Photons and Dileptons from Heavy-Ion Collisions and QCD: Microscopic and Bulk Dynamics

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YUKAWA INSTITUTE FOR THEORETICAL PHYSICS

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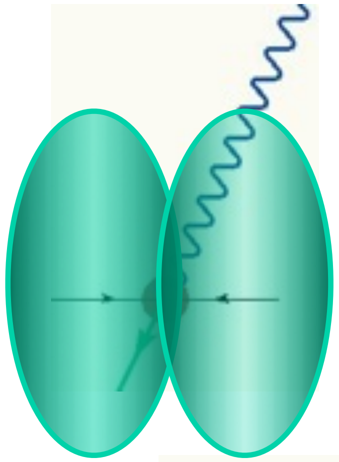
Charles Gale
McGill University

- Sources & EM emissivity: Rates
- Modelling the evolving system:
  - 3D hydro
  - 3D viscous hydro
  - Fluctuating initial states
- How are the photon yields dependent on the dynamics?
- Status of our interpretation of the data
- Photons as a characterization tool

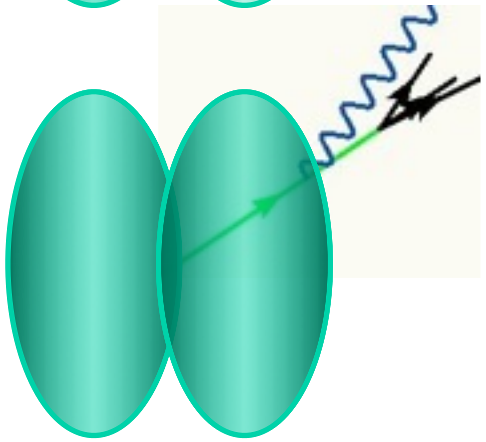
Why study photons and dileptons in relativistic nuclear collisions?

- Penetrating probes: negligible final state effects (α)
- Real and virtual photons are complementary, and they supplement hadronic observables
- Thermal photon emission rate favours hotter zones of the colliding system
- Emitted throughout the collision history
- Low emission rates
- Procedure: Calculate thermal emission rates & use hydrodynamics to model the evolution. Integrate rates over whole history

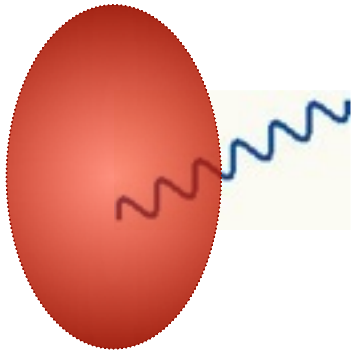
Sources of photons in a relativistic nuclear collision:



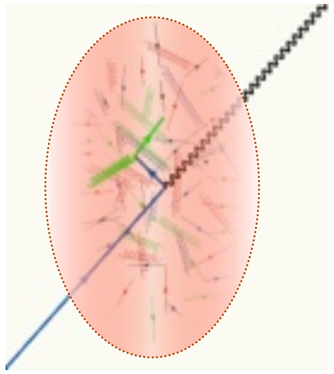
Hard direct photons. pQCD with shadowing
Non-thermal



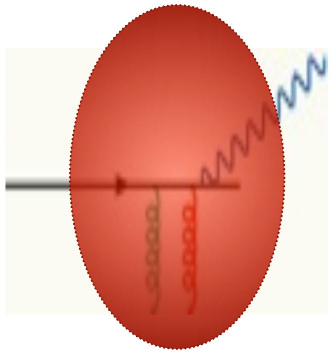
Fragmentation photons. pQCD with shadowing
Non-thermal



Thermal photons
Thermal



Jet-plasma photons
Thermal



Jet in-medium bremsstrahlung
Thermal

Pre-equilibrium?

INFO CARRIED BY THE RADIATION

$$dR = -\frac{g^{\mu\nu}}{2\omega}\frac{d^3k}{(2\pi)^3}\frac{1}{Z}\sum_i e^{-\beta K_i}\sum_f (2\pi)^4\delta(p_i - p_f - k)$$

$$\times \langle j \mid J_\mu \mid i\rangle\langle i \mid J_\nu \mid j\rangle$$

Thermal ensemble average of the current-current correlator

Emission rates:

$$\omega\frac{d^3R}{d^3k} = -\frac{g^{\mu\nu}}{(2\pi)^3}\mathrm{Im}\Pi^R_{\mu\nu}(\omega,k)\frac{1}{e^{\beta\omega}-1} \qquad \text{(photons)}$$

$$E_+E_-\frac{d^6R}{d^3p_+d^3p_-} = \frac{2e^2}{(2\pi)^6}\frac{1}{k^4}L^{\mu\nu}\,\mathrm{Im}\,\Pi^R_{\mu\nu}(\omega,k)\frac{1}{e^{\beta\omega}-1} \qquad \text{(dileptons)}$$

McLerran, Toimela (85), Weldon (90), Gale, Kapusta (91)

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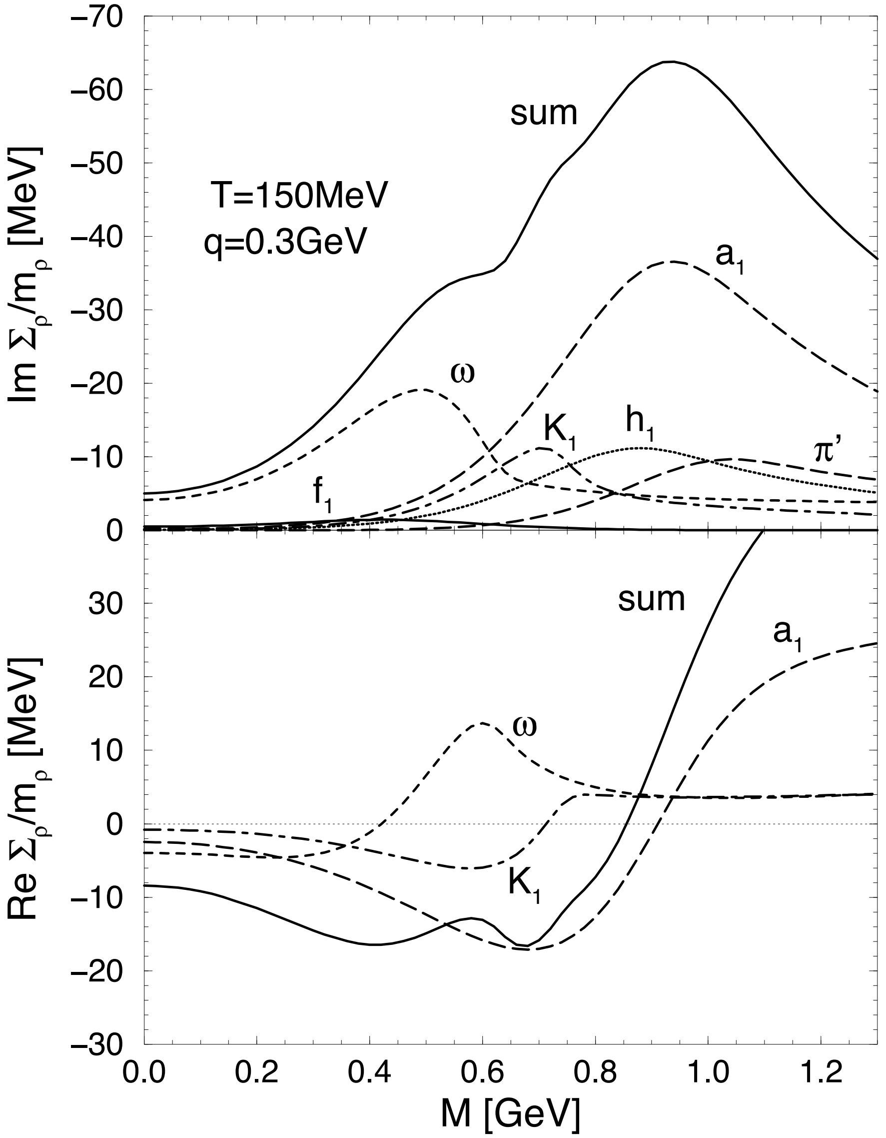
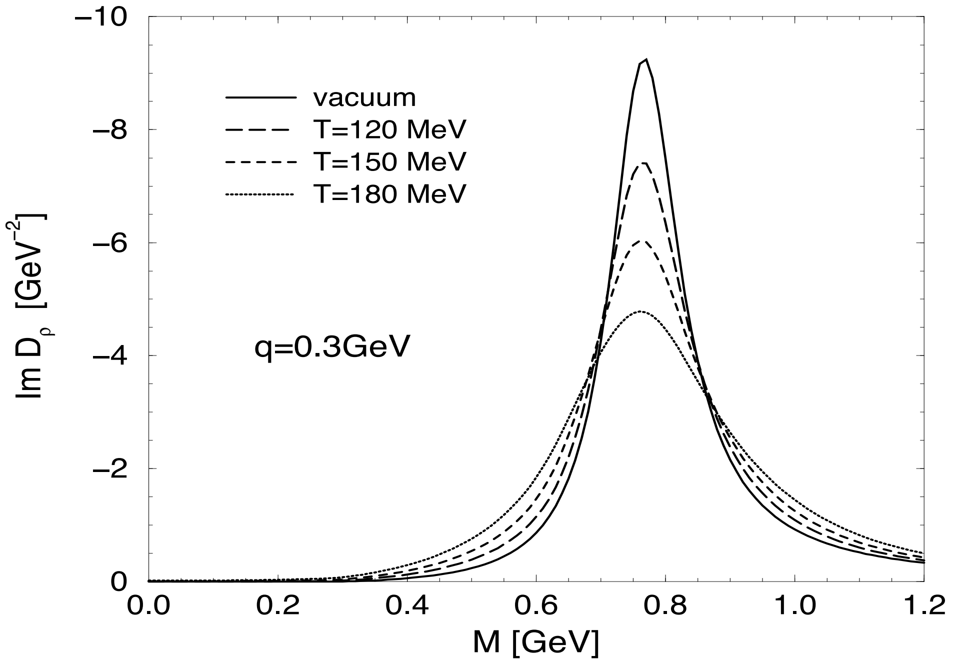
(FOR DILEPTONS:)

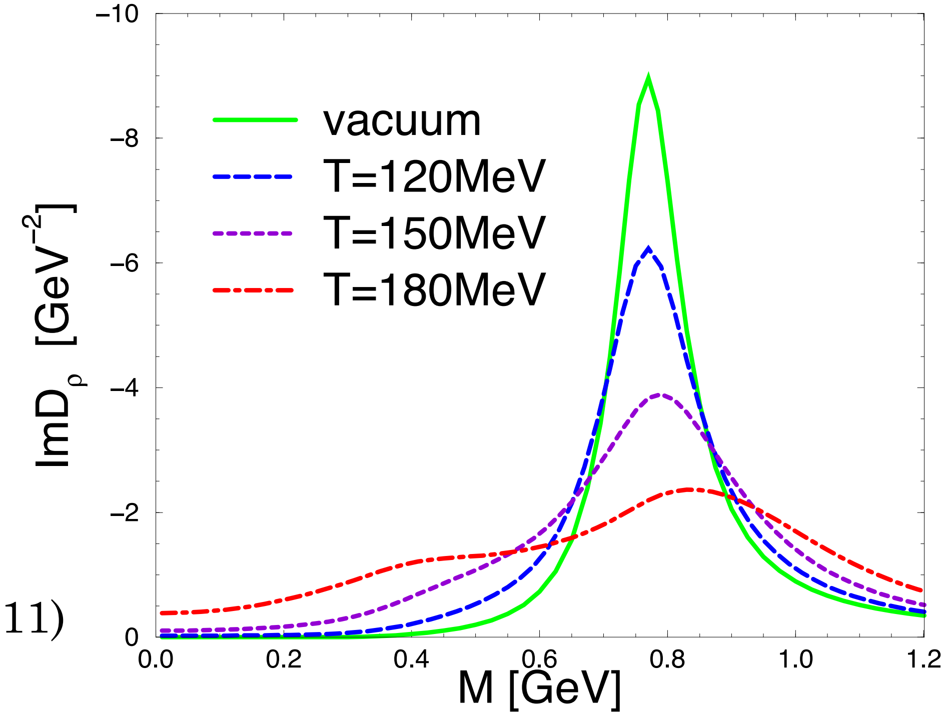
$$\text{Im} < J_\mu J_\nu >_T \overset{\text{VMD}}{\Rightarrow} \text{Im} < \rho_\mu \rho_\nu >_T \Rightarrow \text{Im}\, D^T_{\mu\nu} \Rightarrow \text{Vector spectral density}$$



R. Rapp, Act. Phys. Pol. (2011)

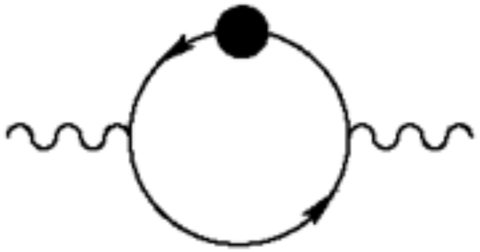


Rapp and Gale, PRC (1999)



$\rho_B / \rho_0 = 0.1, 0.7, 2.6$

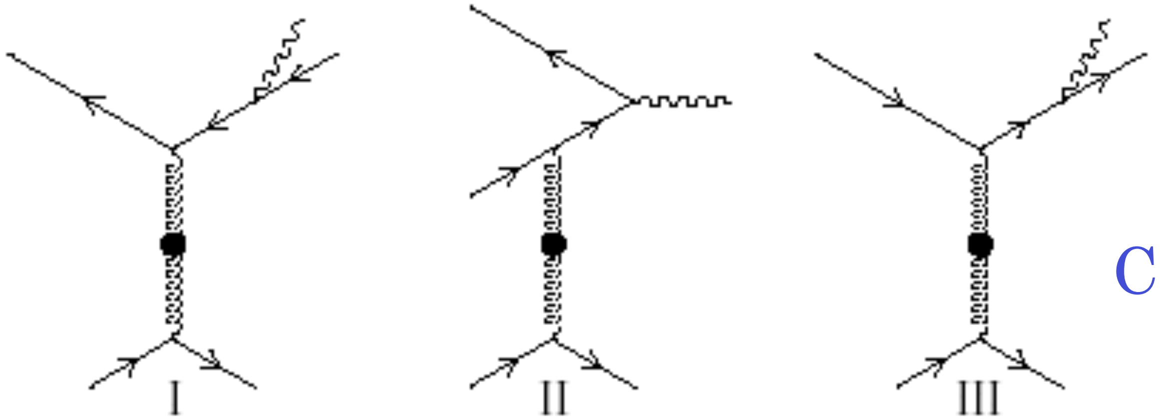
Thermal Photons from hot QCD: HTL program (Klimov (1981), Weldon (1982), Braaten & Pisarski (1990); Frenkel & Taylor (1990))



$$\mathrm{Im}\,\Pi^{\mu}_{R\mu} \sim \ln\left(\frac{\varpi T}{\left(m_{th}\left(\sim gT\right)\right)^{2}}\right)$$

Kapusta, Lichard, Seibert (1991)
Baier, Nakkagawa, Niegawa, Redlich (1992)

Going to two loops: Aurenche, Kobes, Gélis, Petitgirard (1996)
Aurenche, Gélis, Kobes, Zaraket (1998)



Co-linear singularities: $\alpha_s^{2}\left(\frac{T^{2}}{m_{th}^{2}}\right) \sim \alpha_s$

2001: Results complete at $O(\alpha_s)$

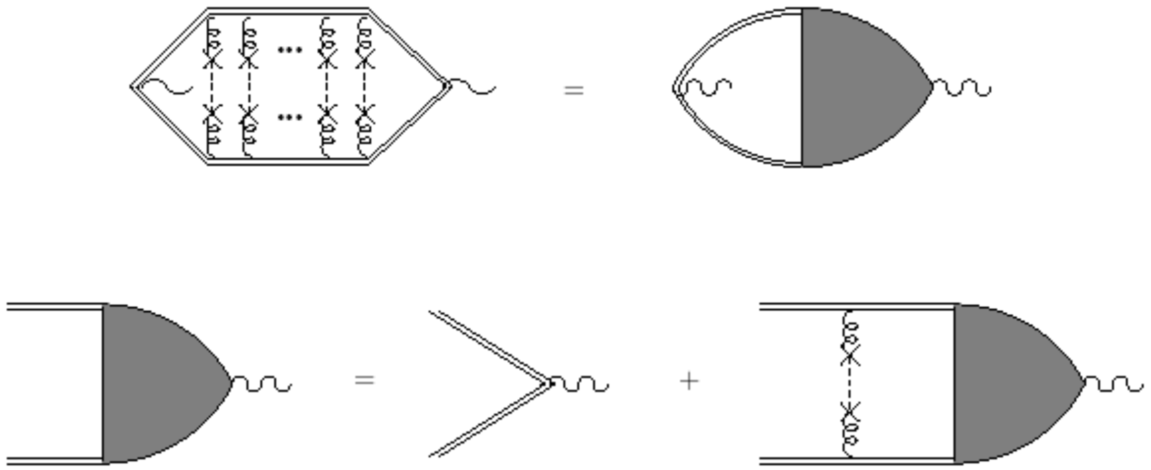
Arnold, Moore, and Yaffe JHEP **12**, 009 (2001); JHEP **11**, 057 (2001)
Incorporate LPM; Inclusive treatment of collinear enhancement, photon and gluon emission

Going beyond LO AMY rates?

- Approach is LO, but

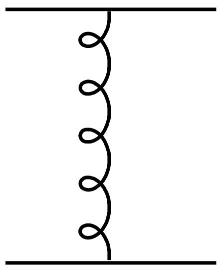
$$\alpha_s \sim 0.2 - 0.3$$

- Integral equation can be written in terms of a Dyson-Schwinger type iteration…



which contains a scattering kernel:

$$C(q_\perp)\big|_{\mathrm{LO}} = g^2 C_R T \frac{m_D^2}{q_\perp^2 (q_\perp^2 + m_D^2)}$$



Aurenche, Gélis, Zaraket (2002)

The techniques used to derive this - and all results in perturbative, finite-temperature field theory - rely on the scale separation:

$$gT \ll T$$

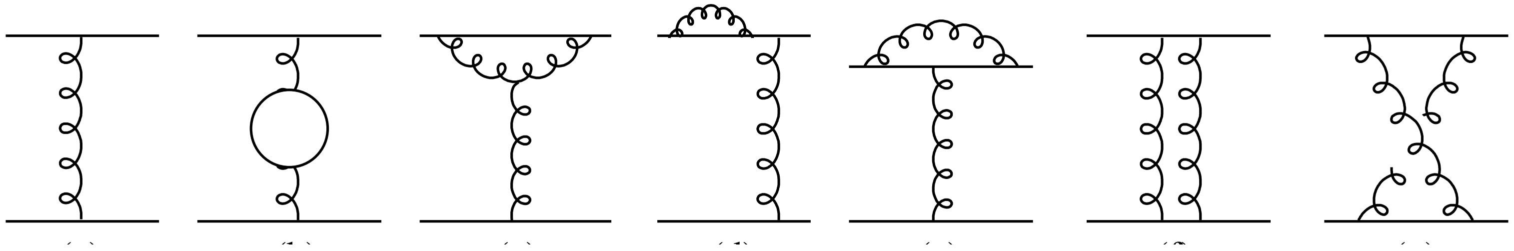
$$\text{soft} \ll \text{hard}$$

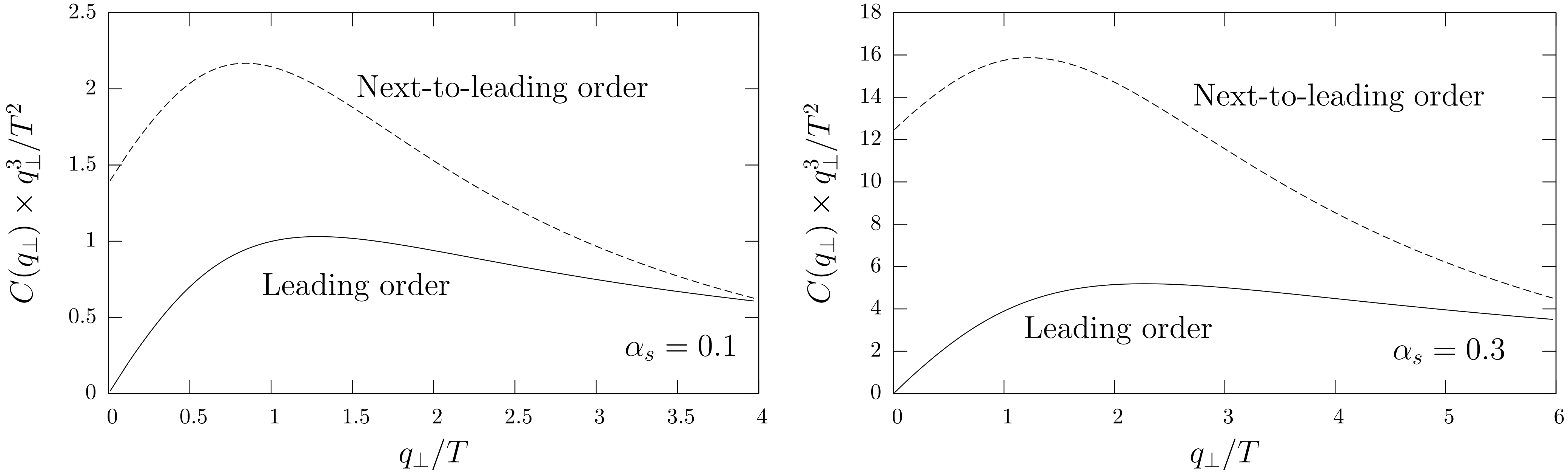
The LO-NLO scattering kernel(s)

Clue that NLO effects might be important: Heavy quark diffusion

$$C(q_\perp)\Big|_{\mathrm{LO}} \rightarrow C(q_\perp)\Big|_{\mathrm{NLO}}$$

Simon Caron-Huot PRD (2009)



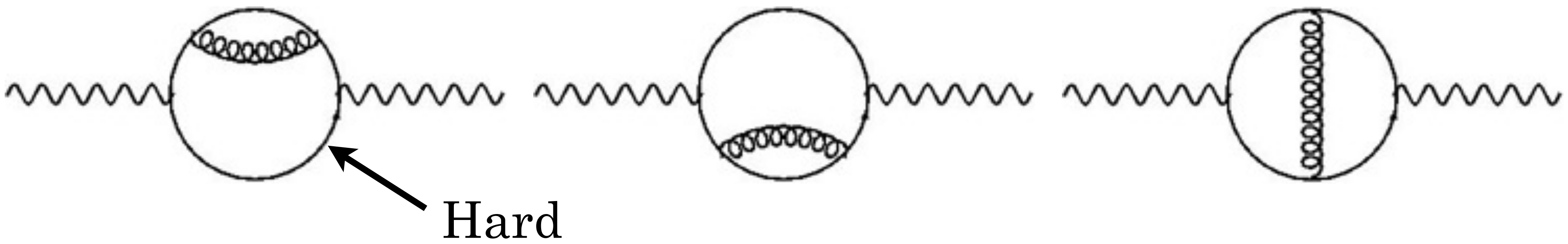


Possible large effects on photon production!?

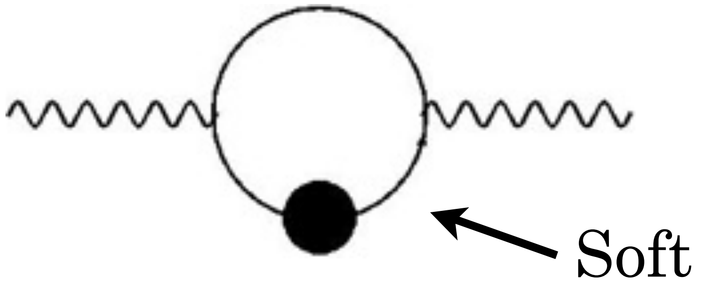
Photon emission at LO

$$\left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{LO}} = \left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{hard}} + \left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{soft}} + \left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{coll}}$$

$$\left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{hard}} \rightarrow$$

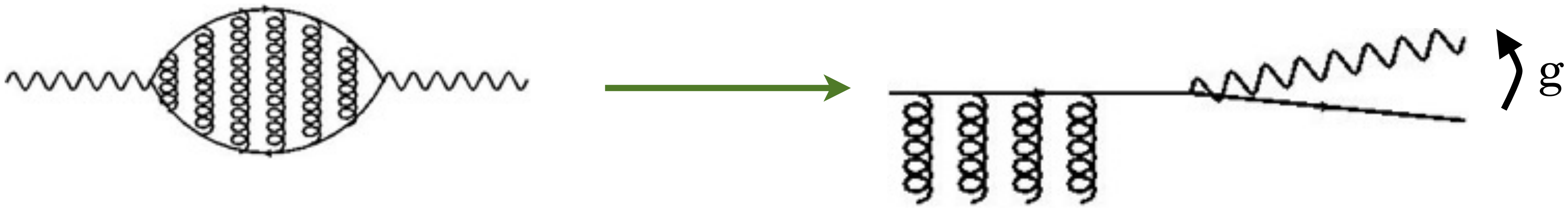


$$\left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{soft}} \rightarrow$$



(Conversion)

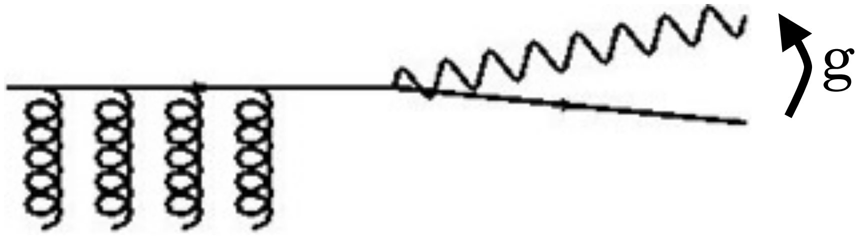
$$\left.\frac{d\Gamma_\gamma}{d^3k}\right|_{\mathrm{coll}} \rightarrow$$



The LO-NLO scattering kernels

Ghiglieri, Hong, Kurkela, Lu, Moore, Teaney, JHEP (2013)

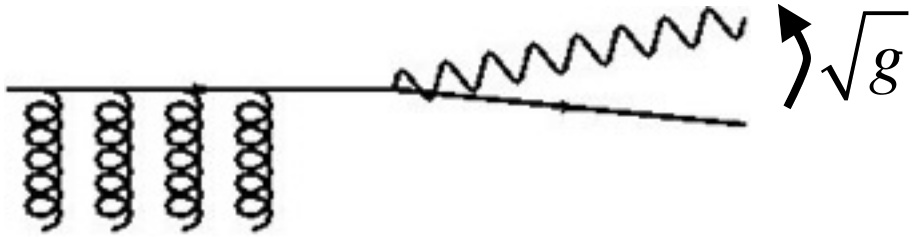
The two main contributions:



$$C(q_T)_{\mathrm{LO}} = \frac{Tg^2 m_D}{q_T(q_T + m_D)} \Rightarrow \mathrm{NLO}$$

Simon Caron-Huot PRD (2009)

Enhanced at NLO



Larger angle bremmstrahlung

Suppressed at NLO

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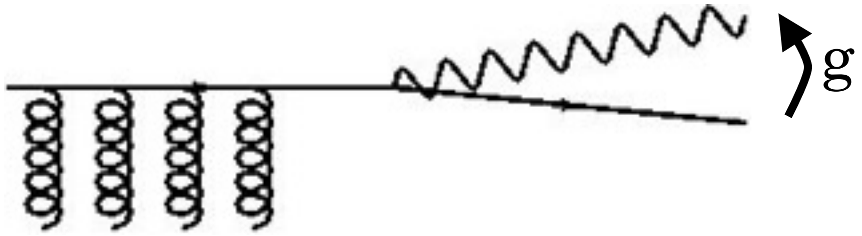

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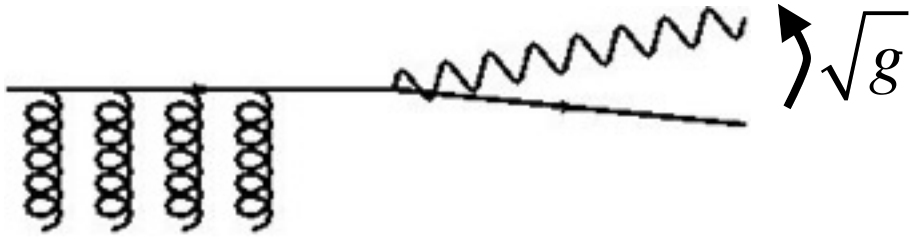

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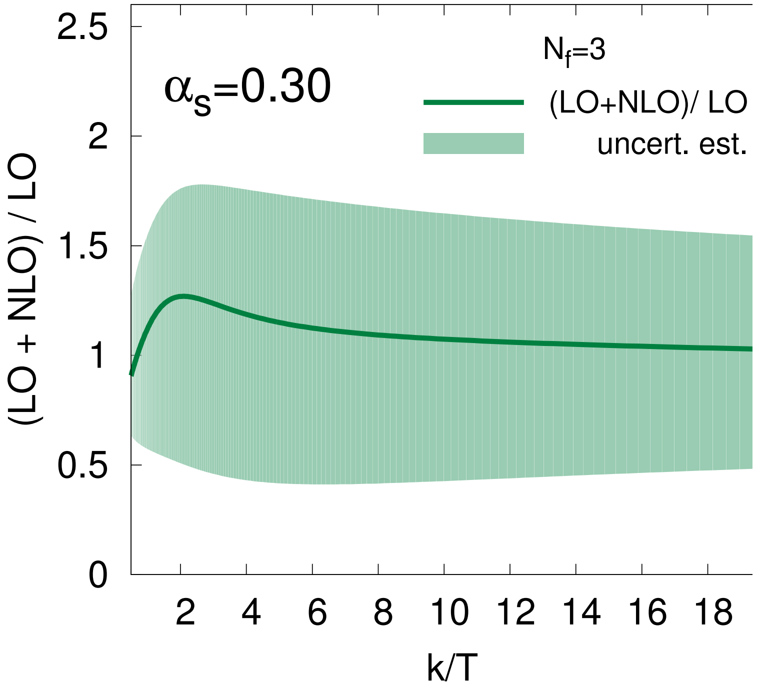
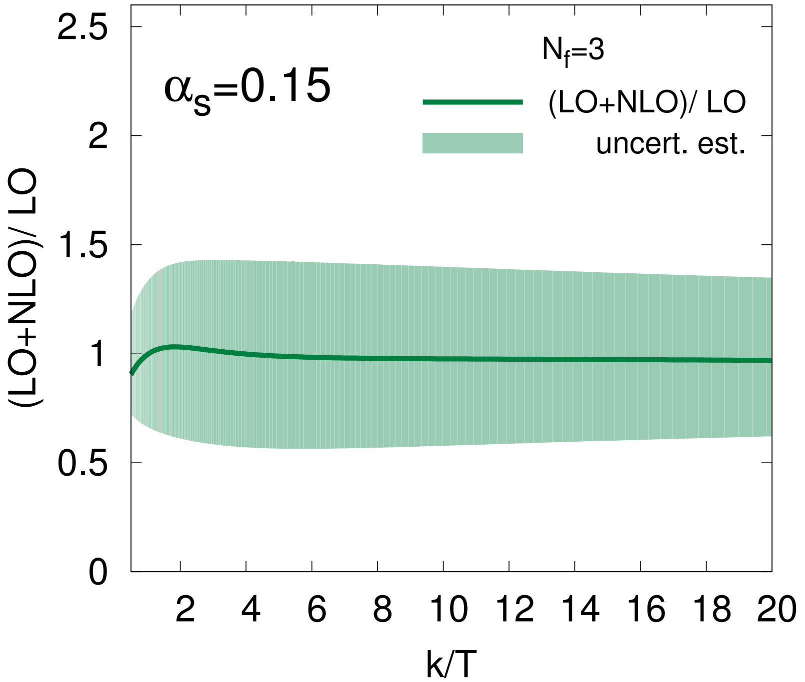




- Net correction to photon production rate is modest up to high k/T
- Techniques developed here have many more applications in FTFT

ELECTROMAGNETIC RADIATION FROM HADRONS

Chiral, Massive Yang-Mills:

O. Kaymakcalan, S. Rajeev, J. Schechter, PRD 30, 594 (1984)

$$L = \frac{1}{8} F_\pi^2 \,\mathrm{Tr}\, D_\mu U D^\mu U^\dagger + \frac{1}{8} F_\pi^2 \,\mathrm{Tr}\, \mathrm{M}\left(\mathrm{U} + \mathrm{U}^\dagger\right)$$

$$-\frac{1}{2}\mathrm{Tr}\left(F_{\mu\nu}^L F^{L\mu\nu} + F_{\mu\nu}^R F^{R\mu\nu}\right) + m_0^2 \mathrm{Tr}\left(A_\mu^\mathrm{L} A^{L\mu} + A_\mu^R A^{R\mu}\right)$$

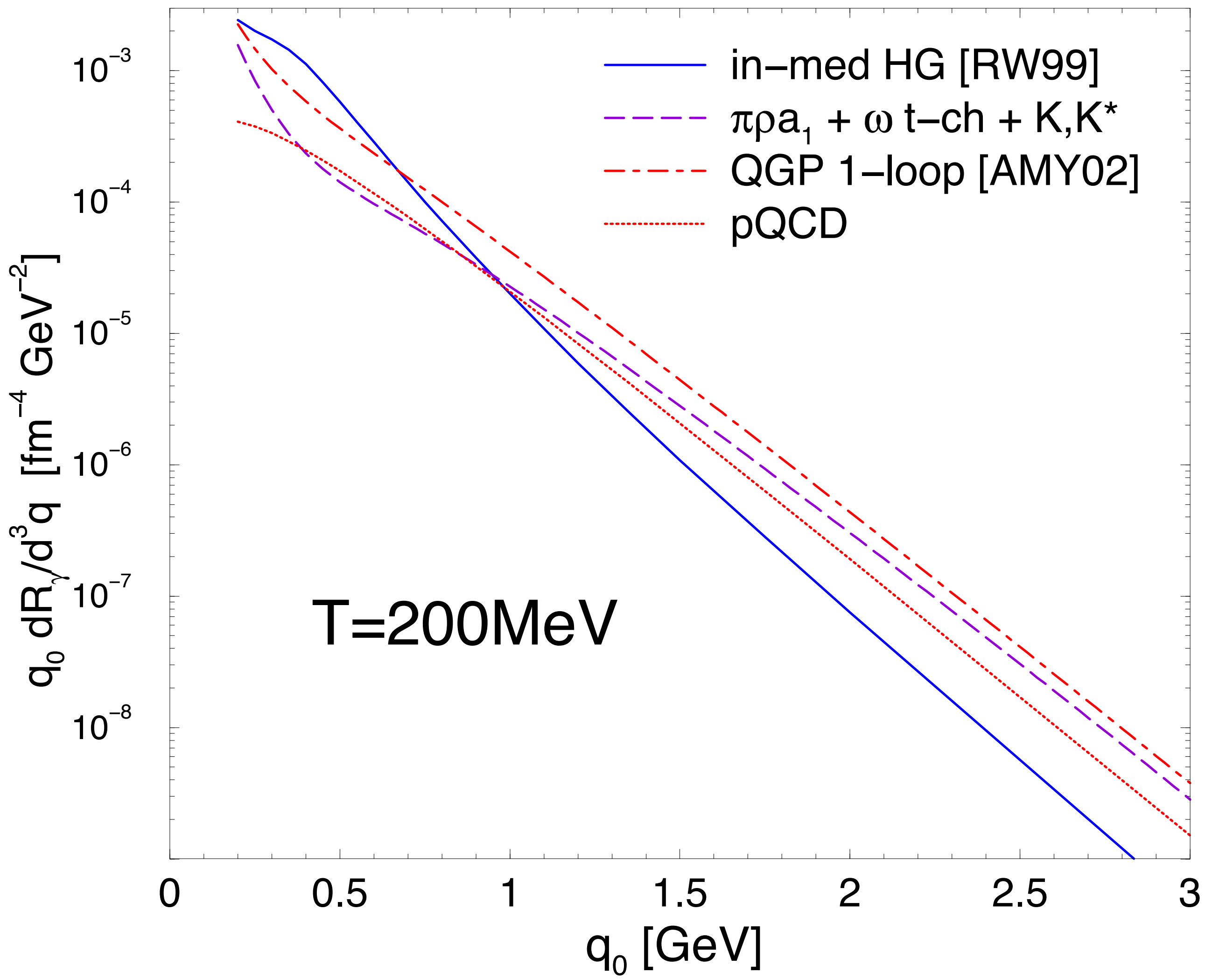
$$+ \text{ non-minimal terms}$$

Parameters and form factors are constrained by hadronic phenomenology:

