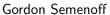

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Scattering with partial information







Laurent Chaurette

arxiv:1606.xxxxx

How can we understand (relativistic) scattering in quantum measurement theory?	the language

$$|0\rangle_A \otimes \sum_i c_i |i\rangle_S \xrightarrow{U} \sum_i c_i |i\rangle_A \otimes |i\rangle_S$$

$$|0\rangle_A \otimes \sum_i c_i \, |i\rangle_S \xrightarrow{U} \sum_i c_i \, |i\rangle_A \otimes |i\rangle_S$$

Scattering:

$$A = \phi_A = \text{apparatus field(s)}$$

 $S = \phi_S = \text{system field(s)}$

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$$U = S : \mathcal{H}^- \to \mathcal{H}^+ = S$$
-matrix

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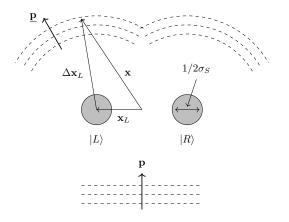
$$U = S \cdot \mathcal{H}^- \rightarrow \mathcal{H}^+ = S$$
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$$\mathcal{H}^{\pm}=\mathcal{H}^{\pm}_{m{A}}\otimes\mathcal{H}^{\pm}_{m{S}}$$

Questions

- 1. Given some property of S we want to measure, what do we observe with A?
- 2. How much A S entanglement does scattering generate?

1. Ex: verifying spatial superpositions



$$S = -\int d^4x \; rac{1}{2} (\partial_\mu \phi_S)^2 + rac{1}{2} (\partial_\mu \phi_A)^2 + rac{1}{2} m_S^2 \phi_S^2 + rac{1}{2} m_A^2 \phi_A^2 + rac{\lambda}{4} \phi_S^2 \phi_A^2$$

$$\rho = |\mathbf{p}^{-}\rangle \left\langle \mathbf{p}^{-}|_{A} \otimes \begin{pmatrix} \frac{1}{2} & \alpha \\ \alpha & \frac{1}{2} \end{pmatrix}_{S}$$

Convex family $0 < \alpha < 1$.

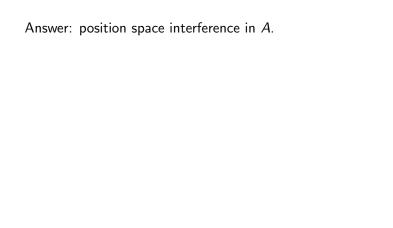
$$\alpha \to 1$$
: coherent superposition of $|L\rangle$, $|R\rangle$ $\alpha \to 0$: classical ensemble $|L\rangle$, $|R\rangle$

$$ho = \left| \mathbf{p}^- \right\rangle \left\langle \mathbf{p}^- \right|_A \otimes \begin{pmatrix} \frac{1}{2} & \alpha \\ \alpha & \frac{1}{2} \end{pmatrix}_{\mathbf{S}}$$

Convex family $0 \le \alpha \le 1$.

$$\alpha \to 1$$
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Question: How do we measure α (in perturbation theory in λ ?)



Answer: position space interference in A.

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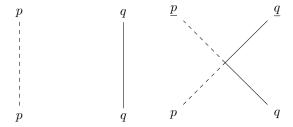
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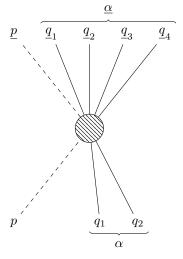
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2. Entanglement entropy generated by scattering



$$|\psi\rangle = |\mathbf{p}\rangle_{\mathbf{A}}^{-} \otimes |\alpha = \mathbf{q}_{1}, \mathbf{q}_{2}, \ldots\rangle_{\mathbf{S}}^{-}$$

$$\rho_{A}^{+}=\operatorname{tr}_{\mathcal{H}_{S}^{+}}\rho$$

$$ho_{\mathsf{A}}^+ = \operatorname{tr}_{\mathcal{H}_{\mathsf{S}}^+}
ho$$

$$I_0 = -2T \operatorname{Im} M_{\mathbf{p}\alpha\mathbf{p}\alpha}$$

$$F(\underline{\mathbf{p}}) = 2\pi T \sum_{\alpha} \left| M_{\underline{\mathbf{p}}\underline{\alpha}\mathbf{p}\alpha} \right|^{2} \delta_{\underline{\mathbf{p}}+\mathbf{p}_{\underline{\alpha}},\mathbf{p}+\mathbf{p}_{\alpha}} \delta(E_{\underline{\mathbf{p}}}^{A} + E_{\underline{\alpha}}^{S} - E_{\mathbf{p}}^{A} - E_{\alpha}^{S})$$

$$\rho_{A}^{+}=\operatorname{tr}_{\mathcal{H}_{S}^{+}}\rho$$

$$ho_A^+ = egin{pmatrix} 1 + I_0 + F(\mathbf{p}) & & & & & \\ & & F(\underline{\mathbf{p}}_1) & & & & \\ & & & F(\underline{\mathbf{p}}_2) & & & \\ & & & \ddots \end{pmatrix}$$

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 $F(\mathbf{p})$: "conditional total cross-section"

So, entanglement entropy:

$$S_A = -(1 + I_0 + F(\mathbf{p})) \ln(1 + I_0 + F(\mathbf{p})) - \sum F(\underline{\mathbf{p}}) \ln F(\underline{\mathbf{p}})$$

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Exact, non-perturbative. Due entirely to symmetry and assumption of single A particle.

Assuming perturbation theory holds,

$$S_A \approx -\sum_{\mathbf{p}} F(\underline{\mathbf{p}}) \ln F(\underline{\mathbf{p}})$$

Ex: $2 \rightarrow 2$ scattering in $\lambda \phi_A^2 \phi_S^2$ theory at lowest order.

(cf. Seki-Park-Sin 1412.7894, Peschanski-Seki 1602.00720)

$$iM_{\underline{pqpq}} = \frac{\underline{p}}{q} = \frac{i\lambda}{(2\pi)^3 \sqrt{16E_{\underline{p}}^A E_{\underline{q}}^S E_{\underline{p}}^A E_{\underline{q}}^S}}$$

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$$S_A = -\frac{T}{V} \frac{\lambda^2}{16\pi} \frac{p_{cm}(E^A + E^S)}{(E^A E^S)^2} \ln \left[\frac{T^2}{V^2} \frac{\lambda^2}{16(E^A E^S)^2} \right]$$

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- 3. Black holes: entropy in soft modes? String scattering?