



Heavy quarks from CGC in p+A collisions

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with K. Watanabe

HF and K. Watanabe NPA915 (2013) 1 [arXiv:1304.2221[hep-ph]]
NPA920 (2013) 78 [arXiv:1308.1258[hep-ph]]

Motivation

- Heavy quarks produced from gluons only in initial hard interactions
- Heavy quarks in pA collisions at the LHC
 - probe small-x gluons
 - crucial baseline for calibration of AA collisions

Outline

- Introduction:

Glueon distribution at small x in CGC (rcBK)

- Quark pair production in pA

- quarkonium

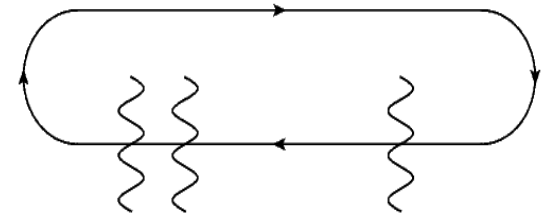
- open charm

- Summary & Outlook

Gluon distribution at small x

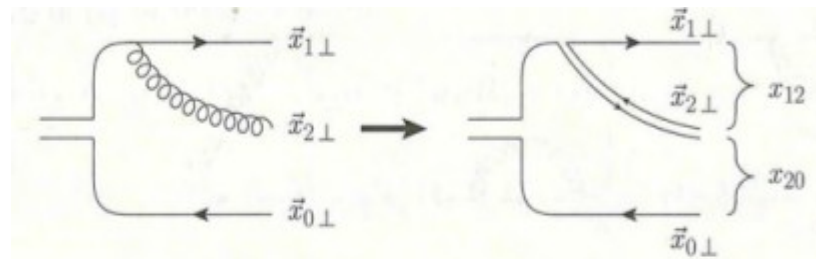
- “Gluon distribution” is related to $1 - \mathcal{N}(r, y) = \frac{1}{N_c} \text{tr} \langle \tilde{U}(r_1) \tilde{U}^\dagger(r_2) \rangle$

$$\phi(\mathbf{k}, y) \sim -\mathbf{k}^2 \text{F.T.} \mathcal{N}(r, y)$$



- x-dependence is described by BK eqn in large N_c
Kovchegov-Weigert, Balitsky-Chirilli

$$\frac{\partial \mathcal{N}(r, y)}{\partial y} = \int d^2 \mathbf{r}_1 K^{\text{run}} [\mathcal{N}(r_1, y) + \mathcal{N}(r_2, y) - \mathcal{N}(r, y) - \mathcal{N}(r_1, y) \mathcal{N}(r_2, y)]$$



Constraining N with DIS data (AAMQS)

Initial condition and evolution

Albacete et al. PRD80

$$\mathcal{N}_F(r, x=x_0) = 1 - \exp \left[-\frac{(r^2 Q_{s0, \text{proton}}^2)^\gamma}{4} \ln \left(\frac{1}{\Lambda r} + e \right) \right]$$

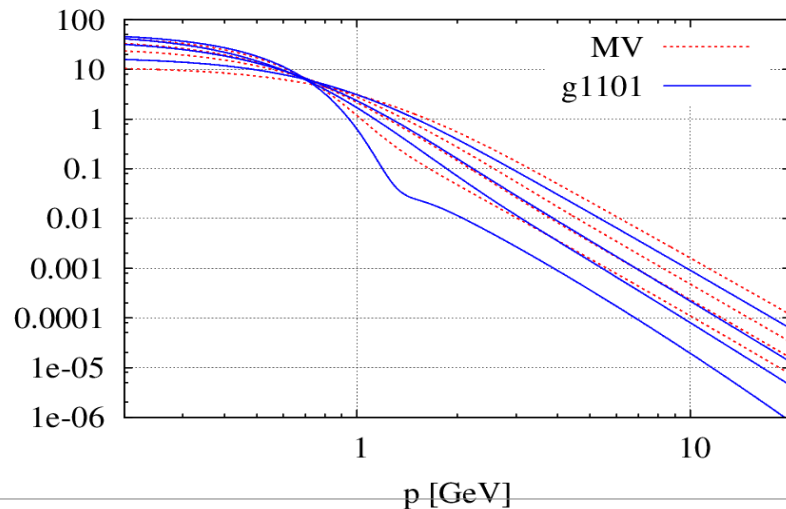
$$\alpha_s(r) = 1/[b_0 \ln(4C/r\Lambda + a)]$$

other parametrization, Lappi et al. arXiv:1309.6963

Parameter set

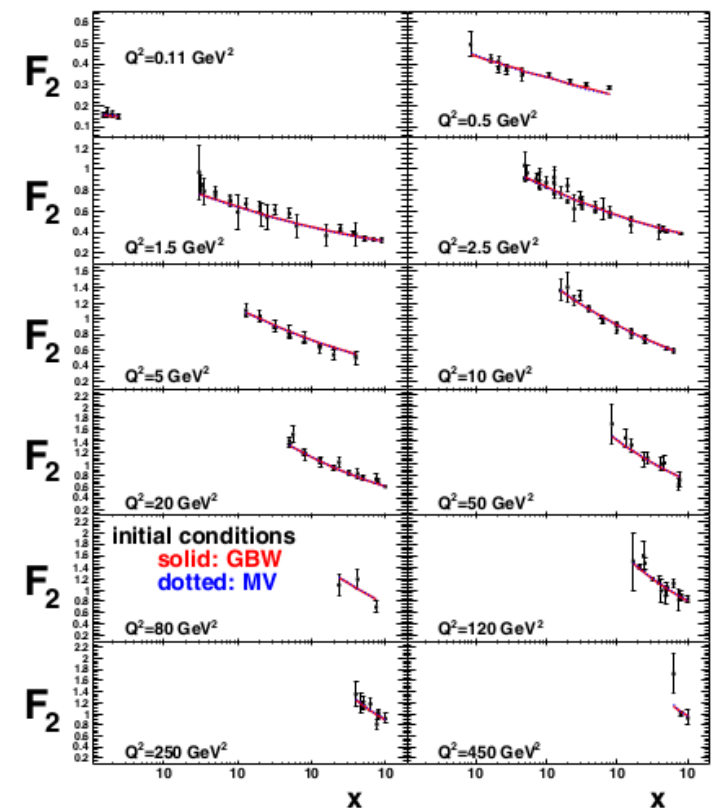
UGD Set	$Q_{s0, \text{proton}}^2$ (GeV ²)	γ	α_{fr}	C
MV	0.2	1	0.5	1
g1.119	0.168	1.119	1.0	2.47
g1.101	0.157	1.101	0.8	1