- Masses & strong decay widths
- Electromagnetic decay widths
- Other hadronic observables:
  - *e.g.* $a_1 \to \pi\rho \quad D/S$ (See also, Lichard and Vojik, Nucl. Phys. (2010); Lichard and Juran, PRD (2008))

EM emissivities computed: Turbide, Rapp, Gale, PRC (2004); Turbide, McGill PhD (2006)

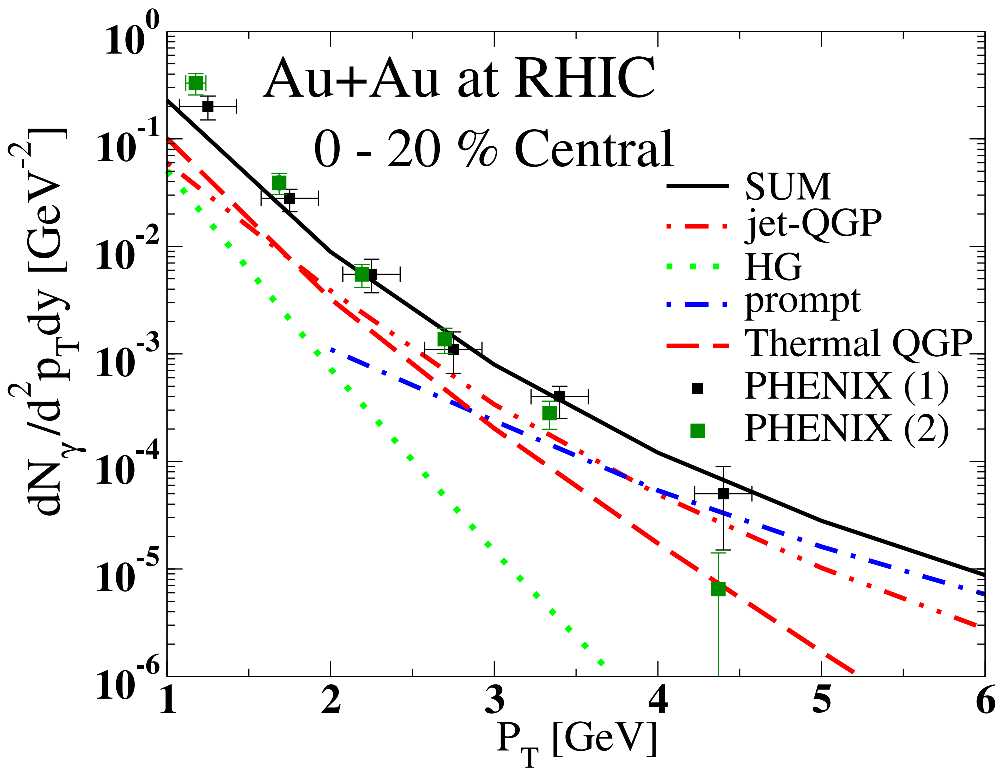
COMPARING RATES

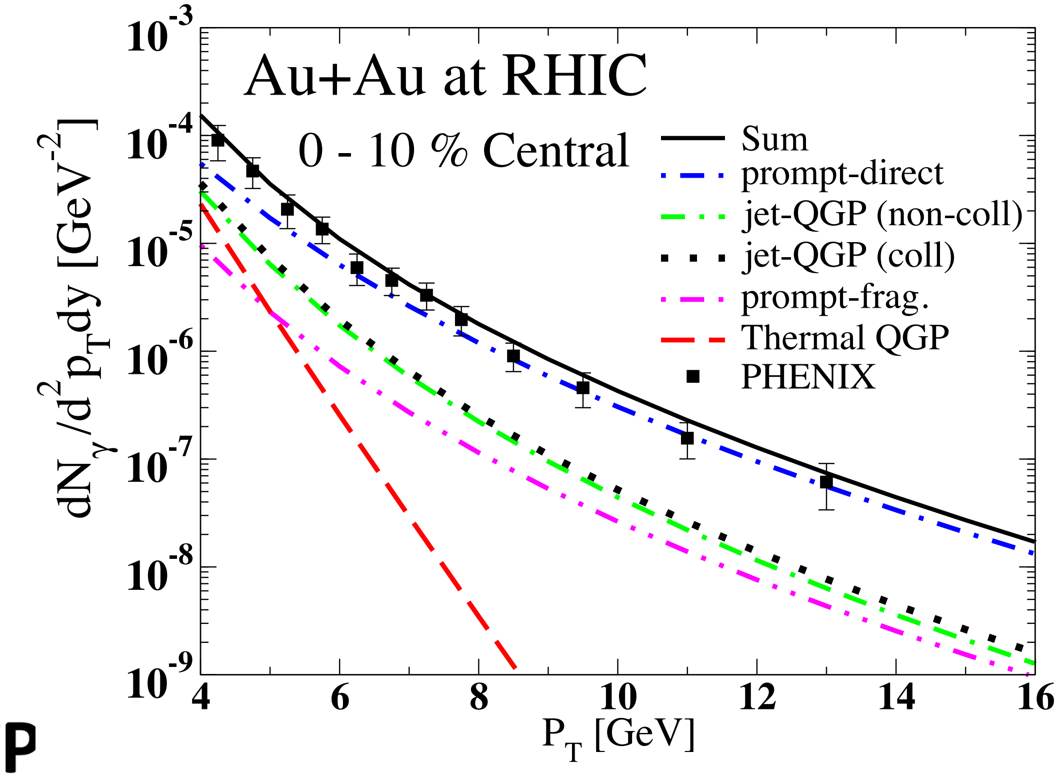


Turbide, Rapp, and Gale, PRC (2011)

Integrate rates with hydro evolution

APPLYING THIS TO INTERPRET PHOTONS MEASURED @ RHIC: RATES ARE INTEGRATED USING RELATIVISTIC HYDRODYNAMIC MODELLING

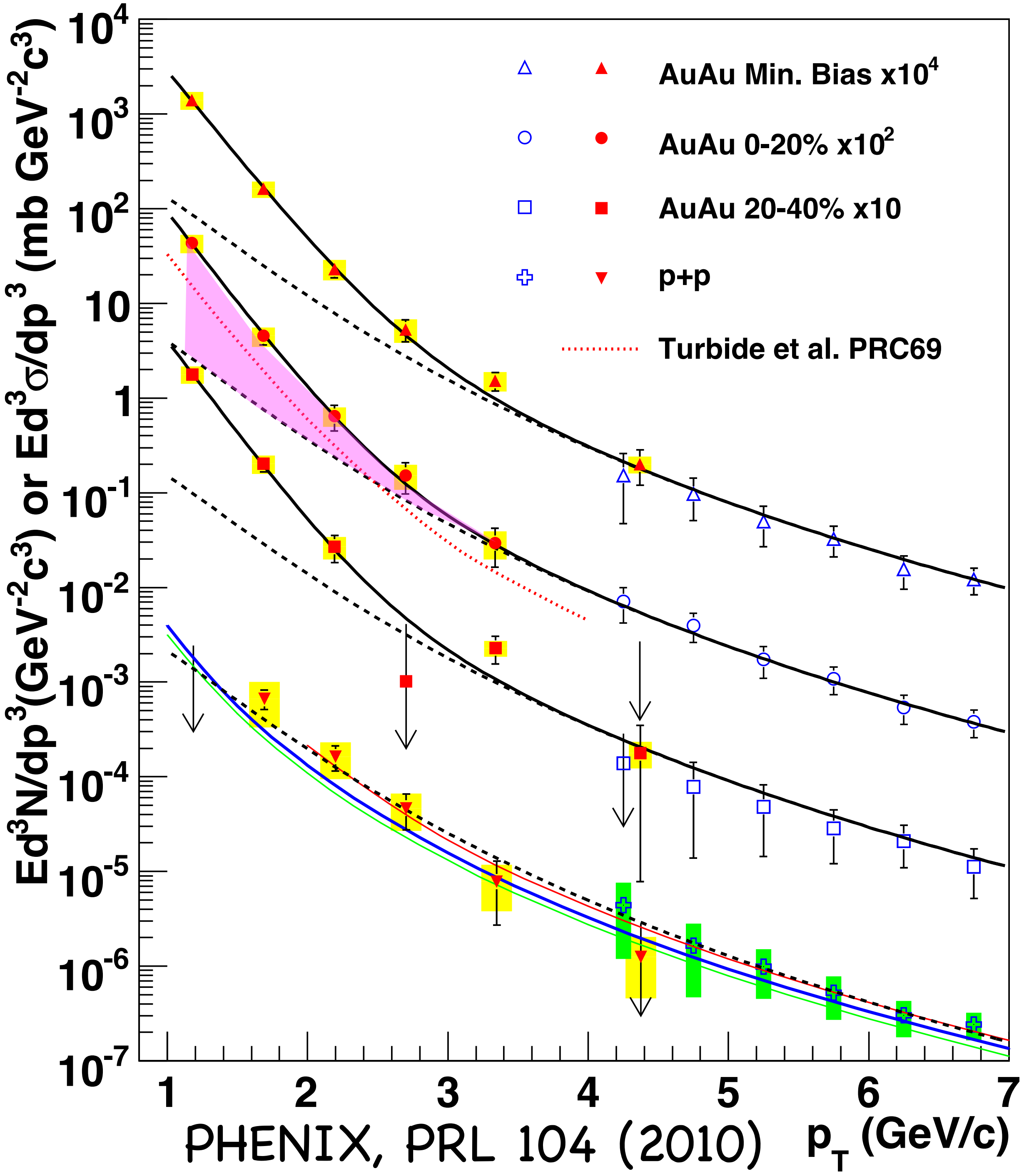




- At low p_T, spectrum dominated by thermal components (HG, QGP)
- At high p_T, spectrum dominated by pQCD
- Window for jet-QPG contributions at mid-p_T?

Turbide, Gale, Frodermann, Heinz, PRC (2008);
Higher p_T: G. Qin et al., PRC (2009)

ONE OF THE USES OF PHOTONS: CHARACTERIZING THE HOT MATTER CREATED AT RHIC



PHENIX, PRL 104 (2010)

$$T_{\text{excess}} = 221 \pm 19 \pm 19\,\text{MeV}$$

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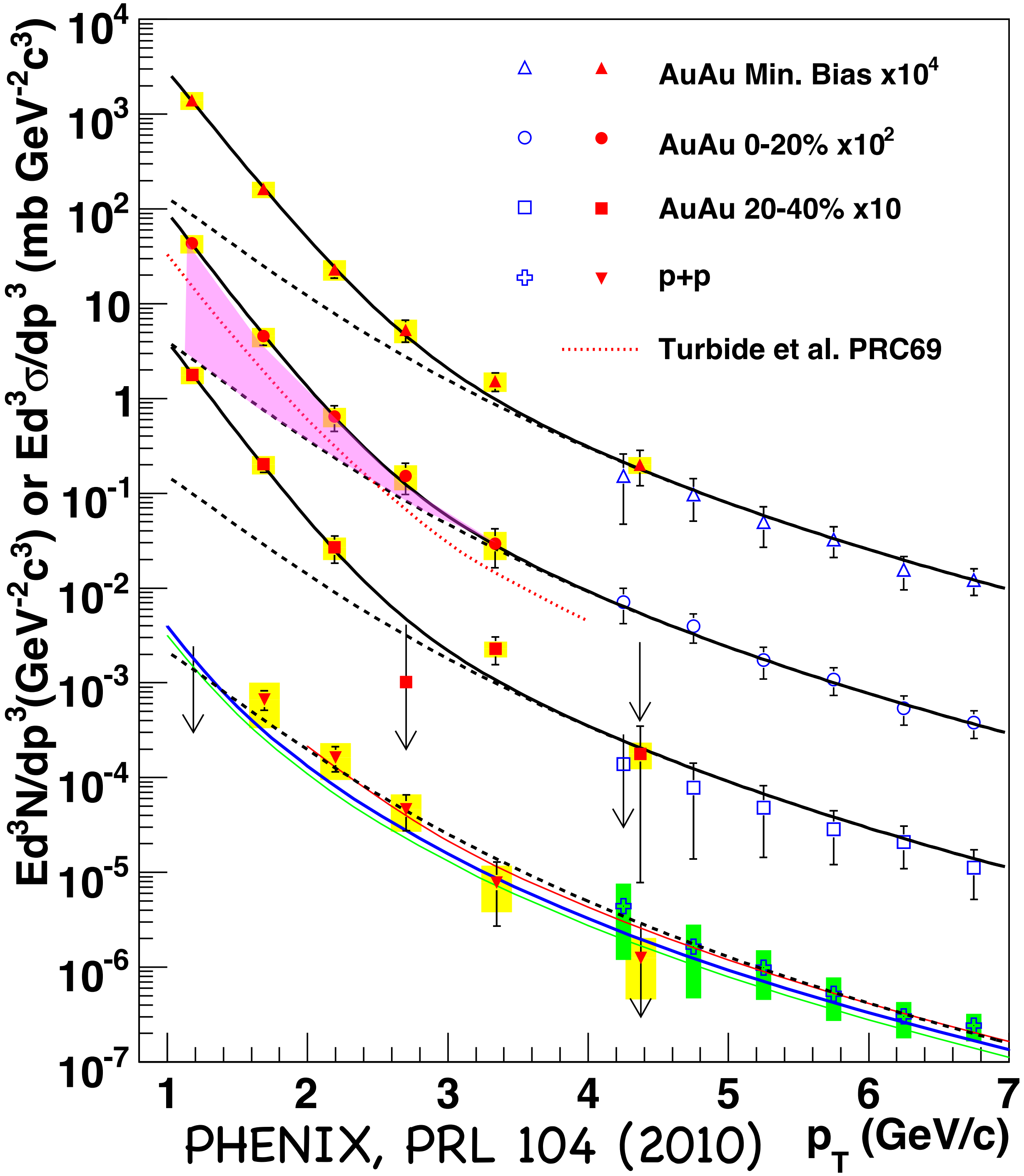


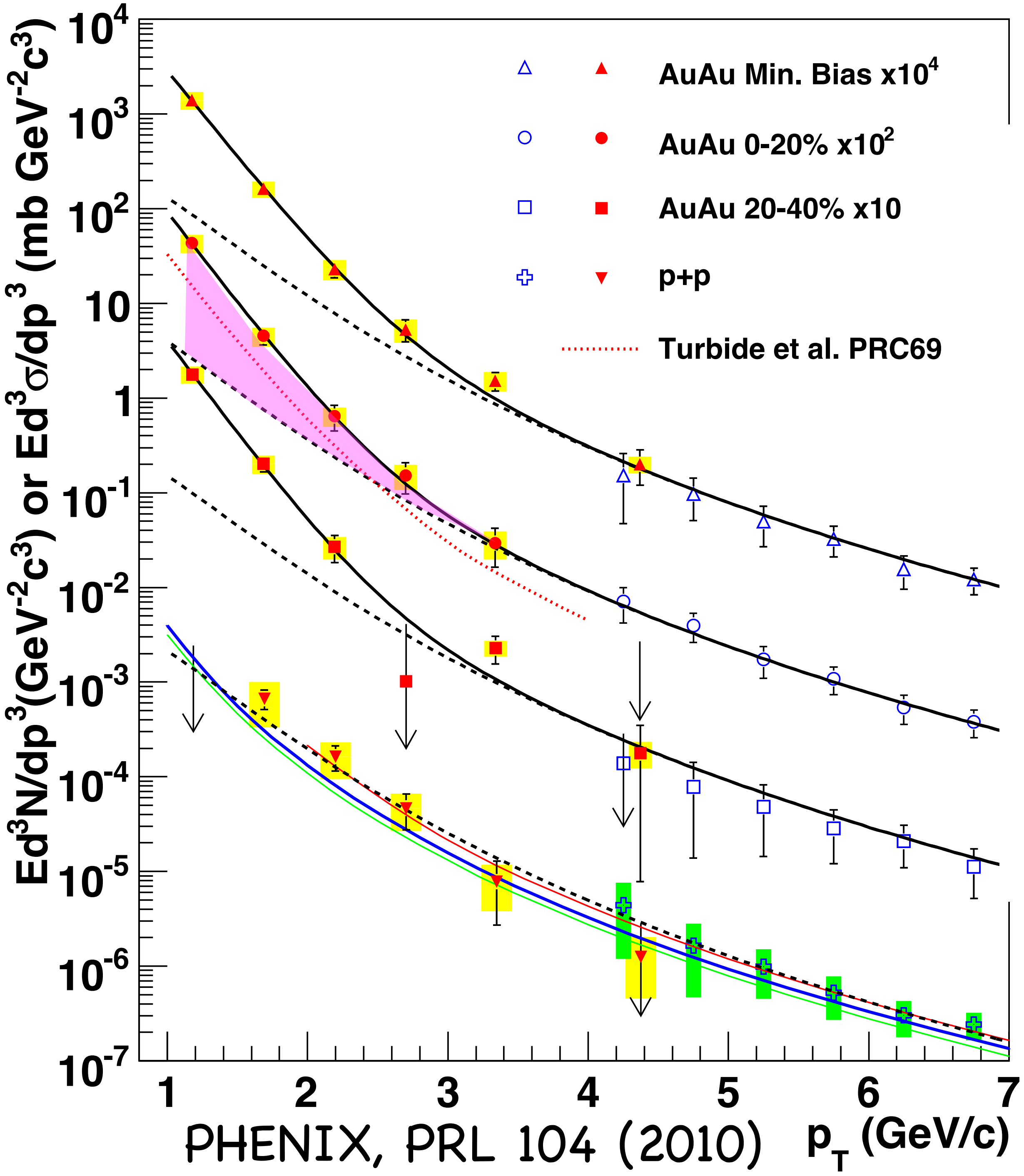
PHENIX, PRL 104 (2010)

$$T_{\mathrm{excess}} = 221 \pm 19 \pm 19\,\mathrm{MeV}$$

Flow effects will be important

- van Hees, Gale Rapp,PRC (2011)
- Shen, Heinz, Paquet, Gale, arXiv:1308.2440

ONE OF THE USES OF PHOTONS: CHARACTERIZING THE HOT MATTER CREATED AT RHIC



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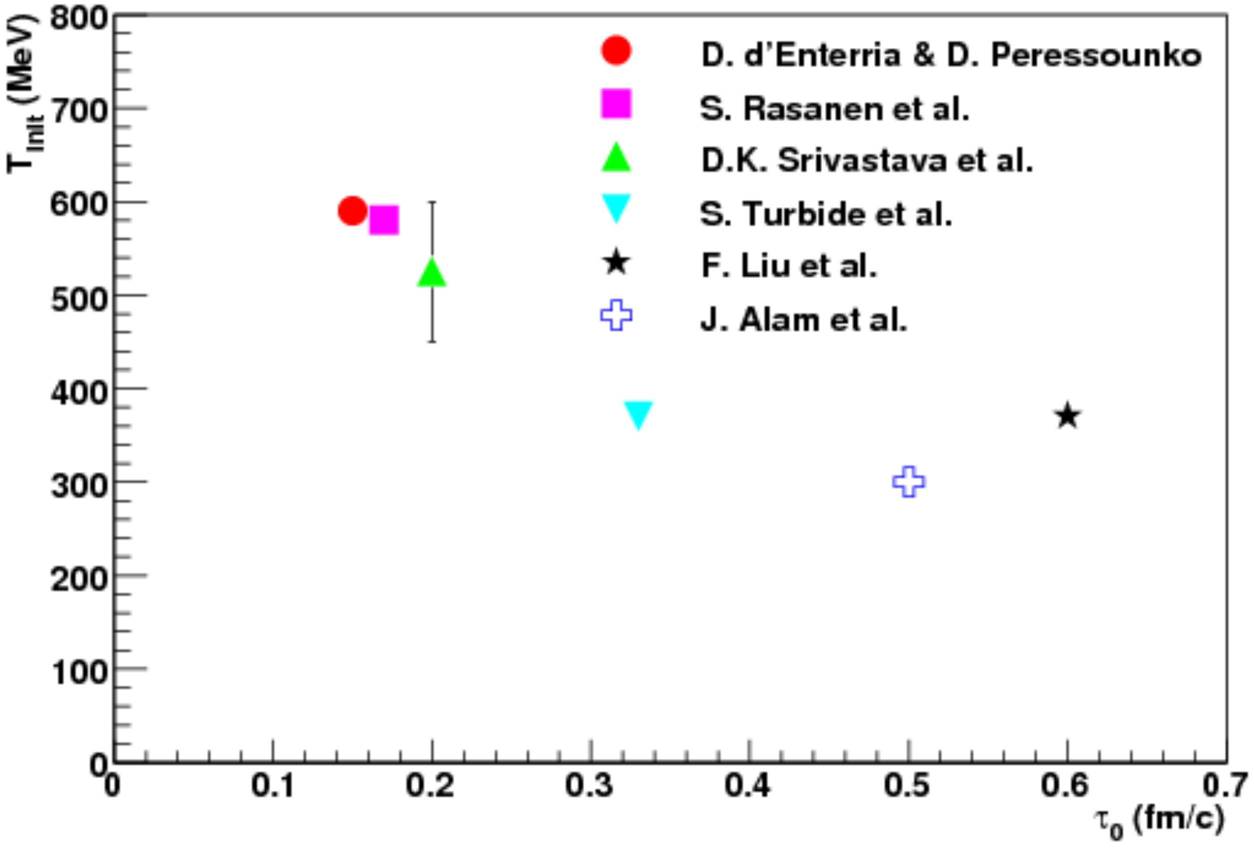
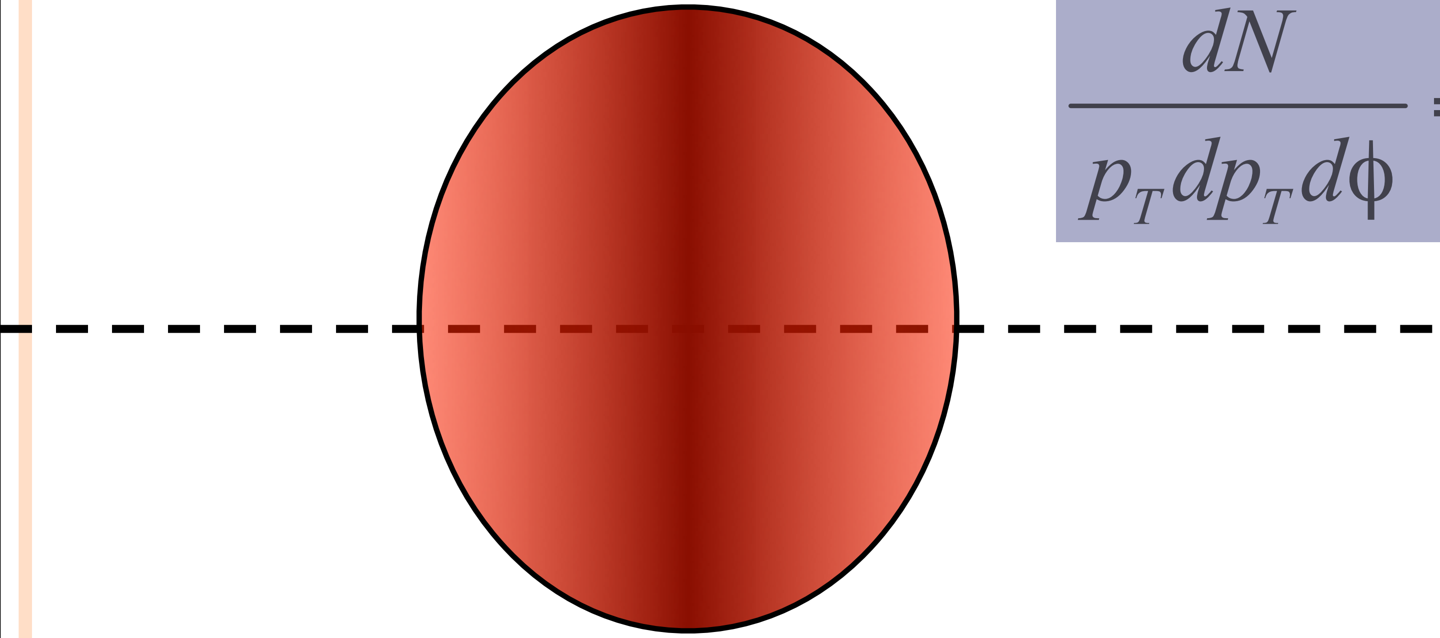
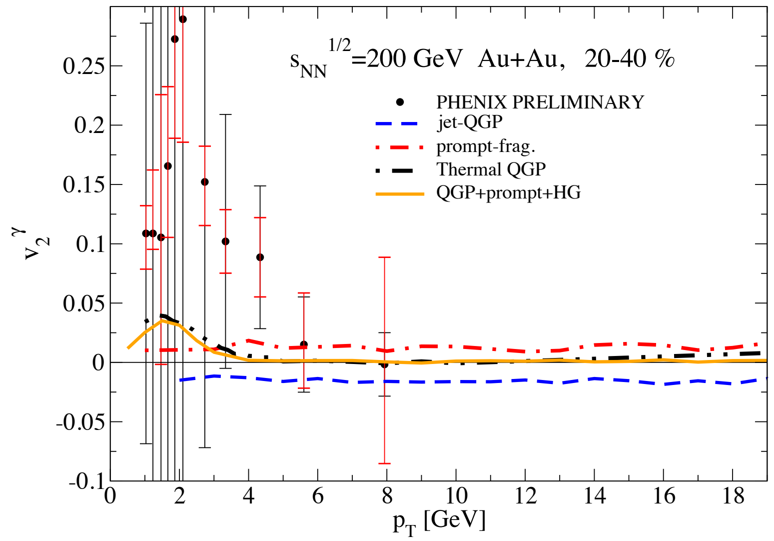


T_{ini} = 300 to 600 MeV
t_0 = 0.15 to 0.5 fm/c

BEYOND SIMPLE SPECTRA: FLOW AND CORRELATIONS

$$\frac{dN}{p_T\,dp_T\,d\phi} = \frac{dN}{2\pi\, p_T\,dp_T}\left[1+\sum_n 2v_n\cos(n\phi)\right]$$



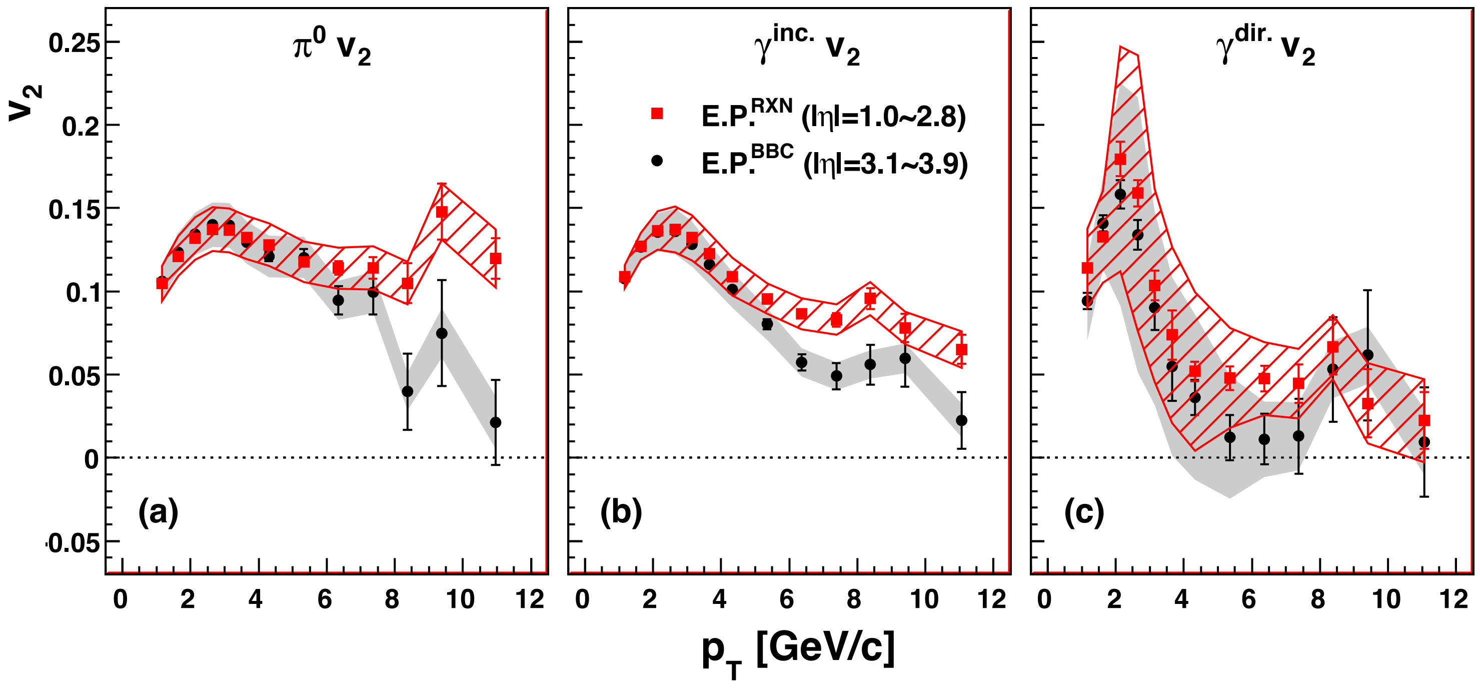


(2008)

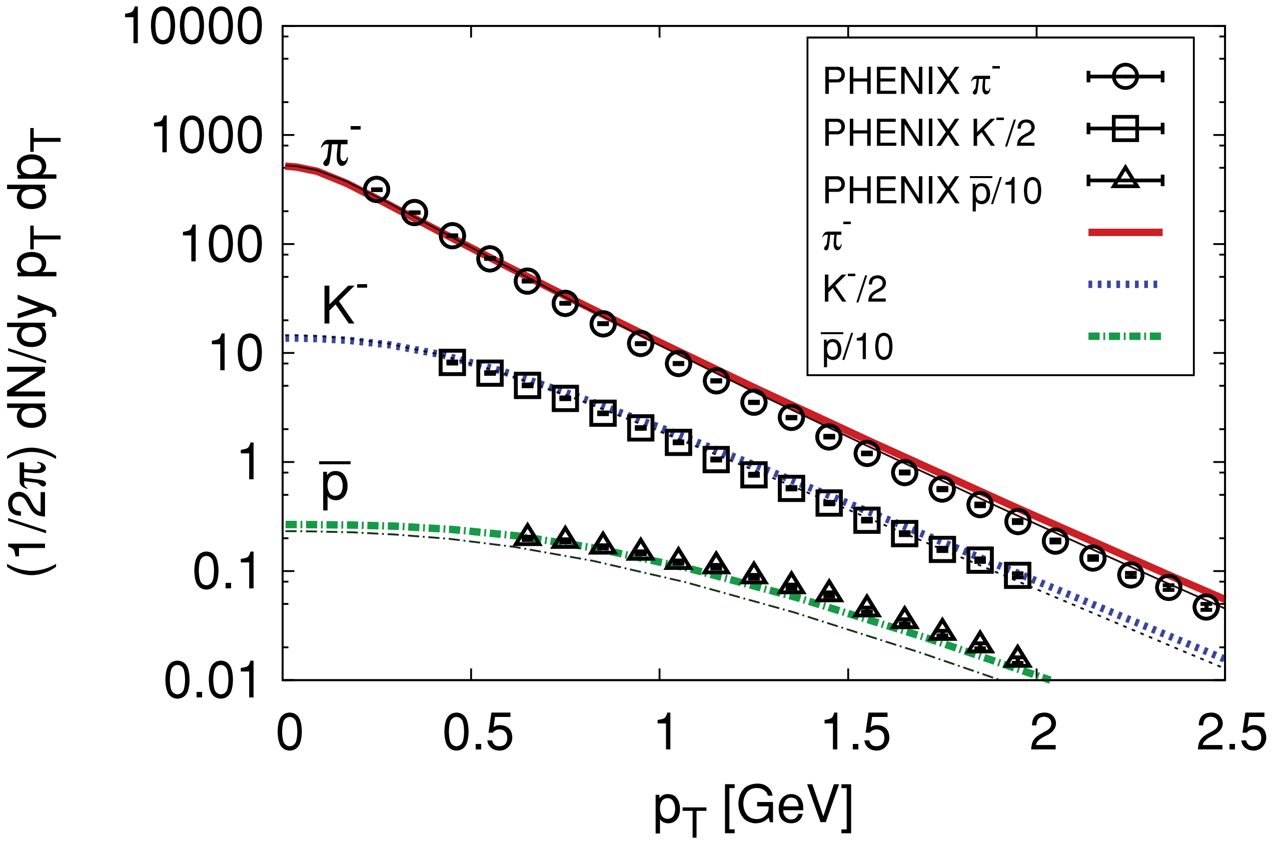
- Soft photons will go with the flow
- Jet-plasma photons: a negative v_2
- Details will matter: flow, T(t). . .

