

カイラル場の揺らぎを取り入れた 強結合極限での QCD 相図

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- Introduction
- 強結合極限でのカイラル場有効作用
- 補助場 Monte-Carlo による QCD 相図 (強結合極限)
- Summary

AO, T. Z. Nakano, in prep.

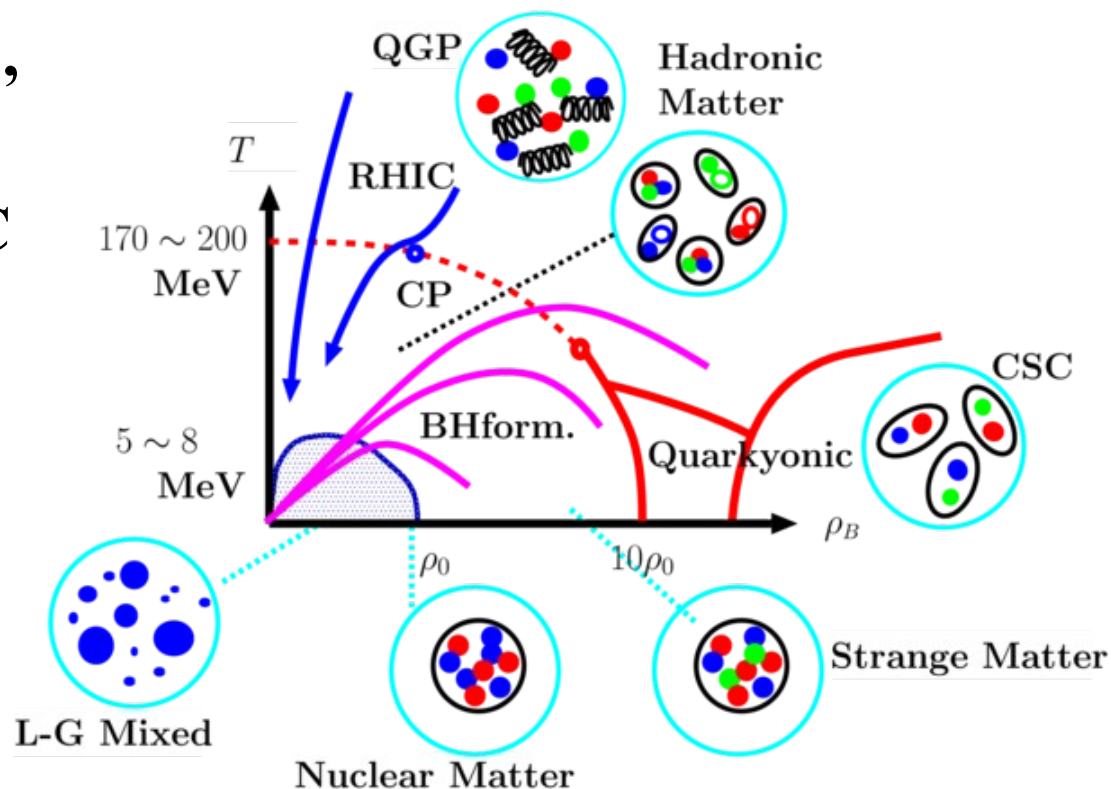
QCD Phase diagram

■ Phase transition at high T

- Early universe / RHIC, LHC / Lattice MC, pQCD, ...

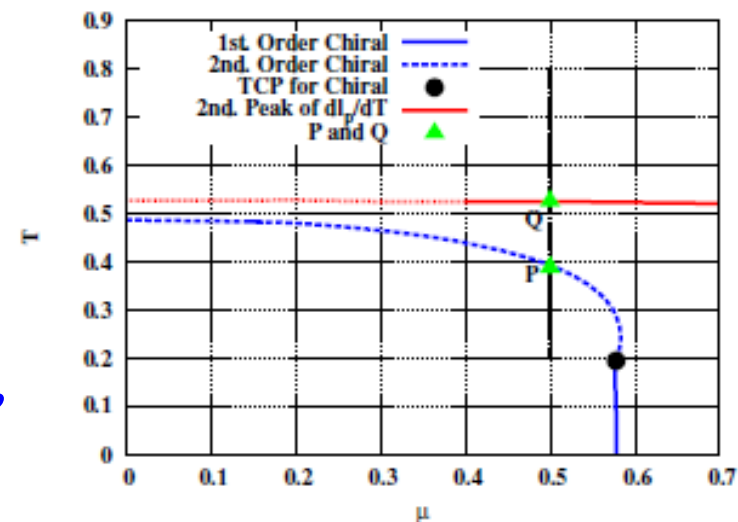
■ High μ transition

- Compact Astrophysical Objects (Neutron stars, Supernovae, Black hole formation, ...)
- RHIC-BES, FAIR, J-PARC, Astro-H, Grav. Wave, ...
- Sign problem in Lattice MC

