## Computational Advances in Nuclear and Hadron Physics Workshop

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**Relativistic Electro-Magneto-Fluid Dynamicsss** 

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#### **Relativistic Electro-Magneto-Fluid Dynamics**

#### **REMFD in Astrophysics & Cosmology**

**SKA Science – infer EoS and transport properties** 

#### NICA & FAIR Science – measure the eos and transport properies

# Pulsar Science Highlights

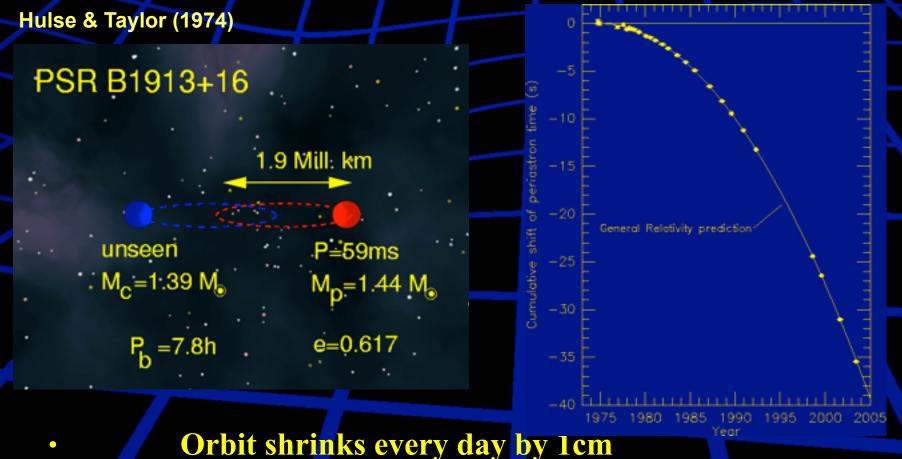
Key Science: Strong-field Tests of Gravity Was Einstein Right? Cosmic Censorship, "No-Hair" Theorem Cosmic Gravitational Wave Background

Variety of Other Major Astrophysical Topics:
Milky Way Structure, ISM
Intergalactic Medium
Relativistic Plasma Physics
Extreme Densities

# SKA

## Noted GR Laboratories

#### Weisberg & Taylor (priv. comm)



Confirmation of existence of gravitational waves

# Pulsars...

- embody physics of the EXTREME
  - surface speed ~0.1c
  - 10x nuclear density in centre
  - some have  $B > B_{q} = 4.4 \times 10^{13} G$
  - Voltage drops ~ 10<sup>12</sup> volts
  - $-F_{EM} = 10^9 F_g = 10^{11} F_{gEarth}$
  - Tsurf ~ million K
- …relativistic plasma physics in action
  - ...probes of turbulent and magnetized ISM ...precision tools, e.g.
    - Period of B1937+21:

0.0015578064924327±0.0000000000000004 s - Orbital eccentricity of J1012+5307:

# Was Einstein right?

General Relativity vs Alternative Theories



Strong Equivalence Principle
Violation of Lorentz-Invariance
Violation of Positional Invariance
Violation of Conservation Laws etc.

Solar System tests provide constraints but only in weak field!

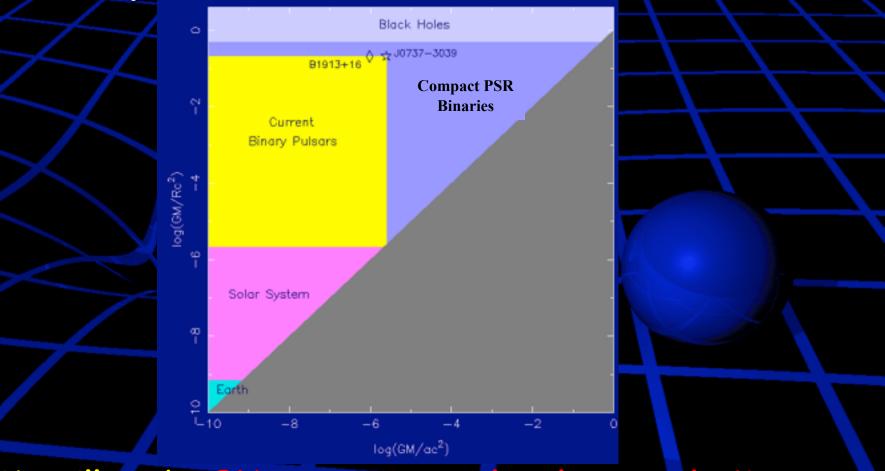


No test of any theory of gravity is complete, if only done in solar system, i.e. strong field limit and radiative aspects need to be tested, too! ≈10<sup>-4</sup>

Jorbit

This is and will be done best with radio pulsars!

