

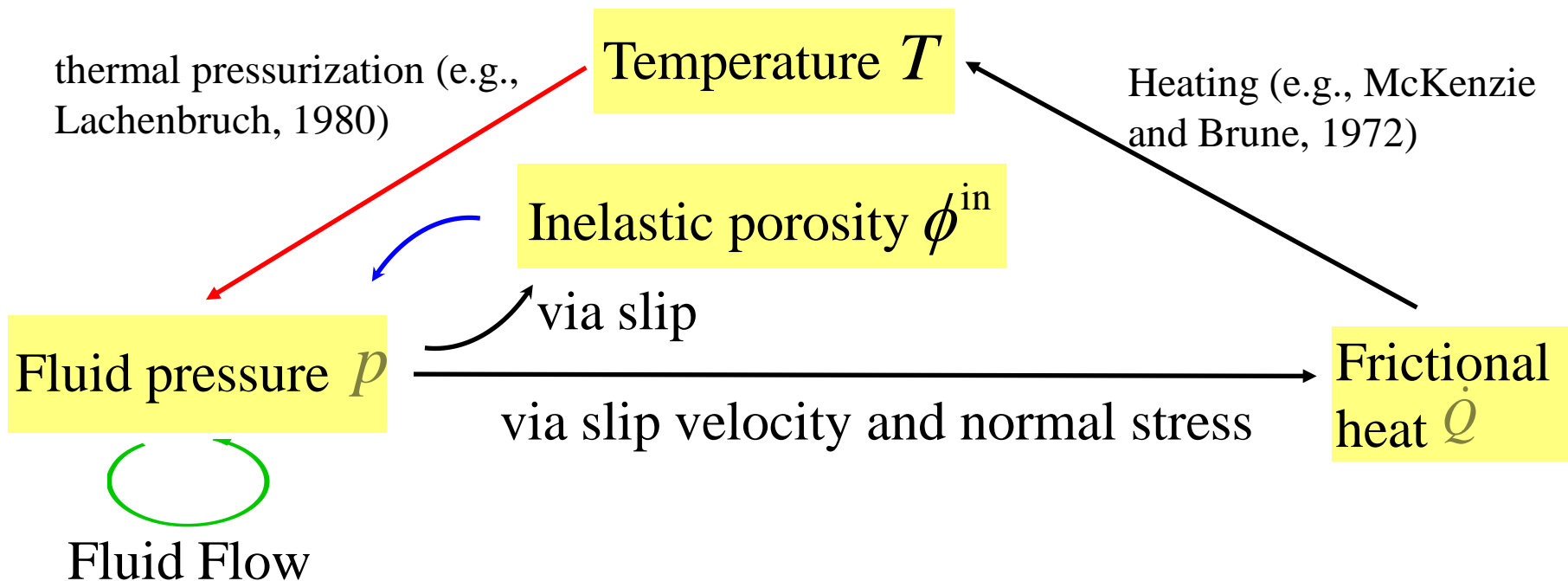


Phase Transition and Power Law
Appearing in Friction Behavior in the
Medium with Heat, Fluid Pressure and
Dilatancy

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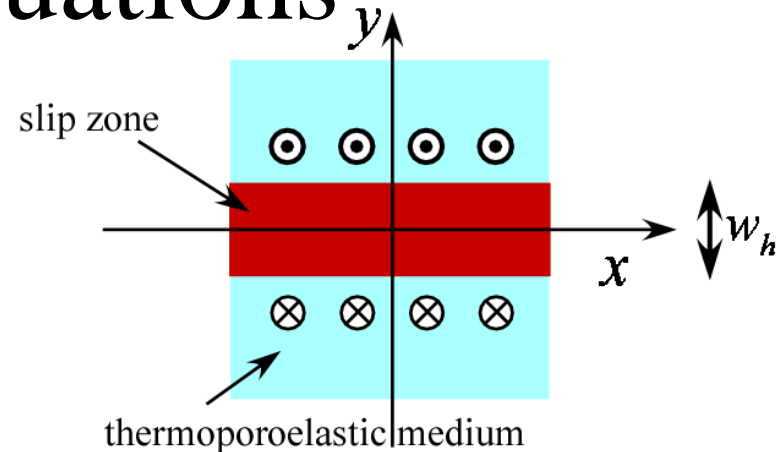
Model Setup

- A one-dimensional (1-D) mode III fault is assumed. The interaction among heat, fluid pressure and inelastic pore creation is investigated. See details in Suzuki and Yamashita (2010, 2014).
- In particular, temporal evolution of the porosity is investigated because it produces an important universality and analytical study of such a behavior gives some implications for understanding the slip behavior.



