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# Cosmological Structure from Quantum Fluctuations

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# outline of talk

#### Cosmic structure

- mapping our Universe

#### • Primordial sound waves

- classical waves an expanding Universe
- problem of initial conditions

#### Inflation in the very early Universe

- quantum fluctuations

# Hubble's law

#### Edwin Hubble made the most remarkable scientific discovery of the 20<sup>th</sup> century



Distant galaxies are moving away from us

and their *speed* is proportional to their *distance* 

M31

Andromeda

=> our universe is expanding

# Sloan Digital Sky Survey

today we use Hubble's law to build 3-dimensional maps of galaxies in our universe

the largest galaxy survey to date is the SDSS

over 200 scientists from 14 institutions (including Japan Participation Group)

it has mapped one million galaxies

see www.galaxyzoo.org

movies from COSMUS University of Chicago



#### SDSS: a map of one million galaxies

movie from COSMUS University of Chicago



cosmologists look back to a simpler past...

## how far back can we look?



The cosmic microwave background Radiation's "surface of last scatter" is analogous to the light coming through the clouds to our eye on a cloudy day. We can only see the surface of the cloud where light was last scattered

#### Cosmic microwave background radiation



- discovered by Penzias and Wilson 1965
- relic thermal radiation from the hot big bang







#### 2.7 K everywhere

+/- 3.3 mK redshift due local motion (at 1 million miles per hour)

+/- 18  $\mu\text{K}$  intrinsic anisotropies

COBE data release 1994

# WMAP first data released February 2003

latest data released February 2010



the surface of last scattering

© NASA

#### WMAP last scattering sphere

© NASA

quicktime movie from WMA http://map.gsfc.nasa.gov

# cosmic fluid equations:

energy density  $\rho$ pressure Pmomentum density  $\vec{q} = (\rho + P)\vec{v}$   $\dot{\rho} = -\vec{\nabla}.\vec{q} - 3H\rho$  $\dot{\vec{q}} = -\vec{\nabla}p - 3H\vec{q}$ 

Hubble expansion, H, dilutes density and momentum

eliminate pressure  $\vec{\nabla}P = c_s^2 \vec{\nabla}\rho$ where  $c_s^2$  = adiabatic sound speed = 1/3 in hot big bang

to obtain second - order wave equation

 $\ddot{\rho} + 3H\dot{\rho} = c_s^2 \nabla^2 \rho$ 

#### cosmic sound waves:

wave equation in an expanding spacetime:

Characteristic timescales for wavelength  $\lambda$ 

- oscillation period/wavelength  $\lambda / c_s$
- Hubble damping time-scale H<sup>-1</sup> -
- small-scales  $\lambda < c_s H^{-1}$  under-damped oscillator

 $\frac{d^2}{dt^2}\rho + 3H\frac{d}{dt}\rho = c_s^2\frac{d^2}{dr^2}\rho$ 

• large-scales  $\lambda > c_s H^{-1}$  over-damped ("frozen-in")

#### cosmo-seismology:

in general relativity should consider both

- effect of curved spacetime on matter
- effect of matter on curved spacetime

...leading to non-linear coupled equations

but often the simple picture works for linearised perturbations

two independent modes:

i. density waves (sound waves)

= scalar metric perturbations

- ii. gravitational waves (spacetime waves)
  - = tensor metric perturbations

study the detailed response of the system (our Universe) to inhomogeneous perturbations

#### Characteristic scales on last-scattering surface



coherent oscillations of a relativistic plasma (photons + electrons + protons) much more at background.uchicago.edu/~whu

#### WMAP angular power spectrum



#### Baryon Acoustic Oscillations also seen in galaxy distribution





# Precision cosmology

e.g., angular scale of first peak indicates flat space geometry, but also depends on nature of energy density in the universe



#### one question leads to another:

where do the primordial perturbations come from?

