

Investigation of the low-energy kaons hadronic interactions in light nuclei by AMADEUS

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On behalf of the AMADEUS collaboration

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MIN16

31/7 – 2/8 2016, YITP, Kyoto University

AMADEUS proposal

Antikaon Matter At DAΦNE: Experiments with Unraveling Spectroscopy

AMADEUS collaboration

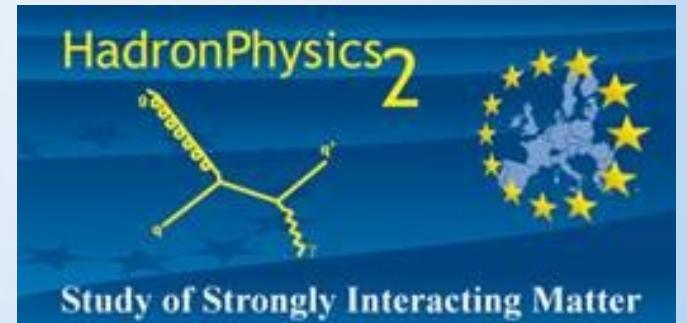
116 scientists from 14 Countries and 34 Institutes

lnf.infn.it/esperimenti/siddharta
and

LNF-07/24(IR) Report on lnf.infn.it web-page (Library)

AMADEUS started in 2005 and
was presented and discussed in all the LNF Scientific Committees

EU Fundings FP7 – I3HP2:
Network WP9 – LEANNIS;
WP24 (SiPM JRA);
WP28 (GEM JRA)



Why AMADEUS & ΔΑΦΝΕ?



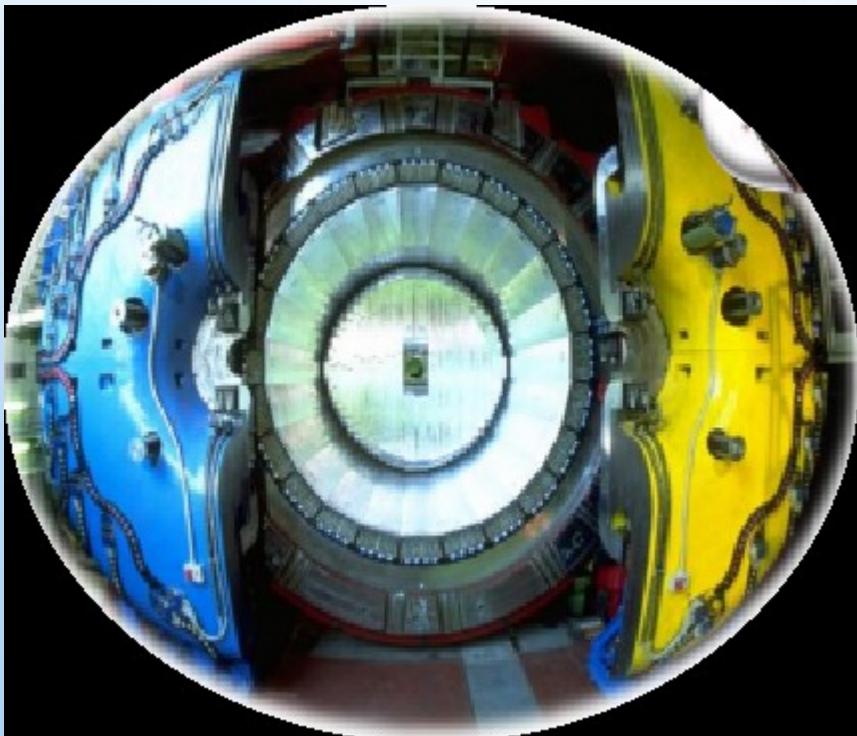
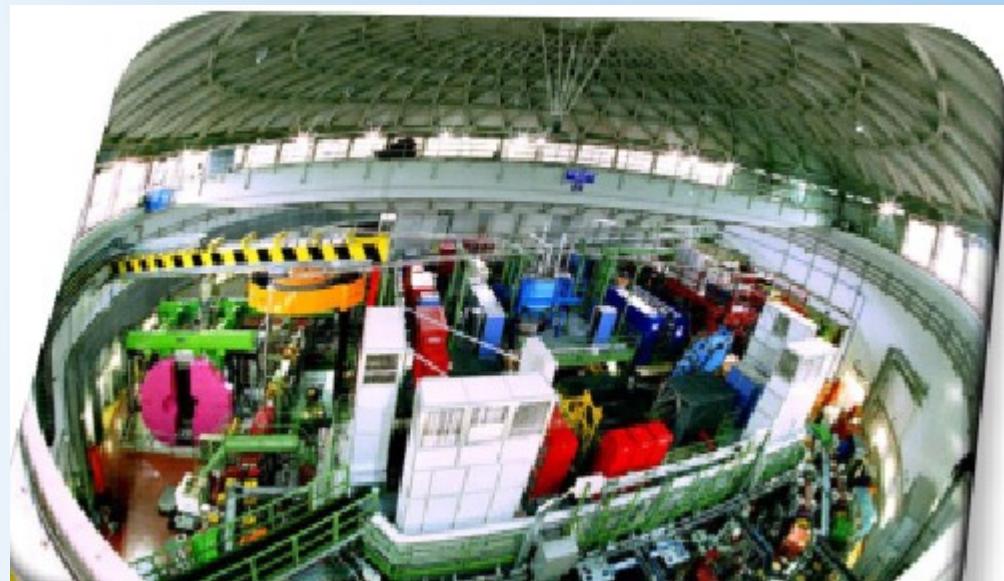
Why AMADEUS & DAΦNE?

DAΦNE

Double ring e^+e^- collider working in C. M. energy of ϕ , producing $\approx 600 K^+K^-$ /s

$\phi \rightarrow K^+K^-$ (BR = $(49.2 \pm 0.6)\%$)

- **low momentum** Kaons
 ≈ 127 Mev/c
- **back to back** K^+K^- topology



KLOE

- 96% acceptance,
- optimized in the energy range of all charged particles involved
- good performance in detecting photons (and neutrons checked by kloNe group (M. Anelli et al., Nucl Inst. Meth. A 581, 368 (2007)))

Why AMADEUS & DAΦNE?

Neutron detection efficiency

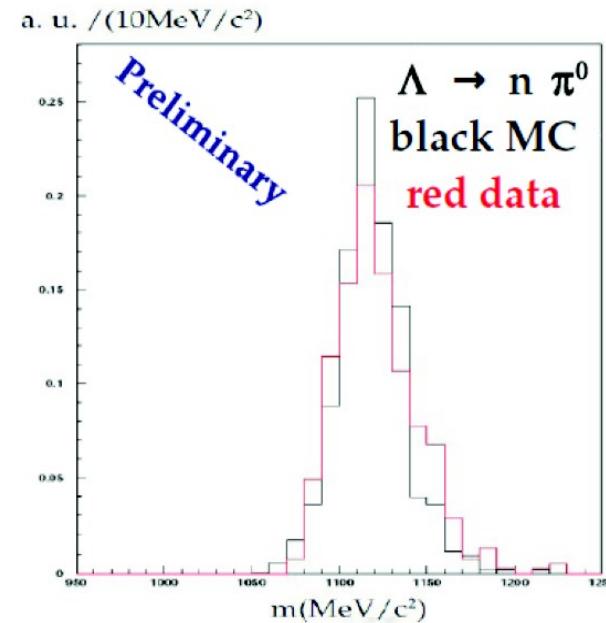
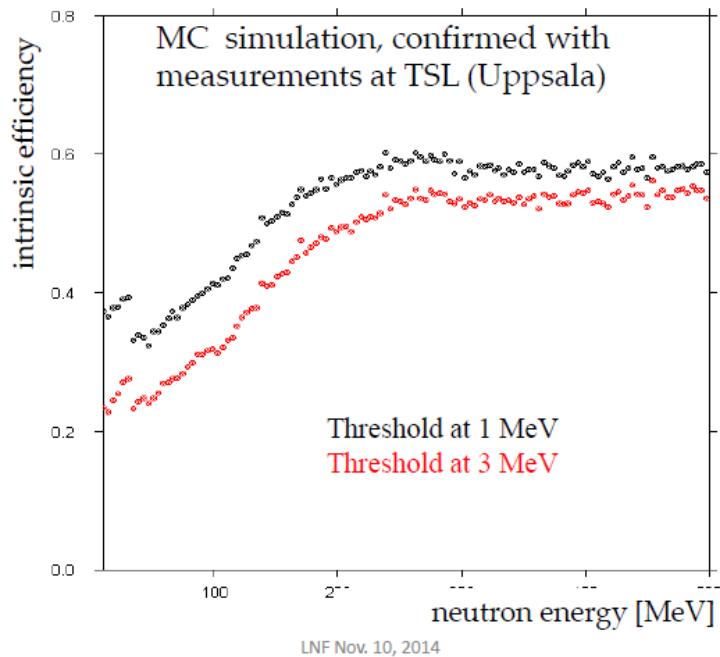


Fig. 1. $n\pi^0$ invariant mass spectrum measured by the KLOE EMC, the red line corresponds to data, the black one corresponds to a Monte Carlo simulation of the $\Lambda \rightarrow n\pi^0$ decay, reconstructed in the KLOE calorimenter.

KLOE

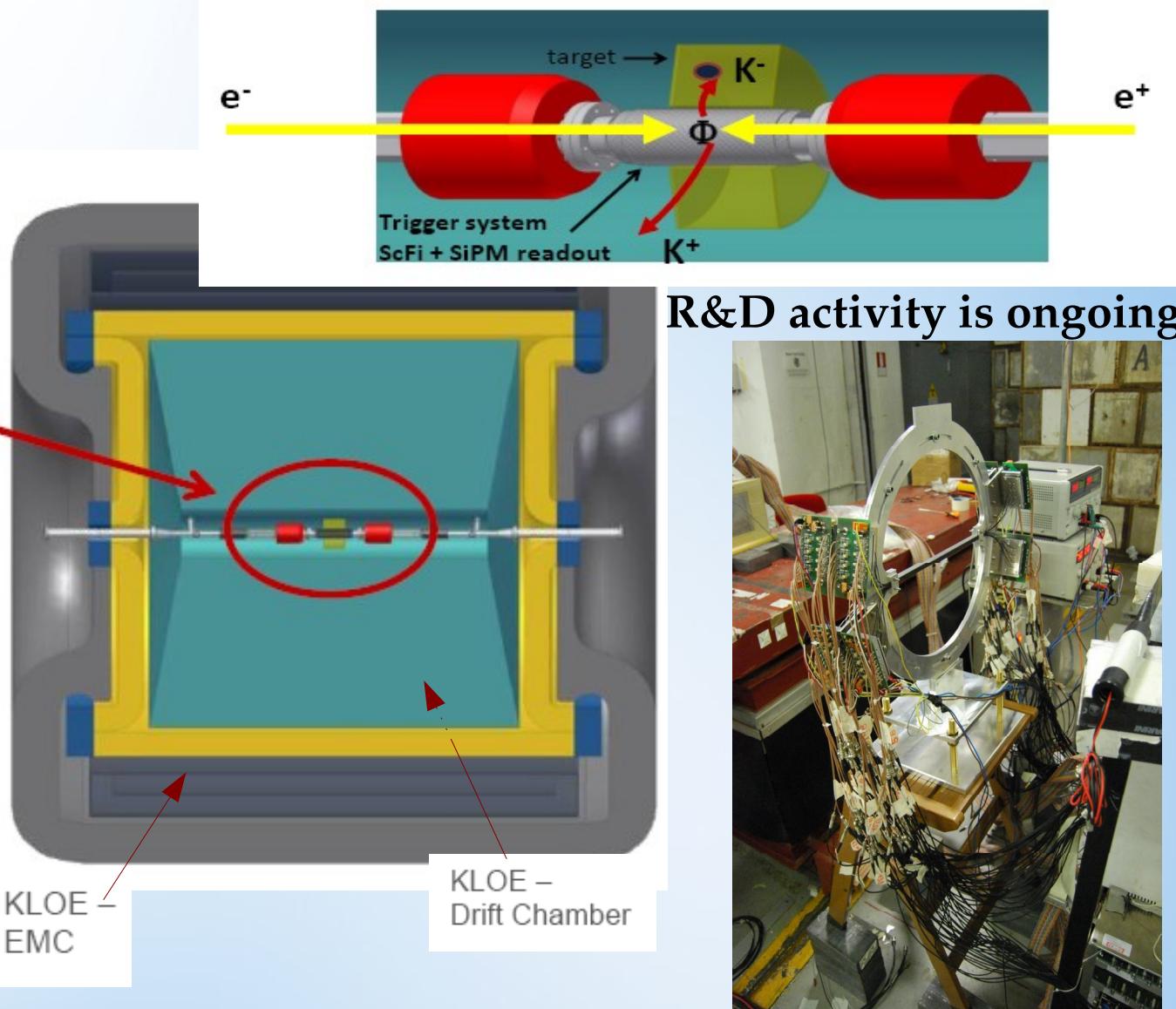
- 96% acceptance,
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AMADEUS target + tracking inside KLOE spectrometer

Implementation of dedicated solid targets & cryogenic gaseous targets (H, d, ^3He , ^4He) inside the KLOE DC.

Nucl. Instrum. Meth. A671 (2012) 125-128

JINST 8 (2013) T05006



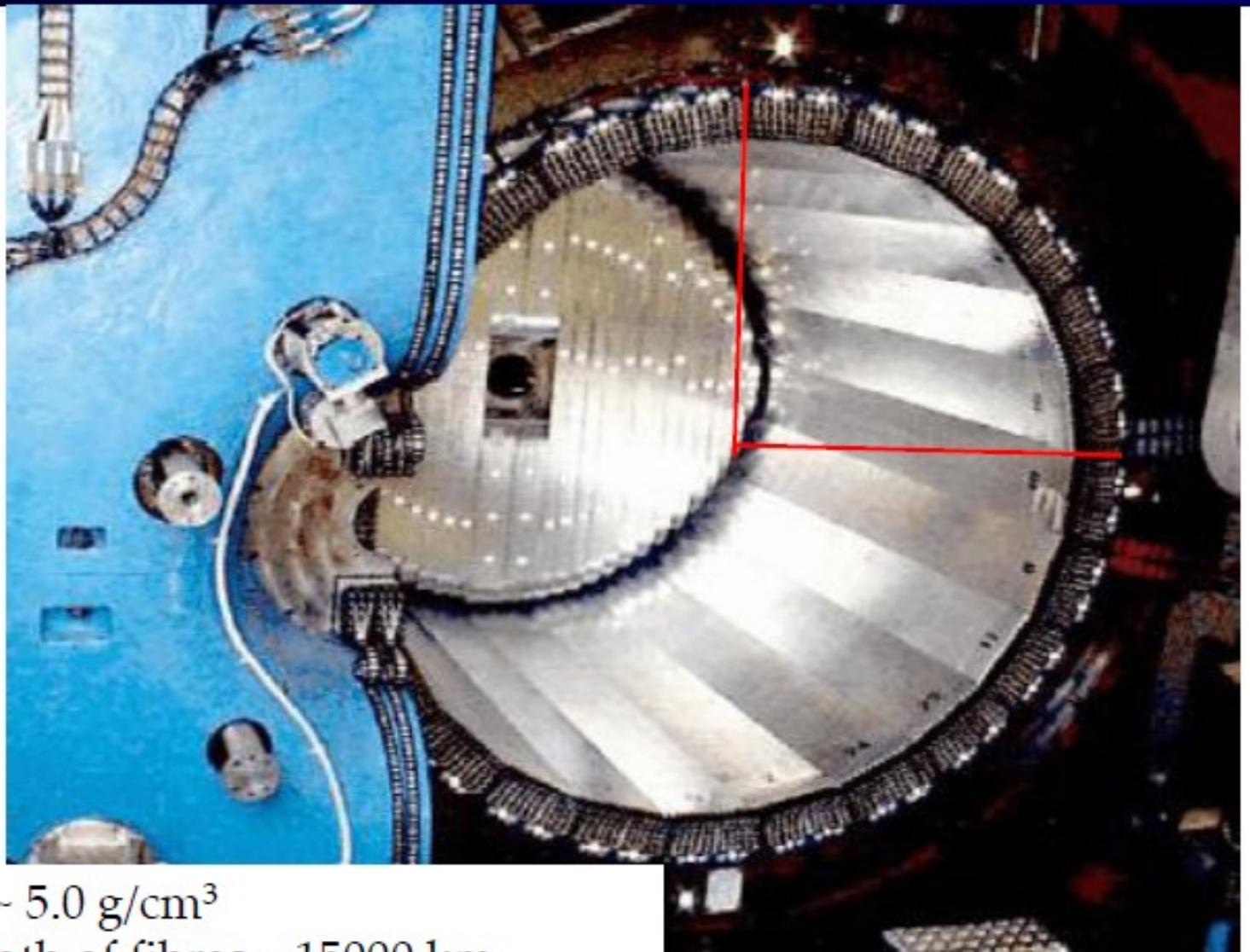
- The AMADEUS setup will be implemented in the 50 cm. gap in KLOE DC around the beam pipe:

- Target** (A gaseous He target for a first phase of study)

- Trigger** (1 or 2 layers of ScFi surrounding the interaction point)

KLOE electromagnetic calorimeter

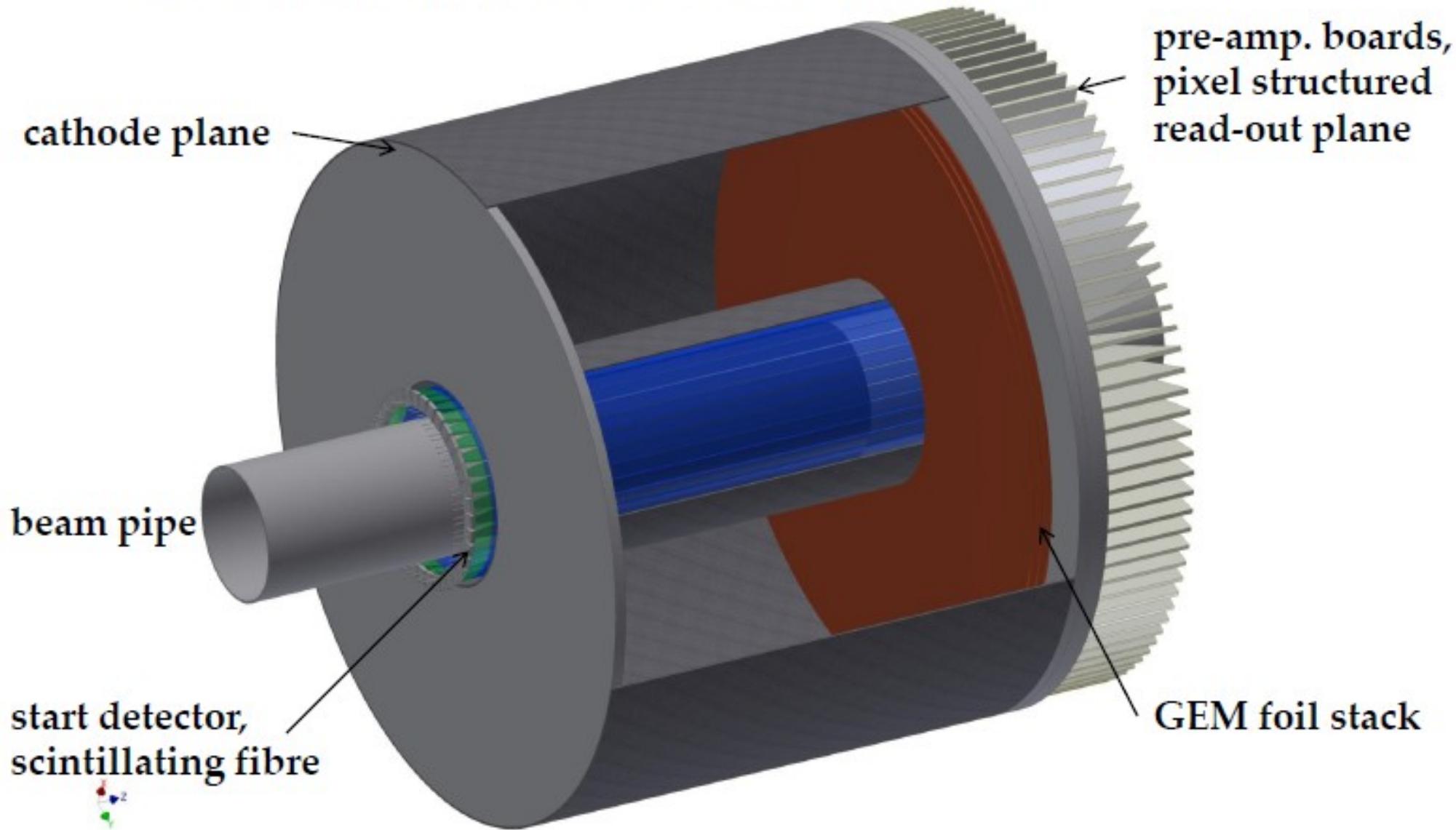
or re-using the
calorimeter only,
with a new
tracking detector



