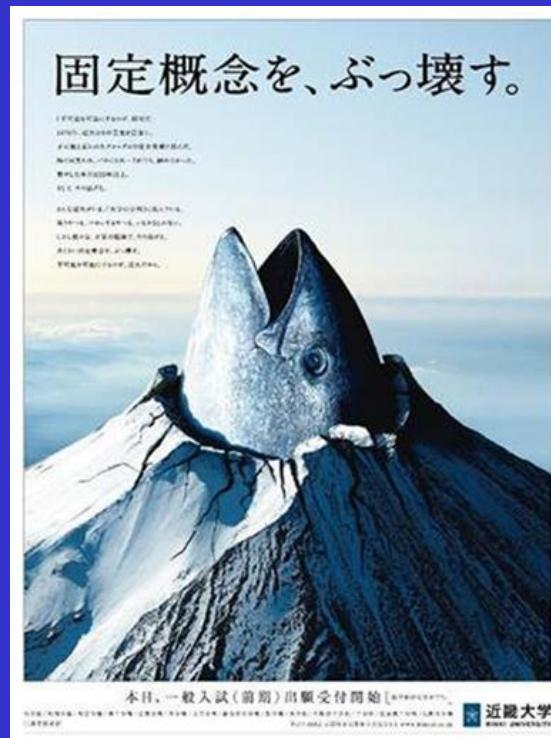


Dynamical Approach to Synthesis of Superheavy Elements

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*Long-term workshop on
“Computational Advances in Nuclear
and Hadron Physics” (CANHP2015)
YITP, Kyoto, Japan, 1st October 2015*

Topics

Structure

Light Nuclei

Reaction

Superheavy Nuclei

Key Words

1. Shell Correction Energy
Two-Center Shell Model
2. Dynamical Approach

Contents

1. Introduction

Super Heavy Elements

Stability of Nuclei, Shell effects

Synthesis of SHE

2. Experimental methods

3. Theoretical calculation

Dynamical model (Fluctuation Dissipation model)

Langevin equation

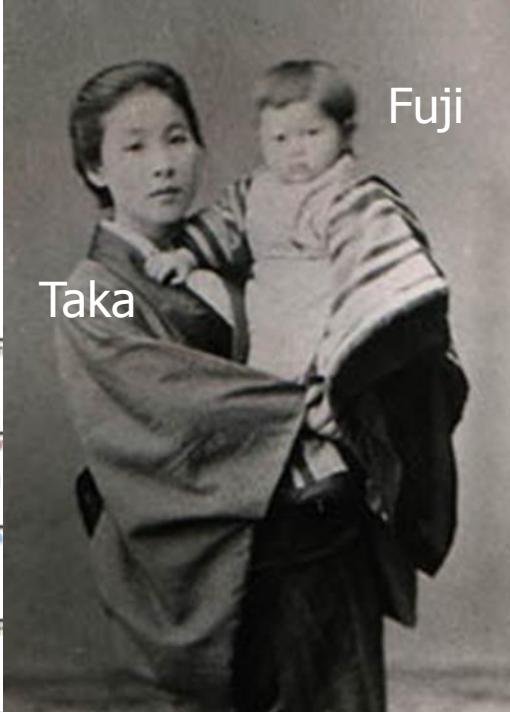
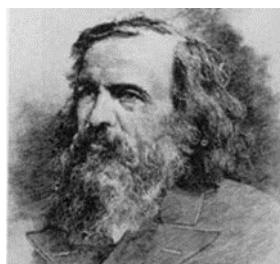
4. Calculation results

fusion-fission process

5. Further study

Periodic Table

1	Hydrogen 1	2
2	Lithium 3	Beryllium 4
3	Sodium 11	Magnesium 12
4	Potassium 19	Calcium 20
5	Rubidium 37	Strontium 38
6	Cesium 55	Barium 56
7	Franclium Fr	Radium Ra
	Actinium Ac**	Rutherfordium Rf
	Thorium 104	Dubnium 105
	Seaborgium Sg	Singapore 106



Fuji

Taka

13	14	15	16	17	18
Boron 5	Carbon 6	Nitrogen 7	Oxygen 8	Fluorine 9	Neon 10
Aluminum 13	Silicon 14	Phosphorus 15	Sulfur 16	Chlorine 17	Argon 18
In 31	Gallium 32	Germanium 33	Arsenic 34	Selenium 35	Bromine 36
Zn 31	Ga 32	Ge 33	As 34	Se 35	Kr 36
Cadmium 49	In 50	Tin 51	Antimony 52	Tellurium 53	Iodine 54
Ce 58	Hafnium 72	Tantalum 73	Bismuth 83	Polonium 84	Xenon 86
Curium 90	Thallium 81	Lead 82	Astatine 85	Radon 86	
Db 96	Tl 81	Pb 82	Po 84	At 85	Rn 86
Sg 106	113	114	115	116	117 118



Менделеев(1834-1907)
1869

May 2012 IUPAC

Fl

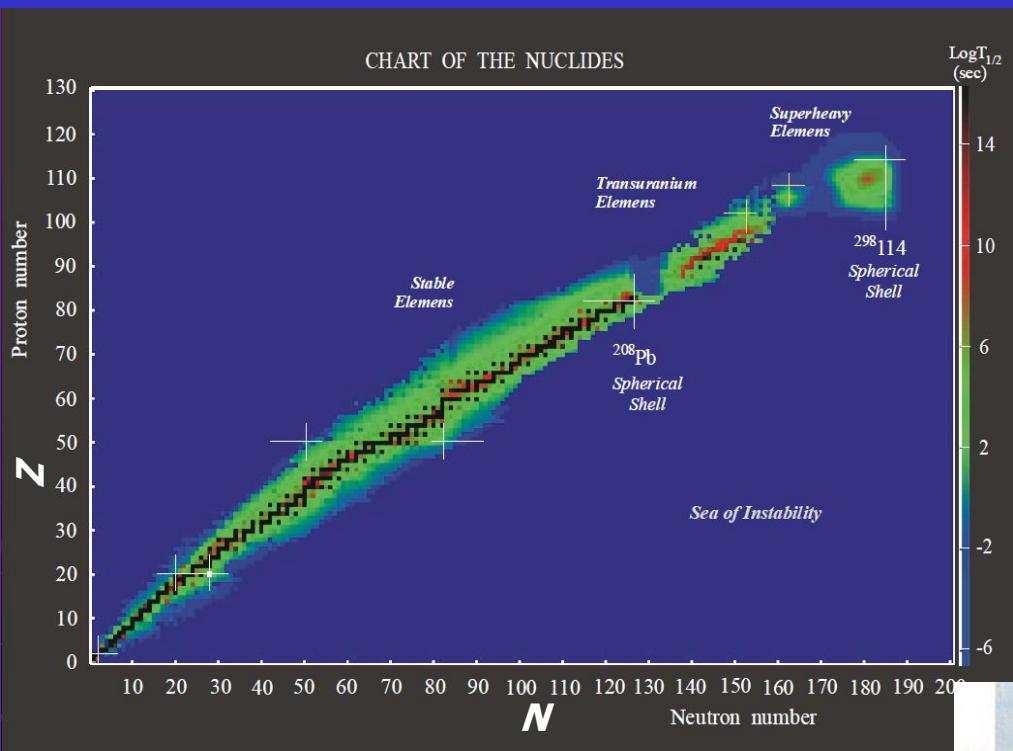
Lv

flerovium livermorium

La*	Lanthanides	Lanthanum 57	Cerium 58	Praseodymium 59	Neodymium 60	Promethium 61	Samarium 62	Europium 63	Gadolinium 64	Terbium 65	Dysprosium 66	Holmium 67	Erbium 68	Thulium 69	Ytterbium 70	Lutetium 71
Ac**	Actinides	Actinium 89	Thorium 90	Protactinium 91	Uranium 92	Neptrium 93	Plutonium 94	Americium 95	Curium 96	Berkelium 97	Californium 98	Einsteinium 99	Fermium 100	Mendeleovium 101	Nobelium 102	Lawrencium 103

Super Heavy Elements → less stable

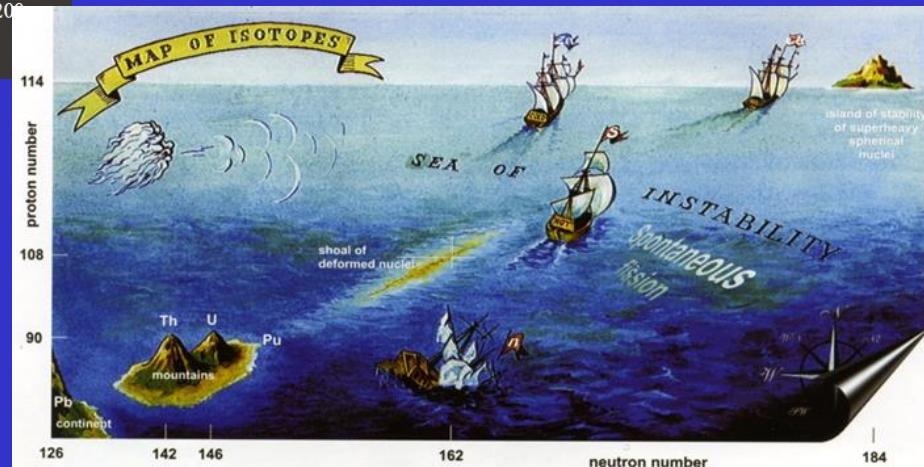
1. Introduction

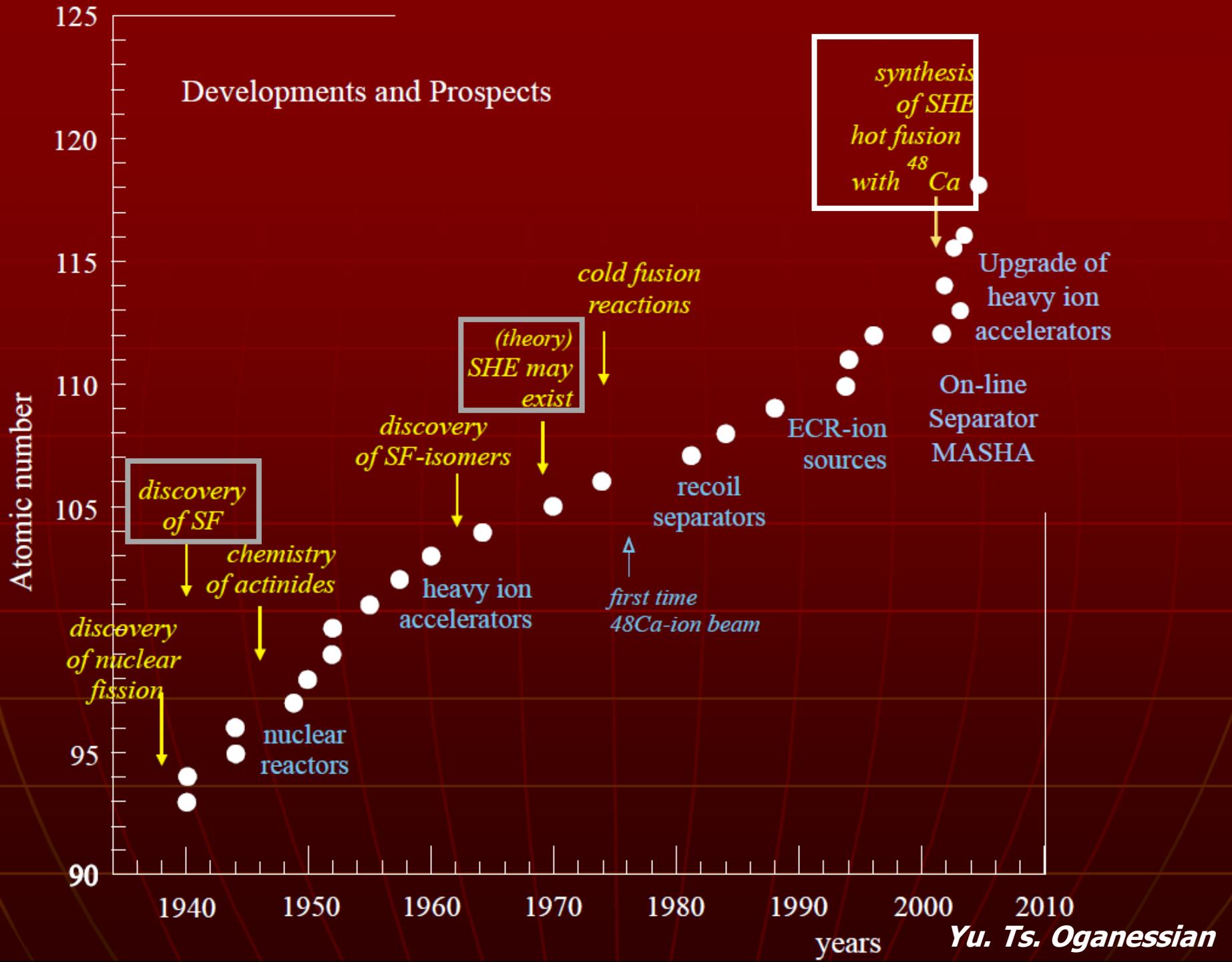


Nuclear Chart
Stability of nuclei

Our Interest

- Next magic number $\leftarrow Z=82, N=126$
- Verification of 'Island of Stability'
(predicted by macroscopic-microscopic model in 1960's)
- Synthesis of new elements





Experimental setup for synthesis of SHE

Lab	Country	City	Accelerator	Separator
FLNR	Russia	Dubna	U400 U400M	DGFRS VASSILISSA
GSI	Germany	Darmstadt	UNILAC	SHIP TASCA
RIKEN	Japan	Wako	RILAC	GALIS
LBNL	USA	Berkeley	88-inch Cyclotron	BGS
GANIL	France	Caen	<i>SPIRAL2's LINAC accelerator</i>	<i>S3 (Super Separator Spectrometer)</i>



FLNR (Russia)

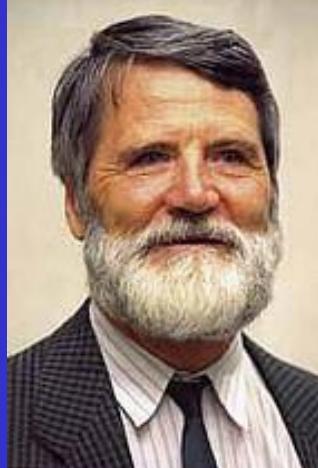


G.N. Flerov
(1913 -1990)



