

# Universal quantum dynamics near the two-dimensional Superfluid- to-Insulator transition

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# Collaborators



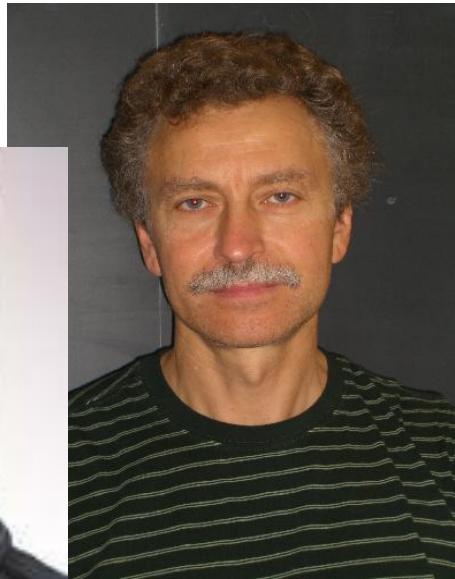
**Longxiang Liu**



**Youjin Deng**



**Lode Pollet**



**Nikolay  
Prokofiev**

# 3D XY universality

- Field theory: (2+1)-Dimensional *relativistic* O(2) model

$$S = \int_0^{1/T} dt \int d^2 r \left[ |\partial_t \Psi|^2 - c^2 |\nabla \Psi|^2 - V(\Psi) \right]$$

$$V(\Psi) = \lambda |\Psi|^2 + u(|\Psi|^2)^2$$

- Complex field  $\Psi$
- Quantum critical point : Superfluid-to-Insulator transition

# Mean field scenario and Higgs mode

$$V(\Psi) = (\lambda - \lambda_c) |\Psi|^2 + u(|\Psi|^2)^2$$

- Ordered phase (Superfluid phase)

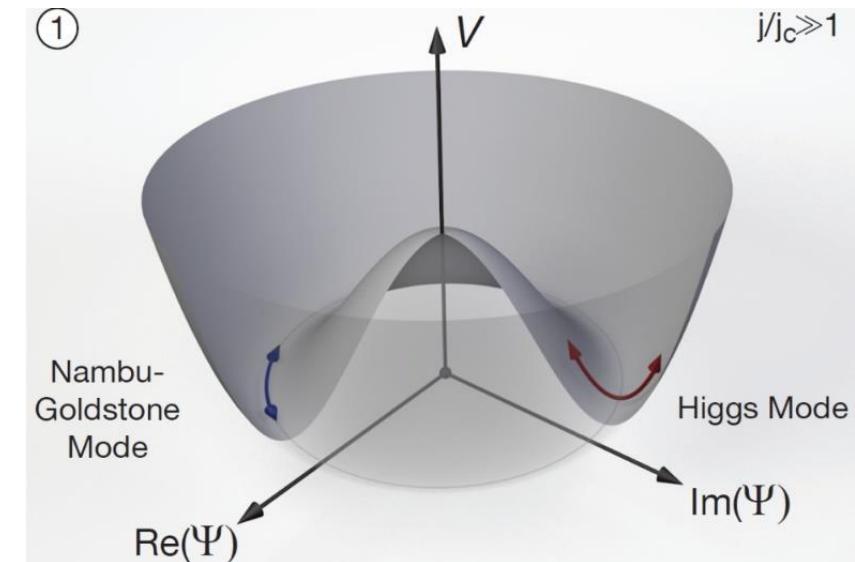
- Two decoupled excitation modes

Nambu-Goldstone mode

fluctuation of phase  
massless (gapless)

Massive Goldstone mode (Higgs mode)

fluctuation of amplitude  
massive (gapped)



# Why 2D?

D=3 Asymptotically exact mean-field → Higgs mode is well-defined.

P. Merchant, B. Normand, K. W. Kramer, M. Boehm, D. F. McMorrow, and Ch. Ruegg

Nature Physics 10, 373-379 (2014)

Ch. Ruegg, et.al, Phys. Rev. Lett. 100, 205701 (2008)

U. Bissbort, et.al, Phys. Rev. Lett. 106, 205303 (2011)

D=2 Destroyed by strong decay into two Goldstone modes.

Chubukov, Sachdev, Ye '93

Altman, Auerbach '02

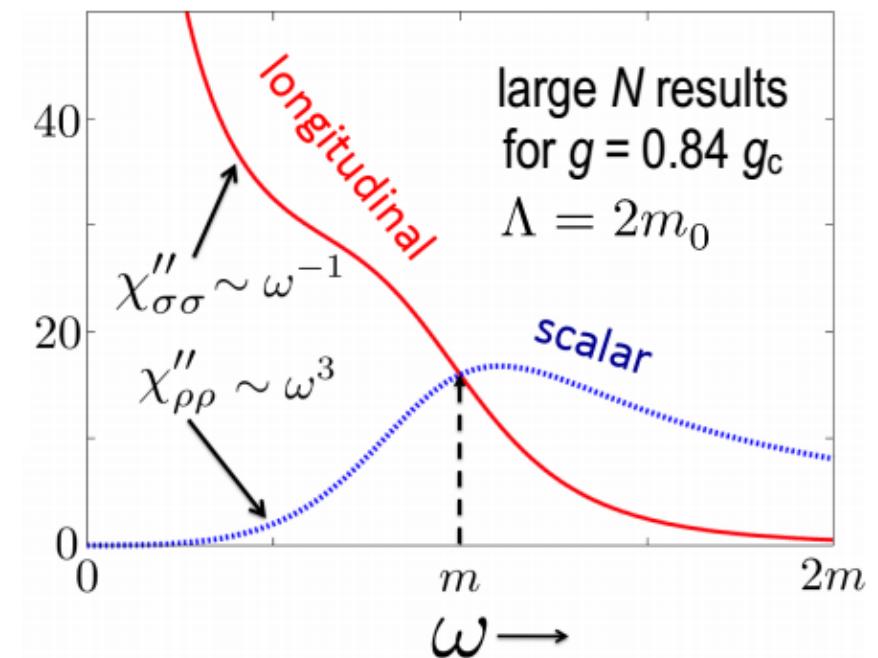
Zwerger '04

Just in case the theory is wrong, the type of the probe matters!

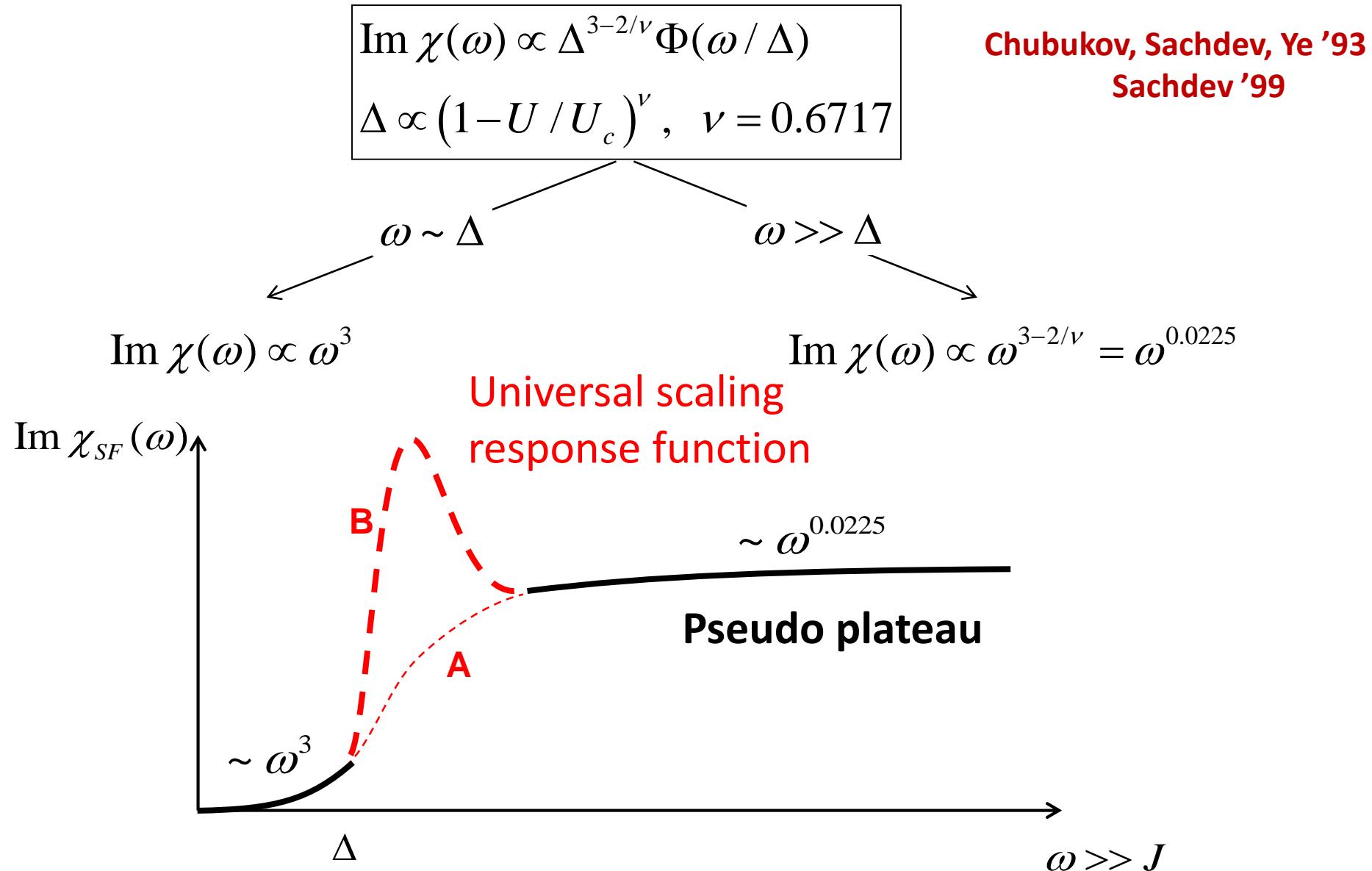
Longitudinal:  $S(\omega) = \langle \psi(0) \psi(t) \rangle_{\omega}$

Scalar:  $S(\omega) = \langle |\psi(0)|^2 |\psi(t)|^2 \rangle_{\omega}$   
is the best candidate