Governing Equations

- From the interaction mentioned above, ...



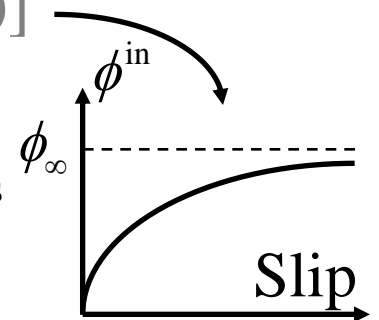
Fluid pressure $\frac{1}{M} \frac{\partial \bar{p}_f}{\partial t} = ((b - \phi_t) \alpha_s + \phi_t \alpha_f) \frac{\partial \bar{T}}{\partial t} + \frac{k}{\eta} \nabla^2 \bar{p}_f - \frac{\partial \phi^{\text{in}}}{\partial t}$

Temperature $\left[(1 - \phi_t) \rho_s C_s + \phi_t \rho_f C_f \right] \frac{\partial \bar{T}}{\partial t} = \sigma_{\text{fric}} \frac{v}{w_h} \left[H\left(y + \frac{w_h}{2}\right) - H\left(y - \frac{w_h}{2}\right) \right]$

Porosity $\frac{\partial \phi^{\text{in}}}{\partial t} = \alpha_0 v \left(1 - \frac{\phi^{\text{in}}}{\phi_\infty} \right) \left[H\left(y + \frac{w_h}{2}\right) - H\left(y - \frac{w_h}{2}\right) \right]$

e.o.m. $\rho_B \frac{\partial^2}{\partial t^2} (\bar{u}_s)_z = \mu_v \nabla^2 (\bar{u}_s)_z$

The fluid pressure affects the boundary condition for e.o.m.



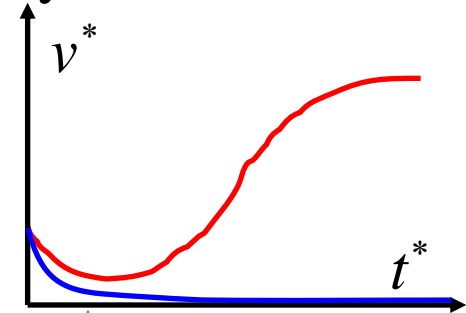
Two Qualitatively Different Slip Behaviors

- If the fluid flow is neglected, governing equations can be rewritten in terms of the normalized slip velocity and inelastic porosity:

$$\dot{v}^* = v^* (1 - v^*) - S_u (1 - \phi^*) v^*$$

$$\dot{\phi}^* = T_a (1 - \phi^*) v^*$$

S_u, T_a : Nondimensional numbers



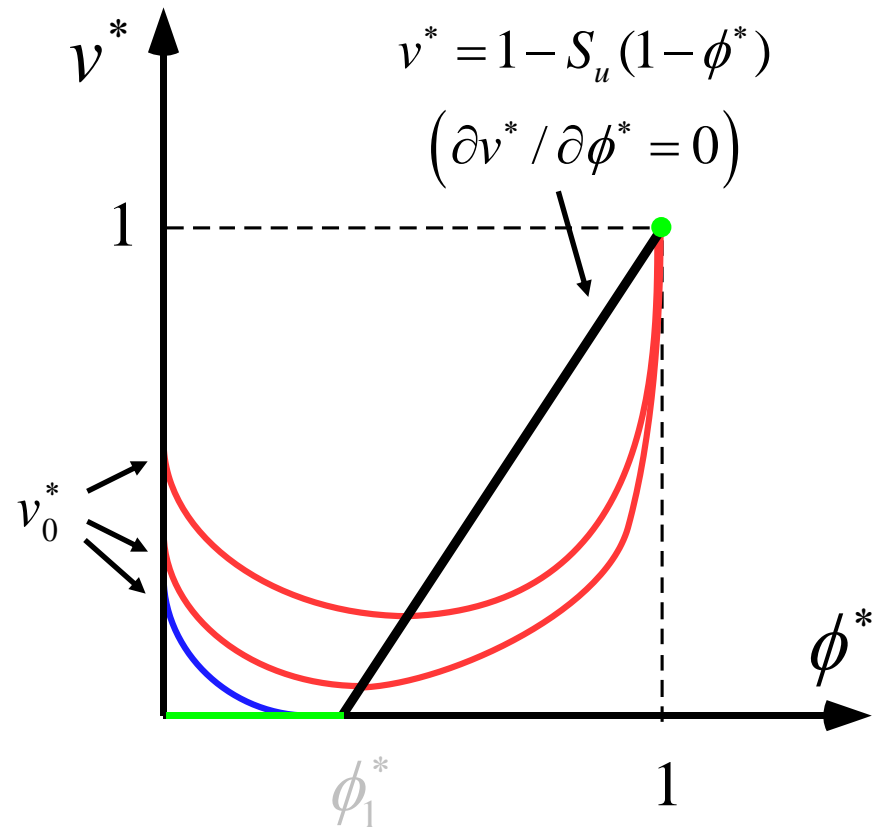
- We assume $S_u > 1 - v_0^*$ (v_0^* is the initial value of v^*). At the slip onset ($t^* = 0$), we can easily show $\dot{v}^* < 0$.
- From brief consideration, we clearly find two qualitatively different slip behaviors.
 - First case: $S_u (1 - \phi^*)$ decreases with increasing time and \dot{v}^* can become positive. → Acceleration
 - Second case: \dot{v}^* is always negative. → Spontaneous slip cessation
- The steady state solutions for the governing equations are $(v_s^*, \phi_s^*) = (0, \text{indeterminate})$ and $(1, 1)$.

