

Light nuclei and nucleon form factors  
in  $N_f = 2 + 1$  lattice QCD

Takeshi Yamazaki



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for PACS Collaboration

## 1. Light nuclei

in collaboration with

K.-I. Ishikawa, Y. Kuramashi, and A. Ukawa for PACS Collaboration

Refs: PRD81:111504(R)(2010); PRD84:054506(2011); PRD86:074514(2012)

PRD92:014501(2015)

## 2. Nucleon form factors

in collaboration with

K.-I. Ishikawa, Y. Kuramashi, S. Sasaki and A. Ukawa

for PACS Collaboration

# Outline

- Introduction
- Calculation method of nuclei in lattice QCD
- Simulation parameters
- Results of light nuclei
  - $^4\text{He}$  and  $^3\text{He}$  channels
  - NN channels
- Preliminary result at  $m_\pi \sim 0.145$  GeV
  - Light nuclei binding energy
  - nucleon form factors
- Summary and future work

# Introduction

Binding force  $\left\{ \begin{array}{l} \text{protons and neutrons} \rightarrow \text{nuclei} \\ \text{quarks and gluons} \rightarrow \text{protons and neutrons} \end{array} \right.$

both from fundamental strong interaction of quark and gluon  
well known, but hard to prove

quark and gluon  $\rightarrow$  proton and neutron  $\rightarrow$  nucleus

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Spectrum of proton and neutron (nucleons)

success of non-perturbative calculation of QCD

degrees of freedom of quarks and gluons

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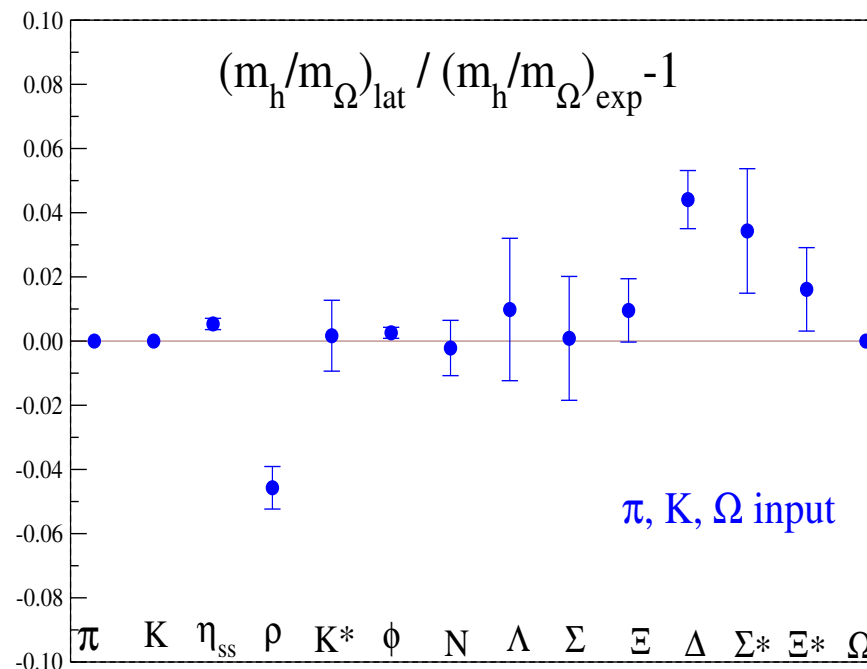
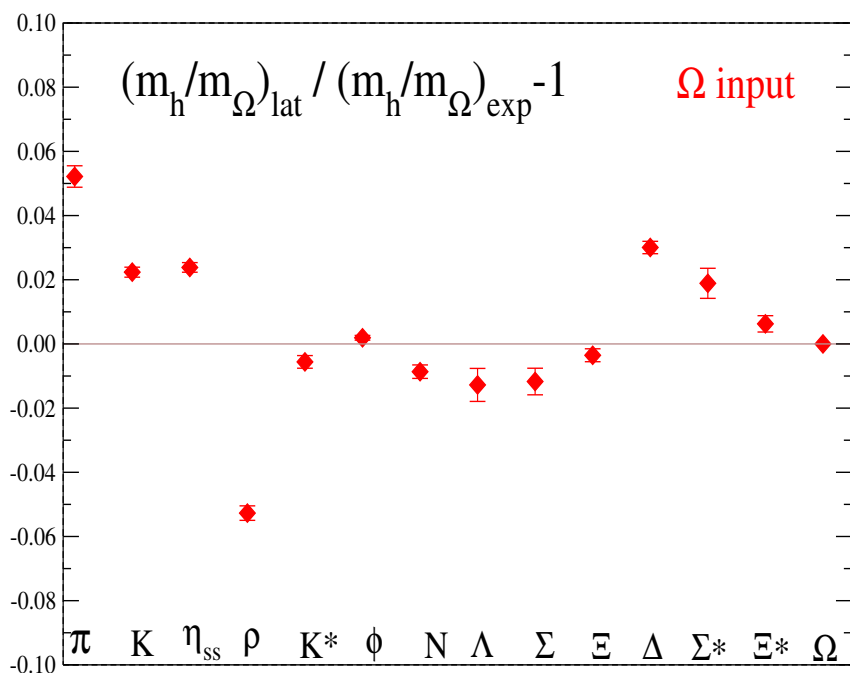
# Hadron spectrum in $N_f = 2 + 1$ QCD

Lattice 2015, Ukita for PACS Collaboration

$m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm (K computer, SPIRE Field 5)  
using reweighting  $m_{u,d}, m_s$  + extrapolation  $\rightarrow$  physical  $m_\pi$  and  $m_K$

$m_\pi \sim 0.145$  GeV

physical point



Stable hadron mass: well reproduced from lattice QCD

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$\underbrace{\text{quark and gluon} \rightarrow \text{proton and neutron}}_{\rightarrow \text{nucleus}}$

goal: quantitatively understand property of nucleus from QCD

So far not many studies for multi-baryon bound states

→ Can we reproduce binding energy of known light nuclei?



# Introduction

Binding force  $\left\{ \begin{array}{l} \text{protons and neutrons} \rightarrow \text{nuclei} \\ \text{quarks and gluons} \rightarrow \text{protons and neutrons} \end{array} \right.$

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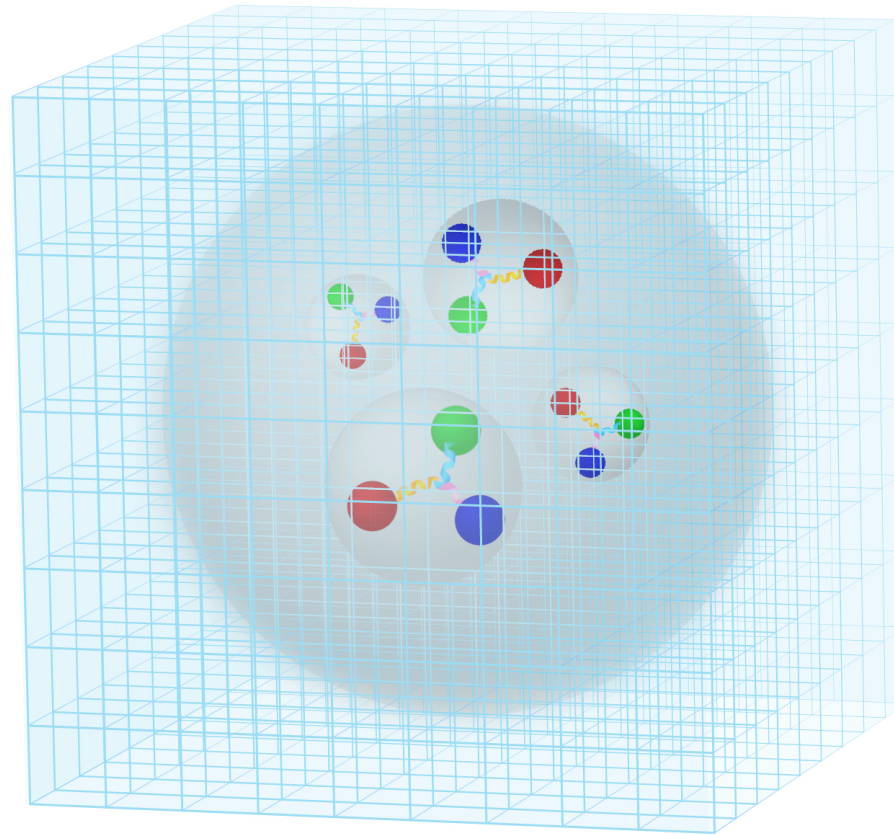
2nd motivation: Nucleon form factors not well understood

→ 2nd part of talk

$\underbrace{\text{quark and gluon} \rightarrow \text{proton and neutron}}_{\text{quark and gluon} \rightarrow \text{proton and neutron}} \rightarrow \text{nucleus}$

goal: quantitatively understand property of nucleus from QCD

# Ultimate goal of lattice QCD



<http://www.jicfus.jp/jp/promotion/pr/mj/2014-1/>

quantitatively understand property of nuclei from QCD

# Multi-baryon system from lattice QCD at '09

## 1. $\Lambda\Lambda$ system (Quenched QCD)

'85 Mackenzie & Thacker    '00 Wetzorke *et al.*

'88 Iwasaki *et al.*                '02 Wetzorke & Karsch

'99 Pochinsky *et al.*            '09 NPLQCD ( $N_f = 2 + 1$ )

**H dibaryon: unbound** except Iwasaki *et al.*

## 2. NN system $^3S_1$ and $^1S_0$

'95 Fukugita *et al.* : Quenched QCD

'06 NPLQCD :  $N_f = 2 + 1$  QCD

'08 Ishii *et al.* : Quenched and  $N_f = 2 + 1$  QCD

'09 NPLQCD :  $N_f = 2 + 1$  QCD

**Deuteron: unbound** due to  $m_\pi \gtrsim 0.3$  GeV

## 3. NNN system

'09 NPLQCD :  $N_f = 2 + 1$  QCD

**Triton: likely unbound**

# Multi-baryon bound state from lattice QCD

Not observed before '09 (except H-dibaryon '88 Iwasaki *et al.*)

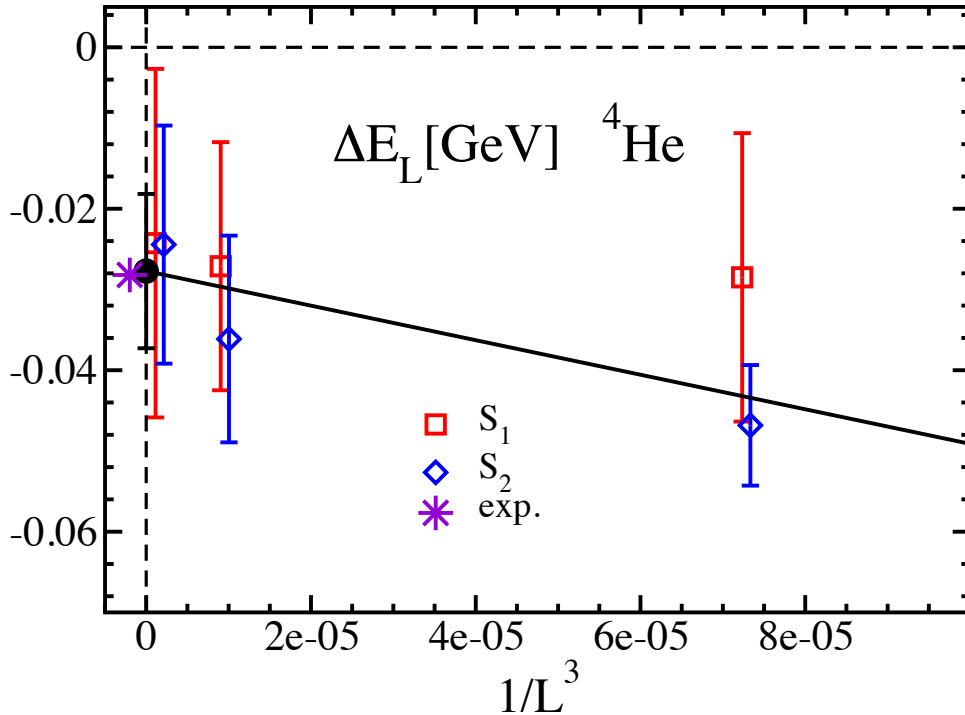
1.  ${}^4\text{He}$  and  ${}^3\text{He}$

'10 PACS-CS  $N_f = 0$   $m_\pi = 0.8$  GeV PRD81:111504(R)(2010)

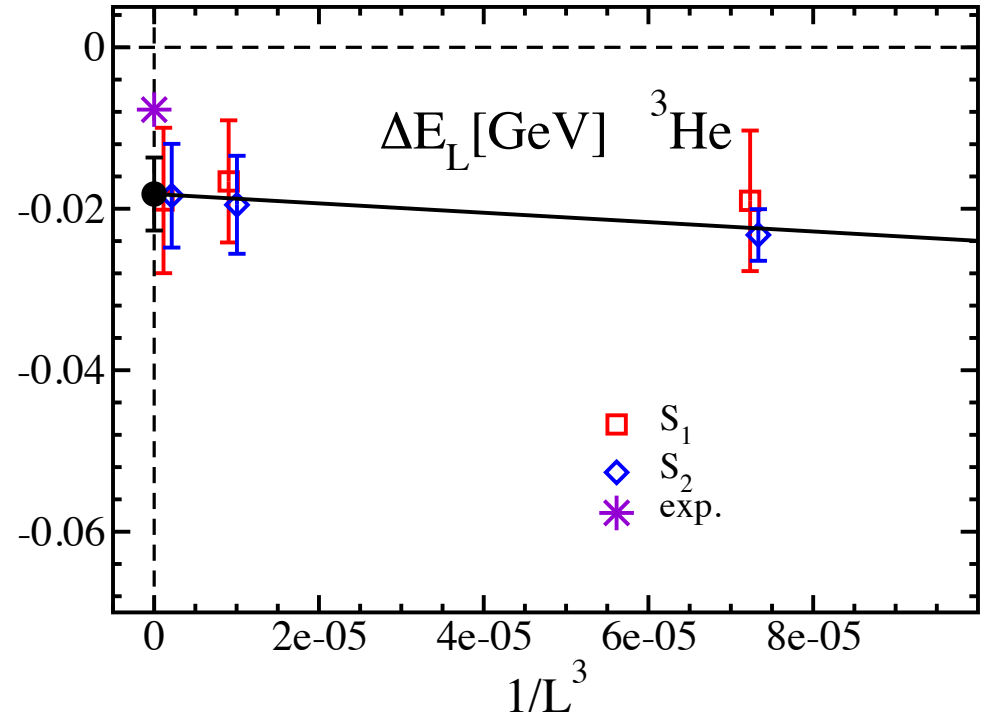
# Exploratory study of three- and four-nucleon systems