Yu.Ts. Oganessian
(1933-)

GSI (Germany)



P. Armbruster **G. Muenzenberg**
(1931-) **(1940-)**



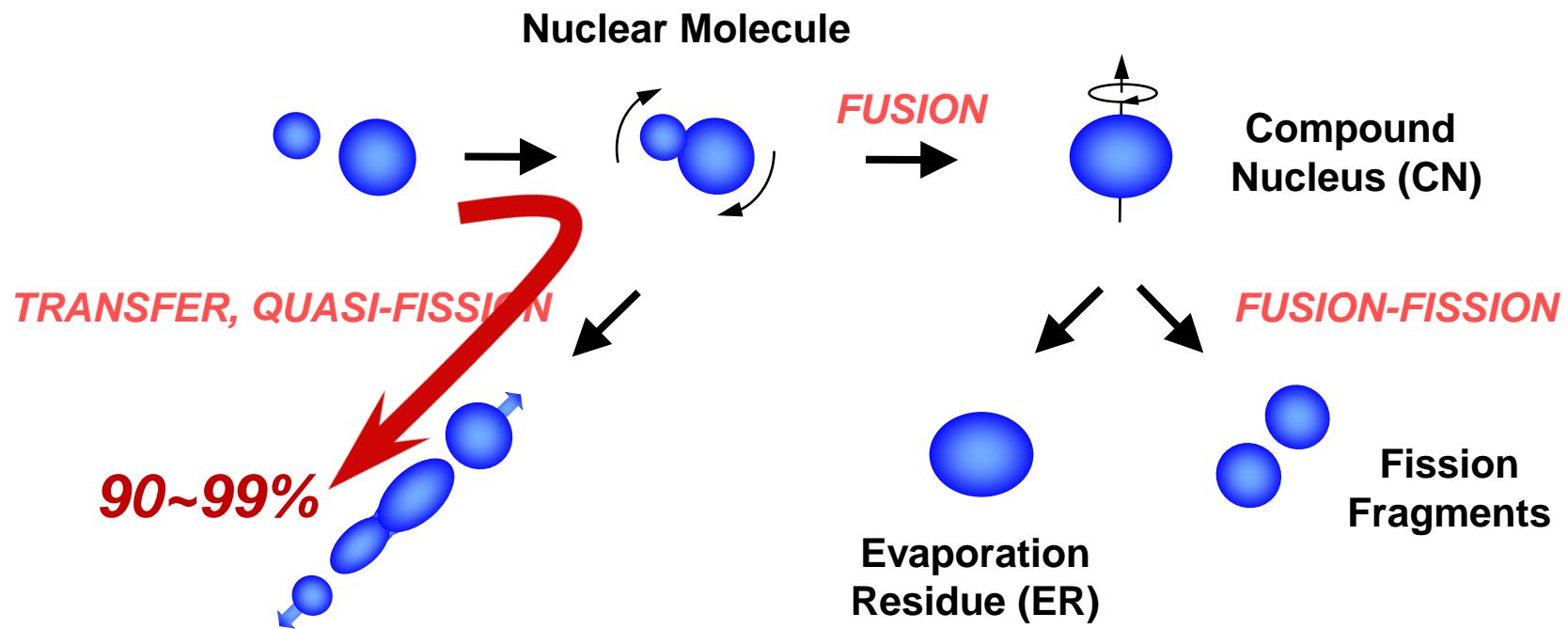
S. Hofmann
(1943-)

RIKEN (Japan)



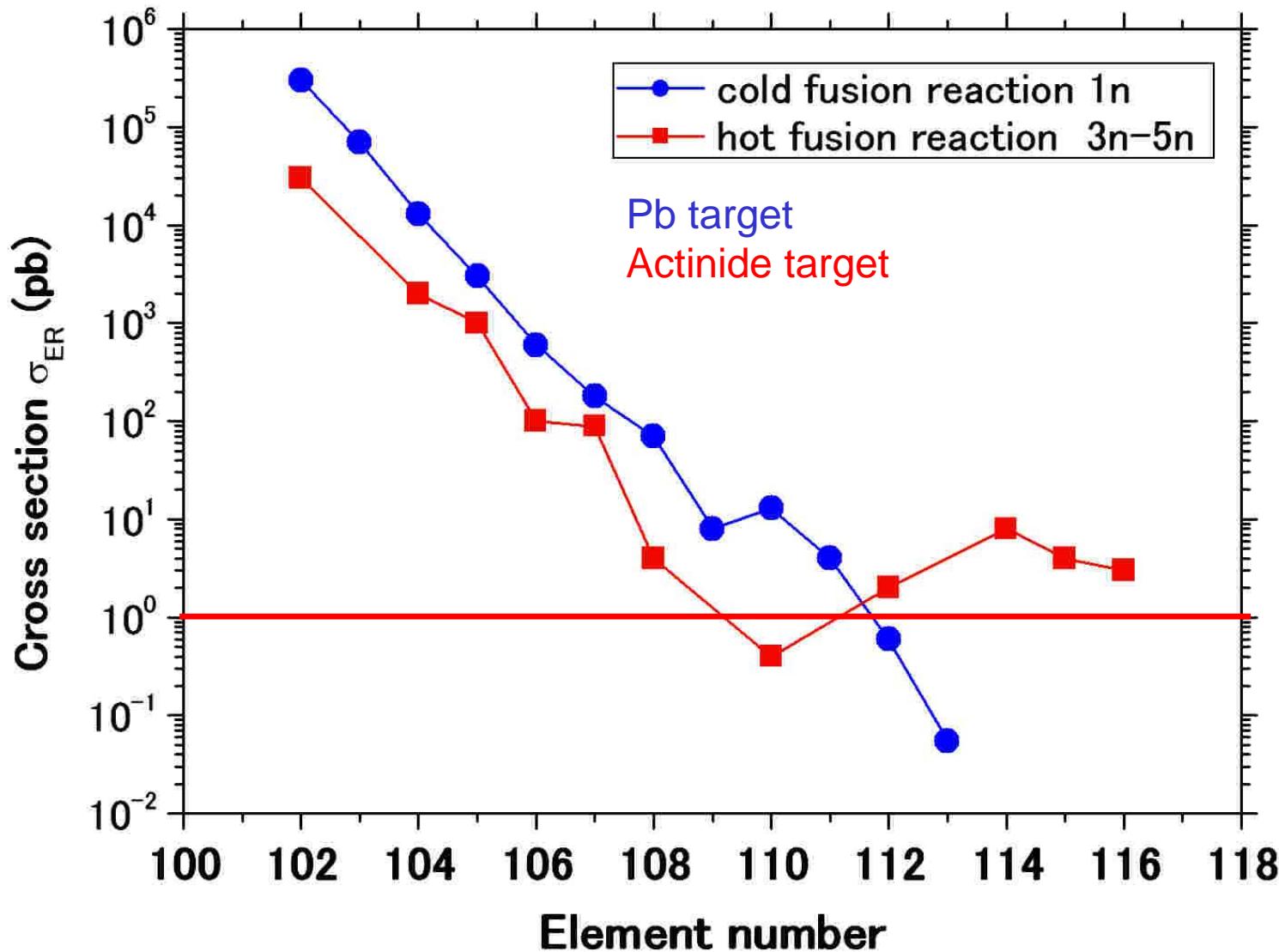
K. Morita
(1957-)

Fusion process in Superheavy mass region

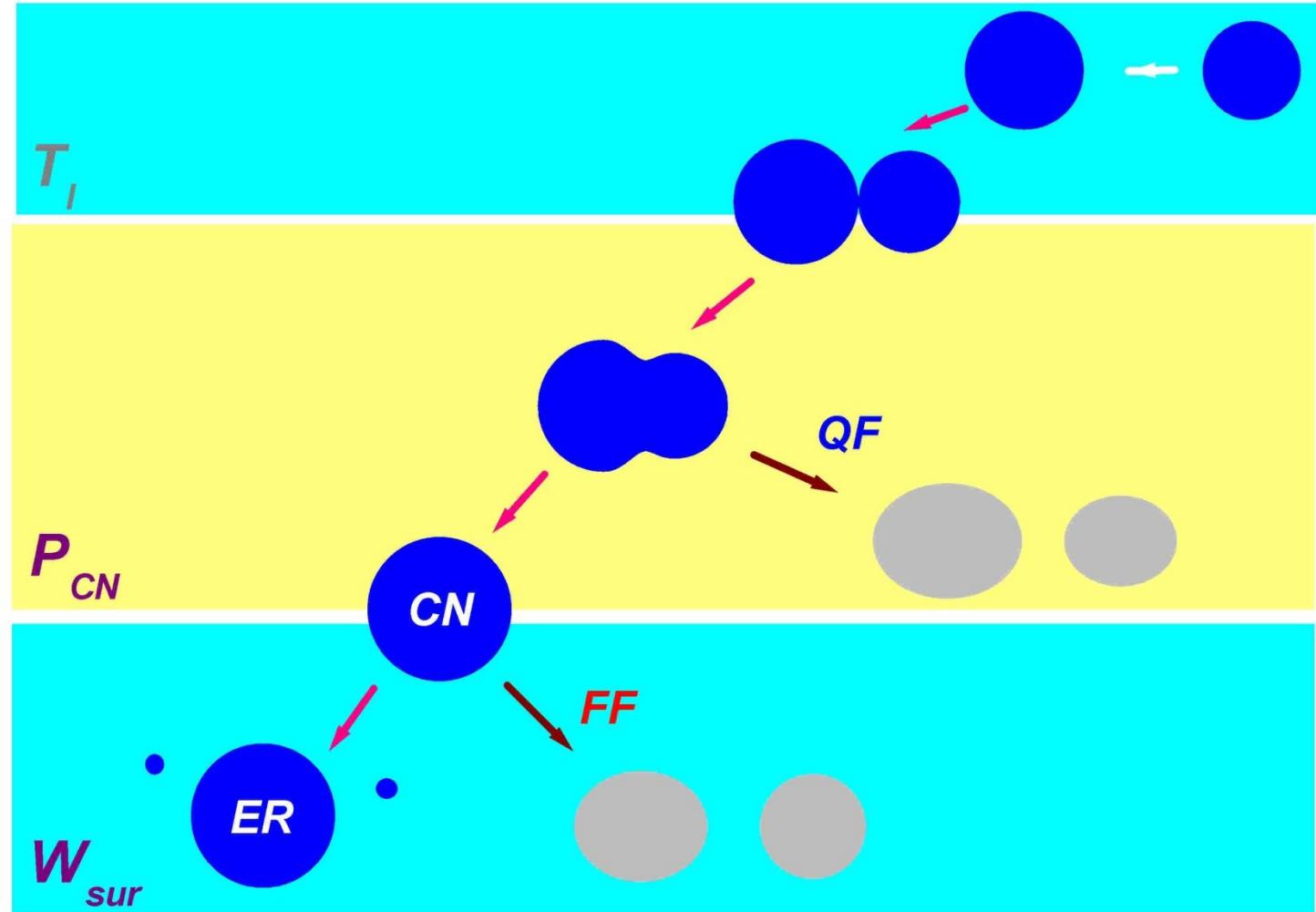


Experimental data

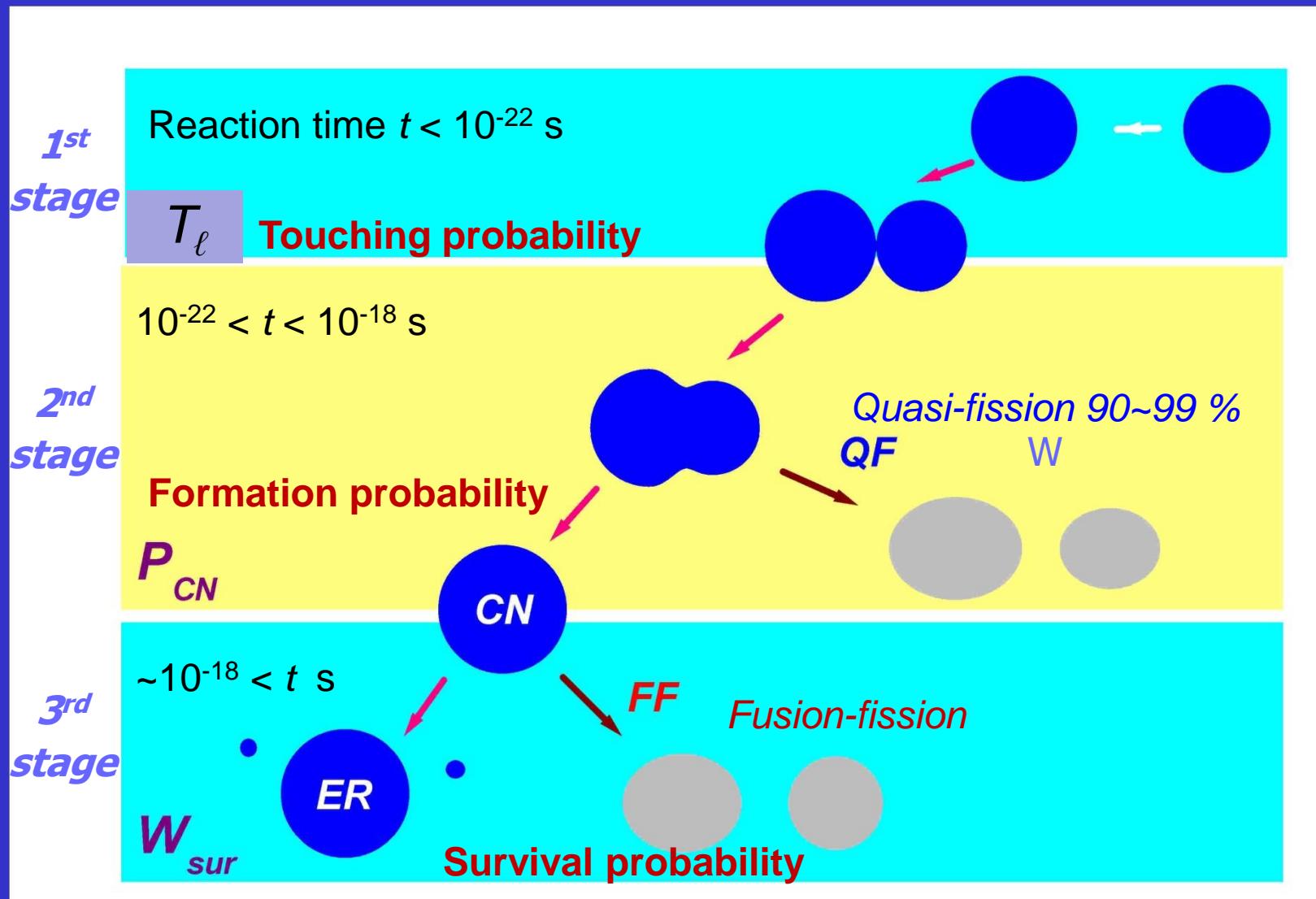
Evaporation residue cross sections



$$\sigma_{ER} = \frac{\pi \hbar^2}{2\mu_0 E_{cm}} \sum_{\ell=0}^{\infty} (2\ell+1) T_{\ell}(E_{cm}, \ell) P_{CN}(E^*, \ell) W(E^*, \ell)$$



$$\sigma_{ER} = \frac{\pi \hbar^2}{2 \mu_0 E_{cm}} \sum_{\ell=0}^{\infty} (2\ell+1) T_{\ell}(E_{cm}, \ell) P_{CN}(E^*, \ell) W(E^*, \ell)$$



