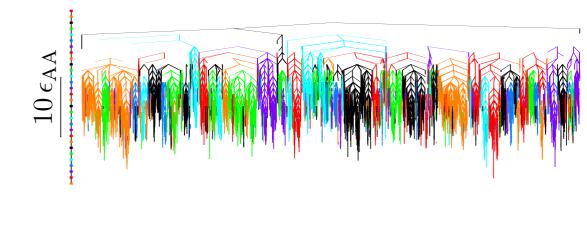
Glassy Dynamics in the Potential Energy Landscape

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Overview

Introduction

Strong and Fragile Glasses Potential Energy Landscape Visualising the Potential Energy Landscape

Glassy Dynamics

Coarse-graining the Landscape - Metabasins Cage-breaking Reversed and Productive Cagebreaks Calculating Diffusion Constants

Cage-break Metabasins

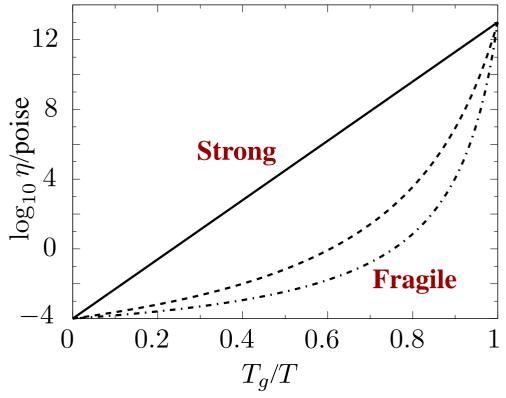
Random Walk Metabasins vs. Cagebreaks

'Super-Arrhenius' behaviour

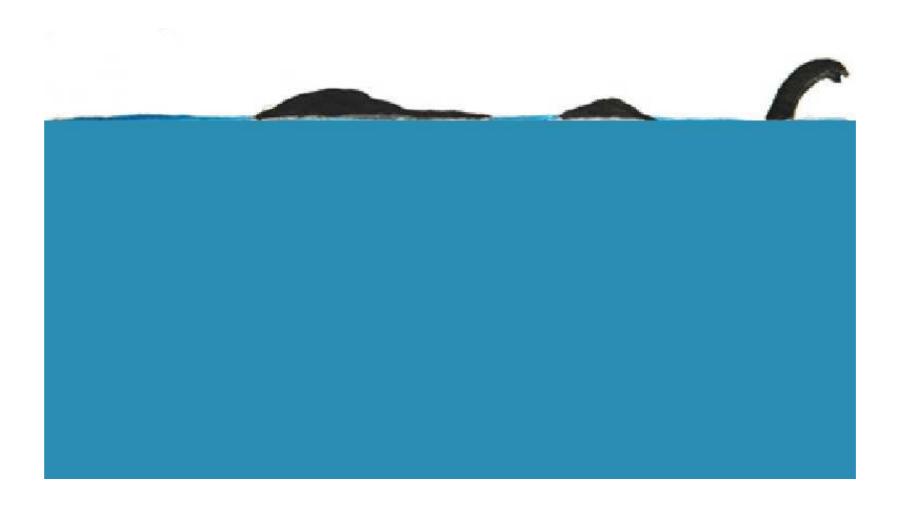
For some supercooled liquids, the temperature dependence of relaxation times or transport properties such as the diffusion constant, D, is stronger than predicted by the Arrhenius law.

	Arrhenius	Super-Arrhenius
Temperature dependence	Arrhenius Law	VTF equation
	$\eta = \eta_0 \exp[A/T]$	$\eta = \eta_0 \exp[A/(T-T_0)]$
Angell's classification	Strong	Fragile

Strong and Fragile Glasses



	Arrhenius	Super-Arrhenius
Temperature dependence	Arrhenius Law	VTF equation
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Angell's classification	Strong	Fragile