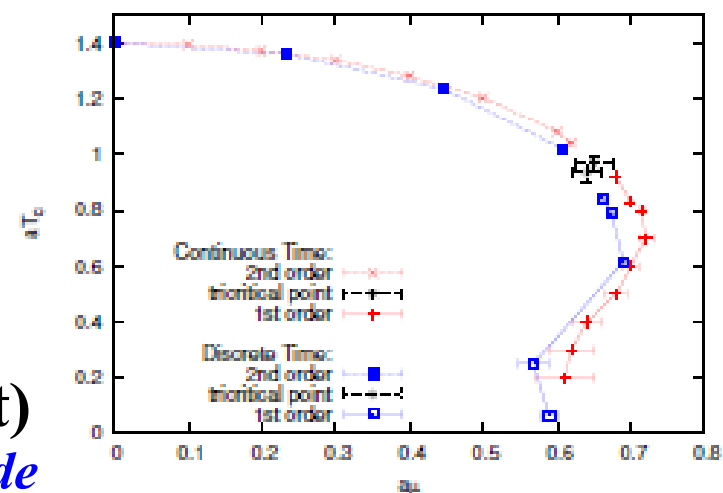


Strong Coupling Lattice QCD for finite μ

- **Effective Models (PNJL, PQM), FRG**
Fukushima ('11); Fujii, Sano ('10); T. K. Herbst, J. M. Pawłowski, B. J. Schaefer, ('11); ...
- **Lattice MC**
Fodor, Katz ('02); de Forcrand, Philipsen ('02); D'Elia, M. Lombardo ('03); Allton et al. ('04); Ejiri ('08); Nagata, Nakamura ('10); Nakagawa, Ejiri, ..(WHOT, '11), ..
- **Strong Coupling Lattice QCD**
 - **Mean Field approaches**
Damagaard, Hochberg, Kawamoto ('85); Bilic, Karsch, Redlich ('92); Fukushima ('04); Nishida ('04); Kawamoto, Miura, AO, Ohnuma ('07).
 - **MDP simulation (Strong Coupling Limit)**
Karsch, Mutter ('89); de Forcrand, Fromm ('10); de Forcrand, Unger ('11)



Miura, Nakano, AO, Kawamoto, arXiv:1106.1219



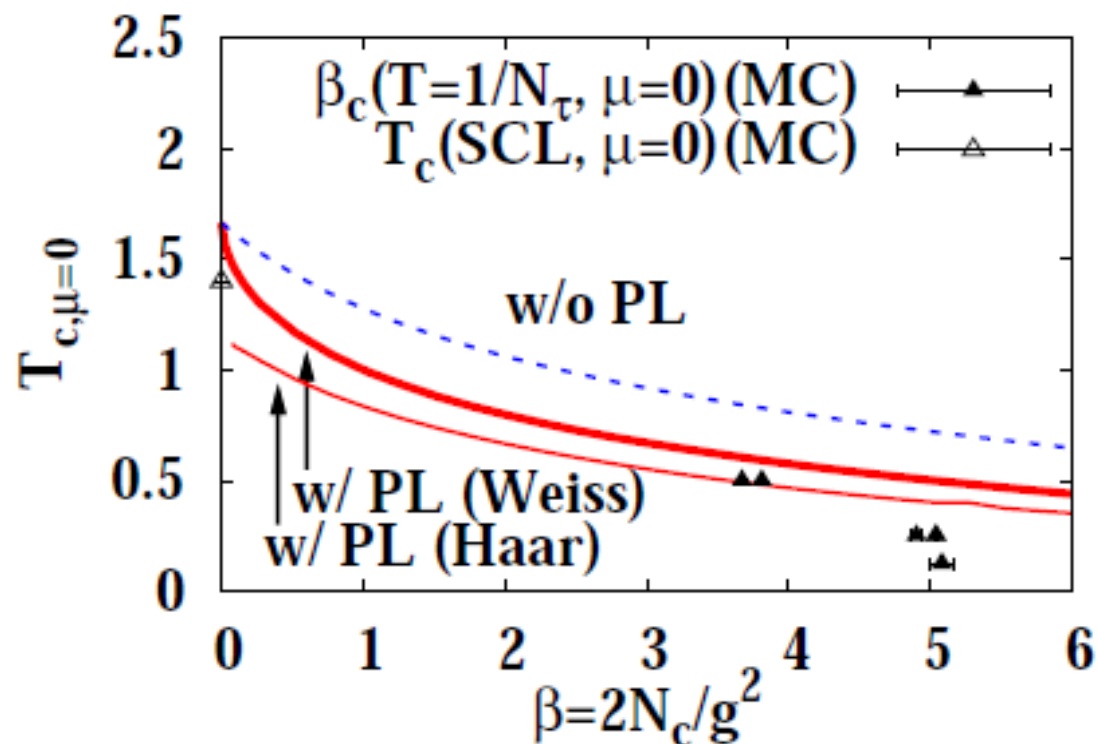
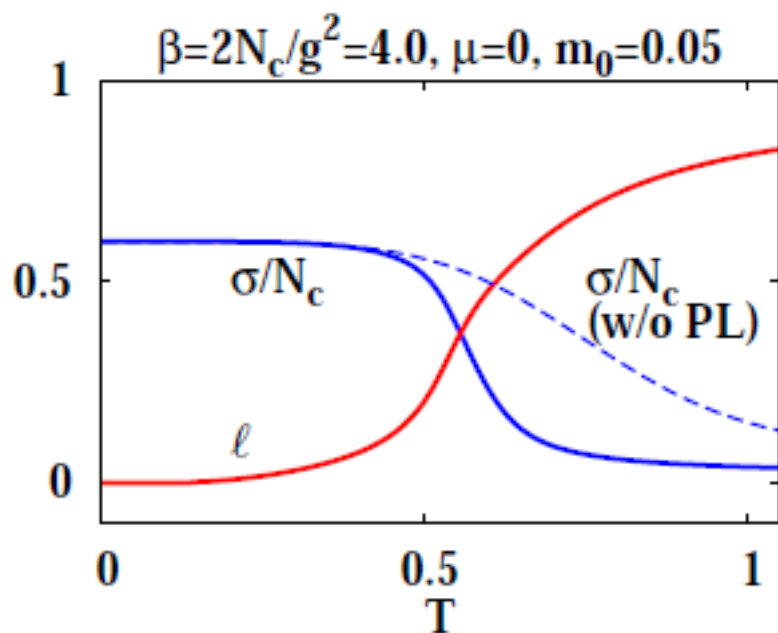
de Forcrand, Unger ('11)

