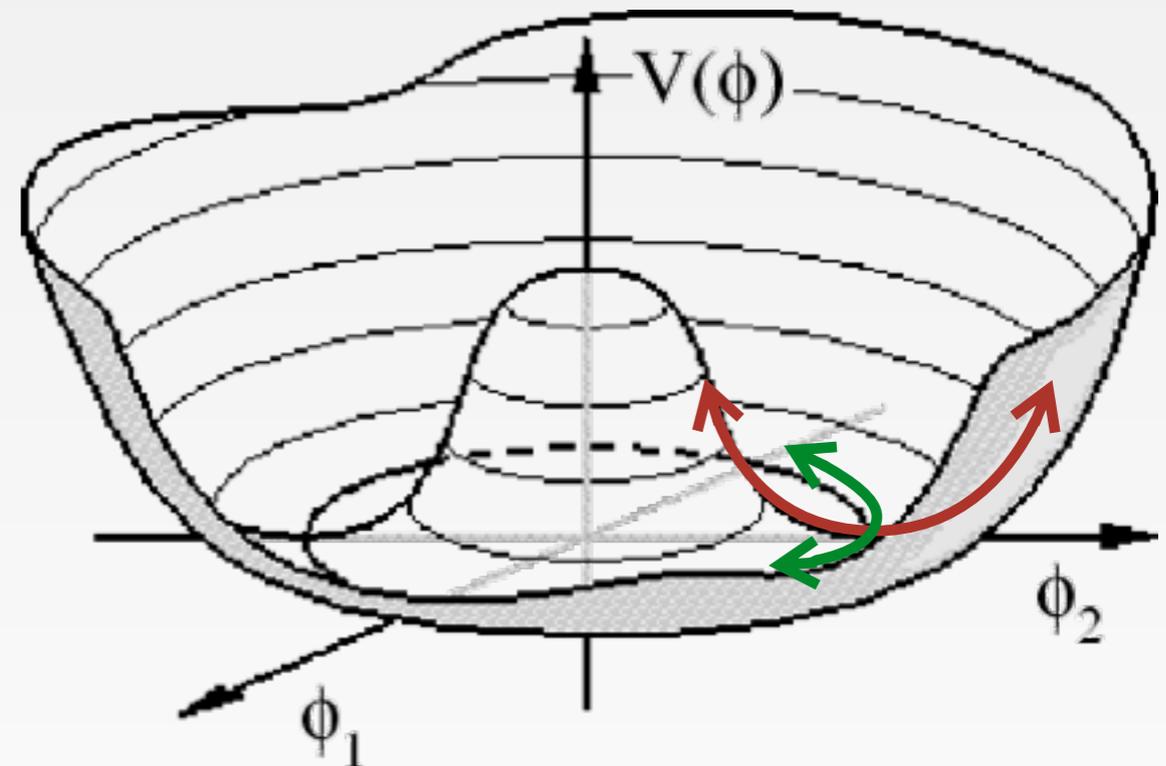


Higgs mode and universal dynamics near quantum criticality

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June 2014



Podolsky, Auerbach, Arovas, PRB 84, 174522 (2011)

Podolsky and Sachdev, PRB 86, 054508 (2012)

Gazit, Podolsky, Auerbach, PRL 110, 140401 (2013)

Gazit, Podolsky, Auerbach, Arovas PRB 88, 235108 (2013)

Tenenbaum Katan and Podolsky, to appear (2014)

Gazit, Podolsky, Auerbach, to appear (2014)



Technion
Israel Institute
of Technology

Outline

- The Higgs mode in condensed matter
- How to see the Higgs mode in 2d?
- Universal dynamics near quantum criticality
- Can the Higgs be seen through conductivity measurements?
- Amplitude mode in solid Helium 4?

Spontaneous Symmetry Breaking

SSB - Ground state has less symmetry than the Hamiltonian

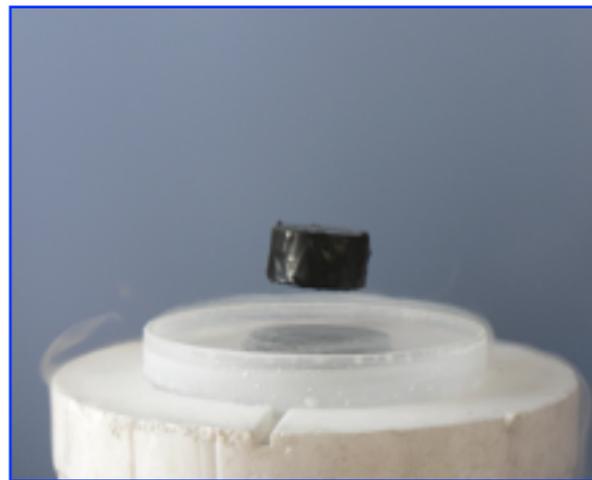
Ubiquitous phenomenon in physics. Order parameter $\phi(x,t)$ gets an expectation value. Usually results in massless Goldstone bosons.

physical system	broken symmetry	Goldstone bosons
ferromagnets	spin rotational invariance	spin waves
crystals, density waves	translational invariance	phonons and phasons
superfluids	phase rotation invariance	phonons



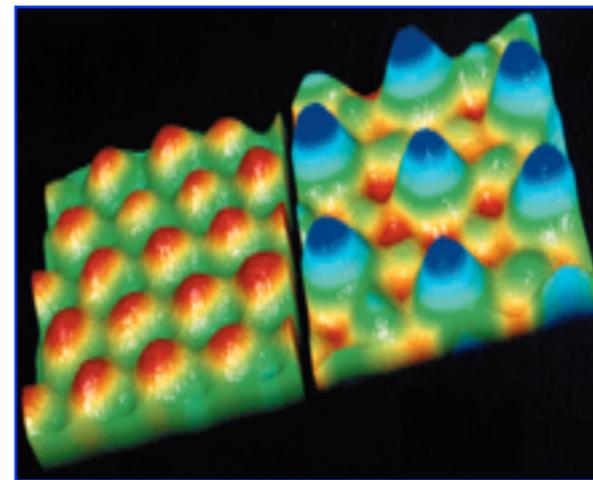
MAGNETS

magnetization M



SUPERCONDUCTORS

pair field $\Psi = |\Psi| e^{i\phi}$



DENSITY WAVES

Fourier mode $|\varrho_G| e^{i\phi_G}$



THE UNIVERSE

Higgs field $\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$

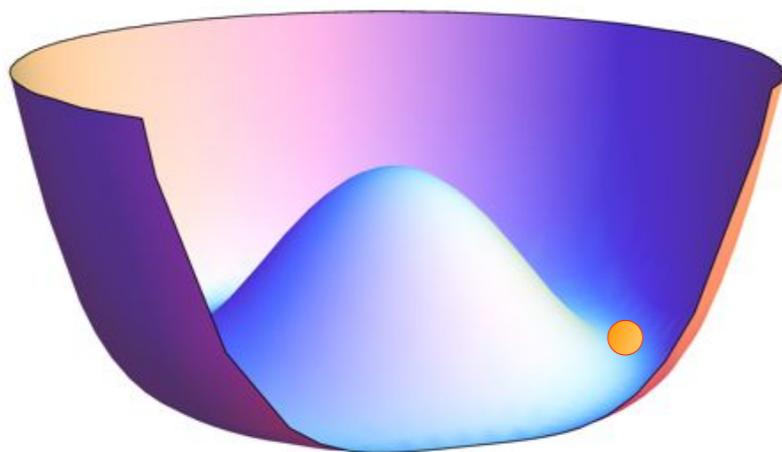
Spontaneous Symmetry Breaking

N -component order parameter:

$$\phi = \begin{pmatrix} \phi_1 \\ \phi_2 \\ \dots \\ \phi_N \end{pmatrix}$$

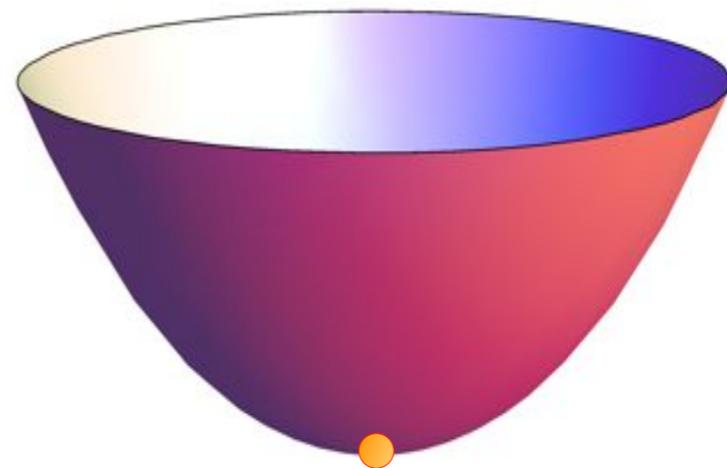
$$V(\phi) = g\phi^2 + u(\phi^2)^2$$

$g < 0$

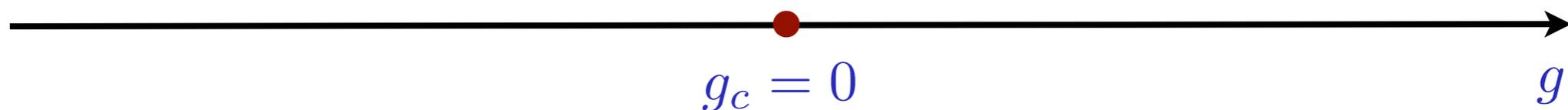


$\langle \phi \rangle \neq 0$

$g > 0$



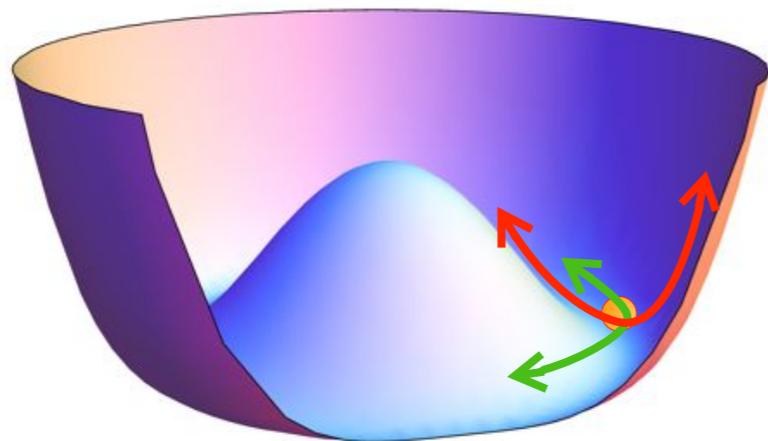
$\langle \phi \rangle = 0$



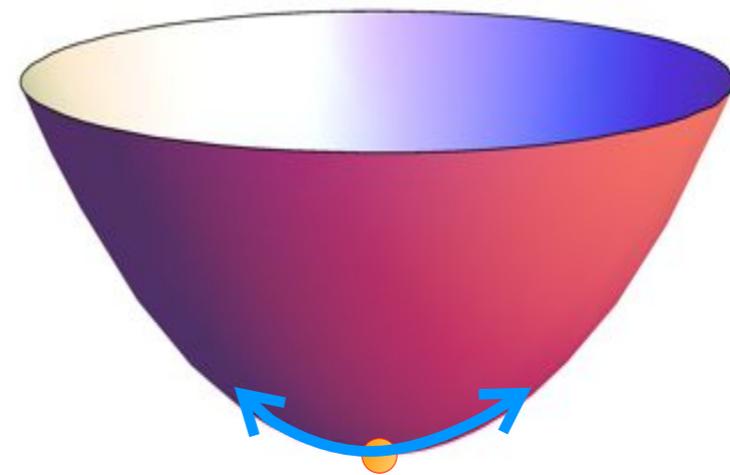
Collective excitations

Action:

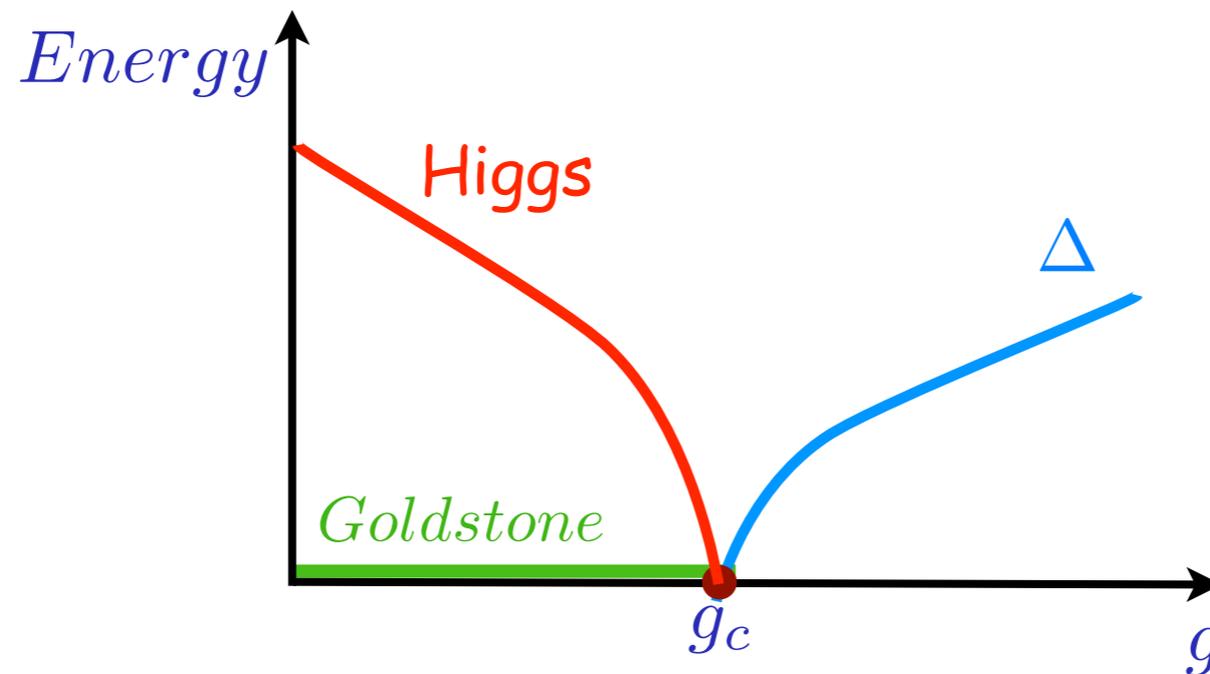
$$\mathcal{S} = \int d^d x dt [(\partial_t \phi)^2 - (\nabla \phi)^2 - g\phi^2 - u(\phi^2)^2]$$



$$\langle \phi \rangle \neq 0$$



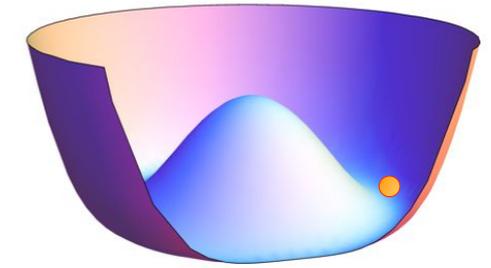
$$\langle \phi \rangle = 0$$



Higgs mechanism vs Higgs boson

Gauged O(2) theory :

$$\mathcal{L} = \frac{1}{2} |(\partial_\mu + ieA_\mu) \psi|^2 - m^2 (\psi^* \psi - 1)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$



Amplitude and phase: $\psi = (1 + \eta) e^{i\xi}$

Gauge invariance: $\psi \rightarrow e^{-i\xi} \psi$, $\tilde{A}_\mu \equiv A_\mu - e^{-1} \partial_\mu \xi$

Then:

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \eta)^2 + m^2 (\eta^2 + 2\eta^3 + 2\eta^4) + e^2 (1 + \eta)^2 \tilde{A}_\mu \tilde{A}^\mu - \frac{1}{4} \tilde{F}_{\mu\nu} \tilde{F}^{\mu\nu}$$

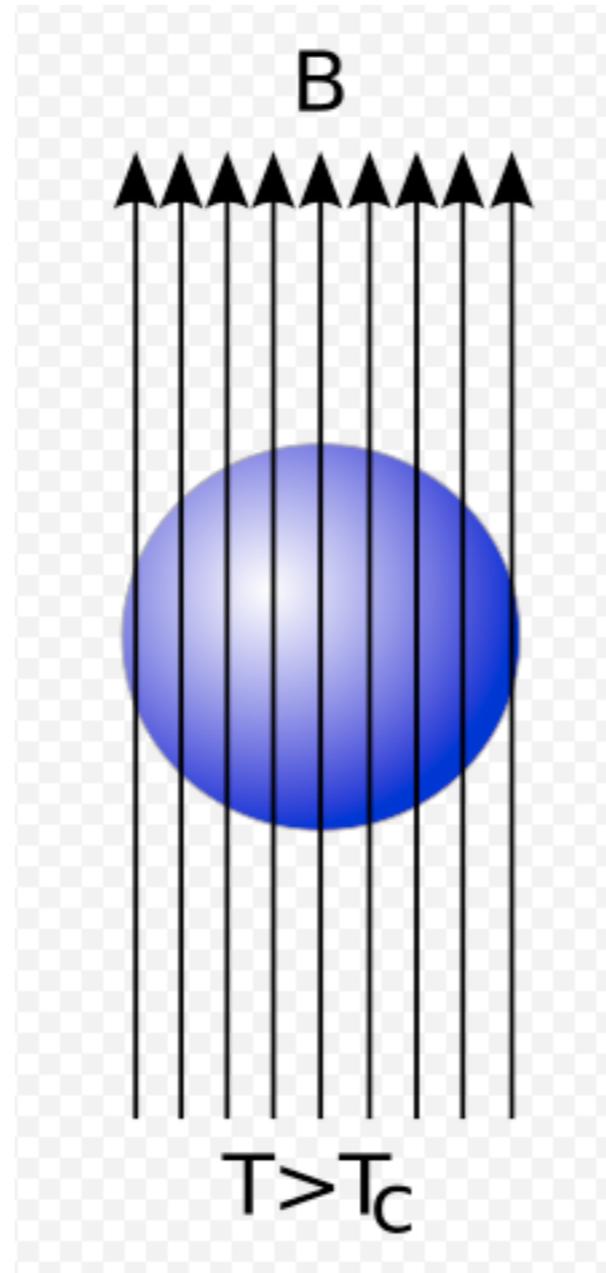
- ★ There is no ξ particle !
- ★ Photon becomes massive triplet \tilde{A}_μ
- ★ The Higgs boson (η) has mass, $m_H = \sqrt{2}m$, independent of e

Higgs mechanism in condensed matter

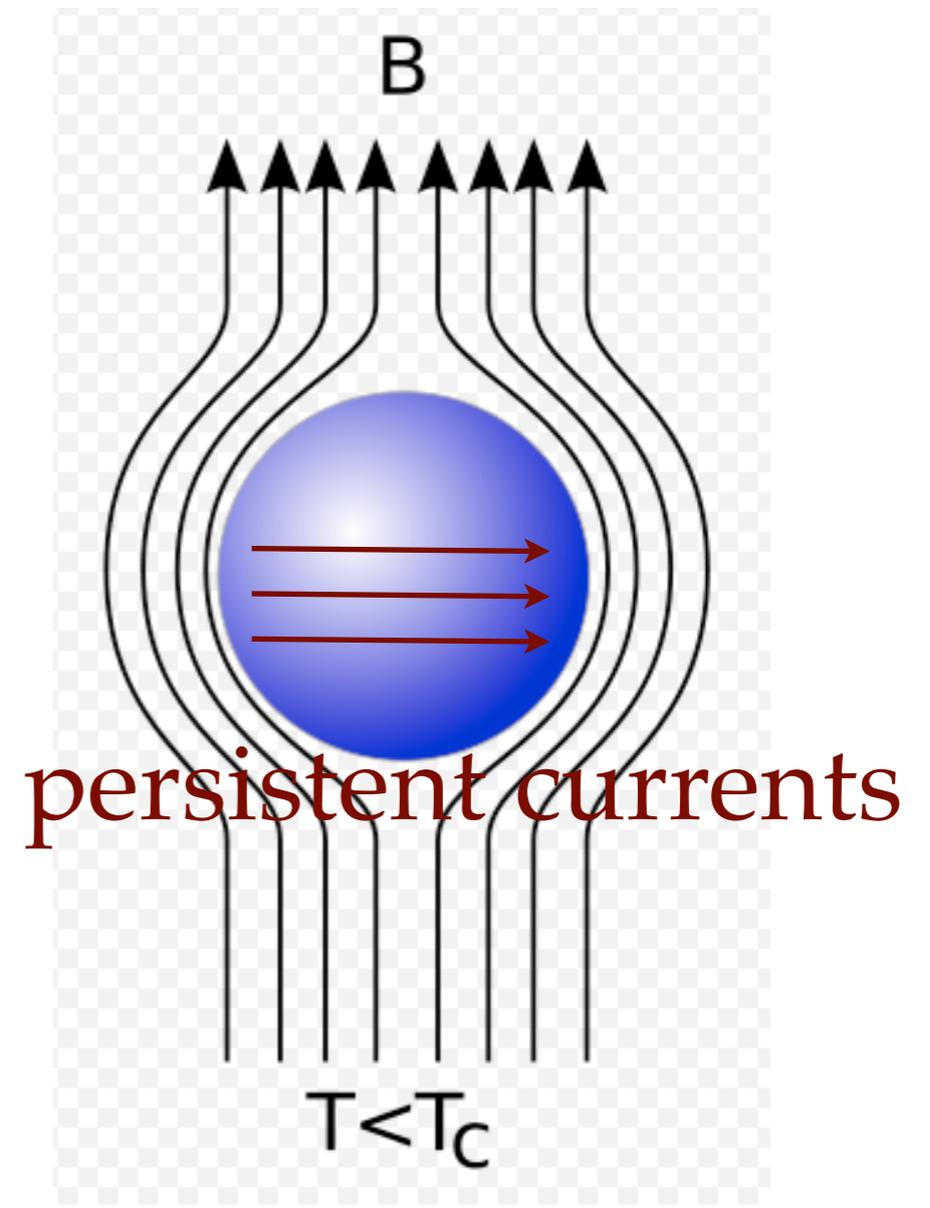


Meissner Effect, 1933

Metal



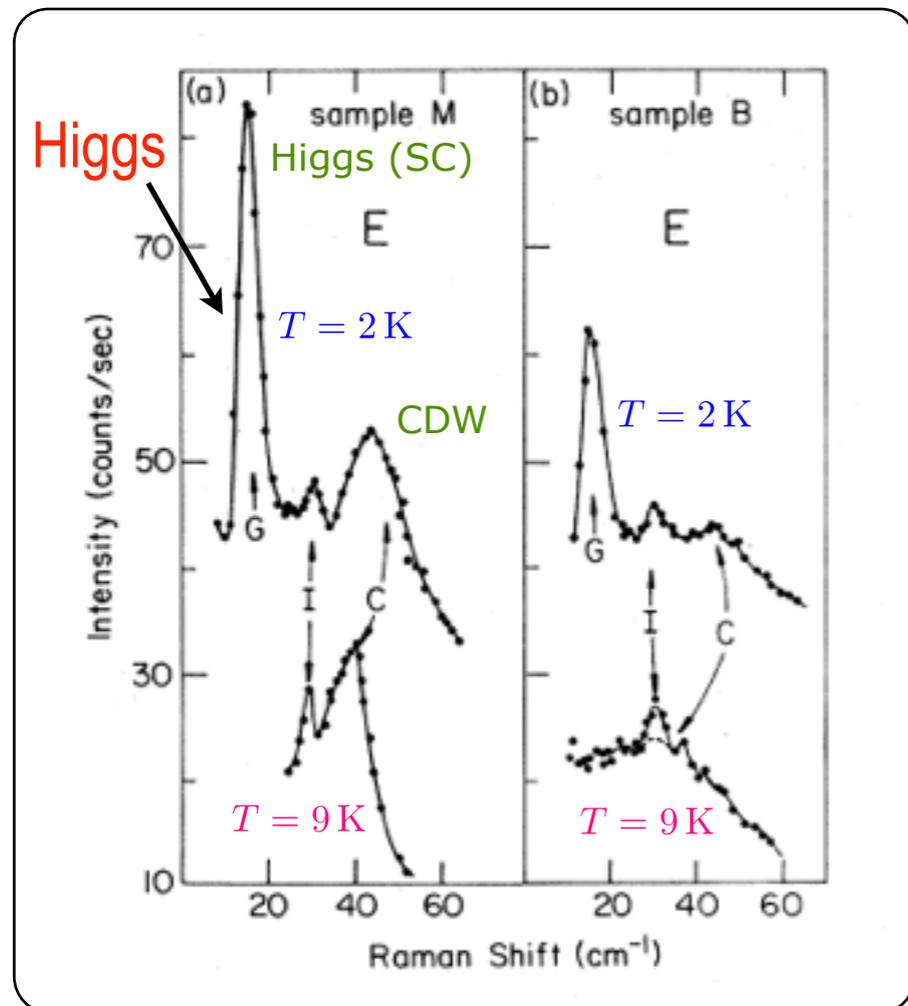
Superconductor



Higgs mechanism, but where is the "Higgs boson"?

