

# Limits on **Statistical Anisotropy** from **Large-scale structure** of the Universe

arXiv:1612.02645 (Phys.Rev.D95, 063508)

arXiv:1704.02868 (**accepted** by MNRAS last night!)

**Naonori Sugiyama** (Kavli IPMU)

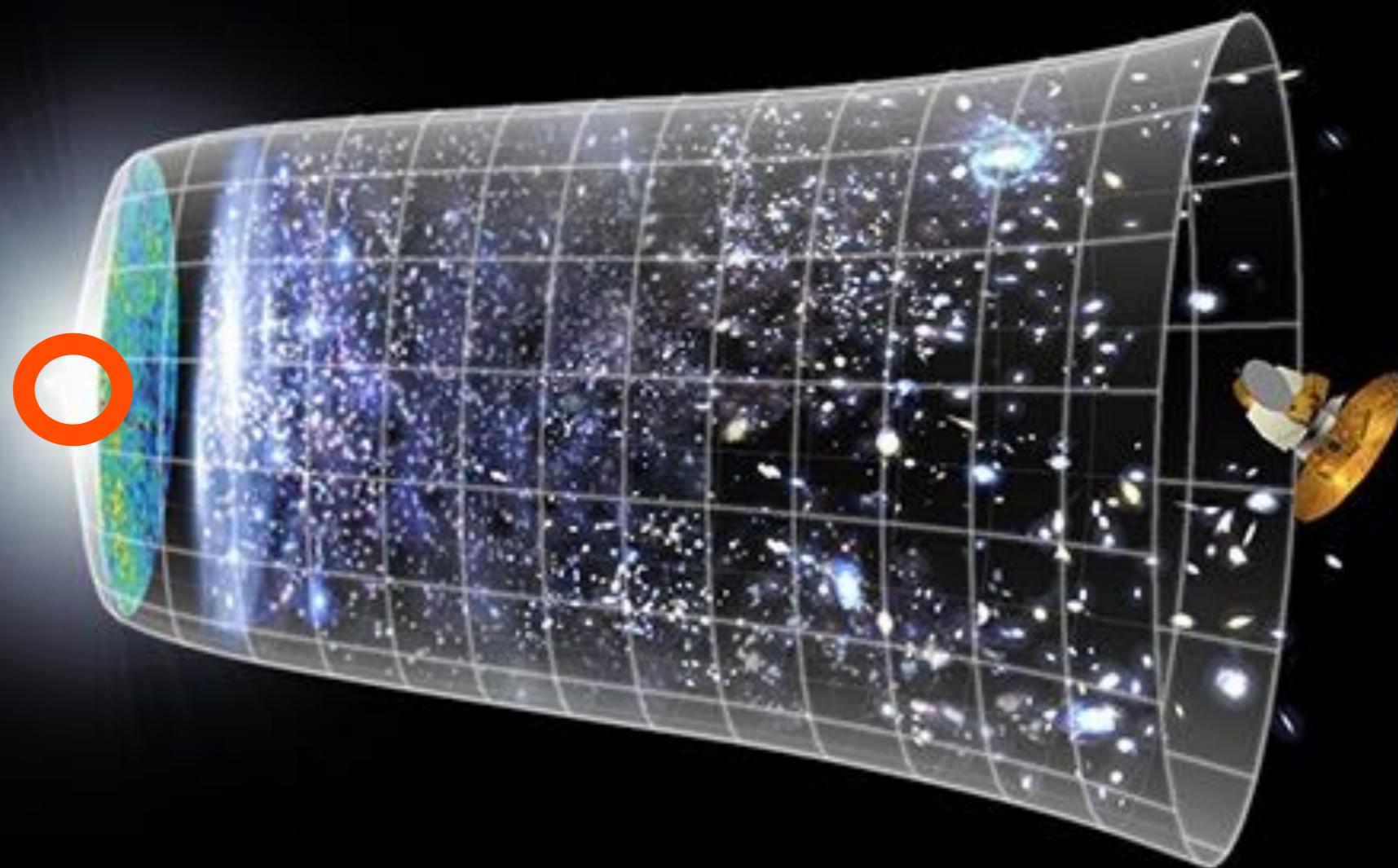
Collaborators:

**Maresuke Shiraishi and Teppei Okumura**

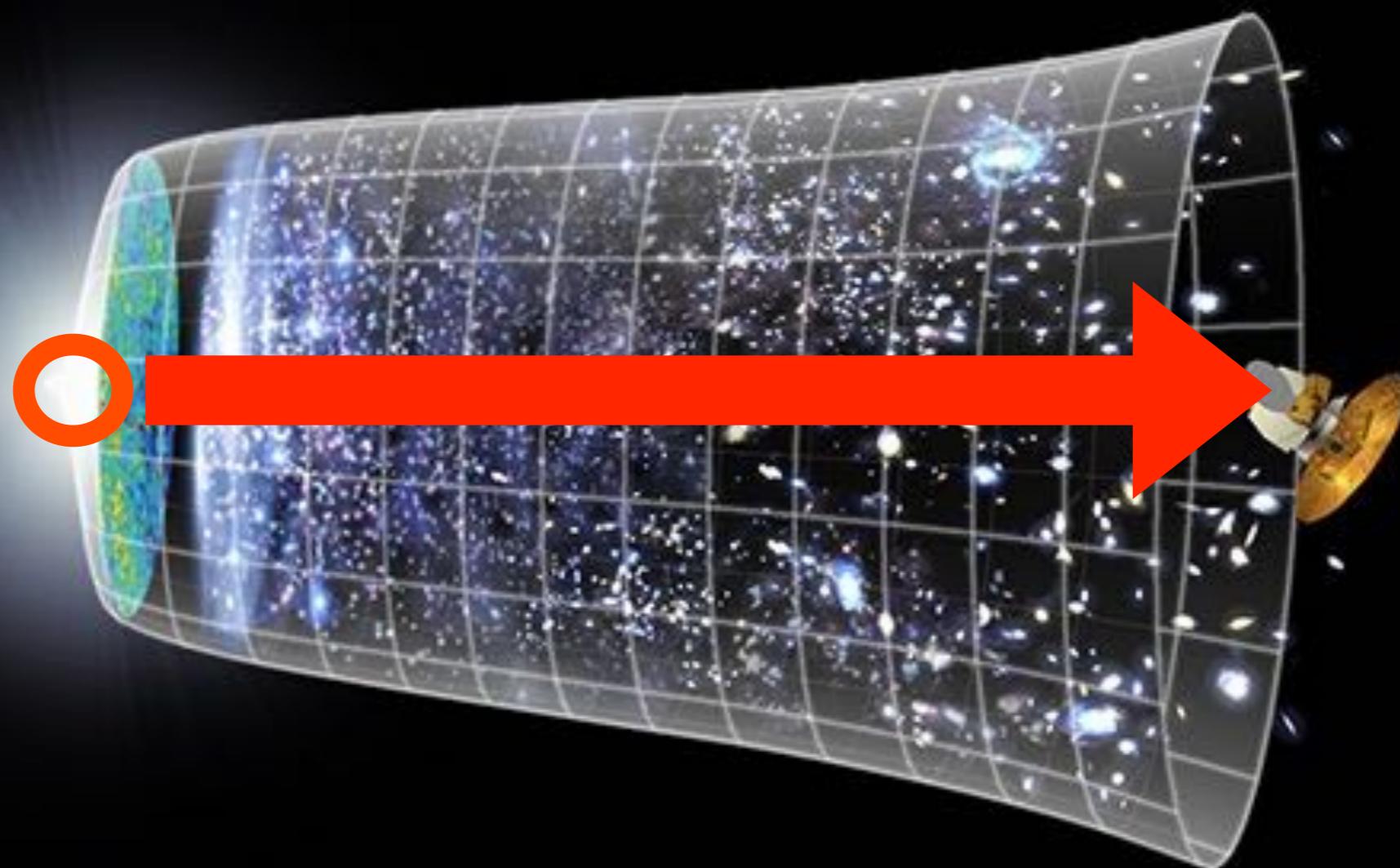
Next generation cosmology with large-scale structure:

CosKASI-ICG-NAOC-YITP joint workshop @ Kyoto Univ., Sep.7-8, 2017

# The beginning of the Universe



# The beginning of the Universe



# The standard inflation scenario **(Slow-roll single scalar field inflation)**

## Lagrangian

$$\mathcal{L} = -\frac{1}{2}g^{\mu\nu}\partial_\mu\varphi\partial_\nu\varphi - V(\varphi)$$

## Energy and pressure

$$\bar{\rho} = \frac{1}{2}\dot{\bar{\varphi}}^2 + V(\bar{\varphi}),$$

$$\bar{p} = \frac{1}{2}\dot{\bar{\varphi}}^2 - V(\bar{\varphi}),$$

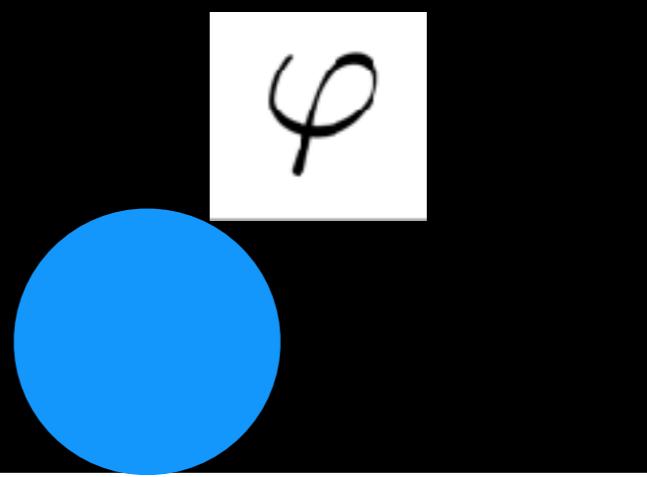
# The standard inflation scenario (Slow-roll single scalar field inflation)

If potential is flat (constant):

$$\bar{\rho} = \frac{1}{2}\dot{\bar{\varphi}}^2 + V(\bar{\varphi}),$$

$$\bar{p} = \frac{1}{2}\dot{\bar{\varphi}}^2 - V(\bar{\varphi}),$$

$$V(\varphi) = \text{const.}$$



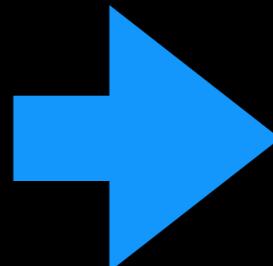
# The standard inflation scenario (Slow-roll single scalar field inflation)

If potential is flat (constant):

$$\bar{\rho} = \frac{1}{2}\dot{\bar{\varphi}}^2 + V(\bar{\varphi}),$$

$$\bar{p} = \frac{1}{2}\dot{\bar{\varphi}}^2 - V(\bar{\varphi}),$$

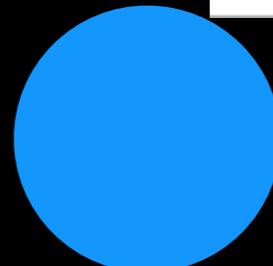
Cosmological constant.