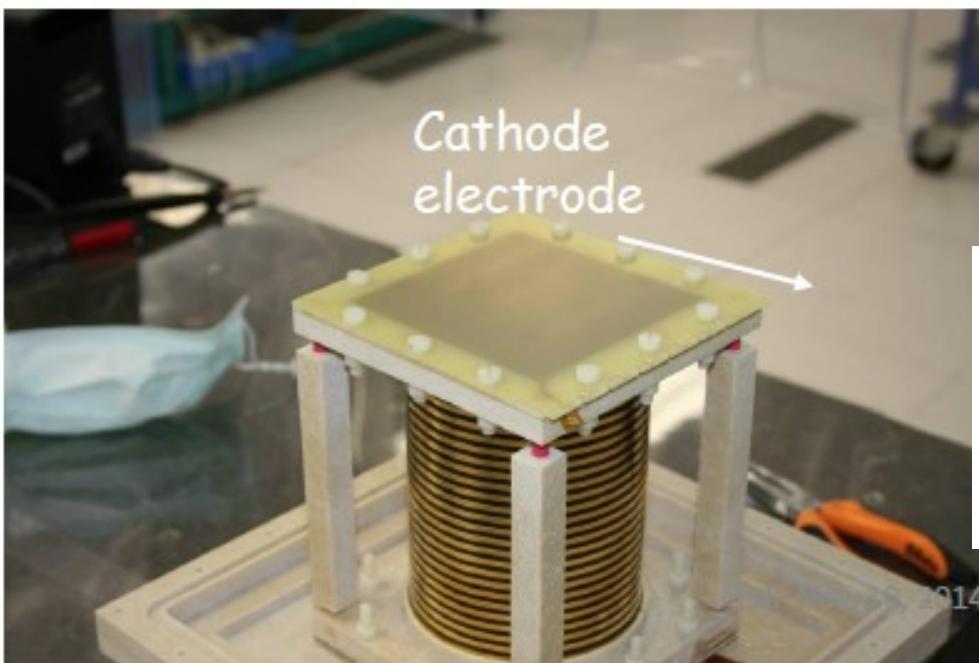
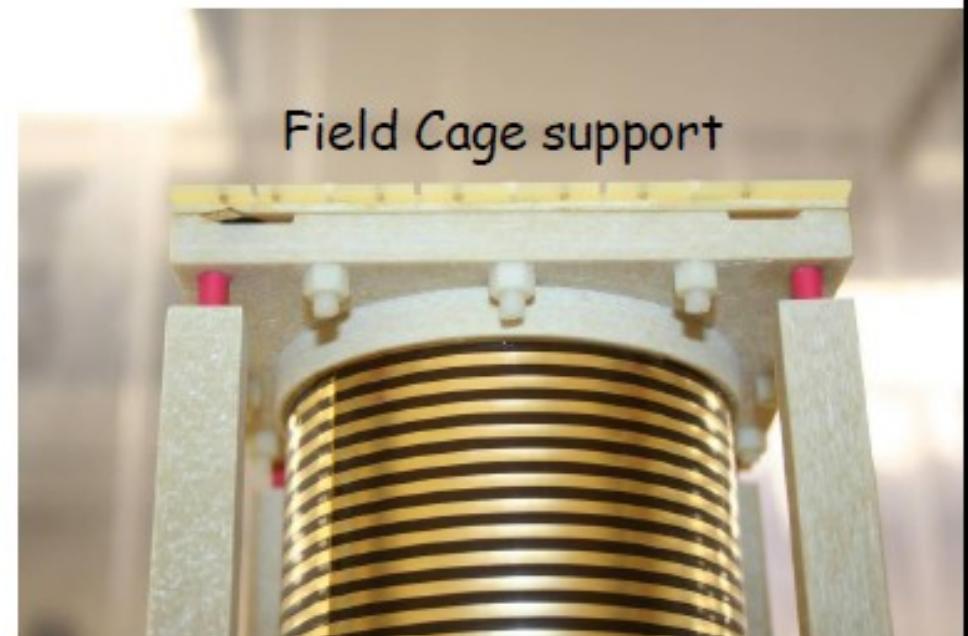
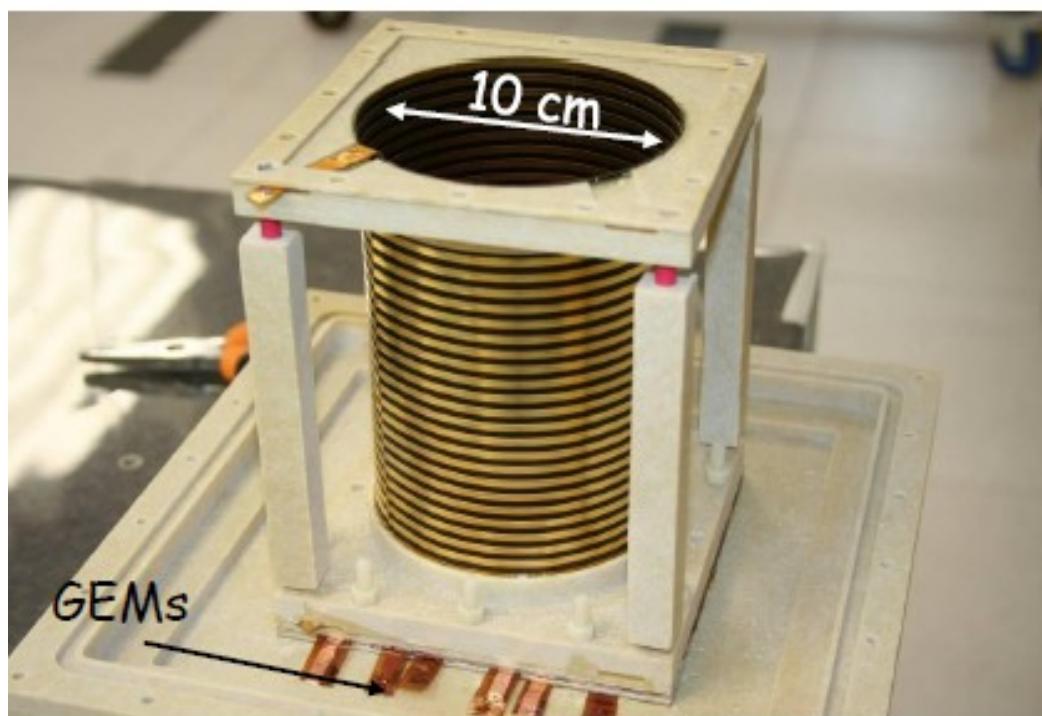
density ~ 5.0 g/cm³
total length of fibres ~ 15000 km
read out by ~ 5000 mesh PM

R&D – advanced setup

- active target TPC with GEM technology, with 6000 pads
 - R&D work within EU-FP7 HadronPhysics3



“active” TPC-GEM test setup at LNF



Tests at PSI

Modern Instrumentation, 2015, 4, 32-41
Published Online July 2015 in SciRes. <http://www.scirp.org/journal/mi>
<http://dx.doi.org/10.4236/mi.2015.43004>



Performances of an Active Target GEM-Based TPC for the AMADEUS Experiment

M.Poli Lener

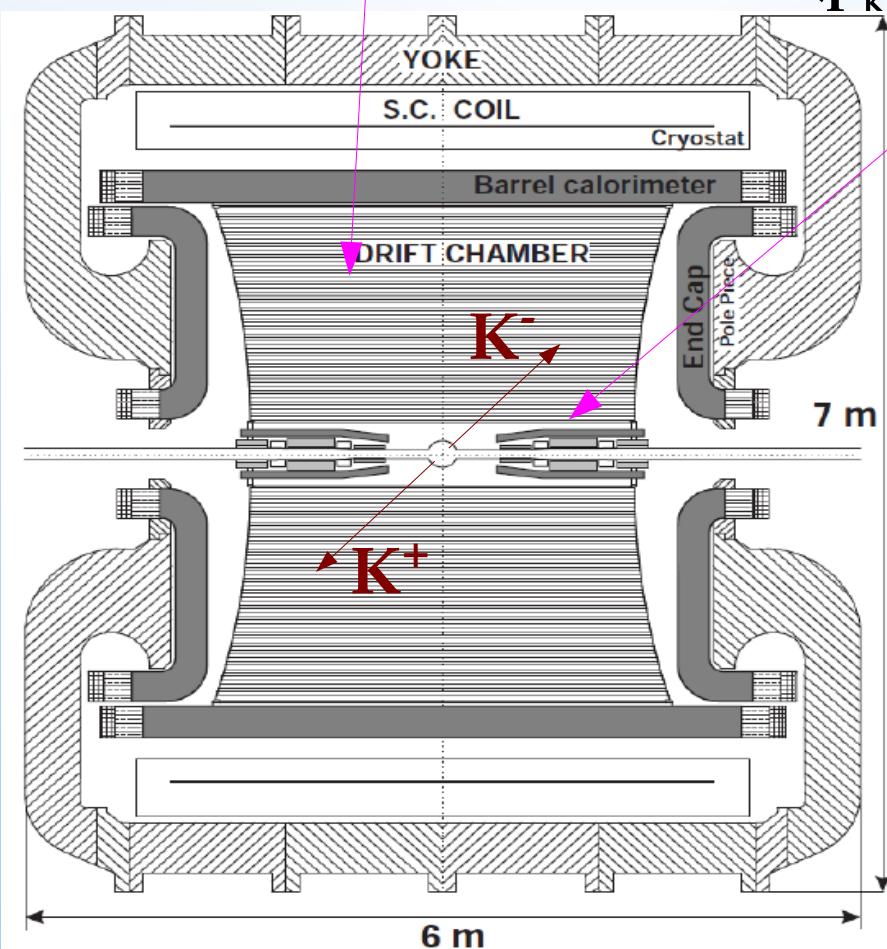
What to do meanwhile?

K⁻ absorption on light nuclei

from the materials of the KLOE detector

DC gas (90% He, 10% C₄H₁₀) & DC wall (C + H)

AT-REST (K⁻ absorbed from atomic orbit) or IN-FLIGHT
(p_K ~100MeV)



Advantage:

excellent resolution ..

$$\sigma_{p\Lambda} = 0.49 \pm 0.01 \text{ MeV/c in DC gas}$$

$$\sigma_{m\gamma\gamma} = 18.3 \pm 0.6 \text{ MeV/c}^2$$

Disadvantage:

Not dedicated target → different nuclei
contamination → complex interpretation .. but
→ new features .. K⁻ in flight absorption.

The scientific goal of AMADEUS

Low energy QCD in strangeness sector is still waiting for experimental conclusive constrains on:

1) **\bar{K} -N potential** → how deep can an antikaon be bound in a nucleus?

- U_{KN} strongly affects the position of the $\Lambda(1405)$ state → we investigate it through $(\Sigma-\pi)^0$ decay --- $Y\pi$ CORRELATION

- if U_{KN} is strongly attractive then $K^- NN$ bound states should appear → we investigate through $(\Lambda/\Sigma-N)$ decay --- YN CORRELATION

2) **$Y-N$ potential** → extremely poor experimental information from scattering data

- U_{YN} determines the strength of the final state YN (elastic & inelastic) scattering in nuclear environment → could be tested by YN CORRELATION

K⁻ - multi nucleon absorption and K⁻ pp bound state search

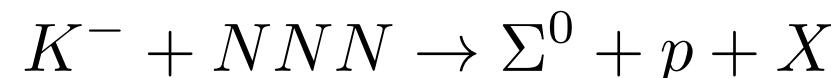
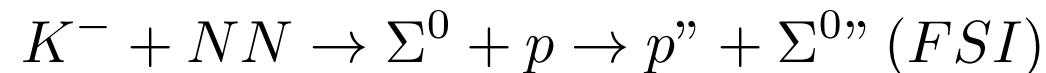
Σ^0 p correlated production, goals of this analysis

K- Absorption

- Pin down the contribution of the process:



with respect to processes as:



Kaonic Bound States



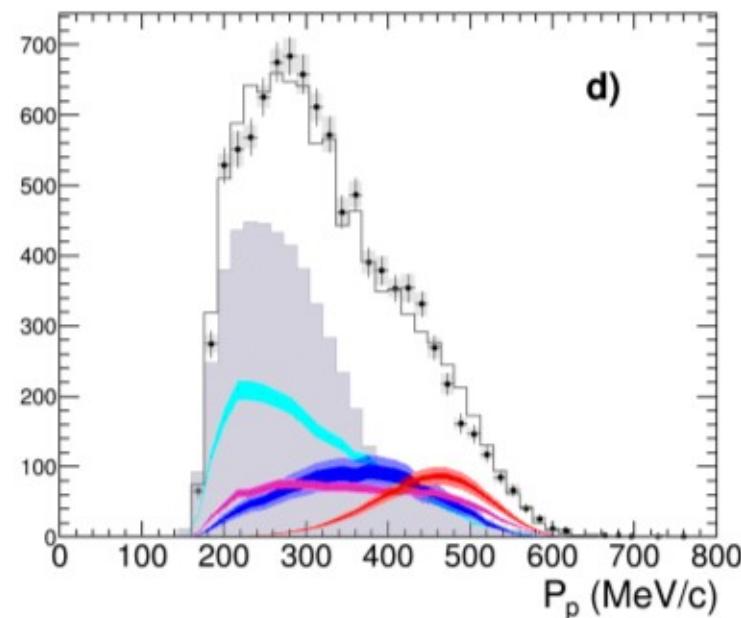
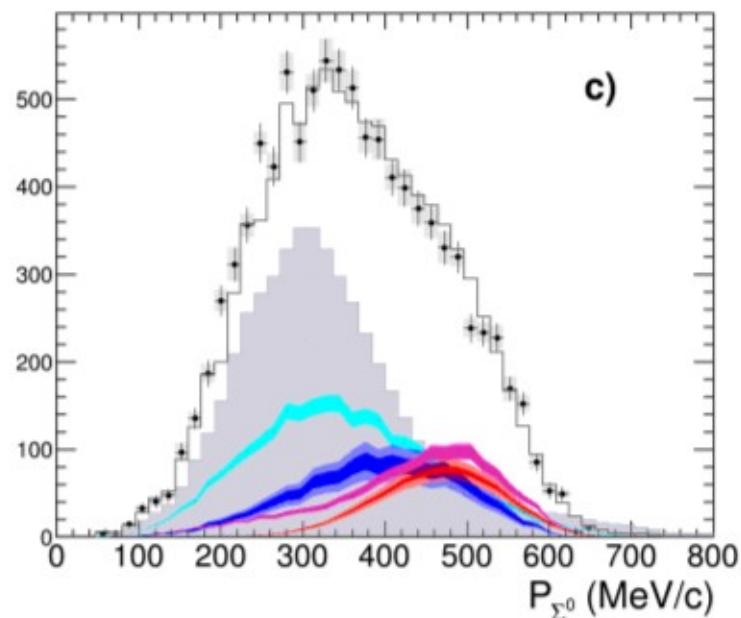
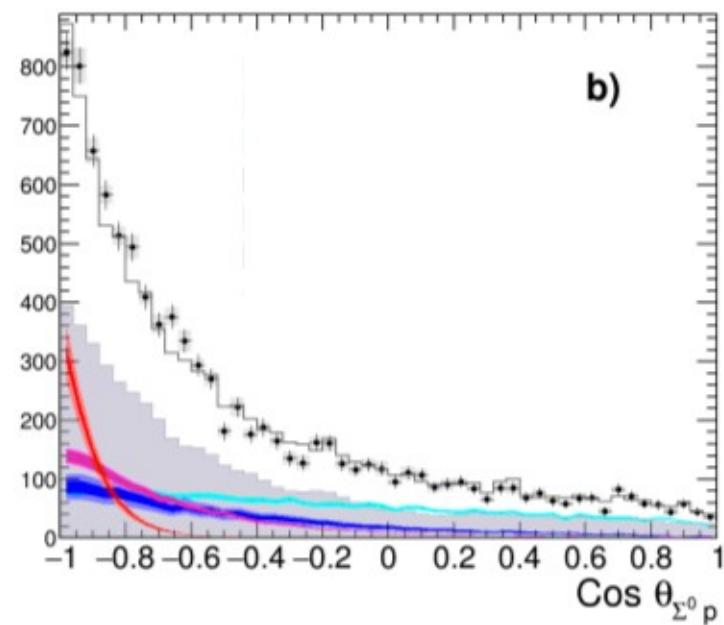
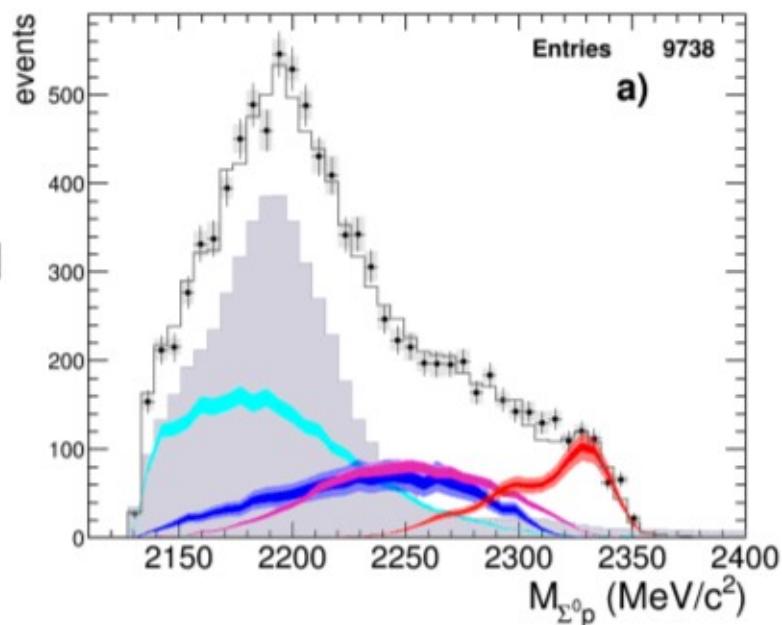
Yield Extraction and Significance

Final fit

- data
- π^0 background
- 4NA+Uncorr.
- 3NA
- 2NA FSI
- 2NA QF
- Total fit

$$\chi^2 = 0.85$$

2NA-QF clearly separated From other processes



From the contributions to the fit, the yields are extracted for K-stop

Absorption results

	yield / $K_{stop}^- \cdot 10^{-2}$	$\sigma_{stat} \cdot 10^{-2}$	$\sigma_{syst} \cdot 10^{-2}$
2NA-QF	0.127	± 0.019	$+0.004$ -0.008
2NA-FSI	0.272	± 0.028	$+0.022$ -0.023
Tot 2NA	0.376	± 0.033	$+0.023$ -0.032
3NA	0.274	± 0.069	$+0.044$ -0.021
Tot 3body	0.546	± 0.074	$+0.048$ -0.033
4NA + bkg.	0.773	± 0.053	$+0.025$ -0.076