PACS-CS Collaboration, PRD81:111504(R)(2010)

## Identification of bound state from volume dependence of $\Delta E$



$$\Delta E_{4\text{He}} = 27.7(7.8)(5.5) \text{ MeV}$$



$$\Delta E_{3\text{He}} = 18.2(3.5)(2.9) \text{ MeV}$$

1. Observe bound state in both channels
2. Same order of  $\Delta E$  to experiment

Several systematic errors included, e.g.,  $N_f = 0$ ,  $m_\pi = 0.8$  GeV

# Multi-baryon bound state from lattice QCD

## 1. ${}^4\text{He}$ and ${}^3\text{He}$

'10 PACS-CS  $N_f = 0$   $m_\pi = 0.8$  GeV PRD81:111504(R)(2010)

'12 HALQCD  $N_f = 3$   $m_\pi = 0.47$  GeV,  $m_\pi > 1$  GeV  ${}^4\text{He}$

'12 NPLQCD  $N_f = 3$   $m_\pi = 0.81$  GeV

'12 TY *et al.*  $N_f = 2 + 1$   $m_\pi = 0.51$  GeV PRD86:074514(2012)

'15 TY *et al.*  $N_f = 2 + 1$   $m_\pi = 0.30$  GeV PRD92:014501(2015)

## 2. H dibaryon in $\Lambda\Lambda$ channel ( $S=-2$ , $I=0$ )

'11, '12 NPLQCD  $N_f = 2 + 1$   $m_\pi = 0.39$  GeV,  $N_f = 3$   $m_\pi = 0.81$  GeV

'11, '12 HALQCD  $N_f = 3$   $m_\pi = 0.47-1.02$  GeV

'11 Luo *et al.*  $N_f = 0$   $m_\pi = 0.5-1.3$  GeV

'14 Mainz  $N_f = 2$   $m_\pi = 0.45, 1.0$  GeV

## 3. NN

'11 PACS-CS  $N_f = 0$   $m_\pi = 0.8$  GeV PRD84:054506(2011)

'12 NPLQCD  $N_f = 2 + 1$   $m_\pi = 0.39$  GeV (Possibility)

'12 NPLQCD, '15 CalLat  $N_f = 3$   $m_\pi = 0.81$  GeV

'12 TY *et al.*  $N_f = 2 + 1$   $m_\pi = 0.51$  GeV PRD86:074514(2012)

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'15 NPLQCD  $N_f = 2 + 1$   $m_\pi = 0.45$  GeV

Other states:  $\Xi\Xi$ , '12 NPLQCD; spin-2  $N\Omega$ ,  ${}^{16}\text{O}$  and  ${}^{40}\text{Ca}$ , '14 HALQCD, ...

Calculation method

# Calculation method of multi-nucleon bound state

Traditional method for example  ${}^4\text{He}$  channel

$$\langle 0 | O_{4\text{He}}(t) O_{4\text{He}}^\dagger(0) | 0 \rangle = \sum_n \langle 0 | O_{4\text{He}} | n \rangle \langle n | O_{4\text{He}}^\dagger | 0 \rangle e^{-E_n t} \xrightarrow{t \gg 1} A_0 e^{-E_0 t}$$

## Difficulties for multi-nucleon calculation

### 1. Statistical error

$$\text{Statistical error} \propto \exp\left(N_N \left[m_N - \frac{3}{2}m_\pi\right] t\right)$$

### 2. Calculation cost

$$\begin{aligned} \text{Wick contraction for } {}^4\text{He} &= p^2 n^2 = (udu)^2 (dud)^2: 518400 \\ \text{proton} &= p = (udu): 2 \end{aligned}$$

### 3. Identification of bound state on finite volume

Finite volume effect of attractive scattering state

$$\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0 \leftrightarrow \text{binding energy}$$



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Most severe problem before '09: (every  $t$ )  $\times N_{\text{meas}} \sim O(10^6)$

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→ heavy quark  $m_\pi = 0.8-0.3$  GeV + large # of measurements

### 2. Calculation cost PACS-CS PRD81:111504(R)(2010)

Wick contraction for  ${}^4\text{He} = p^2 n^2 = (udu)^2 (dud)^2: 518400 \rightarrow 1107$

→ reduction using  $p(n) \leftrightarrow p(n)$   $p \leftrightarrow n$ ,  $u(d) \leftrightarrow u(d)$  in  $p(n)$

+ block of 3 quark props(parallel) and contraction(workstation)

'12 Doi and Endres; Detmold and Orginos; '13 Günther et al.; '15 Nemura

### 3. Identification of bound state on finite volume

attractive scattering state  $\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0$

'86,'91 Lüscher, '07 Beane *et al.*

→ Volume dependence of  $\Delta E_L \rightarrow \Delta E_\infty \neq 0 \rightarrow$  bound state

Spectral weight: '04 Mathur *et al.*, Anti-PBC '05 Ishii *et al.*

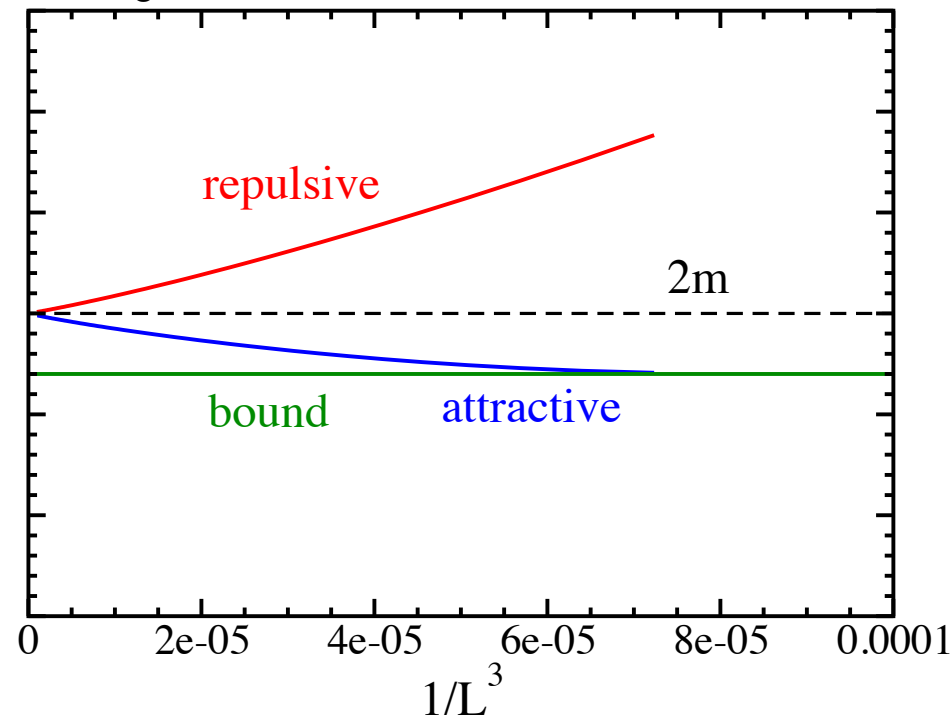
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Ground state energy  $E_0$  in two particle system on finite volume



Hard to distinguish attractive scattering from bound  
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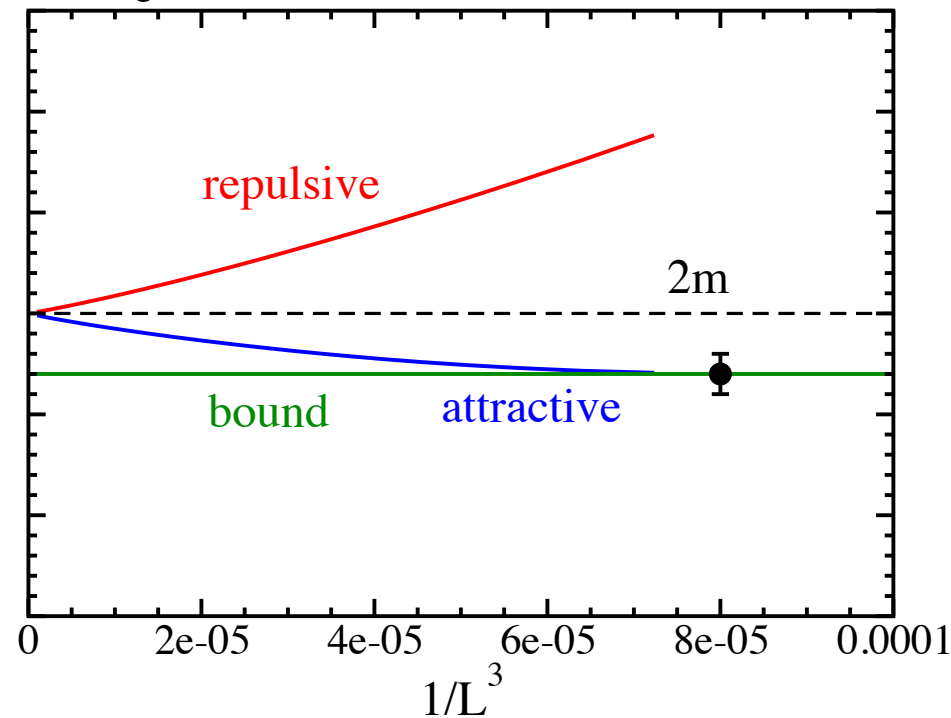
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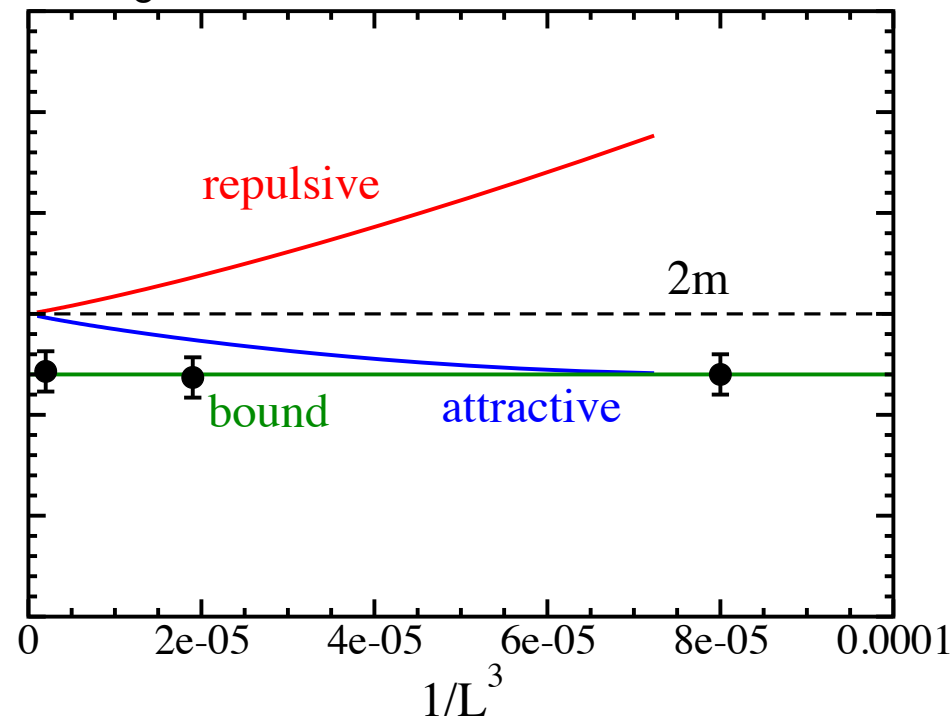
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Simplest way: extrapolation of  $E_0(\Delta E_L)$  to infinite volume limit

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Most severe problem at present

### 2. Calculation cost PACS-CS PRD81:111504(R)(2010)

Wick contraction for  ${}^4\text{He} = p^2 n^2 = (udu)^2 (dud)^2$ : 518400

proton =  $p = (udu)$ : 2

Used to be most severe problem

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Finite volume effect of attractive scattering state

$$\Delta E_L = E_0 - N_N m_N = O(L^{-3}) < 0 \leftrightarrow \text{binding energy}$$

# Simulation parameters

$N_f = 2 + 1$  QCD

Iwasaki gauge action at  $\beta = 1.90$

$a^{-1} = 2.194$  GeV with  $m_\Omega = 1.6725$  GeV '10 PACS-CS

non-perturbative  $O(a)$ -improved Wilson fermion action

$m_\pi = 0.51$  GeV and  $m_N = 1.32$  GeV PRD86:074514(2012)

$m_\pi = 0.30$  GeV and  $m_N = 1.05$  GeV PRD92:014501(2015)

$m_s \sim$  physical strange quark mass

${}^4\text{He}$ ,  ${}^3\text{He}$ , NN( ${}^3\text{S}_1$  and  ${}^1\text{S}_0$ )

		$m_\pi = 0.5$ GeV		$m_\pi = 0.3$ GeV		$R$
$L$	$L$ [fm]	$N_{\text{conf}}$	$N_{\text{meas}}$	$N_{\text{conf}}$	$N_{\text{meas}}$	
32	2.9	200	192			
40	3.6	200	192			
48	4.3	200	192	400	1152	12
64	5.8	190	256	160	1536	5

$$R = (N_{\text{conf}} \cdot N_{\text{meas}})_{0.3\text{GeV}} / (N_{\text{conf}} \cdot N_{\text{meas}})_{0.5\text{GeV}}$$

