



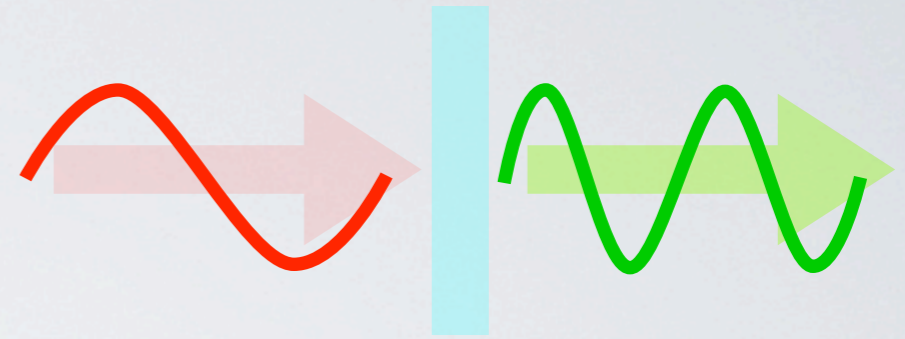
TERAHP

Here !

Quest for Extreme Terahertz Nonlinear Optics : on Dirac Fermions in Graphene

Shuntaro Tani, Kyoto University

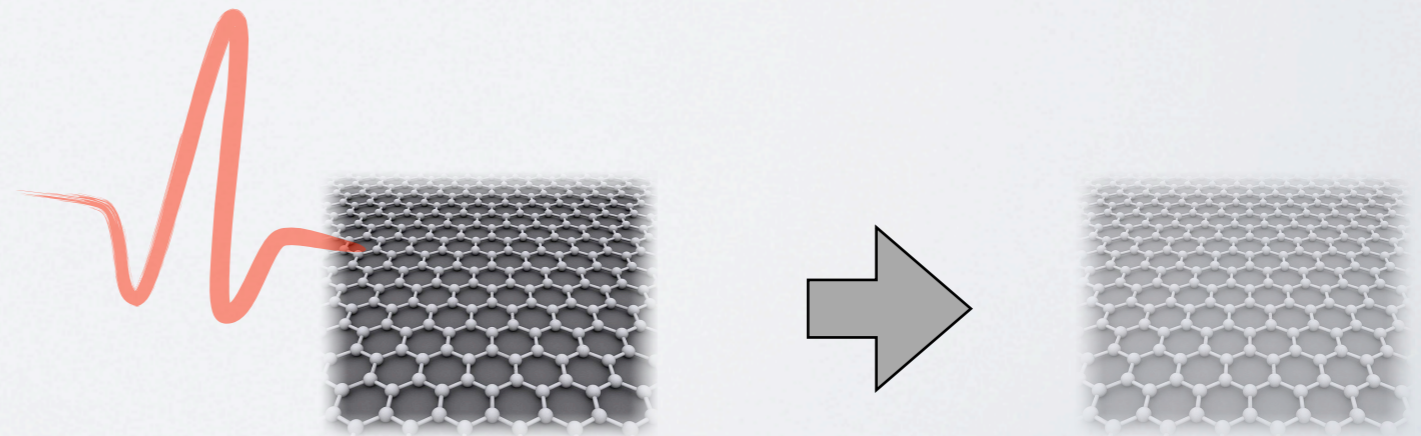
- Nonlinear Optics and Extreme Nonlinear Optics



- Terahertz Nonlinear Optics

$$U_p = \frac{e^2 E^2}{4m\omega^2}$$

- Terahertz Nonlinear Response in Graphene



What is Nonlinear Optics ?

Maxwell equations in medium

$$\nabla \cdot E = -\frac{1}{\epsilon_0} \nabla \cdot P$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\nabla \cdot B = 0$$

$$c^2 \nabla \times B = \frac{\partial E}{\partial t} + \frac{1}{\epsilon_0} \frac{\partial P}{\partial t}$$

Polarization $P \propto \chi E$ (~~$E \ll$ interatomic electric fields~~)

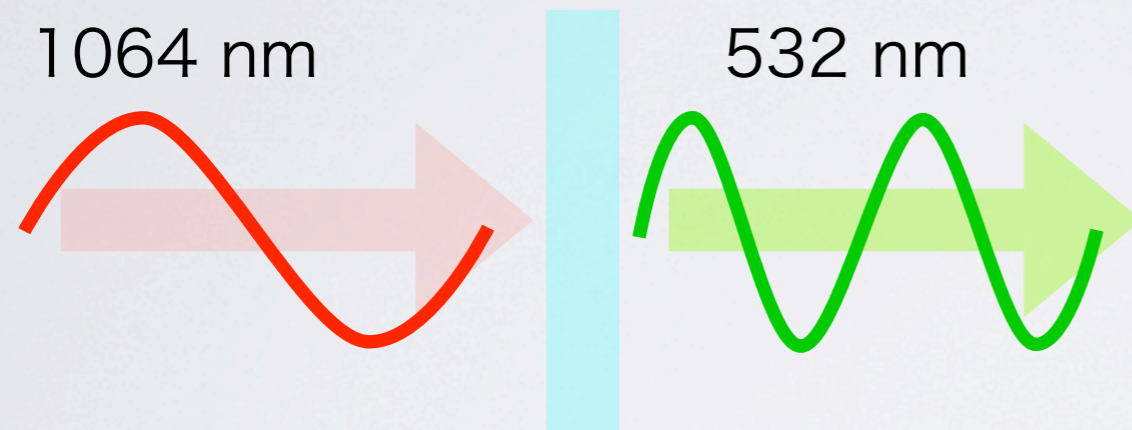


$$P \propto \chi^{(1)} E + \chi^{(2)} EE + \chi^{(3)} EEE + \dots$$

What is Nonlinear Optics ?

$$P \propto \chi^{(1)} E + \chi^{(2)} EE + \chi^{(3)} EEE + \dots$$

$$E \propto e^{-i\omega t} \rightarrow P \propto e^{-i \underline{2\omega t}} \rightarrow E \propto e^{-i \underline{2\omega t}}$$



What happens when the “energy” of light becomes comparable or larger than a characteristic energy of the system.

The energy associated to the light intensity

Ponderomotive energy $U_p \propto I$

Rabi energy $\hbar\Omega_R \propto \sqrt{I}$

Bloch energy $\hbar\Omega_B \propto \sqrt{I}$

Cyclotron energy $\hbar\omega_c \propto \sqrt{I}$

Tunneling energy $\hbar\Omega_{tun} \propto \sqrt{I}$

The characteristic energy of the system

Carrier photon energy $\hbar\omega_0$

Binding energy E_b

Rest Energy m_0c^2

The energy associated to the light intensity

Ponderomotive energy $U_p \propto I$

► An average kinetic energy of a electron in the oscillating electric field

$$U_p = \frac{e^2 E^2}{4m\omega^2} \propto I$$

The characteristic energy of the system

Carrier photon energy $\hbar\omega_0$

Binding energy E_b

Rest Energy $m_0 c^2$

Nonlinear Optical Phenomena with Visible Laser

Electron-positron pairs generation from vacuum

Nonlinear optics of the vacuum

Photo-nuclear fission

Relativistic nonlinear optics of electrons

Tunneling of electrons from atoms

SHG in solid

Linear regime

10^{-10}

10^0

10^{10}

10^{20}

10^{30}

Schwinger intensity

The highest laser intensity available

Laser intensity in 1961

The sunlight on the earth

Cosmic background radiation

W/cm^2

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W/cm^2

Links Among

the Hierarchies

Vacuum → "Vacuum" in a semiconductor
Electron-positron pairs → Electron-hole pairs

Visible Laser → Terahertz-wave

1 THz \sim 4.14 meV \ll visible light (\sim 2 eV)

Elementary excitation in solid

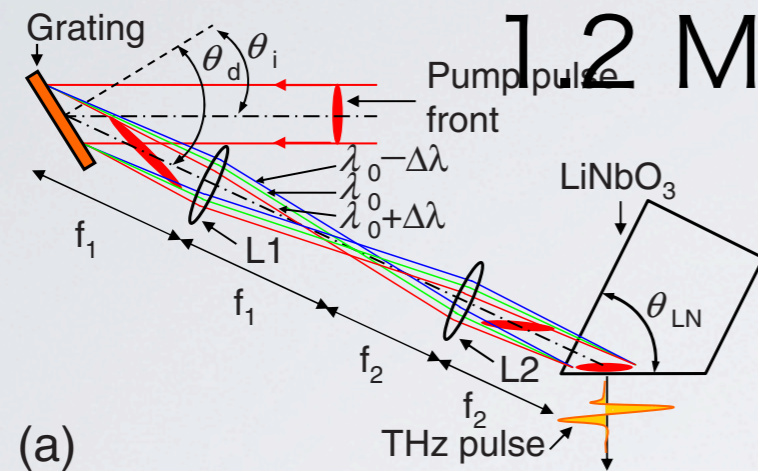
- Exciton binding energy
 - Lattice vibration energy
 - Superconductor gap
- etc...

Temporal waveform measurement is possible!

Coherent single cycle pulse is available!

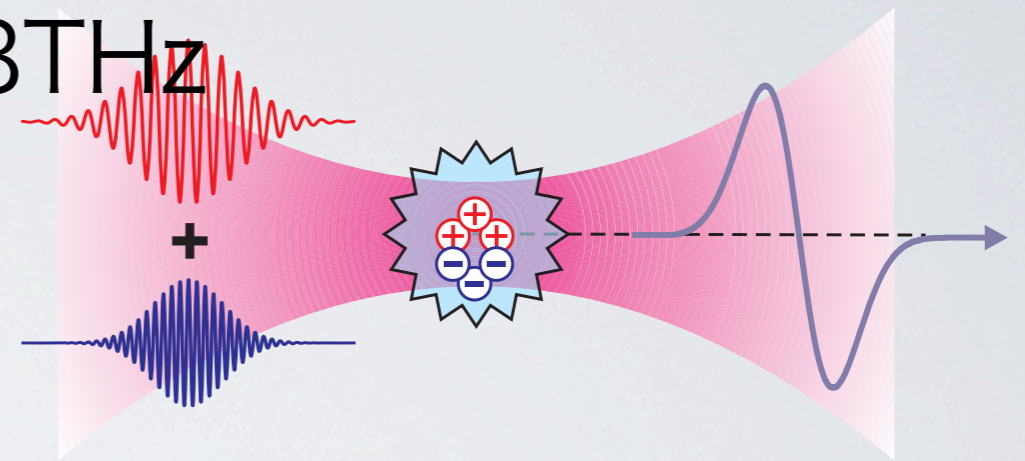
High Power Terahertz-wave Generation

Tilted wavefront method



1.2 MV/cm @ 0.8THz

Air plasma method



H.Hirori, K.Tanaka, et al., APL **98**, 091106 (2011) M.D.Thomson, et al., Laser & Photon. Rev. **1**, 349 (2007)

Terahertz Nonlinear Phenomena

- Coherent control of intra-excitonic state

S.Leinß, et al., Phys. Rev. Lett. **101**, 246401 (2008)

- Light-induced superconductivity in cuprate

D.Fausti, et al., Science **331**, 189 (2011)

- Extraordinary carrier multiplication

H.Hirori, K.Tanaka, et al., Nature Comm. **2**, 594 (2011)

Links Among the Hierarchies

Ponderomotive energy

$$U_p = \frac{e^2 E^2}{4m\omega^2}$$

Visible → **THz**

10^6 times larger

10 eV with 1 MV/cm
@ 1 THz

Carrier photon energy $\hbar\omega_0$

10^3 times smaller

Binding energy E_b

10^3 times smaller

Rest Energy m_0c^2

10^5 times smaller

Nonlinear Optical Phenomena with Visible Laser

Electron-positron pairs generation from vacuum

Nonlinear optics of the vacuum

Photo-nuclear fission

Relativistic nonlinear optics of electrons

Tunneling of electrons from atoms

SHG in solid

Linear regime

10^{-10}

Cosmic background radiation

10^0

The sunlight on the earth

Laser intensity in 1961

10^{10}

Our Terahertz Pulse

10^{20}

The highest laser intensity available

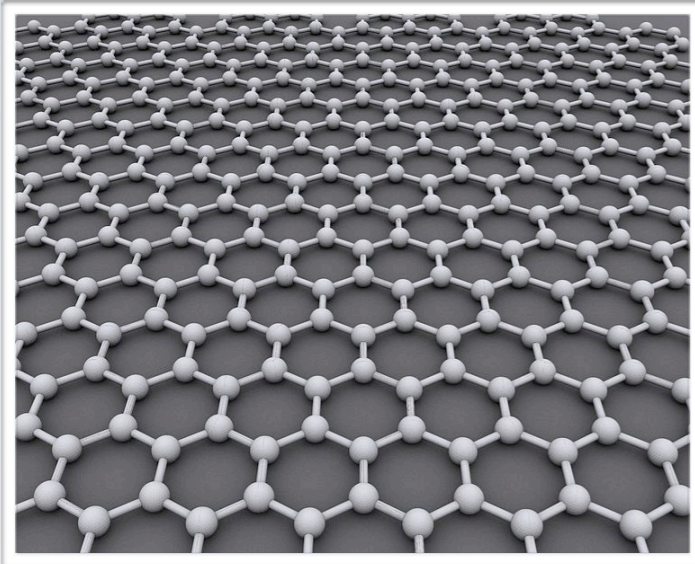
Schwinger intensity

10^{30}

W/cm²

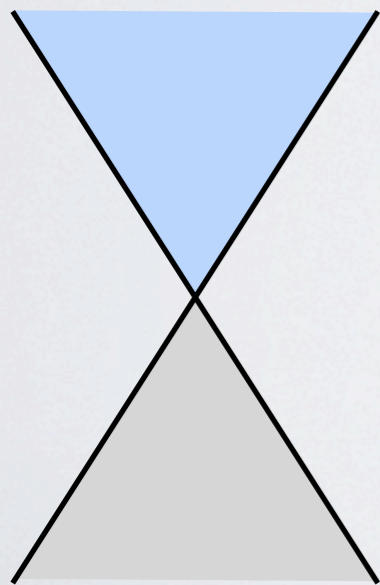
Relativistic Massless Electrons in Solid : Graphene

