

# quantum information metrics in QFT

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distinguishability between

$$\rho = \frac{1}{Z} e^{-\beta H}$$

$$\rho_\lambda = \frac{1}{Z_\lambda} e^{-\beta H + \lambda V}$$

result depends on **cutoff**

cutoff  
*vs*  
resolution

# QFT

local quantum theory in a continuous space

QFT

local quantum theory in a continuous space

REF

Q

R









$$\partial_x^2 \phi(x) + m^2 \phi(x) = 0$$

**finite resolution**

**arbitrary**  
**finite resolution**



$$\pi = \left\{ 3, \frac{31}{10}, \frac{314}{100}, \dots \right\}$$



$x_1, x_2, \dots \in \mathbb{Q}$  is Cauchy if

$\forall \epsilon > 0 \quad \exists N$  such that  $\forall i, j > N, \|x_i - x_j\| < \epsilon$

$\Lambda \mapsto \psi_\Lambda$  is “Cauchy” if

$\forall \sigma, \epsilon \quad \exists \Lambda_0$  such that  $\forall \Lambda, \Lambda' > \Lambda_0, d_\sigma(\psi_\Lambda, \psi_{\Lambda'}) < \epsilon$


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$\Lambda_0(\sigma, \epsilon)$

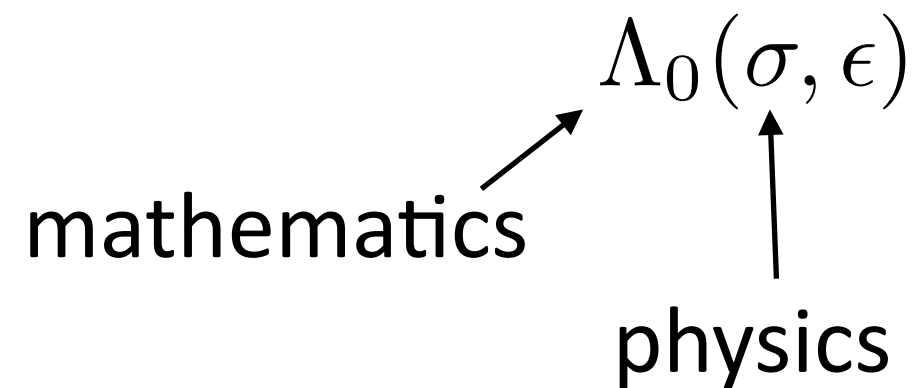
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mathematics   $\Lambda_0(\sigma, \epsilon)$

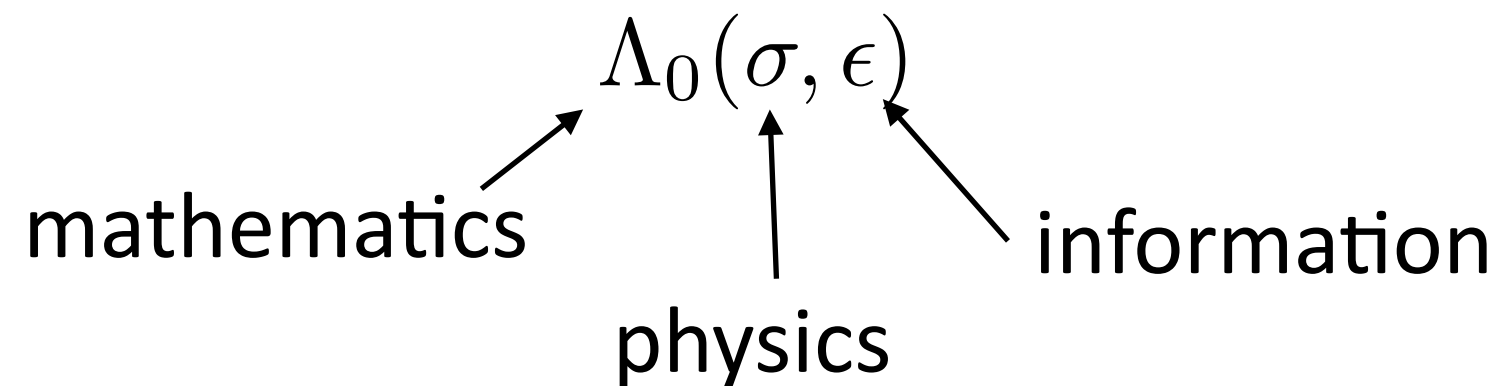
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$$d_{\sigma}(\psi_{\Lambda}, \psi_{\Lambda'})$$