$$\rho = -p$$

$$V(\varphi) = \text{const.}$$

$$\varphi$$

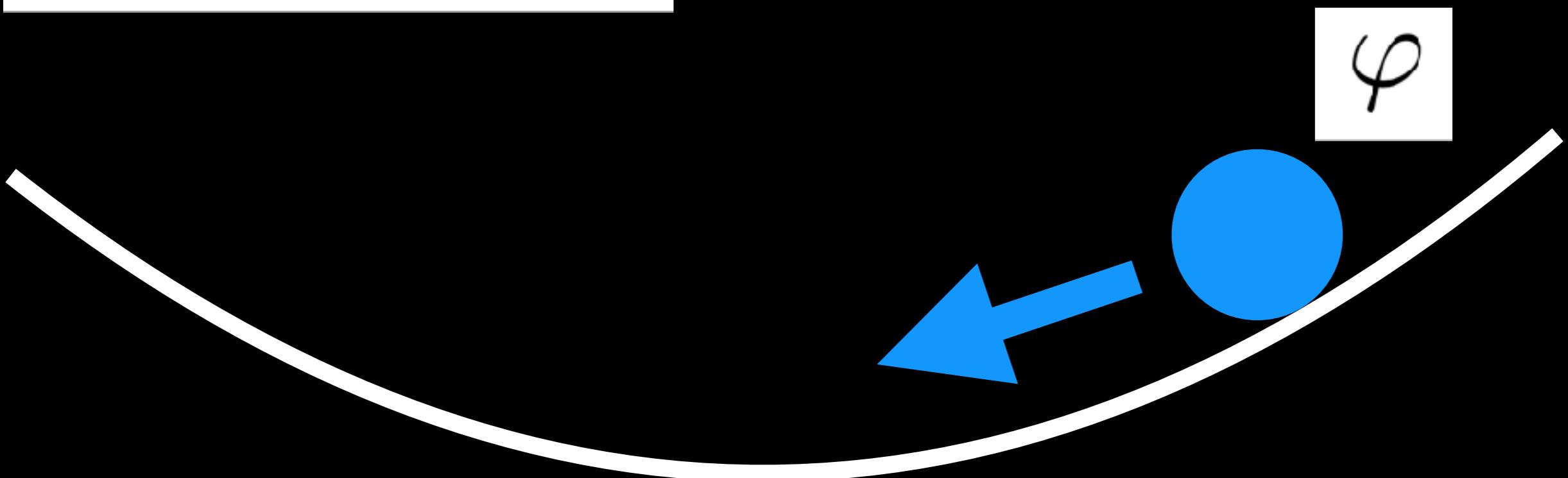


# The standard inflation scenario (Slow-roll single scalar field inflation)

If potential is **nearly flat** (constant):

$$\bar{\rho} = \frac{1}{2}\dot{\bar{\varphi}}^2 + V(\bar{\varphi}),$$

$$\bar{p} = \frac{1}{2}\dot{\bar{\varphi}}^2 - V(\bar{\varphi}),$$



# The standard inflation scenario **(Slow-roll single scalar field inflation)**

- Nearly scale-free power spectrum

$$\frac{k^3}{2\pi^2} P_{\delta\varphi}(k) \propto k^{n_s - 1}$$

- The spectral index is characterized by the derivative of potential:

$$n_s - 1 = 2\eta - 6\varepsilon \quad \eta = \frac{V''}{V} \quad \varepsilon = \frac{1}{2} \left( \frac{V'}{V} \right)^2$$

- Departure from the scale-free:

Detection!

$n_s = 0.9652 \pm 0.0047$

Planck 2015

# The standard inflation scenario (Slow-roll single scalar field inflation)

- Nearly Gaussian:

$$P(\delta\varphi) \sim e^{iS(\delta\varphi)}$$

$$S = S(\delta\varphi^2) + S(\delta\varphi^3) + \dots$$

Linear

Non-Linear

- Departure from the Gaussian:

No detection

$$f_{\text{NL}}^{\text{local}} = 0.8 \pm 5.0$$

# The standard inflation scenario **(Slow-roll single scalar field inflation)**

- Adiabatic condition

$$S_{m\gamma} = \frac{1}{3}\delta_m - \frac{1}{4}\delta_\gamma = 0$$

- Multi-field inflation scenarios can break this
- Departure from adiabatic condition  
(Existence of isogurvature modes):

No detection

# The standard inflation scenario **(Slow-roll single scalar field inflation)**

- Parity-symmetry

$$P_{\delta\varphi}(\vec{k}) = P_{\delta\varphi}(-\vec{k})$$

- Departure from the parity-symmetry:

No detection

# The standard inflation scenario **(Slow-roll single scalar field inflation)**

- Translational invariance:

$$\langle \delta\varphi(\vec{k})\delta\varphi(\vec{k}') \rangle = (2\pi)^3 \delta_D(\vec{k} + \vec{k}') P_{\delta\varphi}(\vec{k})$$

- Departure from the translational invariance:

No detection

# The standard inflation scenario (Slow-roll single scalar field inflation)

- **Rotational invariance:**

$$P_{\delta\varphi}(\vec{k}) = P_{\delta\varphi}(|\vec{k}|)$$

- Vector field inflation theories can break this.
- Departure from the rotational invariance:

No detection

# Quadrupolar anisotropy

$$P_\zeta(\vec{k}) = P_{\text{iso}}(k) \left[ 1 + g_* \underbrace{\left( \hat{k} \cdot \hat{p} \right)^2}_{\text{---}} \right]$$

The simplest model  
**breaking statistical isotropy**  
with preserving  
**party-symmetry** and  
**translational invariance.**

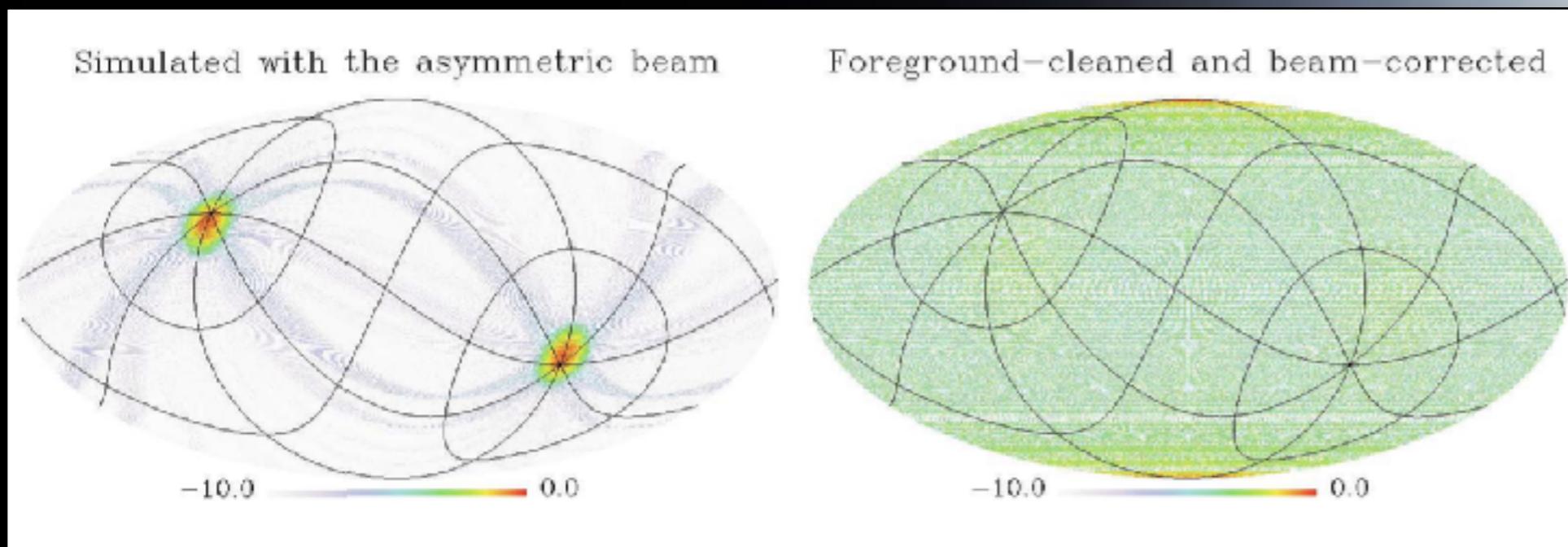


# CMB experiments

Planck provides

$$g_* = 0.002 \pm 0.016 \text{ (68%CL)}$$

(Kim and Komatsu 2014)