Smear source and point sink ( $N$  with  $p = 0$ ) operators

→ after some tests in  $N_f = 0$ , consider large overlap to ground state



# Results at $m_\pi = 0.5$ and $0.3$ GeV

Computational resources

PACS-CS, T2K-Tsukuba, HA-PACS, COMA at Univ. of Tsukuba

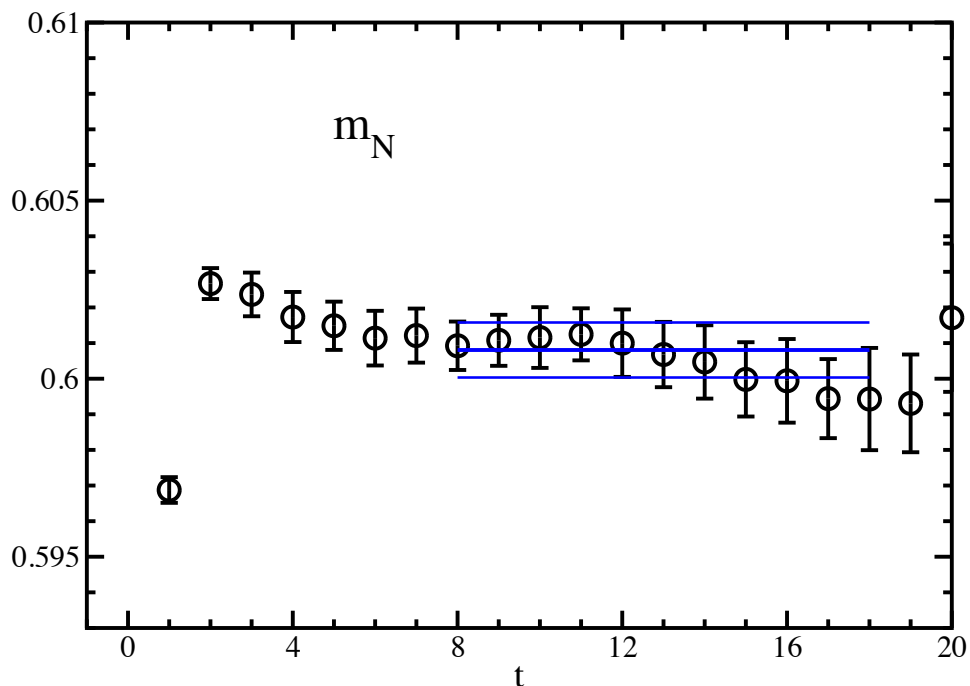
T2K-Tokyo and FX10 at Univ. of Tokyo, and K at AICS

# Results

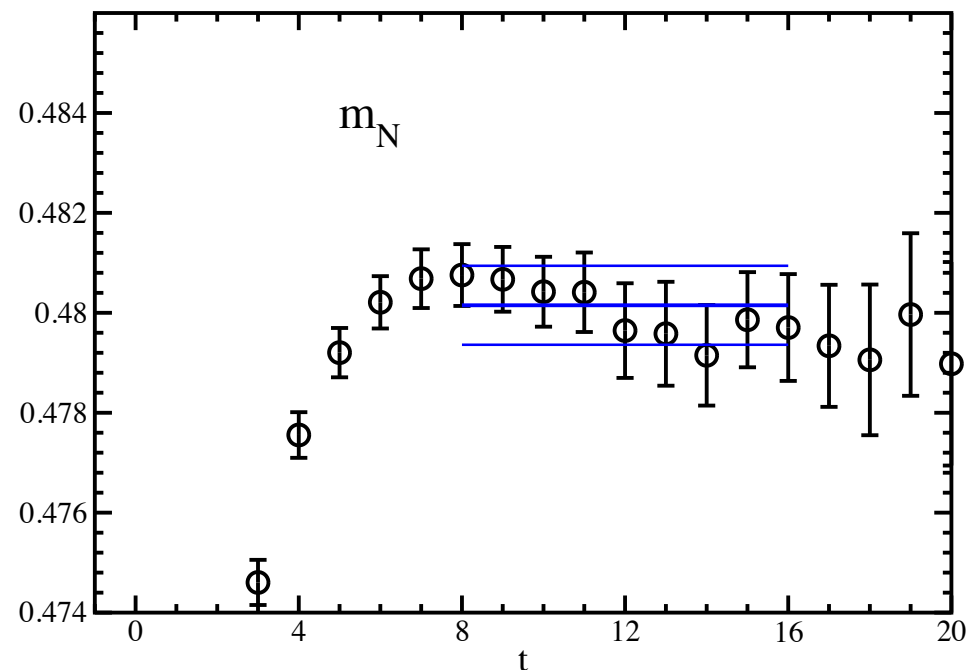
Effective mass of nucleon on  $L = 5.8$  fm

$$\text{Effective } m_N = \log \left( \frac{C_N(t)}{C_N(t+1)} \right)$$

$m_\pi = 0.5$  GeV



$m_\pi = 0.3$  GeV



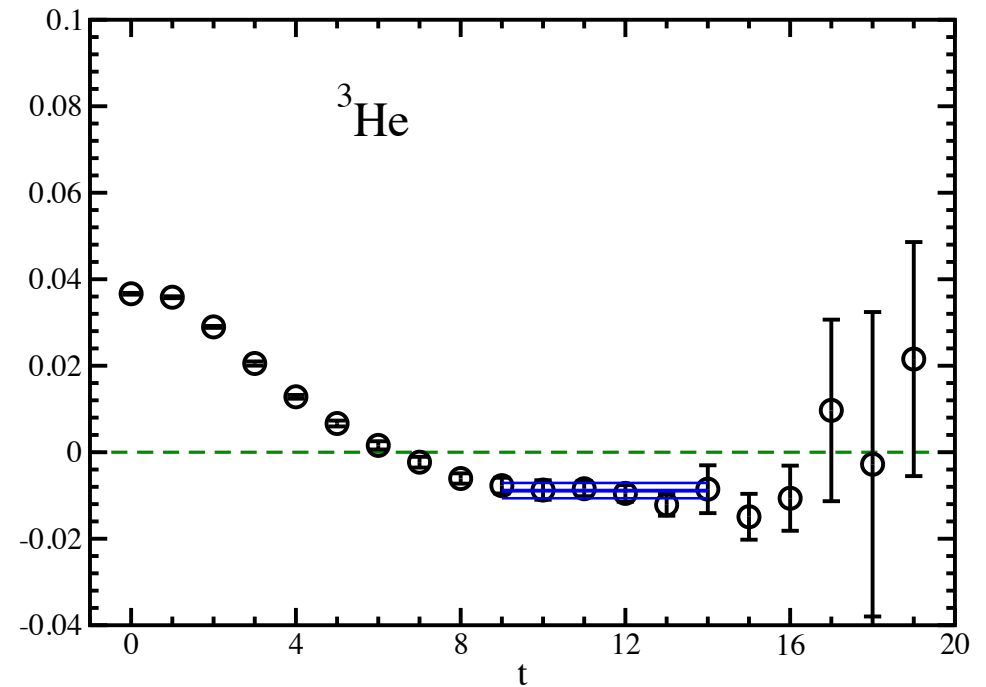
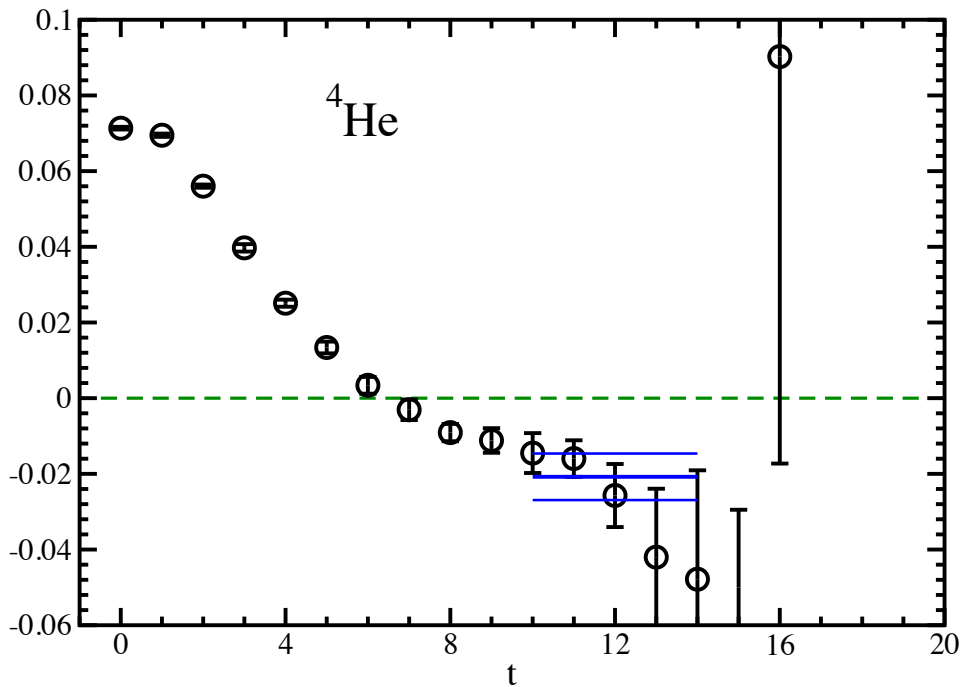
- Good plateau  $t \gtrsim 7$
- Statistical error  $< 0.2\%$

$\Delta E_L = E_0 - N_N m_N$  in  $^4\text{He}$  and  $^3\text{He}$  channels

at  $m_\pi = 0.5$  GeV on  $L = 5.8$  fm

TY *et al.*, PRD86:074514(2012)

$$\Delta E_L = \log \left( \frac{R_{4\text{He}}(t)}{R_{4\text{He}}(t+1)} \right) \text{ with } R_{4\text{He}}(t) = \frac{C_{4\text{He}}(t)}{(C_N(t))^4}$$

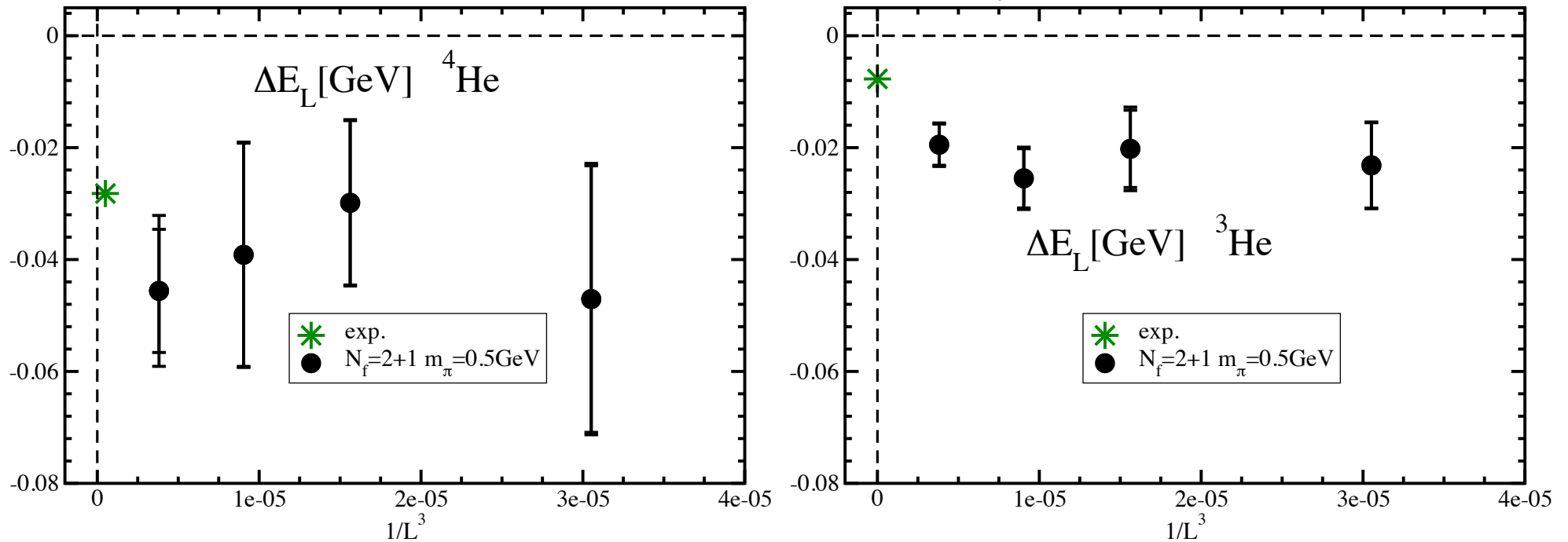


- Statistical error under control in  $t < 12$
- Negative  $\Delta E_L$  in both channels
- Plateau region  $\sim$  plateau region of  $m_N$

$^4\text{He}$  and  $^3\text{He}$  channels  $\Delta E_L = E_0 - N_N m_N$  at  $m_\pi = 0.5 \text{ GeV}$

TY *et al.*, PRD86:074514(2012)

Identification of bound state from volume dependence of  $\Delta E$

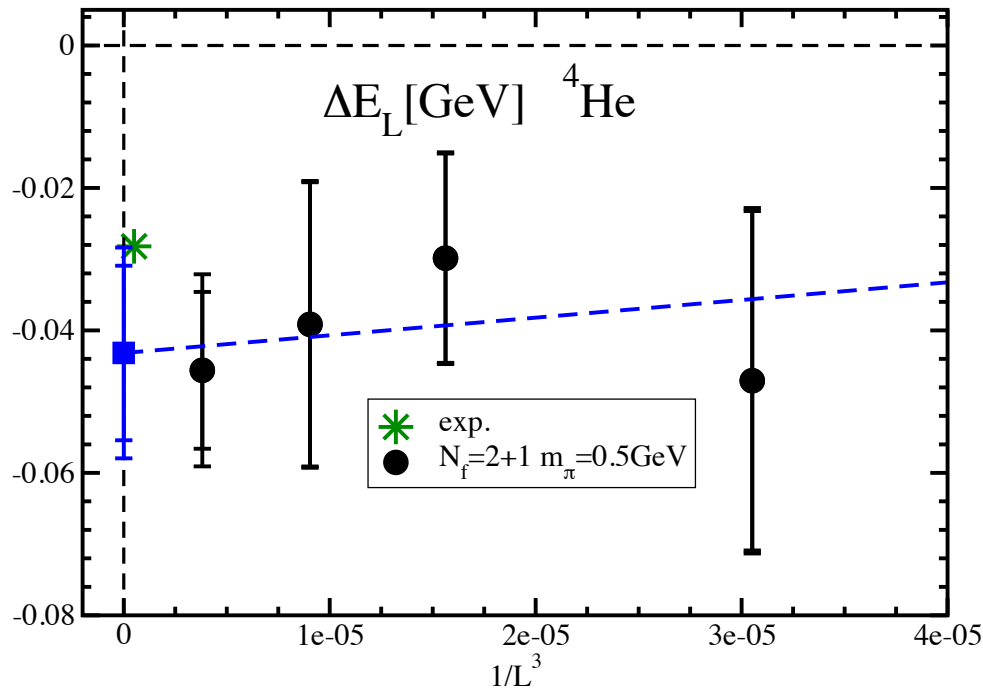


- $\Delta E_L < 0$  and mild volume dependence
- Infinite volume extrapolation with  $\Delta E_L = -\Delta E_{\text{bind}} + C/L^3$

