

# Toward solving the sign problem with path optimization method

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We propose a new approach to circumvent the sign problem in which the integral path is optimized to control the sign problem. We give a trial function specifying the integration path in the complex plane and tune it to optimize the cost function which represents the seriousness of the sign problem. We call it path optimization method. In this method, we do not need to solve the gradient flow required in the Lefschetz-thimble method and then the construction of the integration-path contour arrives at the optimization problem where several efficient methods can be applied. In a simple model with a serious sign problem, the path optimization method is demonstrated to work well; the residual sign problem is resolved and precise results can be obtained even in the region where the global sign problem is serious.

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*Introduction* — The sign problem induced by the oscillating Boltzmann weight of the partition function in the numerical integration for various quantum systems is serious obstruction in the computational science; see Ref. [1] for a review. Particularly, the sign problem attracts much more attention recently in QCD because some new approaches to circumvent the sign problem have been proposed and applied.

Recent promising approaches to evade the sign problem include the complex Langevin method [2, 3] and the Lefschetz-thimble path-integral method [4–6]. The complex Langevin method is based on the stochastic quantization and then we are free from the usual Monte-Carlo sampling. Therefore, the sign problem does not appear, but it is well known that the complex Langevin method sometimes provides us wrong results when the drift term shows singular behavior in the Langevin-time evolution [7, 8]. In comparison, the Lefschetz-thimble path-integral method is based on the Picard-Lefschetz theory [9] and thus it is within the standard path-integral formulation. In this method, we construct the new integration-path contour which is so called the Lefschetz thimbles by solving the gradient flow starting from fixed points. Then, the partition function can be decomposed into the summation over contributions of relevant Lefschetz thimbles which can be determined from the crossing behavior of conjugate gradient flows with the original integration-path contour. On each Lefschetz thimble, the imaginary part of the action is constant and thus the sign problem seems to be resolved, but there are two remnants of the original sign problem. First one is the global sign problem: In the summation process of relevant Lefschetz thimbles, the cancellation can appear because imaginary parts of the action on each Lefschetz thimble have constant but different values. The other is the residual sign

problem; it comes from the Jacobian generated by the bending structure of the new integration-path. Recently, one more serious problem in the Lefschetz-thimble path-integral method has been discussed which is so called the singularity problem: There are singular points and cuts on the complexified variables of integration if the action has the square root and/or the logarithm, explicitly and implicitly. These singularities obstruct to draw continuous Lefschetz-thimbles in the numerical calculation of gradient flows.

In this article, we propose a new method which we call the *path optimization method (POM)* to attack the sign and singularity problem. This method is based on the Lefschetz-thimble path-integral method. The main idea is the modification of the path-integral contour by minimizing the suitable *cost function* which reflects the seriousness of the sign problem. This means that the evading the sign problem arrives at the optimization problem. This fact becomes the strong advantage of this method because the optimization problem is well explored in the computational science and thus we may use several efficient methods such as the machine learning to the optimization process. The path optimization method is demonstrated in the simple model with the serious sign problem where the complex Langevin method can fail.

*Cost and trial functions* — In the path optimization method, the function which is so called the *cost function* plays a crucial role to construct the new integration-path contour which has the controllable sign problem. The cost function is related with the seriousness of the sign problem with weakened weight cancellation by minimizing the function. In this article, we use the following cost function;

$$\mathcal{F}[z(t)] = \frac{1}{2} \int dt |e^{i\theta(t)} - e^{i\theta_0}|^2 \times |J(t)e^{-S(z(t))}|, \quad (1)$$

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with

$$e^{i\theta(t)} = \frac{J(t)e^{-S(z(t))}}{|J(t)e^{-S(z(t))}|}, \quad e^{i\theta_0} = \frac{\mathcal{Z}}{|\mathcal{Z}|}, \quad (2)$$

where  $z$  is the complexified variables of integration,  $\mathcal{Z}$  is the partition function and  $J(t) = dz/dt$ . This function can be expressed by using the reweighting factor as

$$\frac{\mathcal{F}}{|\mathcal{Z}|} = |\langle e^{i\theta} \rangle_{pq}|^{-1} - 1, \quad (3)$$

where

$$\langle \mathcal{O} \rangle_{pq} \equiv \frac{\int dt \mathcal{O} |J e^{-S}|}{\int dt |J e^{-S}|}. \quad (4)$$

It should be noted that the choice of the cost function is not unique and thus we can freely change or extend it as long as the function reflects the seriousness of the sign problem.