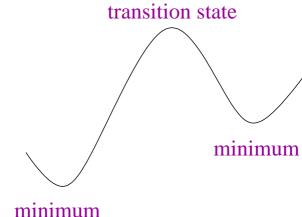


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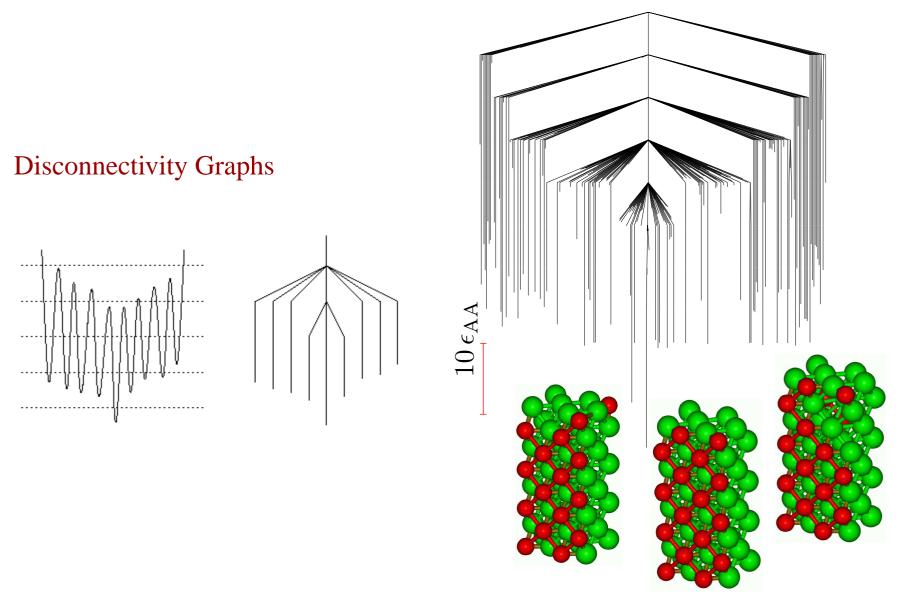
Potential Energy Landscape (PEL): the potential energy as a function of all the relevant particle coordinates.

- Any structure can be minimised to find its inherent structure, a minimum on the PEL.
- Discretisation and simplification of configuration space.



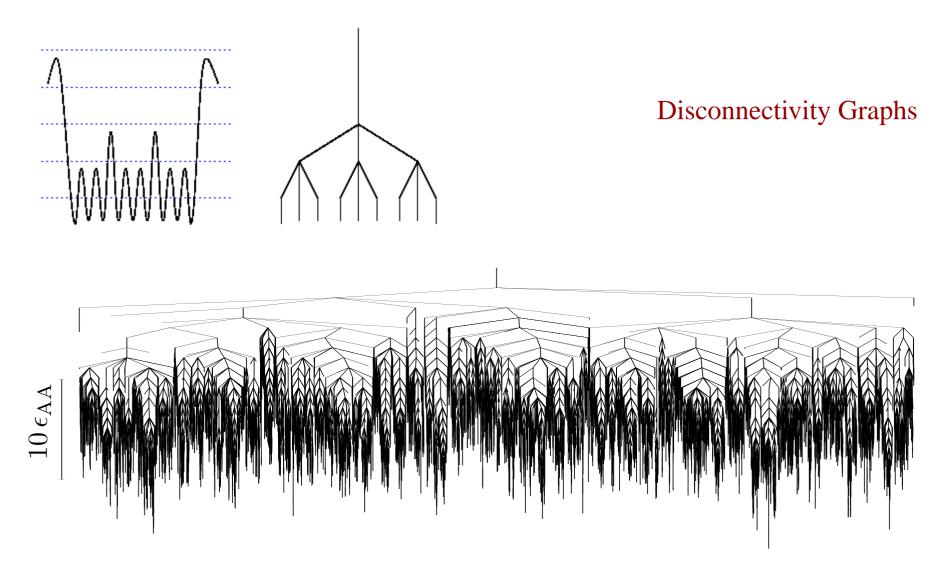
Dynamics requires information about transition states, the highest point on the lowest-energy pathway between two minima.