P-SC-LQCD at $\mu=0$

T. Z. Nakano, K. Miura, AO, PRD 83 (2011), 016014 [arXiv:1009.1518 [hep-lat]]

- P-SC-LQCD reproduces $T_c(\mu=0)$ in the strong coupling region ($\beta = 2N_c/g^2 \leq 4$)

MC data: SCL (Karsch et al. (MDP), de Forcrand, Fromm (MDP)), $N_\tau=2$ (de Forcrand, private), $N_\tau=4$ (Gottlieb et al.('87), Fodor-Katz ('02)), $N_\tau=8$ (Gavai et al.('90))



Lattice Unit

*How can we include
both fluctuation and finite coupling effects ?
→ Auxiliary field MC*

Strong Coupling Limit ($1/g^2=0$) Lattice QCD

Lattice QCD action (staggered fermion)

$$S_{\text{LQCD}} = \frac{1}{2} \sum_{x,y} \bar{\chi}_x D_{x,y}^{(0)} \chi_y + \frac{m_0}{\gamma} \sum_x M_x + \frac{1}{2\gamma} \sum_{x,j} (V_x^{+j} - V_x^{-j}) - \frac{1}{g^2} \sum_P \text{tr}(U_P + U_P^+)$$

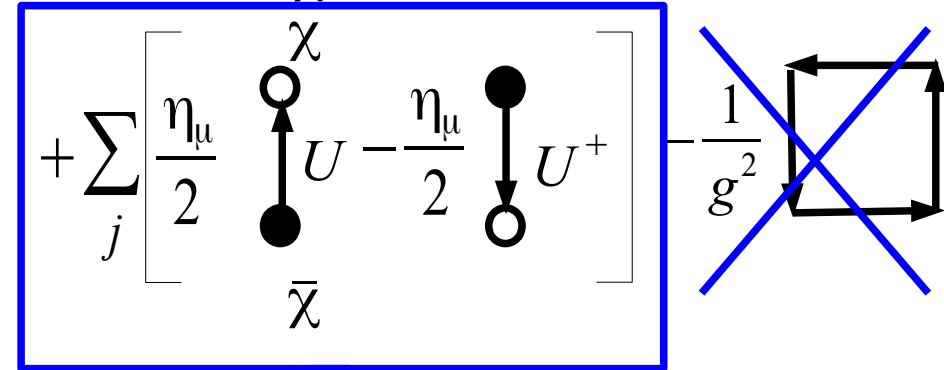
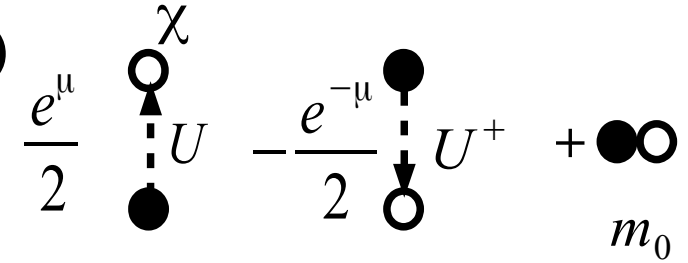
Spatial link integral in SCL

$$S_{\text{eff}} = \frac{1}{2} \sum_{x,y} \bar{\chi}_x D_{x,y}^{(0)} \chi_y + \frac{m_0}{\gamma} \sum_x M_x - \frac{1}{4N_c \gamma^2} \sum_{x,j} M_x M_{x+\hat{j}}$$

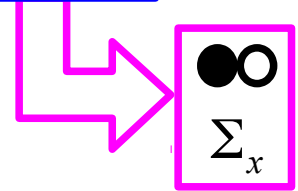
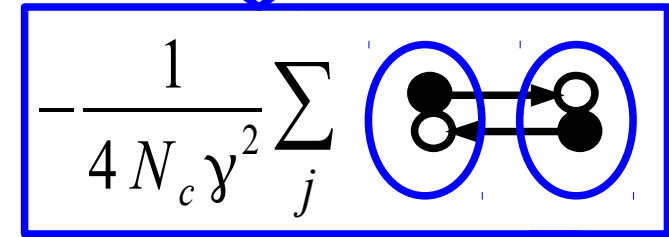
Bosonization

$$S_{\text{eff}} = \frac{1}{2} \sum_{x,y} \bar{\chi}_x D_{x,y}^{(0)} \chi_y + \sum_x \left(\Sigma_x + \frac{m_0}{\gamma} \right) M_x + \frac{\Omega}{4N_c \gamma^2} \sum_{k, f_M(\mathbf{k}) > 0} f_M(\mathbf{k}) [\sigma_k^* \sigma_k + \pi_k^* \pi_k]$$

$$D_{x,y}^{(0)} = \delta_{x+\hat{0},y} \delta_{x,y} e^{\mu/\gamma^2} U_{x,0} - \delta_{x,y+\hat{0}} \delta_{x,y} e^{-\mu/\gamma^2} U_{y,0}^+, V_x^{+\mu} = \eta_{\mu,x} \bar{\chi}_x U_{\mu,x} \chi_{x+\hat{\mu}}, V_x^{-\mu} = \eta_{\mu,x}^{-1} \bar{\chi}_{x+\hat{\mu}} U_{\mu,x}^+ \chi_x, M_x = \bar{\chi}_x \chi_x, \eta_{j,x} = (-1)^{x_0 + \dots + x_{j-1}}$$



$$\int DU U_{ab} U_{cd}^+ = \delta_{ad} \delta_{bc} / N$$



Bosonized Effective Action

■ Bosonization of MM term (Four Fermi (two-body) interaction)

$$-\alpha \sum_{j, x} M_x M_{x+\hat{j}} \rightarrow \sum_x M_x \Sigma_x + \alpha L^3 N_\tau \sum_{k, f_M(\mathbf{k}) > 0} f_M(\mathbf{k}) (\sigma_k^* \sigma_k + \pi_k^* \pi_k)$$

$$f_M(\mathbf{k}) = \sum_j \cos k_j, \quad \Sigma_x = 2\alpha \sum_{k, f_M(\mathbf{k}) > 0} f_M(\mathbf{k}) e^{ikx} (\sigma_k + i\boldsymbol{\varepsilon}(\mathbf{x}) \pi_k)$$

- Negative negative eigen values of meson matrix in “High” k
 \rightarrow Involves a factor $i\boldsymbol{\varepsilon}_x = i(-1)^{*(x_0+x_1+x_2+x_3)}$ in \mathbf{x} repr.