To perform the optimization of Eq. (1), we need the trial function to specify the integration-path contour. One simple way to prepare the trial function is using the complete set,  $\mathcal{H}_m$ , as

$$z(t) = x(t) + iy(t), \quad \begin{cases} x(t) = \sum_m c_{x,m} \mathcal{H}_m(t), \\ y(t) = \sum_m c_{y,m} \mathcal{H}_m(t), \end{cases} \quad (5)$$

with imposing the conditions,  $x(\pm\infty) = \pm\infty$ ,  $|y(\pm\infty)| < \infty$ . We can extend this trial function to more complicated form by performing the feature engineering or the machine learning.

*Example* — In this article, we use the following partition function as an example to demonstrate the path optimization method. The actual form of partition function [8] is

$$\mathcal{Z}_p = \int dx (x + i\alpha)^p e^{-\frac{x^2}{2}}, \quad (6)$$

where  $\alpha$  and  $p$  are input parameters. The analytic result of  $\mathcal{Z}_p$  can be obtained from the recurrence formula;

$$\mathcal{Z}_p = i\alpha \mathcal{Z}_{p-1} + (p-1) \mathcal{Z}_{p-2}, \quad (7)$$

and the expectation value of  $x^2$  is expressed as

$$\langle x^2 \rangle_p = \frac{\mathcal{Z}_{p+2} - 2i\alpha \mathcal{Z}_{p+1} - \alpha^2 \mathcal{Z}_p}{\mathcal{Z}_p}. \quad (8)$$

In the actual optimization, we use simplified version of Eq. (5) based on the Gaussian function;

$$x(t) = t, \quad (9)$$

$$y(t) = c_1 \exp\left(-\frac{c_2^2 t^2}{2}\right) + c_3. \quad (10)$$

This function should be nonlinear to construct the complicated structure of the integration-path contour. The

optimization is numerically performed using the steepest descent method,  $dc_i/d\tau = -\partial\mathcal{F}/\partial c_i$ , and the integration is performed using the double exponential formula.

The optimized integration-path in comparison with the Lefschetz thimble is shown in Fig. 1. It can be seen that the two contours overlap in the vicinity of the fixed point. However, there are qualitative differences on the thimble

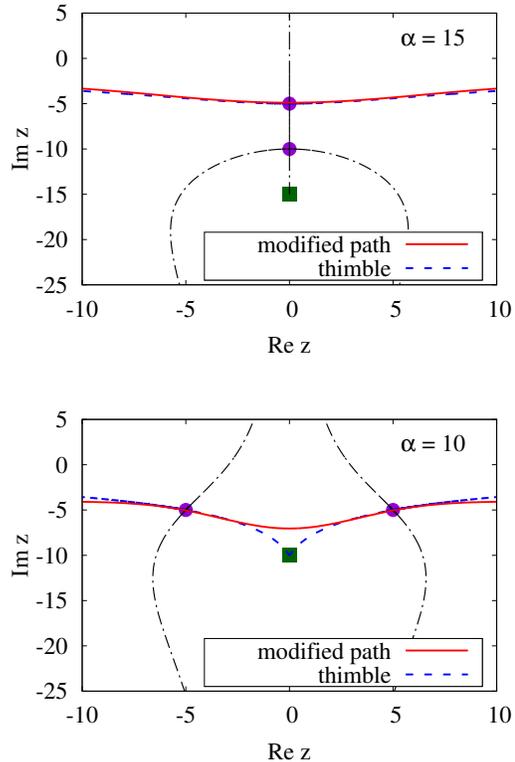


FIG. 1: Modified integration-path by optimizing Eq. (1) and the Lefschetz thimbles for  $p = 50, \alpha = 15, 10$ . Closed circle (square) point shows the fixed (singular) point. Dot-dashed lines are steepest ascent paths.

structure with varying  $\alpha$ . In the case with  $\alpha = 10$ , Lefschetz thimbles are ended at the singular point unlike the case with  $\alpha = 15$  and then the optimized integration-path approaches to the singular point.