Chubukov, Sachdev, Ye '93  
Podolsky, Auerbach, Arovas '11

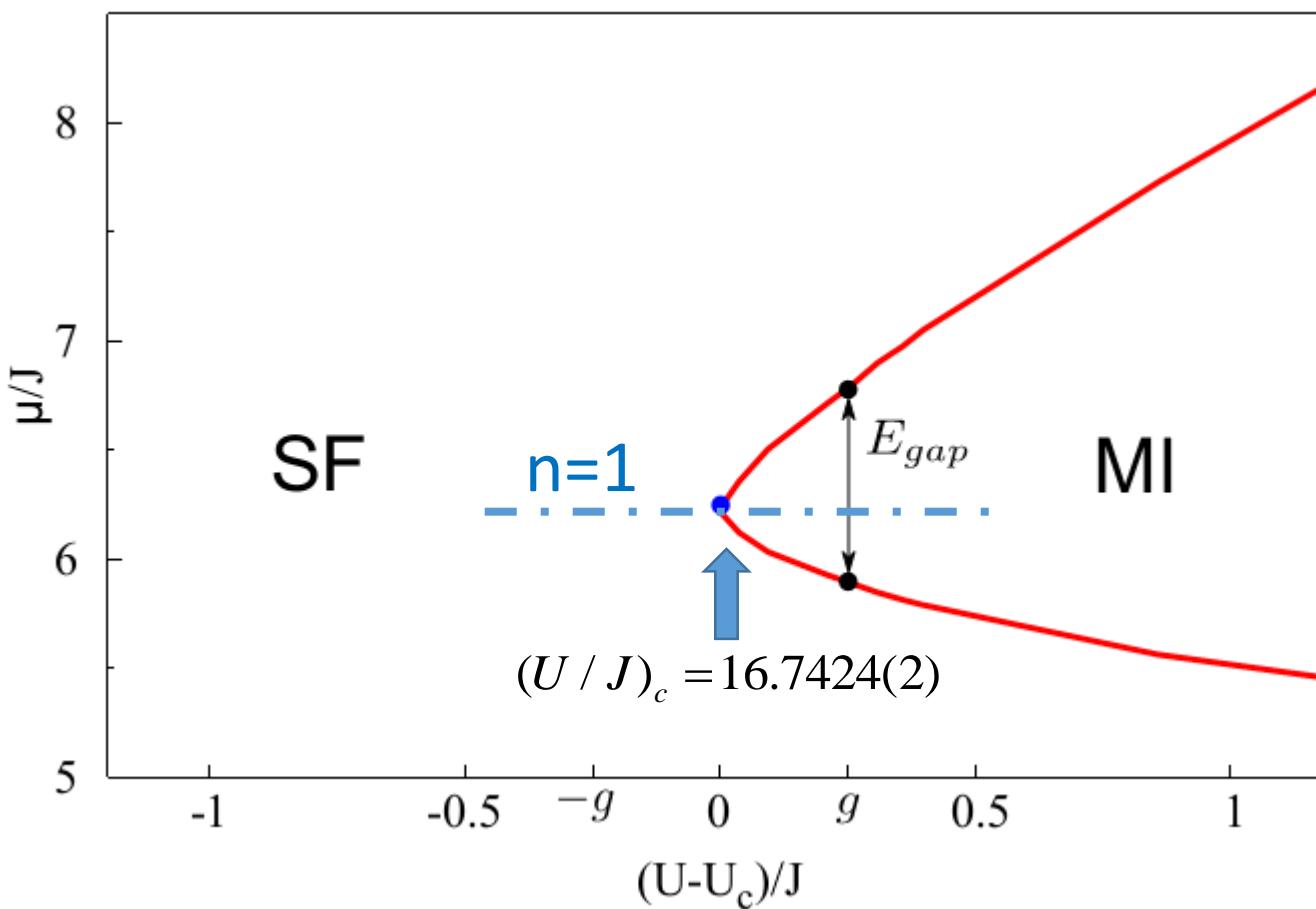


# What to expect for the scalar response function?



# Bose Hubbard model: a testbed

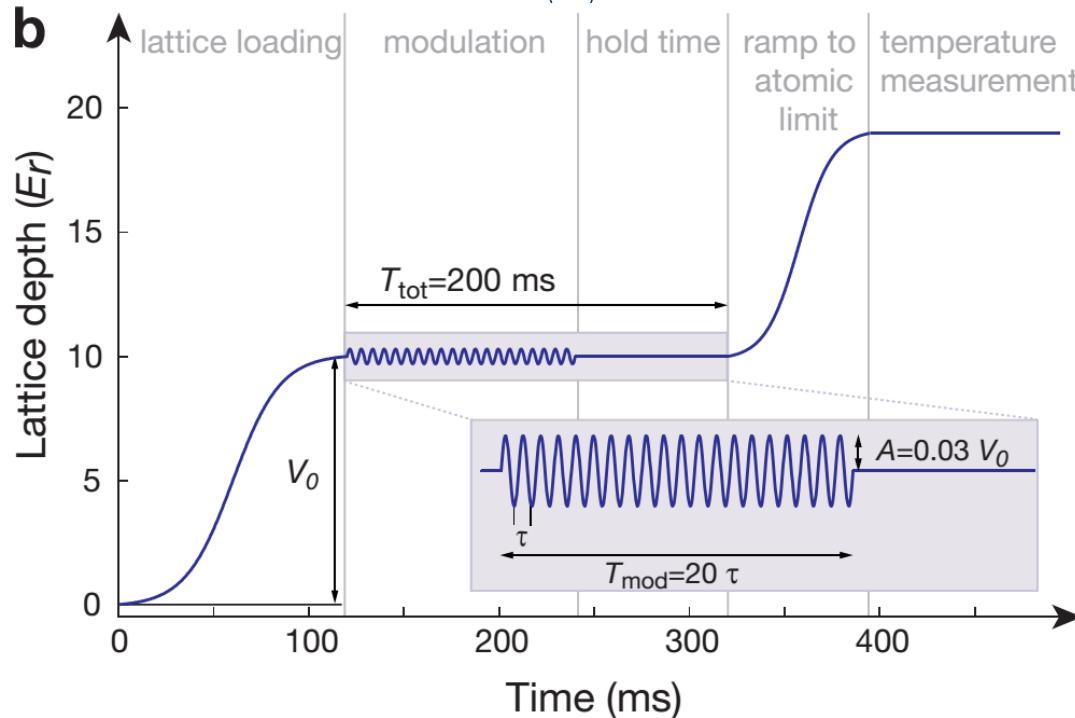
$$\hat{H}_{BH} = -J \sum_{\langle i,j \rangle} a_i^\dagger a_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) - (\mu - V_i) \sum_i \hat{n}_i$$



# Optical Lattice Emulator for BH model

$$\hat{H} = \hat{H}_{BH} + \frac{\delta J}{J} \hat{K} e^{i\omega t} , \quad \hat{K} = -J \sum_{\langle i,j \rangle} b_i^\dagger b_j$$

Linear response for small  $\delta J / J$



Endres, Fukuhara, Pekker, Cheneau, Schau  
Gross, Demler, Kuhr, Bloch, Nature '12

$$\chi(t) = i \left\langle \left[ \hat{K}(t), \hat{K}(0) \right] \right\rangle \Theta(t) \rightarrow \text{Energy dissipation rate} \propto \omega \operatorname{Im} \chi(\omega)$$

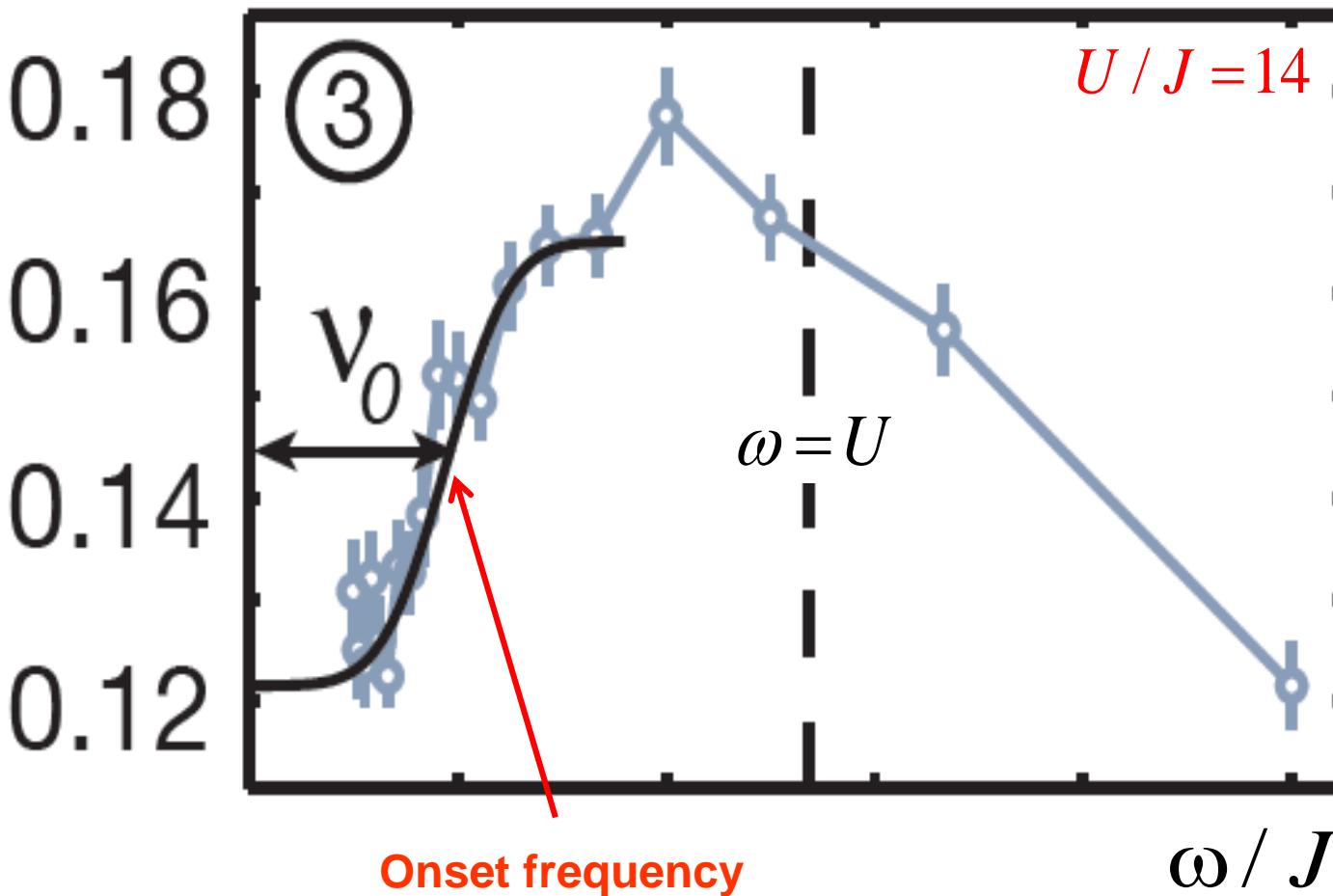
Total energy absorbed:  $\Delta E \propto \omega \operatorname{Im} \chi(\omega) (2\pi M / \omega) \propto \operatorname{Im} \chi(\omega)$

Temperature increase:  $\Delta T = \Delta E / C(T)$

# Recent experiment @ Munich: Onset of quantum critical response

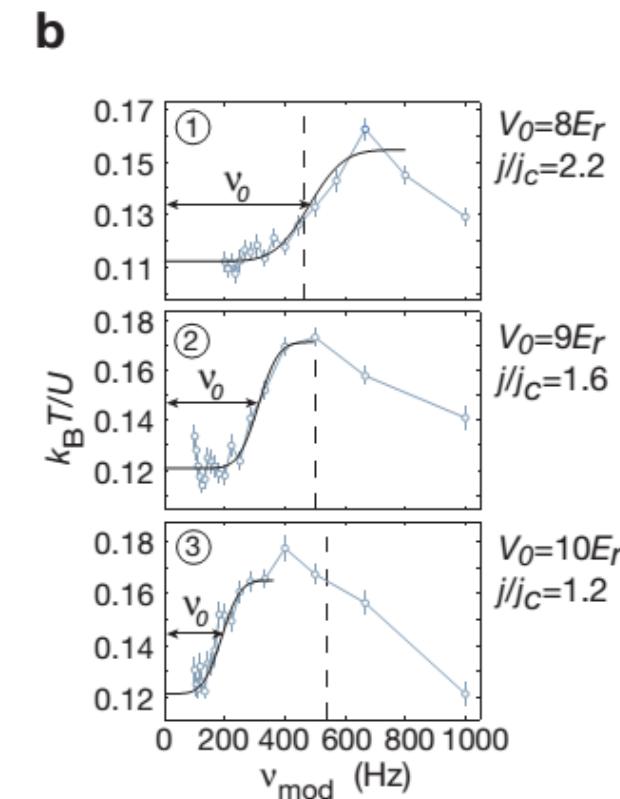
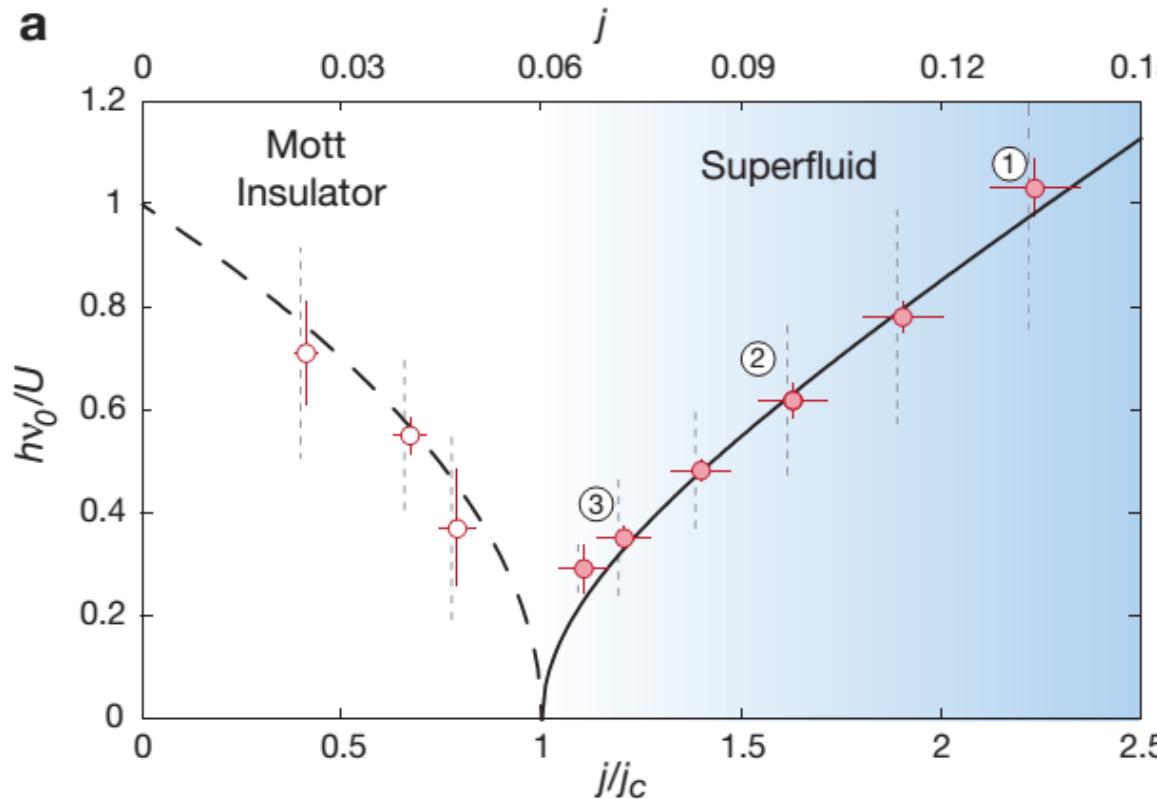
$$\Delta T / U \propto \text{Im } \chi(\omega)$$

Endres, Fukuhara, Pekker, Cheneau, Schau  
Gross, Demler, Kuhr, Bloch, Nature '12



# Nice scaling behavior, but where is Higgs resonance peak?

$$\nu_0 \propto \Delta \propto \left(1 - J / J_c\right)^\nu$$



Time for WORM !



