

Radiative Transfer Calculation for the Thermal Radiation from GRB Jets

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Outline

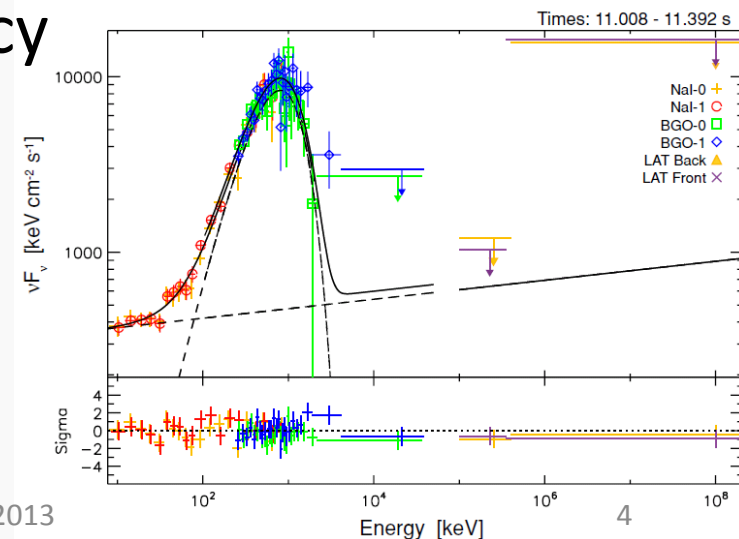
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Introduction

Models for the prompt emission

- Internal shock model
 - A standard scenario for a long time.
 - Low Radiative efficiency, line of death problem
- Photospheric (thermal emission) model
 - Thermal emission from relativistic jets
 - (possibly) high radiative efficiency
 - Some GRBs exhibit blackbody like feature (e.g., GRB090902B).

(Ryde et al 2010)

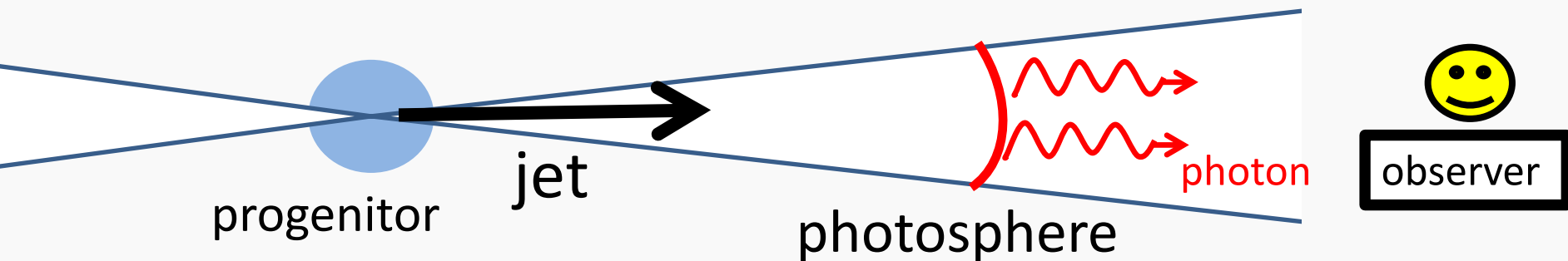


Thermal emission from GRB jets

- Thermal emission from GRB jets have been investigated by performing hydrodynamical simulations.

(Lazzati+2009, Mizuta+11, Nagakura+11)

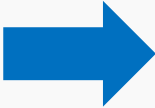
- They calculated the light curves and spectra by superposing blackbody radiation **emitted from the photosphere** with $\tau = 1$.



Thermal emission from GRB jets

However

- The observed photons should be generated in the inner layer with $\tau \gg 1$. (e.g., Beloborodov 13)
- Radiation intensity can be anisotropic even in the comoving frame at $\tau \sim 1$. (Beloborodov 11, Aksenov+ 13)

 In order to treat the thermal radiation from GRB jets properly, both **the radiative transfer in the jet and complex structure of the jet and cocoon** should be taken into account.

We calculate the radiative transfer in the jet and cocoon.

Method

Hydrodynamical simulation

✓ 2D relativistic hydrodynamics (Tominaga 2009)

✓ Setup

– Progenitor: $15M_{\text{sun}}$ WR star ($R_{\text{prog}} \sim 2.3 \times 10^{10} \text{cm}$)