Graphene



Monoatomic layer carbon material

Conduction band



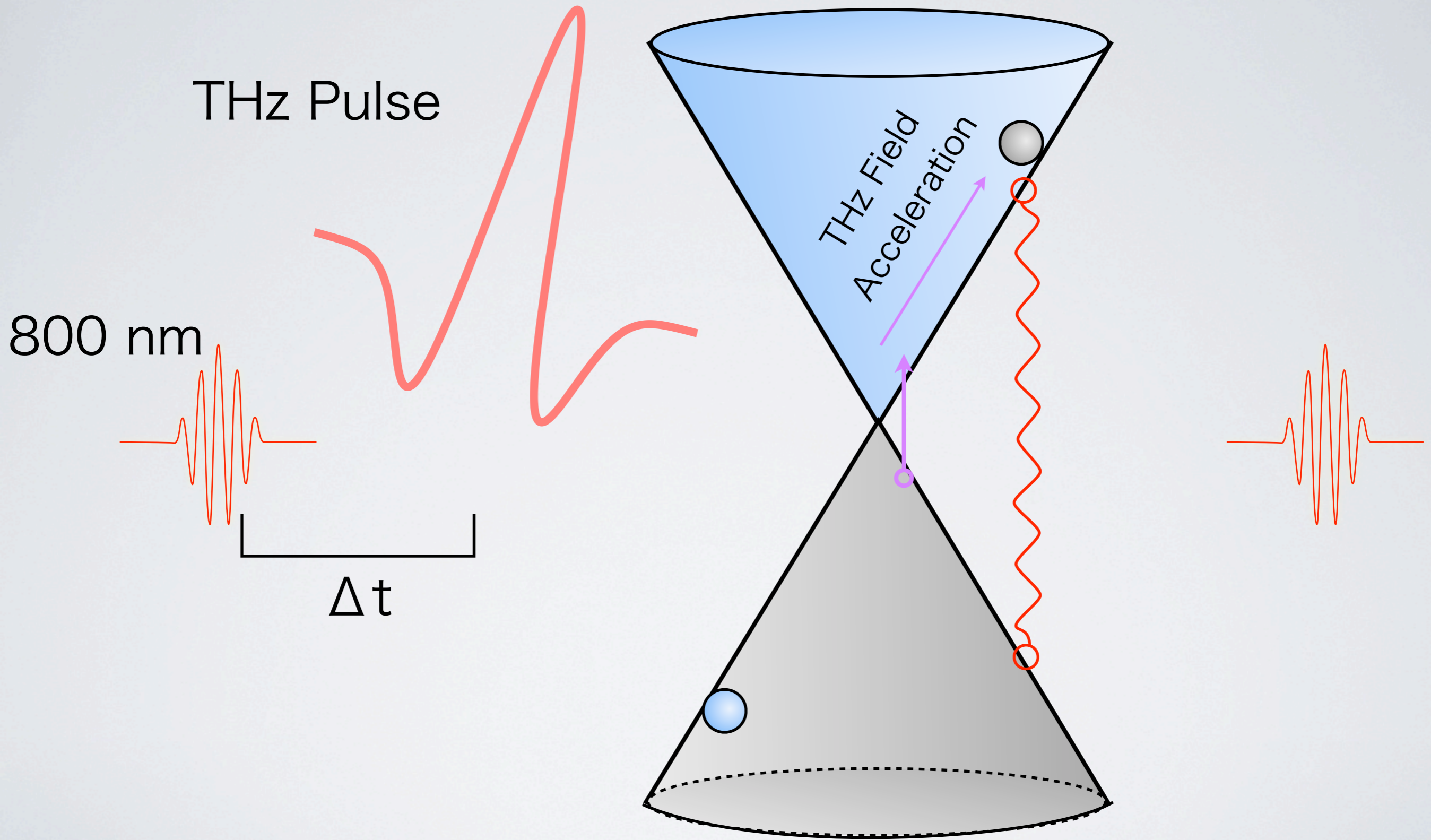
$$E = \pm \hbar v_F k$$

Valence band

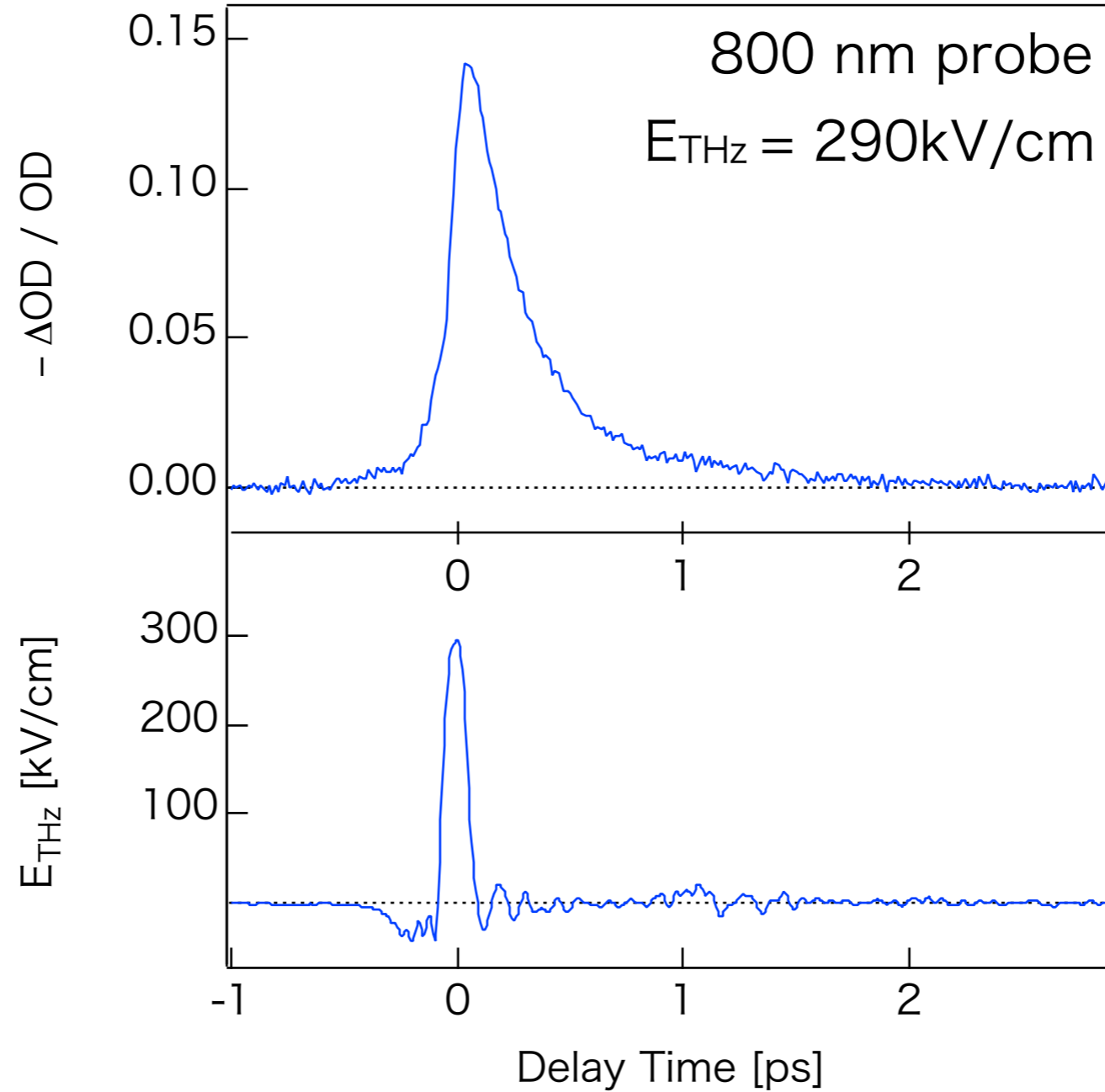
Dirac cone

- ▶ Linear dispersion
- ▶ Relativistic massless electrons

Experimental Scheme



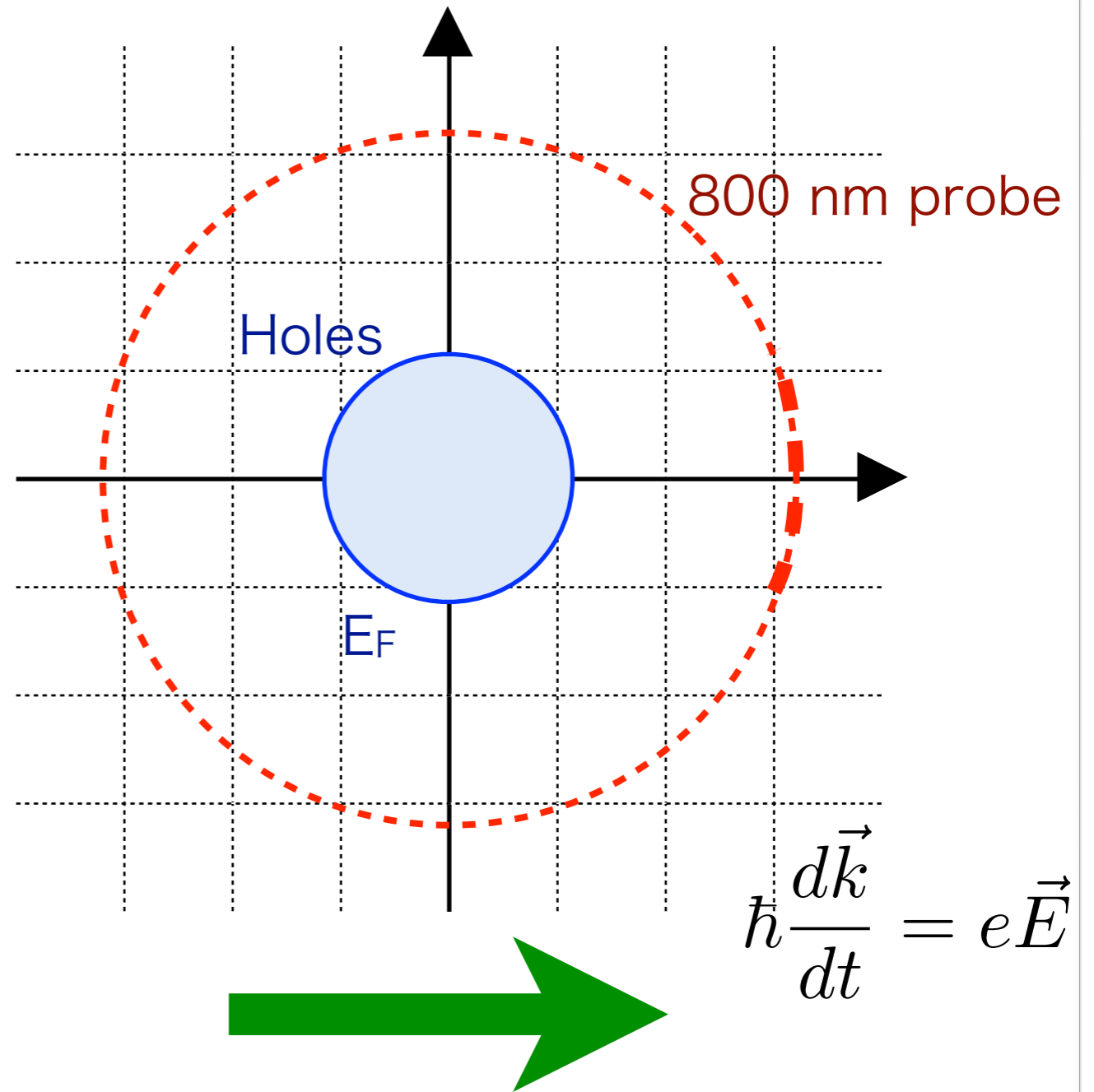
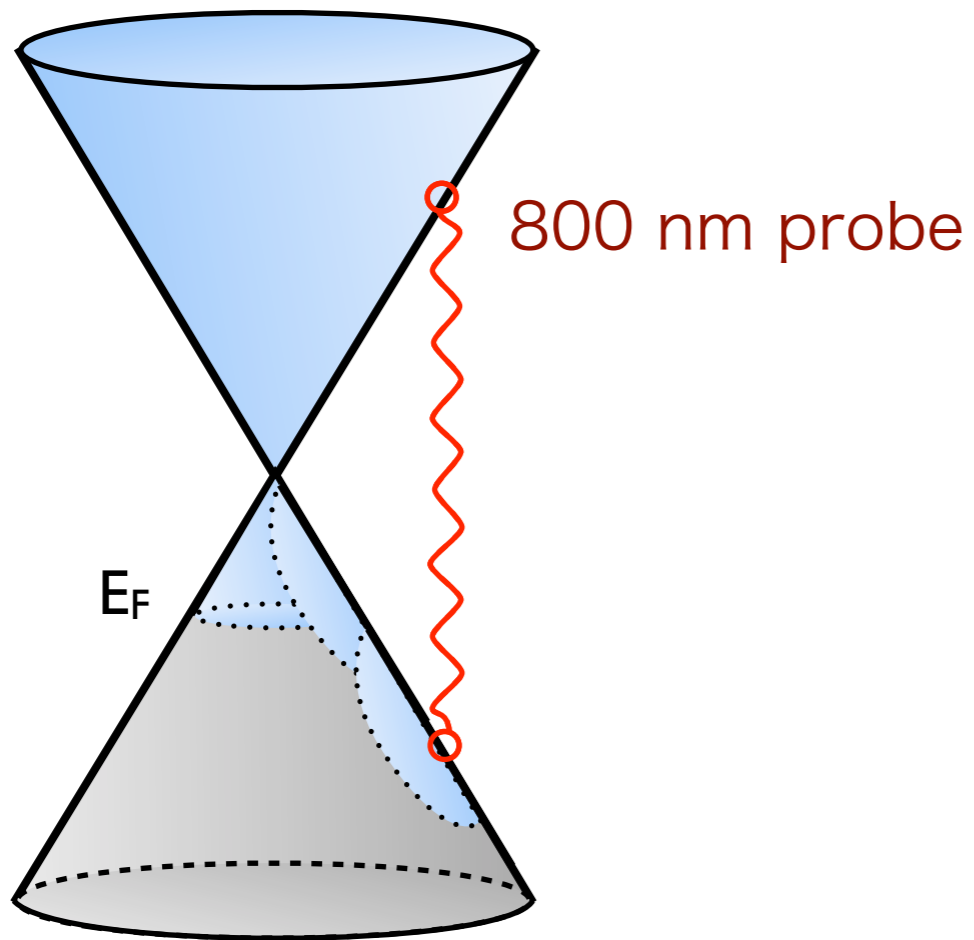
Results : Typical THz-induced Transparency in Near-infrared Region



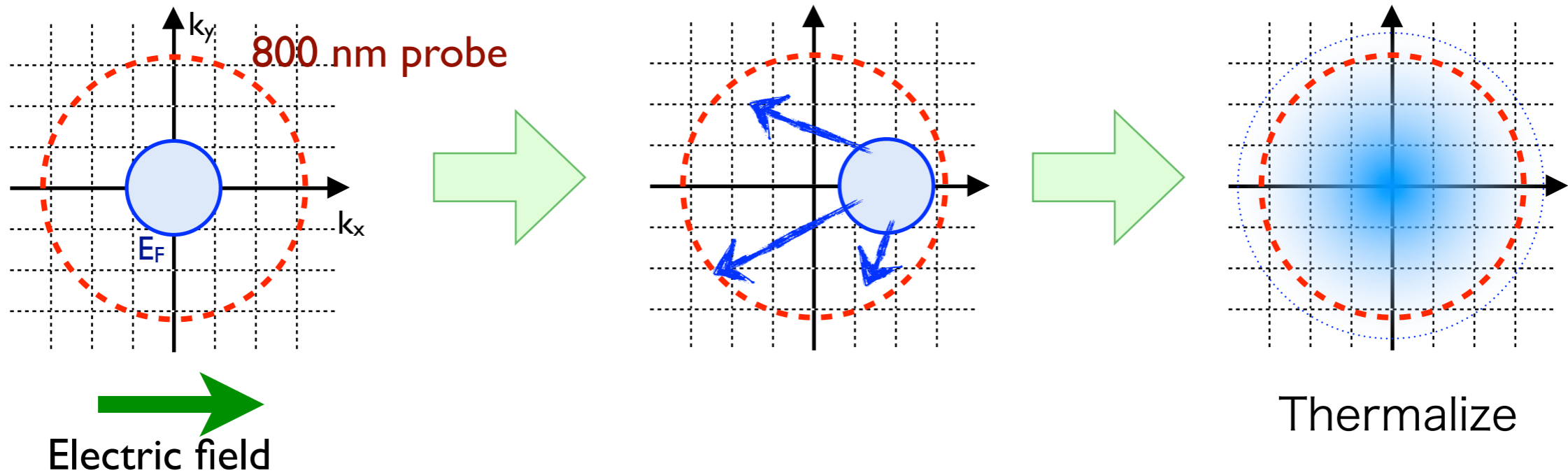
THz induced transparency over 14% at 800 nm

Origin of THz Induced Transparency

Hole distribution in the **momentum space**



Carrier Scattering Process



Boltzmann equation for $f(\vec{k}, t)$

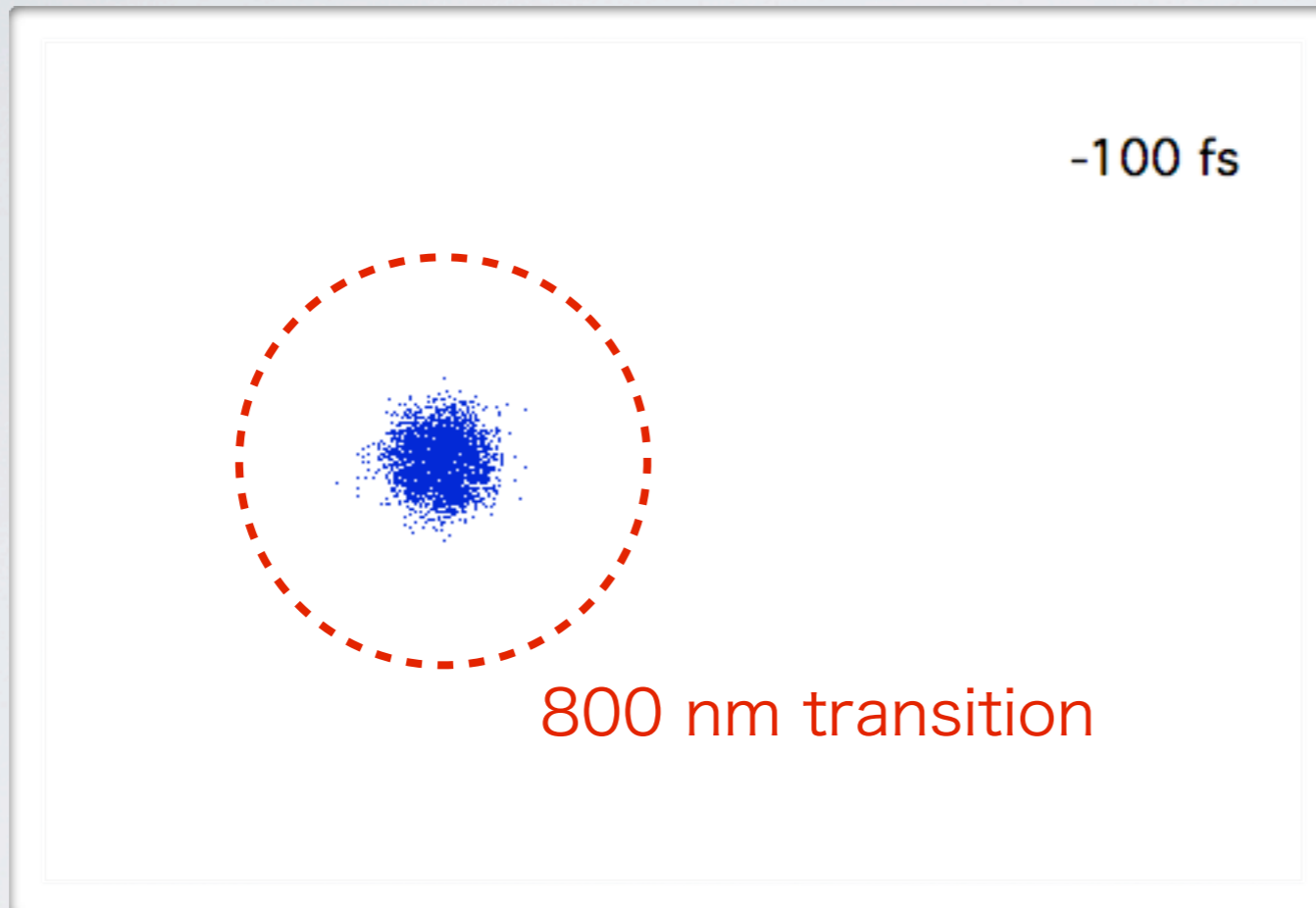
$$\frac{\partial f}{\partial t} + \frac{e\vec{E}}{\hbar} \cdot \frac{\partial f}{\partial \vec{k}} = \frac{\partial f}{\partial t} \Big|_{col}$$

$$= \frac{\partial f}{\partial t} \Big|_{e-e} + \frac{\partial f}{\partial t} \Big|_{h-h} + \frac{\partial f}{\partial t} \Big|_{OPE}$$