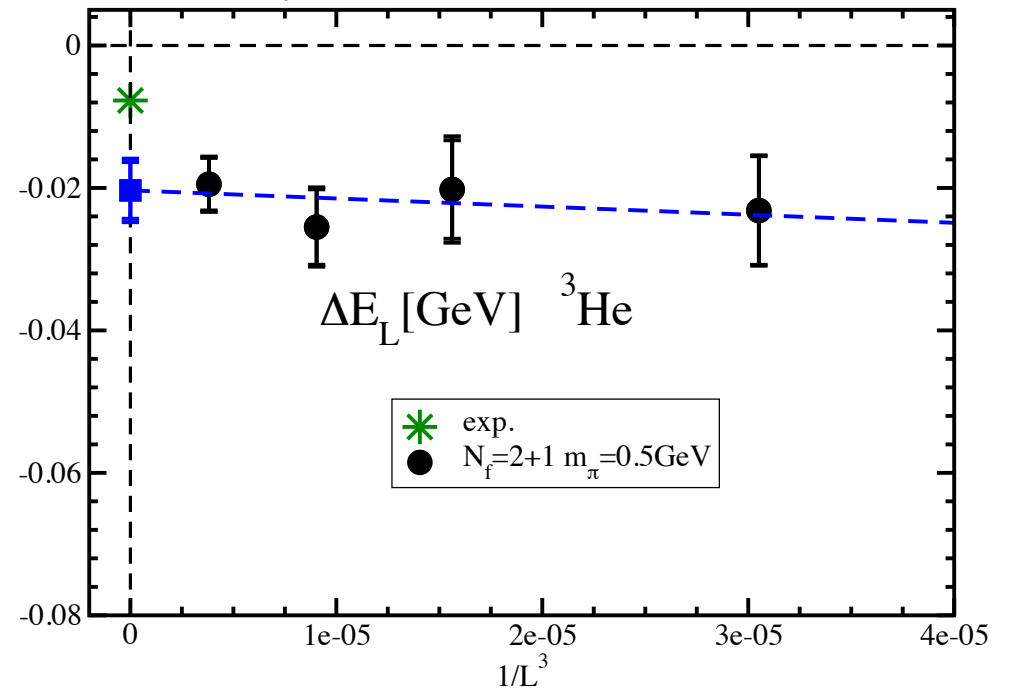
$^4\text{He}$  and  $^3\text{He}$  channels  $\Delta E_L = E_0 - N_N m_N$  at  $m_\pi = 0.5 \text{ GeV}$

TY *et al.*, PRD86:074514(2012)

Identification of bound state from volume dependence of  $\Delta E$



$$\Delta E_{4\text{He}} = 43(12)(8) \text{ MeV}$$



$$\Delta E_{3\text{He}} = 20.3(4.0)(2.0) \text{ MeV}$$

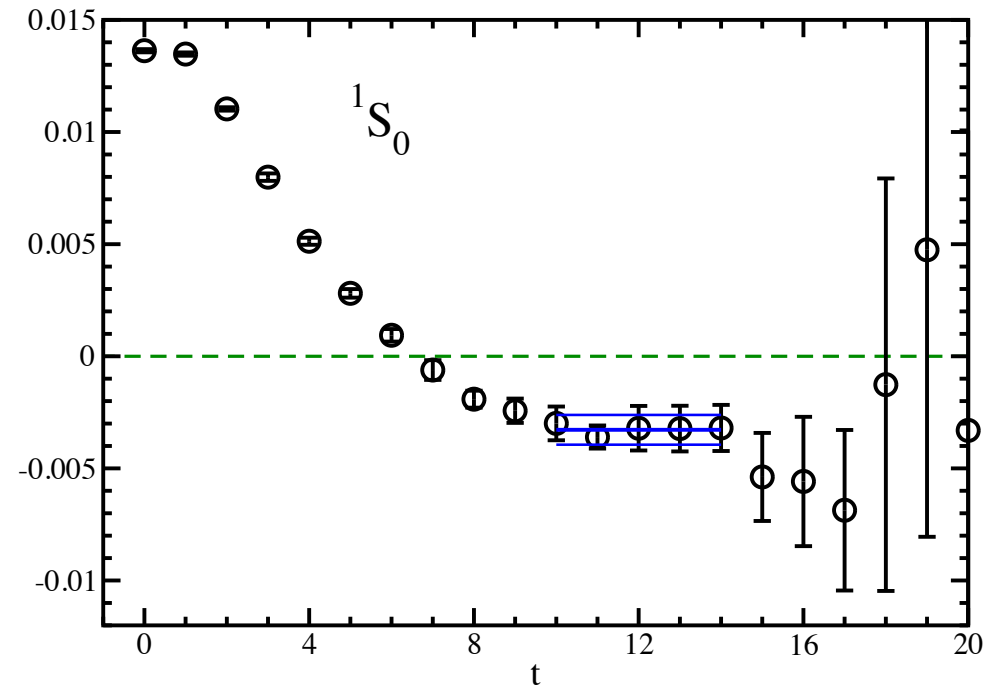
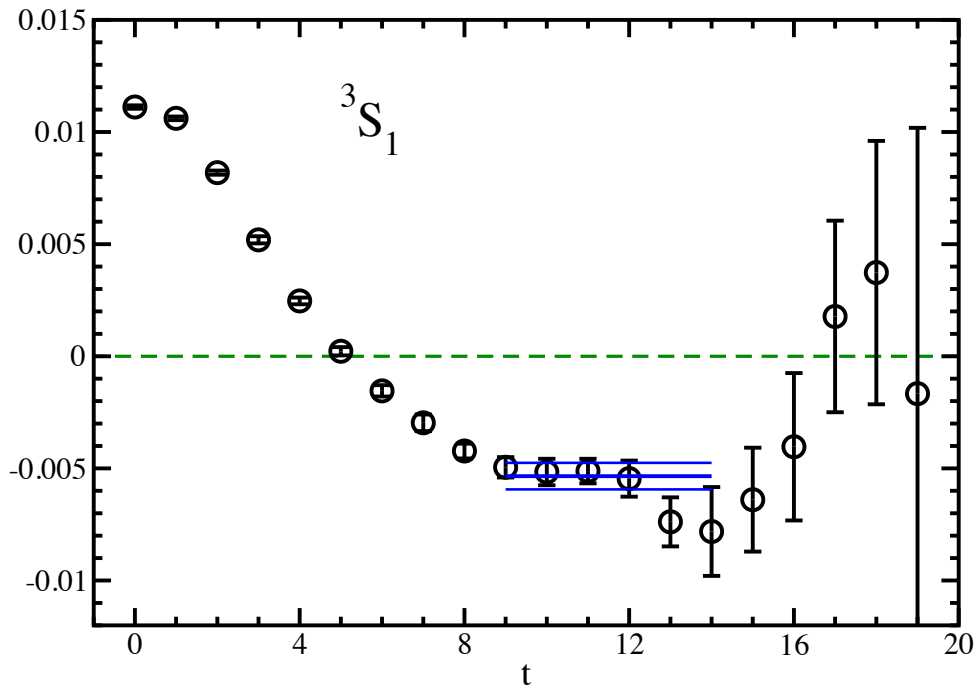
Observe bound state in both channels

$\Delta E$ : small difference from  $\exp(-cL)$  fit due to large error

$\Delta E_L$  in 2-nucleon channels at  $m_\pi = 0.5$  GeV on  $L = 5.8$  fm

TY *et al.*, PRD86:074514(2012)

$$\Delta E_L = \log \left( \frac{R_{NN}(t)}{R_{NN}(t+1)} \right) \text{ with } R_{NN}(t) = \frac{C_{NN}(t)}{(C_N(t))^2}$$

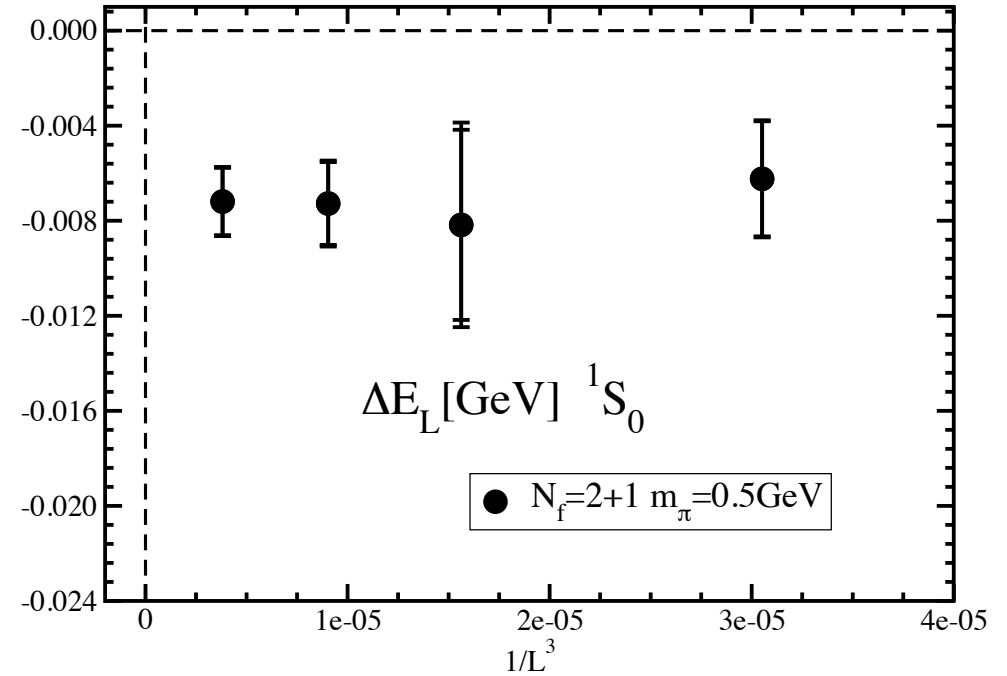
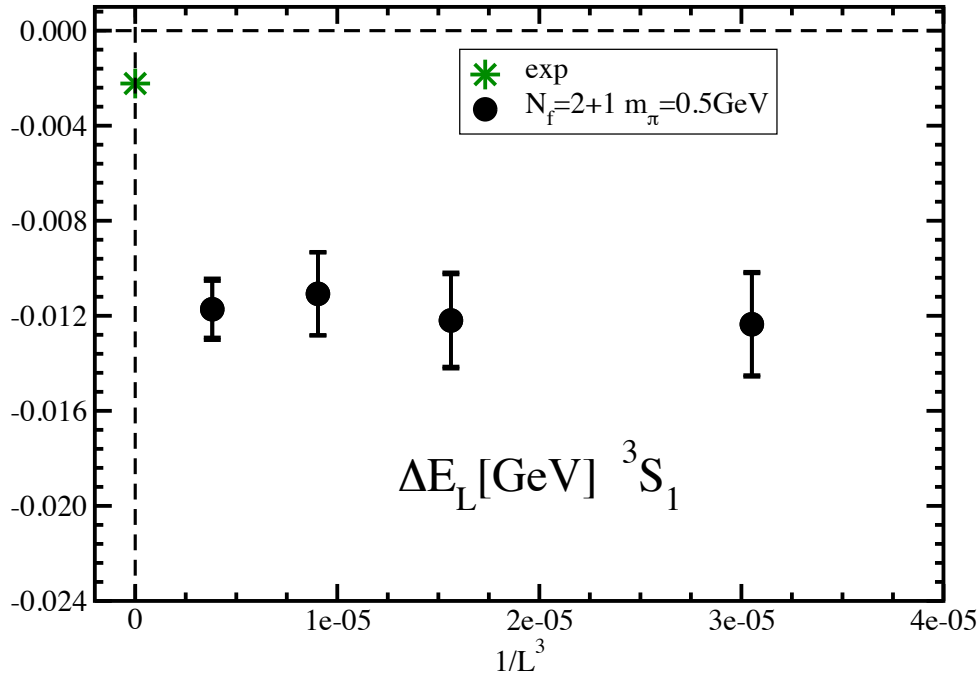


- Statistical error under control in  $t \leq 12$
- Smaller error than  $^4\text{He}$  and  $^3\text{He}$  channels
- Negative  $\Delta E_L$  in both channels
- Plateau region  $\sim$  plateau region of  $m_N$

NN ( $^3S_1$  and  $^1S_0$ ) channels  $\Delta E_L = E_0 - 2m_N$  at  $m_\pi = 0.5$  GeV

TY *et al.*, PRD86:074514(2012)