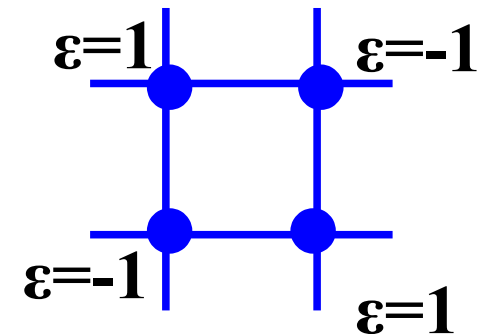
Miura, Nakano, AO (09), Miura, Nakano, AO, Kawamoto (09)

■ Bosonized Effective Action

$$S_{\text{eff}} = \frac{L^3 N_\tau}{4 N_c \gamma^2} \sum_{k, f_M(\mathbf{k}) > 0} f_M(\mathbf{k}) [\sigma_k^* \sigma_k + \pi_k^* \pi_k] - \sum_x \log R(x)$$

$$R(x) = X_N(\mathbf{x})^3 - 2 X_N(\mathbf{x}) + 2 \cosh(3 N_\tau \mu)$$

- $X_N(\mathbf{x}) =$ easily calculated from $\sigma(\mathbf{x})$ and $\pi(\mathbf{x})$.
- Imaginary part (π) involves $\boldsymbol{\varepsilon}_x$
 \rightarrow Phase cancellation for low k .

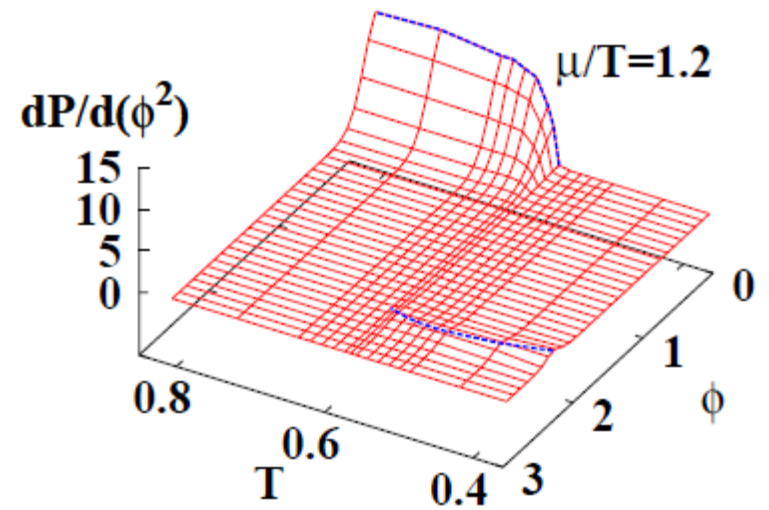
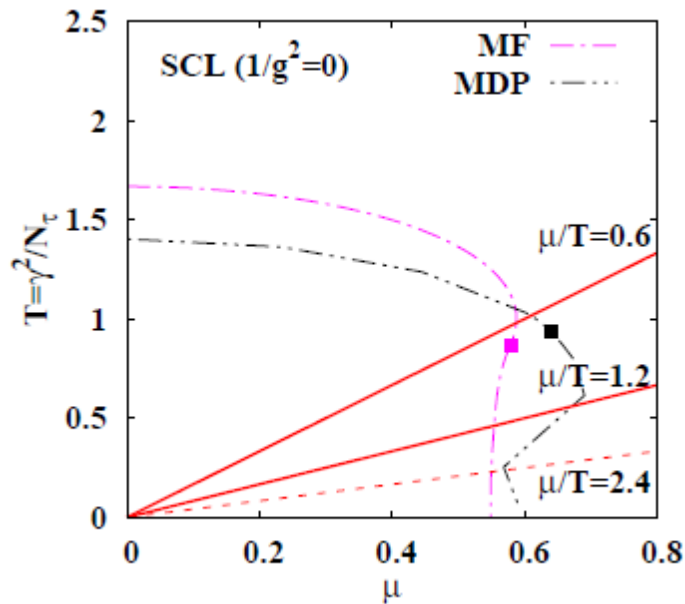
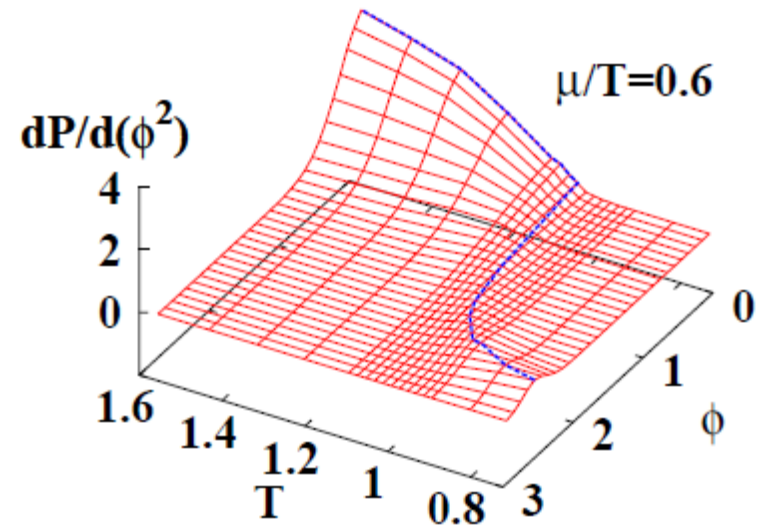


Let's Try !