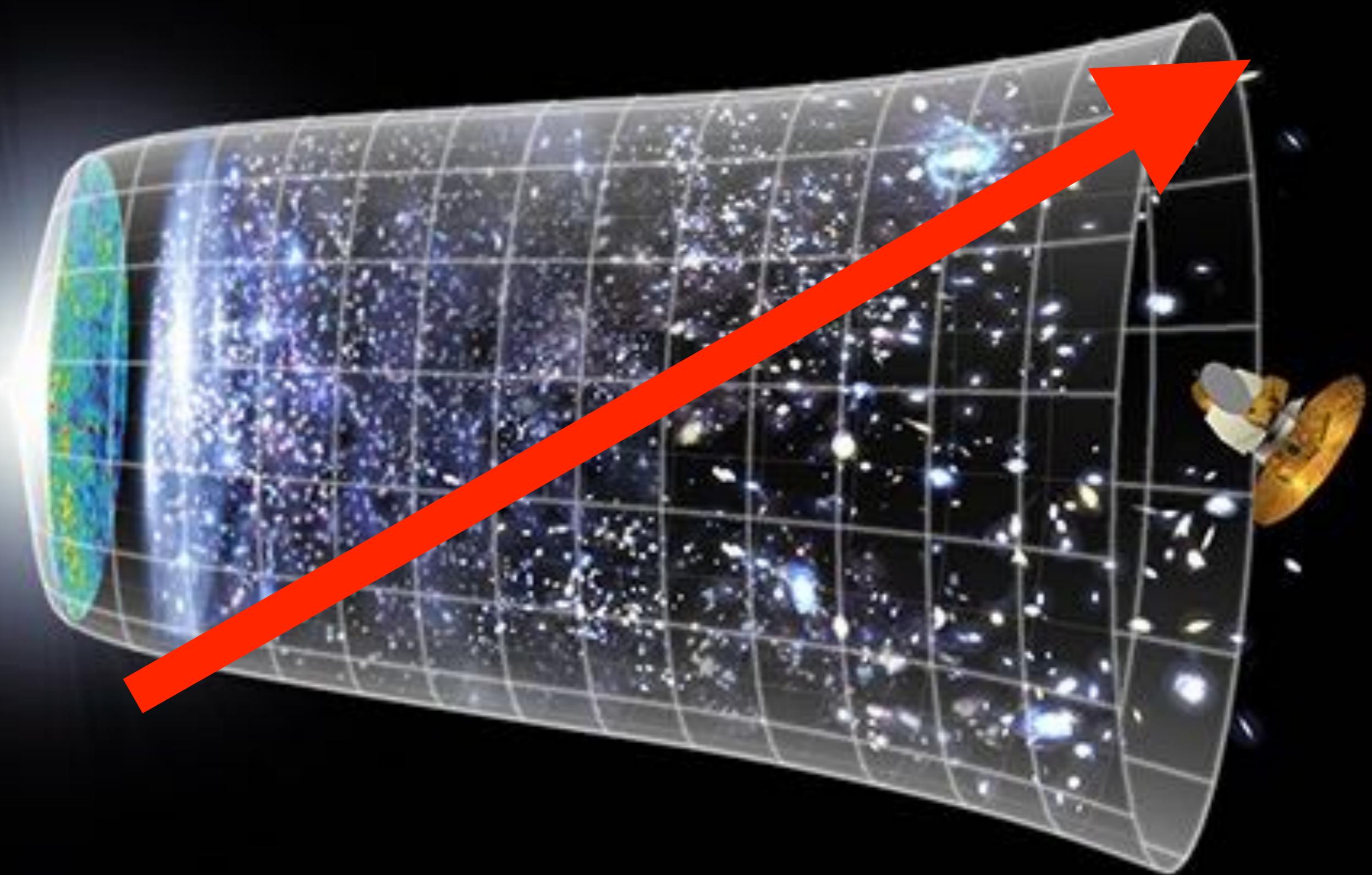


Before Beam Correction

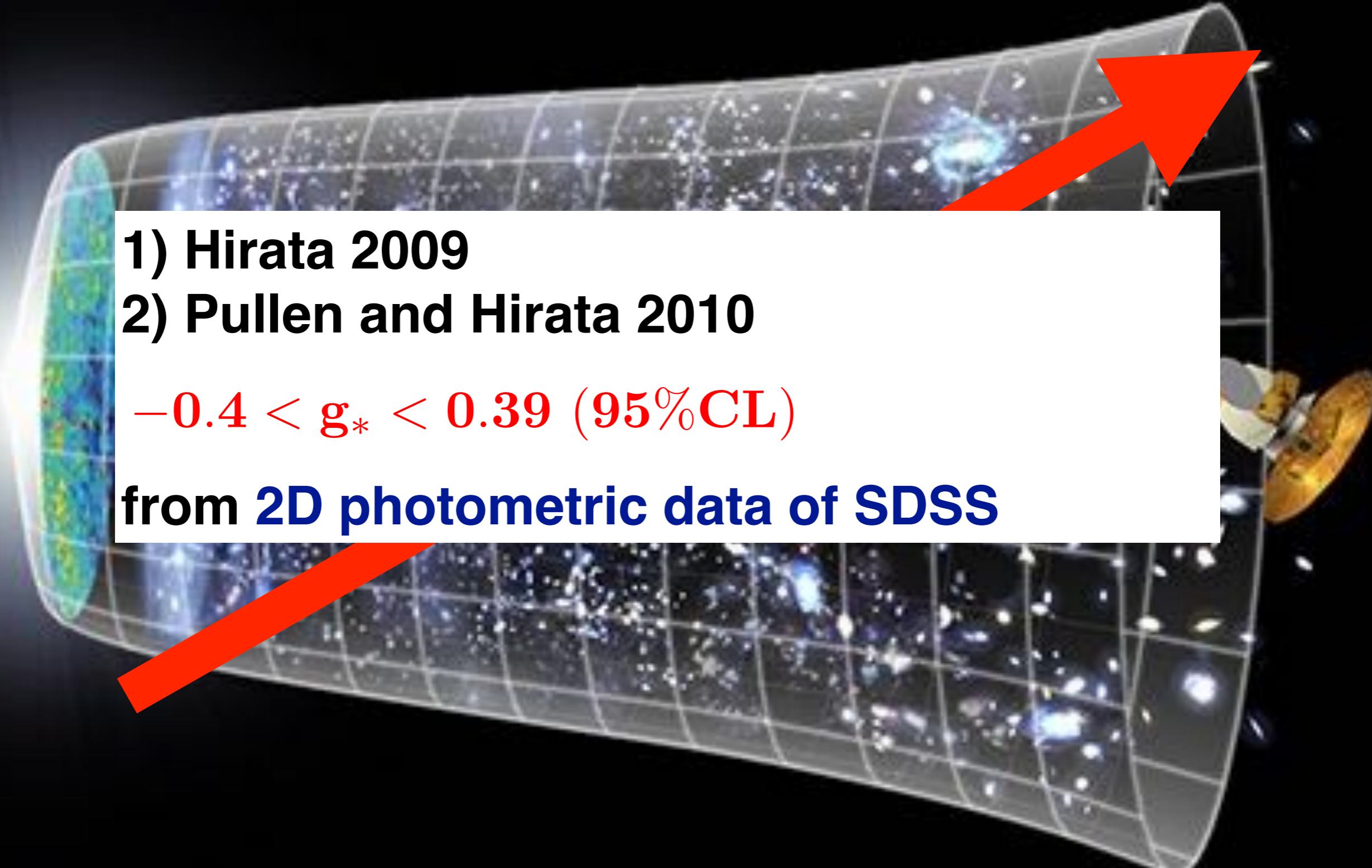
After Beam Correction



# Large-scale structure



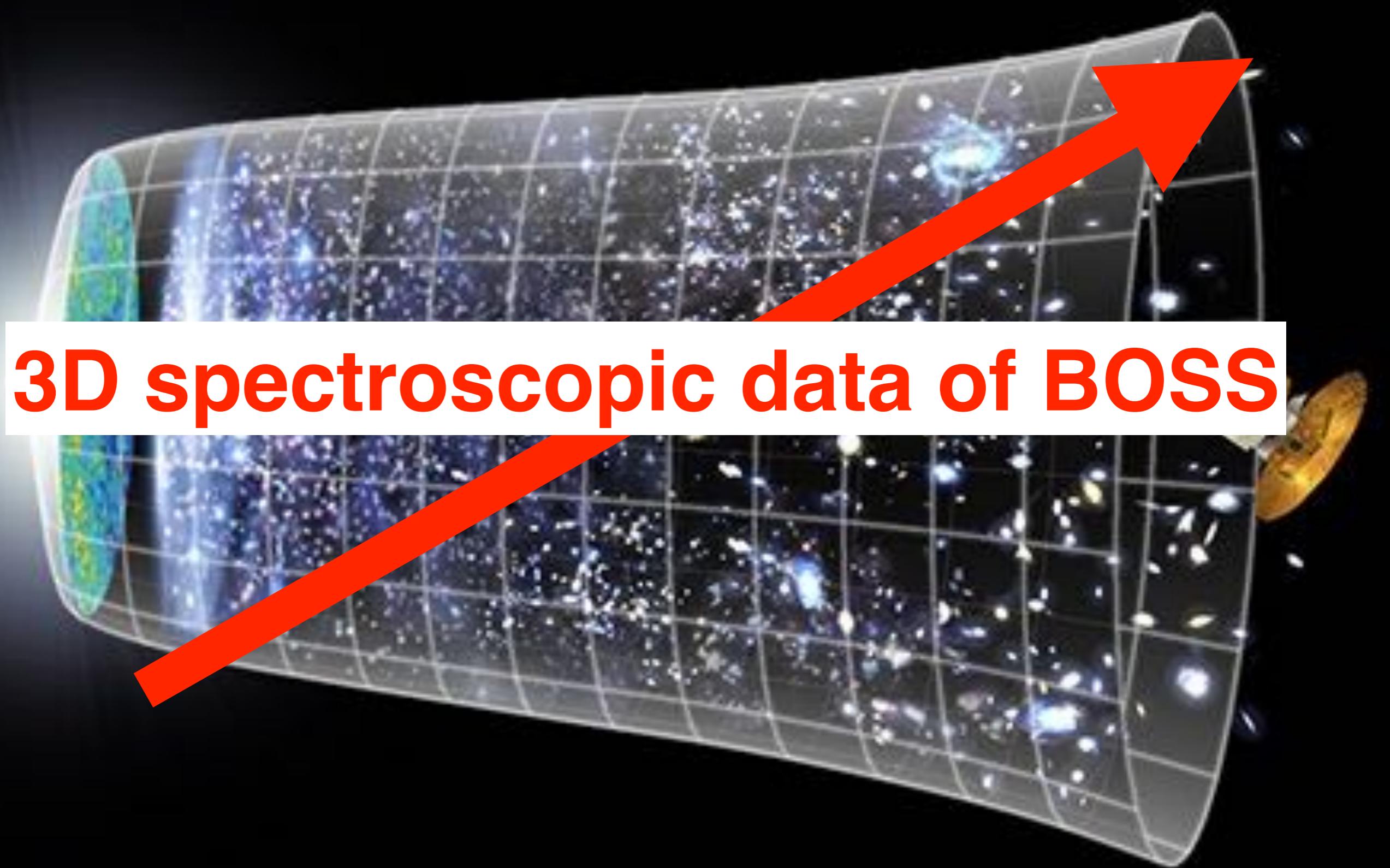
# Large-scale structure

- 
- 1) Hirata 2009
  - 2) Pullen and Hirata 2010

$$-0.4 < g_* < 0.39 \text{ (95%CL)}$$

from 2D photometric data of SDSS

# Large-scale structure



# Large-scale structure

**1, Theory** (M. Shiraishi, NS and T.Okumura, 2017)

- Bipolar spherical harmonics

**2, Measurement** (NS, M.Shiraishi and T.Okumura, 2017)

- BOSS
- Survey window corrections

**3, Future surveys**

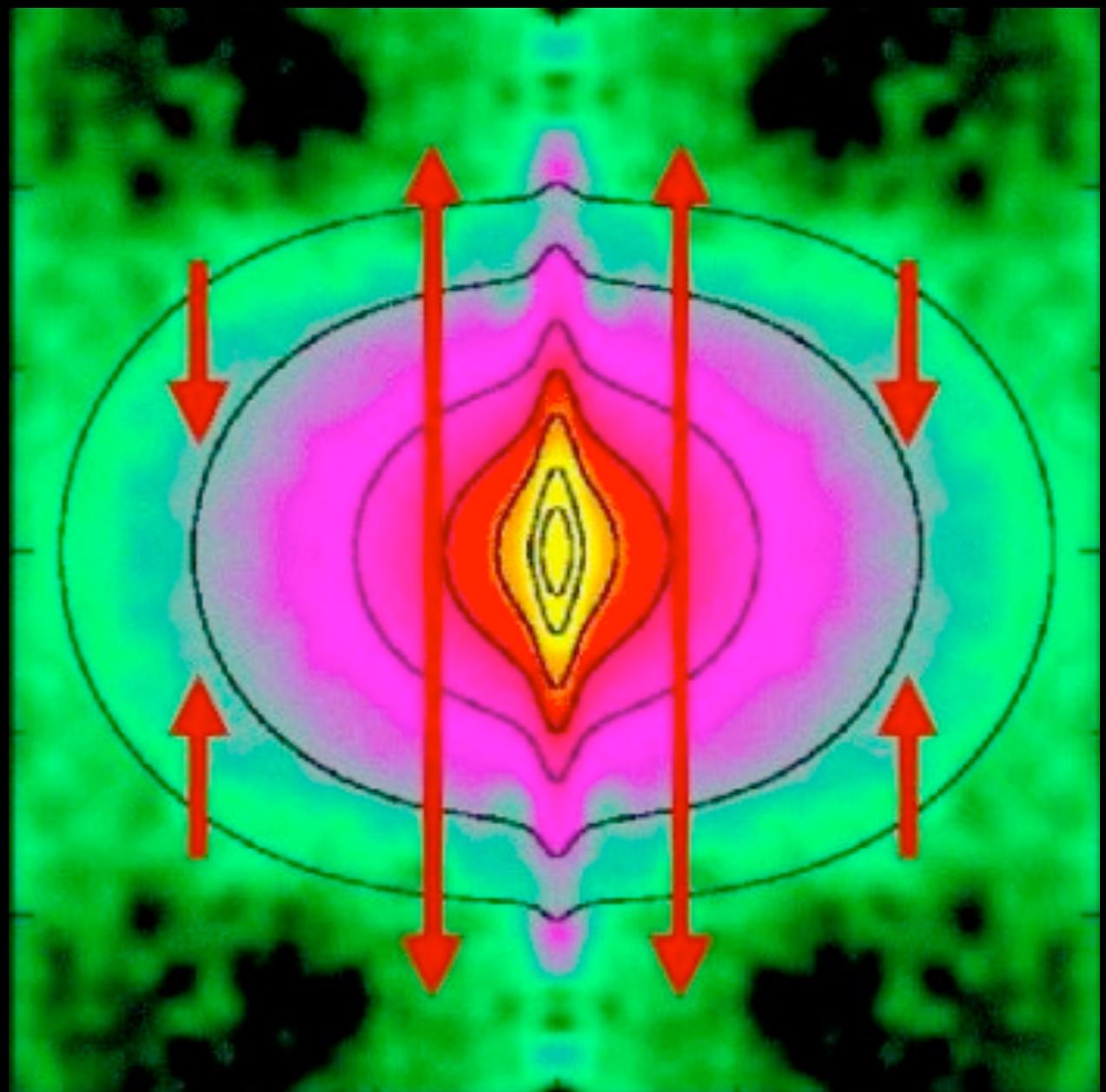
# Spectroscopic galaxy surveys

We observe:

- 1) Right ascension
- 2) Declination
- 3) Redshift

Redshift space distortions:

$$\begin{aligned}1 + z &= \frac{E_e}{E_o} \\&\approx 1 + \bar{z} + \frac{\vec{v} \cdot \hat{n}}{c}\end{aligned}$$



The observed radial distance of galaxies  
is distorted by peculiar velocities.

$$P_{\rm g}(\vec{k},\,\hat{n})$$

**Under the assumption of statistical isotropy,  
Legendre polynomials decomposition:**

$$P_g(\vec{k}, \hat{n}) = \sum_{\ell} P_{\ell}(k) \mathcal{L}_{\ell}(\hat{k} \cdot \hat{n})$$

# Preferred direction

$$P_g(\vec{k}, \hat{n}, \hat{p})$$



# Quadrupolar anisotropy

Preferred direction

Linear theory:



$$P_g(\vec{k}, \hat{n}, \hat{p}) = \underbrace{\left( b + f(\hat{k} \cdot \hat{n})^2 \right)^2}_{\text{Kaiser factor}} \underbrace{\left\{ 1 + g_* \left[ (\hat{k} \cdot \hat{p})^2 - \frac{1}{3} \right] \right\}}_{\text{Legendre decomposition}} P_m(k)$$

Kaiser factor

The Legendre decomposition is not sufficient any more.