Synthesis of New Elements

Reports of new elements

Heavy ion reaction

Cold fusion reaction Hot fusion reaction

1994



1996



1999



2000



2002



2003

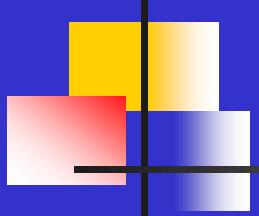


2004



2010





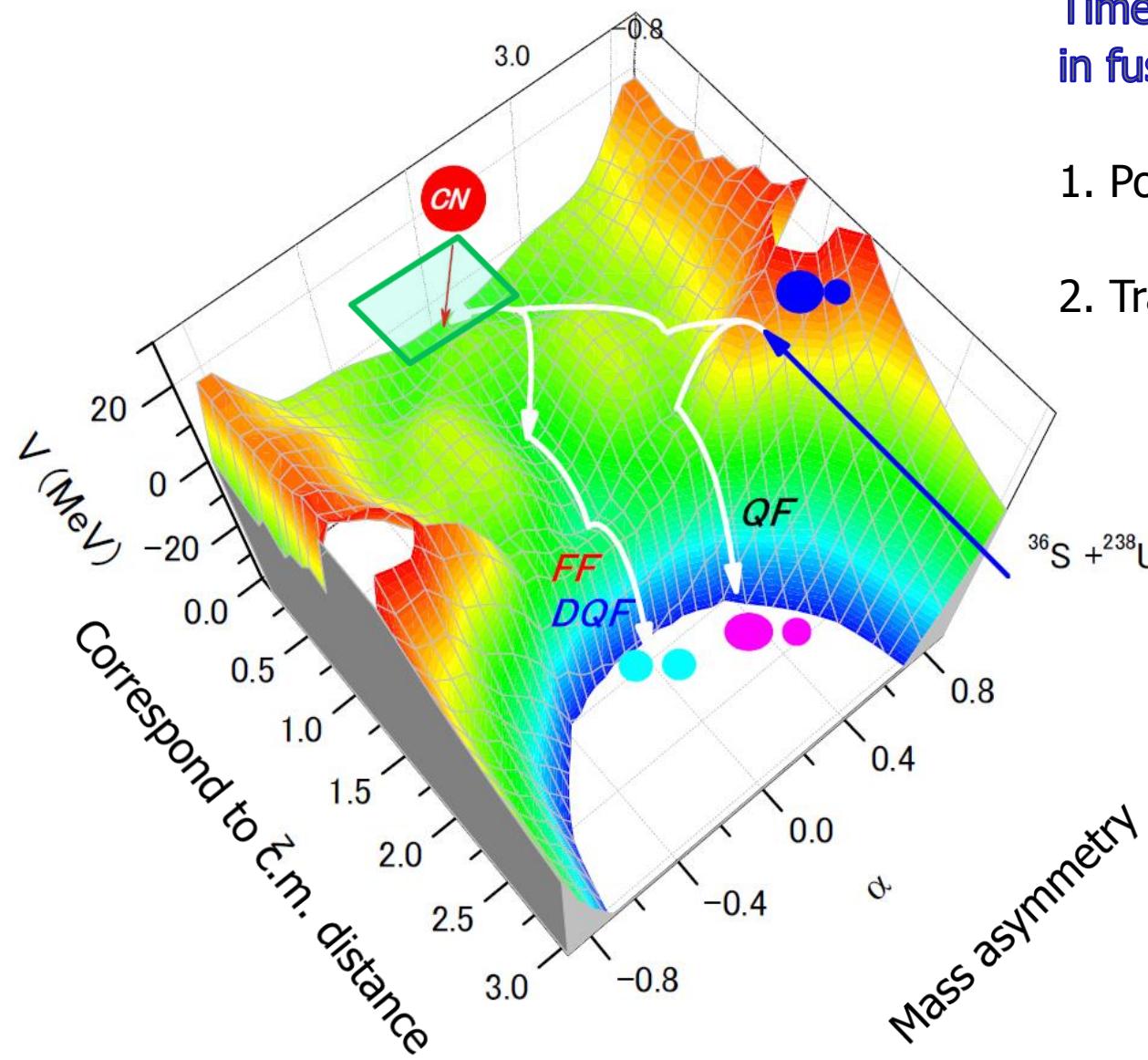
2. Model

2-1. Estimation of cross sections

2-2. Dynamical Equation

Overview of Dynamical Process in reaction $^{36}\text{S} + ^{238}\text{U}$

Time-evolution of nuclear shape
in fusion-fission process



1. Potential energy surface
2. Trajectory → described by equations

Nuclear shape

two-center parametrization (z, δ, α)

(Maruhn and Greiner,
Z. Phys. 251(1972) 431)

$q(z, \delta, \alpha)$

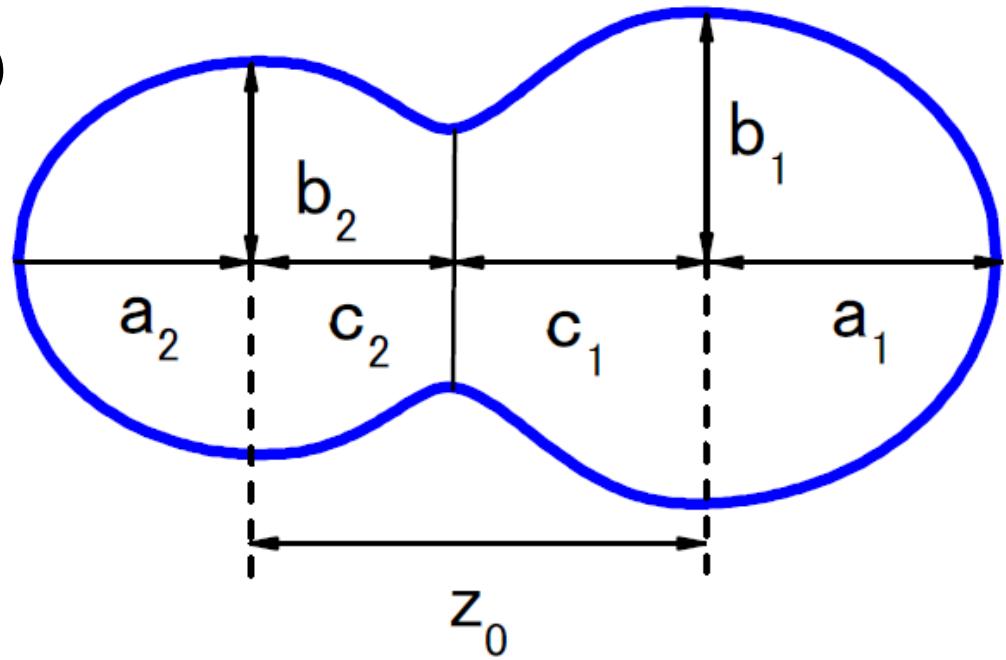
$$z = \frac{z_0}{BR}$$

$$B = \frac{3+\delta}{3-2\delta}$$

R : Radius of the spherical compound nucleus

$$\delta = \frac{3(a-b)}{2a+b} \quad (\delta_1 = \delta_2)$$

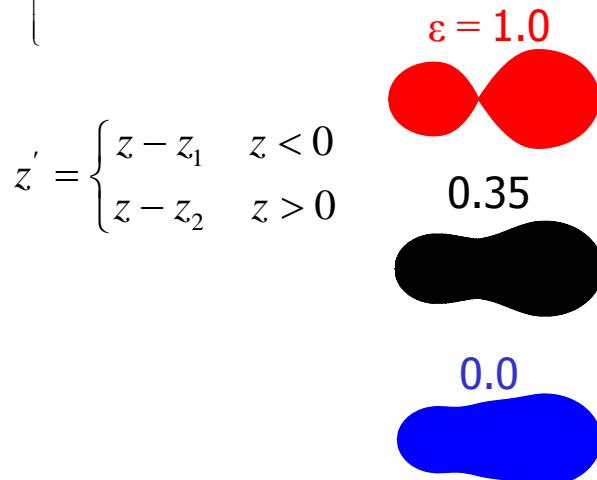
$$\alpha = \frac{A_1 - A_2}{A_{CN}}$$



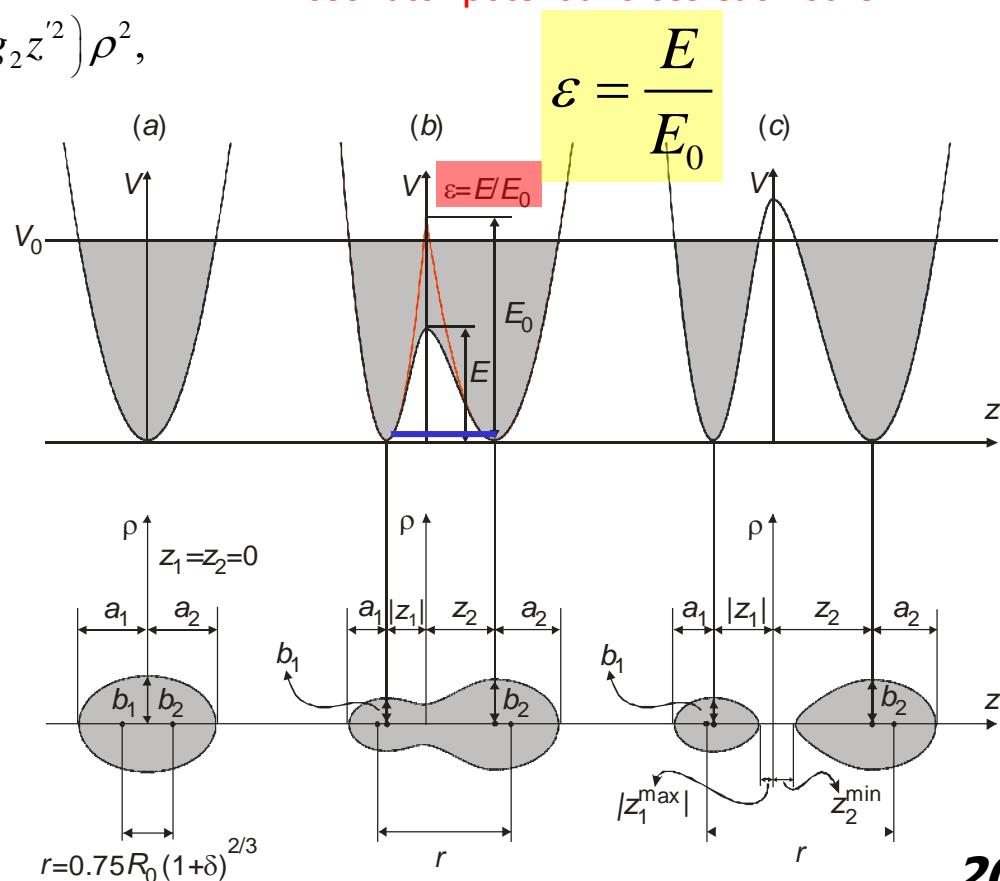
Two Center Shell Model

$$\hat{H} = -\frac{\hbar^2}{2m_0} \nabla^2 + V(\mathbf{r}) + V_{LS}(\mathbf{r}, \mathbf{p}, \mathbf{s}) + V_{L^2}(\mathbf{r}, \mathbf{p}).$$

$$V(\rho, z) = \frac{1}{2} m_0 \begin{cases} \omega_{z1}^2 z'^2 + \omega_{\rho 1}^2 \rho^2, & z < z_1 \\ \omega_{z1}^2 z'^2 (1 + c_1 z' + d_1 z'^2) + \omega_{\rho 1}^2 (1 + g_1 z'^2) \rho^2, & z_1 < z < 0 \\ \omega_{z2}^2 z'^2 (1 + c_2 z' + d_2 z'^2) + \omega_{\rho 2}^2 (1 + g_2 z'^2) \rho^2, & 0 < z < z_2 \\ \omega_{z2}^2 z'^2 + \omega_{\rho 2}^2 \rho^2, & z > z_2, \end{cases}$$



Neck parameter is the ratio of smoothed potential height to the original one where two harmonic oscillator potential cross each other



Potential Energy

$$V(q, \ell, T) = V_{DM}(q) + \frac{\hbar^2 \ell(\ell+1)}{2I(q)} + V_{SH}(q, T)$$

$$V_{DM}(q) = E_S(q) + E_C(q)$$

$$V_{SH}(q, T) = E_{shell}^0(q) \Phi(T)$$

T : nuclear temperature

$$E^* = aT^2 \quad a : \text{level density parameter}$$

Toke and Swiatecki

E_S : Generalized surface energy (finite range effect)