- No screening
- $\omega_{TO} = 160 \text{ meV}$

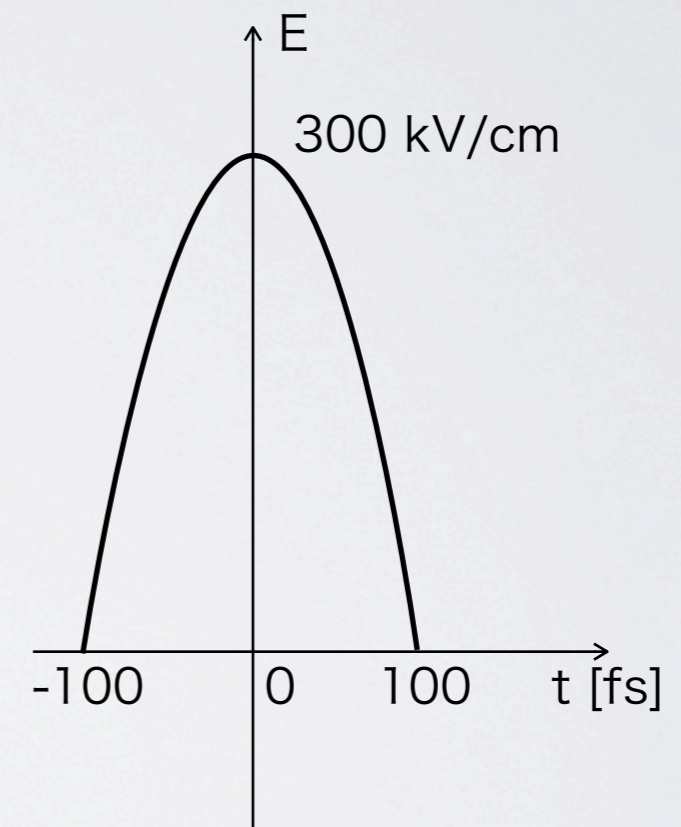
Time Development of Hole Distribution

Hole distribution in the momentum space



$$E_{\text{THz}} = 300 \text{ kV/cm}$$

$$E_F = -200 \text{ meV}$$

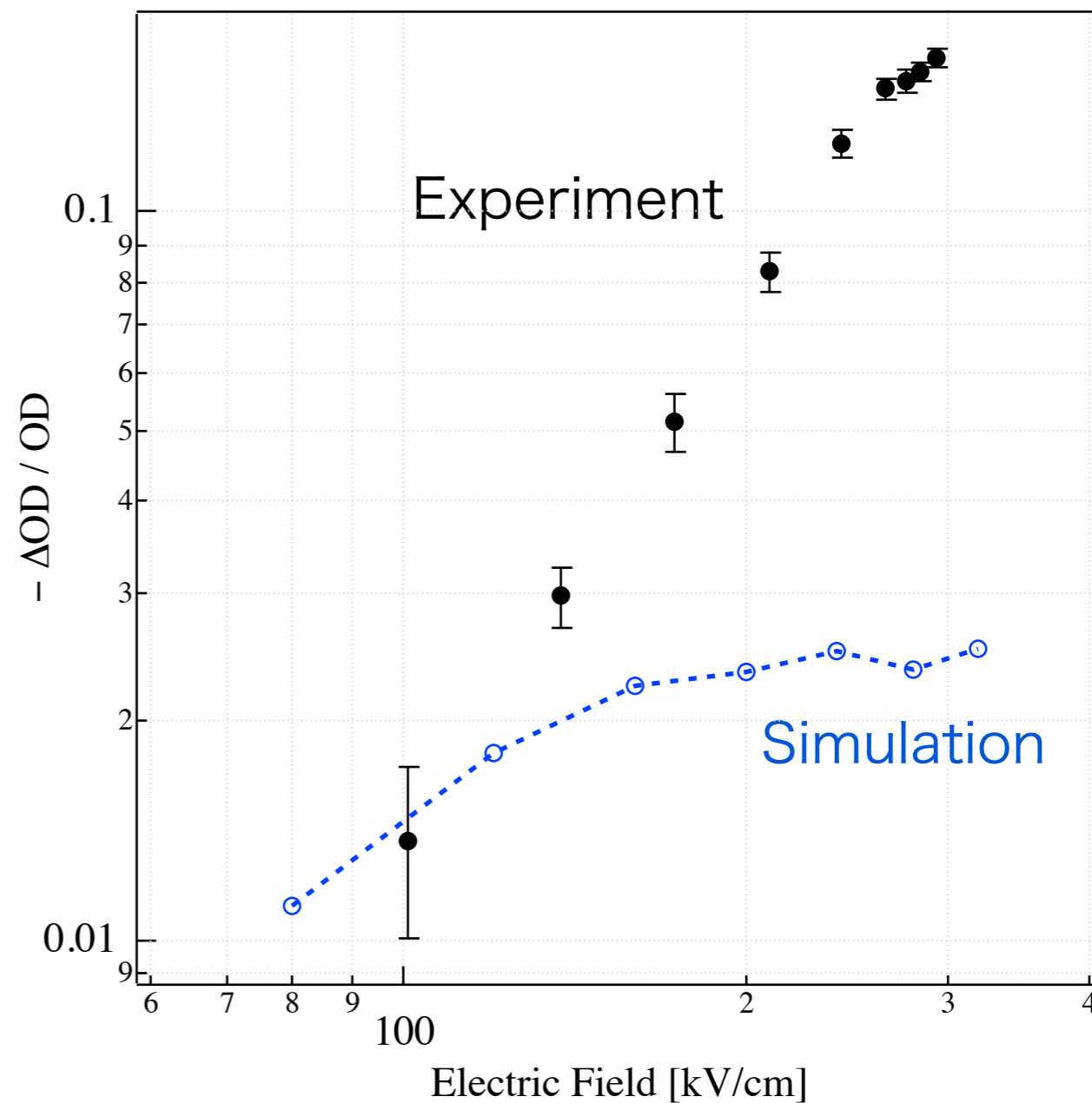
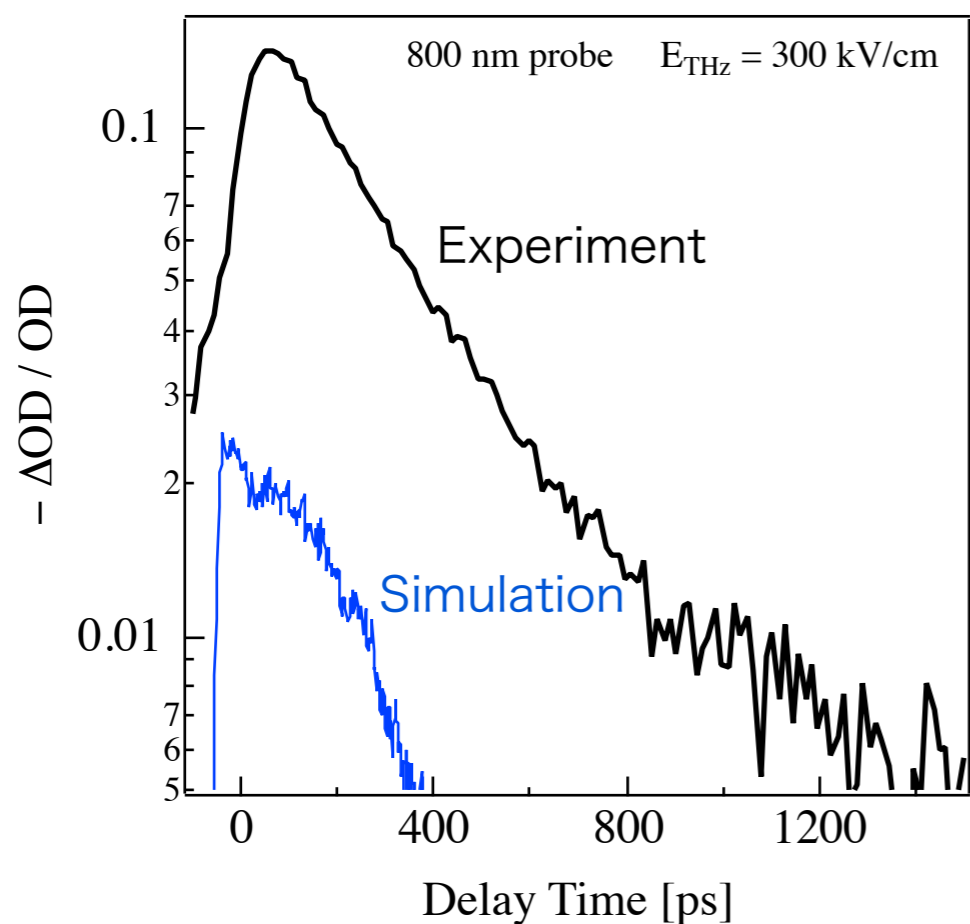
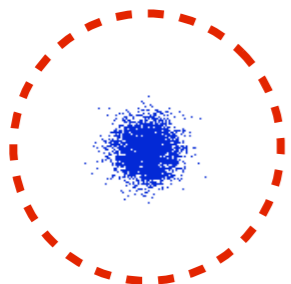


Parameters are chosen to reproduce experimental results.

Simulation Results

$E_{\text{THz}} = 300 \text{ kV/cm}$
 $E_{\text{F}} = -200 \text{ meV}$

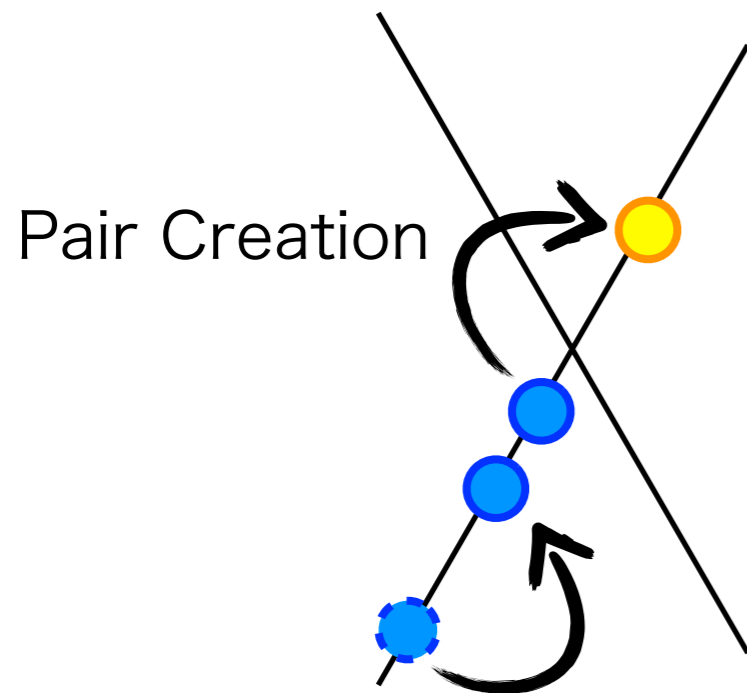
-100 fs



Carriers may not be enough.

Interband Carrier Scattering

Impact Ionization

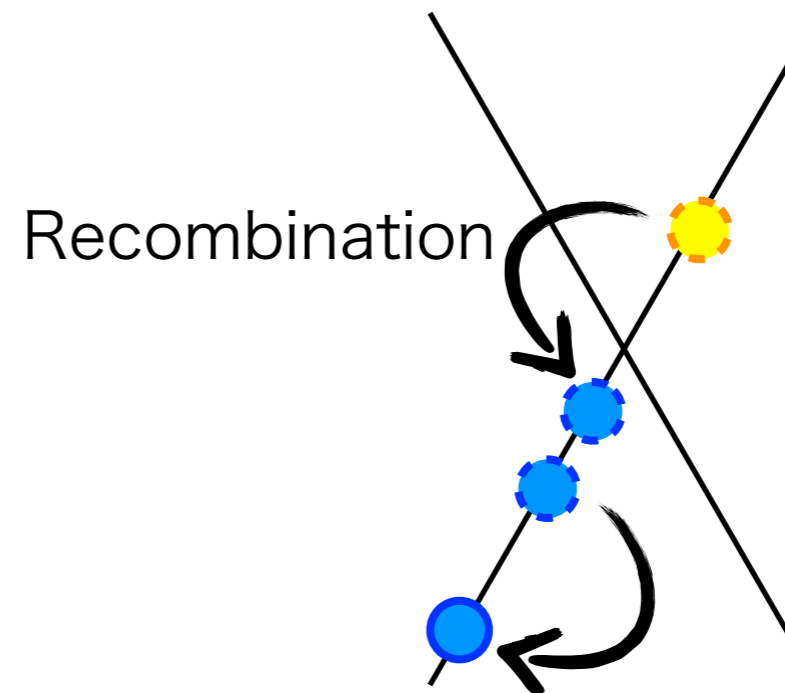


$$h \rightarrow h + h + e$$

$$e \rightarrow e + e + h$$

⇒ Carrier multiplication

Auger Recombination



$$h + h + e \rightarrow h$$

$$e + e + h \rightarrow e$$

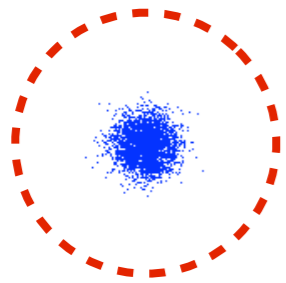
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F.Rana, Phys. Rev. B 76, 155431 (2007)

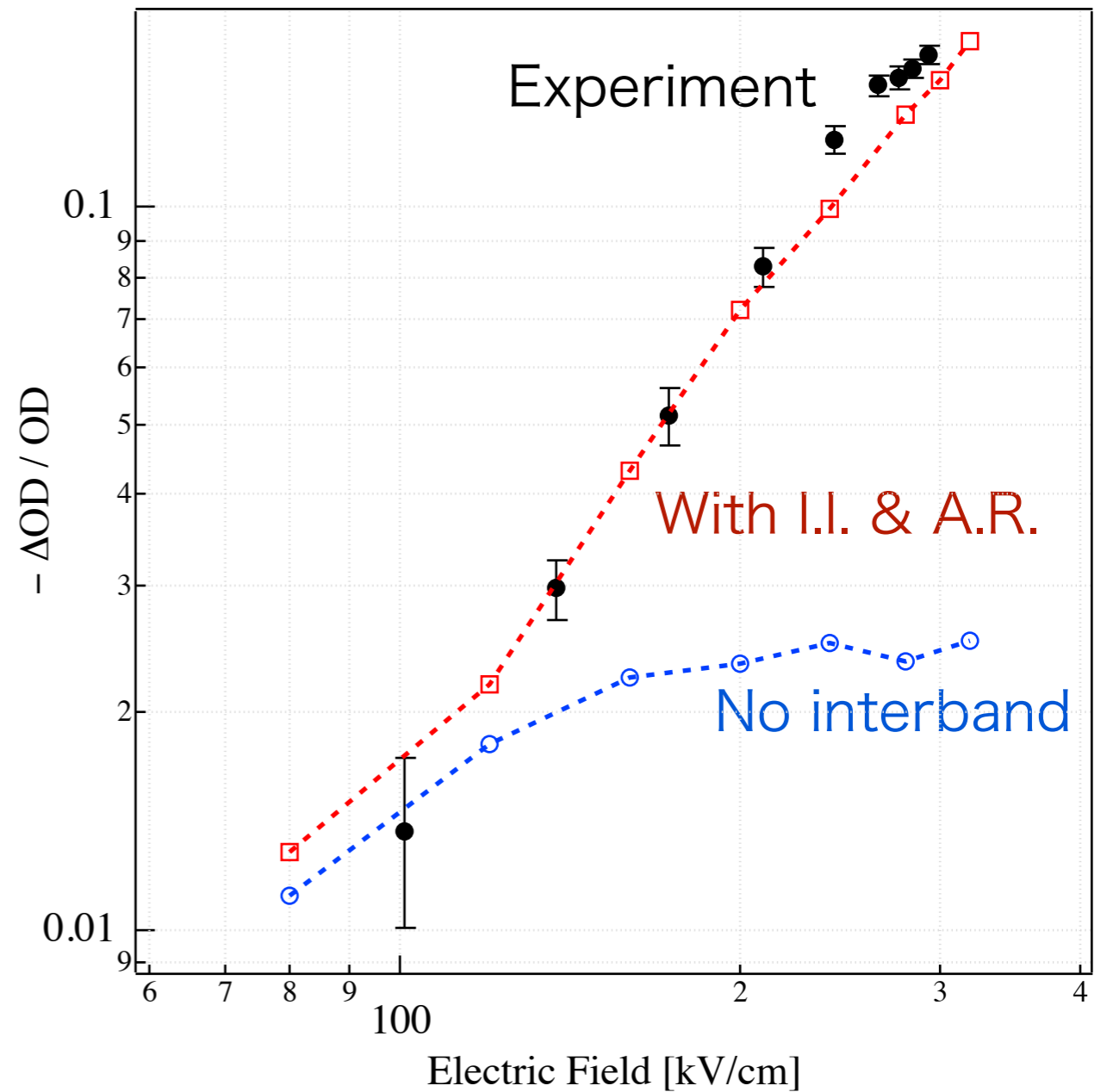
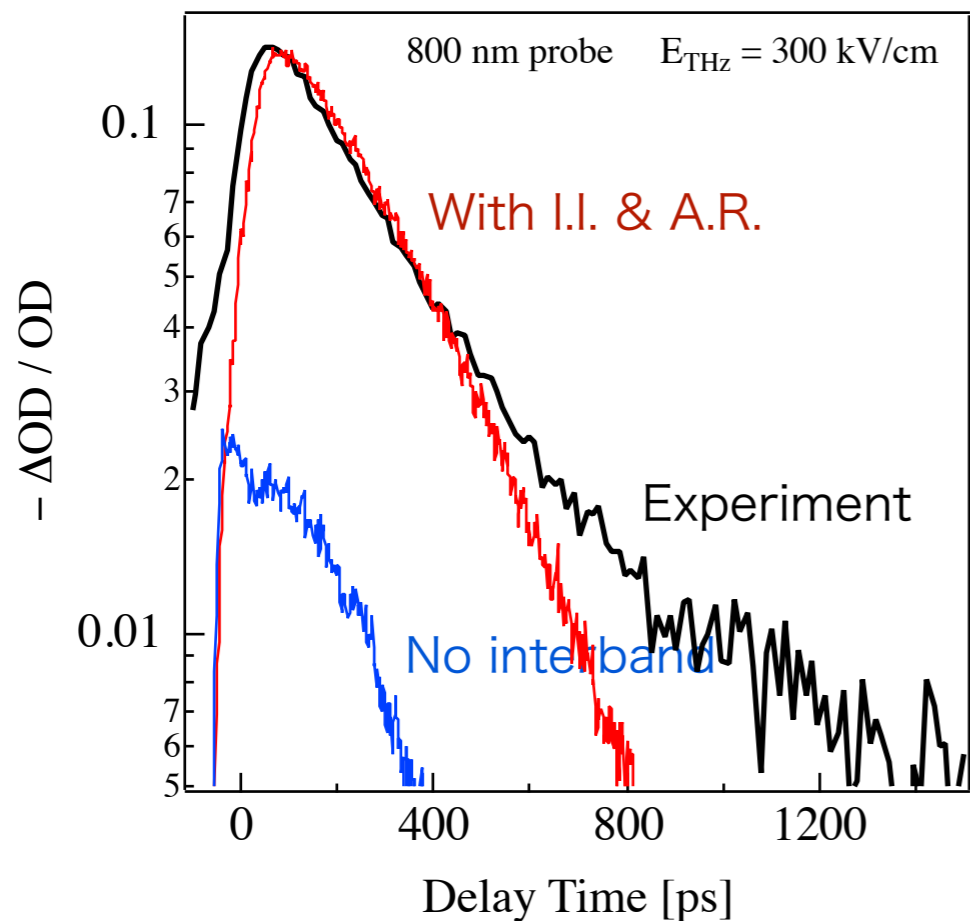
Simulation With Interband Carrier Scattering

$E_{\text{THz}} = 300 \text{ kV/cm}$
 $E_{\text{F}} = -200 \text{ meV}$

-100 fs

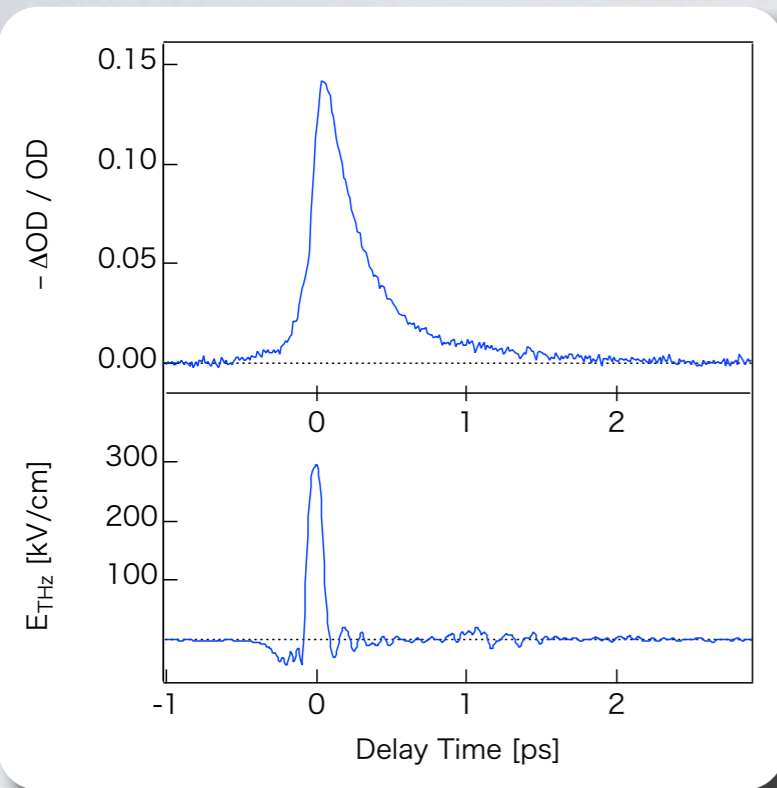


— New hole
— Initial hole



Carrier multiplication is important in extreme nonlinear regime.

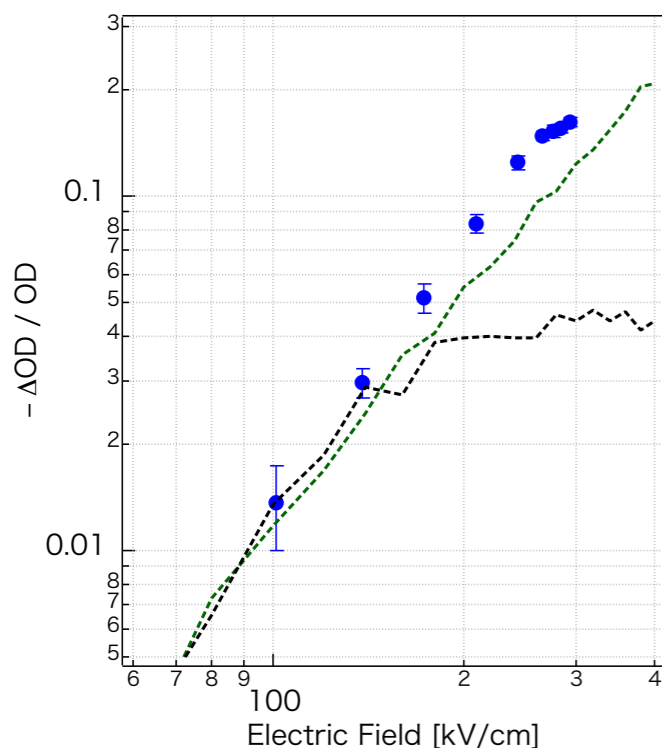
Summary : Thank you for your attention



THz-field driven carrier dynamics is investigated with transient absorption measurement in NIR region.

Huge THz induced transparency over 14% is observed at 800 nm.

Carrier multiplication with interband carrier scattering is essential in high field dynamics.