# Legendre decomposition

$$\int \frac{d^2\hat{n}}{4\pi} \int \frac{d^2\hat{k}}{4\pi} \underline{\mathcal{L}_\ell(\hat{k} \cdot \hat{n}) P_g(\vec{k}, \hat{n}, \hat{p})} =$$

**Linear theory:**

$$\begin{aligned} P_0^s(k) &= \left( b^2 + \frac{2}{3}bf + \frac{1}{5}f^2 \right) P_{\delta\delta}^{\text{lin}}(k), \\ P_2^s(k) &= \left( \frac{4}{3}bf + \frac{4}{7}f^2 \right) P_{\delta\delta}^{\text{lin}}(k), \\ P_4^s(k) &= \frac{8}{35}f^2 P_{\delta\delta}^{\text{lin}}(k). \end{aligned}$$

**No primordial anisotropic contributions!!**

# Spherical harmonic decomposition

$$P_g(\vec{k}, \hat{n}, \hat{p}) = \sum_{\ell m} \sum_{\ell' m'} P_{\ell \ell'}^{mm'}(k, \hat{p}) Y_{\ell m}(\hat{k}) Y_{\ell' m'}(\hat{n})$$

# Bipolar Spherical harmonic decomposition

$$P_g(\vec{k}, \hat{n}, \hat{p}) = \sum_{LM} \sum_{\ell\ell'} P_{\ell\ell'}^{LM}(k, \hat{p}) (Y_\ell(\hat{k}) \otimes Y_{\ell'}(\hat{n}))_{LM}$$

Angular  
momentum  
coupling

# Bipolar Spherical harmonic decomposition

$$P_{\ell\ell'}^{LM}(k) = \sum_{mm'} \begin{pmatrix} \ell & \ell' & L \\ m & m' & M \end{pmatrix} P_{\ell\ell'}^{mm'}(k)$$

Wigner 3j symbol

Angular momentum coupling

# Bipolar Spherical harmonic decomposition

$$P_{\ell\ell'}^{LM}(k) = \sum_{mm'} \begin{pmatrix} \ell & \ell' & L \\ m & m' & M \end{pmatrix} P_{\ell\ell'}^{mm'}(k)$$

Total angular momentum

$L = 0$  : Legendre polynomials

$L > 0$  : Statistical anisotropy

# Quadrupolar anisotropy

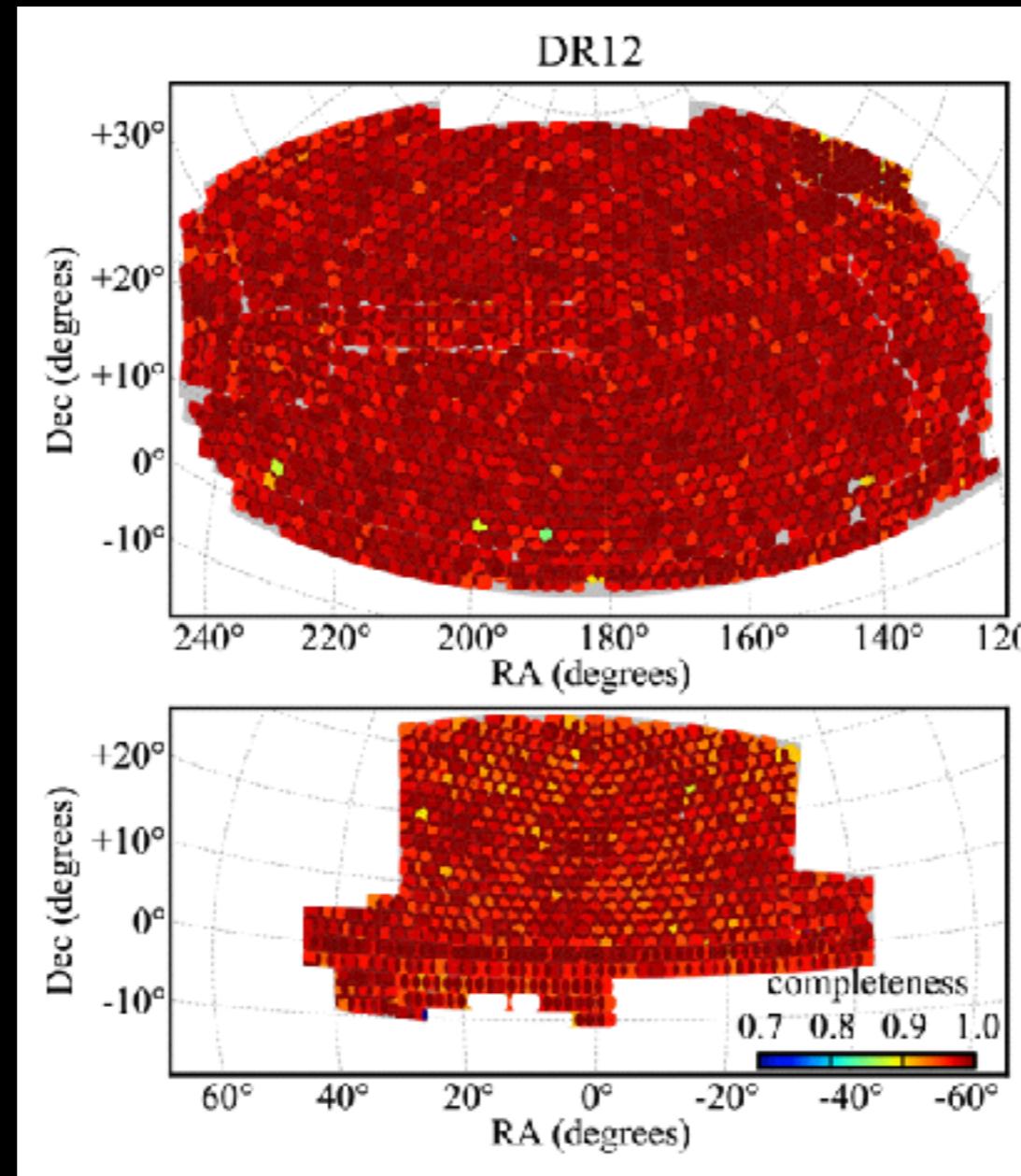
No primordial  
anisotropic contributions

$$P_{\ell\ell'}^{L=0,M=0}(k) = \delta_{\ell\ell'} P_\ell(k)$$

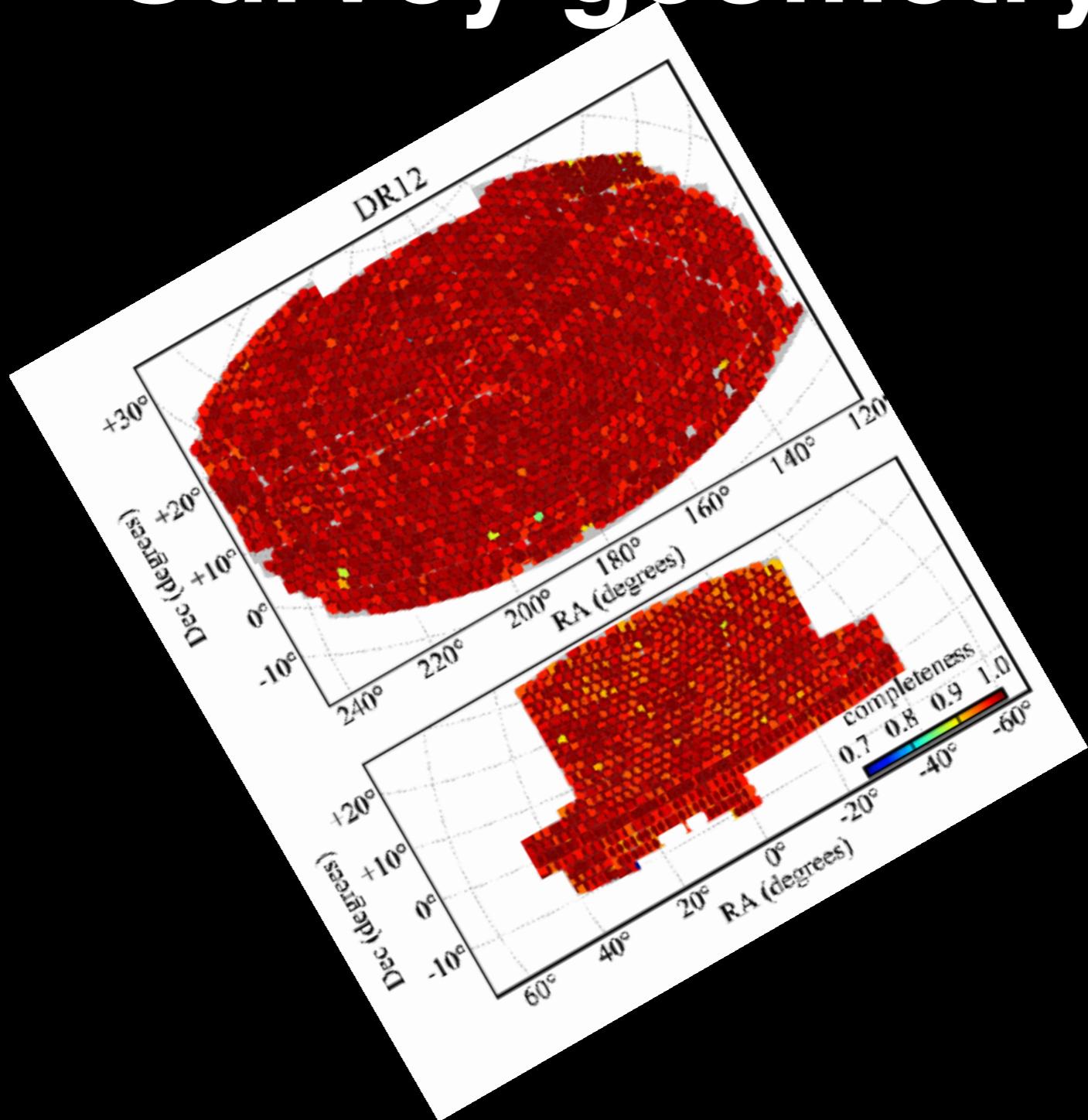


$$P_{\ell\ell'}^{L=2,M}(k) \propto g_* Y_{2M}(\hat{p})$$

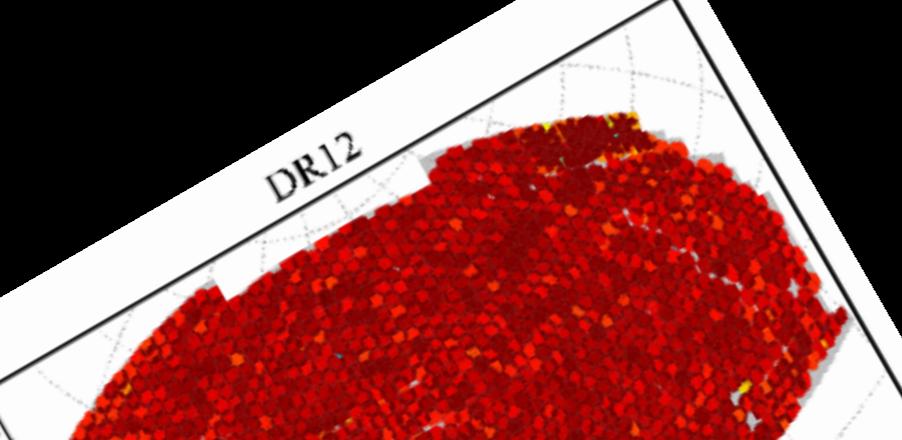
# Survey geometry



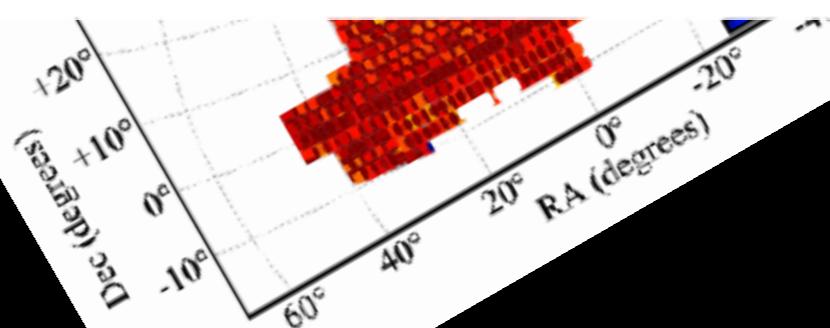
# Survey geometry



# Survey geometry

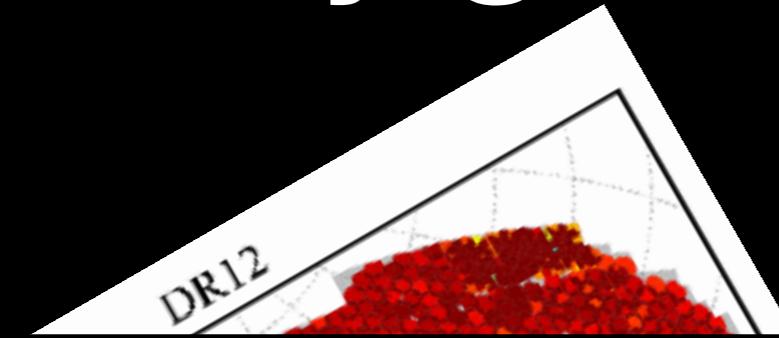


$$P_{\text{obs}} = Q * P_{\text{theory}}$$



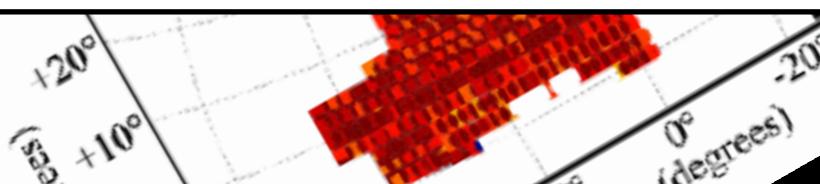
Q: window function

# Survey geometry



DR12

$$\begin{aligned}\langle \hat{\xi}_{20}^{2M}(r)|_A \rangle &= Q_0(r) \xi_{20}^{2M}(r) + Q_{20}^{2M}(r) [\xi_0(r) - \langle \bar{\delta}^2 \rangle] \\ &+ \frac{1}{5} [Q_{02}^{2M}(r) + Q_{22}^{2M}(r) + Q_{42}^{2M}(r)] \xi_2(r) \\ &+ \frac{1}{9} [Q_{24}^{2M}(r) + Q_{44}^{2M}(r)] \xi_4(r) \\ &+ \dots,\end{aligned}$$