Some experimental candidates

Superconducting $2H-NbSe_2$

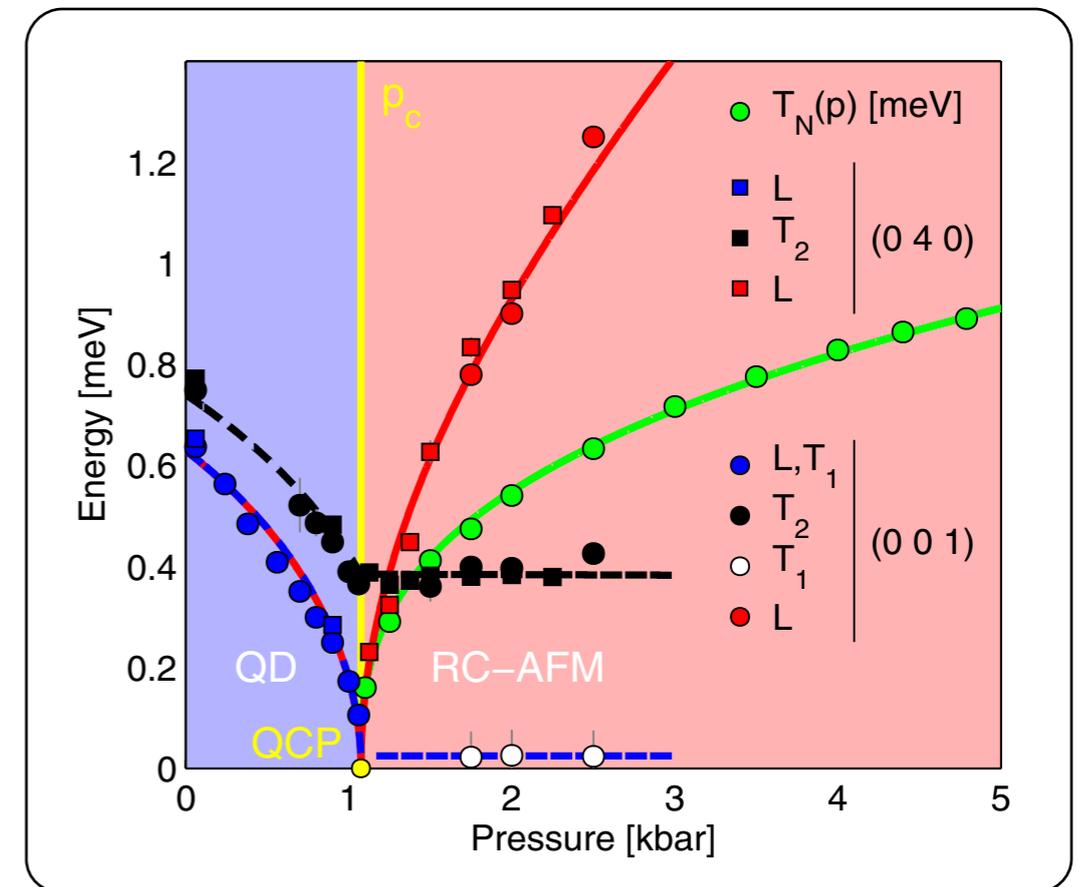


R. Sooryakumar et al., PRB (1981)

First observation of the Higgs mode?

P. B. Littlewood and C. M. Varma, PRB (1982)

Antiferromagnetic $TlCuCl_3$



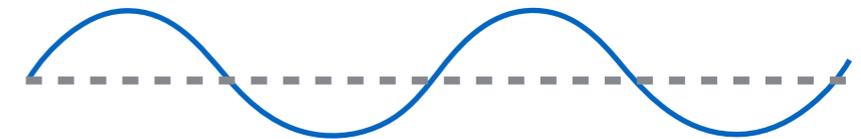
Ch. Rüegg et al., PRL (2008)

Higgs mode softens at quantum phase transition

Charge density waves

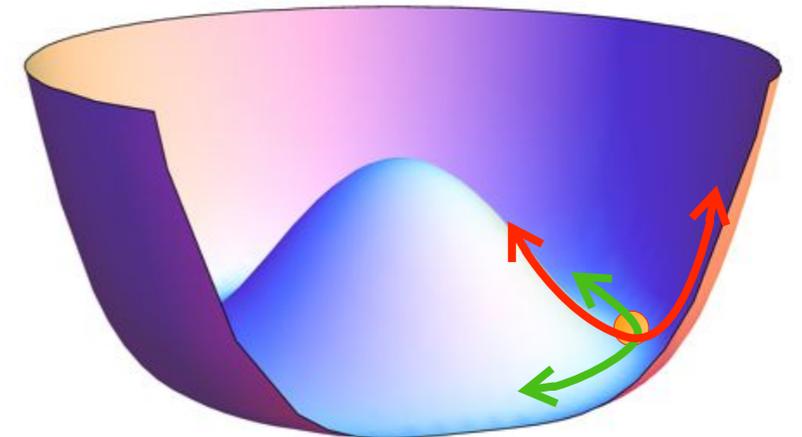
Density modulation:

$$\delta\rho(x) = \text{Re} [\psi e^{iQx}] = |\psi| \cos(Qx + \varphi)$$

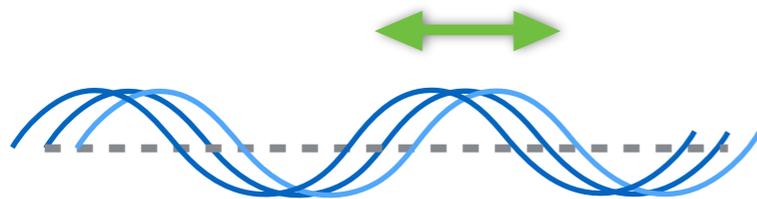


Complex order parameter:

$$\psi = |\psi| e^{i\varphi}$$

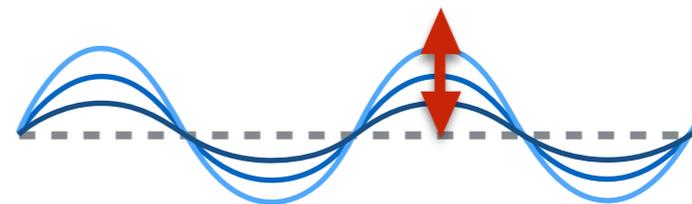


"phason"



$$\omega \sim cq$$

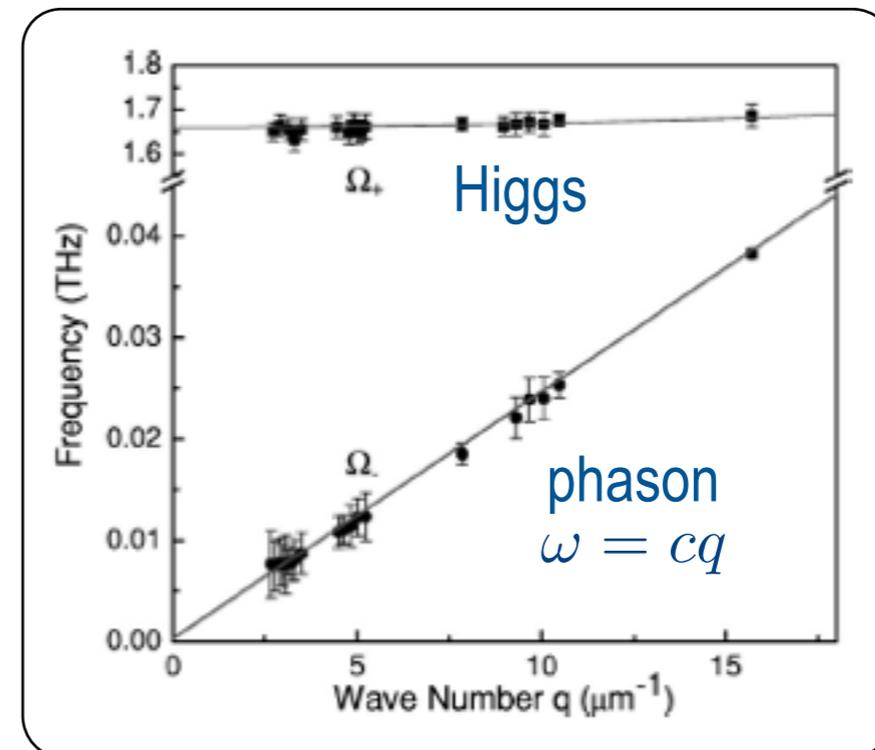
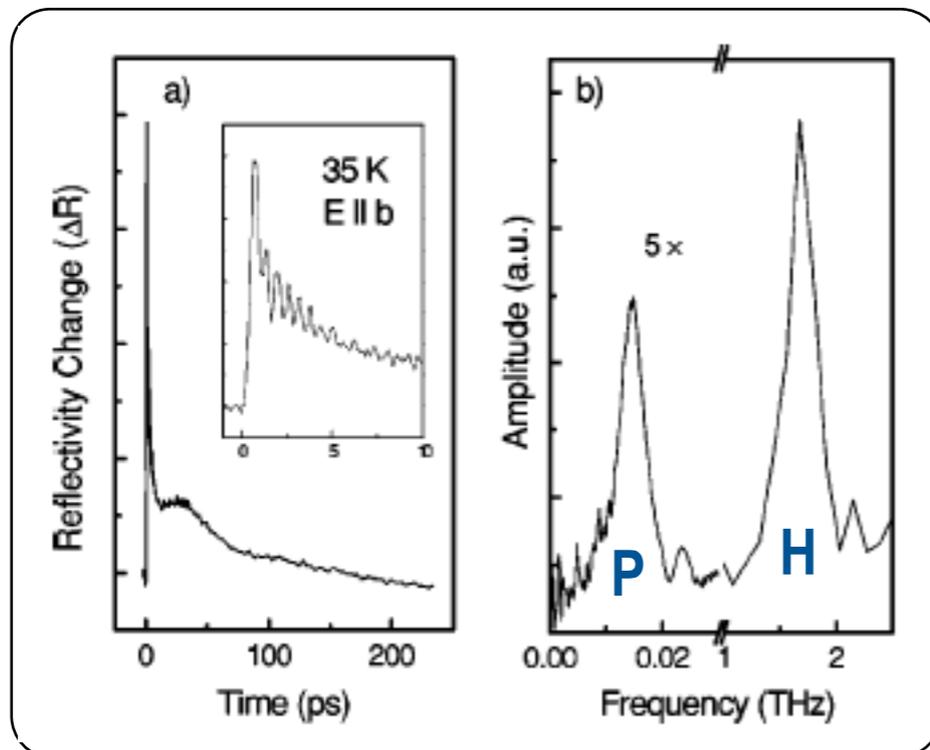
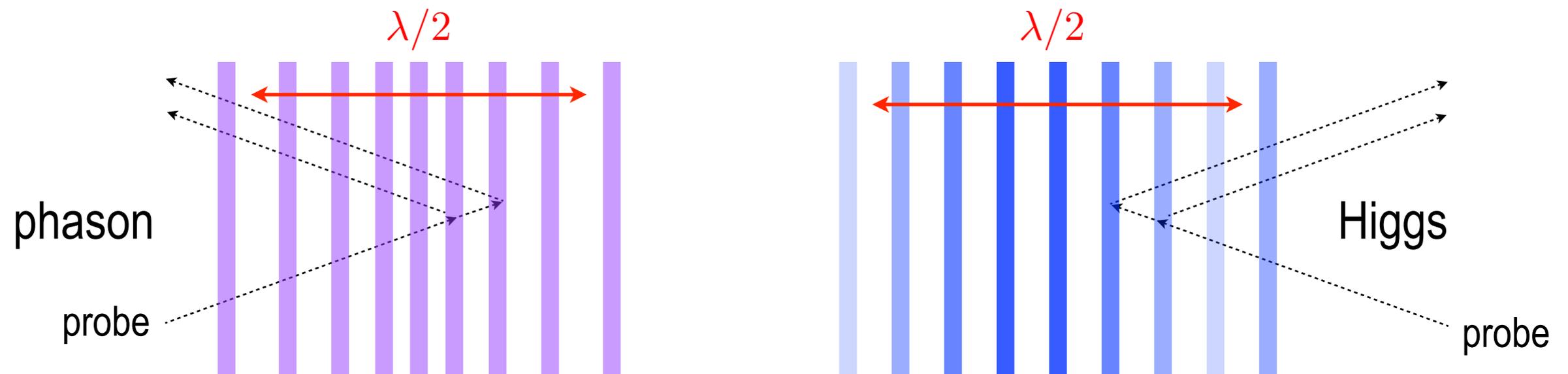
"amplitudon" (Higgs)



$$\omega \sim \sqrt{m^2 + c^2 q^2}$$

FS pump-probe spectroscopy

One-dimensional CDW conductor $K_{0.3}MoO_3$:

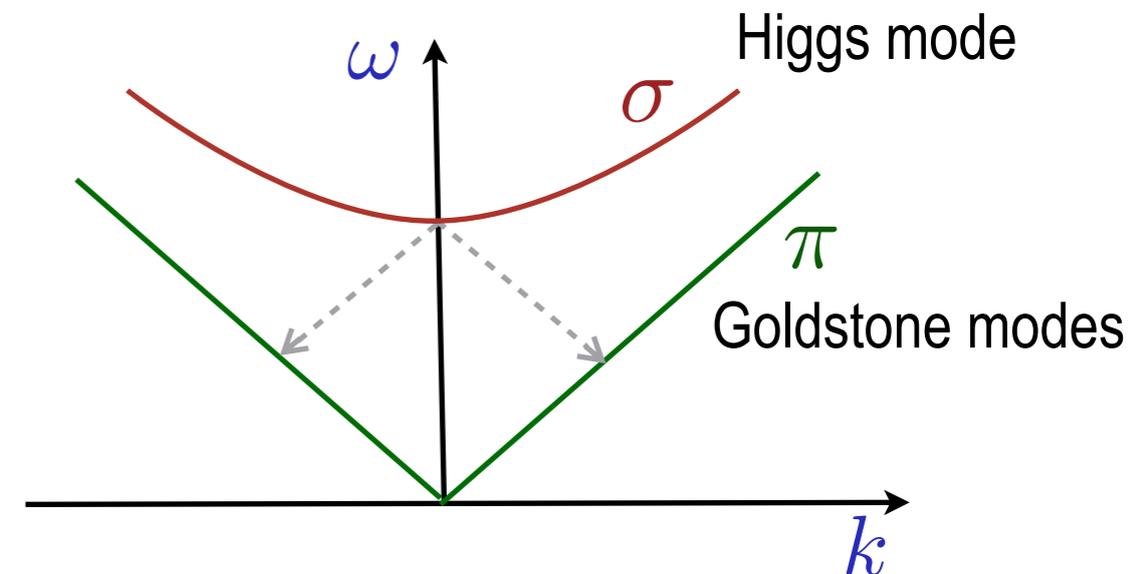


Y. Ren, Z. Xu, and G. Lüpke, J. Chem. Phys. 120, 4755 (2004)

**How to see the Higgs mode
in 2d?**

The Higgs decay

The Higgs mode can decay into a pair of Goldstone bosons:



d=3 Higgs decay rate is **finite**. As $g \rightarrow g_c$, Higgs mode becomes sharper and sharper
Affleck and Wellman (1992)

d=2 Longitudinal response **diverges** at low frequency, **even at weak coupling!**

(Nepomnyaschii)² (1978)
Sachdev (1999), Zwerger (2004)

Behavior of different response functions

longitudinal susceptibility

$$\chi_{\text{long}}(\omega) = \langle \phi_1(\omega) \phi_1(-\omega) \rangle \sim \omega^{-1}$$

infrared divergent in d=2

(Nepomnyaschii)² (1978)

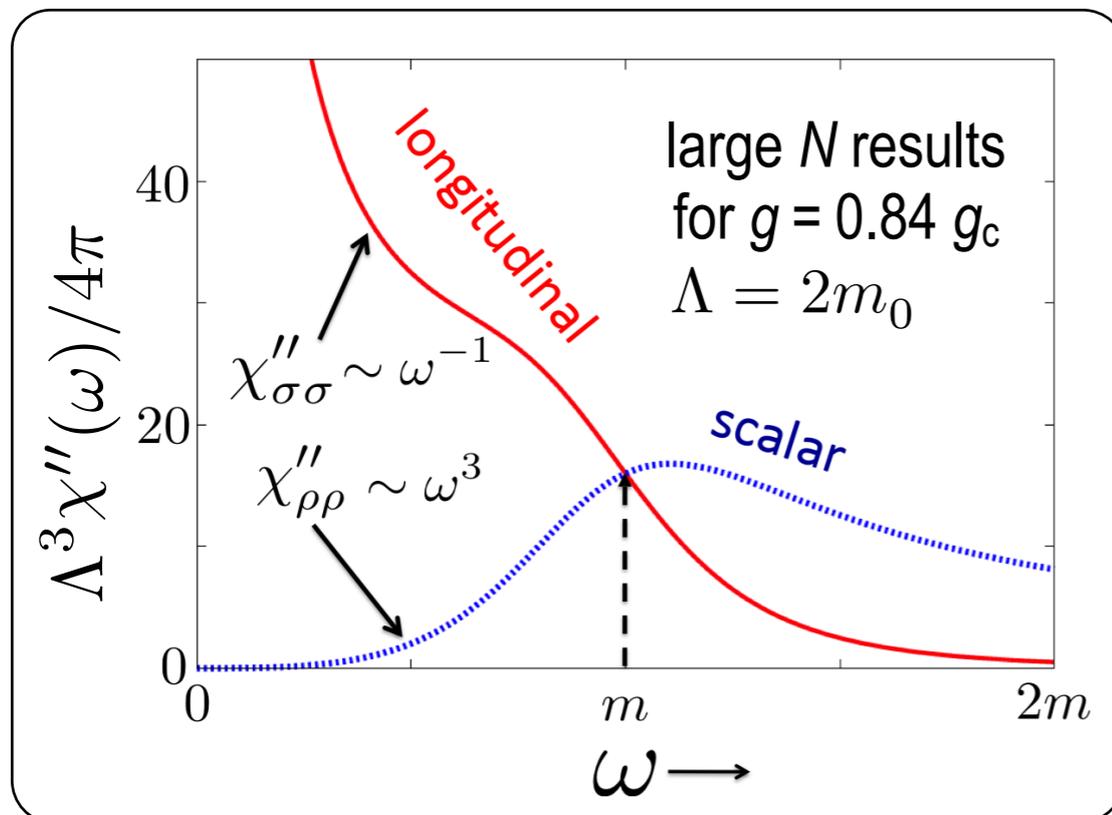
Sachdev (1999), Zwirger (2004)

scalar susceptibility

$$\chi_{\text{scalar}}(\omega) = \langle |\vec{\phi}|^2(\omega) |\vec{\phi}|^2(-\omega) \rangle \sim \omega^3$$

infrared regular in d=2

Podolsky, Auerbach and Arovas, PRB (2011)

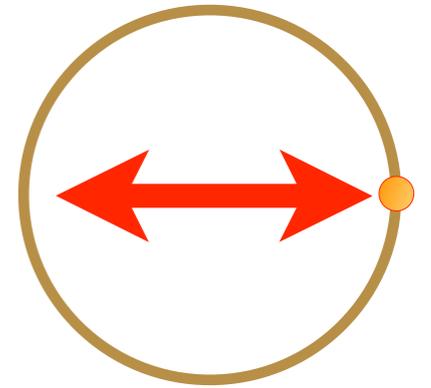


Longitudinal versus scalar measurements

Longitudinal: couples to order parameter as a vector

$$\mathcal{H}_{\text{probe}} = \vec{h}_{\text{ext}} \cdot \vec{\phi}$$

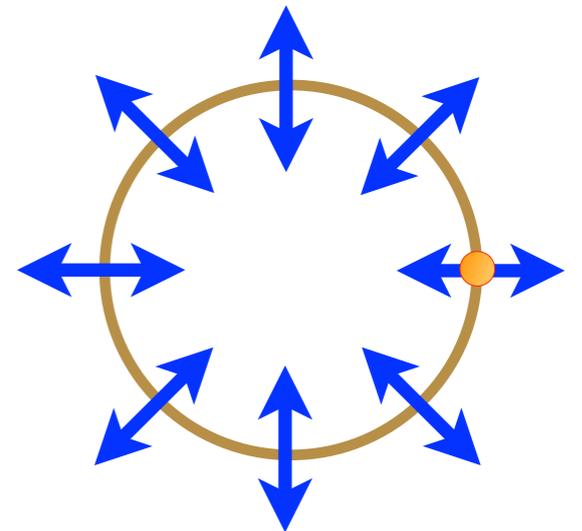
Example: neutron scattering in an antiferromagnet.