Identification of bound state from volume dependence of  $\Delta E$



- Negative  $\Delta E_L$
- Infinite volume extrapolation of  $\Delta E_L$

'04 Beane *et al.*, '06 Sasaki & TY

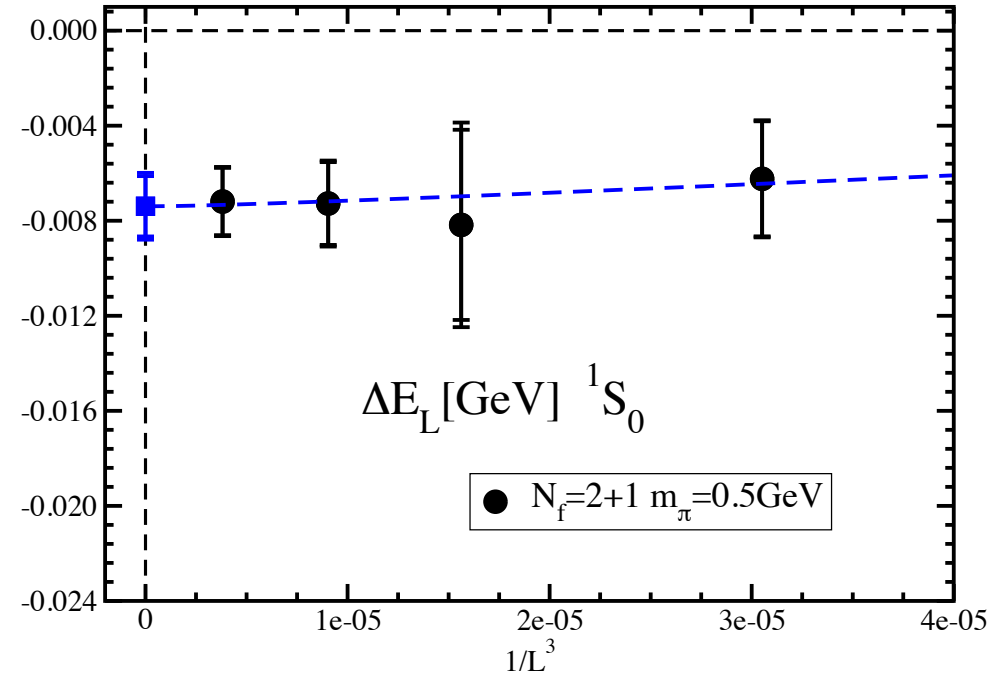
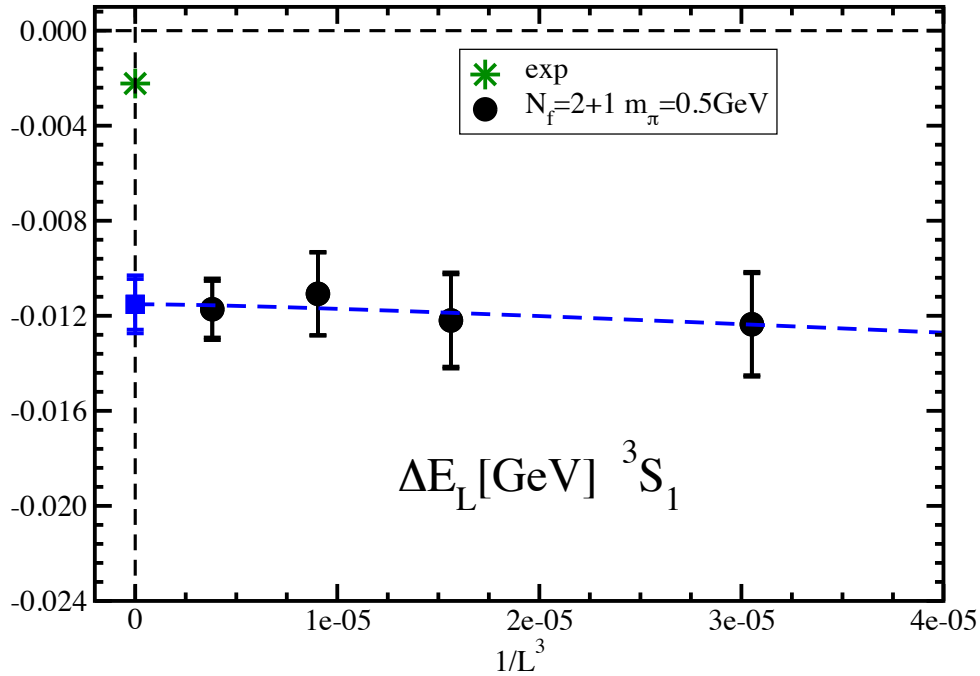
$$\Delta E_L = -\frac{\gamma^2}{m_N} \left\{ 1 + \frac{C_\gamma}{\gamma L} \sum_{\vec{n}}' \frac{\exp(-\gamma L \sqrt{\vec{n}^2})}{\sqrt{\vec{n}^2}} \right\}, \quad \Delta E_{\text{bind}} = \frac{\gamma^2}{m_N}$$

based on Lüscher's finite volume formula

NN ( $^3S_1$  and  $^1S_0$ ) channels  $\Delta E_L = E_0 - 2m_N$  at  $m_\pi = 0.5$  GeV

TY *et al.*, PRD86:074514(2012)

Identification of bound state from volume dependence of  $\Delta E$



Bound state in both channels ← different from experiment

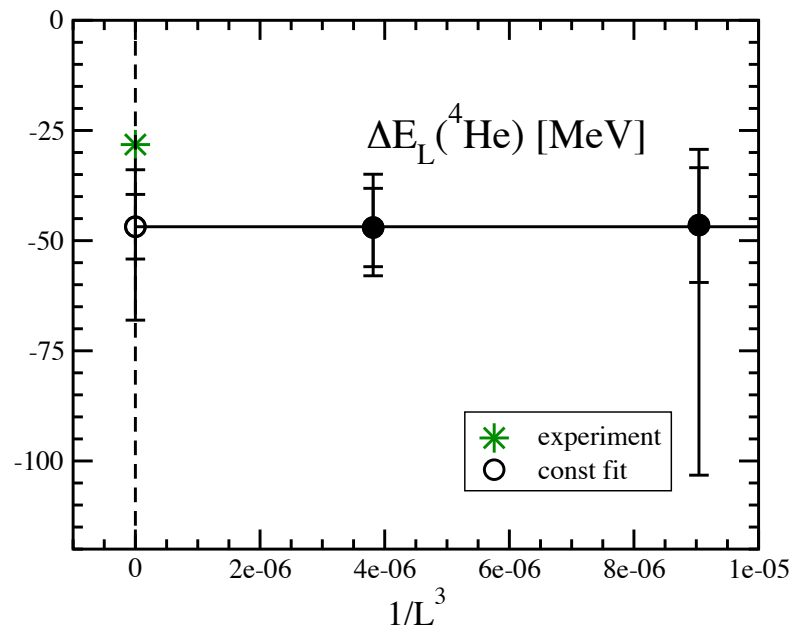
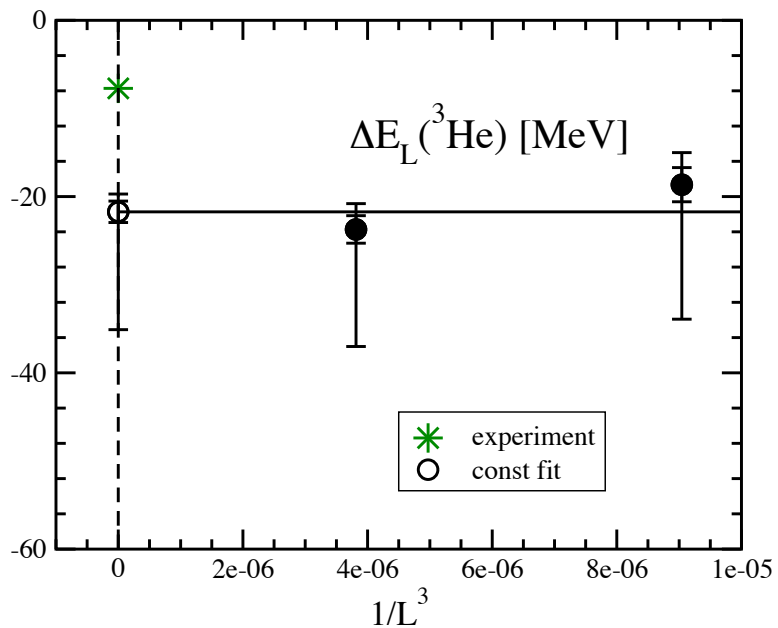
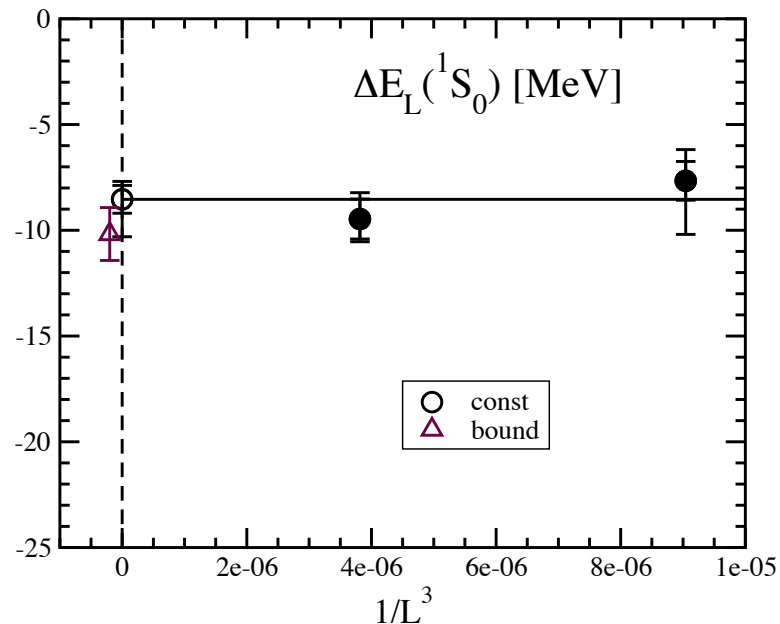
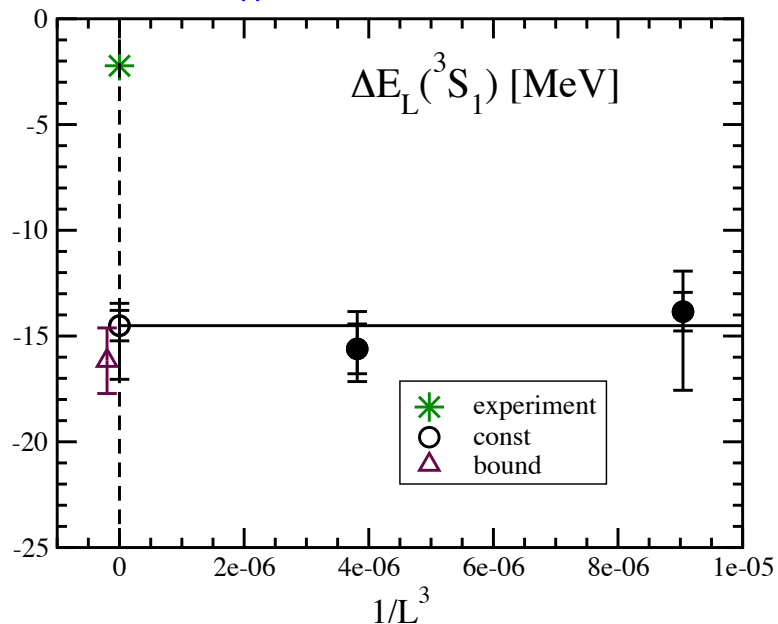
$$\Delta E_{^3S_1} = 11.5(1.1)(0.6) \text{ MeV}$$

$$\Delta E_{^1S_0} = 7.4(1.3)(0.6) \text{ MeV}$$

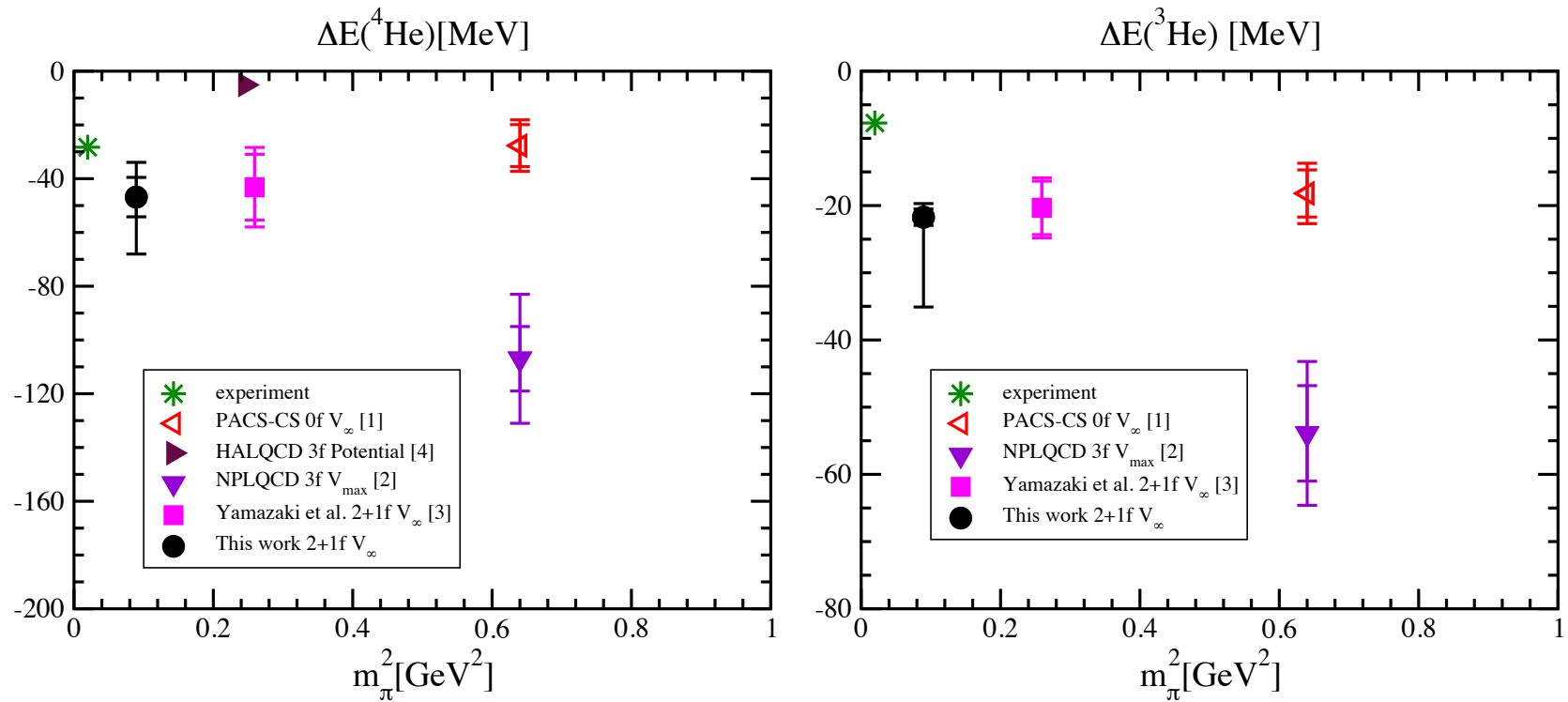


# Results at $m_\pi = 0.3$ GeV

TY *et al.*, PRD92:014501(2015)