E_C : Coulomb repulsion for diffused surface

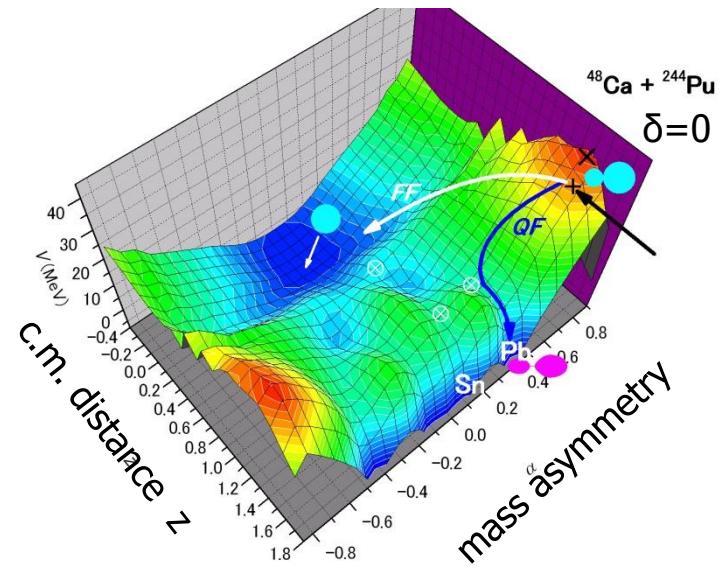
E_{shell}^0 : Shell correction energy at $T=0$

I : Moment of inertia for rigid body

$\Phi(T)$: Temperature dependent factor

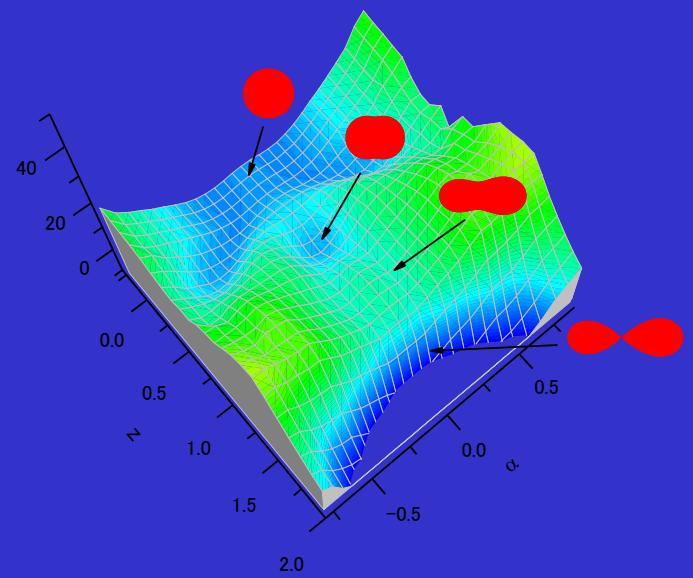
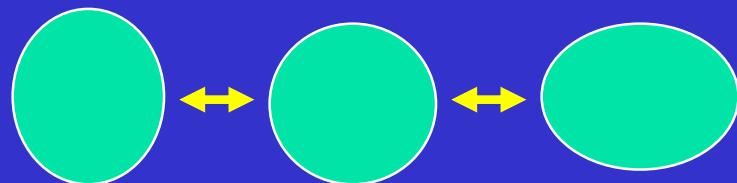
$$\Phi(T) = \exp \left\{ -\frac{aT^2}{E_d} \right\}$$

$$E_d = 20 \text{ MeV}$$



Taking into account the fluctuation around the mean trajectory

Thermal fluctuation of nuclear shape
→ thermal fluctuation of collective motion



Multi-dimensional Langevin Equation

$$\frac{dq_i}{dt} = (m^{-1})_{ij} p_j$$

Friction
dissipation

$$\frac{dp_i}{dt} = -\frac{\partial V}{\partial q_i} - \frac{1}{2} \frac{\partial}{\partial q_i} (m^{-1})_{jk} p_j p_k - \gamma_{ij} (m^{-1})_{jk} p_k + g_{ij} R_j(t)$$

Random force
fluctuation

Newton equation

ordinary differential equation

$\langle R_i(t) \rangle = 0, \langle R_i(t_1)R_j(t_2) \rangle = 2\delta_{ij}\delta(t_1-t_2)$: white noise (Markovian process)

$$\sum_k g_{ik} g_{jk} = T\gamma_{ij}$$

Einstein relation

Fluctuation-dissipation theorem

q_i : deformation coordinate

(nuclear shape)

two-center parametrization (z, δ, α)

(Maruhn and Greiner, Z. Phys. 251(1972) 431)

p_i : momentum

m_{ij} : Hydrodynamical mass

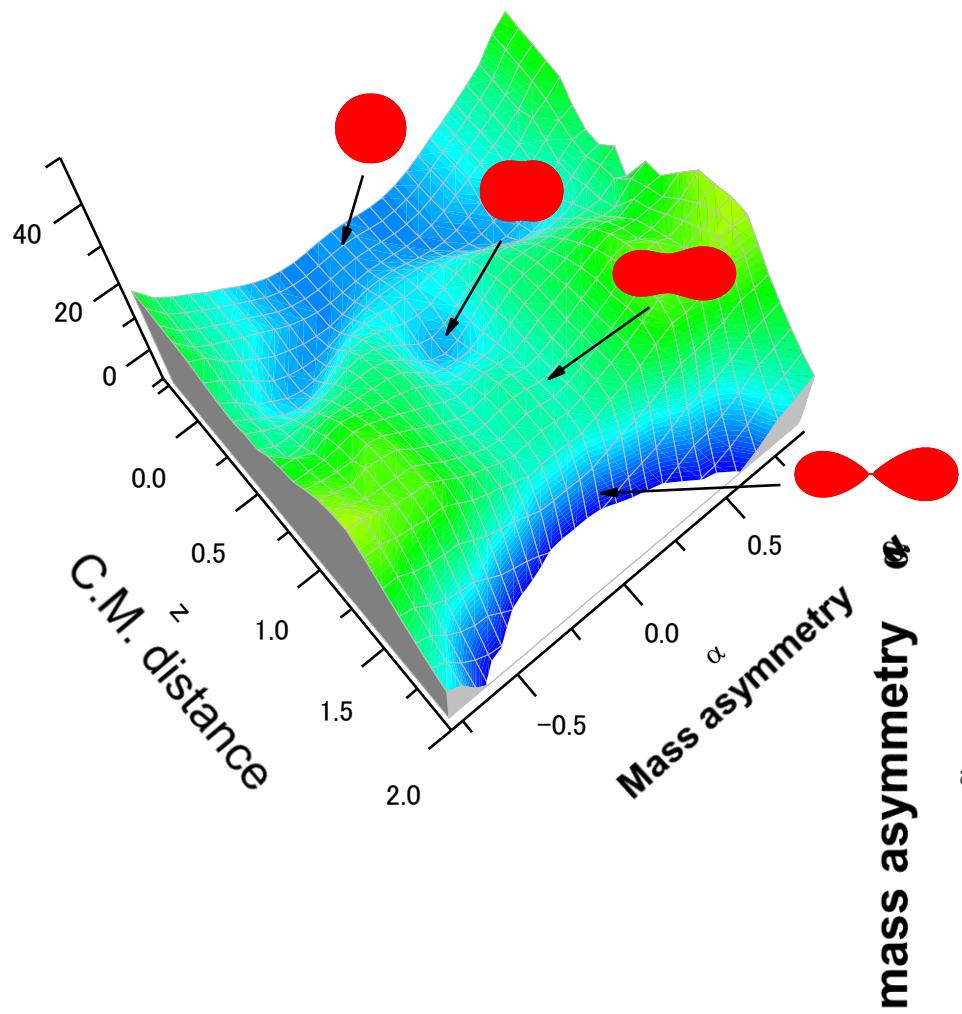
(inertia mass)

γ_{ij} : Wall and Window (one-body) dissipation (friction)

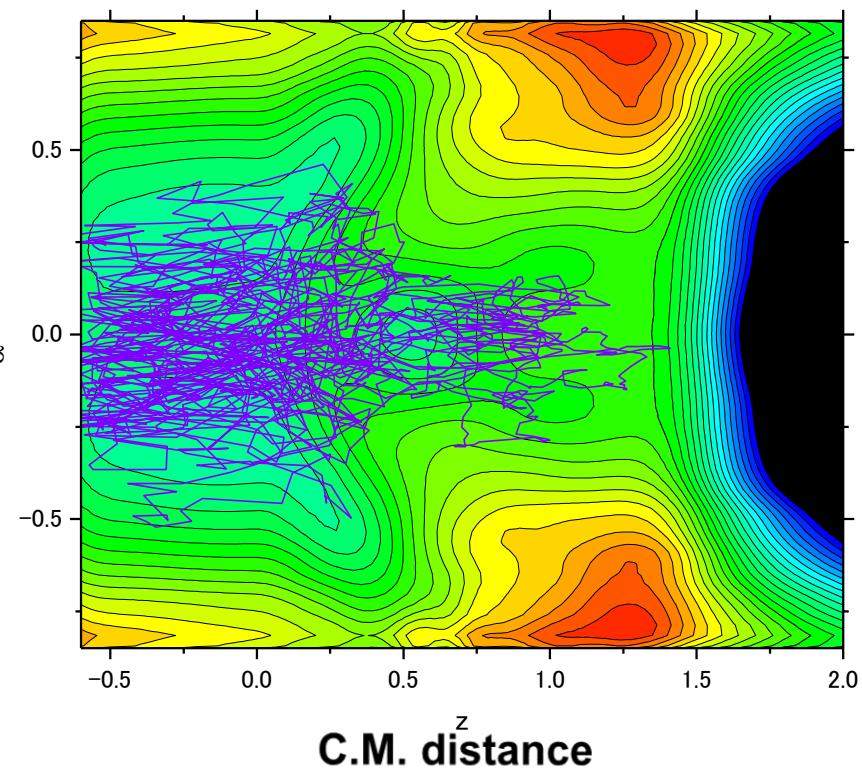
$$E_{\text{int}} = E^* - \frac{1}{2} (m^{-1})_{ij} p_i p_j - V(q)$$

E_{int} : intrinsic energy, E^* : excitation energy

Fission process ^{240}U $E^* < 20$ MeV



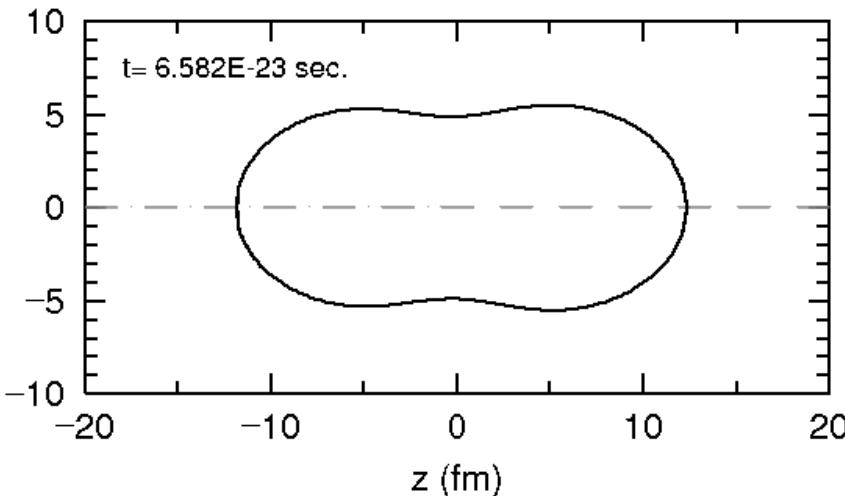
Trajectory on potential energy surface



236U $E^*=20$ MeV

揺動項なし Newton Eq.

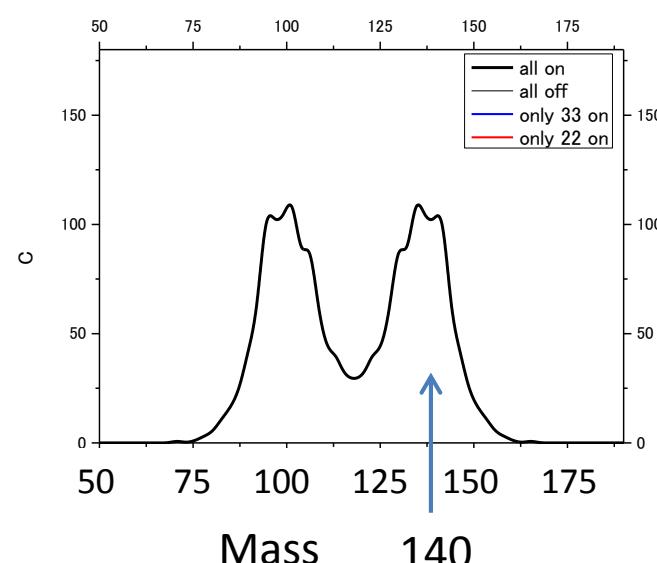
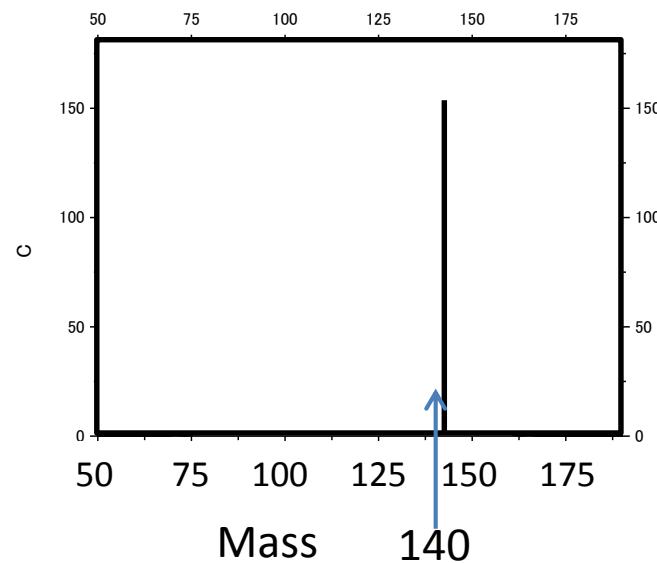
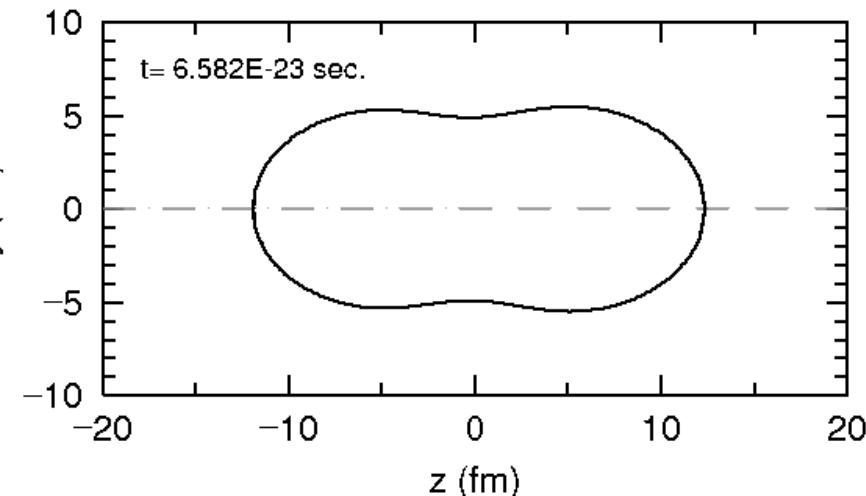
質量分布



