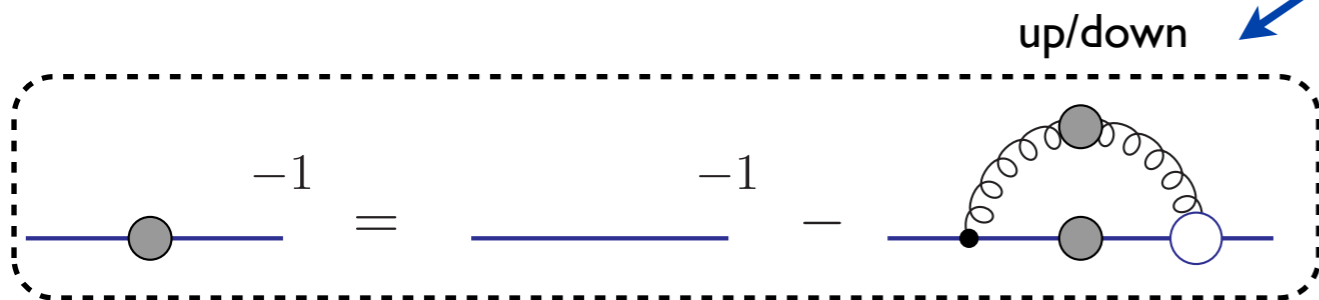
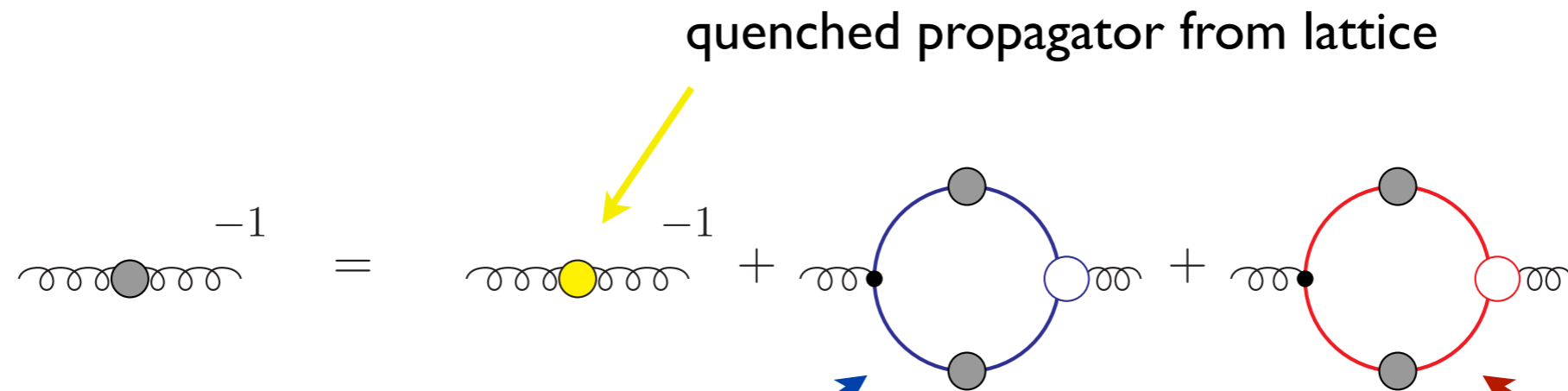
Infrared ansatz:

- d2 fixed to match gluon input
- d1 fixed via quark condensate (see later)
- **correct UV and IR-behavior**

- **crosscheck:** $f_\pi(T = 0) = 88 \text{ MeV}$

Alkofer, CF, Llanes-Estrada, Schwenzer, Annals Phys. 324 (2009)
CF, Pawłowski, PRD 80 (2009) 025023

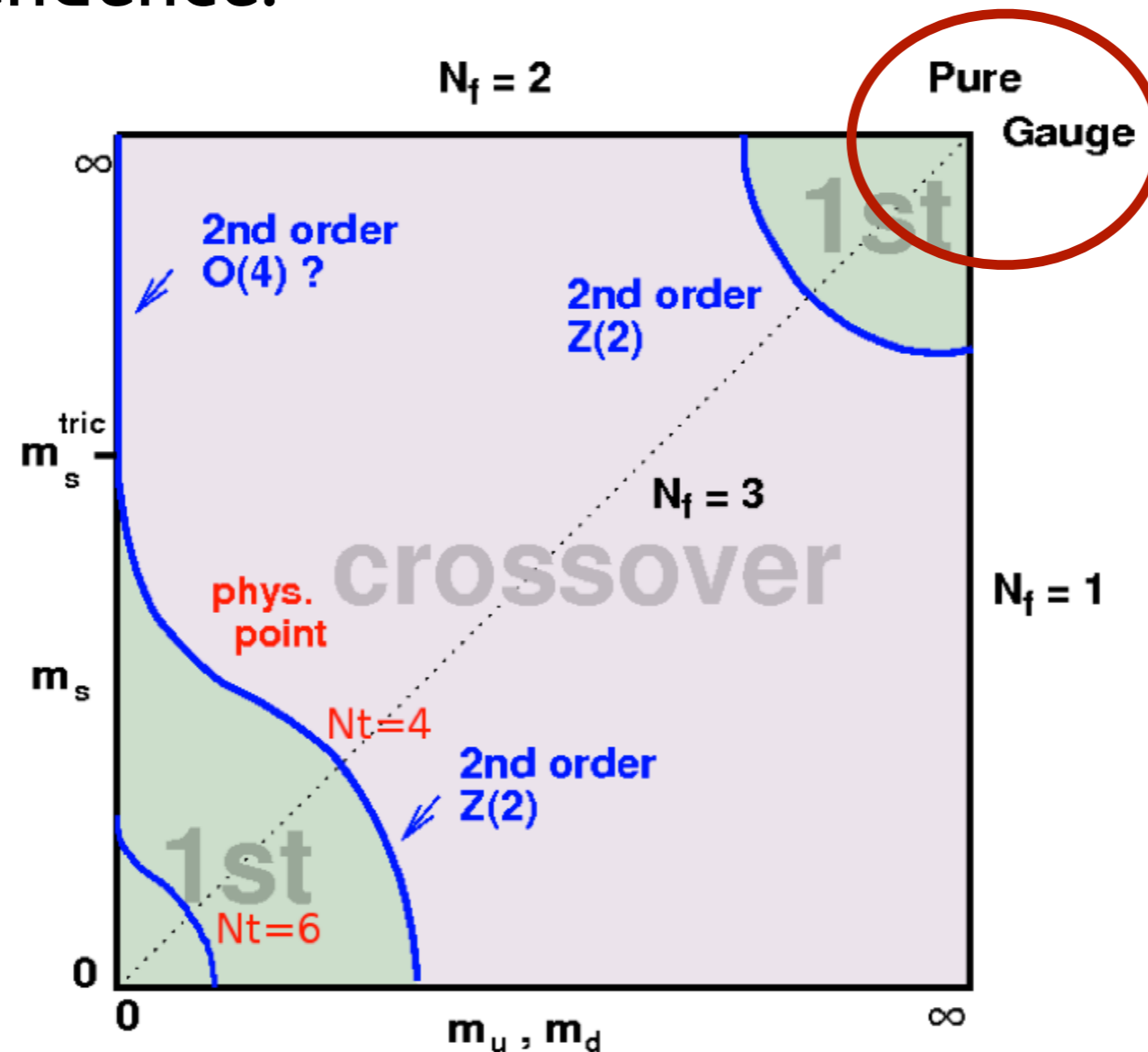
$N_f=2+1$ -QCD with DSEs



- quenched: without quark-loop
- $N_f=2$: isospin symmetry
- $N_f=2+1$: solve coupled system of $2+3+3$ equations

QCD phase transition: heavy quark limit/quenched

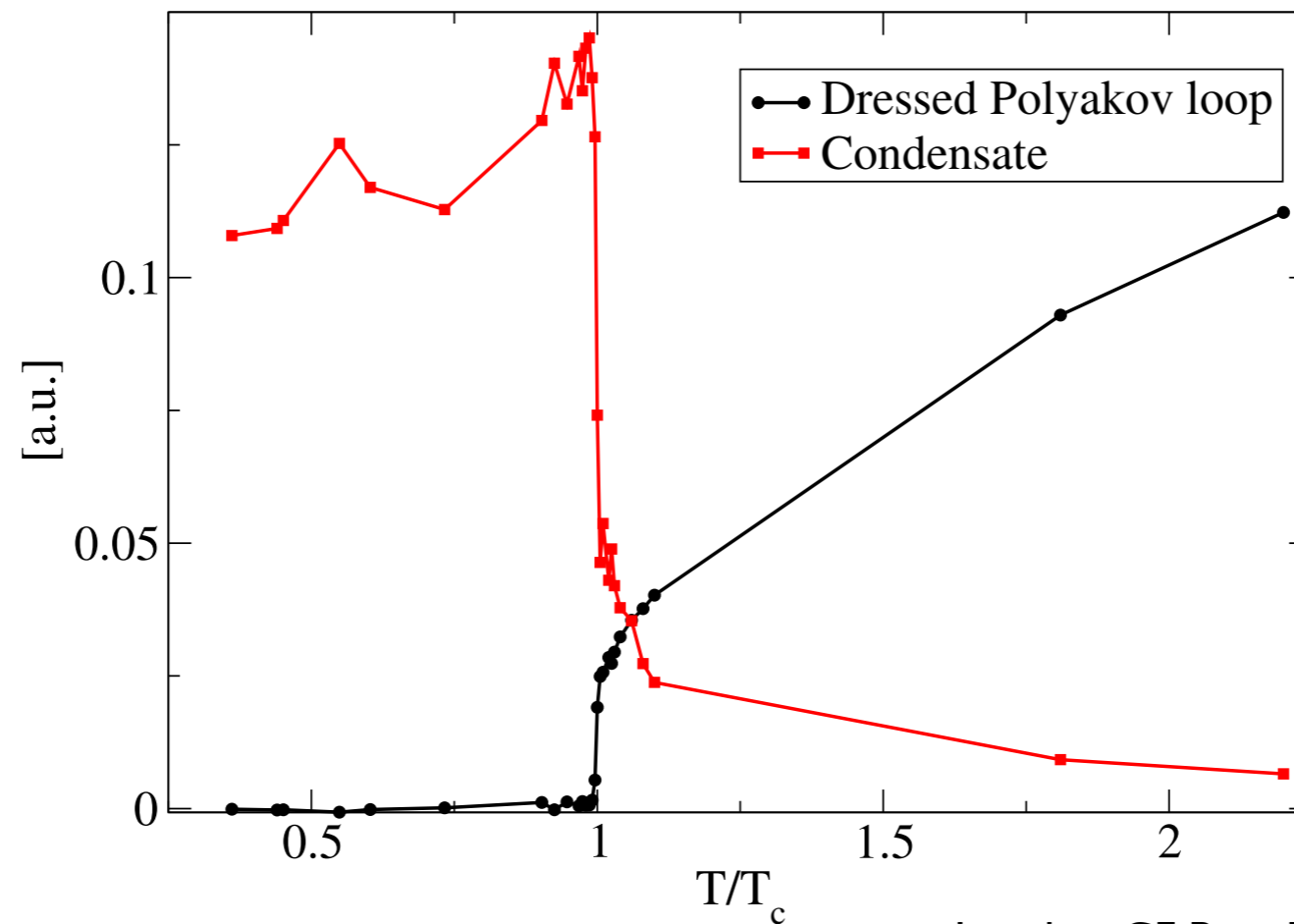
Quark mass dependence:



- Expect: Transitions controlled by deconfinement
- $SU(2)$ second order, $SU(3)$ first order

Transition temperatures, quenched

quenched DSE: SU(3)



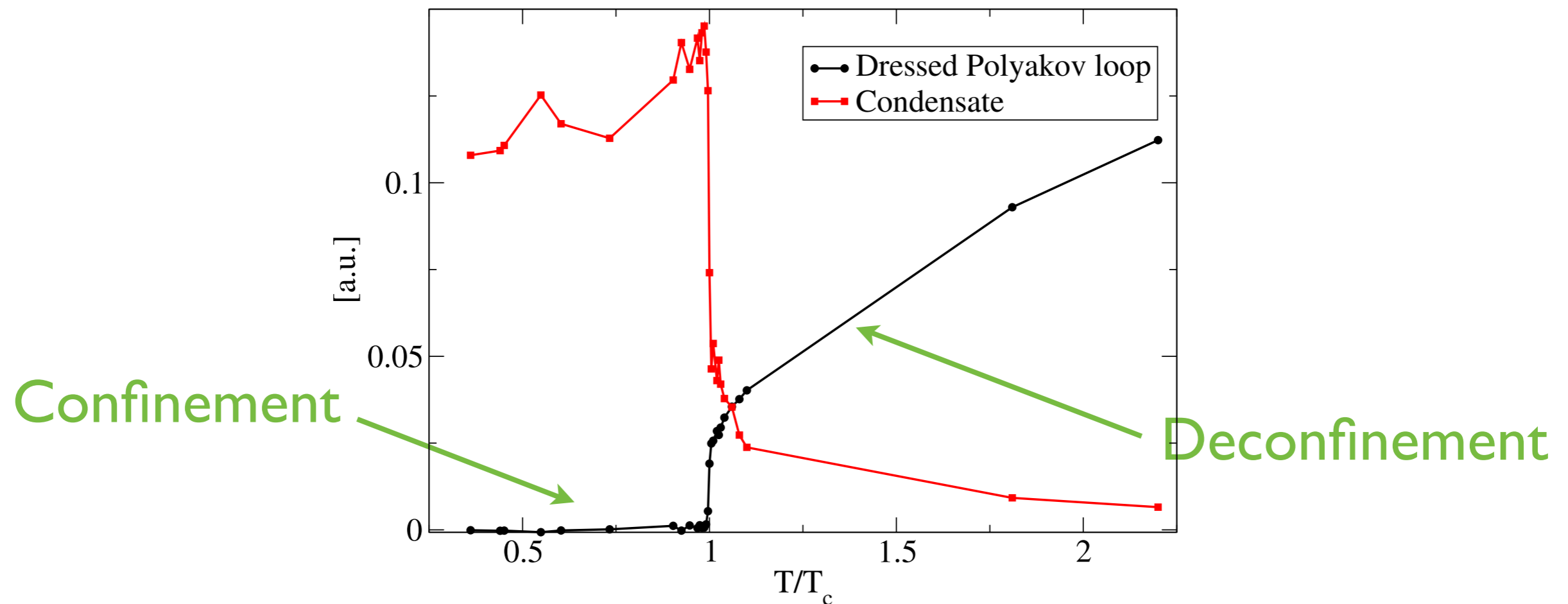
Luecker, CF, Prog.Part.Nucl.Phys. 67 (2012) 200-205
CF, Maas, Mueller EPJC 68 (2010)

- SU(2): $T_c \approx 305$ MeV
- SU(3): $T_c \approx 270$ MeV
- $T \leq T_c$: increasing condensate due to electric part of gluon

cf. Buividovich, Lushevskaya, Polikarpov, PRD 78 (2008) 074505
cf. Braun, Gies, Pawłowski, PLB 684 (2010) 262.

Transition temperatures, quenched

quenched DSE: SU(3)



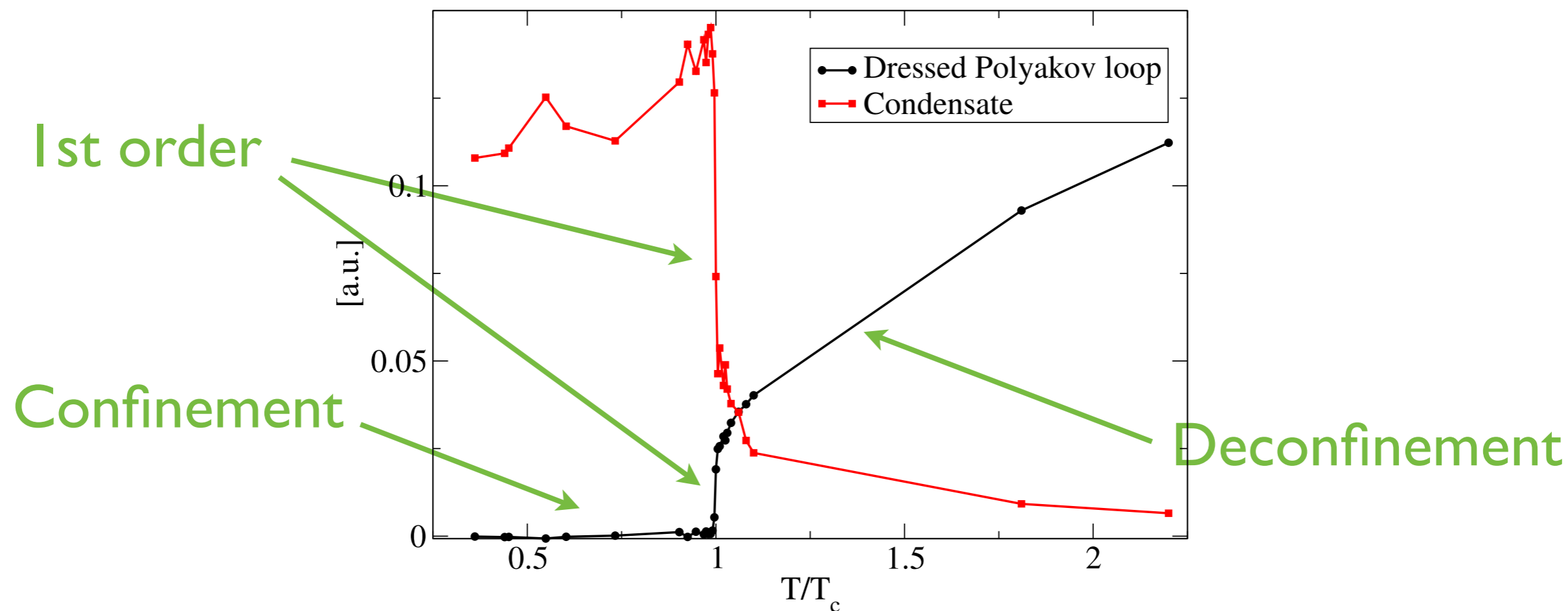
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Transition temperatures, quenched

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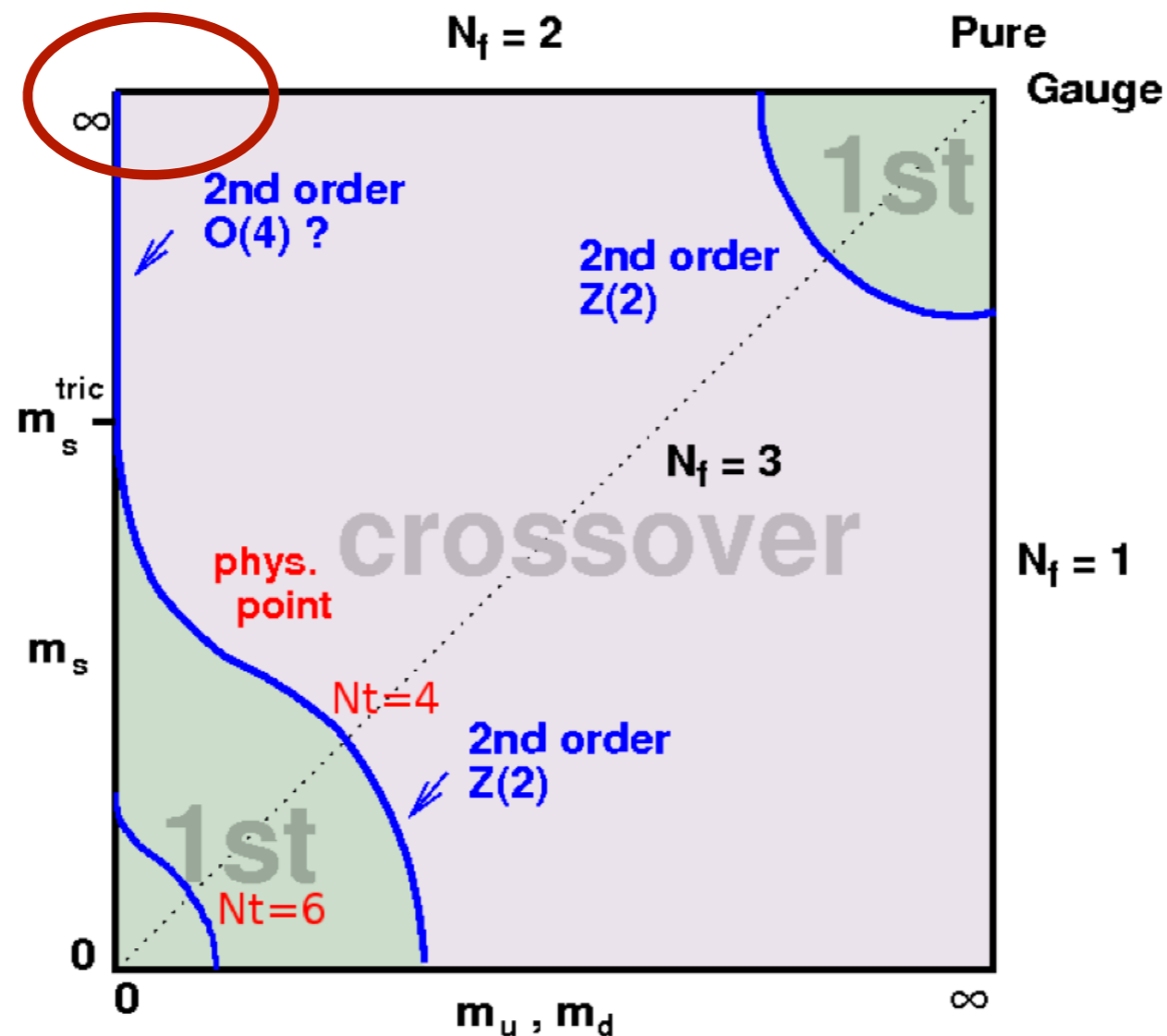