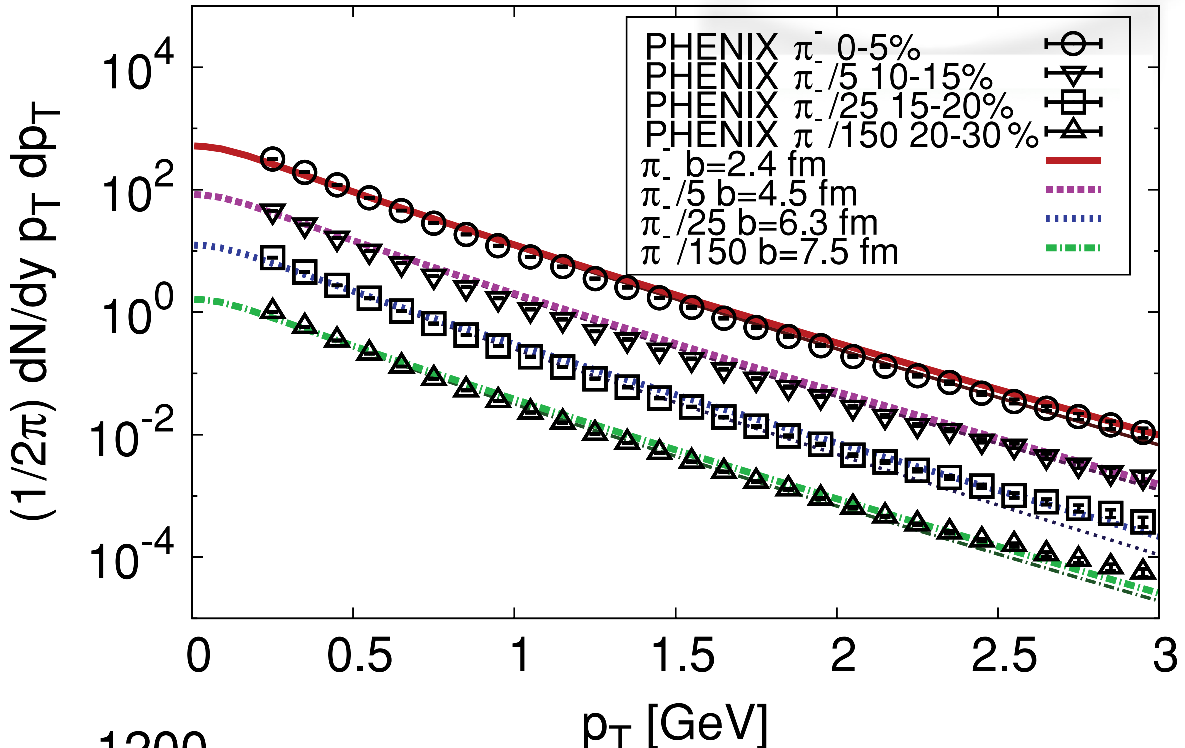
Turbide, Gale, Fries PRL (2006)
Low p_T: Chatterjee *et al.*, PRL (2006)
All p_T: Turbide *et al.*, PRC (2008)

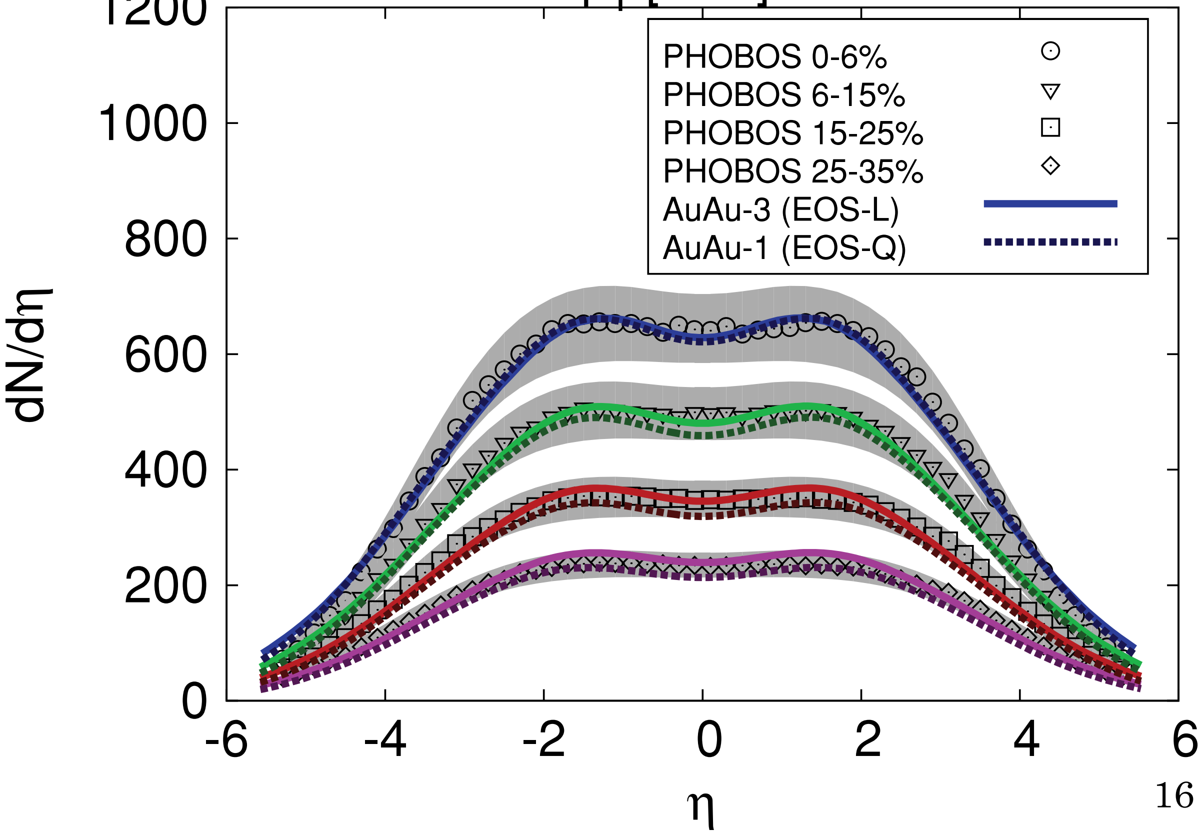
(2011)



PROGRESS IN CHARACTERIZATION TOOL: MUSIC:
3D VISCOUS RELATIVISTIC HYDRODYNAMICS







- MUSIC: 3D relativistic hydro
  - Ideal: Schenke, Jeon, and Gale, PRC (2010)
  - FIC and Viscous: Schenke, Jeon, Gale, PRL (2011)

Viscosity & FIC effects on EM observables?

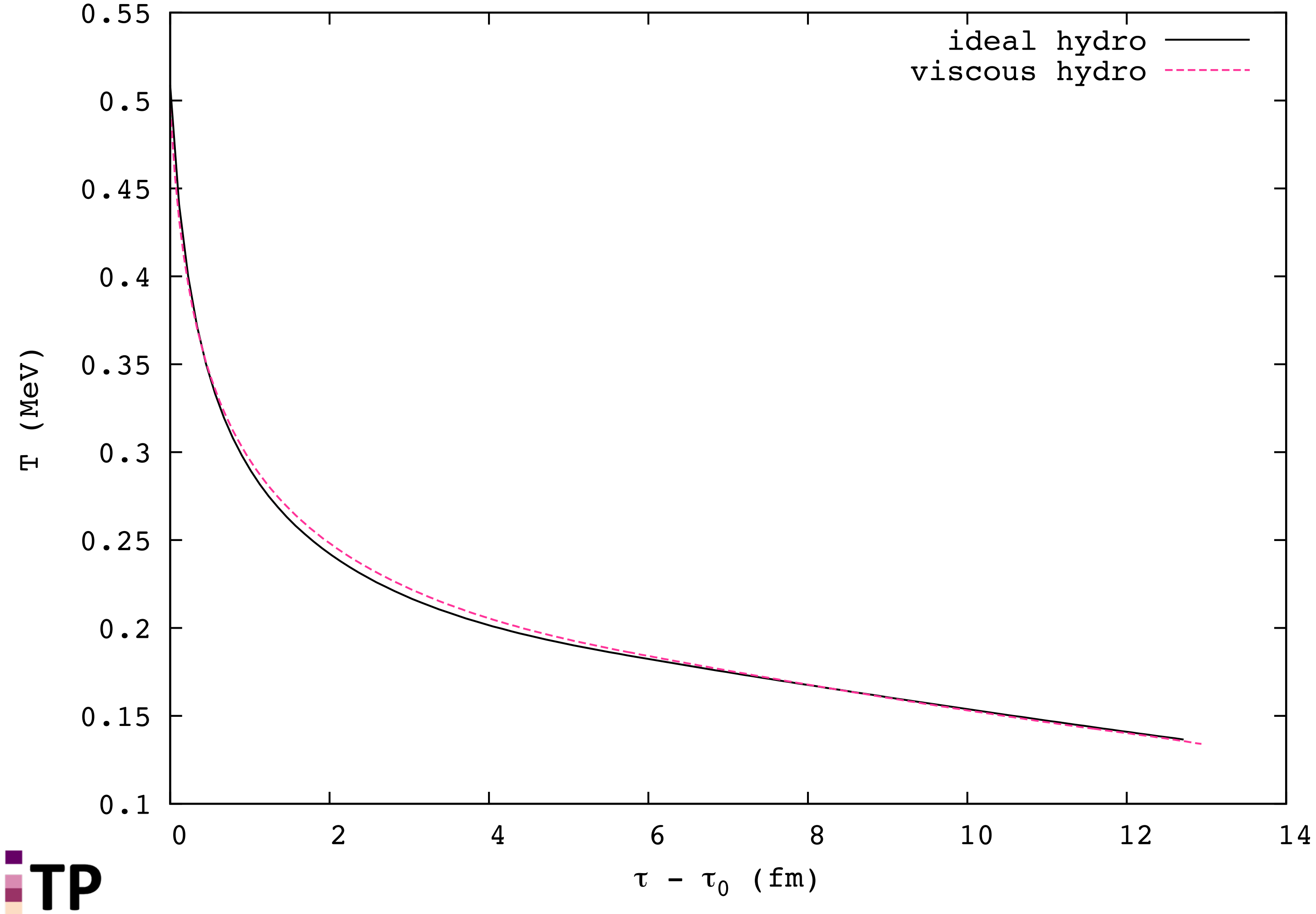
THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS

$$T^{\mu\nu}_{\text{ideal}} = (\varepsilon + P)u^\mu u^\nu - P g^{\mu\nu}$$

$$T^{\mu\nu} = T^{\mu\nu}_{\text{ideal}} + \pi^{\mu\nu}$$

Israël & Stewart, Ann. Phys. (1979), Baier et al., JHEP (2008), Luzum and Romatschke, PRC (2008)

$$\partial_\mu (s u^\mu) \propto \eta$$



- Viscous evolution starts with a lower T
- T drop is slower than ideal case

THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS

$$T^{\mu\nu}_{\text{ideal}} = (\varepsilon + P)u^{\mu}u^{\nu} - Pg^{\mu\nu}$$

$$T^{\mu\nu} = T^{\mu\nu}_{\text{ideal}} + \pi^{\mu\nu}$$

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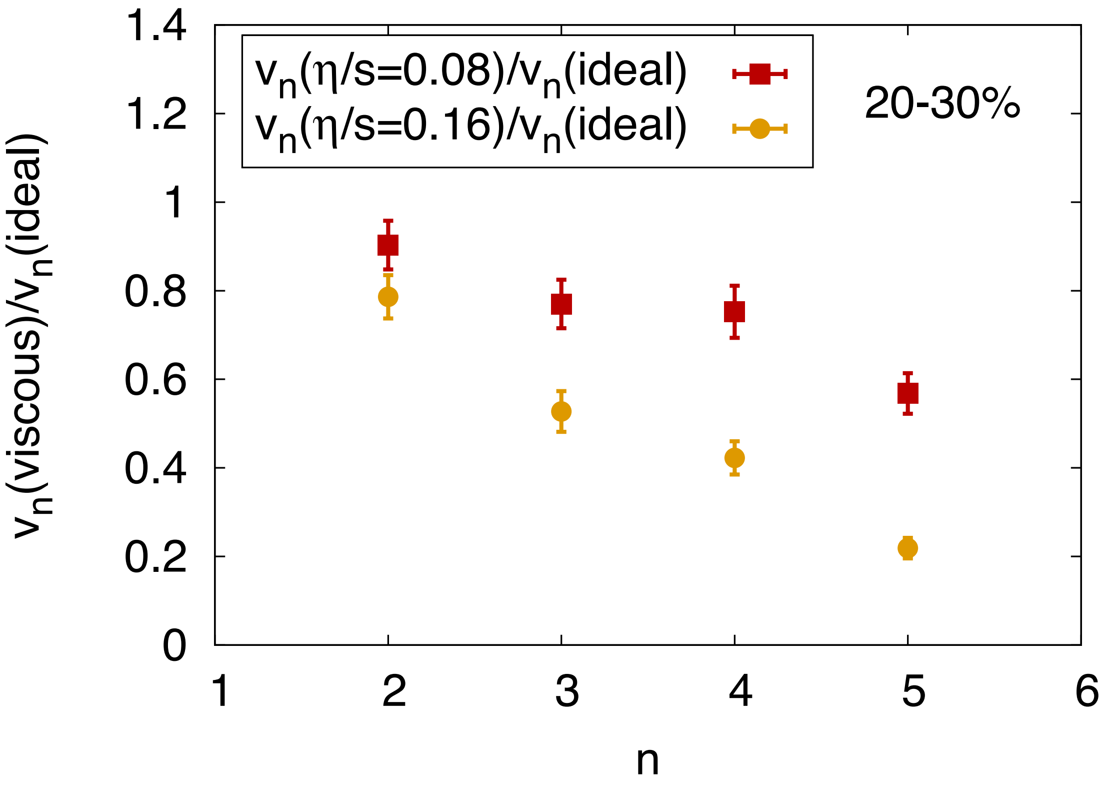
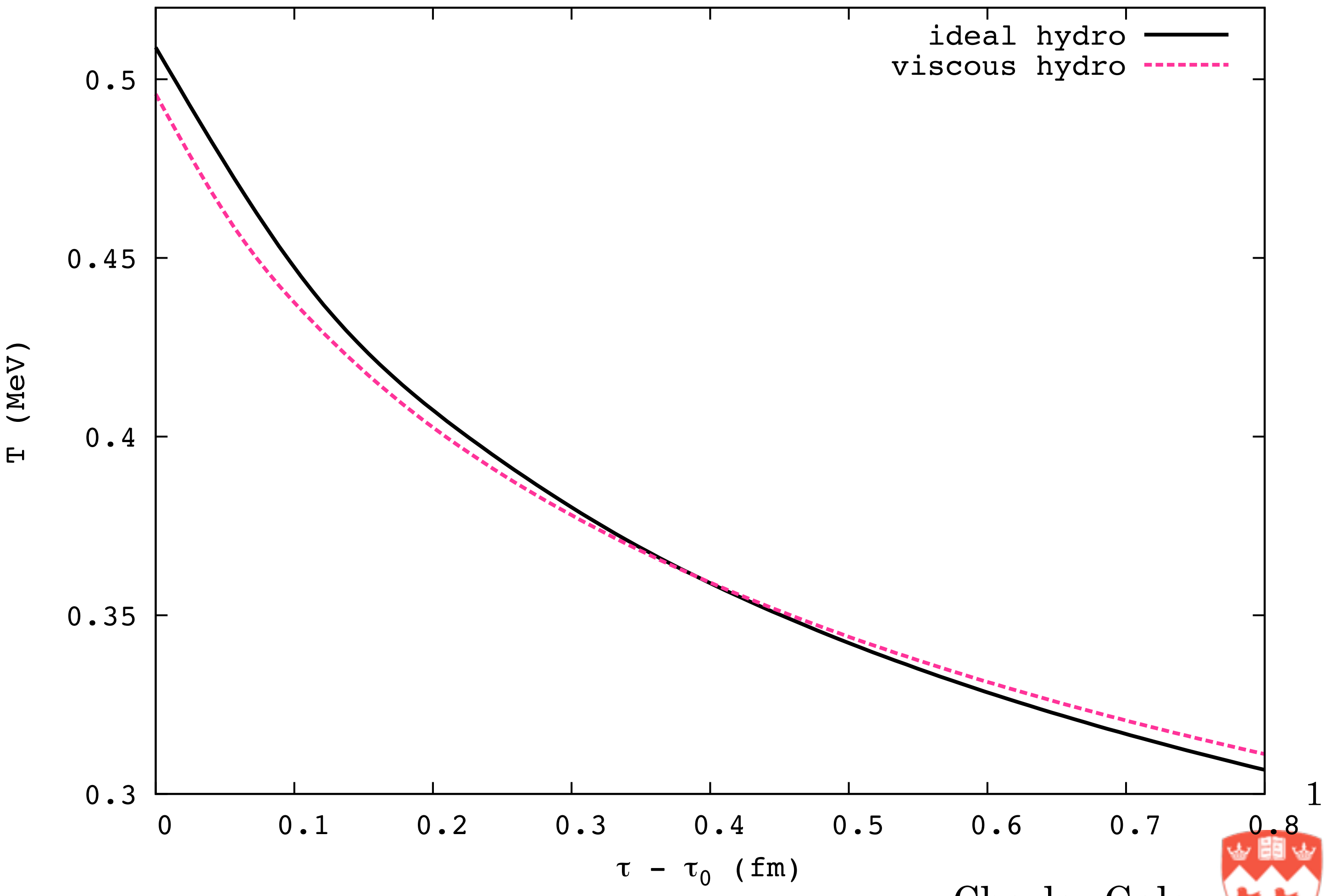
Israël & Stewart, Ann. Phys. (1979), Baier et al., JHEP (2008), Luzum and Romatschke, PRC (2008)





THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION

In-medium **hadrons**:

$$f_0(u^\mu p_\mu) = \frac{1}{(2\pi)^3}\frac{1}{\exp[(u^\mu p_\mu - \mu)/T] \pm 1}$$

$$f \to f_0 + \delta f, \quad \delta f = f_0(1 \pm (2\pi)^3 f_0) p^\alpha p^\beta \pi_{\alpha\beta}\frac{1}{2(\varepsilon + P)T^2}$$

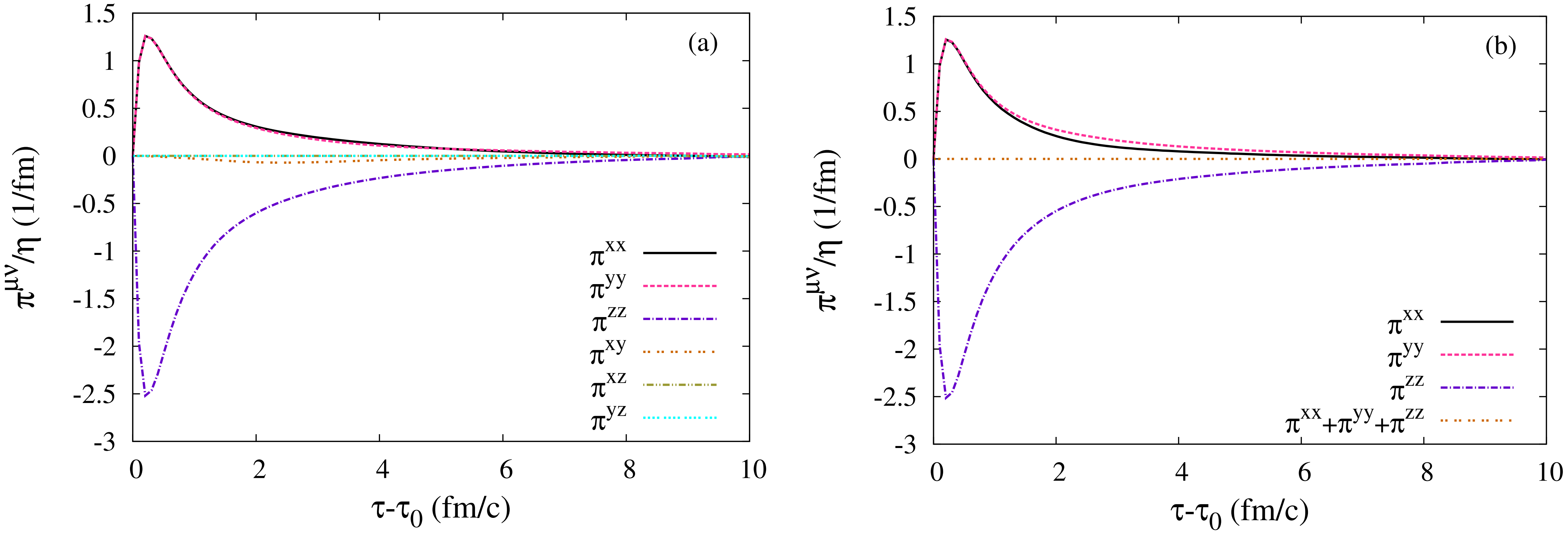
$$q_0\frac{d^3R}{d^3q} = \int \frac{d^3p_1}{2(2\pi)^3E_1}\frac{d^3p_2}{2(2\pi)^3E_2}\frac{d^3p_3}{2(2\pi)^3E_3}(2\pi)^4\left|M\right|^2\delta^4(\ldots)\frac{f(E_1)f(E_2)\left[1 \pm f(E_3)\right]}{2(2\pi)^3}$$

One considers all the reaction and radiative decay channels of external state combinations of:

$\{\pi, K, \rho, K^*, a_1\}$ With hadronic form factors

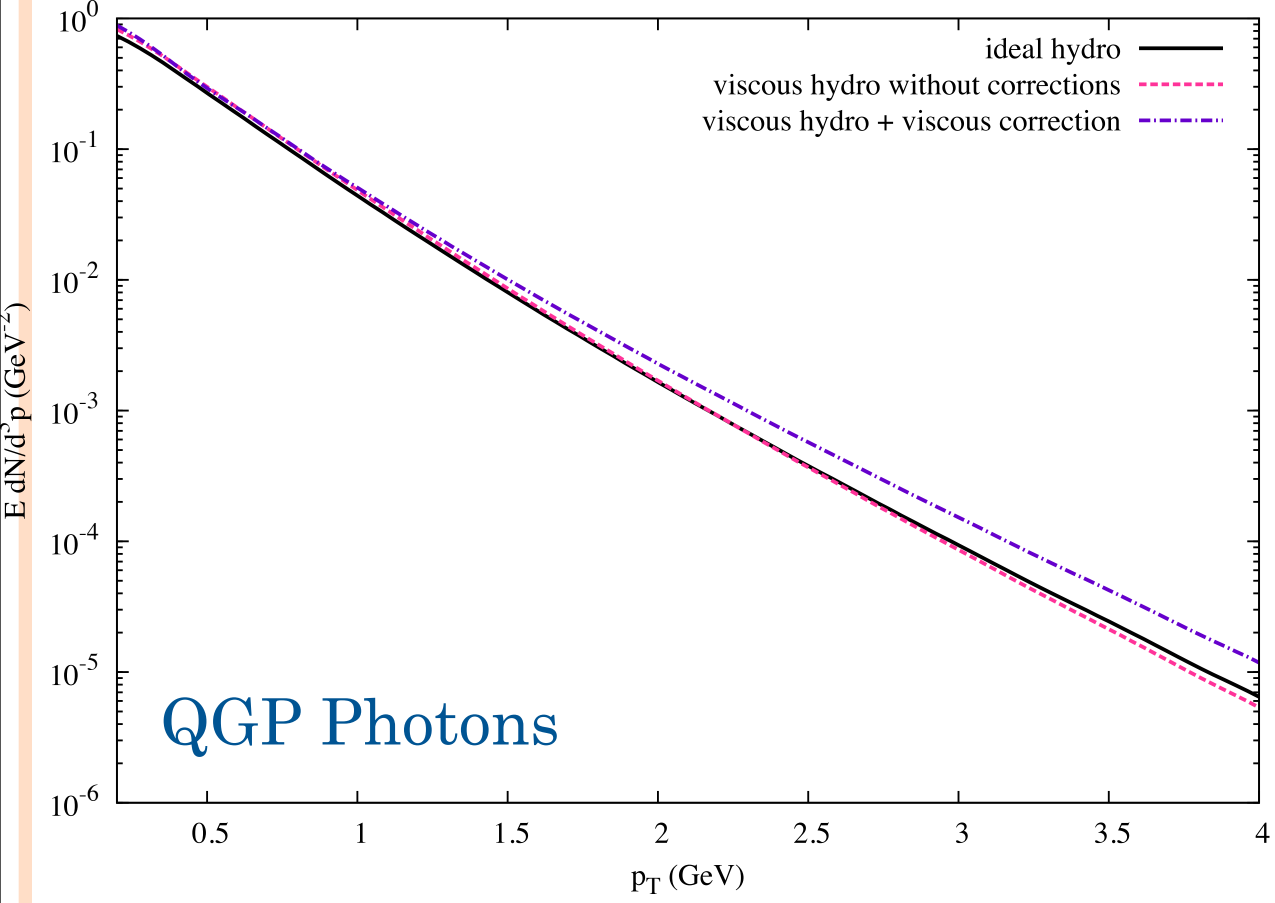
+ QGP Photons

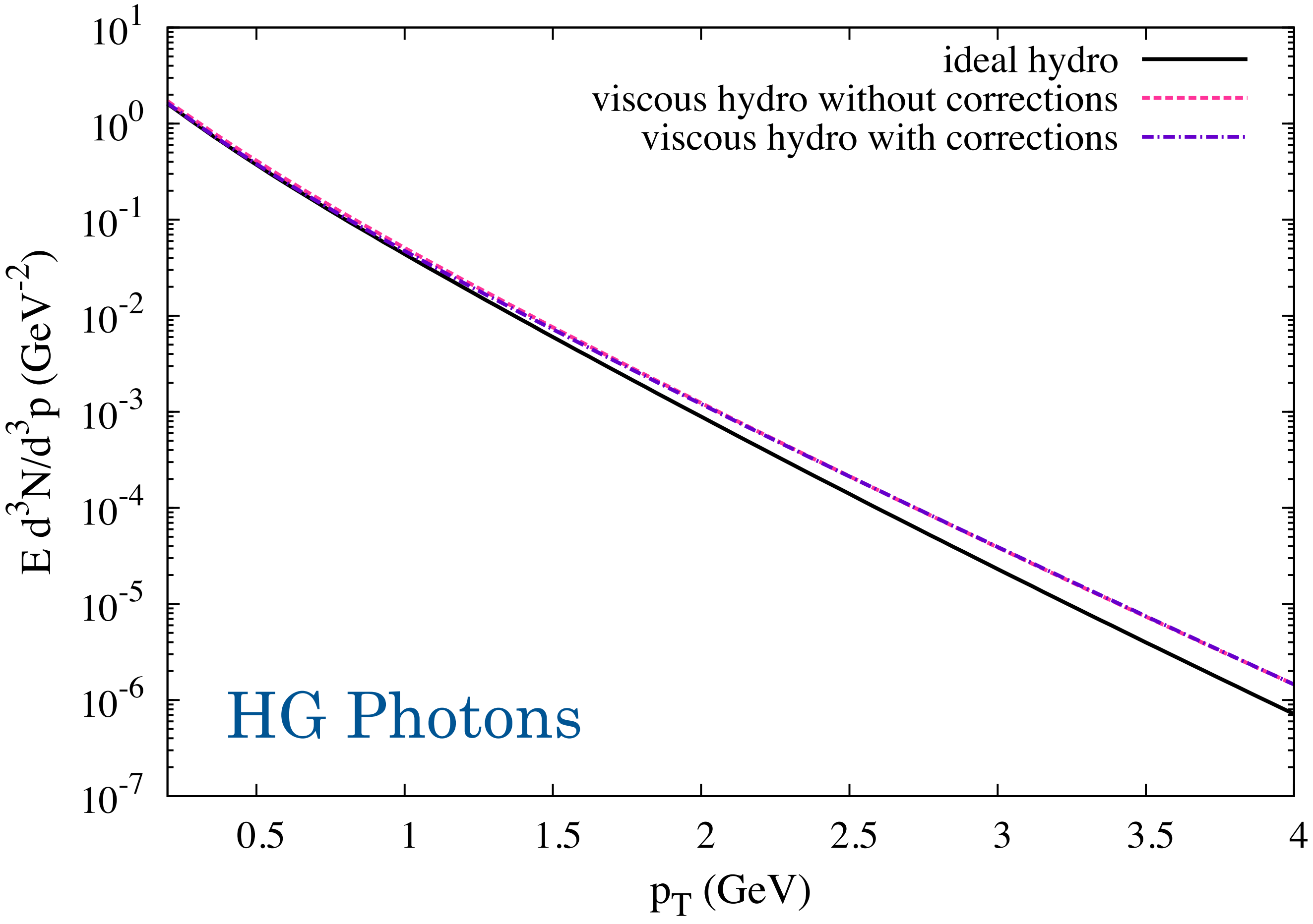
THE EFFECTS OF SHEAR VISCOSITY ON THE BULK DYNAMICS



- Large at early times
- Small at later times: viscosity corrections to the distribution functions will also vanish

THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION



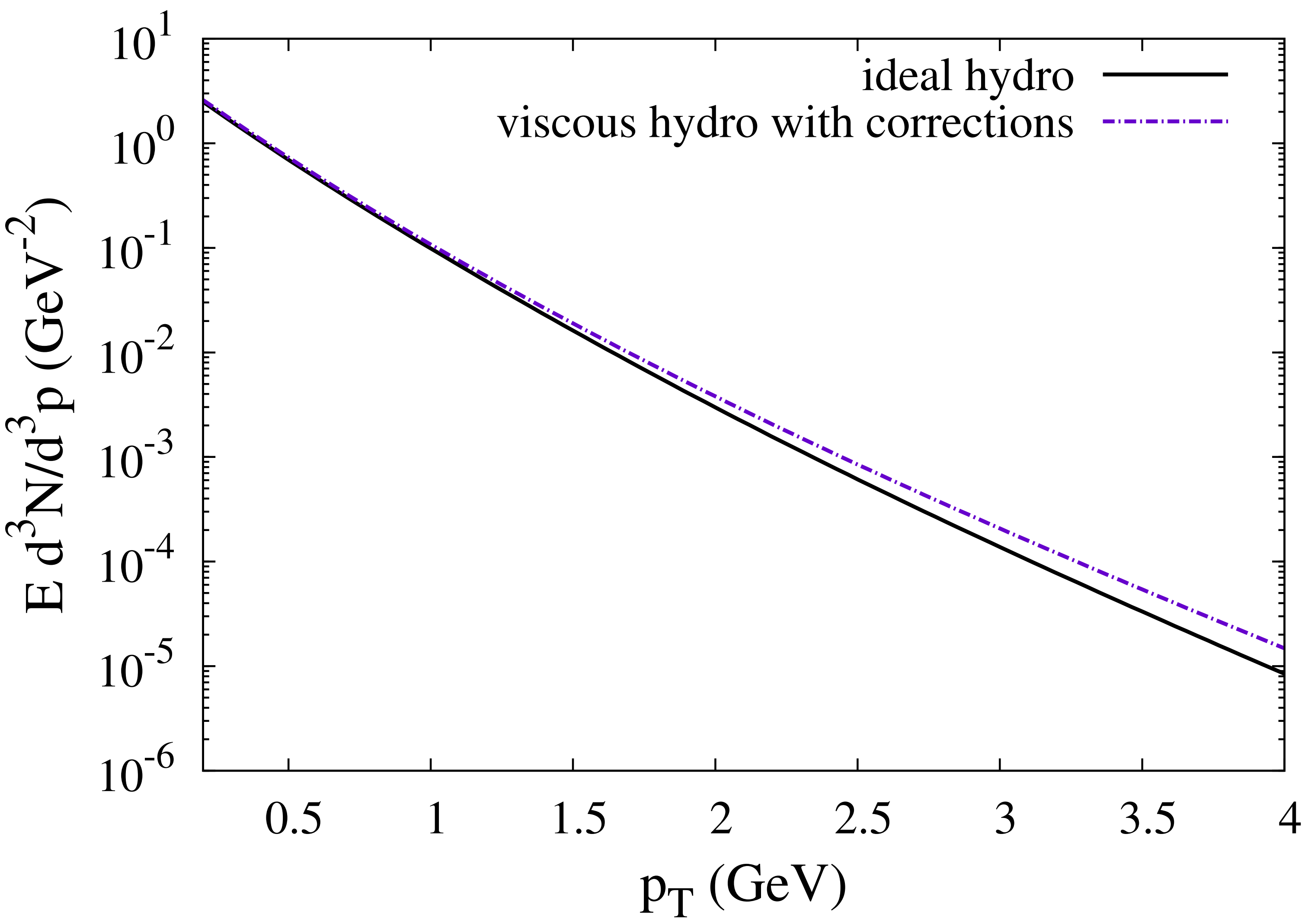


K. Dusling NPA (2010)
Chaudhuri & Sinha, PRC (2011)

Viscous effects harden the photon spectrum

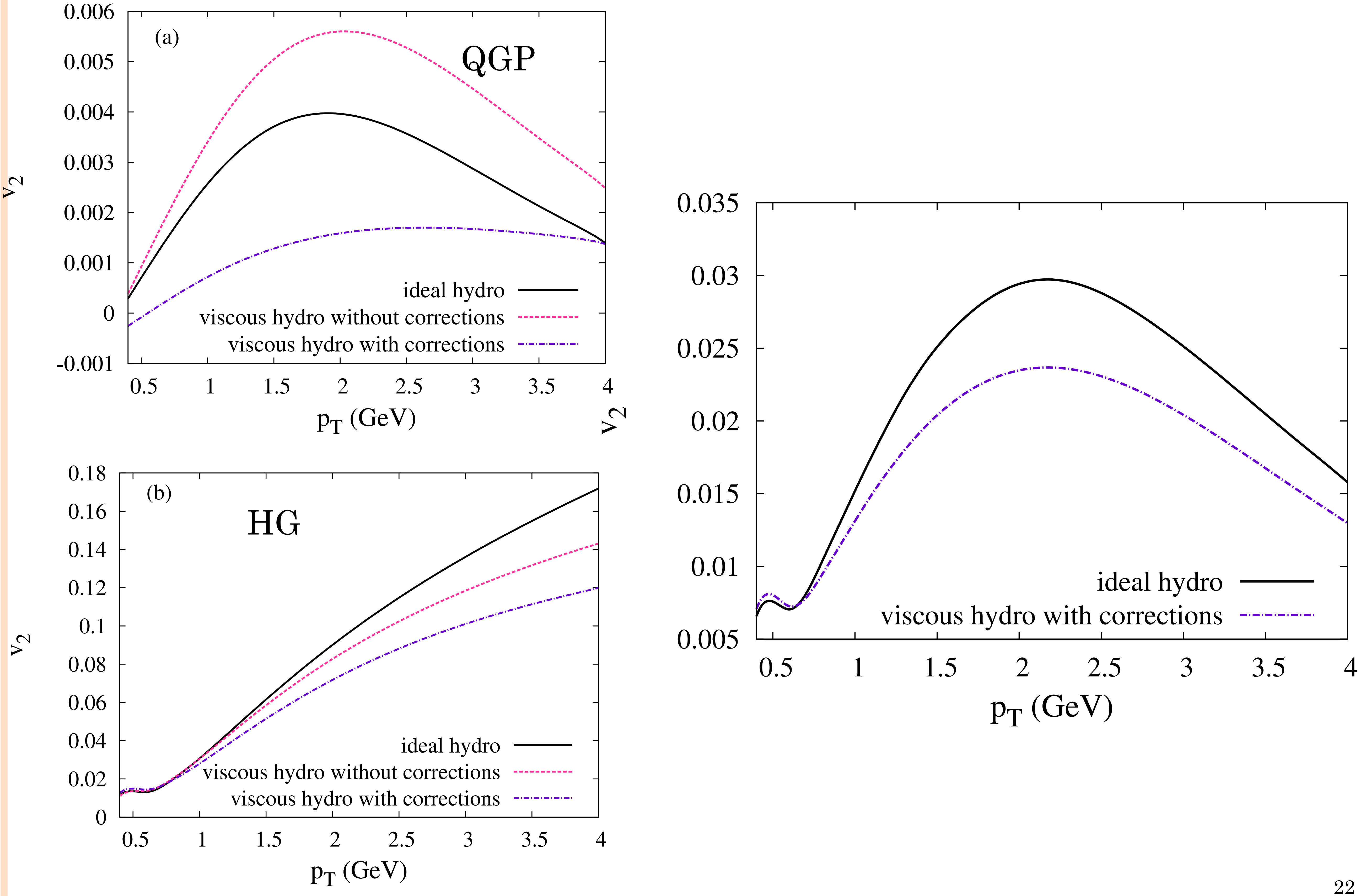
M. Dion et al., PRC (2011)

THE NET THERMAL PHOTON YIELD



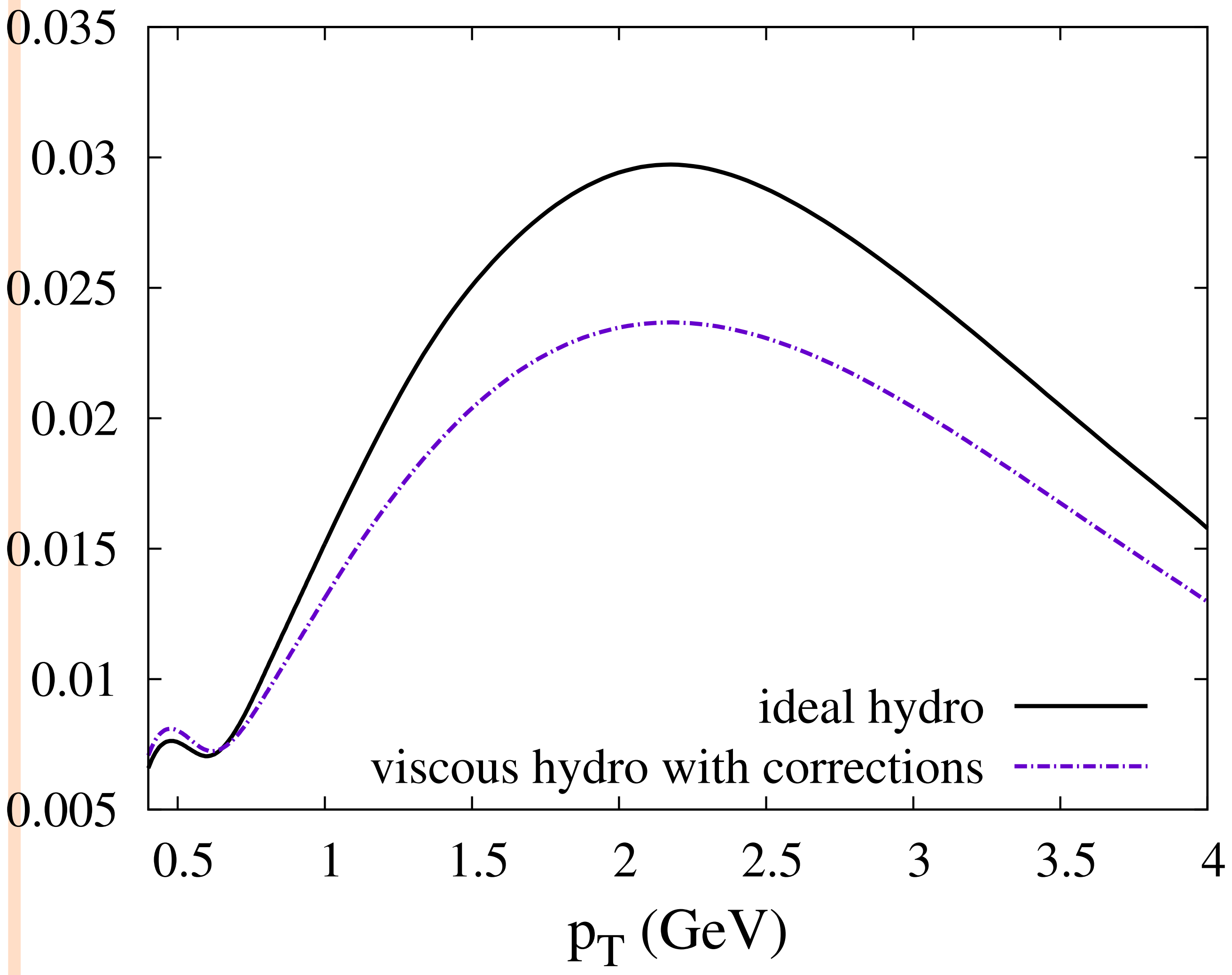
- Viscous corrections make the spectrum harder, ≈100% at $p_T = 4$ GeV.
- Increase in the slope of ≈15% at $p_T = 2$ GeV.
- Extracting the viscosity from the photon spectra will be challenging
- Once pQCD photons are included: a few % effect from viscosity
- More work is still needed to properly include all photon sources in a consistent way