Backup Slides



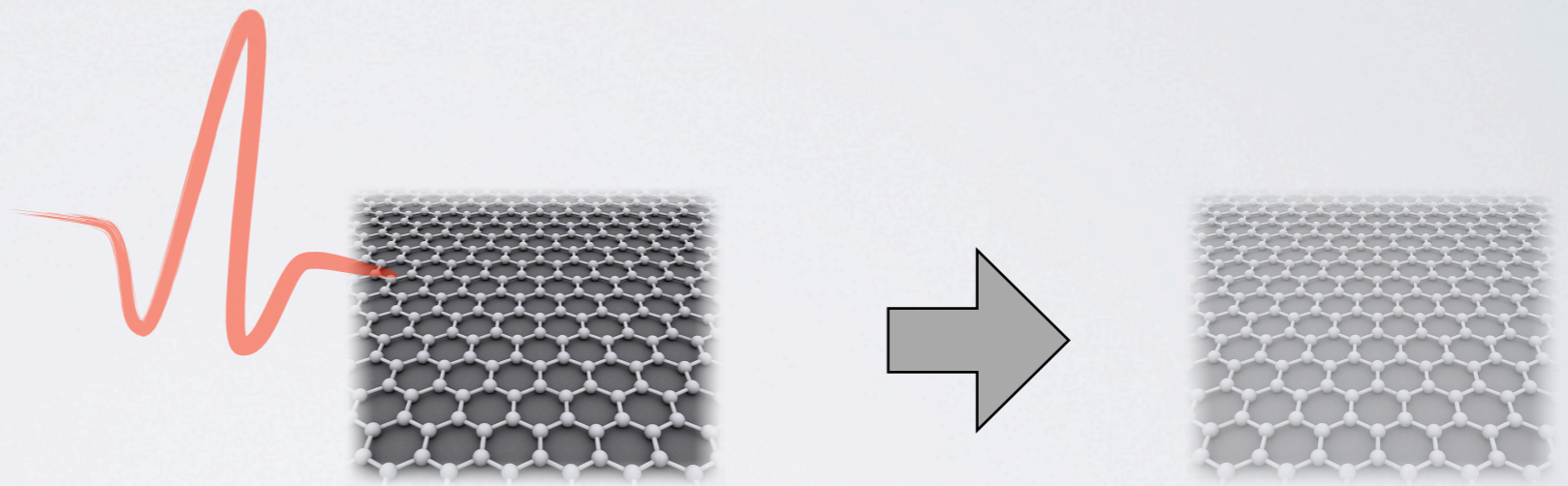
Conduction Band

The diagram illustrates the energy bands of a semiconductor. It features two thick black parabolic curves. The upper curve, representing the conduction band, opens upwards and is labeled 'Conduction Band'. The lower curve, representing the valence band, opens downwards and is labeled 'Valence Band'. There is a clear vertical gap between the minimum of the conduction band and the maximum of the valence band, representing the band gap.

Valence Band

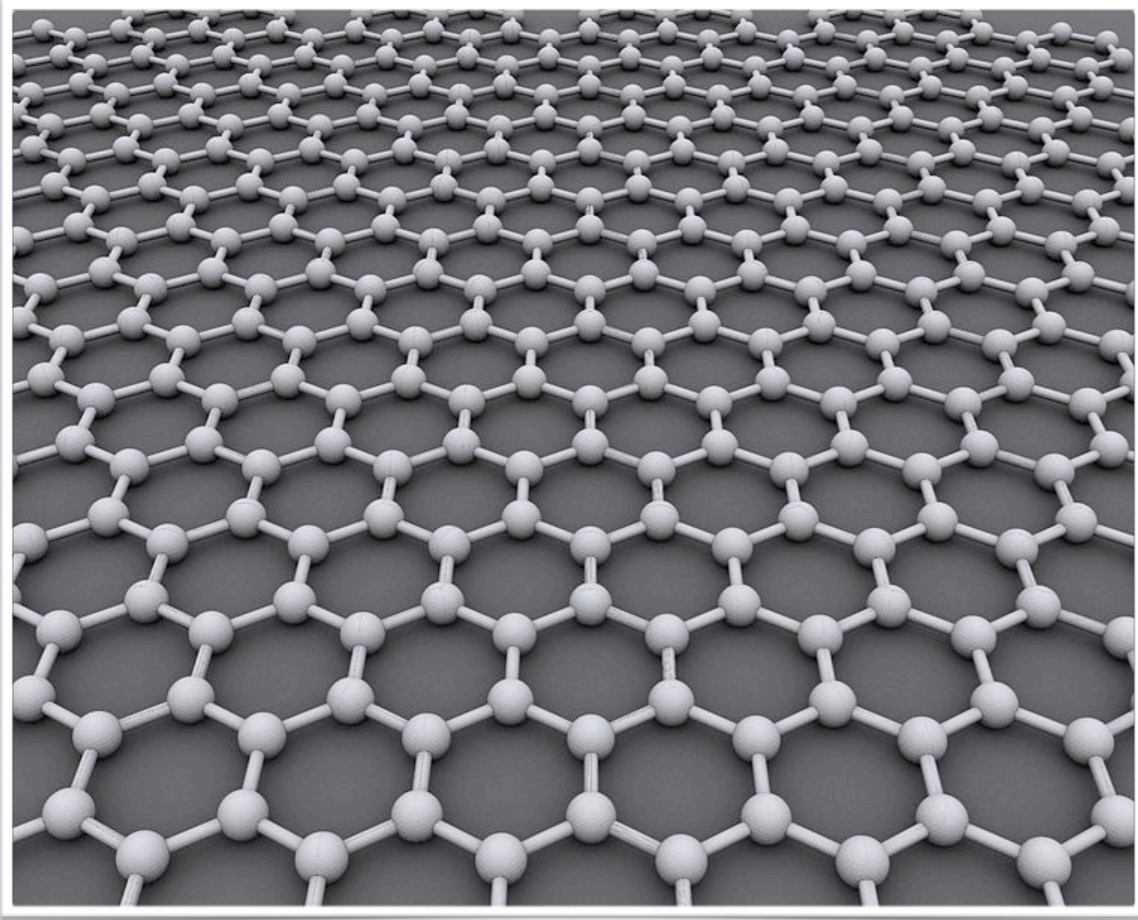
Experiment

- THz pump - NIR transient absorption measurement.
- Induced transparency over **14 %** at 800 nm by THz wave.



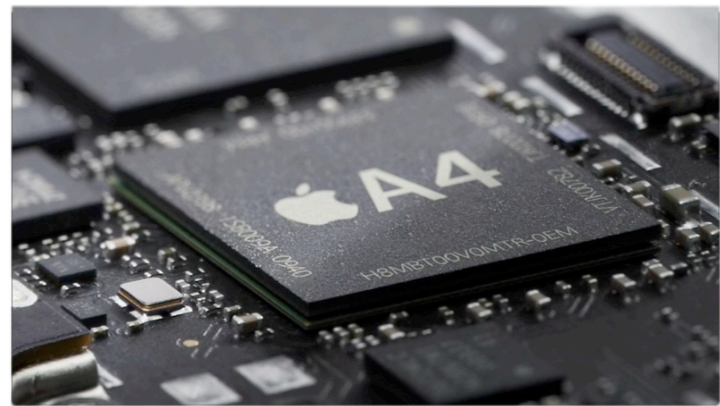
Simulation

- Monte Carlo simulation including carrier-carrier and optical phonon scattering.



- Monoatomic layer carbon material
- Characteristic band structure
 - ➔ Massless electrons
- High conductivity at $T = 300\text{ K}$
 - ➔ Promising material for ultra high-speed nano devices
Graphene transistor, graphene LSI, etc...

High Field Carrier Dynamics in Graphene



Typical Nano Device

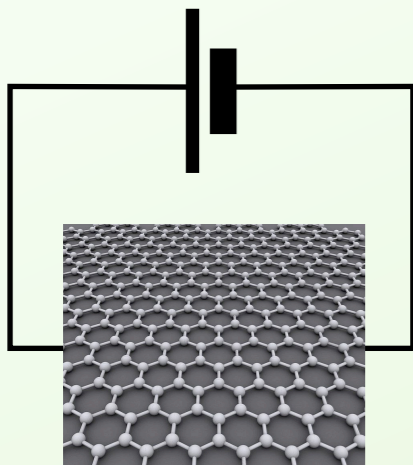
Voltage ~ 1 V

Length ~ 10 nm

→ Electric Field ~ 0.1 V/nm = 1 MV/cm

High electric field \Rightarrow Extremely non-equilibrium state

For future applications such as high-speed devices, ultrafast carrier dynamics under high field is important.



Under high DC field

J. Phys. Cond. Mat. 21 (2009) 344201

• Drift velocity saturates by phonon scattering.

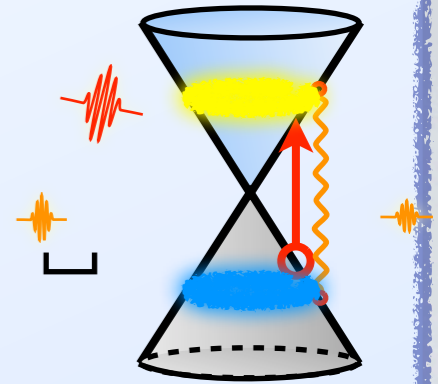
→ Averaged dynamics

→ Difficult to understand ultrafast response

High Field Carrier Dynamics in Graphene

According to optical-pump optical-probe measurement

- Carrier-carrier scattering : 2 ~ 10 fs
- Optical phonon emission : 100 ~ 400 fs



Requirement

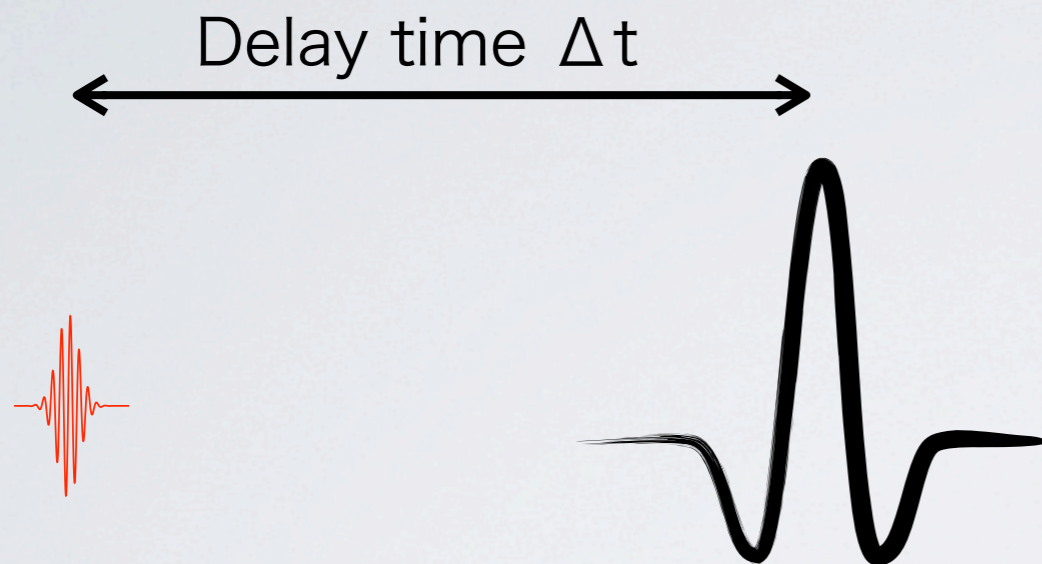
To investigate ultrafast carrier dynamics under high field, time resolution better than 100 fs is required.

Solution

Near-infrared transient absorption measurement
with **single cycle THz pulse** excitation

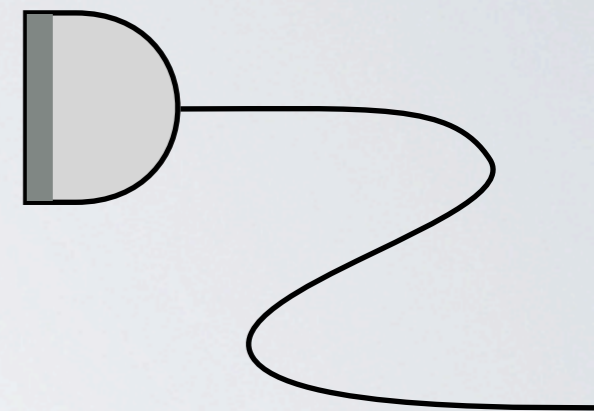
- Apply high electric field in very short time duration
- High time resolution ~ 50 fs