# Monte Carlo probe: Kink-kink correlation

Kinetic energy :

$$\beta K_{MC} = - \sum_{kinks}$$

MC measurement:  
(Matsubara frequency)

$$K_{MC}(i\omega_n) = \int_0^b dt e^{i\omega_n t} K_{MC}(t) = - \sum_{k=kinks} e^{i\omega_n t_k}$$



kink-kink correlation function

$$C(i\omega_n) = \langle |K_{MC}(i\omega_n)|^2 \rangle + \langle K \rangle$$

Analytical continuation:  
(real frequency)



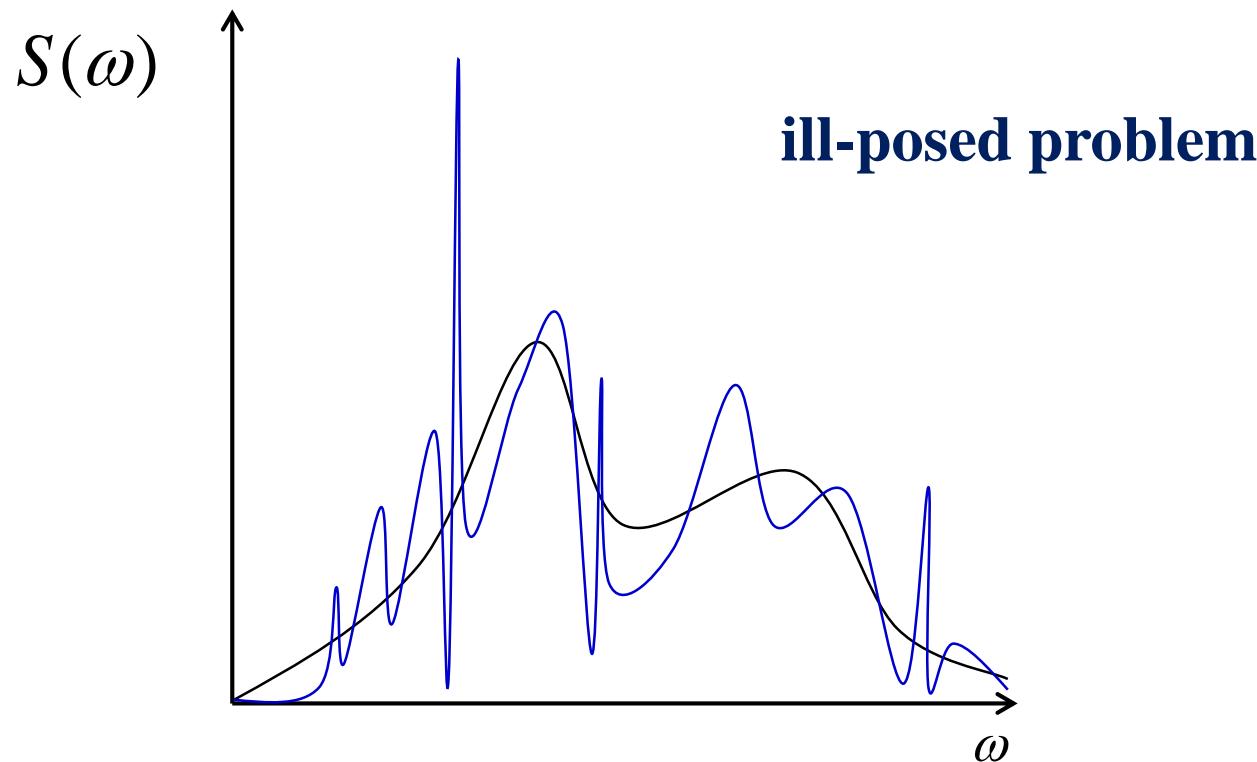
$$C(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \langle K(t) K(0) \rangle e^{-i\omega t} dt$$

But how to do the analytical continuation in  
numeric sense?

# Numerical analytical-continuation method

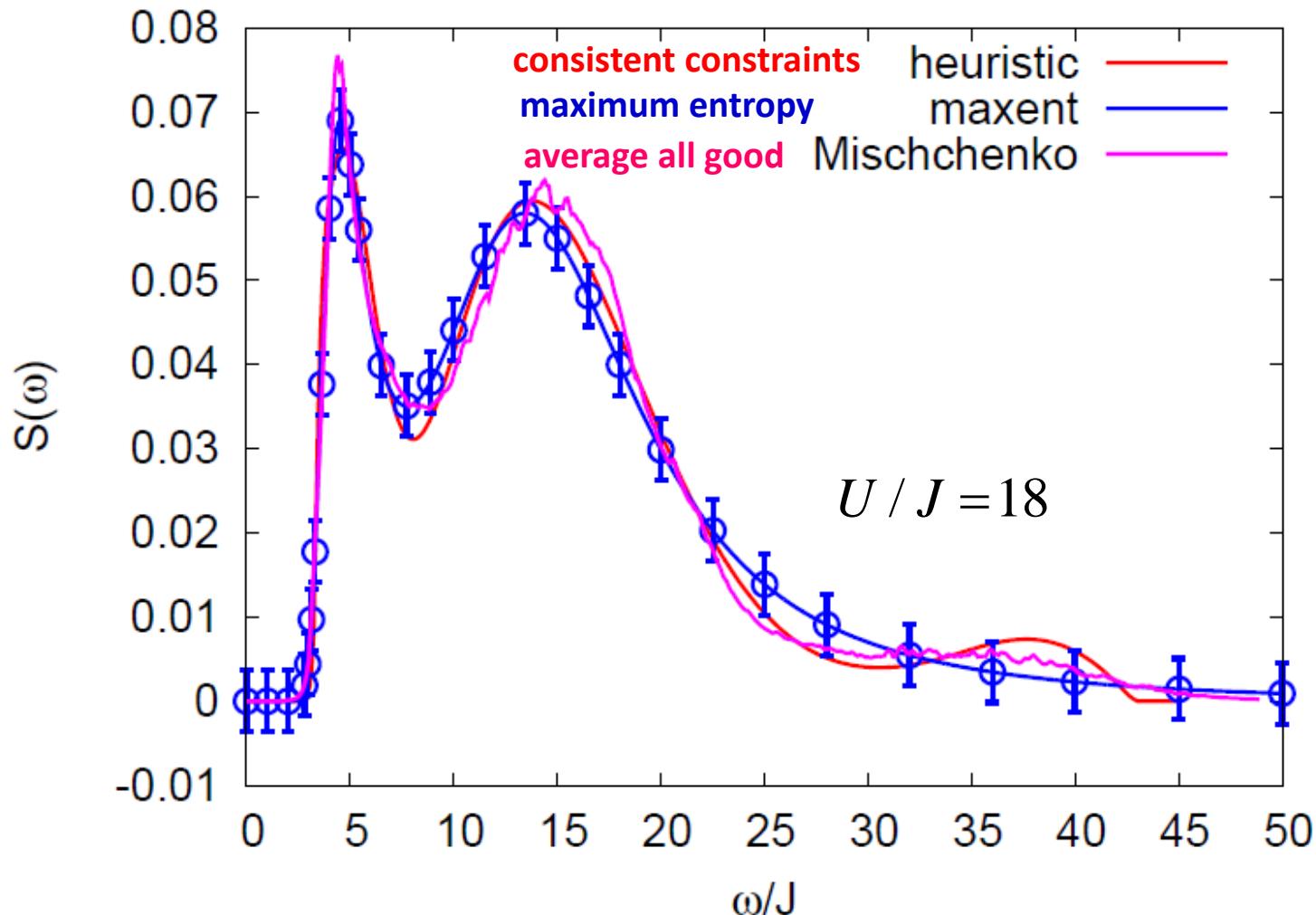
$$C(\tau) = \frac{1}{2\pi} \int_0^\infty e^{-\omega\tau} S(\omega) d\omega$$

$$S(\omega) = 2 \operatorname{Im} \chi(\omega) / (1 - e^{-\beta\omega})$$



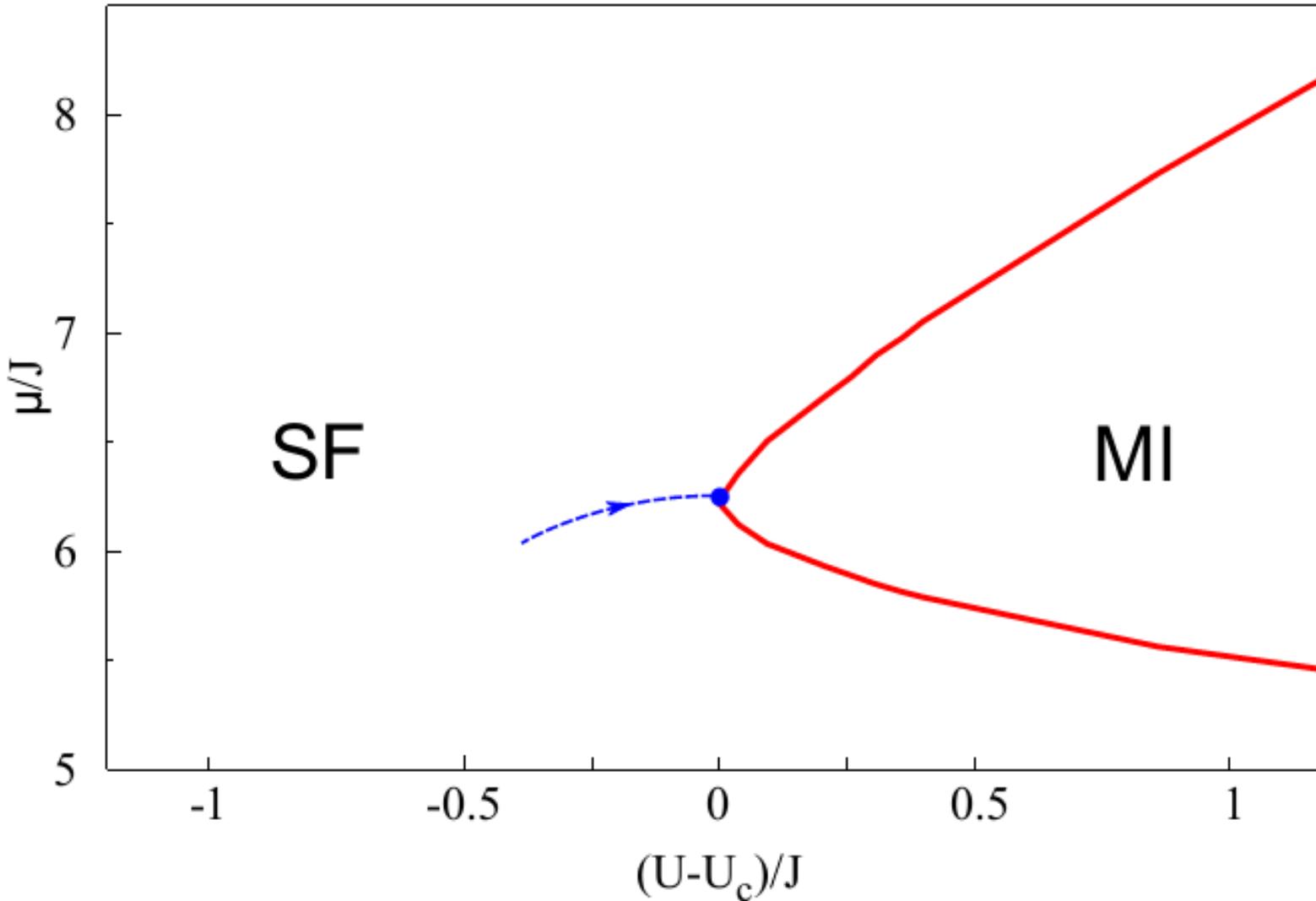
Hopefully, we know what we are doing ...