**Bipolar Spherical harmonic  
decomposition of the window function**

# Survey geometry

## Anisotropic signal

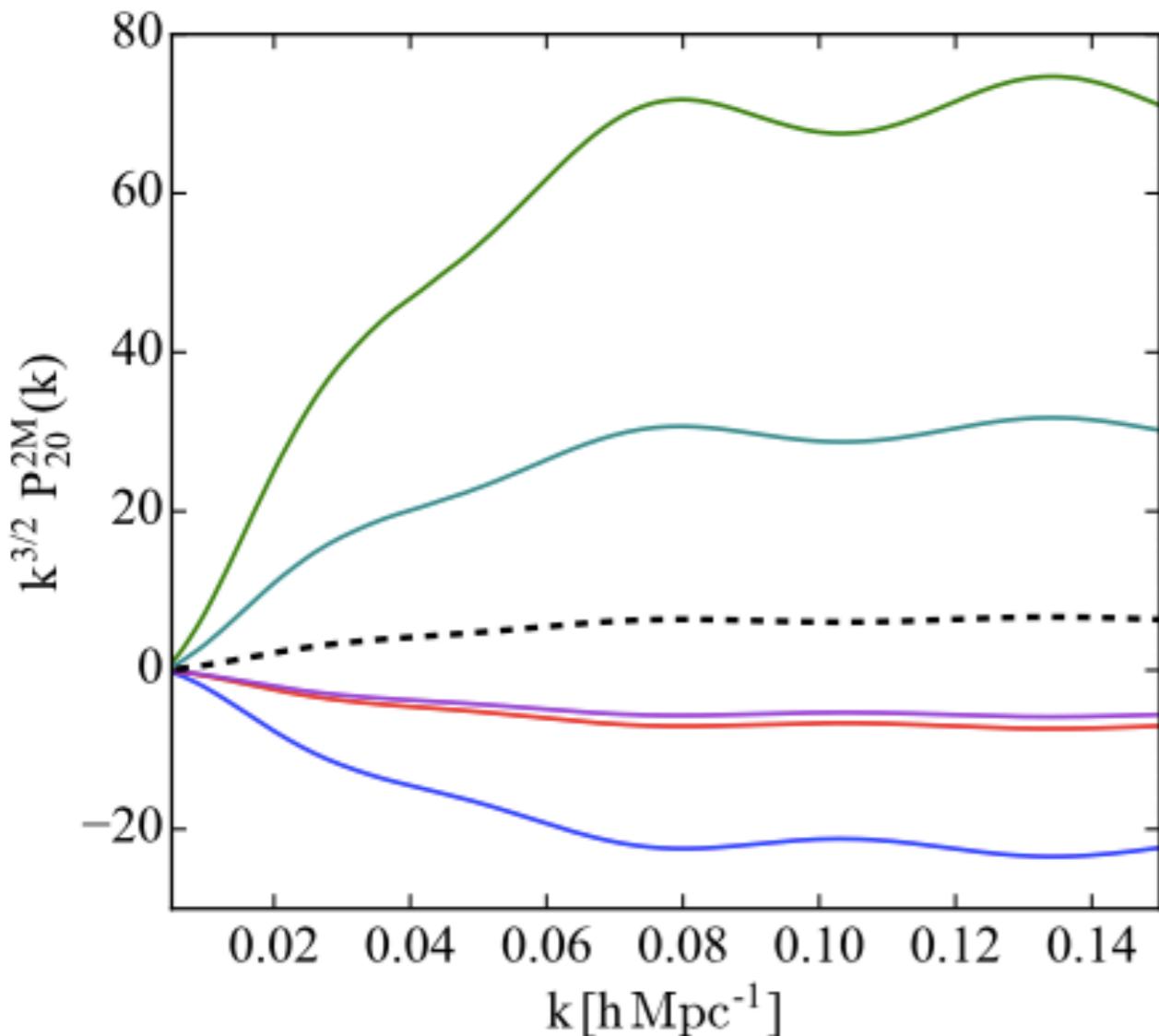
## Window effects

$$\begin{aligned}\langle \hat{\xi}_{20}^{2M}(r)|_A \rangle &= Q_0(r) \xi_{20}^{2M}(r) + Q_{20}^{2M}(r) [\xi_0(r) - \langle \bar{\delta}^2 \rangle] \\ &+ \frac{1}{5} [Q_{02}^{2M}(r) + Q_{22}^{2M}(r) + Q_{42}^{2M}(r)] \xi_2(r) \\ &+ \frac{1}{9} [Q_{24}^{2M}(r) + Q_{44}^{2M}(r)] \xi_4(r) \\ &+ \dots,\end{aligned}$$

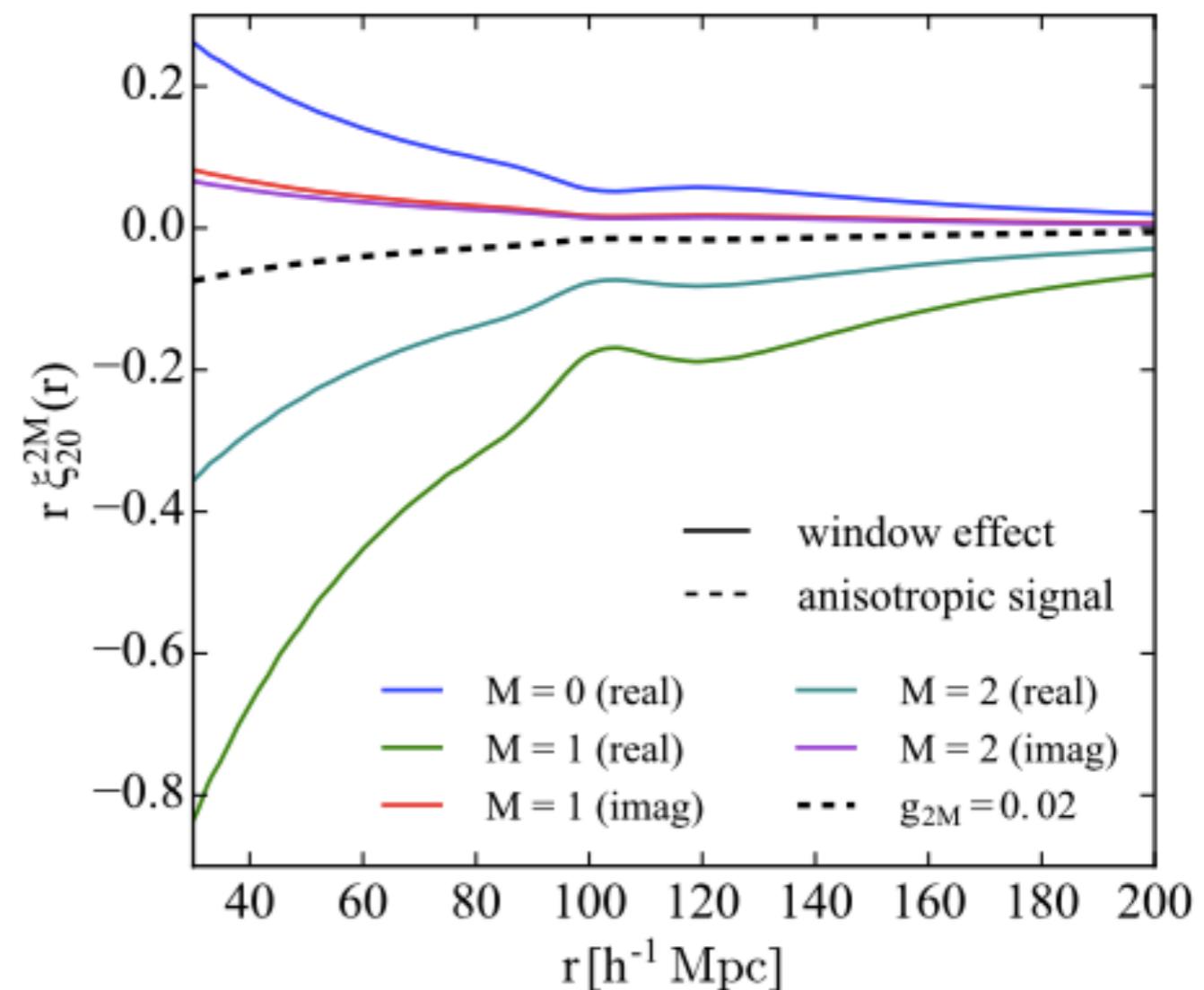
Bipolar Spherical harmonic  
decomposition of the window function