Figure 2 shows  $J e^{-S}$  on the optimized integration-path. We can see that there is the large probability distribution ( $W(t) \sim |J e^{-S}|$ ) with almost the constant phase near the fixed point. Therefore, Monte-Carlo sampling works with  $\alpha = 15$ . In comparison,  $W(t)$  has two peaks in the case with  $\alpha = 10$ . The sign of  $\text{Im } J e^{-S}$  are opposite at both peaks and thus there are serious cancellations between them when we take into account both peaks to the integration. If we take into account only one peak, the wrong result comes up. The cancellations are induced by the singular point when the optimized integration-path approaches to it: The Boltzmann weight becomes zero at singular point and thus the sign of the Boltzmann weight can be easily flipped near the singular point. In

the present case, there is exact parity symmetry between  $\text{Re } z$  and  $-\text{Re } z$  and thus the cancellation is very serious. This cancellation reflects the hidden Lefschetz thimble structure behind the path optimization method.

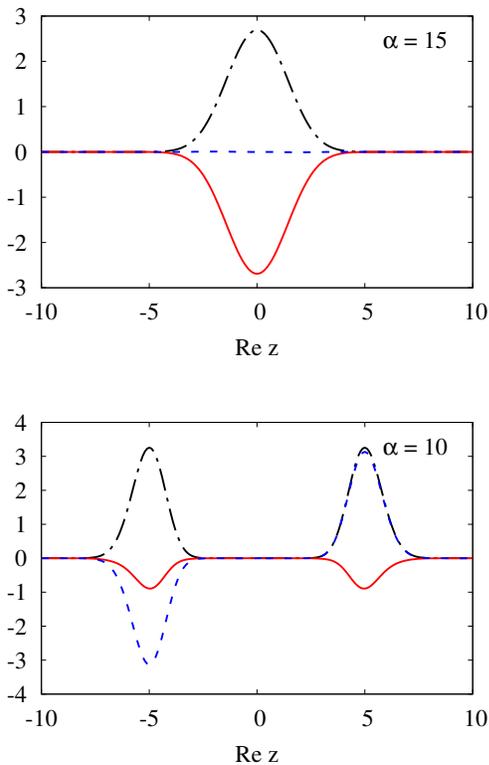


FIG. 2: Boltzmann weight on the modified integration-path with  $p = 50$ ,  $\alpha = 15, 10$ . Solid (dashed) line represents real (imaginary) part of  $J e^{-S}$ . The dot-dashed line indicates  $|J e^{-S}|$ . Where, the amplitudes is normalized by factor  $10^{55}$  with  $\alpha = 15$ , and  $10^{42}$  with  $\alpha = 10$ .

The optimized average phase factor is shown in Fig. 3. From the difference between the full calculation and the calculation in the  $\text{Re } z \in [0, \infty)$  range, we can clarify the seriousness of the global sign problem. In the case with  $\alpha \gtrsim 14$ , we can see that the sign problem can be solved because the path is represented by one thimble. In the case with  $\alpha \lesssim 14$ , contributions from the two thimbles cancel with each other. In the path optimization method, we can resolve the residual sign problem, but not the global sign problem. This problem also exists in the ordinary and generalized Lefschetz thimble methods.

Figure 4 shows the expectation value of  $x^2$  by hybrid Monte-Carlo method on modified integration-path. We calculate the expectation value in Eq. (4) by using the reflection symmetry of the Boltzmann weight,  $W(t) = W(-t)$ , in this setting. This treatment replicates the parallel tempering algorithm which has been applied to the generalized Lefschetz-thimble path-integral method [10, 11]. The results well agree with the analytic results (8). Readers can find the calculation of this model

by using the complex Langevin method in Ref. [8].