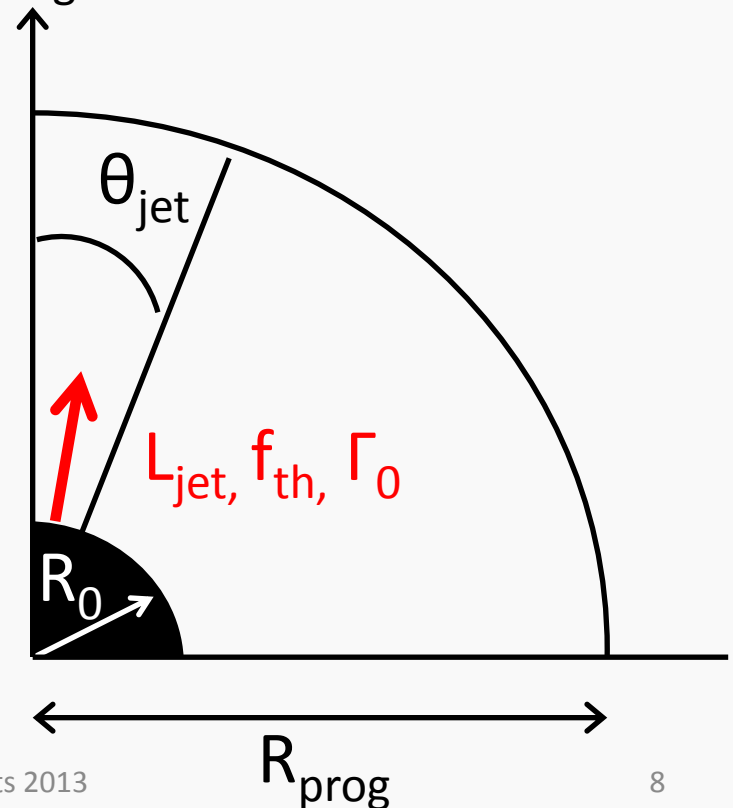
– $\Gamma_0 = 5$

– $\Theta_{\text{jet}} = 10^\circ$

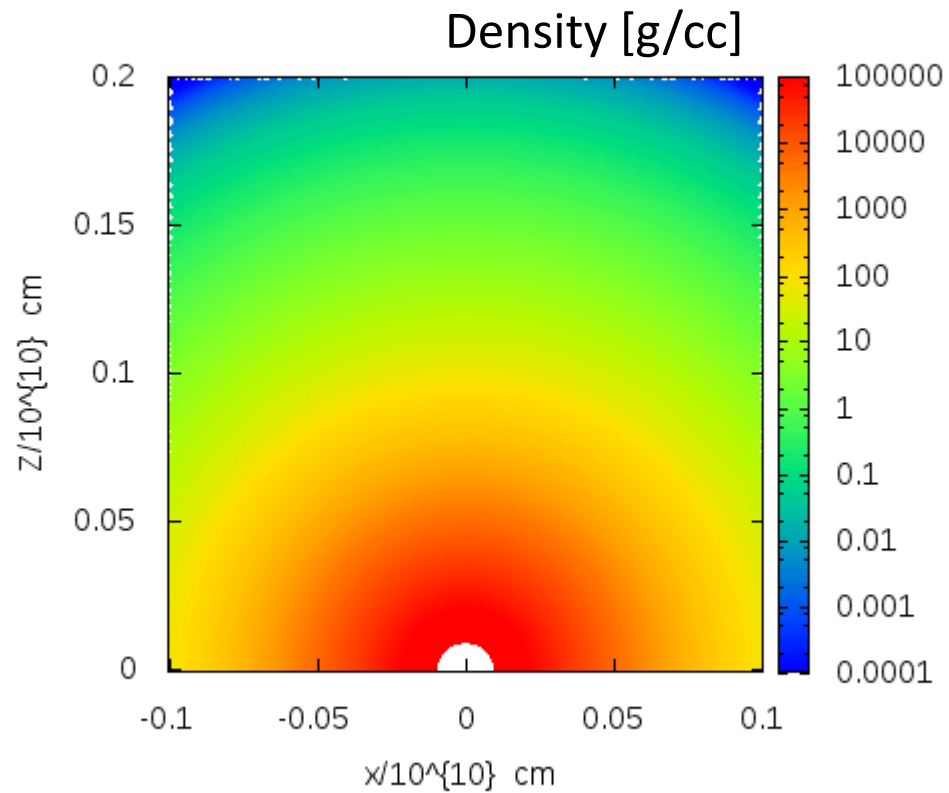
– $L_{\text{jet}} = 5.3 \times 10^{50} \text{erg s}^{-1}$

– $f_{\text{th}} = 0.9925$ ($e_{\text{int}}/\rho c^2 = 80$)