- 4^4 asymmetric lattice + Metropolis sampling of σ_k and π_k .
- Metropolis sampling of full configuration (σ_k and π_k) at a time.
(efficient for small lattice)
- Chiral limit ($m=0$) simulation \rightarrow Symmetry in $\sigma \leftrightarrow -\sigma$
- Computer: My PC (Core i7).
(Partially in SR 16000 (single core))

Results (1): σ distribution

- Fixed μ/T simulation: $\mu/T = 0 \sim 2.4$
- Low μ region: Second order (Single peak: finite $\sigma \rightarrow$ zero)
- High μ region: First order (Dist. func. has two peaks)



Lattice Unit

Results (2): Susceptibility and Quark density

Weight factor $\langle \cos \theta \rangle$

$$\langle \cos \theta \rangle = Z / Z_{\text{abs}}$$

$$Z = \int D\sigma_k D\pi_k \exp(-S_{\text{eff}})$$

$$= \int D\sigma_k D\pi_k \exp(-\text{Re} S_{\text{eff}}) e^{i\theta}$$

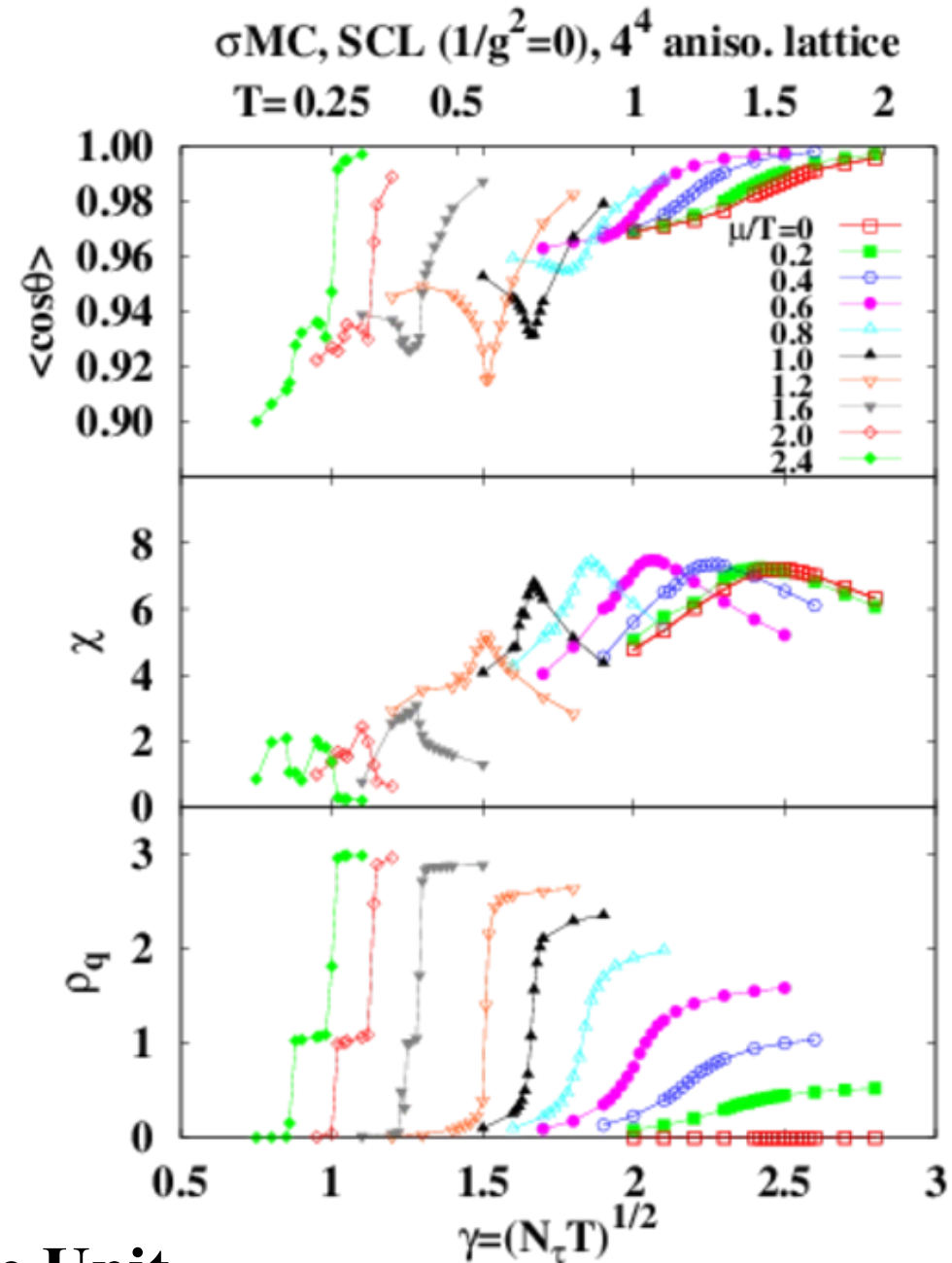
$$Z_{\text{abs}} = \int D\sigma_k D\pi_k \exp(-\text{Re} S_{\text{eff}})$$