Visualising the Landscape - Crystal Landscapes



Calvo, Bogdan, de Souza and Wales, JCP 127, 044508 (2007)

Visualising the Landscape - Glassy Landscapes



de Souza and Wales, JCP 129, 164507 (2008)

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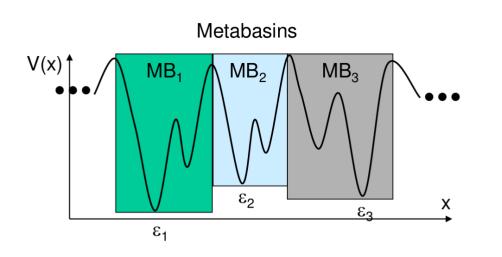
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Coarse-graining the landscape



- Transitions between metabasins follow a random walk
- Metabasins are well-characterised by an energy and waiting time
- Diffusion constants can be calculated

Doliwa and Heuer, PRE (2003)

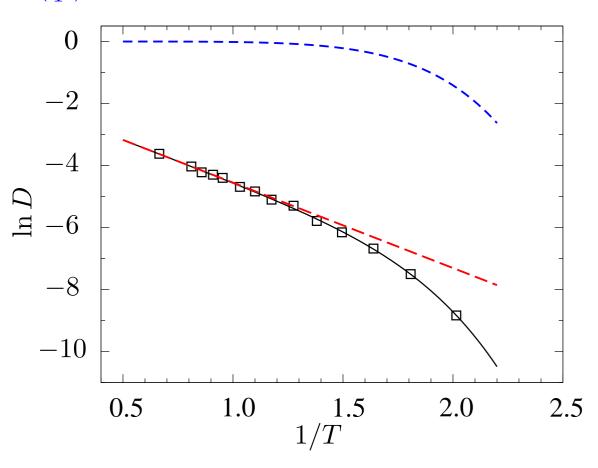
Problems with this approach:

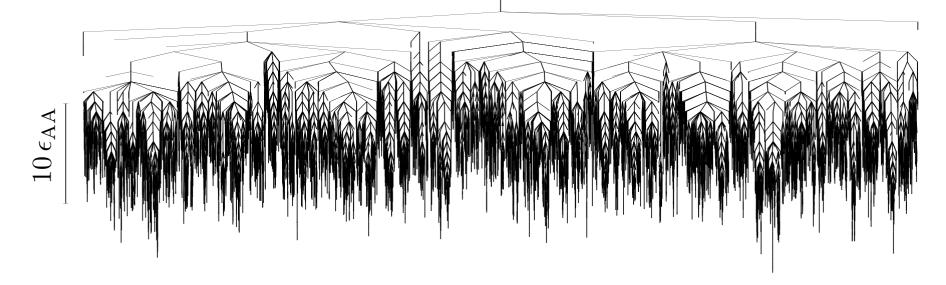
- How but not Why.
- No information about microscopic mechanisms, within metabasins or for transitions between metabasins.
- Identify minima by total system energy, the method cannot be scaled for larger system sizes, restricted to around 65 atoms.

Fitting to Super-Arrhenius Behaviour

- $\ln D_{\rm erg}(T) = -\left(\frac{m}{T}\right)^n \frac{c}{T} + \ln D_0$
- Arrhenius component: $-\frac{c}{T} + \ln D_0$
- Correction: $-\left(\frac{m}{T}\right)^n$

de Souza and Wales PRB 74, 134202 (2006) PRL 96, 057802 (2006)





Negative correlation in Minima-to-Minima Transitions

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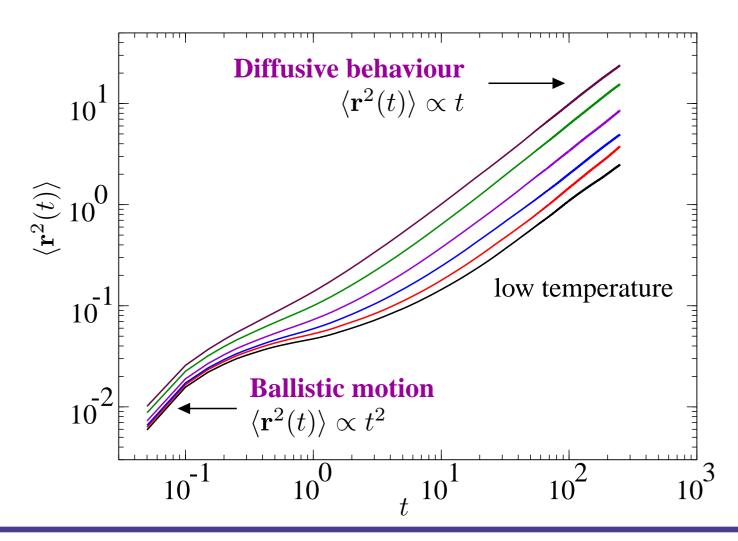
Negatively correlated Diffusive Processes