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distinguishability

$$\rho_1 \otimes \cdots \otimes \rho_1$$

or

$$\rho_2 \otimes \cdots \otimes \rho_2$$

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$$\text{minimize} \quad \begin{cases} p_1 = \text{Tr}(P \rho_2^{\otimes N}) \\ p_2 = 1 - \text{Tr}(P \rho_1^{\otimes N}) \end{cases} \quad \text{over } P$$

$$\min_{p_2 = \epsilon} p_1 \approx e^{-NS(\rho_2 \parallel \rho_1)}$$

relative entropy  
(Kullback–Leibler divergence)

$$\min_{p_2 = \epsilon} p_1 \simeq e^{-NS(\rho_2 || \rho_1)}$$

$$S(\rho_2 || \rho_1) = \text{Tr}[\rho_2(\ln \rho_2 - \ln \rho_1)]$$

entropy

$$S(\rho) = \ln d - S(\rho \| \frac{1}{d} I)$$

entropy

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mutual information

$$I(A : B) = S(\rho_{AB} \| \rho_A \otimes \rho_B)$$

quantum channels



$\mathcal{N}(\rho)$

$$\mathcal{N}(\rho) = \rho$$

$$\mathcal{N}(\rho) = \rho \otimes \rho'$$

$$\mathcal{N}(\rho) = U (\rho \otimes \rho') U^\dagger$$

$$\mathcal{N}(\rho) = \text{Tr}' U (\rho \otimes \rho') U^\dagger$$

distinguishability  
must be monotonic

$$S(\mathcal{N}(\rho_2) \parallel \mathcal{N}(\rho_1)) \leq S(\rho_2 \parallel \rho_1)$$

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$$\text{Tr}(P\mathcal{N}(\rho)) = \text{Tr}(\mathcal{N}^\dagger(P)\rho)$$

# Information metrics



## Riemannian metrics on manifold of states

$$\mathcal{M} = \{\rho : \rho \geq 0, \text{Tr}\rho = 1\}$$

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$$\mathcal{M} = \{\rho : \rho \geq 0, \text{Tr}\rho = 1\}$$

which are monotonic

$$d(\mathcal{N}(\rho_1), \mathcal{N}(\rho_2)) \leq d(\rho_1, \rho_2)$$

for all channels  $\mathcal{N}$

classically: **unique**

$$\langle X, Y \rangle_\rho = \int \frac{X(x)Y(x)}{\rho(x)} dx$$

$$\rho + \epsilon X$$

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quantumly: classified by Petz

$$\langle X, Y \rangle_\rho = \text{Tr}(X \Omega_\rho^{-1}(Y))$$

$$\rho + \epsilon X$$

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$$\langle X, Y \rangle_\rho = \text{Tr}(X \Omega_\rho^{-1}(Y))$$

$$\Omega_\rho(Y) \sim \rho Y$$

## Kubo-Mori metric

$$S(\rho + \epsilon X \parallel \rho) = \epsilon^2 \text{Tr}(X \Omega_\rho^{-1}(X)) + \mathcal{O}(\epsilon^4)$$

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$$\langle X, X \rangle_\rho = \int_0^1 \text{Tr}(A \rho^s A \rho^{1-s}) ds = \int_0^1 \langle A(0) A(is) \rangle ds$$