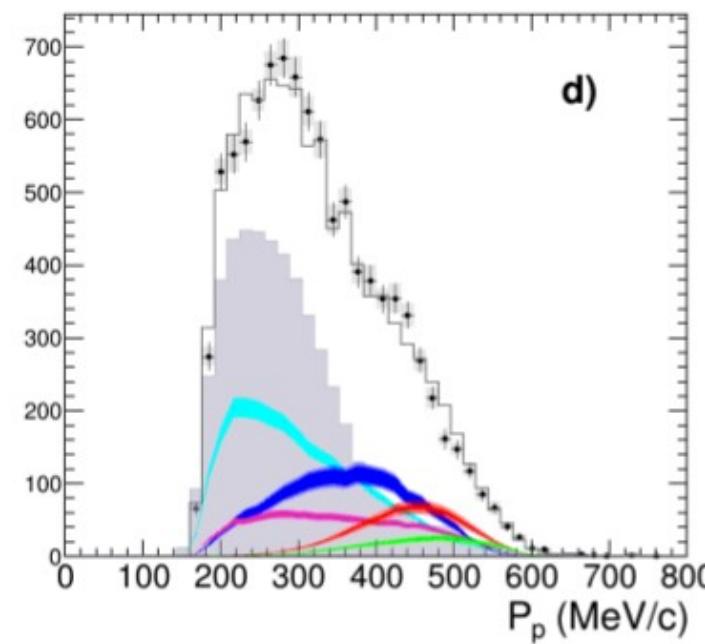
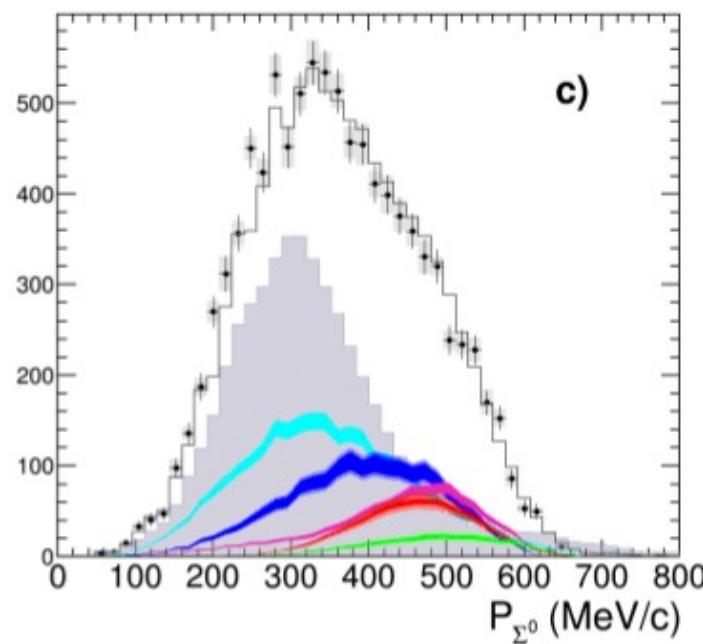
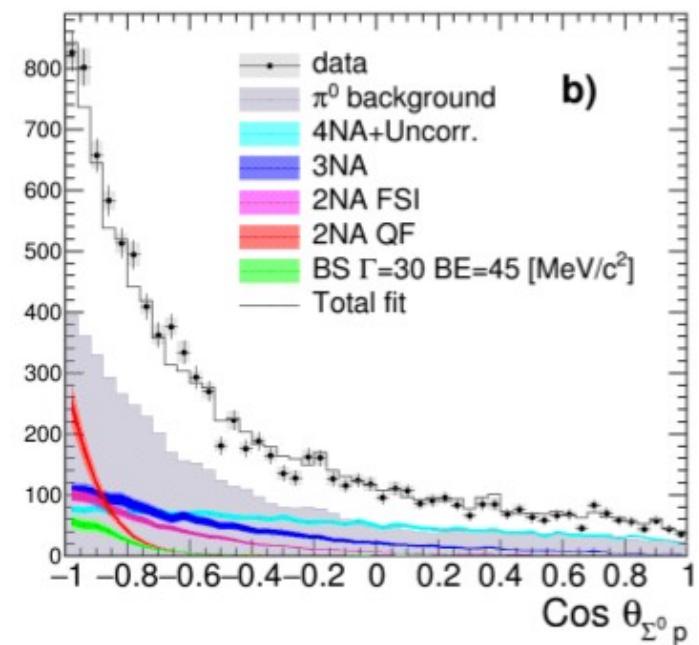
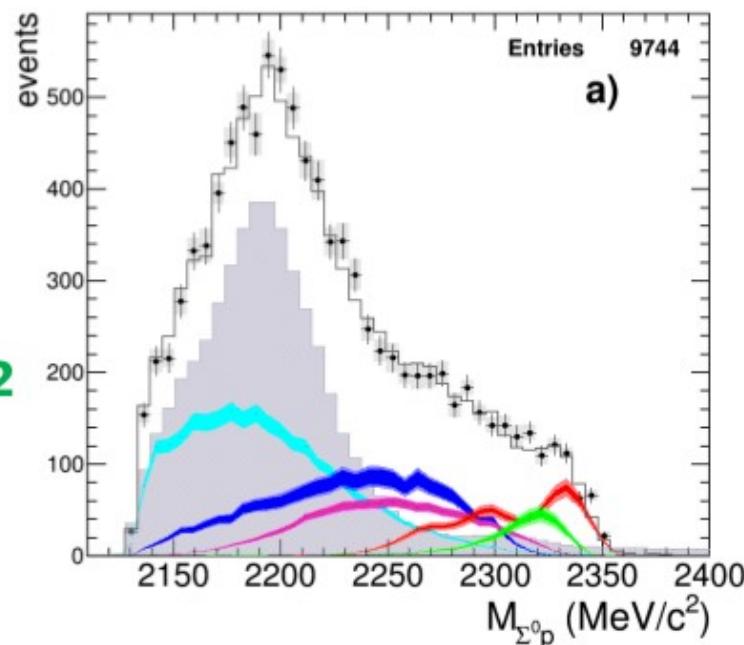
...is there room for the signal of a **ppK- bound state?**

Fit with ppK-

Best solution:
 (best χ^2 and higher yield)
- B.E. = 45 MeV/c²
- Width = 30 MeV/c²

$$\chi^2 = 0.807$$

- data
- π^0 background
- 4NA+Uncorr.
- 3NA
- 2NA FSI
- 2NA QF
- BS $\Gamma=30$ BE=45 [MeV/c²]
- Total fit



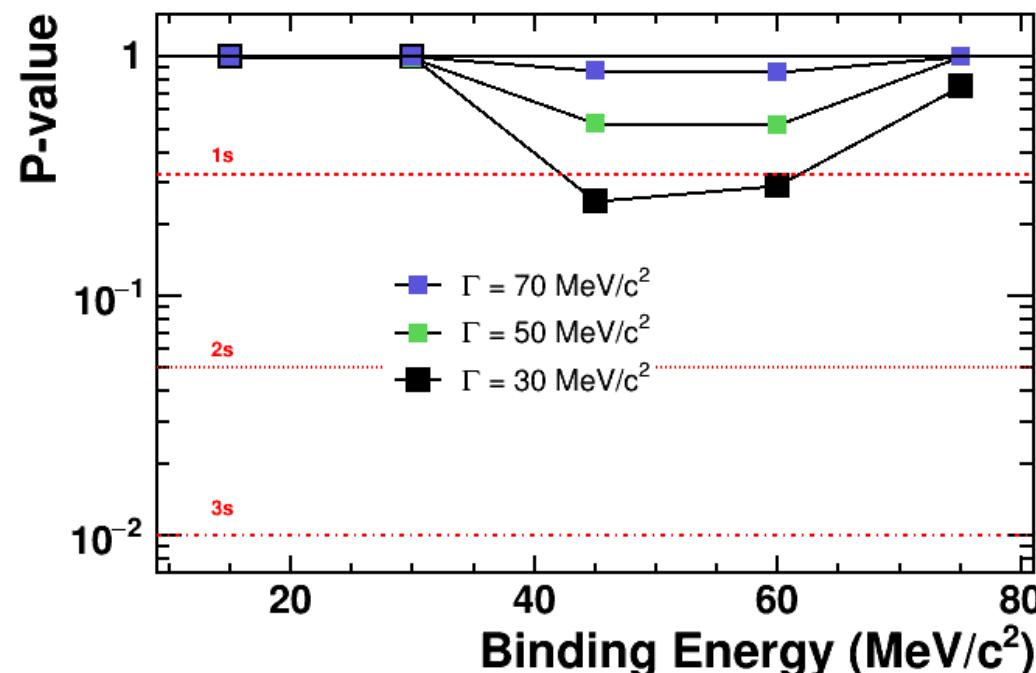
Evaluation of the significance of the ppK⁻ signal

For B.E. = 45 MeV/c², Width = 30 MeV/c²

$$Yield/K_{stop}^- = (0.044 \pm 0.009 stat)^{+0.004}_{-0.005 syst} \cdot 10^{-2}$$

F-test to evaluate the addition of an extra parameter to the fit:

Significance of “signal” hypothesis w.r.t
“Null-Hypothesis” (no bound state)



Conclusions

- 2NA-QF yield

	yield / $K_{stop}^- \cdot 10^{-2}$	$\sigma_{stat} \cdot 10^{-2}$	$\sigma_{syst} \cdot 10^{-2}$
2NA-QF	0.127	± 0.019	$^{+0.004}_{-0.008}$

- Bound state ppK- yield for B.E. 45 MeV/c2 and Width 30 MeV/c2

$$Yield/K_{stop}^- = (0.044 \pm 0.009 stat)^{+0.004}_{-0.005 syst} \cdot 10^{-2}$$

- the significance of the ppK- signal is of 1σ according to F-test

O. Vazquez Doce et al., Physics Letters B 758 (2016) 134



4NA cross section and yield

Λt available data

Available data:

- in Helium :

- bubble chamber experiment

[M.Roosen, J.H. Wickens, Il Nuovo Cimento 66, (1981), 101]

K^- stopped in liquid helium, Λ dn/t search. **3 events** compatible with the Λt kinematics were found

$$BR(K^{-4}\text{He} \rightarrow \Lambda t) = (3 \pm 2) \times 10^{-4} / K_{\text{stop}} \quad \underline{\text{global, no 4NA}}$$

- Solid targets

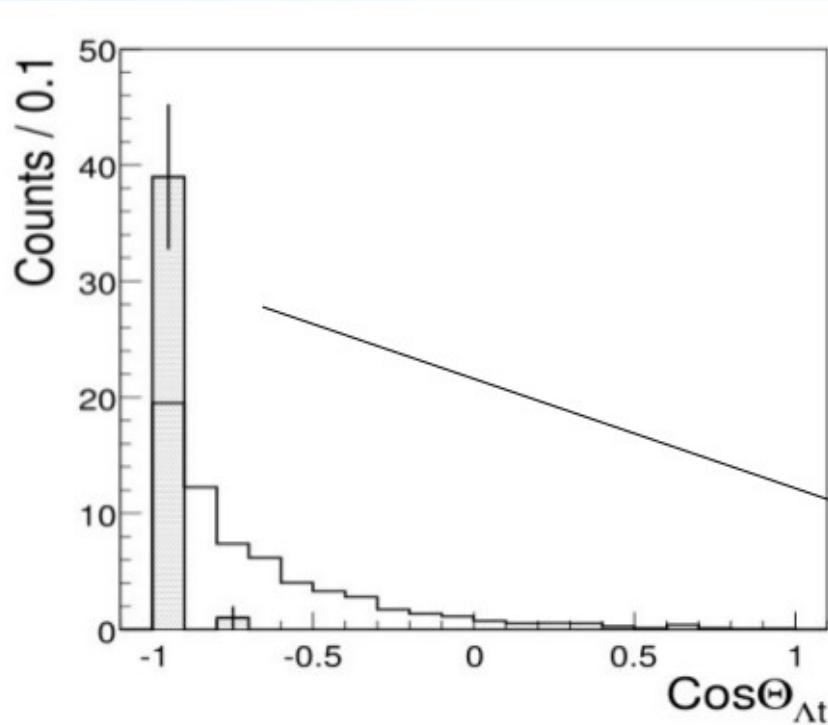
- FINUDA [Phys.Lett. B669 (2008) 229]
(**40 events** in different solid targets)

- T.Suzuki et al. Three- and four-nucleon absorption processes observed in the $K^{-4}\text{He}$ reaction at rest , arXiv:1009.5082v1 [nucl-ex] 26 Sep 2010 .

Λt available data

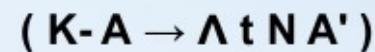
FINUDA presented [Phys.Lett.B (2008) 229]:

- a study of Λ vs t momentum correlation and an opening angle distribution
- **40 events** collected and added together coming from different targets ($^{6,7}\text{Li}$, ^9Be)



Filled histogram = data

Open histogram = Phase space simulation



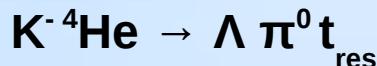
Unclear back to back topology

Λt emission yield $\rightarrow 10^{-3} - 10^{-4} / K_{\text{stop}}$
global, no 4NA

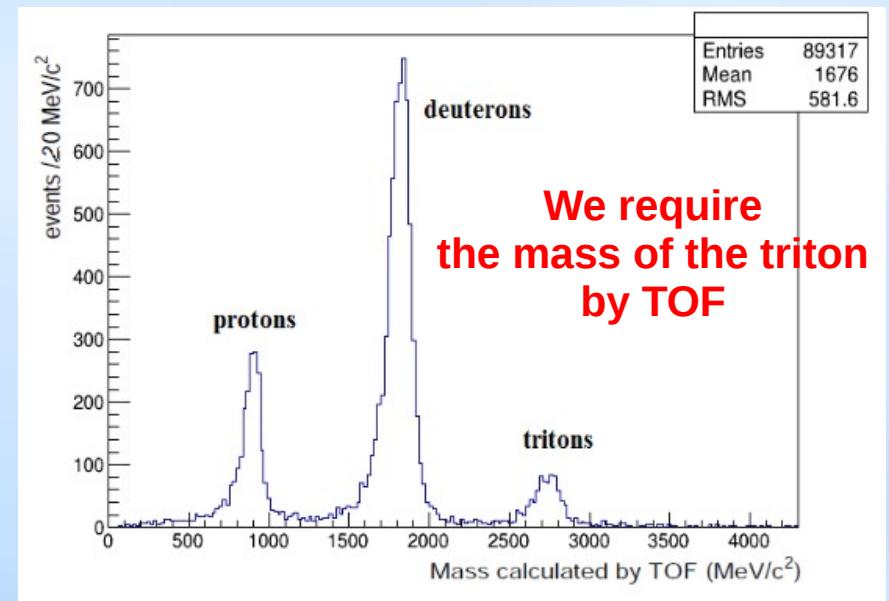
Experimental data only back-to-back

At correlation studies in ${}^4\text{He}$ from the DC gas : contributing processes

single nucleon absorption (1NA)



conversion on triton:



Tritons are spectators, **too low momentum**: $p_t \sim$ Fermi momentum

lower than the calorimeter threshold ($p_t \sim 500 \text{ MeV}/c$)

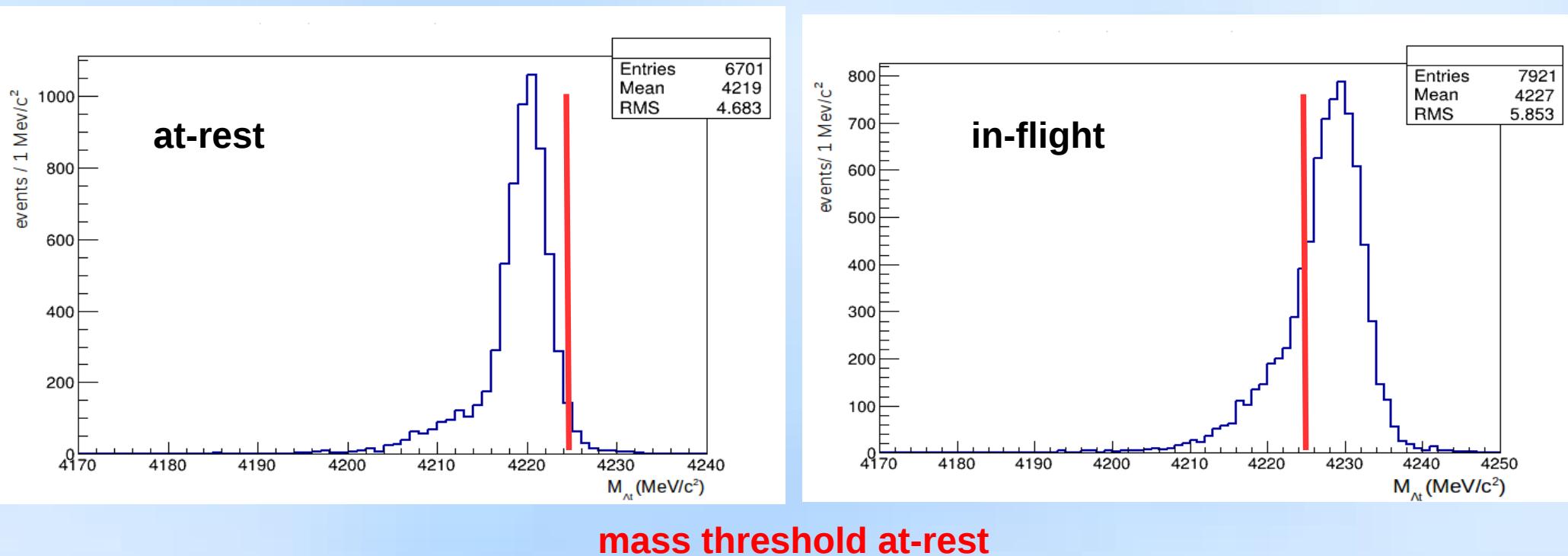
checked by MC simulations

4NA processes – K^- absorbed by the α particle:



conversion is suppressed
By the
 Σ^0 -t
Back to back topology!