$N(k, y)$, $y=0, 1.5, 3, 6$



Fit to HERA-DIS data

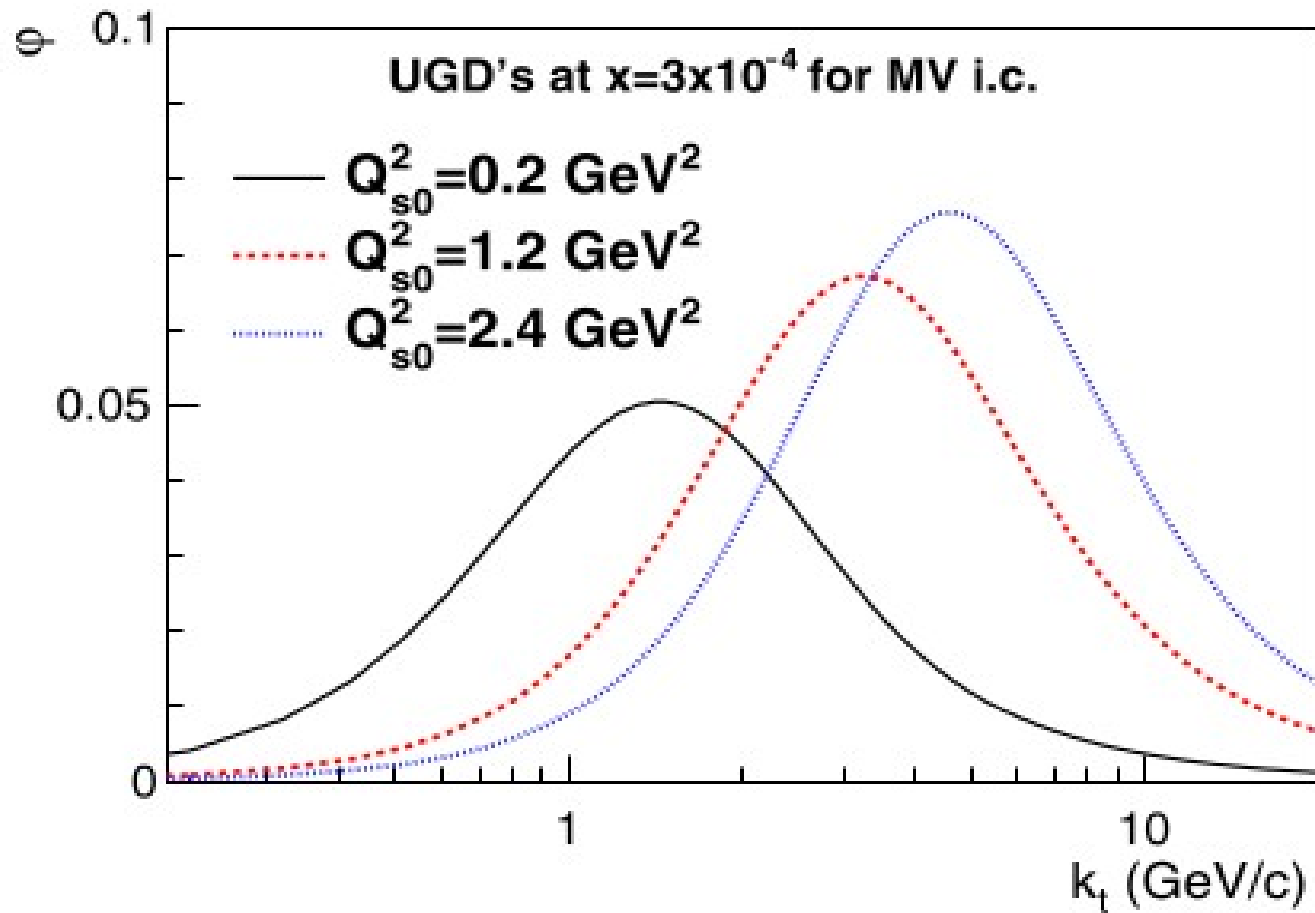
$$\sigma_{T,L}(x, Q^2) = \sigma_0 \int_0^1 dz \int dr |\Psi_{T,L}(z, Q^2, \mathbf{r})|^2 \mathcal{N}(r, Y)$$