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Random Walk between Metabasins

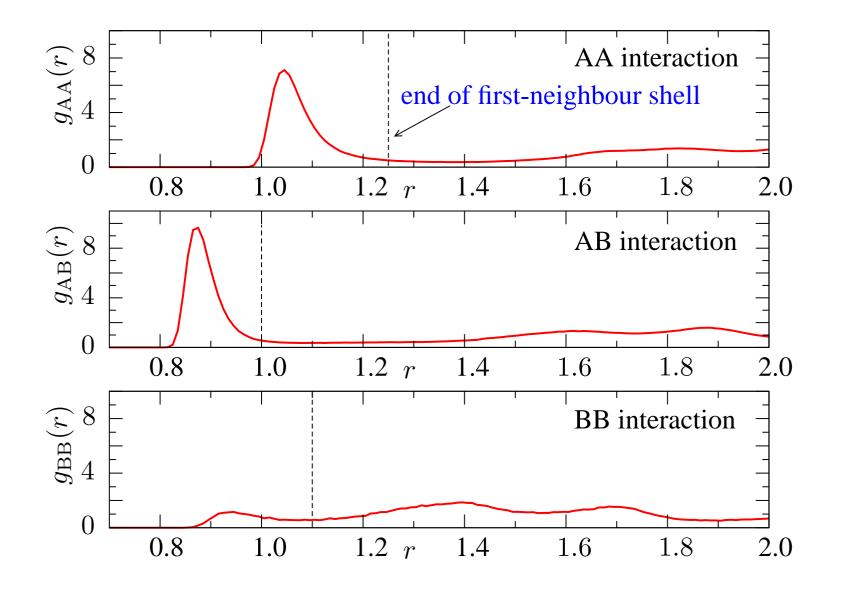
Mean square displacement \rightarrow Diffusion

Einstein relation: $D = \lim_{t \to \infty} \frac{1}{6t} \langle \Delta \mathbf{r}^2(t) \rangle$