SHEAR VISCOSITY AND PHOTON V2



M. Dion *et al.*, PRC 2011

SHEAR VISCOSITY AND PHOTON v_2

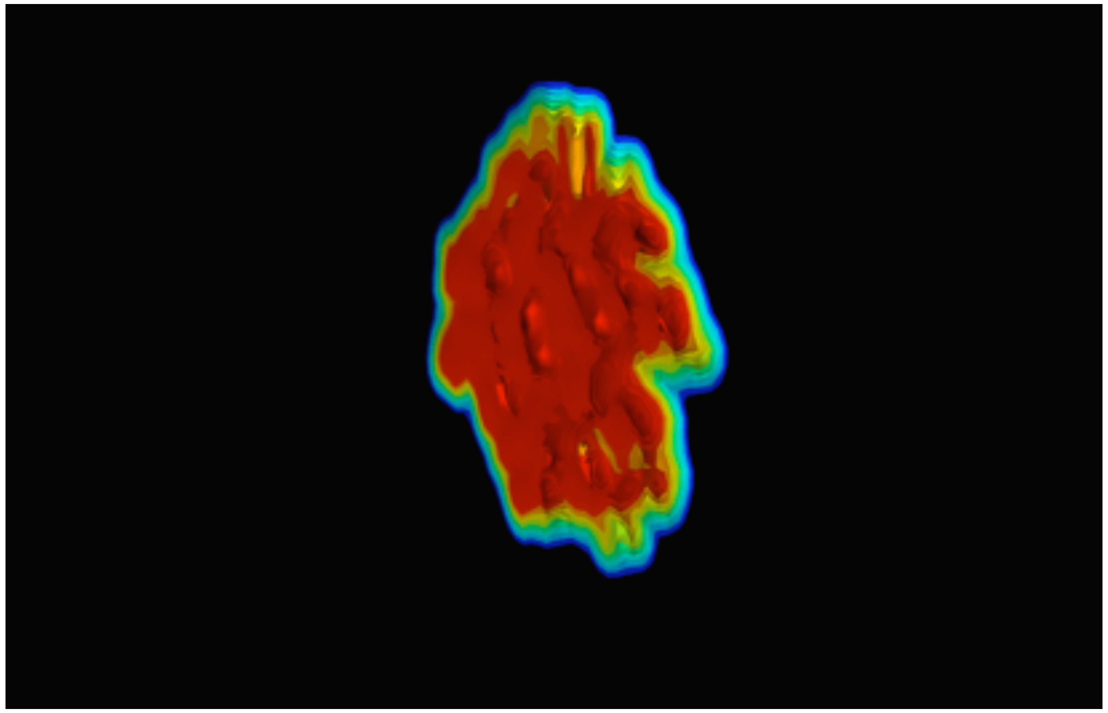


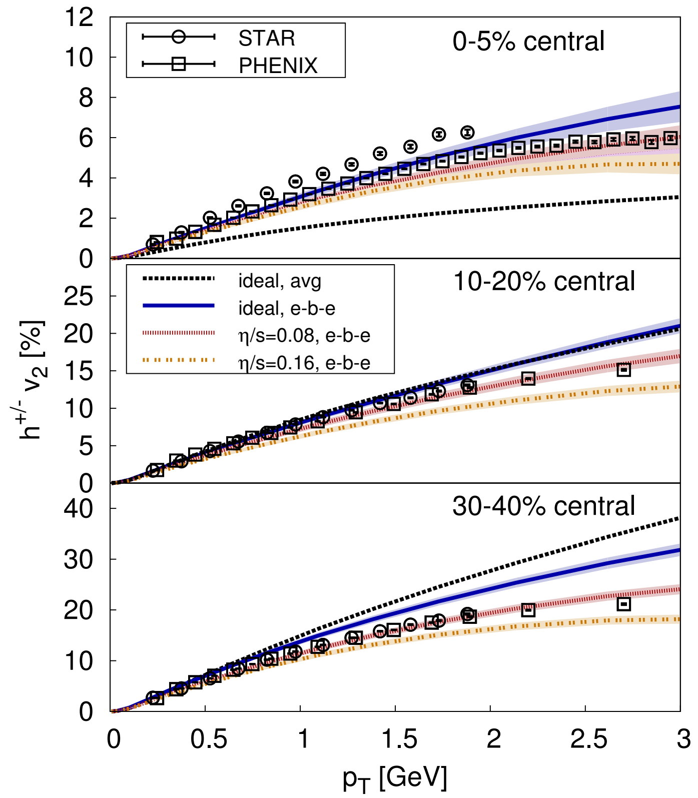
- The net elliptic flow is a *weighted average*. A larger QGP yield will yield a smaller v_2. Same story - *mutatis mutandis* - for the HG
- The turnover at $p_T \approx 2$ GeV could be QGP-driven and/or pQCD-driven
- The net effect of viscous corrections makes the photon elliptic flow *smaller*, as it does for hadrons

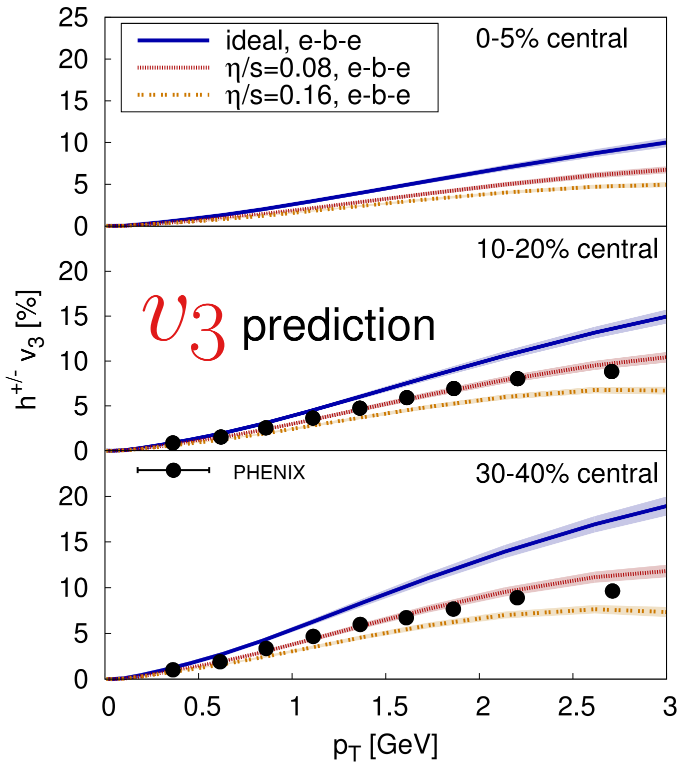
M. Dion *et al.*, PRC 2011

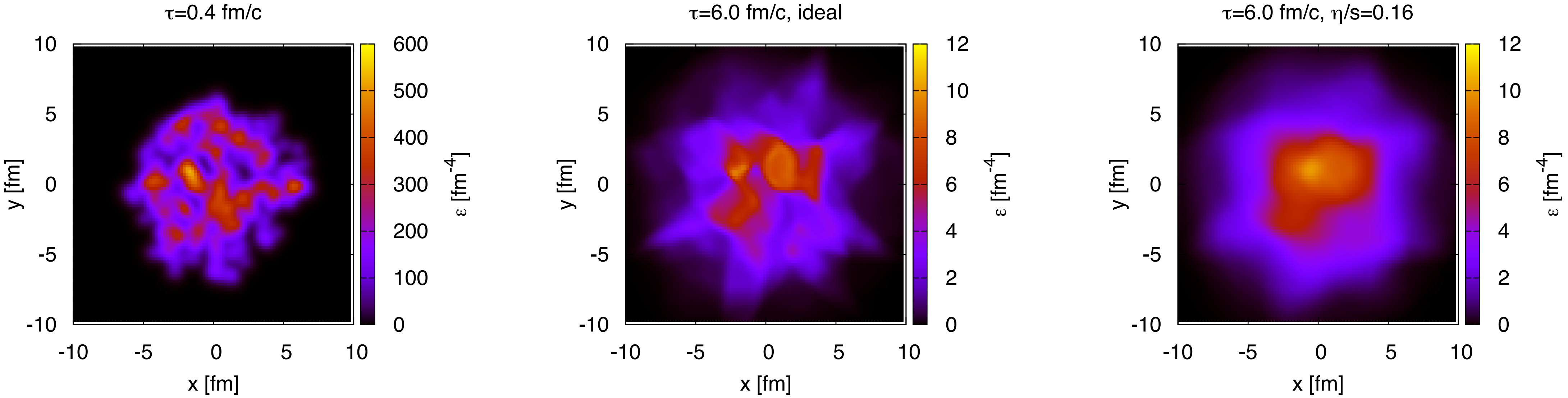
INITIAL STATE FLUCTUATIONS: A PARADIGM SHIFT IN HEAVY ION ANALYSES

Lumpy
MUSIC









Schenke, Jeon, Gale, PRL (2011)

INITIAL STATE FLUCTUATIONS: A PARADIGM SHIFT IN HEAVY ION ANALYSES

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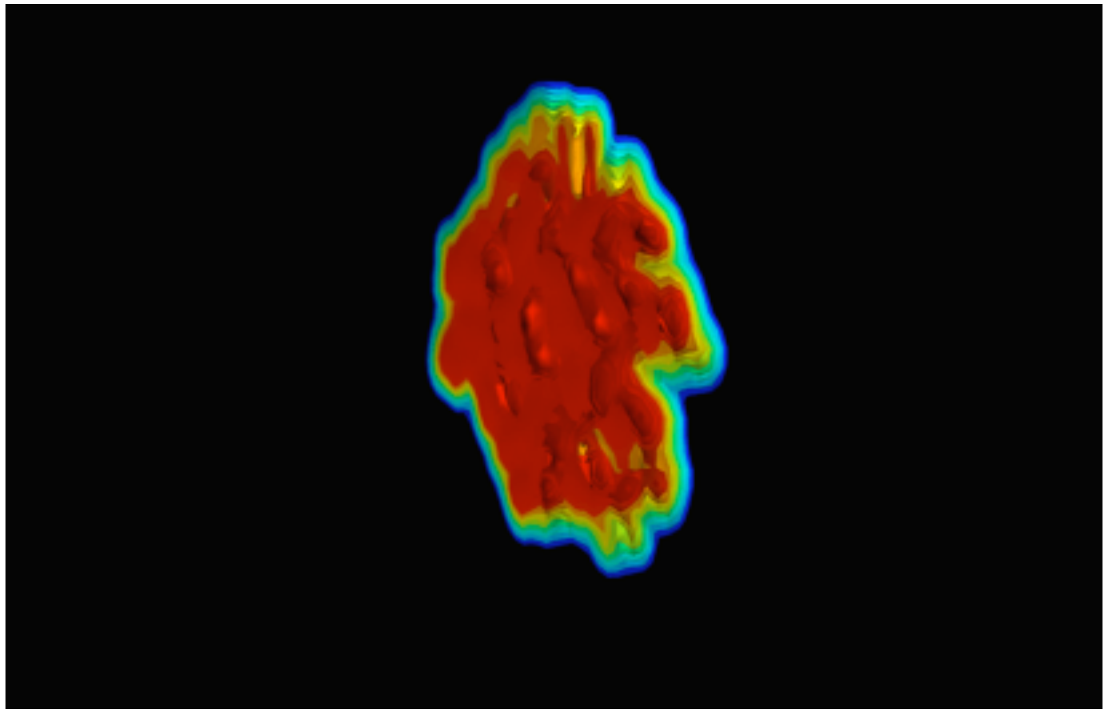
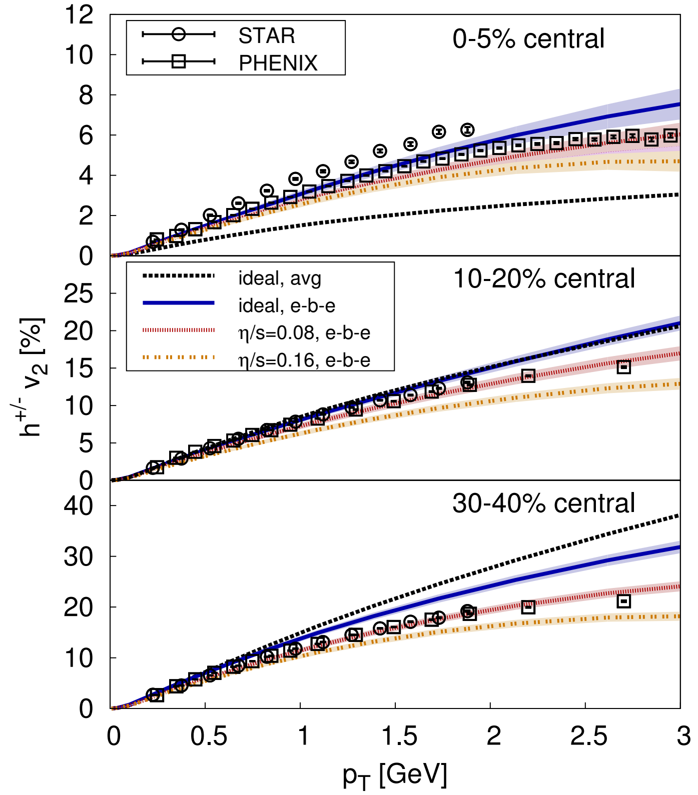
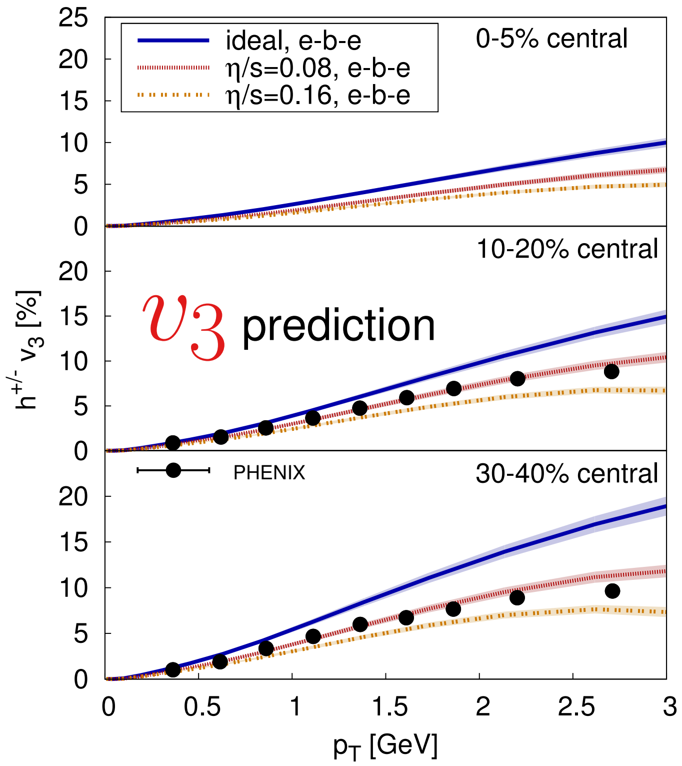
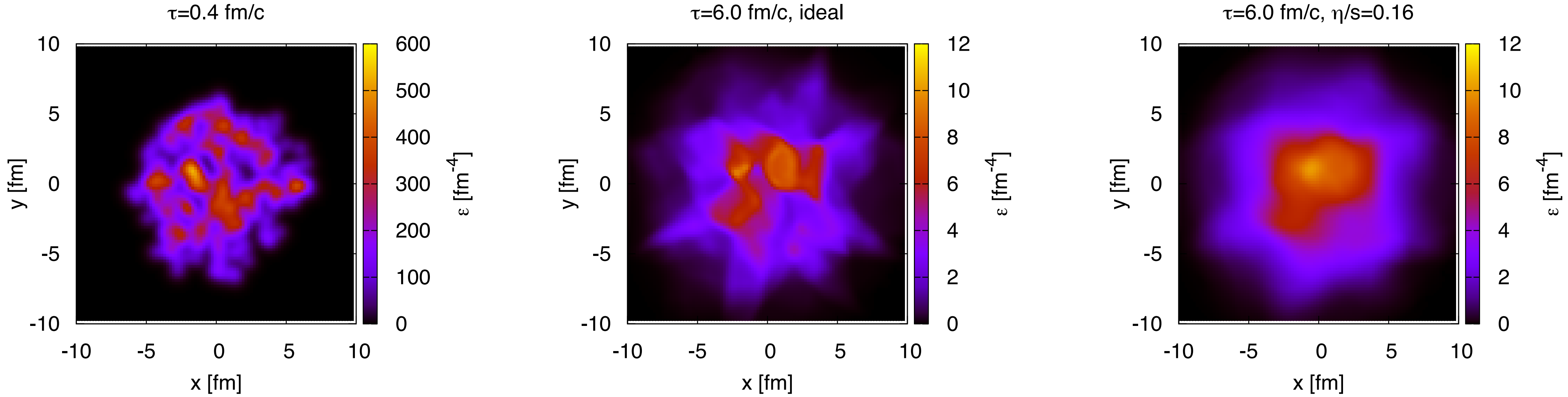


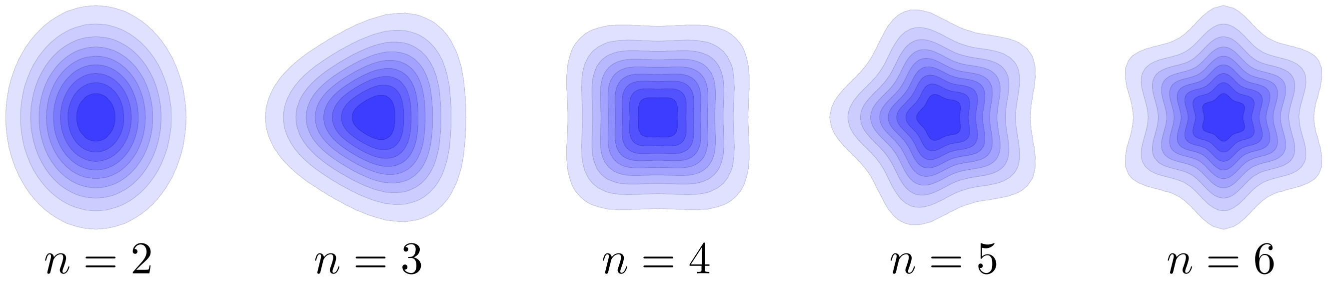


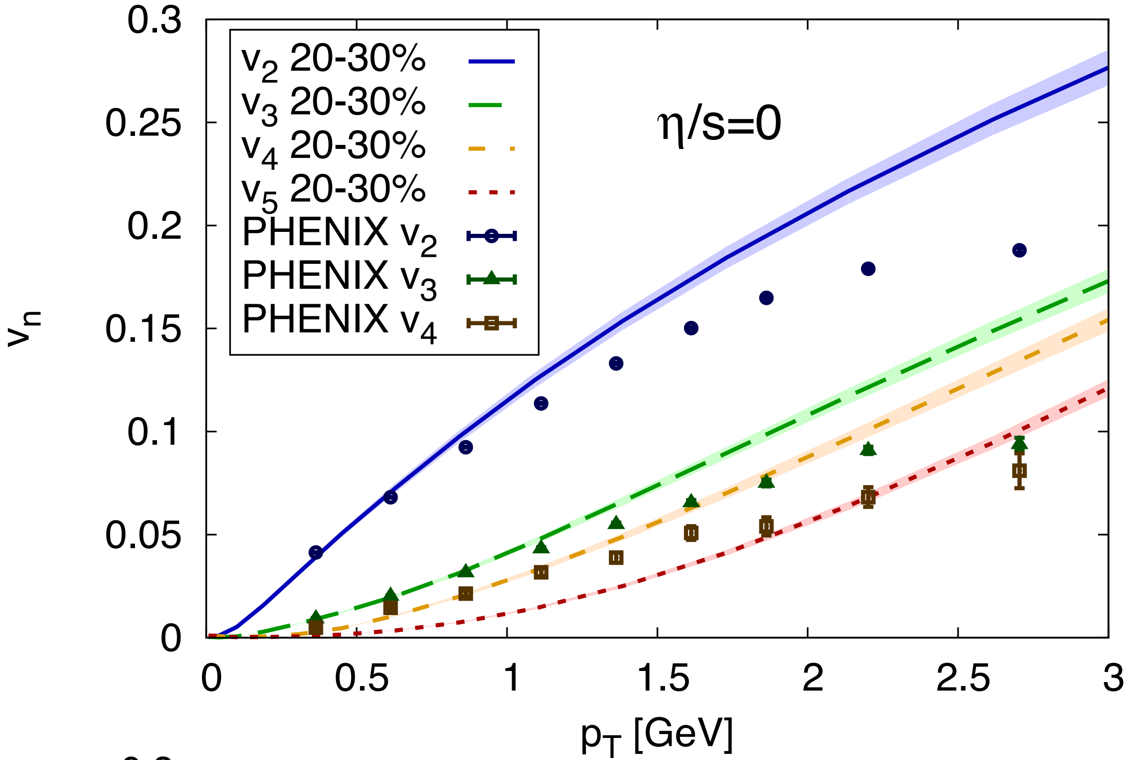


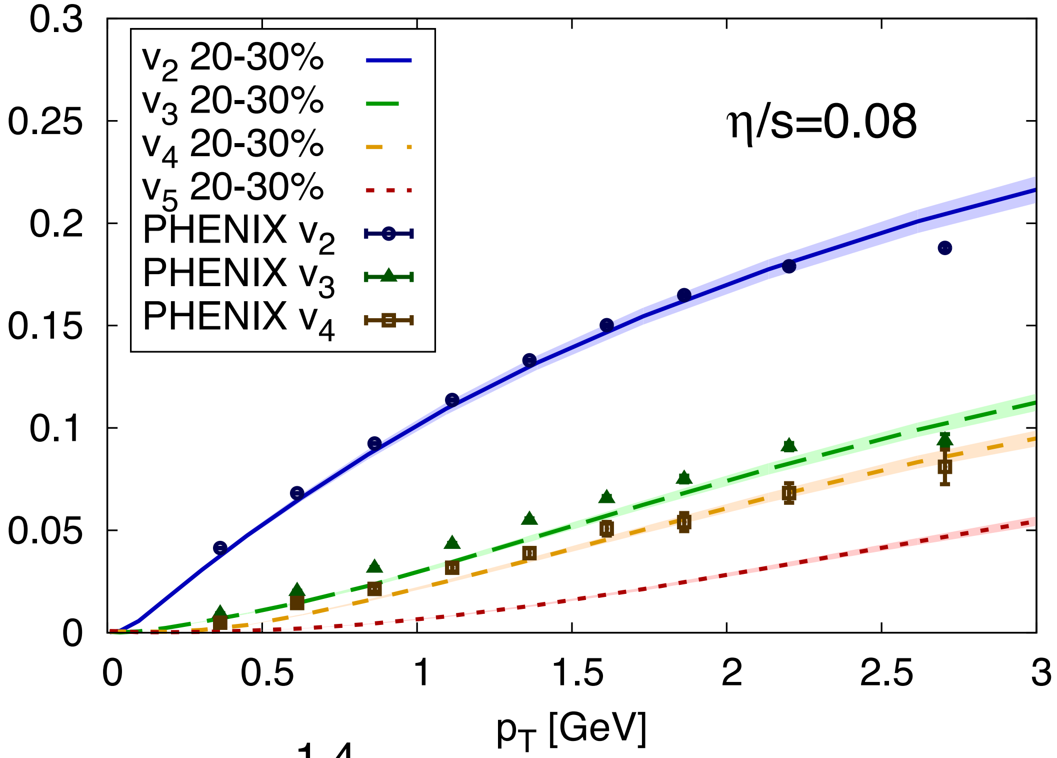


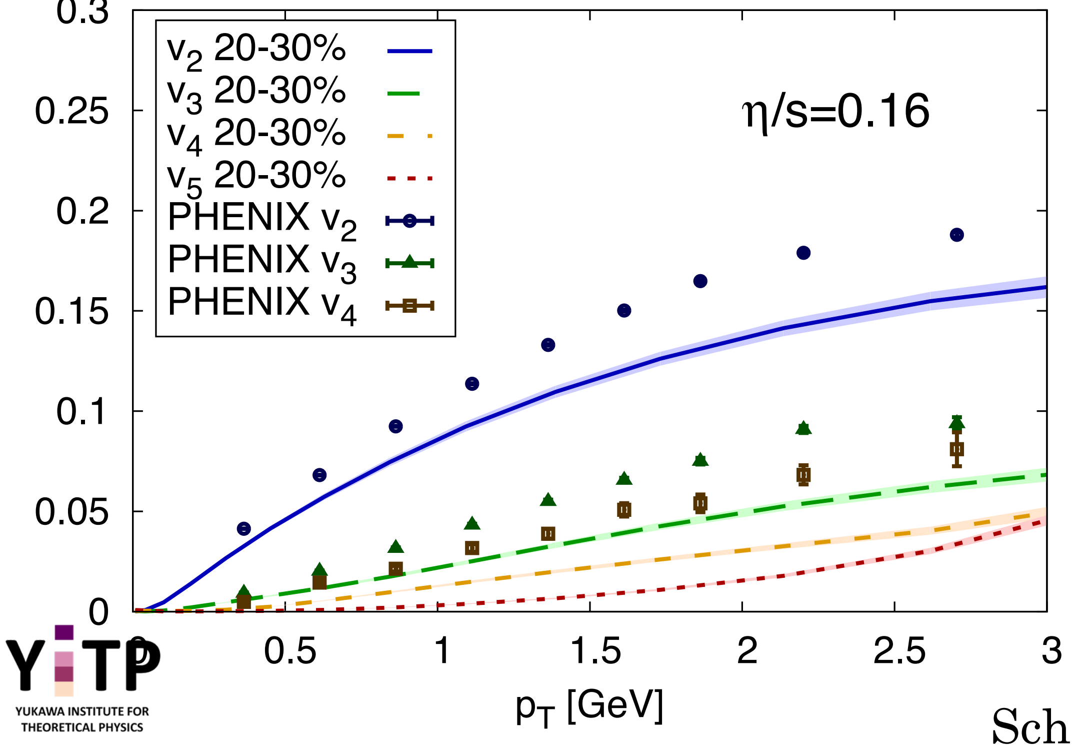
Schenke, Jeon, Gale, PRL (2011)

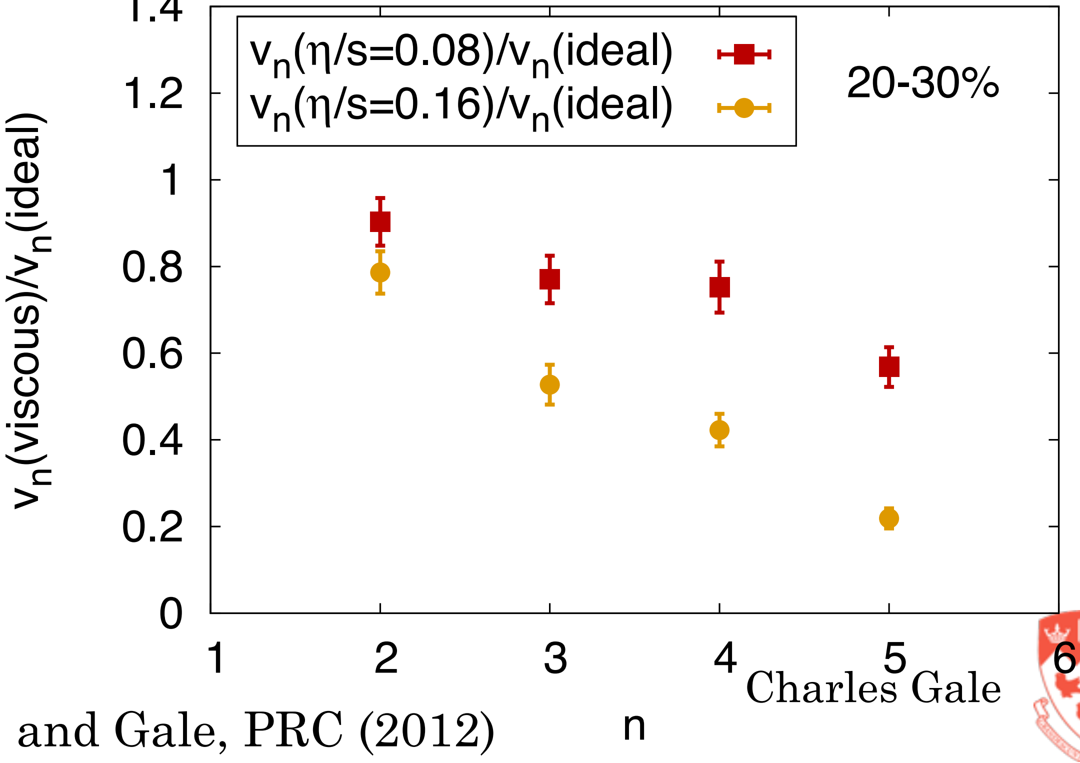
MOVING INTO THE "CHARACTERIZATION" PHASE…





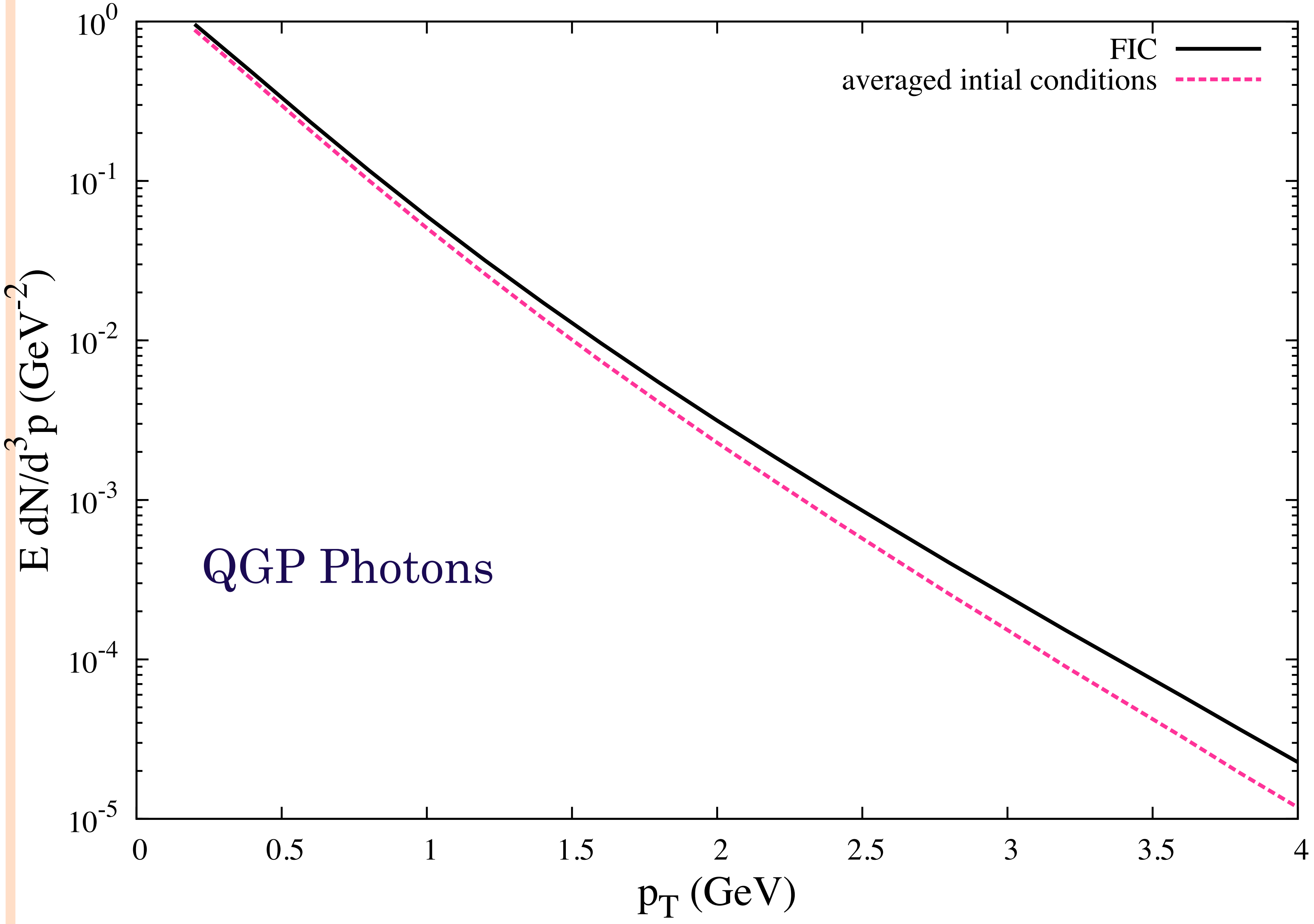






Schenke, Jeon, and Gale, PRC (2012)

THE EFFECT OF FIC ON THE THERMAL PHOTON SPECTRUM



- FIC produces higher initial T (hot spots), and higher initial gradients
- FIC conditions are demanded by hadronic data (v_{odd})
- These lead to a harder spectrum, *as for hadrons*

Dion et al., PRC (2011)
Chatterjee et al., PRC (2011)

ALL TOGETHER: FICS + VISCOSITY

- Combined with viscous corrections, FIC yield an enhancement by ≈5 @ 4 GeV, and ≈2 @ 2 GeV
- Temperature estimated by slopes can vary considerably
- A combination of hot spots and blue shift hardens spectra
- Once pQCD photons are included: only modest changes from viscous corrections + FICs