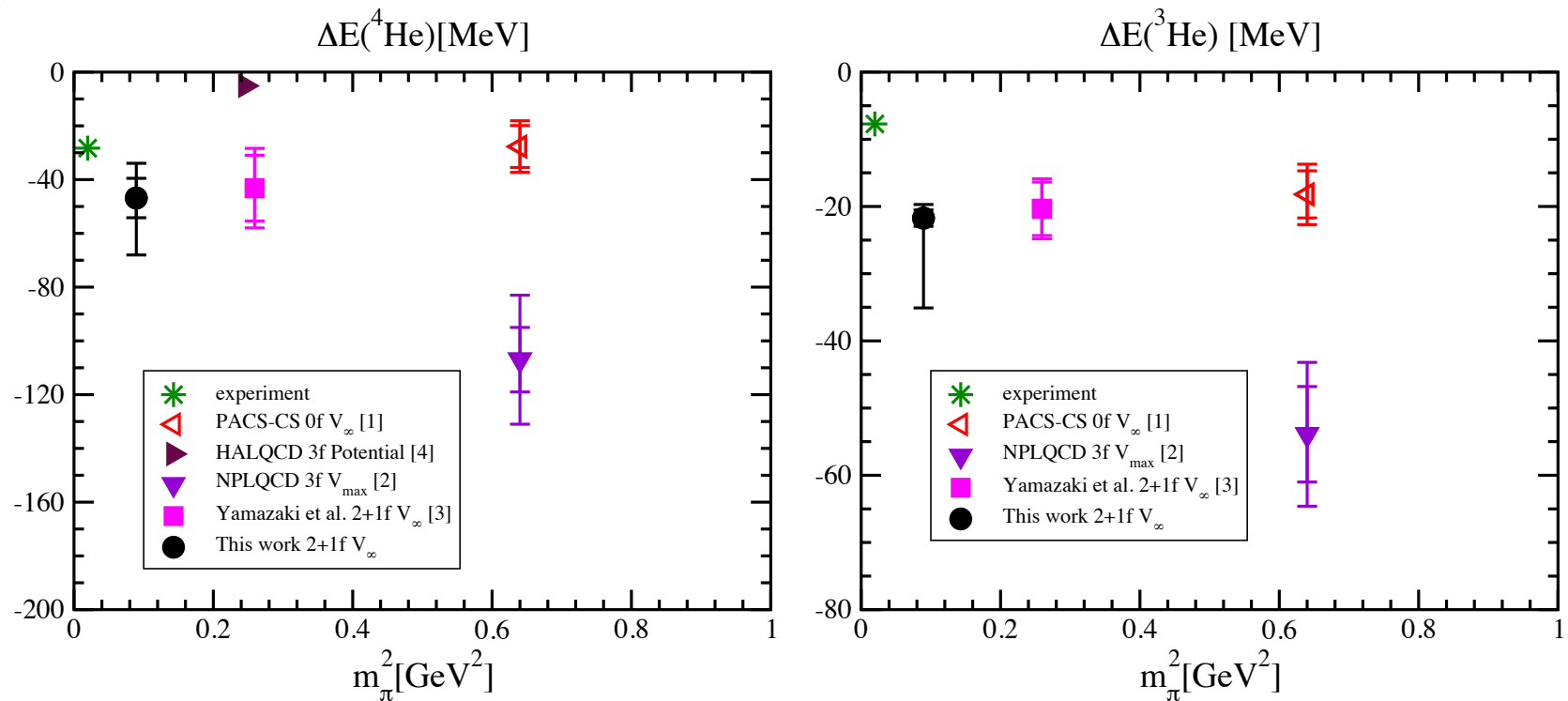


## Comparison of ${}^4\text{He}$ and ${}^3\text{He}$ channels



Light nuclei likely formed in  $0.3 \text{ GeV} \leq m_\pi \leq 0.8 \text{ GeV}$   
 Same order of  $\Delta E$  to experiments

## Comparison of ${}^4\text{He}$ and ${}^3\text{He}$ channels



Light nuclei likely formed in  $0.3 \text{ GeV} \leq m_\pi \leq 0.8 \text{ GeV}$

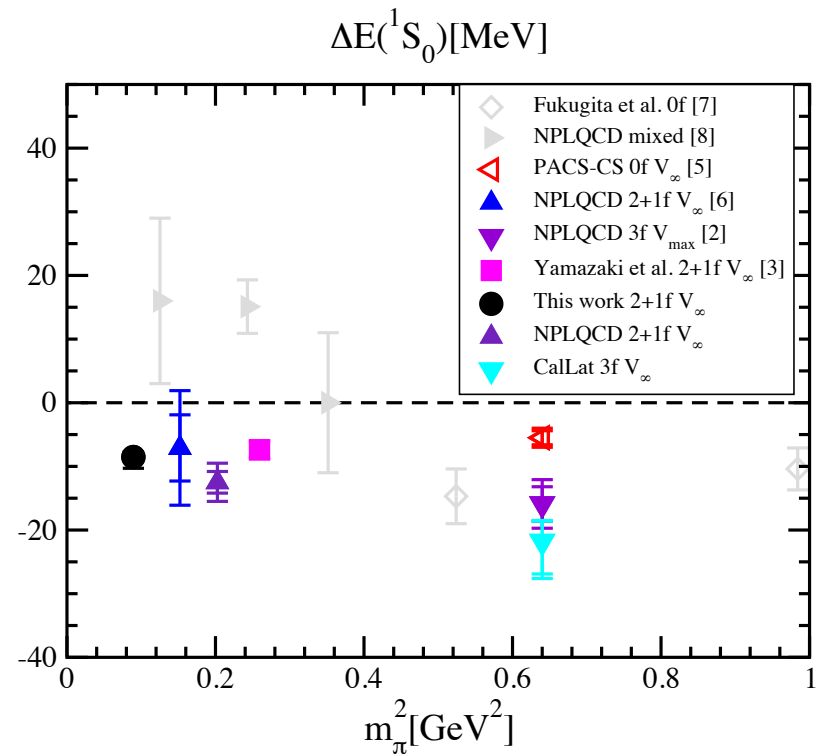
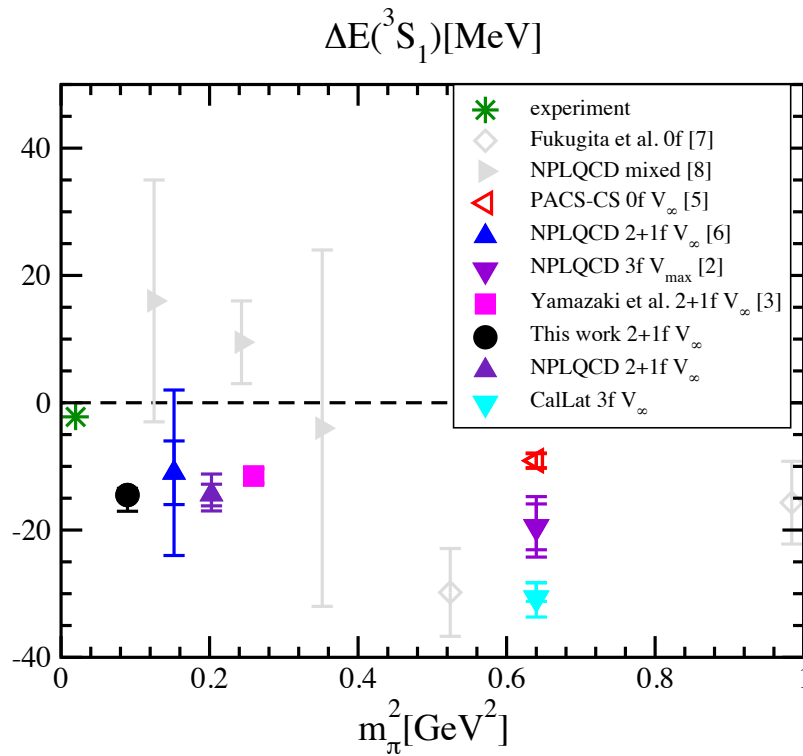
Same order of  $\Delta E$  to experiments  $\rightarrow$  relatively easier than  $NN$

large  $|\Delta E|$  makes less  $V$  dependence at physical  $m_\pi$

touchstone of quantitative understanding of nuclei from lattice QCD

Investigations of  $m_\pi$  dependence  $\rightarrow m_\pi \sim 0.145 \text{ GeV}$  on  $L \sim 8 \text{ fm}$

# Comparison of $NN$ channels



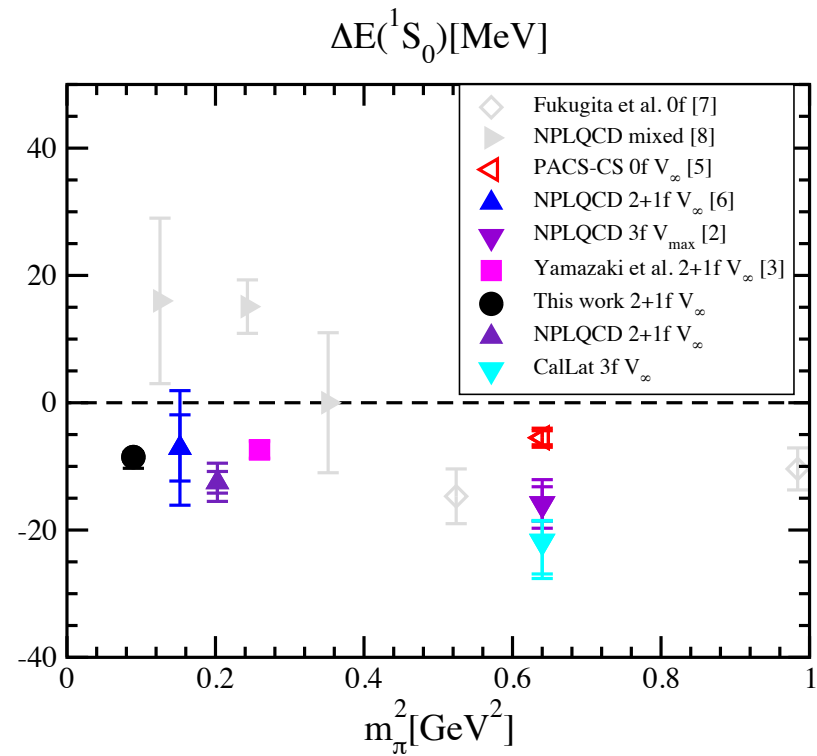
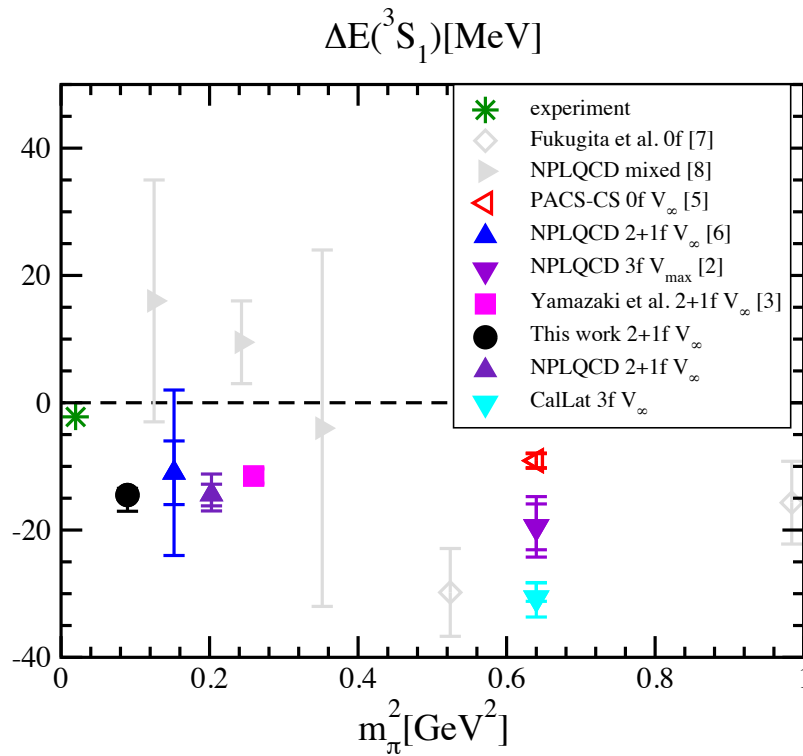
gray data: single volume calculation

$L^3 \rightarrow \infty$  data: **existence of bound states in  $^3S_1$  and  $^1S_0$**

**inconsistent with experiment due to larger  $m_\pi$  (?)**

Investigation of  $m_\pi$  dependence  $\rightarrow m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

