# Holographic Model of the Dark Universe & the Late Time Evolution



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$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} - \frac{1}{L}\left(\mathcal{K}_{\mu\nu} - \mathcal{K}g_{\mu\nu}\right) = \kappa_4 T_{\mu\nu}$$



Ref: JHEP 1810 (2018) 009 [arXiv: 1712.09326] by: Rong-Gen Cai (ITP/ Beijing) Sichun Sun (NTU/ Taipei) Yun-Long Zhang (APCTP/ Pohang)

 & Fitting with SNIa [arXiv: <u>1812.11105</u>]
 with Bum-Hoon Lee (Sogang U./ Seoul) Sunly Khimphun (Phnom Penh)
 <u>Gansukh Tumurtushaa</u> (IBS/Daejeon)



## Motivations — Holographic Hydrodynamics



Figures Credit: RHIC & Google

### Toy Duality: Field & Surface Matter



Figures Credit: searched from Google

Yun-Long Zhang Holographic Dark Fluid

### Membrane Duality: Black Hole & Surface Fluid

T. Doumer & K. Thorne, (1980s-)



**Effective Description** 

$$\mathcal{T}_{ab} = -2(K_{ab} - K\gamma_{ab})$$

Membrane on Stretched horizon

Viscosity & Conductivity



### Echoes from Compact Objects [1706.06155 PRD'17]



Echoes from the Abyss [1612.00266 PRD'17]

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# AdS/CMT Duality: Geometry & Quantum Matters



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Entangled state

(1)))

 $|\psi\rangle \neq |\psi\rangle_A \otimes |\psi\rangle_B$ 

Mixed

B

Trace

## Fluid/Gravity Duality (2008-)

### AdS Black Brane

$$ds^{2} = -2 u_{\mu} dx^{\mu} dr + r^{2} (\eta_{\mu\nu} + [1 - f(r/\pi T)] u_{\mu} u_{\nu}) dx^{\mu} dx^{\nu} + 2r \left[ \frac{r}{\pi T} F(r/\pi T) \sigma_{\mu\nu} + \frac{1}{3} u_{\mu} u_{\nu} \partial_{\lambda} u^{\lambda} - \frac{1}{2} u^{\lambda} \partial_{\lambda} (u_{\nu} u_{\mu}) \right] dx^{\mu} dx^{\nu},$$

$$g_{ab} = \sum_{k=0}^{\infty} \epsilon^{k} g_{ab}^{(k)}, \quad T = \sum_{k=0}^{\infty} \epsilon^{k} T^{(k)}, \quad u_{\mu} = \sum_{k=0}^{\infty} \epsilon^{k} u_{\mu}^{(k)}.$$

$$T^{\mu\nu} = \lim_{r_{\nu} \to \infty} \frac{r_{c}^{d-2}}{16\pi G_{N}} \left[ K^{\mu\nu} - K \gamma^{\mu\nu} - (d-1) \gamma^{\mu\nu} - \frac{1}{d-2} \left( R^{\mu\nu} - \frac{1}{2} R \gamma^{\mu\nu} \right) \right]$$

$$P = \frac{1}{16\pi G_{N}} \left( \frac{4\pi T}{d} \right)^{d}$$

$$\eta = \frac{s}{4\pi} = \frac{1}{16\pi G_{N}} \left( \frac{4\pi T}{d} \right)^{d-1}$$

$$\eta = \frac{s}{4\pi T} \int_{1}^{\infty} dy \frac{y^{d-2} - 1}{y(y^{d-1})}$$

$$\kappa = \frac{d}{4\pi T} \int_{1}^{\infty} dy \frac{y^{d-2} - 1}{y(y^{d-1})}$$

[Bhattacharyya, Hubeny, Minwalla, Rangamani et al]

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 $\xi_{\omega} = 0 \ .$ 

# Moving the Holographic Screen to the Finite Cutoff



Wilsonian Approach to Fluid/Gravity Duality [Bredberg, Keeler, Lysov, Strominger, JHEP 1103(2011)141]

Non-Relativistic Fluid Dual to Asymptotically AdS Gravity at Finite Cutoff Surface [Cai, Li, Zhang, JHEP 1107(2011)027]