Gluon distributions for p and A

$$Q_{sA}^2 = A^{1/3} Q_{sp}^2$$

or MC

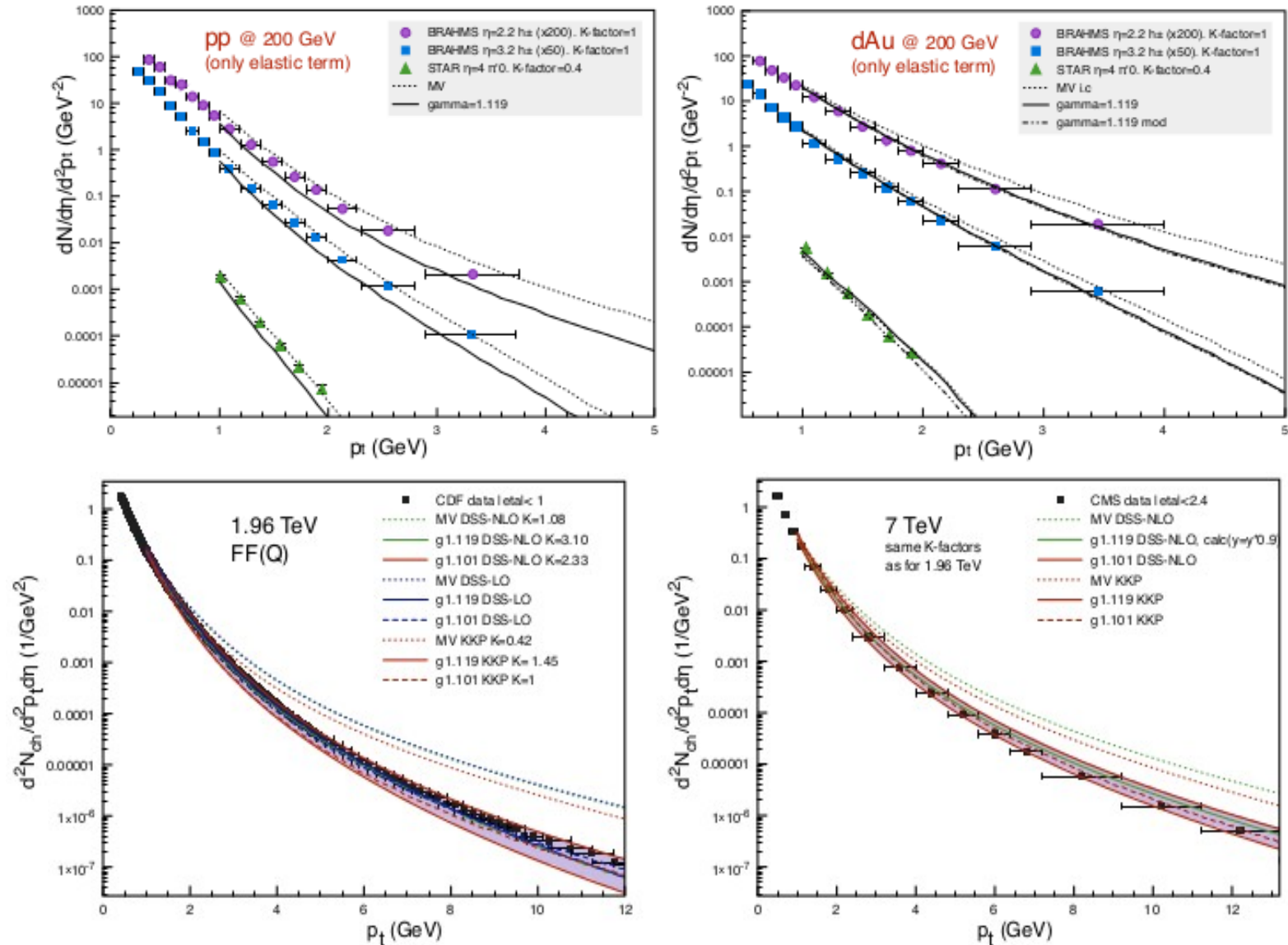


Hadron spectrum at RHIC, Tevatron, LHC

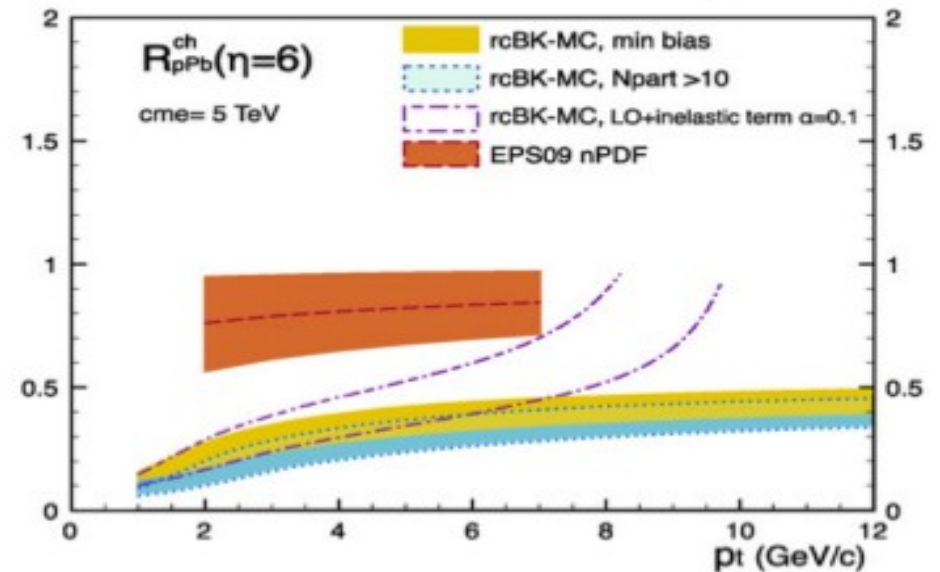
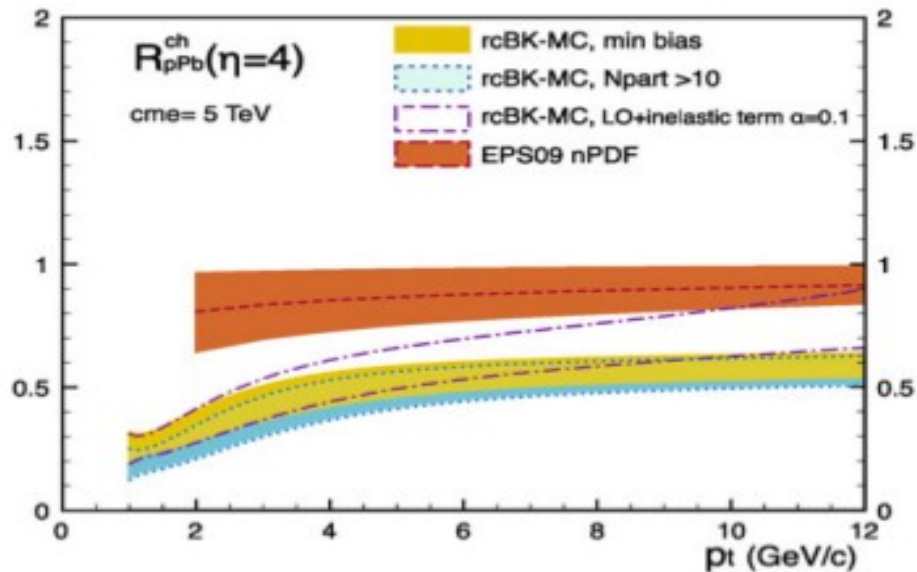
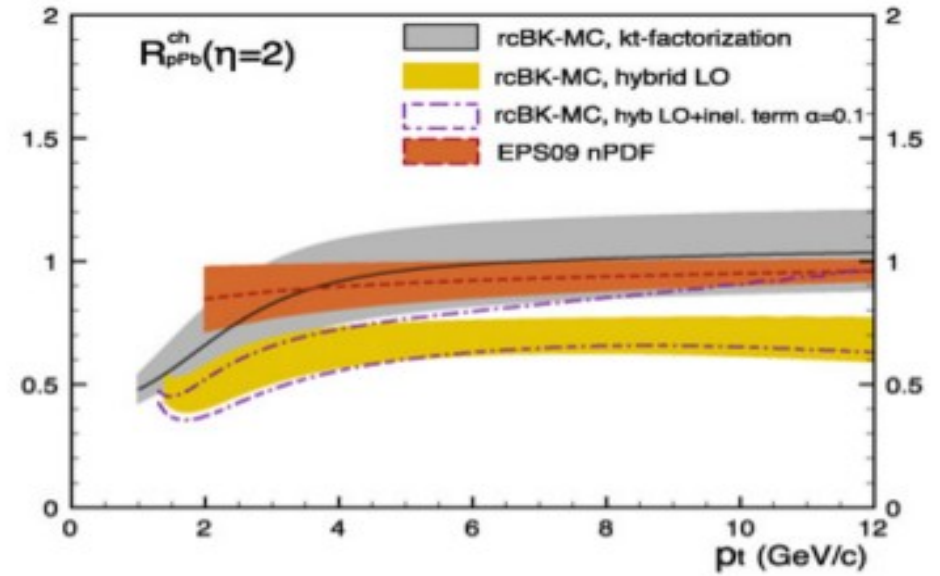
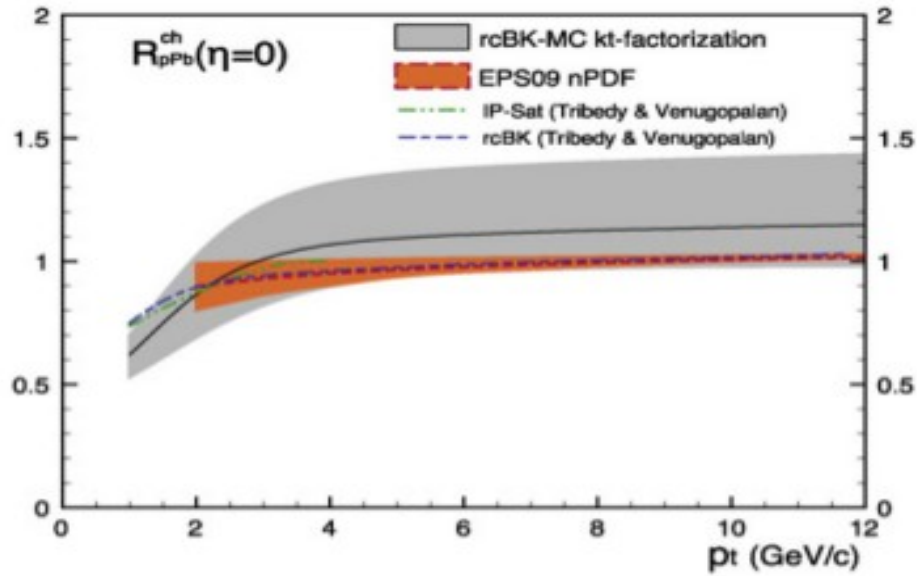
k_T factorization ($y \sim 0$), DHJ ($y > 0$) formulae

$$Q_{sA}^2 = A^{1/3} Q_{sp}^2$$

or MC



RpA(pT) of single hadron at $y=0,2,4,6$ at the LHC



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- Quark pair production in pA