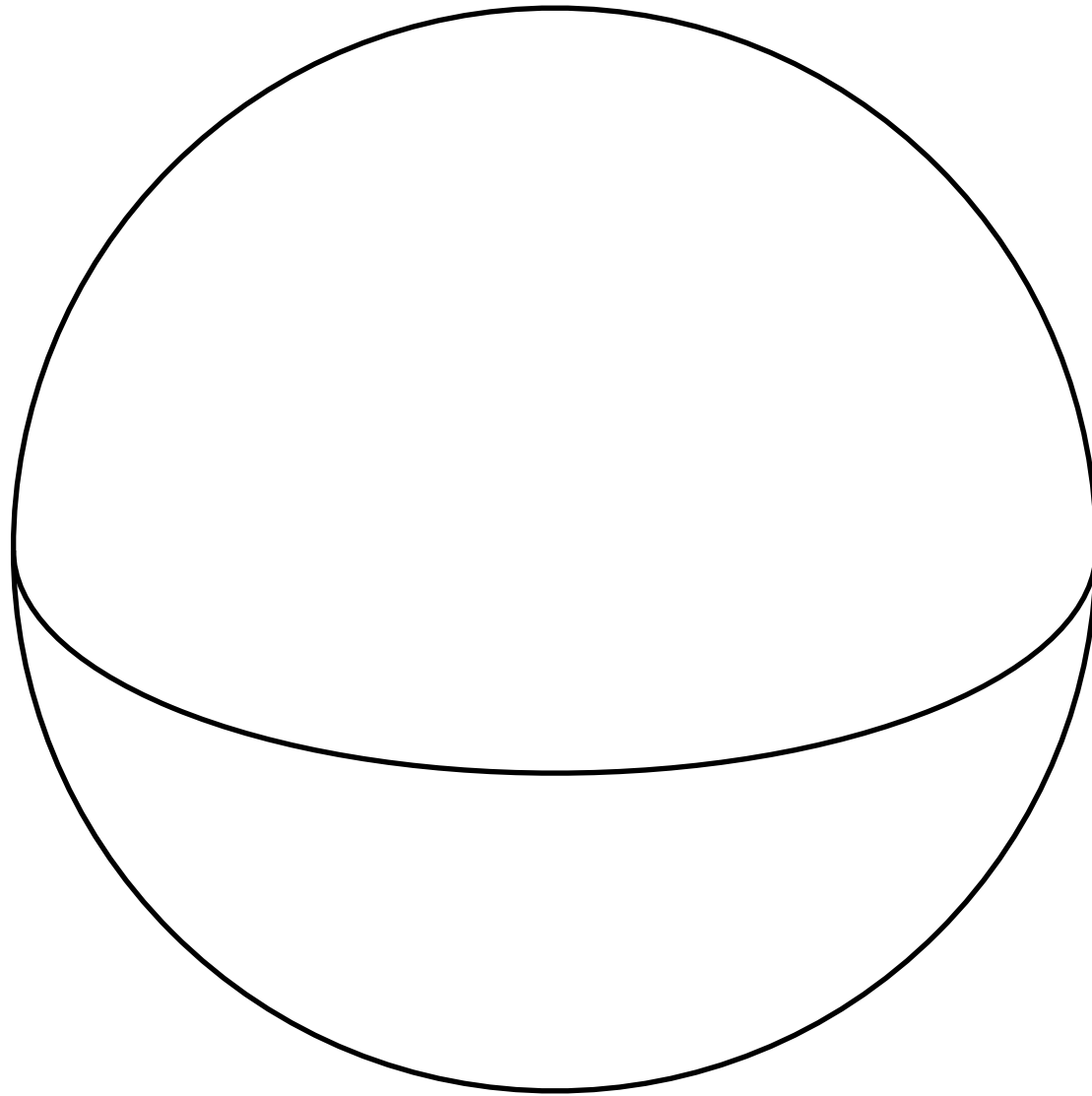
but  
this all depends on the cutoff...

coarse-grained  
distinguishability

$$d_{\sigma}(\rho_2, \rho_1) = d(\mathcal{N}_{\sigma}(\rho_2), \mathcal{N}_{\sigma}(\rho_1))$$