Experimental Condition



Transmission
measurement

$$\Delta OD = OD_{\text{THz}} - OD_{\text{ref}}$$



Probe pulse

- 800 nm
- Pulse width : 50 fs

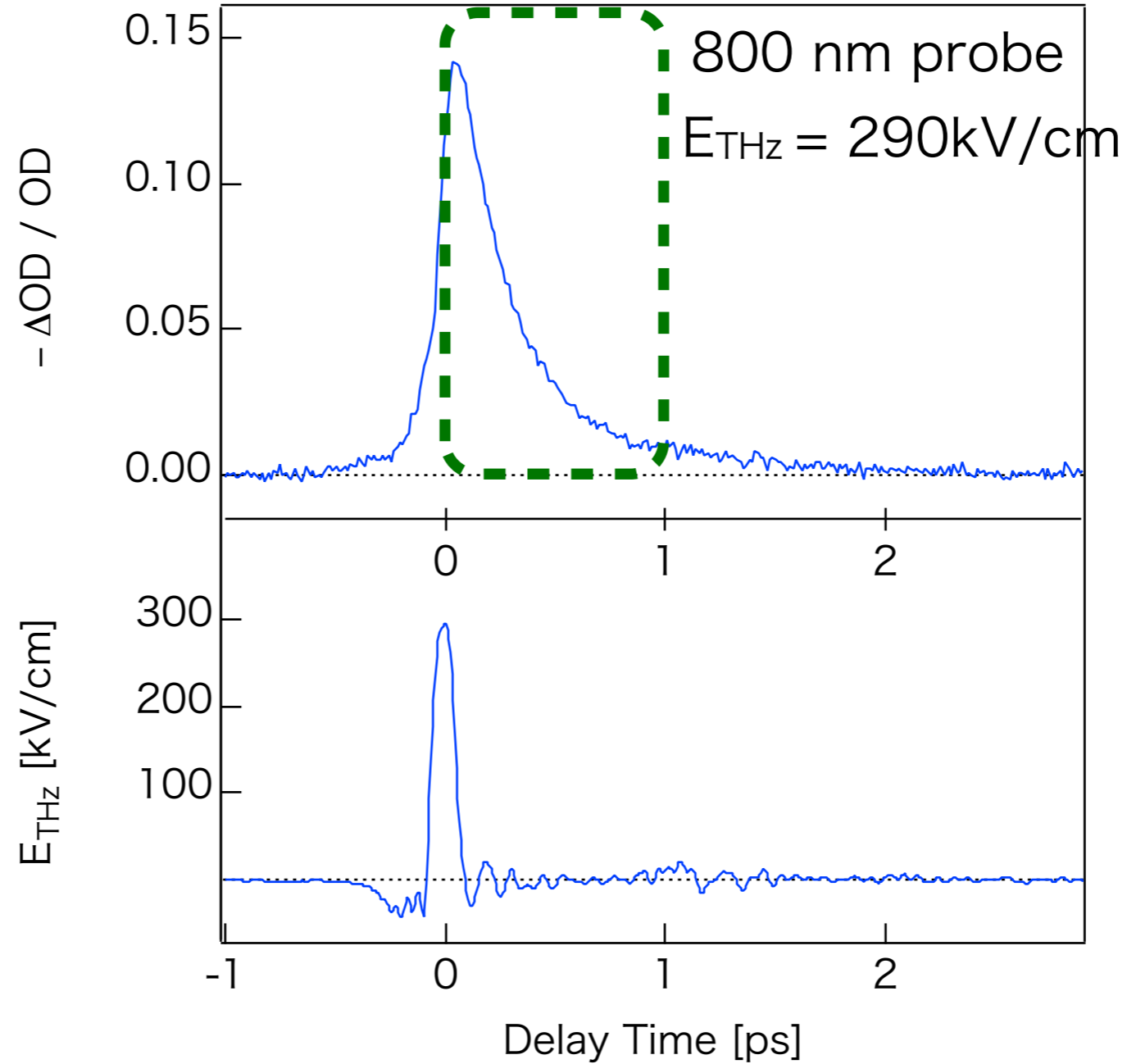
THz pulse

- Air plasma method
- Field : ~ 300 kV/cm
- Pulse width : 100 fs

Sample

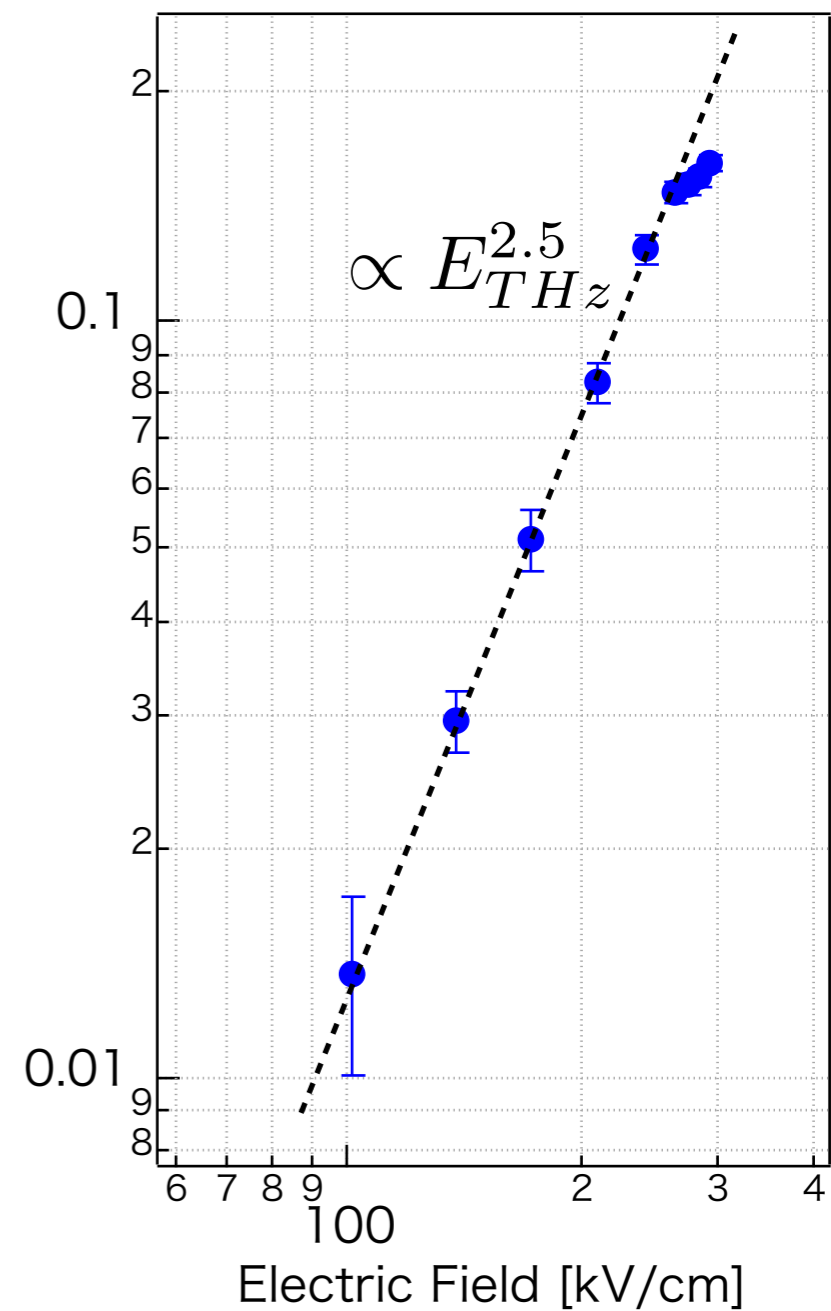
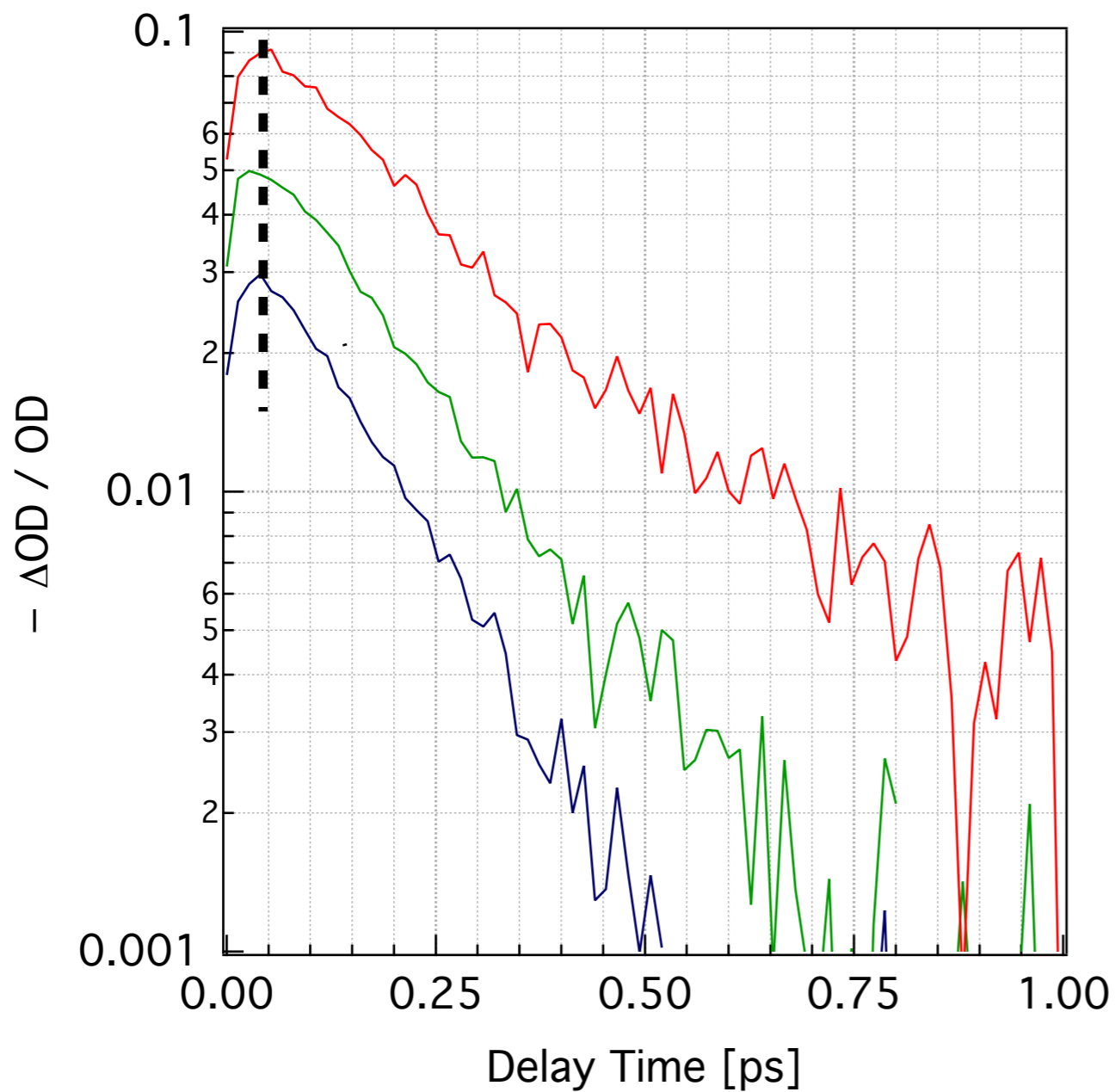
- CVD growth graphene
- On SiO₂ substrate
- Transmission 97 % at 800 nm
- T = 294 K

Results : Typical THz-induced Transparency in Near-infrared Region

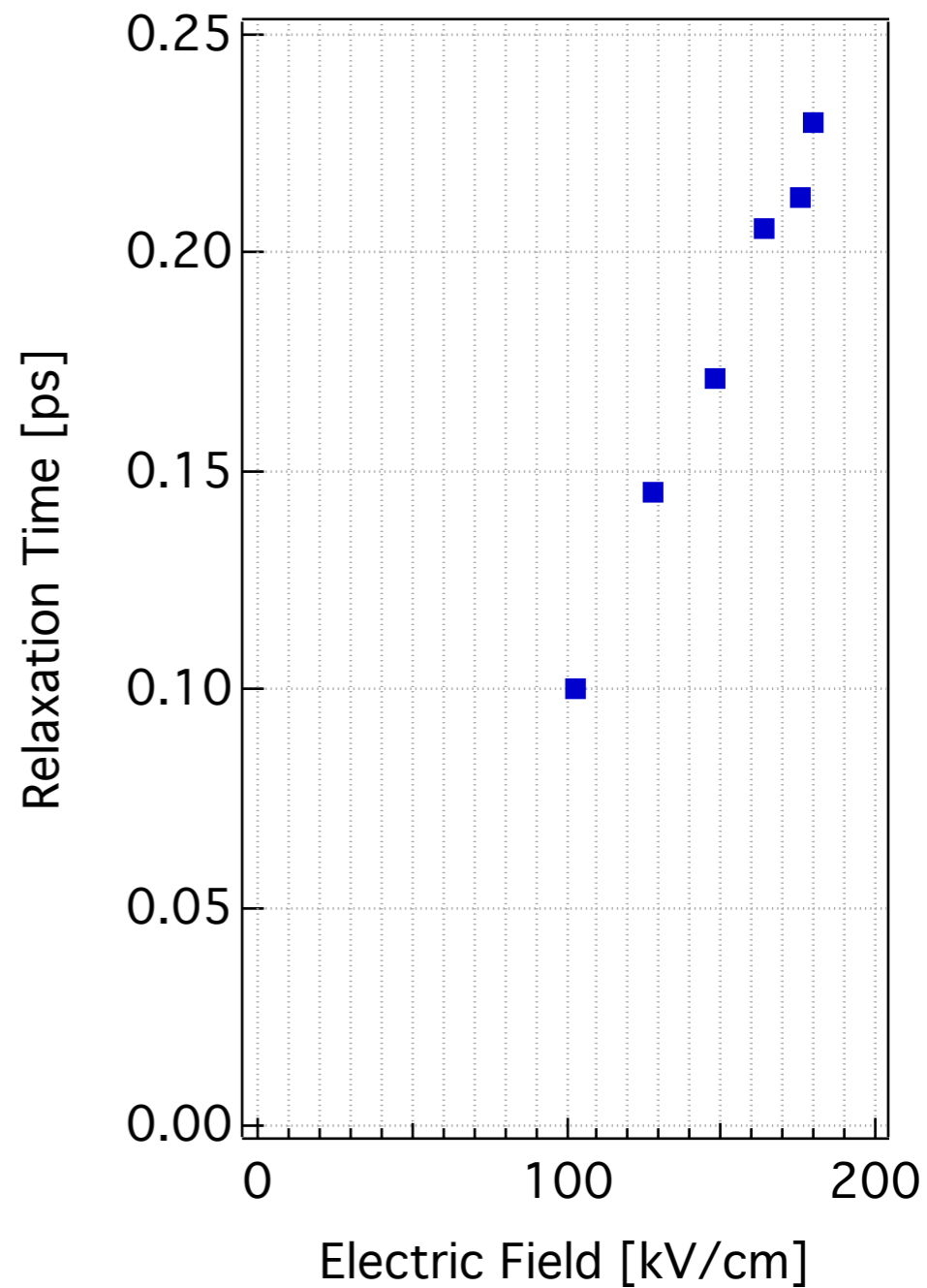
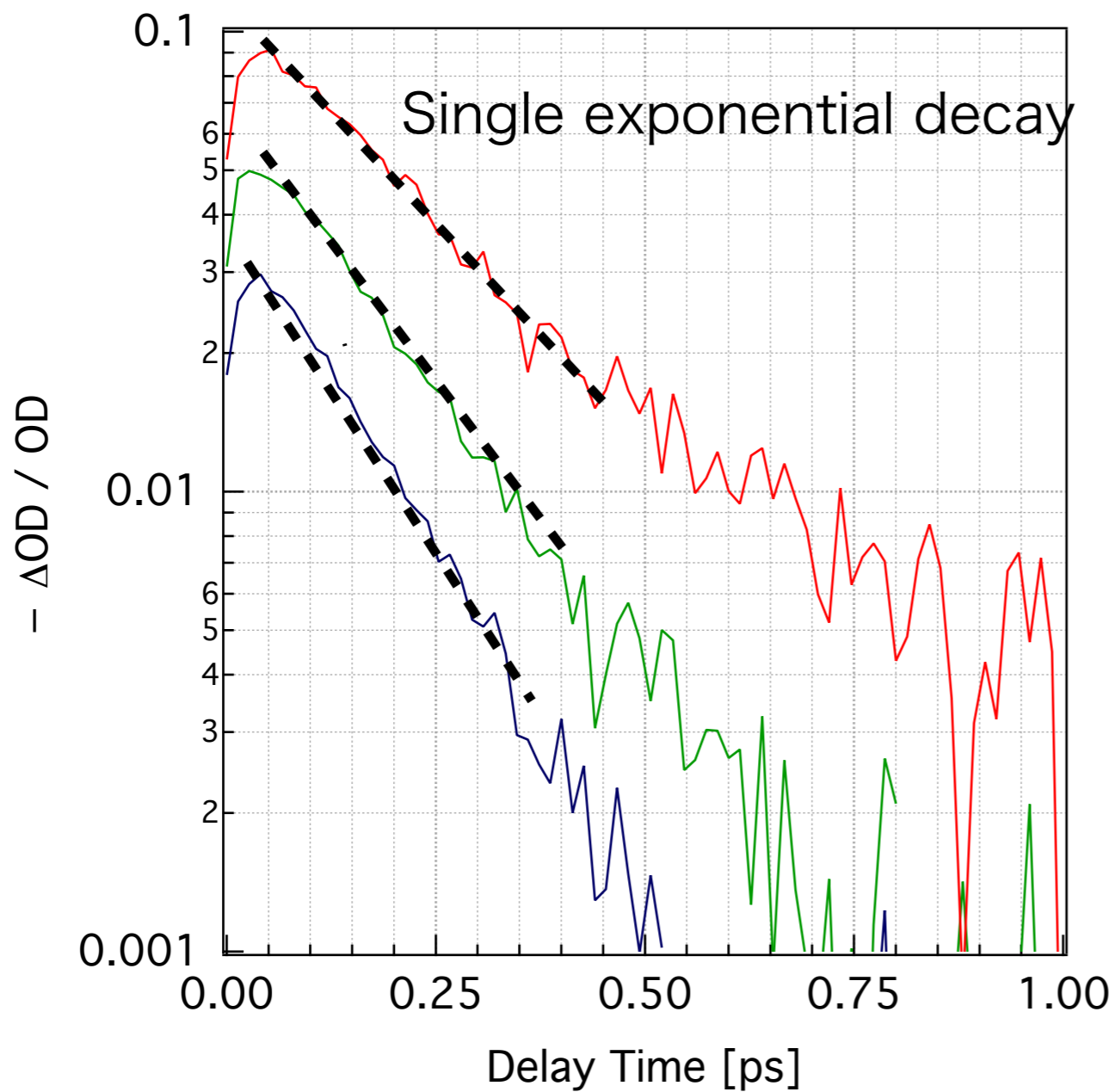


