# Higgs amplitude mode in the vicinity of a (2+1)-dimensional quantum critical point

Adam Rançon UChicago



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# Outline

- Introduction
- Non-Perturbative Renormalization Group
- Fate of the amplitude mode close to the QCP
- Conclusion

# Quantum O(N) model

$$S = \int_{0}^{\beta} d\tau \int d^{d}x \left\{ \frac{1}{2} \psi_{i} (-c^{-2} \partial_{\tau}^{2} - \nabla^{2} + r_{0}) \psi_{i} + \frac{u_{0}}{4! N} (\psi_{i}^{2})^{2} \right\}$$

- Generalization of classical O(N) model
- at T=0, Lorentz symmetry
- Describes critical regime of a number of systems :
  - bosons in optical lattices
  - antiferromagnets
  - Josephson junction arrays
  - granular superconductors, ...

### Typical phase diagram in 2D



Quantum critical point described by the classical O(N) model in 2+1 D

Non trivial critical exponent :  $\eta, 
u$ 

One (quantum) exponent : dynamical exponent z = 1 (Lorentz symmetry)

### Amplitude mode : Mean Field



Broken symmetry phase  $\langle \phi_i \rangle = \delta_{i1} \phi_0$ 

 $p^2 = \omega^2 + c^2 \mathbf{q}^2$ 

Mean-field picture : - N-1 transverse (massless) Goldstone modes  $G_{\pi}(f)$ 

- 1 gapped longitudinal mode

es 
$$G_{\pi}(p) = \frac{1}{p^2}$$
  
 $G_{\sigma}(p) = \frac{1}{p^2 + 2\Delta^2}$ 

MF spectral function of one gapped mode arbitrarily close to the QCP (  $\Delta \rightarrow 0$  )

$$\chi_{\sigma}(\omega) \propto \delta(\omega - \sqrt{2}|\Delta|)$$



Strong coupling between longitudinal and transverse fluctuations :

$$G_{\sigma}(p) \propto rac{1}{|p|}$$
 at low energy

Amplitude fluctuations seem not to be a well defined mode close to criticality: the 1/p divergence might the resonance as  $\Delta$  goes to zero.

Zwerger 2004, Dupuis 2011

### Amplitude mode : scalar fluctuation

Podolsky et al. 2011 : it might depend on the correlation function !

$$\phi = \sqrt{\rho} \mathbf{n} \qquad \qquad G_{\rho} = \left< \rho \rho \right>$$
 magnitude

Not the same correlation function ! Example for antiferromagnets: - longitudinal = neutron scattering - scalar (amplitude) = Raman scattering

At low energy, coupling for decay of an amplitude fluctuation into two Goldstone is small (derivative coupling – proportional to the energy square).

We expect the spectral function to be of order  $\omega^3$  at small  $\omega$  (no divergence), so the resonance might still be visible close to the QCP.

But is the amplitude mode still well defined close to the quantum critical point ?

Spectral functions : far from criticality



Spectral functions : far from criticality



Spectral functions : closer to criticality



Spectral functions : closer to criticality



Spectral functions : critical regime – no amplitude mode



### Amplitude mode : Small N Monte Carlo

Recent Monte Carlo simulations : Polet et al. (2012), Gazit et al. (2013), Chen et al. (2013) for N=2 and 3

Also measured in cold atoms experiments : ``Higgs" mode (I. Bloch group)



Amplitude mode defined close to criticality, but how does it disappear as N increases ?

## Amplitude mode : disordered phase ?

Kun Chen et al. : amplitude mode in the symmetric (disordered) phase !?



Gave an RG argument, but is it true ?

**Non-Pertubative Renormalization Group** 

#### Family of actions indexed by momentum scale k



Exact Flow equation (Wetterich `93):  $\partial_k \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \left\{ \partial_k R_k \left( \Gamma_k^{(2)} + R_k \right)^{-1} \right\}$ 

# Average Effective Action ("Gibbs" free energy)

Bare action :

$$S = \int_0^\beta d\tau \int d^d x \left\{ \frac{1}{2} \psi_i (-c^{-2} \partial_\tau^2 - \nabla^2 + r_0) \psi_i + \frac{u_0}{4!N} (\psi_i^2)^2 \right\}$$

"Functional free energy" gives access to all the physics :

$$\begin{split} W_k[J,h] &= \ln \int D\psi_i e^{-S - \Delta S_k} + J_i \psi_i + h\rho \\ \text{modified Legendre transform} & \phi_i(\tau) &= \langle \psi_i(\tau) \rangle \\ & \Gamma_k[\phi,h] \\ \text{for h=0 in constant field} \\ \rho &= \sum_i \phi_i^2 \\ V_k(\rho)^i \\ \text{thermodynamics :} \\ \text{pressure, order parameter, etc.} \\ \end{split}$$

### **Approximations**

Write down an Ansatz for  $\,\Gamma_k[\phi,h]\,$  (expansion in  $\phi$  and h)



Propagator computed within a derivative expansion but keep the Matsubara frequency dependence. In the end, analytic continuation.

Technical details in AR and N. Dupuis, PRB 2014

### Amplitude mode : N=2

N=2, m≈2.4 Δ



AR & N. Dupuis, PRB 89 180501(R) (2014)

### Amplitude mode : N=3

![](_page_18_Figure_1.jpeg)

AR & N. Dupuis, PRB 89 180501(R) (2014)

### Amplitude mode : all N

No resonance for N>2

![](_page_19_Figure_2.jpeg)

## T=0 disordered phase - 1

-"Resonance" close to gap at 2Δ : related to amplitude mode ?

-If so : resonance must appear in the critical regime (otherwise : not related)

-But in critical regime : scaling  $\chi_{\rm crit}''(\omega) \propto \omega^{3-2/\nu}$ 

-Implies no resonance in critical regime

- Thus resonance in ordered/ disordered regime are not related – not the same physics

![](_page_20_Figure_6.jpeg)

![](_page_20_Figure_7.jpeg)

# Quantum disordered phase - 2

No resonance in the disordered phase : no amplitude mode precursor in the renormalization flow when the peak appears in the disordered phase.

![](_page_21_Figure_2.jpeg)

### Amplitude mode : Kosterlitz-Thouless phase

![](_page_22_Figure_1.jpeg)

AR & N. Dupuis, PRB 89 180501(R) (2014)

# **Conclusion and Perspectives**

- NPRG study of amplitude mode close to a non-trivial critical point. Available approaches : Monte-Carlo, NPRG, others ?
- Study of the amplitude mode for all N. Nice agreement with Monte-Carlo for N=2. Case N=3 needs further study.
- No amplitude mode resonance in the quantum disordered phase.
- Resonance exists in the BKT phase (vanishing order parameter !). Up to what temperature ? (T<sub>BKT</sub> ≈ m<sub>H</sub>/5)
- Future calculation : Conductivity / viscosity within NPRG.

# Typical energy scales

$$S = \int_0^\beta d\tau \int d^d x \left\{ \frac{1}{2} \psi_i (-c^{-2} \partial_\tau^2 - \nabla^2 + r_0) \psi_i^2 + \frac{u_0}{4!N} (\psi_i^2)^2 \right\}$$

Four energy scales : T ,  $\Delta$  ,  $\, u_0$  ,  $\omega_c \propto \sqrt{r_0 - r_{0,c}}\,$  (not all independent)