– $(\log r, \theta) = (600, 150)$ grids
from $R_0 = 10^9 \text{cm}$

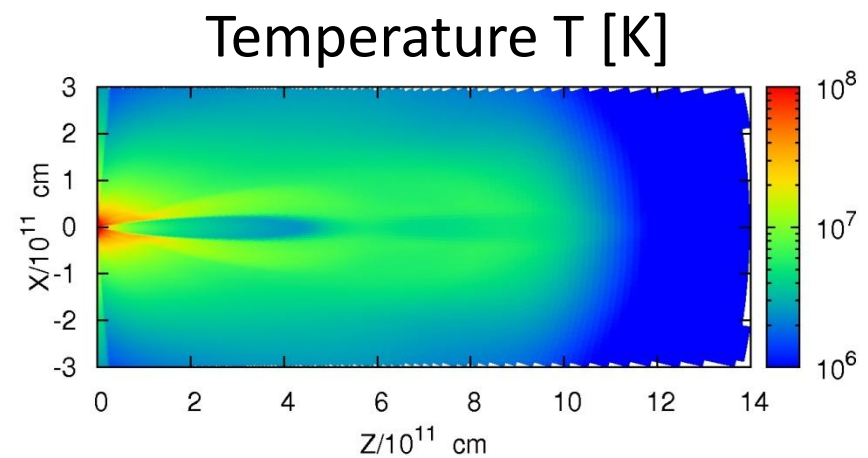
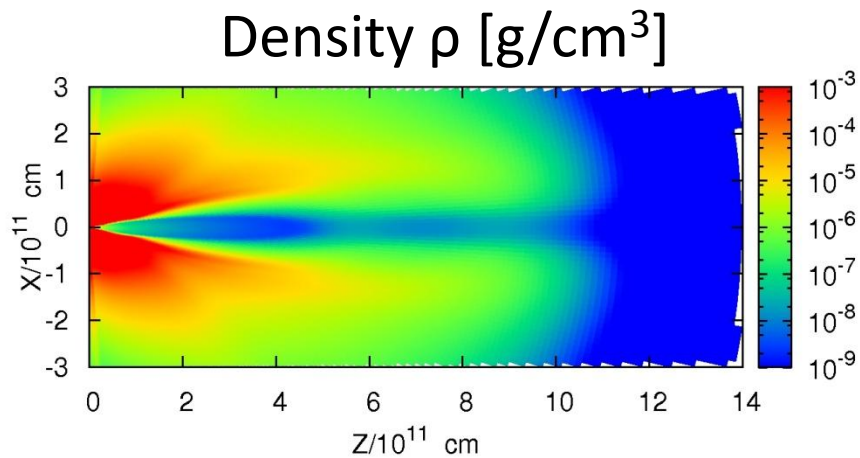


Hydrodynamical simulation



Snapshot at 40s

- We use a snapshot at 40s for the structures of the jet and cocoon.



The site of photon production

- The effective optical depth τ_*

For the static medium (Rybicki & Lightman 79)

$$\tau_*^{\text{NR}} \sim \sqrt{\tau_a(\tau_a + \tau_s)}$$

For the relativistic medium

$$\tau_*^{\text{R}} = \left\{ \frac{\Gamma^2}{3}(\beta^2 + 3) + (\Gamma\beta)^2 \frac{\tau_s}{\tau_a} \right\}^{-1/2} \frac{\sqrt{\tau_a(\tau_a + \tau_s)}}{\Gamma(1 - \beta \cos \theta_v)}$$

$$\tau_a = \Gamma(1 - \beta \cos \theta_v) \alpha' L, \quad \tau_s = \Gamma(1 - \beta \cos \theta_v) \sigma' L$$

In the non-relativistic limit, $\tau_*^{\text{R}} \rightarrow \tau_*^{\text{NR}}$

In the relativistic limit, $\tau_*^{\text{R}} \rightarrow 2 \tau_a$ for $\Theta=0$

Absorption processes

$$\alpha'(x) = \alpha'_{\text{ff}}(x) + \alpha'_{\text{DC}}(x)$$

- Free-free absorption ($e + p + \gamma \rightarrow e + p$)

$$\alpha'_{\text{ff}}(x) = \frac{\alpha_{\text{fin}} \lambda_c^3 \sigma_T}{\sqrt{6\pi}} \theta^{-1/2} Z^2 n_e n_i x^{-3} (1 - e^{-x/\theta}) g_{\text{ff}} \quad (\text{Rybicki\&Lightman 79})$$

- Double Compton absorption ($\gamma + \gamma + e \rightarrow \gamma + e$)

$$\alpha'_{\text{DC}}(x) = \frac{2\alpha_{\text{fin}} \lambda_c^3 \sigma_T}{\pi^2} \theta^2 n_e n_\gamma x^{-3} (e^{x/\theta} - 1) g_{\text{DC}} \quad (\text{Svensson 84})$$