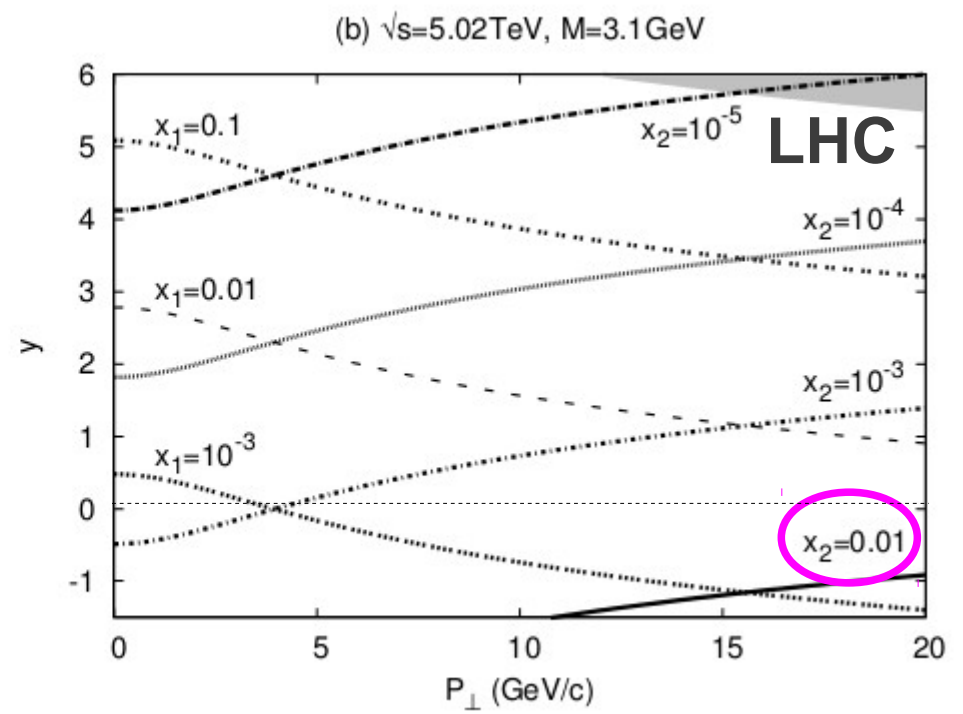
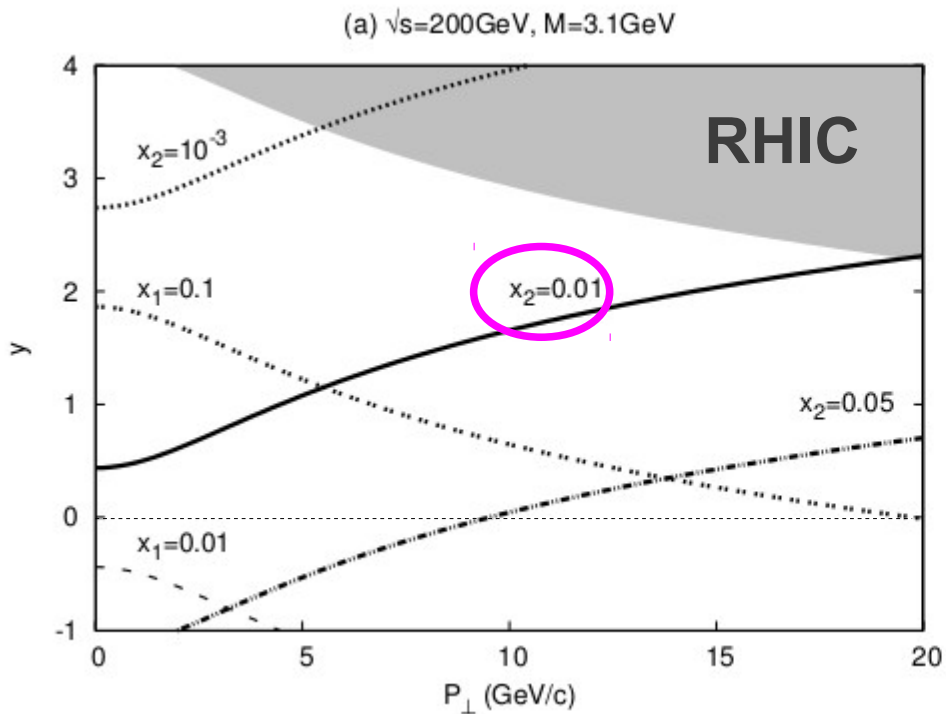
- quakonium

- open charm

- Summary & Outlook

Heavy quarks sensitive small x?

- Quarkonium; $gg \rightarrow J/\psi$
- maybe at RHIC, while must be at the LHC

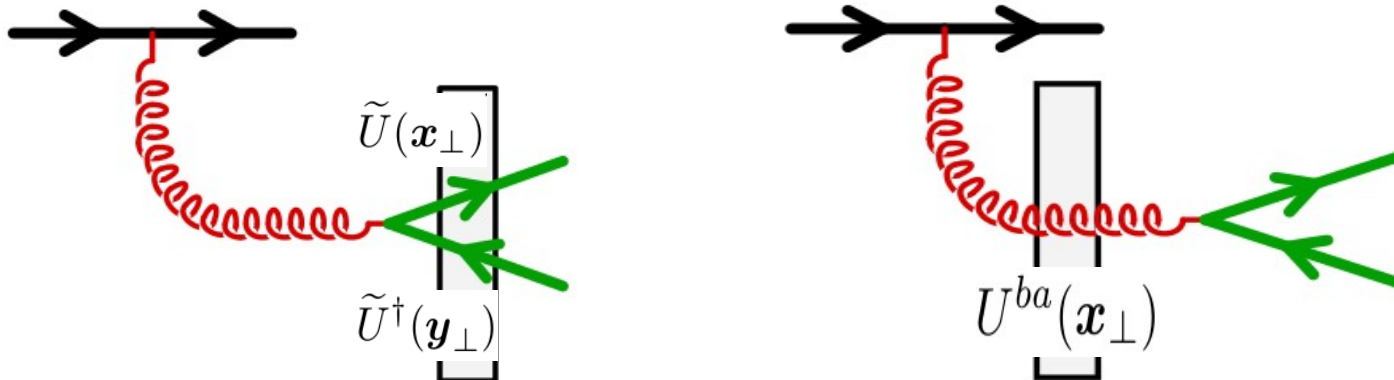


Quark pair production amplitude

Blaizot-Gelis-Venugopalan (2004)

$$\mathcal{M}_F(\mathbf{q}, \mathbf{p}) = g^2 \int \frac{d^2 \mathbf{k}_{1\perp}}{(2\pi)^2} \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} \frac{\rho_{p,a}(\mathbf{k}_{1\perp})}{k_{1\perp}^2} \int d^2 \mathbf{x}_\perp d^2 \mathbf{y}_\perp e^{i\mathbf{k}_\perp \cdot \mathbf{x}_\perp} e^{i(\mathbf{p}_\perp + \mathbf{q}_\perp - \mathbf{k}_\perp - \mathbf{k}_{1\perp}) \cdot \mathbf{y}_\perp}$$

$$\times \bar{u}(\mathbf{q}) \left\{ T_{q\bar{q}}(\mathbf{k}_{1\perp}, \mathbf{k}_\perp) [\tilde{U}(\mathbf{x}_\perp) t^a \tilde{U}^\dagger(\mathbf{y}_\perp)] + T_g(\mathbf{k}_{1\perp}) [t^b U^{ba}(\mathbf{x}_\perp)] \right\} v(\mathbf{p}),$$

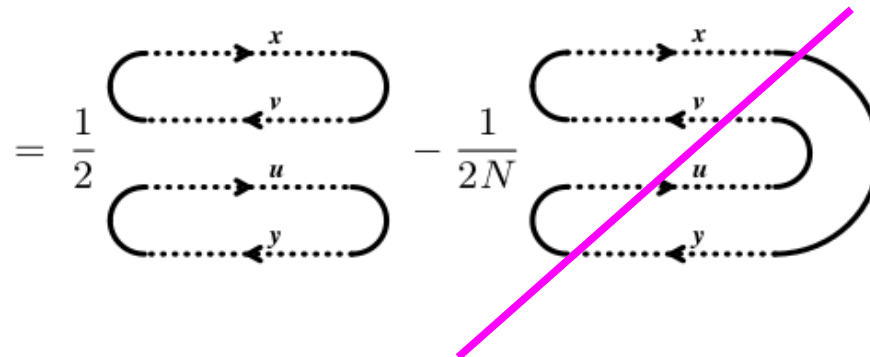


Cross-section in large N

$$\frac{dN_{q\bar{q}}}{d^2\mathbf{p}_\perp d^2\mathbf{q}_\perp dy_p dy_q} = \frac{1}{\pi R_A^2} \frac{\alpha_s^2 N}{8\pi^4 d_A} \frac{1}{(2\pi)^2} \int_{\mathbf{k}_{2\perp}, \mathbf{k}_\perp} \frac{\Xi(\mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}, \mathbf{k}_\perp)}{k_{1\perp}^2 k_{2\perp}^2} \phi_{A, y_2}^{q\bar{q}, g}(\mathbf{k}_{2\perp}, \mathbf{k}_\perp) \varphi_{p, y_1}(\mathbf{k}_{1\perp})$$