# Survey geometry

Power spectrum



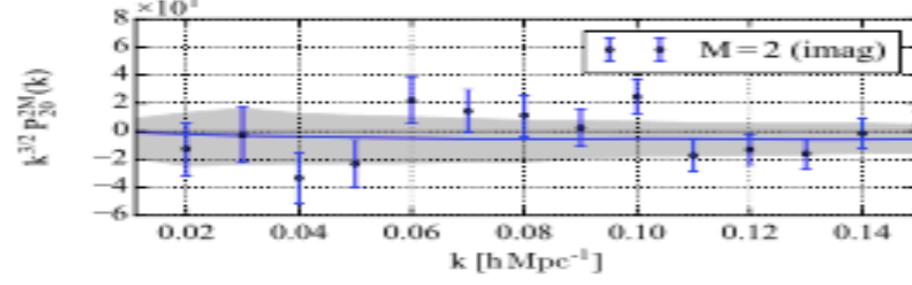
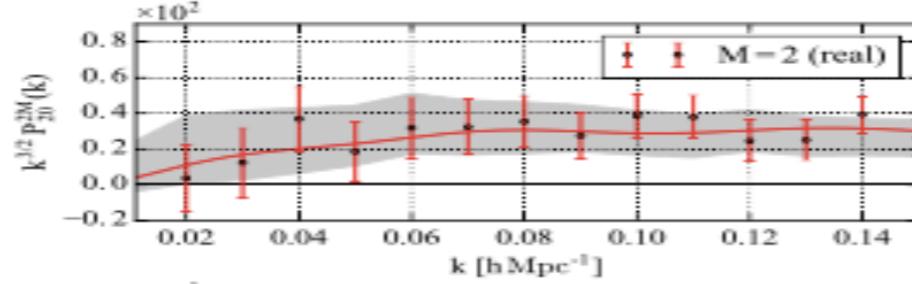
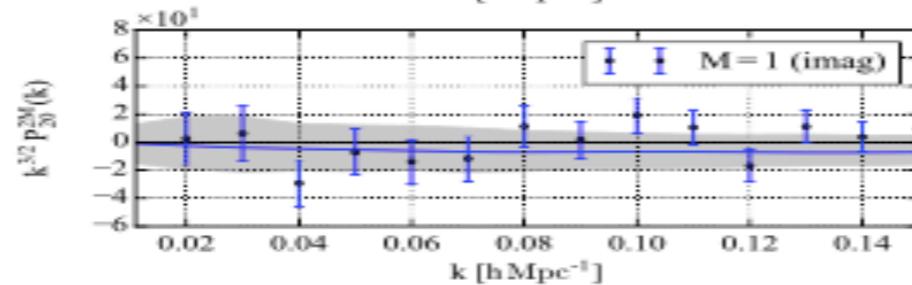
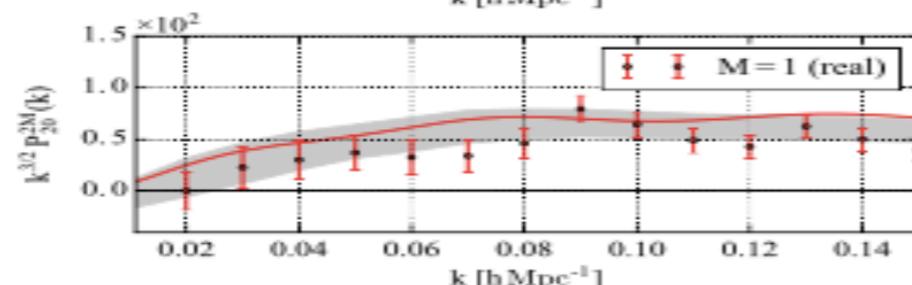
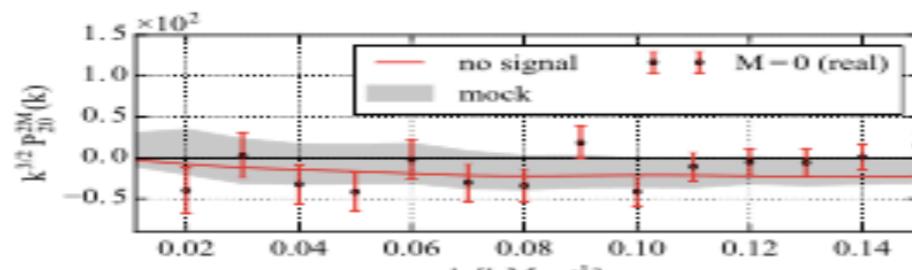
2pt correlation function



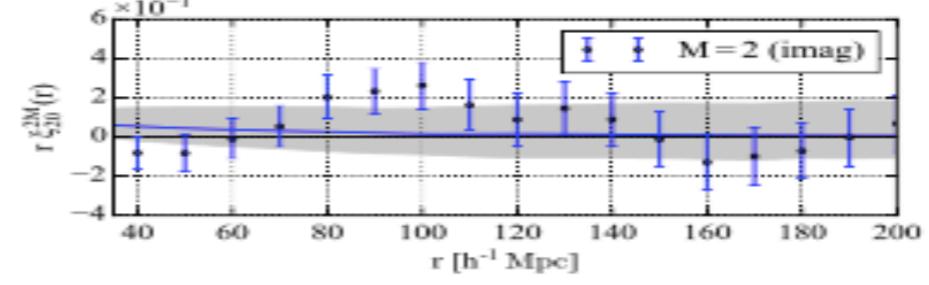
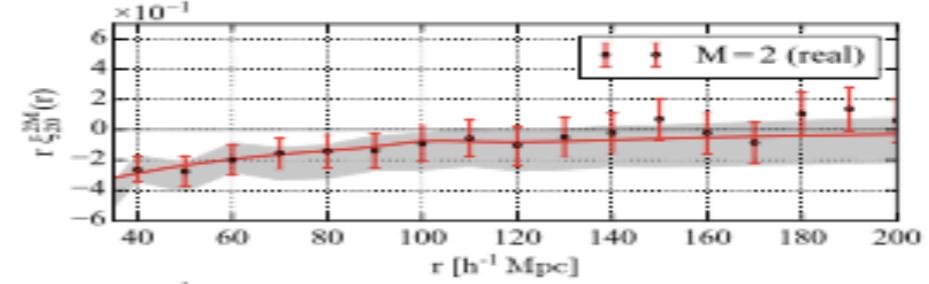
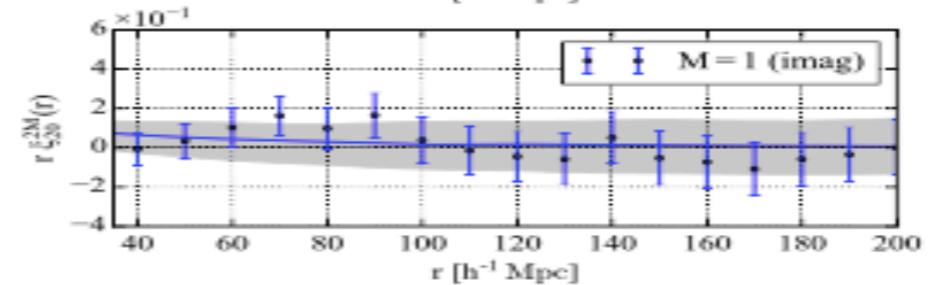
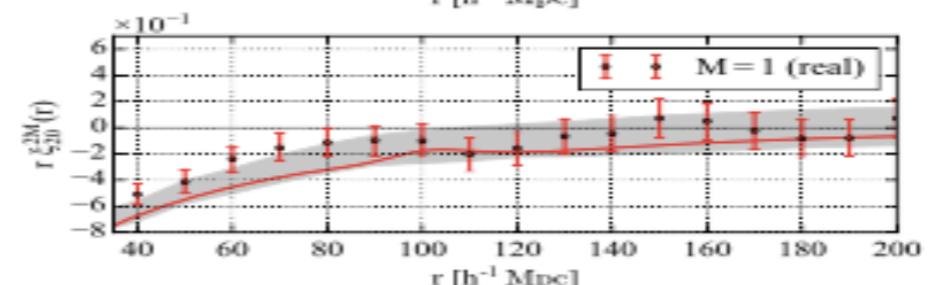
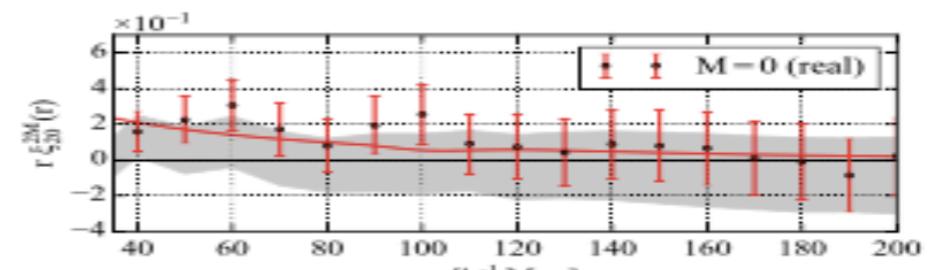
# L=2

# M=-2,1,0,1,2

## Power spectrum



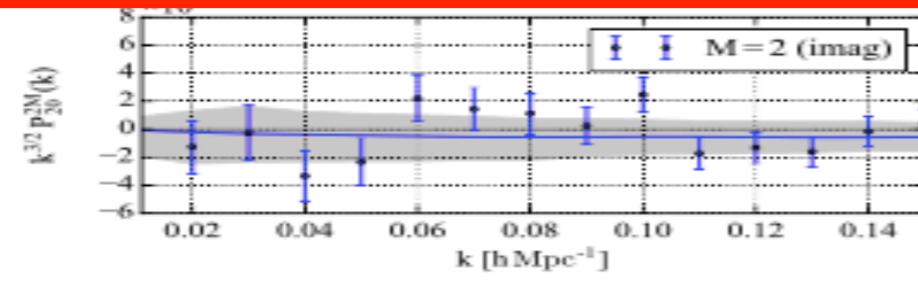
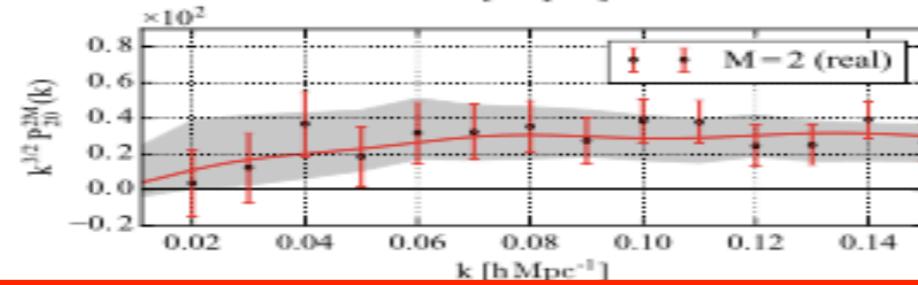
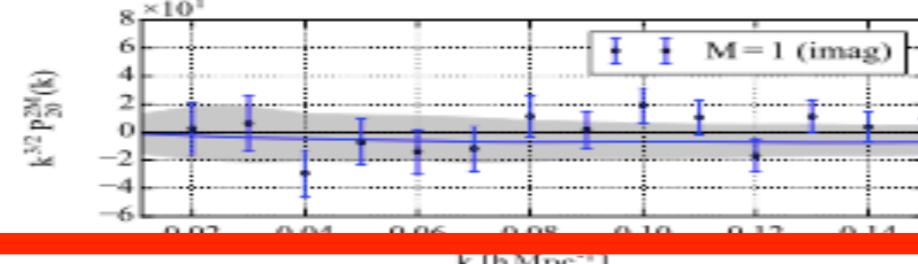
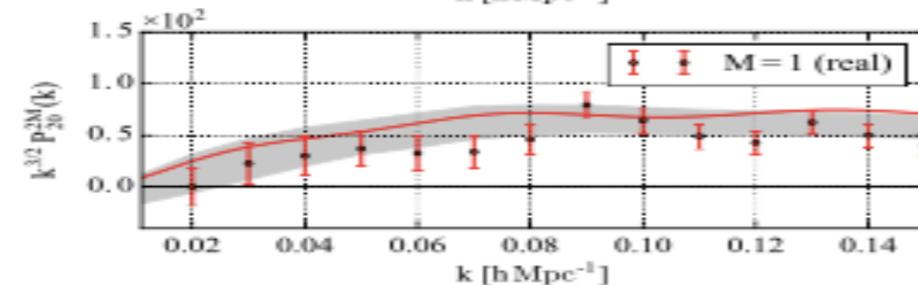
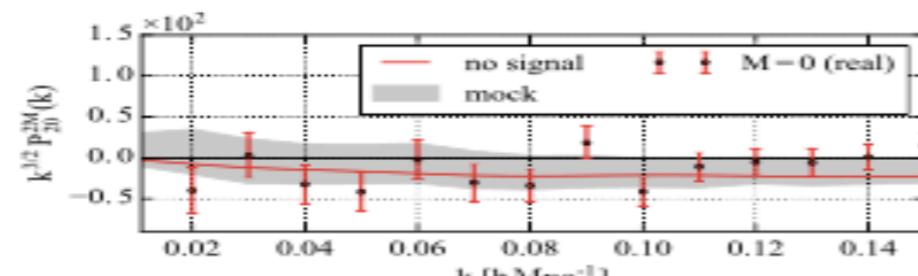
## 2pt function