- We assume $h\nu = kT$

$$x = \frac{h\nu}{m_e c^2}$$

$$\theta = \frac{k_B T}{m_e c^2}$$

Photosphere

- τ_* to a radius R_*

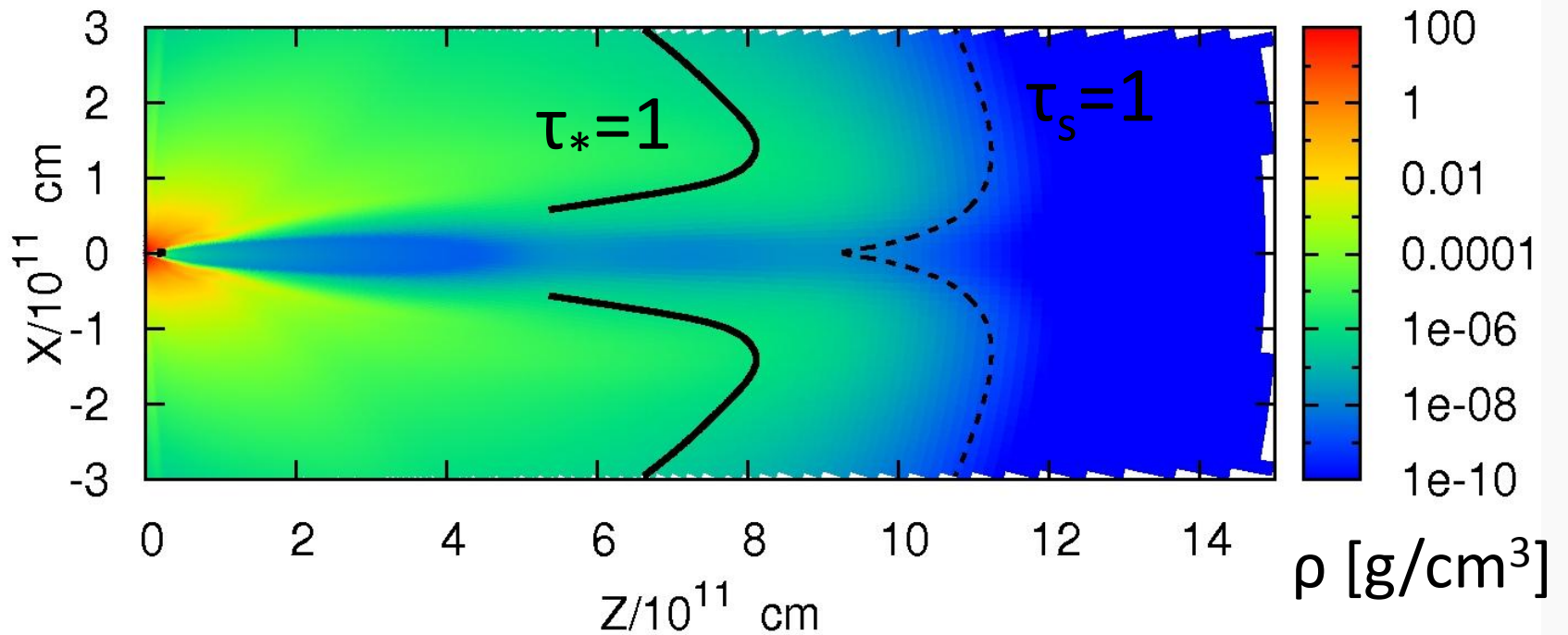
$$\tau_* = \int_{R_*}^{\infty} \left\{ \frac{\Gamma^2}{3} (\beta^2 + 3) + (\Gamma\beta)^2 \frac{\sigma'}{\alpha'} \right\}^{-1/2} \sqrt{\alpha'(\alpha' + \sigma')} dr$$