- Need 4-pt function (kT fact breaking)
- but, in large N limit, it reduces to a product of dipole amplitudes

$$C(\mathbf{x}_\perp, \mathbf{y}_\perp; \mathbf{u}_\perp, \mathbf{v}_\perp) \equiv \text{tr}_c \langle \tilde{U}(\mathbf{x}_\perp) t^a \tilde{U}^\dagger(\mathbf{y}_\perp) \tilde{U}(\mathbf{u}_\perp) t^a \tilde{U}^\dagger(\mathbf{v}_\perp) \rangle$$



Quarkonium in Color Evaporation Model (CEM)

$$\frac{dN_{J/\psi}}{d^2\mathbf{P}_\perp dy} = F_{J/\psi} \int_{4m_c^2}^{4M_D^2} dM^2 \frac{dN_{c\bar{c}}}{d^2\mathbf{P}_\perp dM^2 dy}$$

- Assumption:

Pair bounds to J/psi non-perturbatively, irrespective of its color, outside the nucleus

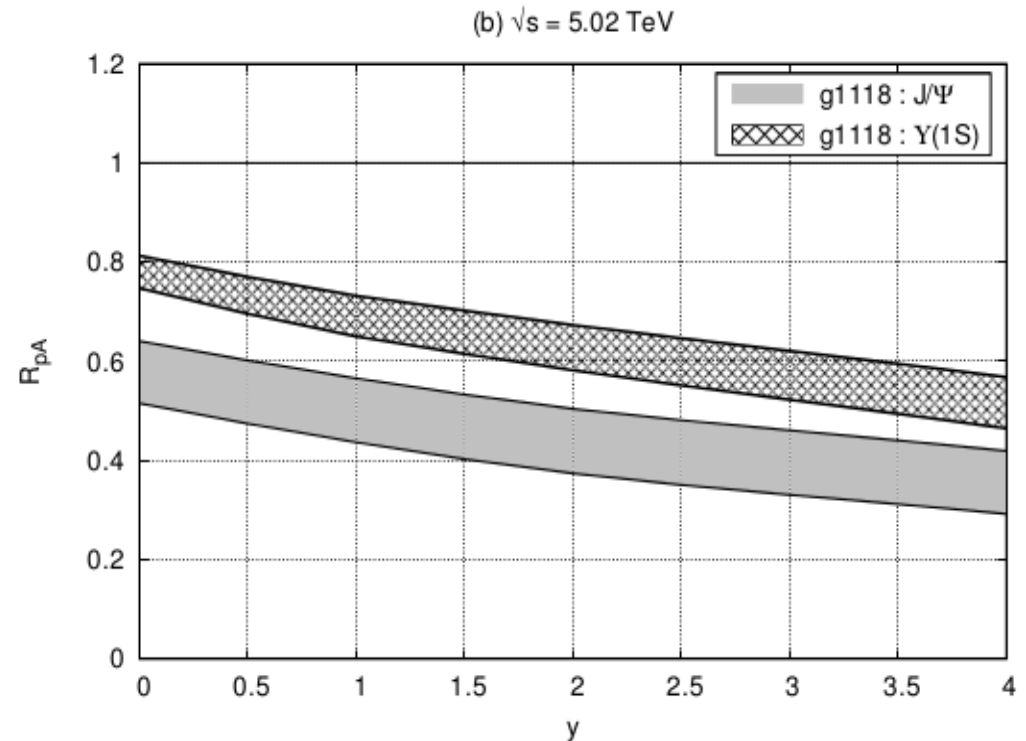
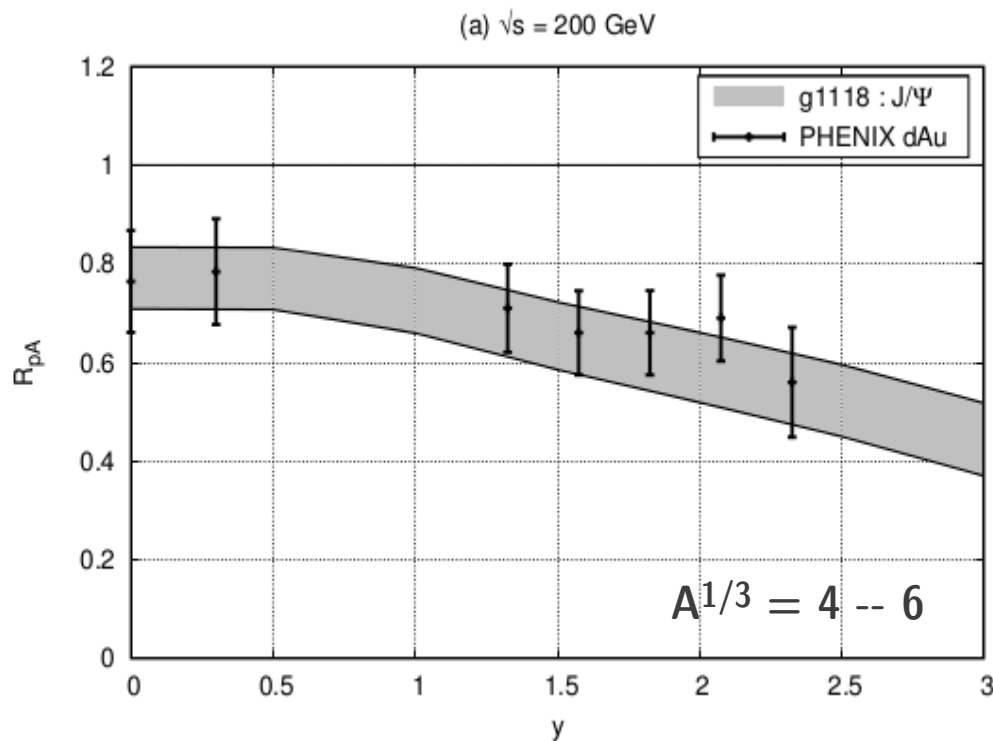
- Cold Nuclear Effect = Initial gluon saturation + multiple scatterings of pair