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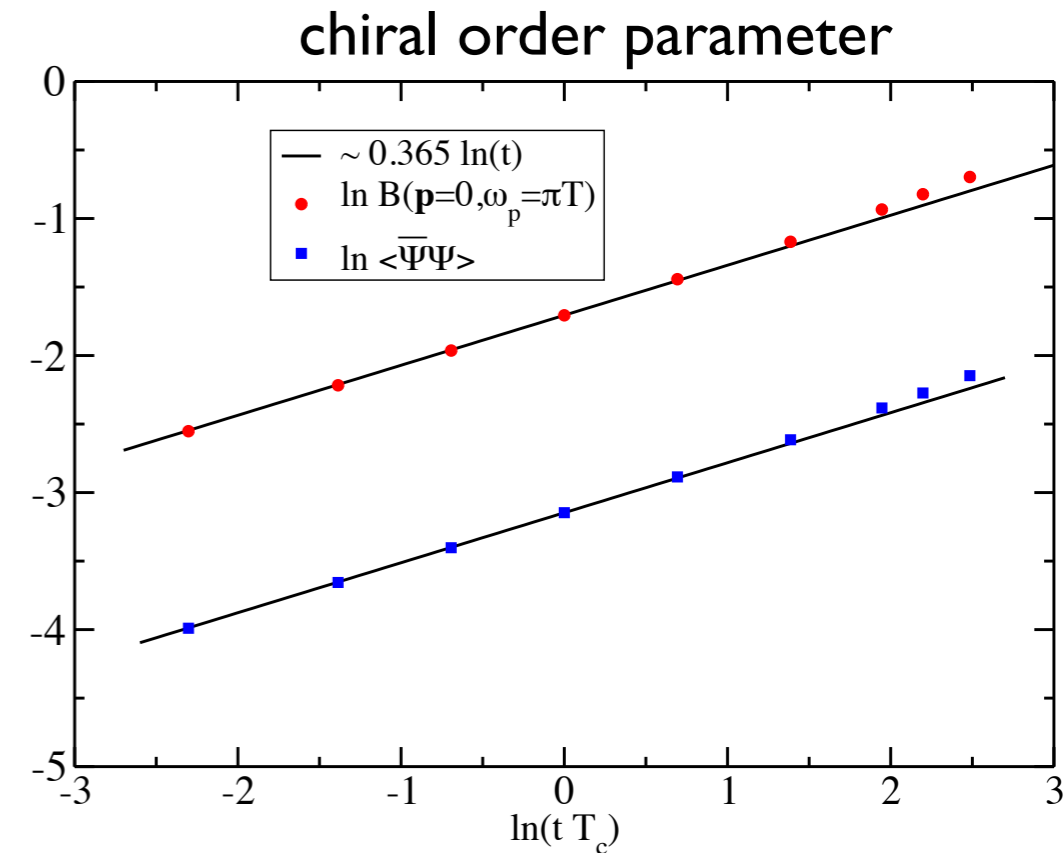
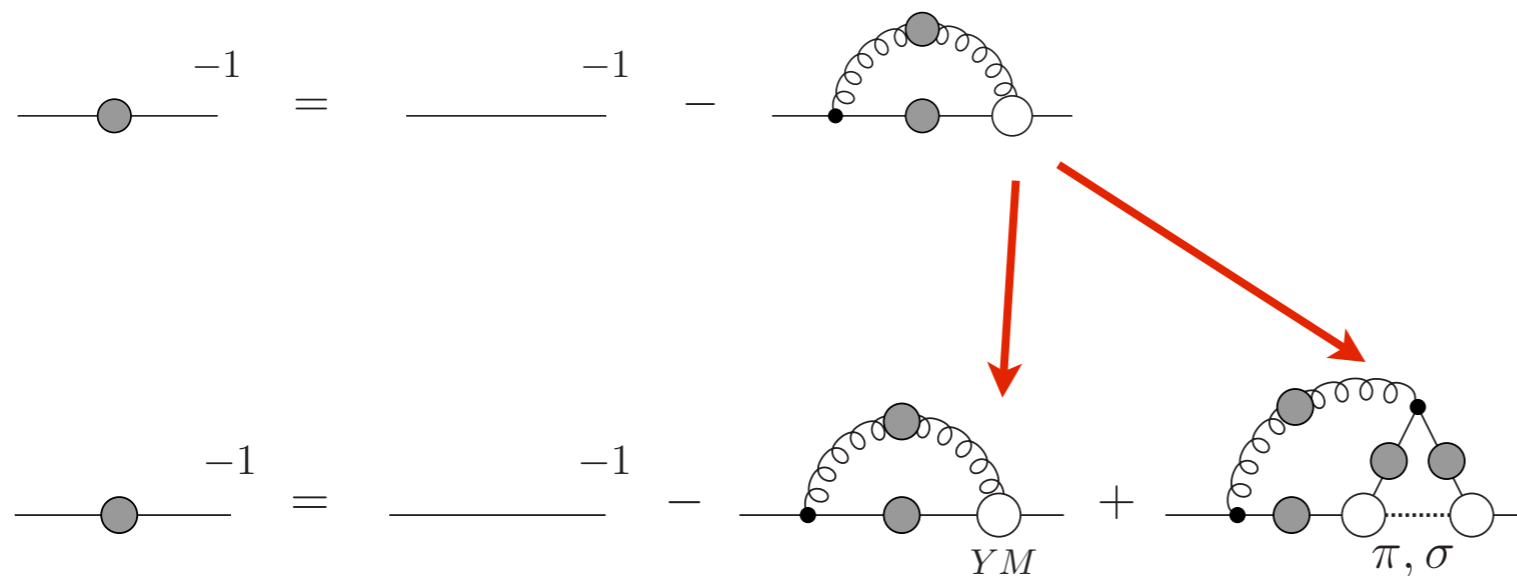
QCD phase transitions: chiral limit



- $N_f=2$, chiral limit: phase transition dominated by Goldstone boson physics \rightarrow Quark-Meson (QM) model
- $SU(2) \times SU(2) \cong O(4)$ -second order vs. $O(2) \times O(4)$ -first order

Pisarski and Wilczek, PRD 29 (1984) 338

$N_f=2$, chiral limit: Critical scaling from DSEs



- Crucial: take meson part of vertex explicitly into account

- $T \neq T_c$: meson corrections of order of 10-20 %

CF, Williams, PRD 78 (2008) 074006

- $T = T_c$: 2nd order phase transition,

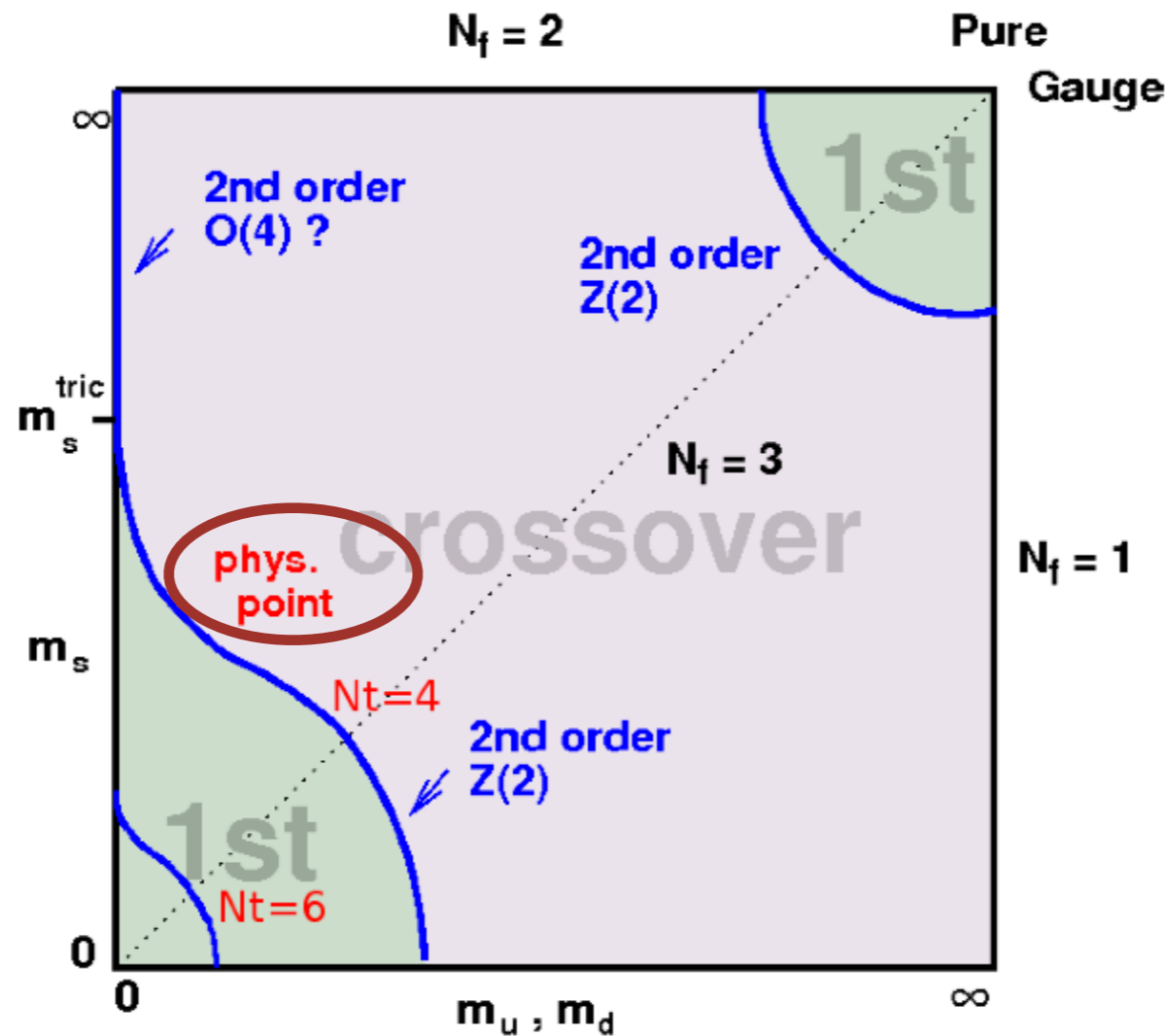
meson contributions are dominant ! (\rightarrow universality)