Chiral susceptibility

$$\chi = -\frac{1}{L^3} \frac{\partial^2 \log Z}{\partial m_0^2}$$

Quark number density

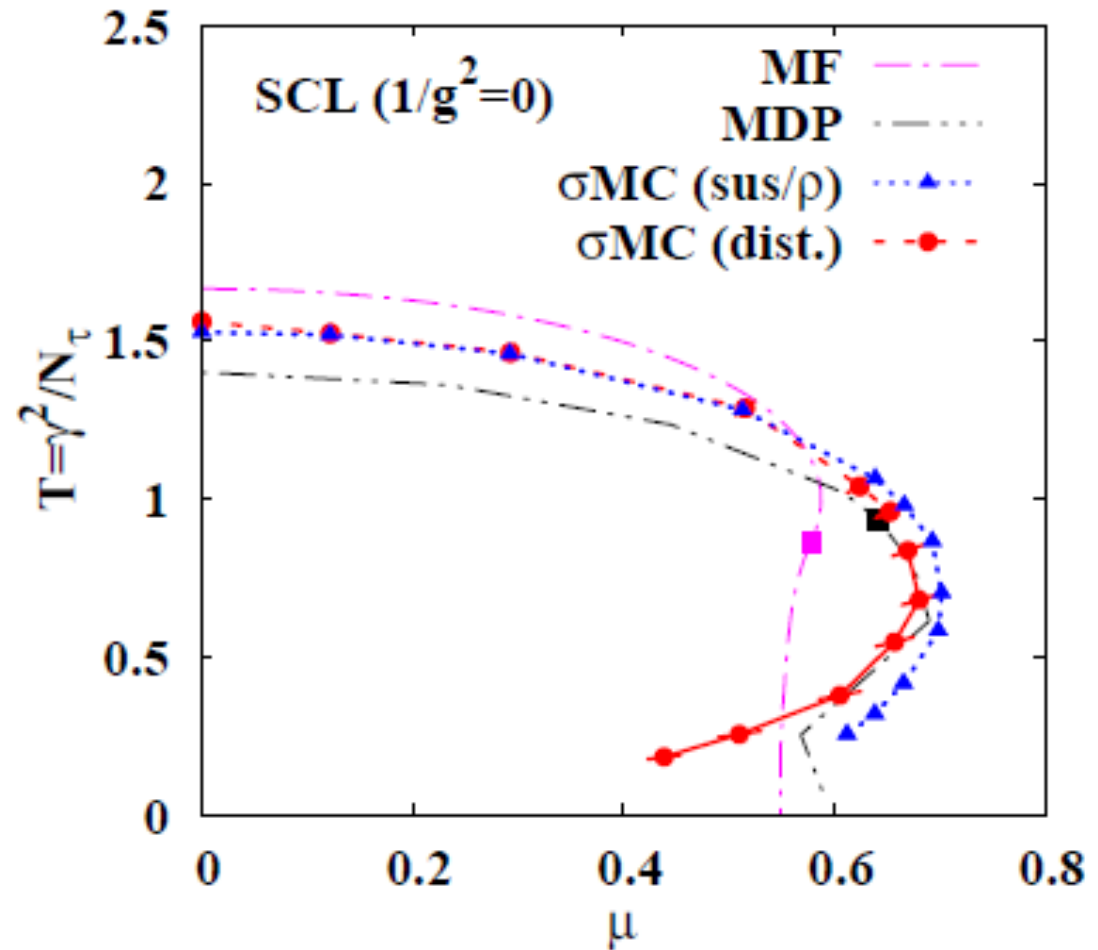
$$\rho_q = -\frac{1}{L^3} \frac{\gamma^2}{N_\tau} \frac{\partial \log Z}{\partial \mu}$$



Lattice Unit

Results (3): Phase diagram

- By taking $T = \gamma^2/N_\tau$,
 γ dep. of the phase boundary becomes small. *Bilic et al. ('92)*
- Phase boundary
 - σ dist.(red) & χ peak (blue)
- Fluctuation effects
 - T_c reduction at $\mu=0$
 μ_c enh. at medium T
 - Consistent with MDP
de Forcrand, Fromm ('09);
de Forcrand, Unger ('11)
- Bending at low T : Small size
de Forcrand, Wenger ('06)