## Hydrodynamics of the Cutoff AdS Fluid

$$\begin{array}{ll} \mbox{AdS Metric} & ds_{p+2}^2 = -r^2 f(r) d\tau^2 + \frac{1}{r^2 f(r)} dr^2 + r^2 dx_i dx^i \\ \\ \mbox{Induced Metric} & ds_{p+1}^2 = -r_c^2 f(r_c) d\tau^2 + r_c^2 dx_i dx^i \\ \\ \mbox{Dual Tensor} & \overline{\mathcal{T}_{ab} = -2(K_{ab} - K\gamma_{ab} + C\gamma_{ab})} \\ \\ \mbox{Constraint equations} & 2G_{\mu b} n^{\mu}|_{r_c} = 2\partial^a (K_{ab} - \gamma_{ab} K) = 0 \Rightarrow \partial^a T_{ab} = 0 \end{array}$$

### Our contributions in the Holographic Cutoff AdS Fluid

| Holographic Cutoff AdS Fluid in Non-relativistic limit<br>[Cai, Li, Zhang, JHEP 1107(2011)027]                 | $\overset{\mathbf{t}}{\partial_r} \sim \epsilon^0,  \partial_i \sim v_i \sim \partial_i \phi \sim \epsilon^1,  \partial_\tau \sim P \sim \epsilon^2$  |
|--|---|
| <u>Holographic Forced Fluid Dynamics on finite Cutoff</u><br>[Cai, Li, Nie, <b>Zhang</b> , NPB 864 (2012) 260] | $\partial_i v^i = 0,  \partial_\tau v_i + v^j \partial_j v_i + \partial_i P - \nu \partial^2 v_i = f_i^{\phi} + f_i^q$  |
| Incompressible Navier-Stokes from Chern-Simons Mc Cai, Li, Qi, Zhang, PRD 86 (2012) 086008] $\partial_	au v$   | <u>odified Gravity</u><br>$v_i + v^j \partial_j v_i + \partial_i P - \nu \partial^2 v_i - (\tilde{\nu} \epsilon_{ij} \partial^2 v_j + \tilde{\zeta} \epsilon^{jk} \partial_i \partial_j v_k) = f_i$ |
| Holographic Charged Fluid with Anomalous Current a<br>[Bai, Hu, Lee, Zhang, JHEP 1211 (2012) 054]              | $\frac{\text{at Finite Cutoff}}{\xi_B} = c\left(\mu - \frac{1}{2}\frac{n\mu^2}{\rho + p}\right),  \xi_V = c\left(\mu^2 - \frac{2}{3}\frac{n\mu^3}{\rho + p}\right)$                                 |

# **Cutoff AdS Fluid with Momentum Relaxation**



Ref: by Khimphun, Lee, Park, Zhang JHEP 1801 (2018) 058

# Running From Conformal Fluid to Rindler Fluid



Momentum Relaxation Rate







Ref: by Khimphun, Lee, Park, Zhang JHEP 1801 (2018) 058

## From AdS/CFT to Holographic Rindler Fluid

Holographic Screen the Time-like boundary

$$\mathcal{T}_{ab} = -2(K_{ab} - K\gamma_{ab})$$

What is Rindler Fluid?

Fluid dual to Rindler spacetime



AdS Spacetime

Navier-Stokes Equations:BredFluid/Gravity Expansion:CompetendeEntropy Current and Constraint:Comparison with AdS/Fluid:Comparison with AdS/Fluid:MRindler Fluid and Recurrence RelationRindler Fluid with Momentum Relaxation

Flat Spacetime

Bredberg, Keeler, Lysov, Strominger ['10,'11] Compere, McFadden, Skenderis, Taylor ['11,'12] t: Chirco, Eling, Liberati, Meyer, Oz ['12,'13] Matsuo, Natsuume, Ohta, Okamura ['12,13] ation Cai, Li, Yang, Zhang ['13,'14] xation Khimphun, Lee, Park, Zhang ['17]