- Critical scaling: $\langle \bar{\Psi} \Psi \rangle(t) \sim B(t) \sim t^{\nu/2}$

$$f_{\pi,s}^2 \sim t^{\nu} \quad (t = (T_c - T)/T_c)$$

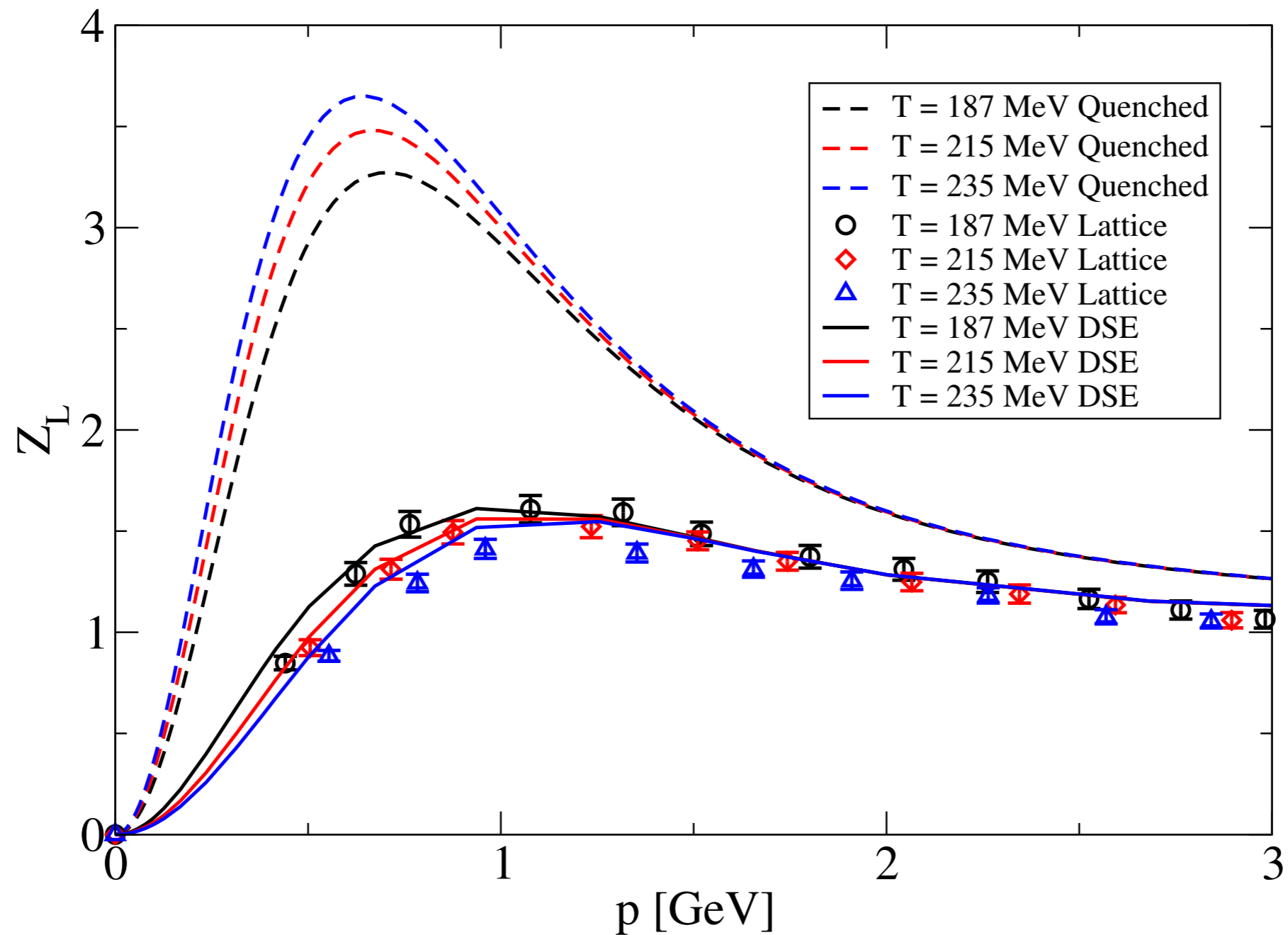
CF, Mueller, PRD 84 (2011) 054013

QCD phase transitions: $N_f=2+1$



- Physical up/down and strange quark masses
- Transition controlled by chiral dynamics
- at $\mu=0$: compare to available lattice results

Unquenched Gluon DSE vs Lattice

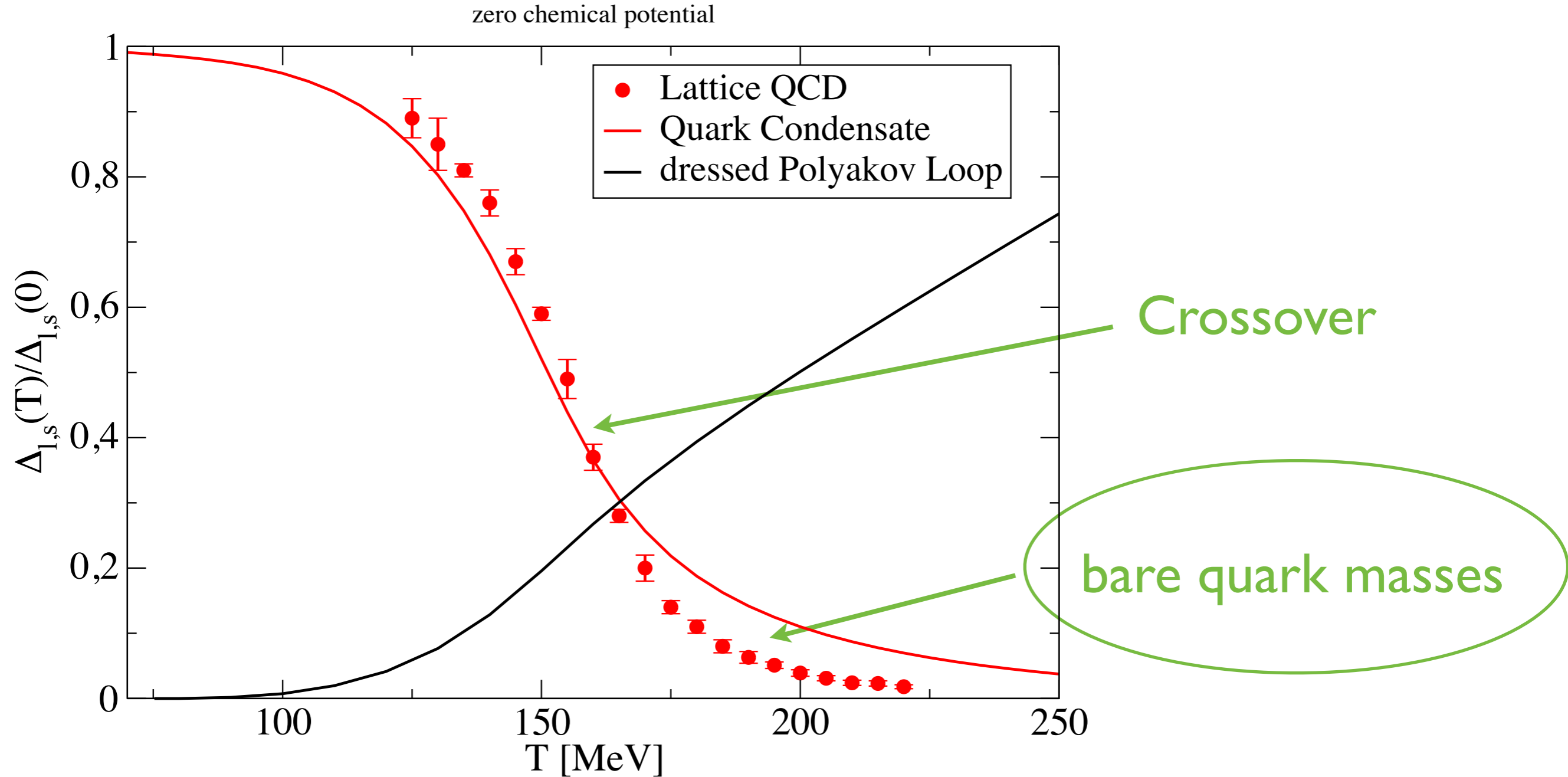


● quantitative agreement: **DSE prediction verified by lattice**

DSE: CF, Luecker, PLB 718 (2013) 1036 [arXiv:1206.5191]

Lattice: Aouane, Burger, Ilgenfritz, Muller-Preussker and Sternbeck, arXiv:1212.1102