- $\sigma' = n_e \sigma_T$
- α' depends on n_e, n_γ, T .
- We assume

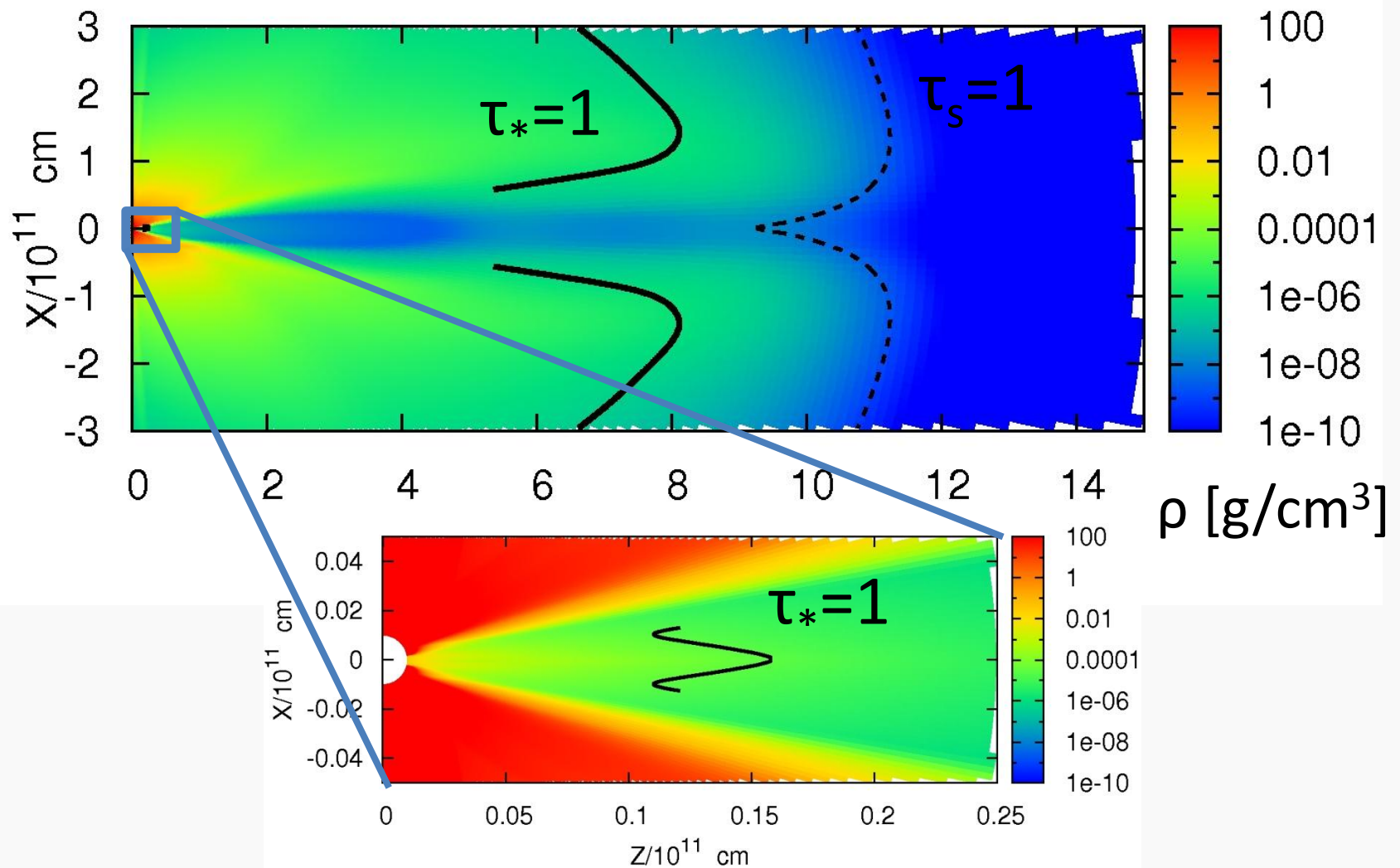
$$n_\gamma = n_{\gamma*} \left(\frac{R_*}{r} \right)^2 \quad \text{and} \quad n_{\gamma*} \equiv n_{\text{bb}}(R_*)$$

We find the R_* which satisfies $\tau_* = 1$

Photosphere



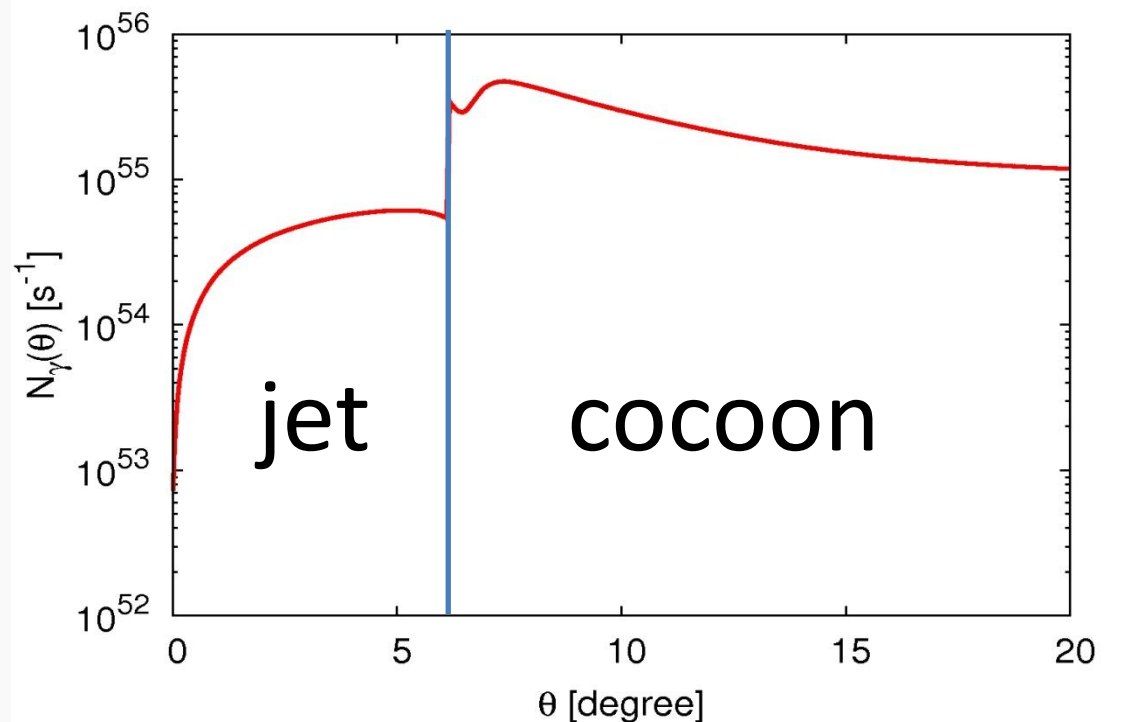
Photosphere



Photosphere

- The number of photons emitted at the photosphere:

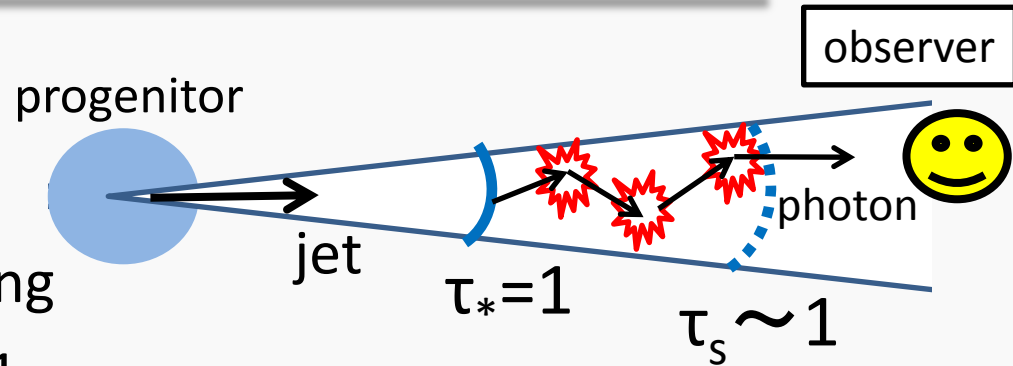
$$N_{\gamma}(\Theta) = 16\pi^2 \Gamma(3) \zeta(3) (kT_*/hc)^3 R_*^2 \sin\Theta_*$$



Radiative transfer

✓ Numerical code

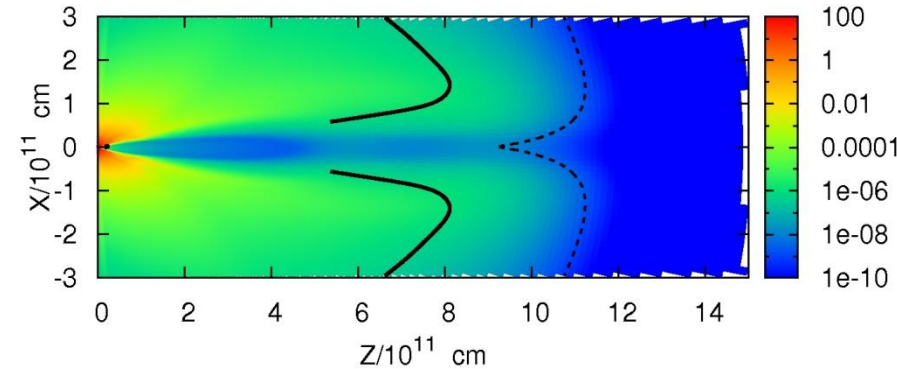
- Monte Carlo method
- Calculate Compton scattering
- Photons are injected at $\tau_* = 1$



✓ Photon injection

- Spatial distribution: $N_\nu(\Theta)$
- Planck distribution with local plasma temperatures
- Isotropic in the comoving frame

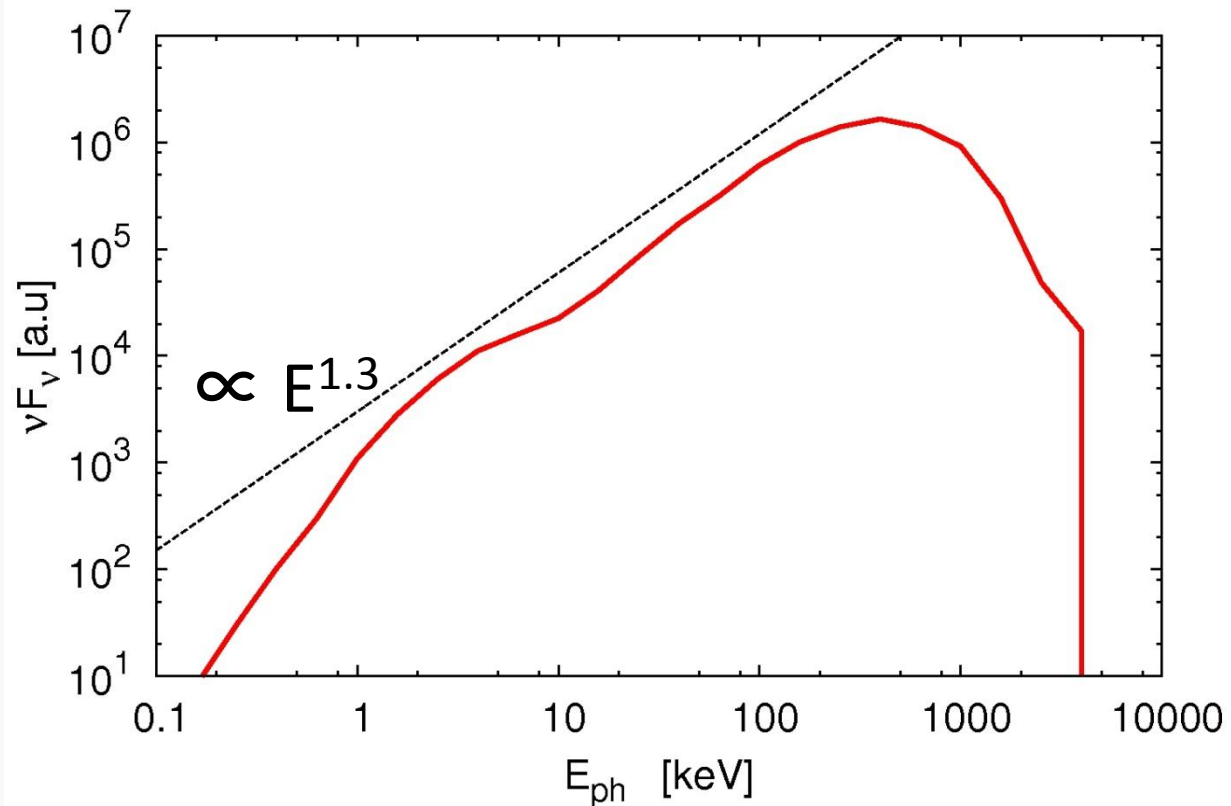
We use a snapshot at $t=40s$ for the jet and cocoon structure.