# Exploration of Black Holes



#### We will probe BH properties with pulsars and SKA:

- precise measurements
- no assumptions about EoS or accretion physics
  - test masses well separated not deformed

# 

ullet

Result is relativistic & classical spin-orbit coupling Visible as a precession of the orbit: Measure higher order derivatives of secular changes in semi-major axis and longitude of periastron (relativistic) or transient TOA perturbations (classical) Not easy! It is not possible today!

Wide range of applications:

 Galactic probes:
 Interstellar
 medium/magnetic
 field

 Magnetic field
 Star
 formation
 history

Dynamics

Population via distances (ISM, VLBI)

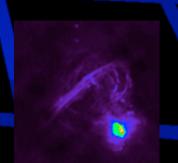
distribution

4.85 GHz PSR 81758-23

1.41 GHz

and the second second second second

50 100 150 200 250 300 350 400 450 Time (msec)



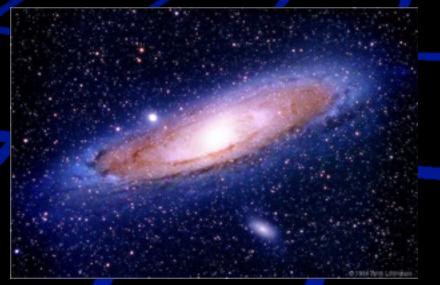
#### Galactic Centre

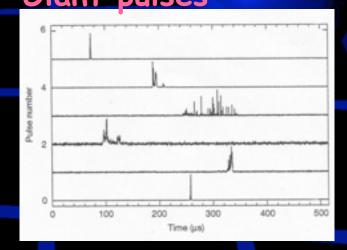
Movement in potential

Wide range of applications:

# Galactic probes Extragalactic pulsars: Missing Baryon Problem Formation & Population Turbulent magnetized IGM Giant pulses

#### Search nearby galaxies!





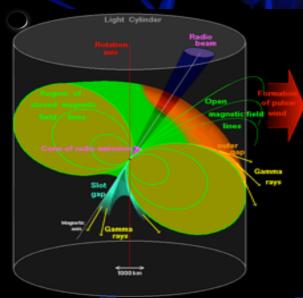
Reach the local group!

Wide range of applications:

#### Extragalactic pulsar

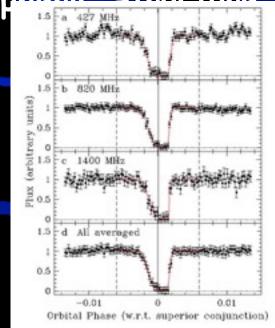
Galactic probes

#### Relativistic plasma physics: Emission Processes



## Magnetosphenic Structure

Pulsar Wind Nebulae



Mass [solar masses]

Wide range of applications:

Extragalactic pulsars Relativistic plasma physics

Galactic probes

#### Matter Physics: Ultra-strong B-fields

10

Atmosphere Superhot plasma

> Outer crust Starquakes / Crystal lattice: 200 m deep nuclei + electrons

Inner crust Starquakes Crystal lattice: 1 km deep nuclei + electrons + neutron drip

🚽 20 km (12 mi) diameter

Outer core Atomic particle fluid

nner core Solid block of subatomic particles? Equation-of-State



10<sup>17</sup> kg/m<sup>3</sup>

Radius [km]

1000

10000

100

Wide range of applications:

Galactic probes Extragalactic pulsa



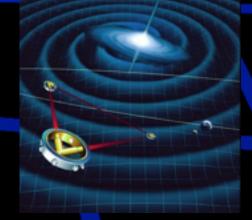
Extreme Dense Matter Physi

Multi wavelength studies: Photonic windows

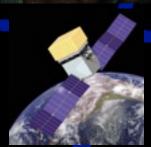
#### Non-photonic windows











## Pulsar Astrophysics with SKA Wide range of applications: Holy Grail: PSR-BH Galactic probe Extragalactic pulsars Relativistic plasma physics Extreme Dense Matter Physics Multi-wavelength studies 'stems: planets pulsar/MS binaries millisecond pulsars relativistic binaries double pulsars Double Pulsars PSR-BH systems Planets

Cosmological Gravitational Wave Background

 stochastic gravitational wave background expected on theoretical grounds

Possible Sources:

Inflation String cosmology Cosmic strings phase transitions



and also: merging massive BH binaries  $\hbar_0^2 \Omega_{GW}(1) \approx 1^{2/3}$ in early galaxy evolution

## Cosmological Gravitational Wave Background

Pulsars discovered in Galactic Census also provide network of arms of a huge cosmic gravitational wave detector