# Comparison of $NN$ channels



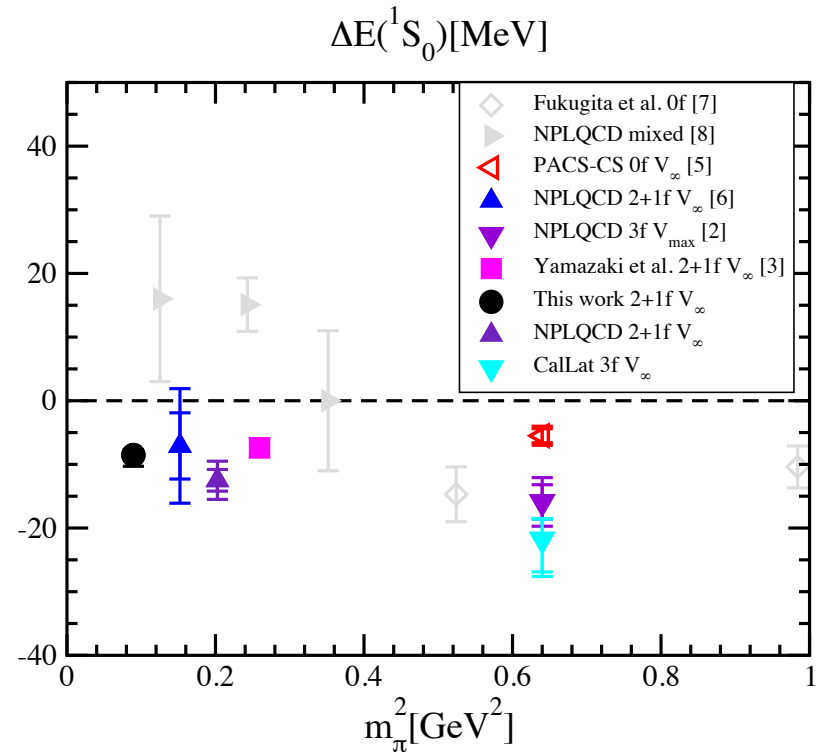
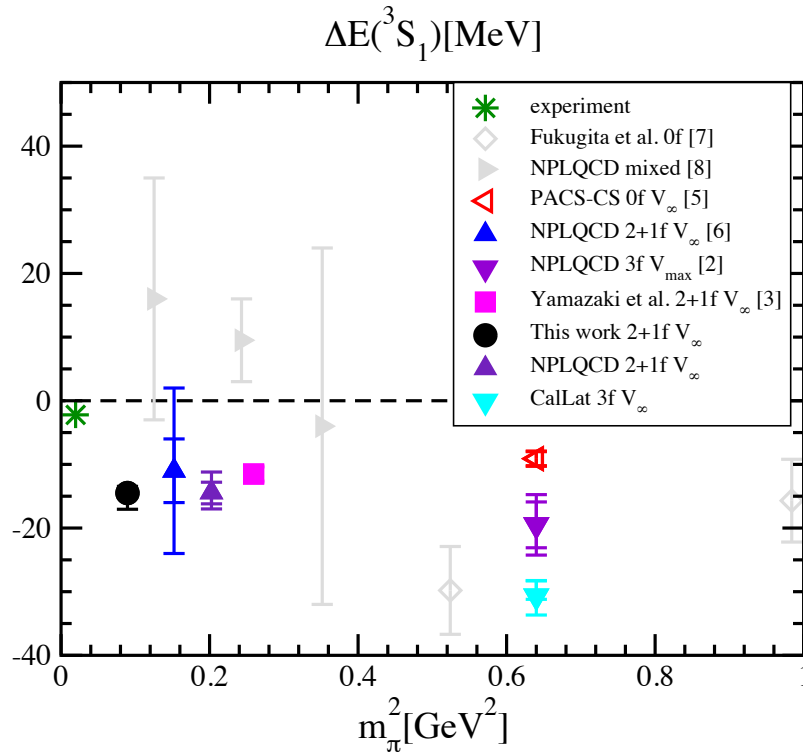
gray data: single volume calculation

Investigations of  $m_\pi$  dependence  $\rightarrow m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

Large finite volume effect expected even on  $L \sim 8$  fm

'86 Lüscher, '04 Beane *et al.*, '14 Briceño *et al.*

# Comparison of $NN$ channels



gray data: single volume calculation

Investigations of  $m_\pi$  dependence  $\rightarrow m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

Large finite volume effect expected even on  $L \sim 8$  fm

$${}^3S_1: \Delta E_{\text{exp}} = 2.2 \text{ MeV}$$

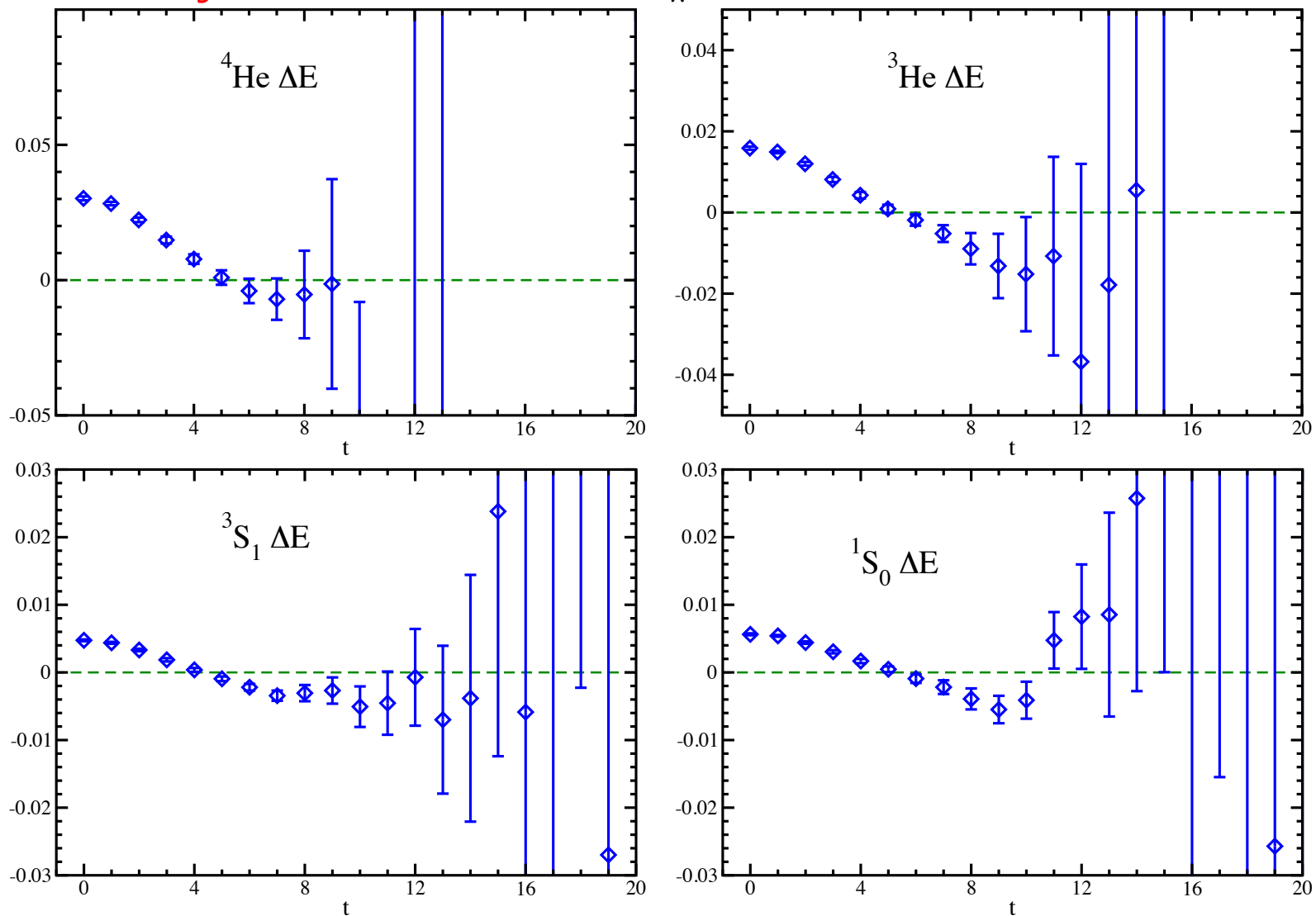
'86 Lüscher, '04 Beane *et al.*, '14 Briceño *et al.*

$$\Delta E_L = -(\Delta E_{\text{exp}} + \mathcal{O}(\exp(-L\sqrt{m_N \Delta E_{\text{exp}}})) \lesssim -4 \text{ MeV}$$

$${}^1S_0: a_0^{\text{exp}} = 23.7 \text{ fm}$$

$$\Delta E_L = -\frac{4\pi a_0^{\text{exp}}}{m_N L^3} + \mathcal{O}(1/L^4) \lesssim -2 \text{ MeV}$$

Very preliminary results of  $\Delta E$  at  $m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm



Computational resources

HA-PACS, COMA @Univ. of Tsukuba, K @AICS, FX100 @RIKEN

# Nucleon form factors at almost physical $m_\pi$

in collaboration with

K.-I. Ishikawa, Y. Kuramashi, S. Sasaki, and A. Ukawa  
for PACS Collaboration

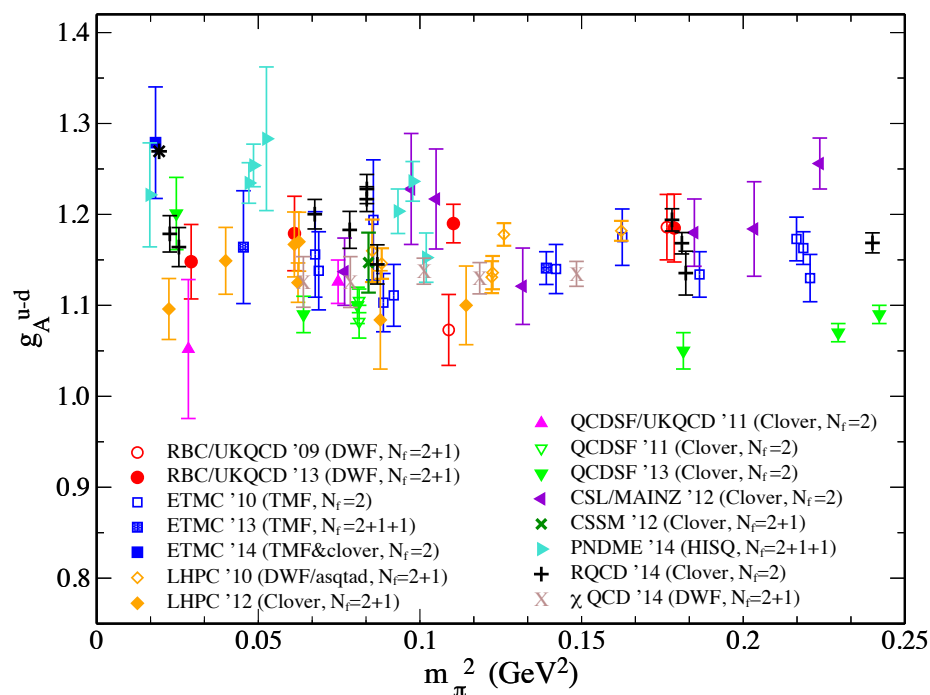
Computational resources (the HPCI System Research Project: hp140155, hp150135)  
COMA @Univ. of Tsukuba, FX10 @Univ. of Tokyo,  
FX100 @RIKEN, System E @Kyoto Univ., FX100 @Nagoya Univ.



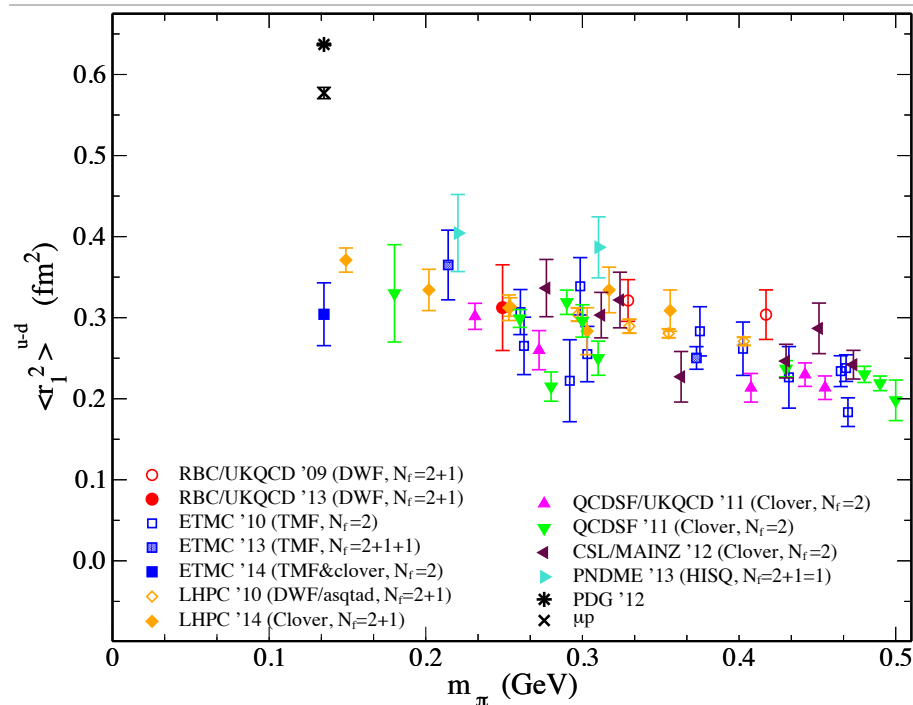
# Example of large quark mass dependence near $m_\pi \rightarrow 0$

## Isvector radii from form factors $F_1$ and $F_2$

Constantinou, Lat14 plenary  
axial charge



charge radius  $\langle r_1^2 \rangle$



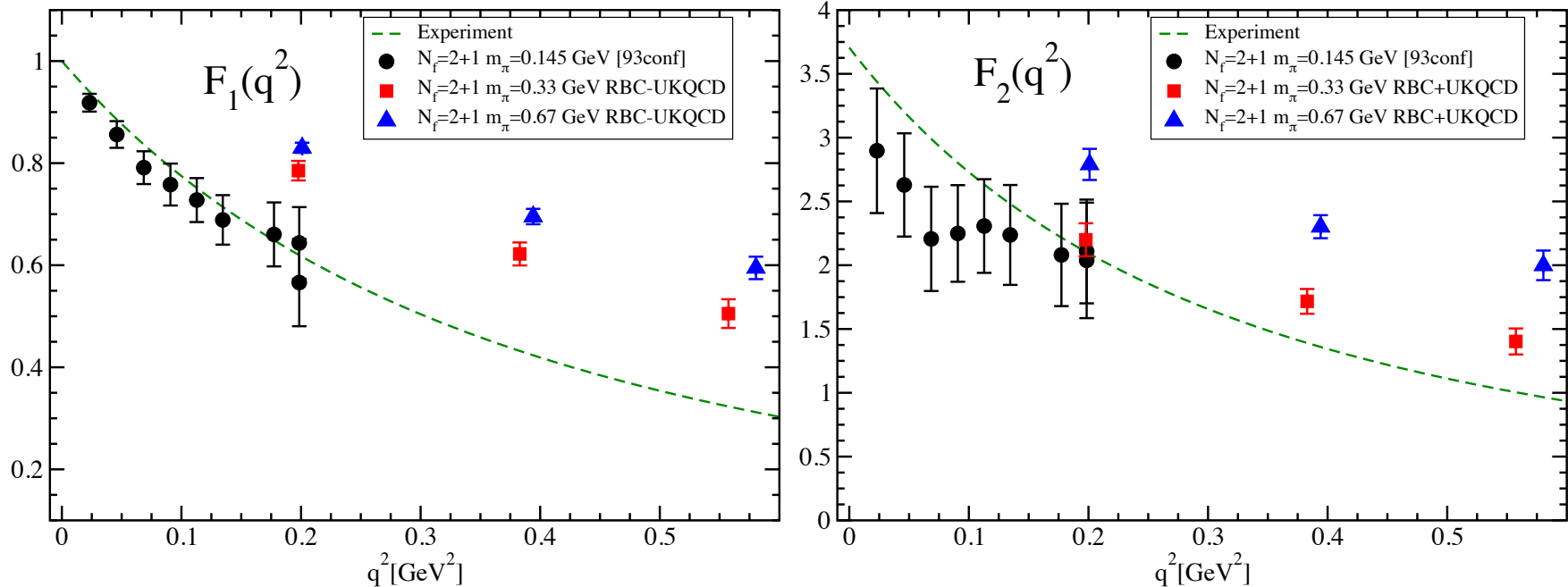
c.f.) '14 LHP, '15 Capitani *et al.*, '15 ETM, see also James's Lat15 plenary

important for understanding of nucleus property

Can we reproduce experiment at physical  $m_\pi$ ?