- Open charm production is evaluated with Fragmentation Func

R_{pA}(y) for J/ψ at RHIC & LHC

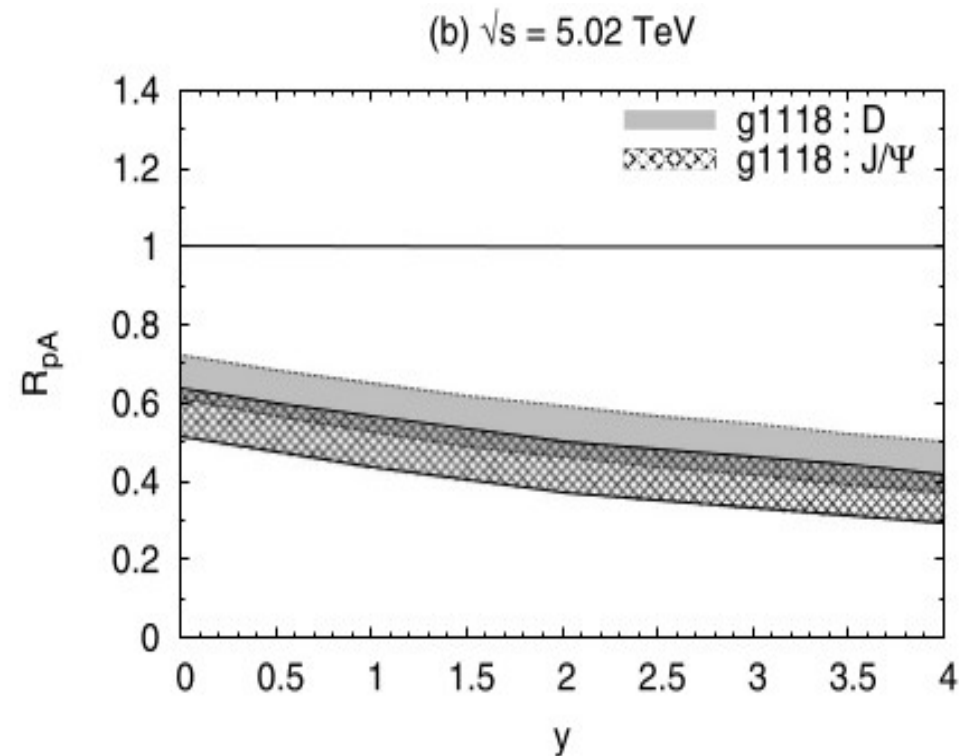
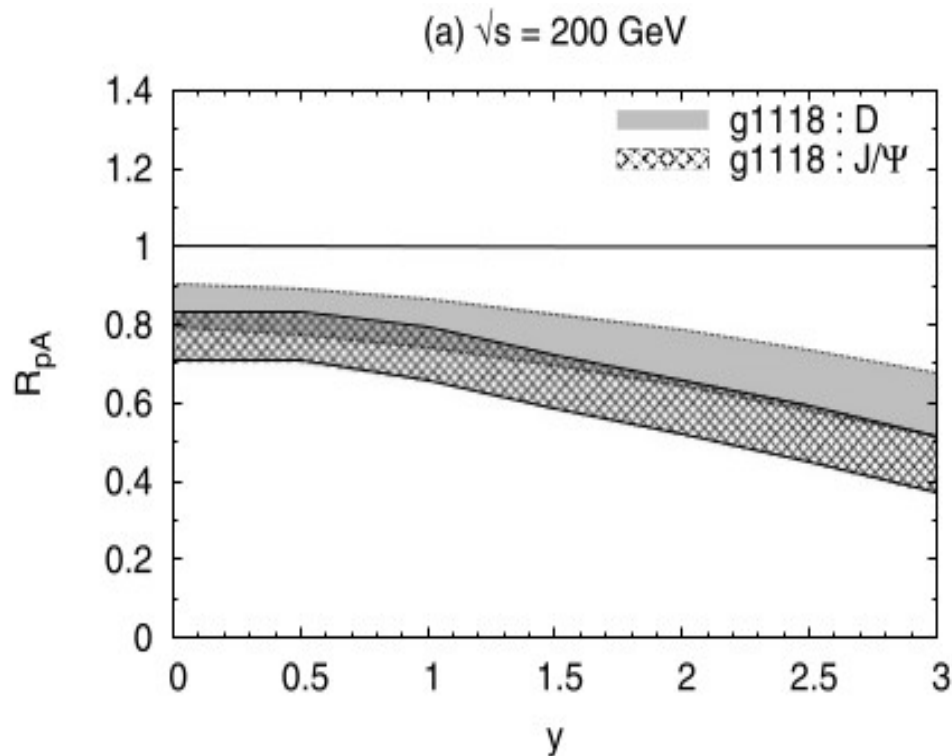
$$R_{pA} = \frac{dN_{J/\psi}/d^2P_{\perp} dy|_{pA}}{N_{\text{coll}} dN_{J/\psi}/d^2P_{\perp} dy|_{pp}} \quad N_{\text{coll}} = A^{1/3}$$

- expect stronger suppression at LHC than at RHIC



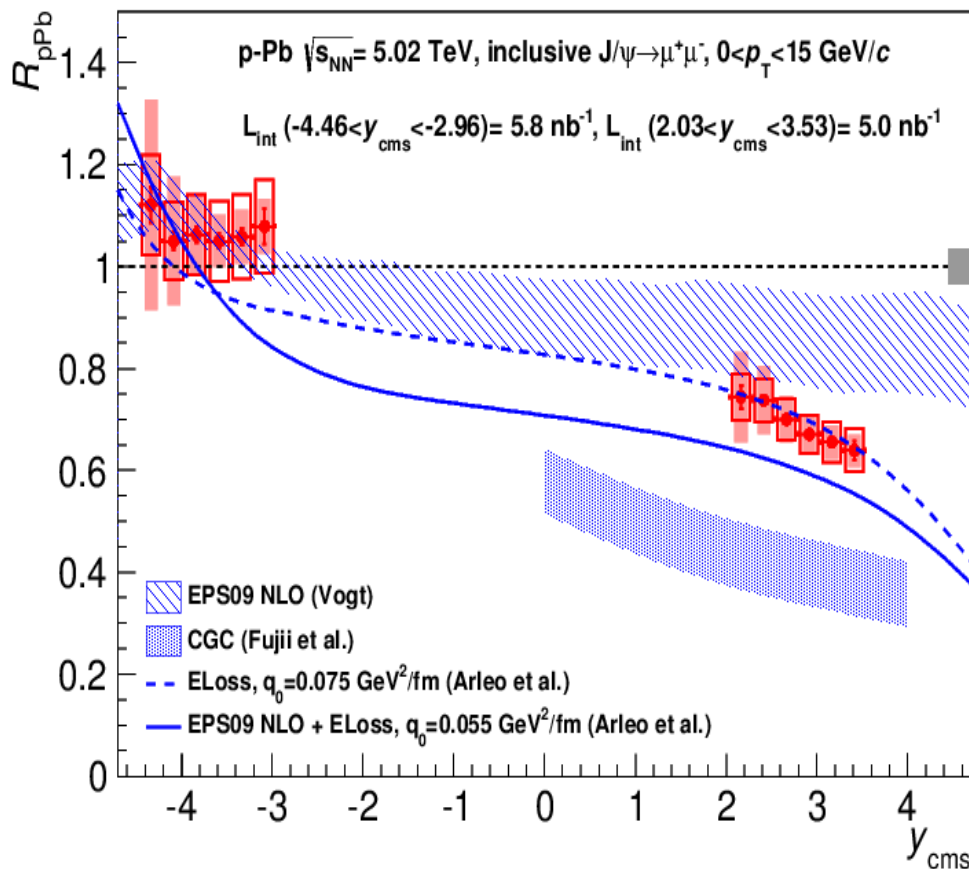
Comparison: quarkonium vs open charm

- Saturation is stronger at forward rapidity (smaller x)
- Multiple scattering increases the pair's invariant mass, which more suppresses quarkonium formation in CEM