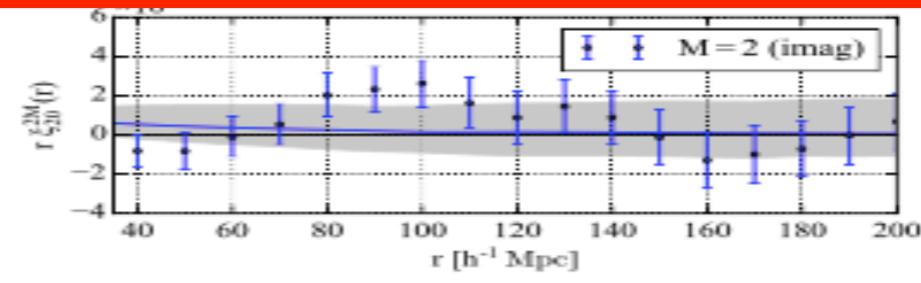
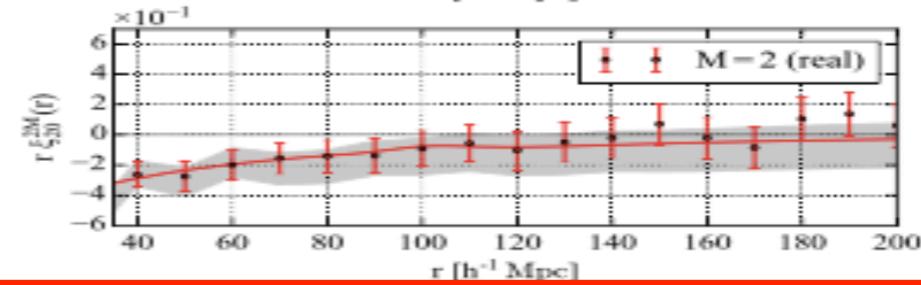
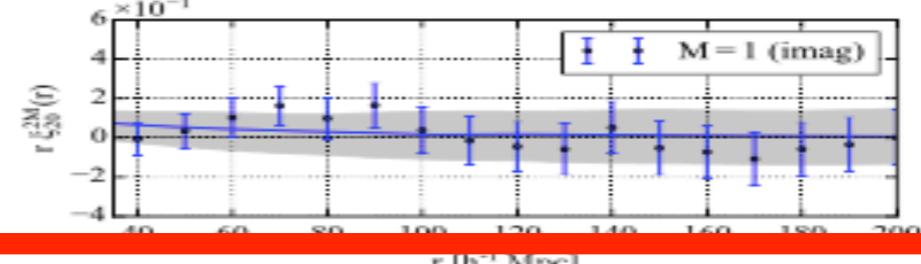
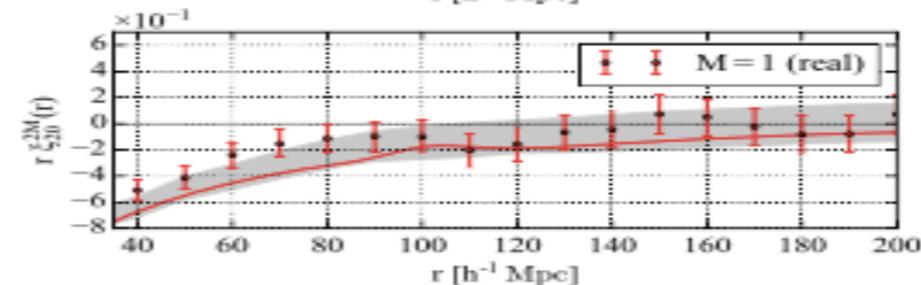
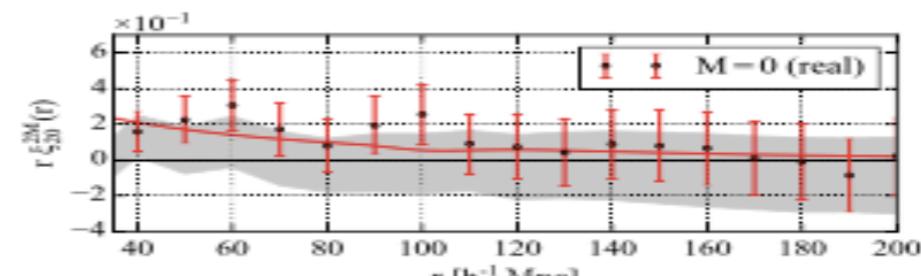
# L=2