THz induced transparency over 14% at 800 nm

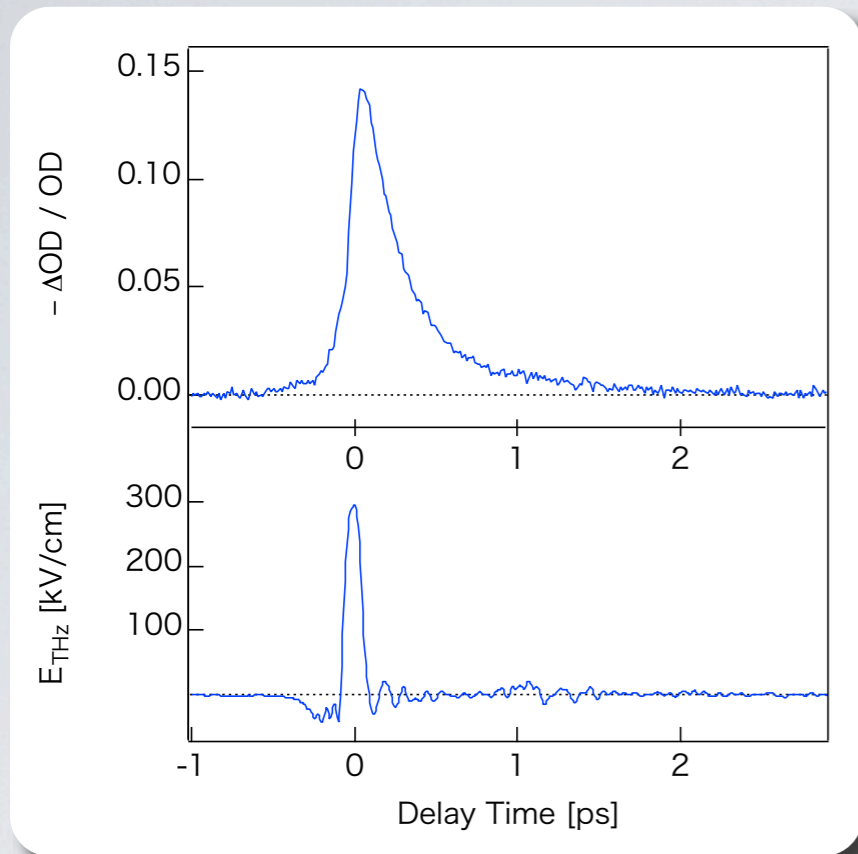
Results : Field Dependence



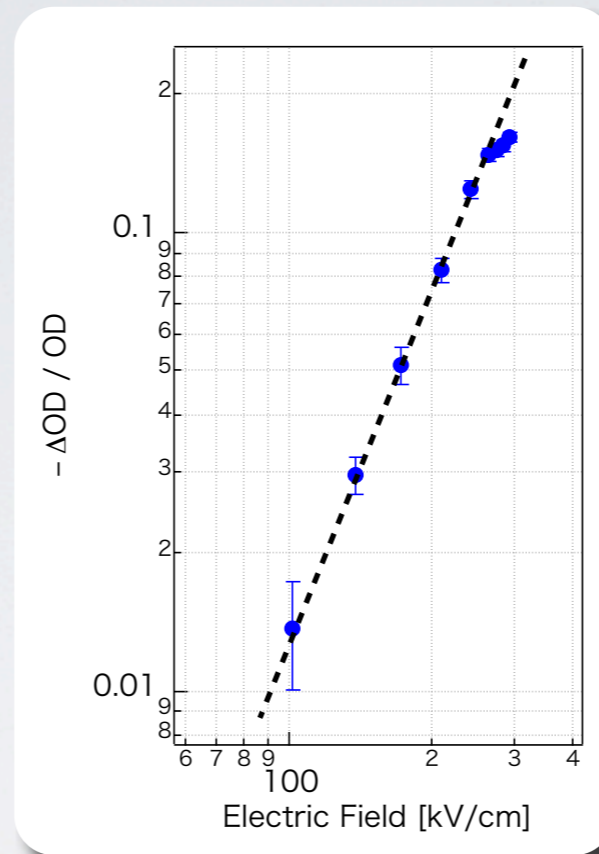
Results : Field Dependence



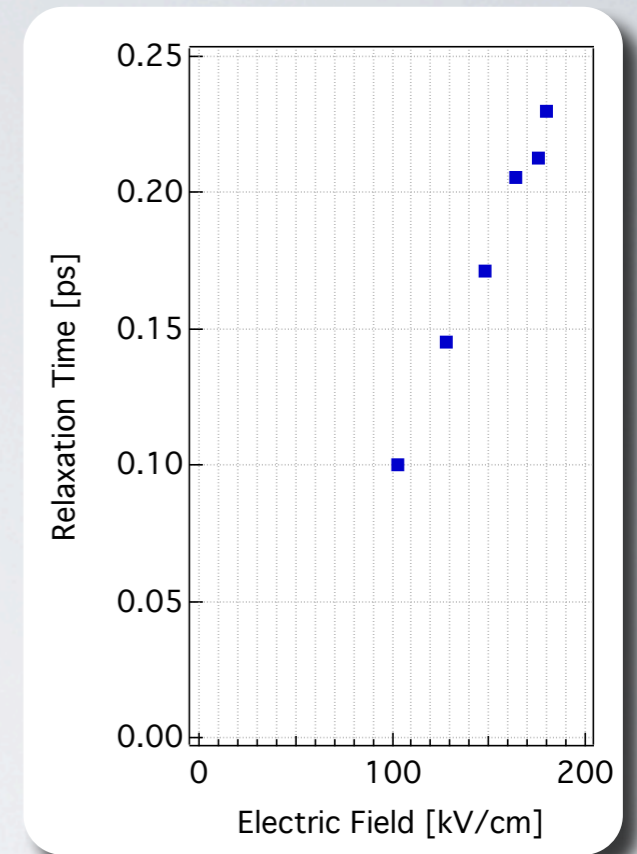
Summary : Experimental Results



THz induced
transparency over
14% at 800 nm



$$-\Delta OD / OD \propto E_{THz}^{2.5}$$

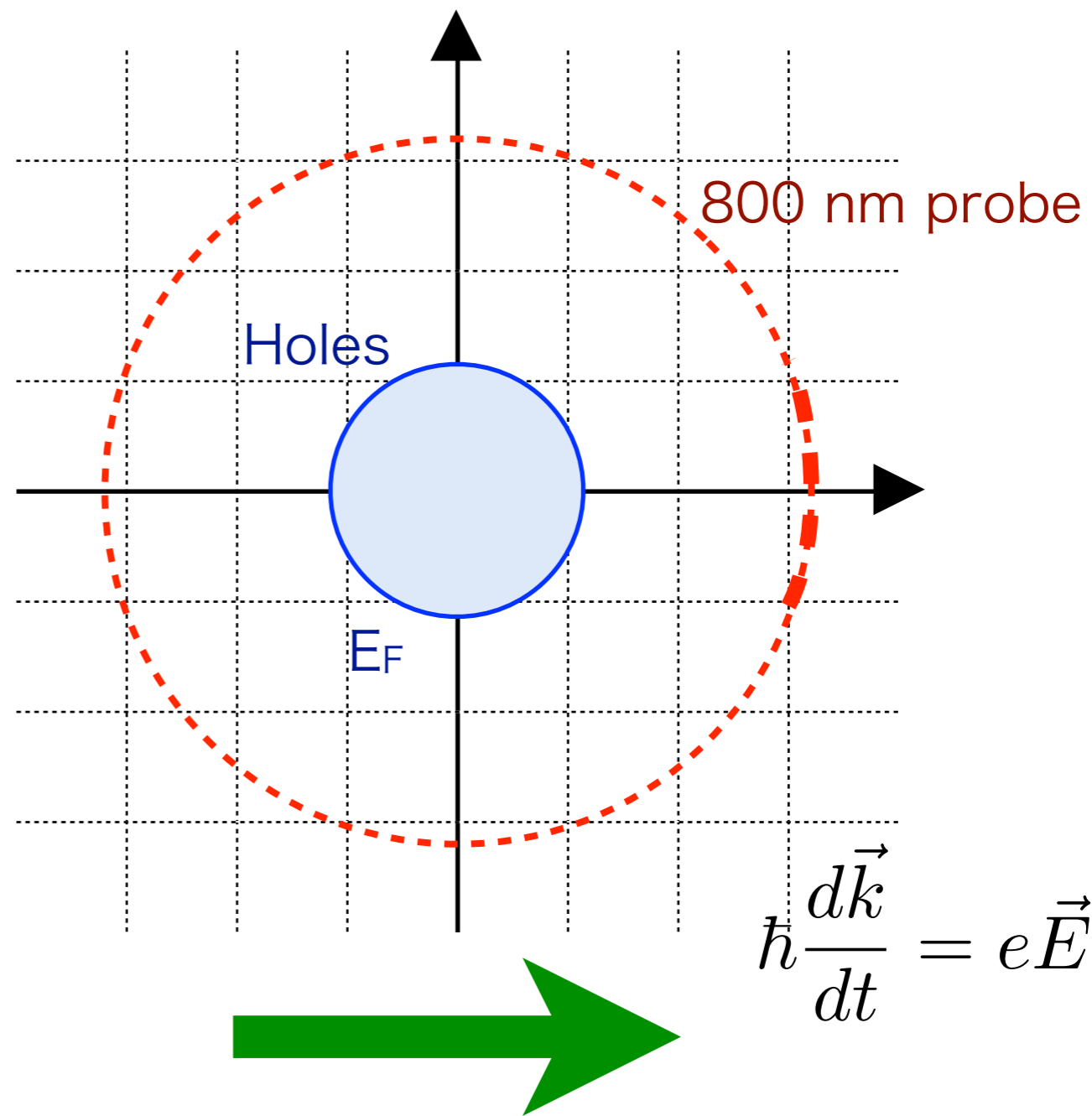
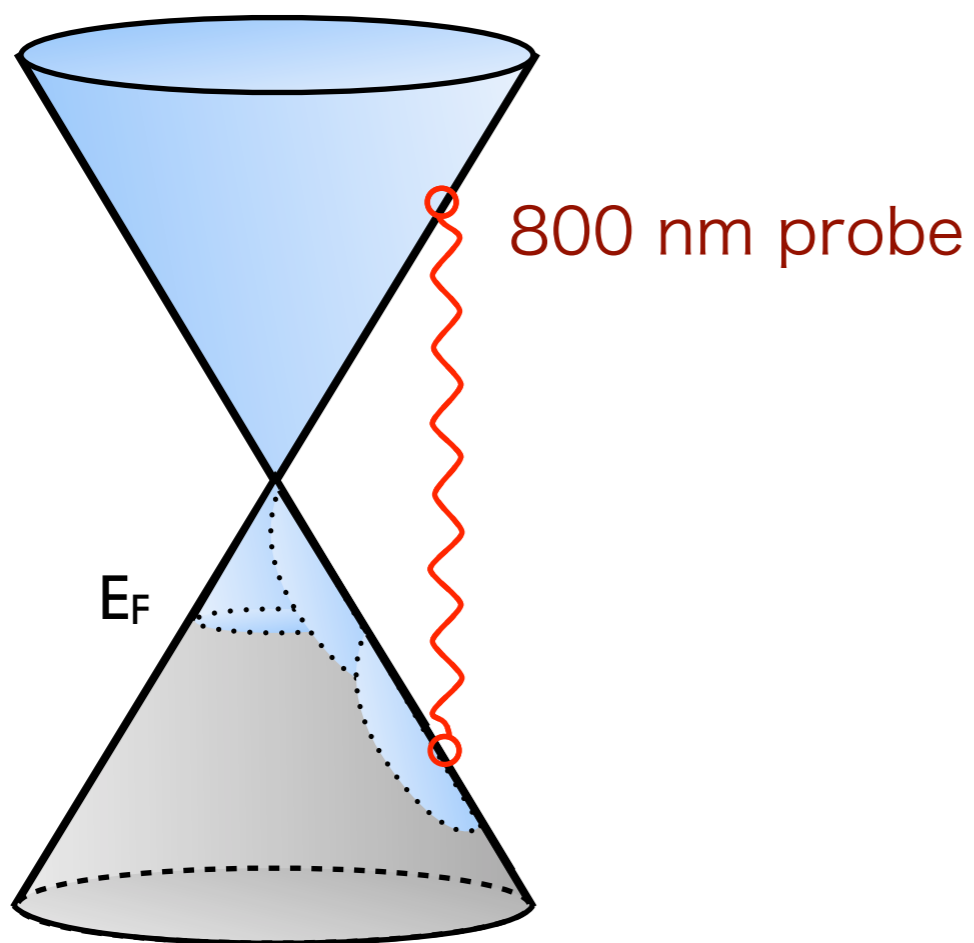


High field
↓
Longer
relaxation time

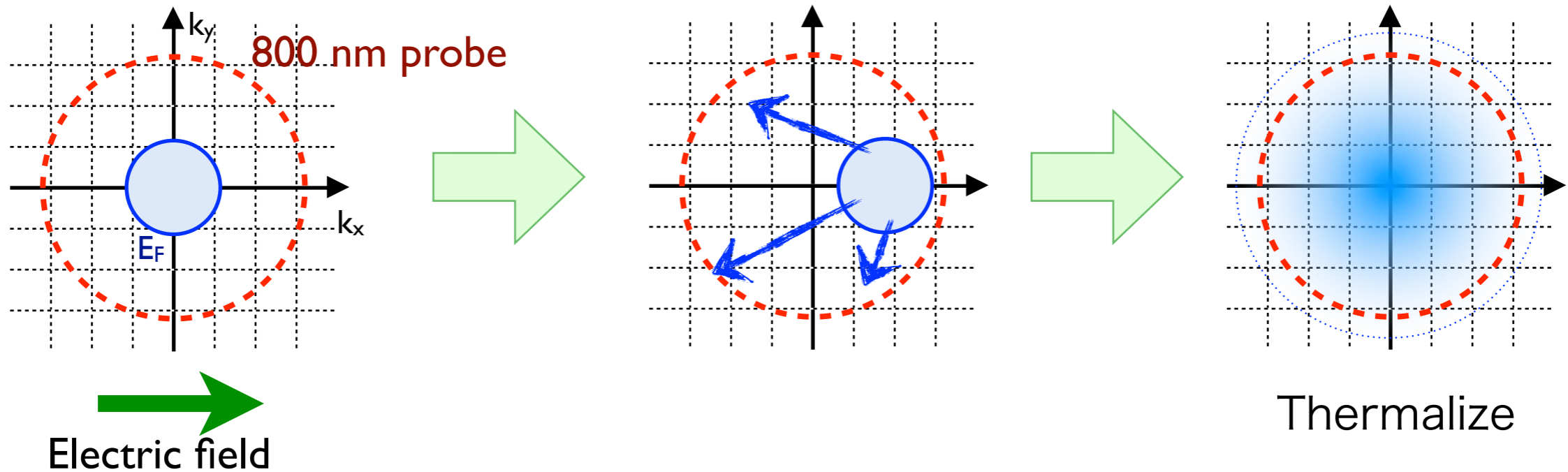
→ Monte Carlo simulation

Origin of THz Induced Transparency

Hole distribution in the **momentum space**



Carrier Scattering Process



Boltzmann equation for $f(\vec{k}, t)$

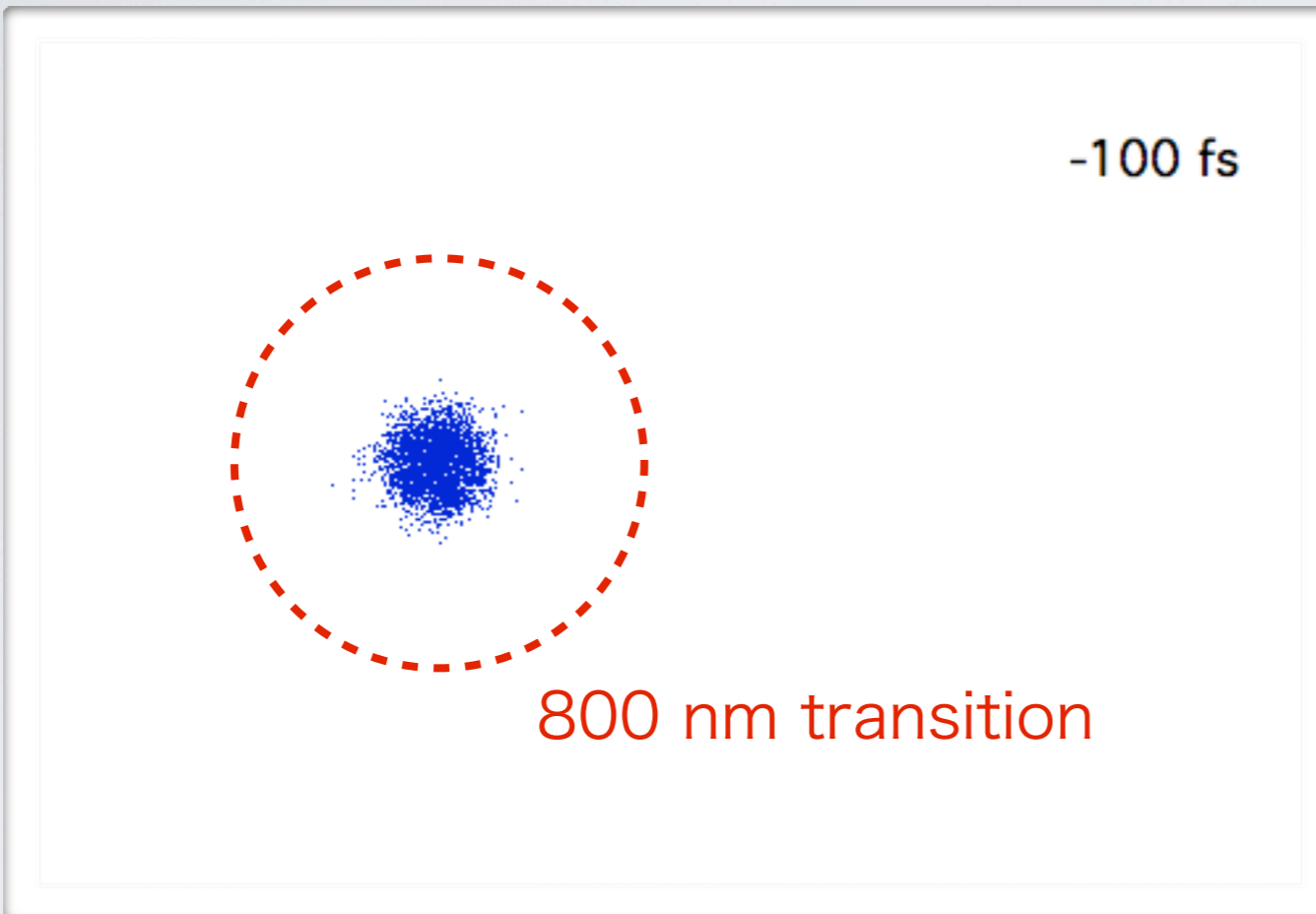
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- No screening
- $\omega_{TO} = 160 \text{ meV}$

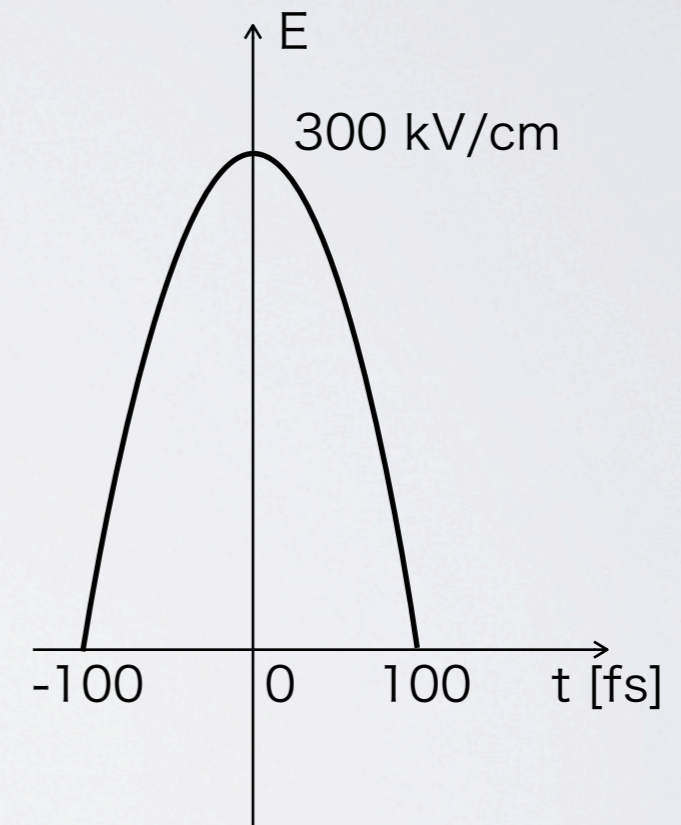
Time Development of Hole Distribution

Hole distribution in the momentum space



$$E_{\text{THz}} = 300 \text{ kV/cm}$$

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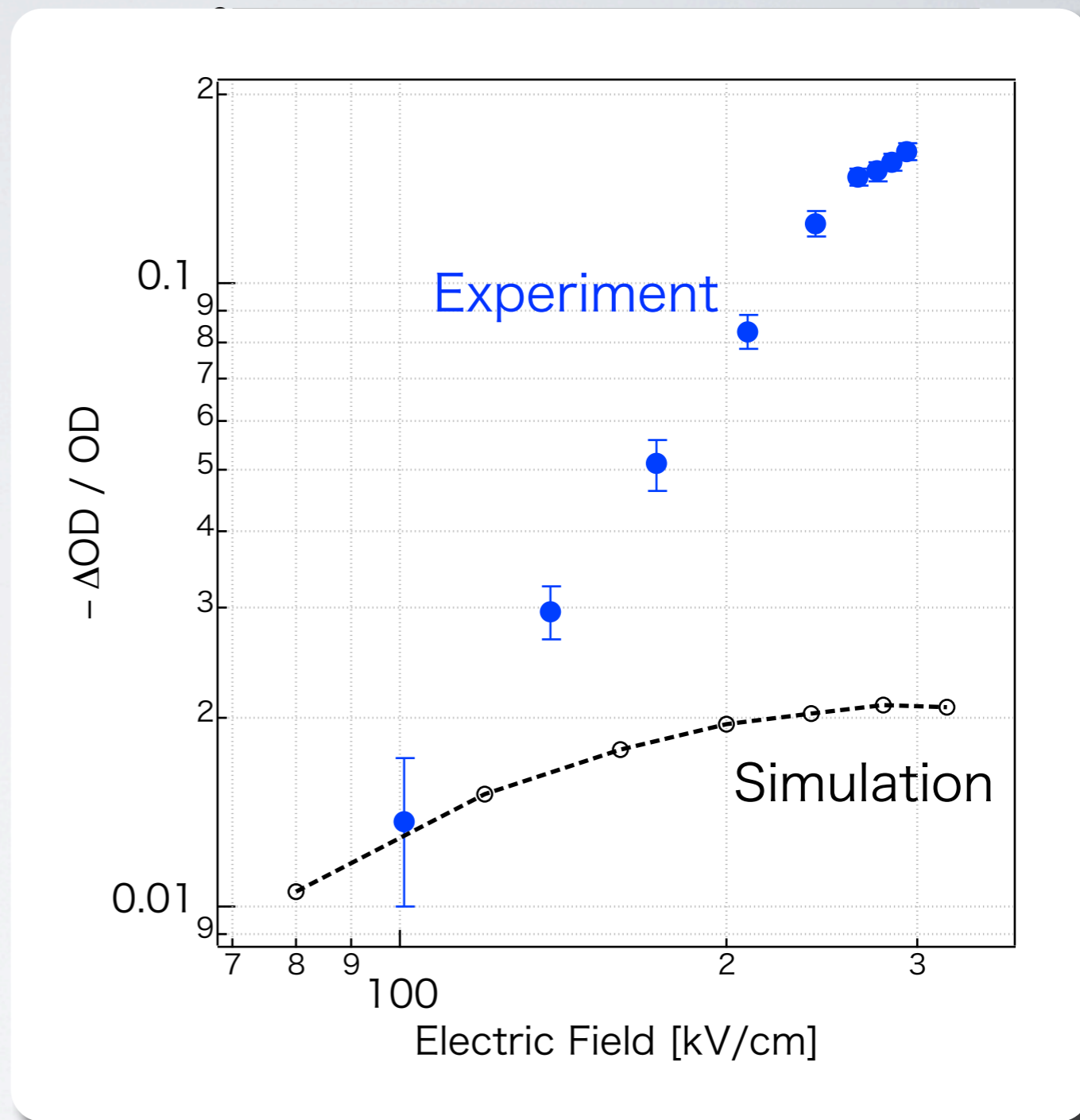
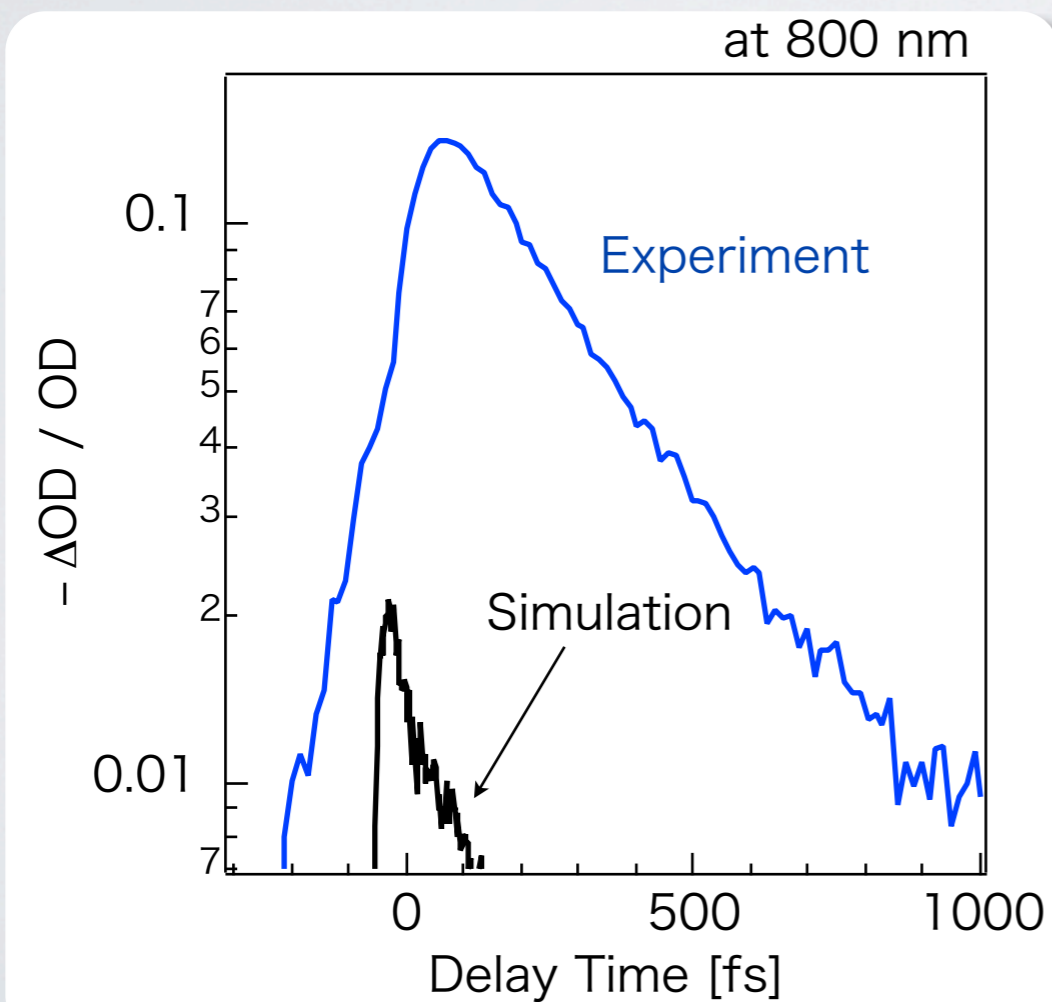
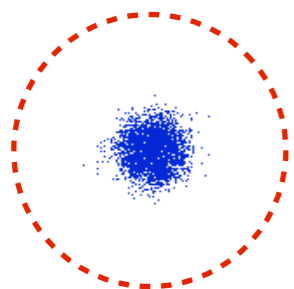


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Simulation Results

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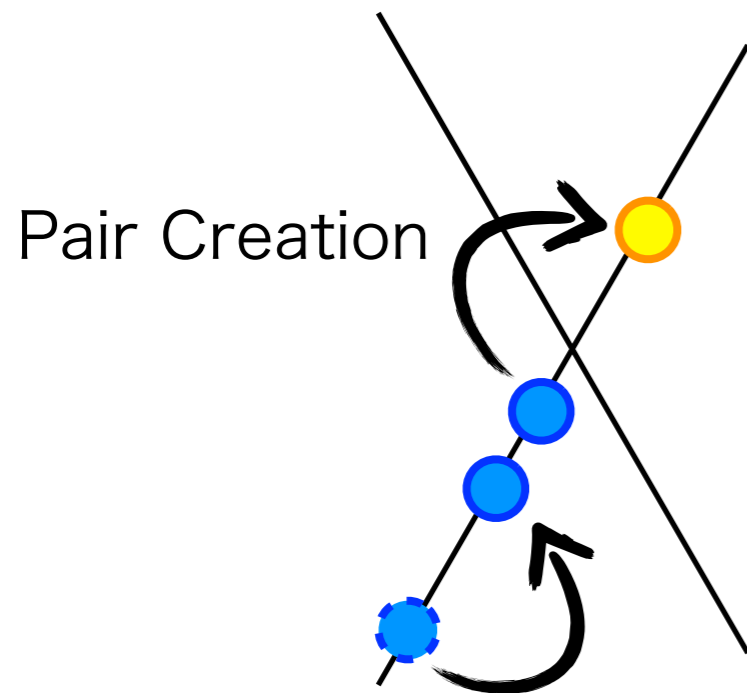
-100 fs



Carriers may not be enough.