...quantum fluctuations during a period of inflation in the very early universe

## problem of initial conditions: $\delta\ddot{\rho} + 3H\delta\dot{\rho} + (c_s/\lambda)^2\delta\rho = 0$

Characteristic timescales for waves, fixed comoving wavelength

small-scales = late times,  $\lambda / c_s < H^1$ , under-damped oscillator 

large-scales = early times ,  $\lambda / c_s > H^1$ , "frozen-in"



inflation = accelerated expansion with almost constant Hubble rate

#### Cosmological inflation:

- period of accelerated expansion in the very early universe
- driven by vacuum energy

e.g. self-interacting scalar field  $V(\phi)$ 

speculative and uncertain physics

 just the kind of peculiar cosmological behaviour we observe today Starobinsky (1980) Guth (1981)



## Vacuum fluctuations

Hawking '82, Starobinsky '82, Guth & Pi '82

 $\delta \phi_k \approx$ 



- small-scale/underdamped zero-point fluctuations
- large-scale/overdamped perturbations in growing mode linear evolution ⇒ Gaussian random field

$$\left\langle \delta \phi^2 \right\rangle_{k=aH} \approx \frac{4\pi k^3 \left| \delta \phi_k^2 \right|}{\left(2\pi\right)^3} = \left(\frac{H}{2\pi}\right)^2$$

fluctuations of any light fields (m<3H/2) `frozen-in' on large scales

# cosmological Hawking radiation

semi-classical gravity - quantum fields in curved spacetime

from a black-hole event horizon





Hawking '73

or from a cosmological event horizon during a period of inflation in the very early universe



Hawking '82 Starobinsky '82 Guth and Pi '82

# scalar metric/density perturbations from inflation the $\delta N$ formalism for primordial perturbations



on large scales, neglect spatial gradients, treat as "separate universes"



$$\zeta = N(\phi_{initial}) - \overline{N} \approx \sum_{I} \frac{\partial N}{\partial \phi_{I}} \delta \phi_{I}$$

Starobinsky `85; Sasaki & Stewart `96 Lyth & Rodriguez '05 – works to any order

#### density perturbations from inflaton field

• quantum field fluctuations on unperturbed (flat) hypersurfaces during inflation leads to scalar metric perturbation

$$\zeta = \frac{dN}{d\phi} \,\delta\phi = \left(-\frac{H}{\dot{\phi}} \,\delta\phi\right)_{k=aH}$$

• produce primordial density perturbations in radiation-dominated era  $\Rightarrow \left\langle \frac{\delta T^2}{T^2} \right\rangle \approx \frac{1}{25} \left\langle \zeta^2 \right\rangle \approx \frac{1}{25} \left( \frac{H^2}{2\pi \dot{\phi}} \right)_{h=aH}^2$ 

slow time-dependence during inflation ⇒ weak scale-dependence

## WMAP 7 year data February 2010





scale-invariant *Harrison-Zel'dovich (n=1)* spectrum **excluded** by WMAP+BAO at 3sigma (Komatsu et al 2010)

#### what next?



#### WMAP 7-year data (Komatsu et al 2010)

running spectral index -0.086 < d n / d ln k < 0.018 higher-order in slow-roll

gravitational waves tensor-scalar ratio r < 0.36

non-Gaussianity -10 < f<sub>NL</sub> < 74



#### gravitational waves



• *transverse, traceless "tensor" metric perturbations* 

$$\delta g_{ij}(t,x) \approx \int d^3k \ h_k(t) \ e_{ij}^{(+,\times)}(x)$$

- amplitude, h(t), obeys same wave equation for massless field
  early time quantum fluctuation also frozen-in at late times
- *decoupled from matter perturbations*

$$\Rightarrow \langle T^2 \rangle_{today} \approx \left( \frac{64\pi}{M_{Pl}^2} \right) \left( \frac{H}{2\pi} \right)_{k=aH}^2$$

#### constrain the Hubble rate during inflation relative to Planck scale

#### ESA Planck satellite launched! 14th May 2009



#### next all-sky survey

50 million pixels improved polarisation sensitivity

data release 2012

Gravitational waves  $r \approx 0.1?$ 

Non-Gaussianity  $f_{NL} = 8?$ 

#### summary:

- *large-scale structure of our Universe can* come from small scale *quantum fluctuations*
- "golden era" for cosmological discovery, and ambitious plans
  - cosmic microwave background (CMBPol, B-pol, LiteBIRD)
  - large-scale structure in galaxy surveys (dark energy survey...)
  - radio surveys (Square Kilometre Array)
  - gravitational wave detectors (LIGO, LISA...)
- can go beyond linear theory and study nonlinear interactions in early universe
  - non-Gaussianity of density perturbations
  - gravitational waves from early universe