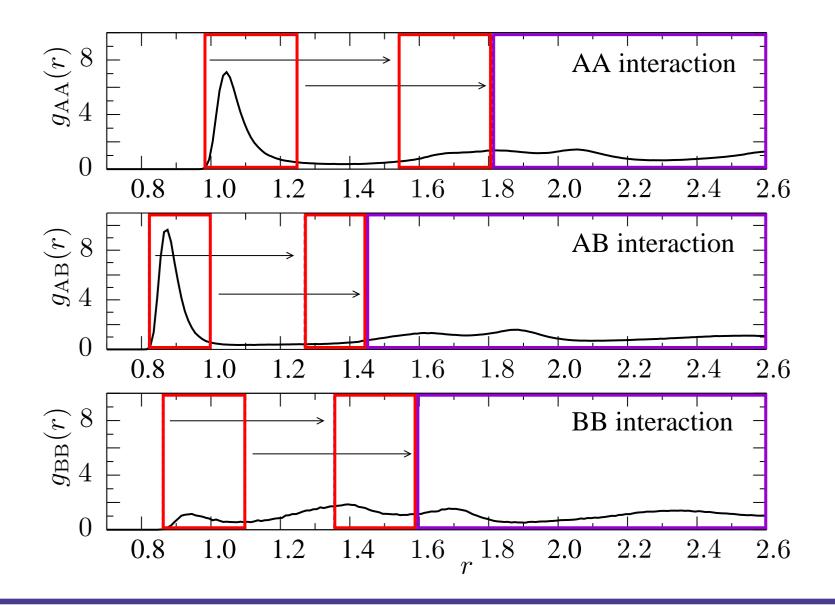


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Nearest Neighbours



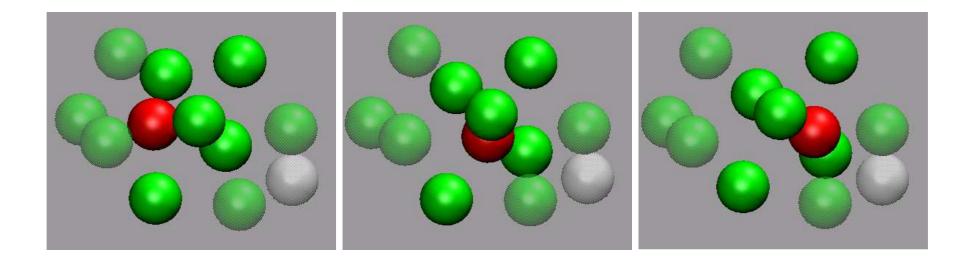
Glassy Dynamics in the Potential Energy Landscape - p. 13



Glassy Dynamics in the Potential Energy Landscape - p. 13,

Cage-Breaking Criteria

- Nearest neighbours are within a distance of 1.25 for an AA interaction.
- For the loss of a neighbour, relative distance changes by more than 0.561, which corresponds to half the equilibrium pair separation.
- A cage-break occurs with the loss/gain of at least two neighbours.



● Sequence of minimum – transition state – minimum for a cagebreak.

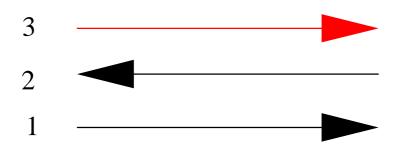
de Souza and Wales, JCP 129, 164507 (2008)

Reversed Cage-Breaks

- Identified when the net displacement squared is less than 10^{-5} .
- Chains of repeatedly reversed cage-breaks are found.
- Determine cage-breaks which are Productive towards long-term diffusion:

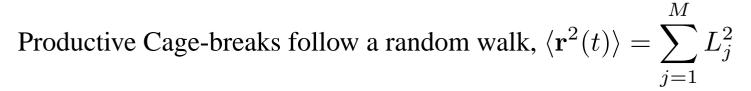
The cage-break is not followed by the reverse event. The cage-break is not part of a reversal chain OR

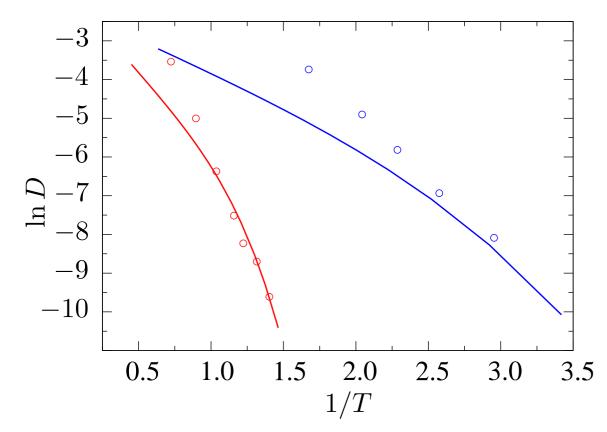
ends a chain with an even number of reversals.



3 cage-breaks2 reversalsLast cage-break is Productive

Diffusion from Productive Cage-Breaks





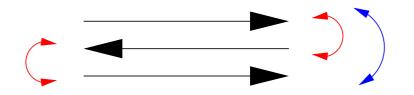
60-atom binary Lennard-Jones at number densities of 1.3 and 1.1Landscape-influenced regime (1/T):0.78 and 1.78Landscape-dominanced regime (1/T):1.56 and 3.56

Accounting for correlation

The following simplifications are suggested by our studies of diffusion using Molecular Dynamics trajectories:

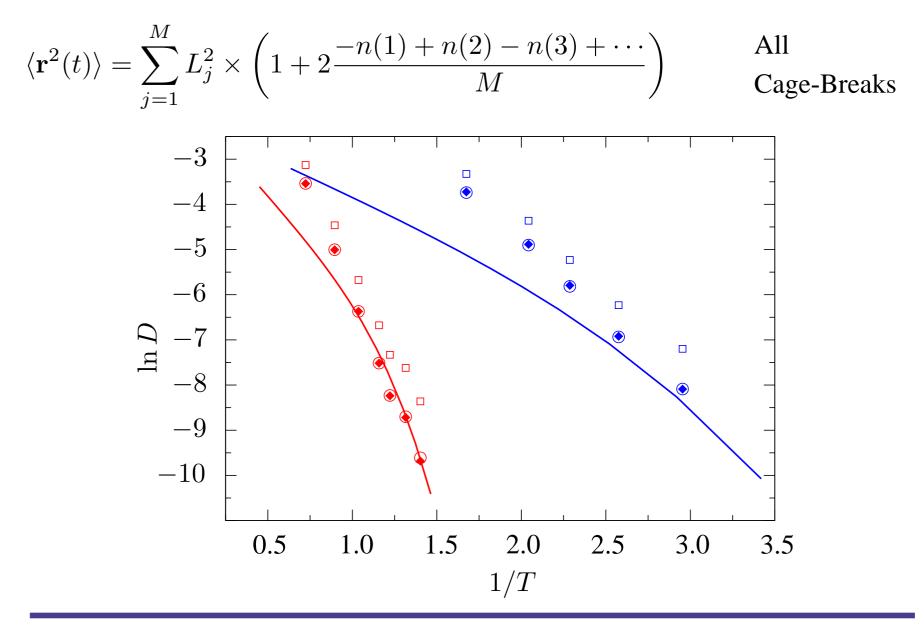
- The displacements of cage-breaks are similar and can be represented by a constant, L.
- Correlation arises from direct return events.
- We can account for correlation effects using a count of reversal chains of length z, n(z).

$$\langle \mathbf{r}^{2}(t) \rangle = ML^{2} \left(1 + 2 \frac{-n(1) + n(2) - n(3) + \cdots}{M} \right)$$



Reversal chain, z=2. Two reversal chains, z=1. n(1) = 2 and n(2) = 1

Diffusion from All Cage-Breaks



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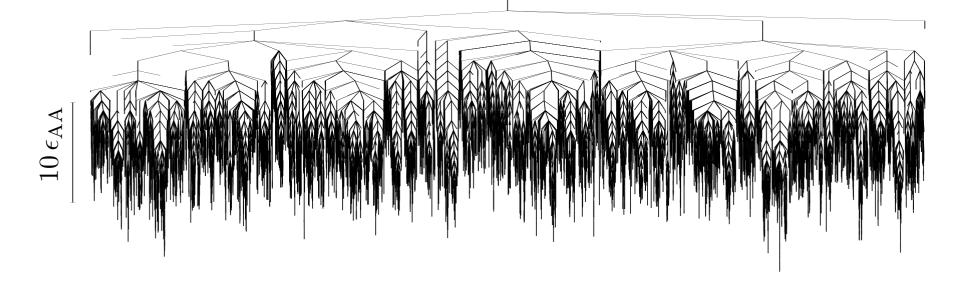
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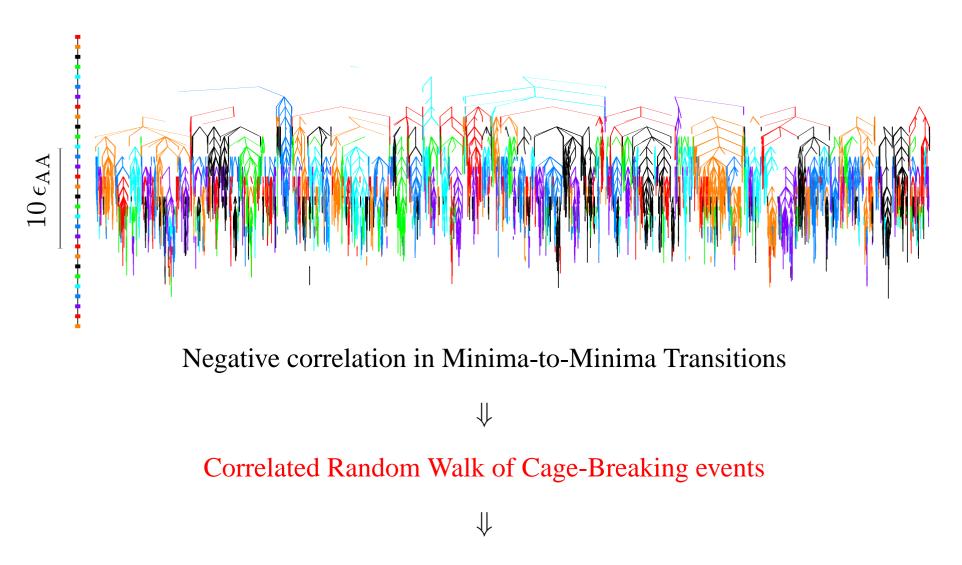
Negative correlation in Minima-to-Minima Transitions