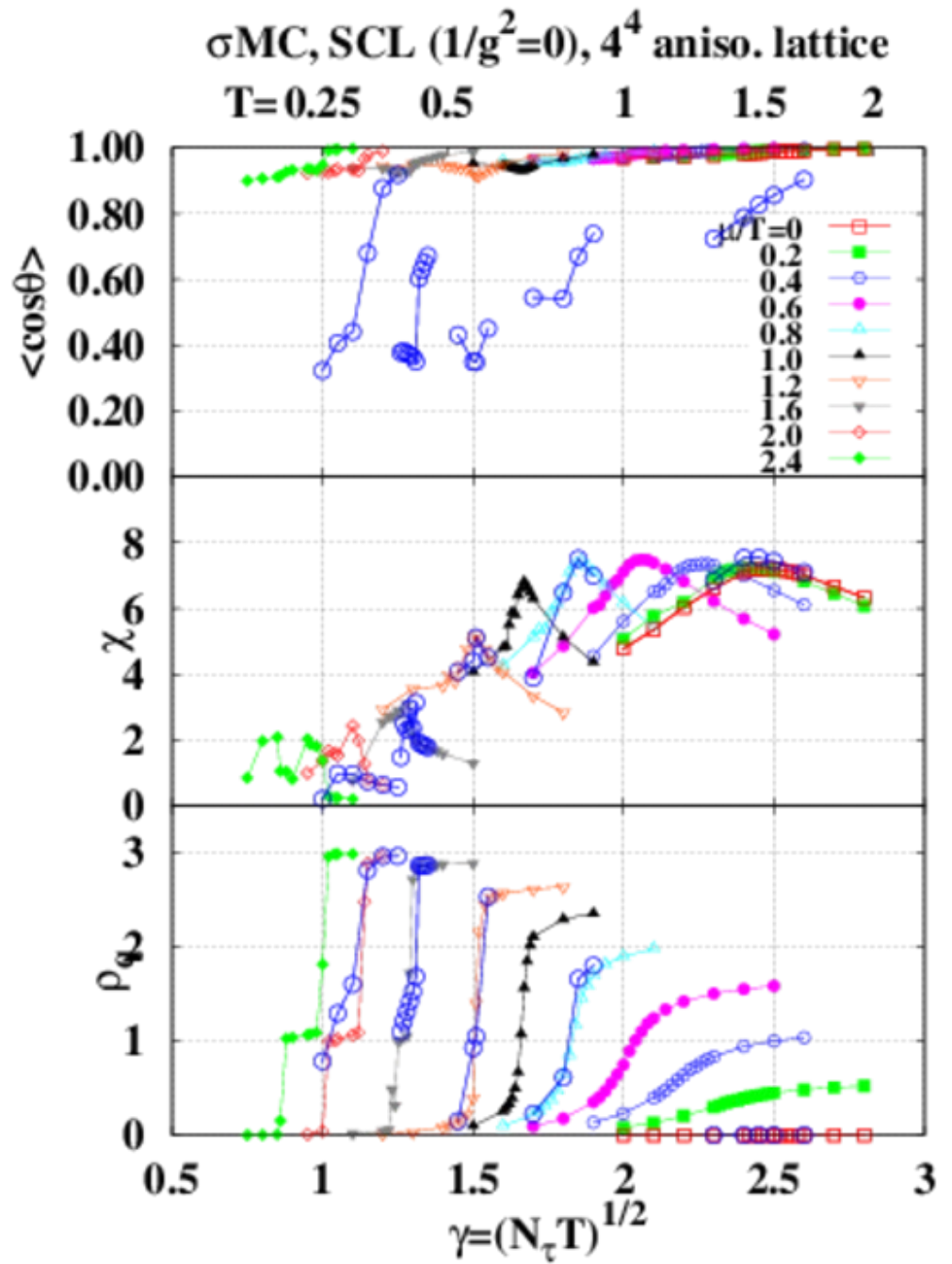


Summary

- We have proposed an auxiliary field MC method (σ MC), to simulate the SCL quark- U_0 action (LO in strong coupling ($1/g^0$) and $1/d$ ($1/d^0$) expansion) without further approximation.
c.f. Determinantal MC by Abe, Seki
- **Sign problem is mild** in small lattice ($\langle \cos \theta \rangle \sim (0.9-1)$ for 4^4), because of **the phase cancellation mechanism** coming from nearest neighbor interaction.
- Phase boundary is moderately modified from MF results by fluctuations, as in MDP simulation.
- Future work
 - Extension to NLO SC-LQCD is straightforward.
 - Larger lattice, finite coupling, different Fermion, higher $1/d$ terms including baryonic action and chiral Polyakov coupling.

Thank you

Preliminary Results with $8^3 \times 4$ Lattice



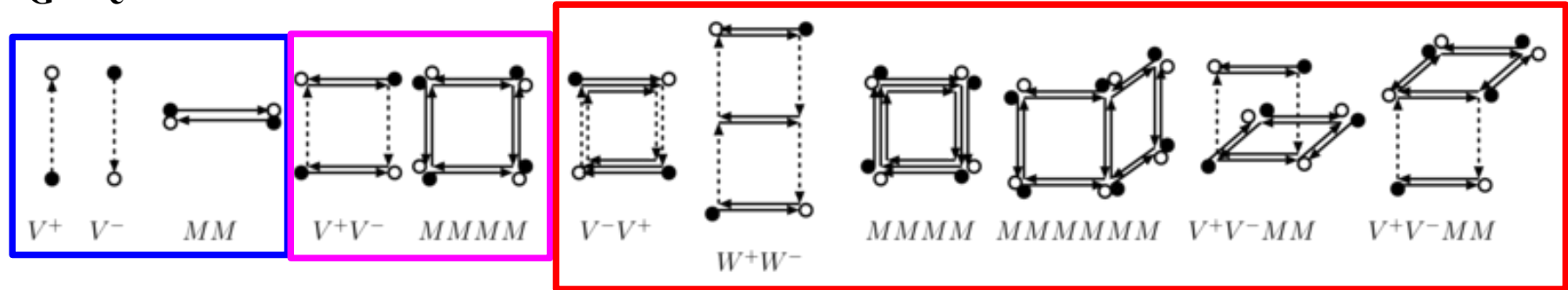
SC-LQCD with Fermions

Effective Action with finite coupling corrections

Integral of $\exp(-S_G)$ over spatial links with $\exp(-S_F)$ weight $\rightarrow S_{\text{eff}}$

$$S_{\text{eff}} = S_{\text{SCL}} - \log \langle \exp(-S_G) \rangle = S_{\text{SCL}} - \sum_{n=1} \frac{(-1)^n}{n!} \langle S_G^n \rangle_c$$

$\langle S_G^n \rangle_c = \text{Cumulant (connected diagram contr.)}$ *c.f. R.Kubo('62)*



$$S_{\text{eff}} = \frac{1}{2} \sum_x (V_x^+ - V_x^-) - \frac{b_\sigma}{2d} \sum_{x,j>0} [MM]_{j,x}$$

SCL (Kawamoto-Smit, '81)

$$+ \frac{1}{2} \frac{\beta_\tau}{2d} \sum_{x,j>0} [V^+V^- + V^-V^+]_{j,x} - \frac{1}{2} \frac{\beta_s}{d(d-1)} \sum_{x,j>0,k>0,k \neq j} [MMMM]_{jk,x}$$