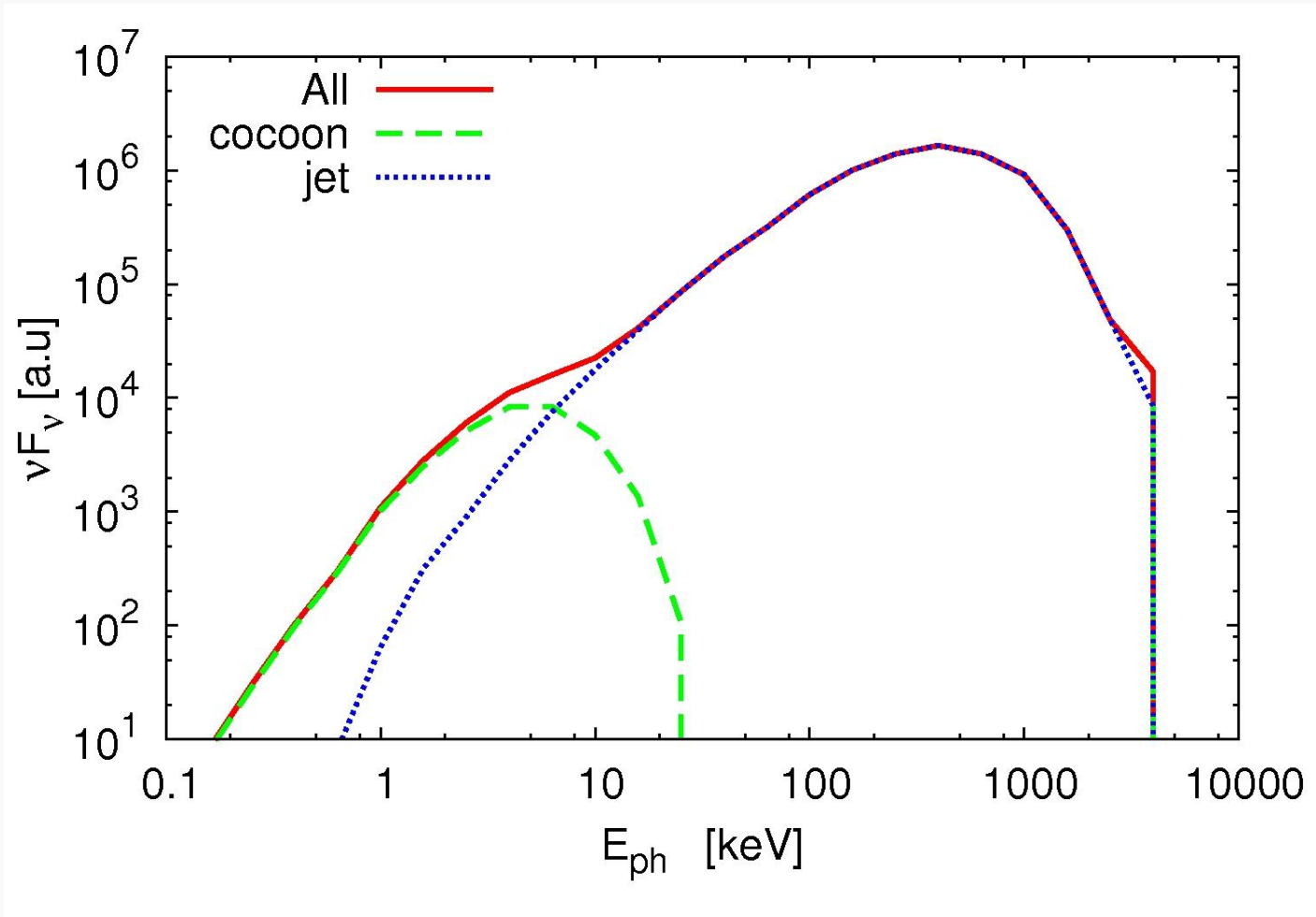
Results

Observed spectrum

- $E_{\text{peak}} \sim 450\text{keV}$
- A bump like feature at low energies
- At the low energy,
 $\nu F_{\nu} \propto E^{1.3}$
 $\rightarrow N_{\nu} \propto E^{-0.7}$
- No high energy PL

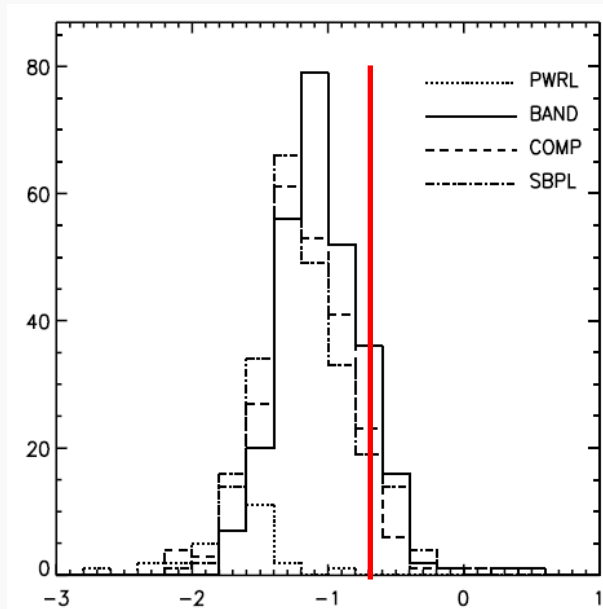


Origin of the bump?

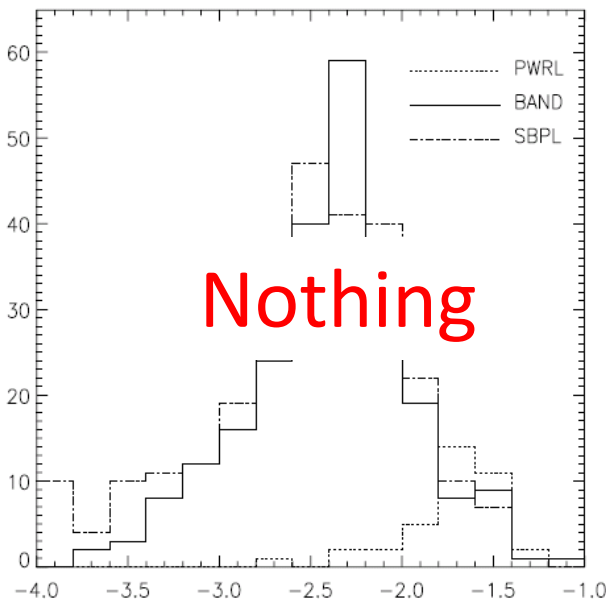


Comparison with the observations