# Reliability of Analytical continuation

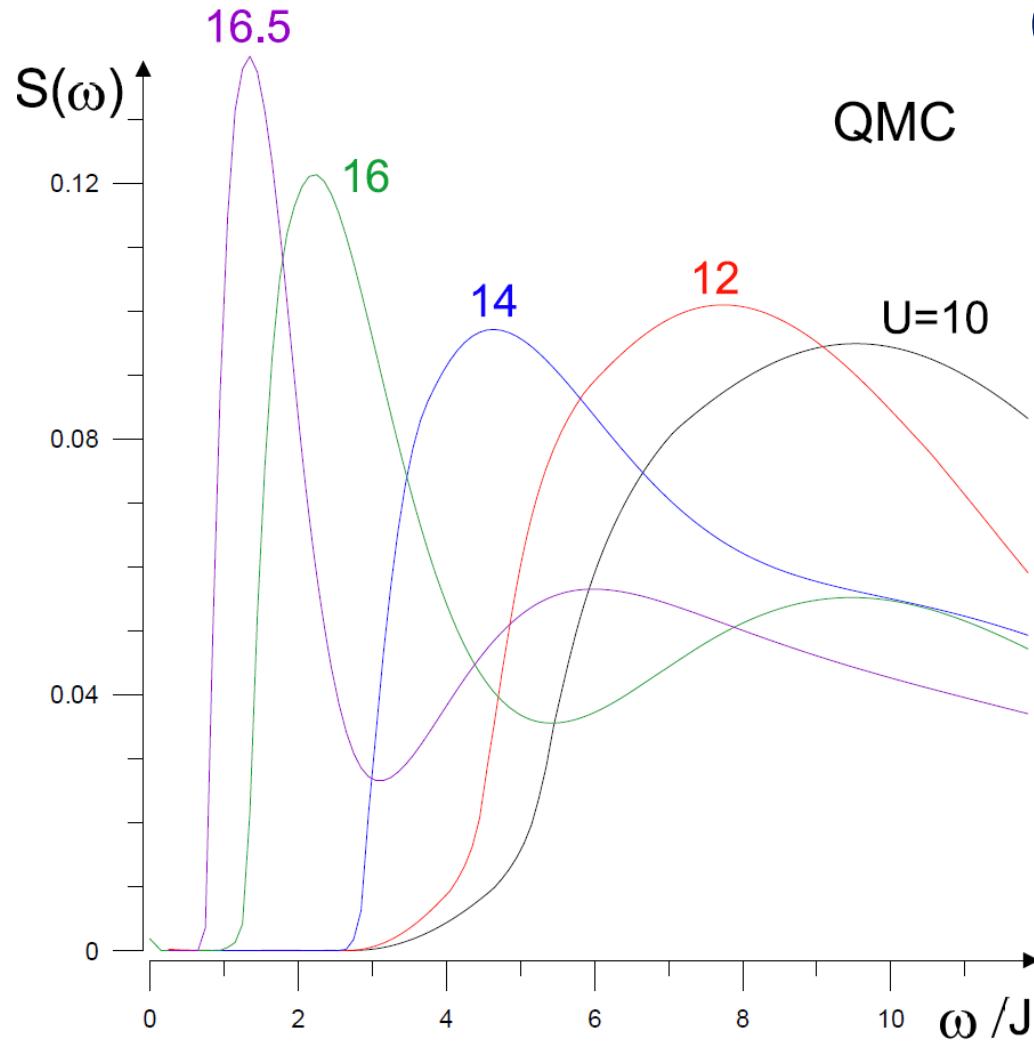


L. Pollet and N. Prokof'ev, PRL 109, 010401 (2012).

# In Superfluid



# In Superfluid



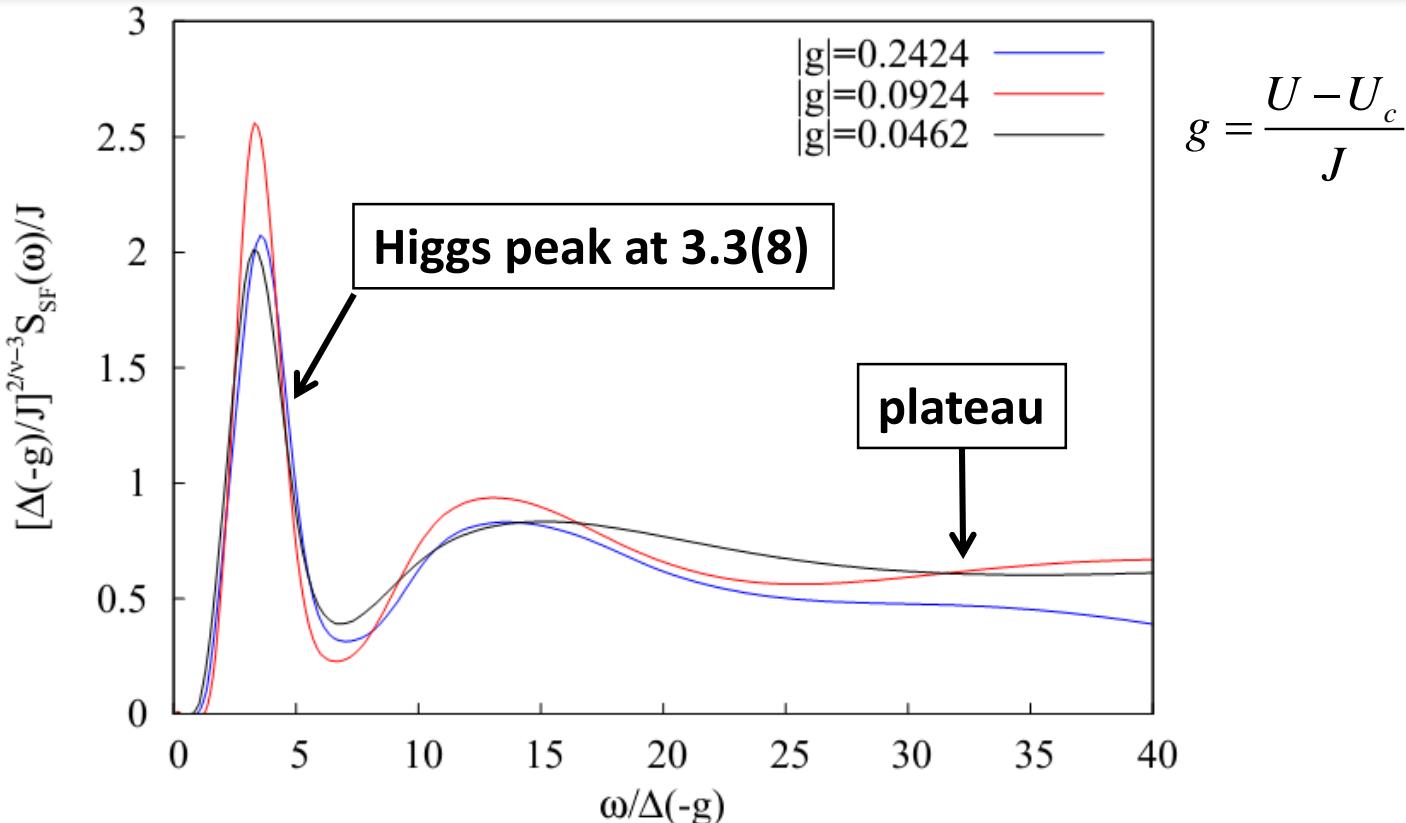
Oops, Higgs resonance is there!

$$(U / J)_c = 16.7424(2)$$

There is a resonance at low frequency which

- emerges at  $U \geq 14$
- softens as  $U \rightarrow U_c$
- gets more narrow as  $U \rightarrow U_c$

# In Superfluid

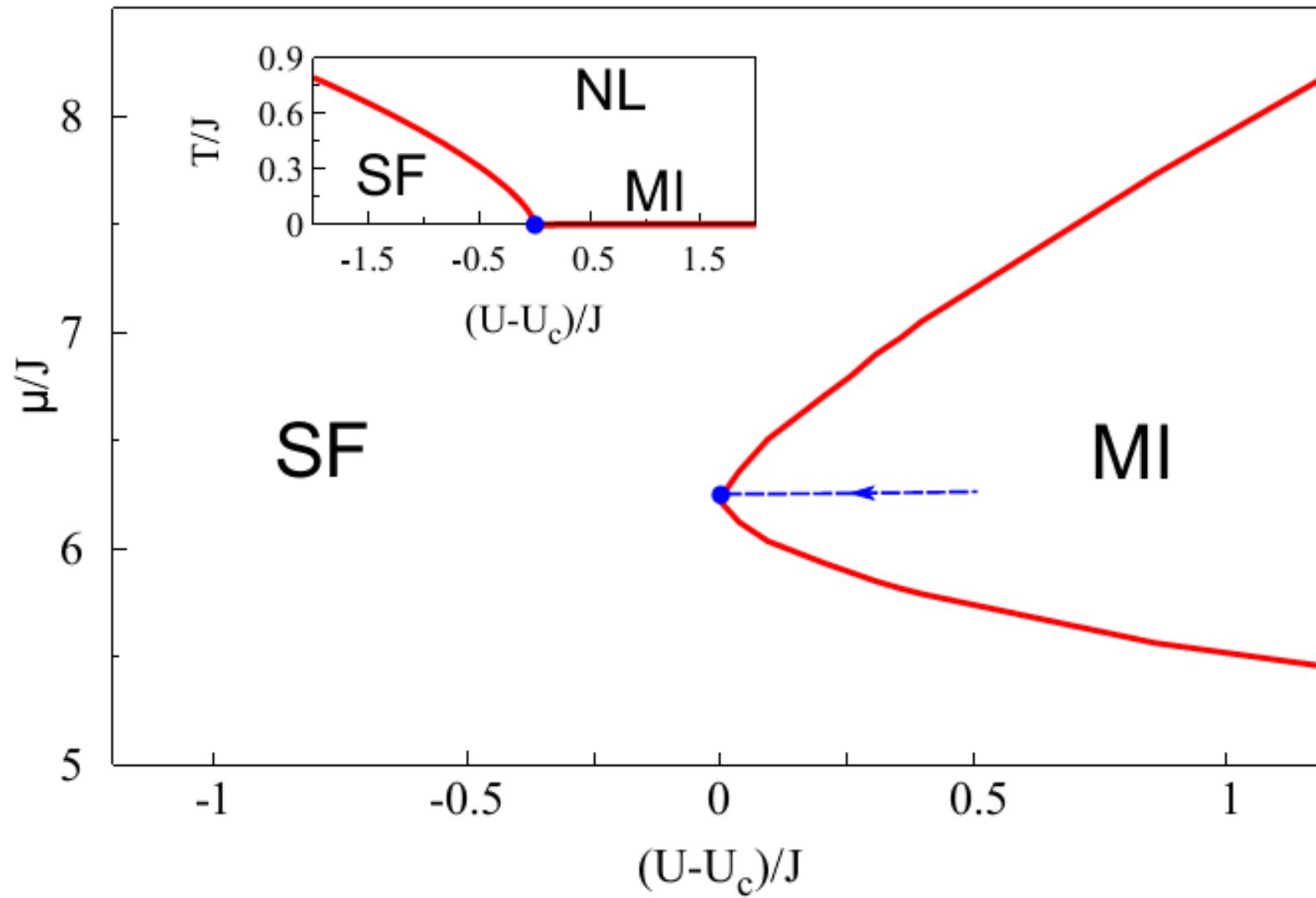


Chen, Liu, Deng, Pollet, and Prokof'ev, PRL 110, 170403 (2013)

★Nice Collapse: universal response function!

★Clear existence of plateau—a must physics condition

# In Mott Insulator



# Universal scaling spectral function

Chubukov, Sachdev, Ye '93

- Similar form for SF and MI phases:

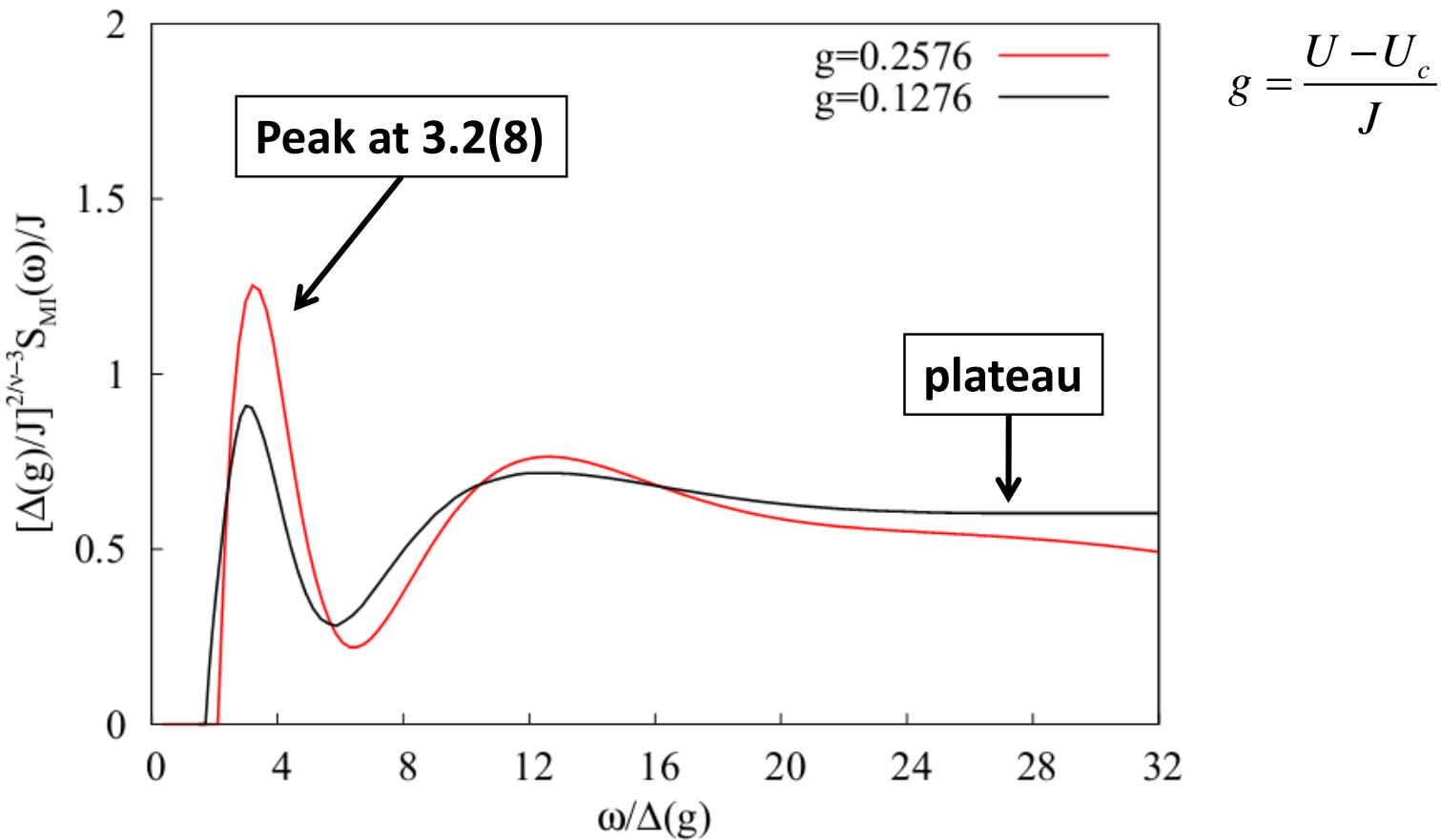
Sachdev '99

$$S_{SF,MI}(\omega) \propto \Delta^{3-2/\nu} \Phi_{SF,MI}(\omega / \Delta), \nu = 0.6717$$

$g = \frac{U - U_c}{J}$	Superfluid $\Delta(-g) = E_{gap} / 2$	MI $\Delta(g) = E_{gap} / 2$
$\omega \ll \Delta$	$\sim \omega^3$	$\sim \Theta(\omega - 2\Delta)$
$\omega \gg \Delta$	$\sim \omega^{0.0225}$	$\sim \omega^{0.0225}$

*A plateau!*

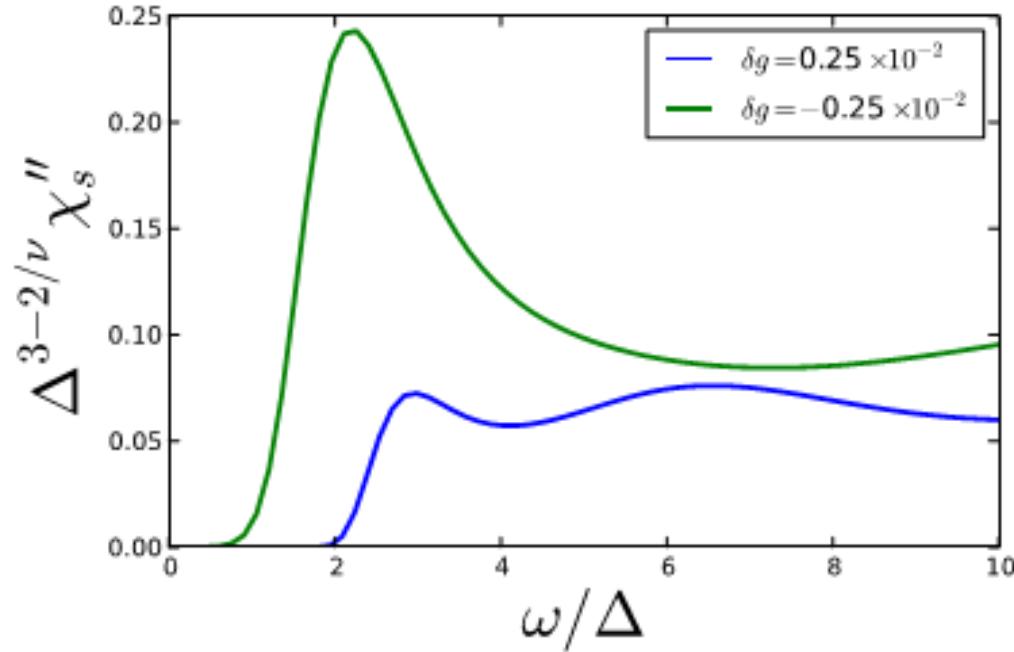
# In Mott Insulator



Chen, Liu, Deng, Pollet, and Prokof'ev, PRL 110, 170403 (2013)

- ★ What is the physics of the resonance peak?
- ★ Looks so much as in the superfluid phase!