Phase Diagram

- All orbits of solutions are first apparently absorbed into $v^* = 0$.
- If the orbits do not cross the straight line $\partial v^* / \partial \phi^* = 0$, they are finally absorbed to the stable steady state

$(v_s^*, \phi_s^*) = (0, \text{indeterminate})$.

- On the other hand, if the orbits cross the line $\partial v^* / \partial \phi^* = 0$, the relationship $\partial v^* / \partial \phi^* > 0$ is satisfied and they are finally absorbed into the stable steady state $(v_s^*, \phi_s^*) = (1, 1)$.



(1) The Function G Determining Slip Behavior Completely

- The function $G(S_u, T_a, v_0^*)$ determining the system behavior can be given by (Suzuki and Yamashita, 2014)

$$G(S_u, T_a, v_0^*) = 1 - \left(\frac{1}{S_u T_a} \left((1 - v_0^*) T_a + S_u - 1 + v_0^* \right) \right)^{\frac{T_a}{T_a - 1}} \cdot S_u$$

- If $G > 0$,
 - \dot{v}^* becomes positive. $\rightarrow (v_s^*, \phi_s^*) = (1, 1)$
- If $G < 0$,
 - \dot{v}^* does not become positive and v^* approaches zero.
 $\rightarrow (v_s^*, \phi_s^*) = (0, \text{indeterminate})$

(2) Critical Porosity ϕ_1^*

- For the acceleration case, the deceleration changes to the acceleration at $\dot{v}^* = 0$, and we define v_a^* and ϕ_a^* as the values of v^* and ϕ^* at that time:

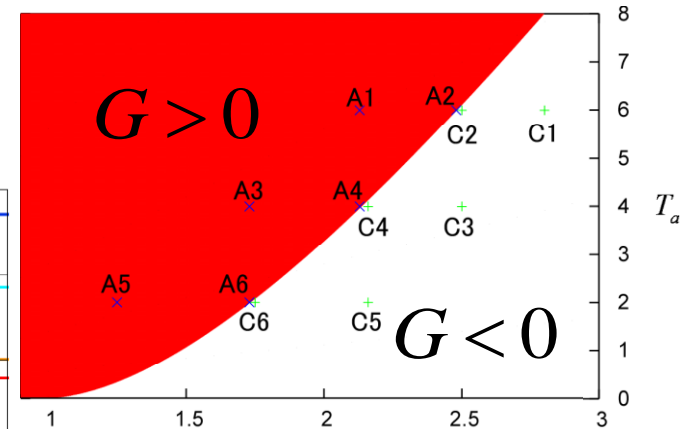
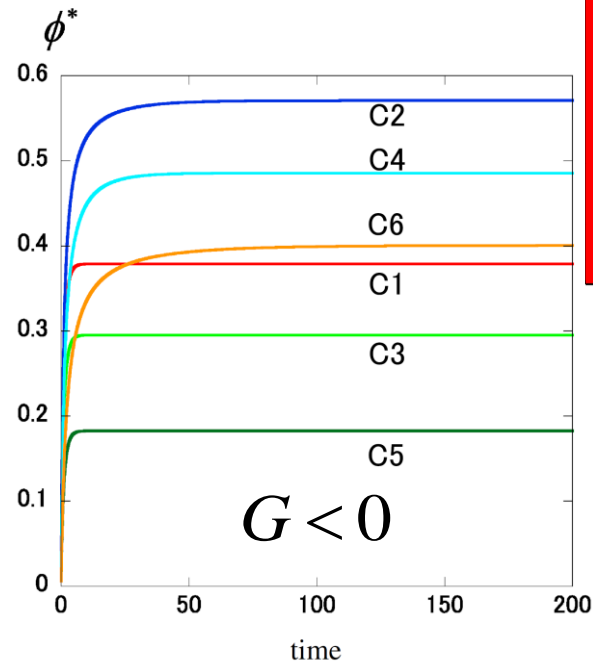
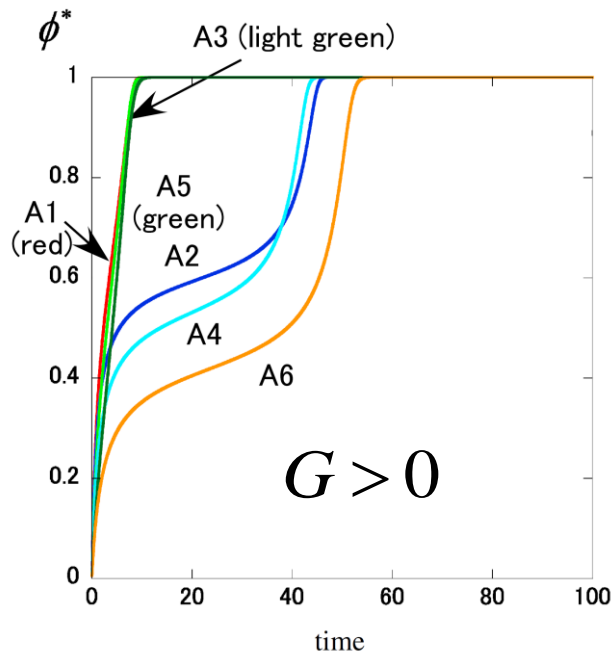
$$\dot{v}^* = (1 - v_a^* - S_u(1 - \phi_a^*))v_a^* = 0$$

- Since the condition $v_a^* > 0$ must be satisfied, we obtain

$$\phi_a^* > \frac{S_u - 1}{S_u} (\equiv \phi_1^*)$$

- We can conclude that if $\phi^* > \phi_1^*$ is satisfied, $(v_s^*, \phi_s^*) = (1, 1)$ can emerge.
- In other words, it is the upper value of ϕ^* for the case $(v_s^*, \phi_s^*) = (0, \text{indeterminate})$: $\phi_s^* \leq \phi_1^*$

Temporal Changes in Porosity



$$(v_0^* = 0.1)$$

$$\phi_1^*$$

C1:0.64, C2:0.6

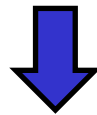
C3:0.6, C4:0.54

C5:0.54, C6:0.43

- If $G > 0$, $\phi^* \rightarrow 1$.
- If $G < 0$, ϕ^* approaches the constant values.
 - The constant values are below ϕ_1^* .

Phase Transition and Power Law

- We introduce the variables $w^* \equiv 1 - v^*$ and $\psi^* \equiv 1 - \phi^*$ and change $w_0^* \equiv 1 - v_0^*$ as a parameter.
 - If $G > 0$, $(w^*, \psi^*) \rightarrow (0, 0)$
 - If $G < 0$, $(w^*, \psi^*) \rightarrow (1, 1 - \phi_s^*)$
- For $G < 0$, the value ϕ_s^* changes continuously with changing v_0^* .



- The power law near (w_c^*, ψ_c^*) is expected for $G < 0$:

$$\psi_\infty^* - \psi_c^* \sim (w_0^* - w_c^*)^\alpha \quad \left(G(S_u, T_a, 1 - w_c^*) = 0, \quad w_c^* = \frac{S_u^{1/T_a} T_a - S_u}{T_a - 1} \right)$$