NLO (Faldt-Petersson, '86)

$$- \frac{\beta_{\tau\tau}}{2d} \sum_{x,j>0} [W^+W^- + W^-W^+]_{j,x} - \frac{\beta_{ss}}{4d(d-1)(d-2)} \sum_{x,j>0, |k|>0, |l|>0, |k| \neq j, |l| \neq j, |l| \neq |k|} [MMMM]_{jk,x} [MM]_{j,x+\hat{l}}$$

$$+ \frac{\beta_{\tau s}}{8d(d-1)} \sum_{x,j>0, |k| \neq j} [V^+V^- + V^-V^+]_{j,x} \left([MM]_{j,x+\hat{k}} + [MM]_{j,x+\hat{k}+\hat{0}} \right)$$

NNLO (Nakano, Miura, AO, '09)

Approximations in Pol. loop extended SC-LQCD

■ Strong coupling expansion

- Fermion terms: LO($1/g^0$, SCL), NLO($1/g^2$), NNLO ($1/g^4$)
- Plaquette action: LO ($1/g^{2N\tau}$)

■ Large dimensional approximation

- 1/d expansion (d=spatial dim.)
→ Smaller quark # configs. are preferred.
 $\sum_j M_x M_{x+j} = O(1/d^0) \rightarrow M \propto d^{-1/2} \rightarrow \chi \propto d^{-1/4}$
- Only LO ($1/d^0$) terms are mainly evaluated.

■ (Unrooted) staggered Fermion

- Nf=4 in the continuum limit.

■ Mean field approximation

- Auxiliary fields are assumed to be constant.

*This work:
Auxiliary Field MC
in SCL*

Fermion Determinant

Faldt, Petersson, 1986

- Fermion action is separated to each spatial point and bi-linear
 → Determinant of $N_\tau \times N_c$ matrix

$$\exp(-V_{\text{eff}}/T) = \int dU_0 \begin{vmatrix} I_1 & e^\mu & 0 & & e^{-\mu} U^+ \\ -e^{-\mu} & I_2 & e^\mu & & \\ 0 & -e^{-\mu} & I_3 & e^\mu & \\ \vdots & & & \ddots & \\ -e^\mu U & & & -e^{-\mu} & I_N \end{vmatrix} \quad \begin{matrix} \uparrow \\ \downarrow \end{matrix} N_c \times N_\tau$$

$$= \int dU_0 \det \left[\underbrace{X_N[\sigma] \otimes \mathbf{1}_c}_{\text{magenta}} + \underbrace{e^{-\mu/T} U^+}_{\text{blue}} + \underbrace{(-1)^{N_\tau} e^{\mu/T} U}_{\text{red}} \right] \quad \begin{matrix} \uparrow \\ \downarrow \end{matrix} N_c$$

$$= X_N^3 - 2 X_N + 2 \cosh(3 N_\tau \mu)$$

$$I_\tau/2 = [\sigma(x) + i\varepsilon(x)\pi(x)]/2 N_c \gamma^2 + m_0/\gamma$$

$$X_N = B_N + B_{N-2} (2; N-1)$$

$$B_N = I_N B_{N-1} + B_{N-2}$$

$$B_N = \begin{vmatrix} I_1 & e^\mu & 0 & & \\ -e^{-\mu} & I_2 & e^\mu & & \\ 0 & -e^{-\mu} & I_3 & e^\mu & \\ \vdots & & & \ddots & \\ -e^{-\mu} & & & & I_N \end{vmatrix}$$

Clausius-Clapeyron Relation

- First order phase boundary \rightarrow two phases coexist

$$P_h = P_q \rightarrow dP_h = dP_q \rightarrow \frac{d\mu}{dT} = -\frac{s_q - s_h}{\rho_q - \rho_h}$$

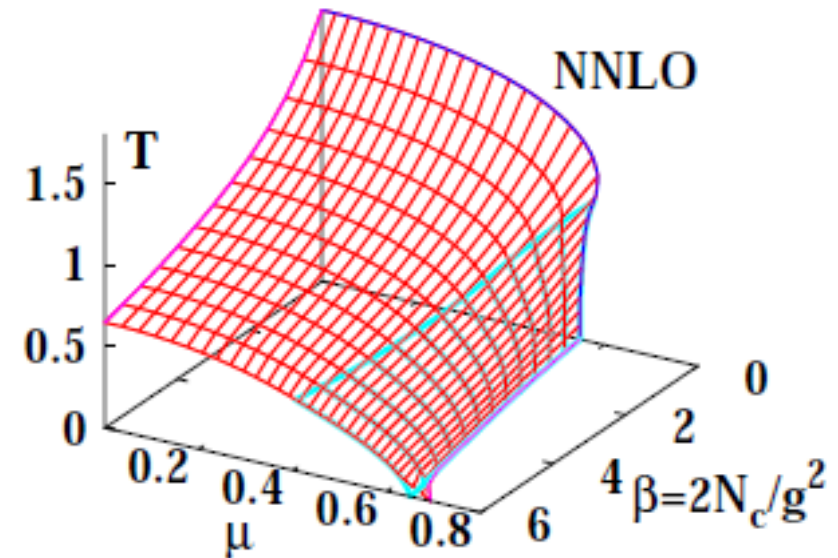
$$dP_h = \rho_h d\mu + s_h dT, \quad dP_q = \rho_q d\mu + s_q dT$$

- Continuum theory

\rightarrow Quark matter has larger entropy and density ($d\mu/dT < 0$)

- Strong coupling lattice

- ◆ SCL: Quark density is larger than half-filling, and “Quark hole” carries entropy $\rightarrow d\mu/dT > 0$
- ◆ NLO, NNLO $\rightarrow d\mu/dT < 0$



AO, Miura, Nakano, Kawamoto ('09)