# Higgs mode in Relativistic O(2) model



Discrete version of relativistic O(2) model

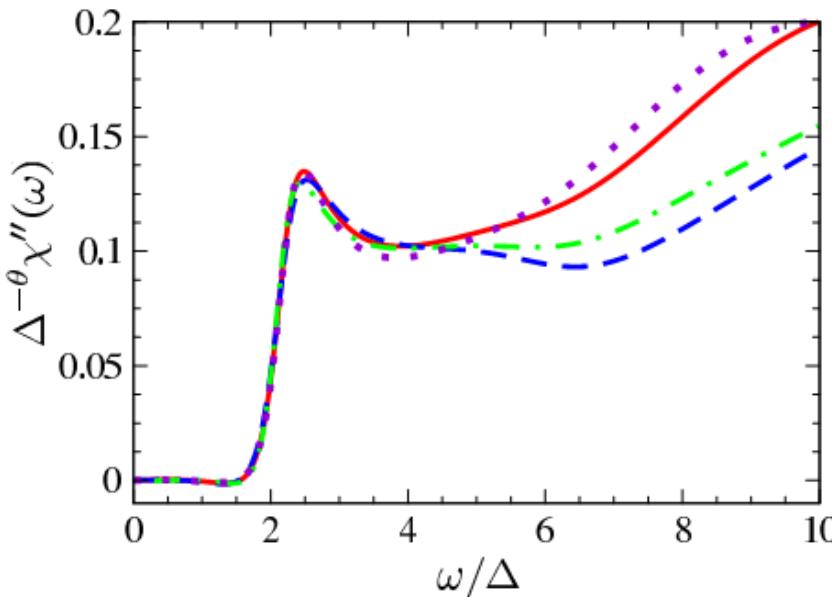
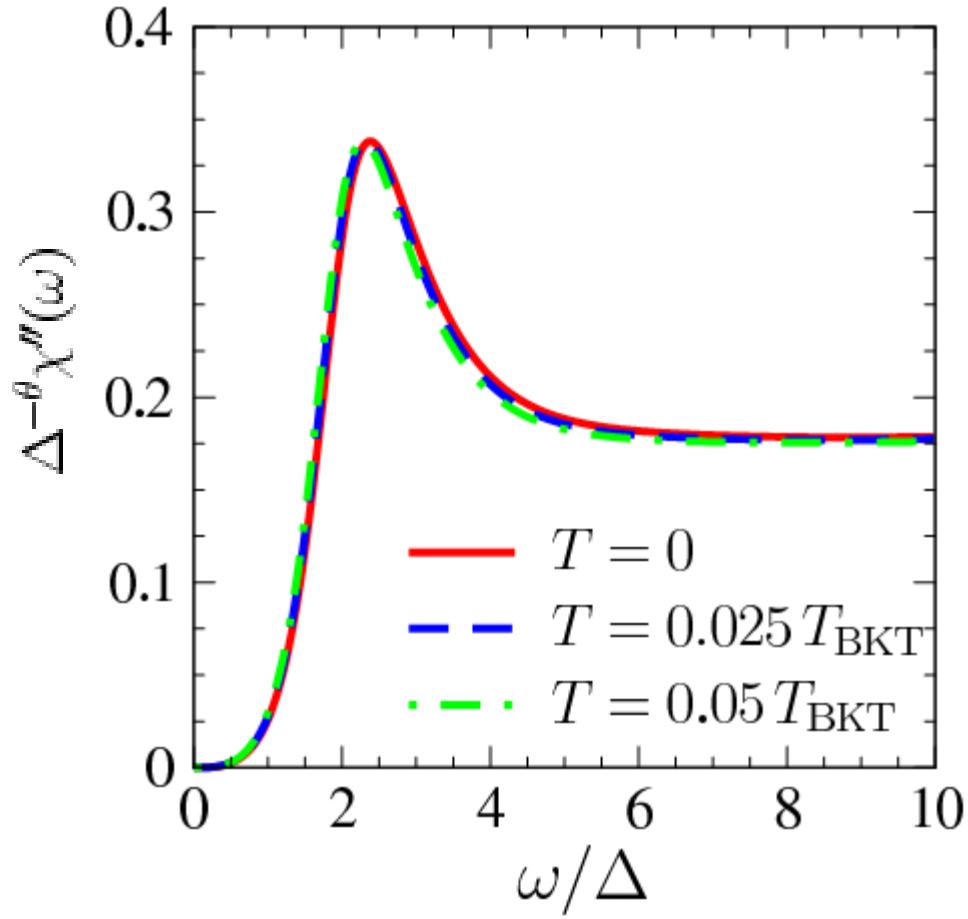
$$\chi(\omega) = \left\langle |\psi(t)|^2 |\psi(0)|^2 \right\rangle_{\omega}$$

$$\omega_H / \Delta \approx 2.1(3)$$

S. Gazit, D. Podolsky, A. Auerbach, PRL 110, 140401 (2013)

S. Gazit, D. Podolsky, A. Auerbach, and D.P. Arovas, Phys. Rev. B 88, 235108

# Nonperturbative renormalization group study



$$\chi(\omega) = \left\langle |\psi(t)|^2 |\psi(0)|^2 \right\rangle_\omega$$

$$\omega_H / \Delta \approx 2.4$$

# Numerics vs Experiment

With a harmonic trap term:  $V_i = \frac{1}{2} \alpha \omega^2 r_i^2$  in  $\hat{H}_{BH}$

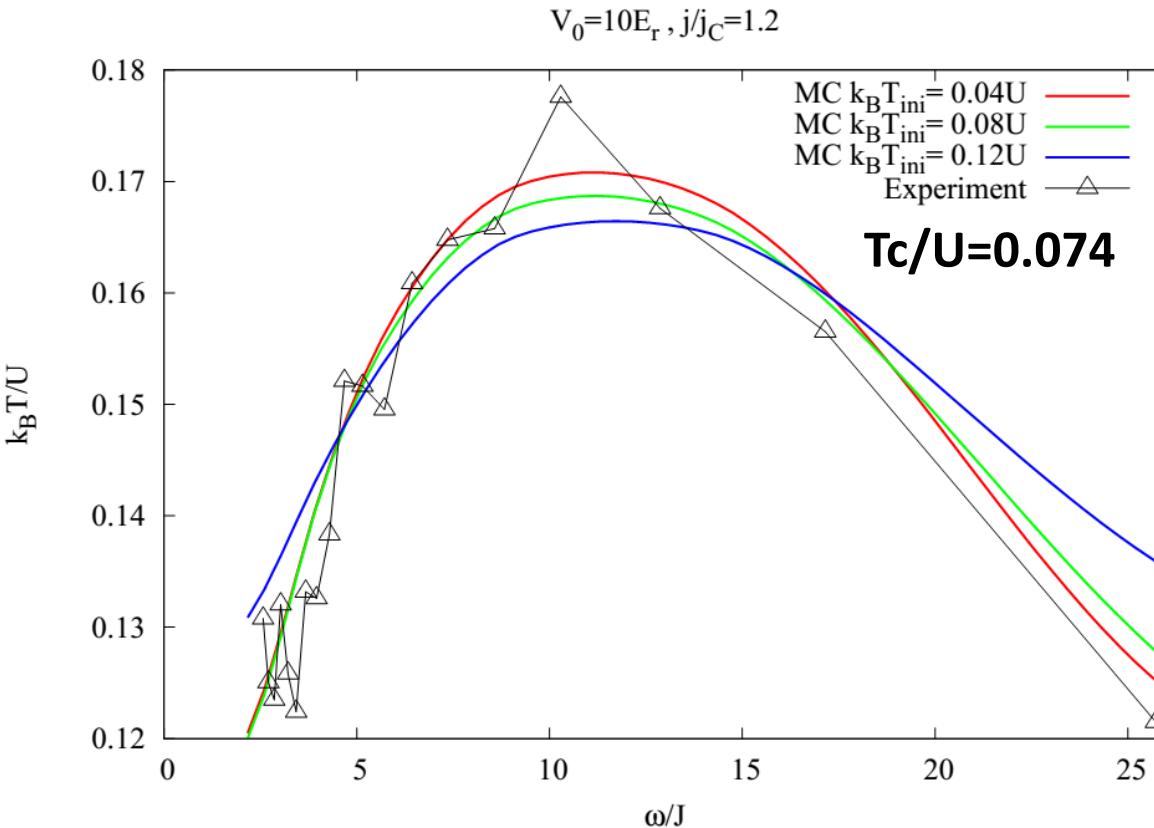
$T_0 = 0.12U$

$\dot{E} = \frac{1}{2} \omega S(\omega, T_i) \left( \frac{\delta J}{J} \right)^2$

$\dot{T}_i = (\dot{E} + Q) / C(T_i) = f(T_i, \omega)$

$T_{i+1} = T_i + f(T_i, \omega) \cdot dt$

$T_{SF}$

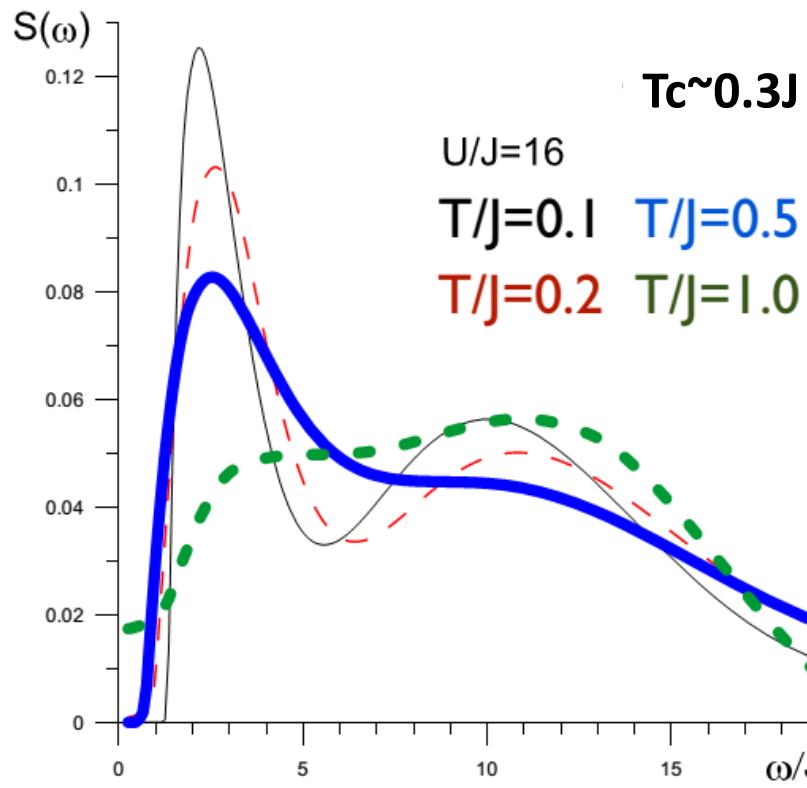


K. Chen et.al, unpublished

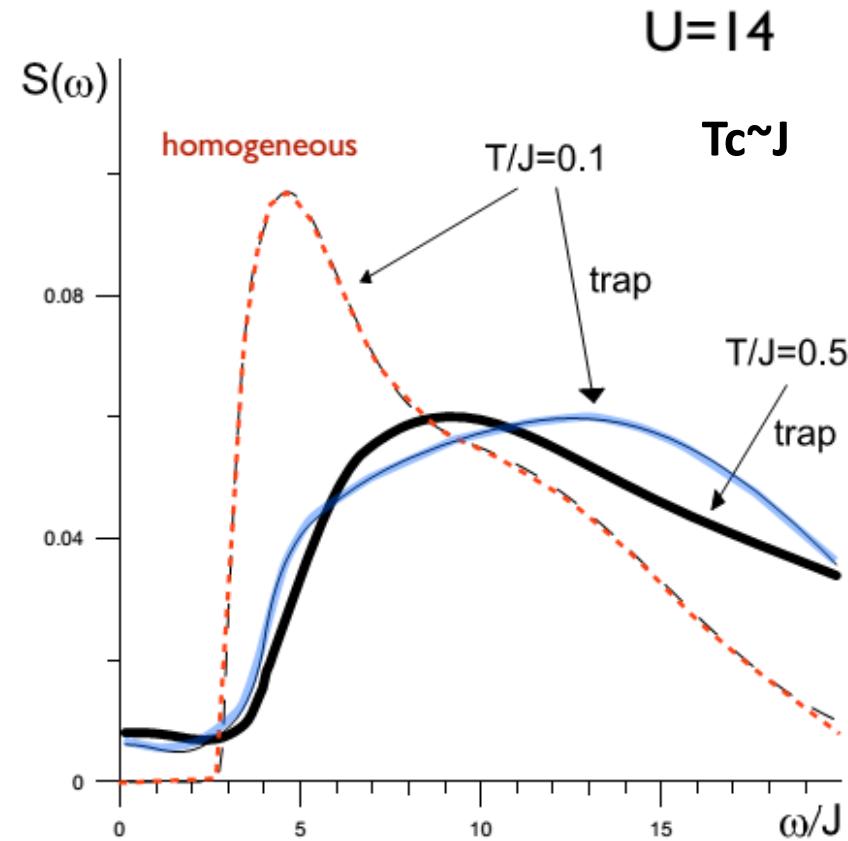
Consistent with experimental result!