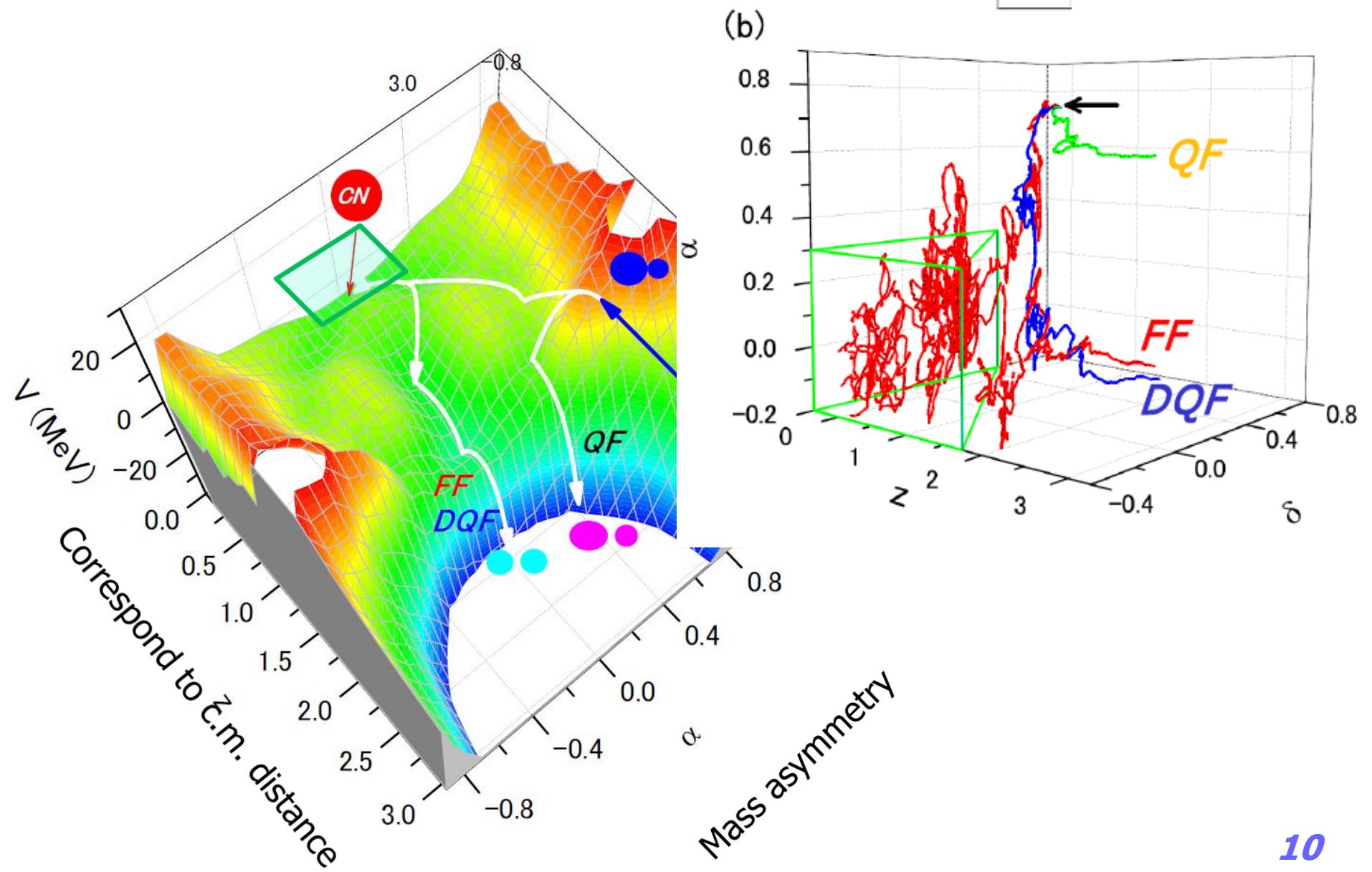
分裂過程

質量分布

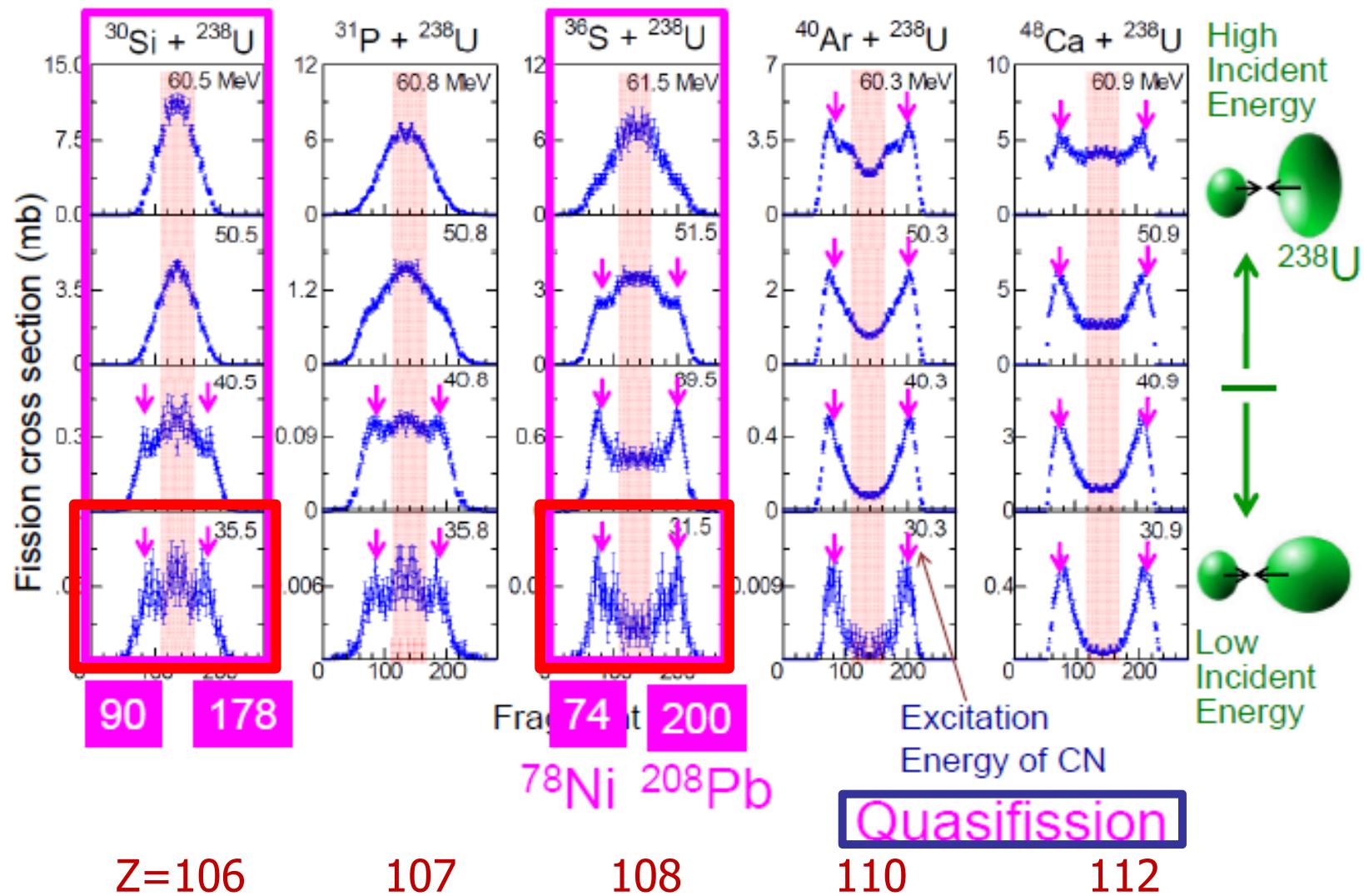
揺動項あり Langevin Eq.



Overview of Dynamical Process in reaction $^{36}\text{S} + ^{238}\text{U}$



Projectile dependence of fragment mass distributions



Z=106

107

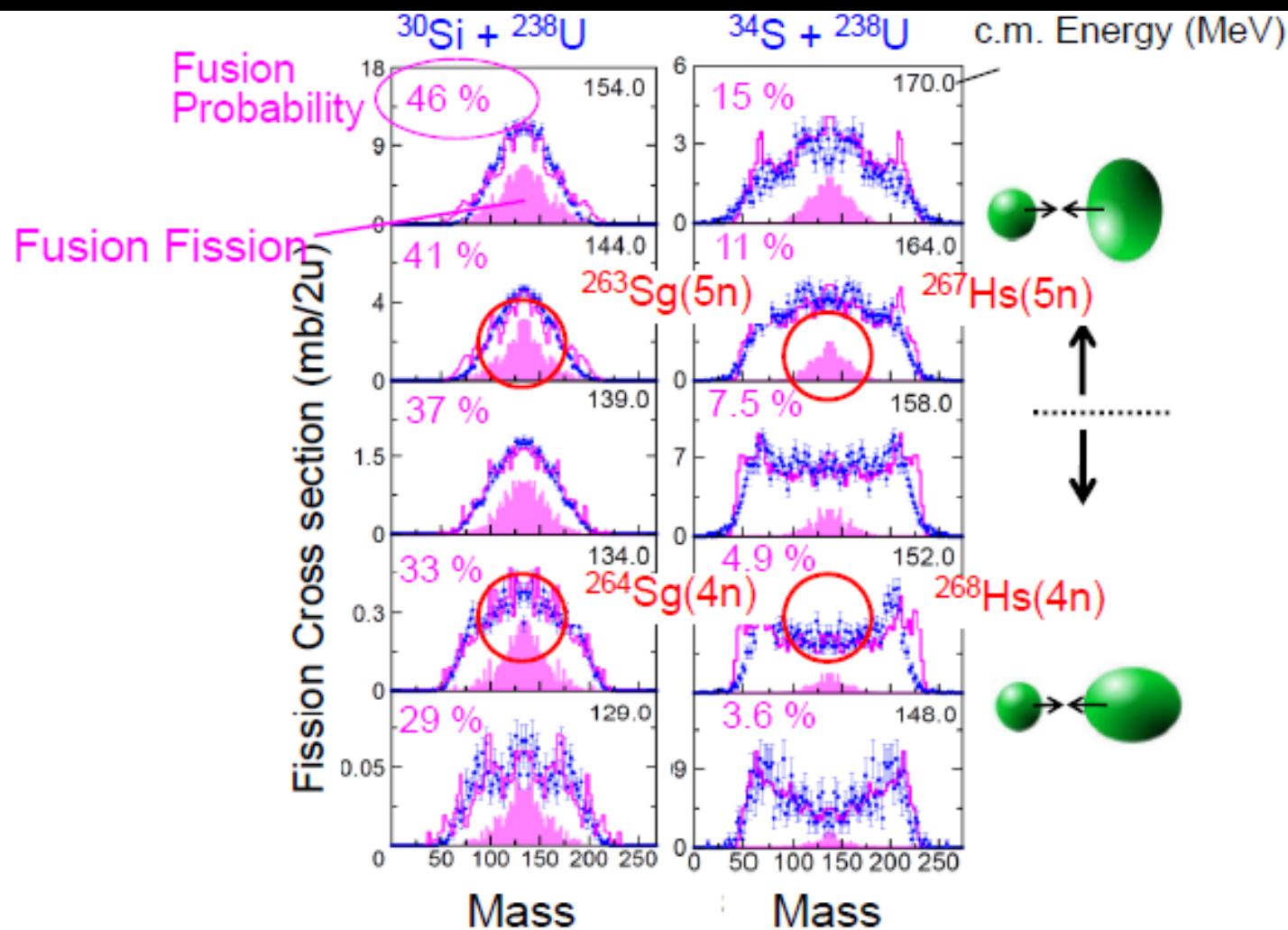
108

110

112

Experiments by K. Nishio et al. (JAEA)

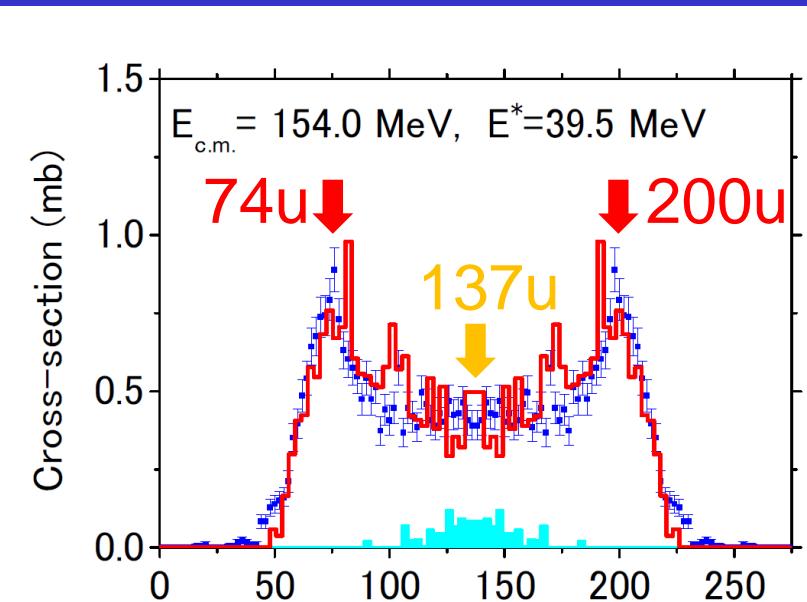
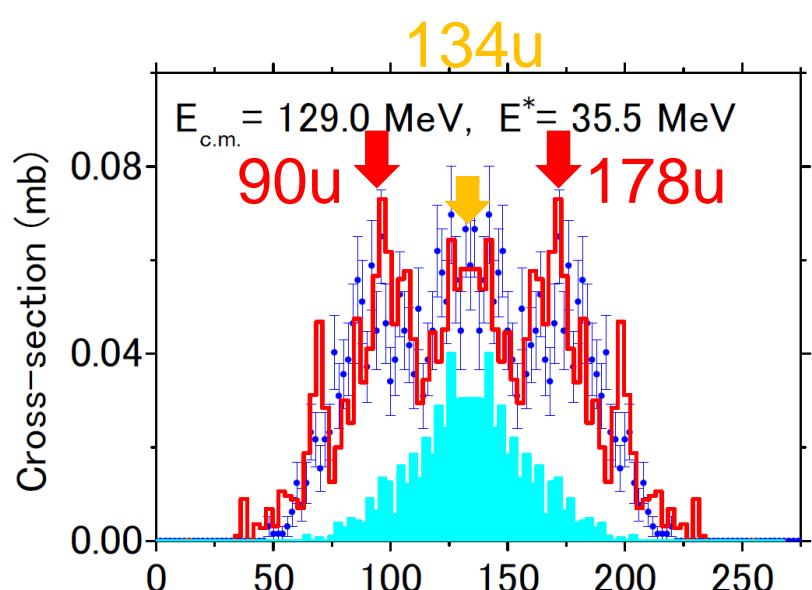
Calculated spectra for fusion-fission and quasi-fission