Interband Carrier Scattering

Impact Ionization

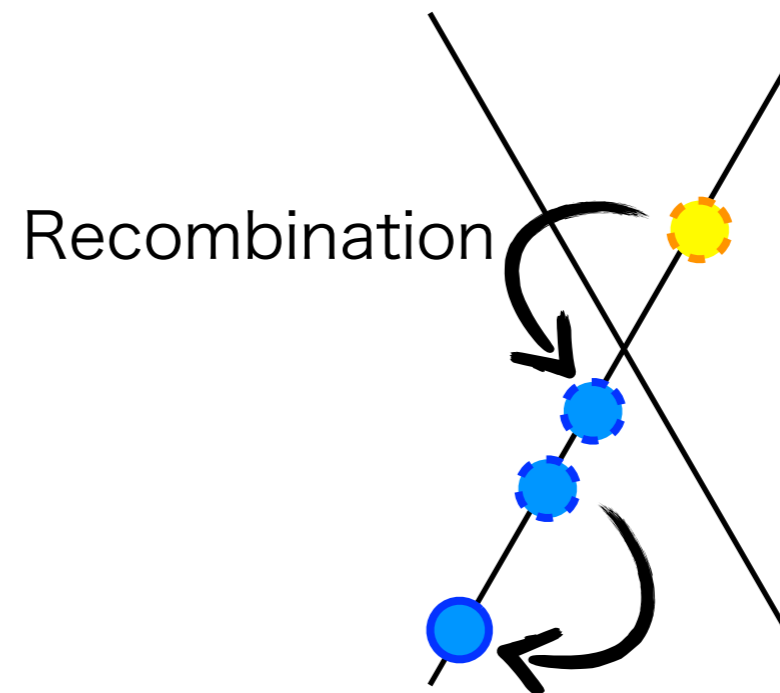


$$h \rightarrow h + h + e$$

$$e \rightarrow e + e + h$$

⇒ Carrier multiplication

Auger Recombination



$$h + h + e \rightarrow h$$

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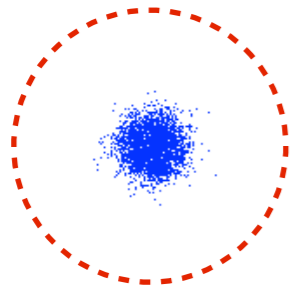
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F.Rana, Phys. Rev. B 76, 155431 (2007)

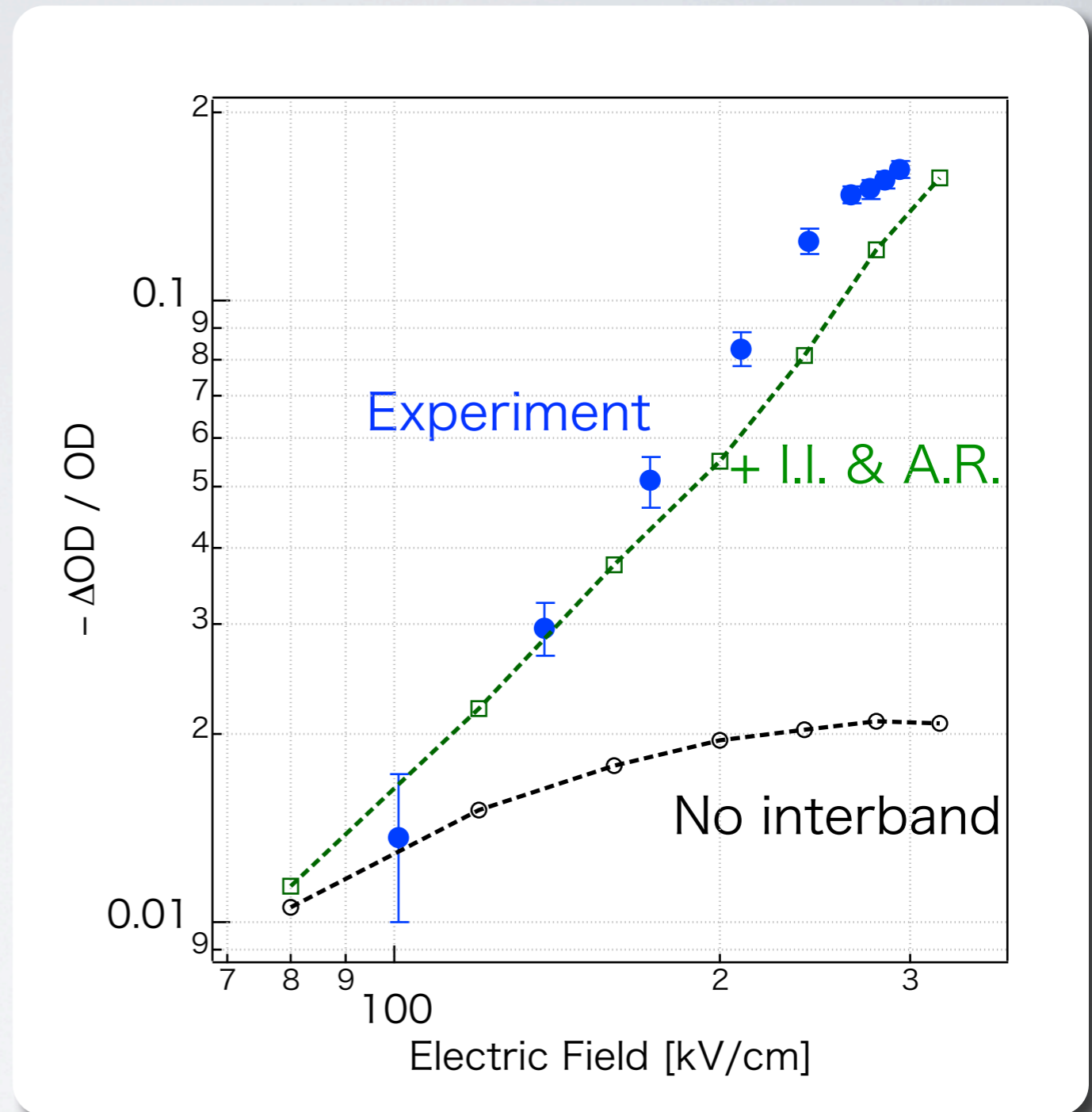
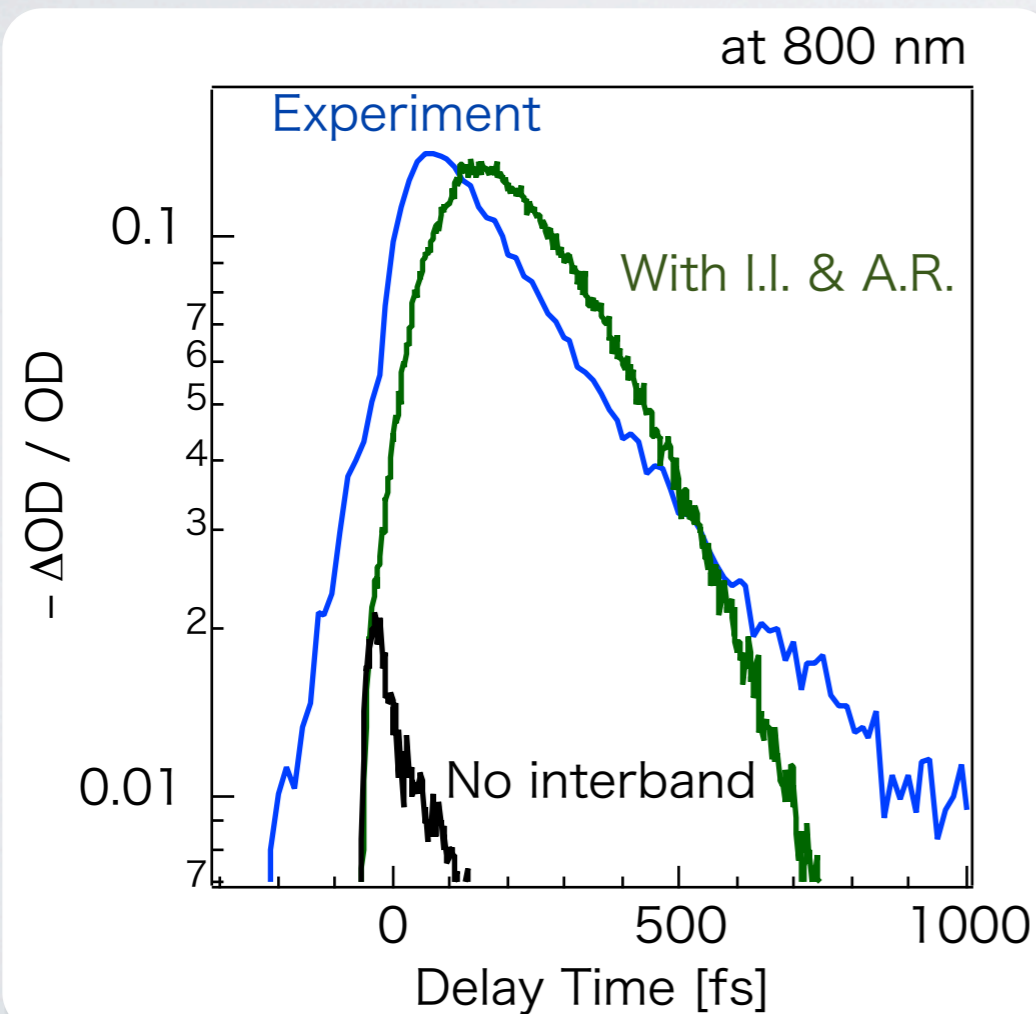
Simulation With Interband Carrier Scattering

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-100 fs

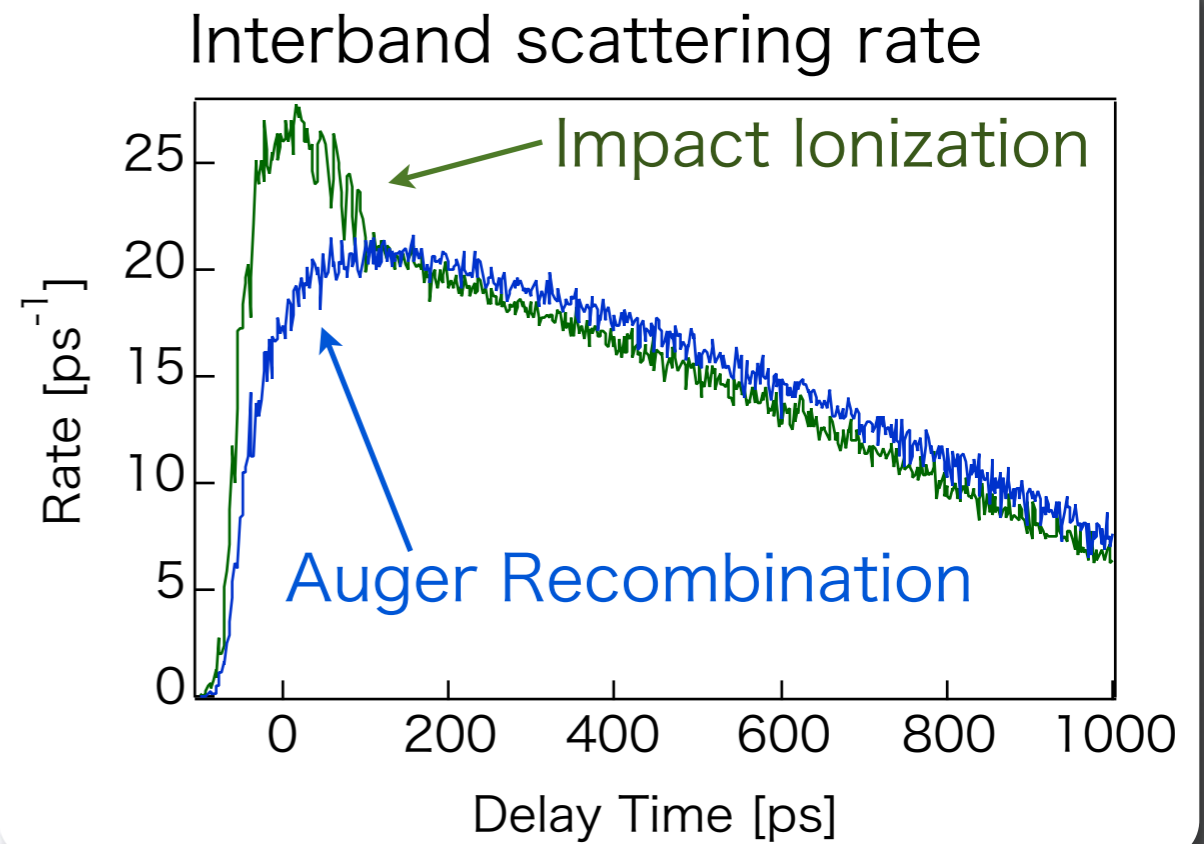
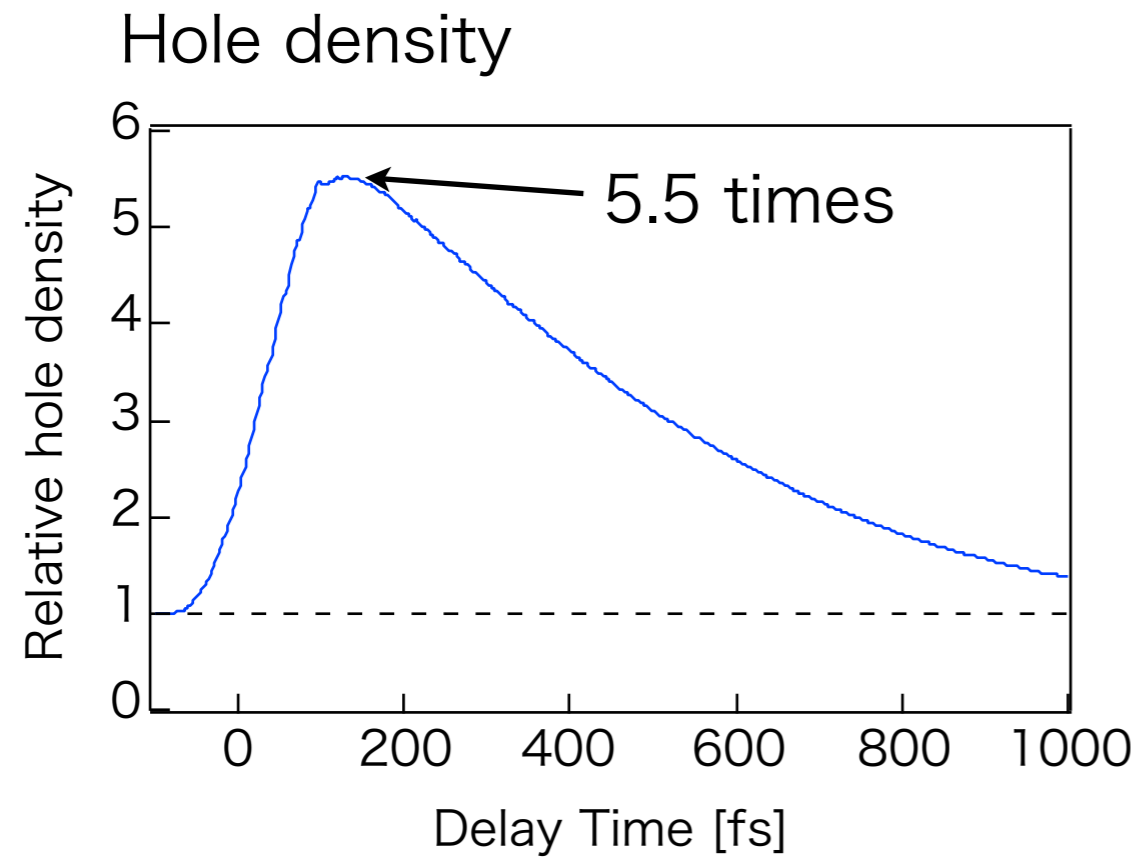


— New hole
— Initial hole

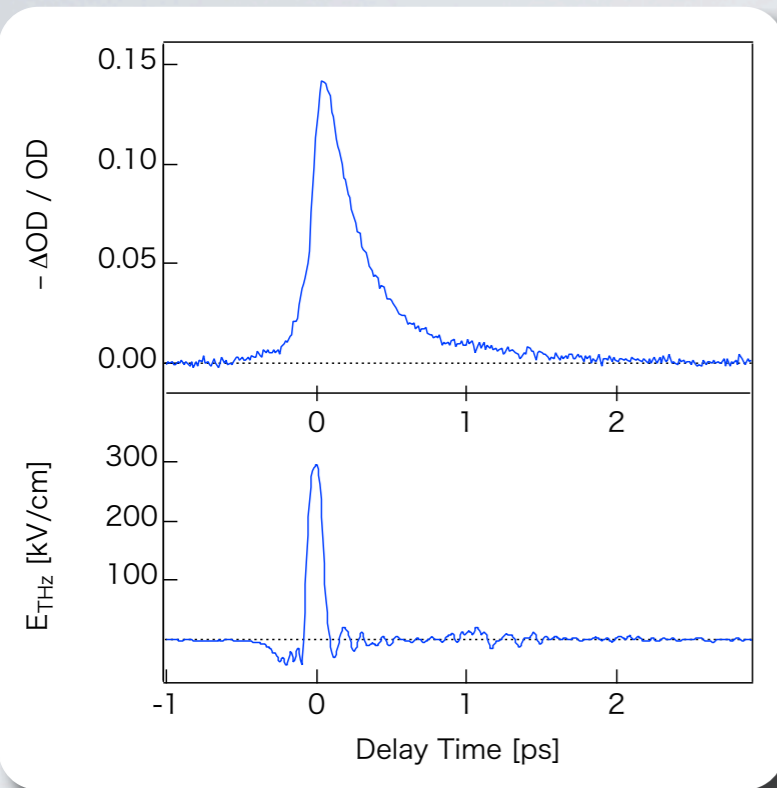


Carrier multiplication is important in high field dynamics.