# Why no resonance peak?

Temperature effect

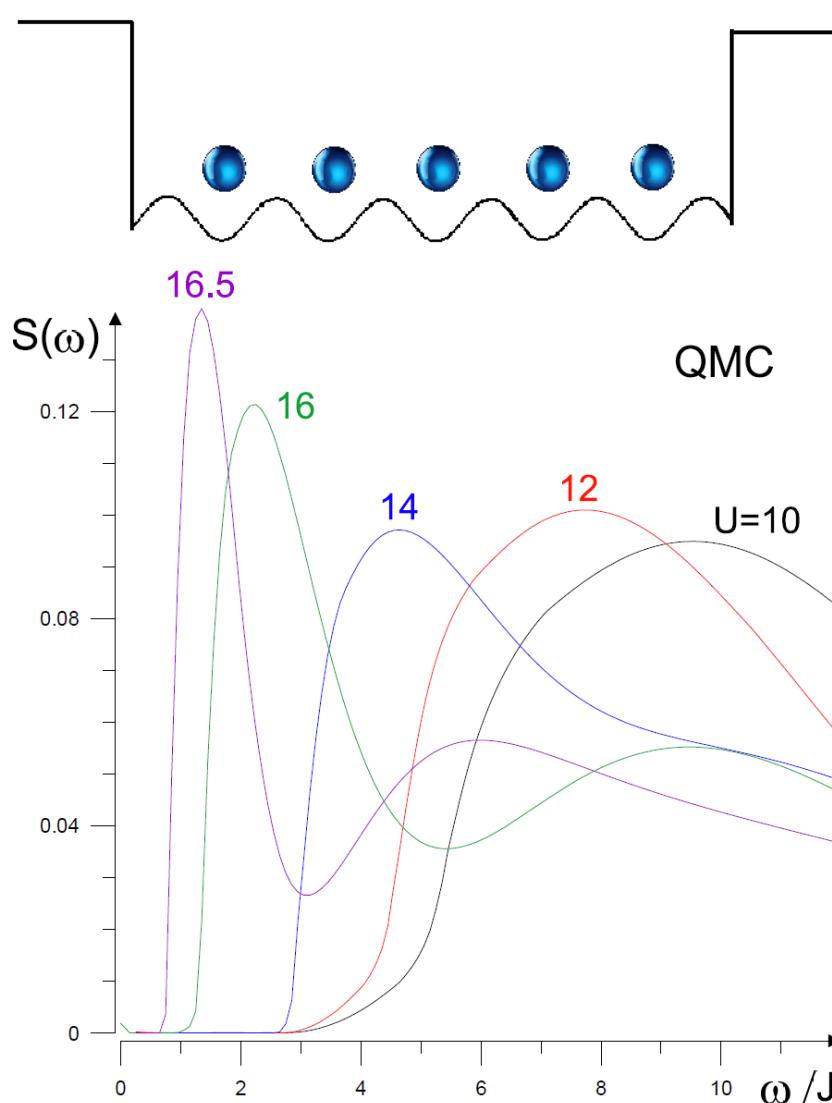


Trap effect



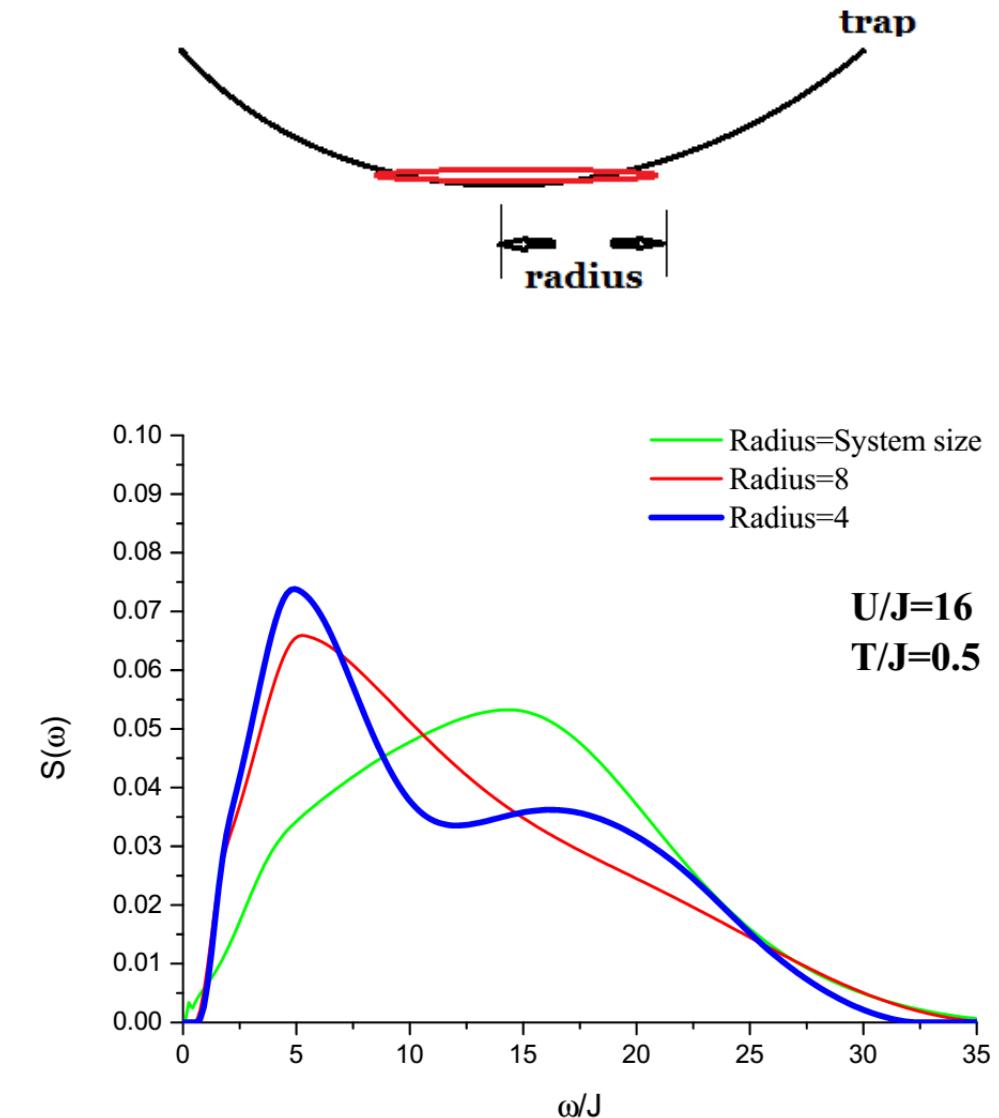
Pollet and Prokof'ev, PRL 109, 010401 (2012).

# How to reproduce resonance peak?



QMC

OR



# Summary

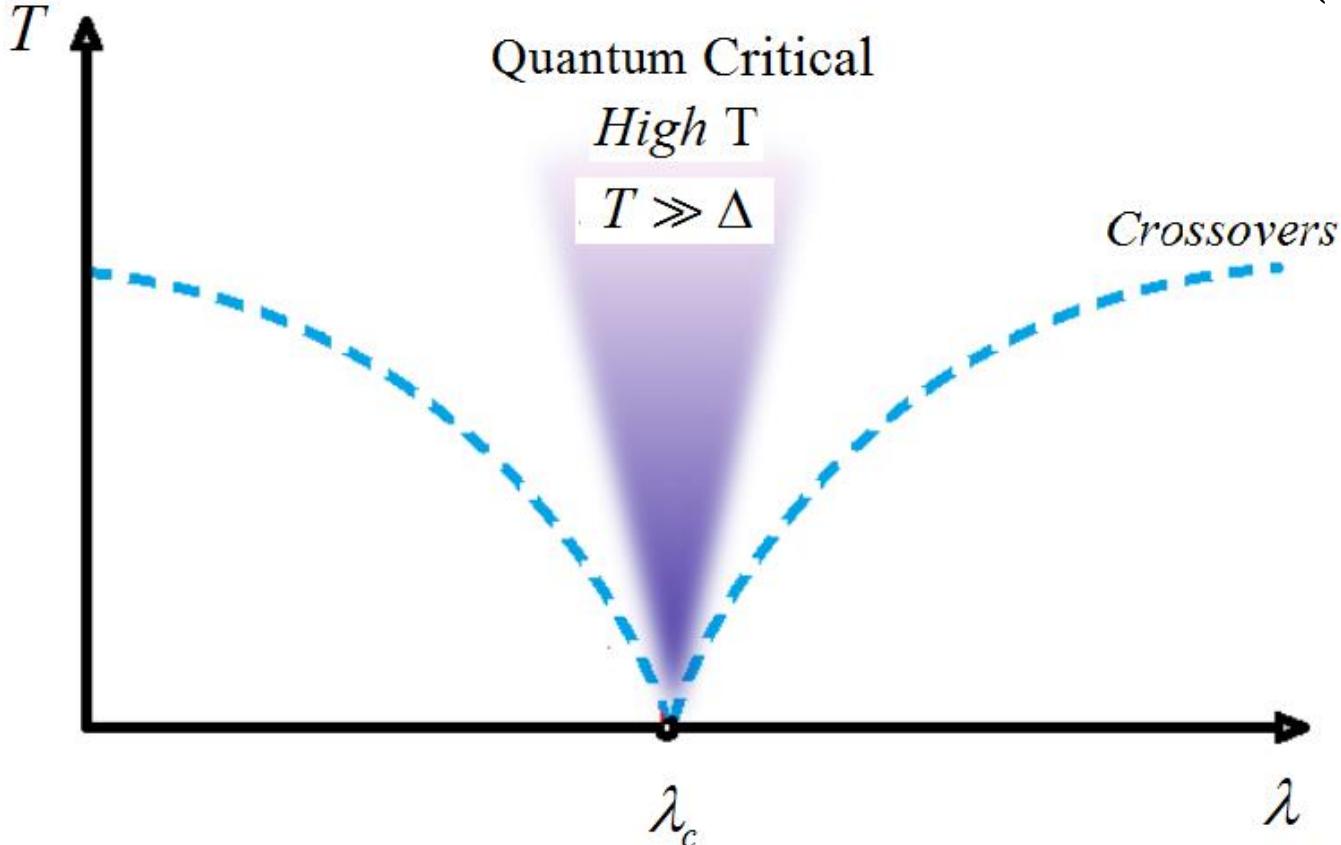
- Higgs(-like) mode exists in SF and MI.
- Universal spectral functions:
  - Sharp, low-frequency resonance
  - Damped second oscillation
  - Saturation to quantum critical plateau
- A direct observation of Higgs resonance peak in experiment still need to be done.

# When dynamics meets quantum criticality

High-T regime:  $T > 0, \lambda \rightarrow \lambda_c$

➤ Only one energy scale T !