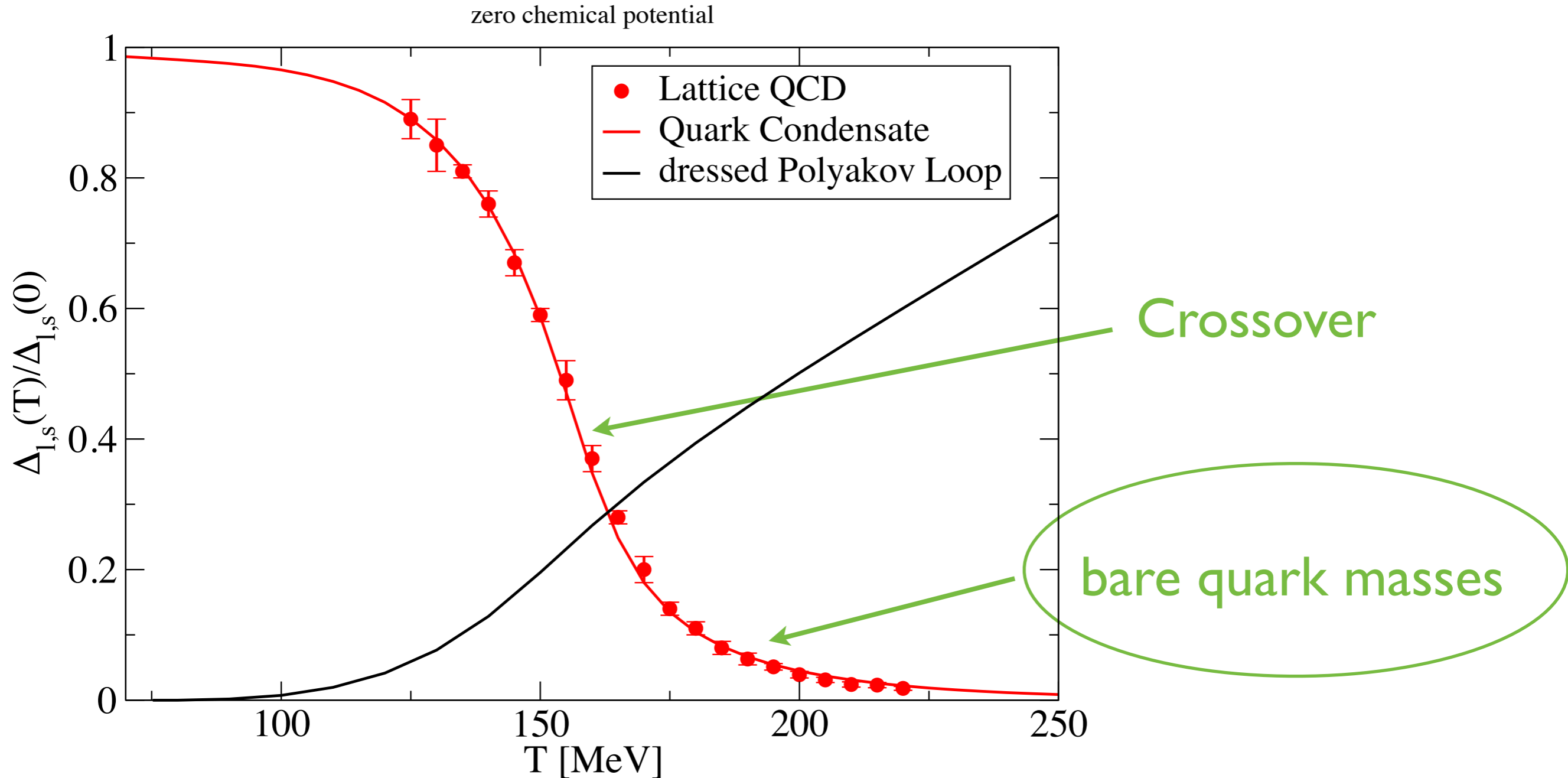
$N_f=2+1$, zero chemical potential



Lattice: Borsanyi *et al.* [Wuppertal-Budapest Collaboration], JHEP 1009(2010) 073

DSE: CF, Luecker, PLB 718 (2013) 1036, CF, Luecker, Welzbacher, in prep.

$N_f=2+1$, zero chemical potential

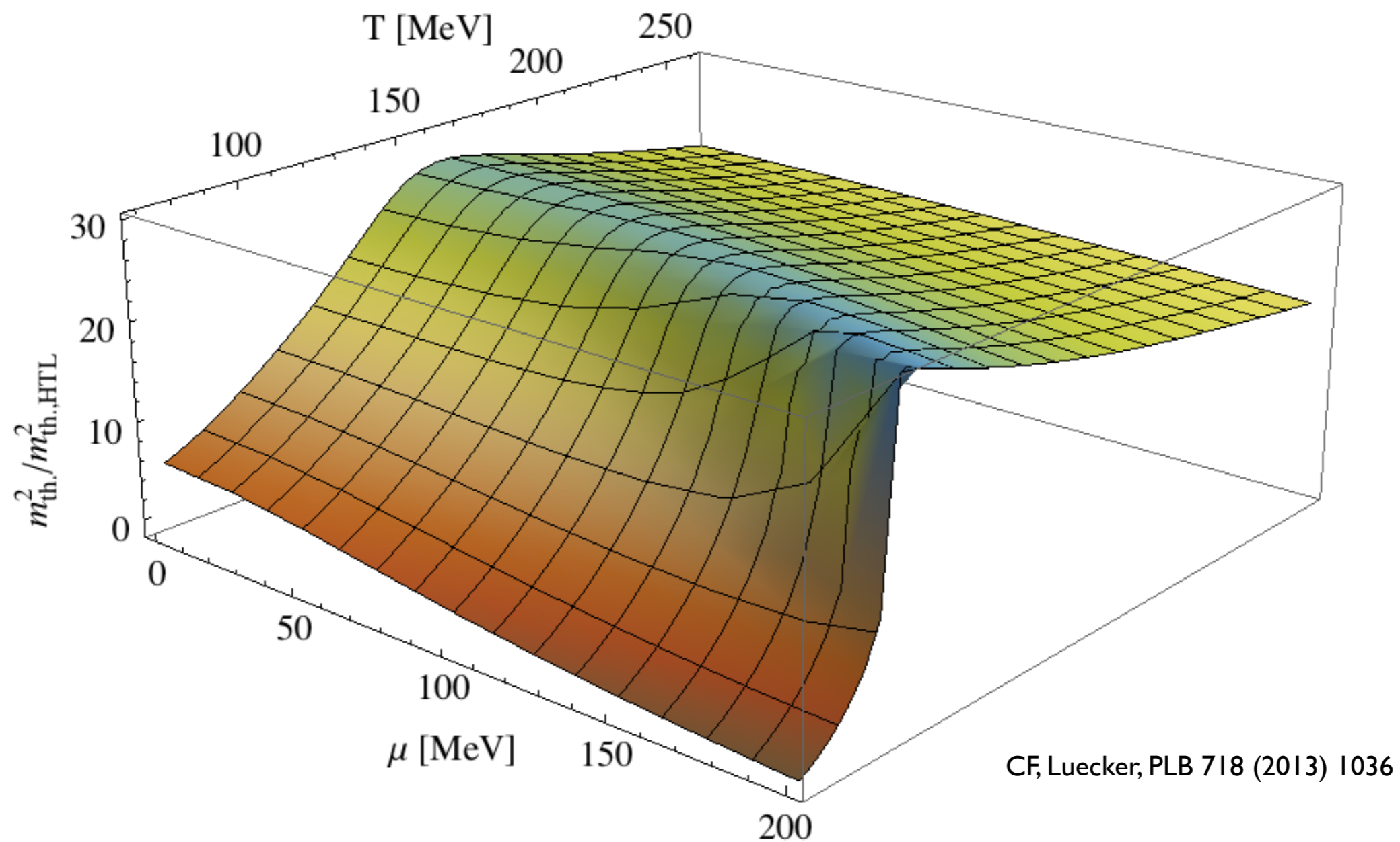


Lattice: Borsanyi *et al.* [Wuppertal-Budapest Collaboration], JHEP 1009(2010) 073

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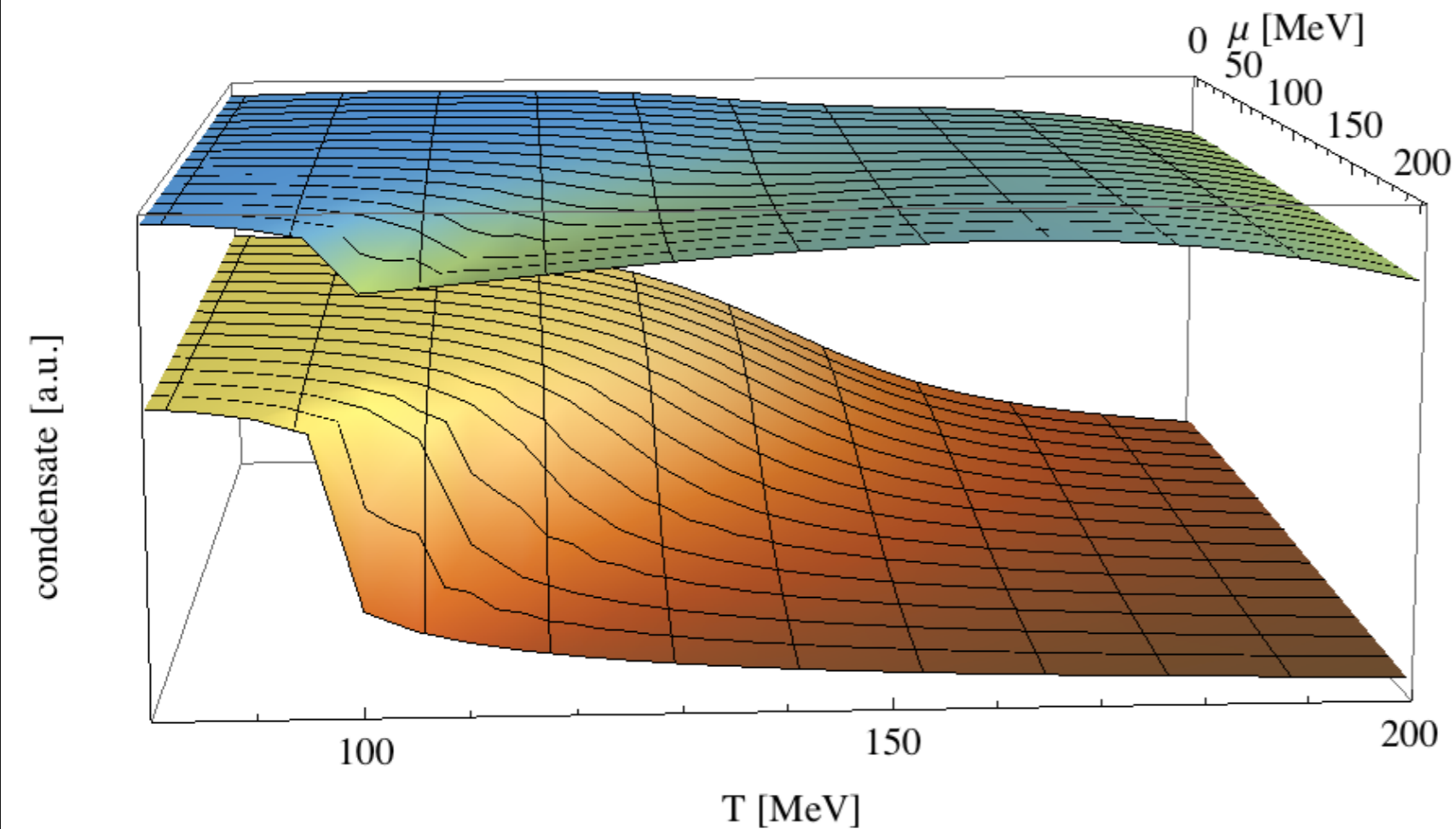
● quantitative agreement