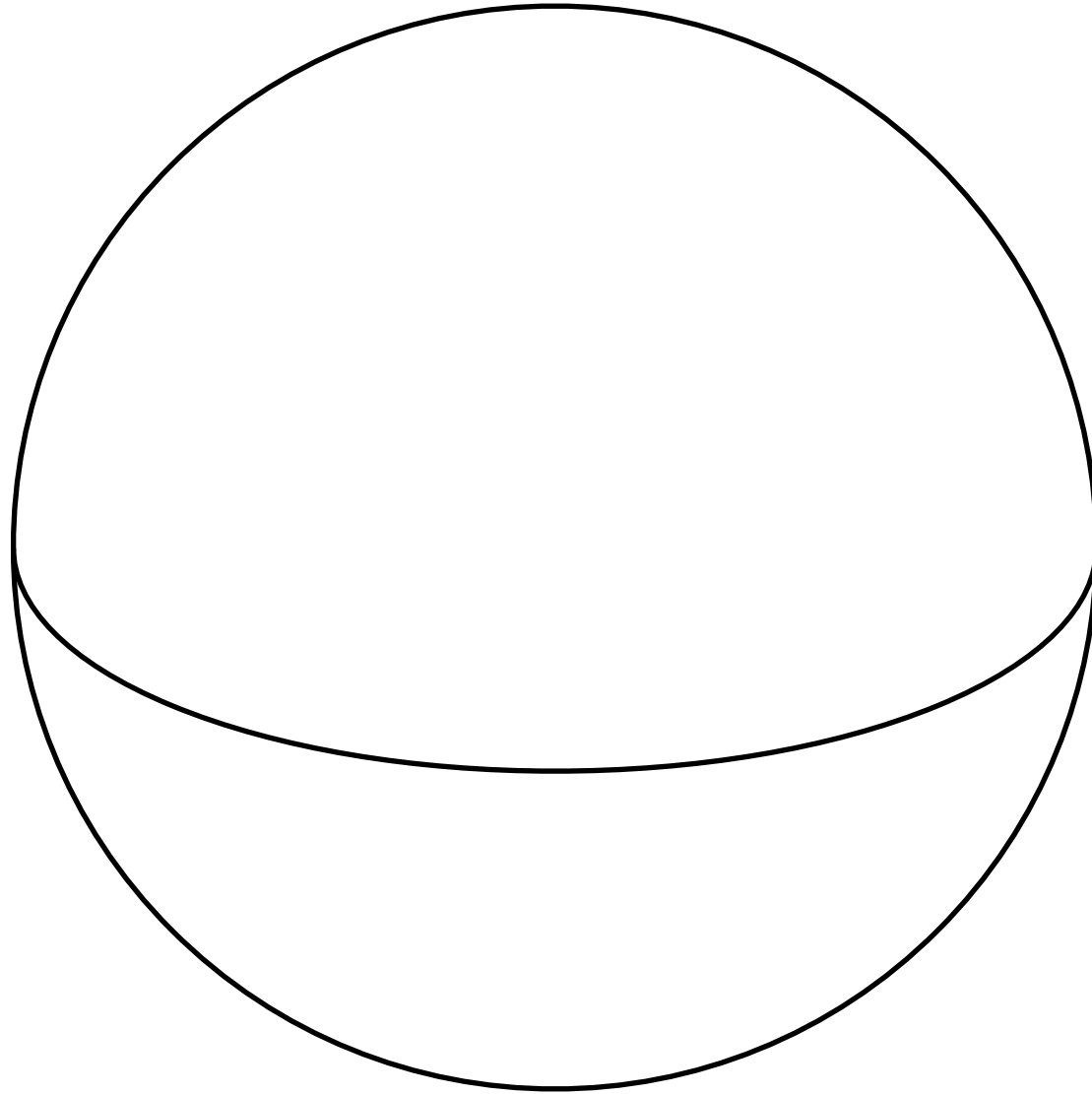
$$d_\sigma(\rho_2, \rho_1) = d(\mathcal{N}_\sigma(\rho_2), \mathcal{N}_\sigma(\rho_1))$$

$$\langle\langle X, Y \rangle\rangle_\rho^\sigma = \langle \mathcal{N}_\sigma(X), \mathcal{N}_\sigma(Y) \rangle_{\mathcal{N}_\sigma(\rho)}$$