# Isvector $F_1$ and $F_2$ form factors

Preliminary results at  $m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm



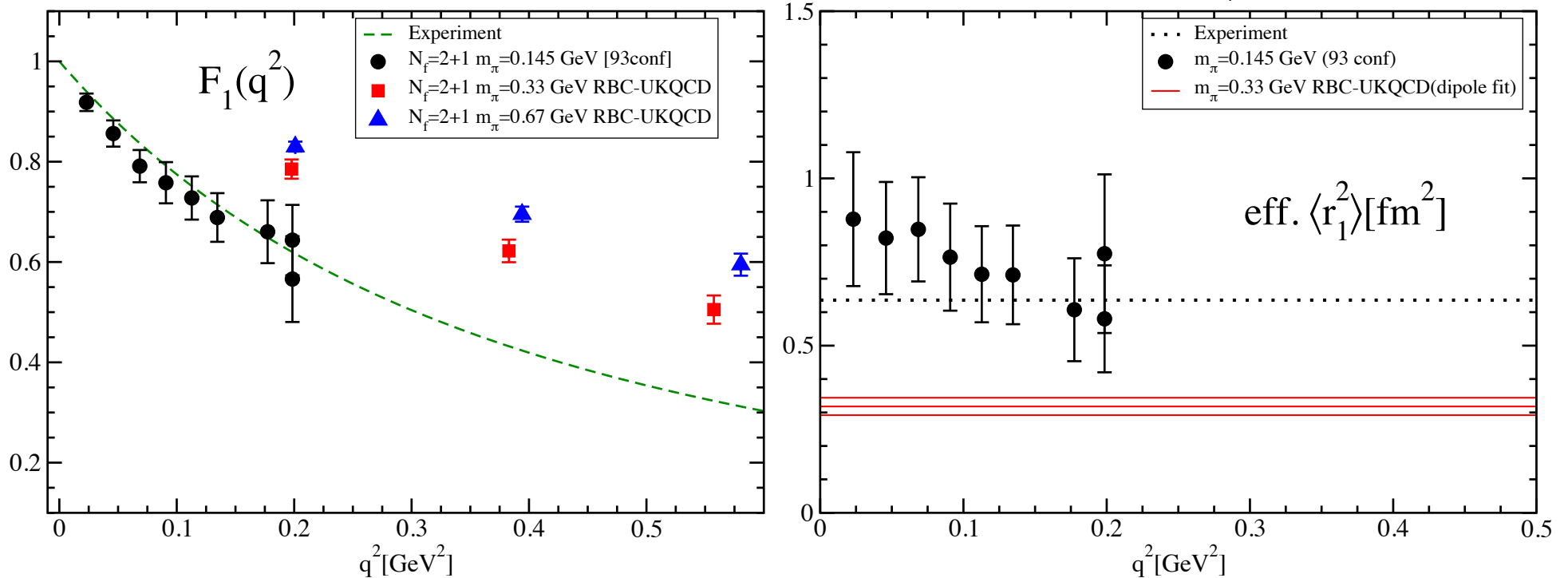
Need much more statistics  
but encouraging signal in  $G_E$

# Charge radius $\langle r_1^2 \rangle$

Preliminary results at  $m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

$$\text{Dipole form } F_1(q^2) = \left( 1 + \frac{q^2}{12} \langle r_1^2 \rangle \right)^{-2}$$

$$\text{Eff. } \langle r_1^2 \rangle = \frac{12}{q^2} \left( \sqrt{\frac{1}{F_1(q^2)}} - 1 \right)$$



Statistical error is large

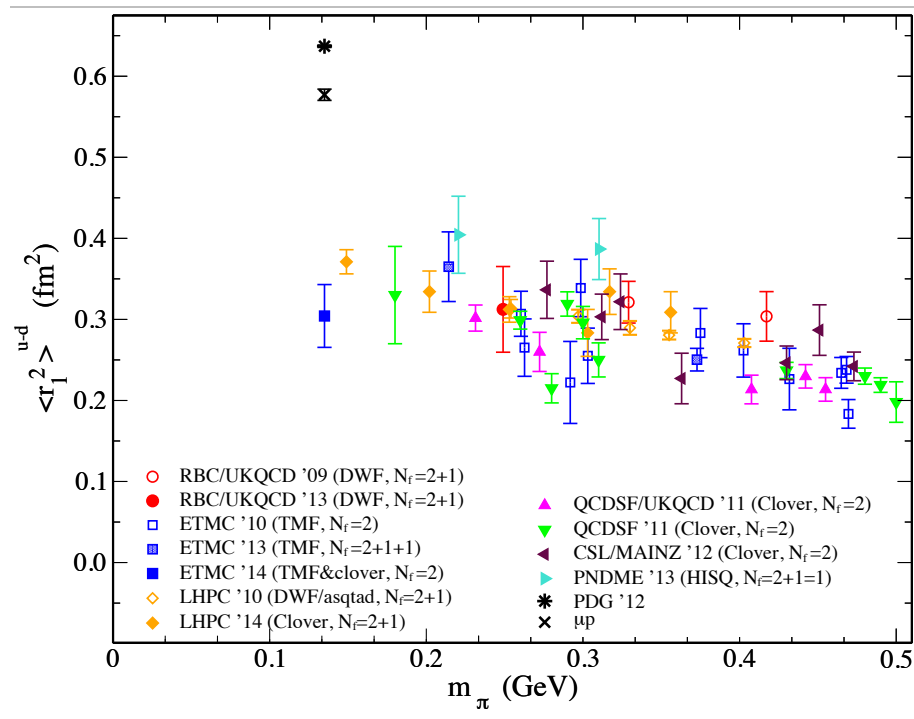
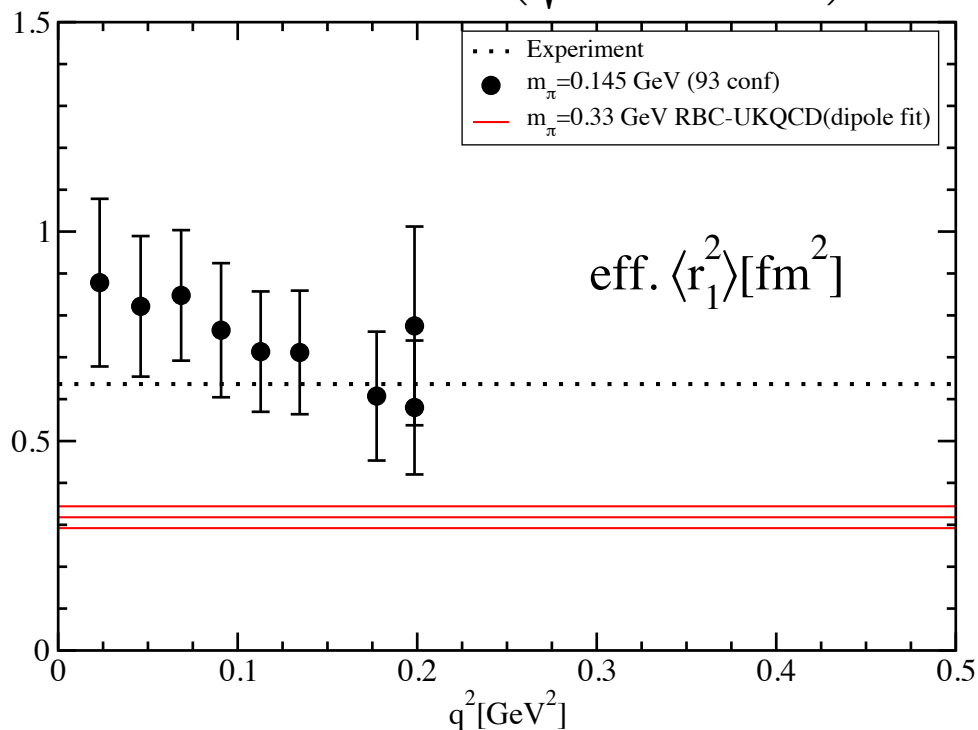
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Constantinou, Lat14 plenary



Need much more statistics  
but encouraging signal

Axial charge  $g_A = Z_A g_A^{\text{bare}}$

Preliminary results at  $m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

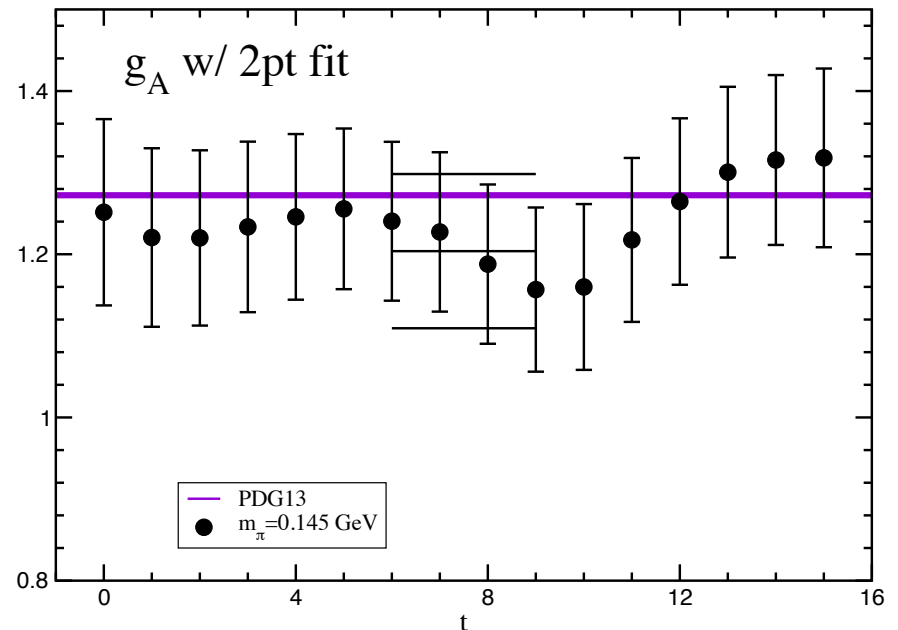
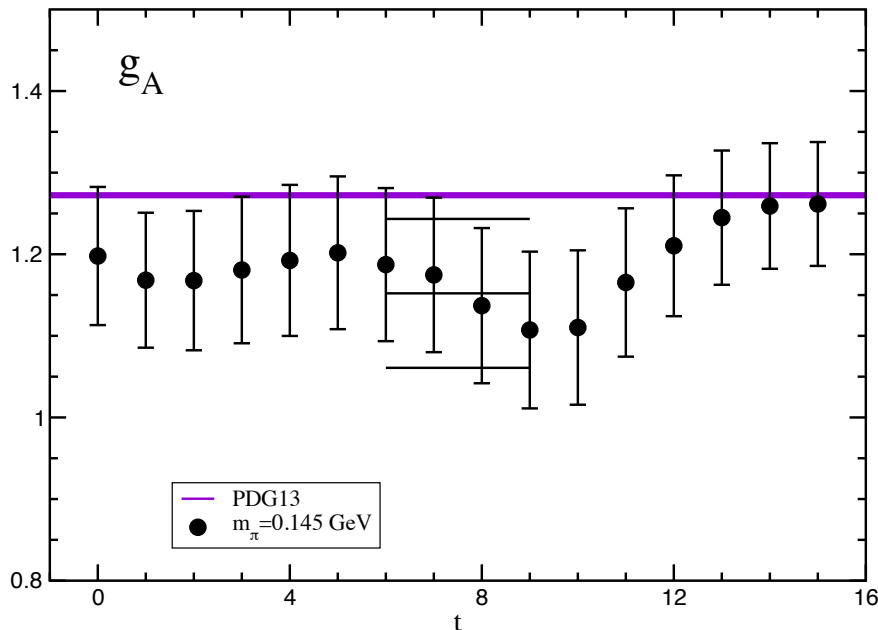
$Z_A$  from SF scheme (Lattice 2015, Ishikawa for PACS Collaboration)

consistent with  $Z_V = 1/G_E(0)$  within 1–2%

$$g_A^{\text{bare}} = C_{A_3}(t)/C_N(t_{\text{sink}})$$

$$g_A^{\text{bare}} = C_{A_3}(t)/(Z_N^2 \exp(-M_N t_{\text{sink}}))$$

$Z_N$  and  $M_N$  from fit of  $C_N(t)$



Discrepancy of two results  $\rightarrow$  systematic error

roughly consistent with experiment,

but need much more statistics for stringent test

# Summary

$N_f = 2 + 1$  lattice QCD at  $m_\pi = 0.5$  and  $0.3$  GeV

- Volume dependence of  $\Delta E$

$\Delta E \neq 0$  of 0th state in infinite volume limit

→ bound state in  ${}^4\text{He}$ ,  ${}^3\text{He}$ ,  ${}^3\text{S}_1$  and  ${}^1\text{S}_0$   
at  $m_\pi = 0.5$  and  $0.3$  GeV

- $\Delta E$  larger than experiment and small  $m_\pi$  dependence
- Bound state in  ${}^1\text{S}_0$  not observed in experiment

$N_f = 3$  at  $m_\pi = 0.8$  GeV by NPLQCD  
and CalLat with sophisticated sources

$N_f = 2 + 1$   $m_\pi = 0.45$  GeV by NPLQCD

No bound state in HALQCD method

variational method could give hint to solve the difference

## Need further investigations

e.g. systematic error from large  $m_\pi$  and finite lattice spacing

$N_f = 2 + 1$   $m_\pi \sim 0.145$  GeV on  $L \sim 8$  fm

$\Delta E$  for nuclei and Isovector form factors of nucleon