MC simulations: efficiency & resolution



M_{Λ_t} invariant mass resolution = 2.2 MeV/c²

overall detection + reconstruction efficiency for 4NA direct Λ_t production :

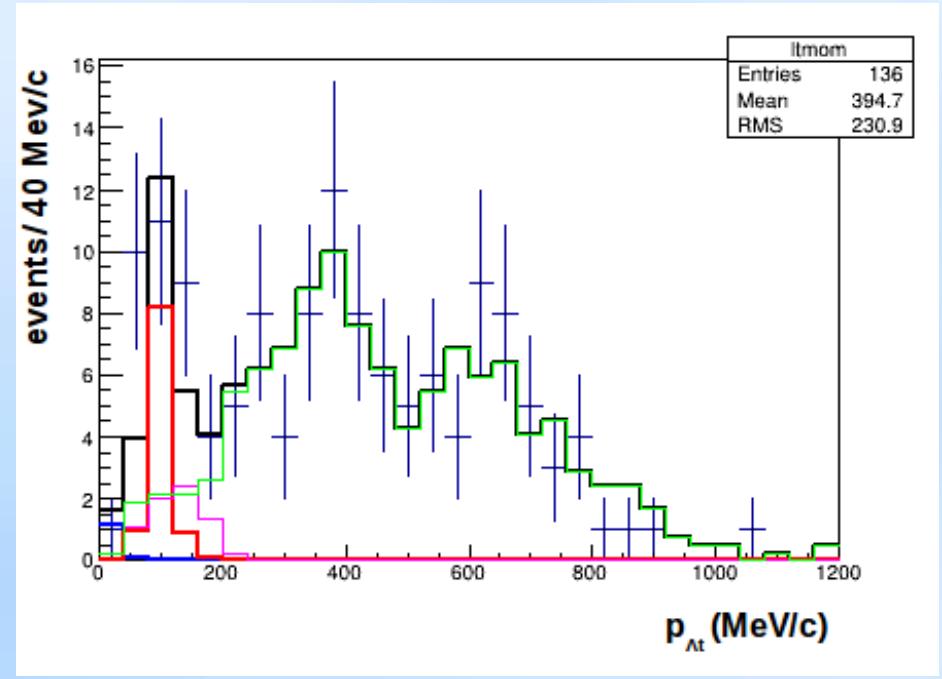
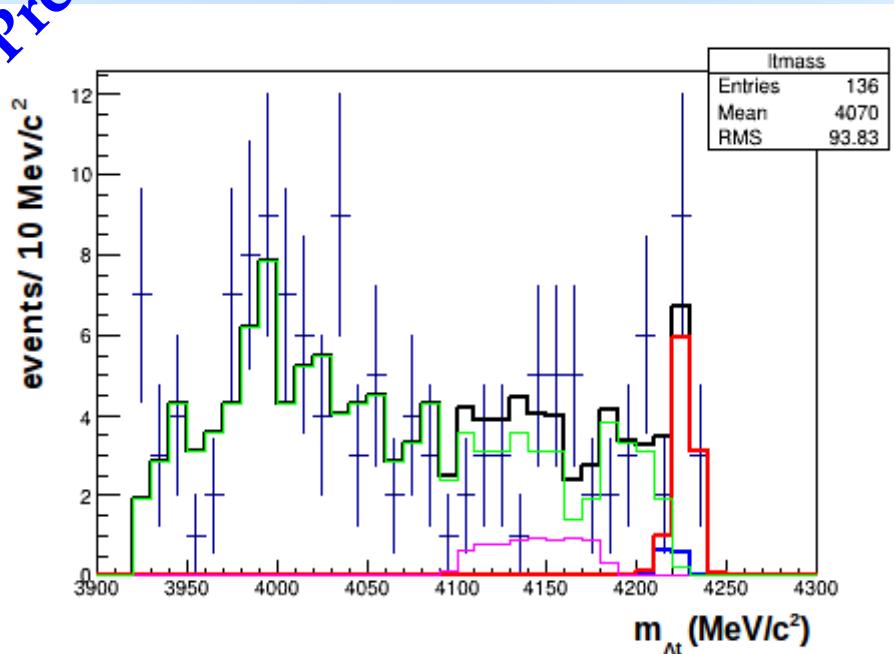
$$\epsilon_{4NA,ar,\Lambda_t} = 0.0493 \pm 0.0006 \quad ; \quad \epsilon_{4NA,if,\Lambda_t} = 0.0578 \pm 0.0006,$$

at-rest

in-flight

At correlation studies in ${}^4\text{He}$: mass, momentum and angle simultaneous fit

Preliminary



+ data

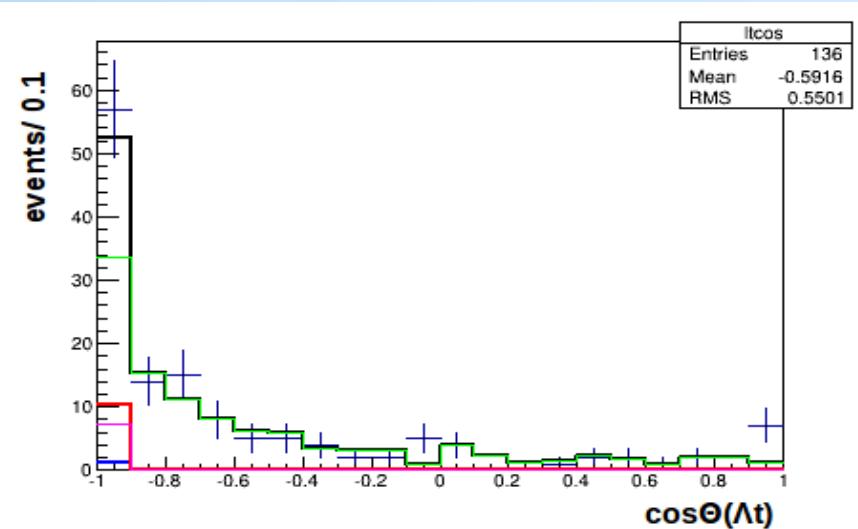
--- carbon data from DC wall

--- 4NA K⁻ ${}^4\text{He} \rightarrow \Lambda t$ in flight MC

--- 4NA K⁻ ${}^4\text{He} \rightarrow \Lambda t$ at rest MC

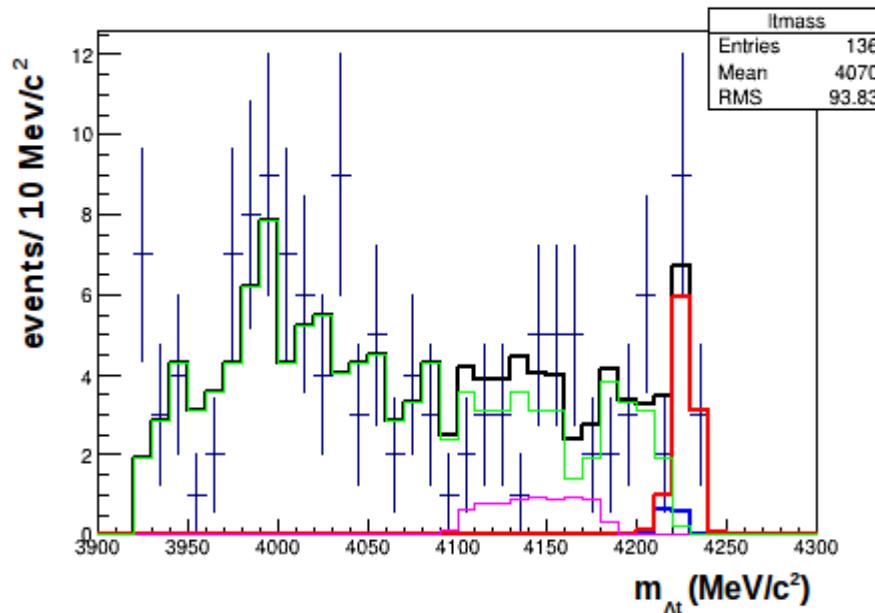
--- 4NA K⁻ ${}^4\text{He} \rightarrow \Sigma^0 t$, $\Sigma^0 \rightarrow \Lambda \gamma$ MC

--- 4NA K⁻ ${}^4\text{He} \rightarrow \Sigma^0 t$, $\Sigma^0 \rightarrow \Lambda \gamma$ MC



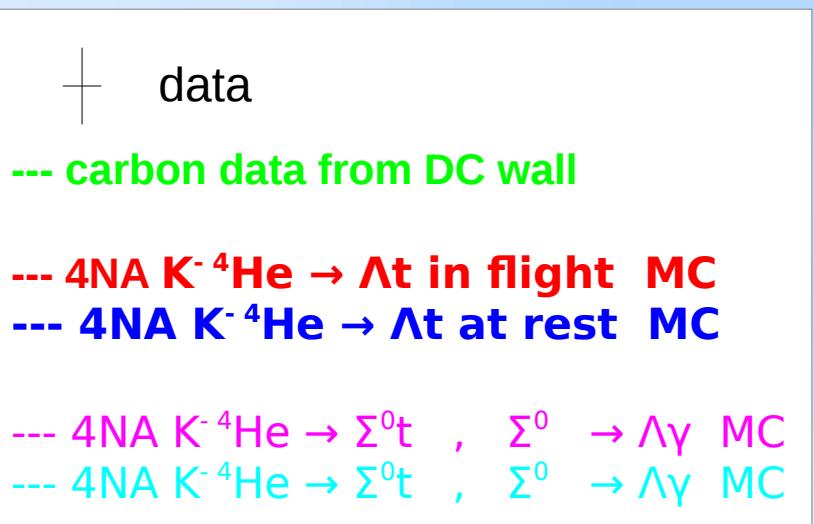
At correlation studies in ${}^4\text{He}$: preliminary mass and angle momentum simultaneous fit

Preliminary

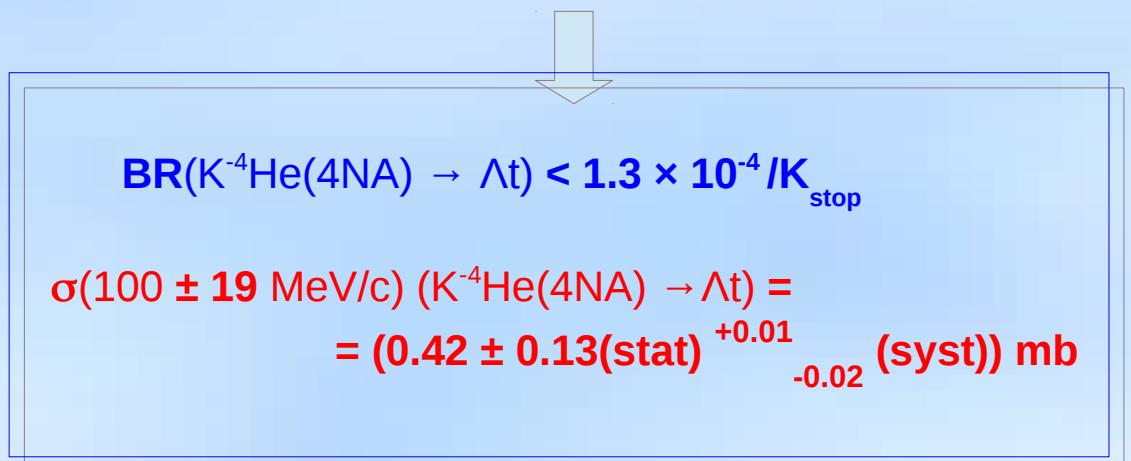


Contribution to the spectra	Parameter value
$K^- {}^4\text{He} \rightarrow \Lambda t$ at rest	0.01 ± 0.01
$K^- {}^4\text{He} \rightarrow \Lambda t$ in-flight	0.09 ± 0.02
$K^- {}^4\text{He} \rightarrow \Sigma^0 t$ in-flight	0.05 ± 0.03
$K^- {}^{12}\text{C} \rightarrow \Lambda t$ experimental distribution from the carbon DC wall	0.85 ± 0.06
χ^2 / ndf	0.654

parameters giving the contribution of the each process



Total number of events = 136
 4NA $K^- {}^4\text{He} \rightarrow \Lambda t$ at rest $\rightarrow 1 \pm 1$ events
 4NA $K^- {}^4\text{He} \rightarrow \Lambda t$ in flight $\rightarrow 12 \pm 3$ events



K⁻ - N single nucleon absorption processes

$\Lambda(1405)$ case

- Chiral unitary models: $\Lambda(1405)$ is an $I = 0$ quasibound state emerging from the coupling between the $\bar{K}\bar{N}$ and the $\Sigma\pi$ channels. Two poles in the neighborhood of the $\Lambda(1405)$:

two poles: about 1420 ; about = 1380) MeV

Phys. Lett. B 500 (2001), Phys. Rev. C 66 (2002), (Nucl. Phys.

A 725(2003) 181) .. many others .. (Nucl. Phys. A881, 98 (2012)) .. others

mainly coupled to $\bar{K}\bar{N}$

mainly coupled to $\Sigma\pi$

→ line-shape depends on production mechanism

- Akaishi-Esmaili-Yamazaki phenomenological potential

Phys. Lett. B 686 (2010) 23-28 Confirmation of single pole ansatz?

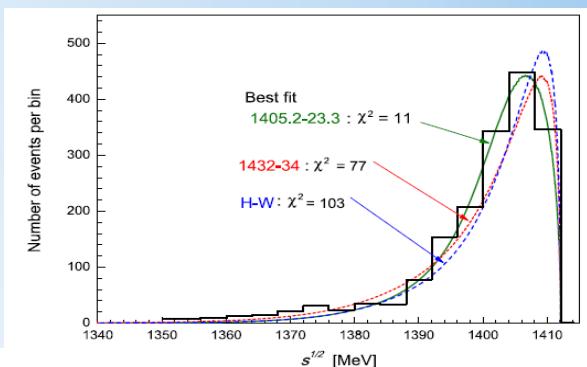
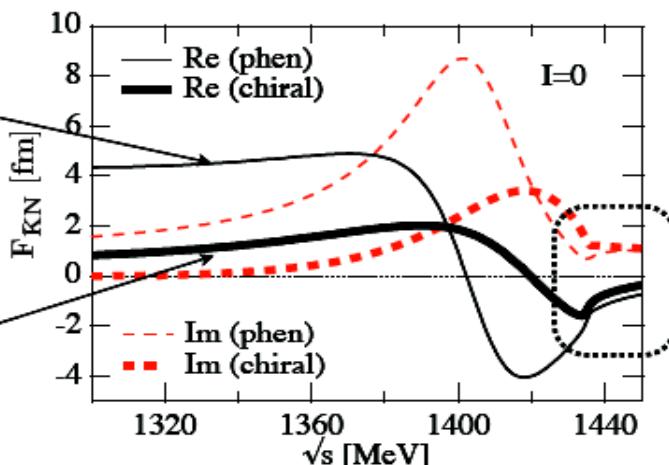


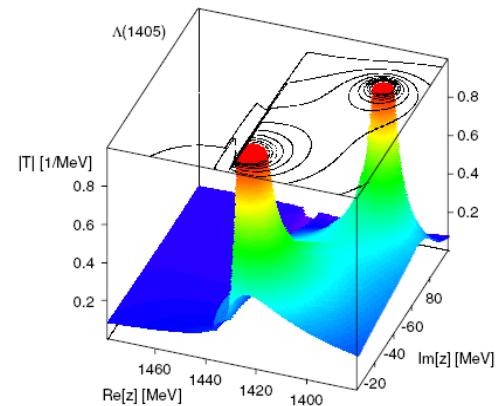
Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

AY
phenom.
potential

chiral
SU(3)
dynamics



large differences
in
subthreshold
extrapolations



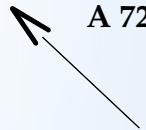
- Chiral dynamics predicts significantly weaker attraction than AY (local, energy independent) potential in far-subthreshold region

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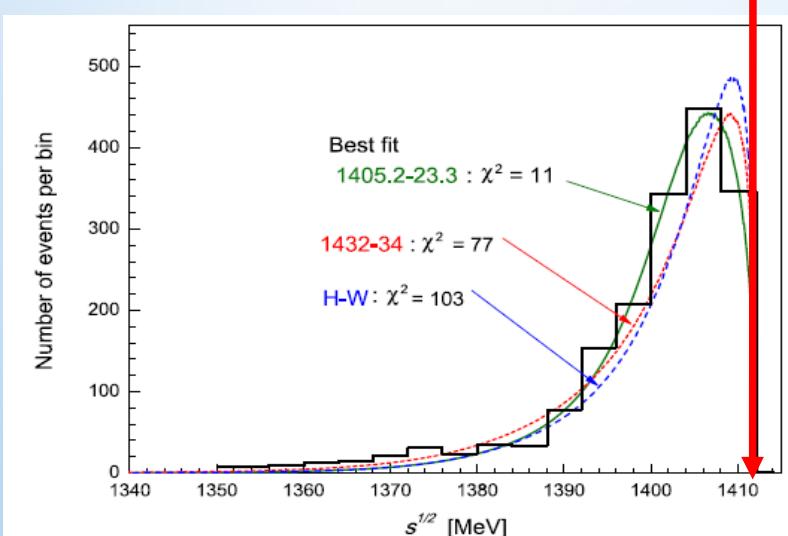
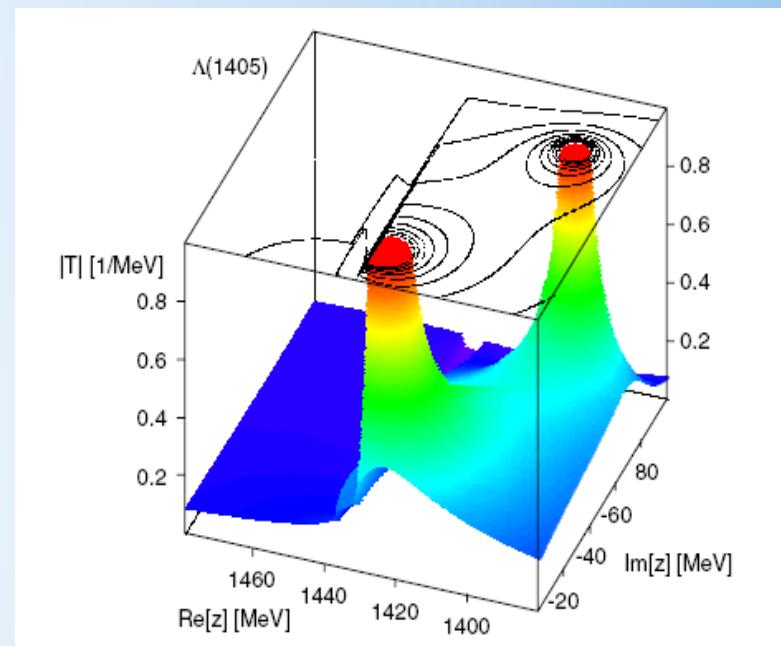


Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.



CUT AT THE ENERGY LIMIT AT-REST ?

NON RESONANT SHAPE ?

$\Lambda(1405)$ case

Phys.Rev.Lett.95:052301,2005

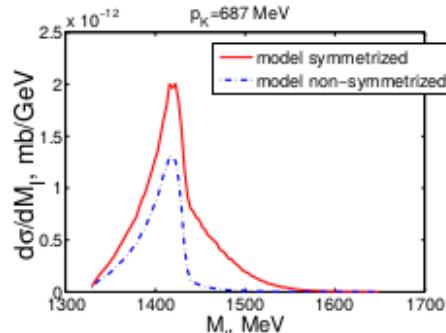


FIG. 4: Theoretical ($\pi^0 \Sigma^0$) invariant mass distribution for an initial kaon lab momenta of 687 MeV. The non-symmetrized distribution also contains the factor 1/2 in the cross section.