$N_f=2+1$: thermal electric gluon mass

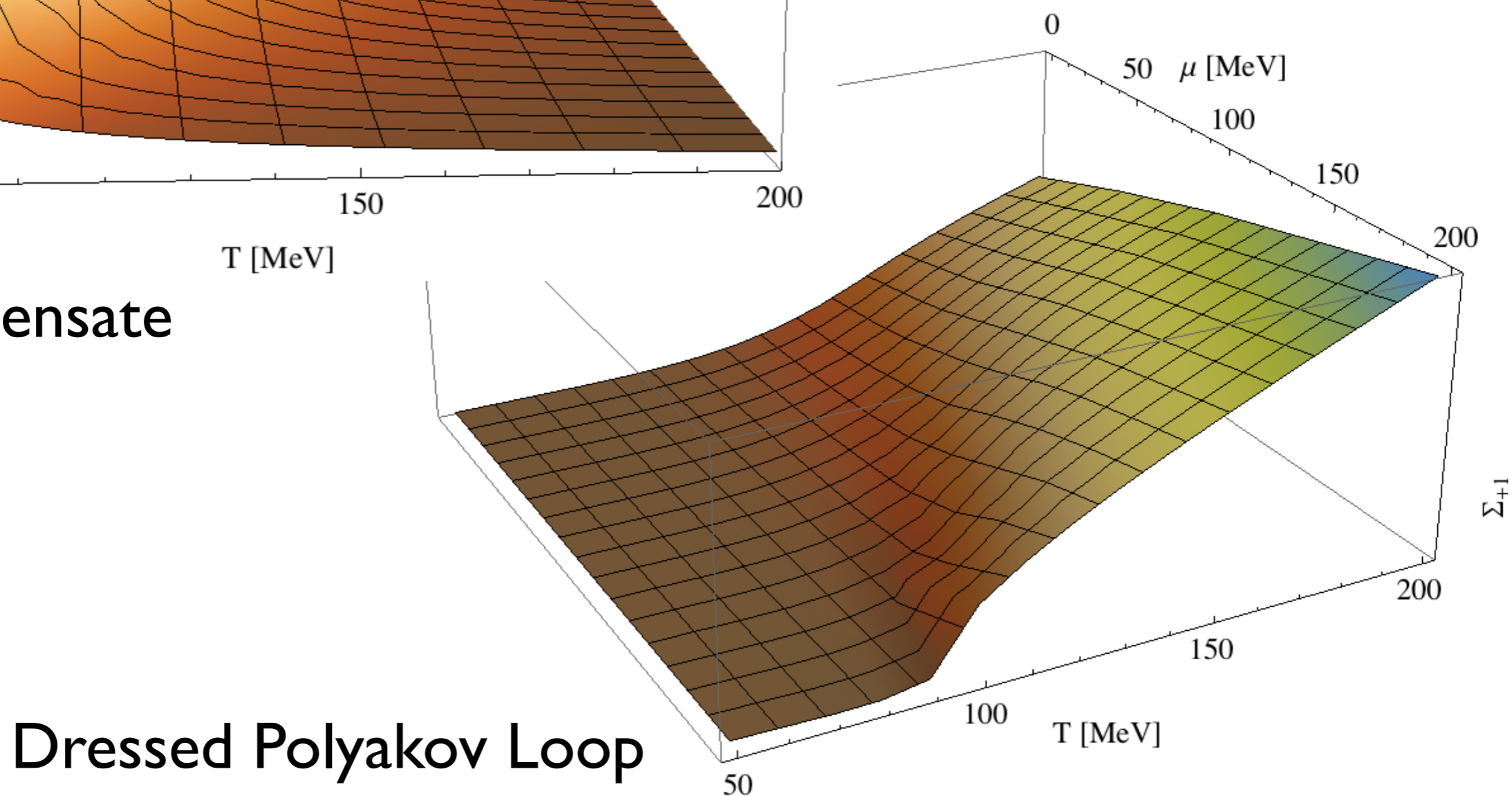


- large temperatures: behavior as expected from HTL
- first order transition at large chemical potential

$N_f=2+1$: Condensate and dressed Polyakov Loop



Quark condensate



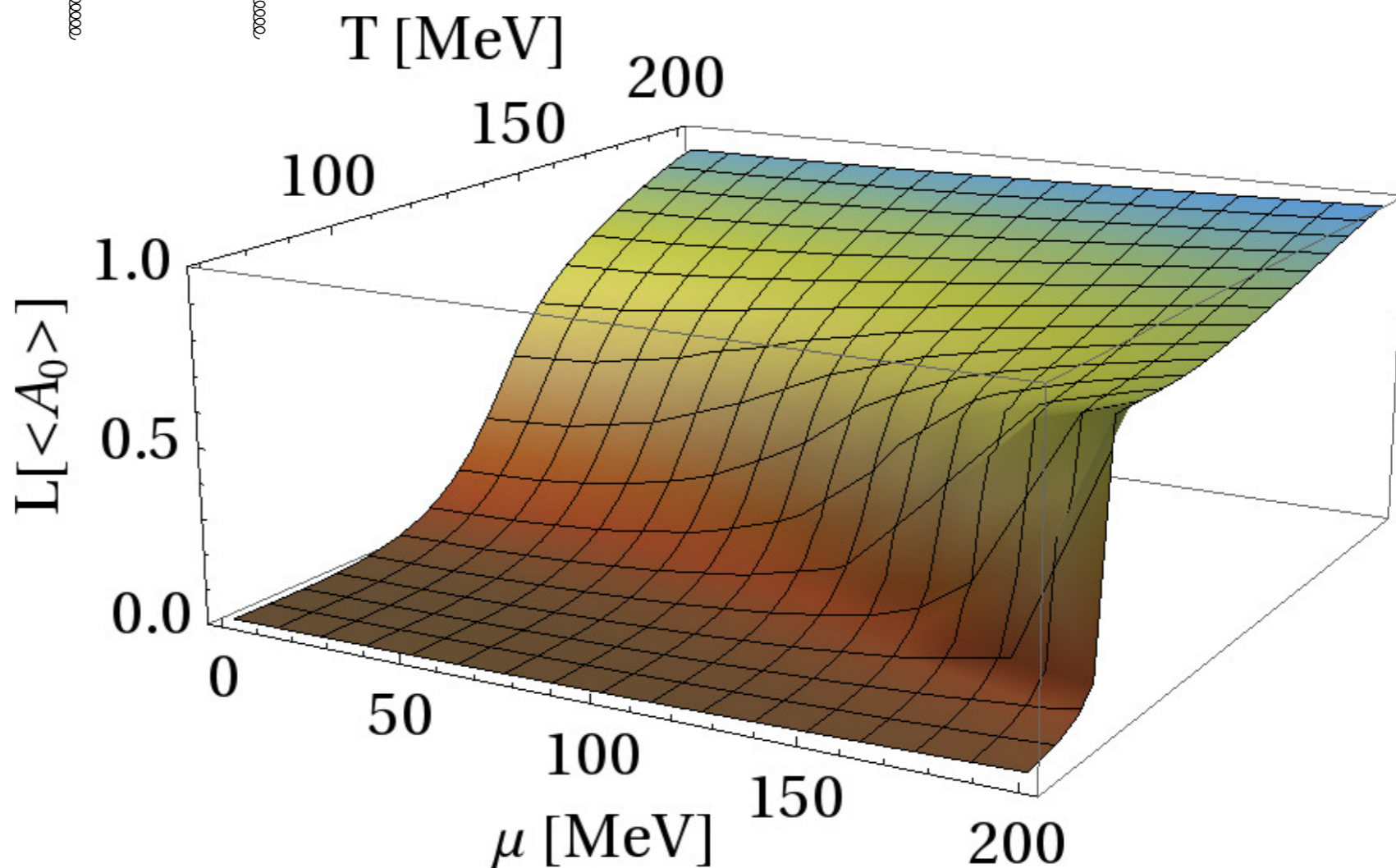
Dressed Polyakov Loop

$N_f=2+1$: Polyakov loop potential at finite μ

$$\frac{\delta(\Gamma - S)}{\delta A_0} = \frac{1}{2} \left[\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} - \frac{1}{6} \left(\text{Diagram 4} + \text{Diagram 5} \right) \right]$$

Polyakov-Loop

$$L = \frac{1}{N_c} \text{tr} e^{ig \int A_0}$$

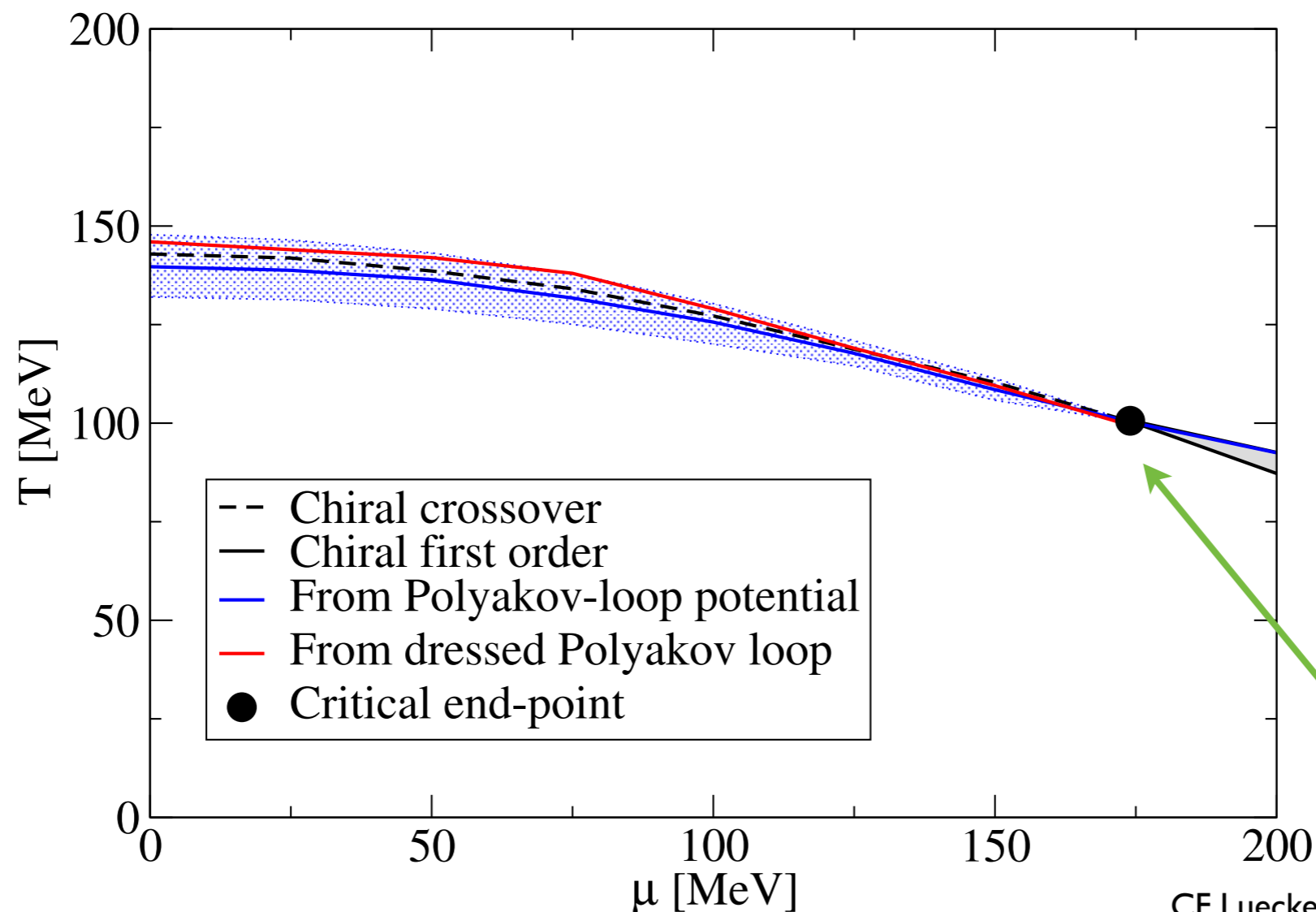


CF, Fister, Luecker, Pawłowski, arXiv:1306.6022

- evaluated from Polyakov-Loop potential
- important input for P-models: PQM, PNJL !