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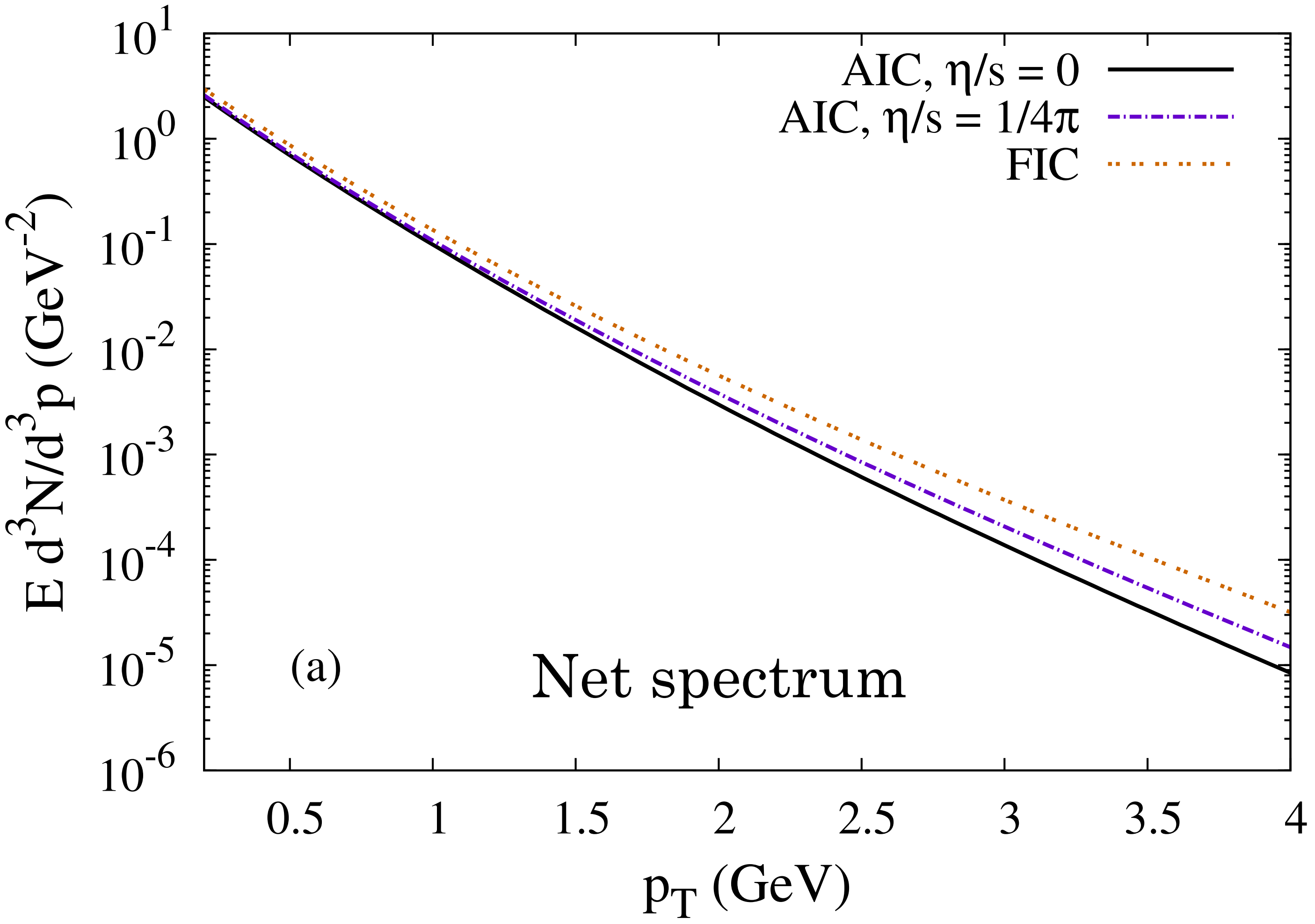

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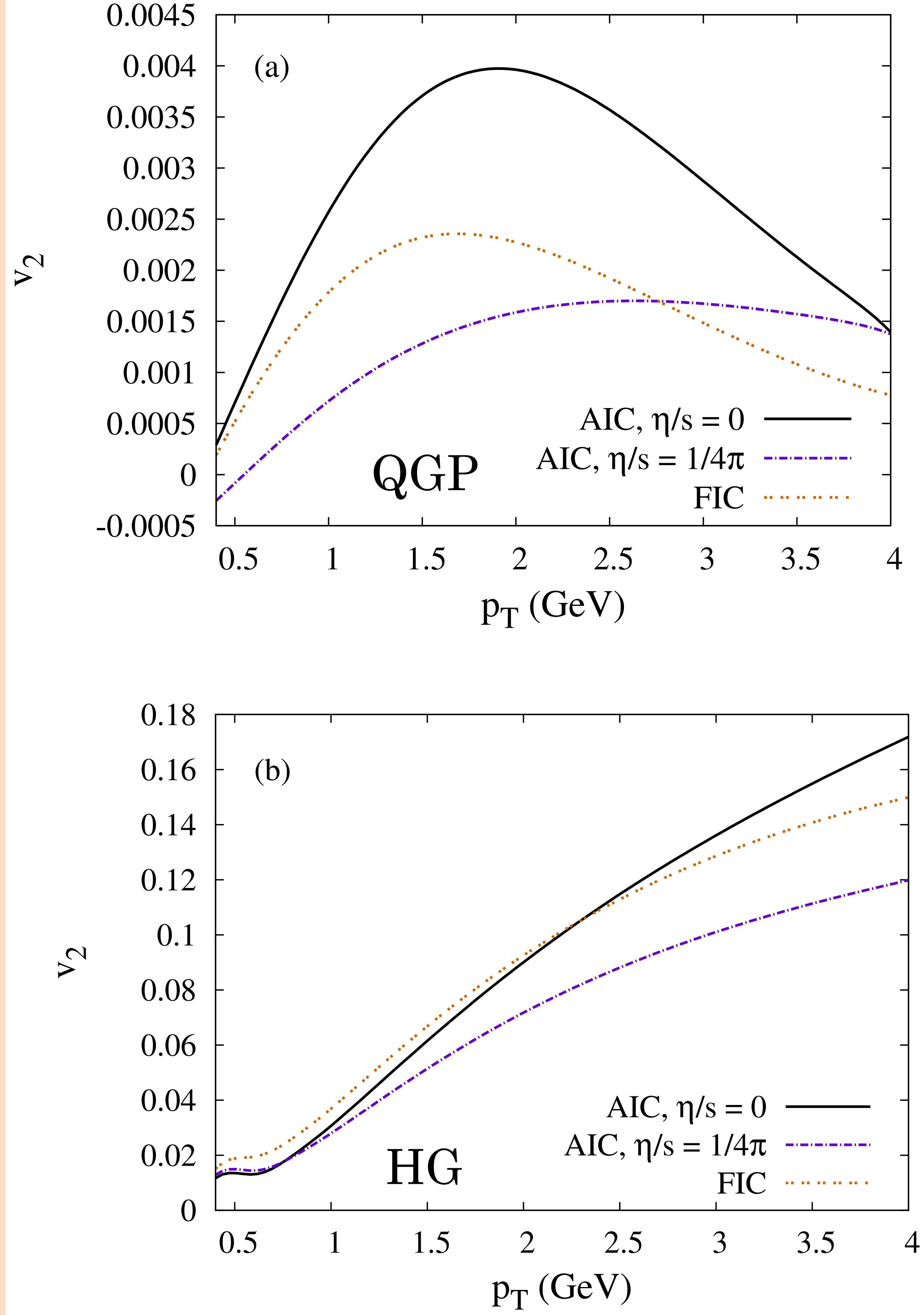
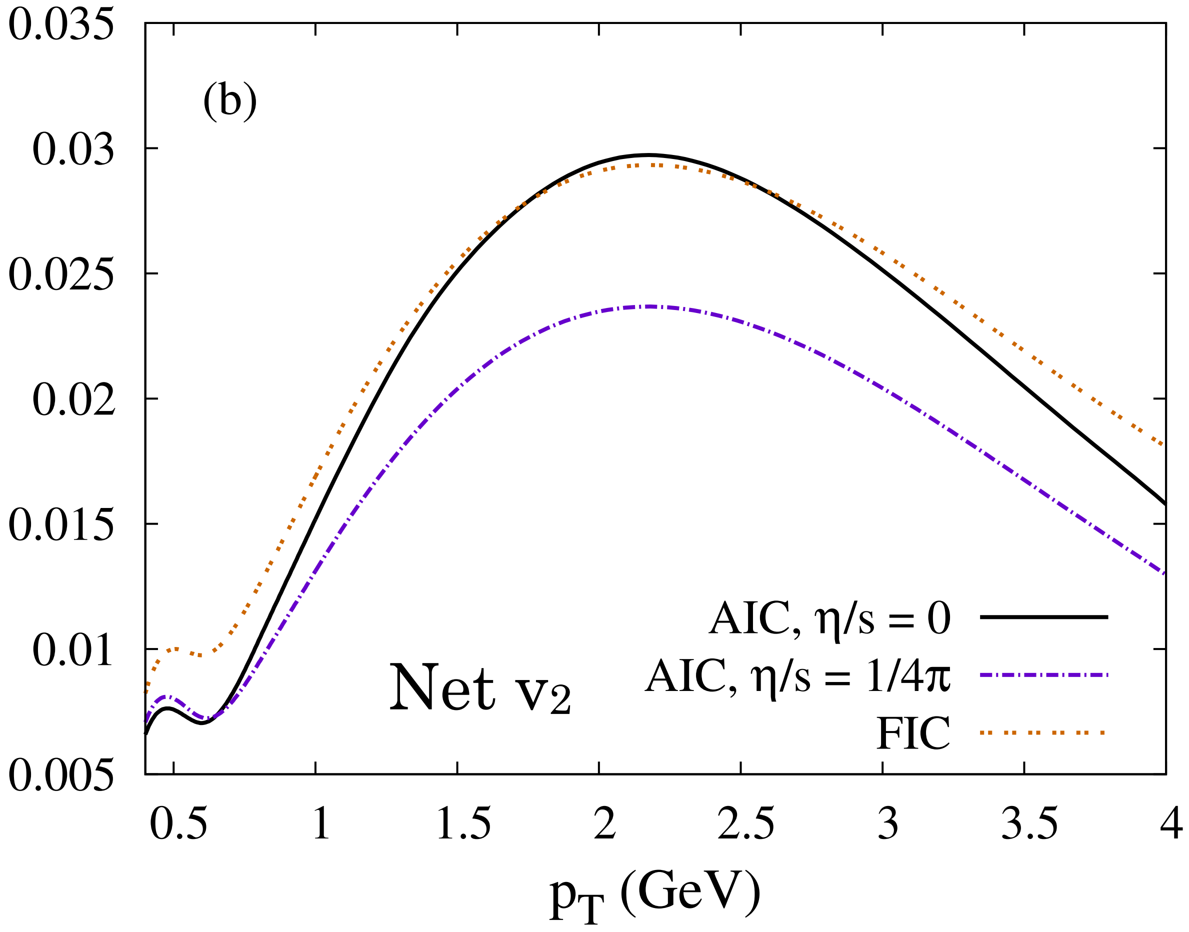
FICs AND THERMAL PHOTON V_2



- FICs enhance v_2 in this centrality class (0-20%), as for hadrons
- For hadrons measured in events belonging to large centrality, FICs will *decrease* v_2
- HG elliptic flow is much larger than QGP elliptic flow, but remember net v_2 is a weighted average

FICs AND THERMAL PHOTON v_2



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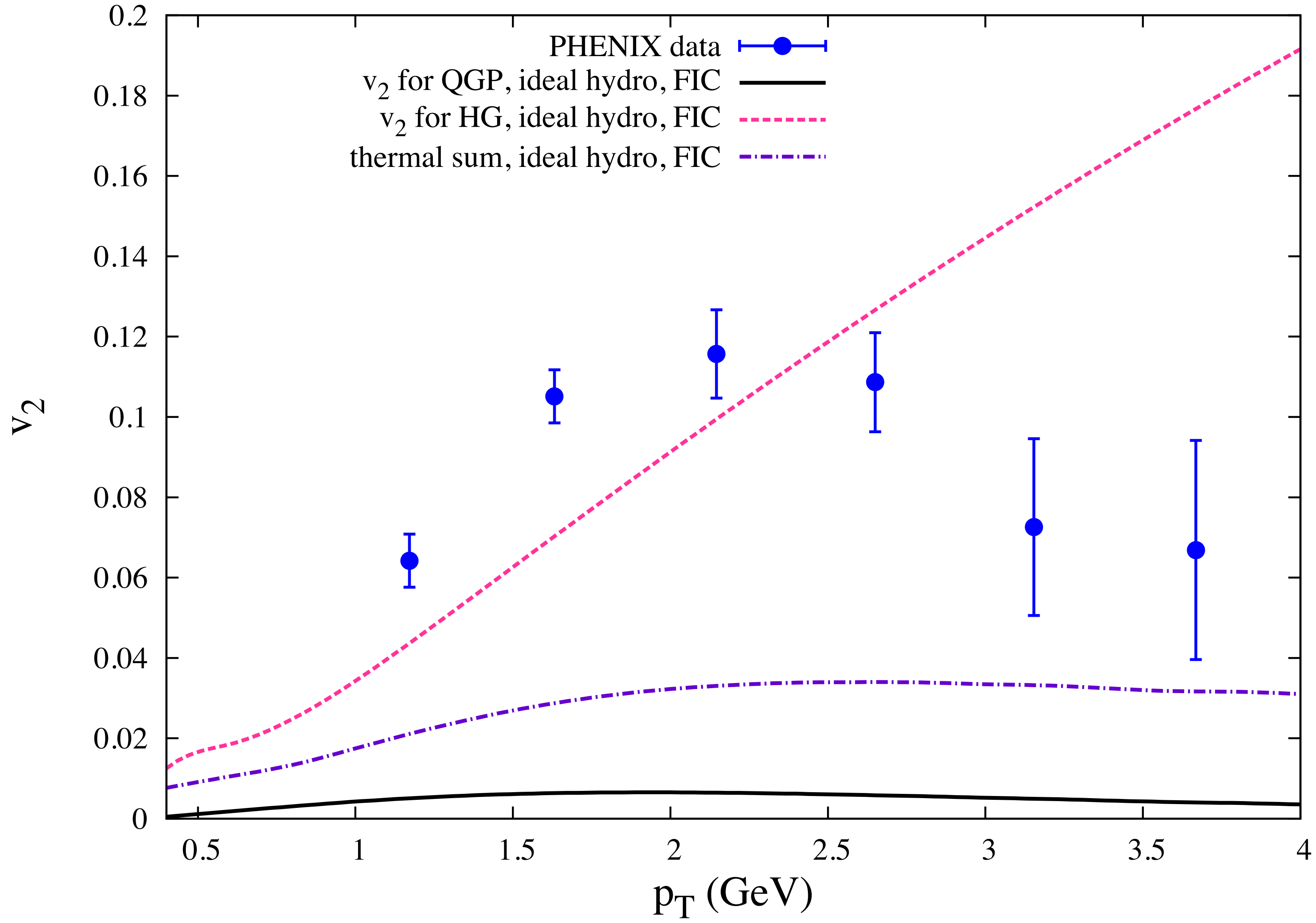
- For hadrons measured in events belonging to large centrality, FICs will *decrease* v_2
- HG elliptic flow is much larger than QGP elliptic flow, but remember net v_2 is a weighted average
- Net v_2 is comparable in size to that with ideal medium.

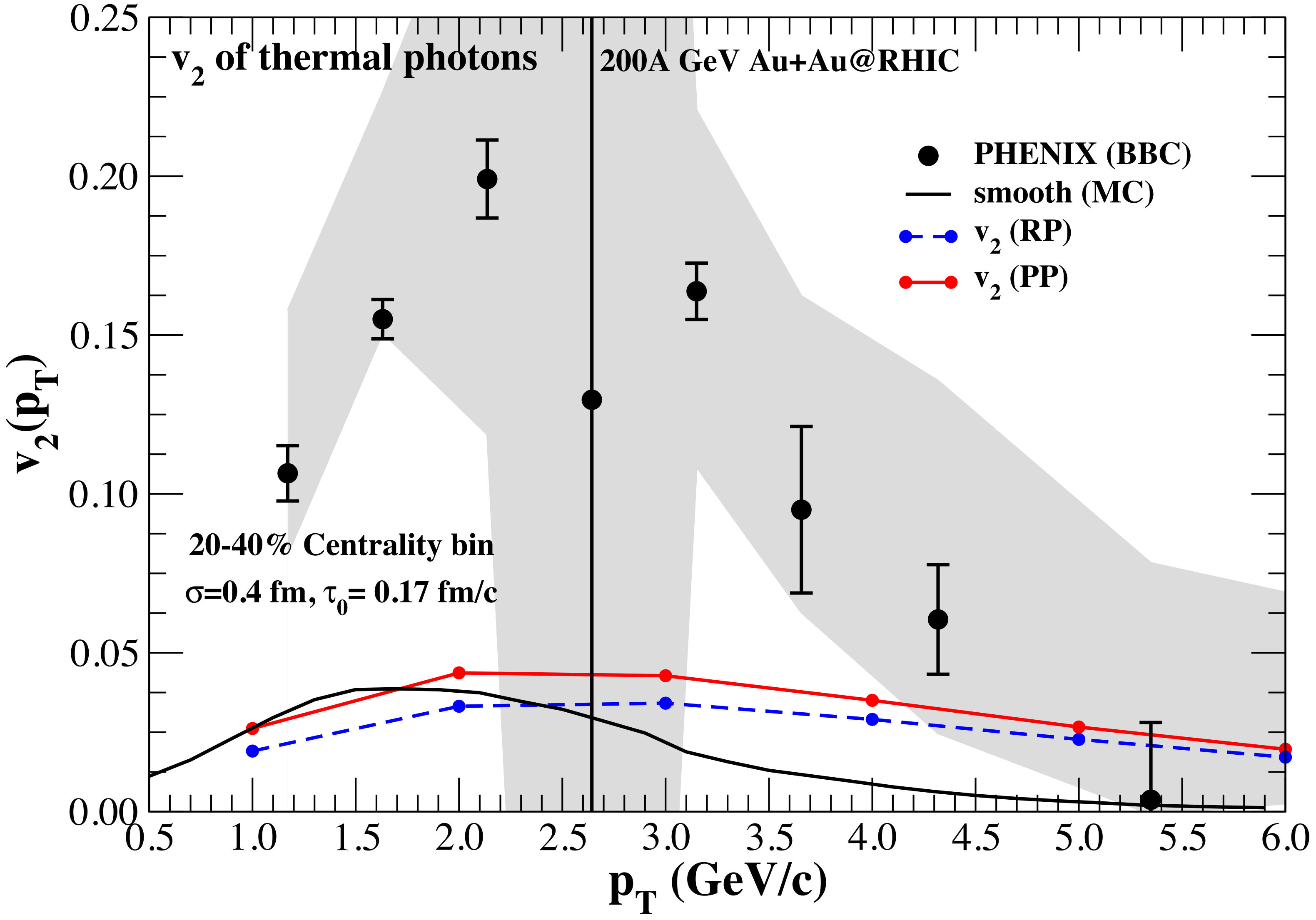
PHOTON v_2 DATA?



Dion et al. (2011)

- Data is higher than calculation, even with e-b-e initial state fluctuations, and ideal hydro
- Size comparable with HG v_2

PHOTON v_2 DATA?



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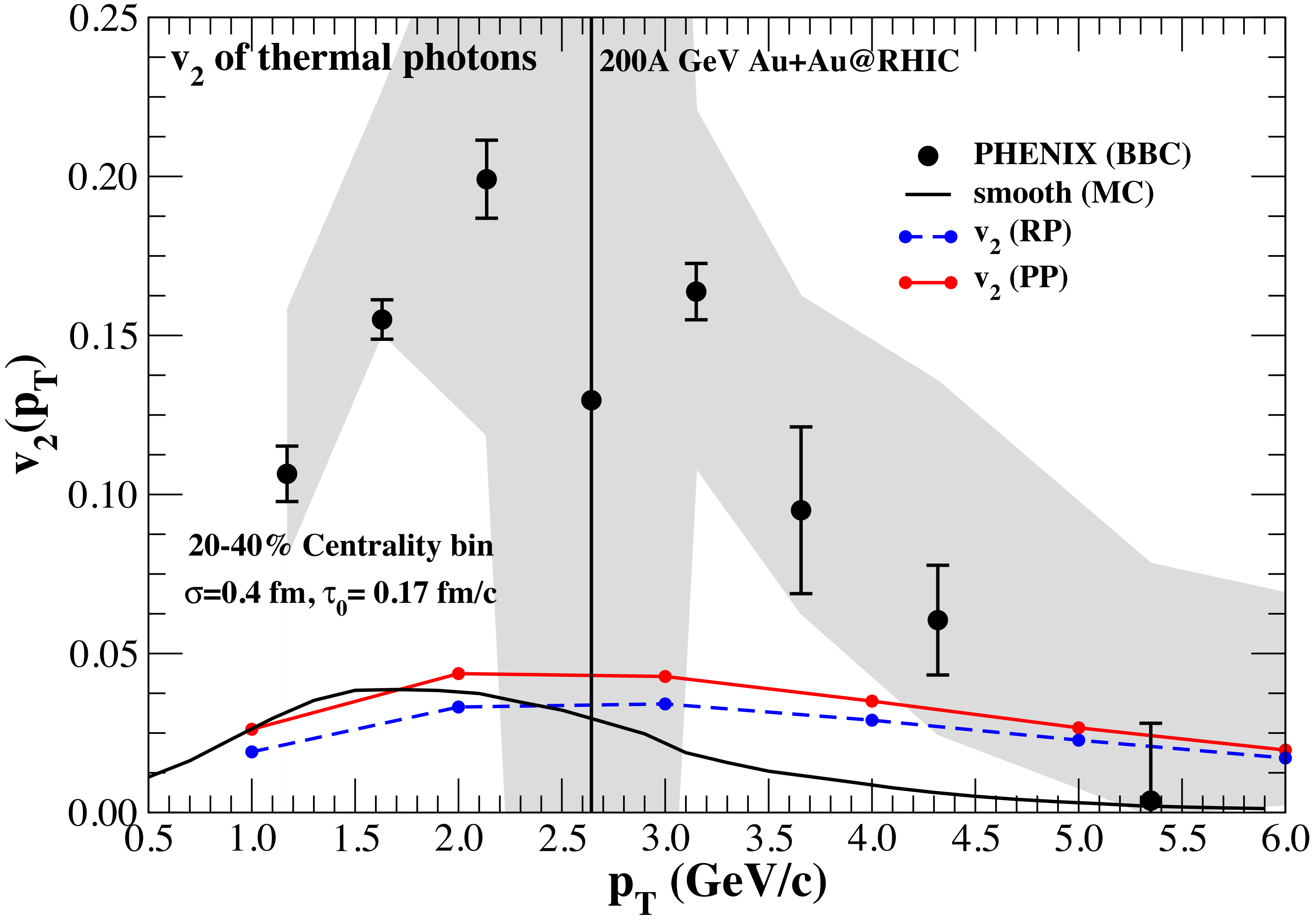

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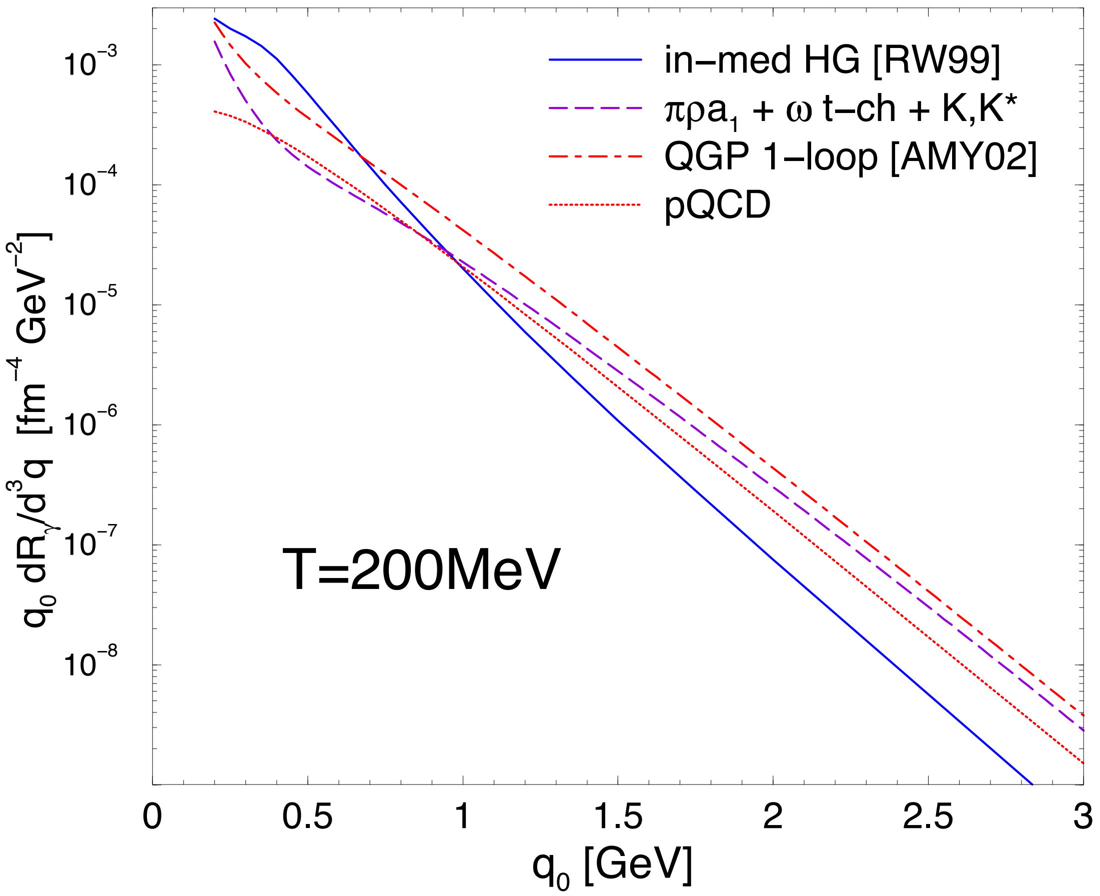
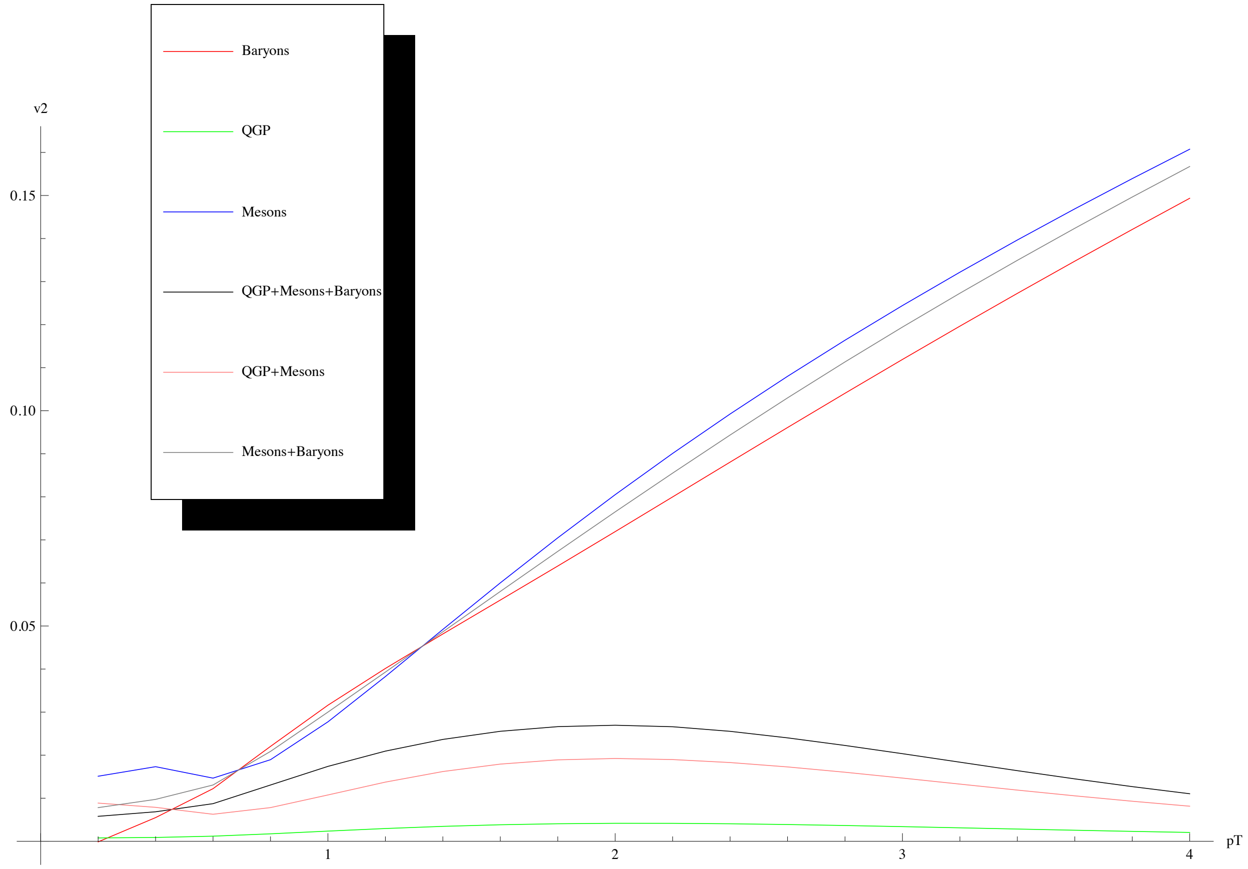
SOME FACTS AND SOME LEADS

- FICs are here to stay. The meaning of "initial temperature" is altered.
- (Some?) Room to explore systematically hydro initialization and parameters. This requires consistency with the hadronic data.
- Making the QGP signal larger will *decrease* the v_2. The T=0 photons, *decrease* v_2.
- Early-times magnetic field effects? (Basar, Kharzeev, Skokov, PRL (2012))
- Is the large photon elliptic flow telling us about the dynamics? Baryons?
- Non-zero initial shear tensor? Primordial flow? Can we improve on the hydro initial states?
- Can we improve on the hydrodynamic evolution?

EFFECTS OF BARYONS ON PHOTON V_2?



Turbide, Rapp, Gale, PRC (2004)

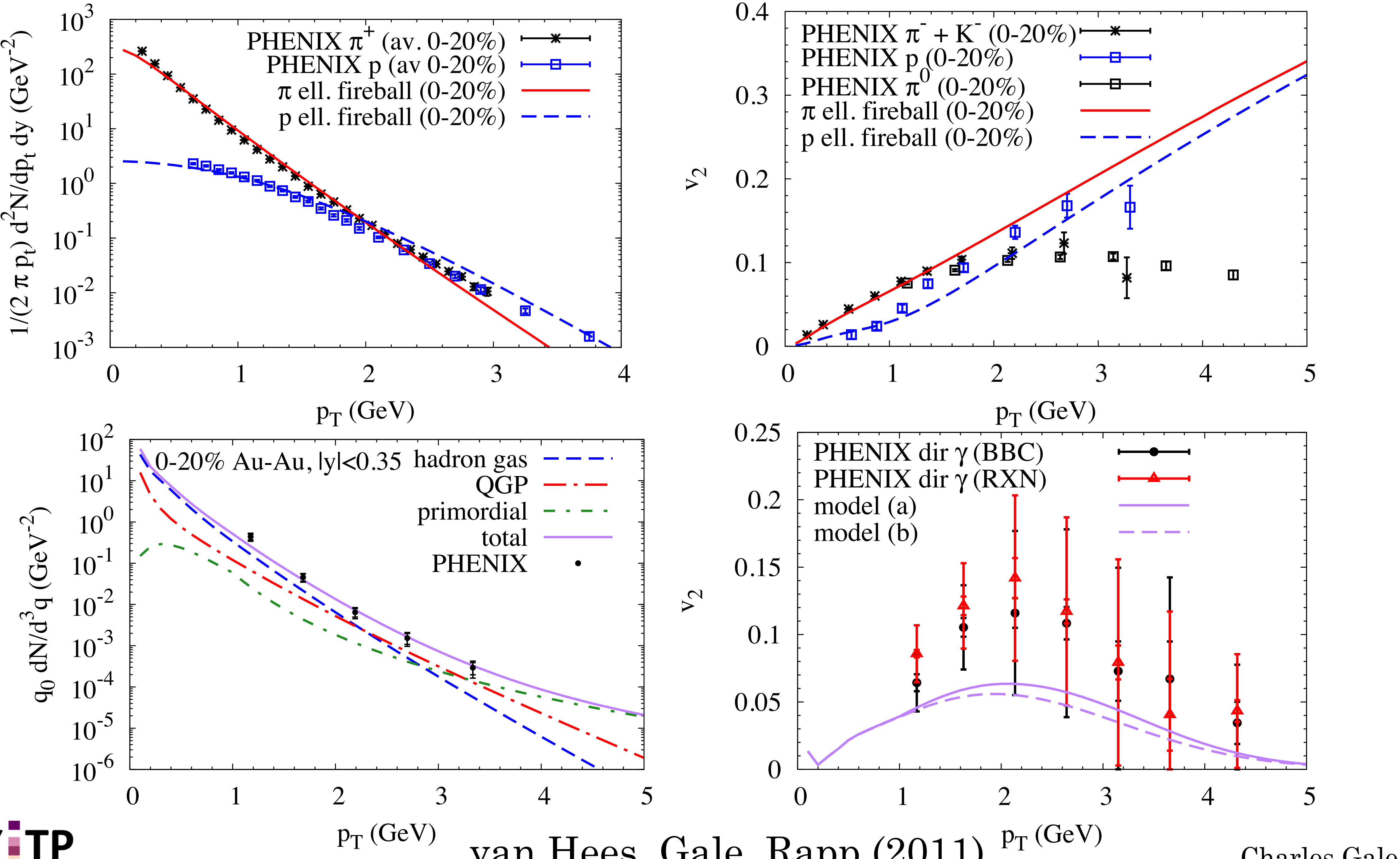


J.-F. Paquet et al., (2013)

- Adds to meson elliptic flow (~2)
- Promising. Still needs a proper, consistent calculation

ELLIPTIC FLOW AND SPACE-TIME DYNAMICS

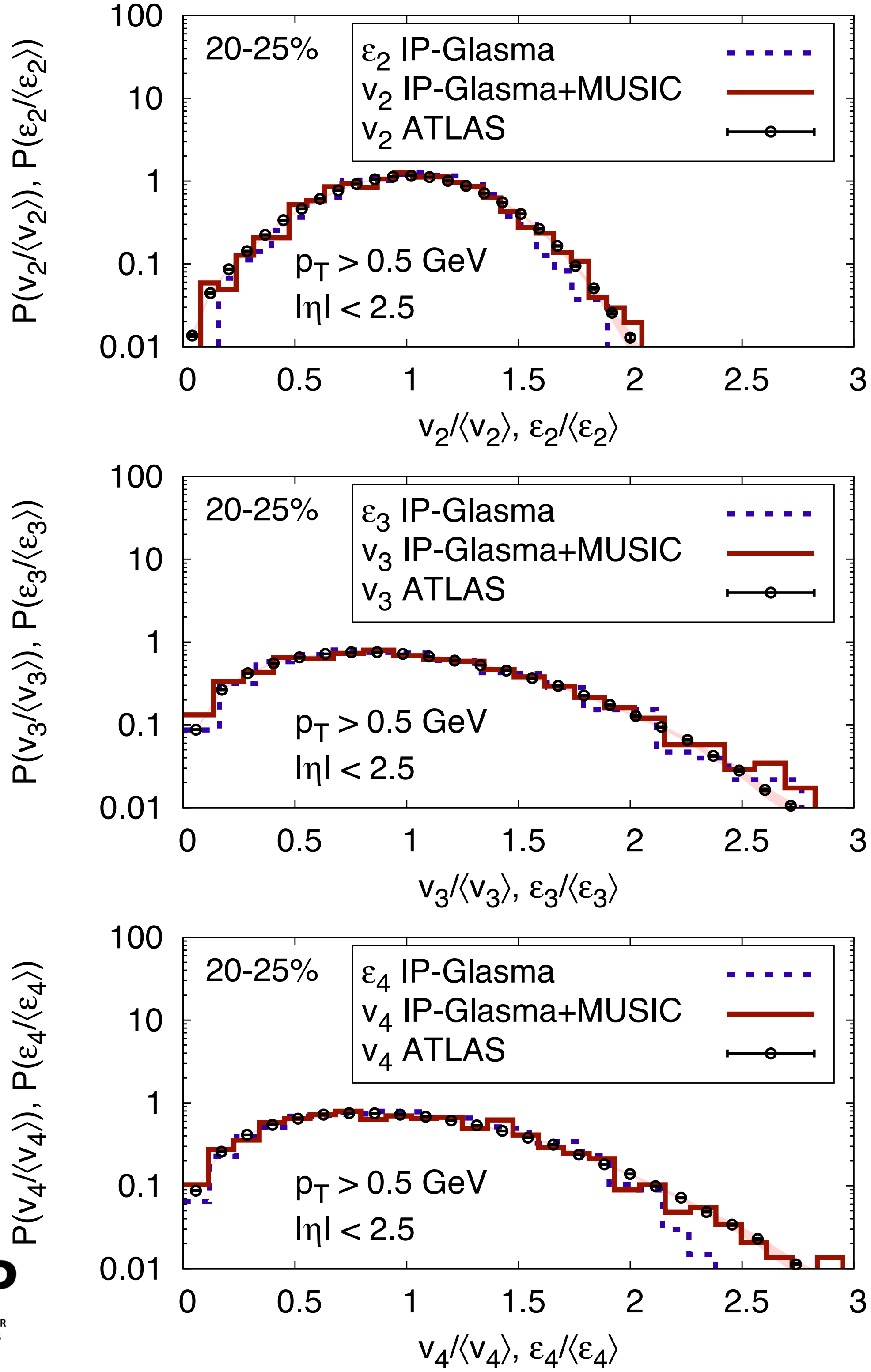
- In a thermal fireball picture, the net photon yield is sensitive to the value of the acceleration parameter, and to details of the initial state. The photons **do** report on the details of the dynamics.
- How uniquely determined are these? How unique is the entire evolution?