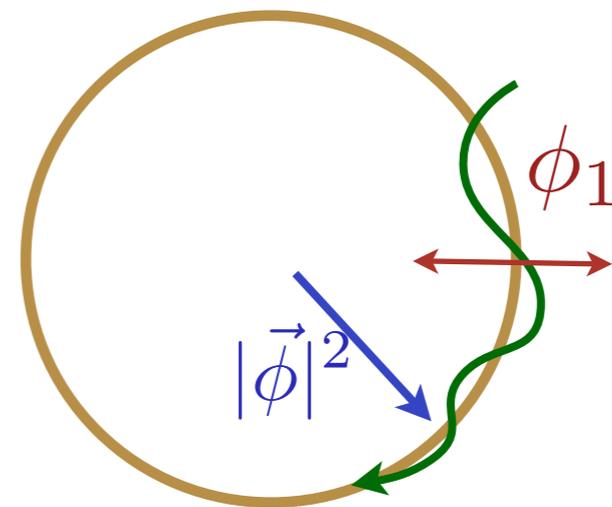
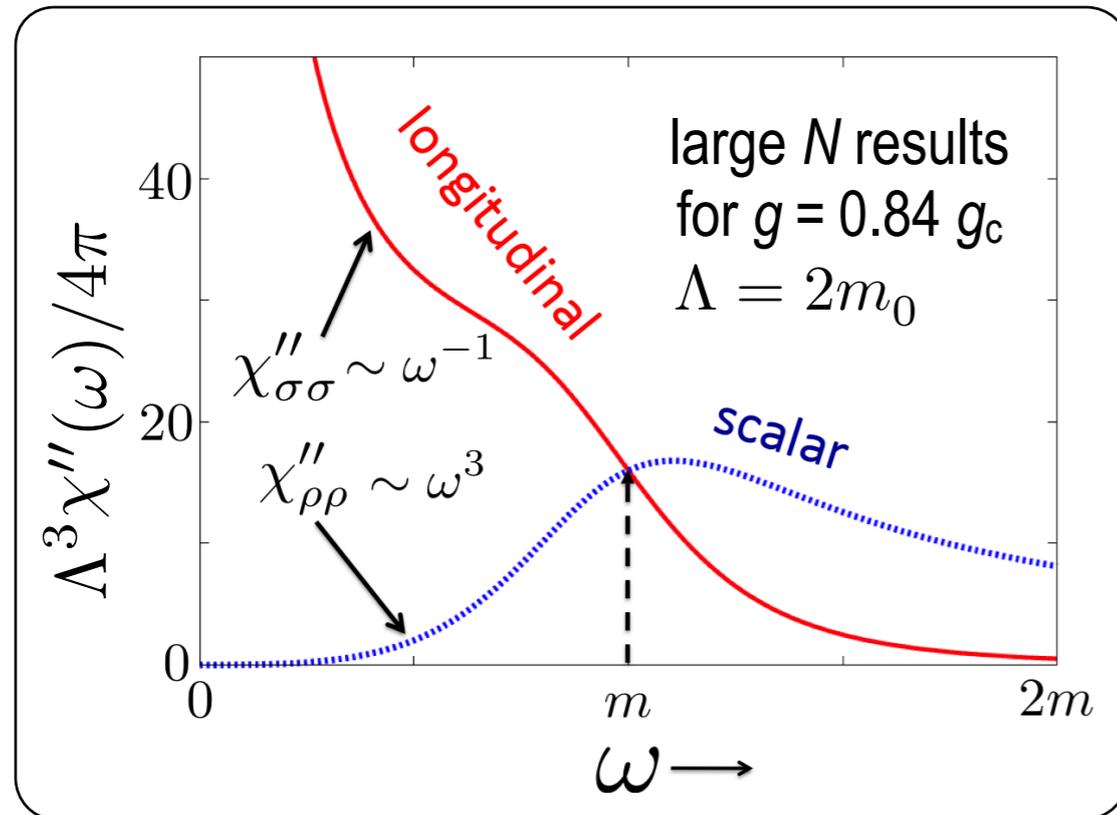
Scalar: couples to the magnitude of the order parameter

$$\mathcal{H}_{\text{probe}} = u_{\text{ext}} |\vec{\phi}|^2$$

Example: lattice depth modulation of bosons



Why is the **scalar** response sharper?



Radial motion is less damped since it is not effected by azimuthal meandering.

Bosons in an optical lattice

Bose-Hubbard model $H = -t \sum_{\langle ij \rangle} b_i^\dagger b_j + U \sum_i n_i^2 - \mu \sum_i n_i$

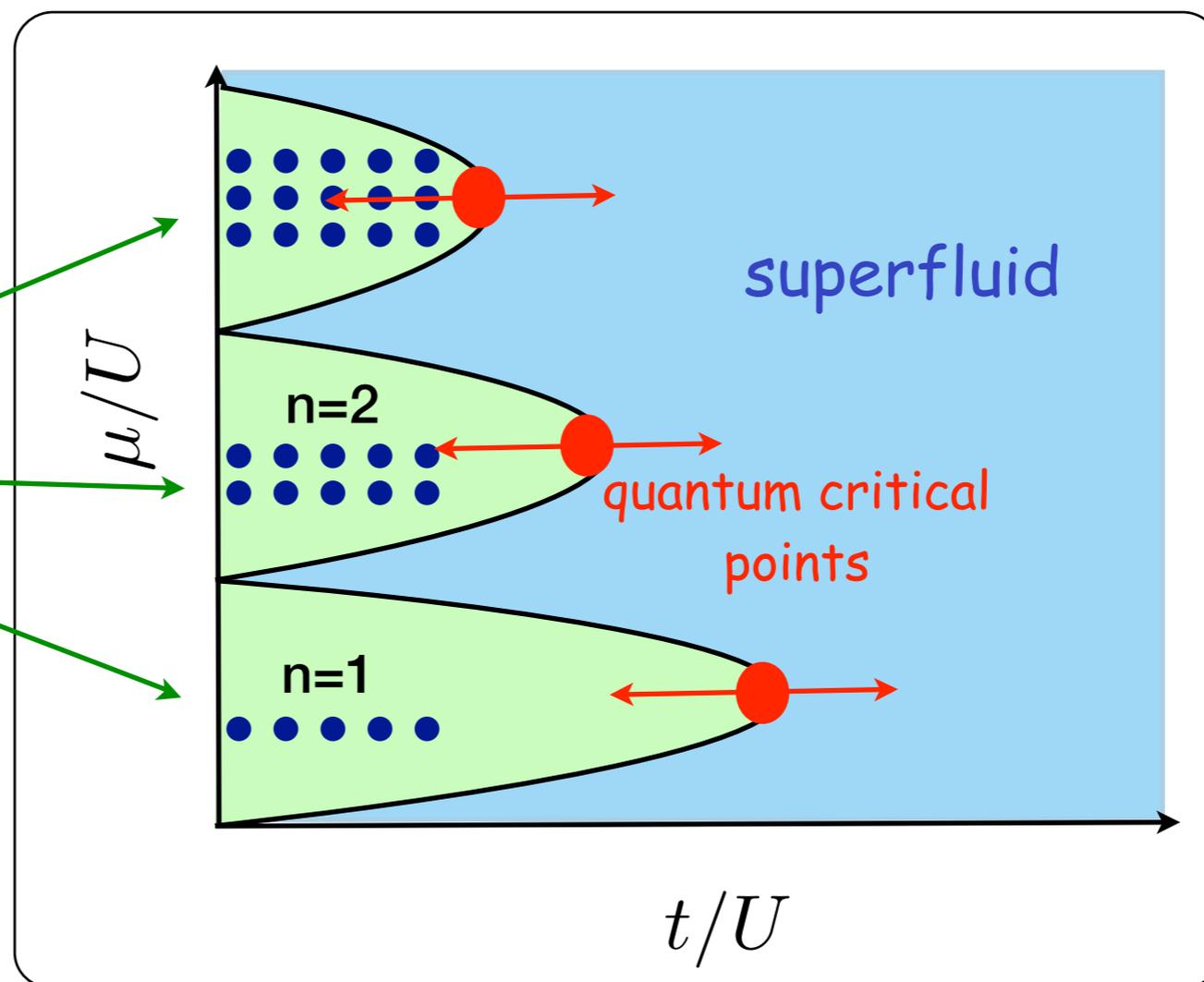
$t \gg U$: **superfluid** (Bose condensate)

$U \gg t$: **Mott insulator** (gapped charge fluctuations)

Mott insulator
incompressible

$$\langle b \rangle = 0$$

$$\frac{\partial \langle n_i \rangle}{\partial \mu} = 0$$



Bosons in an optical lattice

Bose-Hubbard model
$$H = -t \sum_{\langle ij \rangle} b_i^\dagger b_j + U \sum_i n_i^2 - \mu \sum_i n_i$$

Dynamics:

1) **Far from Mott lobe**, Gross-Pitaevskii model,

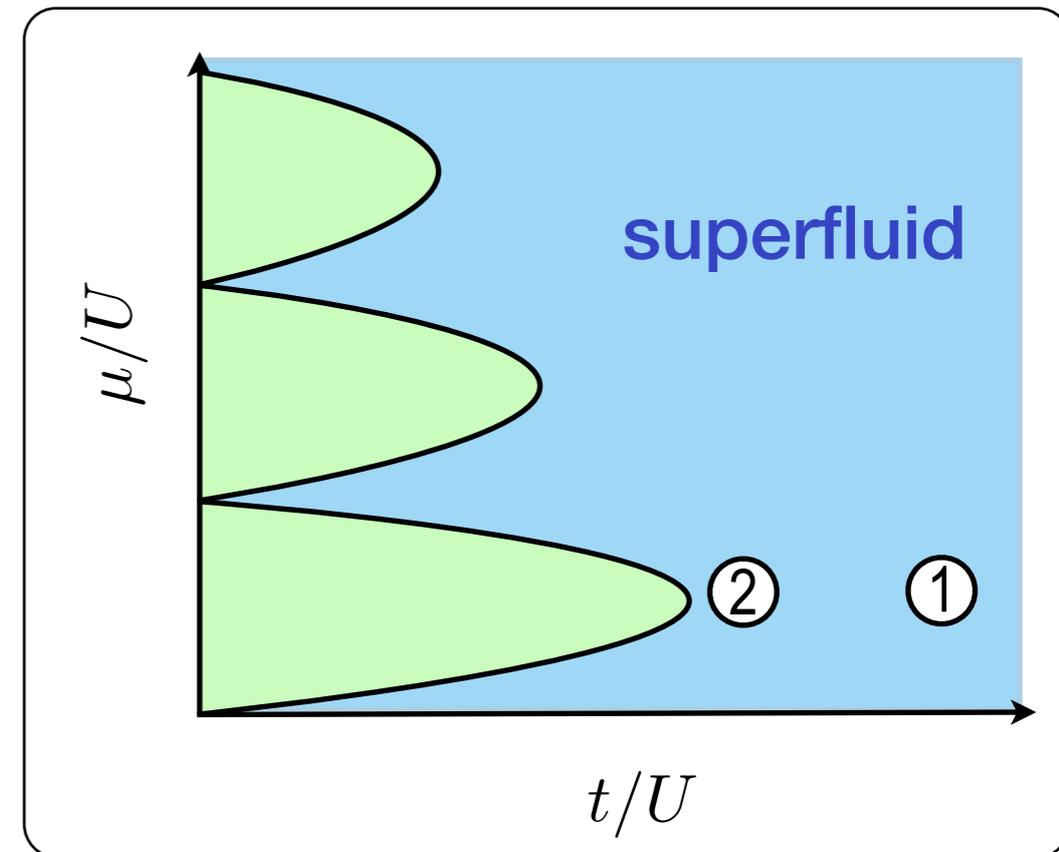
$$\mathcal{L} = -i\psi^* \partial_t \psi - \frac{\hbar^2}{2m^*} |\nabla \psi|^2 + \mu |\psi|^2 - g |\psi|^4$$

Gapless Goldstone mode, but no Higgs.

Varma, J. Low Temp. Phys. (2002)

Huber et al, PRB (2008)

Huber and Lindner, PNAS (2011)



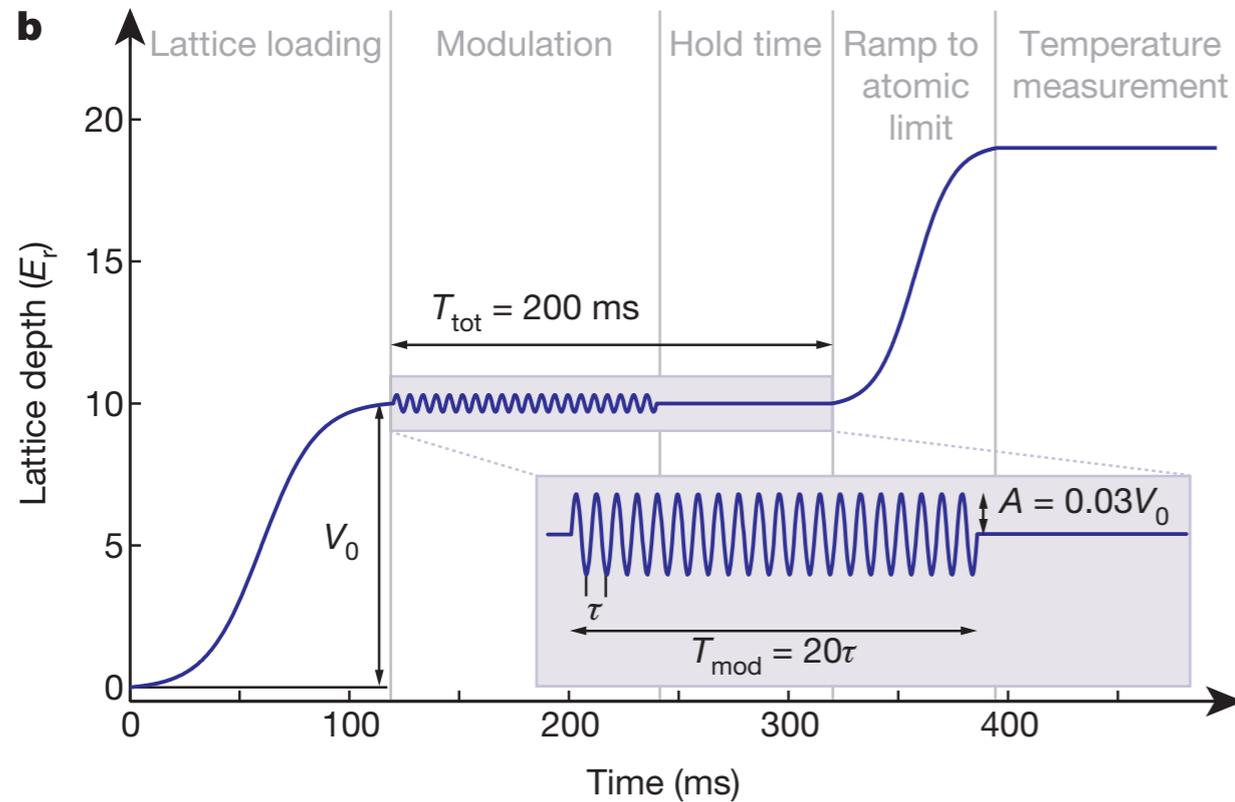
2) **Close to Mott lobe**, relativistic model,

$$\mathcal{L} = |\partial_t \psi|^2 - c^2 |\nabla \psi|^2 + r |\psi|^2 - u |\psi|^4$$

Goldstone and Higgs.

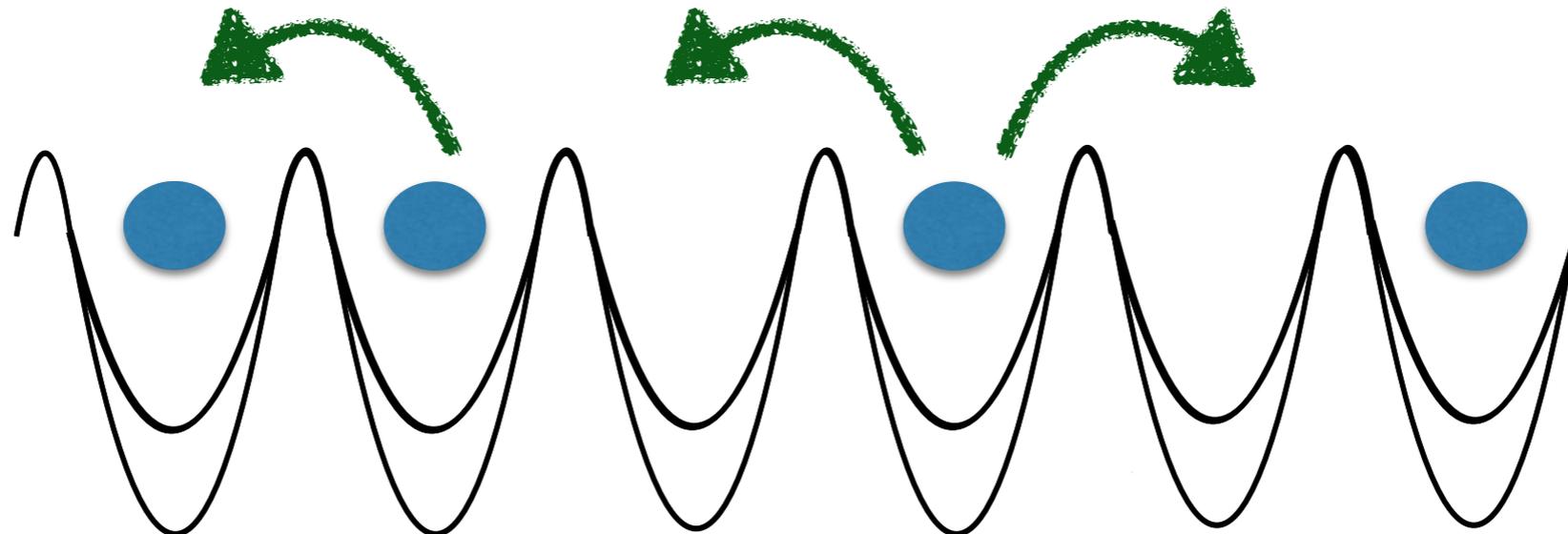
The ‘Higgs’ amplitude mode at the two-dimensional superfluid/Mott insulator transition

Manuel Endres¹, Takeshi Fukuhara¹, David Pekker², Marc Cheneau¹, Peter Schauß¹, Christian Gross¹, Eugene Demler³, Stefan Kuhr^{1,4} & Immanuel Bloch^{1,5}



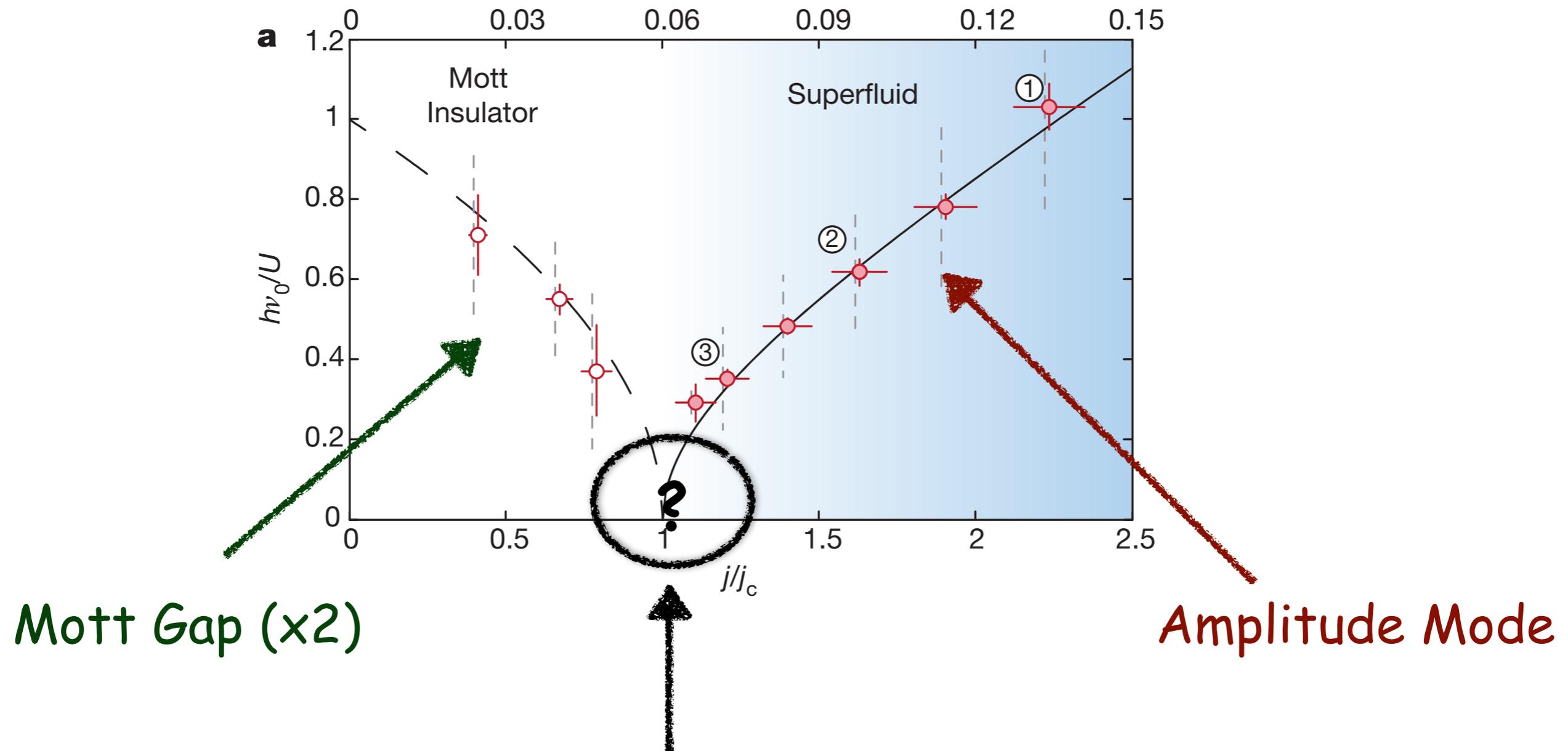
Energy absorption rate of periodically modulated lattice

$$\propto \omega \chi''_{\text{scalar}}(\omega)$$



The 'Higgs' amplitude mode at the two-dimensional superfluid/Mott insulator transition

Manuel Endres¹, Takeshi Fukuhara¹, David Pekker², Marc Cheneau¹, Peter Schauß¹, Christian Gross¹, Eugene Demler³, Stefan Kuhr^{1,4} & Immanuel Bloch^{1,5}



What happens near the quantum critical point???