Perturbation in space-time can be detected in timing residuals Pulsar Timing Array

**Sensitivity: dimensionless strain** 



#### **Relativistic Electro-Magneto-Fluid Dynamics**

#### **Combined EMFD**

- **Relativistic dissipative fluids relaxation equations**
- Maxwell-Lorentz general equations
- **Einstein GR equations** 
  - **Computational techniques**

#### **Thermodynamics of Polarized Media**

The Maxwell-Lorentz field equations are

$$F_{\alpha\beta} = 2\partial_{[\alpha}A_{\beta]} \quad H^{\alpha\beta}_{\ \ |\beta} = 4\pi\vartheta^{\alpha} \quad H^{\alpha\beta} = \Phi^{\alpha\beta} - 4\pi M^{\alpha\beta}$$

The skew tensors may be decomposed into pairs of spatial vectors as

$$F_{\alpha\beta} = (B,E)$$
  $H^{\alpha\beta} = (H,\Delta)$   $M^{\alpha\beta} = (M,-\Pi)$ 

so that, for example,  $E_{\alpha} = \Phi_{\alpha\beta} \upsilon^{\beta}$ 

where  $u_{\alpha}$  is the 4-predominant of the dome dome defined ( $\upsilon_{\alpha}\upsilon^{\alpha} = -1$ )

# The energy-momentum tensor of the electromagnetic field is given by

$$4\pi T^{\lambda\mu}_{(\epsilon,\mu)} = \Phi^{\lambda}_{\ \mu} H^{\mu\alpha} - \Pi^{\lambda} \upsilon^{\mu} - \Pi^{[\lambda} E^{\mu]} - M^{[\lambda} B^{\mu]} - \frac{1}{4} \gamma^{\lambda\mu} \Phi^{\alpha\beta} \Phi_{\alpha\beta}$$
with

$$\Pi^{\lambda} = 2u_{\mu}F^{[\lambda}{}_{\alpha}M^{\mu}{}^{]\alpha} \quad \Pi = \mathbf{P} \times \mathbf{B} - \mathbf{M} \times \mathbf{E}$$

#### **Thermodynamics of Polarized Media**

The primary variables  $N^{\mu}$ ,  $T^{\lambda\mu}_{(mat)}$  for an arbitrary state are decomposed in the form

$$N^{\mu} = \nu \upsilon^{\mu}$$

$$T^{\lambda\mu}_{\ (\mu \circ \eta)} = \rho \upsilon^{\lambda} \upsilon^{\mu} + 2\theta^{(\lambda} \upsilon^{\mu)} + \Pi \Delta^{\lambda\mu} + \pi^{\lambda\mu}$$

$$\upsilon_{\lambda} \Pi^{\lambda} = \upsilon_{\lambda} \theta^{\lambda} = \upsilon_{\lambda} \pi^{\lambda\mu} = \pi^{\lambda\mu} \upsilon_{\mu} = \pi^{\lambda\mu} \Delta_{\lambda\mu} = 0$$

#### **Total energy-momentum tensor**

The defination of the electromagnetic energy momentum tensor is purely formal.

The physically significant quantity is the total energy-momentum tensor , which couples to Einstein tensor in the gravitational field equations and satisfies

$$T^{[\lambda\mu]} = T^{\lambda\mu}_{\ |\mu} = 0$$

## **Entropy production and transport equations**

The fundamental reation between the primary variables for an arbitrary state close to equilibrium

$$S^{\mu} = \prod_{(0)} \beta^{\mu} - \alpha N^{\mu} - \beta^{\lambda} T_{\lambda}^{(\mu^{-\lambda})} - \Theta^{\mu}$$

From which we get the entropy production

$$\begin{split} TS^{\mu}_{|\mu} &= \phi^{\mu} E_{\mu} - \theta^{\mu} \left( \beta^{-1} \partial_{\mu} \beta - \dot{\upsilon}^{}_{\mu} \right) + \pi^{\lambda \mu}_{(1)} \upsilon_{(\lambda|\mu)} \\ &+ \Pi^{\mu}_{(1)} \left( \Lambda_{\mu} E_{\mu} - \upsilon_{(\lambda|\mu)} E^{\lambda} \right) + M^{\mu}_{(1)} \left( \Lambda_{\mu} B_{\mu} - \upsilon_{(\lambda|\mu)} B^{\lambda} \right) \\ &- T\Theta^{\mu}_{|\mu} \end{split}$$

## **Relativistic equations for stellar structure**

**Static** and **sphericaly symmetric** self-gravitating mass distribution

$$ds^{2} = g_{\mu} v dx^{\mu} dx^{\nu} = e^{2\Phi(r)}c^{2} d^{2} - e^{2\lambda(r)} d^{2} - r^{2} (d\theta^{2} + \sin^{2}\theta d\phi^{2})$$
$$e^{\lambda(r)} i \left[1 - \frac{2Gm(r)}{c^{2}r}\right]^{1/2}$$