Herbst, Mitter, Pawłowski, Schaefer, Stiele, arXiv:1308.3621

$N_f=2+1$: Polyakov loop and phase diagram



CF, Luecker, PLB 718 (2013) 1036,
CF, Fister, Luecker, Pawłowski, arXiv:1306.6022

CEP at large μ

- no CEP at $(\mu_B)_c/T_c < 2$ in agreement with lattice and FRG

de Forcrand, Philipsen, JHEP 0811 (2008) 012; Nucl Phys. B642 (2002) 290-306

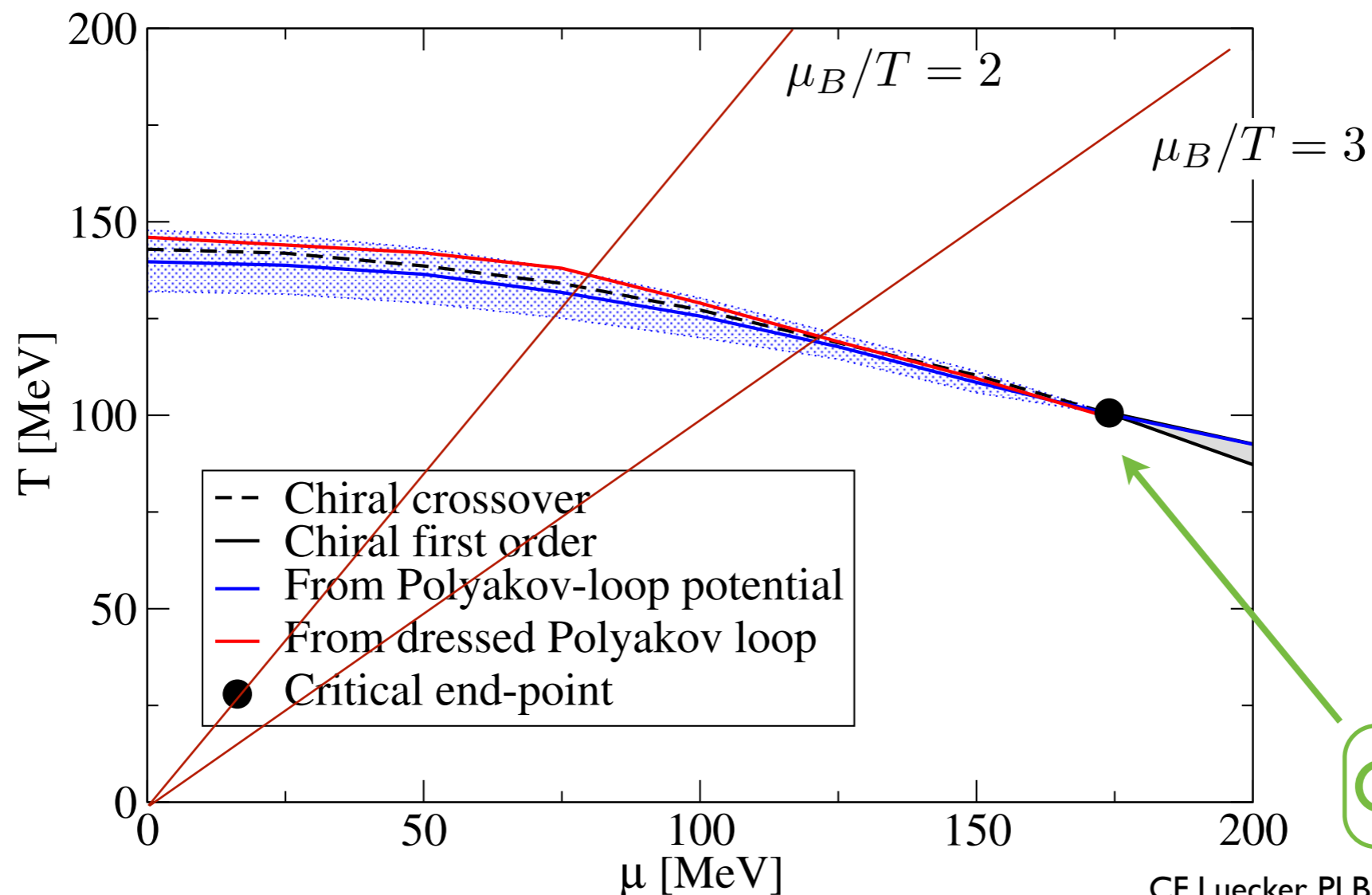
Endrodi, Fodor, Katz, Szabo, JHEP 1104 (2011) 001

Herbst, Pawłowski, Schaefer, PRD 88 (2013) 014007

Caveat: baryon effects missing...

$N_c=2$: Brauner, Fukushima and Hidaka, PRD 80 (2009) 74035
Strodthoff, Schaefer and Smekal, PRD 85 (2012) 074007