**Does the Higgs survive near
quantum criticality?**

Scaling near criticality

gap: $\Delta \sim |g - g_c|^\nu$ $\nu = 0.6717(1)$ ($N = 2$)

$$\chi_{scalar}(\omega) = \Delta^{3-2/\nu} \Phi_s \left(\frac{\omega}{\Delta} \right) + \dots$$

universal function



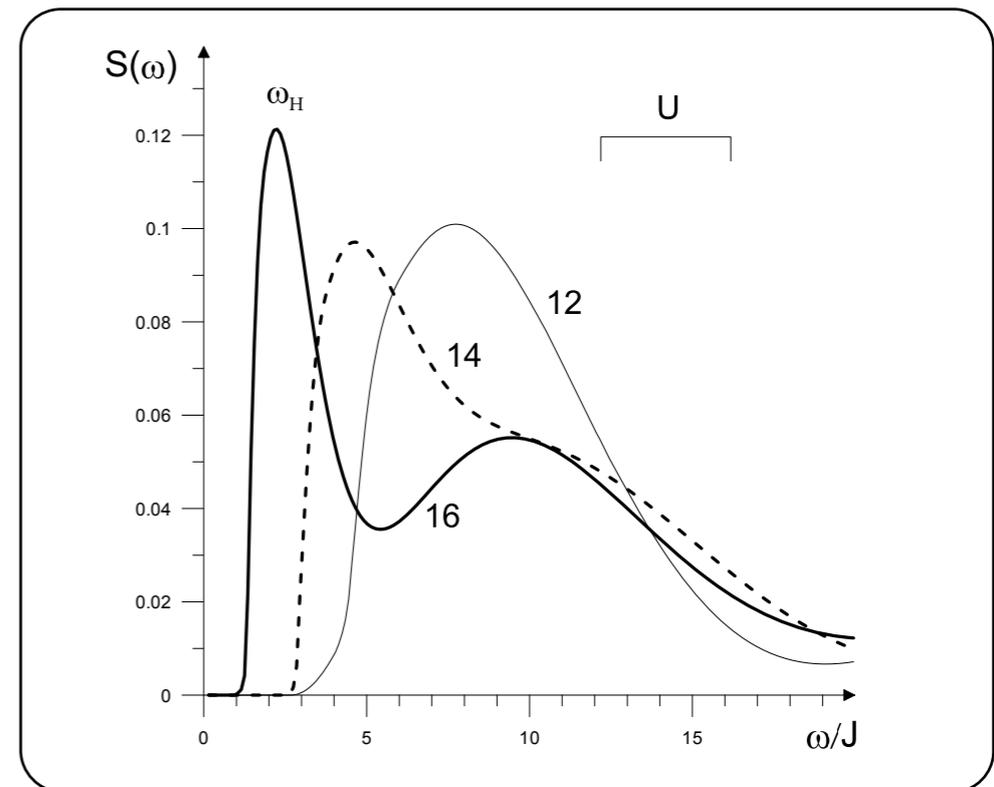
Podolsky and Sachdev, PRB (2012)

Does it have a peak?

First indications

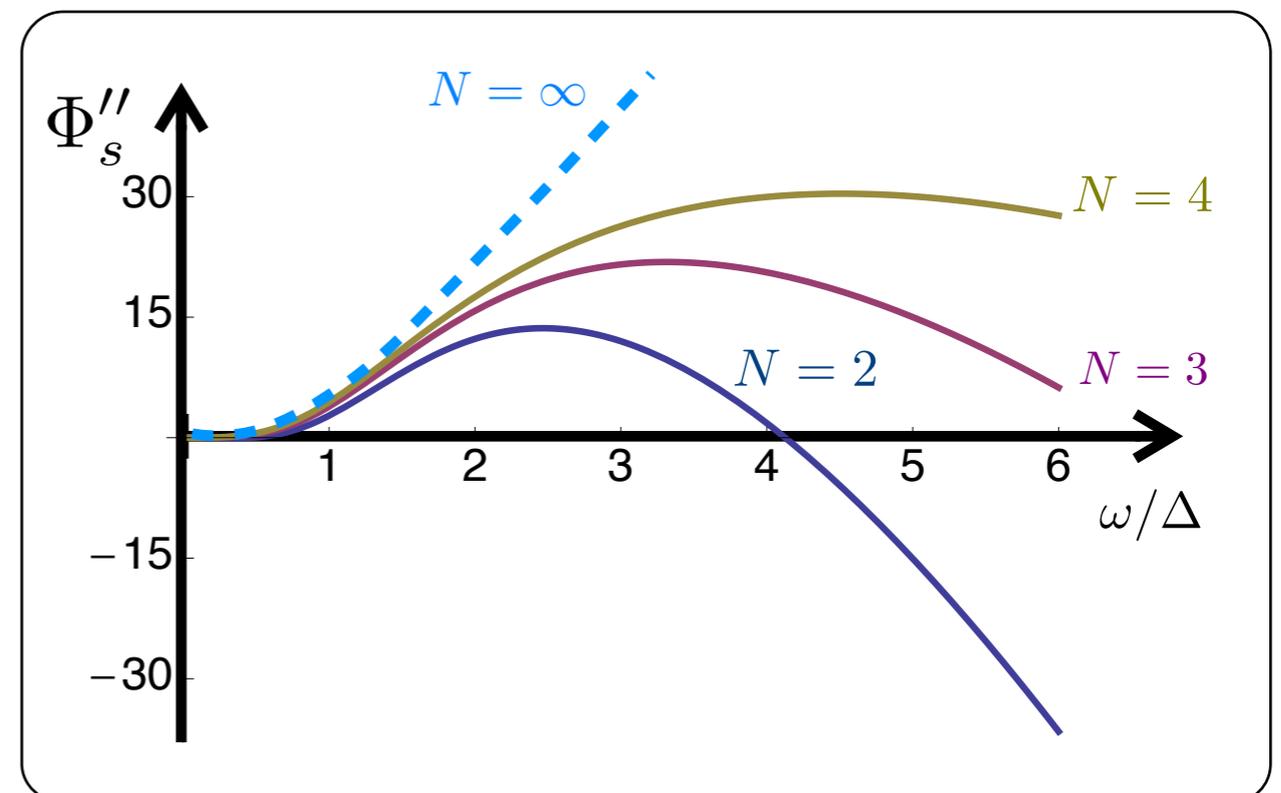
Numerics on Bose-Hubbard model

L. Pollet and N. Prokof'ev, PRL (2012)



Scaling function to $O(1/N)$

Podolsky and Sachdev, PRB (2012)

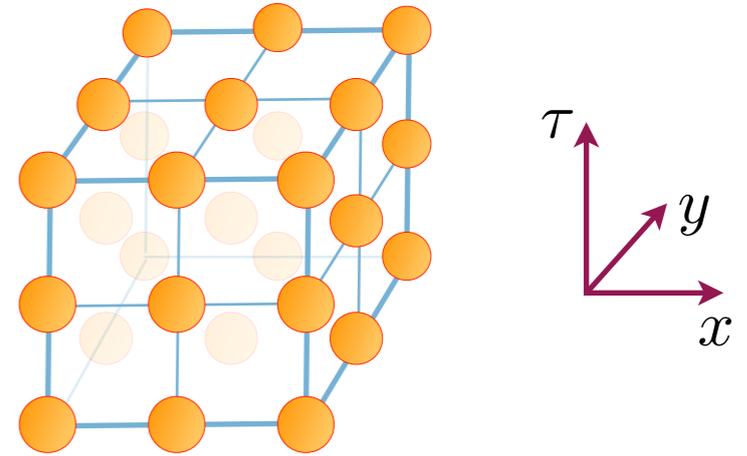


Monte Carlo Simulations

Discrete model:

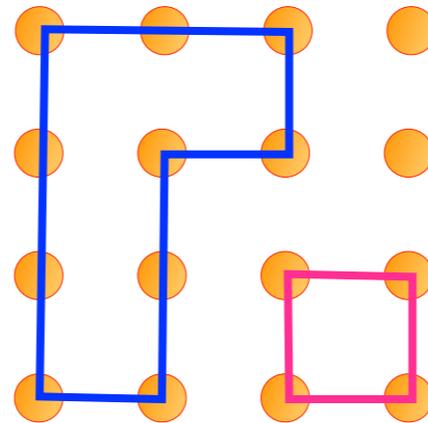
$$\mathcal{Z} = \int \mathcal{D}\vec{\phi} e^{-S[\vec{\phi}]}$$

$$S = - \sum_{\langle ij \rangle} \vec{\phi}_i \cdot \vec{\phi}_j + \mu \sum_i |\vec{\phi}_i|^2 + g \sum_i |\vec{\phi}_i|^4$$



Worm algorithm:

Dual loop model
with N flavors:

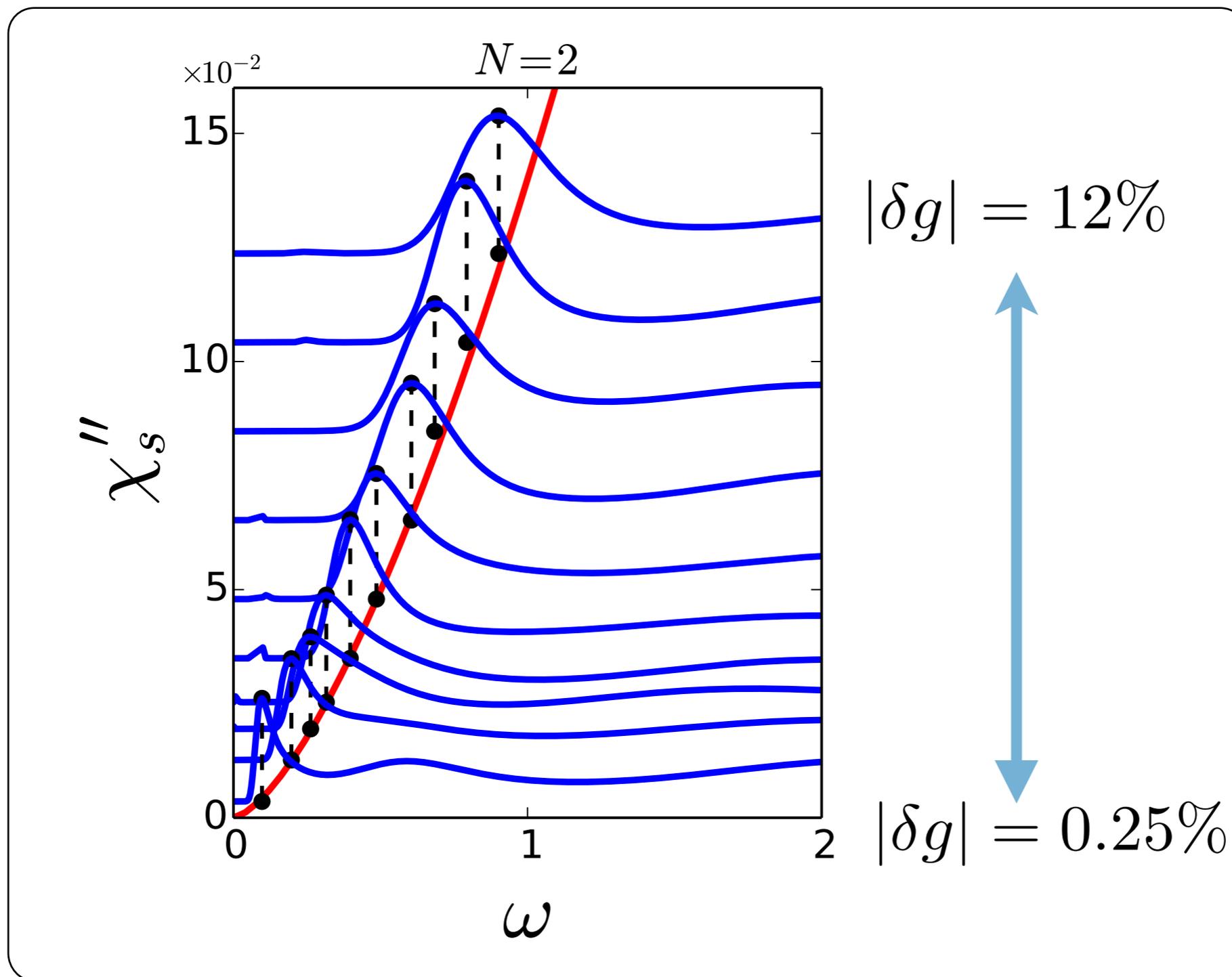


System size: $1 \ll \xi \ll L$ ($1 \ll 30 \ll 200$)

Numerical analytical continuation from Matsubara to real frequencies

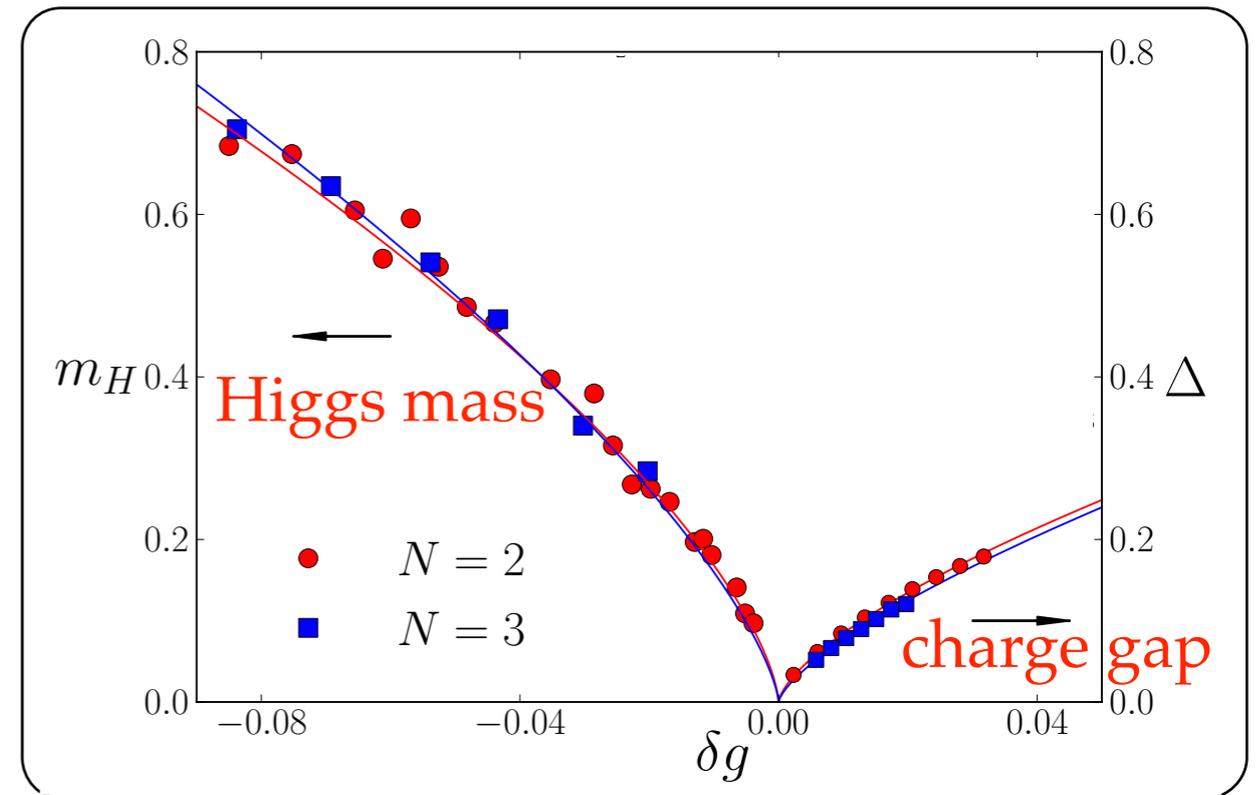
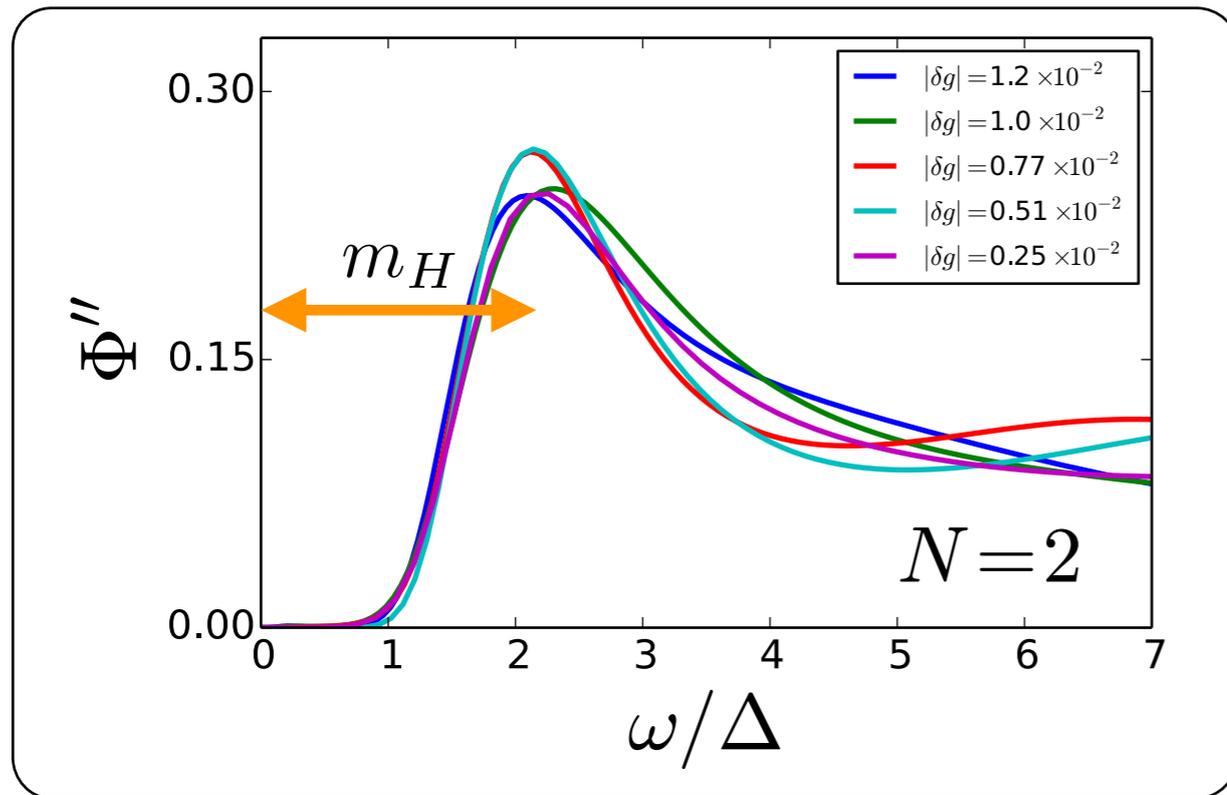
Tracking the Higgs peak