$$\langle X, Y \rangle_{\rho}$$

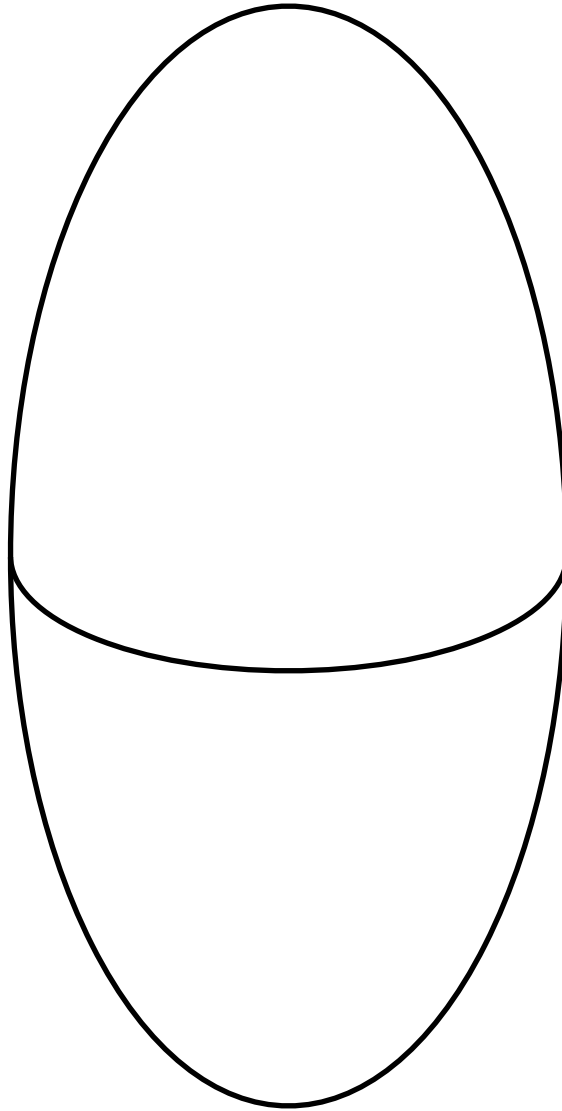
$$\sigma = 0 \mu m$$



$$\langle\langle X, Y \rangle\rangle_{\rho}^{\sigma}$$

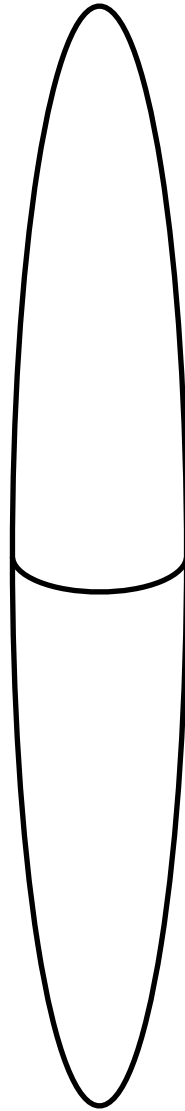


$$\sigma = 10 \mu m$$



$$\langle\langle X, Y \rangle\rangle_{\rho}^{\sigma}$$

$$\sigma = 100 \mu m$$



$$\langle\langle X, Y \rangle\rangle_{\rho}^{\sigma}$$

$$\sigma = 1000 \mu m$$



$$\langle\langle X, Y \rangle\rangle_{\rho}^{\sigma}$$

$$\rho_t = \frac{1}{Z_t} e^{-\beta(H+tV)}$$

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$$D(V) := \langle \partial_t \mathcal{N}_\sigma(\rho_t), \partial_t \mathcal{N}_\sigma(\rho_t) \rangle_{\mathcal{N}_\sigma(\rho_t)} \Big|_{t=0}$$

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$$d(V) = \lim_{\text{vol} \rightarrow \infty} \frac{1}{\text{vol}} D(V)$$

example

$f \in$  phase space

$$\text{Tr}(\rho e^{i\phi(f)}) = e^{-\frac{1}{2}(f, Af)}$$



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$$\text{Tr}(\rho e^{i\phi(f)}) = e^{-\frac{1}{2}(f, Af)}$$

$$\text{Tr}(\mathcal{N}(\rho) e^{i\phi(f)}) = e^{-\frac{1}{2}(f, (X^\dagger AX + Y)f)}$$

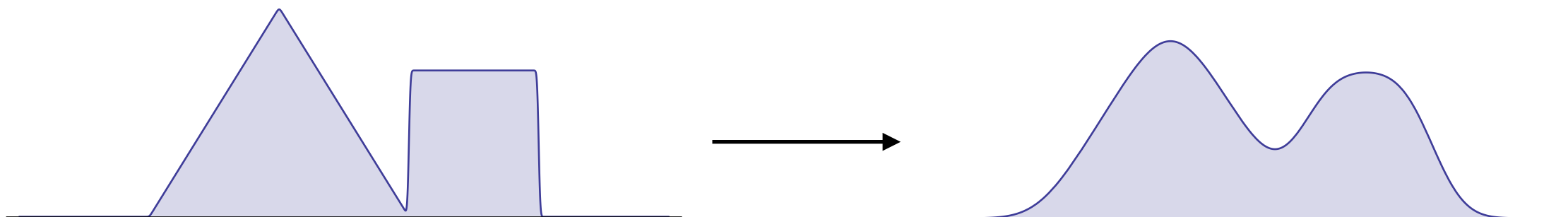
$f \in$  phase space

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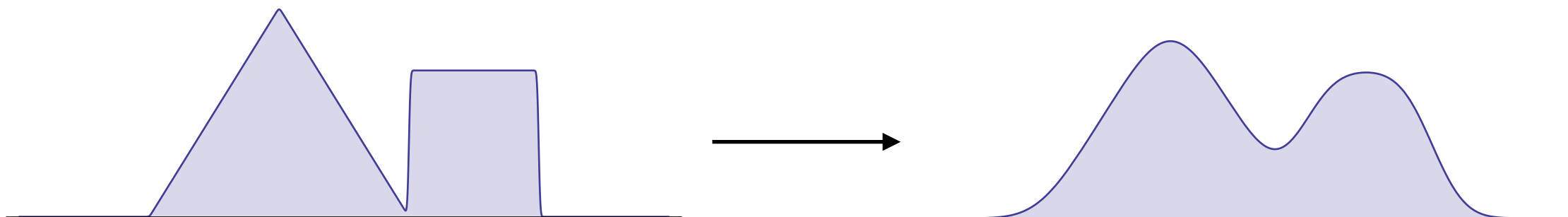
$$\text{Tr}(\mathcal{N}(\rho) e^{i\phi(f)}) = e^{-\frac{1}{2}(f, (X^\dagger AX + Y)f)}$$