$N_f=2+1$: Polyakov loop and phase diagram



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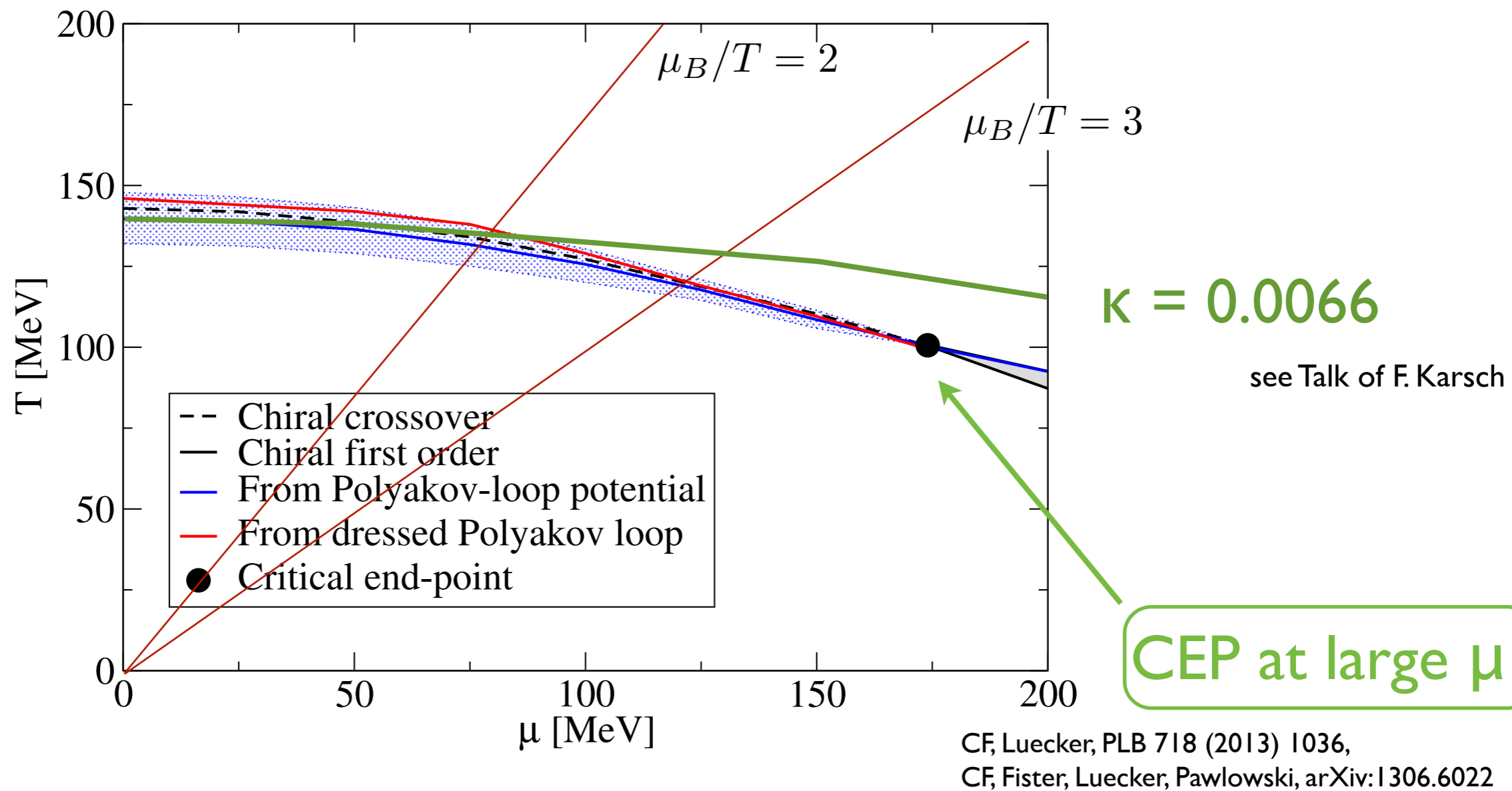
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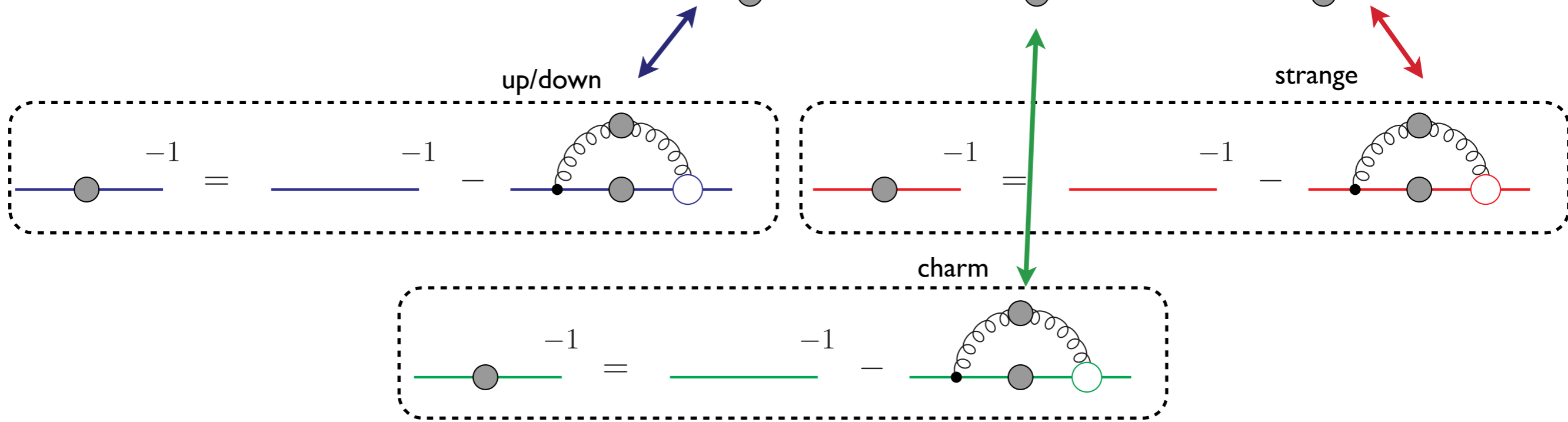
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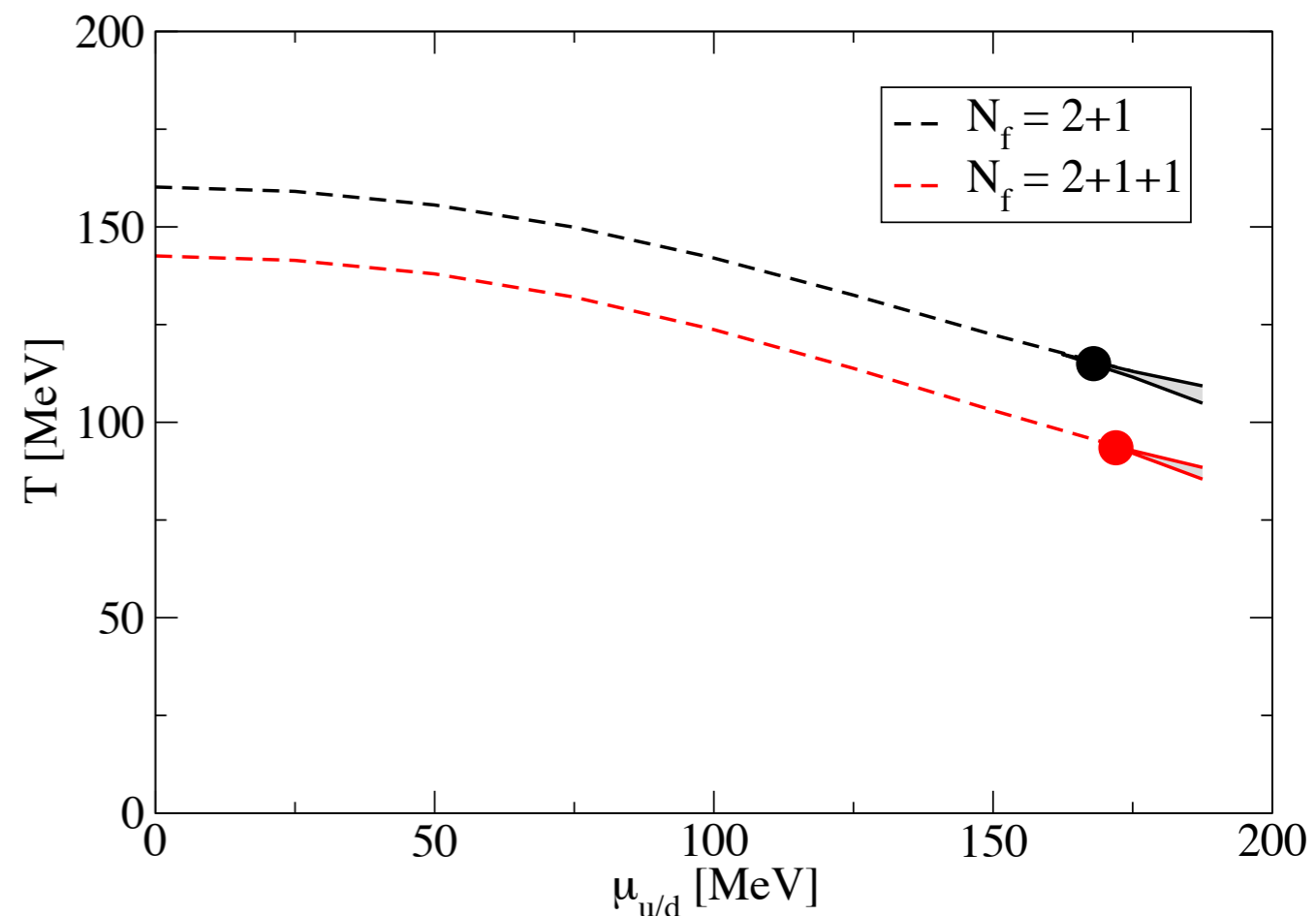
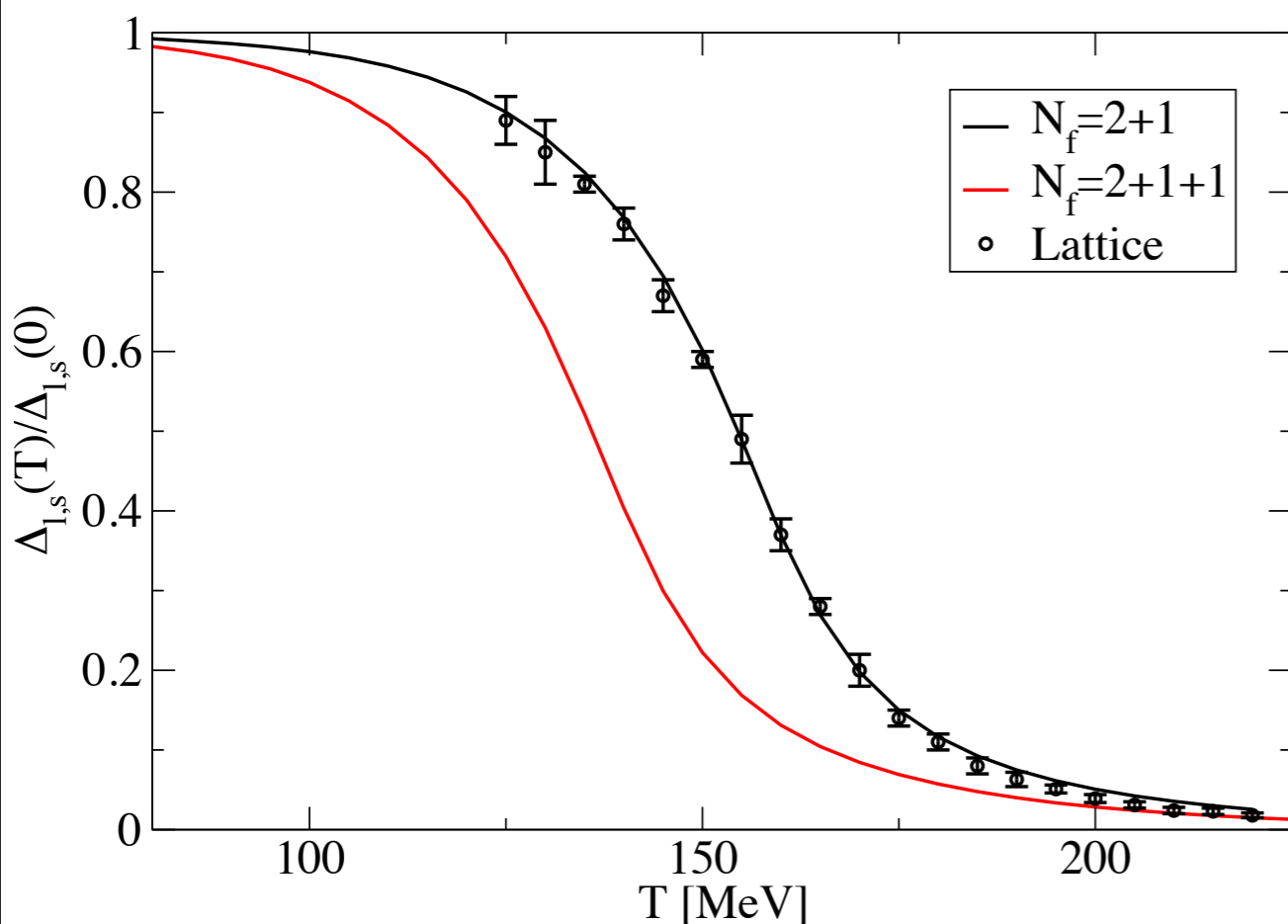
$N_c=2$: Brauner, Fukushima and Hidaka, PRD 80 (2009) 74035
 Strodthoff, Schaefer and Smekal, PRD 85 (2012) 074007

$N_f=2+1+1$ -QCD with DSEs

$$\begin{array}{c} \text{wavy line with grey blob} \end{array}^{-1} = \begin{array}{c} \text{wavy line with yellow blob} \end{array}^{-1} + \begin{array}{c} \text{blue loop} \end{array} + \begin{array}{c} \text{green loop} \end{array} + \begin{array}{c} \text{red loop} \end{array}$$



$N_f=2+1+1$ -QCD with DSEs (preliminary)



CF, Luecker, Welzbacher, in preparation

- Interaction fixed: T_{PC} decreases by $O(10 \text{ MeV})$
- Physics fixed (m_π, f_π): T_{PC} similar (not shown)

→ see also Talk by C. Sasaki, 2nd week

Summary

- Gluon spectral functions at $T=0$: **positivity violation**
- Temperature dependent gluon propagator
 - characteristic behavior of electric gluon
 - 'melting' of magnetic gluon with temperature
- Deconf. T_{pc} from **dressed Polyakov-loop/Polyakov-loop potential**
- QCD with finite chemical potential (beyond mean field)
 - **backreaction of quarks onto gluons important**
 - **$N_f=2+1$ and $N_f=2+1+1$: CEP at $\mu_c/T_c > 1$**

Work in progress: include baryons...

include magnetic field...

Mueller, Bonnet, CF in preparation

→ see also Talk by T. Kojo, 2nd week