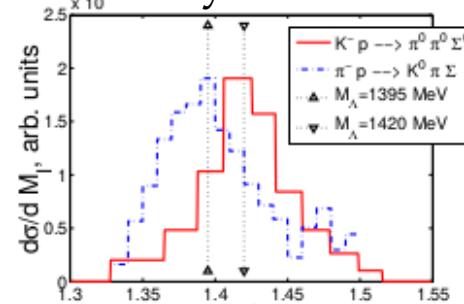
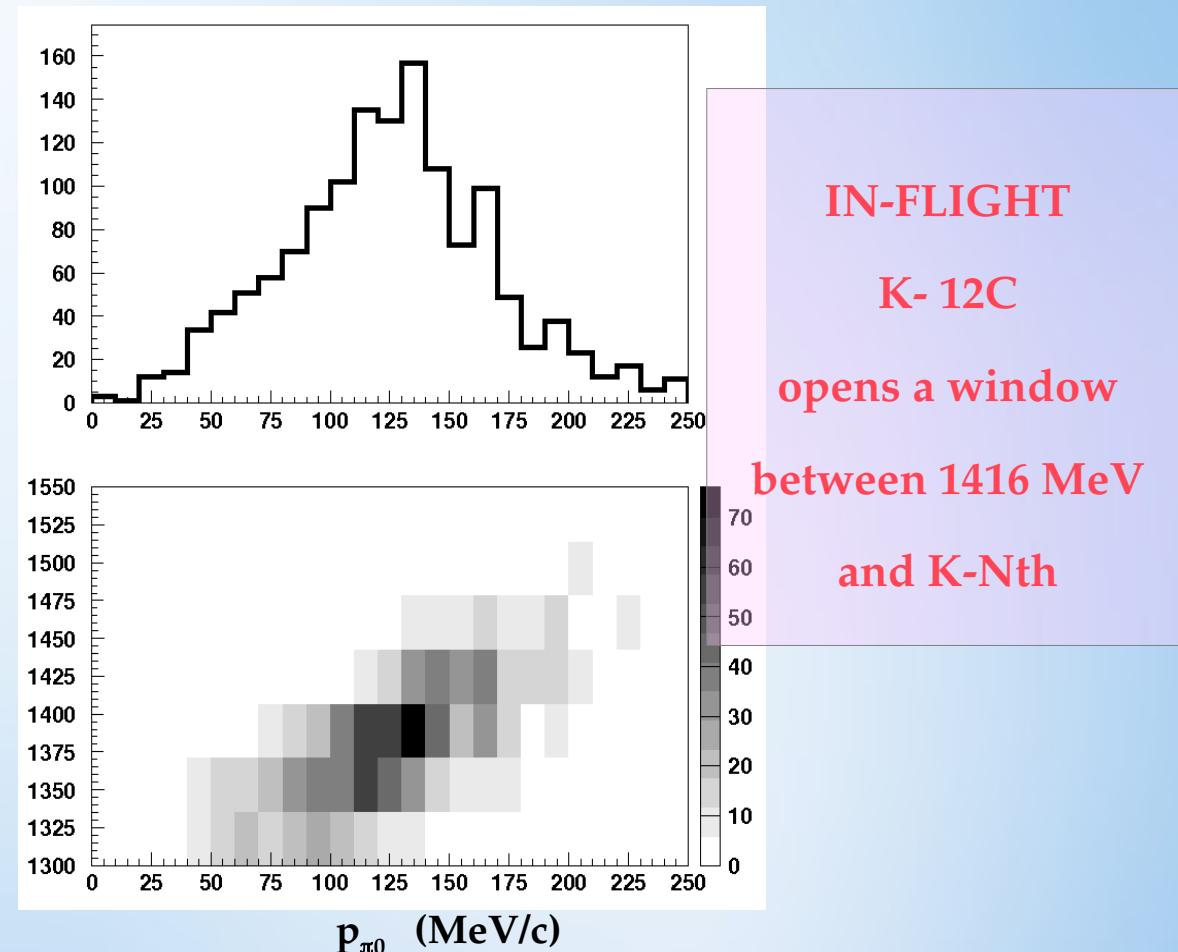
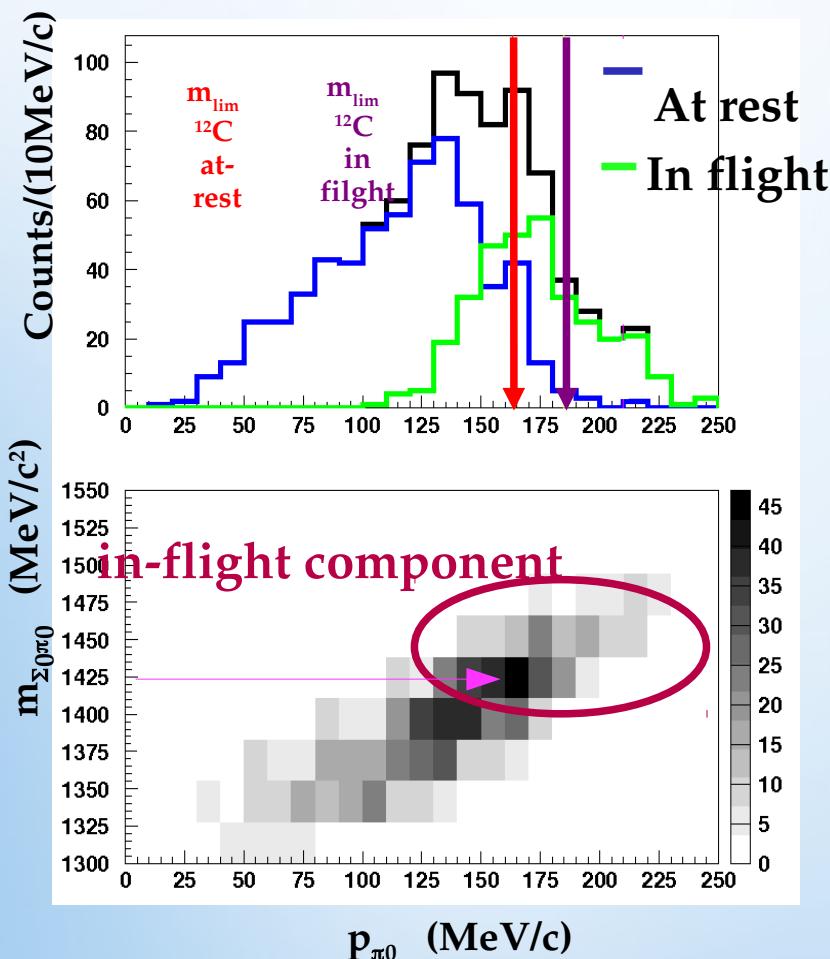


FIG. 5: Two experimental shapes of $\Lambda(1405)$ resonance. See text for more details.

p_{π^0} resolution: $\sigma_p \approx 12$ MeV/c



IN-FLIGHT

K- 12C

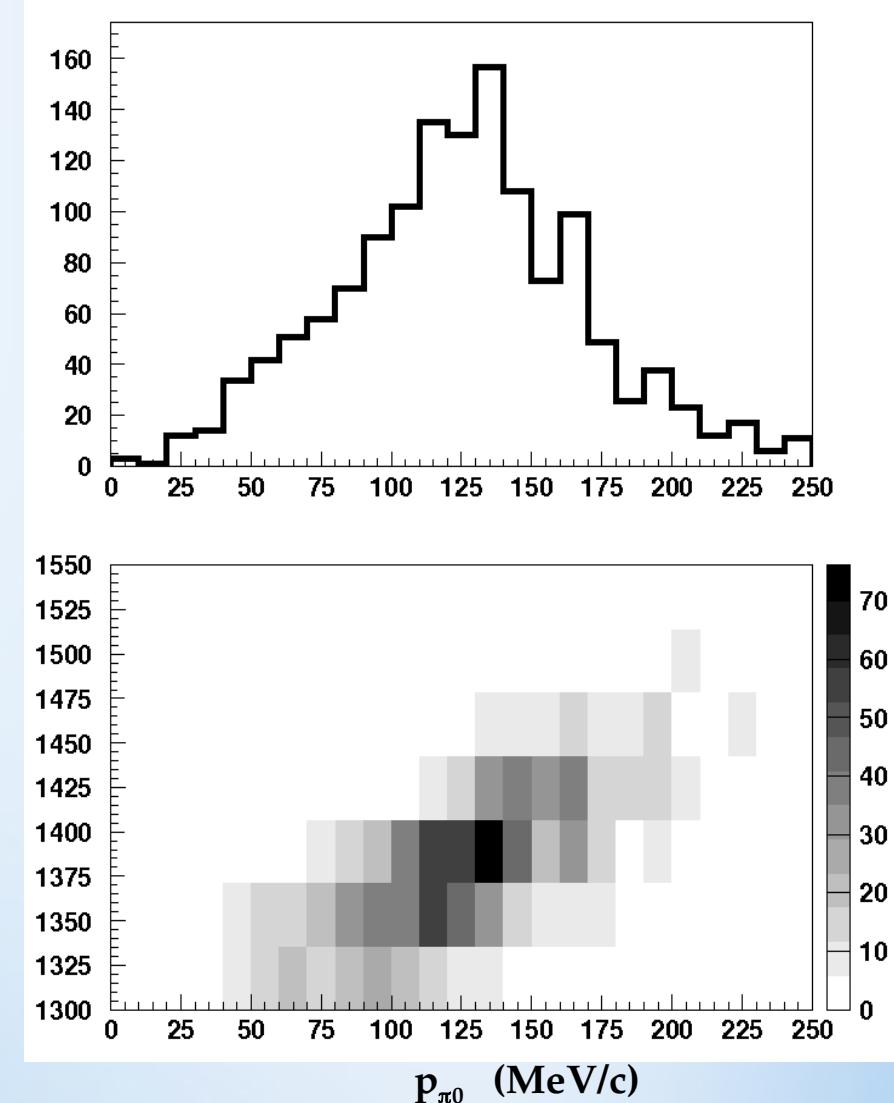
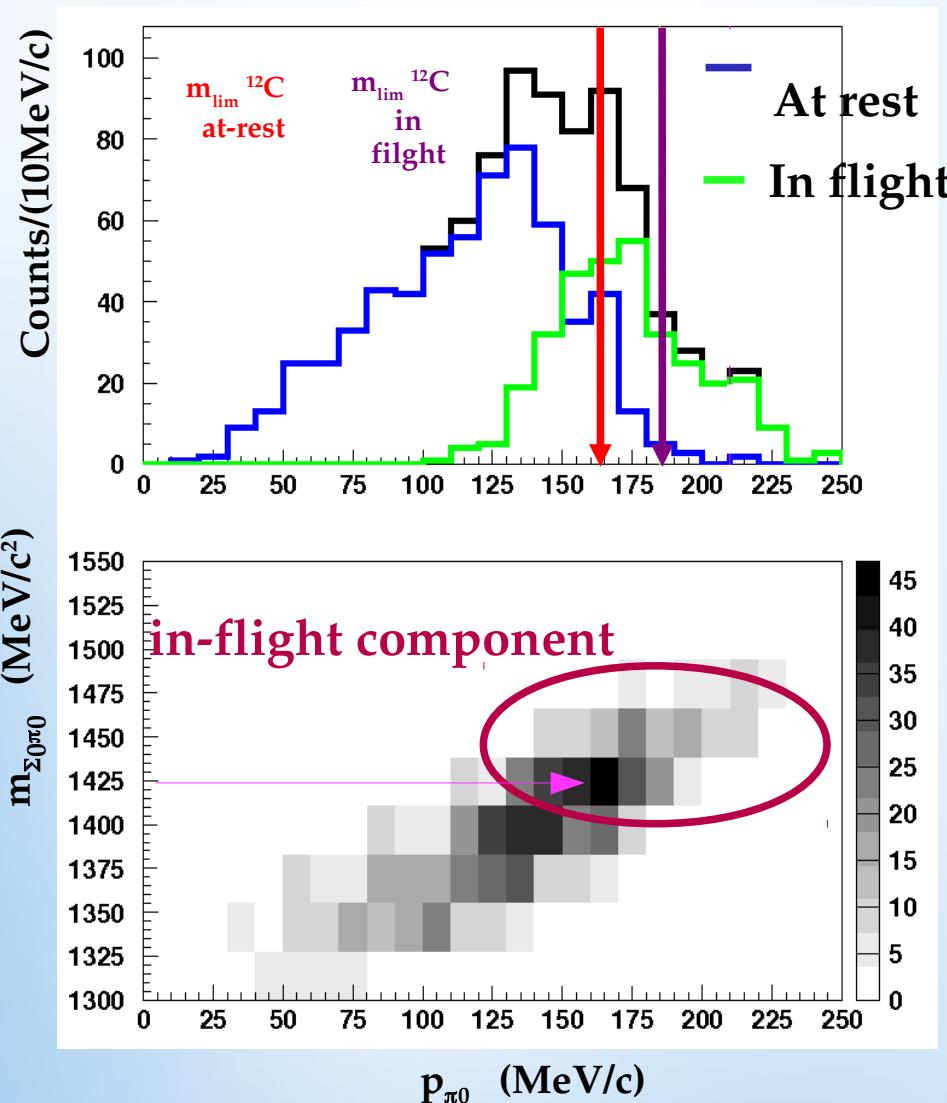
opens a window

between 1416 MeV

and K-Nth

Complex interpretation due to K- H absorptions ongoing with the collaboration of A. Cieply (UJF, Prague)

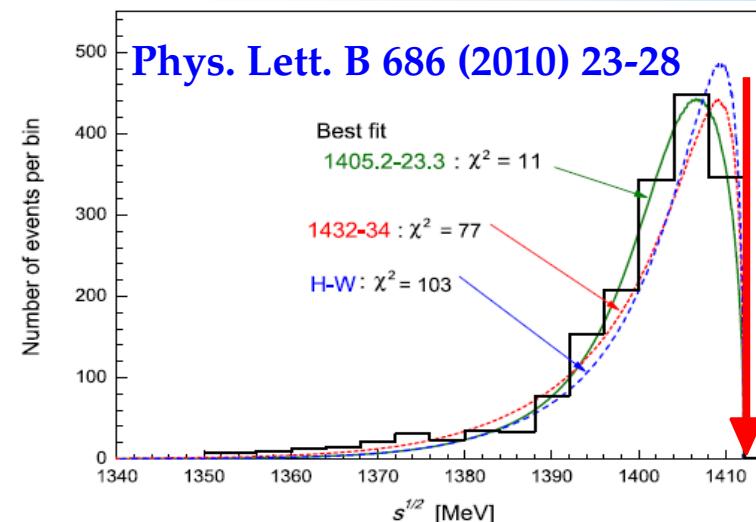
p_{π^0} resolution: $\sigma_p \approx 12 \text{ MeV/c}$



$\Sigma^+ \pi^-$ correlation

$K^- p \rightarrow \Sigma^+ \pi^-$ detected via: $(p\pi^0) \pi^-$

Possibility to disentangle: Hydrogen, in-flight, at-rest, K^- capture



p_{π^-} resolution: $\sigma_p \approx 1$ MeV/c

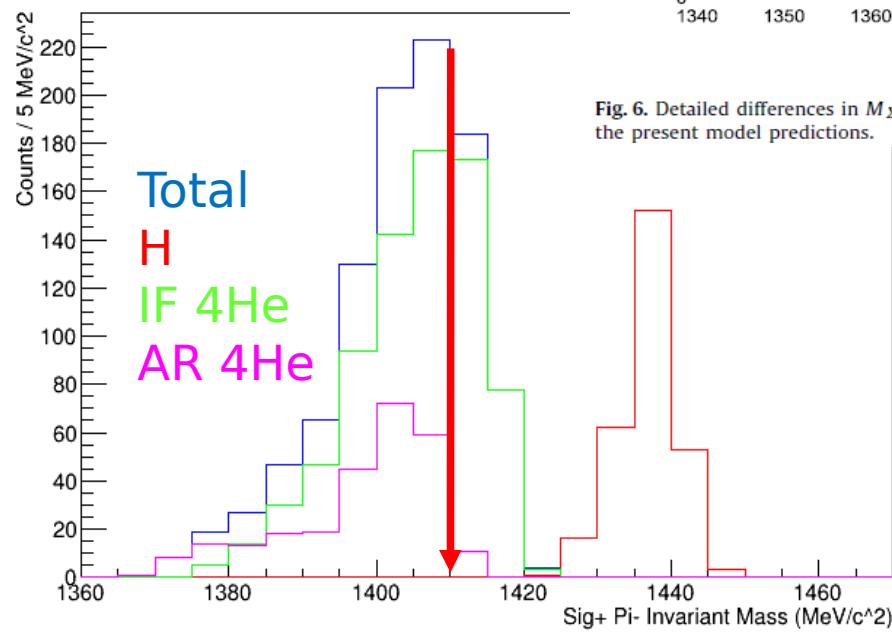
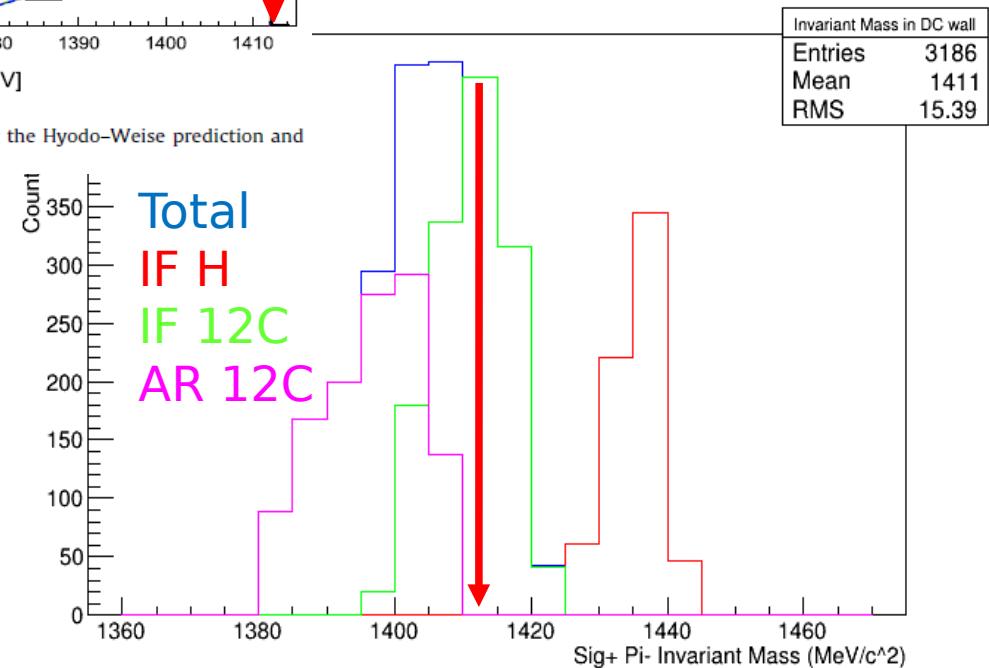


Fig. 6. Detailed differences in $M_{\Sigma\pi}$ spectra among the Hyodo-Weise prediction and the present model predictions.

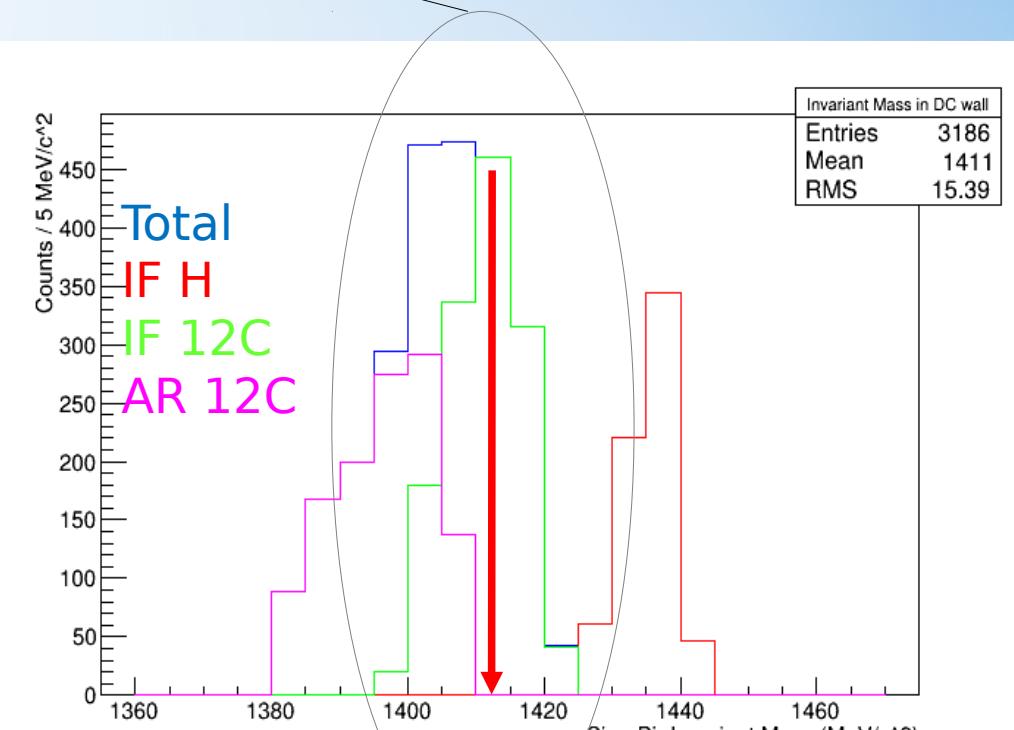
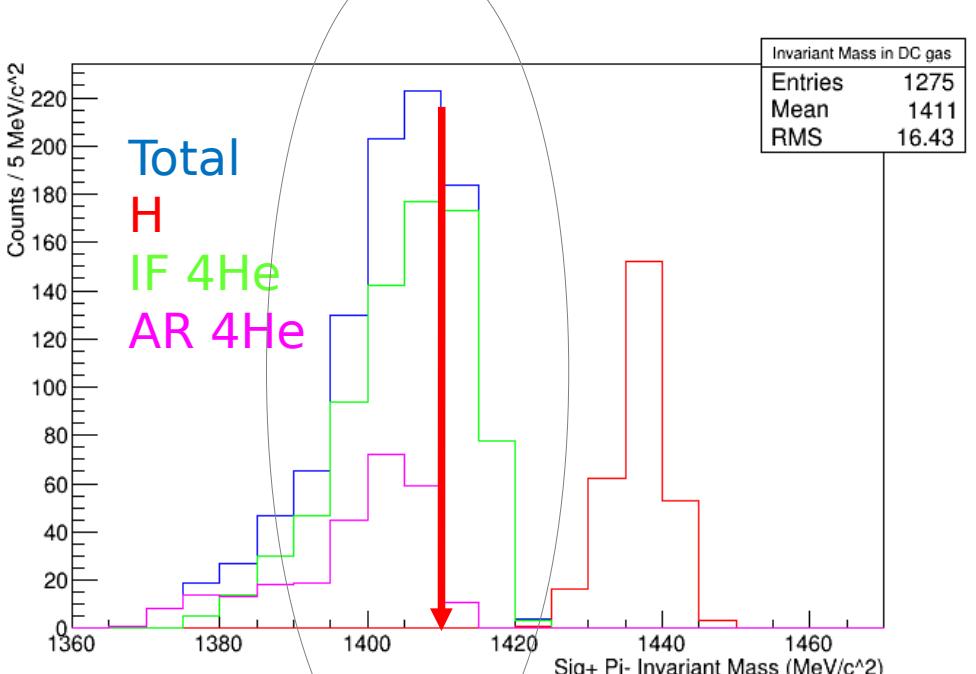


$\Sigma^+ \pi^-$ correlation

$K^- p \rightarrow \Sigma^+ \pi^-$ detected via: $(p\pi^0) \pi^-$

Possibility to disentangle: Hydrogen, in-flight, at-rest, K^- capture

if resonant production contribution is important a high mass component appears!