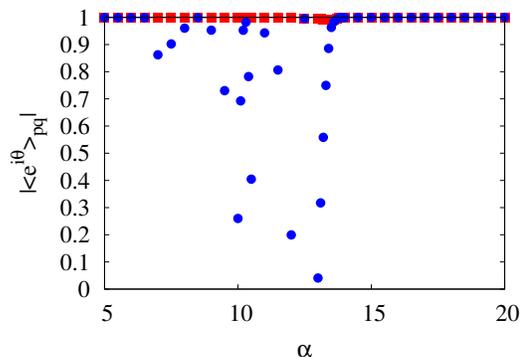


FIG. 3: The phase quenched expectation value of reweighting factor with  $p = 50$ . Closed circle shows the expectation in  $\text{Re } z \in (-\infty, \infty)$ , and square point in  $\text{Re } z \in [0, \infty)$ . The expectation value is calculated by numerical integration.

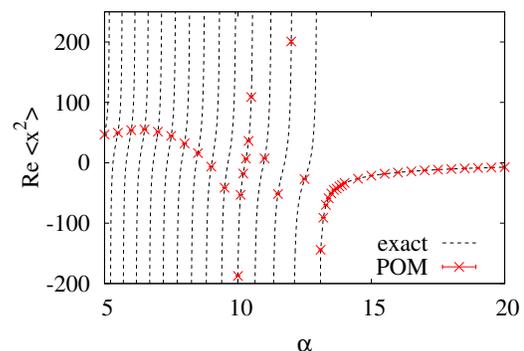


FIG. 4: The expectation value of  $x^2$  with  $p = 50$ . Errors are estimated by Jackknife method.

In the ordinary and generalized Lefschetz-thimble path-integral method, its expansibility is not large because the integration-path is uniquely determined by solving the gradient flows. In comparison, the path optimization method has large expansibility because we can freely replace or extend the cost function as long as it reflects the seriousness of the sign problem. This point is important advantage of the path optimization method. We leave the actual extension of the cost function as a future work since this study is the first attempt to demonstrate our method.

*Summary* — In this article, we have proposed a new approach to circumvent the sign problem which is based on the Lefschetz-thimble path-integral method. We call the *path optimization method*. In the method, the new integration-path is constructed in the plane of complexified variables of integration by minimizing the *cost function*. The cost function is set to reflect the seriousness of the sign problem. The actual optimization of the

integration-path is carried out by using the trial function.

We have demonstrated the path optimization method by using the simple model with the serious sign problem proposed in Ref. [8]. In the path optimization method, we can resolve the residual sign problem which appears in the ordinary and generalized Lefschetz-thimble path-integral methods. But, at least on our present choice of the cost function, the global sign problem cannot be resolved at least in the present setting. However, we can well reproduce the exact results by using the path optimization method in the wide range of the model parameter space.

Finally, we summarize advantages of the path optimization method:

1. No residual sign problem.
2. Applicability of various efficient methods to the op-

timization process.

3. Controllability of the singularity problem.
4. Large extensibility of the cost function.

Possible disadvantage may be the numerical cost. In the complex system, the sign-problem weakened integration-path is expected to have a very complicated shape. Therefore, we should check which optimization method is suitable or not step by step in the future.

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- [1] P. de Forcrand, PoS **LAT2009**, 010 (2009), 1005.0539.  
 [2] G. Parisi and Y.-s. Wu, *Sci.Sin.* **24**, 483 (1981).  
 [3] G. Parisi, *Phys.Lett.* **B131**, 393 (1983).  
 [4] E. Witten, *AMS/IP Stud. Adv. Math.* **50**, 347 (2011), 1001.2933.  
 [5] M. Cristoforetti, F. Di Renzo, and L. Scorzato (Aurora-Science Collaboration), *Phys.Rev.* **D86**, 074506 (2012), 1205.3996.  
 [6] H. Fujii, D. Honda, M. Kato, Y. Kikukawa, S. Komatsu, et al., *JHEP* **1310**, 147 (2013), 1309.4371.  
 [7] G. Aarts, E. Seiler, and I.-O. Stamatescu, *Phys. Rev.* **D81**, 054508 (2010), 0912.3360.  
 [8] J. Nishimura and S. Shimasaki, *Phys. Rev.* **D92**, 011501 (2015), 1504.08359.  
 [9] F. Pham, in *Proc. Symp. Pure Math* (1983), 40, pp. 319–333.  
 [10] M. Fukuma and N. Umeda (2017), 1703.00861.  
 [11] A. Alexandru, G. Basar, P. F. Bedaque, and N. C. Warrington (2017), 1703.02414.