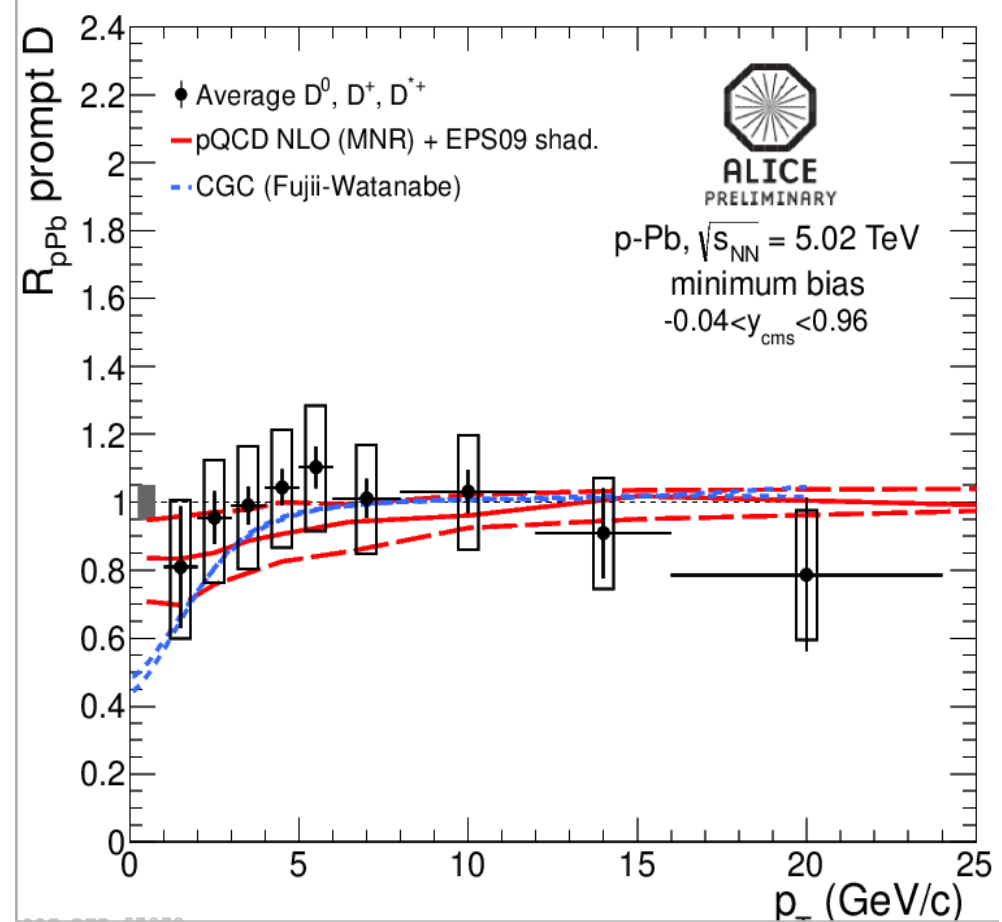


Comparison with data

ALICE, 1308.6726



ALICE, 1310.1714



Hadronization mechanism

- Octet and/or Singlet?
 - COM correlator dominates in CEM in large N
 - CSM picks up a new type of 4pt correlation
 - NRQCD provides a systematic framework

Kang et al, 1309.7337
 Qiu et al, 1310.2230
 HF-Watanabe, work in progress

$$\begin{aligned}
 C_{(xyuv)}^{\text{COM}} &= \frac{1}{4} \left(\begin{array}{c} \text{---} \rightarrow \text{---} \\ \text{---} \leftarrow \text{---} \\ \text{---} \rightarrow \text{---} \\ \text{---} \leftarrow \text{---} \end{array} \right) - \frac{1}{4N} \left(\begin{array}{c} \text{---} \rightarrow \text{---} \\ \text{---} \rightarrow \text{---} \\ \text{---} \leftarrow \text{---} \\ \text{---} \leftarrow \text{---} \end{array} \right) - \frac{1}{4N} \left(\begin{array}{c} \text{---} \leftarrow \text{---} \\ \text{---} \leftarrow \text{---} \\ \text{---} \rightarrow \text{---} \\ \text{---} \rightarrow \text{---} \end{array} \right) + \frac{1}{4N^2} \left(\begin{array}{c} \text{---} \rightarrow \text{---} \\ \text{---} \leftarrow \text{---} \\ \text{---} \rightarrow \text{---} \\ \text{---} \leftarrow \text{---} \end{array} \right) \\
 C_{(xyuv)}^{\text{CSM}} &= \frac{1}{2} \left(\begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} \right) - \frac{1}{2N} \left(\begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} \right)
 \end{aligned}$$

Summary & outlook

- QQbar production is expressed with known gluon distribution, constrained by rcBK phenomenology
- J/psi & D productions are obtained by CEM & FF
- RpA of J/psi is over-suppressed than ALICE data
- RpA of D is almost consistent with data
-

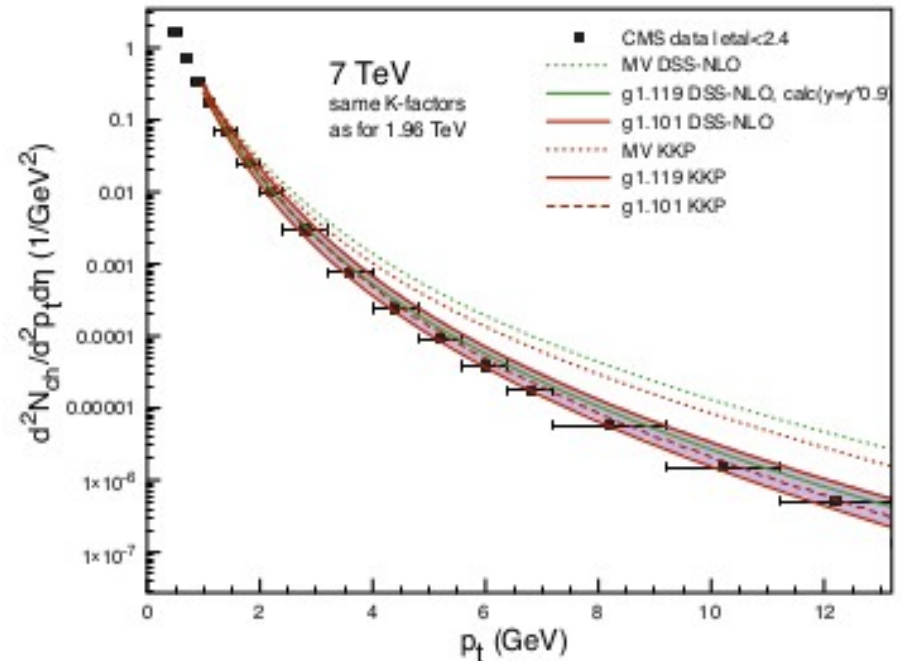
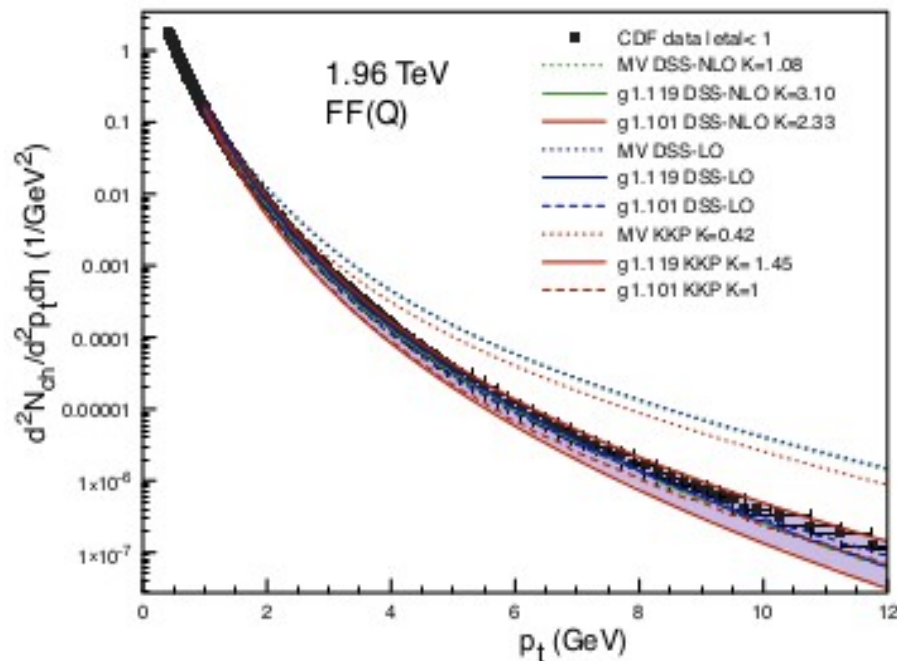
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- RpA of J/psi is over-suppressed than ALICE data
- RpA of D is almost consistent with data
- **Outlook**
 - LHC gives new information on J/psi production process
 - CGC+NRQCD; needs new 4pt gluon correlator
 - **NLO?** See Bowen's talk for NLO factorization in single hadron production