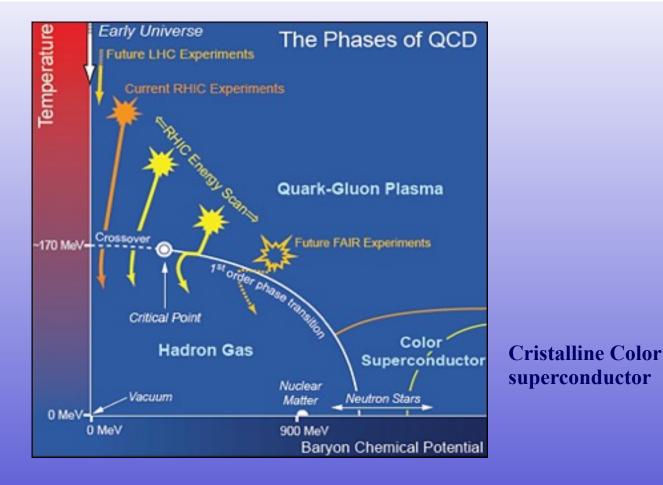
for the present case the Einstein's field equations take the form called the **Tolman – Oppenheimer – Volkov equations (TOV)** 

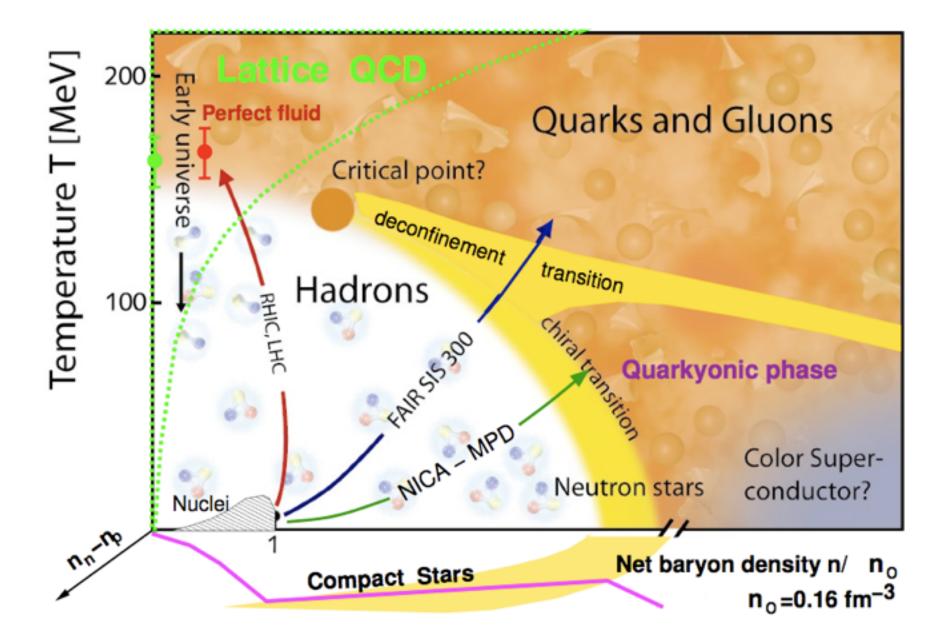
$$\frac{dP}{dr} = -G\frac{m(r)\rho(r)}{r^2} \left(1 + \frac{P(r)}{c^2\rho(r)}\right) \left(1 + 4\pi \frac{r^3 P(r)}{m(r)c^2}\right) \left[1 - \frac{2Gm(r)}{c^2 r}\right]^{-1}$$

$$\frac{dm}{dr} = 4 \pi r^2 \rho(r)$$
$$\frac{d\Phi}{dr} = -\frac{1}{\rho(r)c^2} \frac{dP}{dr} \left(1 + \frac{P(r)}{\rho(r)c^2}\right)^{-1}$$

One needs the equation of state (EOS) of dense matter,  $P = P(\rho)$ , up to very high densities

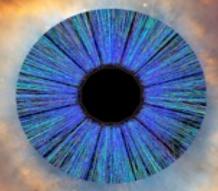
## Neutron Stars in the QCD phase diagram





## 27<sup>th</sup> Chris Engelbrecht Summer School 2016 Hot and Dense Matter in Heavy Ion Collisions and Astrophysics

Pretoria, South Africa, 11 – 22 January 2016



Organizing Committee A Muronga (UJ) F.G. Scholtz (NITheP)

#### Topics