$$A \longrightarrow X^\dagger AX + Y$$

$$(Xf)(x) \propto \int e^{-\frac{(x-y)^2}{2\sigma^2}} f(y) dy$$

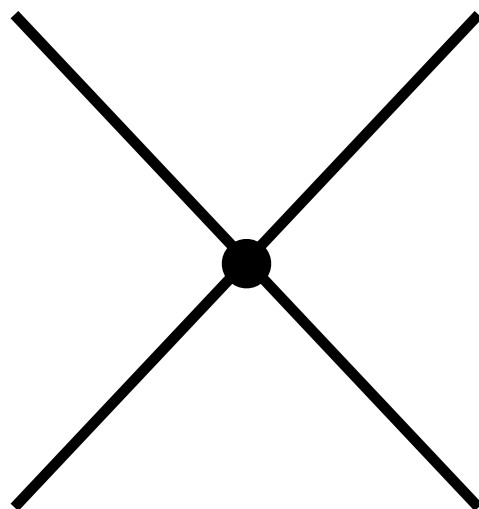


$$(Xf)(x) \propto \int e^{-\frac{(x-y)^2}{2\sigma^2}} f(y) dy$$

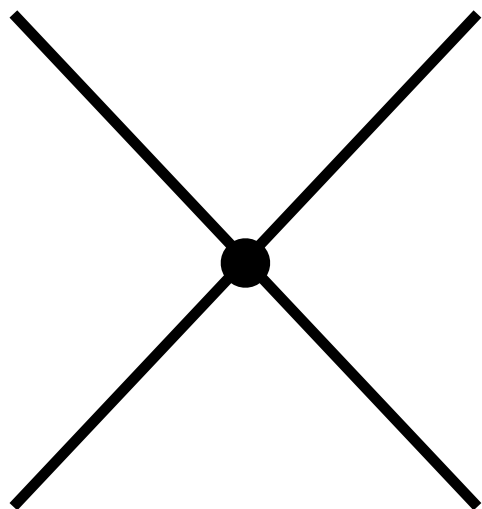


$$Y = \begin{pmatrix} y_{\phi}^2 I & 0 \\ 0 & y_{\pi}^2 I \end{pmatrix}$$

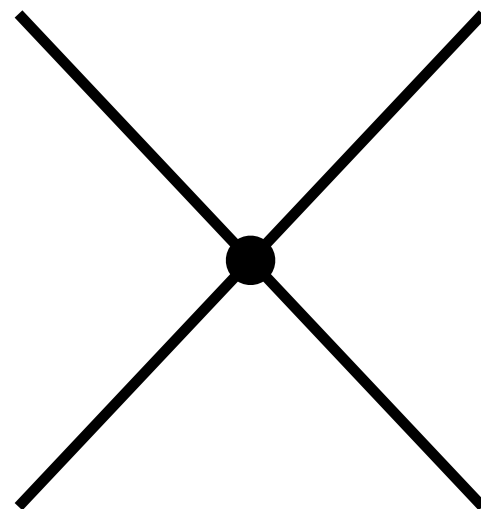
$$V = \int \phi(x)^4 dx$$



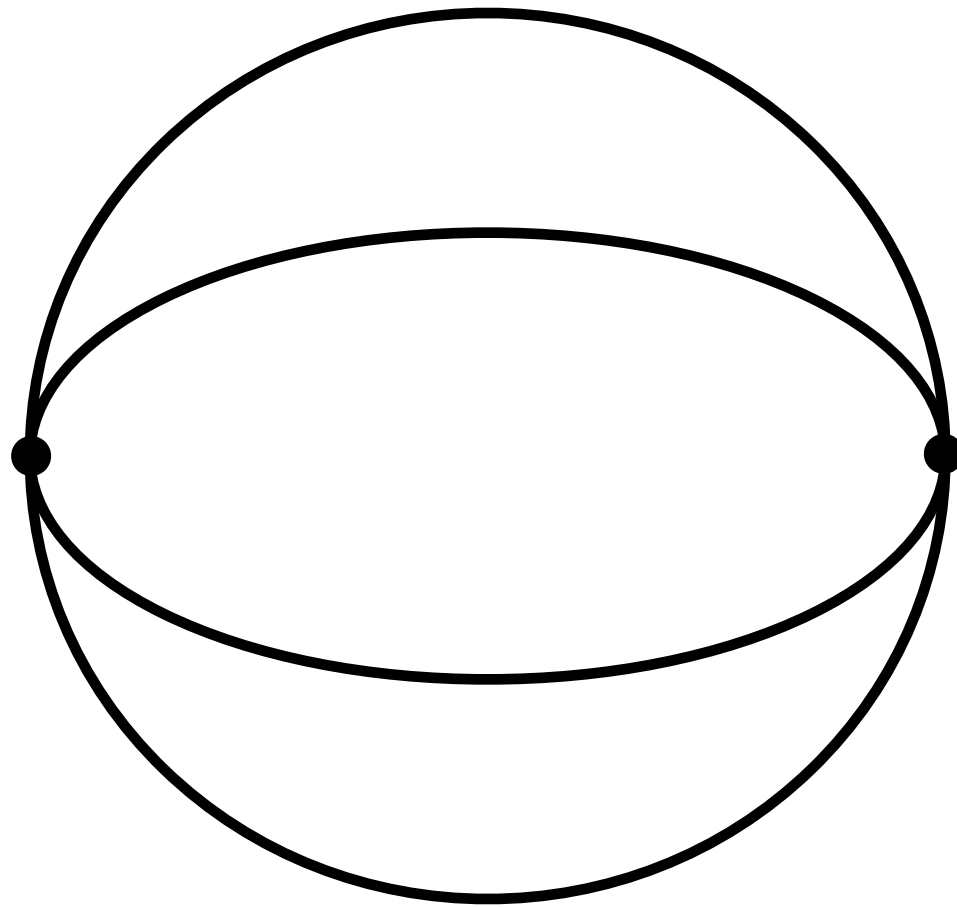
$V$



$V$



$d(V) =$



classical,  $\mathcal{N} = \text{id}$

$$\text{————} = A$$



classical

$$\underline{\hspace{2cm}} = AX(X^\dagger AX + Y)^{-1} X^\dagger A$$

quantum,  $y_\phi y_\pi \gg \hbar$

quantum,  $y_\phi y_\pi \gg \hbar$

$$\text{————} = \left(A + \frac{i}{2}\Delta\right) X B_\beta^{-1} X^\dagger \left(A + \frac{i}{2}\Delta\right)$$

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$$\begin{array}{c} \diagup \\ \diagdown \end{array} \bullet \begin{array}{c} \diagdown \\ \diagup \end{array} = \int_0^\beta dt \begin{array}{c} \diagup \circ \diagdown \\ \diagdown \circ \diagup \end{array} dt$$

quantum,  $y_\phi y_\pi \gg \hbar$

$$\text{---} = \left(A + \frac{i}{2}\Delta\right) X B_\beta^{-1} X^\dagger \left(A + \frac{i}{2}\Delta\right)$$