Simulation : Carrier Multiplication



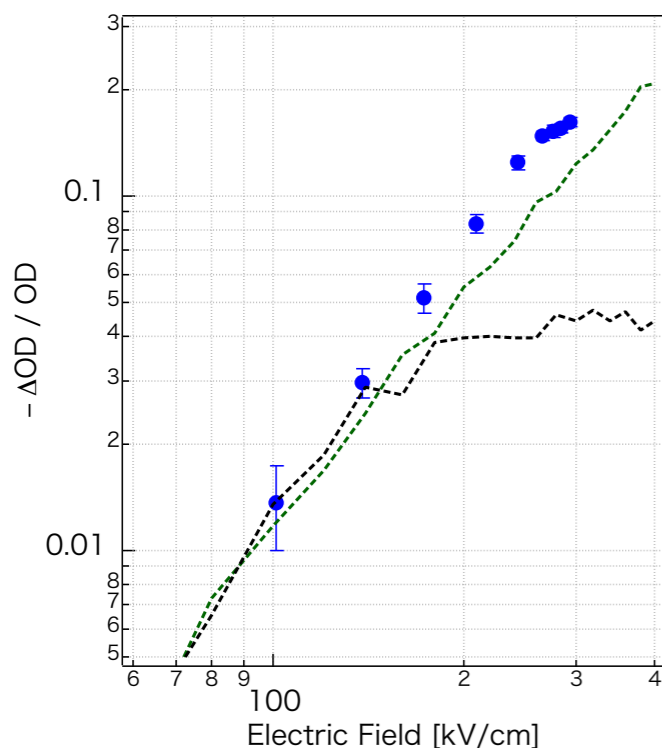
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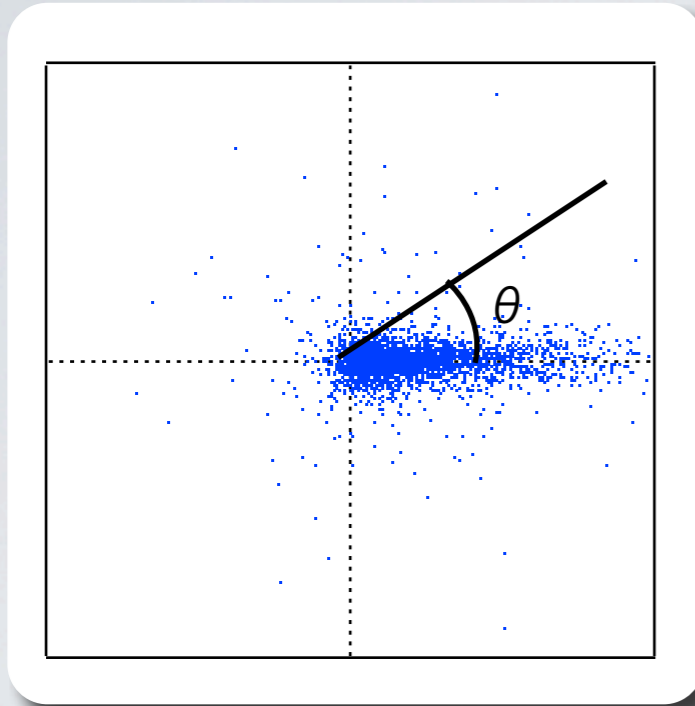
Huge THz induced transparency over 14% is observed at 800 nm.

Carrier multiplication with interband carrier scattering is essential in high field dynamics.

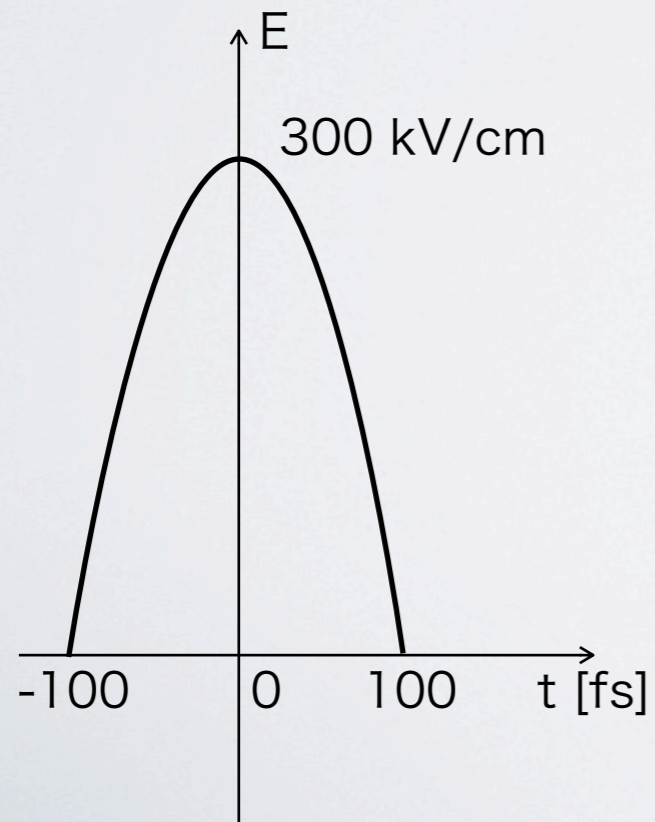
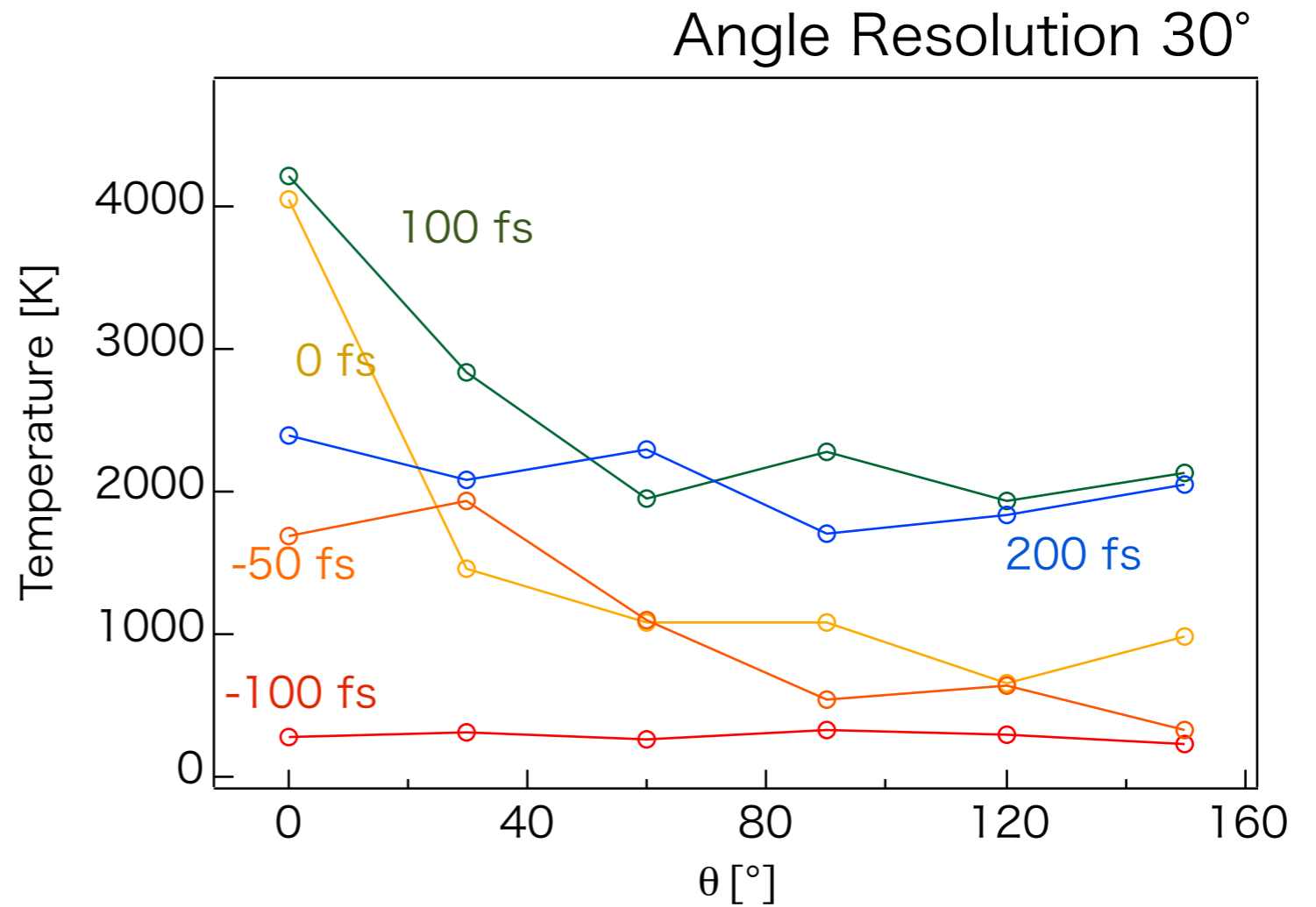


Angle Dependent Temperature

Hole distribution



Angle Dependent Temperature

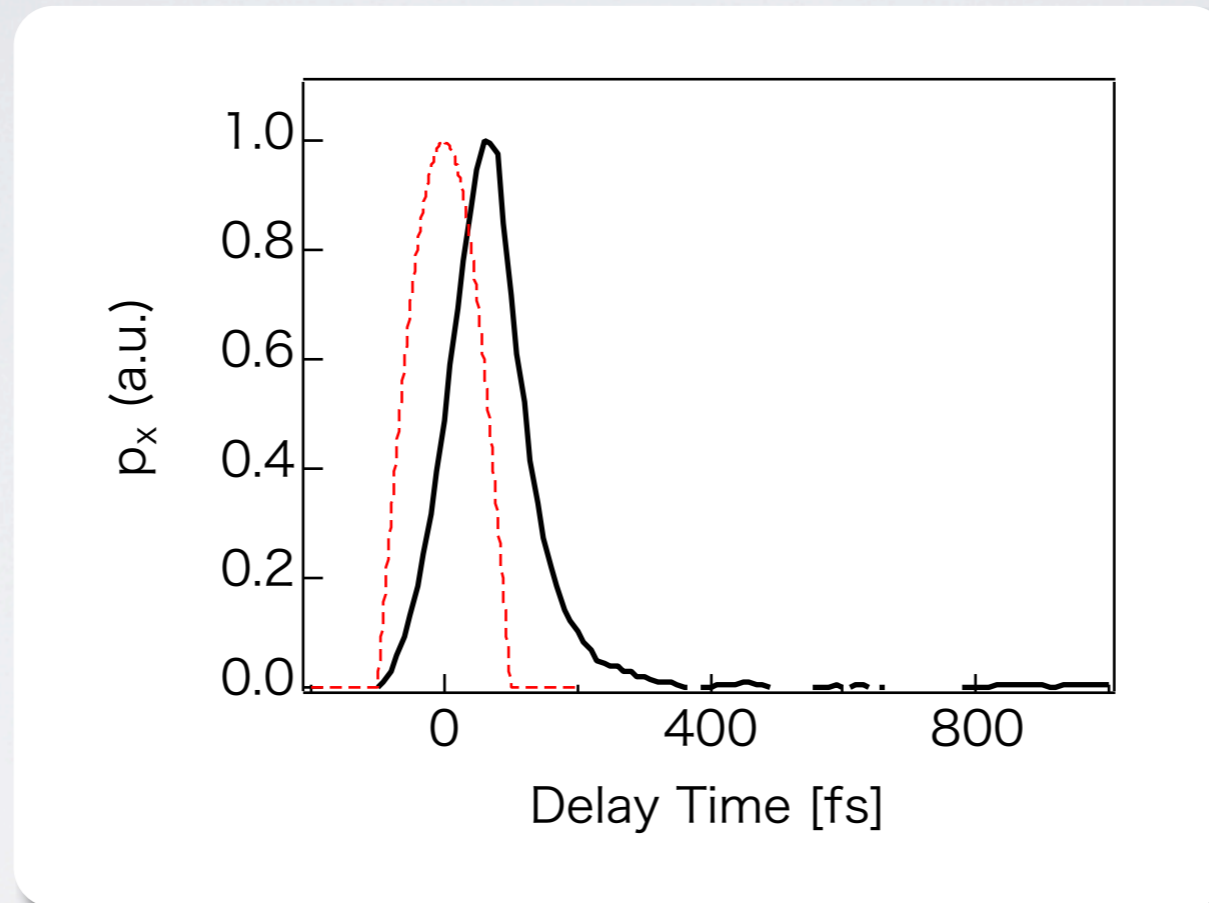


$$E_{\text{THz}} = 300 \text{ kV/cm}$$

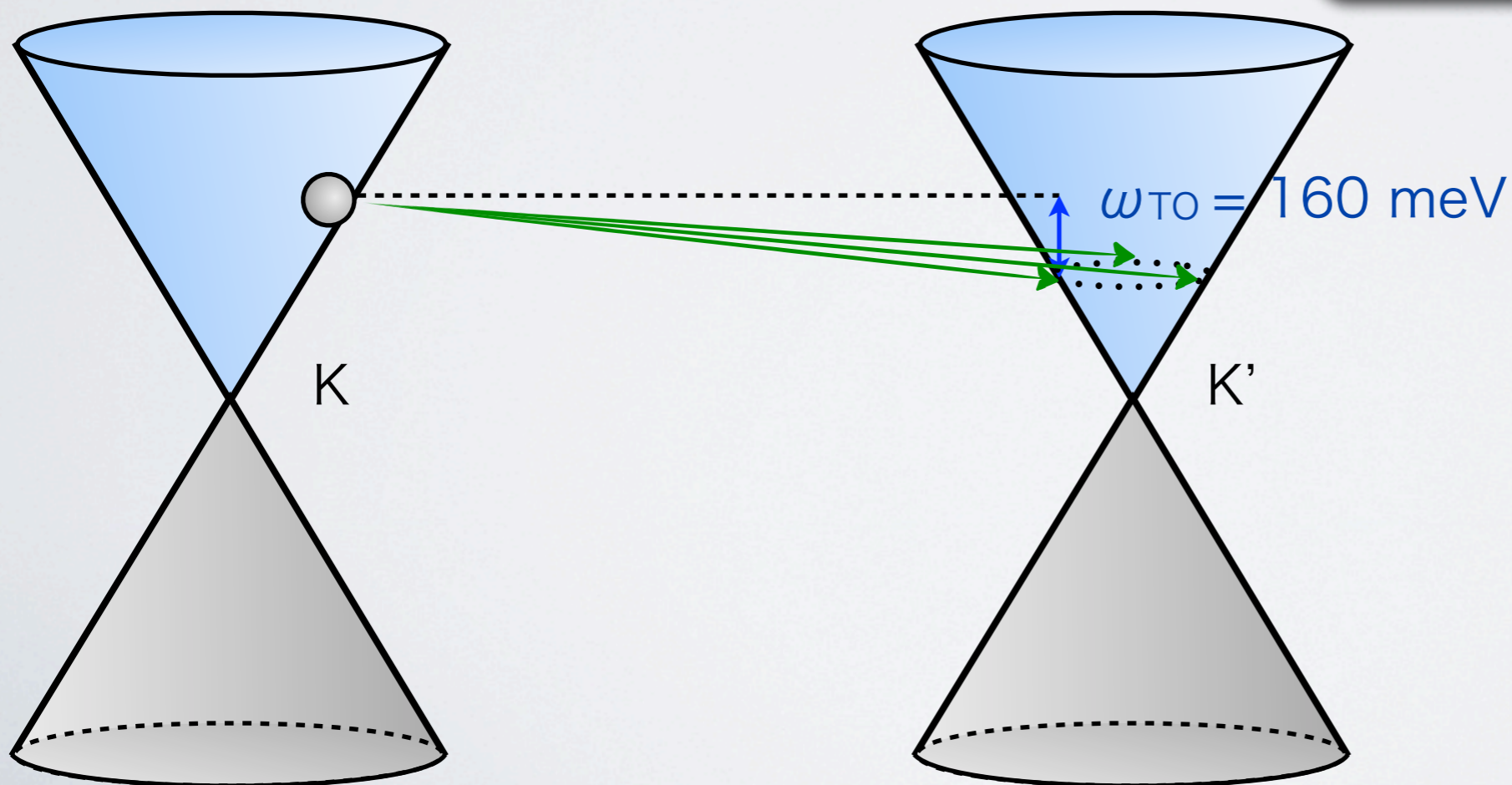
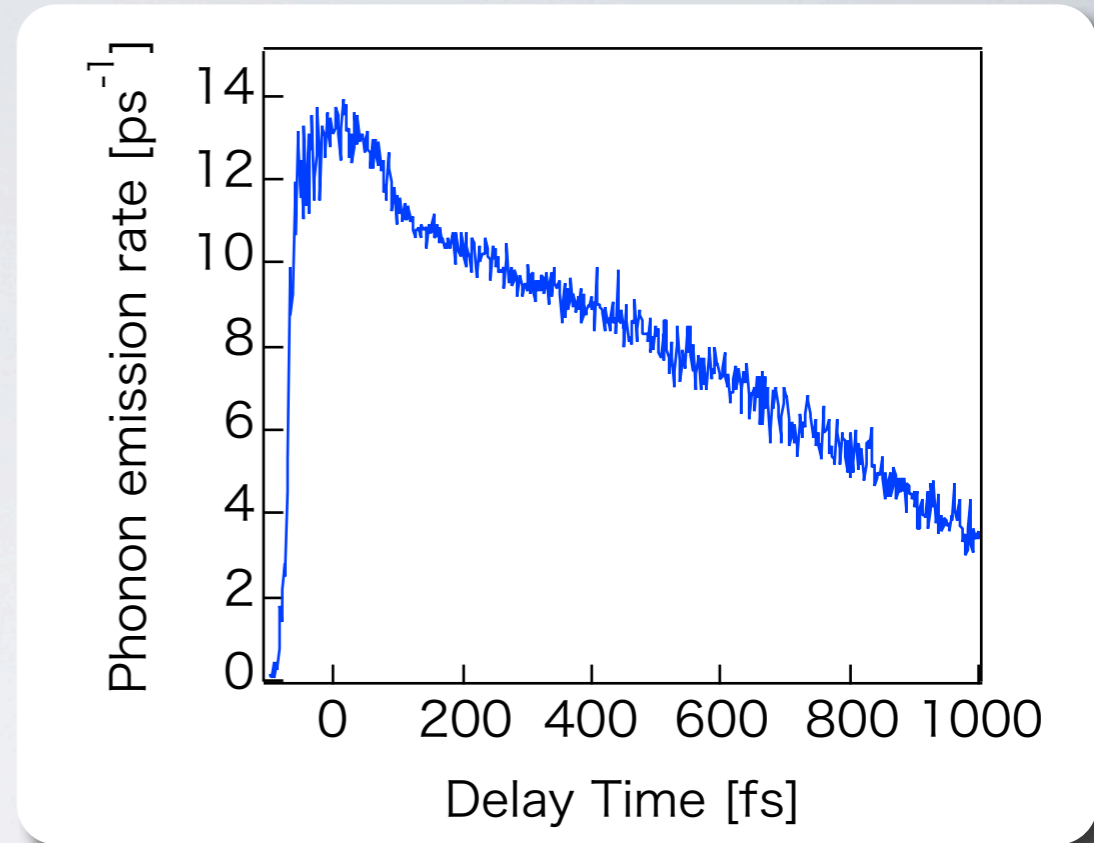
$$E_{\text{F}} = -200 \text{ meV}$$

Monte Carlo Simulation: Total Momentum

Hole total momentum in k_x direction



Monte Carlo Simulation : Phonon Scattering



Carrier Scattering