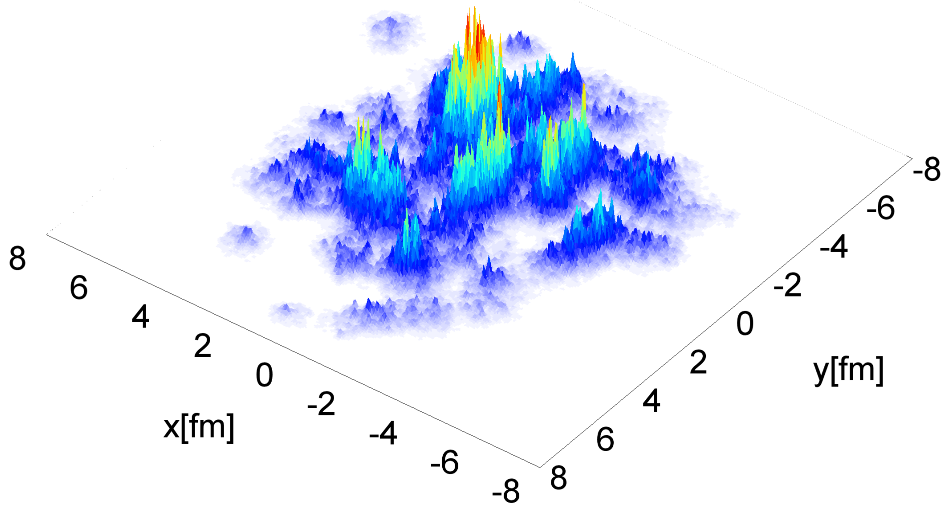


van Hees, Gale, Rapp (2011)

BEYOND GLAUBER: IP-GLASMA + MUSIC
EFFECT ON HADRONIC OBSERVABLES

- Flow harmonics reproduced up to v_5 at RHIC and LHC
- Distributions of v_n at LHC:





- IP-Glasma + MUSIC provides consistent flow systematics at RHIC & LHC
- Contains an initial flow: Investigating the effects on EM variables

Gale, Jeon, Schenke, Tribedy, Venugopalan PRL (2013)

Is the hydrodynamic modelling complete?

- In the last ~5-8 years, relativistic hydrodynamics has undergone a revolution
  - 3D
  - 3D - Shear viscosity
  - 3D - Shear viscosity - Fluctuating initial conditions
  - 3D - Shear viscosity - Fluctuating initial conditions also in y
- What's left?

$$T^{\mu\nu} = -Pg^{\mu\nu} + \omega u^{\mu} u^{\nu} + \Delta T^{\mu\nu}$$

The dissipative terms:

$$\Delta T^{\mu\nu} = \eta\left(\Delta^{\mu} u^{\nu} + \Delta^{\nu} u^{\mu}\right) + \left(\frac{2}{3}\eta - \zeta\right) H^{\mu\nu} \partial_{\rho} u^{\rho} - \chi (H^{\mu\alpha} u^{\nu} + H^{\nu\alpha} u^{\nu}) Q_{\alpha}$$

No simulation incorporates all of these

BULK VISCOSITY?

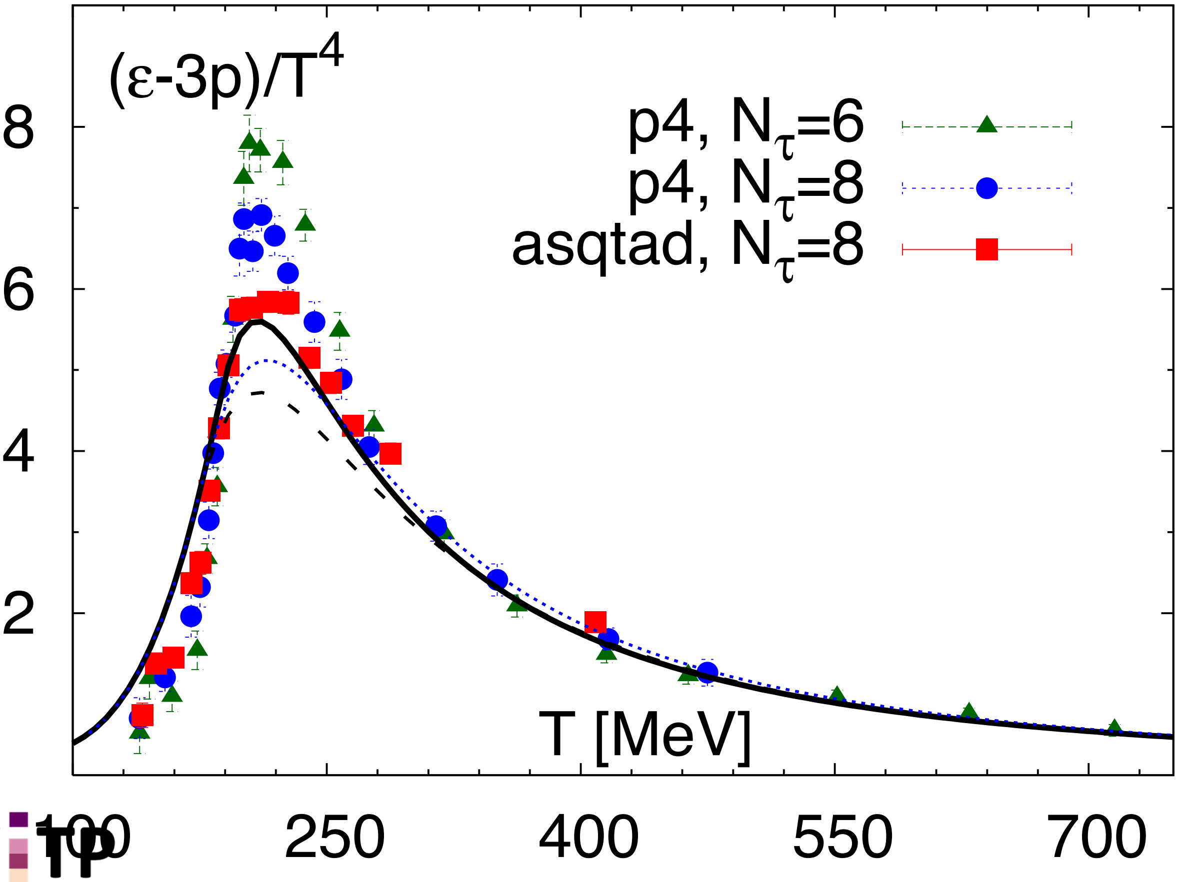
$$\zeta \approx 15\eta\left(\frac{1}{3}-c_s^2\right)^2$$

S. Weinberg, Ap. J (1971)

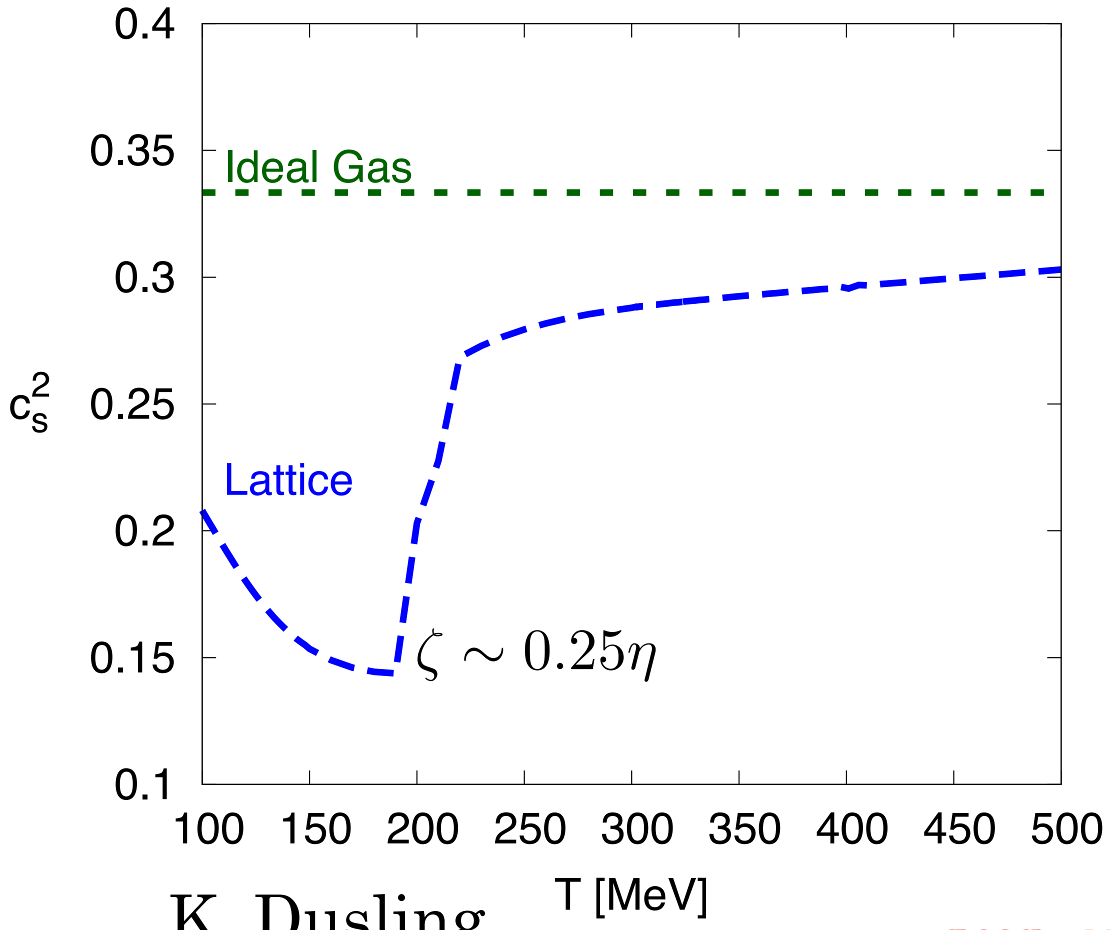
$$\zeta \gtrsim 2\eta\left(\frac{1}{3}-c_s^2\right)$$

A. Buchel, Phys. Lett. (2008)

Bulk viscosity vanishes in conformal fluids. QCD is only very approximately conformal:



Huovinen & Petreczky



K. Dusling

BULK VISCOSITY?

- Quantifies deviations from equilibrium, when the fluids expands or contracts more quickly than the time needed to relax back to equilibrium

$$\Delta T^{\mu\nu} = \eta\left(\Delta^{\mu}u^{\nu} + \Delta^{\nu}u^{\mu}\right) + \left(\frac{2}{3}\eta - \zeta\right)H^{\mu\nu}\partial_{\rho}u^{\rho} - \chi(H^{\mu\alpha}u^{\nu} + H^{\nu\alpha}u^{\mu})Q_{\alpha}$$

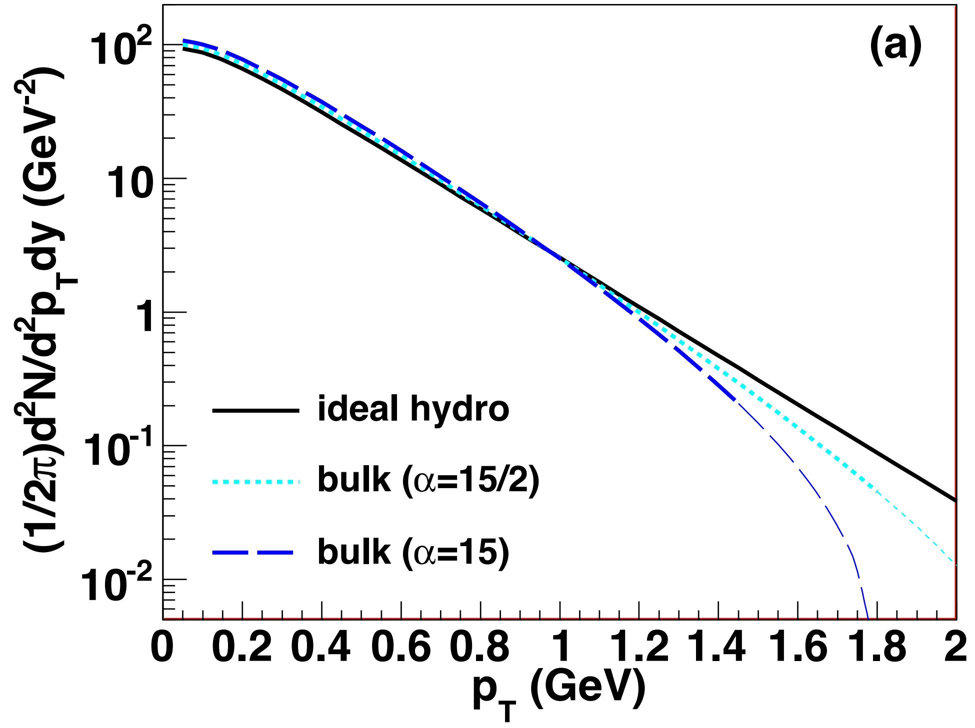
$$\frac{\delta f}{f_0} \sim p_T^2\left(\frac{1}{3} - c_s^2\right)^2(\partial_{\mu}u^{\mu})$$

Relaxation Time Approximation, Dusling & Schäfer (2012)

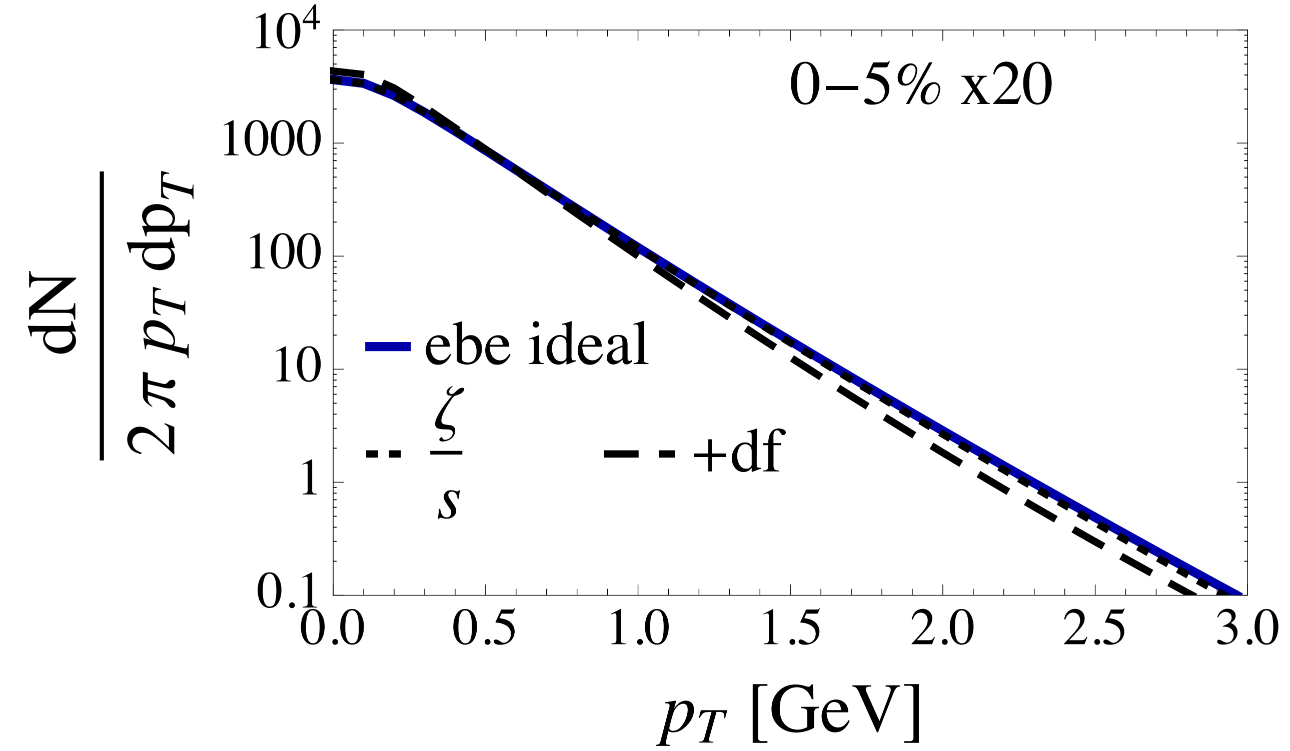
$$\frac{\delta f}{f_0} \sim \zeta\left(\partial_{\mu}u^{\mu}\right)\left[\alpha + \beta u \cdot k + \gamma(u \cdot k)^2\right]$$

Modified Moment Expansion Noronha-Hostler, Denicol, et al., (2013)

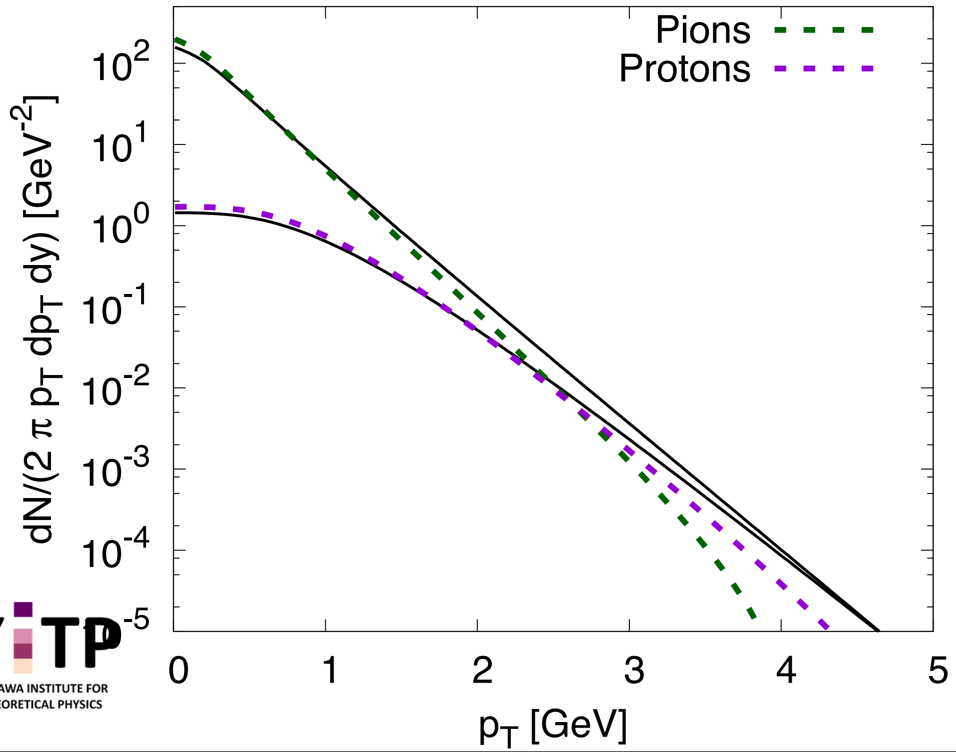
RHIC



Monnai and Hirano (2009)
3+1D Hydro



Noronha-Hostler, Denicol, Andrade, Grassi (2013), 2+1D Hydro



Dusling and Schäfer, (2012)
2+1D Hydro

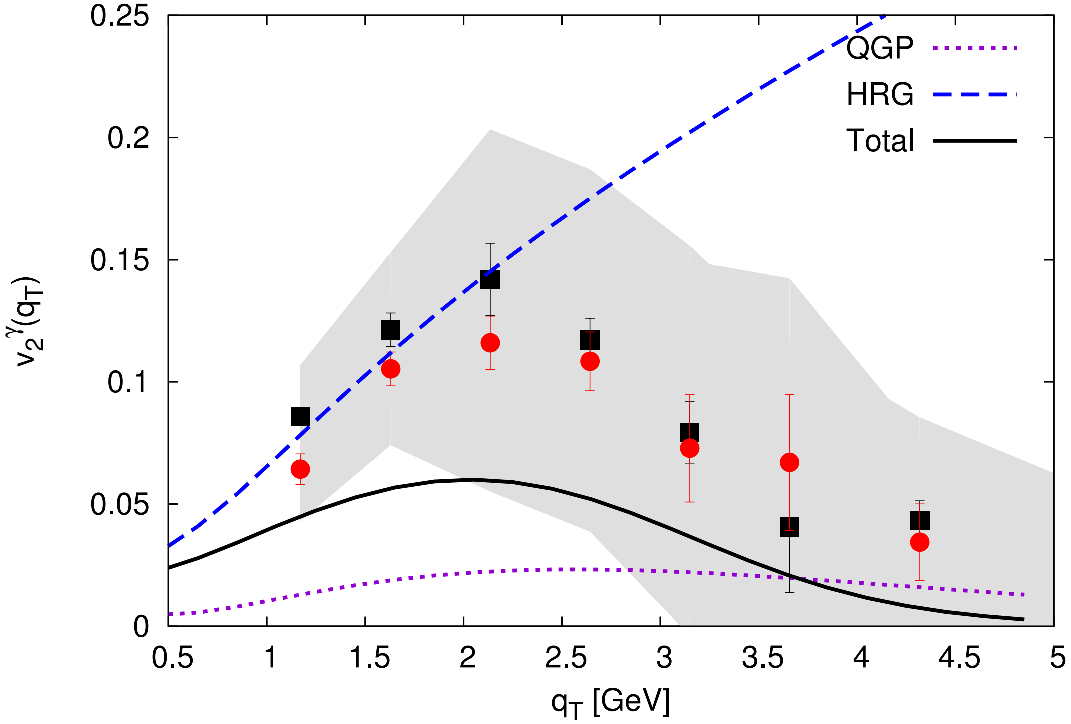
— $\eta/s = 0.16$

_ _ $\eta/s = 0.16, \zeta/s = 0.005$

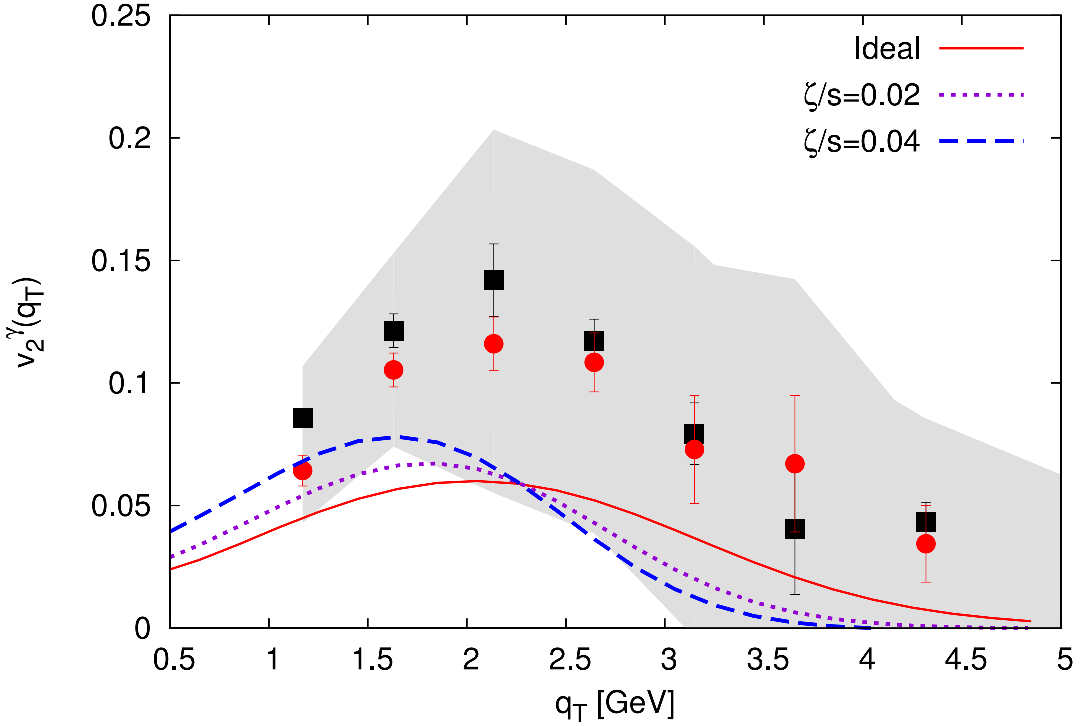
- Spectra are systematically softer
- Details depend on the scheme to implement the viscosity correction(s)
- Some cancellation between shear and bulk effects

BULK VISCOSITY EFFECTS ON PHOTONS?

Ideal photon $v_2(q_T)$



Viscous photon $v_2(q_T)$



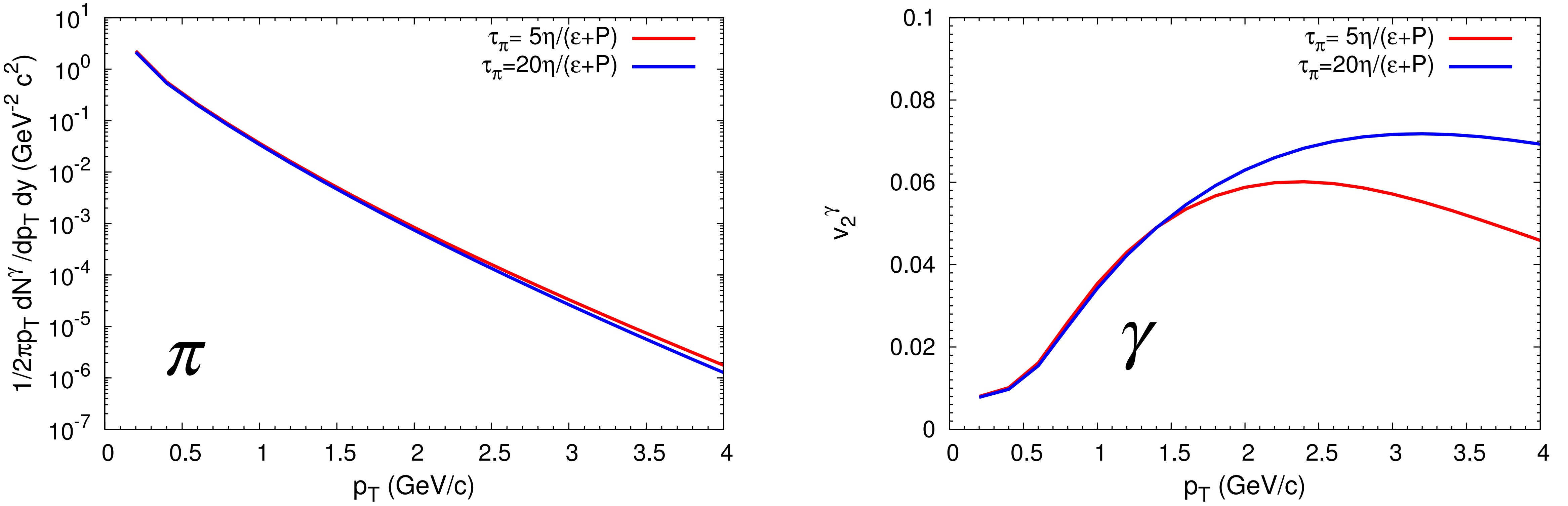
K. Dusling

- Bulk viscosity enhances the elliptic flow
- Effects are however large enough: a <u>consistent</u> inclusion of bulk is warranted.
- $\frac{\zeta}{s}(T)$ etc.. (Gabriel Denicol's talk)

MORE ON THE HYDRO MODELLING AND PHOTON PRODUCTION

$$\tau_\pi \dot{\pi}^{\langle \mu\nu \rangle} + \pi^{\mu\nu} == 2\eta\sigma^{\mu\nu} - \frac{4}{3}\tau_\pi \pi^{\mu\nu}\theta$$

- Can the relaxation time be changed? Does this affect anything?



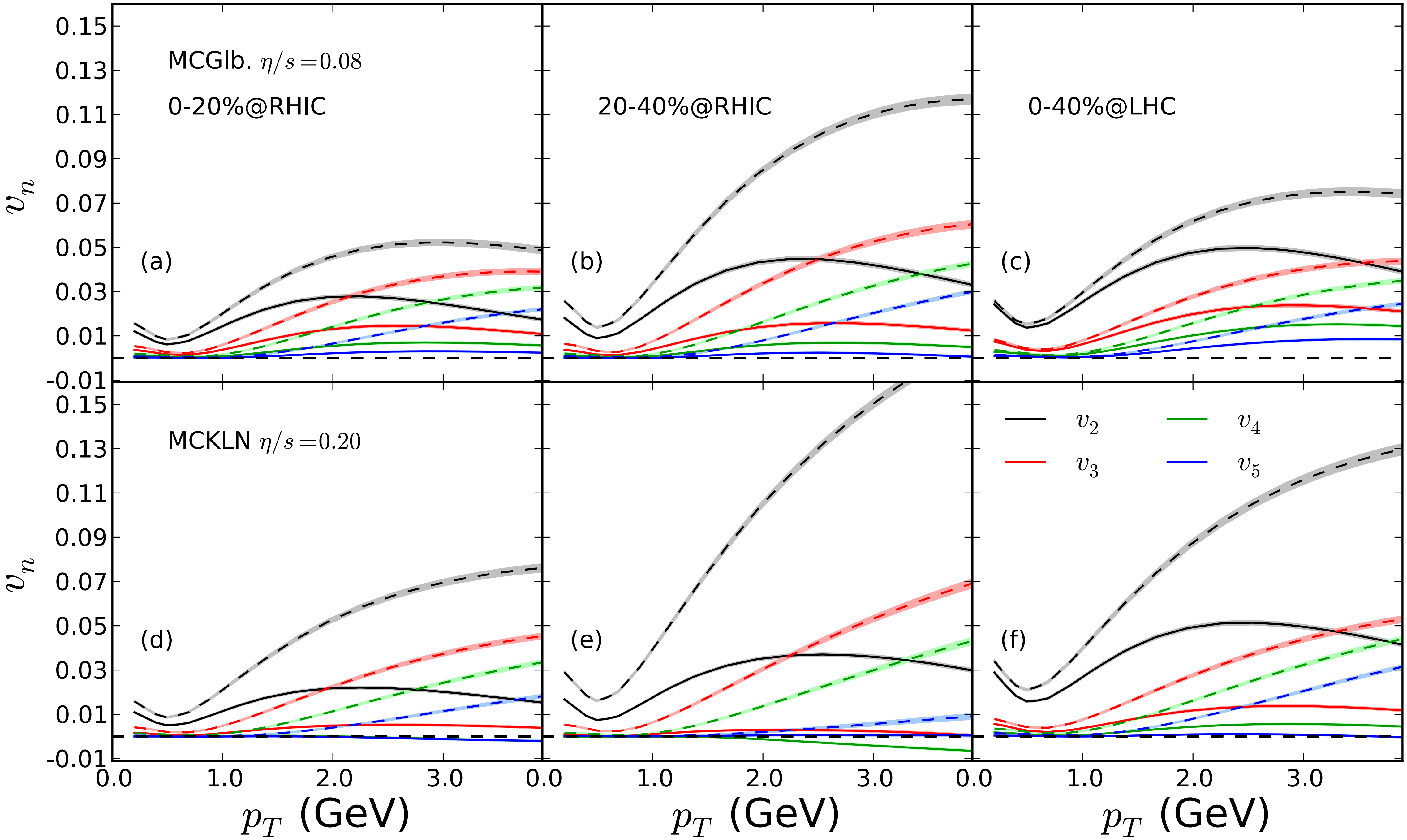
Paquet, Vujanovic, Denicol et al. (2013)

Photons are sensitive to early time dynamics

THERMAL PHOTONS AS A VISCOMETER

$$\omega \frac{dR}{d^3q} = \Gamma_0 + \frac{\pi^{\mu\nu}}{2(\epsilon + P)} \Gamma_{\mu\nu}(p,T)$$

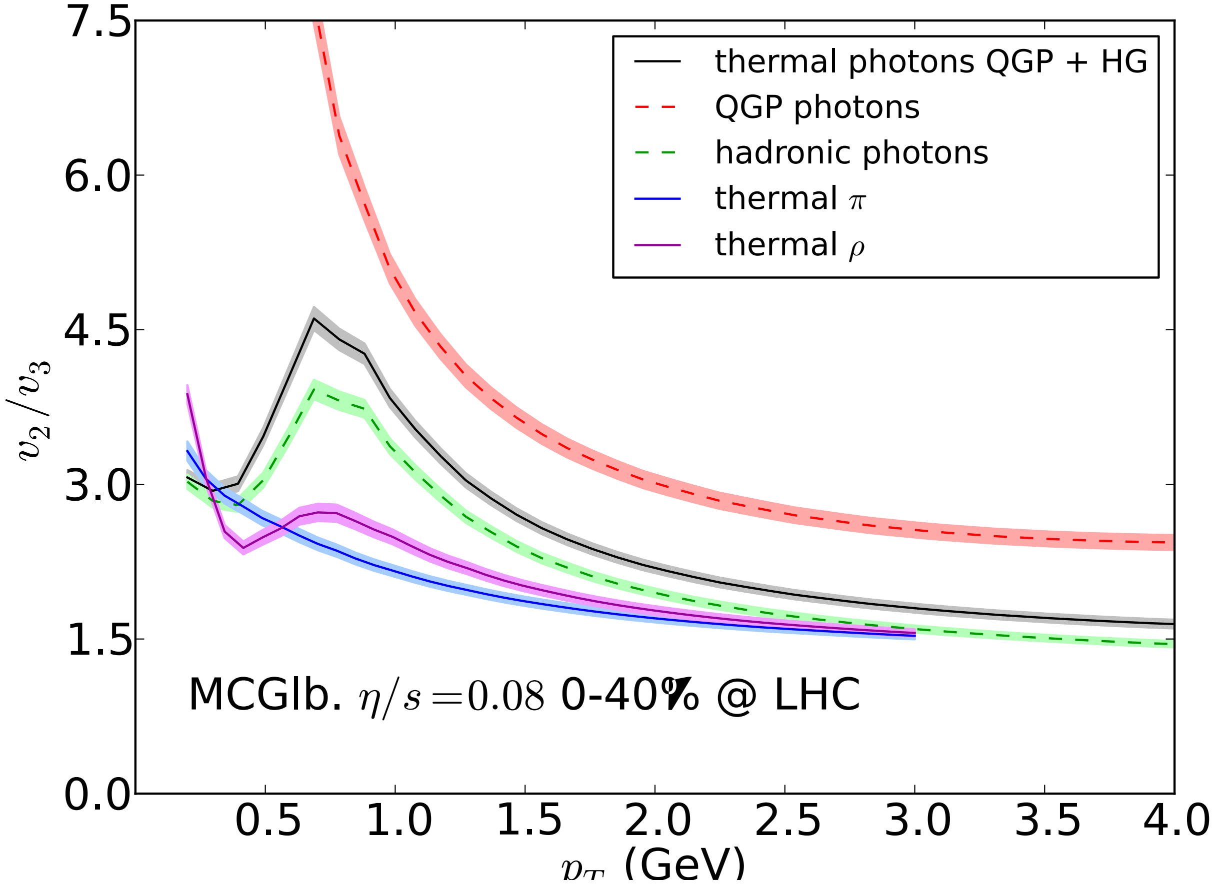
Shen, Heinz, Paquet, Kozlov, Gale, arXiv 1308.2111