➤ Scaling of any dynamic quantity:  $S(\omega) \propto F(\omega/T, T \rightarrow 0)$

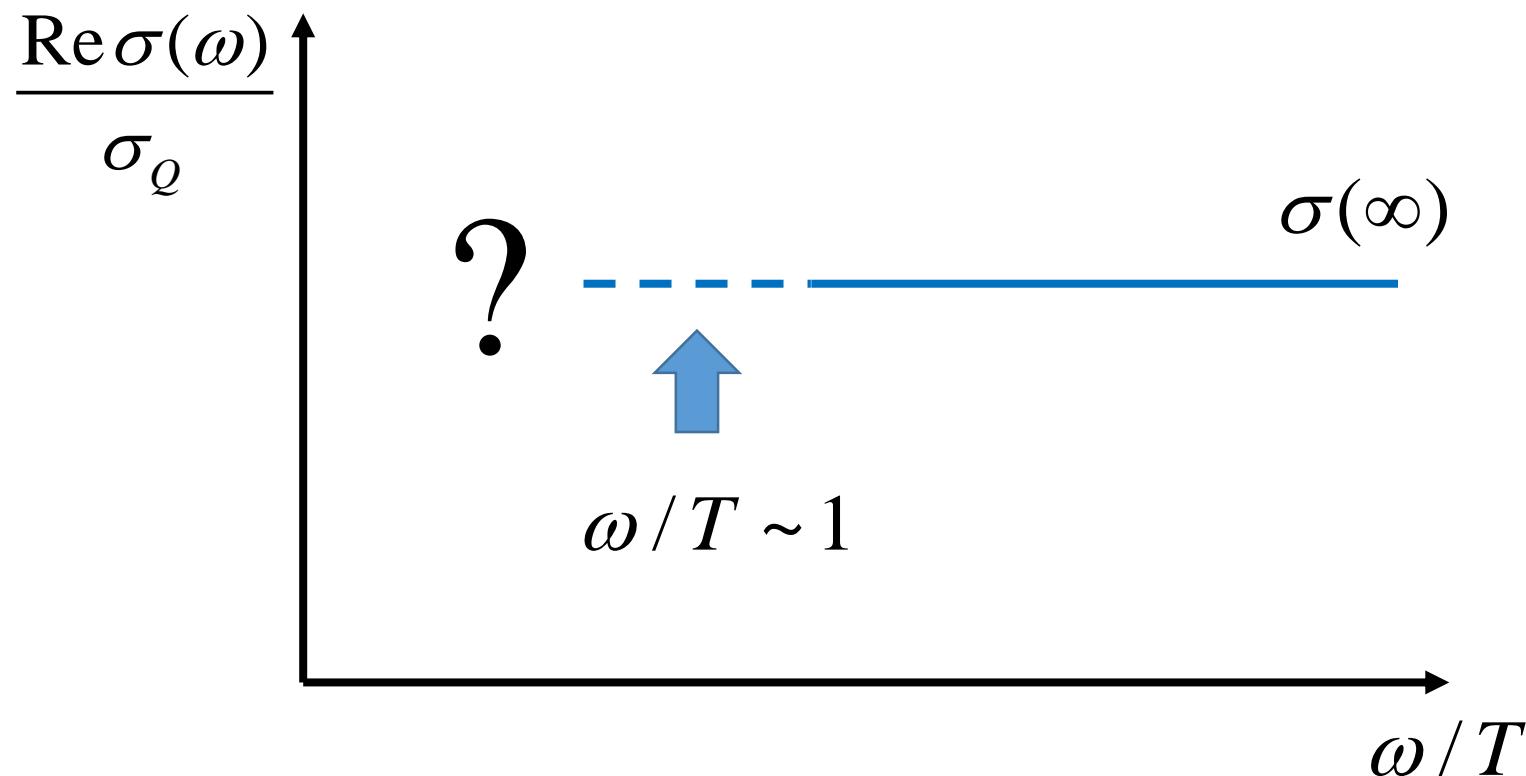


# Universal Conductivity in high-T regime

For a two-dimensional quantum critical system

$$\sigma(\omega) = \sigma_Q \Sigma(\omega/T, T \rightarrow 0) \quad \text{or} \quad \sigma(i\omega_n) = \sigma_Q \Sigma(\omega_n/T, T \rightarrow 0)$$

K. Damle and S. Sachdev, Phys. Rev. B56, 8714 (1997)



# Holographic insight

A holographic theory (AdS4) suggests:

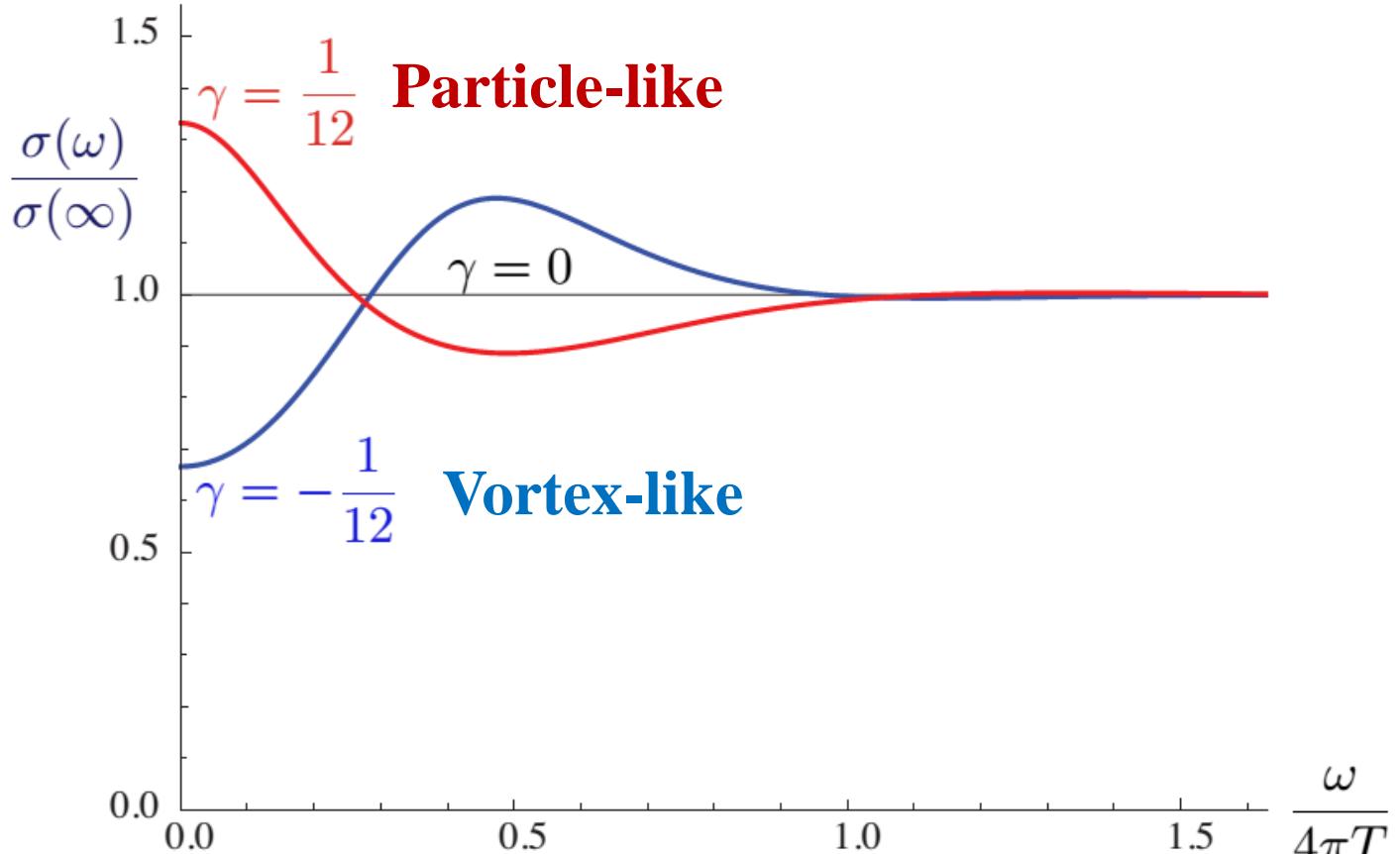
The universal conductivity curve is controlled by two parameters:

$$\sigma(0) \text{ and } \sigma(\infty)$$

So define

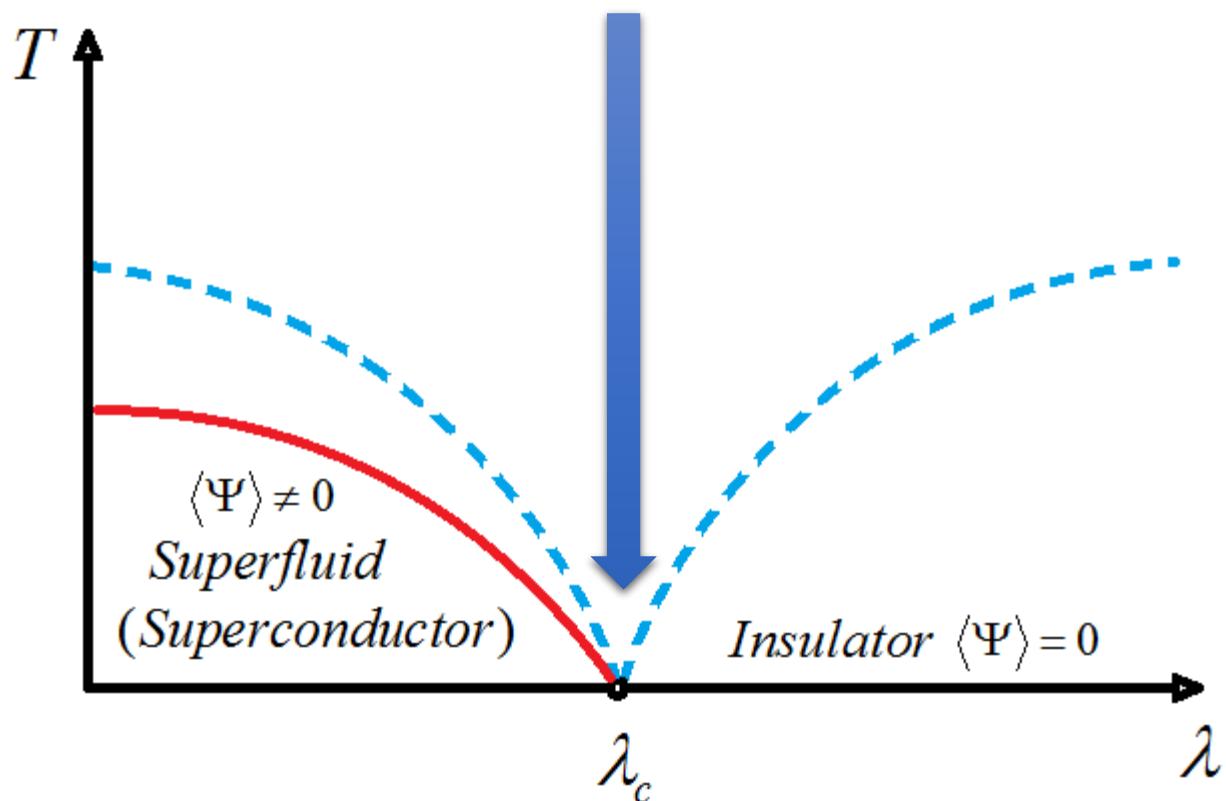
$$\gamma = \frac{\sigma(0) - \sigma(\infty)}{4\sigma(\infty)}$$

$$= \begin{cases} < 0 & \text{vortex-like} \\ > 0 & \text{particle-like} \end{cases}$$



R. C. Myers, S. Sachdev, and A. Singh, *Physical Review D* **83**, 066017 (2011)

# Could it describe the simplest CFT3 --relativistic O(2) QCP?



# Models characterize 3D XY universality

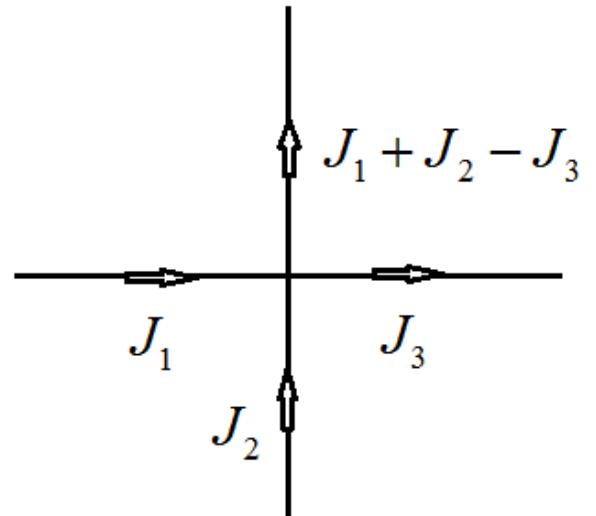
- Bose Hubbard model at quantum critical point

$$\hat{H}_{BH} = -J \sum_{\langle i,j \rangle} b_i^\dagger b_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) - \mu \sum_i \hat{n}_i$$

- J current model (Villain Model, classical) at critical point

$$H = \frac{1}{2K} \sum_{\langle ij \rangle} \nabla \mathbf{J} = 0 J_{\langle i,j \rangle}^2 \quad J_{\langle i,j \rangle} \in (-\infty, \infty)$$

- square lattice
- Monte Carlo with worm algorithm



# Measurement of conductivity

- Kubo formula in Matsubara frequency:

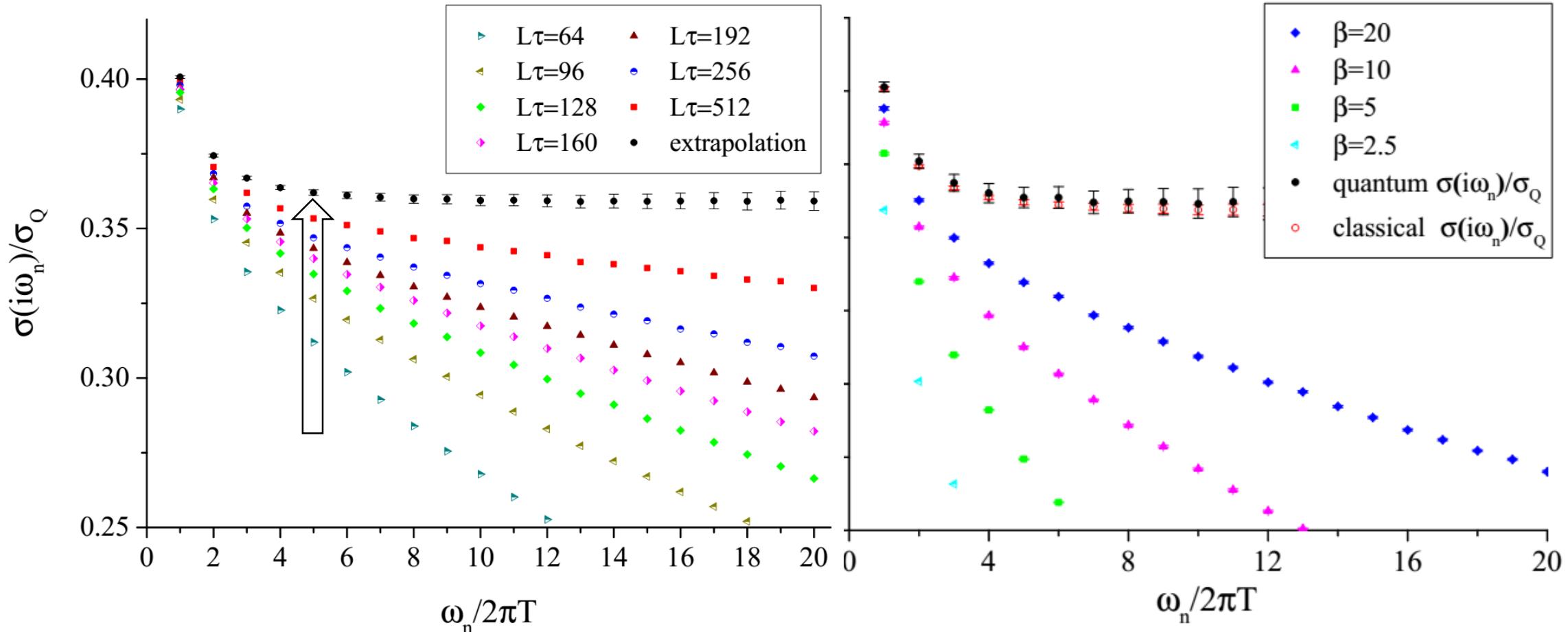
$$\sigma(i\omega_n) \propto \frac{\langle -\mathcal{E}_{kinect} \rangle - \langle j(\tau)j(0) \rangle_{i\omega_n}}{\omega_n}, \quad \omega_n = 2\pi n T, \quad (n=1,2,3\dots)$$