# **Rindler Fluid in Relativistic Derivative Expansion**

### **Rindler Metric with Perturbations**

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = -2pu_{a}dx^{a}dr + g_{ab}dx^{a}dx^{b},$$
where  $g_{ab} = g_{ab}^{(0)} + g_{ab}^{(1)} + g_{ab}^{(2)},$ 
 $g_{ab}^{(0)} = -p^{2}(r - r_{c})u_{a}u_{b} + \gamma_{ab},$ 
 $g_{ab}^{(1)} = 2p(r - r_{c})\left(u^{c}\partial_{c}\ln pu_{a}u_{b} + 2a_{(a}u_{b)}\right),$ 
 $g_{ab}^{(2)} = 2(r - r_{c})\left[\left(\mathcal{K}_{cd}\mathcal{K}^{cd}\right)u_{a}u_{b} - 2u_{(a}h_{b)}^{c}\partial_{d}\mathcal{K}_{c}^{d}\right]$ 
 $-\mathcal{K}_{a}^{c}\mathcal{K}_{cb} + 2\mathcal{K}_{c(a}\Omega_{b)}^{c} - 2h_{a}^{c}h_{b}^{d}u^{e}\partial_{e}\mathcal{K}_{cd}\right]$ 
 $+p^{2}(r - r_{c})^{2}\left\{\left(\frac{1}{2}\mathcal{K}_{cd}\mathcal{K}^{cd} + a_{c}a^{c}\right)u_{a}u_{b}\right.$ 
 $+2u_{(a}h_{b)}^{c}\left[\partial_{d}\mathcal{K}_{c}^{d} - \left(\mathcal{K}_{cd} + \Omega_{cd}\right)a^{d}\right] - \Omega_{ac}\Omega_{b}^{c}\right\}$ 

# Holographic Fluid $T_{ab} = T_{ab}^{(0)} + T_{ab}^{(1)} + T_{ab}^{(2)} + O(\partial^3),$ $T_{ab}^{(0)} = I_{ab}$

$$\begin{split} T_{ab}^{(0)} &= \mathbb{p}h_{ab}, \\ T_{ab}^{(1)} &= \zeta'(u^c\partial_c \mathrm{ln}\mathbb{p})u_a u_b - 2\eta \mathcal{K}_{ab}, \\ T_{ab}^{(2)} &= \mathbb{p}^{-1} \Big\{ \Big[ d_1 \mathcal{K}_{ab} \mathcal{K}^{ab} + d_2 \Omega_{ab} \Omega^{ab} + d_3 (u^c\partial_c \mathrm{ln}\mathbb{p})^2 \\ &+ d_4 u^c \partial_c (u^d\partial_d \mathrm{ln}\mathbb{p}) + d_5 h^{cd} (\partial_c \mathrm{ln}\mathbb{p}) (\partial_d \mathrm{ln}\mathbb{p}) \Big] u_a u_b \\ &+ \Big[ c_1 \mathcal{K}_{ac} \mathcal{K}^c_{\ b} + c_2 \mathcal{K}_{c(a} \Omega^c_{\ b)} + c_3 \Omega_{ac} \Omega^c_{\ b} + c_4 h^c_a h^d_b \partial_c \partial_d \mathrm{ln}\mathbb{p} \\ &+ c_5 \mathcal{K}_{ab} (u^c \partial_c \mathrm{ln}\mathbb{p}) + c_6 (h^c_a \partial_c \mathrm{ln}\mathbb{p}) (h^d_b \partial_d \mathrm{ln}\mathbb{p}) \Big] \Big\}. \end{split}$$

### **Constraint Equations**

 $\partial_a u^a = 2\mathbb{p}^{-1}\mathcal{K}_{ab}\mathcal{K}^{ab} + O(\partial^3),$  $a_a + h^b_a\partial_b \ln\mathbb{p} = 2\mathbb{p}^{-1}h^c_a\partial_b\mathcal{K}^b_c + O(\partial^3).$ 

Ref: Eling, Meyer, Oz , JHEP 1205 (2012) 116 Compère, McFadden, Skenderis ,Taylor, JHEP 1203 (2012) 076

### Transport coefficients

$$\zeta' = 0, \qquad \eta = 1,$$
  
 $d_1 = -2, \quad d_2 = d_3 = d_4 = d_5 = 0,$   
 $c_1 = -2, \quad c_2 = c_3 = c_4 = c_5 = -c_6 = -4.$ 

### **Rindler Fluid and Recurrence Relation**

Petrov Type I Condition

$$\mathbb{P}_{ab} \equiv \boldsymbol{n}^r h_a^c \boldsymbol{n}^r h_b^d C_{rcrd}$$

$$T_{ab}^{(0)} = \mathcal{E}u_a u_b + \mathcal{P}h_{ab} \stackrel{\mathbb{P}_{ab}=0}{\Longrightarrow} T_{ab}^{(1)} = -2\eta\sigma_{ab} + \dots \stackrel{\mathbb{P}_{ab}=0}{\Longrightarrow} T_{ab}^{(2)} = \dots$$

$$4\mathbb{P}_{ab} = h_a^m h_b^n \left[ \left( T_{mc} T_{nd} - T_{mn} T_{cd} \right) u^c u^d - T_{mc} T_n^c \right. \\ \left. - 4u^c \partial_c T_{mn} + 4u^c \partial_{(m} T_{n)c} \right] \\ \left. + p^{-2} \left[ T(T + p T_{cd} u^c u^d) + 4p u^c \partial_c T \right] h_{ab}.$$