Scalar susceptibility in ordered phase: $\chi_s(\omega) = \langle \phi^2(\omega)\phi^2(-\omega) \rangle$



$$m_H = B|\delta g|^\nu$$

Spectral function at the QCP



$$\frac{m_H}{\Delta} = 2.1(3)$$

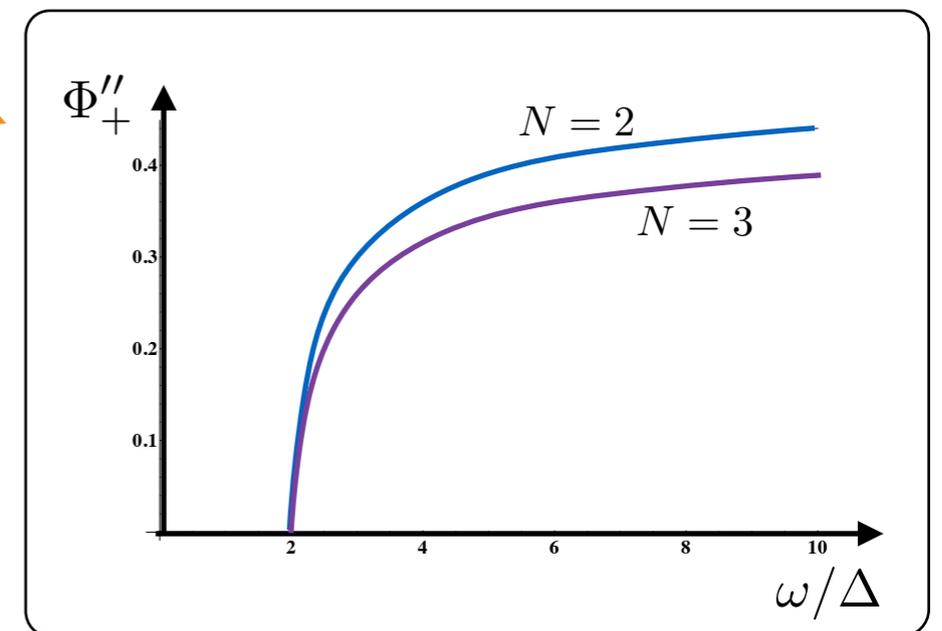
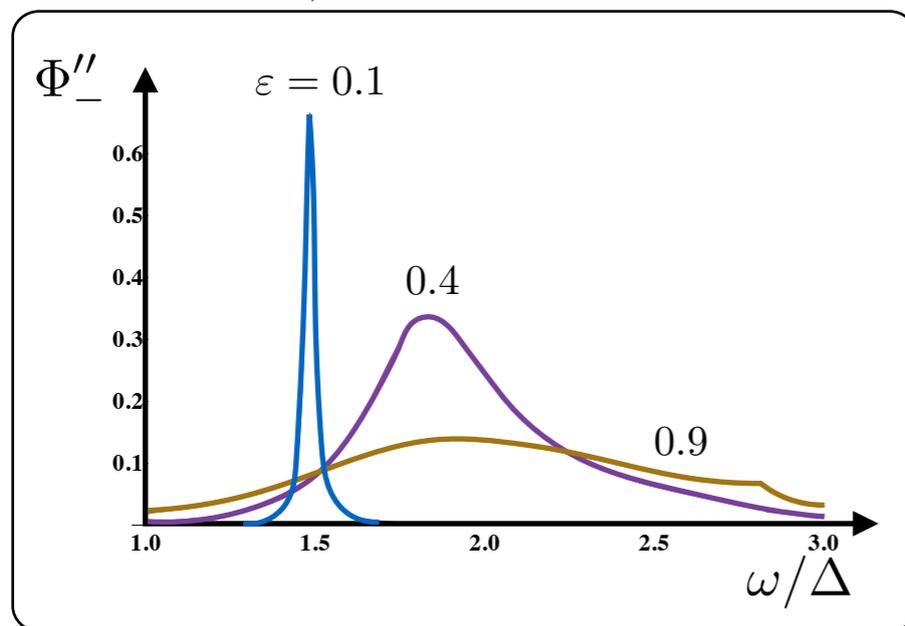
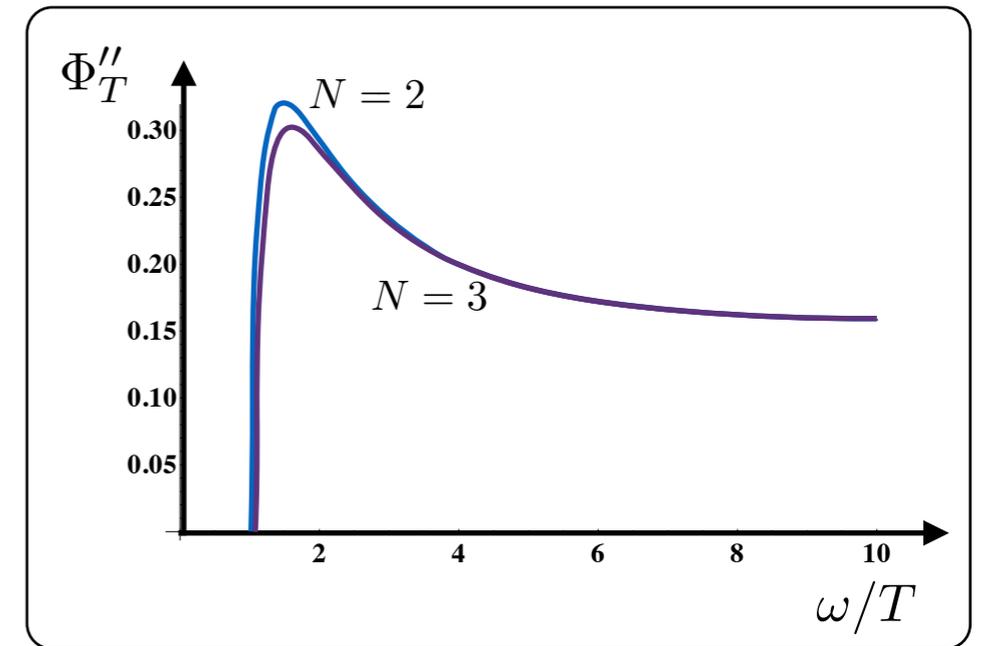
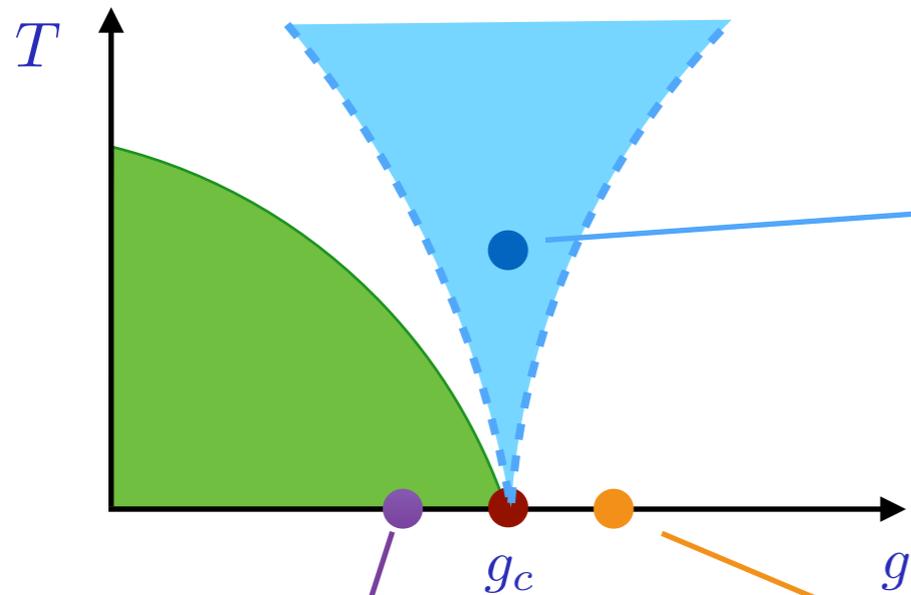
Mean field: $\frac{m_H}{\Delta} = \sqrt{2}$

Conclusion: Higgs resonance **survives** close to criticality in $d=2$

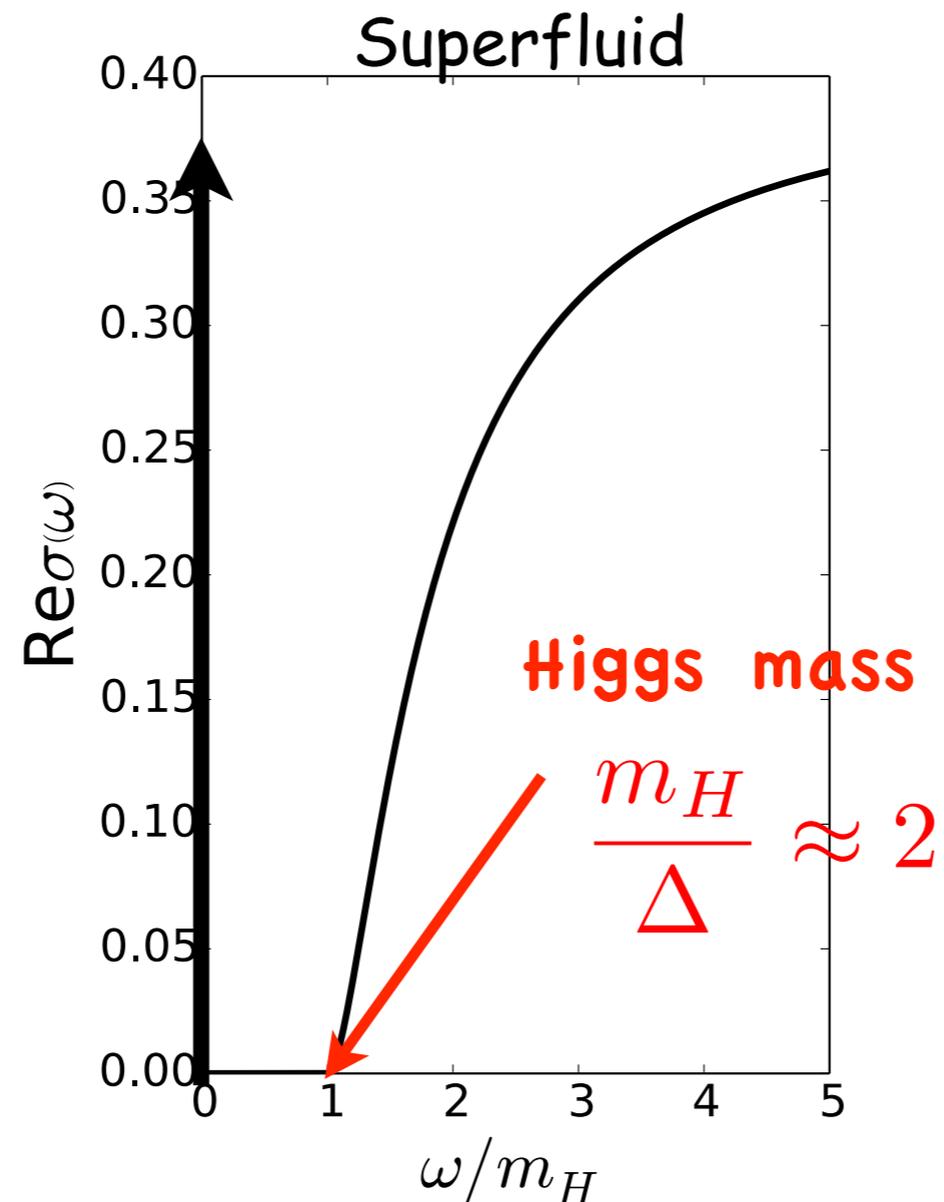
Chen et al, Bose-Hubbard Model (2013): $\frac{m_H}{\Delta} = 3.3(8)$

Rancon and Dupuis (2014): $\frac{m_H}{\Delta} = 2.4$

Scaling functions in $D = 4 - \varepsilon$ dimensions



Higgs in optical conductivity



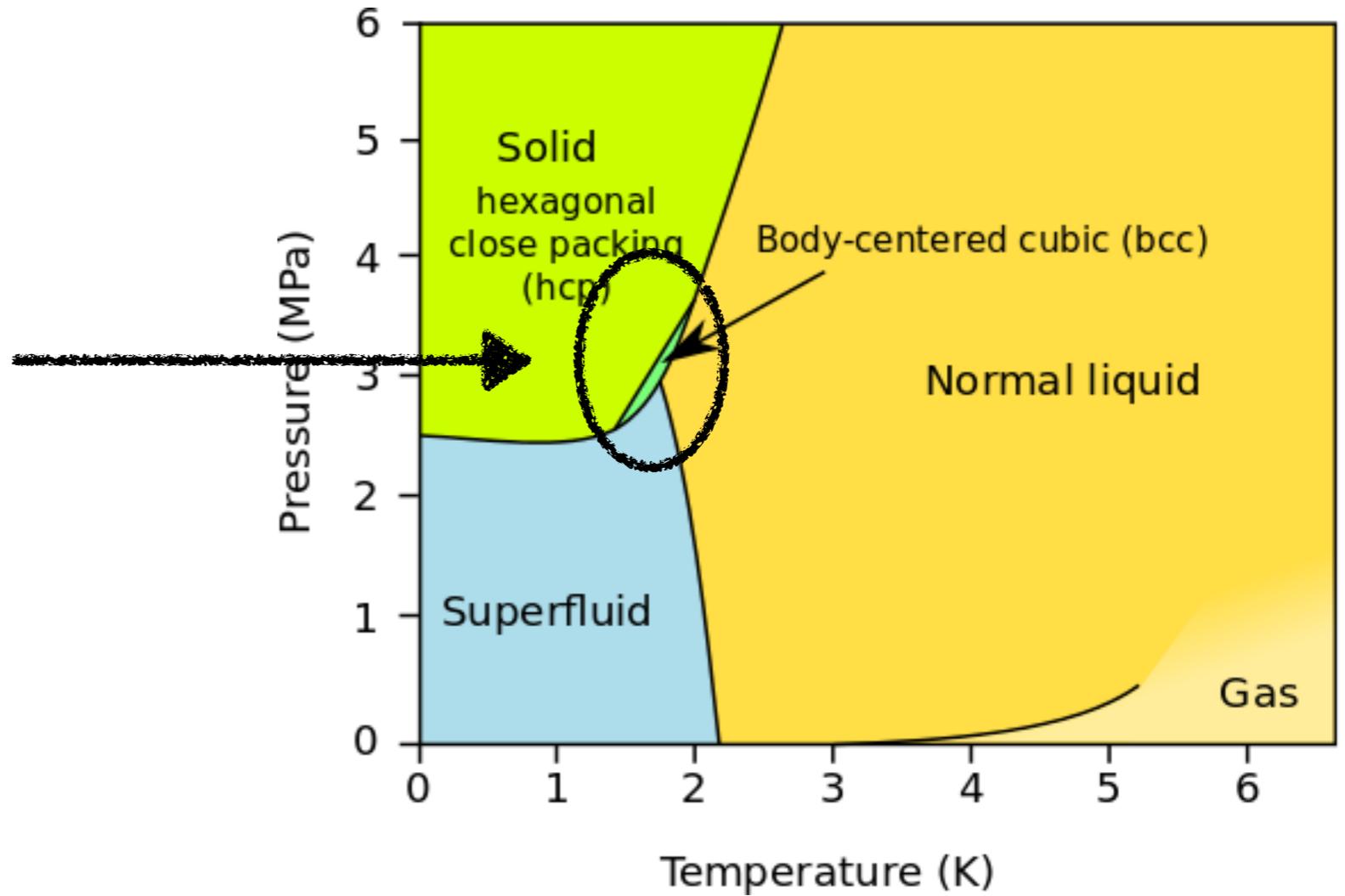
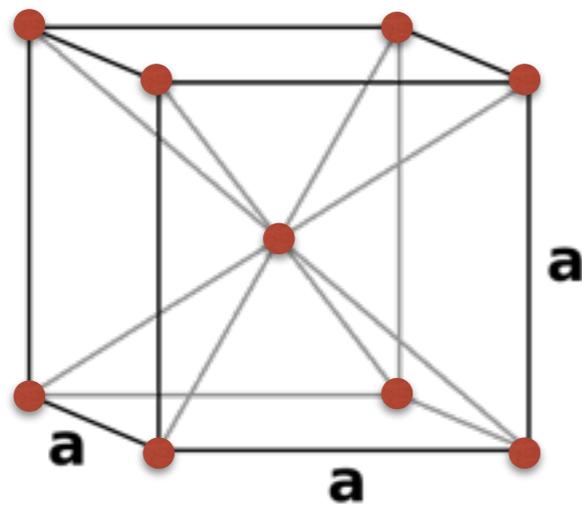
$$\sigma_{\text{SF}}(\omega) = 2\pi\sigma_Q \left(\frac{\omega^2 - m_H^2}{4\omega^2} \right)^2 \Theta(\omega - m_H)$$

Lindner, Auerbach (2010)

Podolsky, Auerbach, Arovas (2011)

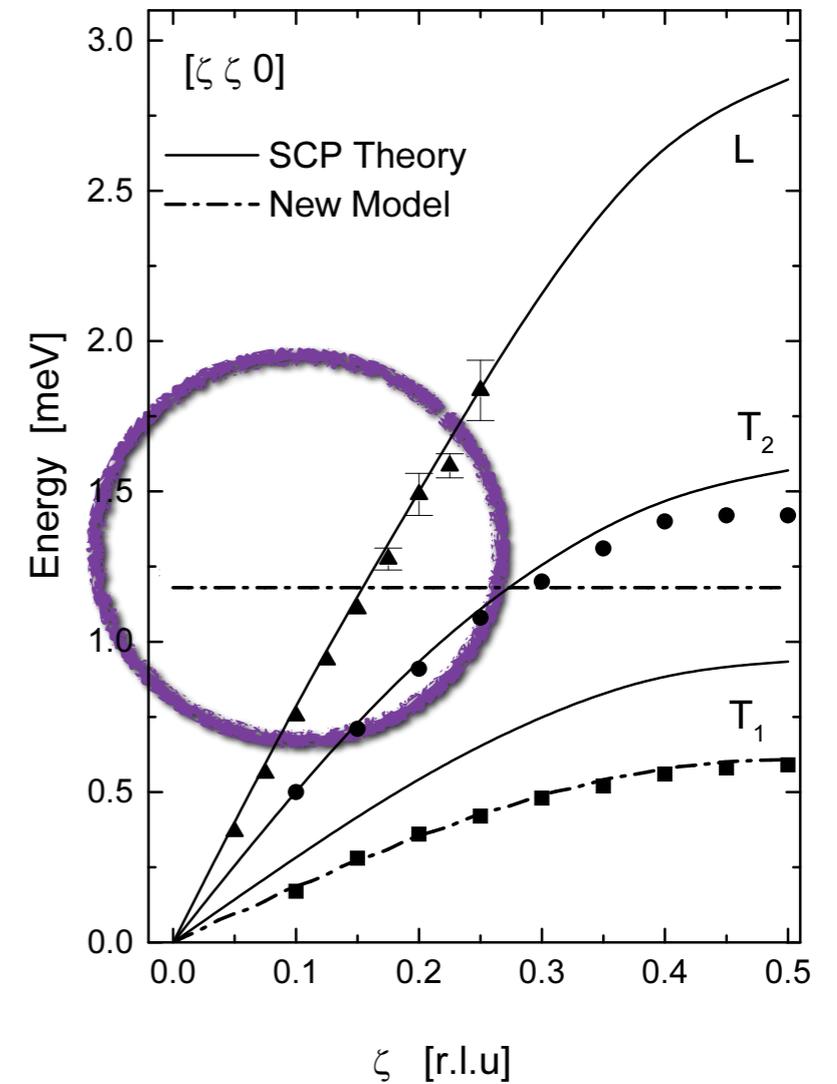
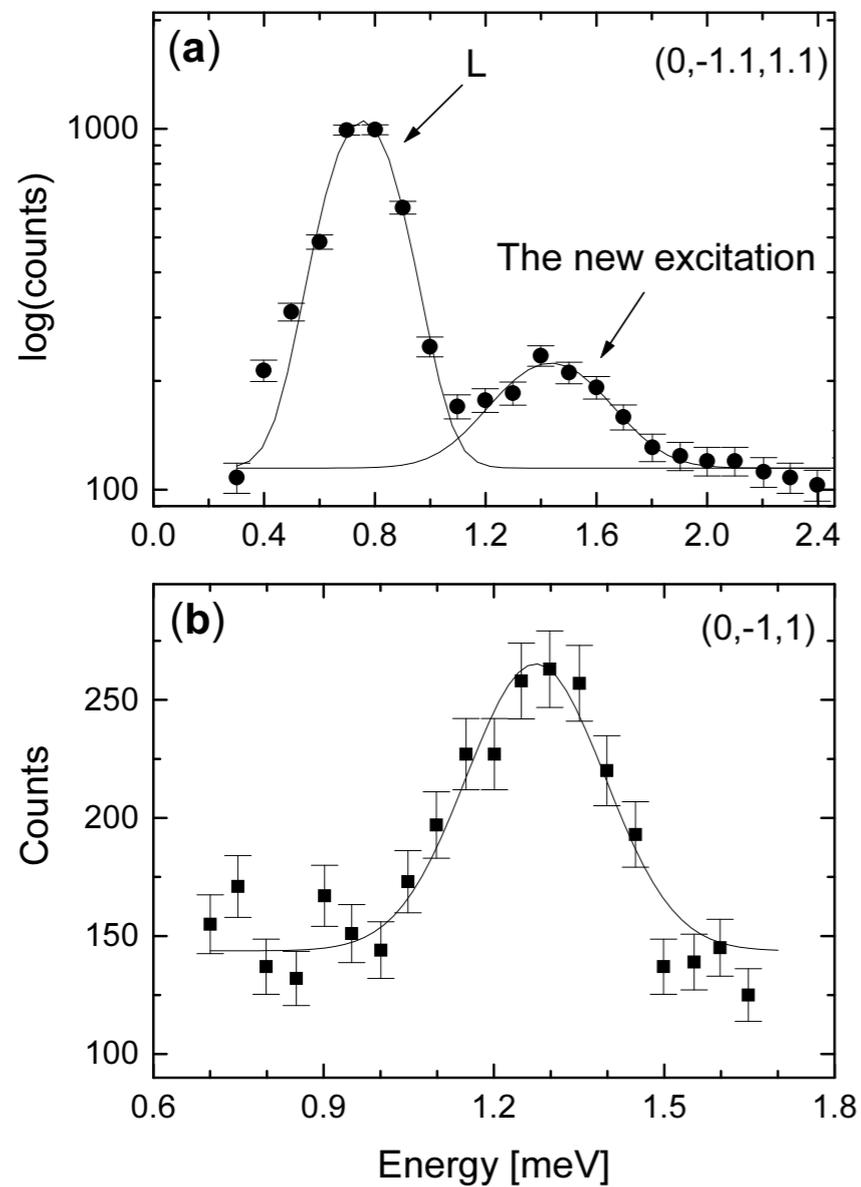
Gapped modes of a quantum solid

Helium 4 - Phase diagram



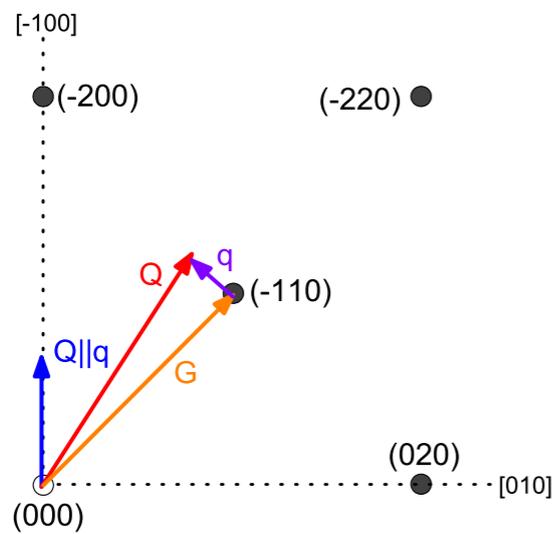
Inelastic neutron scattering

Optical mode observed!

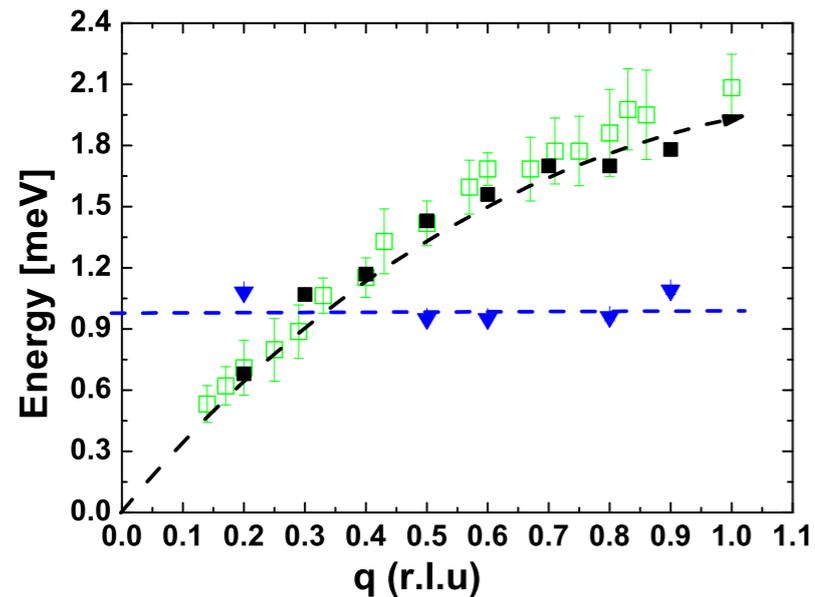


Multiple optical modes?