- Analytical continuation(  $\sigma(i\omega_n) \xrightarrow{i\omega_n \rightarrow \omega} \sigma(\omega)$  )

$$\sigma(i\omega_n) = \frac{2}{\pi} \int_0^\infty \frac{\omega_n}{\omega^2 + \omega_n^2} Re\sigma(\omega) d\omega$$

- ill-defined problem
- Don't expect too much for  $\omega/T < \omega_1/T = 2\pi$

# J current model vs BH model



**Extrapolate all curves to  $T=0$  with finite temperature correction hypothesis:** 
$$\sigma(i\omega_n/T, T) = \sigma(i\omega_n/T, T \rightarrow 0) + A_n T^{0.85}$$

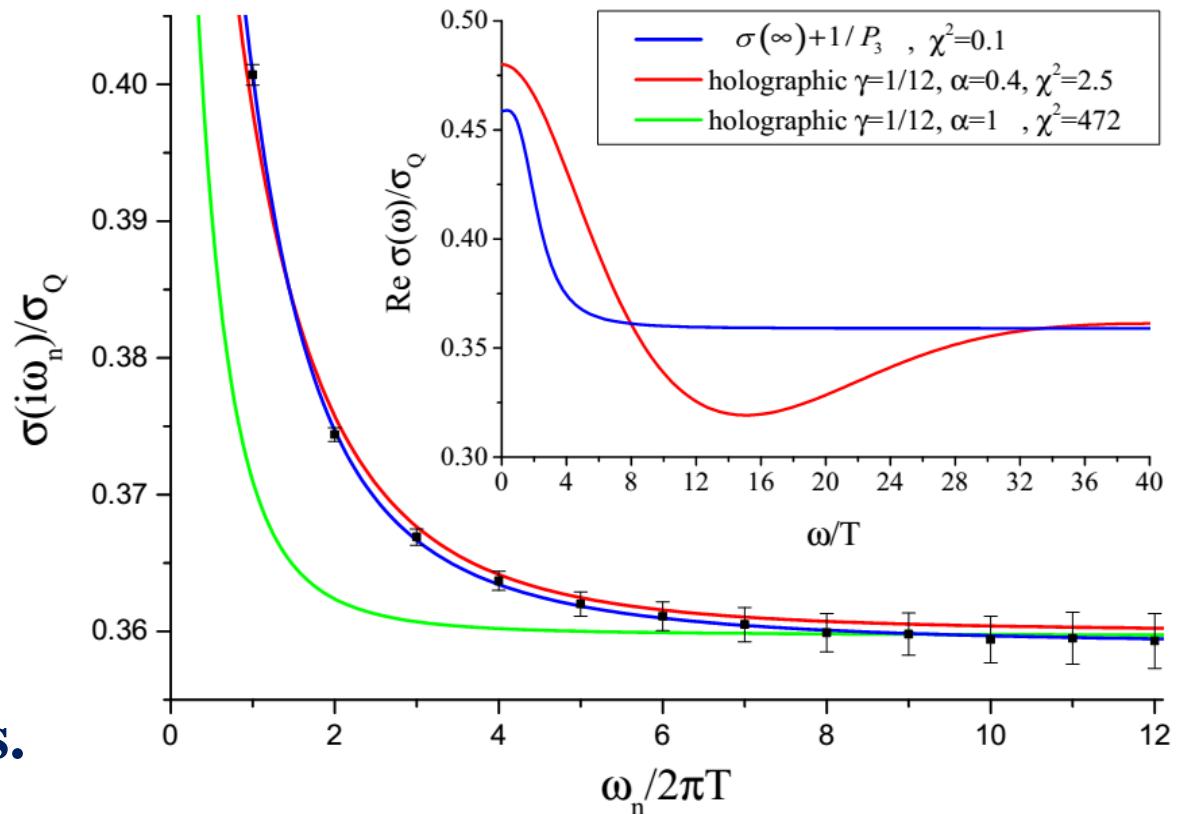
# Main results

- Universal conductivity at large frequency:

$$\sigma(\infty) = 0.359(4)$$

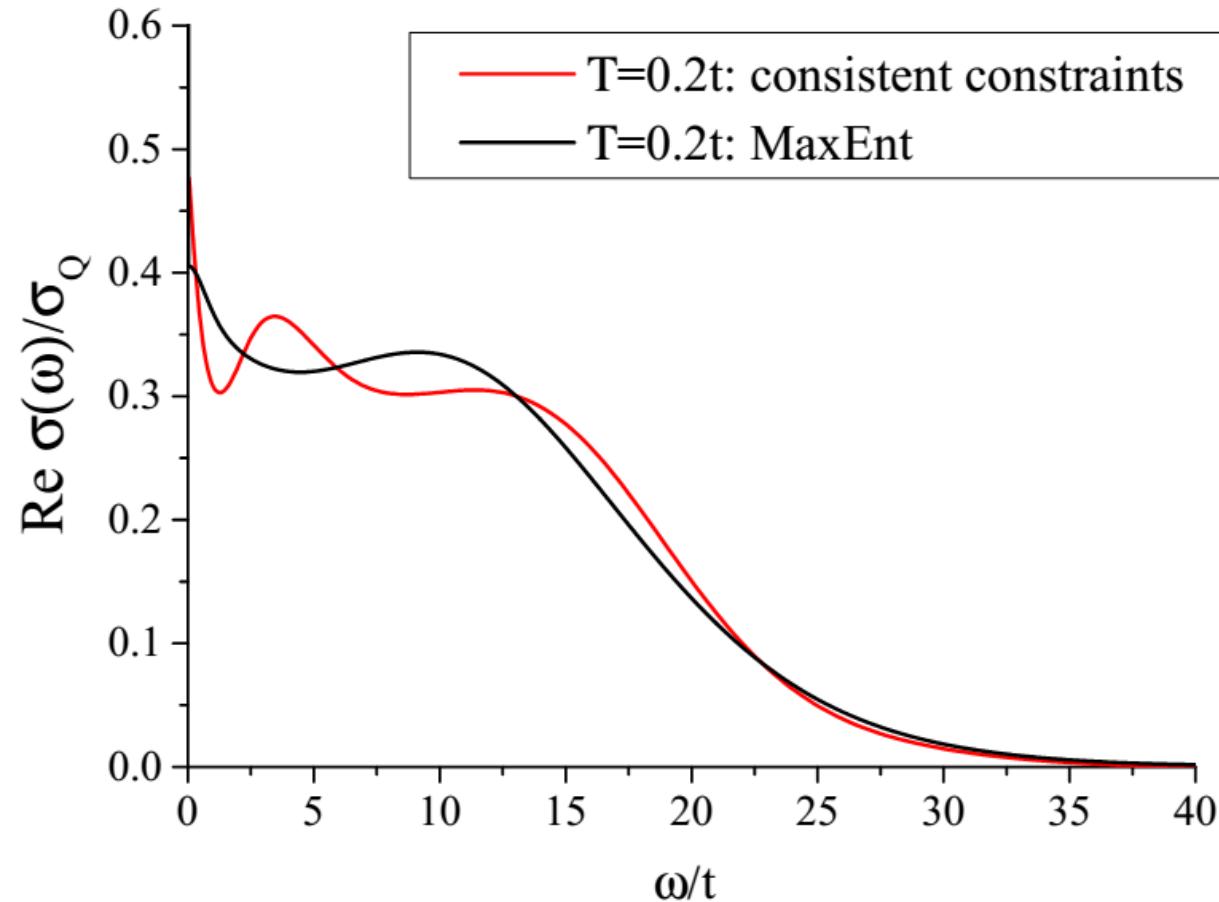
- Particle-like transport.
- The holographic result can not fit our data unless we are allowed to rescale the temperature of black hole by a factor  $\sim 0.4$ .
- Low frequency part of  $\sigma(\omega/T)$  is unstable against different analytical continuation procedures.

S. Gazit et.al. Phys. Rev. B 88, 235108 (2013).  
W. Witczak-Krempa et.al. arXiv:1309.2941 (2013).  
M. Swanson et.al. arXiv. 1310.1073 (2013).



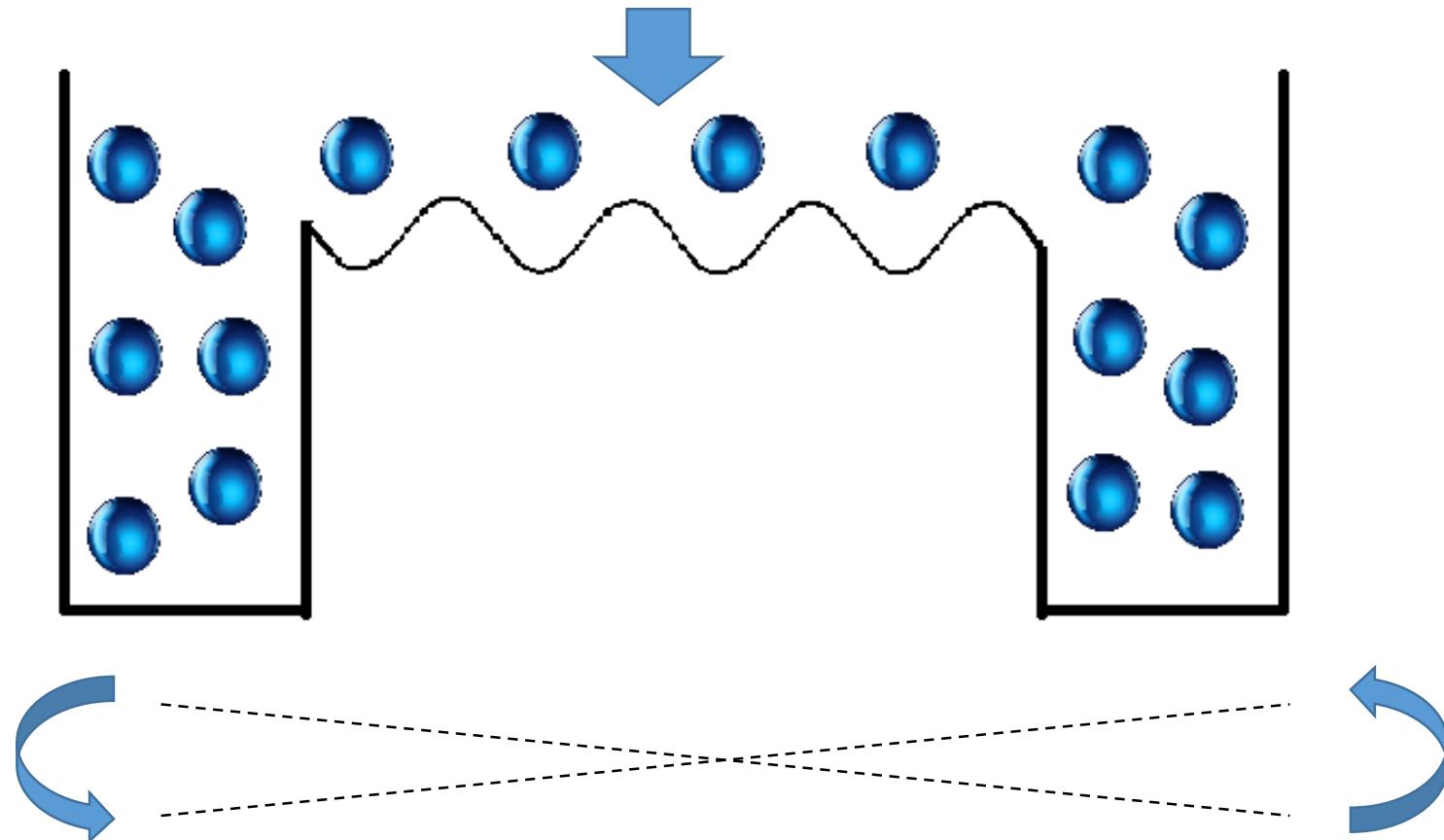
# Experimental accessible conductivity

The plateau value  $\sigma(\infty)$  can also be recovered at experimental accessible temperature



# Ultracold atoms meet quantum gravity

Lower the temperature to  $T/t < 1 \sim 2$



“Sway” the lattice with frequency  $\omega$

# Conclusion

**Phys. Rev. Lett. 110, 170403 (2013). Editors' suggestion**

- Universal conductivity at large frequency:  
 $\sigma(\infty) = 0.359(4)$
- Particle-like transport is found.
- The holographic result can not fit our data unless we are allowed to rescale the temperature of black hole by a factor ~0.4.