Hadron spectrum at pp colliders

- kT-factorized formula

- Normalized at $p_T=1$ GeV for $\sqrt{s}=1.96$ TeV
- uGD set ($\gamma \sim 1.1$) describes energy and p_T dependences



dN/dy in p-Pb collisions

ALICE:
Phys.Rev.Lett. 110 (2013) 032301

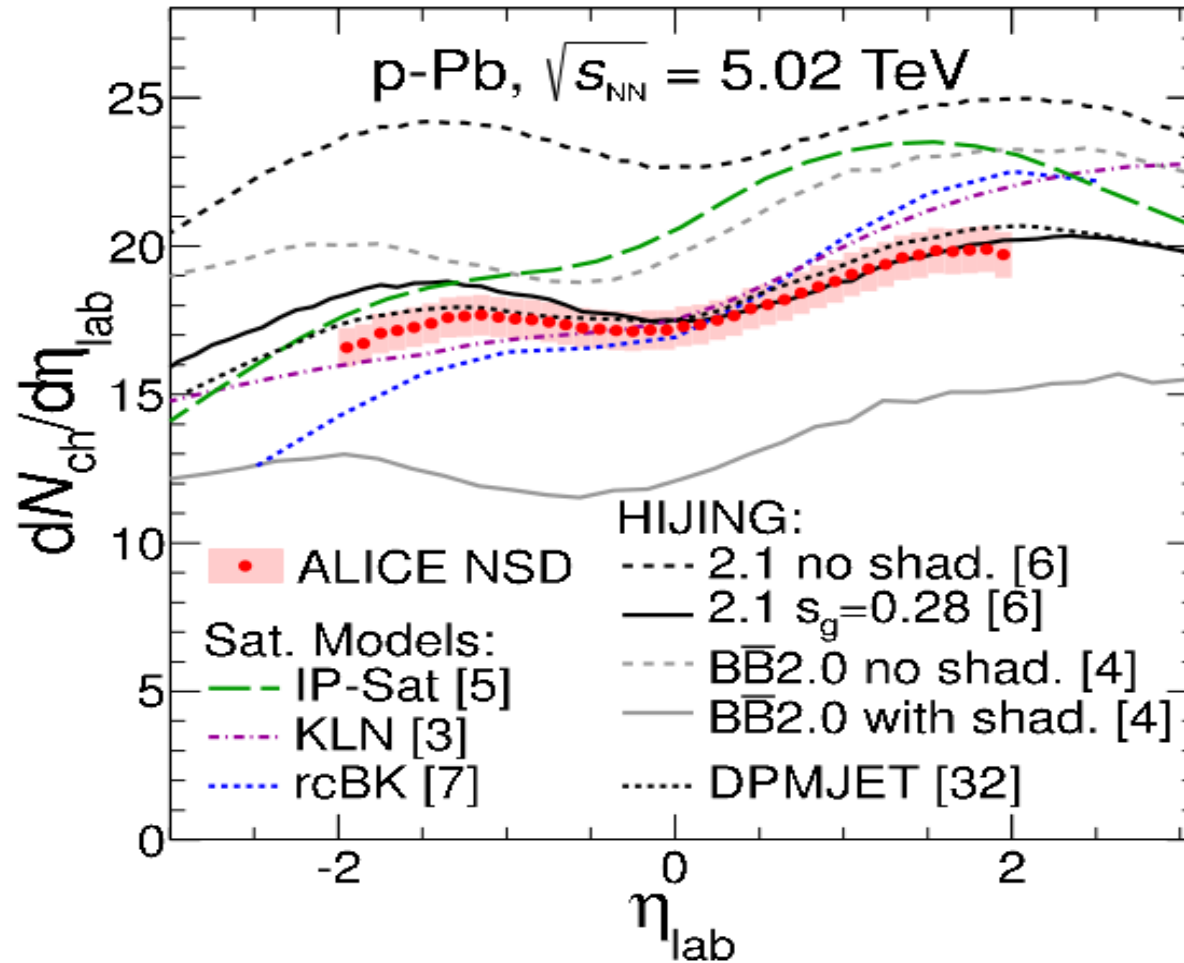


Fig. 1: Pseudorapidity density of charged particles measured in NSD p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV compared to theoretical predictions [3-7]. The calculations [4,5] have been shifted to the laboratory system.

Hadron production

- **INPUT:** gluon dist from rcBK in large N_c

$$\phi(k, y) \sim k^2 N_A(k, y), \quad 1 - N_A = (1 - N_F)^2$$

- **kT-factorization for small $x_{1,2}$ ($y \sim 0$)**

$$\begin{aligned} \frac{d\sigma^{A+B \rightarrow g}}{dy d^2 p_t d^2 R} &= K^k \frac{2}{C_F} \frac{1}{p_t^2} \int \frac{d^2 k_t}{4} \\ &\times \int d^2 b \alpha_s(Q) \varphi_P\left(\frac{|p_t + k_t|}{2}, x_1; b\right) \varphi_T\left(\frac{|p_t - k_t|}{2}, x_2; R - b\right) \end{aligned}$$

- **DHJ *hybrid* formula for small x_2 for $y > 0$**

$$\begin{aligned} \left[\frac{dN_h}{d\eta d^2 k} \right]_{\text{el}} &= \frac{1}{(2\pi)^2} \int_{x_F}^1 \frac{dz}{z^2} \left[\sum_q x_1 f_{q/p}(x_1, Q^2) \tilde{N}_F\left(x_2, \frac{k}{z}\right) D_{h/q}(z, Q^2) \right. \\ &\quad \left. + x_1 f_{g/p}(x_1, Q^2) \tilde{N}_A\left(x_2, \frac{k}{z}\right) D_{h/g}(z, Q^2) \right], \end{aligned}$$

Dumitru-Hayashigaki-Jalilian-Marian,
Nucl. Phys. A 765, 464 (2006).

Quark pair production in pA

- Formula at $O(\rho_p \rho_A^{00})$ in large N_c limit

$$\frac{dN_{q\bar{q}}}{d^2\mathbf{p}_\perp d^2\mathbf{q}_\perp dy_p dy_q} = \frac{1}{\pi R_A^2} \frac{\alpha_s^2 N_c}{8\pi^4 d_A} \frac{1}{(2\pi)^2} \int_{\mathbf{k}_{2\perp}, \mathbf{k}_\perp} \frac{\Xi(\mathbf{k}_{1\perp}, \mathbf{k}_{2\perp}, \mathbf{k}_\perp)}{k_{1\perp}^2 k_{2\perp}^2} \phi_{A,y_2}^{q\bar{q},g}(\mathbf{k}_{2\perp}, \mathbf{k}_\perp) \varphi_{p,y_1}(\mathbf{k}_{1\perp})$$

Gelis-Blaizot-Venugopalan
HF-Gelis-Venugopalan

- Pair production \rightarrow multi-parton correlators

$$\phi_A^{g,g}(\mathbf{l}_\perp) = \text{diagram 1}$$

$$\phi_A^{q\bar{q},g}(\mathbf{l}_\perp; \mathbf{k}_\perp) = \text{diagram 2}$$

$$\phi_A^{q\bar{q},q\bar{q}}(\mathbf{l}_\perp; \mathbf{k}_\perp, \mathbf{k}'_\perp) = \text{diagram 3}$$

- 4-pt & 3-pt functions simplify to a product of fundamental 2-pt funs in large N_c limit