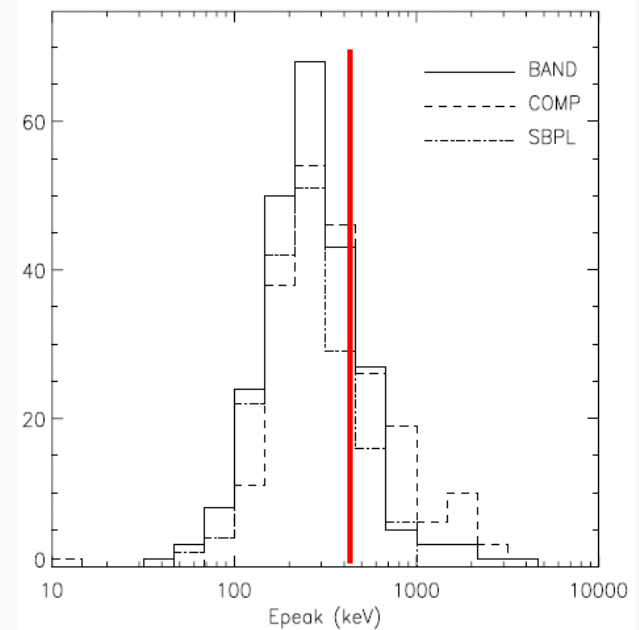
Kaneko et al 2006



low energy index (α)



high energy index (β)



peak energy (E_{peak})

Summary

Summary

- ✓ We calculate radiative transfer for the thermal radiation from GRB jet.
- ✓ Both the jet and cocoon components constitute the observed spectrum.
- ✓ The low energy index may be determined by the relative brightness of these two components.