Look in different directions & polarizations

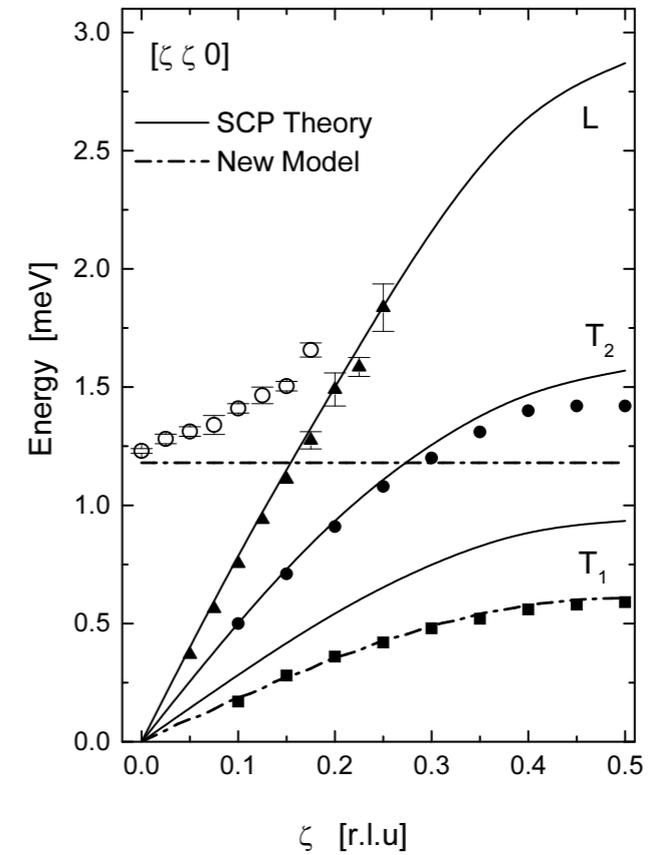


$\tau(100)$



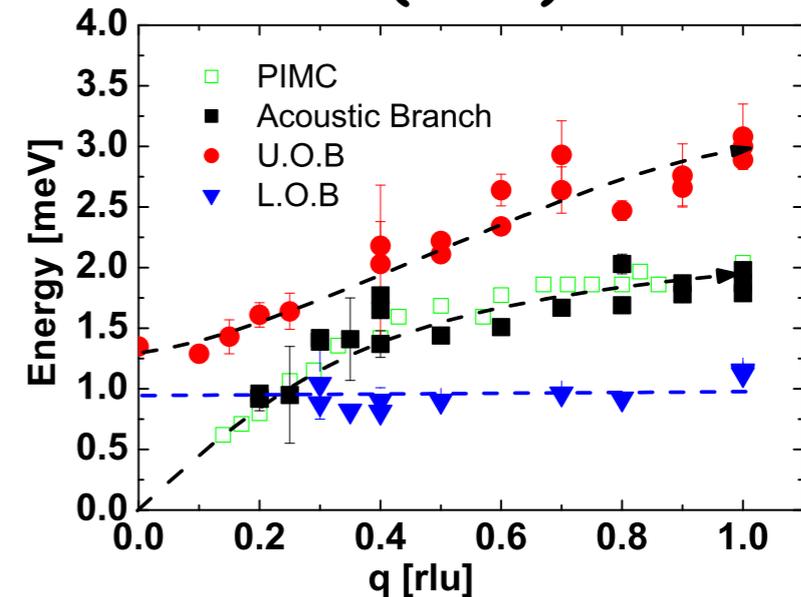
Pelleg et al, PRB 73, 180301R ('06)

(110)



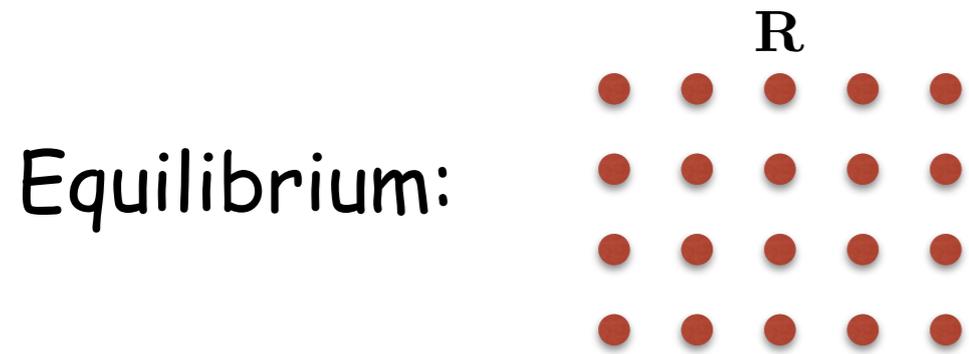
Markovic et al., PRL 88, 195301 ('02)

L(100)

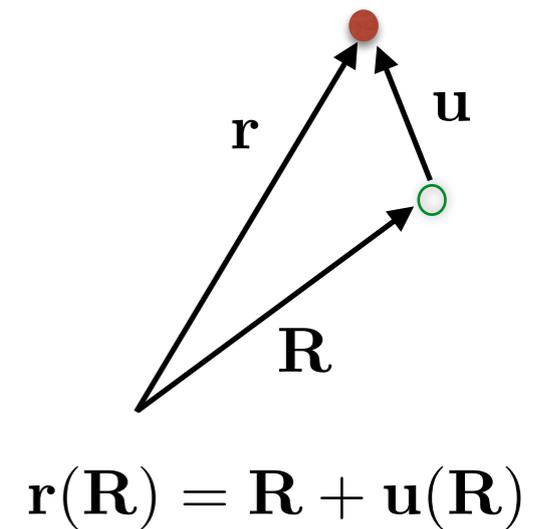


Pelleg et al, JLTP 151, 1164 ('08)

Harmonic theory of solids



Fluctuations:



Small fluctuations $\sqrt{\langle \mathbf{u}^2 \rangle} \ll \Delta R$

$$U_{\text{harm}} = \frac{1}{2} \sum_{\mathbf{R}\mathbf{R}'} \sum_{\mu\nu} u_{\mu}(\mathbf{R}) D_{\mu\nu}(\mathbf{R} - \mathbf{R}') u_{\nu}(\mathbf{R}')$$

Monatomic Bravais lattice \Leftrightarrow acoustic phonons only

Corrections to harmonic theory: $U_{\text{anh}} \sim u^3 + u^4 + \dots$

Lindemann criterion: $\sqrt{\langle \mathbf{u}^2 \rangle} = 0.1 \Delta R \Leftrightarrow$ melting

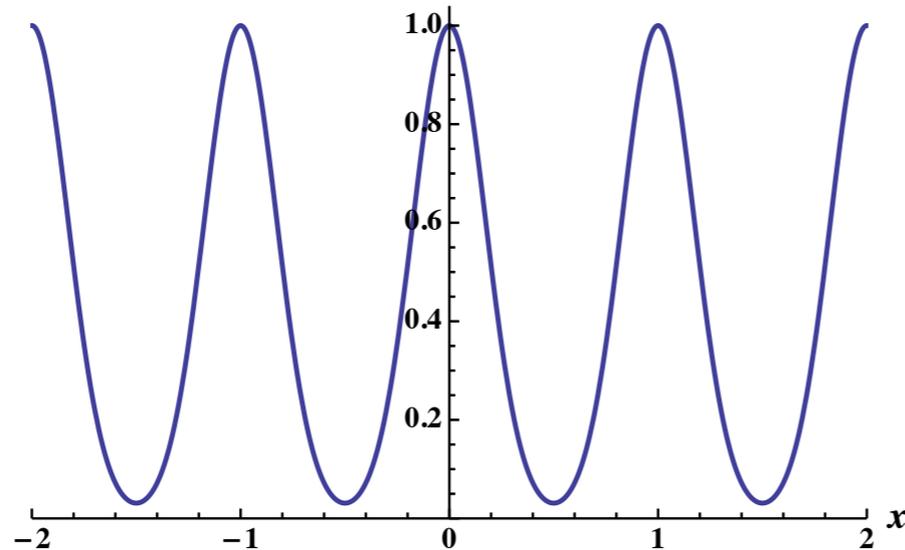
Helium - A quantum solid

Zero point motion

H. Glyde, "Helium, Solid"

Rare-gas crystal	Debye temperature θ_D (K)	Melting temperature T_M (K)	Debye zero point energy $E_{ZD} = \frac{9}{8}\theta_D$	Lindemann parameter $\delta = \langle u^2 \rangle^{1/2}/R$
$^3\text{He}(\text{bcc})$	19	0.65	21	0.368
$^4\text{He}(\text{bcc})$	25	1.6	28	0.292
Ne	66	24.6	74	0.091
Ar	84	83.8	95	0.048
Kr	64	161.4	72	0.036
Xe	55	202.0	62	0.028

Can we think of solid He-4 as a charge density wave (CDW)?



Collective modes of 3d CDW

Apply Ginzburg-Landau analysis to 3d solid Alexander and McTague, PRL (1978)

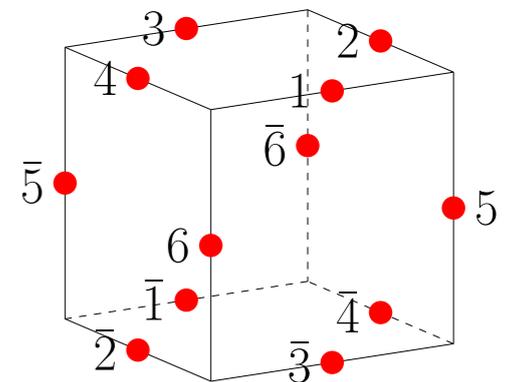
Assume weakly first order transition $\Rightarrow \rho(\mathbf{r}) \equiv n(\mathbf{r}) - n_0$ is small

Dynamical Ginzburg-Landau:
$$L = \frac{1}{2} \int d^3r \left(\frac{\partial \rho}{\partial t} \right)^2 - F_{\text{GL}}$$

Fluctuations about mean-field:
$$\rho(\mathbf{r}, t) = \sum_i (\bar{\rho}_i + \psi_i(\mathbf{r}, t)) e^{i\mathbf{G}_i \cdot \mathbf{r}}$$
$$\psi_i(\mathbf{r}, t) = \psi_{-i}^*(\mathbf{r}, t)$$

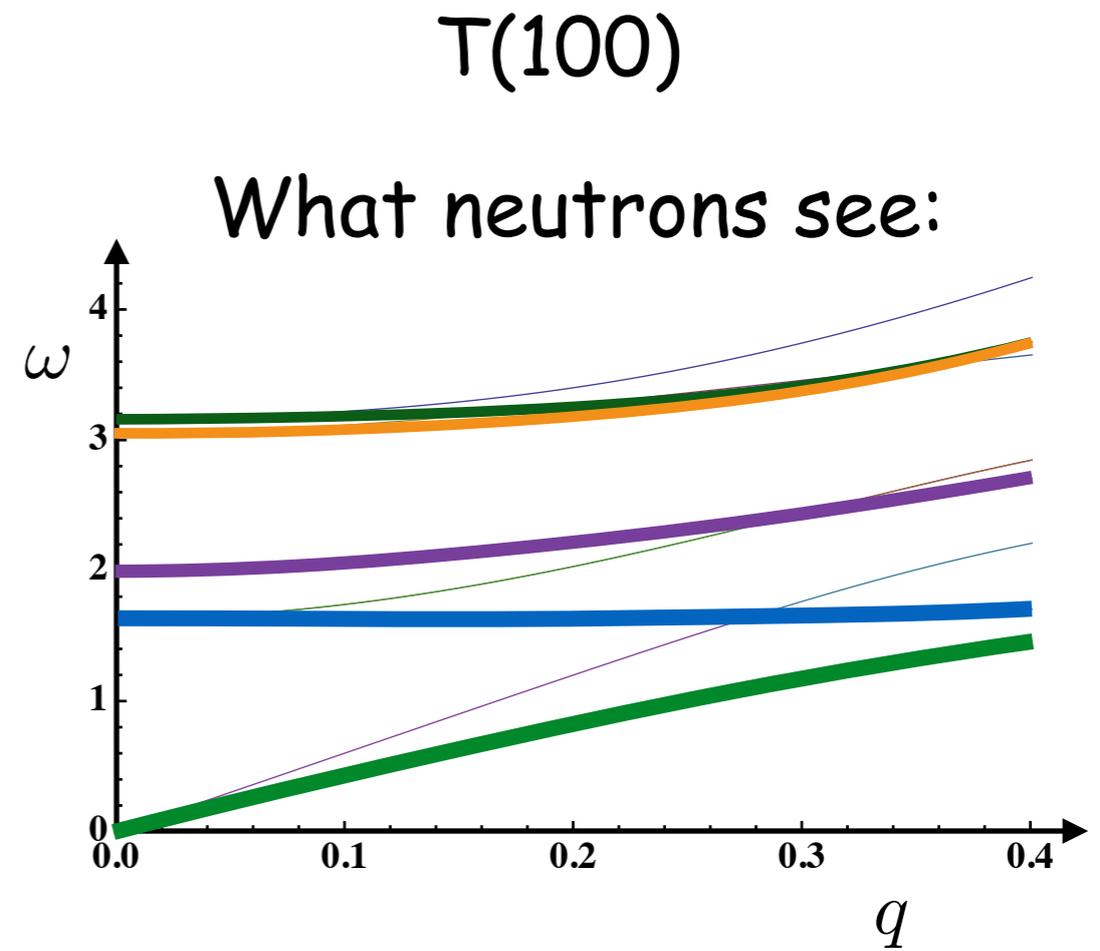
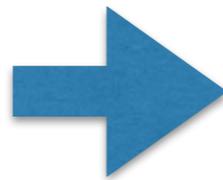
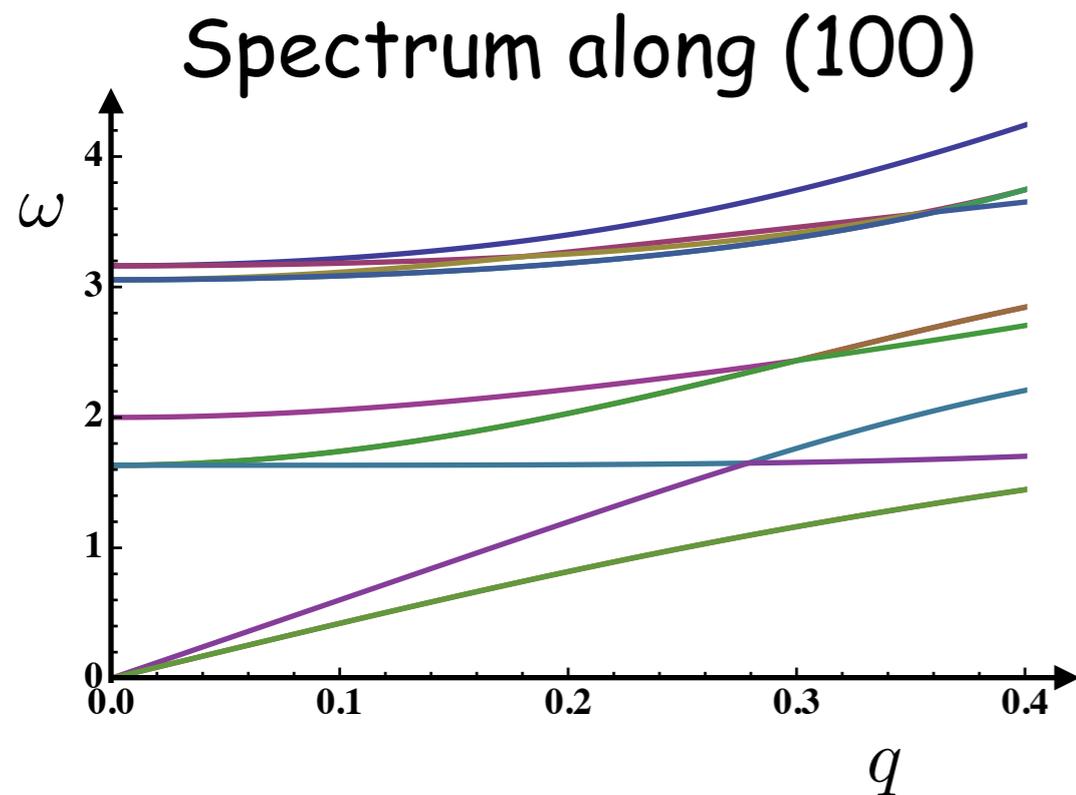
Solve linearized Euler-Lagrange equations

6 pairs of reciprocal lattice vectors \Rightarrow 12 modes!



Spectrum

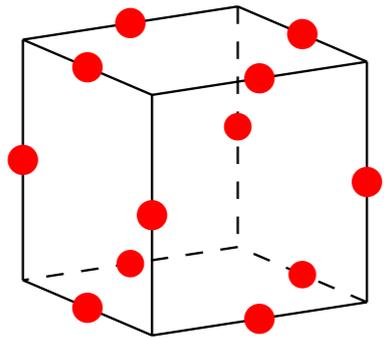
12 modes = 3 acoustic + 9 optical



Structure factor: $S(\mathbf{q}, \omega) = \langle \rho(\mathbf{q}, \omega) \rho(-\mathbf{q}, -\omega) \rangle$

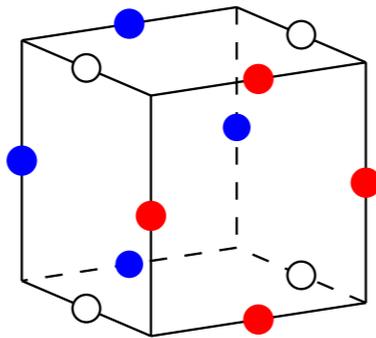
Symmetry of the excitations

“breather” (a_1)

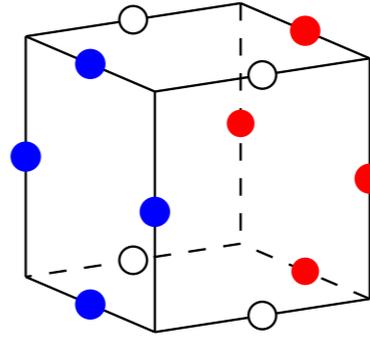


s

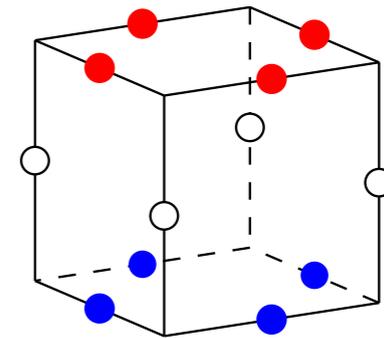
acoustic (t_{1u})



p_x

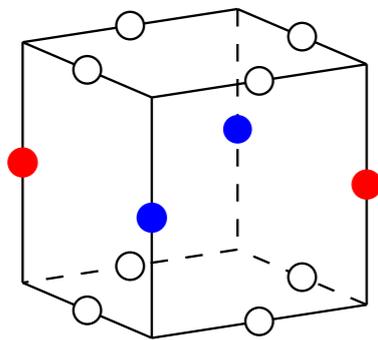


p_y

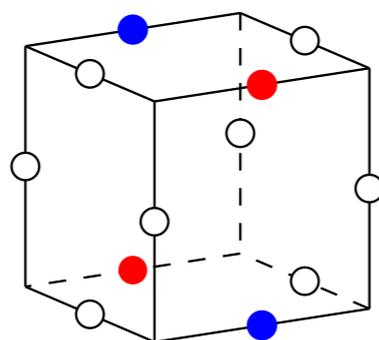


p_z

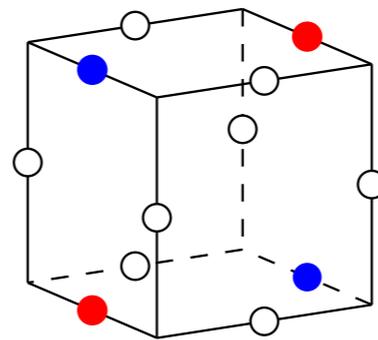
quadrupole (t_{2g})



d_{xy}

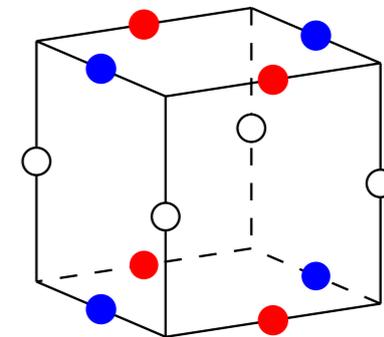


d_{xz}

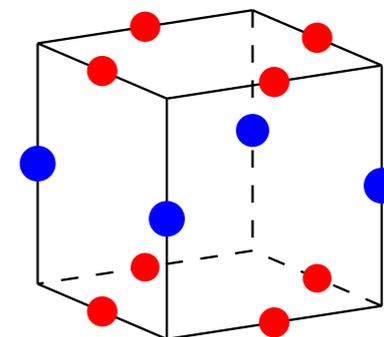


d_{yz}

quadrupole (e_g)

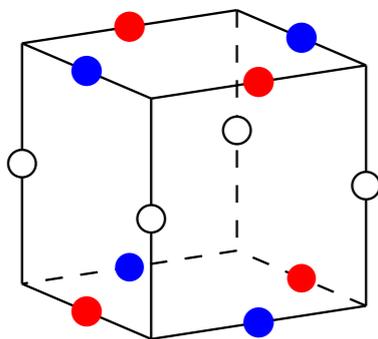


$d_{x^2-y^2}$

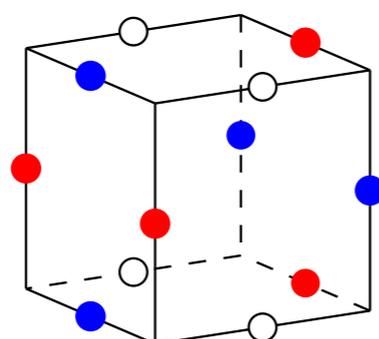


d_{z^2}

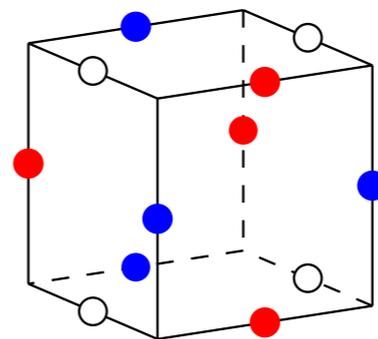
octupole (t_{2u})