$$\text{X} = \int_0^\beta dt \text{X}(t)$$

$$\text{---}(t) = \text{imaginary time evolution in phase space}$$

$$d(\int \phi^4)$$

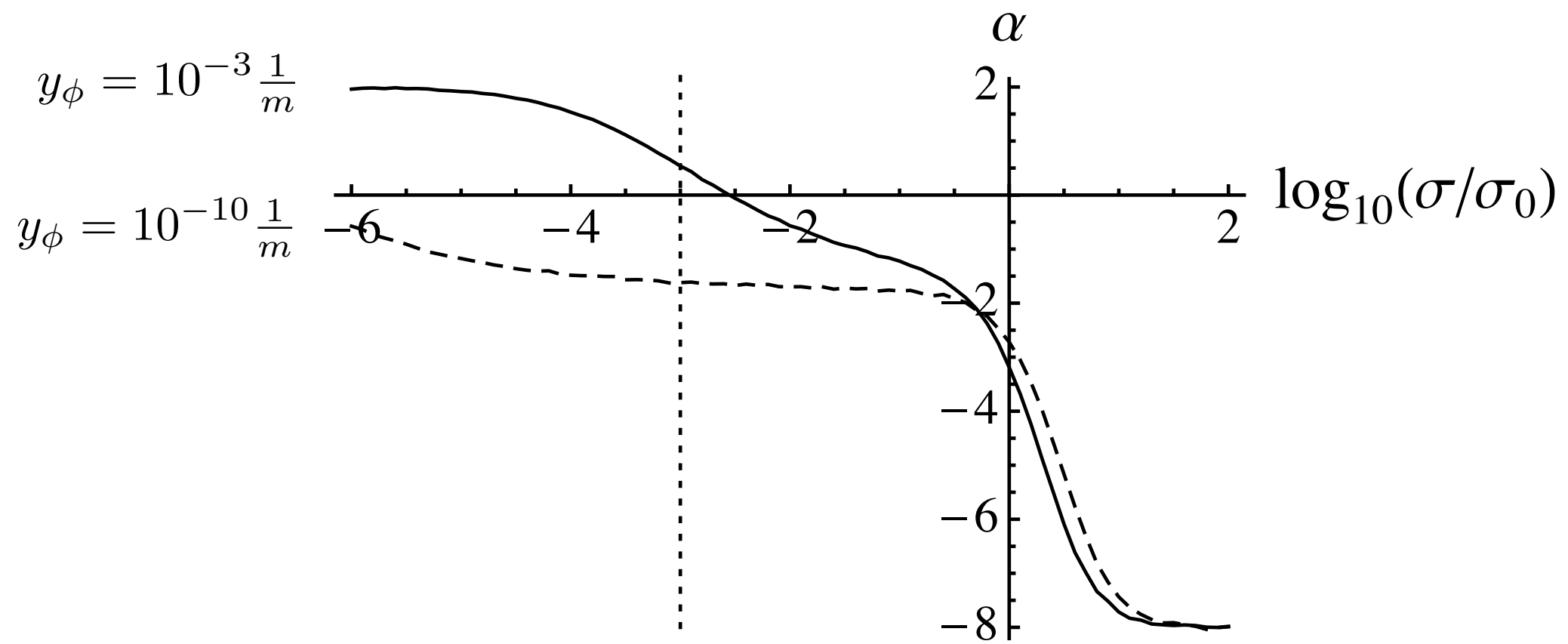
$$\sigma^D d(\int \phi^4)$$

$$\alpha = \frac{\partial \log(\sigma^D d(\int \phi^4))}{\partial \log \sigma}$$



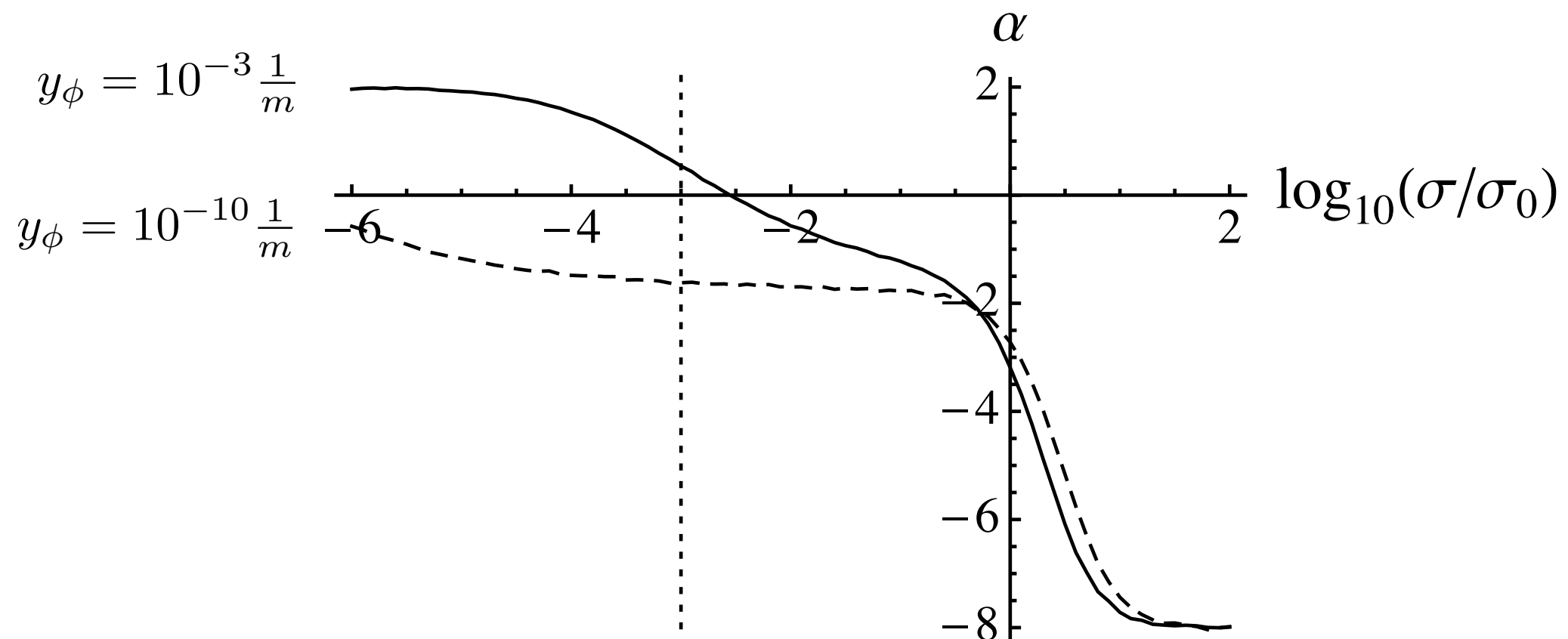
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4+1 dimensions



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*Thanks!*