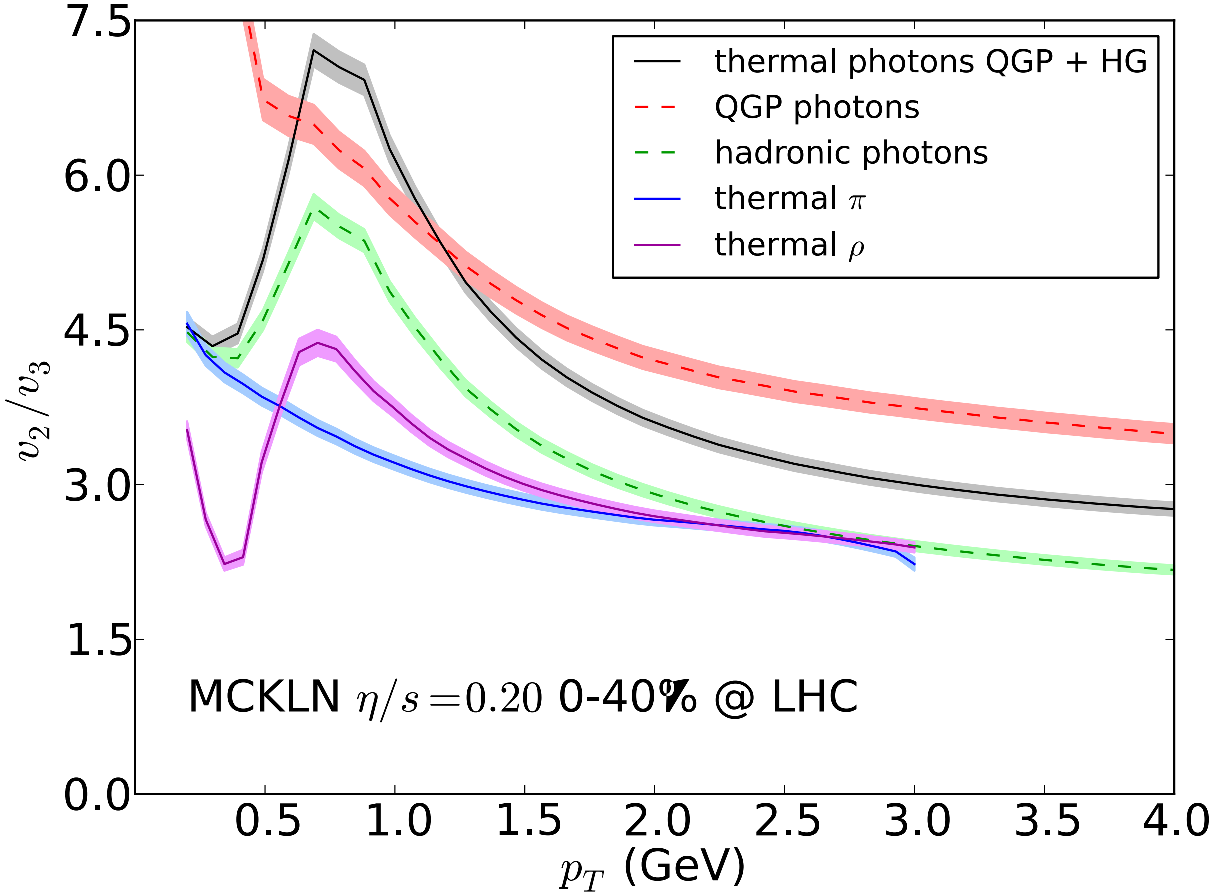


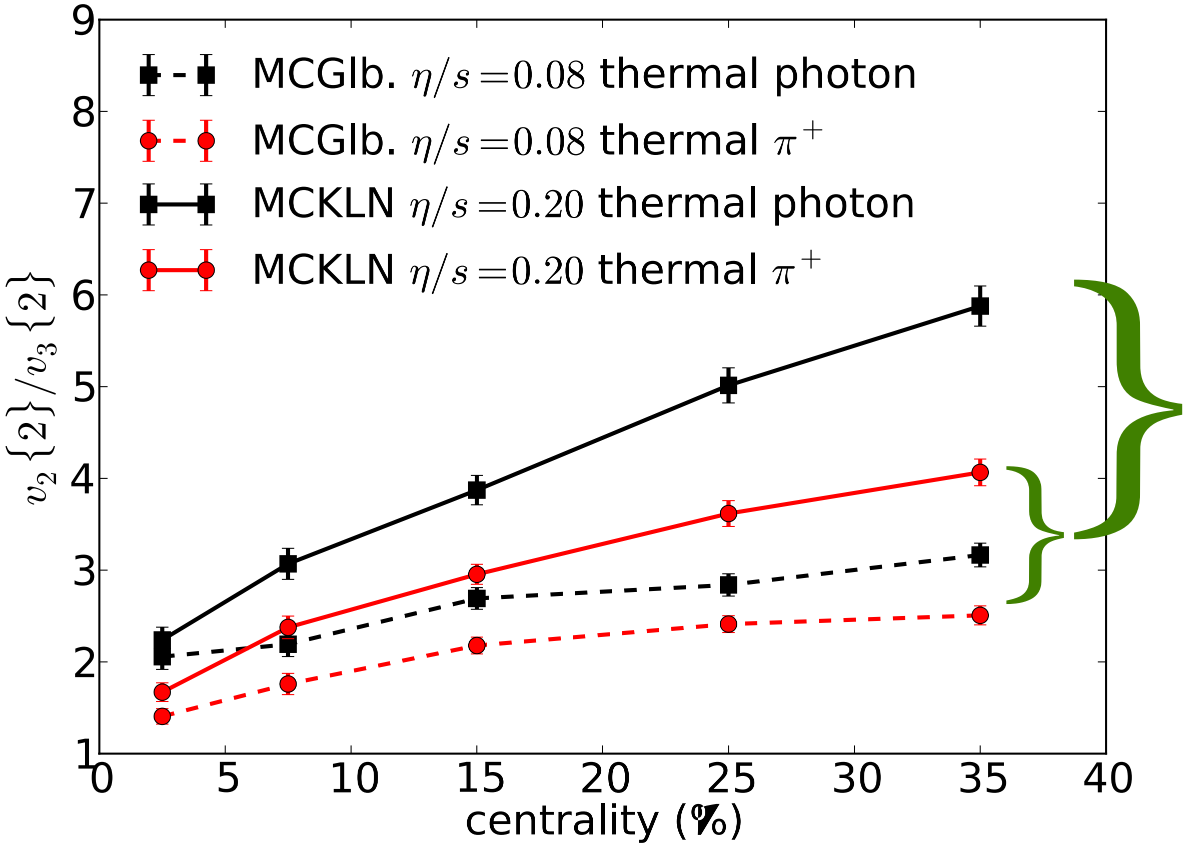
——— Viscous - - - - - Ideal

Charles Gale

MAXIMIZING THE EFFECT



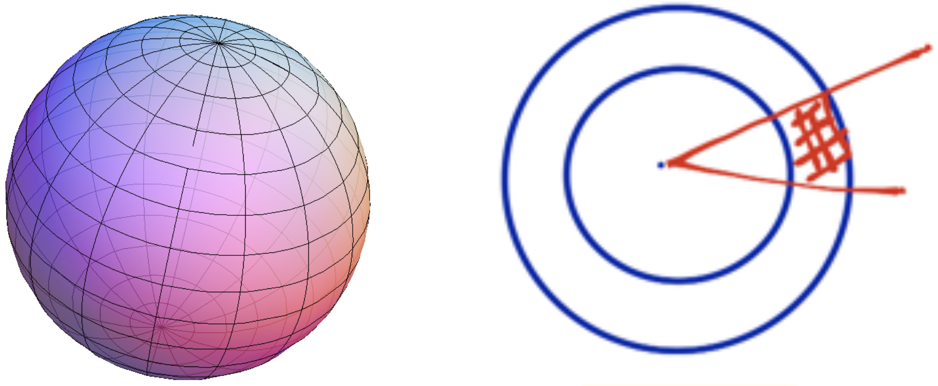




- Slope of ratio vs centrality grows with viscosity
- The ratio has stronger centrality dependence than for hadrons: photons access earlier times with larger viscous tensor

PHOTONS AS A THERMOMETER

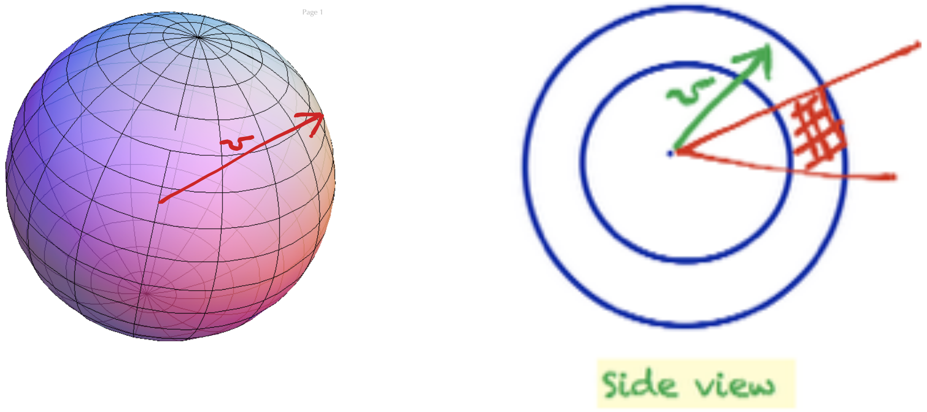
Suppose a static source at temperature T:



$$E\frac{d^3n}{d^3p} = E\,e^{-\beta E}$$

Read off the temperature from the exponent

Suppose an expanding source at local temperature T:

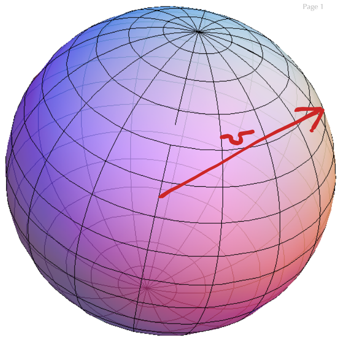


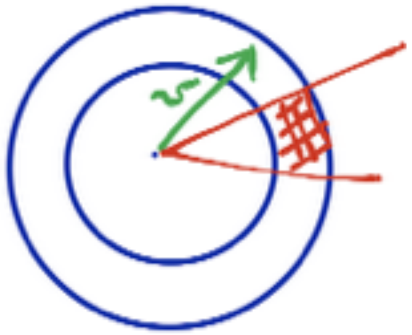
$$E\frac{d^3n}{d^3p} \approx E\,e^{-\beta\gamma E+\beta\gamma v E}$$

$$\boxed{T_e = \sqrt{\frac{1+v}{1-v}}\,T}$$

The effective temperature (deduced from the slope) is not the true temperature

Suppose an expanding source at local temperature T:





Side view

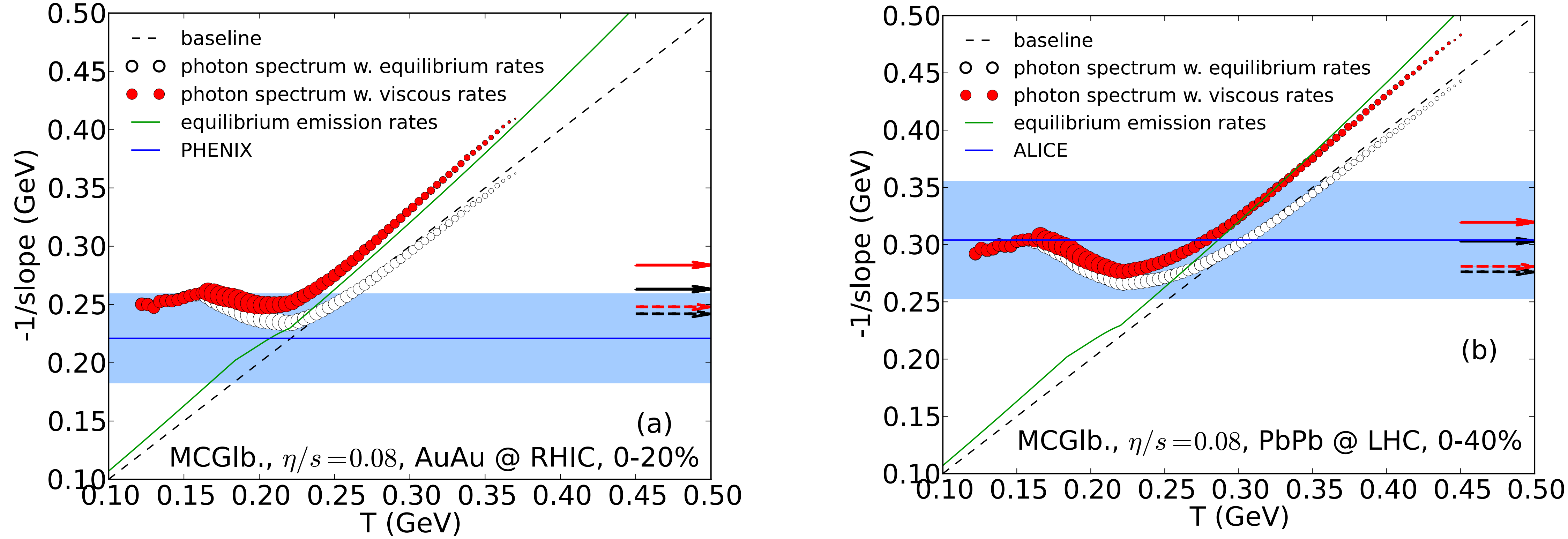
$$E\frac{d^3n}{d^3p} \approx E\,e^{-\beta\gamma E+\beta\gamma v E}$$

$$T_e = \sqrt{\frac{1+v}{1-v}}\,T$$

Doppler shift

The effective temperature (deduced from the slope) is <u>not</u> the true temperature

USING A HYDRO SIMULATION



Shen, Heinz, Paquet, Gale, arXiv 1308.2440
van Hees, Gale, Rapp, PRC (2011)

- The apparent temperature deviates from the true temperature: flow contamination
- The system does go through regions with T>>T$_c$, but a model is needed to extract T

WHAT ABOUT DILEPTONS?

- Additional degree of freedom: M and p_T may be varied independently
- Same approach as for photons: integrate rates with hydro

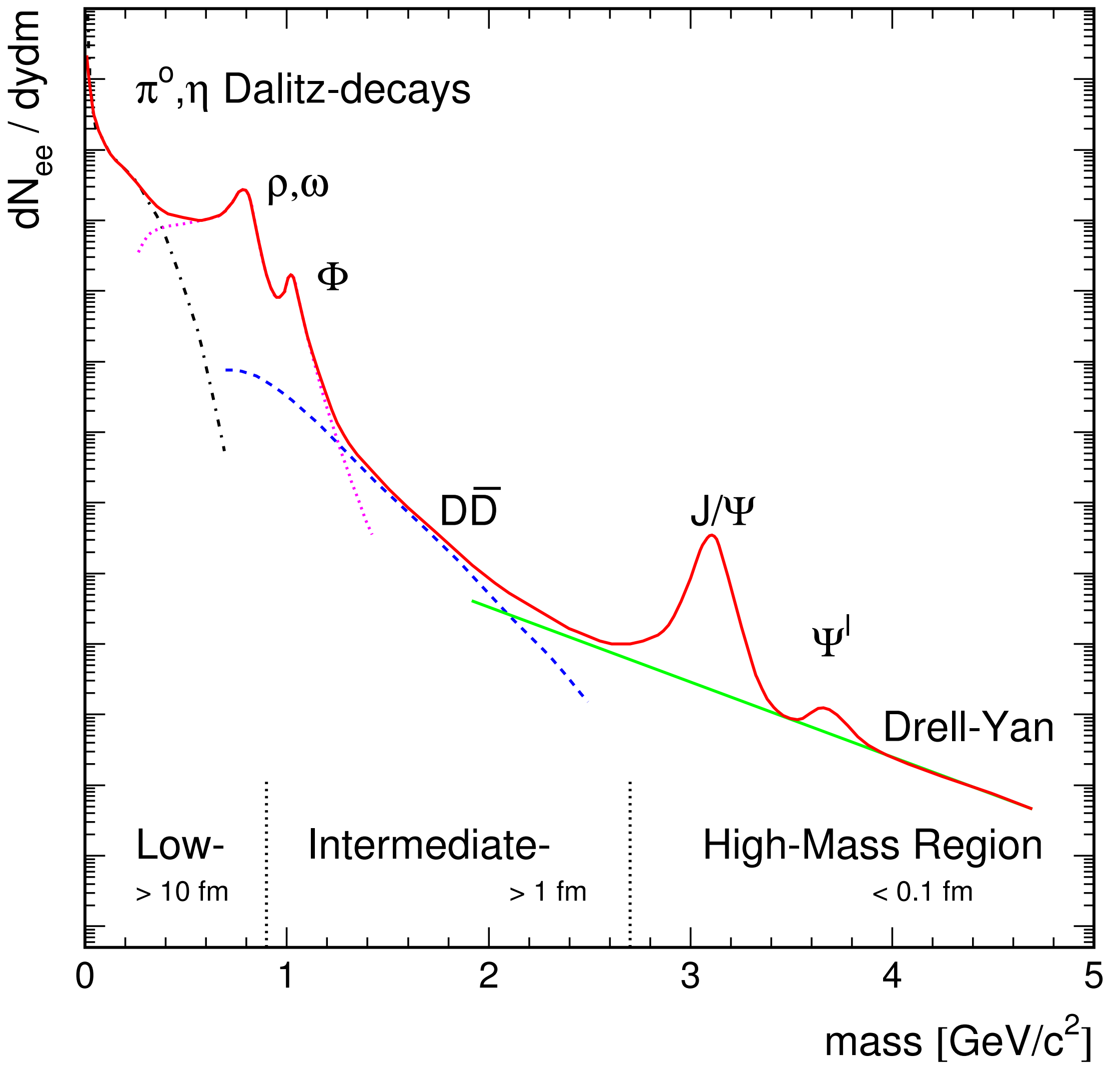
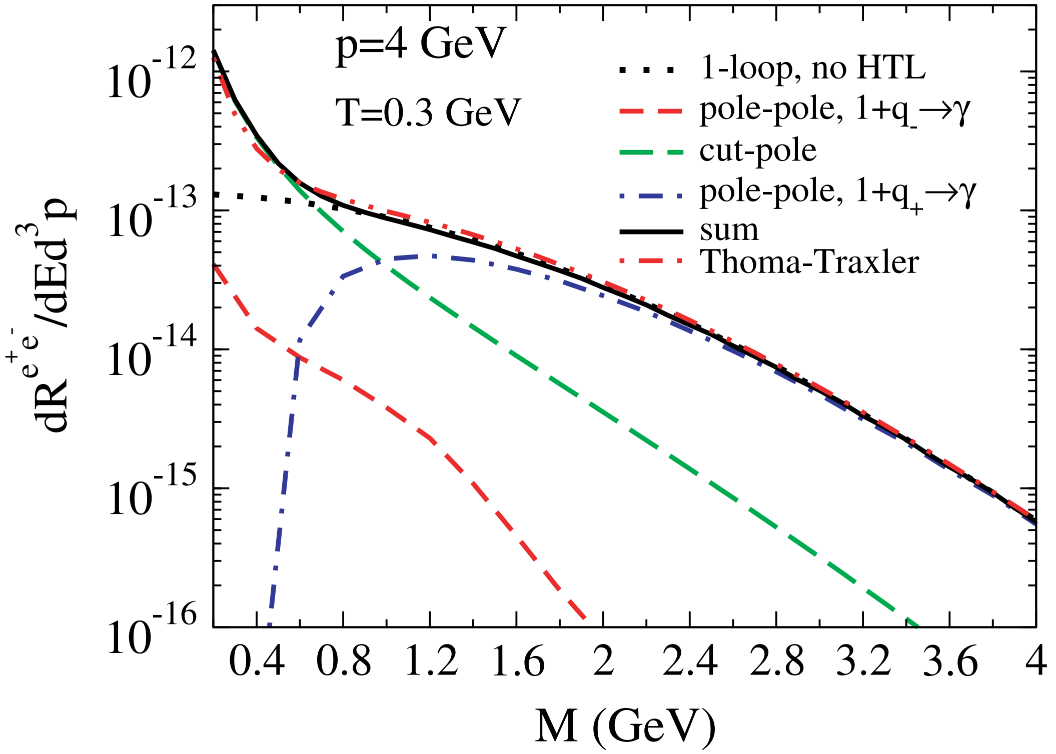


Figure: A. Drees

THERMAL DILEPTON SOURCES, QGP

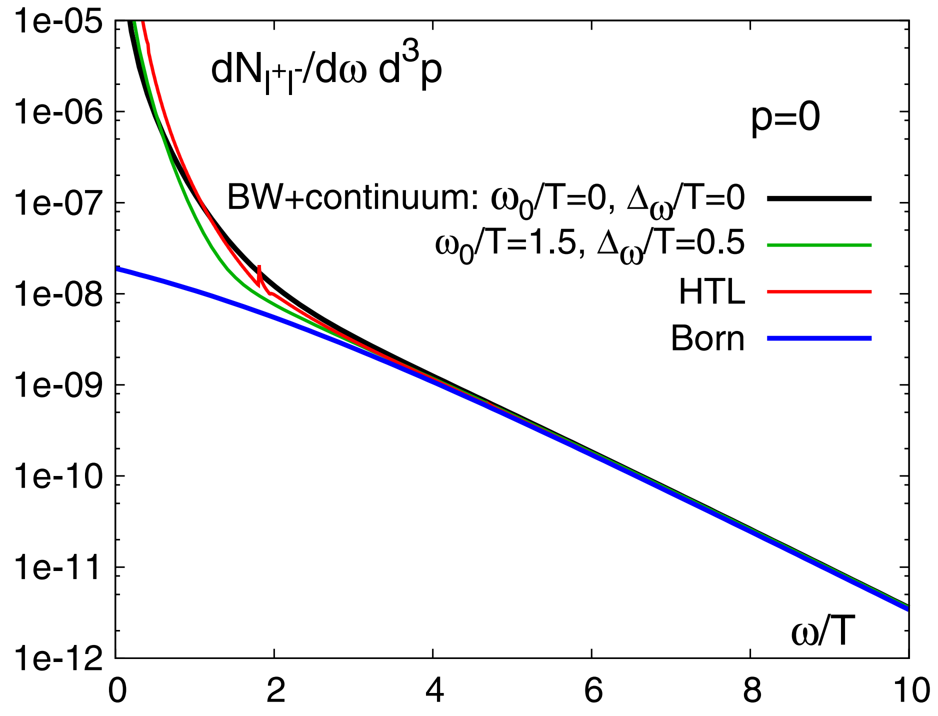
- HTL at zero momentum: Braaten, Pisarski and Yuan, PRL (1990)
- 2-loop, p=0, E>>T: Majumder and Gale, PRC (2002)
- HTL, *M~gT, E>T*: Aurenche, Gélis, Moore, Zaraket, JHEP (2008)

- HTL at finite momentum:



Turbide, Gale, Srivastava, Fries PRC (2006)

- Non-perturbative calculation:



Ding et al., PRD (2011)

No single calculation covers the entire dilepton kinematical phase space

M. Laine, arXiv:1310.0164

$$M^2 \gtrsim (\pi T)^2,\ p \neq 0$$

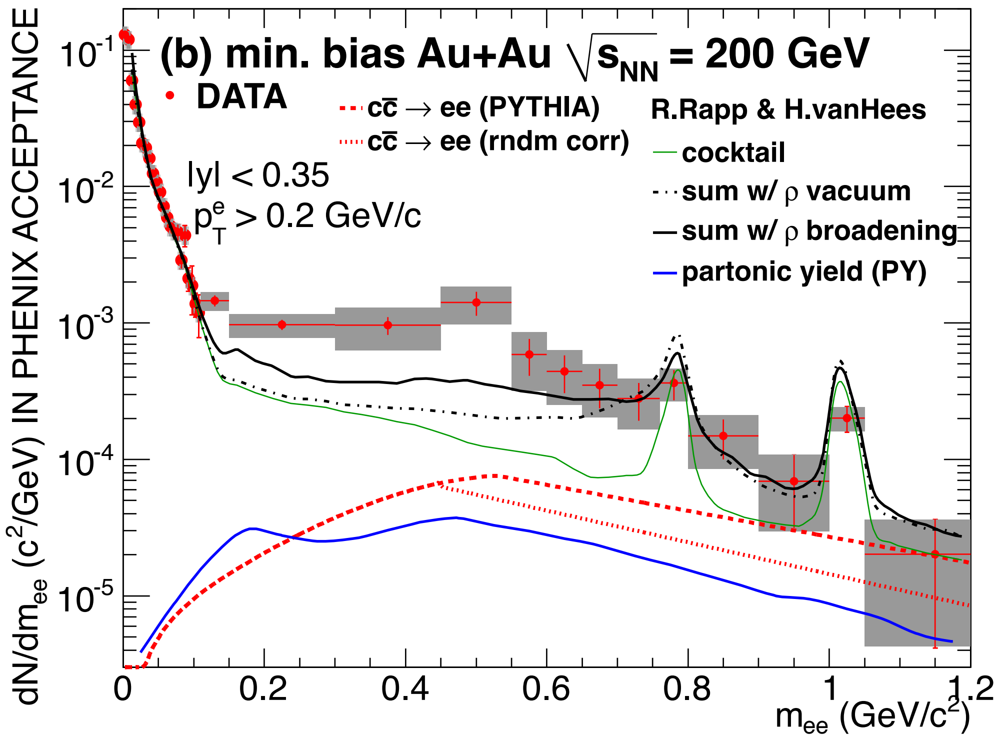
THERMAL DILEPTON SOURCES, HG

- HG contribution: calculate the in-medium vector spectral density
  - Many-Body approach with hadronic effective Lagrangians
    - Rapp and Wambach, ANP (2000)
  - Empirical evaluation of the vector mesons forward-scattering amplitudes

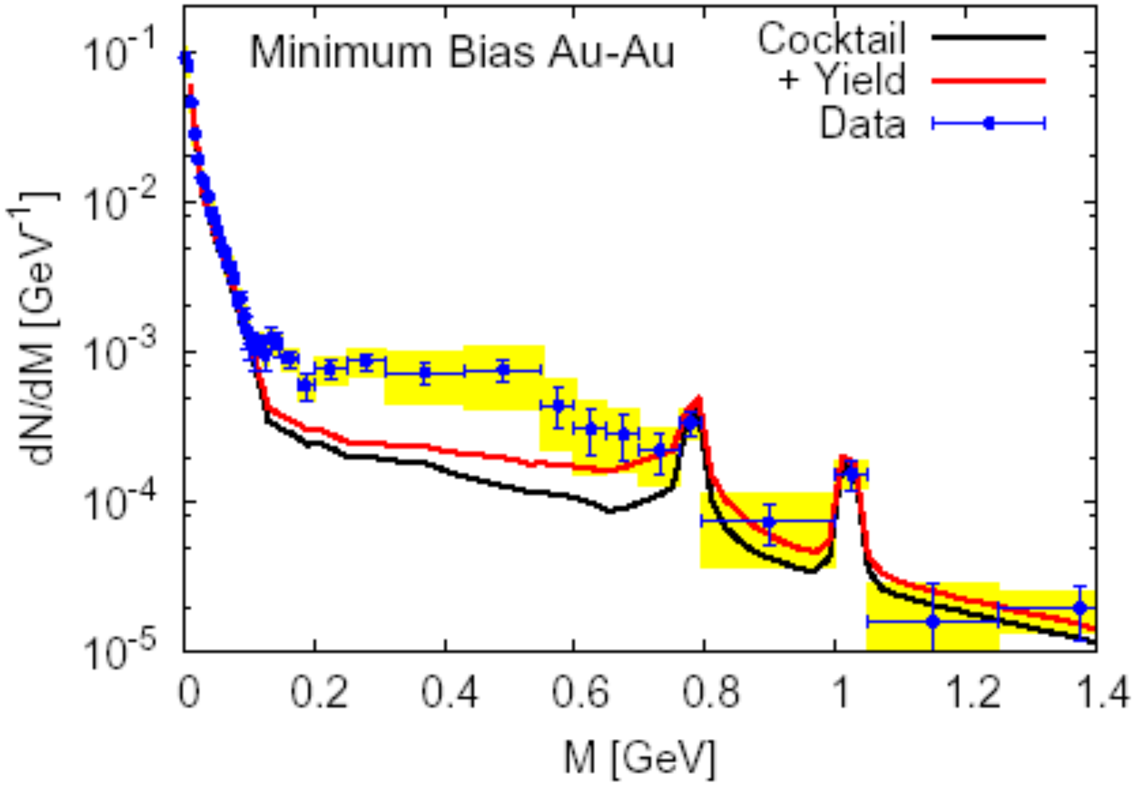
$$\Pi_{ab}(E,p) = -4\pi \int \frac{d^3k}{(2\pi)^3} n_b(\omega) \frac{\sqrt{s}}{\omega} f_{ab}^{\text{c.m.}}(s)$$

    - E. Shuryak, NPA (1991)
    - Eletsky, Ioffe, Kapusta (1999)
    - Vujanovic, Gale (2009)
  - Chiral Reduction formulae
    - Yamagishi, Zahed (1996)

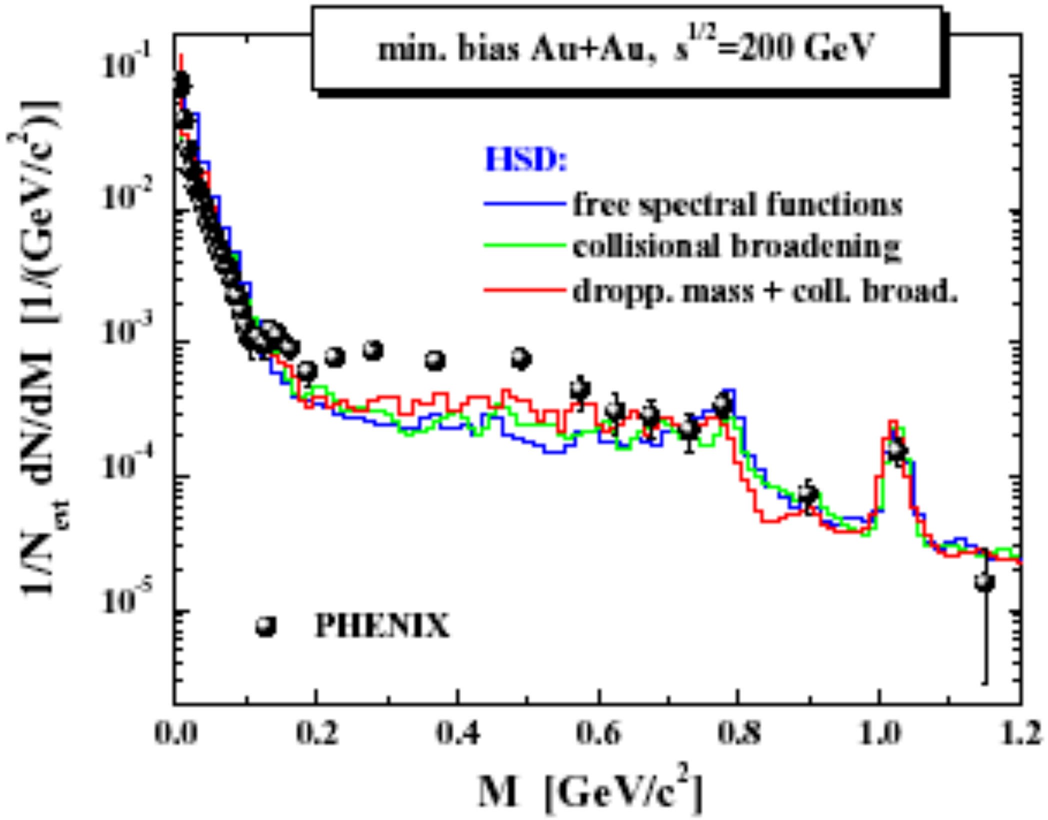
DILEPTONS, THE STORY AS OF A YEAR AGO



van Hees, Rapp (2010)



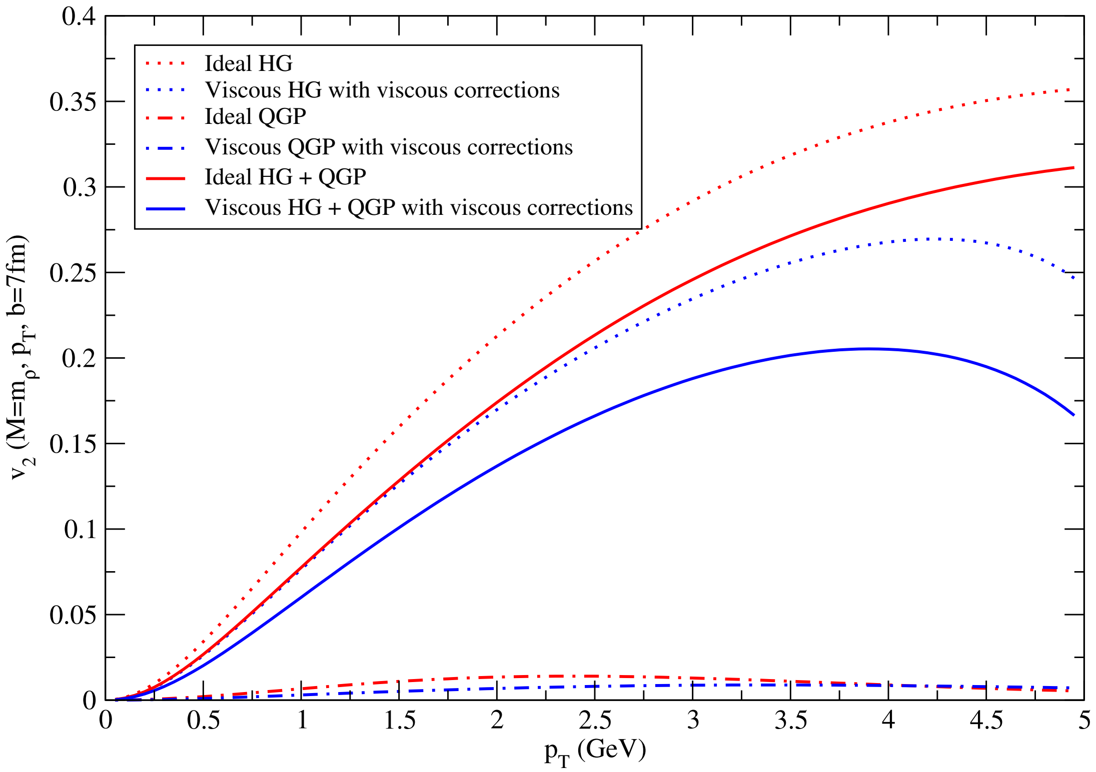
Dusling, Zahed (2009)

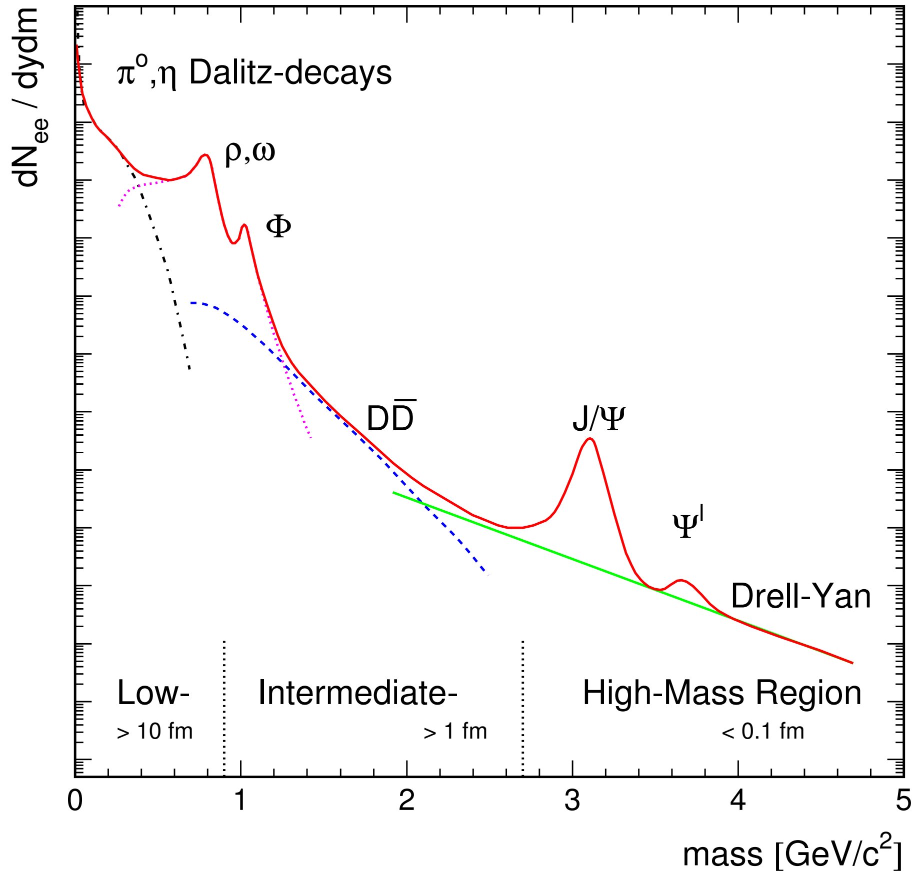


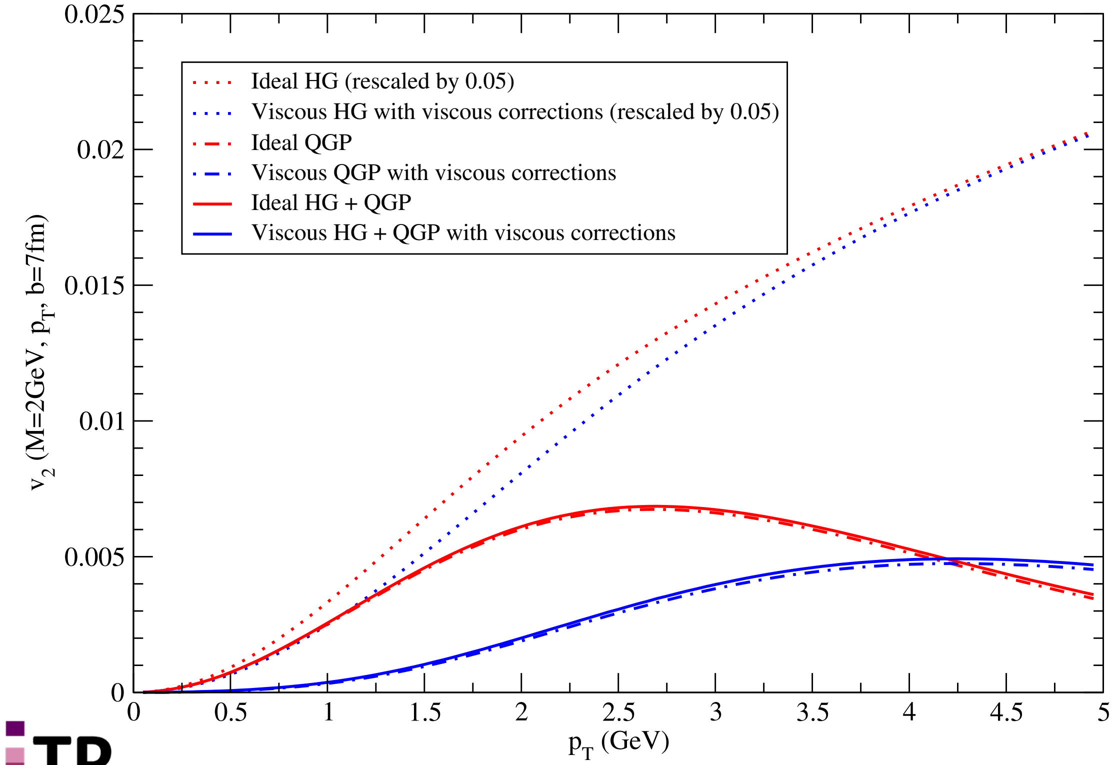
Bratkovskaya, Cassing, Linnyk (2012)