Resonant VS non-resonant



in medium, how much comes from resonance ?

Non resonant transition amplitude:

- Never measured before below threshold

(33 MeV below threshold kinetic energy in the Kn CM system):

$$E_{Kn} = -|B_n| - \frac{p_3^2}{2\mu_{\pi,\Lambda,3He}},$$

- few, old theoretical calculations
(Nucl. Phys. B179 (1981) 33-48)

Resonant VS non-resonant

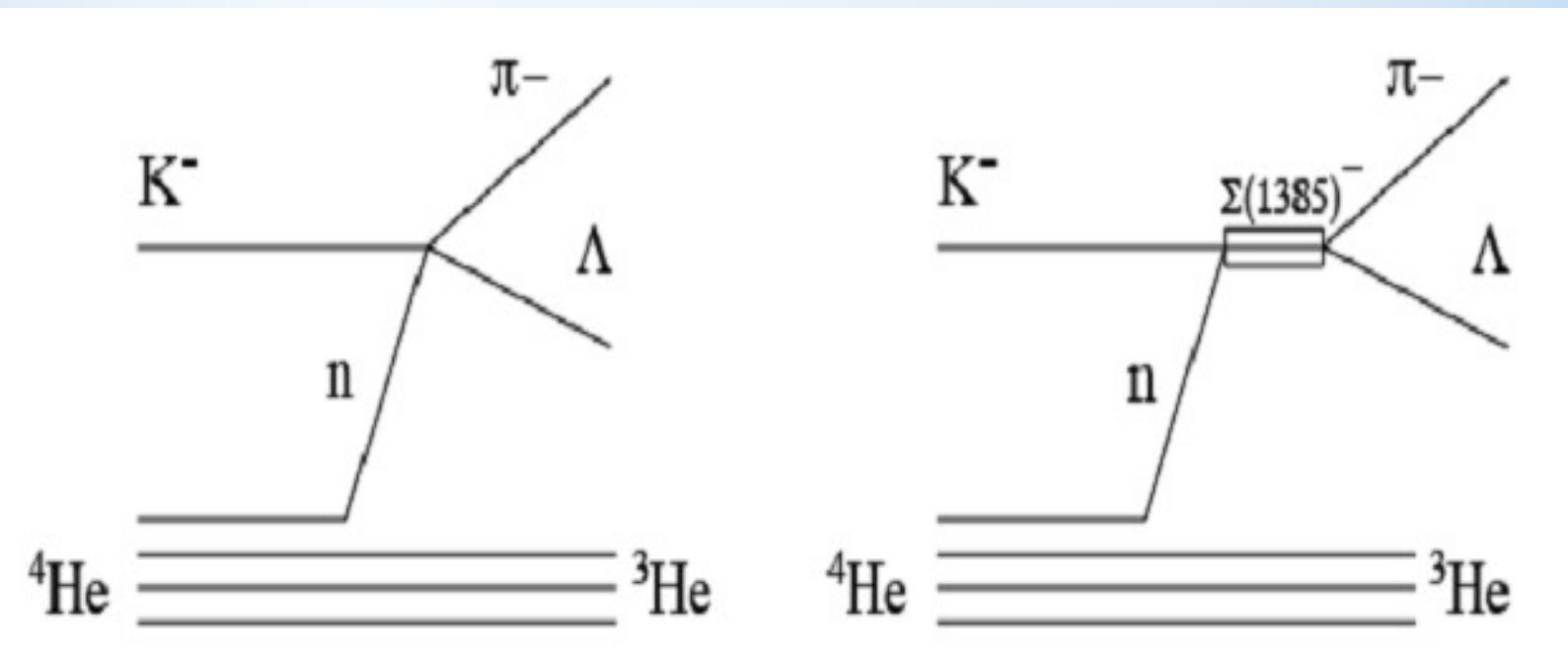
Investigated using:



direct formation in ${}^4\text{He}$

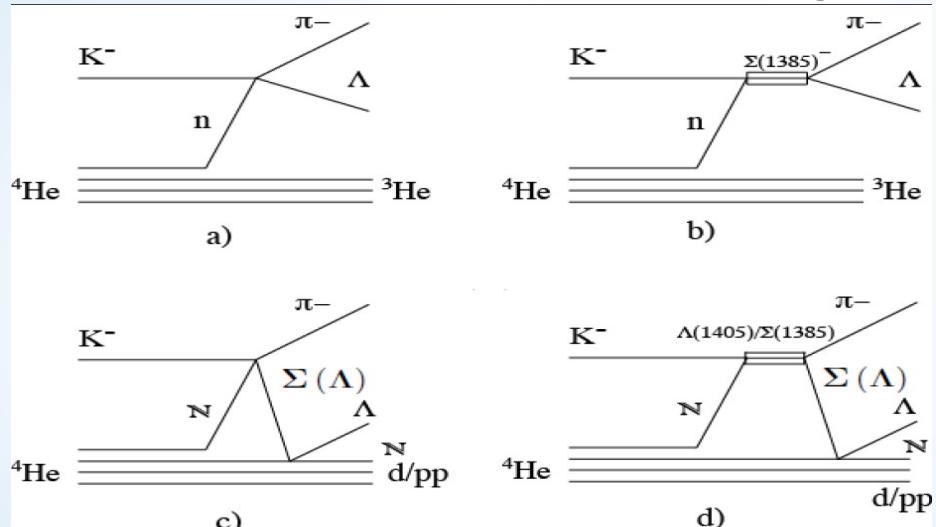
the goal is to measure $|f^{N-R}_{\Lambda\pi^-}(I=1)|$

to get information on $|f^{N-R}_{\Sigma\pi^-}(I=0)|$



$K^- {}^4He \rightarrow \Lambda p^- {}^3He$ resonant and non-resonant processes

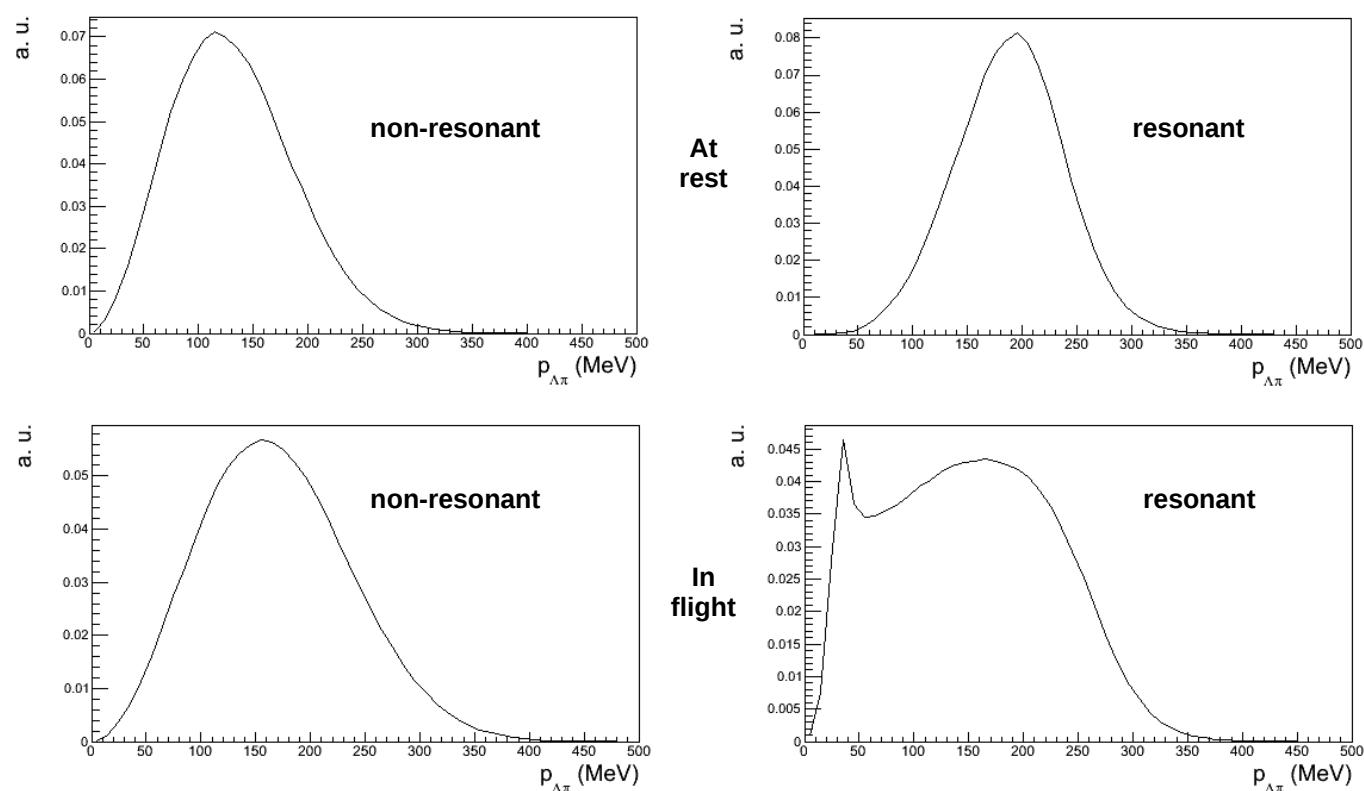
Nucl. Phys. A954 (2016) 75-93



Collaboration with
S. Wycech

Theoretical shapes for :

total $\Lambda\pi^-$ momentum spectra for the resonant (Σ^*) and non-resonant ($I = 1$) processes were calculated, for both S-state and P-state K^- capture at-rest and in-flight. Corrections to the amplitudes due to Λ/π final state interactions were estimated.



How to extract the $K^- n \rightarrow \Lambda\pi^-$ non resonant transition amplitude

simultaneous fit ($p_{\Lambda\pi^-} - m_{\Lambda\pi^-} - \cos(\theta_{\Lambda\pi^-})$) with signal  and background  processes :

- non resonant K^- capture at-rest from S states in ${}^4\text{He}$
- resonant K^- capture at-rest from S states in ${}^4\text{He}$
- non resonant K^- capture in-flight in ${}^4\text{He}$
- resonant K^- capture in-flight in ${}^4\text{He}$

- primary $\Sigma\pi^-$ production followed by the $\Sigma N \rightarrow \Lambda N'$ conversion process
- K^- capture processes in ${}^{12}\text{C}$ giving rise to $\Lambda\pi^-$ in the final state

In order to extract:

NR-ar/RES-ar

&

NR-if/RES-if

Results for the $K^- n \rightarrow \Lambda\pi^-$ non resonant transition amplitude

Preliminary

	percentage $\cdot 10^{-2}$	$\sigma_{stat} \cdot 10^{-2}$	$\sigma_{syst} \cdot 10^{-2}$
NR-ar	12.00	± 1.66	+1.96 -2.77
RES-ar/NR-ar	0.39	± 0.04	+0.18 -0.07
NR-if	19.24	± 4.38	+5.90 -3.33
RES-if//NR-if	0.23	± 0.03	+0.23 -0.22
$\Sigma \rightarrow \Lambda$ conv.	2.16	± 0.30	+1.62 -0.83
$K^- {}^{12}\text{C}$ capture	57.00	± 1.23	+2.21 -3.19

Table 1. Percentages of the different physical contributions to the selected $\Lambda\pi^-$ sample, from top: non-resonant at-rest, resonant at-rest, non-resonant in-flight, resonant in-flight, $\Sigma N \rightarrow \Lambda N'$ conversion processes, K^- absorptions in Carbon. The statistical and systematic errors are shown as well.

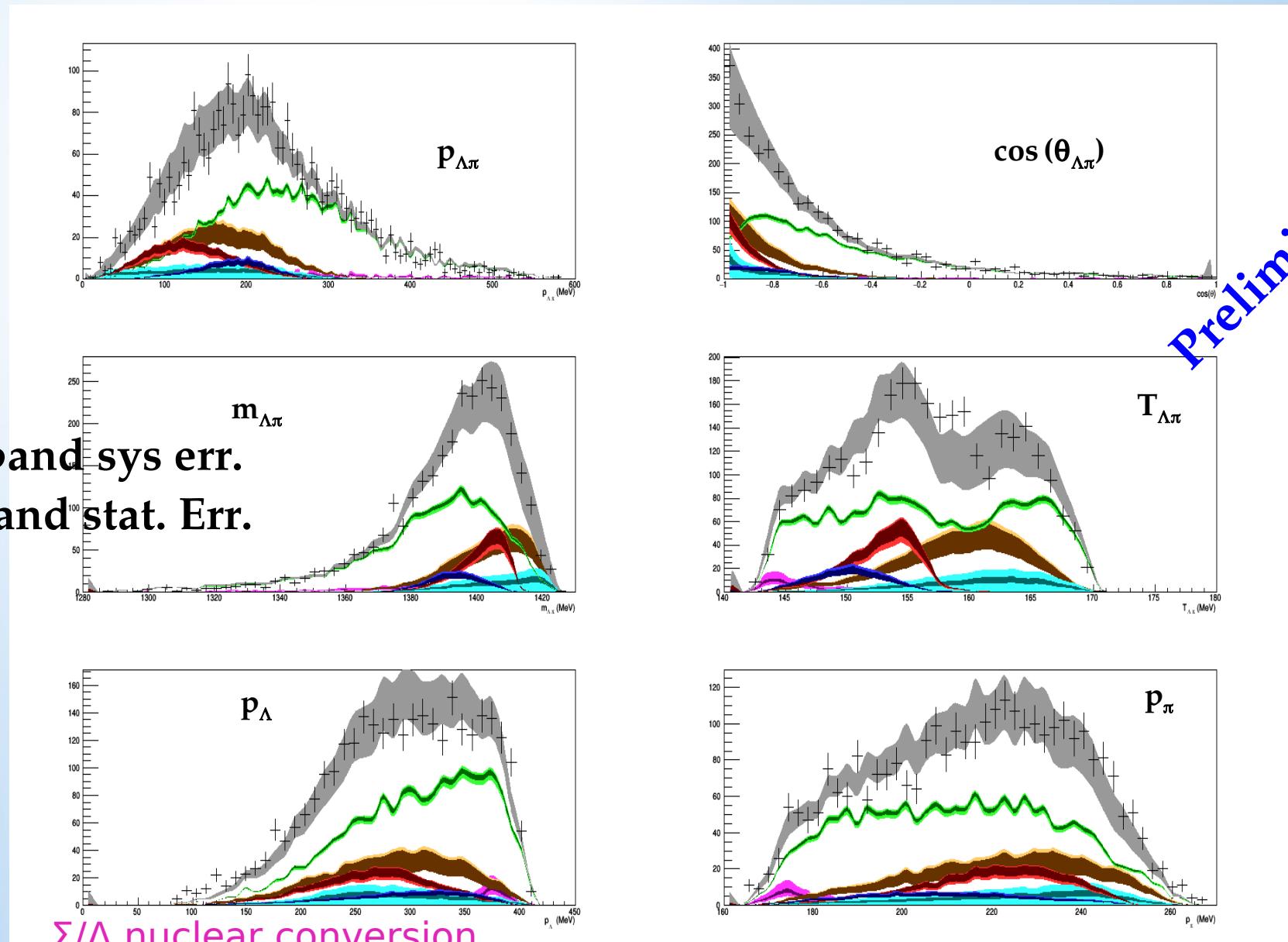
extracted:

NR-ar/RES-ar

&

NR-if/RES-if

Simultaneous momentum – angle – mass fit

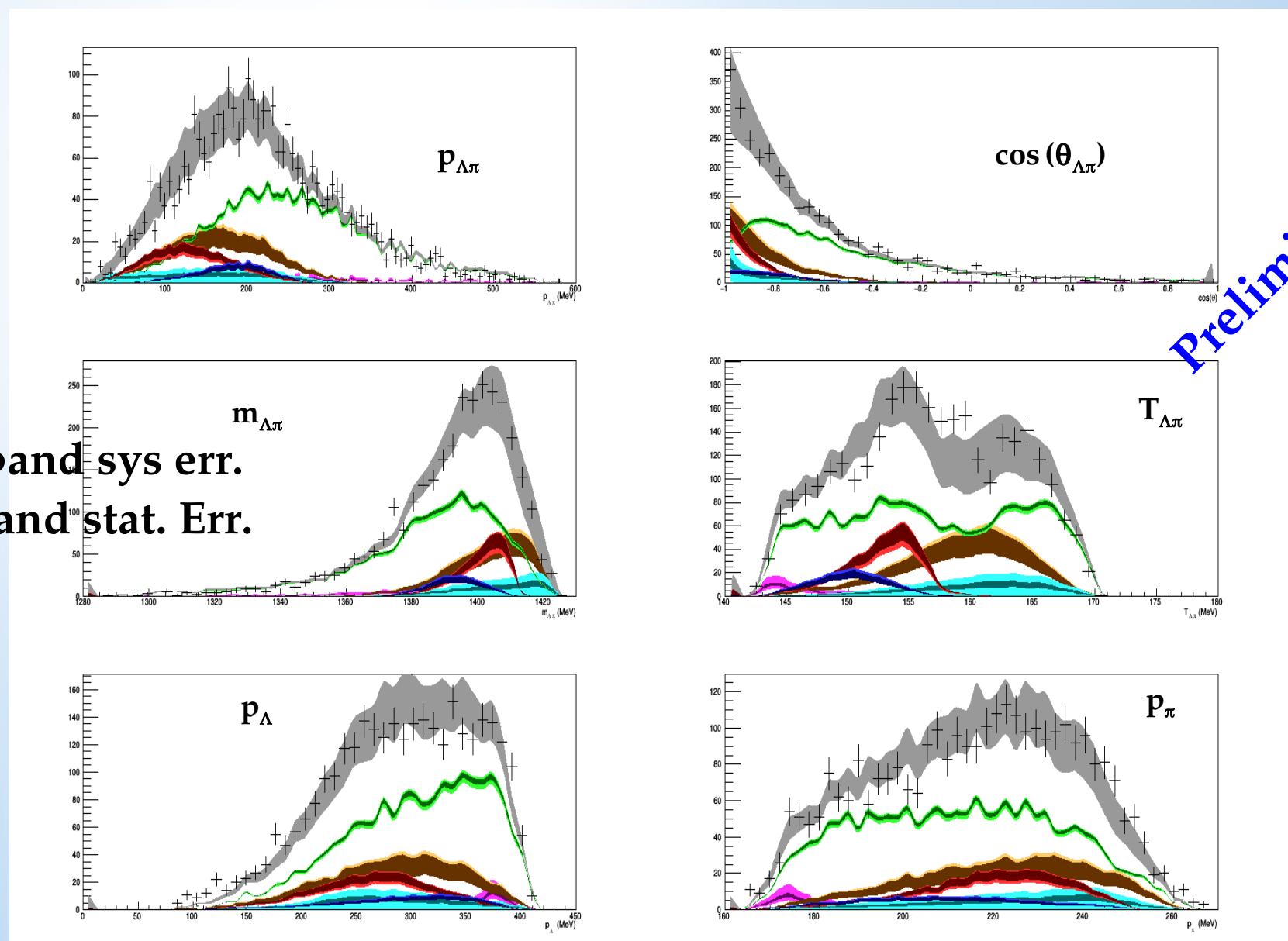


Preliminary

Simultaneous momentum – angle – mass fit

Preliminary

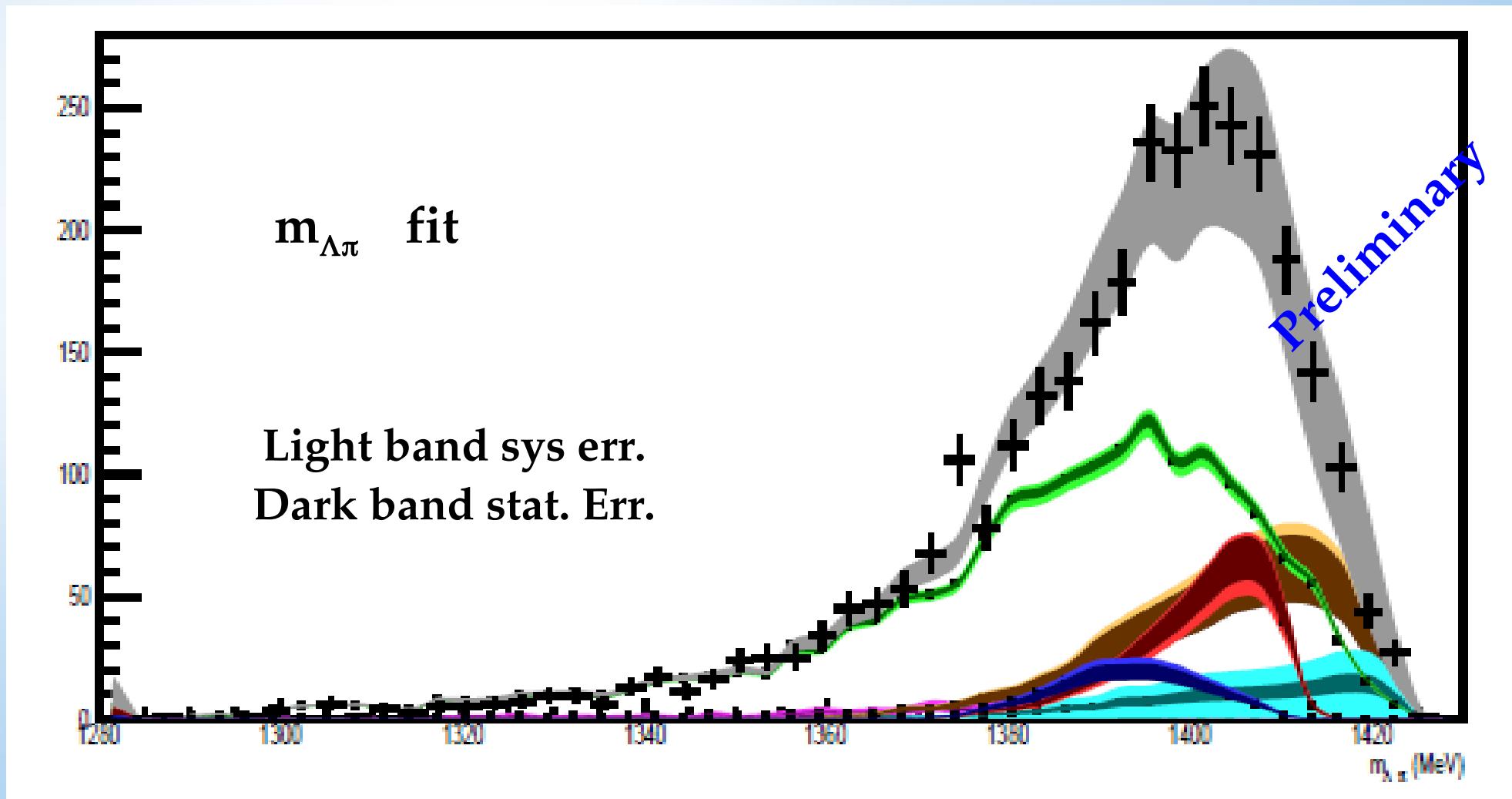
Light band sys err.
Dark band stat. Err.



Non-Resonant
(at-rest)
(in-flight)

Resonant Σ^*
(at-rest)
(in-flight)

Comparison



**Non-Resonant
(at-rest)
(in-flight)**

**Resonant Σ^*
(at-rest)
(in-flight)**

Outcome of the measurement

From the well known Σ^* transition probability:

$$\frac{\int_0^{p_{max}} P_s^s(p_{\Lambda\pi}) dp_{\Lambda\pi}}{\int_0^{p_{max}} P_s^p(p_{\Lambda\pi}) dp_{\Lambda\pi}} = |f_{ar}^s|^2 \cdot 8,94 \cdot 10^5 MeV^2.$$

Using the measured value for the ratio $\frac{NR-ar}{RES-ar}$, the module of the non resonant transition amplitude is found to be:

$$|f_{ar}^s| = (0.334 \pm 0.018 stat^{+0.34}_{-0.58} syst) fm. \quad (27)$$

The sub-threshold result is compatible with corresponding values extracted from $K^- p \rightarrow \Lambda \pi^0$ cross sections above threshold

J. K. Kim, Columbia University Report, Nevis 149 (1966)

J. K. Kim, Phys Rev Lett, 19 (1977) 1074:

$E = -33 MeV$	$p_{lab} = 120 MeV$	160 MeV	200 MeV	245 MeV
$0.334 \pm 0.018 stat^{+0.034}_{-0.058} syst$	0.33(11)	0.29(10)	0.24 (6)	0.28(2)

Preliminary

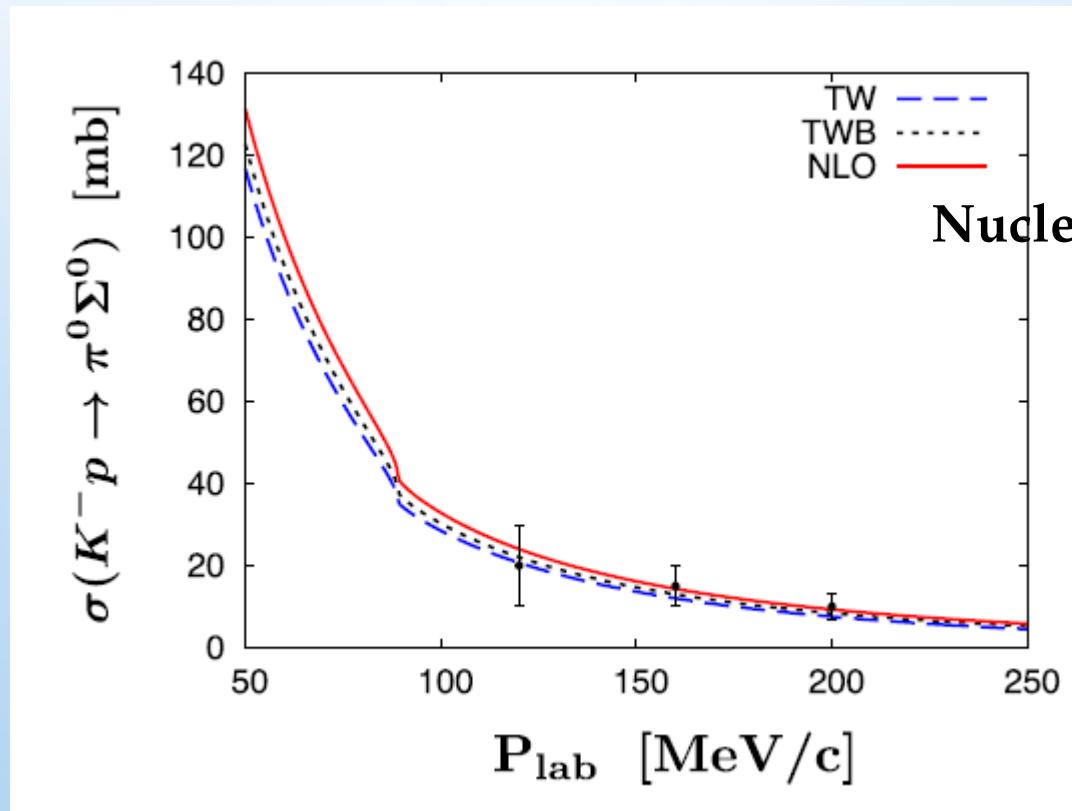
K⁻

future AMADEUS physics case



$K^- N$ cross section measurement below $p_K = 100 \text{ MeV}/c$, example..

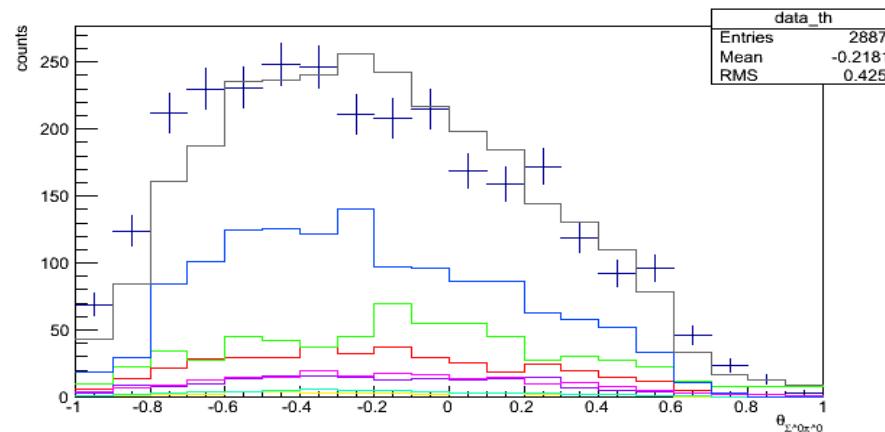
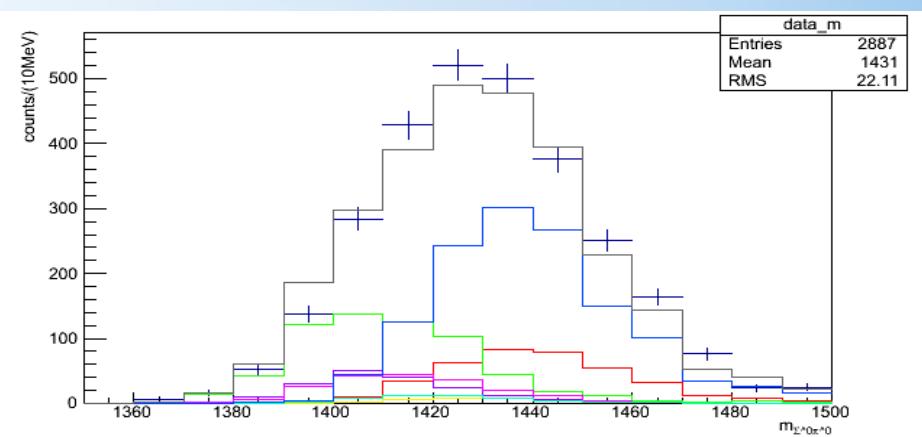
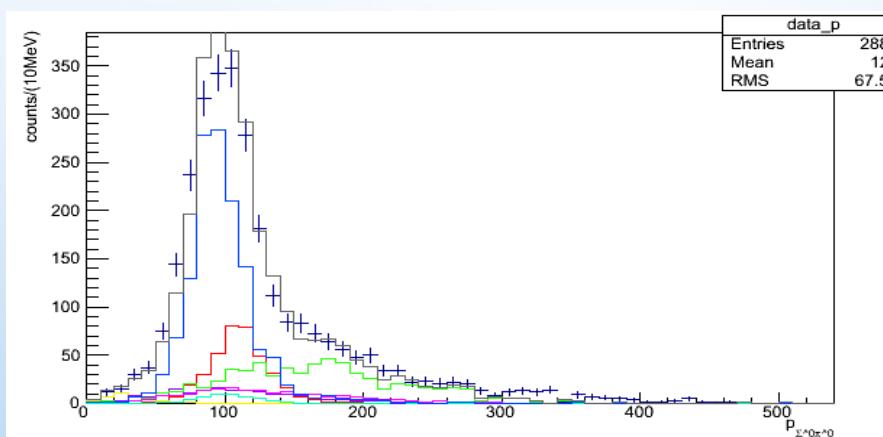
- $K^- p \rightarrow \Sigma^0 \pi^0$ cross section measurement at or below 100 MeV/c missing
- existing data at (120, 160, ..) MeV/c with big relative errors (about 50% & 120 MeV/c)



Nuclear Physics A 881 (2012) 98–114

$K^- p \rightarrow \Sigma^0 \pi^0$ cross section measurement can be done

- K- H capture at-rest → kinematics is closed
 - K- H capture at-rest ($p_K = 90$ MeV) → kinematics is closed
 - K- H capture at-rest ($p_K = 100$ MeV) → kinematics is closed
 - K- 4He capture at-rest + in-flight ($l_K = 1$)
 - K- 12C capture at-rest + in-flight ($l_K = 2$, valence proton)
- 



Preliminary

$\Lambda(1405)$ from K- induced reactions on deuterons

- Extract the $\Lambda(1405)$ resonance shape from **K- induced reactions on deuterons** (resonance is produced by the KN channel) **in-fight** (with $p_K = 130$ MeV/c)
- Calculations performed for these particular reactions based on the Esmaili - Akaishi - Yamazaki framework [Phys. Rev. C 83, (2011)] (left) greatly differ from chiral predictions [Eur. Phys. J. A 42, 257 (2009)] & [Eur. Phys. J. A (2011) 47] (right)

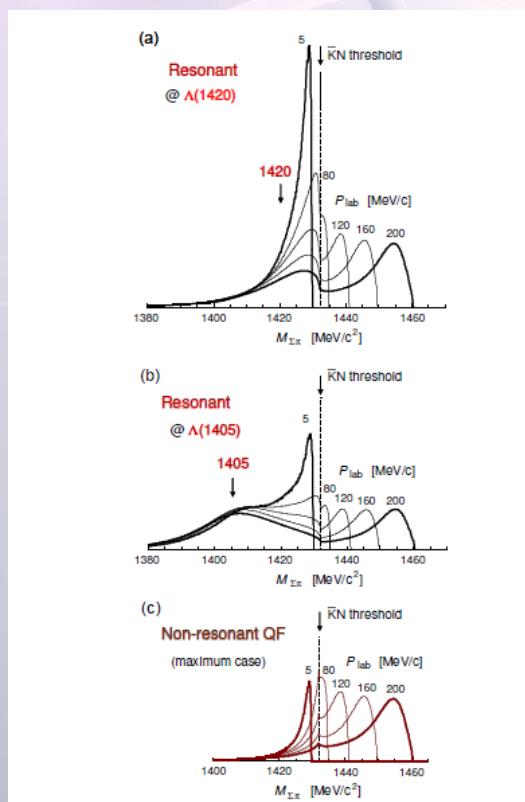


Figure 6: T_{21} invariant-mass spectra from low-momentum $K^- + d$ reactions calculated with AY interactions. The total spectra (resonant plus non-resonant): (a) for AY 1432-34 and (b) for AY 1405-30. The curves (a) and (b) are normalized to unit area to compare with $K^- d$ atomic absorption. (c) The case of non-resonant QF contribution for (a), in which the resonant QF contributions via Λ^* are subtracted by reducing U_{11} to 0.59 U_{11} where the Λ^* quasi-bound state disappears.

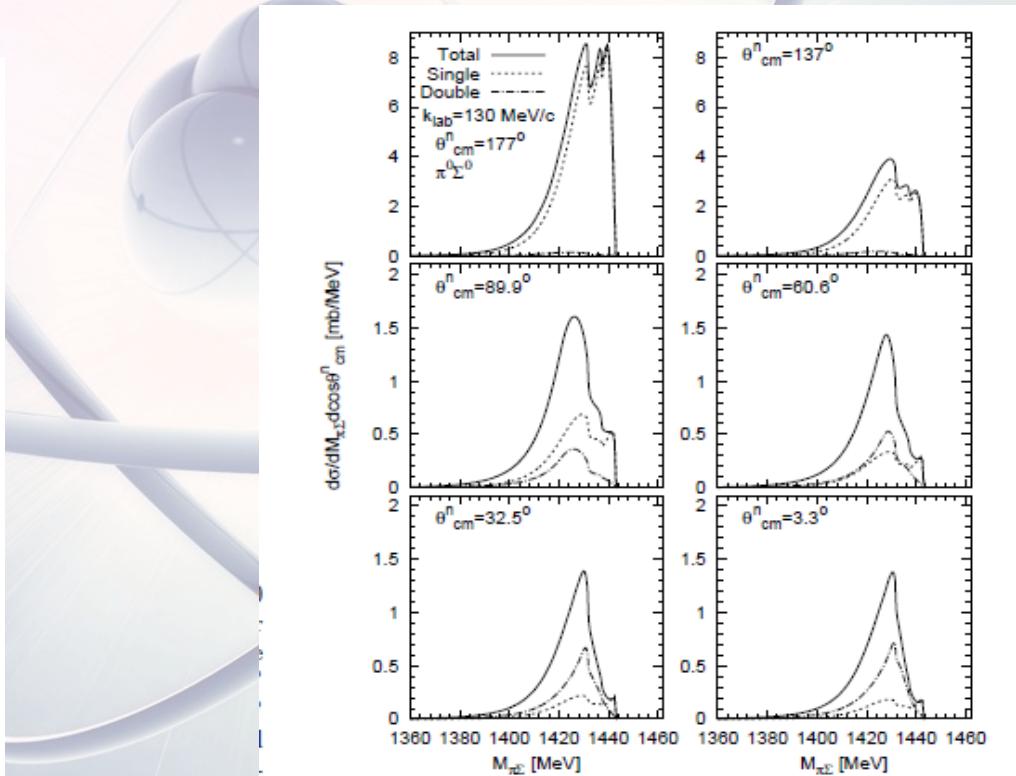


Fig. 5. Differential $\pi\Sigma$ invariant mass spectra of $K^- d \rightarrow \pi^0 \Sigma^0 n$ for 130 MeV/c of incident K^- momentum fixing the angle between the emitted neutron and the incident K^- in the CM frame. In each panel, the solid line denotes the total contributions of the three diagrams, while the dotted and dash-dotted lines show the calculations from the single and double scatterings, respectively.

detection of neutrals:

K^-

- Extract the $\Sigma^0\pi^0$ spectra:

free from $\Sigma(1385)$ background I=1

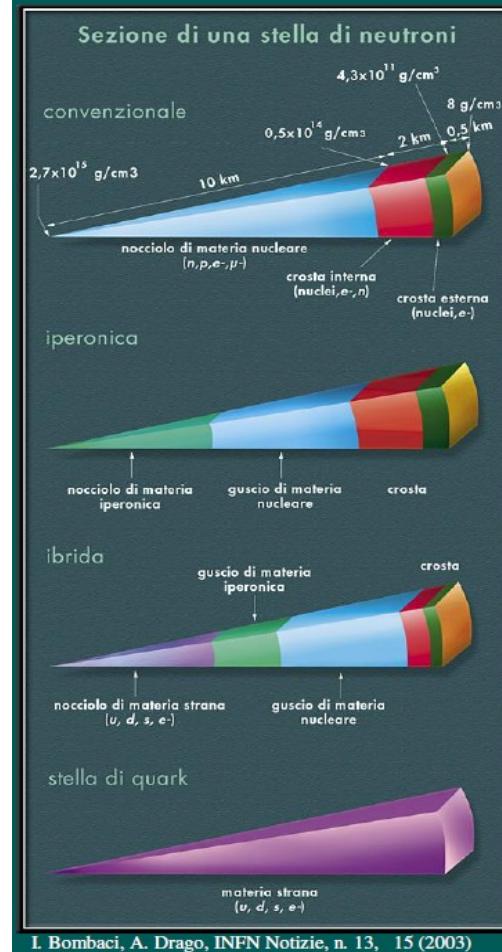
- The kinematics of the $K^- d \rightarrow \Sigma^0\pi^0 n$ reaction would be closed with the neutron detection → **event by event separation of at-rest from in-flight K^- nuclear captures.**
- The selection of forward neutrons (angle between the outgoing neutron and the incident K^- in the CM frame) drastically reduces the contribution of the single scattering, resonant (1405) formation dominates.