Petrov type I Spacetime and Dual Relativistic Fluids [Cai, Yang, Zhang, PRD90 (2014) 041901(R)]



[Cai, Yang, Zhang, JHEP 1412 (2014) 147]

### I. Flat Bulk/Rindler Fluid

Rindler Metric 
$$ds^2 = -rd\tau^2 + \frac{1}{r}dr^2 + dx_i dx^i$$
  
Induced Metric  $ds^2 = -r_c d\tau^2 + dx_i dx^i$   
Dual Tensor  $\mathcal{T}_{ab} = -2(K_{ab} - K\gamma_{ab})$ 

Constraint equations

$$2G_{\mu b}n^{\mu}|_{r_c} = 2\partial^a (K_{ab} - \gamma_{ab}K) = 0 \implies \partial^a T_{ab} = 0$$

$$2G_{\mu\nu}n^{\mu}n^{\nu}|_{r_c} = (K^2 - K_{ab}K^{ab}) = 0 \implies T^2 - pT_{ab}T^{ab} = 0$$

Non-relativistic limit -> Incompressible Navier-Stokes Equations

$$\partial_i v^i = 0$$
  $\partial_i P + \partial_\tau v_i + v^j \partial_j v_i - \eta \partial^2 v_i = 0$ 

Bredberg, Keeler, Lysov, Strominger (JHEP 07 (2012) 146)

## II. Flat Bulk/de-Siter Fluid

## 1) Holographic Stress Tensor — Dark Sectors

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \kappa_4 T_{\mu\nu} + \kappa_4 \langle \mathcal{T} \rangle_{\mu\nu}, \quad \langle \mathcal{T} \rangle_{\mu\nu} \equiv \frac{1}{\kappa_4 L} \left( \mathcal{K}_{\mu\nu} - \mathcal{K} g_{\mu\nu} \right)$$

Modified Einstein equations

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} - \frac{1}{L}\left(\mathcal{K}_{\mu\nu} - \mathcal{K}g_{\mu\nu}\right) = \kappa_4 T_{\mu\nu}$$

Hamiltonian constraints

 $\mathcal{K}^2 - \mathcal{K}_{\mu\nu}\mathcal{K}^{\mu\nu} = R + 2\,G_{MN}^{(d+1)}\mathcal{N}^M\mathcal{N}^N,$ 

$$\langle T_{\mu\nu} 
angle = -\frac{\Lambda}{\kappa_4} g_{\mu\nu}$$
  
 $L = \frac{\kappa_5}{\kappa_4}$ 



2) Embedding in higher dimensions — Brane World Models

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \mathcal{T}^{\mathcal{M}}_{\mu\nu} + T^{B}_{\mu\nu},$$
  
$$\mathcal{T}^{\mathcal{M}}_{\mu\nu} \equiv (\mathcal{K} g_{\mu\sigma} - \mathcal{K}_{\mu\sigma}) \mathcal{K}^{\sigma}_{\ \nu} + \mathcal{M}_{\mu\nu} - \frac{1}{2} \left( \mathcal{K}^{2} - \mathcal{K}_{\rho\sigma} \mathcal{K}^{\rho\sigma} \right) g_{\mu\nu},$$
  
$$\mathcal{M}_{\mu\nu} \equiv g^{\ M}_{\mu} g^{\ N}_{\nu} R^{(d+1)}_{MN} - g^{\ M}_{\mu} \mathcal{N}^{P} g^{\ N}_{\nu} \mathcal{N}^{Q} R^{(d+1)}_{MPNQ}.$$

[Maeda, Mukohyama, Sasaki, Shiromizu, ..., ['99, '10]



### Ref: 1106.2476 [Living Rev. '10]

### III. Flat Bulk/FRW Fluid



The Total Action  $S_{tot} = \int_{\mathcal{H}} d^4 x \sqrt{-g} \left( \frac{1}{2\kappa_4} R + \mathcal{L}_m \right) + S_5 ,$   $S_5 \equiv \int_{\mathcal{M}} d^5 x \sqrt{-\tilde{g}} \left( \frac{1}{2\kappa_5} \mathcal{R} \right) + \int_{\mathcal{H}} d^4 x \sqrt{-g} \frac{1}{\kappa_5} \mathcal{K},$ 