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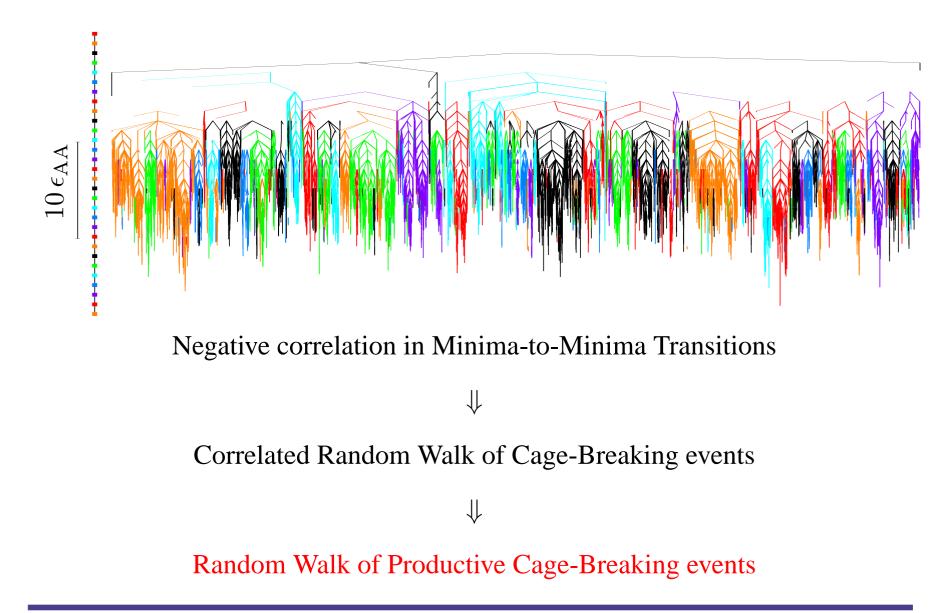
Negatively correlated Diffusive Processes

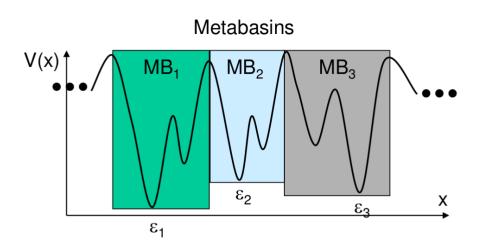
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Random Walk between Metabasins



Random Walk between Metabasins





Metabasins vs. Cagebreaks

- Transitions between metabasins follow a random walk
- Metabasins are well-characterised by an energy and waiting time
- Diffusion constants can be calculated

de Souza, Rehwald and Heuer, in preparation (2013)

Advantages of this method:

- How and Why.
- Information about microscopic mechanisms, within metabasins and for transitions between metabasins.
- Method can be scaled for larger system sizes.

Conclusions

- The Potential Energy Landscape for glass-forming systems is extremely complex.
- The landscape can be coarse-grained into metabasins
- Important transitions such as cagebreaks can be identified
- We have reconciled the two approaches, providing a microscopic description for metabasins within the PEL in the form of productive cagebreaks.
- Microscopic mechanisms <-> Macroscopic properties