# M=-2,1,0,1,2

## Power spectrum

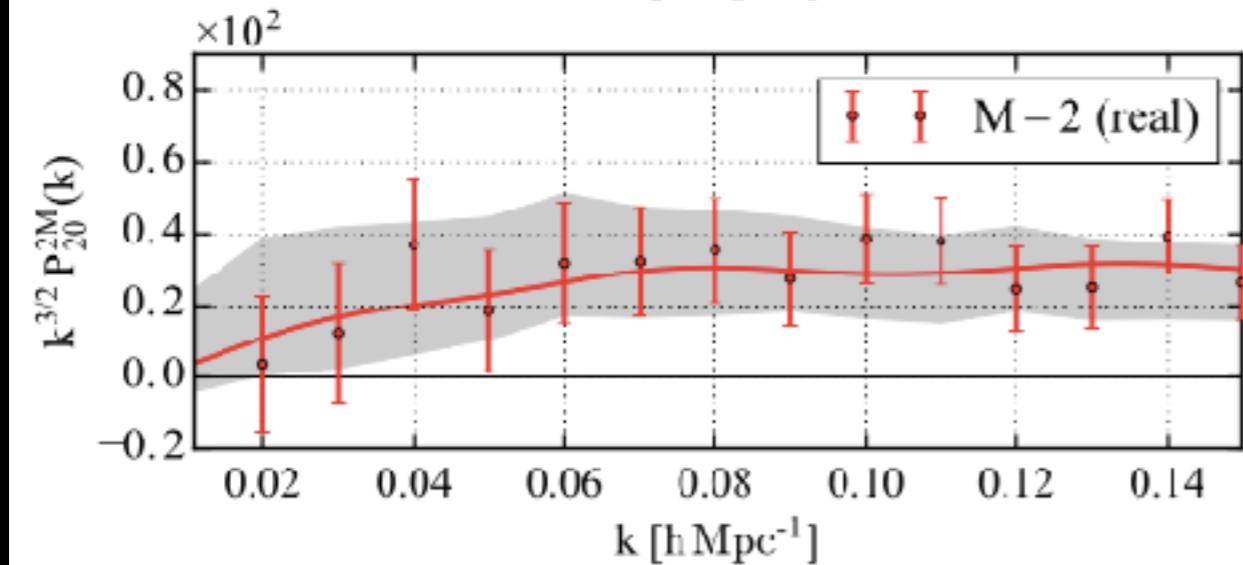


## 2pt function

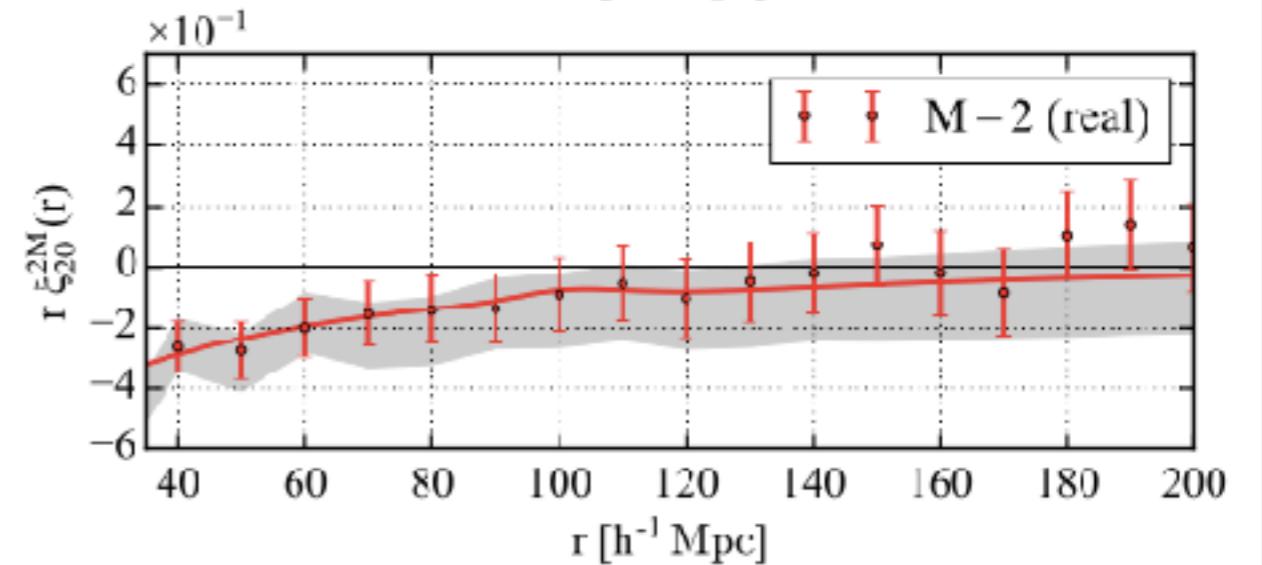


$$(L, M) = (2, 2)$$

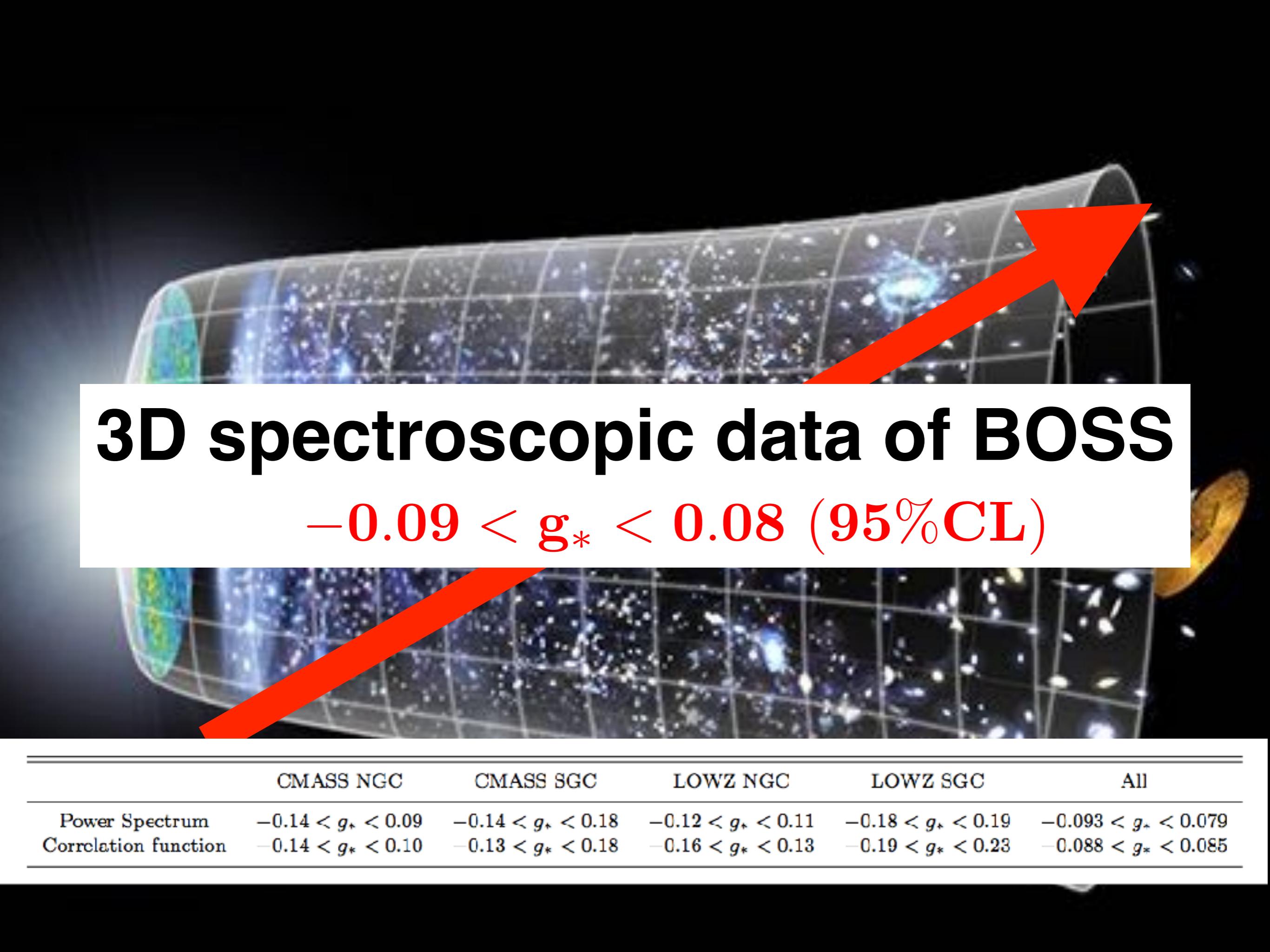
**Power spectrum**



**2pt correlation function**



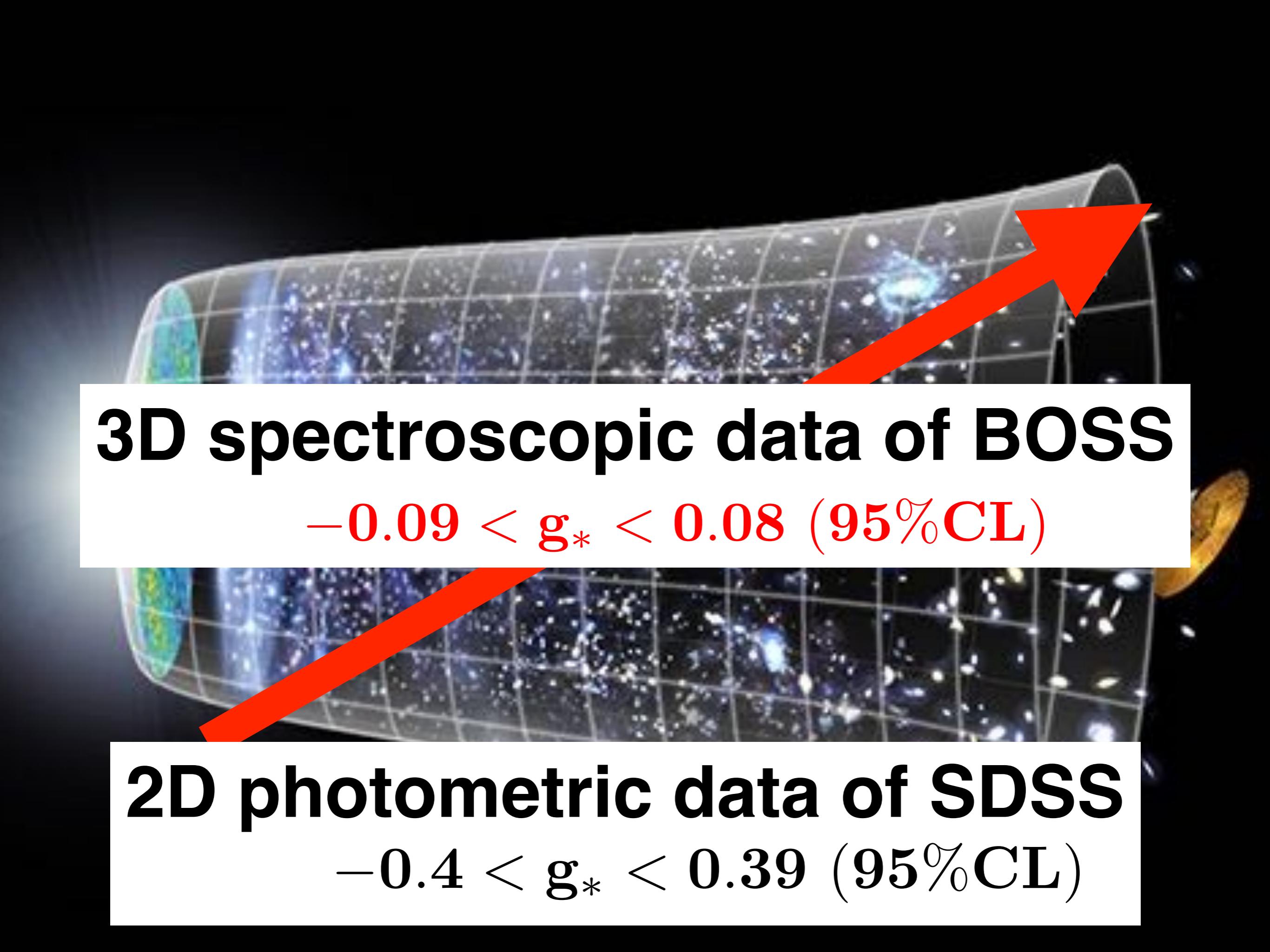
- Shaded regions:**  
mocks without statistical anisotropic signal
- Solid lines:**  
theory without statistical anisotropic signal
- Points:**  
observed data



# 3D spectroscopic data of BOSS

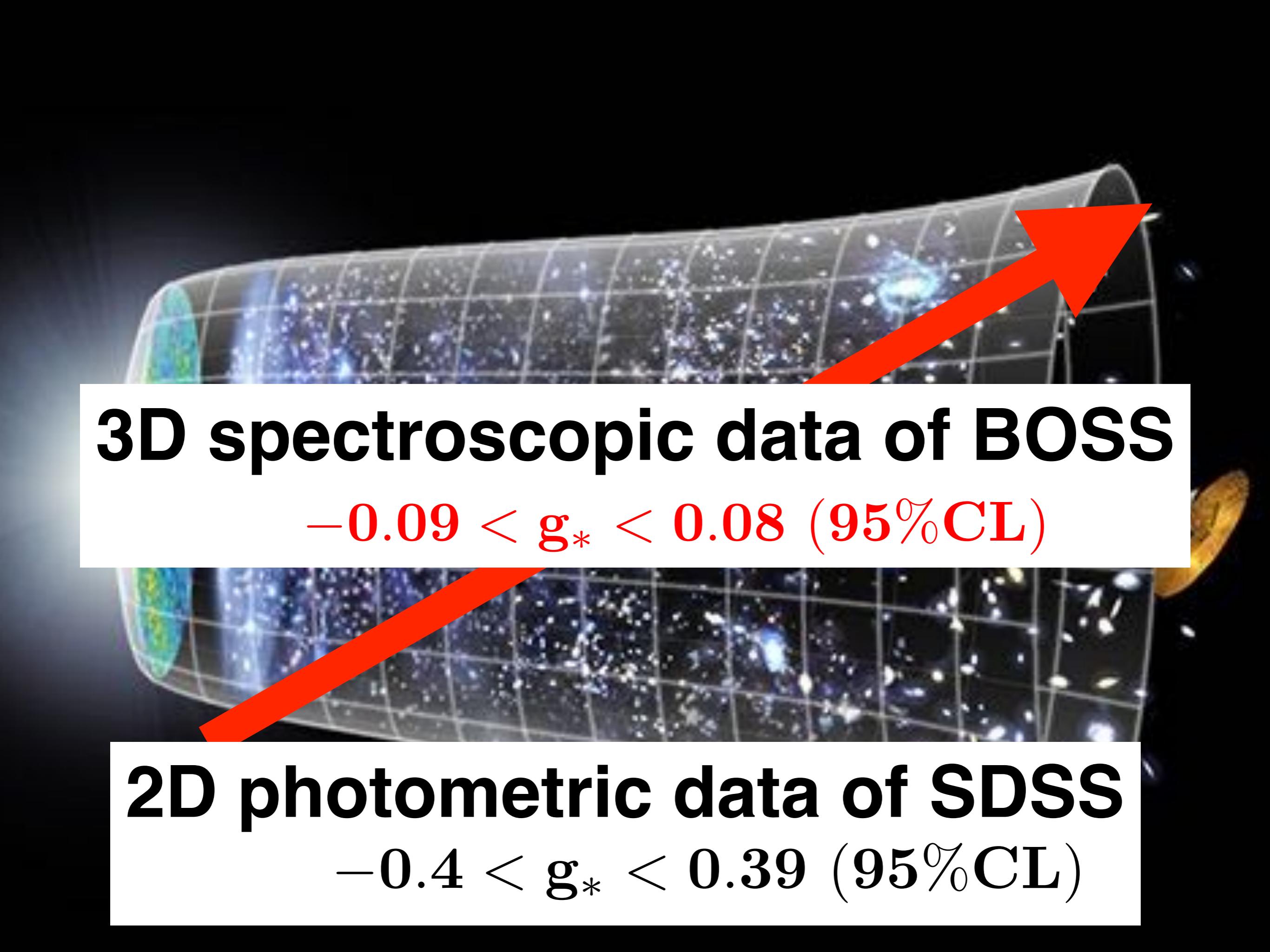
$-0.09 < g_* < 0.08$  (95%CL)

	CMASS NGC	CMASS SGC	LOWZ NGC	LOWZ SGC	All
Power Spectrum	$-0.14 < g_* < 0.09$	$-0.14 < g_* < 0.18$	$-0.12 < g_* < 0.11$	$-0.18 < g_* < 0.19$	$-0.093 < g_* < 0.079$
Correlation function	$-0.14 < g_* < 0.10$	$-0.13 < g_* < 0.18$	$-0.16 < g_* < 0.13$	$-0.19 < g_* < 0.23$	$-0.088 < g_* < 0.085$



## 3D spectroscopic data of BOSS

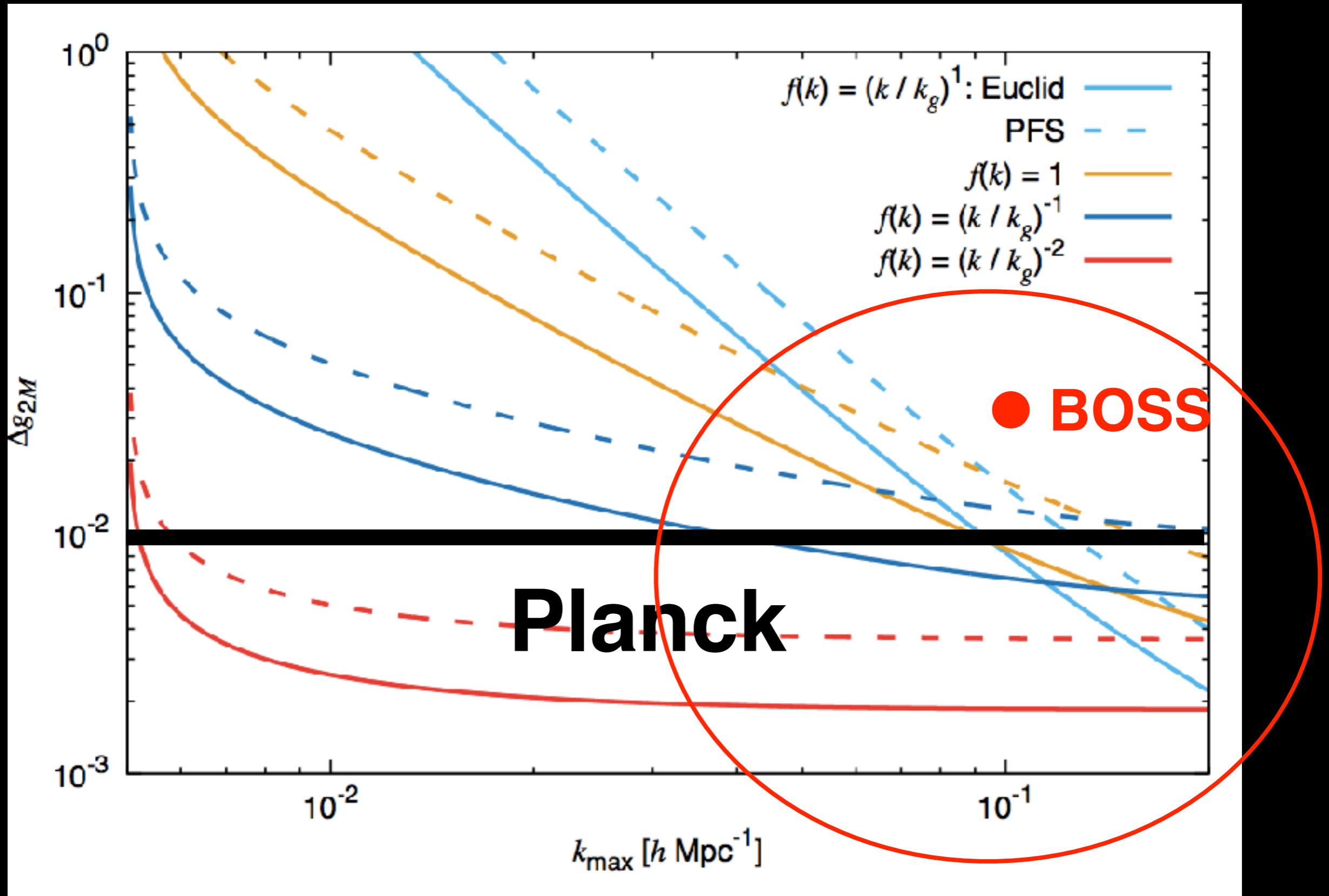
$$-0.09 < g_* < 0.08 \text{ (95\%CL)}$$



## 2D photometric data of SDSS

$$-0.4 < g_* < 0.39 \text{ (95\%CL)}$$

# Future surveys



# Summary

- First limits on statistical anisotropy from BOSS
- Bipolar spherical harmonic decomposition
- Estimator of the bipolar coefficients
- Survey geometry

**3D spectroscopic data of BOSS**

$$-0.09 < g_* < 0.08 \text{ (95\% CL)}$$

# **Extra slides**

# Non-linear information