$f_z(x^2-y^2)$

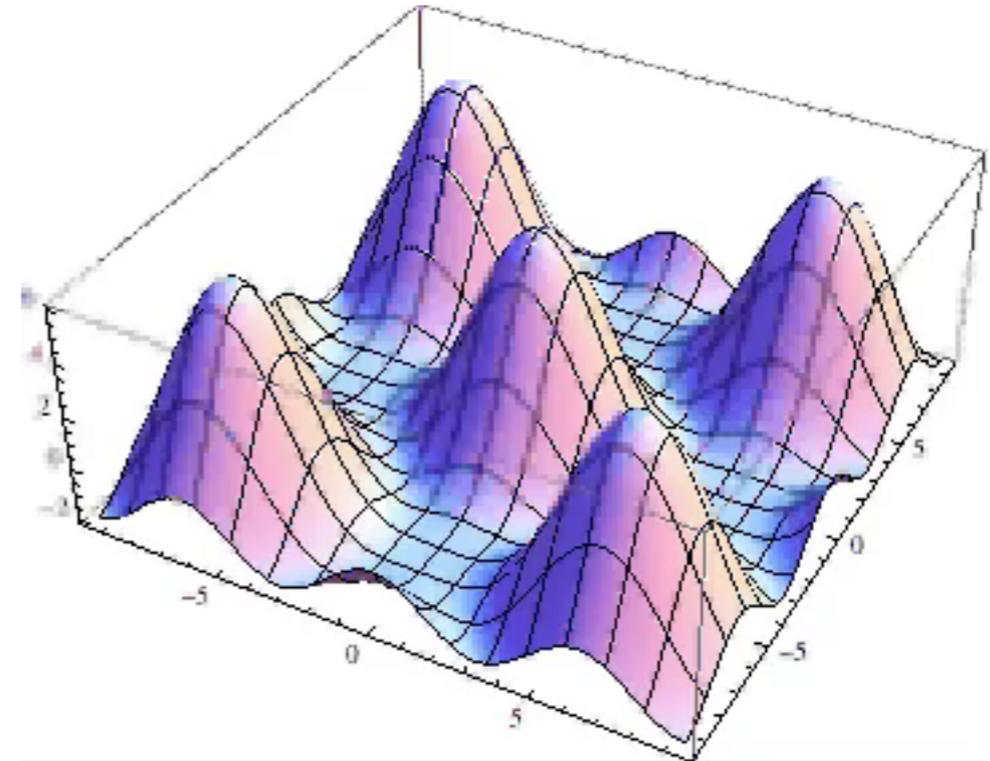
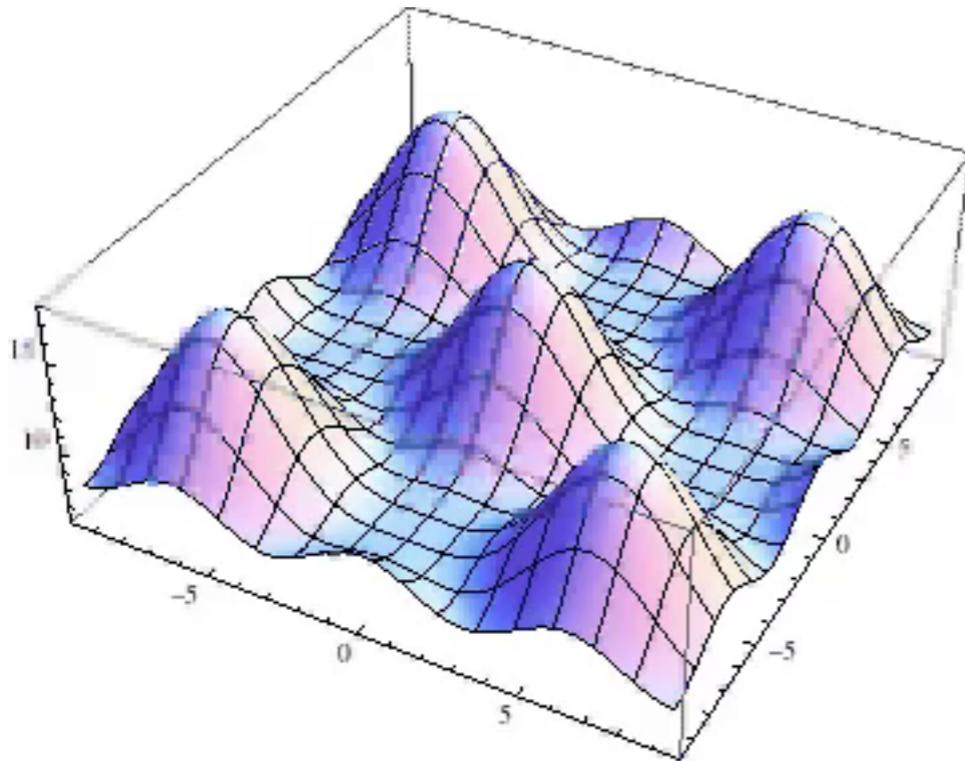
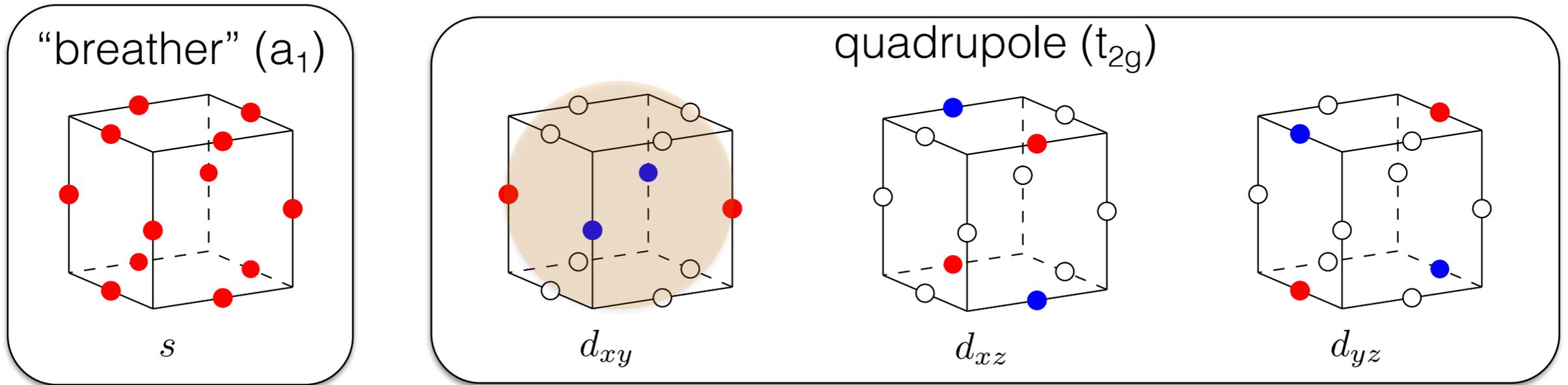


$f_y(z^2-x^2)$



$f_x(z^2-y^2)$

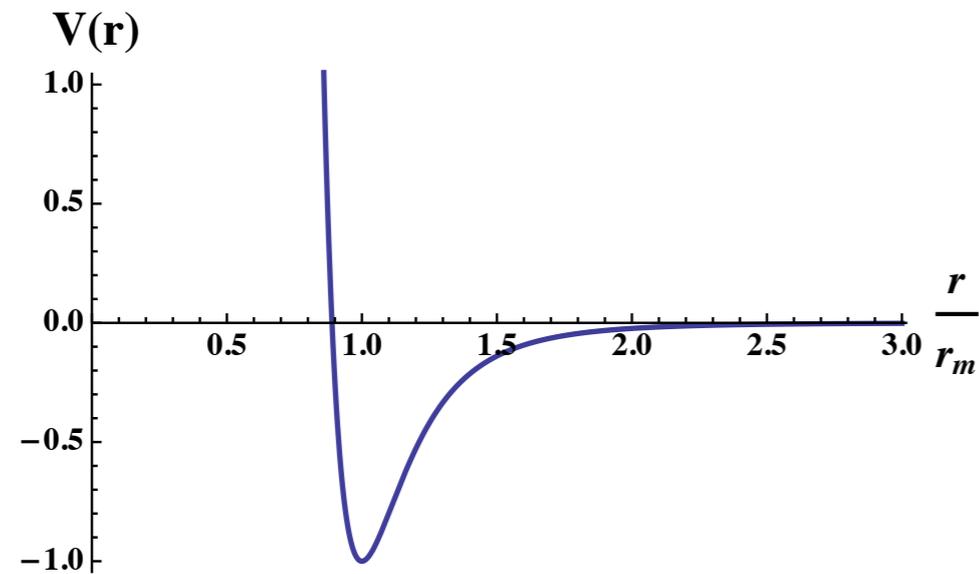
Visualizing the optical modes



d_{xy} “quadrupolon” has vanishing z-axis spring constant \Rightarrow flat band

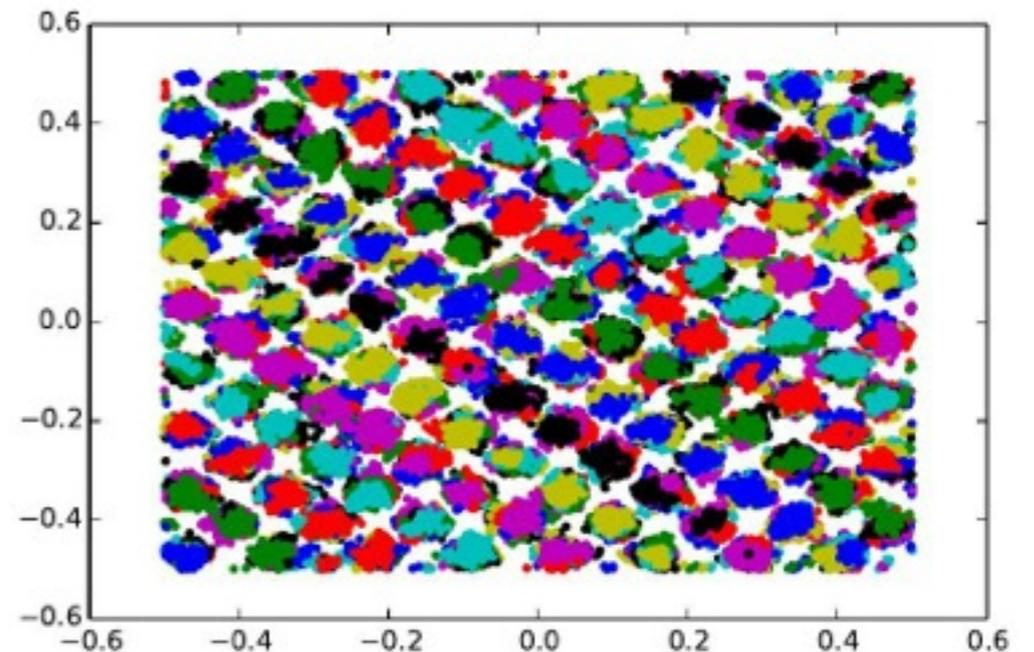
Quantum Monte Carlo

AB-initio simulations (Aziz potential)



Continuous space path integral QMC

2000 He4 Atoms



QMC results

Structure factor:

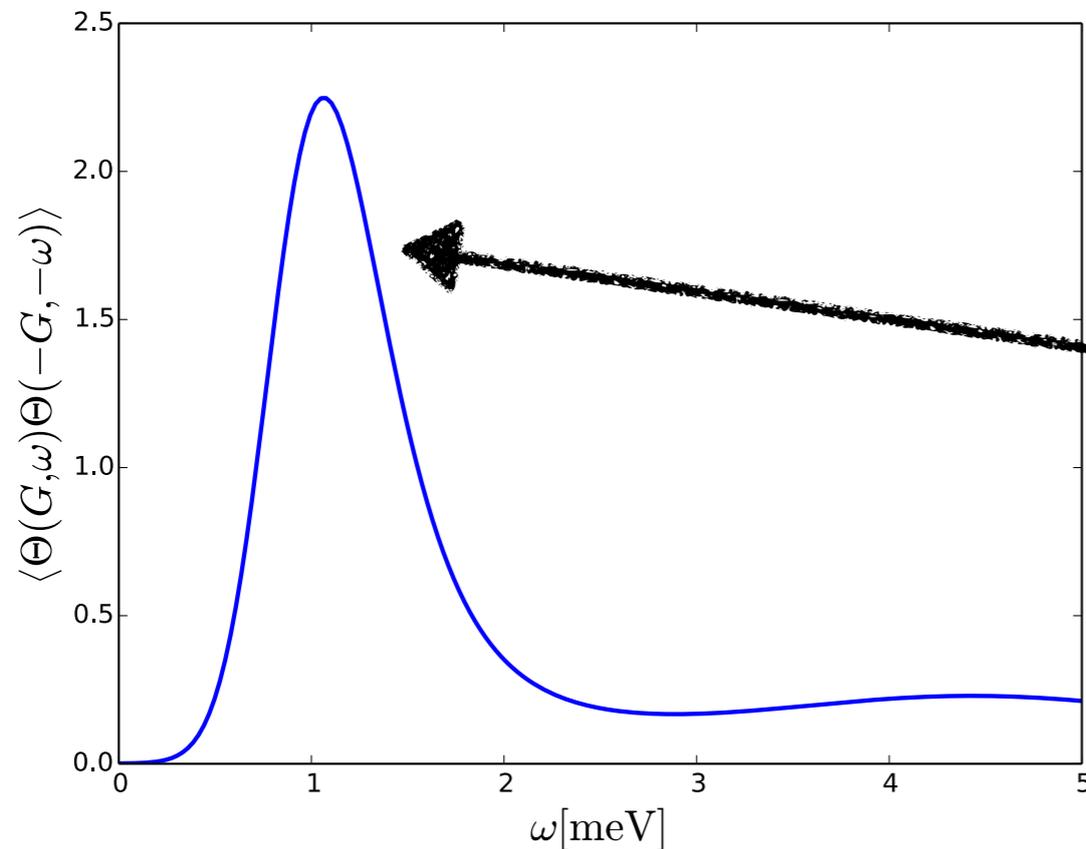
$$S(\mathbf{q}, \omega) = \langle \rho(\mathbf{q}, \omega) \rho(-\mathbf{q}, -\omega) \rangle$$

$$\rho(\mathbf{q}, t) = \sum_n e^{i\mathbf{q} \cdot \mathbf{r}_n(t)}$$

"Scalar susceptibility":

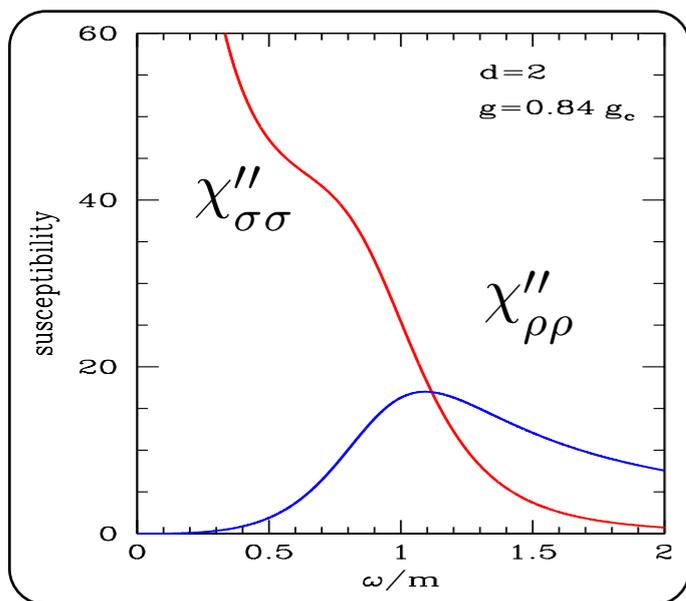
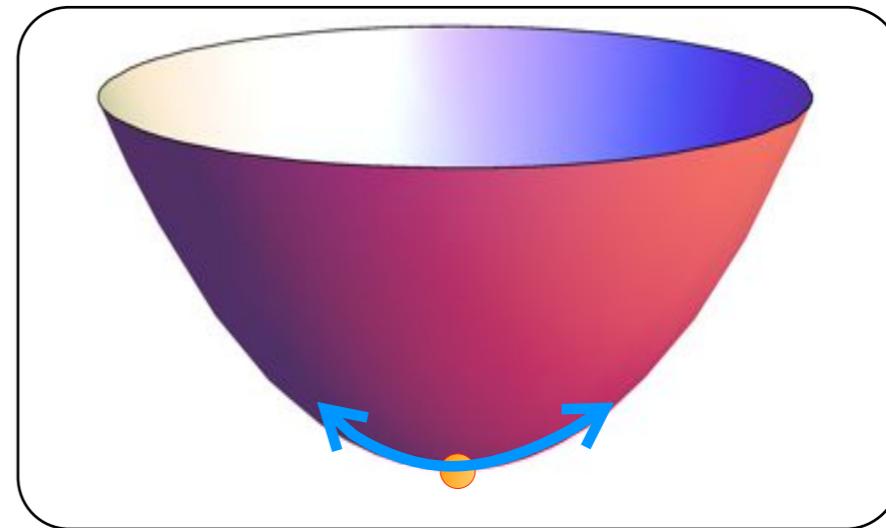
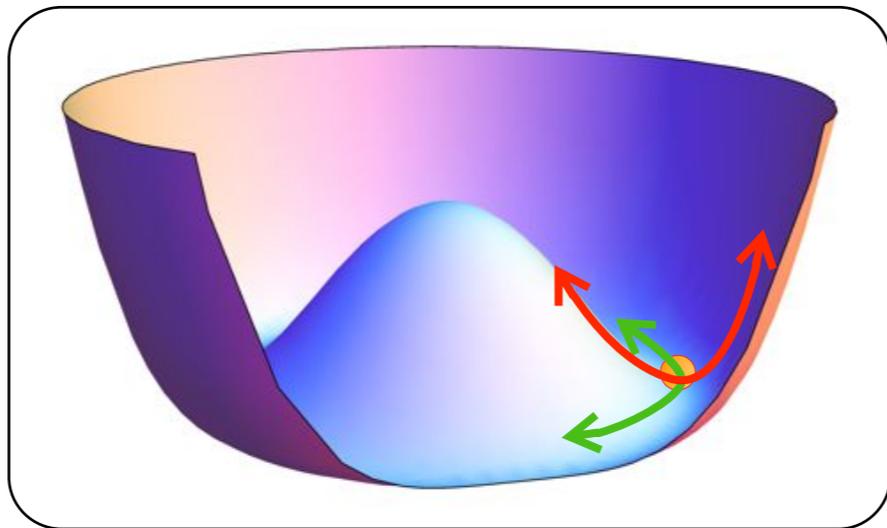
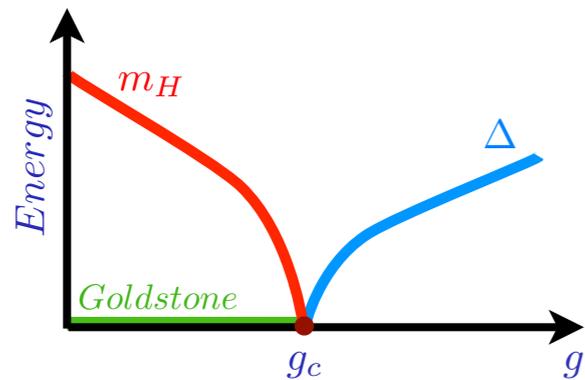
$$S_s(\mathbf{q}, \omega) = \langle \Theta(\mathbf{q}, \omega) \Theta(-\mathbf{q}, -\omega) \rangle$$

$$\Theta(\mathbf{q}, t) = \left| \sum_n e^{i\mathbf{q} \cdot \mathbf{r}_n(t)} \right|^2$$