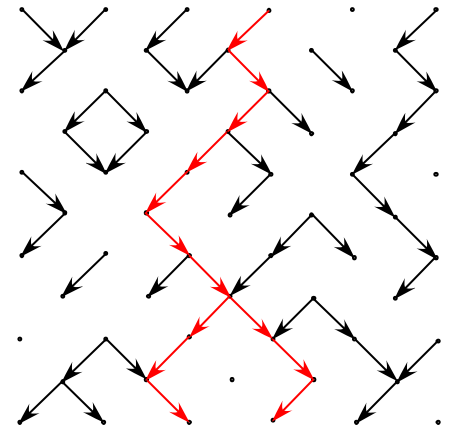
$$\psi_\infty^* = 1 - \phi_s^* \quad \left(\psi_c^* \equiv 1 - \phi_1^* = 1 - \frac{S_u - 1}{S_u} = \frac{1}{S_u} \right)$$

Absorbing Phase Transition

- Once the system enters one state, the system cannot escape the state.
- ex.) Directed percolation
 - Percolation in a lattice with channels
 - Each channel is open with the possibility p , while it is closed with the possibility $1 - p$.
 - Flow occurs from upstream to downstream.
 - If we define P_∞ as the possibility that the infinite percolation occurs, P_∞ is known to show a critical phenomenon:

$$P_\infty \sim \begin{cases} (p - p_c)^\beta & (p \geq p_c) \\ 0 & (p \leq p_c) \end{cases}$$

(p_c : critical value)



Power Law (Universality...?)

- Figure shows the power law near (w_c^*, ψ_c^*) :

$$\psi_\infty^* - \psi_c^* = \beta(w_0^* - w_c^*)^\alpha$$

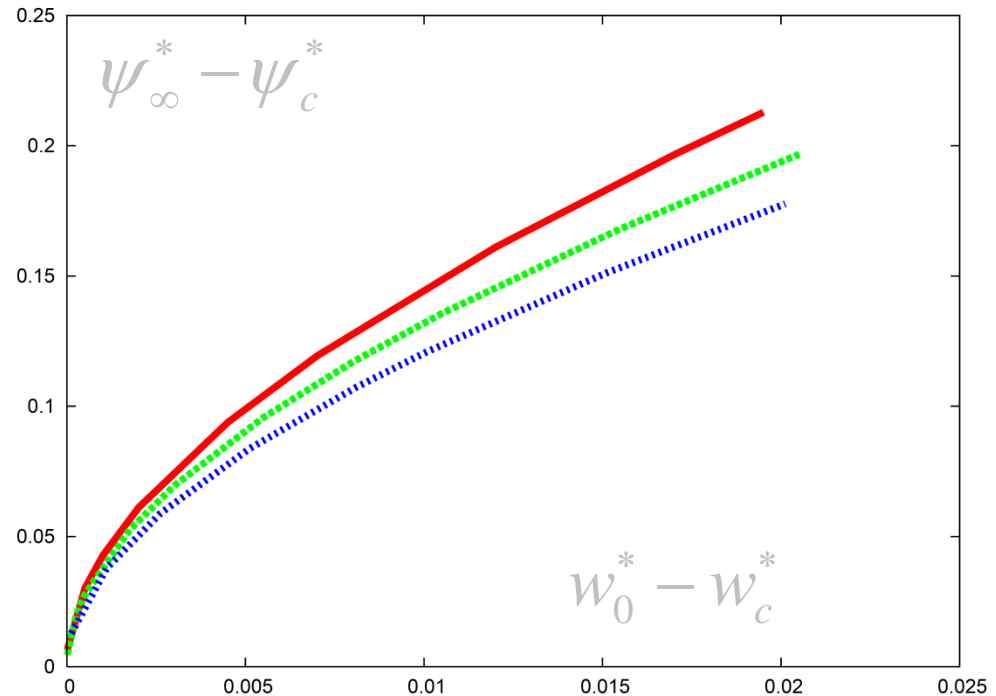
- Values for α are:

$$\alpha = 0.548 \quad (S_u = 2.5, T_a = 6)$$

$$\alpha = 0.538 \quad (S_u = 1.5, T_a = 2)$$

$$\alpha = 0.540 \quad (S_u = 2, T_a = 3)$$

- Universal?
- If this value is universal, porosity values observed in natural faults are related to the initial slip velocity and we may presume the ancient slip behavior from the porosity.



- $S_u = 2.5, T_a = 6$
- ⋯ $S_u = 1.5, T_a = 2$
- ⋯ $S_u = 2, T_a = 3$

Conclusions and Future Works

- Phase transition observed in porosity change within the framework including the heat, fluid pressure and dilatancy
- Universal value for α ?
 - New index characterizing the slip behavior?
- Analytical treatments to show the universality is required.