Experiments by K. Nishio et al. (JAEA)

4. Mechanism of Dynamical process

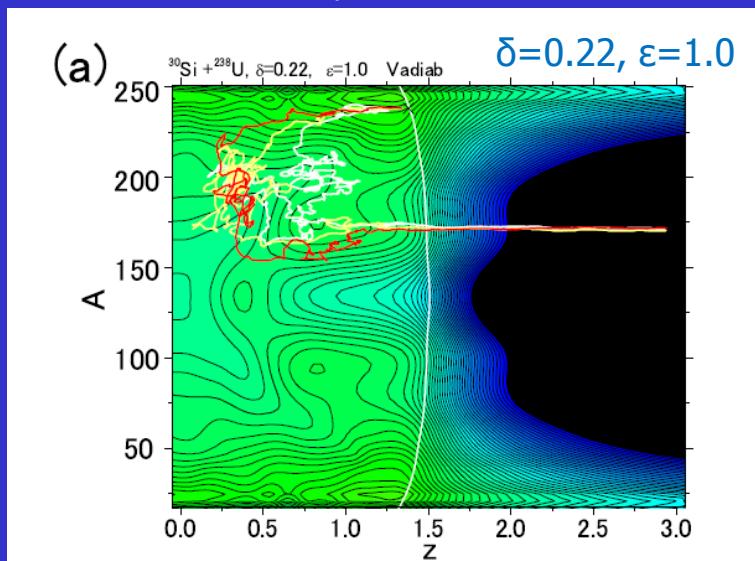
MDFF at Low incident energy



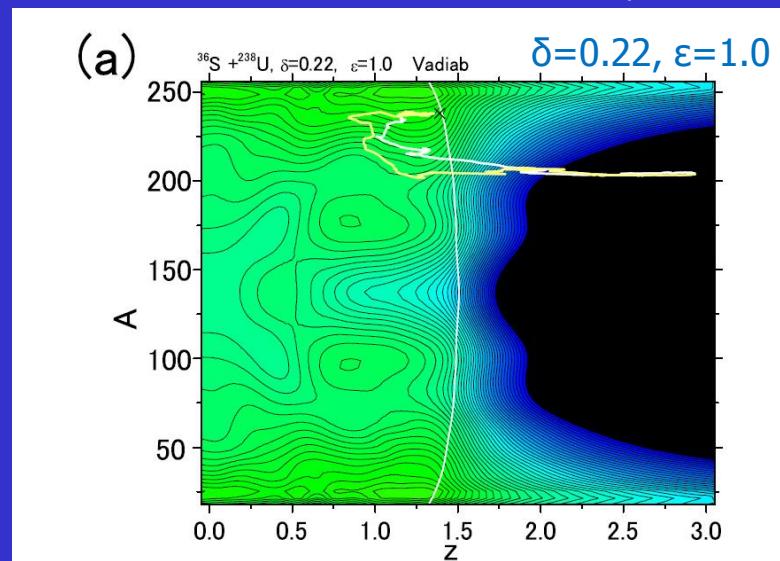
Clarify the origin of the difference

(b) Trajectory Analysis on Potential Energy Surface z-A plane

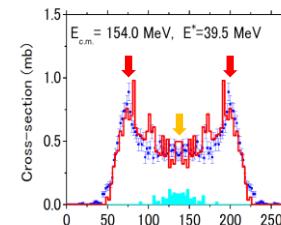
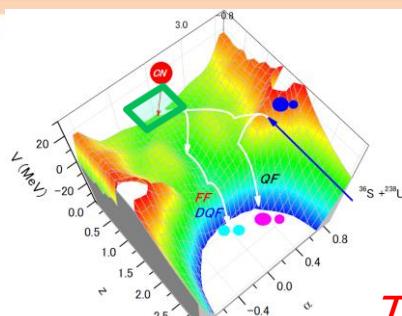
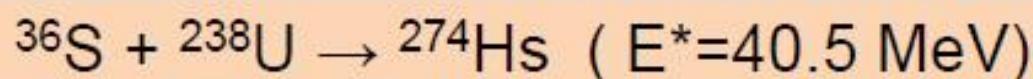
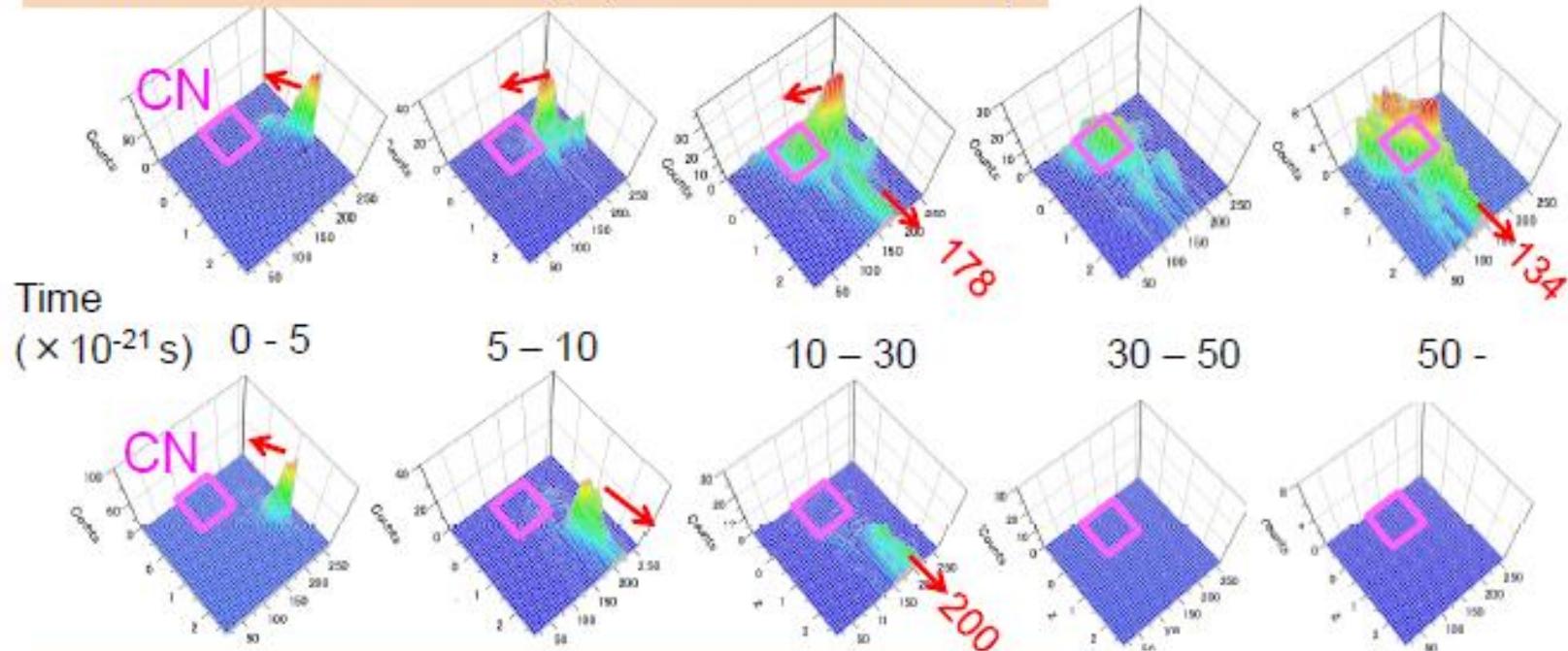
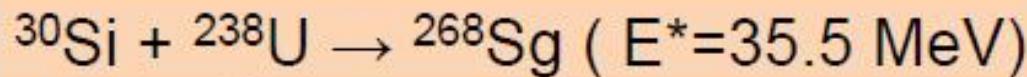
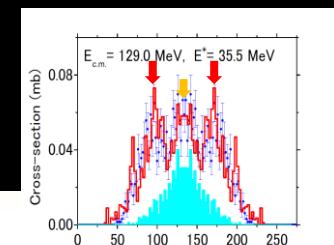
$^{30}\text{Si} + ^{238}\text{U}$ $E^* = 35.5 \text{ MeV}$
 $L=0, \theta=0$



$^{36}\text{S} + ^{238}\text{U}$ $E^* = 39.5 \text{ MeV}$
 $L=0, \theta=0$



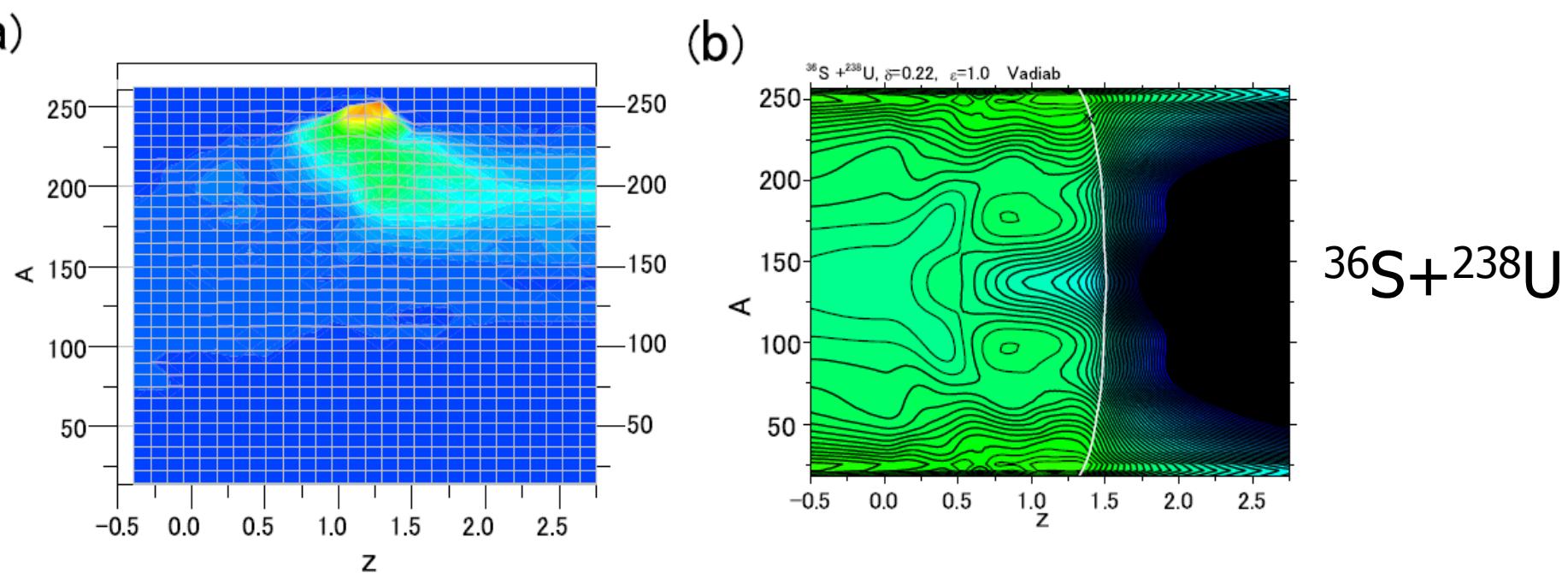
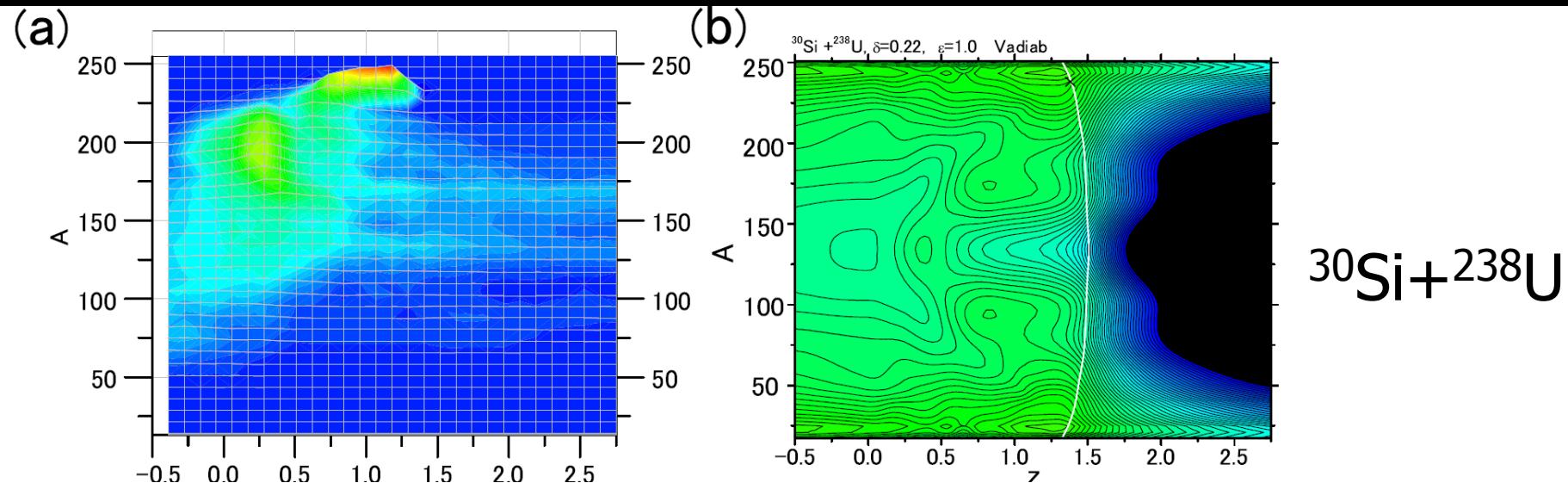
Time evolution of probability distribution