A clear peak at
 $\omega_H \approx 1 meV$

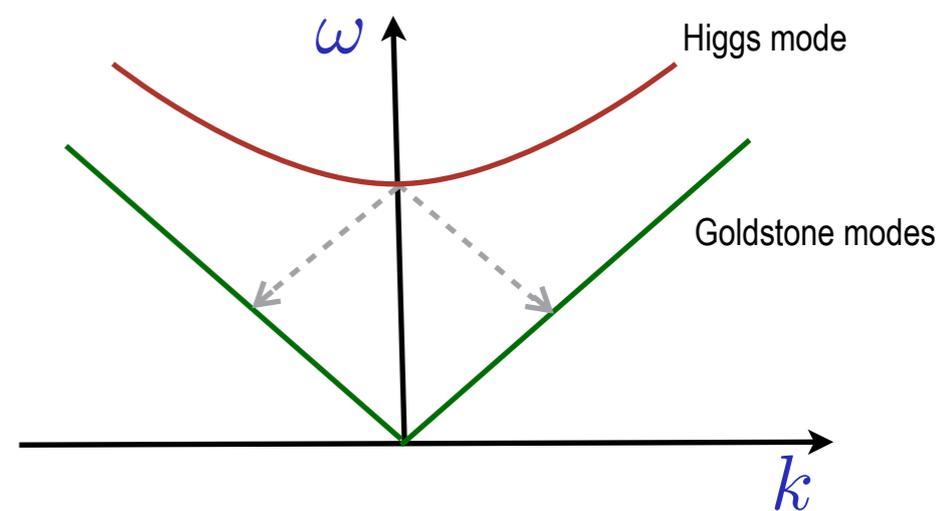
O(N) model



Higgs decays into Goldstone bosons

Summary

scalar $\chi''_{\rho\rho}$ is sharper than longitudinal $\chi''_{\sigma\sigma}$

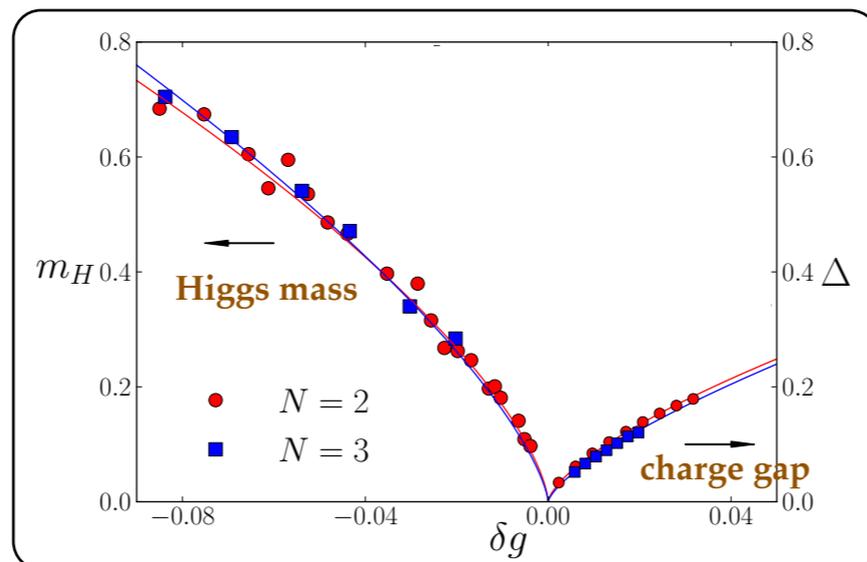


O(1/N), scaling and numerical simulations

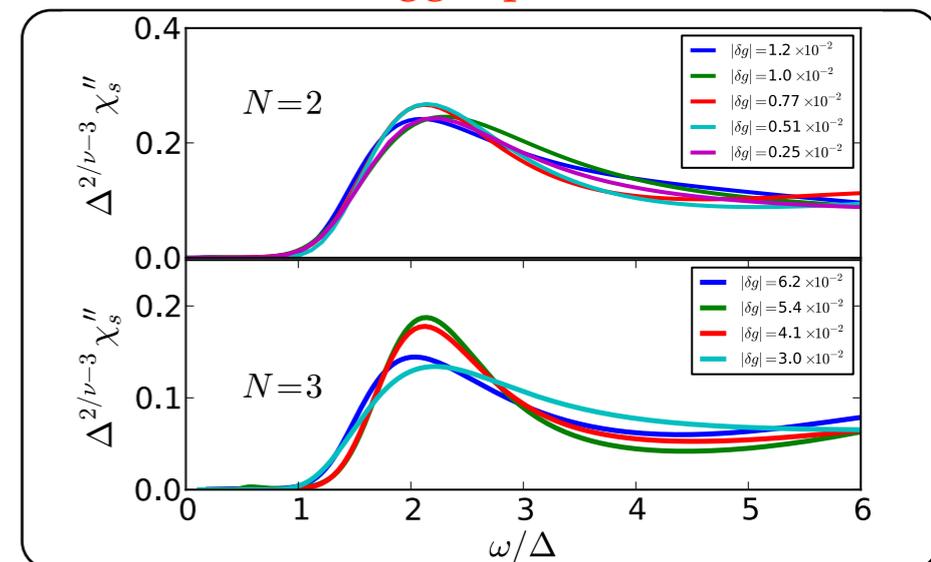
$$\chi_{\rho\rho}(\omega) = \Delta^{3-2/\nu} \Phi_{\rho} \left(\frac{\omega}{\Delta} \right)$$

$$\frac{m_H}{\Delta} = 2.1(3)$$

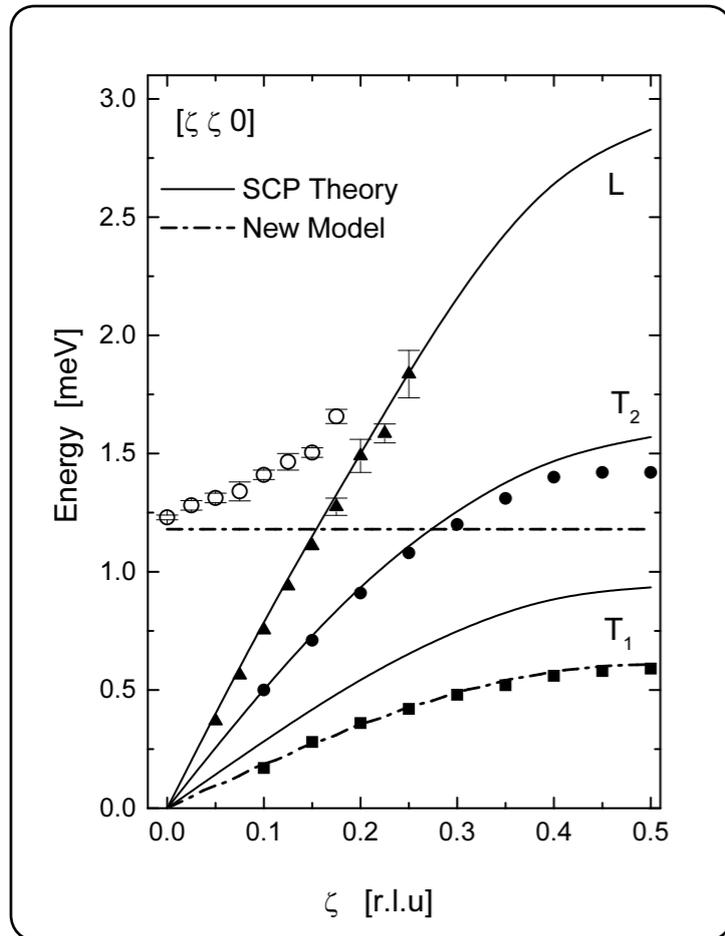
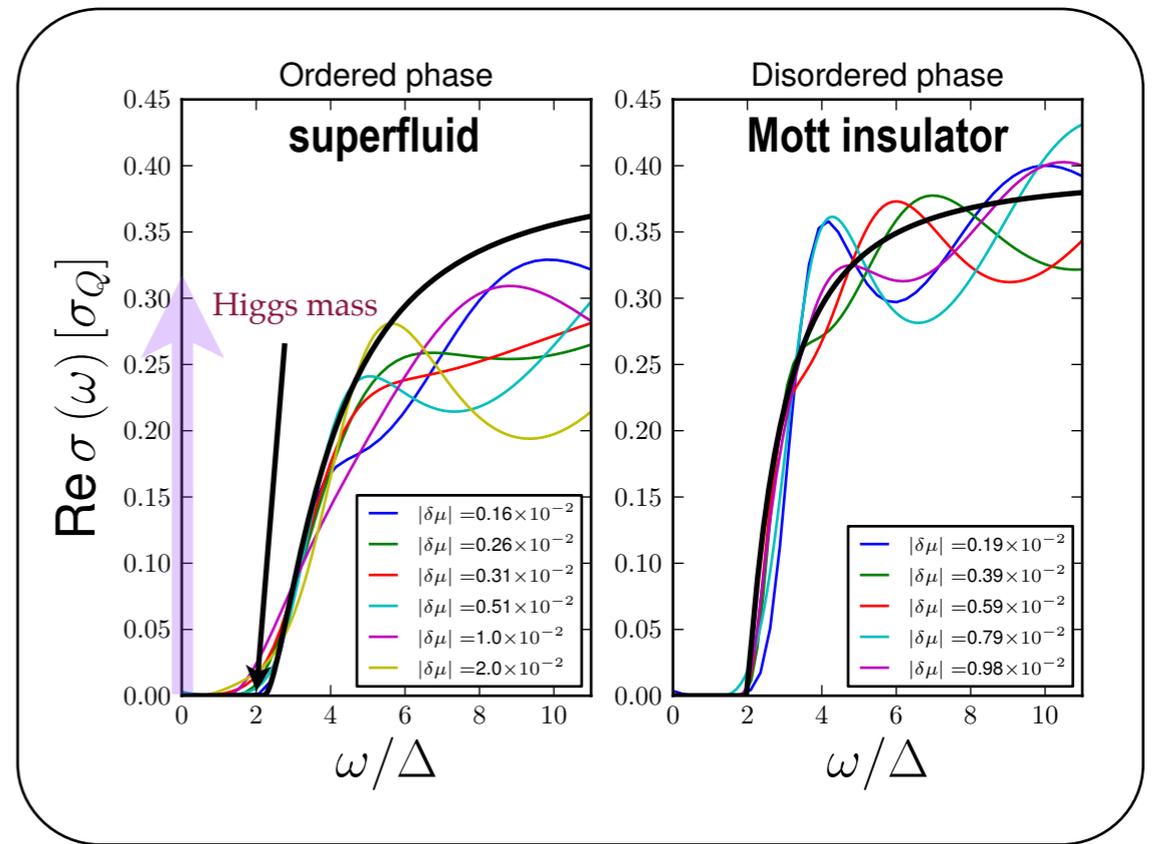
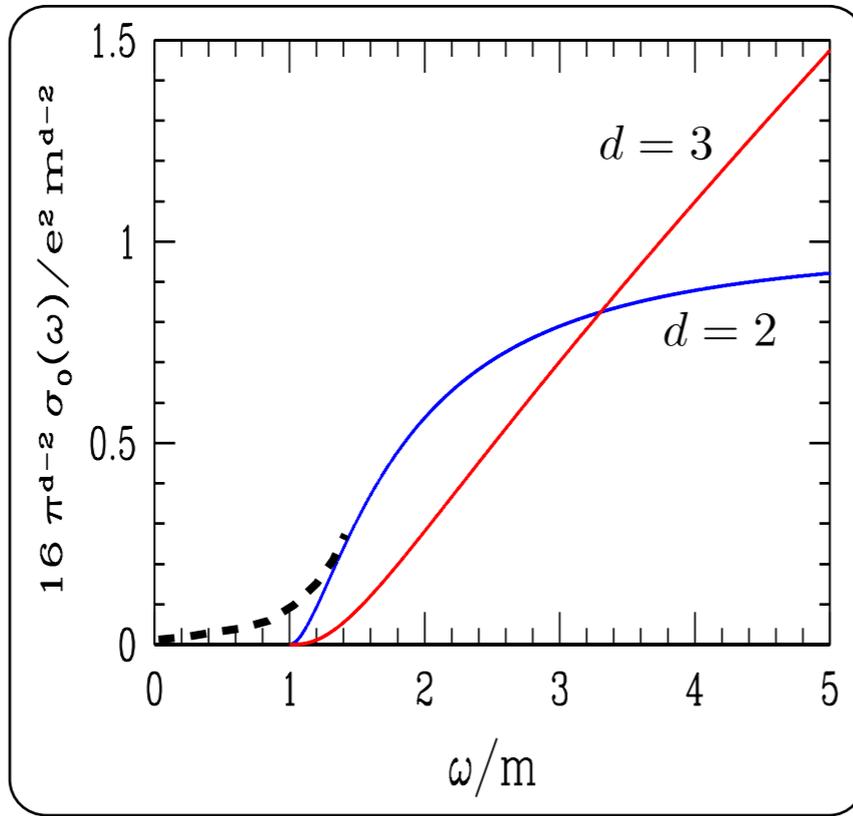
N=2 and N=3 critical behavior



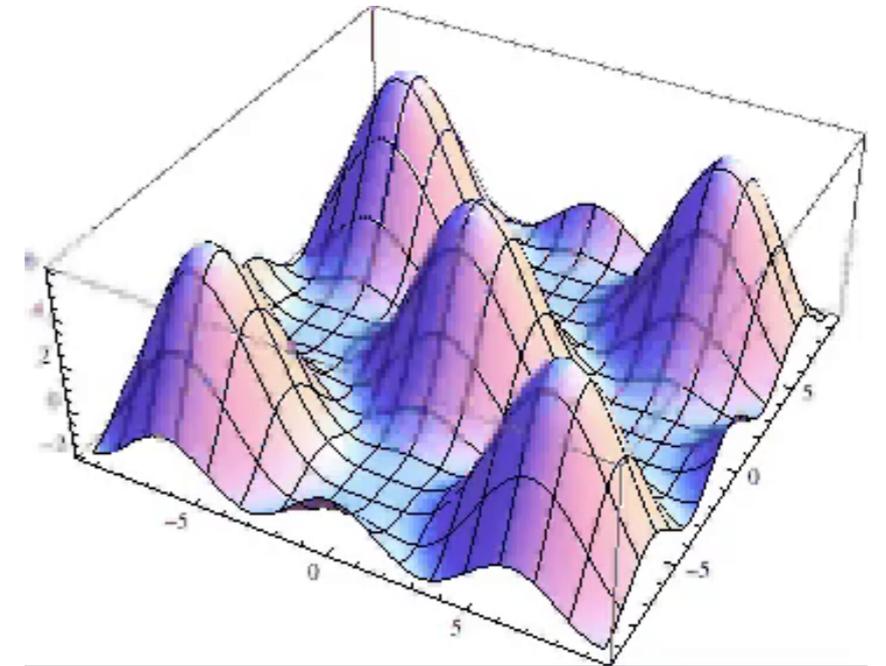
universal Higgs spectral function

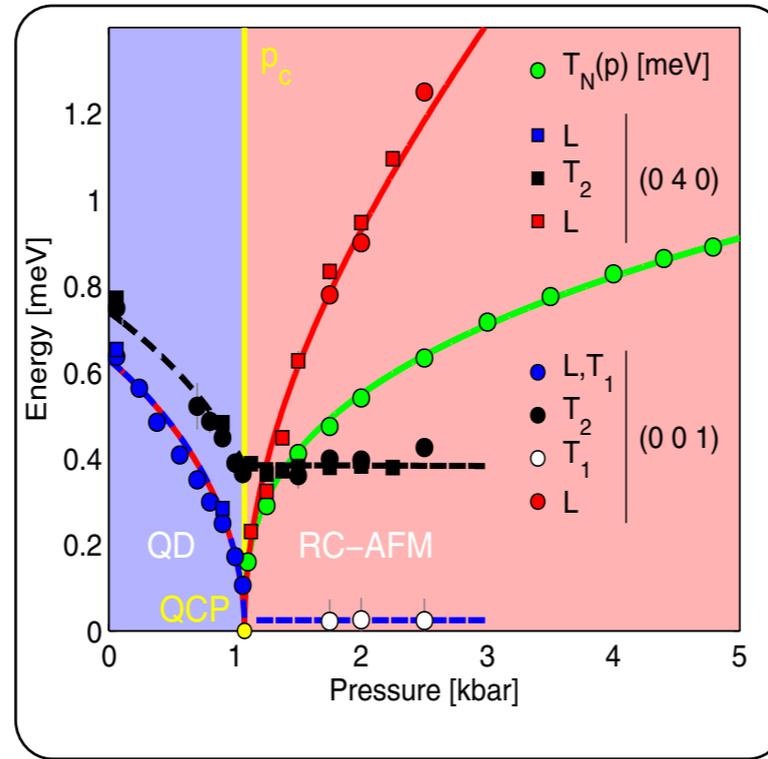
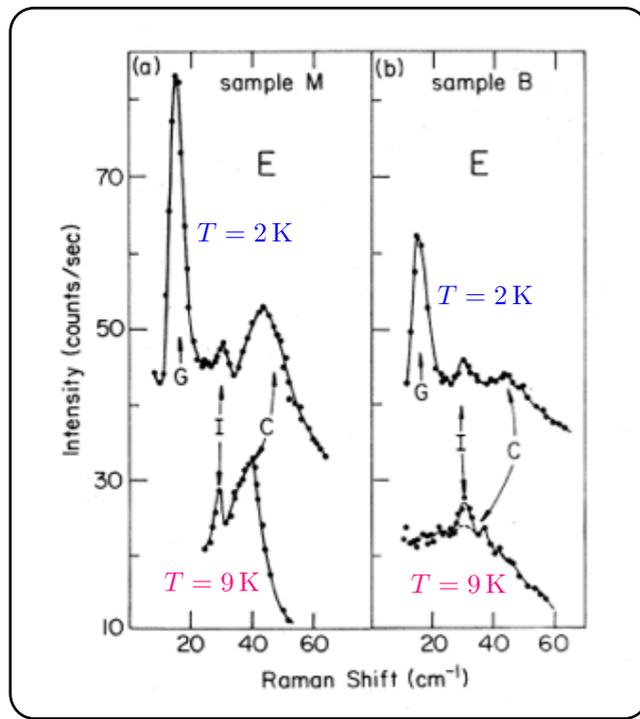


conductivity
pseudogap

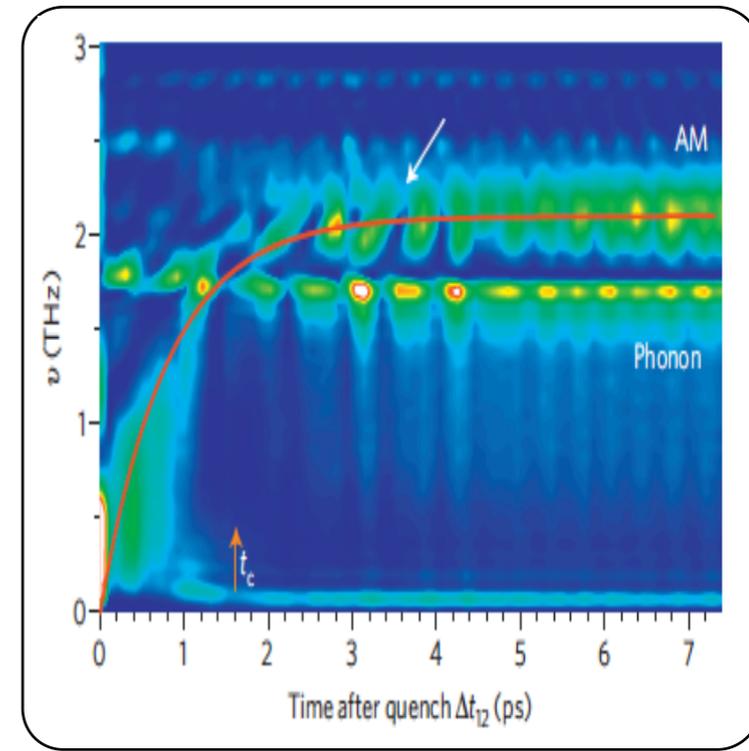


Amplitude mode
in bcc ^4He ?



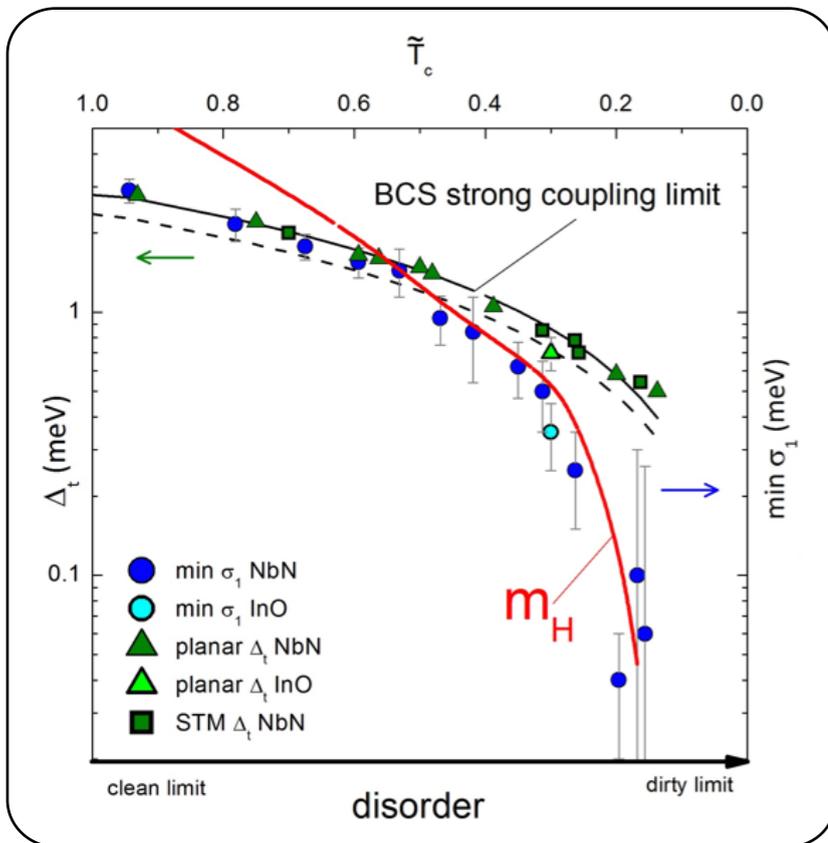


d=3

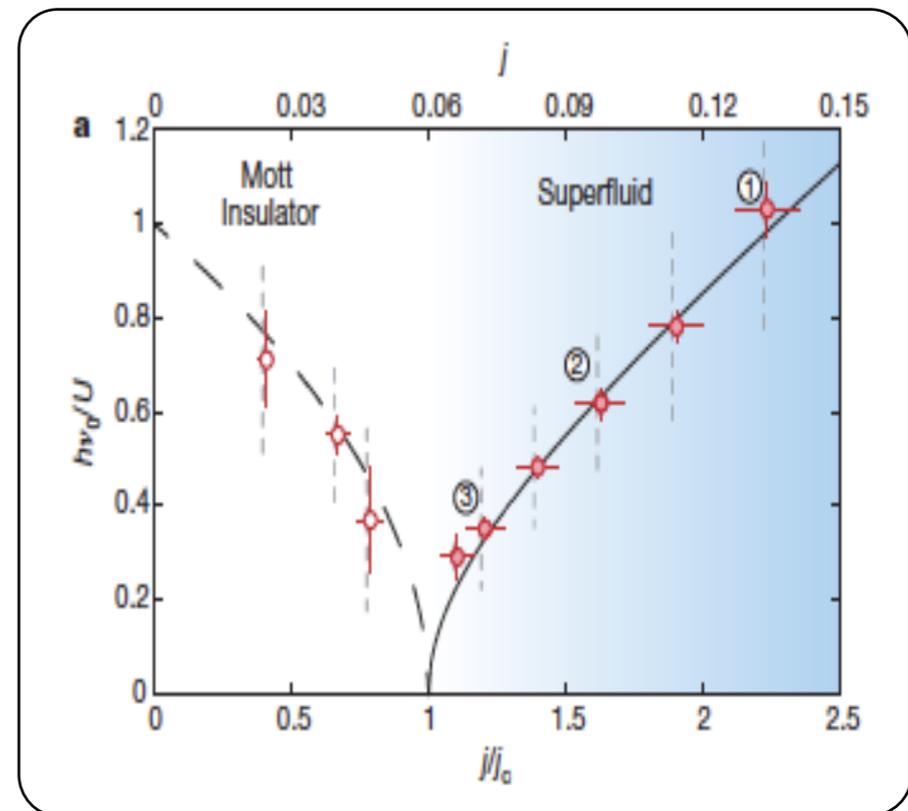


superconductors, antiferromagnets, cold atomic gases, CDWs

Experimental systems



d=2



Higgs Hunters



Dan Arovas
circa 1981



Asa Auerbach



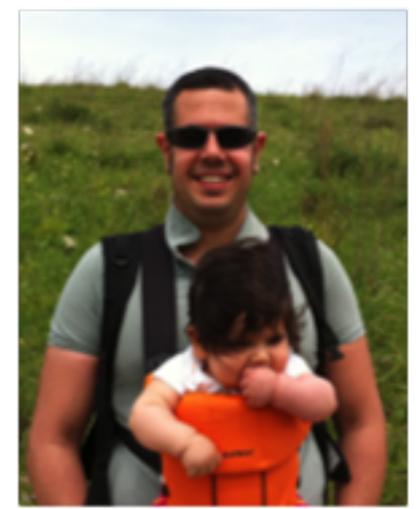
Snir Gazit



Heloise Nonne



Subir Sachdev



Yaniv Tenenbaum Katan

Thank you!