**Einstein Field Equations** 

$$\frac{1}{\kappa_4}G_{\mu\nu} = T^m_{\mu\nu} + \langle \mathcal{T} \rangle^d_{\mu\nu},$$

### Holographic Dark Fluid?

Stress Energy Tensors

$$T^m_{\mu\nu} \equiv -\frac{2}{\sqrt{-g}} \frac{\delta(S_m)}{\delta g^{\mu\nu}}, \qquad \langle \mathcal{T} \rangle^d_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta(\mathcal{S}_5)}{\delta g^{\mu\nu}} = \frac{1}{\kappa_5} \left( \mathcal{K}_{\mu\nu} - \mathcal{K}g_{\mu\nu} \right)$$

In the Bulk => Modified GR

Hamilton Constraint Equation 
$$\delta_{CSZ} \equiv \Omega_D^2 - \frac{1}{2}\Omega_{\Lambda}(\Omega_D - \Omega_B) \simeq -0.003$$
  
& Emergent de-Sitter Universe  $\delta_V \equiv \Omega_D^2 - \frac{4}{3}\Omega_B \simeq 0.004$ .

Ref: [arXiv: JHEP 1810 (2018) 009] by Cai, Sun, Zhang

### holographic Emergent Dark Universe (hEDU) & SNIa Data



Ref: [arXiv: 1812.11105] with Gansukh Tumurtushaa(IBS/Korea) et al

### Late-Time Evolution of Holographic FRW Model



Hamiltonian Constraint from FRW Hypersurface

$$\tilde{\Omega}_D^2 = \frac{\tilde{\Omega}_\Lambda}{2(1+3\tilde{w}_D)} \big[ \tilde{\Omega}_D (1-3\tilde{w}_D) - \tilde{\Omega}_B \big].$$

$$\tilde{\Omega}_{\Lambda} = \Omega_{\Lambda}, \qquad \tilde{\Omega}_{D} = \Omega_{\mathcal{H}}(t) - \Omega_{\Lambda},$$

$$1 \qquad \dot{\Omega}_{\alpha}(t)$$

$$\tilde{w}_D = -1 - \frac{1}{3H(t)} \frac{\Omega_{\mathcal{H}}(t)}{\Omega_{\mathcal{H}}(t) - \Omega_{\Lambda}},$$

$$\tilde{\Omega}_B \equiv \frac{\Omega_M}{a(t)^3} = \frac{H(t)^2}{H_0^2} - \Omega_{\mathcal{H}}(t),$$
$$\Omega_{\mathcal{H}}(t) \equiv \frac{\rho_{\mathcal{H}}}{\rho_c} = \Omega_{\Lambda}^{1/2} \left[ \frac{H(t)^2}{H_0^2} + \frac{\Omega_I}{a(t)^4} \right]^1$$

Late Time Universe 
$$\begin{split} &\delta_{CSZ} \equiv \Omega_D^2 - \frac{1}{2}\Omega_{\Lambda}(\Omega_D - \Omega_B) \simeq -0.003, \\ &\delta_V \equiv \Omega_D^2 - \frac{4}{3}\Omega_B \simeq 0.004. \qquad \Omega_{\Lambda} \simeq 0.685, \quad \Omega_D \simeq 0.265, \quad \Omega_B \simeq 0.050, \end{split}$$

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## **Effective Potential & Swampland Criteria**

$$\begin{split} \text{Effective Action} \qquad S_{tot} &= \int d^4x \sqrt{-g} \Big[ \frac{1}{2\kappa_4} R + \mathcal{L}_m - \frac{1}{2} (\partial \phi)^2 - V(\phi) \Big]. \\ \hline V[\phi(t)] &= \frac{1}{2} [\rho_d(t) - p_d(t)], \\ \dot{\phi}(t) &= -\sqrt{\rho_d(t) + p_d(t)}. \\ \end{split}$$

Ref: [arXiv: 1812.11105] Emergent Dark Universe and the Swampland Criteria Yun-Long Zhang Holographic Dark Fluid 19

### Summary & Outlook



Yun-Long Zhang Holographic Dark Fluid