Mass Asymmetry: α
Charge Center Distance: Z

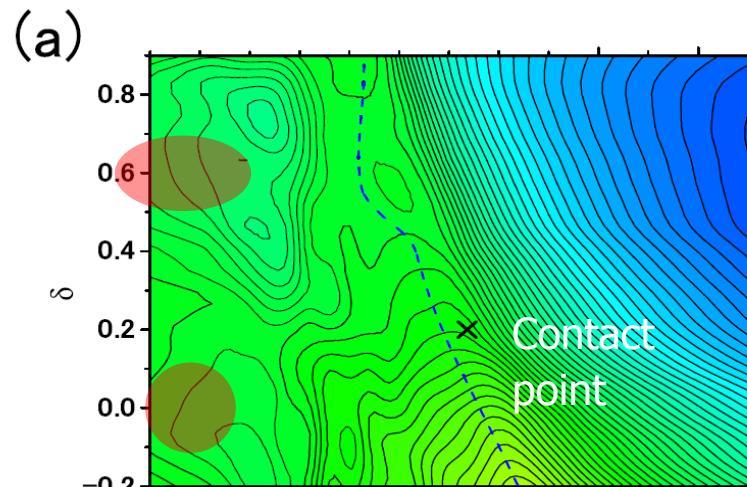
Try to clarify the origin of difference between the both cases →

Probability distribution on the z-A plane

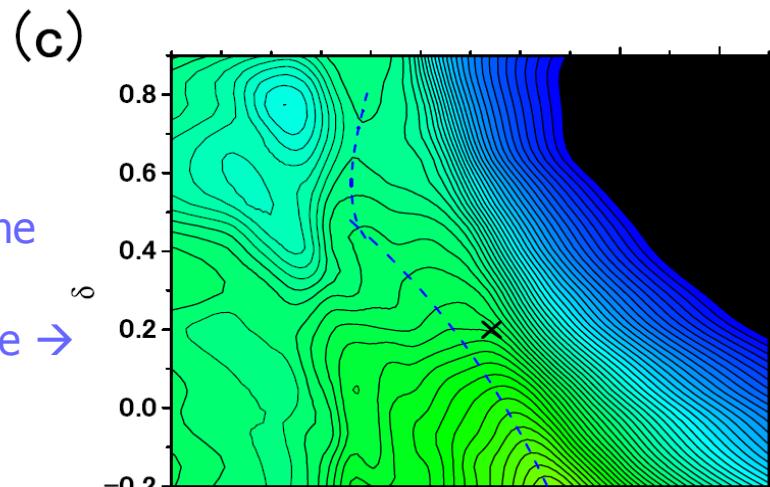


Probability distribution of total time on the z- δ plane

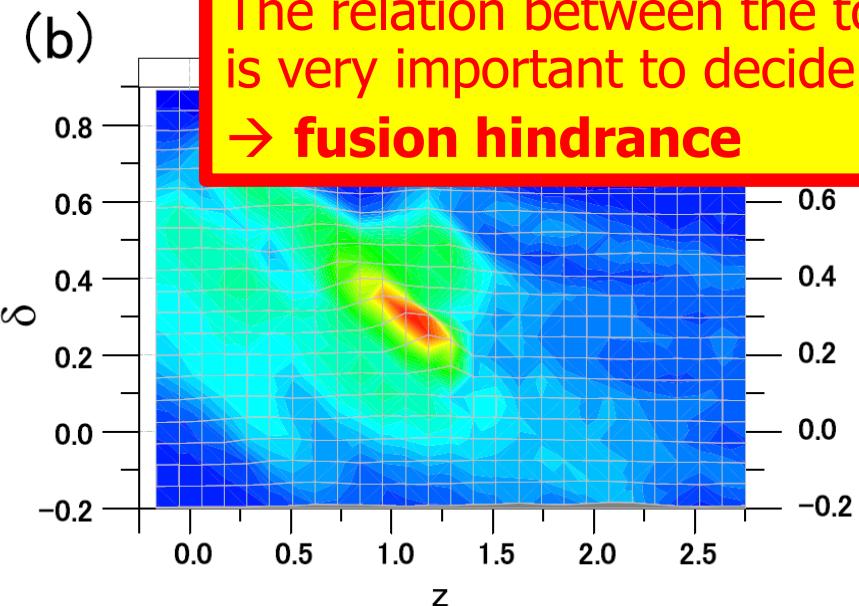
$^{30}\text{Si} + ^{238}\text{U}$, $E^* = 35.5 \text{ MeV}$



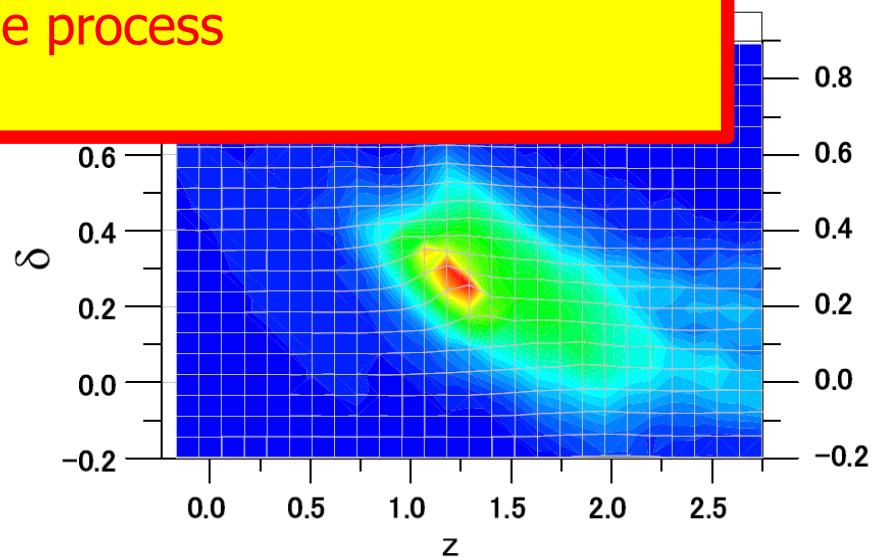
$^{36}\text{S} + ^{238}\text{U}$, $E^* = 39.5 \text{ MeV}$

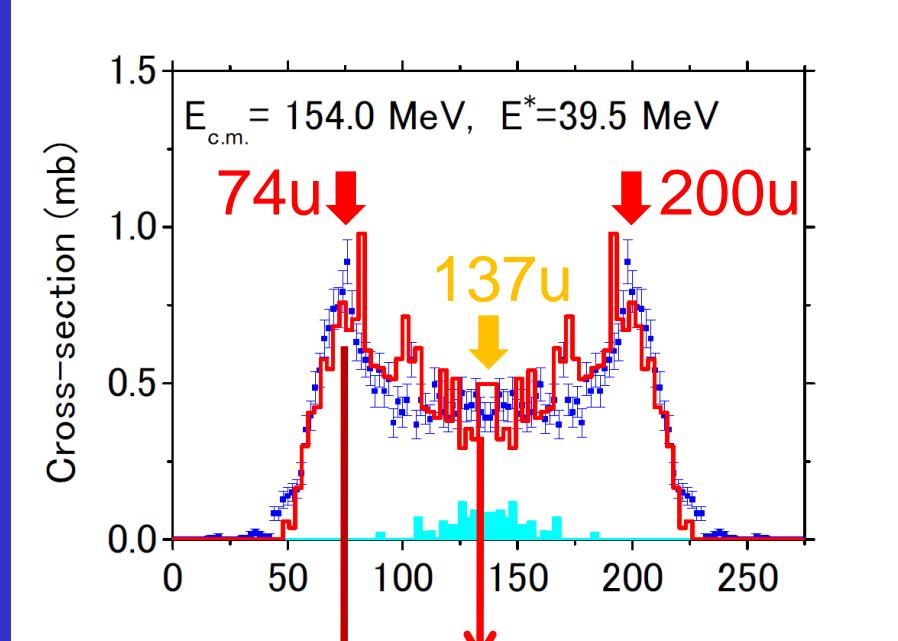
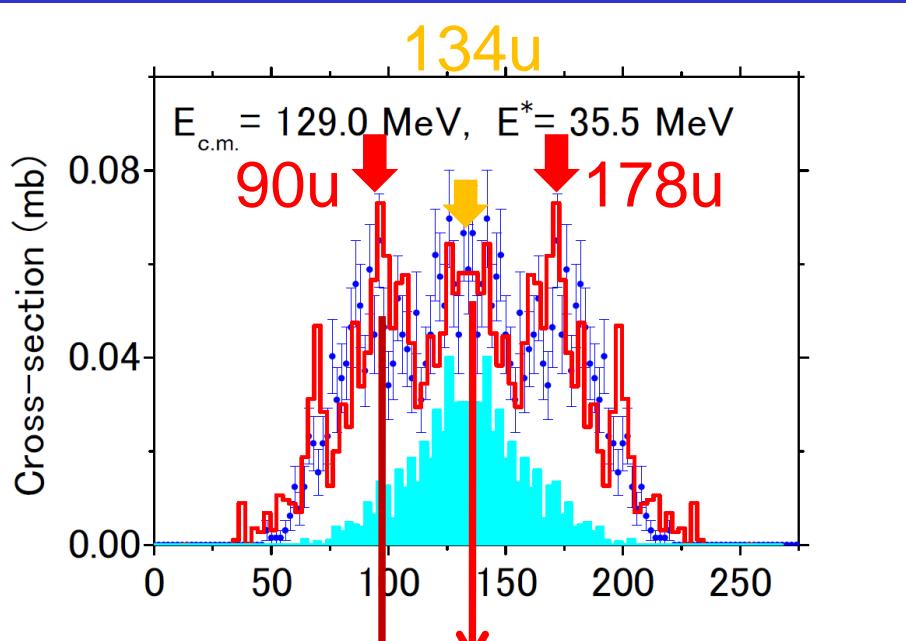


Ridge line
←steep
gentle →



The relation between the touching point and the ridge line
is very important to decide the process
→ **fusion hindrance**





FF and DQF
 $t > 50 \times 10^{-21} \text{ sec}$
 $-0.2 < \delta < 0.2$ (peak 0)

QF via mono-nucleus
 $t < 30 \times 10^{-21} \text{ sec}$
 $0.2 < \delta < 0.5$ (peak 0.4)

FF and DQF
 $t < 30 \times 10^{-21} \text{ sec}$
 $0 < \delta < 0.4$ (peak 0.2)

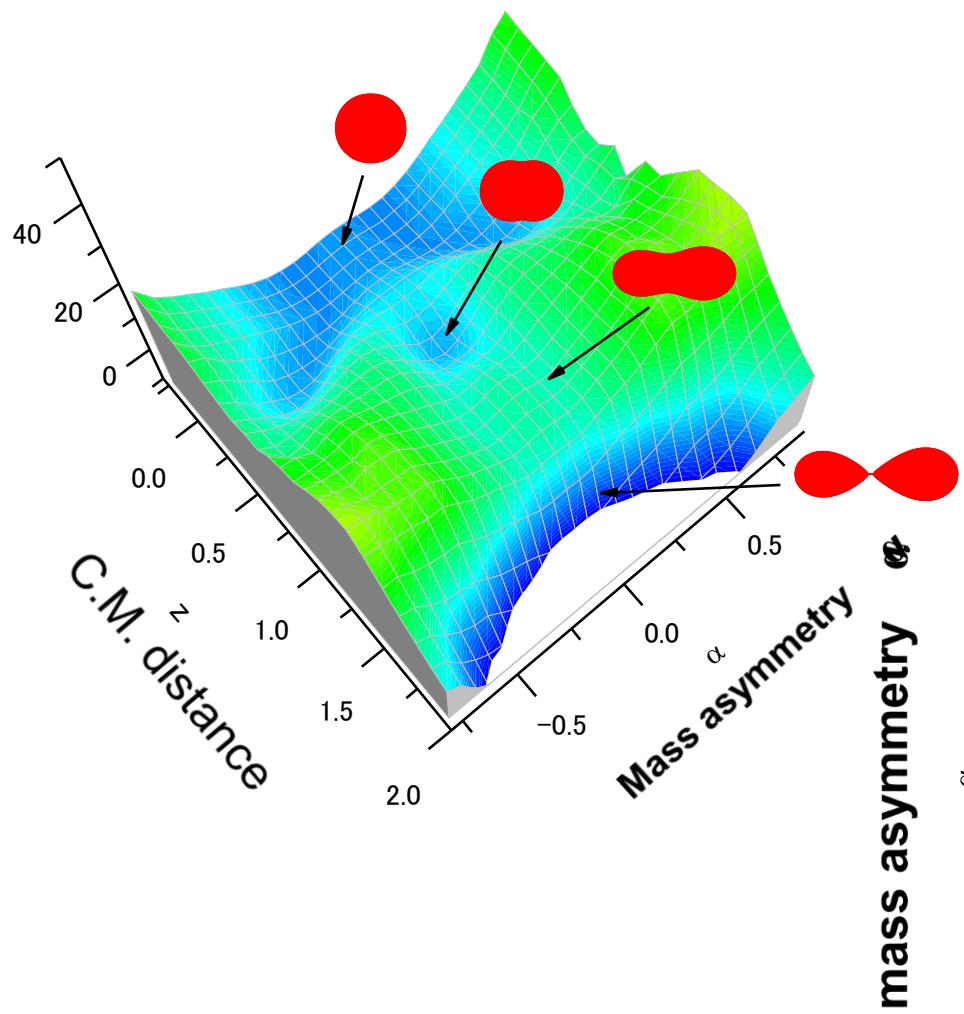
QF
 $t < 10 \times 10^{-21} \text{ sec}$
 $0 < \delta < 0.2$ (peak 0)

- (1) Origin of the reaction process
- (2) Building times
- (3) Deformation of fragments