Y-N/NN interaction essential impact on the case of NEUTRON STARS

ECT*, Trento (Italy), 27 – 31 October 2014

Strangeness in Neutron Stars

Ignazio Bombaci
Dipartimento di Fisica “E. Fermi”, Università di Pisa
INFN Sezione di Pisa



“Neutron

Nucleon Stars

Hyperon Stars

Hybrid Stars

Strange Stars

Microscopic approach to hyperonic matter EOS

input

2BF: **nucleon-nucleon (NN), nucleon-hyperon (NY), hyperon-hyperon (YY)**
e.g. Nijmegen, Julich models

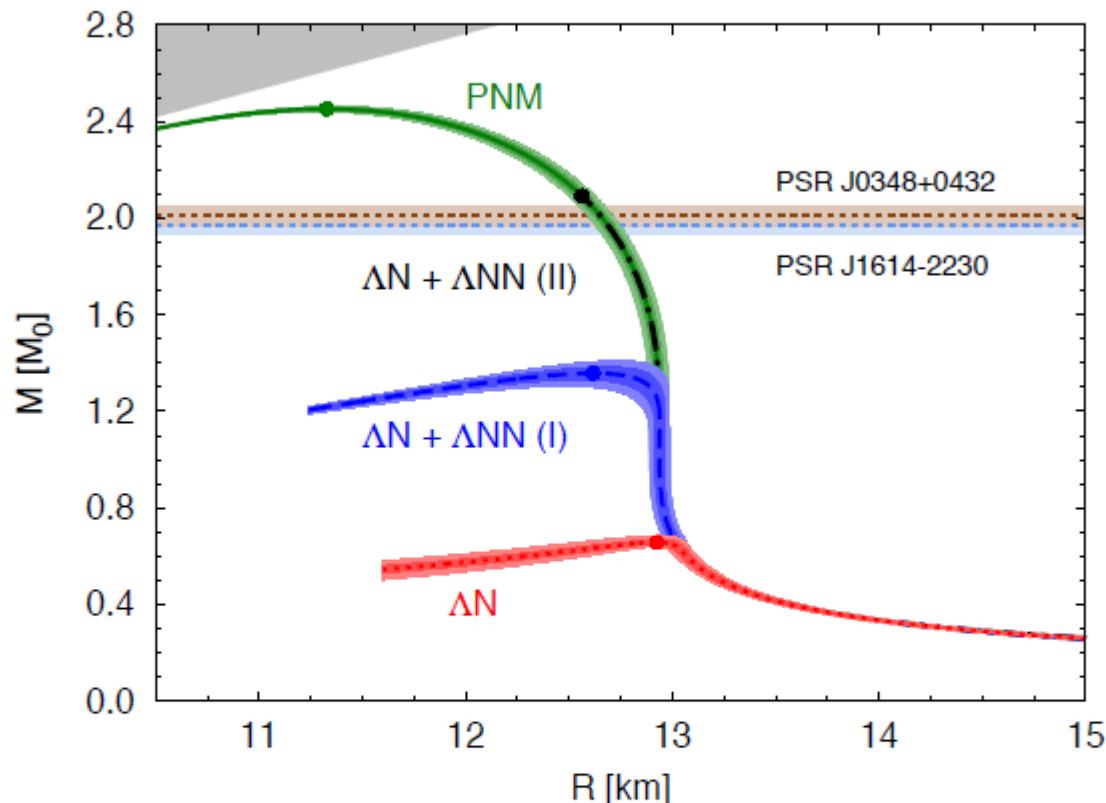
3BF: **NNN, NNY, NYY, YYY**

Hyperonic sector: experimental data

1. **YN scattering** (very few data)
2. **Hypernuclei**

No experimental information on Σ^0 -N/NN interaction

Λ -neutron matter

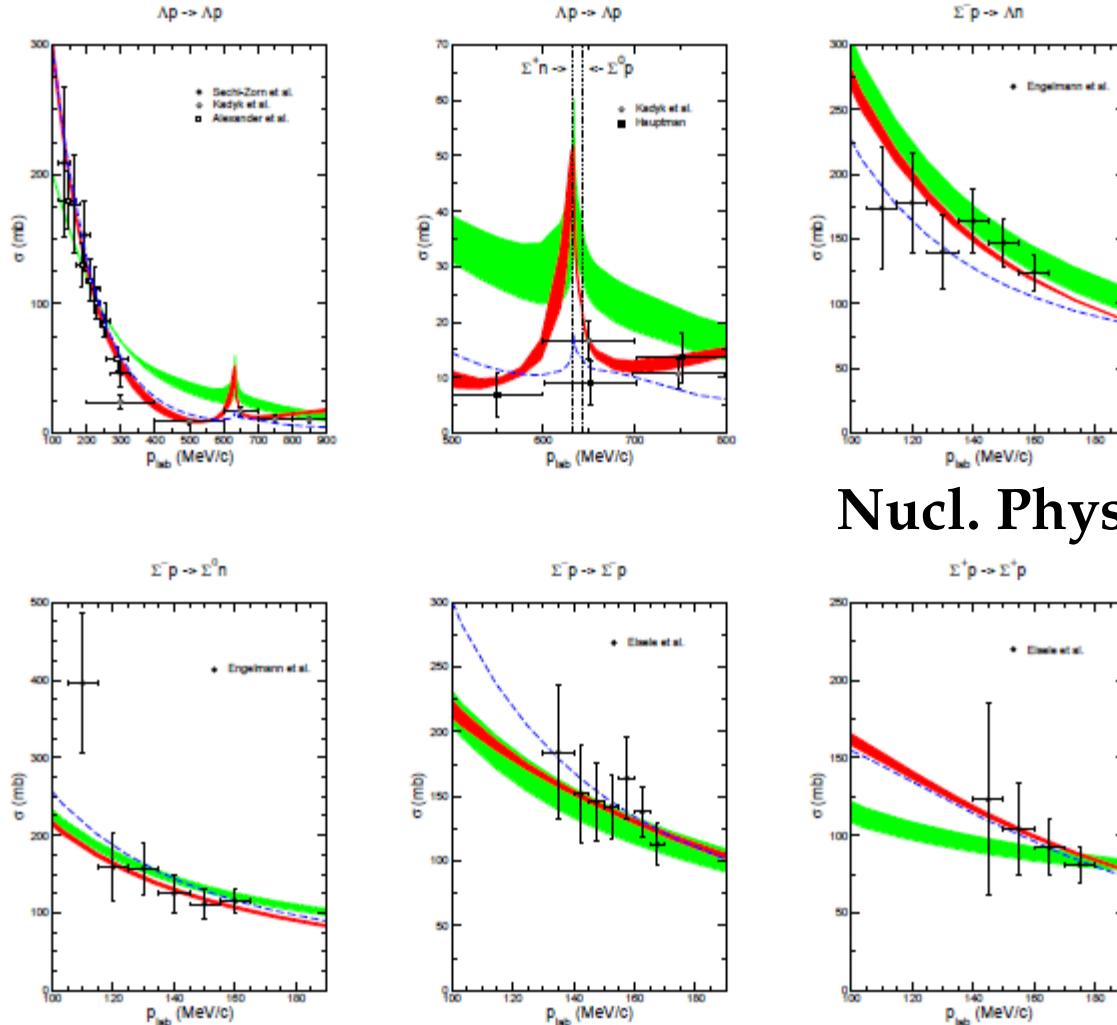


Lonardoni, Lovato, Gandolfi, Pederiva, PRL (2015)

Drastic role played by ΔNN . Calculations can be compatible with neutron star observations.

Note: no $v_{\Lambda\Lambda}$, no protons, and no other hyperons included yet...

No experimental information on Σ^0 -N/NN interaction



Nucl. Phys. A 915 (2013) 24-58

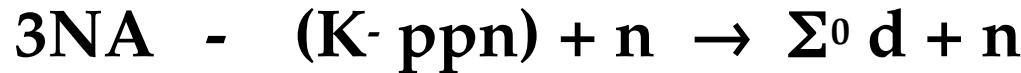
Figure 2: "Total" cross section σ (as defined in Eq. (24)) as a function of p_{lab} . The experimental cross sections are taken from Refs. [52] (filled circles), [53] (open squares), [65] (open circles), and [66] (filled squares) ($\Lambda p \rightarrow \Lambda p$), from [54] ($\Sigma^- p \rightarrow \Lambda n$, $\Sigma^- p \rightarrow \Sigma^0 n$) and from [55] ($\Sigma^- p \rightarrow \Sigma^- p$, $\Sigma^+ p \rightarrow \Sigma^+ p$). The red/dark band shows the chiral EFT results to NLO for variations of the cutoff in the range $\Lambda = 500, \dots, 650$ MeV, while the green/light band are results to LO for $\Lambda = 550, \dots, 700$ MeV. The dashed curve is the result of the Jülich '04 meson-exchange potential [36].



3NA in ${}^4\text{He}$

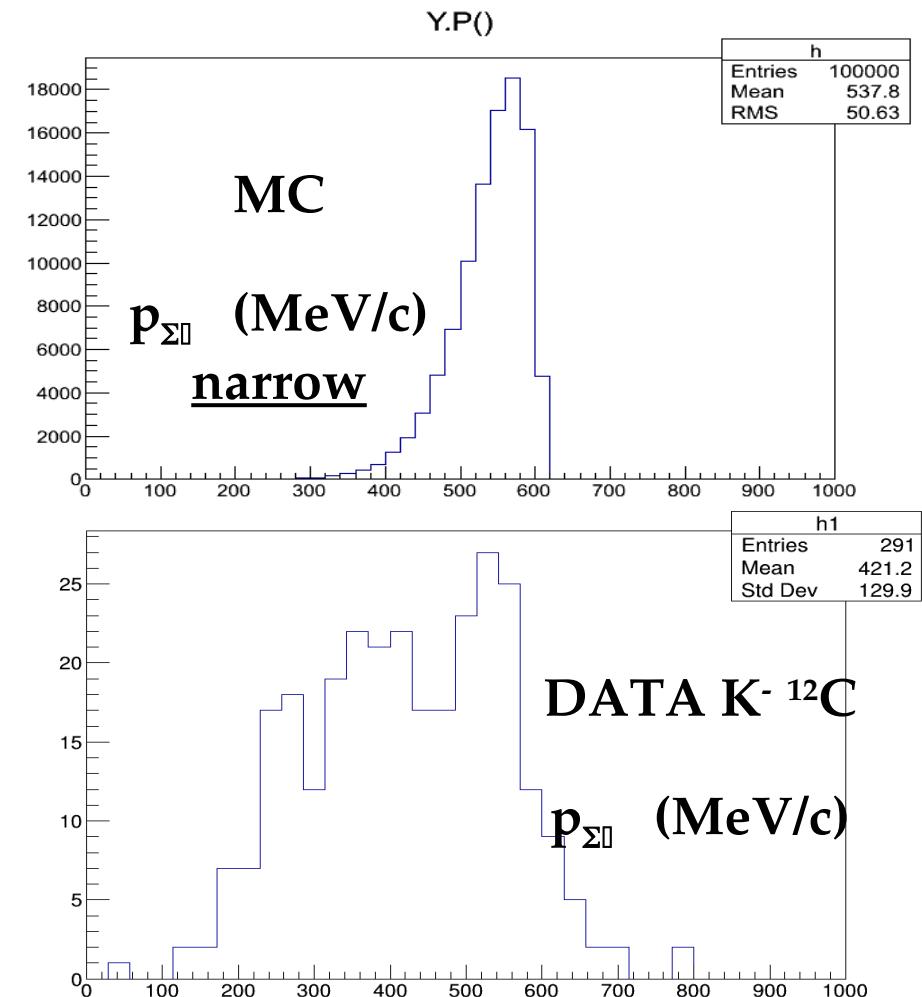
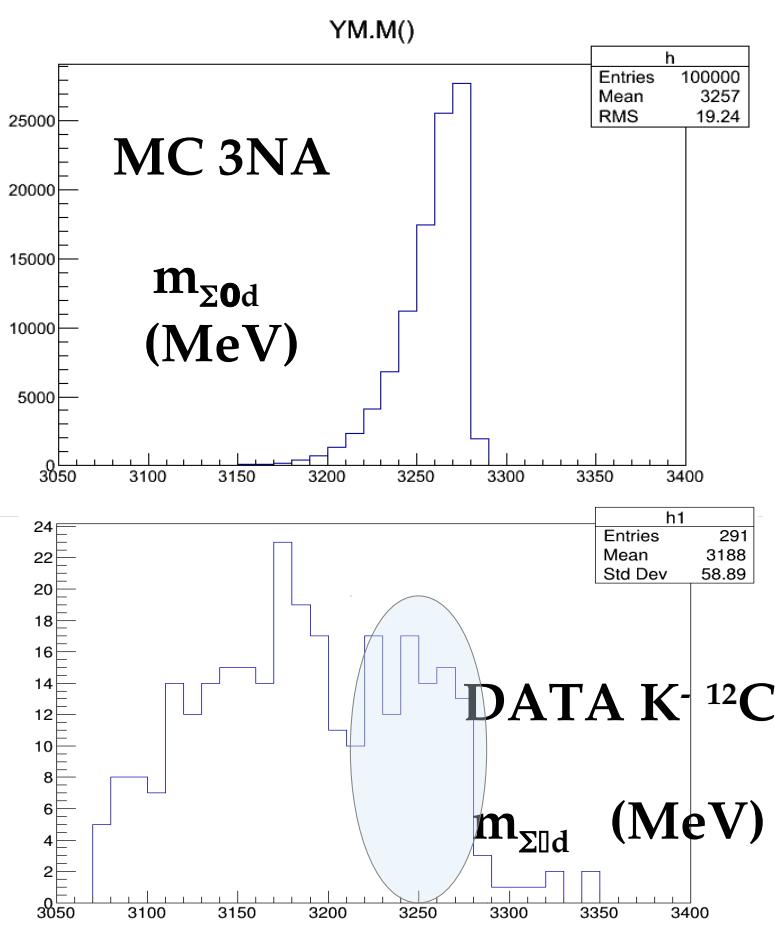
for the investigation of the

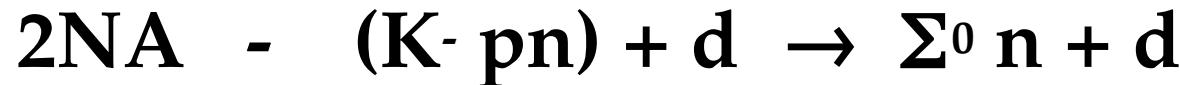
$\Sigma^0\text{-N}$ & $\Sigma^0\text{-(NN)}$ interaction



3NA can be followed by two possible elastic FSI

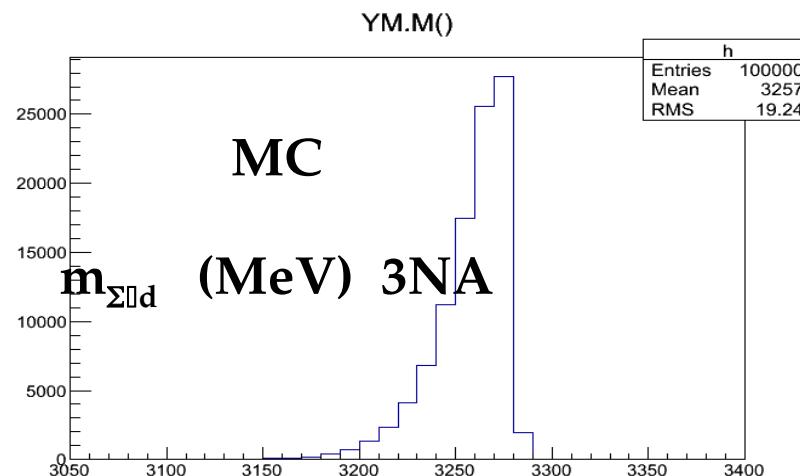
- 1) $\text{n d} \rightarrow \text{n d}$ we may take advantage of the well known σ_{NN} data
- 2) $\Sigma^0 \text{ n/d} \rightarrow \Sigma^0 \text{ n/d}$ from which to extract information on $\Sigma^0\text{-N}$, $\Sigma^0\text{-(NN)}$ interaction.



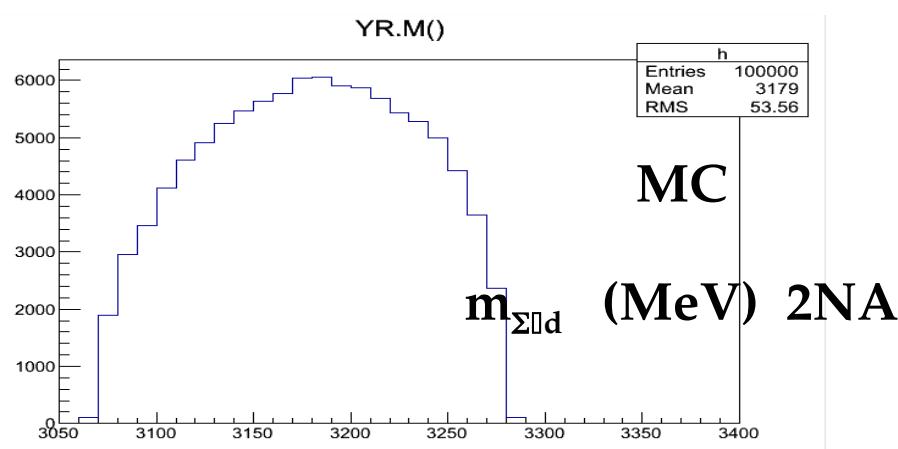


2 possible elastic FSI

- 1) $\text{n d} \rightarrow \text{n d}$ we may take advantage of the well known σ_{NN} data
- 2) $\Sigma^0 \text{ d/n} \rightarrow \Sigma^0 \text{ d/n}$ *hopefully well separated in the lower energy*



part of the final state phase space



Accurate model of the:



3NA in ${}^4\text{He}$

+



**is needed to extract the corresponding
cross sections from the measured shapes.**

K⁻