DILEPTON V2? [R. CHATTERJEE ET AL., PRC (2007)]

THERMAL DILEPTON V2 WITH VISCOUS EFFECTS







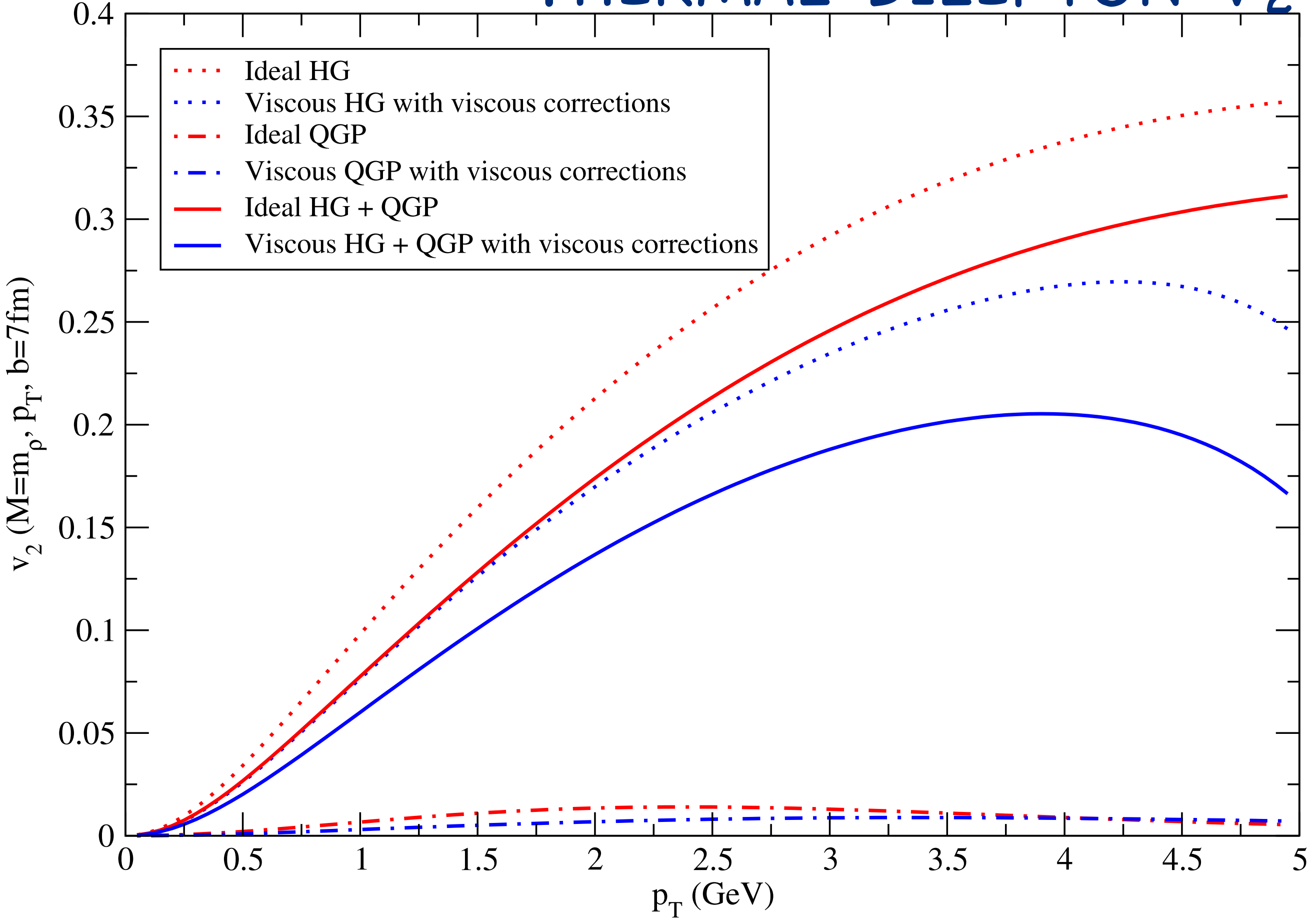
- Low M: HG-dominated
- High-M: QGP dominated

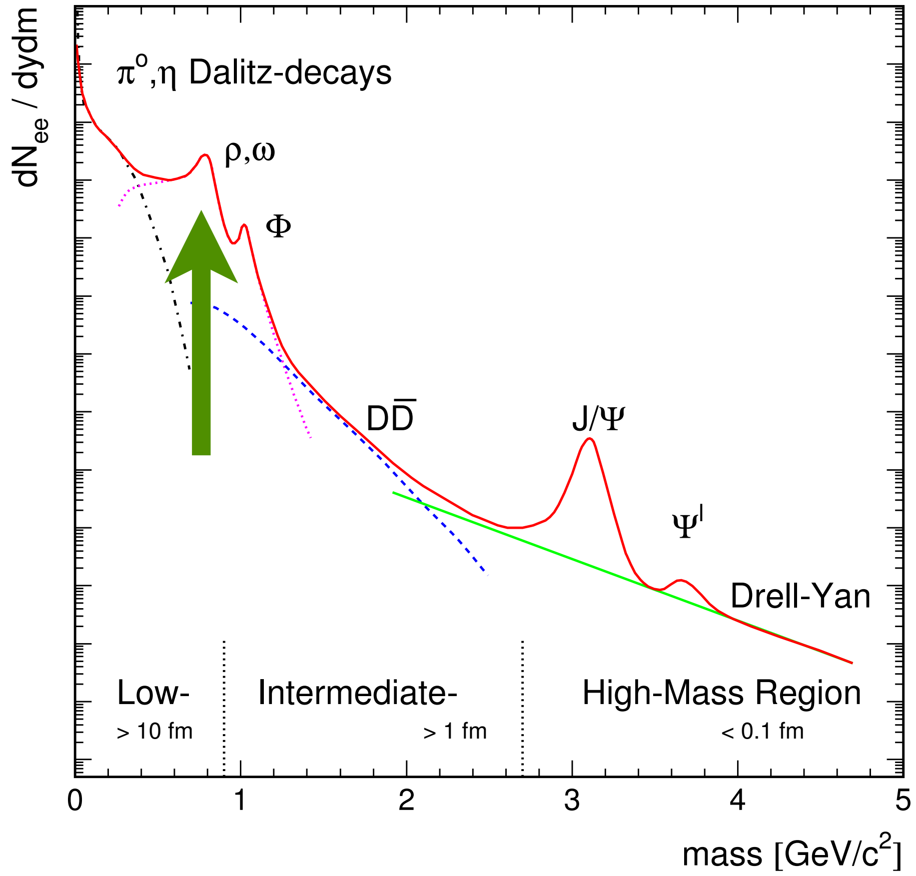
- No open charm here
- $v_2(p_T)$ for different M's contain info on the transition regime
- Viscous effects are moderate

G. Vujanovic et al.,arXiv:1312.0676

DILEPTON V₂? [R. CHATTERJEE ET AL., PRC (2007)]

THERMAL DILEPTON V₂ WITH VISCOUS EFFECTS





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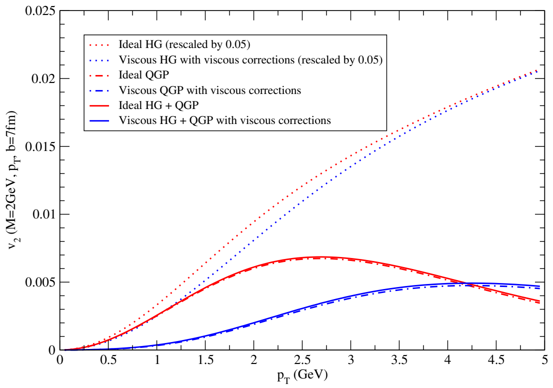

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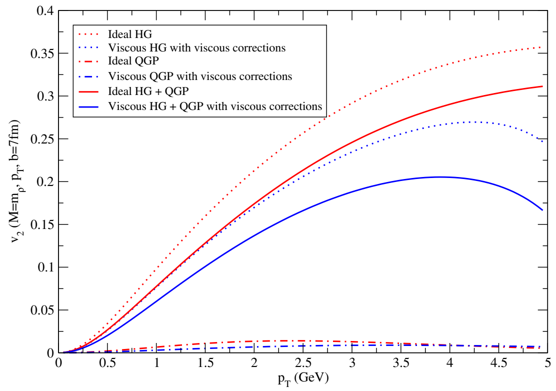
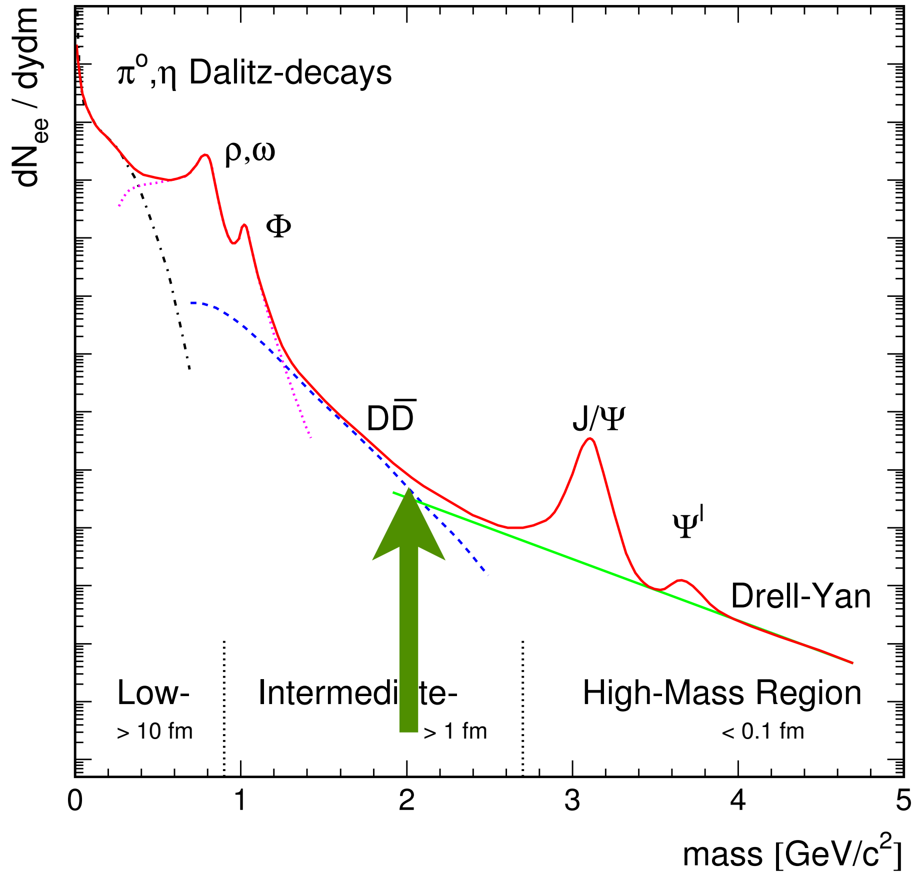
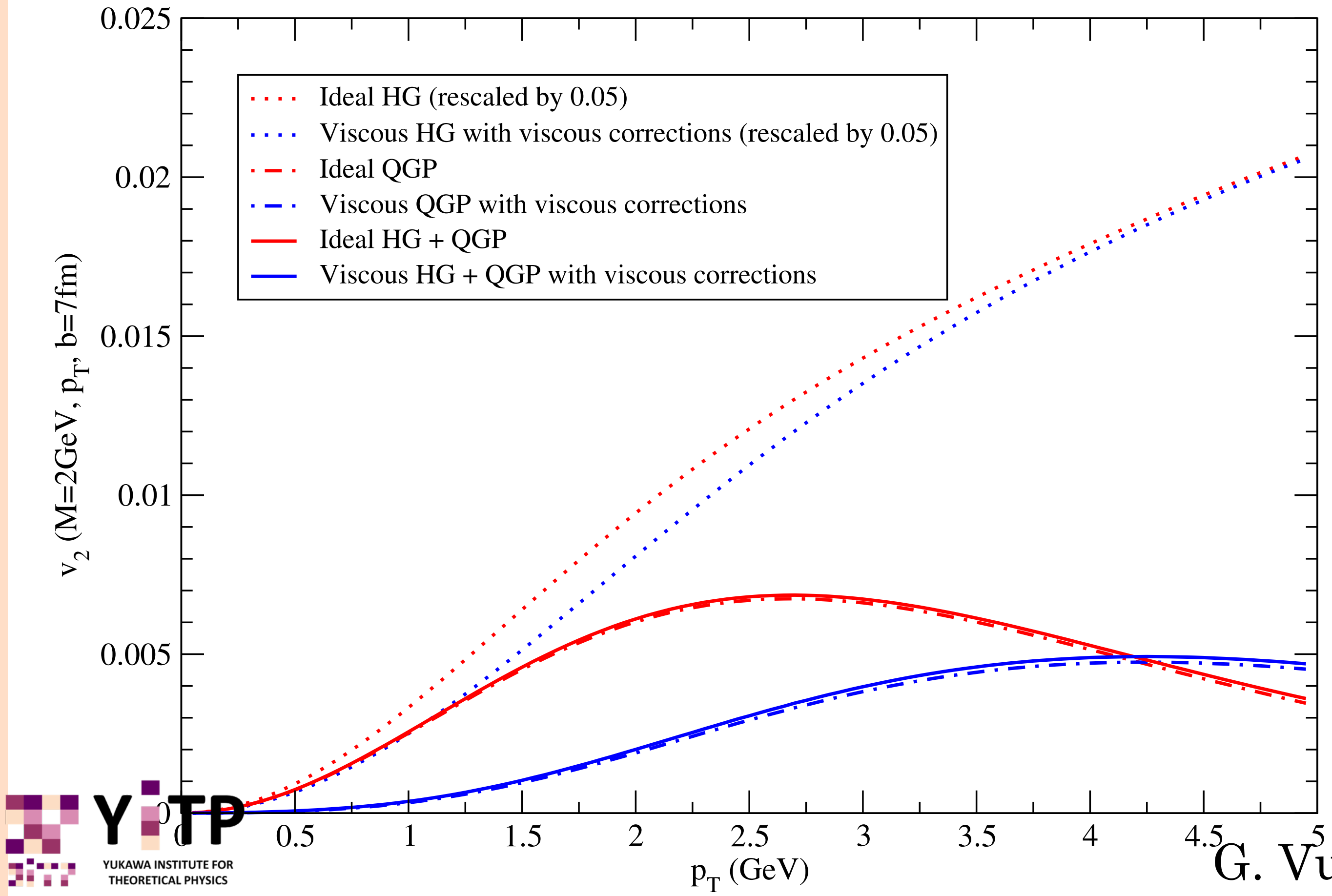
DILEPTON V2? [R. CHATTERJEE ET AL., PRC (2007)]

THERMAL DILEPTON V2 WITH VISCOUS EFFECTS







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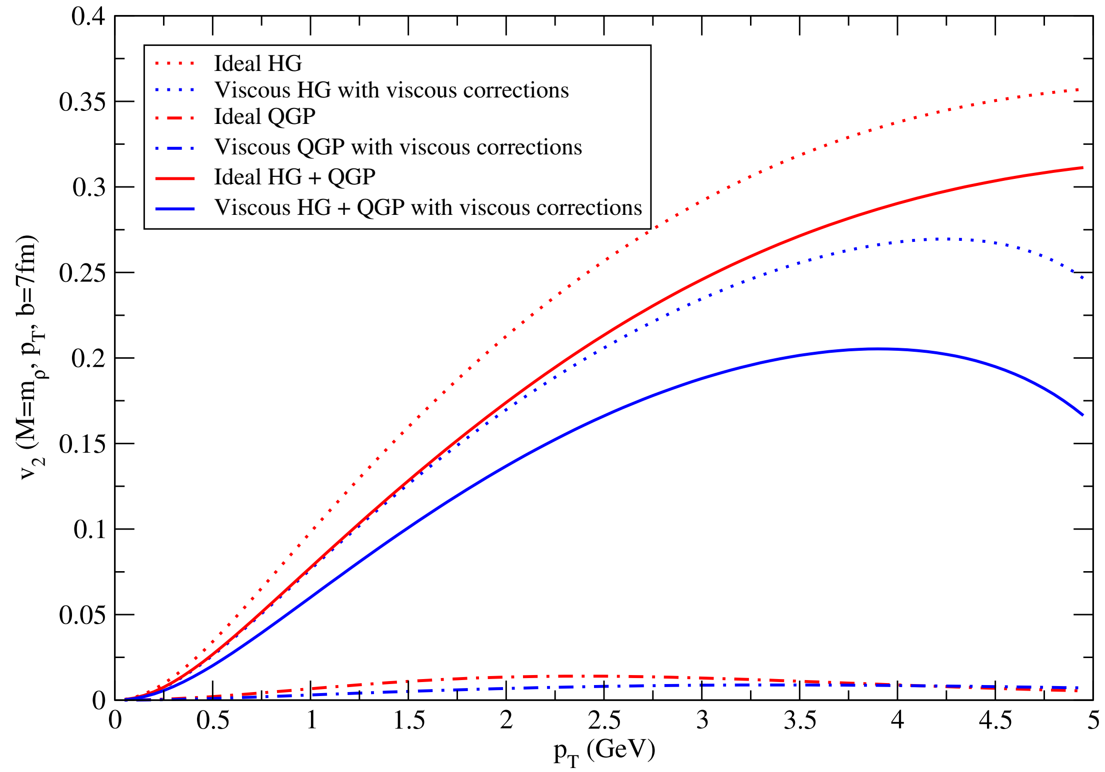
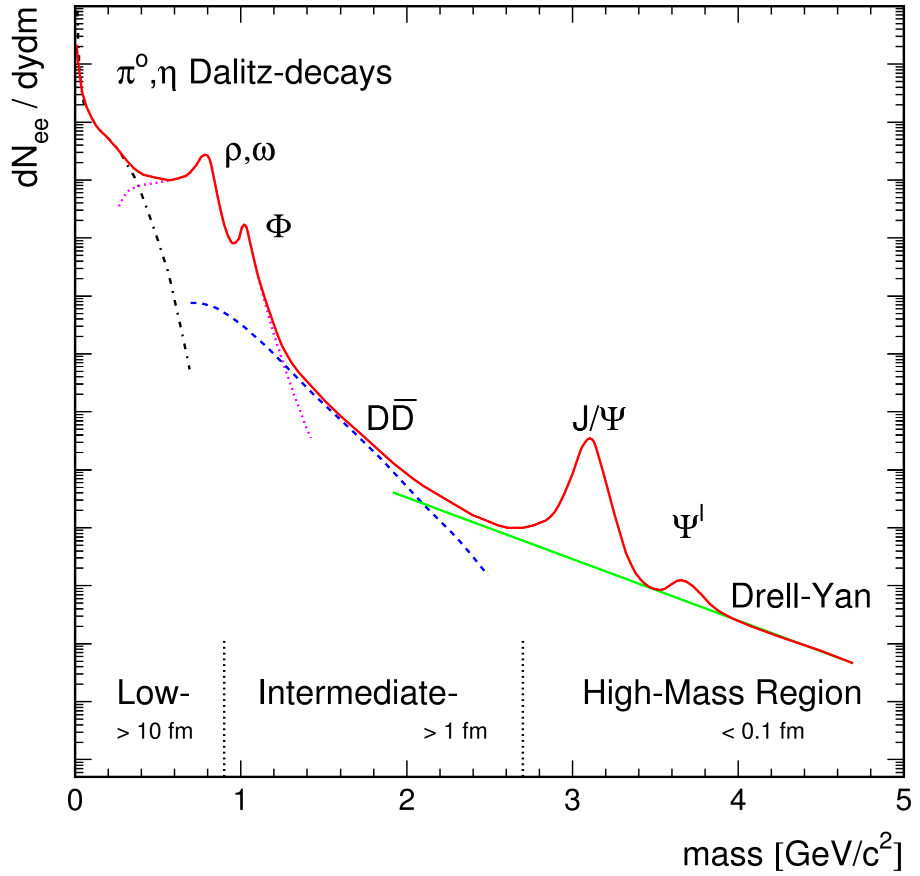
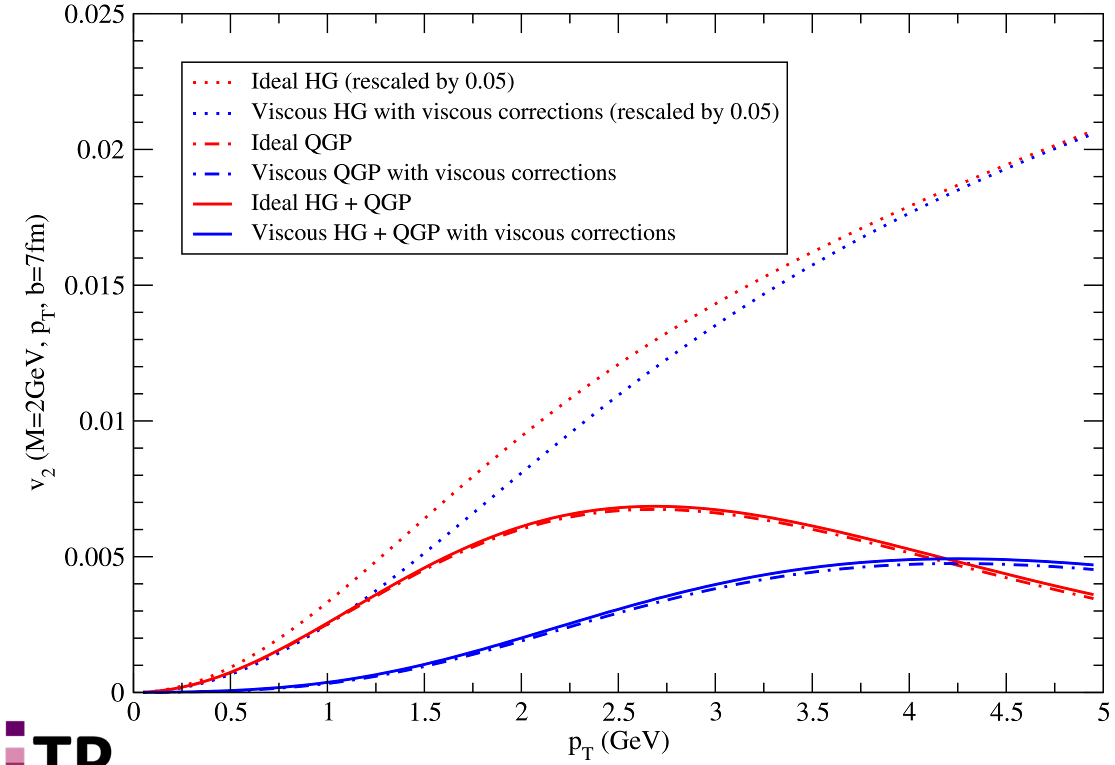
DILEPTON V_2? [R. CHATTERJEE ET AL., PRC (2007)]

THERMAL DILEPTON V_2 WITH VISCOUS EFFECTS







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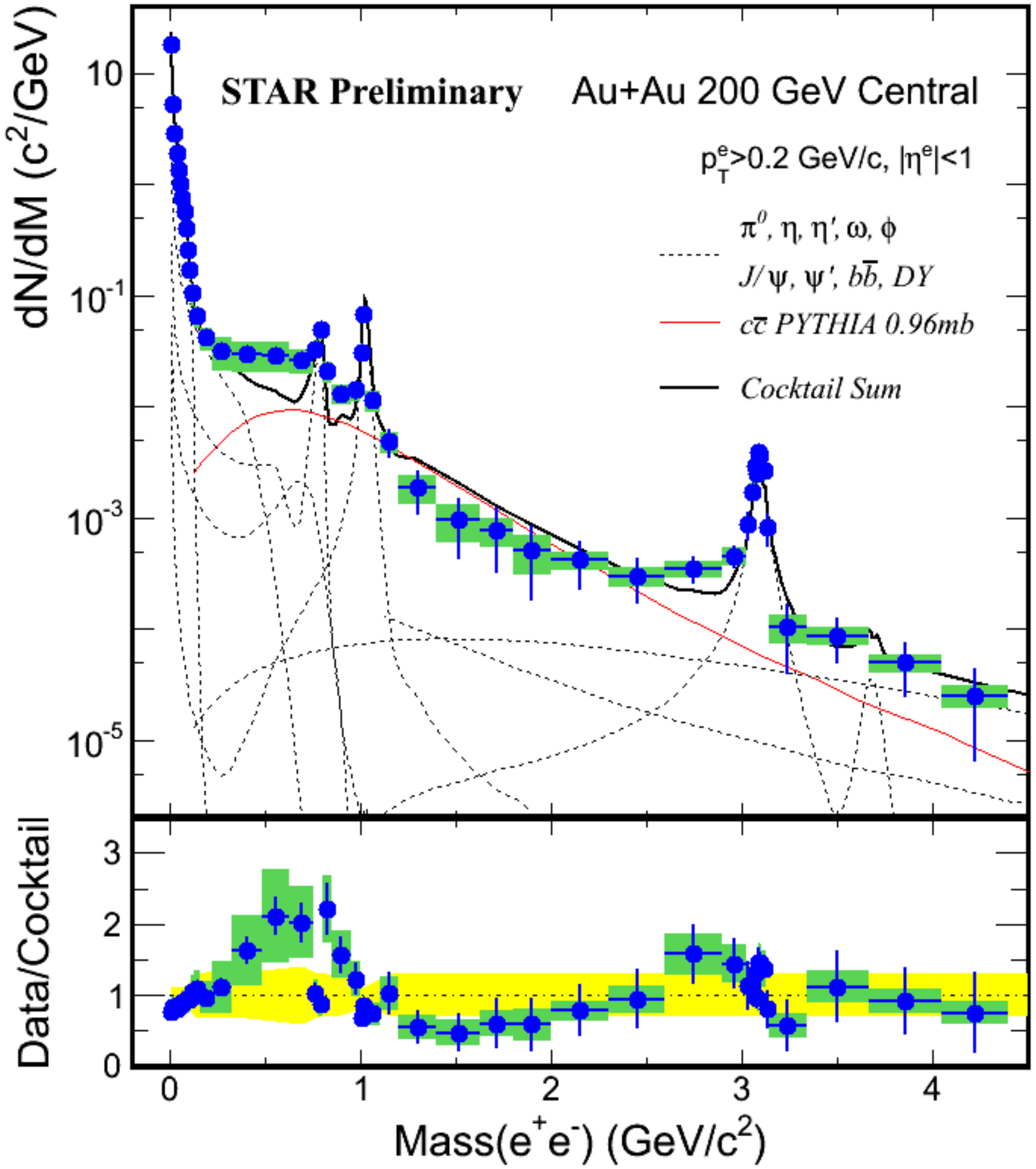
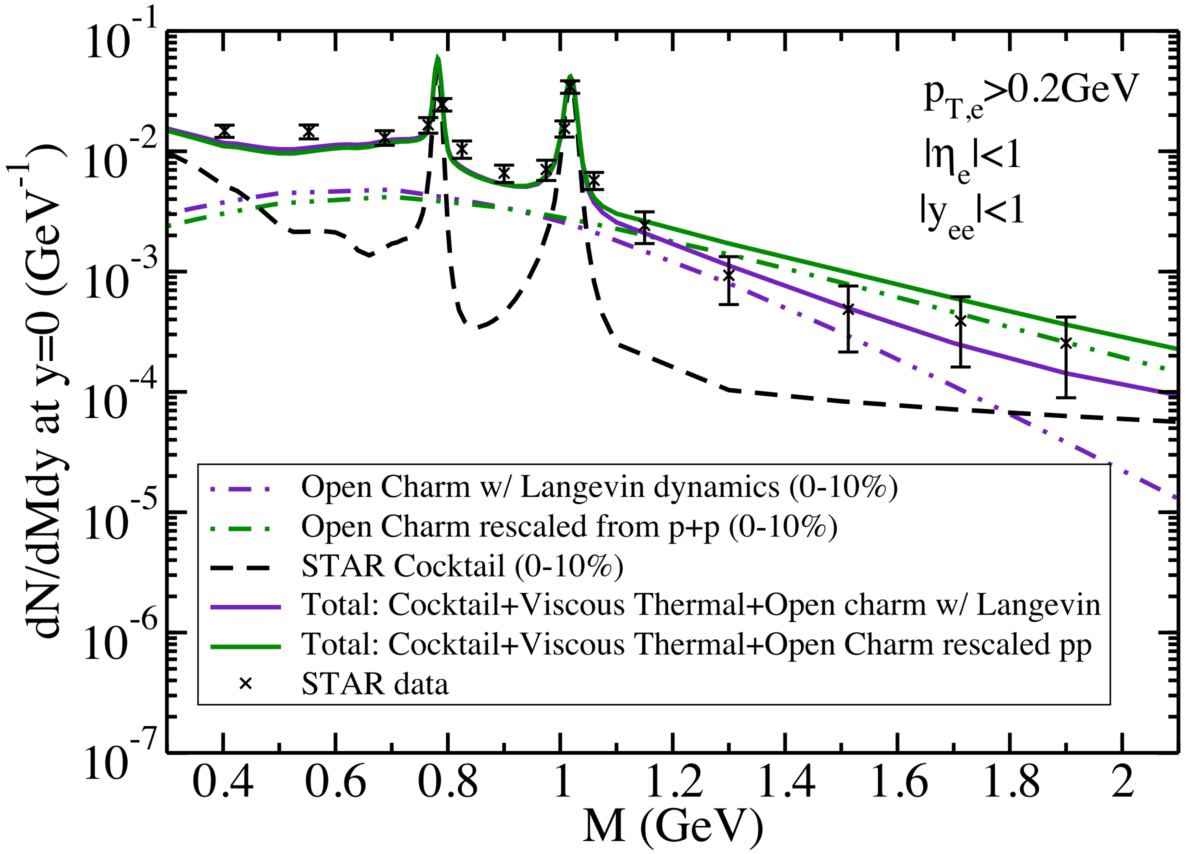
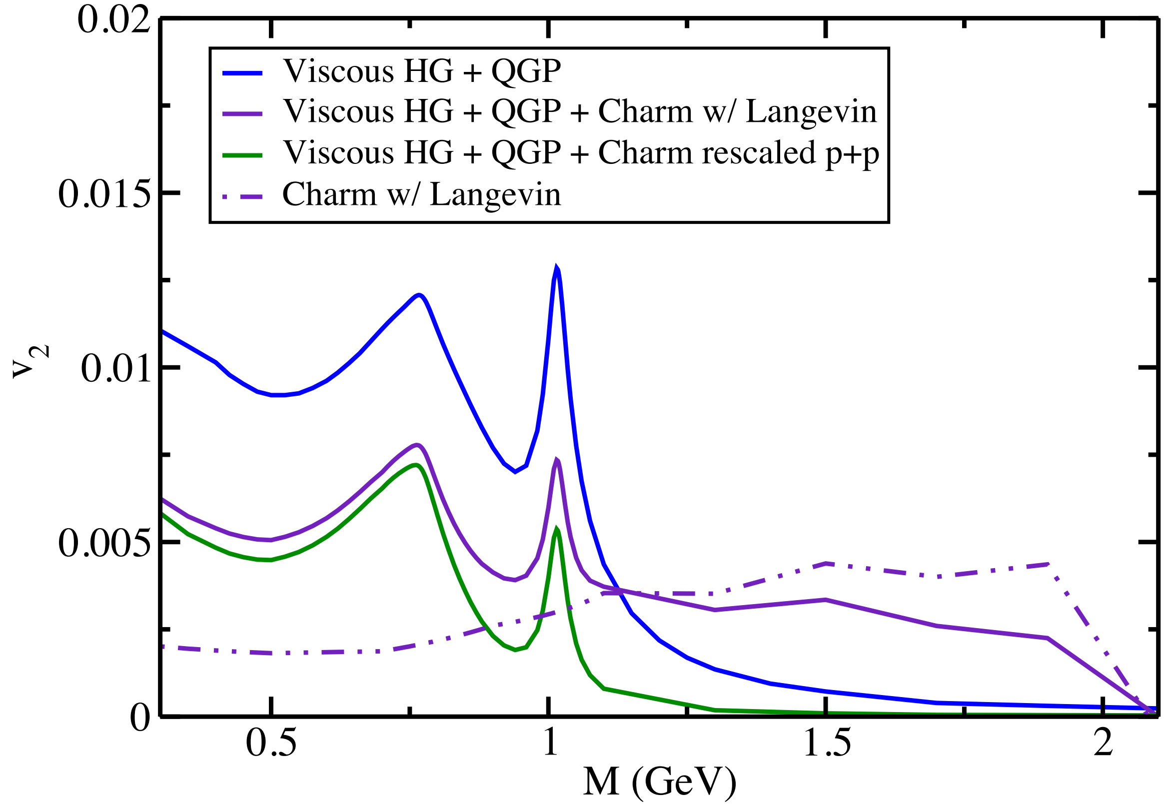
G. Vujanovic et al.,arXiv:1312.0676

Dileptons, some recent results







- High mass region and v_2, sensitive to heavy quark energy loss in the plasma

CONCLUSIONS

- The status of EM rates and their integration in dynamical models is still in flux
- Photon v_2 is sensitive to the EOS, and to various hydro parameters such as viscosity, and initial conditions (time and FICs). Current v_2 data: new physics? Measuring photon v_3, v_n at RHIC and LHC will help complete this picture
- FICs and viscosity(ies) make a difference in photon (and dilepton) characterization of the HICs: one must be consistent with hadronic data
- Jet-plasma photons need to be included: MARTINI
- Known unknowns: pre-equilibrium radiation
- **More work to be done**

Thanks to

- G. Denicol
- U. Heinz (OSU)
- S. Jeon
- I. Kozlov
- M. Luzum
- J.-F. Paquet
- R. Rapp (Texas A&M)
- B. Schenke
- C. Shen (OSU)
- H. van Hees (Frankfurt)
- G. Vujanovic