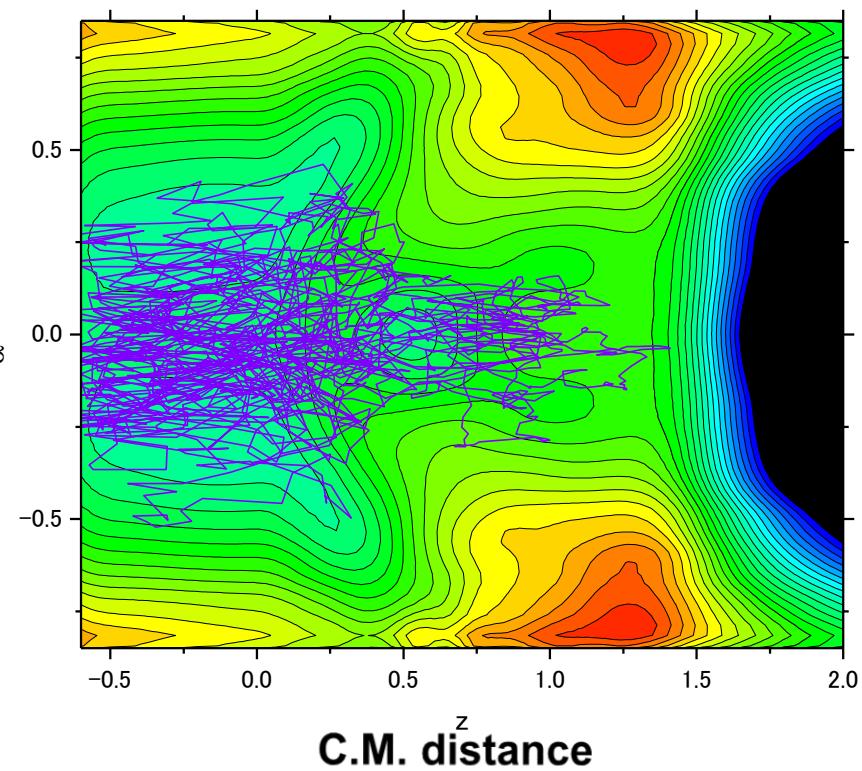


Fission Process

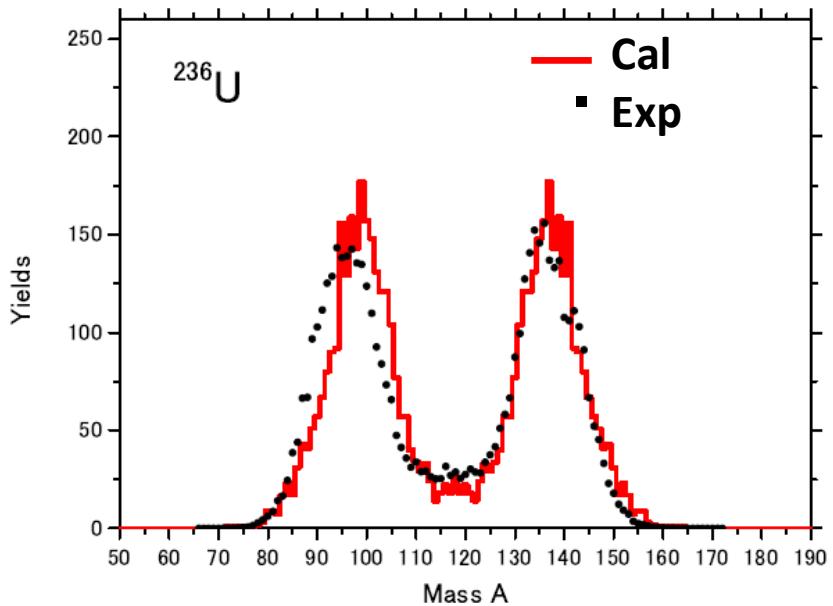
Fission process ^{240}U $E^* < 20$ MeV



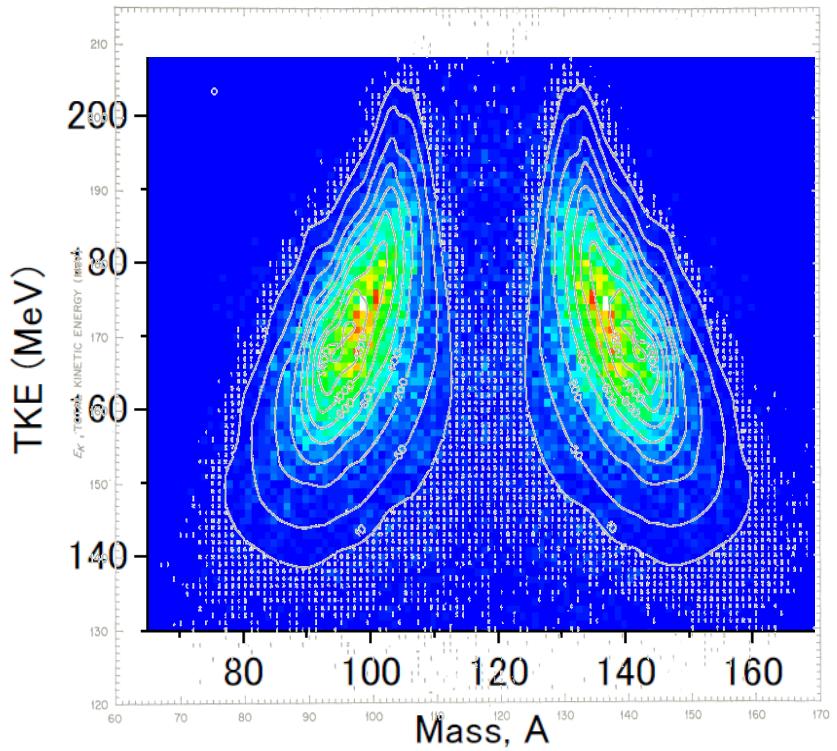
Trajectory on potential energy surface



$^{236}\text{U}^*$ ($E^* = 20\text{MeV}$)



Experiment J.Katakura, JENDL FP Decay Data File 2011
and Fission Yields

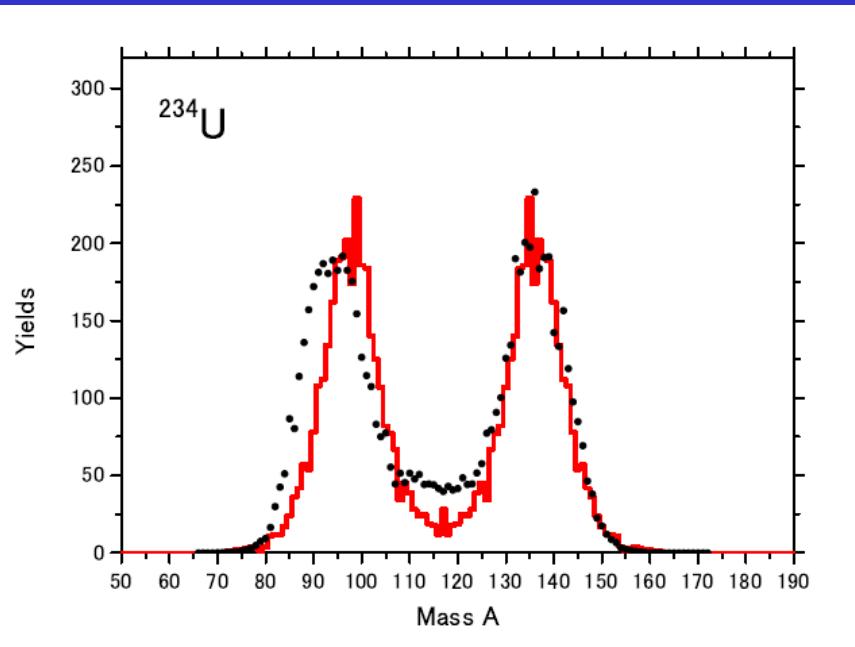


calculation Y. Aritomo and S. Chiba, PRC 88, 044614(2013)

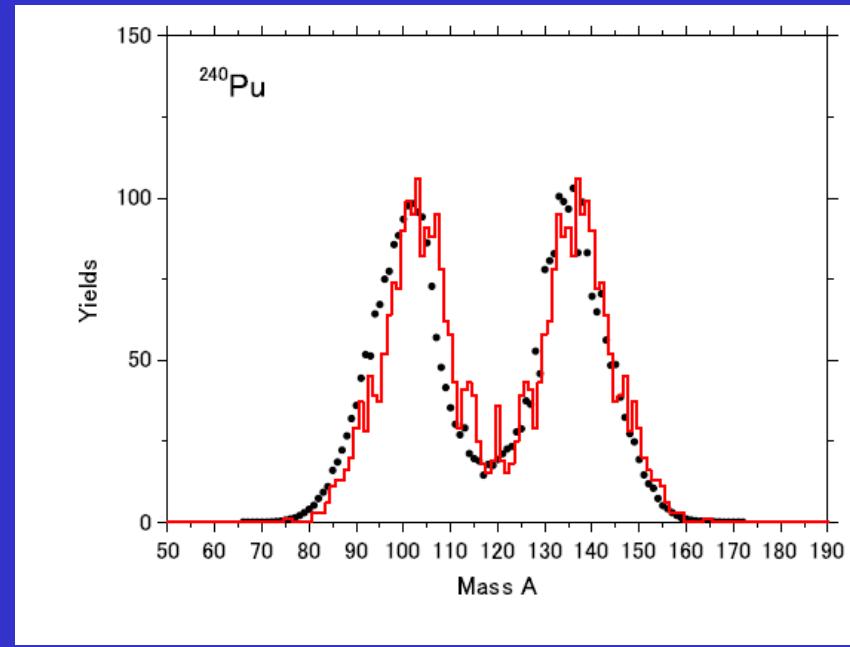
Exp. Phys.Rev. 141(1966)1146

Mass distribution of fission fragments $E^* = 20$ MeV

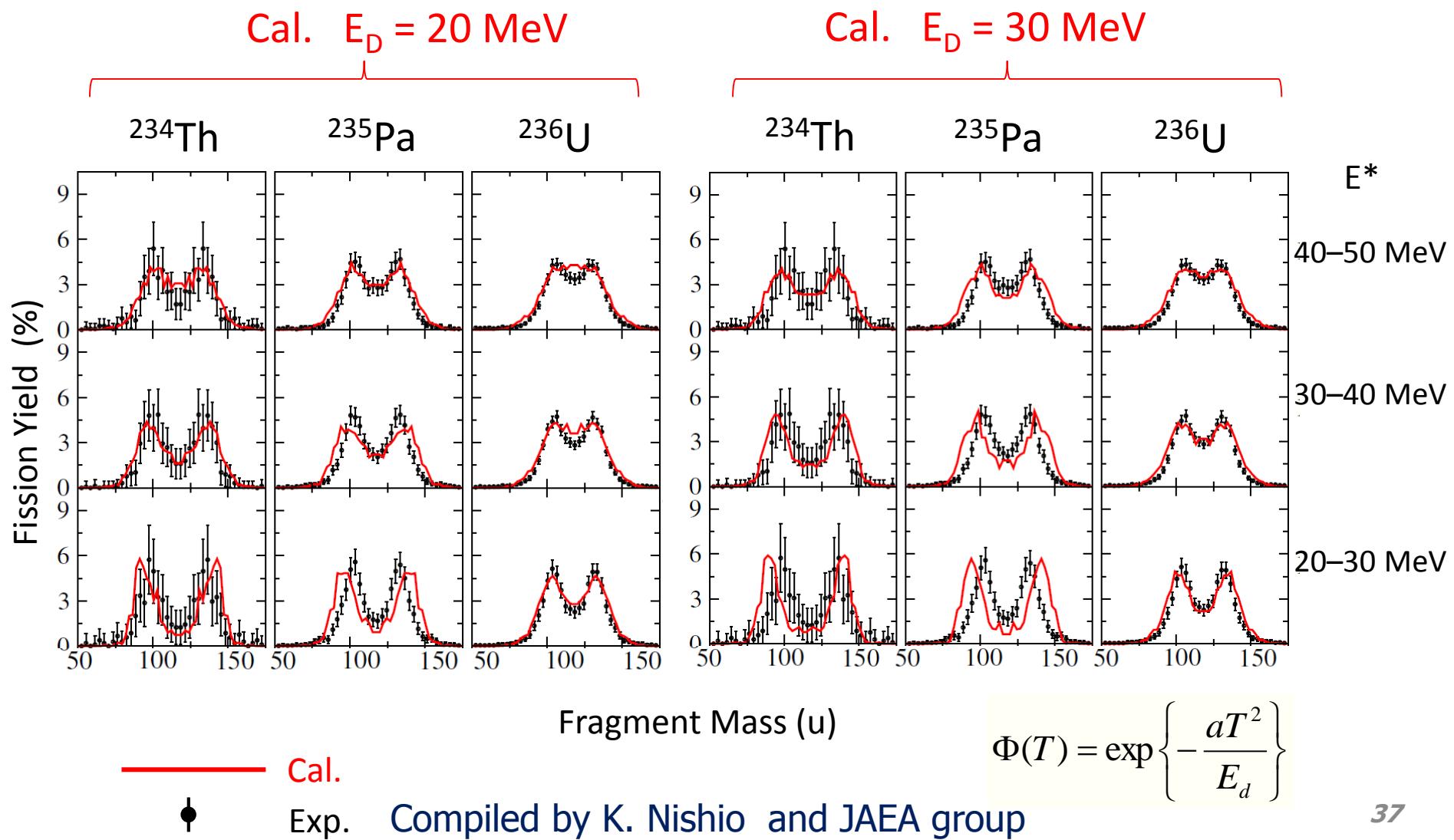
^{234}U



^{240}Pu



Comparison between Cal. and Exp.



5. Summary

1. In order to analyze the fusion-fission process in superheavy mass region, we apply the Couple channels method + Langevin calculation.
2. Incident energy dependence of mass distribution of fission fragments (MDFF) is reproduced in reaction $^{36}\text{S}+^{238}\text{U}$ and $^{30}\text{Si}+^{238}\text{U}$.
3. The shape of the MDFF is analyzed using
probability distribution
4. The relation between the touching point and the ridge line is very important to decide the process → fusion hindrance

And....

Collaborators

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S. Chiba

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V.I. Zagrebaev, A.V. Karpov

Flerov Laboratory of Nuclear Reactions



W. Greiner

Frankfurt Institute for Advanced Studies, J.W. Goethe University

